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Detection of drift-kinetic vortices as a manifestation of Alfvénic turbulence

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<u>Outline</u>

Keywords: turbulence, coherent structures, space plasma measurements

- Phenomenology of turbulence
- Hydrodynamic and MHD turbulence
- MHD turbulence in the solar wind
- Turbulence in the polar cusp
- Drift-wave turbulence and coherent structures
- Analytical solutions, drift-kinetic Alfvén vortices
- In situ detection of drift-kinetic electromagnetic vortices in the cusp

Classical picture of turbulence

Richardson (1922) cascade

- •Turbulent (Richardson) cascade:
- •Energy goes from larger scales to smaller scales



•General classical picture:

•Energy fed at some large scale lo

•Dissipationless transfer towards smaller scales in the inertial range

•Dissipation into heat at the dissipation scale

based on Richardson's cascade picture

•Assumptions: homogeneous, isotropic, incompressible turbulence

Kolmogorov (1941) spectrum

•Statistical and dimensional arguments



Hydrodynamical turbulence





Vortex street

Re = 10,000

Hydrodynamics vs MHD

Fluid turbulence is contained in the Navier-Stokes equation

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla)\mathbf{v} = -\frac{1}{\rho}\nabla p + \frac{\mu}{\rho}\nabla^2 \mathbf{v}.$$

MHD turbulence in the MHD equations

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla)\mathbf{v} = -\frac{1}{\rho}\nabla p + \frac{1}{4\pi\rho}(\nabla \times \mathbf{B}) \times \mathbf{B} + \nu\nabla^2 \mathbf{v}$$

$$\frac{\partial \mathbf{B}}{\partial t} = \boldsymbol{\nabla} \times (\mathbf{v} \times \mathbf{B}) + \mu \nabla^2 \mathbf{B}.$$

Magnetic field complicates the dynamics

In magnetized plasmas, the ambient magnetic field decouples \parallel and \perp dynamics

Kraichnan-Iroshnikov (1965): slower energy transfer rate

quasi-two-dimensional turbulence

suppresses k-spectrum: not Kolmogorov scaling any more

The inverse cascade and coherent structures

Direct cascade: $k_{in} \rightarrow large$ wavenumbers (smaller scales)

Inverse cascade: $k_{in} \rightarrow small$ wavenumbers (larger scales)

General picture: inverse cascade towards larger scales – self-organization into large scale coherent structures

Details: Nonlinearities of vector type cause waves to self-interact with subsequent trapping into vortices

MHD turbulence in the solar wind

Energy injection – solar wind (supersonic) $\beta \approx 5$, $v_{ph} = v_A$ (Alfvénic)

Taylor's hypothesis (Doppler shift): $\omega = kv_{sw}$





Kolmogorov spectrum with index = -5/3 (fluid like)

Close to "dissipation scale", Alfvén waves --> Kinetic Alfvén waves

Kinetic Alfvén waves --> Heating?

Bale et al. *Measurement of the Electric Fluctuation Spectrum of* Magnetohydrodynamic Turbulence, PRL, 94, 215002-1 (2005)

Solar wind Dayside cusp Bow shock

Goldstein Turbulence on a small scale, Nature, 436, 782 (2005).

High-altitude cusp

 $m_e/m_i << \beta << 1$

Strongly inhomogeneous

Kinetic Alfvén waves and ion cyclotron waves close to dissipation scale

MHD inadequate ---> Drift-wave Alfvénic turbulence

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Nykyri et al. Ion cyclotron waves in high altitude cusp: CLUSTER observations at varying spacecraft separations, GRL, 30, 2263 (2003).

Sundkvist et al. *Wave characteristics near proton* gyrofrequency in high-altitude cusp, Ann. Geophys., 23, 983-995 (2005).

Drift-wave turbulence and coherent structures

A model equation for low-frequency drift-wave plasma turbulence is the Charney-Hasegawa-Mima (CHM) equation

$$\frac{\partial}{\partial t} \left(\nabla^2 \phi - \phi \right) - \left[\nabla \phi \times \hat{\boldsymbol{z}} \cdot \nabla \right] \left[\nabla^2 \phi - \ln \left(\frac{n_0}{\omega_{cl}} \right) \right] = 0$$

The CHM-equation admit both wave and coherent (vortex) solutions and is appropriate to describe low-frequency drift-wave turbulence





Horton & Hasegawa Quasi-two-dimensional dynamics of plasmas and fluids, Chaos, 4, 227-251 (1994)

Generalization of CHM-equation

Can be derived from two-fluid equations:

$$\partial_{\mathbf{t}} \mathbf{v}_{\mathbf{j}} + (\mathbf{v}_{\mathbf{j}} \cdot \boldsymbol{\nabla}) \mathbf{v}_{\mathbf{j}} = \frac{q}{m_{\mathbf{j}}} (\mathbf{E} + v_{\mathbf{j}} \times \mathbf{B}) - \frac{1}{m_{\mathbf{j}}n} \boldsymbol{\nabla} p_{\mathbf{j}}$$

For the parameters in the cusp we can assume

$$\mathbf{B} = B_0 \hat{z} + \delta \mathbf{B}_\perp$$
$$\mathbf{E}_\perp = -\boldsymbol{\nabla}_\perp \phi$$

Momentum fluid equations lead to the drift velocities

$$\begin{split} \mathbf{v}_{i,\perp} &\approx \frac{c}{B_0} \hat{z} \times \nabla_{\perp} \phi + \frac{cT_i}{eB_0 n_i} \hat{z} \times \nabla n_i + \frac{c}{B_0 \omega_{ci}} (\partial_t + \mathbf{v}_i \cdot \nabla) \mathbf{E}_{\perp} \\ &\mathbf{v}_e \approx \frac{c}{B_0} \hat{z} \times \nabla_{\perp} \phi - \frac{cT_e}{eB_0 n_e} \hat{z} \times \nabla n_e + v_{ez} \left(\hat{z} + \frac{\mathbf{B}_{\perp}}{B_0} \right) \end{split}$$

which together with quasi-neutrality and the momentum and continuity equations gives the system

$$\begin{pmatrix} \frac{\partial}{\partial t} + \mathbf{v}_{\Xi} \cdot \nabla \end{pmatrix} p_{i} - \frac{\dot{\mathbf{c}}}{\mathrm{d}t} p_{i} - 0$$

$$\frac{\mathrm{d}}{\mathrm{d}t} \left(n - \frac{en_{0}}{m_{i}\omega_{\mathrm{c}i}^{2}} \nabla_{\perp}^{2} \varphi \right) - \frac{e}{m_{i}^{2}\omega_{\mathrm{c}i}^{3}} \nabla \cdot \{ p_{i}, \nabla_{\perp} \varphi \} = 0$$

$$\{ A, B \} = \frac{\partial A}{\partial x} \frac{\partial B}{\partial y} - \frac{\partial A}{\partial y} \frac{\partial B}{\partial x}$$

$$ne\hat{z} \cdot \nabla \phi + \frac{ne}{c} \frac{\partial A_{z}}{\partial t} - \frac{m_{e}c}{4\pi e} \frac{\mathrm{d}}{\mathrm{d}t} \nabla_{\perp}^{2} A_{z} = \hat{z} \cdot \nabla p_{e}$$

$$\frac{\partial n}{\partial t} + \mathbf{v}_{\mathbf{E}} \cdot \nabla n + \frac{c}{4\pi e} \hat{z} \cdot \nabla \nabla_{\perp}^{2} A_{z} = 0$$

This system generalizes and reduce to CHM for $A_z = 0$, p = 0. The system describes both drift-Alfvén waves and vortices

Vortex solutions

Solutions exist and can be found by using the Larichev-Reznik ansatz:

Assume function of x and $\eta = y + \alpha z - ut$

Transformed system have solution

$$\phi = [ur + c_1 J_1(k_1 r) + c_2 J_1(k_2 r)] \cos \theta \qquad (r \le r_0)$$

$$\phi = c_5 K_1(sr) \cos \theta \qquad (r \ge r_0)$$

Dipole vortex

In addition vortex chains (cat-eyes) and monopole vortices can exist

Questions: Do such coherent vortices exist in magnetoplasmas, and if they do, what are their charcteristics?

Shukla, P. K., Yu, M. Y. & Varma, R. K., Phys. Lett. A 109, 322-324 (1985).
Petviashvili, V. I. & Pokhotelov, O. A., JETP Lett. 42, 54-56 (1985).
Shukla, P. K., Yu, M. Y. & Stenflo, L., Phys. Rev. A 34, 3478-3480 (1986).
Liu, J. & Horton, W., J. Plasma Phys. 36, 1-24 (1986).
Kaladze et al. Plasma Phys. and Cont. Fusion 29, 589 (1987)
Chmyrev, V. M. et al., Physica Scripta 38, 841-854 (1988).

Previous single-spacecraft observations of coherent structures

- Chmyrev et al (1988) suggested that they observed drift-kinetic vortices in the plasma sheet boundary layer
- To be sure one needs to separate temporal from spatial phenomena



Fig. 5. Variations of the horizontal component of the electric field vector in the frequency range 0.12-0.8 Hz along the satellite orbit at the second interval in Fig. 1.

Chmyrev et al. *Alfvén vortices and related phenomena in the Ionosphere and the Magnetosphere, Physica Scripta, 38, 841-854 (1988).*

Multi-satellite detection of DKA vortices

Cusp crossing 2002-03-09



Northern outbound crossing

Northward IMF -> high-latitude reconnection

Turbulent boundary layer formed

Hodograms along trajectory suggest vortex motion

Cluster --> Multi-point measurements D. Sundkvist, ICTP, 22 Sept 2005 (davids@irfu.se)



Sundkvist, Krasnoselskikh, Shukla et al. In Situ Multi-Satellite Detection of Coherent Vortices as a Manifestation of Alfvenic Turbulence, Nature, 436, 825-828 (2005).

Detection of DKA vortices – Multipoint timeseries analysis

New point of view: look for coherent structures by examining time-series from several spacecraft simultaneously

Rotation of field vectors spatially coherent

In accord with vortex convected past the spacecraft

Vortex radius approximately 2-6 $\rho_{\rm p}$

Sundkvist et al. In Situ Multi-Satellite Detection of Coherent Vortices as a Manifestation of Alfvenic Turbulence, Nature, 436, 825-828 (2005).





Detection of DKA vortices – Comparison with theory



Model of vortex turbulence in cusp: isolated DKA vortices

In accord with observations

spectra

rotation of field vectors

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Sundkvist et al. In Situ Multi-Satellite Detection of Coherent Vortices as a Manifestation of Alfvenic Turbulence, Nature, 436, 825-828 (2005). Importance of findings:

Drift-wave Alfvénic turbulence exhibit coherent structures

Cross-field anomalous transport

New way of transporting mass and energy through boundary layers

<u>Summary</u>

- High-altitude cusp boundary layer:
- low beta (kinetic regime) drift-wave Alfvénic turbulence
- Appropriate equations are generalizations of Charney-Hasegawa-Mima equation
- Strong drift-turbulence: drift-waves and coherent vortices (owing to selfinteraction among kinetic Alfvén waves via vector nonlinearities)
- Analytical solutions exist: Drift-Kinetic electromagnetic Alfvén vortices (DKA vortices)
- Multi-point measurements by Cluster four spacecraft confirms existence of small(large) scale coherent structures (vortices)
- Observations of coherent structures can be described as DKA vortices and vortex chains
- Cross-field transport