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AUTUMN COLLEGE ON PLASMA PHYSICS

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Plasma Accelerators

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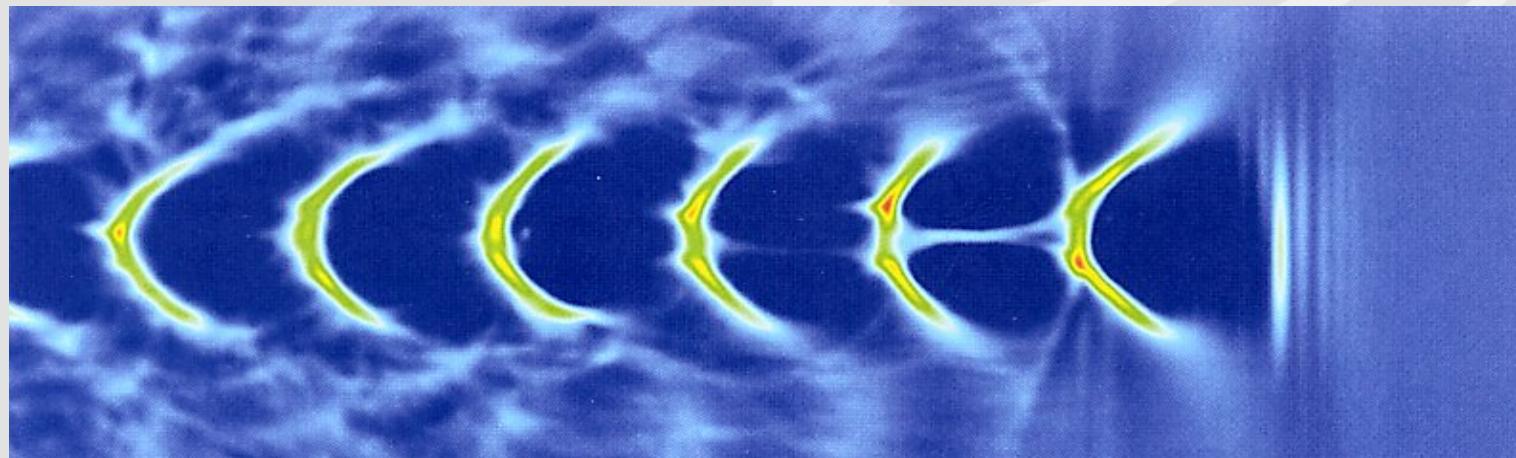
Abdus Salom ICTP Plasma Physics , 2005
Rosenbluth Symposium



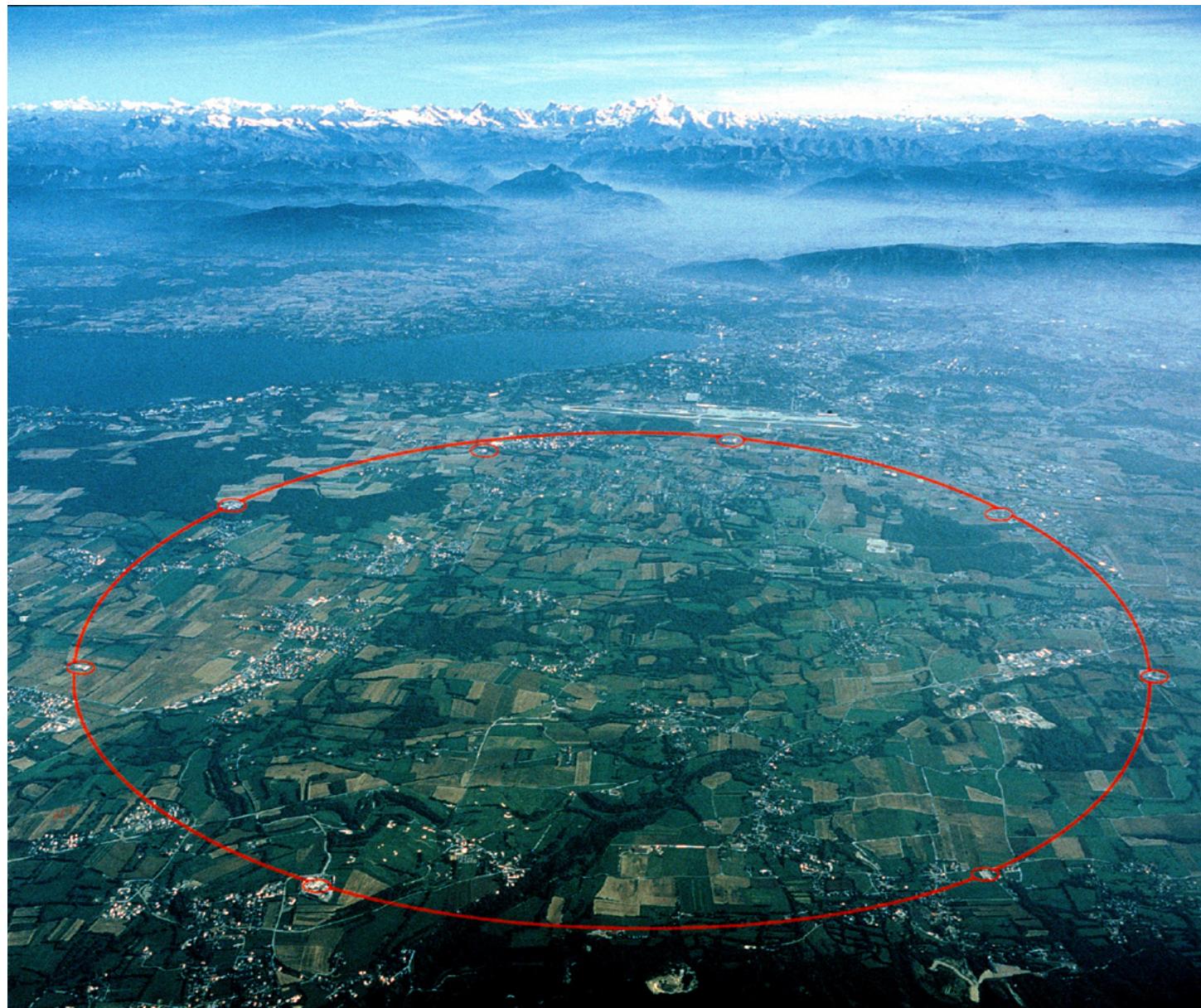
Plasma Accelerators

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Centre for Fundamental Physics



CERN – LEP





Marshall Rosenbluth contribution enormous!

Parametric Instabilities:

Raman, Brillouin, filamentation, self-modulational, scattering.

Seminal papers on role of inhomogeneity PRLs
(All important for laser fusion)

Plasma Accelerators:

Relativistic Plasma Wave

- Saturates by relativistic de-tuning
- Rosenbluth Lui saturation PRL (1972)

Nonlinear Vacuum:

Photon-photon scattering PRL (1970)

Laser Plasma Accelerators



Conventional Accelerators

- Limited by peak power and breakdown
- 20-100 MV/m
- Large Hadron Collider (LHC) -- 27km, 2010
- Plans for “Next” Linear Collider (NLC) -- 100km ?

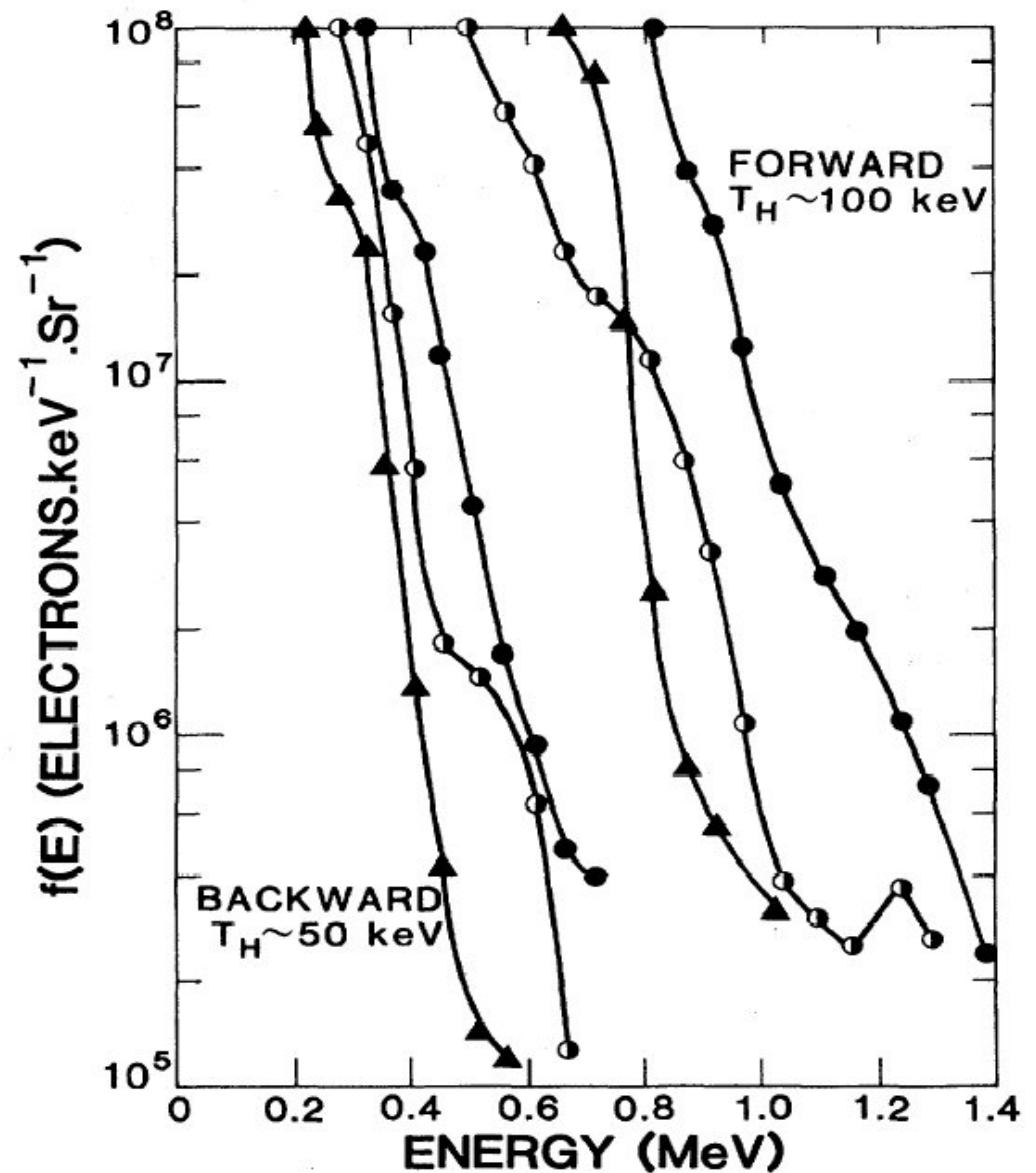
Plasma

- No breakdown limit
- 10-100 GV/m

High Energy Particles

- Relativistic Plasma Wave Acceleration
- The problem is to generate large amplitude plasma wave travelling with a velocity close to the speed of light c
- 4 Approaches
 1. Plasma Beat Wave
 2. Laser Plasma Wakefield
 3. Self-Modulated Laser Wakefield (RFS)
 4. Electron Beam Plasma Wakefield

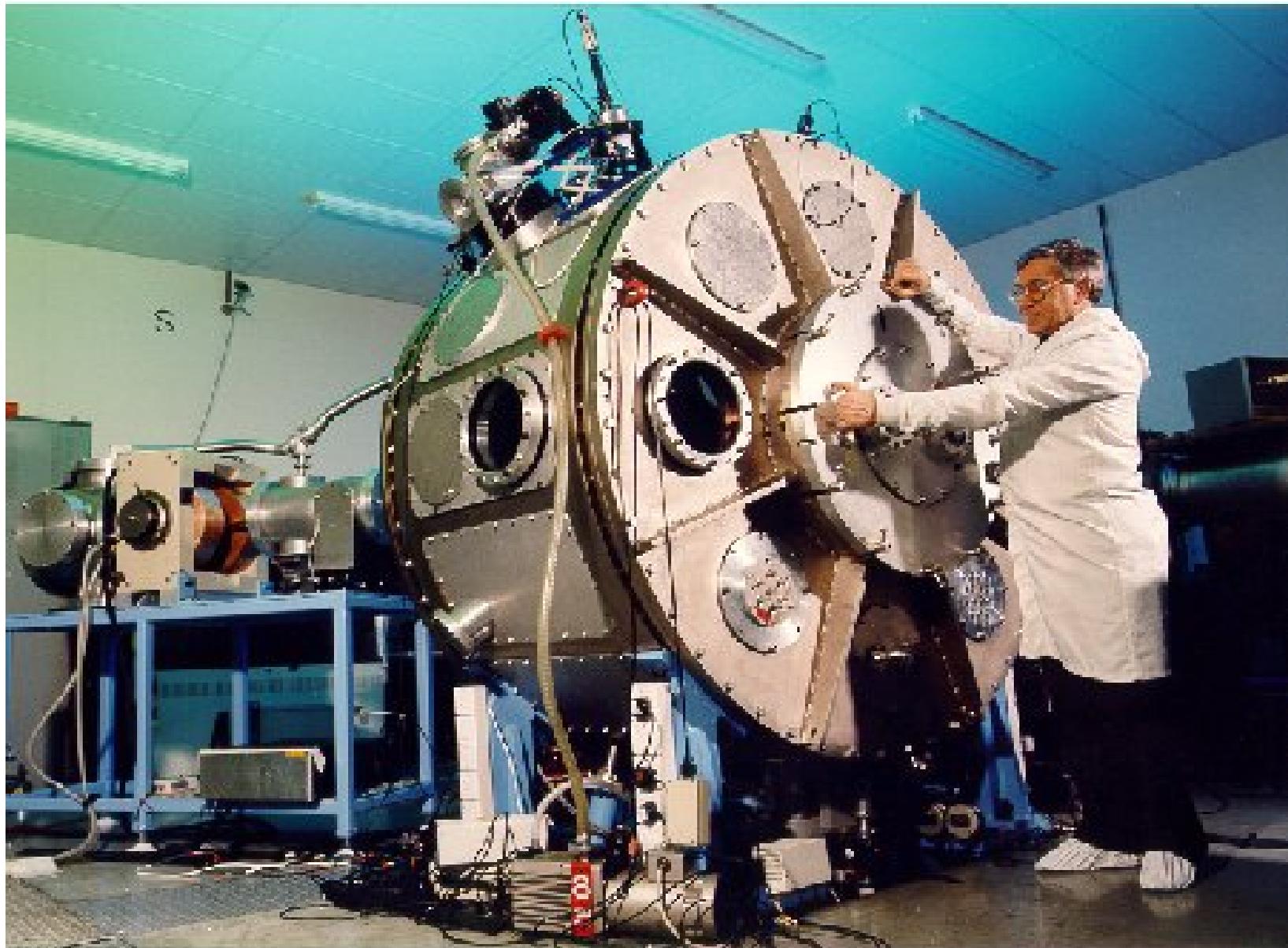
Experimental electron energy distributions in the forward and backward directions. Three different shots are represented.



Drivers for Plasma Based Accelerators

- Lasers – Terawatt, Petawatt Compact Lasers 10^{12} – 10^{15} Watts already exist.
 - Some with high rep. Rates *i.e.* **10 Hz**.
 - Capable of 10^{19} – 10^{21} *Watts/cm²* on target.
 - Future $\sim 10^{23}$ *Watts/cm²* using OPCPA.
- Electrons Beams – Shaped electron beams such as the proposed Stanford/USC/UCLA experiment to generate $1GV/m$ accelerating gradient using the 30 – $50GeV$ beam in a 1 meter long Lithium Plasma.

Vulcan Target Area



Laser Plasma Accelerators

- The electric field of a laser in vacuum is given by

$$E_{\perp} = 30\sqrt{I} \text{ V/cm}$$

- For short pulse intense lasers,

$$P = 10 \text{ TW}, \lambda_0 = 1 \mu\text{m}, I = 1.6 \times 10^{18} \text{ W/cm}^2$$

$$E_{\perp} = 40 \text{ GV/cm}$$

- Unfortunately, this field is perpendicular to the direction of propagation and no significant acceleration takes place.
- The longitudinal electric field associated with electron plasma waves can be extremely large and can accelerate charged particles.

Plasmas

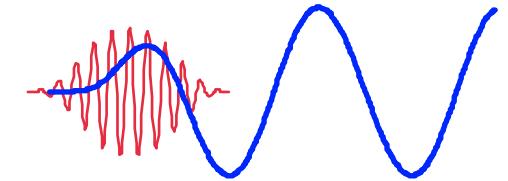
- Conventional accelerators limited by electrical breakdown of accelerating structures
- Plasmas are already broken down.
 - The accelerating fields limited only by plasma density.
- Plasmas can support longitudinal accelerating fields moving close to the speed of light; Relativistic electron plasma waves.
- Lasers easily couple to plasmas and can generate relativistic electron plasma waves.

Laser Plasma Accelerators

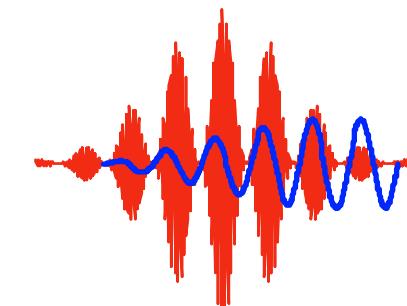
- Large accelerating fields $\sim 1 \text{ GeV/cm}$
- No electrical breakdown limit

Laser Wakefield Acceleration

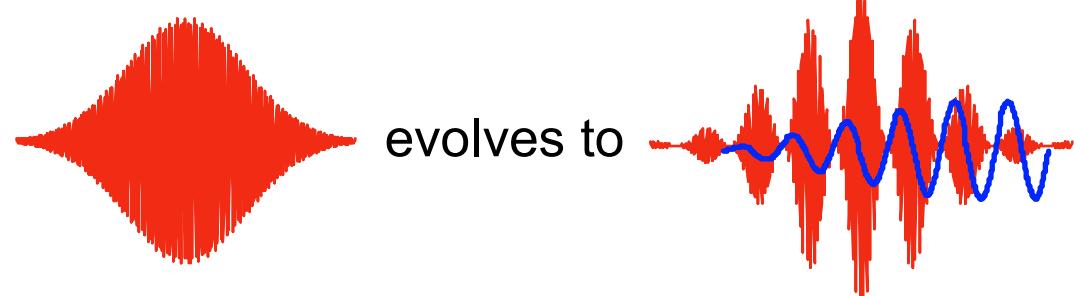
- Laser Wake Field Accelerator(LWFA)
A single short-pulse of photons



- Plasma Beat Wave Accelerator(PBWA)
Two-frequencies, i.e., a train of pulses



- Self Modulated Laser Wake Field Accelerator(SMLWFA)
Raman forward scattering instability

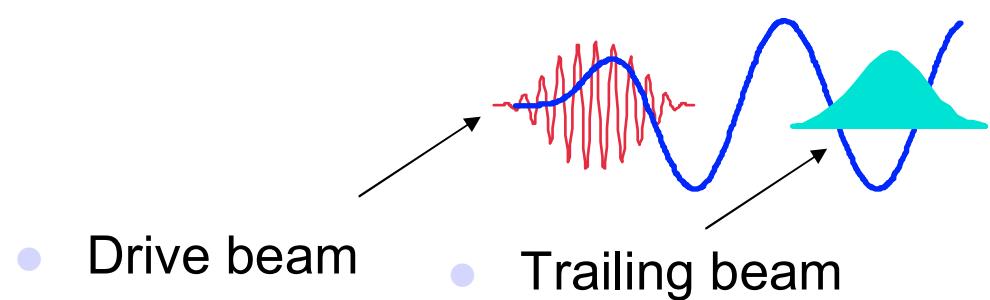


Concepts for Plasma based Accelerators

- Plasma Wake Field Accelerator(PWFA)
A high energy electron bunch

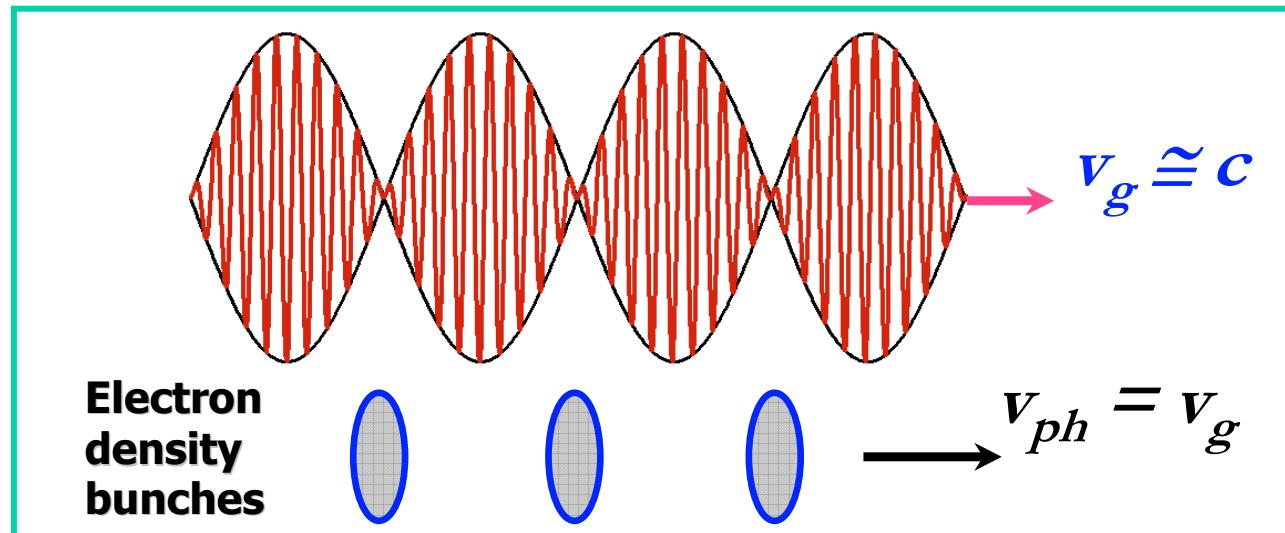


- Laser Wake Field Accelerator(LWFA, SMLWFA, PBWA)
A single short-pulse of photons



Plasma Beat Wave Accelerator (PBWA)

- In the Plasma Beat Wave Accelerator (PBWA) a relativistic plasma wave is resonantly excited by the “ponderomotive” force of two lasers separated by the plasma frequency ω_p .
- The two laser beams beat together forming a modulated beat pattern in the plasma.



- For relativistic plasma wave the accelerating field $E_{||}$ is given by

$$E_{||} = \epsilon \sqrt{n_0} \quad \text{V/cm}$$

ϵ is the fractional electron density bunching, n_0 is the plasma density. For $n_0 = 10^{18} \text{ cm}^{-3}$, $\epsilon = 10\%$ $\Rightarrow E_{||} = 10^8 \text{ V/cm}$

Plasma Beat Wave

Relativistic plasma wave driven by beating 2 lasers in a plasma

$$\omega_1 - \omega_2 \cong \omega_p \quad \text{energy}$$

$$\underline{k}_1 - \underline{k}_2 \cong k_p \quad \text{momentum}$$

For $\omega_1, \omega_2 \gg \omega_p$ i.e. $\omega_1 = 10\omega_p$ $\omega_2 = 9\omega_p$

Then
$$\left. \begin{array}{l} \underline{k}_1 - \underline{k}_2 \sim \Delta k \\ \omega_1 - \omega_2 \sim \Delta\omega \end{array} \right\} \frac{\Delta\omega}{\Delta k} = v_g$$

v_g is the group velocity of the laser beat pattern.

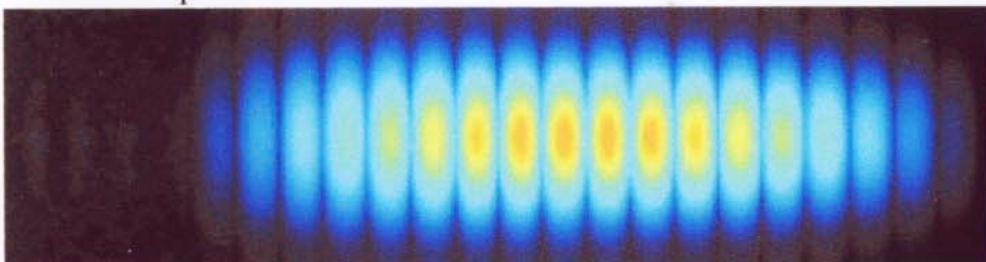
But $k_1 - k_2 \sim k_p$; $\omega_1 - \omega_2 \sim \omega_p$

$$\Rightarrow \frac{\omega_p}{k_p} = v_{ph} \equiv v_g \quad v_g = c \left(1 - \frac{\omega_p^2}{\omega_{1,2}} \right)^{1/2}$$

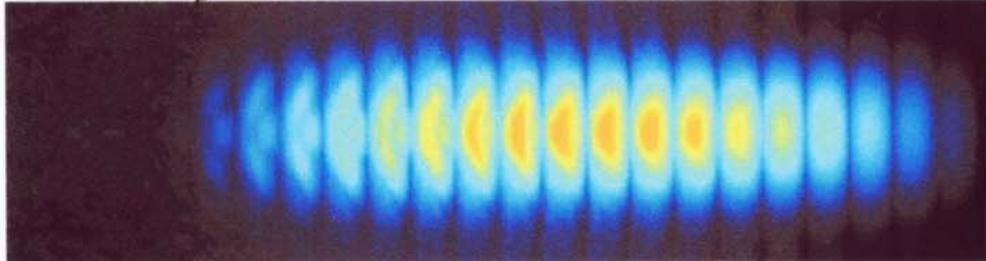
For $\omega_1, \omega_2 \gg \omega_p \Rightarrow v_g \approx c \Rightarrow \text{"Hence relativistic"}$

PBWA: Evolution of Laser Intensity and Accelerating Field

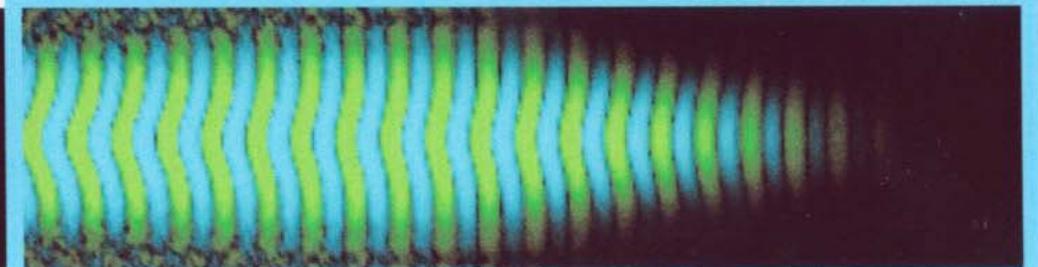
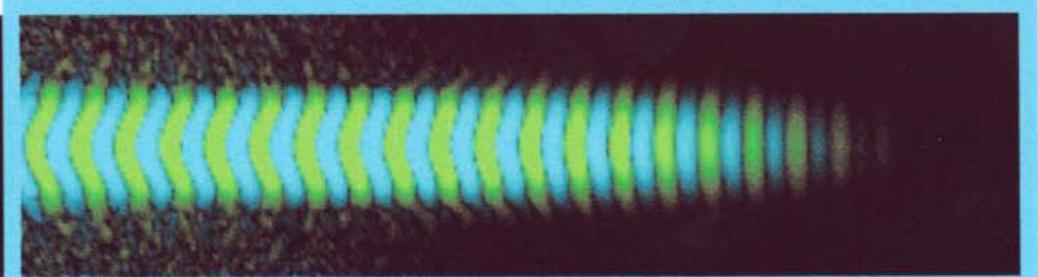
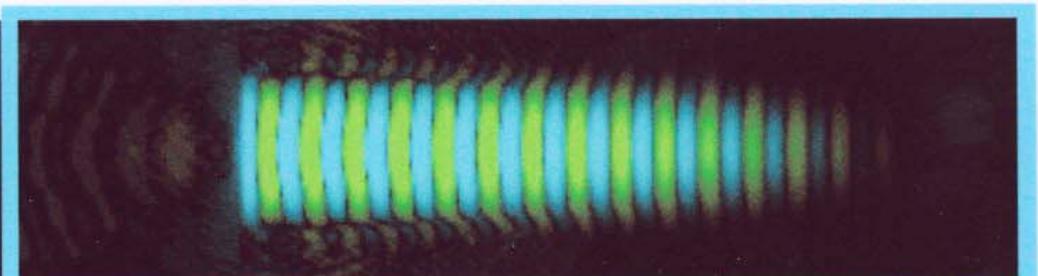
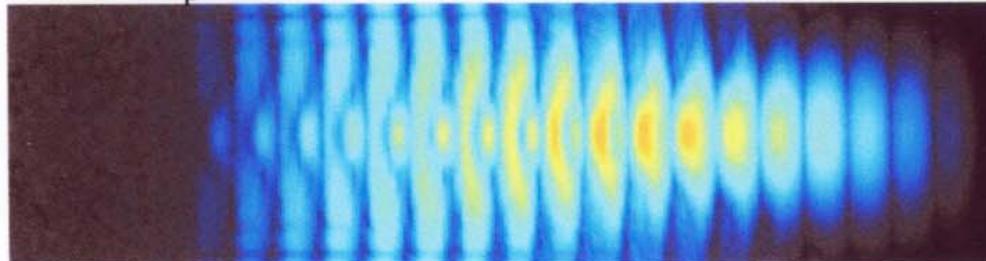
$t=111\omega_p^{-1}$



$t=222\omega_p^{-1}$



$t=489\omega_p^{-1}$



Energy Gain

- For $n_0 = 10^{18} \text{ cm}^{-3}$, $\varepsilon = n_1 / n_0 = 10\%$
- Gain in energy of electron ΔW

$$\Delta W = 2 \varepsilon \gamma^2 m_e c^2$$

- $\gamma = \omega_1 / \omega_p$ is the Lorentz factor
- For a neodymium laser, $\omega_1 / \omega_p \sim 30$, and $n_0 \sim 10^{18} \text{ cm}^{-3}$
- Maximum energy gain $\Delta W \approx eE_p l$
 l dephasing length

$$l = \frac{\lambda_p}{2} \frac{\omega^2}{\omega_p^2} = \frac{\lambda_p}{2} \gamma^2$$

$$eEl = 2\varepsilon\gamma^2 m_e c^2$$

$$\Delta W \approx 100 \text{ MeV}$$

The Relativistic Plasma Wave is described by

$$\left(\frac{\partial^2}{\partial t^2} + \omega_p^2 \right) \delta n = \frac{3}{8} \omega_p^2 \frac{\delta n^2}{n_0^2} \delta n - \frac{n_0}{2} \omega_p^2 \alpha_1 \alpha_2 e^{-i\delta t}$$

$$\alpha_j = \frac{eE}{m_e \omega_j c} \quad ; \quad \delta = \omega_1 - \omega_2$$

For $\alpha_1 = \alpha_2 = \text{constant}$

$$\frac{\delta n}{n_0} = \frac{\delta n_0(0)}{n_0} + \frac{1}{4} \alpha_1 \alpha_2 \omega_p t$$

Linear growth: However due to 1st term on RHS i.e. cubic non-linearity wave saturates before reaching wave breaking limit $\delta n/n \sim 1$; acts as nonlinear frequency shift.

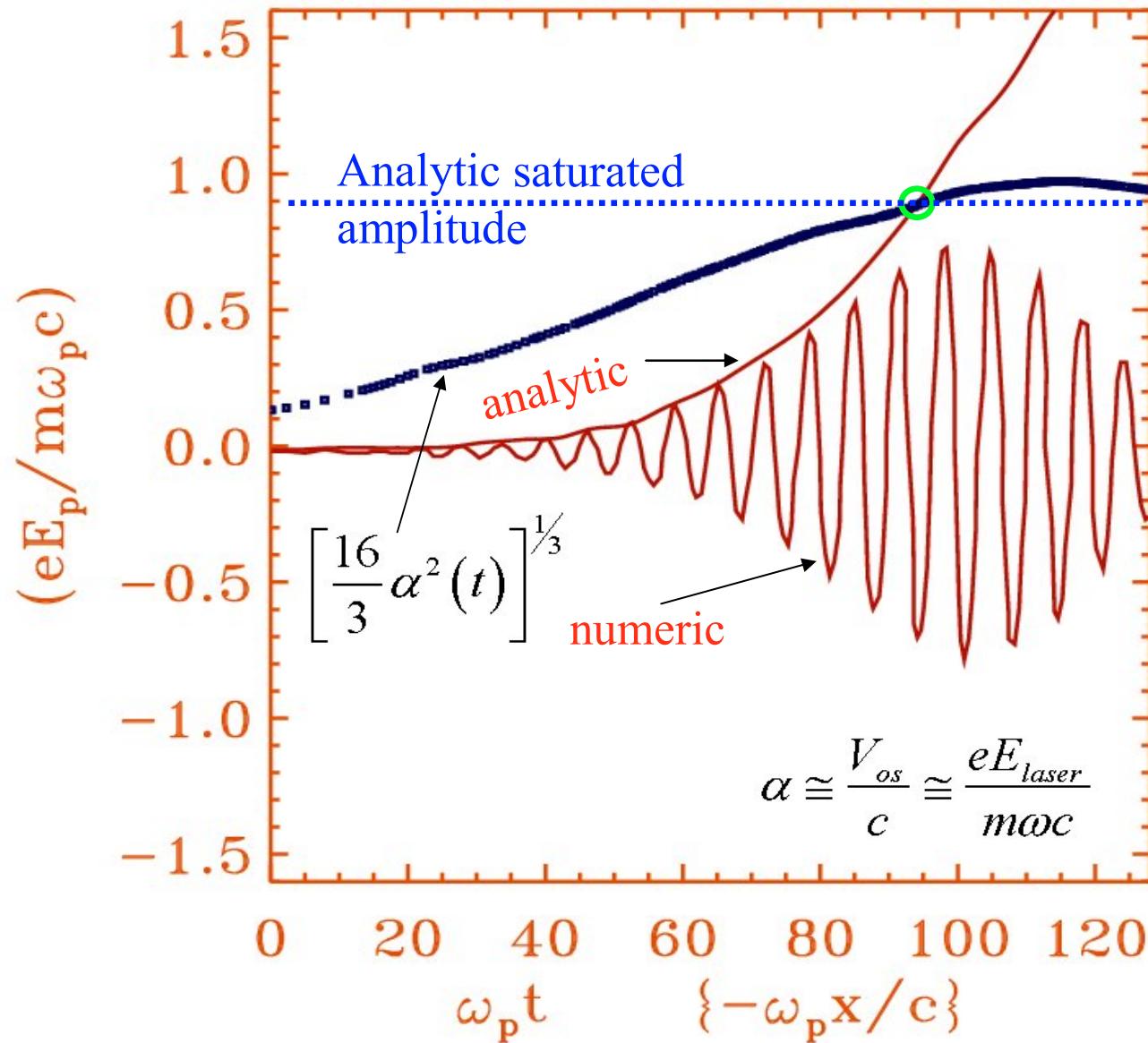
$$\frac{\delta n}{n} \max = \left(\frac{16}{3} \alpha_1 \alpha_2 \right)^{1/3} = \varepsilon$$

Relativistic mass increase of electrons reduces natural frequency

$$\omega_p^1 = \omega_p \left(1 - \frac{3}{8} \frac{v_{osc}^2}{c^2} \right)^{1/2} = \omega_p \left(1 - \frac{3}{8} \frac{\delta n^2}{n^2} \right)^{1/2}$$

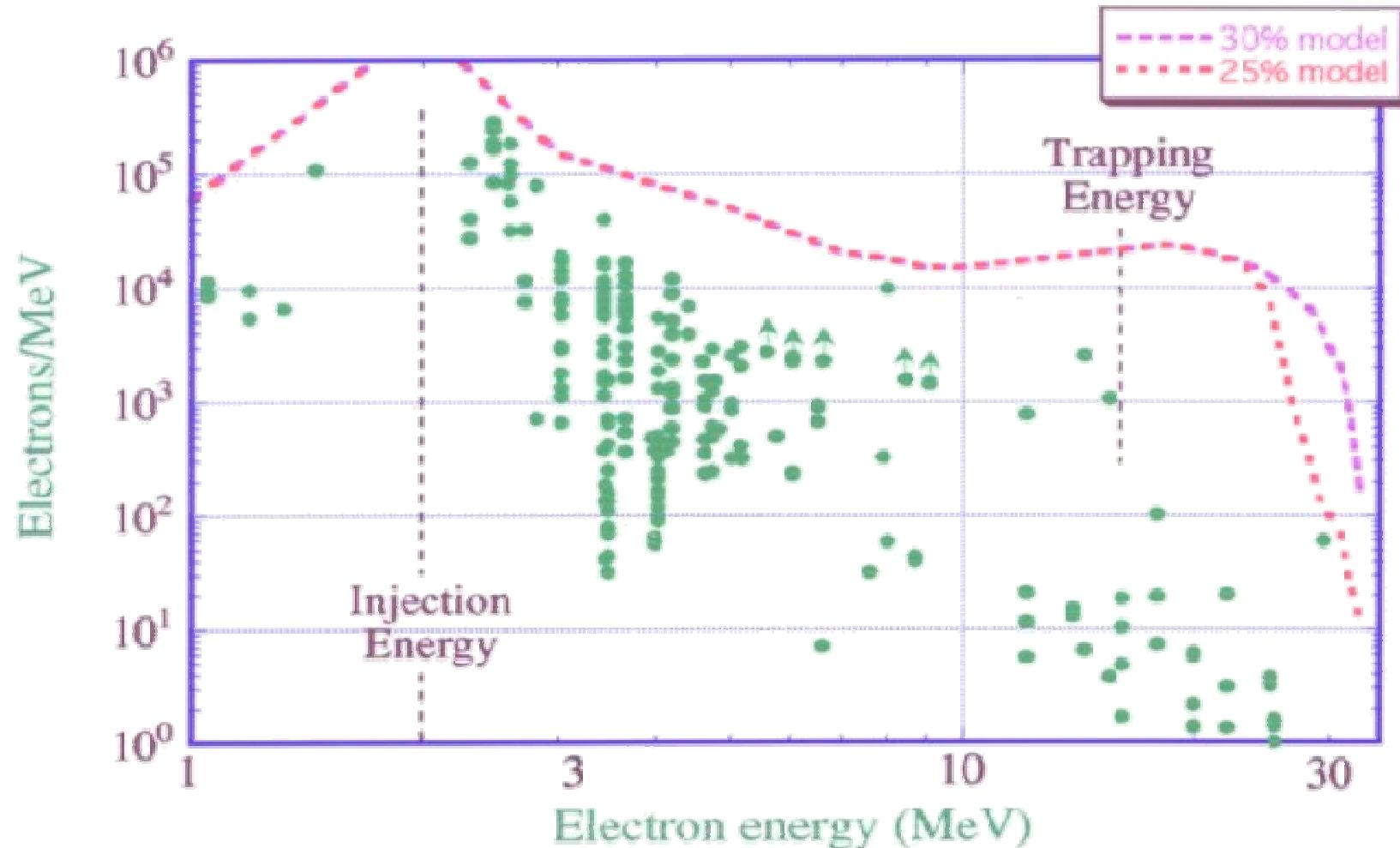
This results in $\omega_1 - \omega_2 \neq \omega_p \Rightarrow$ non-resonant interaction \Rightarrow saturation.

Beat Wave Growth



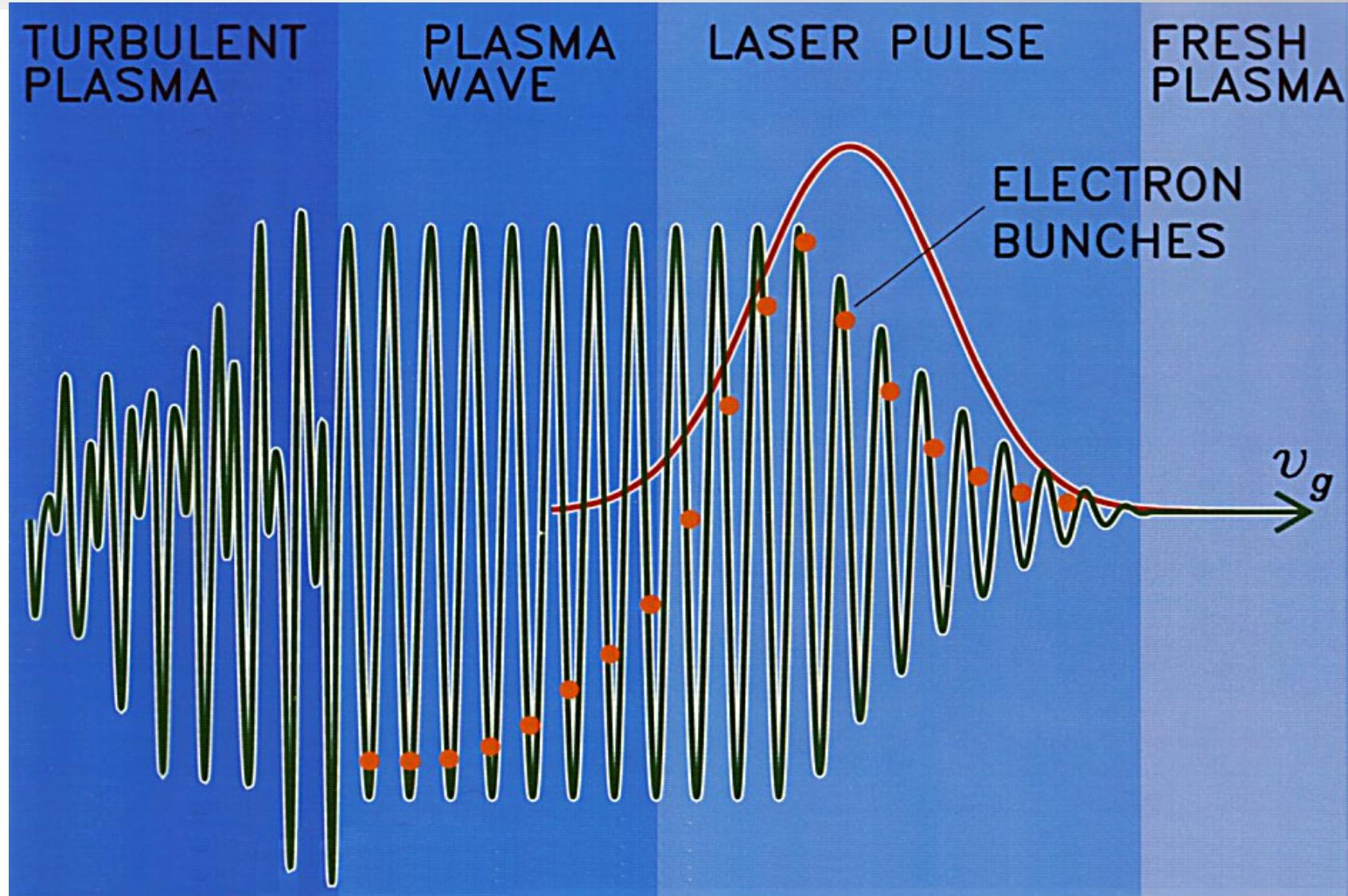
UCLA Beatwave Results

- Electrons injected from 0.3 GHz rf LINAC \Rightarrow train of pulses, <10ps duration



- 1% or 10^5 electrons are accelerated in the diffraction length of $\sim 1\text{cm}$.
- 2 \rightarrow 30 MeV Gradient of 3 GeV/m

Limitations of beatwave scheme



Ion dynamics becomes important: strong plasma turbulence driven by electron plasma waves destroys the accelerating wave structure.

Surfing The Waves!

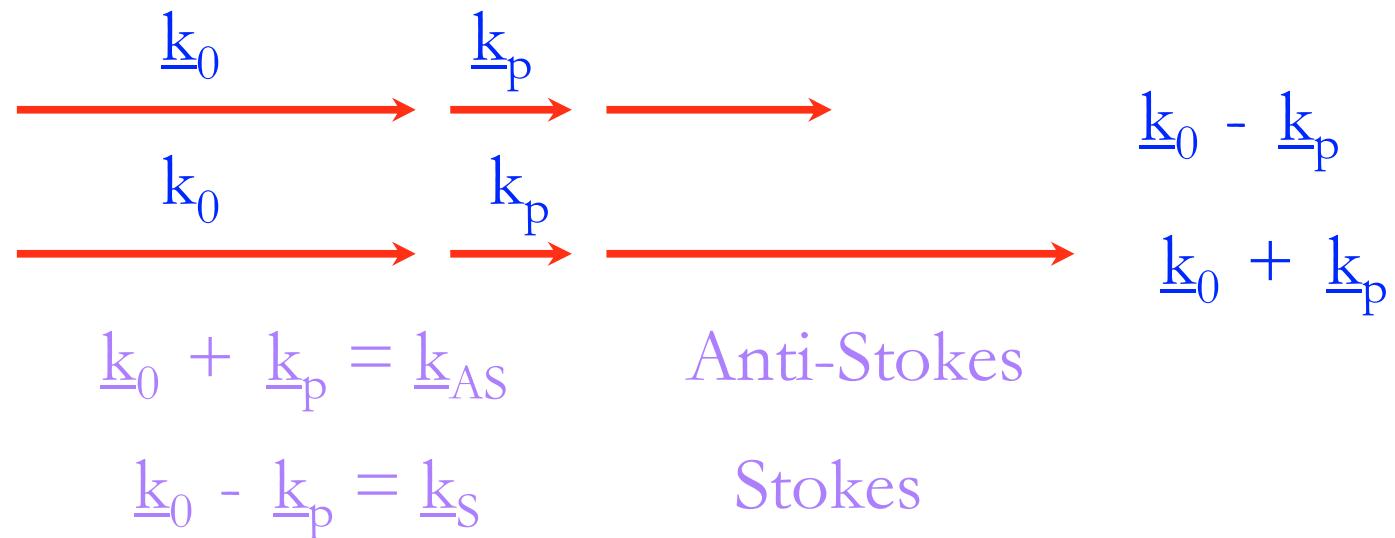
**At 10^{21} W.cm $^{-2}$,
electrons will be
accelerated to
beyond 100 MeV,
generating gamma
rays, proton beams
and exotic isotopes.**



Self-Modulated Wakefield

High intensity long pulse lasers (pulse length greater than plasma wave period) $\tau_p \gg 1/\omega_p$

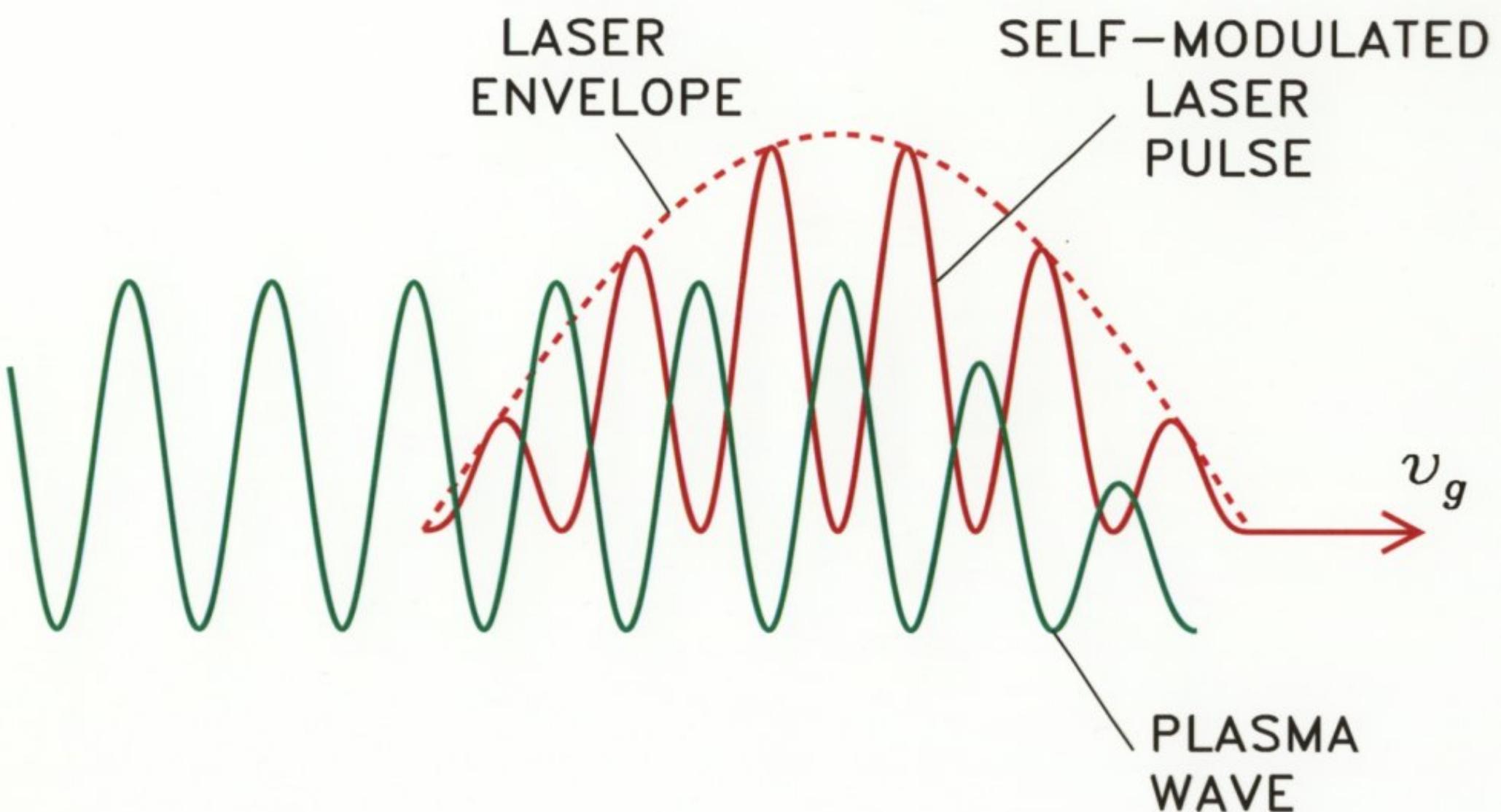
Break-up of a long pulse $L \gg k_p$ via Forward Raman Scattering or an Envelope-Self-Modulation instability.



Pulse Power > relativistic self-focusing

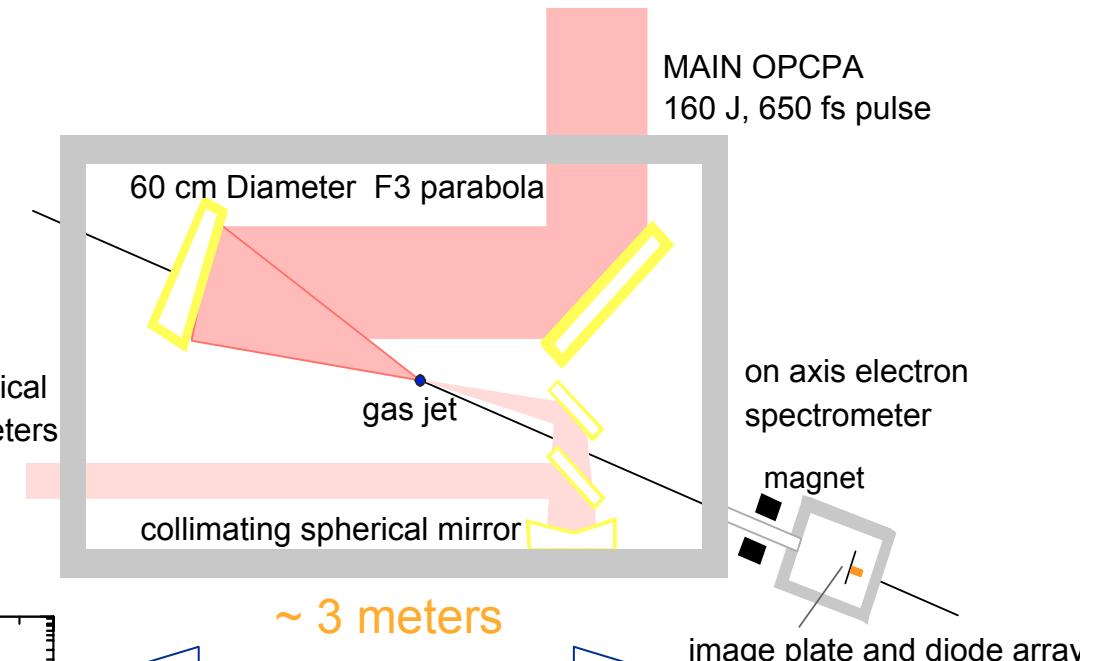
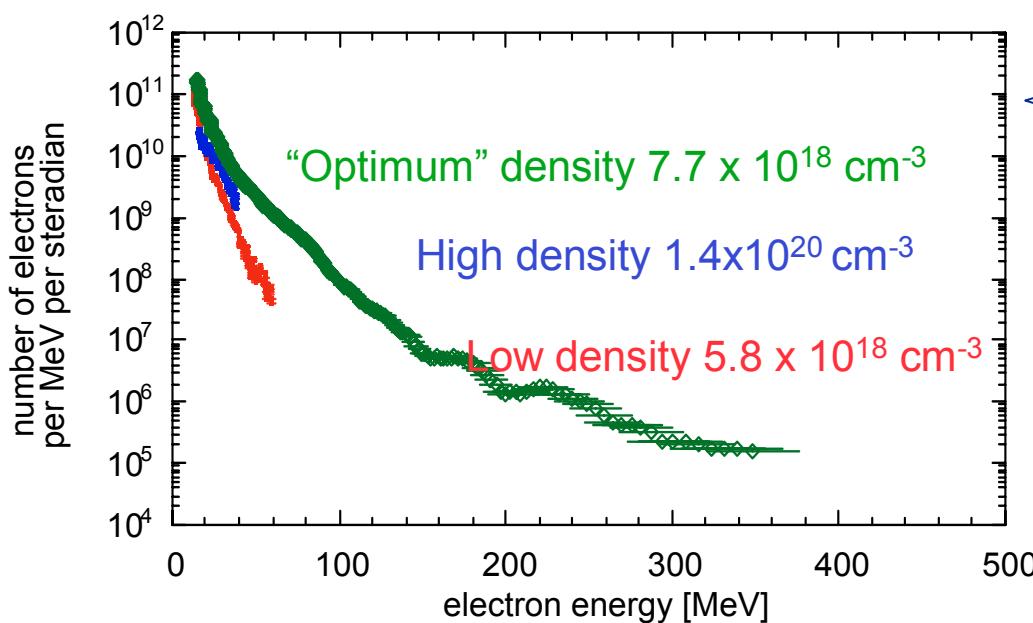
$$P > P_c = 17 \frac{\omega^2}{\omega_p^2} \text{ GW}$$

Self-Modulated Laser Pulse



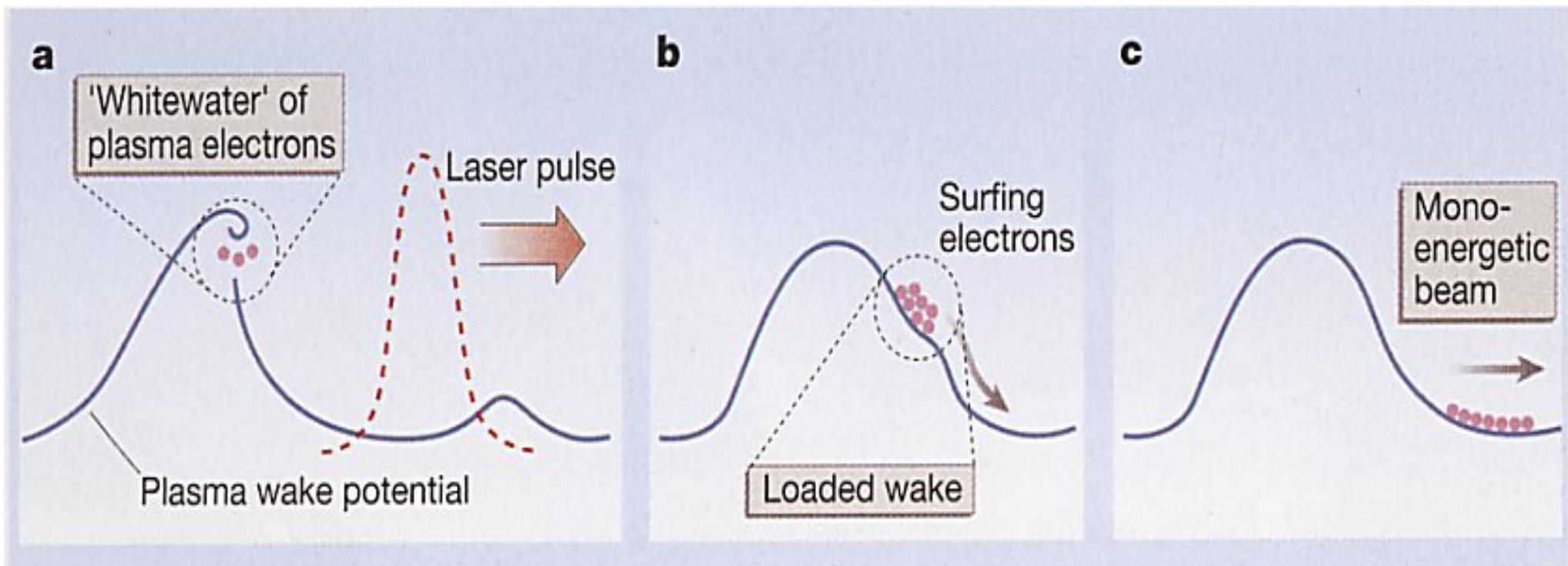
Courtesy of K. Krushelnick et al.

- Vulcan@RAL: 160 J in 650 fs
- Single shot laser



- 350 MeV electrons observed
- Energy spread large

Laser Wakefield



Wakefield acceleration. a, In a plasma excited by a laser pulse, the wake potential rises until it steepens and breaks. Electrons from the plasma are caught in the 'whitewater' and surf the wave. b, The load of the electrons deforms the wake, stopping further trapping of electrons from the plasma. c, As the electrons surf to the bottom of the wake potential, they each arrive bearing a similar amount of energy.

Plasma Wakes – Scaling

- Plasma wake is a relativistic electron plasma wave

$$v_{ph} \leq c$$

- Capable of growing to large values

$$E_W = \varepsilon \sqrt{n_e} \quad \varepsilon < 1$$

- For $n_e \sim 10^{14} \text{ cm}^{-3}$ and $\varepsilon \sim 10\%$

$$E_W \approx 10^8 \text{ V/m} \quad \text{or } 1 \text{ GeV in } \underline{10m}$$

- For $n_e \sim 10^{20} \text{ cm}^{-3}$

$$E_W \approx 10^{11} \text{ V/m} \quad \text{or } 1 \text{ TeV in } \underline{10m}$$

Linear Plasma Wakefield Theory

$$(\partial_t^2 + \omega_p^2) \frac{n_1}{n_o} = -\omega_p^2 \left(\frac{n_b}{n_o} + k_p^2 \nabla^2 \sqrt{1 + a_o^2} \right)$$

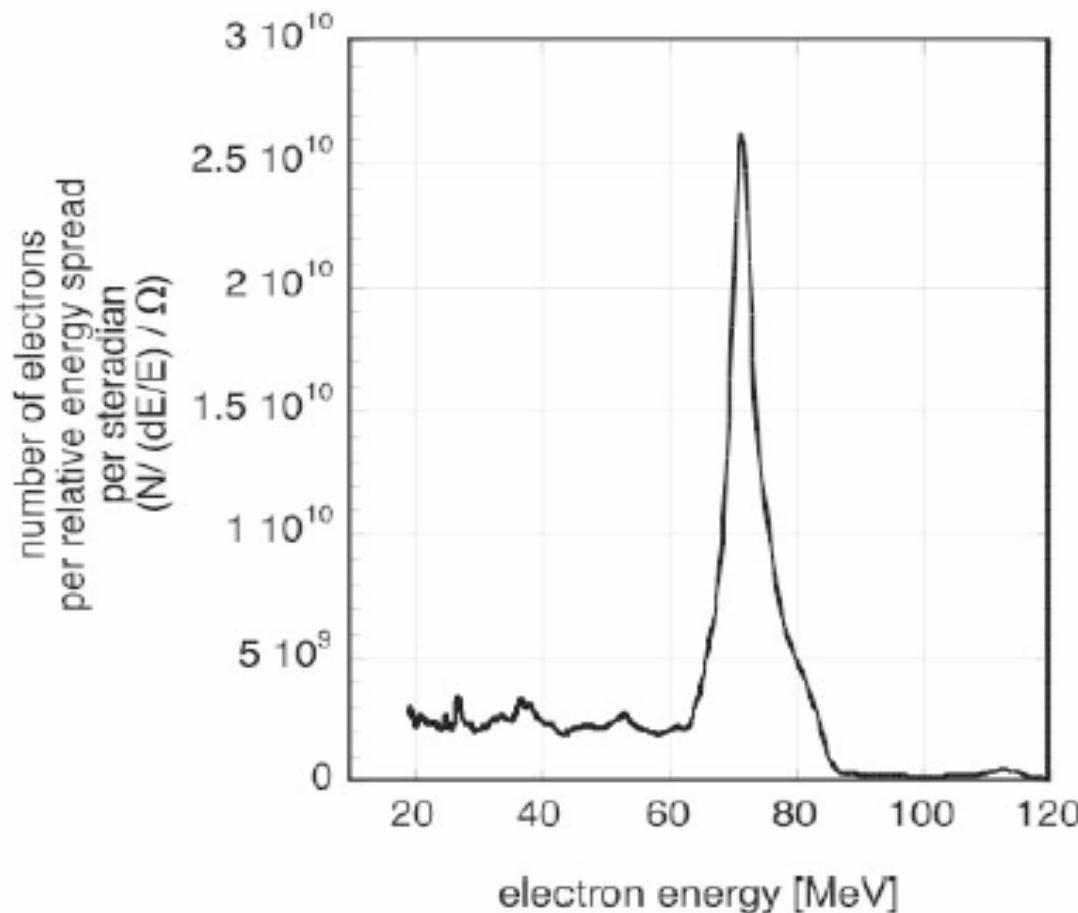
Large wake for a beam density $\mathbf{n}_b \sim \mathbf{n}_o$ or laser amplitude $\mathbf{a}_o = eE_o/m\omega_o c \sim 1$ for τ_{pulse} of order $\omega_p^{-1} \sim 100\text{fs}$ $(10^{16}/n_o)^{1/2}$ and speed $\sim c = \omega_p/k_p$

$$\nabla \bullet E = -4\pi en_1 \Rightarrow eE = \frac{n_1}{n_o} \sqrt{\frac{n_o}{10^{16} cm^{-3}}} 10 GeV/m \cos \omega_p(t - z/c)$$

**But interesting wakes are very nonlinear – e.g. blowout regime
=> PIC simulations**

Recent breakthrough from three laboratories.

Mono-energetic spectra can be observed
at higher power ($\Delta E/E = 6\%$)



$E \sim 500$ mJ,
pulse duration ~ 40 fsec
Focal spot ~ 25 μ m
Density $\sim 2 \times 10^{19}$ cm $^{-3}$

Significant shot-to-shot
fluctuations in
a) energy spread
b) peak energy

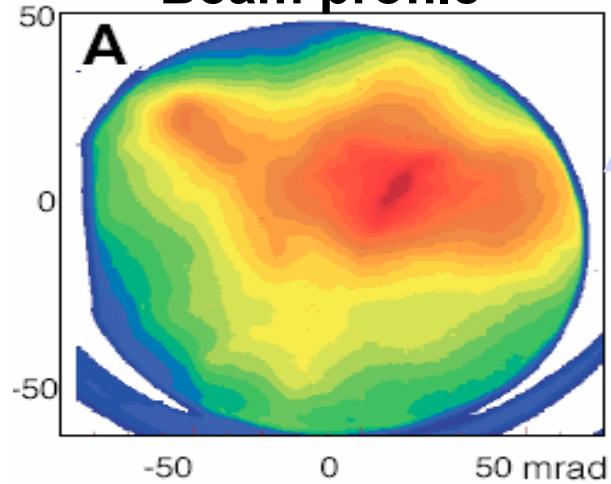
Careful control of laser
and plasma conditions is
necessary

Courtesy: K. Krushelnick, IC

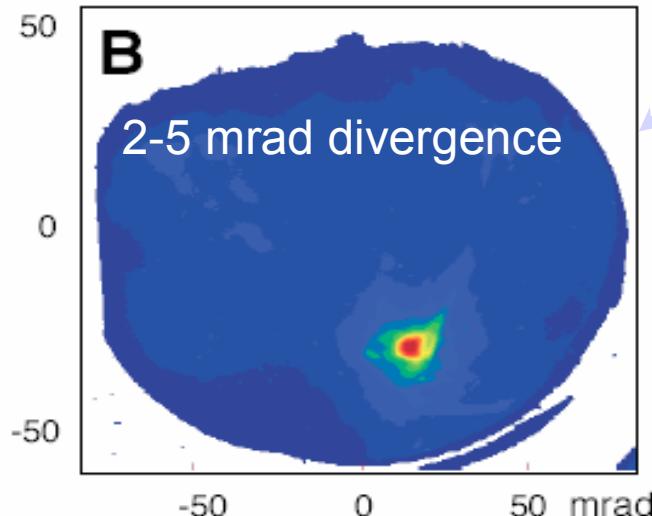
Nature, 2004.

85 MeV e-beam with %-level energy spread observed from laser accelerator

Beam profile

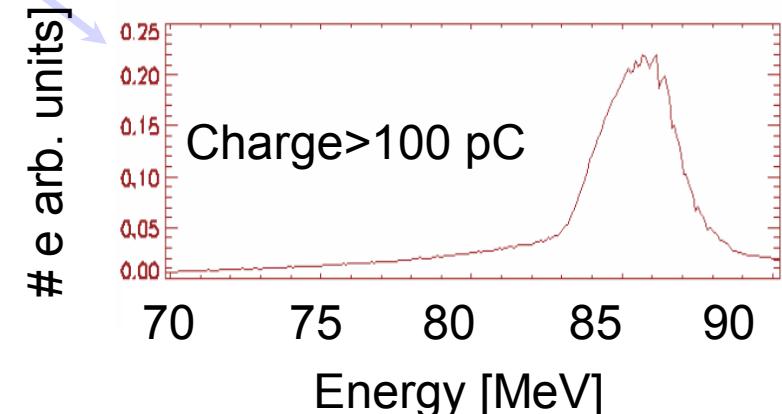
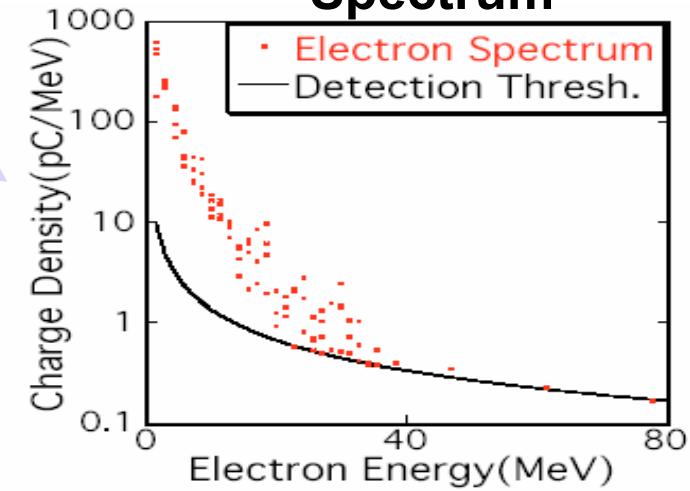


Unguided



Guided

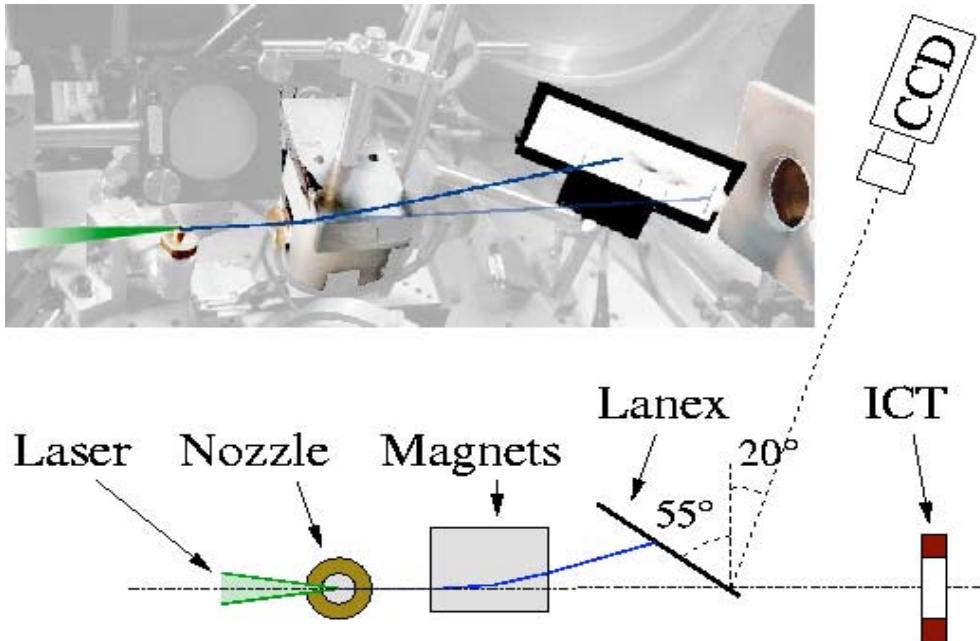
Spectrum



C.Geddes *et al.*, 2004. LBNL

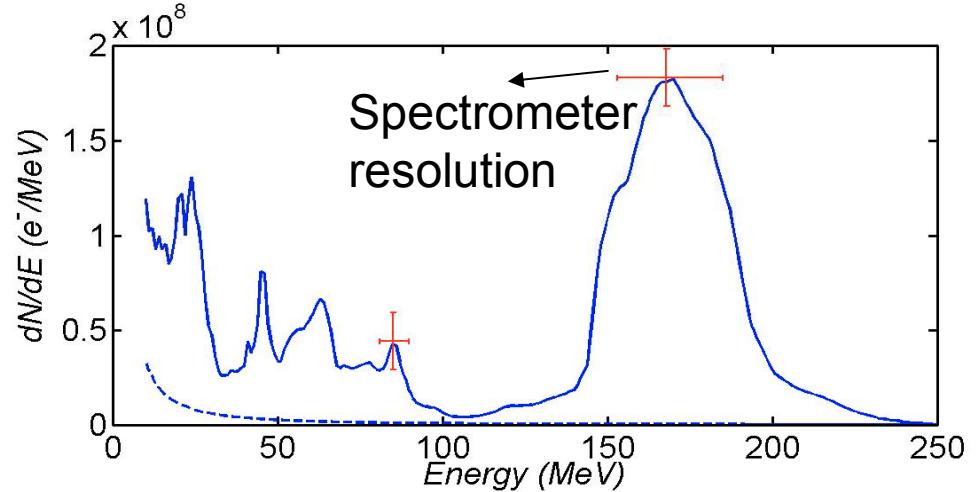
Electrons > 150 MeV observed

Recent Breakthrough -- Mono-energetic Beams! 3 Labs!

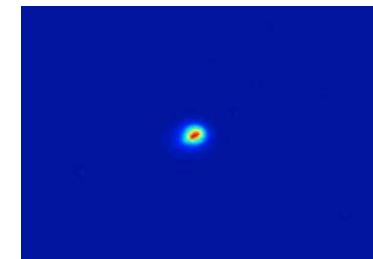


Parameters: $n_e = 6 \times 10^{18} \text{ cm}^{-3}$,
 $a_0 = 1.3$, $\tau = 30 \text{ fs}$ $P = 30 \text{ TW}$
Results obtained with 1 m off-axis parabola:
 $w_0 = 18 \mu\text{m}$, $z_R = 1.25 \text{ mm}$

Quasi-monoenergetic spectrum
Hundreds of pC at 170 MeV +/- 20 MeV



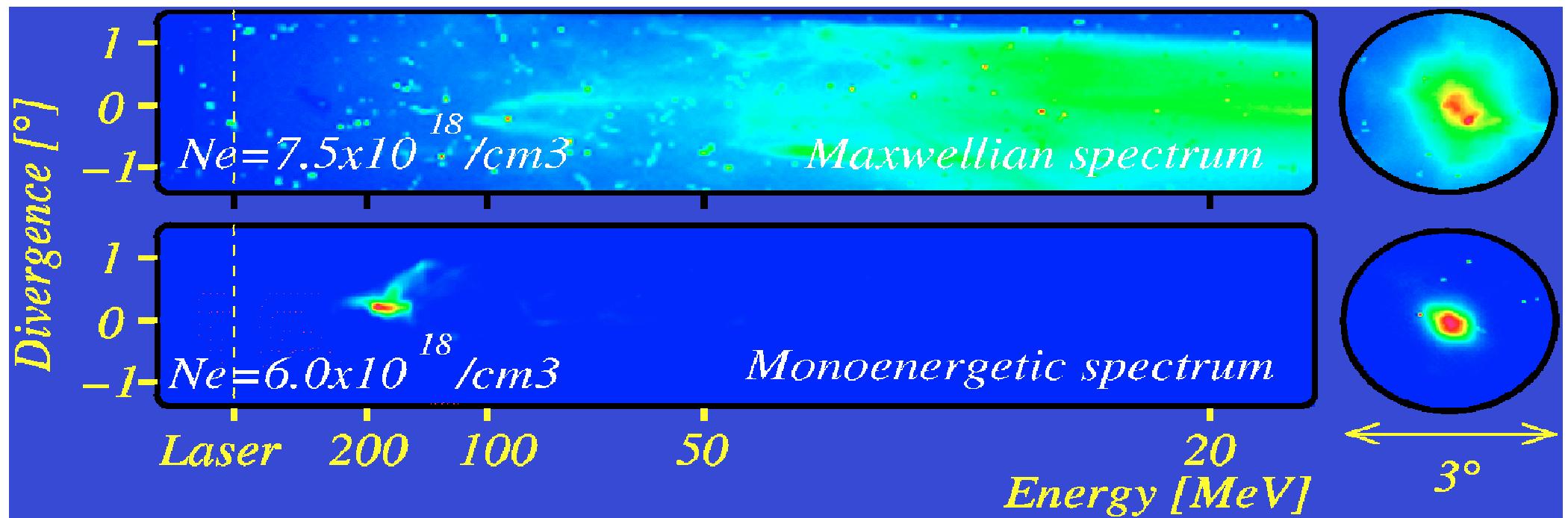
Electron beam
profile on LANEX



Courtesy J. Faure, LOA

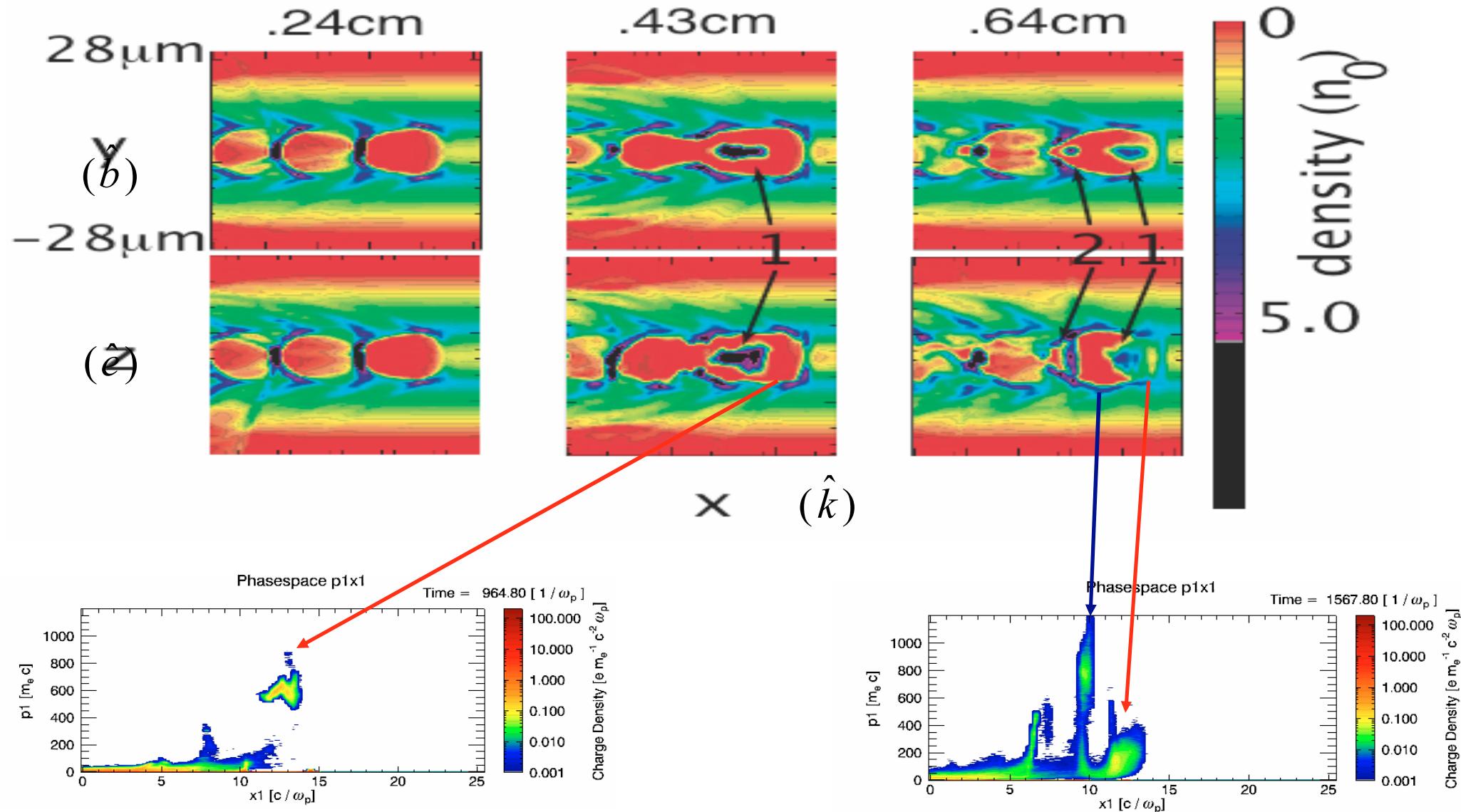
Divergence FWHM = 6 mrad

Recent results on e-beam : Energy distribution improvements



N.B. : color tables are different

Beam loading of first bunch contributes to the generation of a second bunch

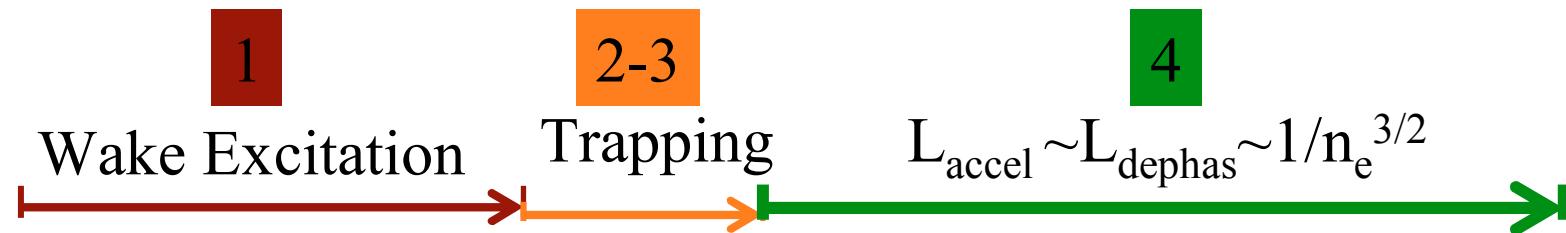


Production of a Monoenergetic Beam

1. Excitation of wake (e.g., self-modulation of laser)
2. Onset of self-trapping (e.g., wavebreaking)
3. Termination of trapping (e.g., beam loading)
4. Acceleration

If $>$ or $<$ dephasing length: large energy spread

If \sim dephasing length: monoenergetic

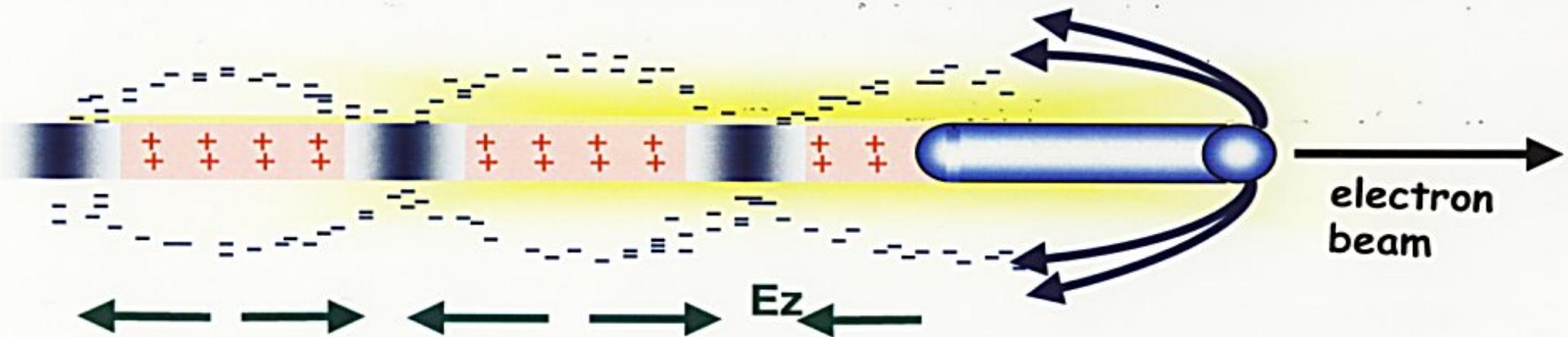


Optimal choice of the plasma density: the smallest possible density
For conditions 1 -4 to be fulfilled.

Physical Principles of the Plasma Wakefield Accelerator

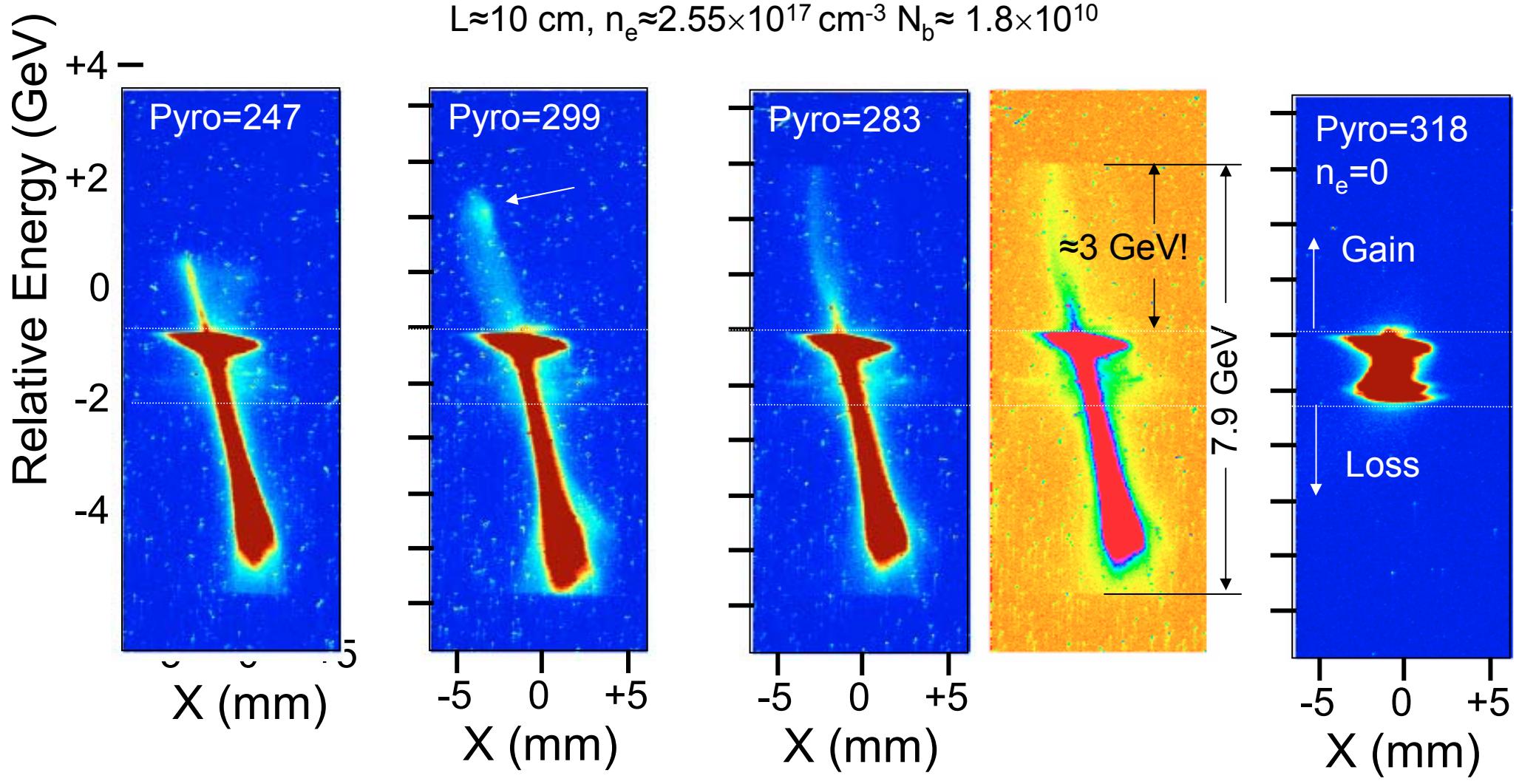
- Plasma is used to create a longitudinal accelerating gradient

- Space charge of drive beam displaces plasma electrons



- Plasma ions exert restoring force => Space charge oscillations
- Wake Phase Velocity = Beam Velocity (like wake on a boat)
- Wake amplitude $\propto N_b / \sigma_z^2$ (*for $2\sigma_z \approx \lambda_p \propto \frac{1}{\sqrt{n_o}}$*)

E-164X Breaks GeV Barrier



Energy gain exceeds $\approx 4 \text{ GeV}$ in 10 cm

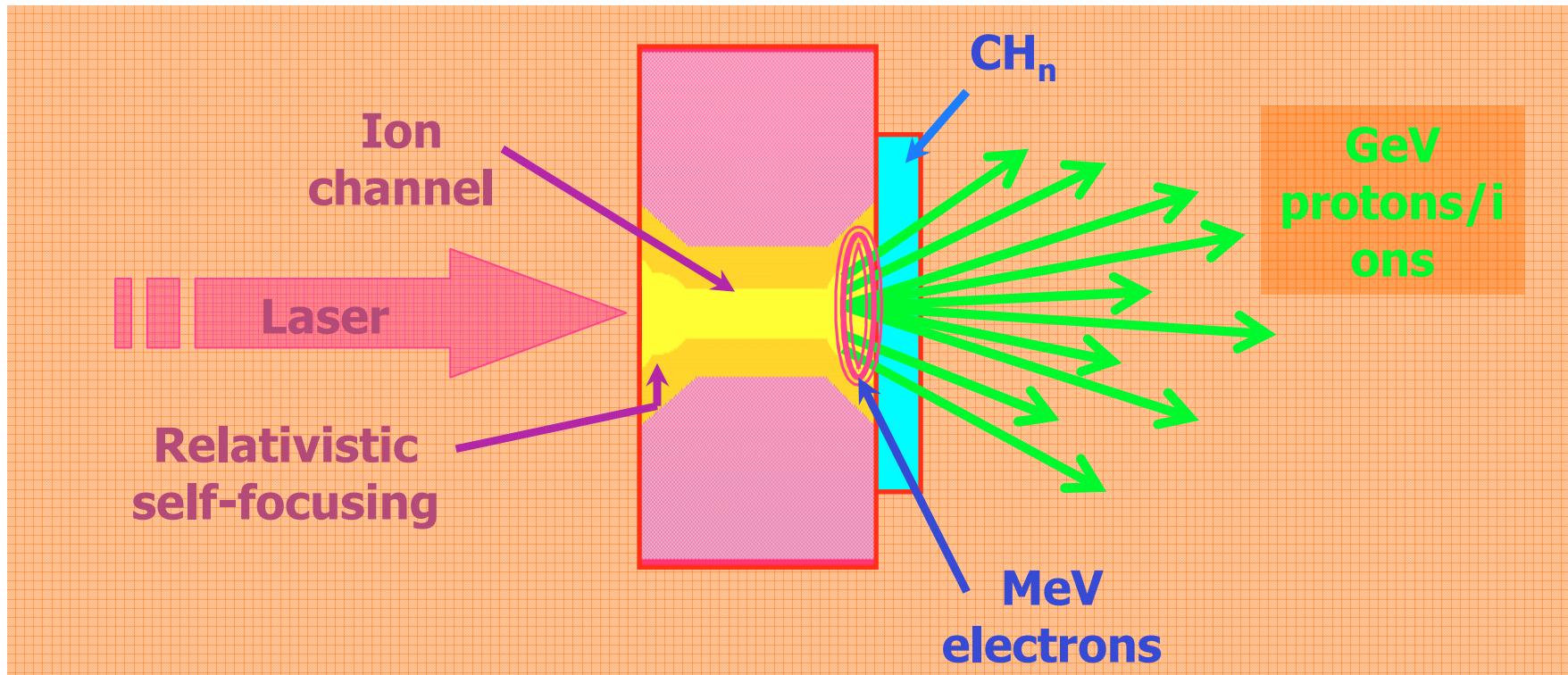
Plasma Wakefield

- The maximum plasma wake amplitude of an electron bunch scales as current over pulse length ($\tau=2\sigma_z/c$)

$$eE_z \cong 1 \text{ GeV/m} \frac{I}{kAmp} \cdot \frac{4 \text{ ps}}{\tau}$$

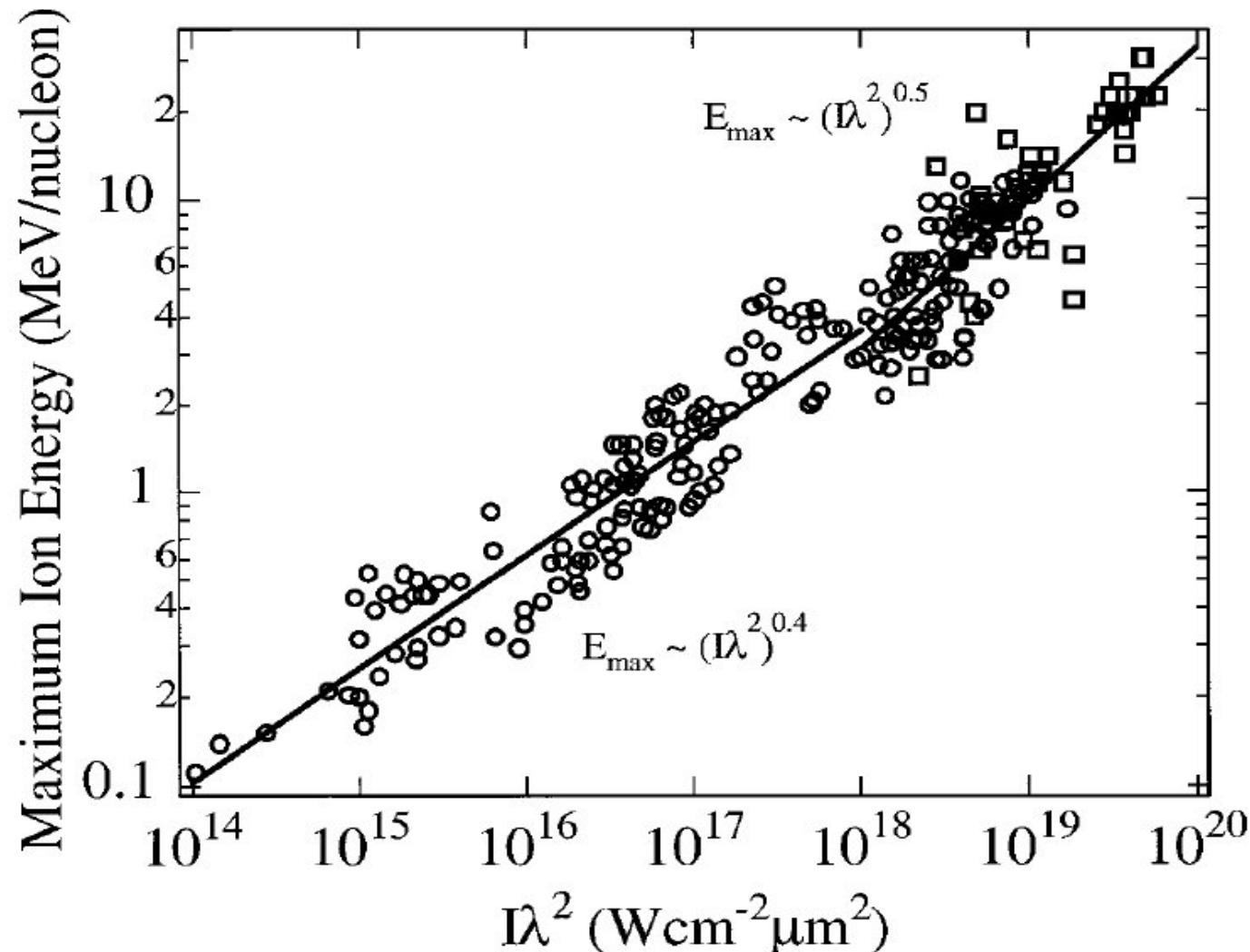
- So, a 1 kAmp pulse, of duration 4 picoseconds is needed to generate a 1 GeV/m wake.

Proton Acceleration



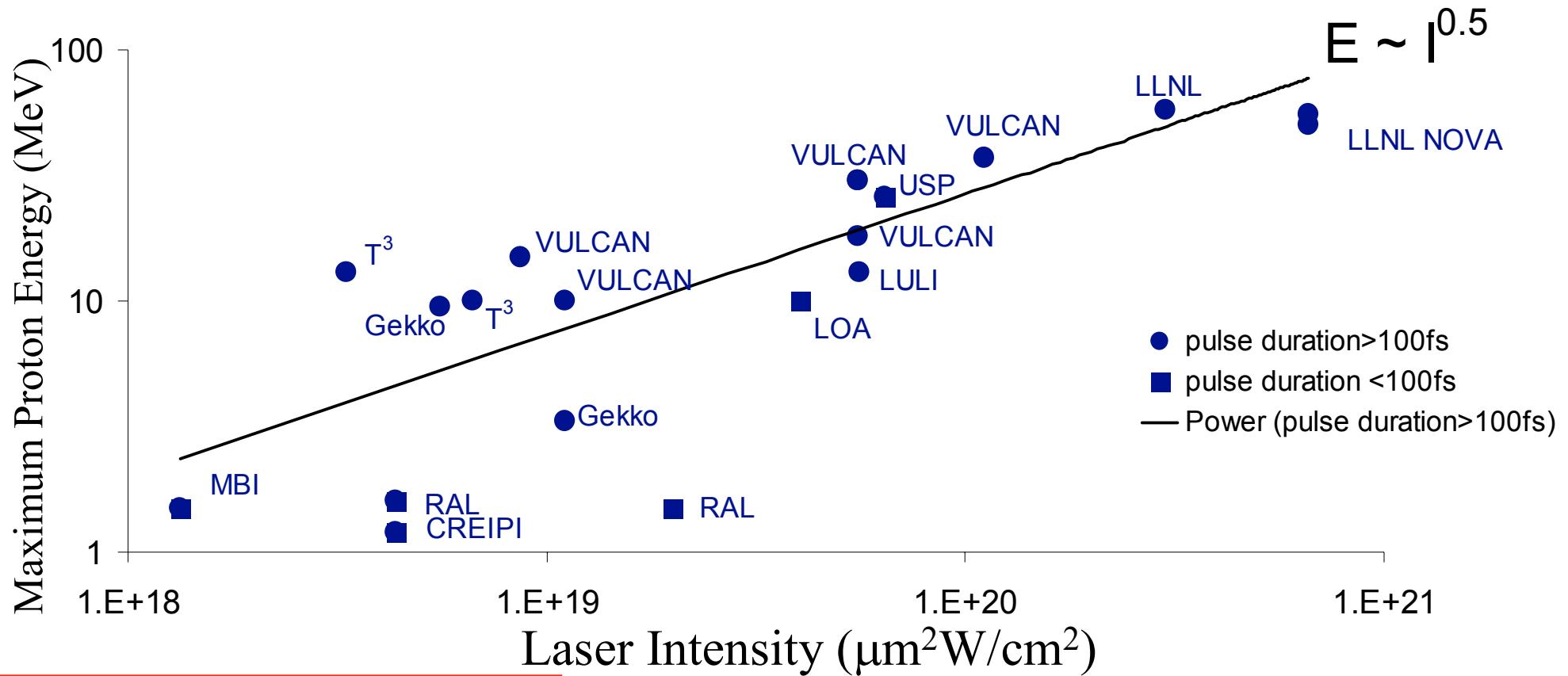
At present, experiments achieve 10s of MeV per nucleon.
Future experiments aim to reach 100s of MeV to GeV.

Heavy Ion/Proton Generation by Ultraintense Lasers



*Clark et al.,
2000, PRL,
85, 1654.*

FIG. 4. Maximum ion energy as a function of $I\lambda^2$. Data from Refs. [7] and [10] are indicated by circles. Squares denote data from experiments discussed here.



- Hi charge: 10^{10} - 10^{13} ions
- Short pulses
- 100's MA/cm²

(Courtesy T. Lin)

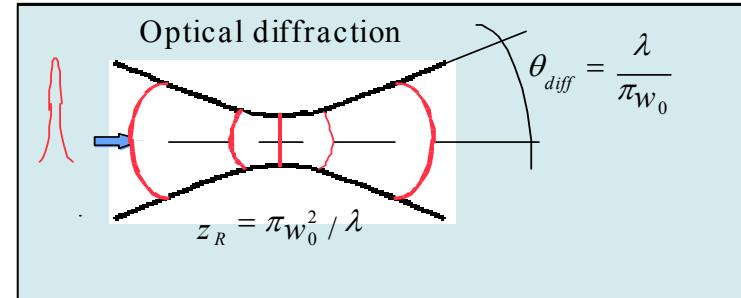
Key Issues

<u>Key Issue</u>	<u>Experiment</u>	<u>Theory/Simulation</u>
Acceler. Length $\text{mm} \rightarrow \text{cm+}$	Channel Formation Plasma Sources	1-to-1 models parallel 3-D hybrid
Beam Quality $\Delta\gamma$ ε N	Injectors 50 fs bunch 50 μm spot Blowout regime	Beam Dynamics matching β injection phase
Efficiency (new)		Drive beam evolution Shaped driver and load Transformer Ratio Energy Spread

3 Limits to Energy gain $DW = eE_z L_{acc}$ (laser driver)

- Diffraction:

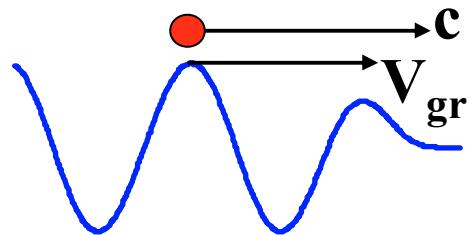
$$L_{dif} \cong \pi L_R = \pi^2 w_0^2 / \lambda$$



order mm!

(but overcome w/ channels or relativistic self-focusing)

- Dephasing:



$$L_{dph} = \frac{\lambda_p / 2}{1 - V_{gr} / c}$$

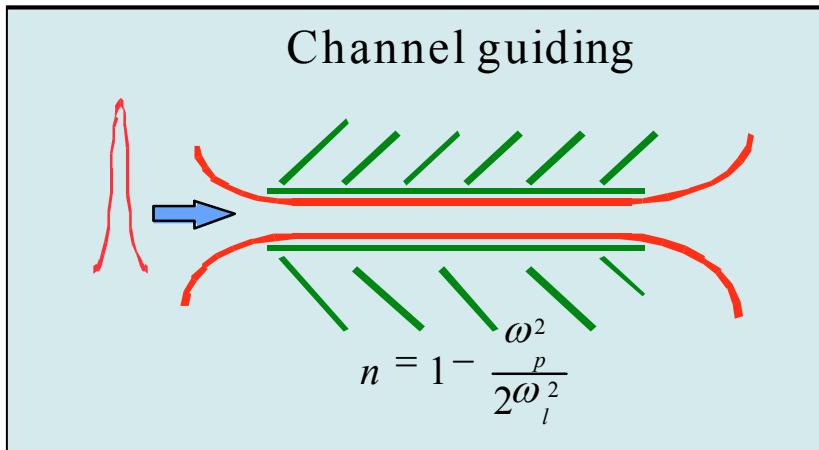
order 10 cm
 $\times 10^{16}/n_o$

- Depletion:

For small a_0 $\gg L_{dph}$
 For $a_0 > \sim 1$ $L_{dph} \sim L_{depl}$

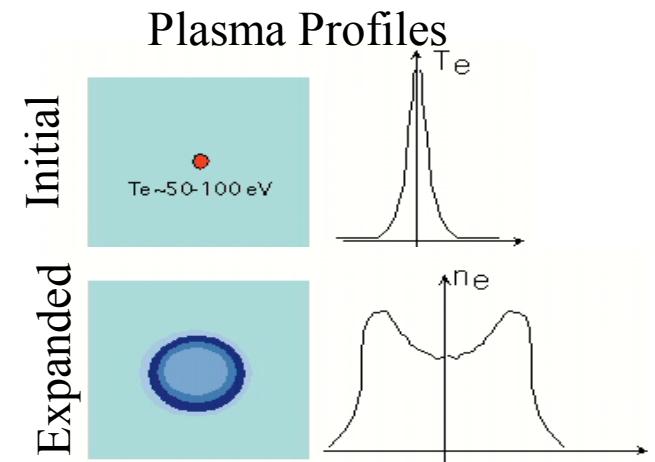
$$\Delta W_{ch}[MeV] \sim 60 \left(\lambda_p / w_0 \right)^2 P[TW]$$

Plasma channel: structure for guiding and acceleration



Step 1: Heat

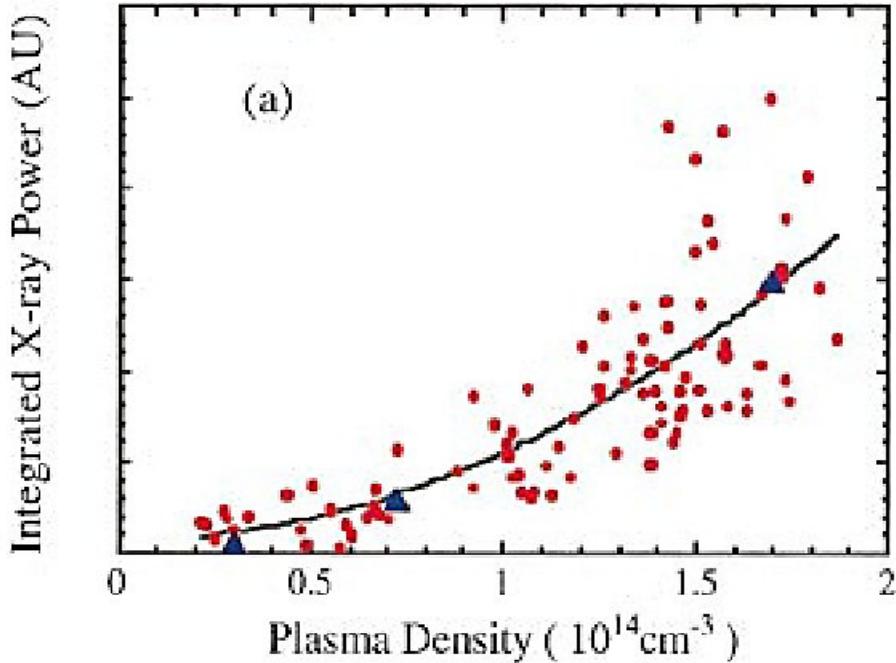
Step 2: expand



- Hydro-dynamically formed plasma channel
 - On-axis axicon (C.G. Durfee and H. Milchberg, PRL 71 (1993))
 - Ignitor-Heater (P. Volfbeyn et al., Phys. Plasmas 6 (1999))
 - Discharge assisted (E. Gaul et al., Appl. Phys. Lett. 77 (2000))
 - Cluster jets (Kim et al., PRL 90 (2003))
- Discharge ablated capillary discharges (Y. Ehrlich et al., PRL 77 (1996))
- Z-pinch discharge (T. Hosokai et al., Opt. Lett. 25 (2000))
- Hydrogen filled capillary discharge (D. Spence and S. M. Hooker, JOSA B (2000))
- Glass capillaries (B. Cross et al., IEEE Trans. PS 28(2000), Y. Kitagawa PRL (2004))

Radiation from Plasma Channels

Betatron X-rays



$I \sim 10^{19} \text{ photons/s}.1\%\text{bw-mm}^2\text{-mr}^2 @ 6 \text{ keV}$

Radiation from energetic electrons in plasma channels.

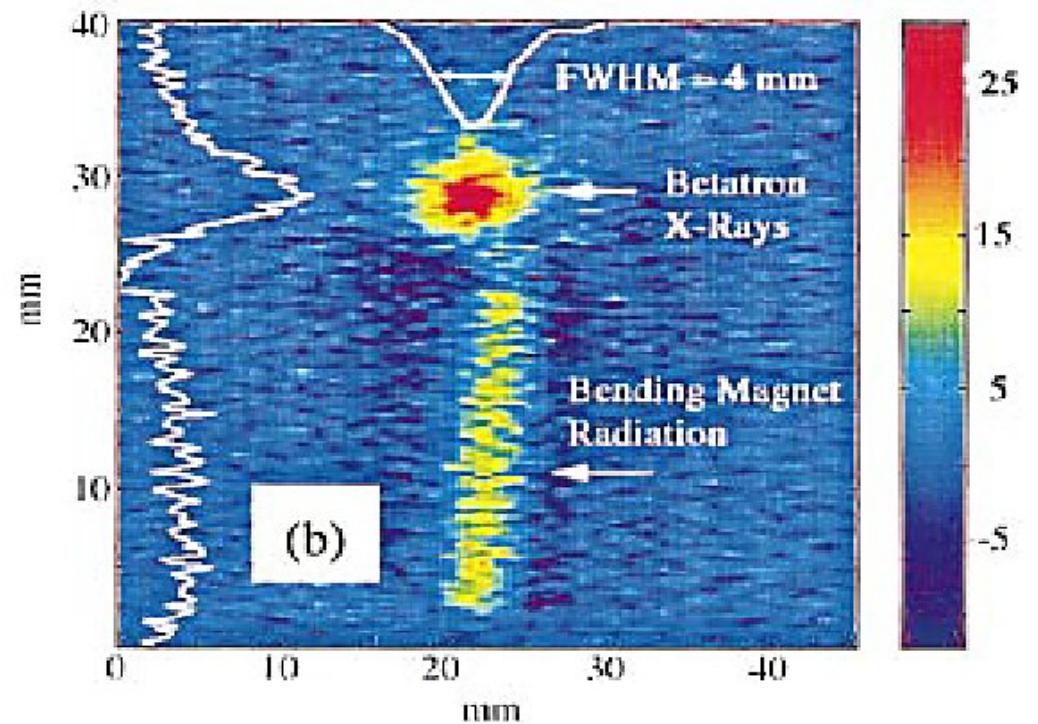


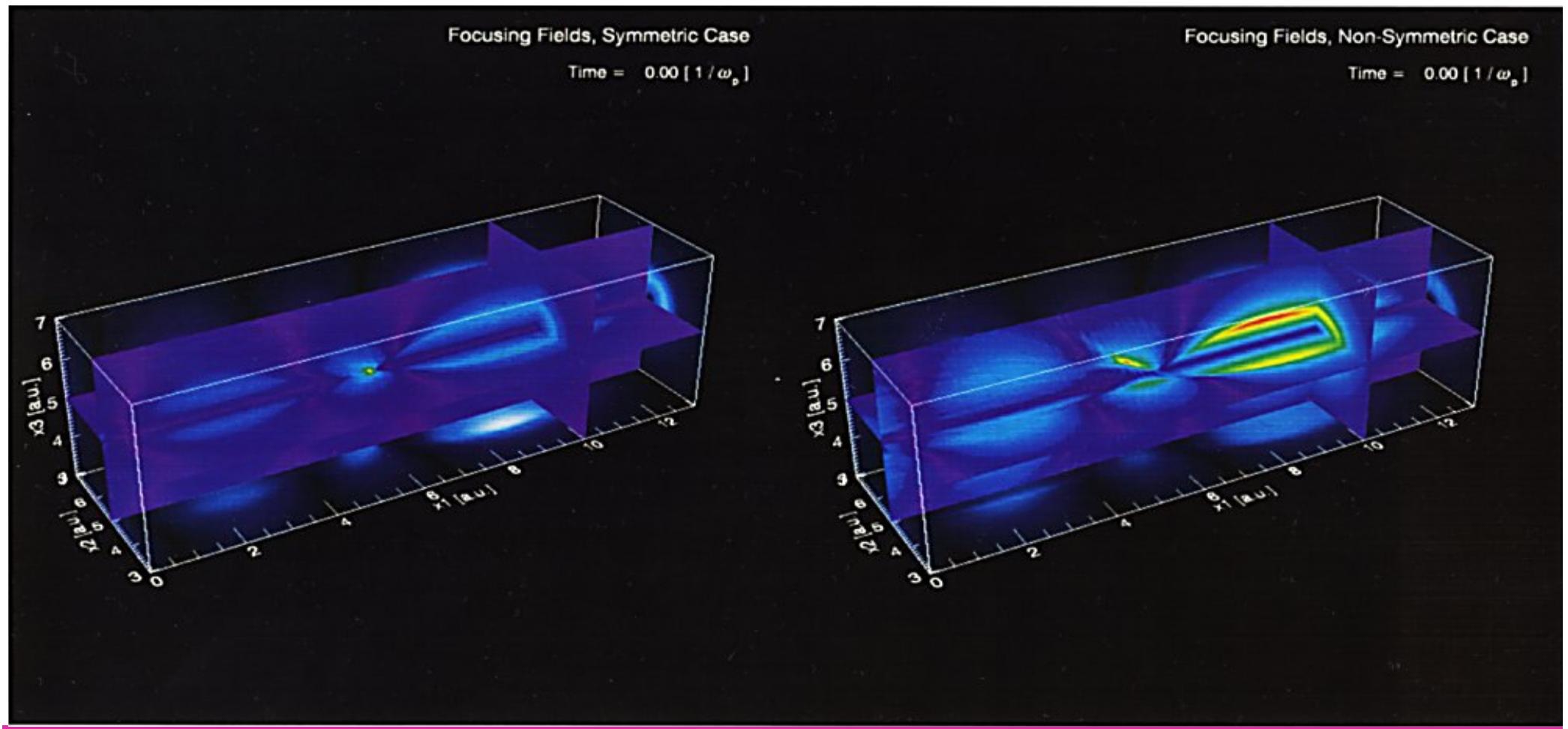
FIG. 8. (Color) (a) The estimated (triangles) and the measured (dots) x-ray energy in the 5–30 keV range as a function of plasma density. The solid line is a quadratic fit to the data. (b) Processed image produced on a fluorescent screen as recorded by a CCD camera showing the betatron x-rays produced by the plasma $n_p = 2 \times 10^{13} \text{ cm}^{-3}$ (circle at the top) and a vertical stripe of remnant synchrotron radiation produced by a dipole bend magnet.

Joshi *et al.*, Physics of Plasmas, 9, 1845, 2002.

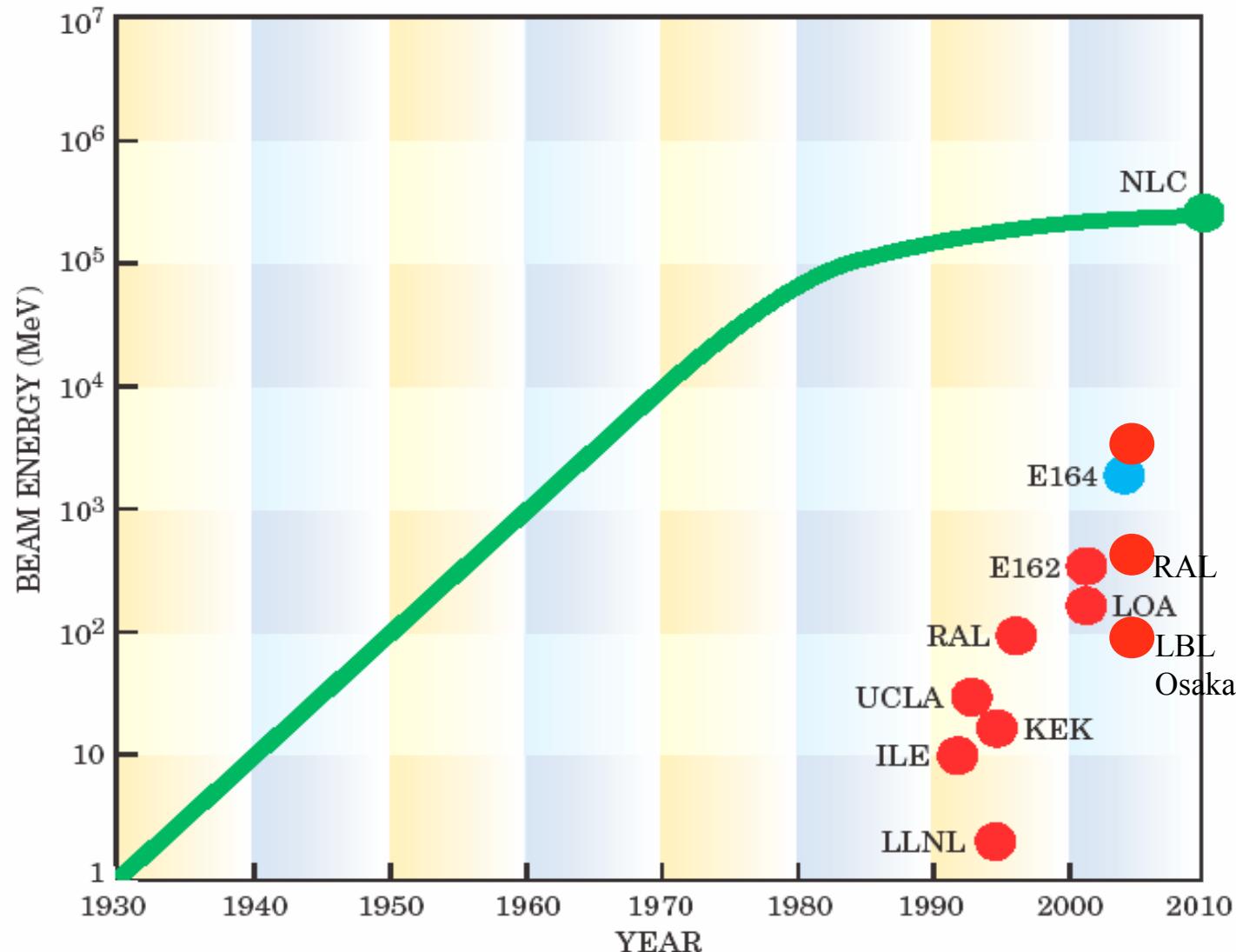
S. Wang *et al.* Phys. Rev. Lett. Vol. 88 Num. 13

Betatron Oscillations

- Different spot sizes lead to different focusing forces and betatron oscillations



Plasma Accelerator Progress and the Livingston Curve

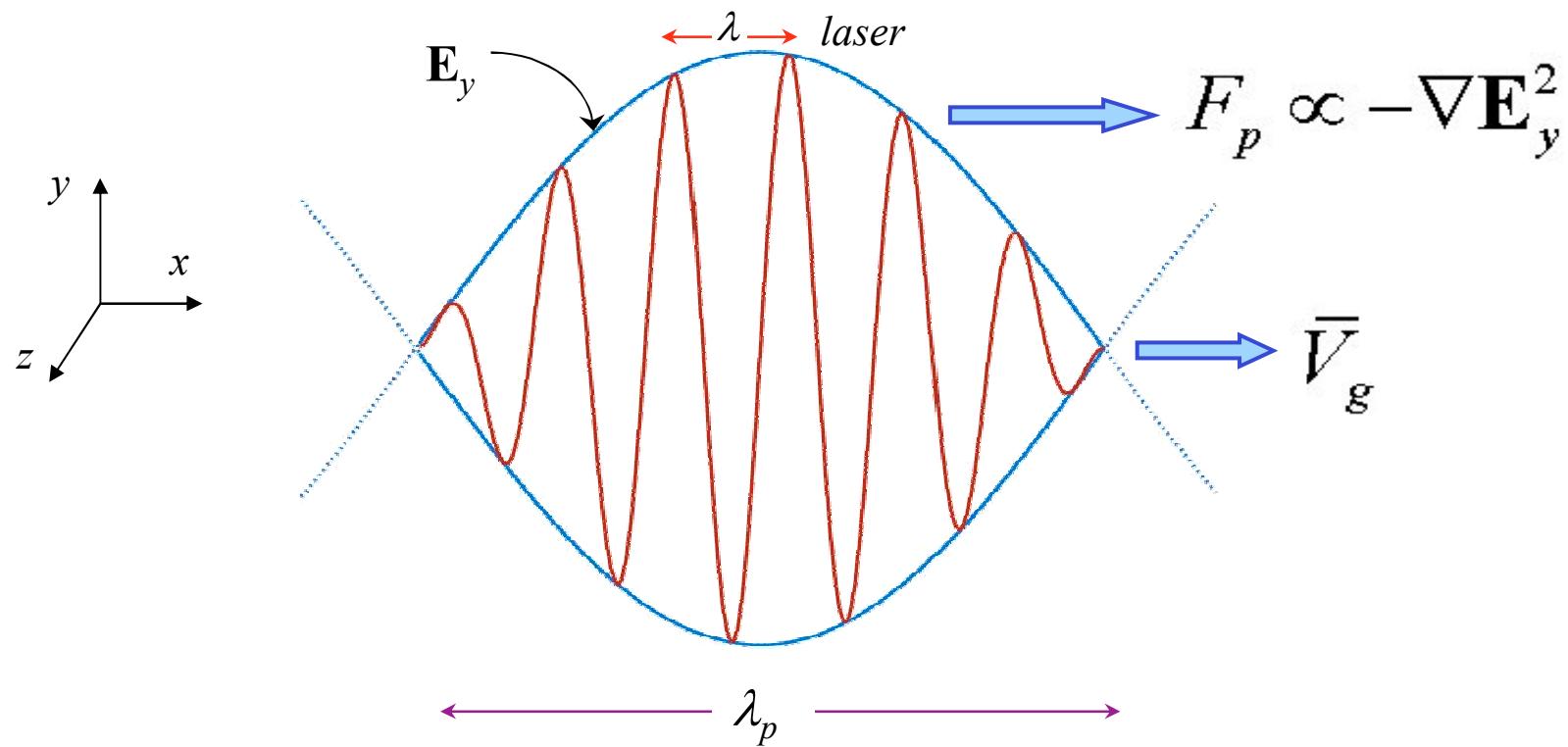


Conclusions

- The dawn of compact particle accelerators is here.
- Laser Plasma Accelerators > 100 MeV
- Numerous applications for 100 MeV-1 GeV beams – Medicine, Light Sources, Industry
- Future goals for laser plasma accelerators are mono-energetic, multi GeV beams.
- With advances in laser technology the TeV energy scale is a long term target.

Particle Accelerators – Requirements for High Energy Physics

- High Energy
- High Luminosity (event rate)
 - $L=fN^2/4\pi\sigma_x\sigma_y$
- High Beam Quality
 - Energy spread $\delta\gamma/\gamma \sim .1 - 10\%$
 - Low emittance: $\varepsilon_n \sim \gamma\sigma_y\theta_y < 1 \text{ mm-mrad}$
- Low Cost (one-tenth of \$6B/TeV)
 - Gradients $> 100 \text{ MeV/m}$
 - Efficiency $> \text{few \%}$



- Envelope of high frequency field moving at group speed \underline{v}_g

$$\left. \begin{aligned} \underline{v}_g &= c \left(1 - \frac{\omega_{pe}^2}{\omega^2} \right)^{1/2} \\ \omega^2 &= \omega_{pe}^2 + c^2 k^2 \end{aligned} \right\} \quad v_g = \frac{d\omega}{dk} = \frac{c^2 k^2}{\omega} = c \left(1 - \frac{\omega_{pe}^2}{\omega^2} \right)^{1/2}$$

- Ponderomotive force

$$F_p = -\nabla \mathbf{E}_y^2 \quad \text{Laser field } \mathbf{E}_y$$

Electron beam

$$(\partial_t^2 + \omega_{pe0}^2) \delta n_e = -\omega_{pe0}^2 n_{e-beam}$$

$\delta n_e \equiv$ Perturbed electron plasma density

Photons

$$(\partial_t^2 + \omega_{pe0}^2) \delta n_e = \frac{\omega_{pe0}^2}{2m_e} \nabla^2 \int \frac{d\mathbf{k}}{(2\pi)^3} \hbar \frac{N_\gamma}{\omega_\mathbf{k}}$$

Neutrinos

$$(\partial_t^2 + \omega_{pe0}^2) \delta n_e = \frac{\sqrt{2} n_{e0} G_F}{m_e} \nabla^2 n_\nu$$

+

Ponderomotive force

physics/9807049

physics/9807050

Kinetic/fluid equations for electron beam, photons, neutrinos coupled with electron density perturbations due to PW

Self-consistent picture of collective e,γ,ν-plasma interactions

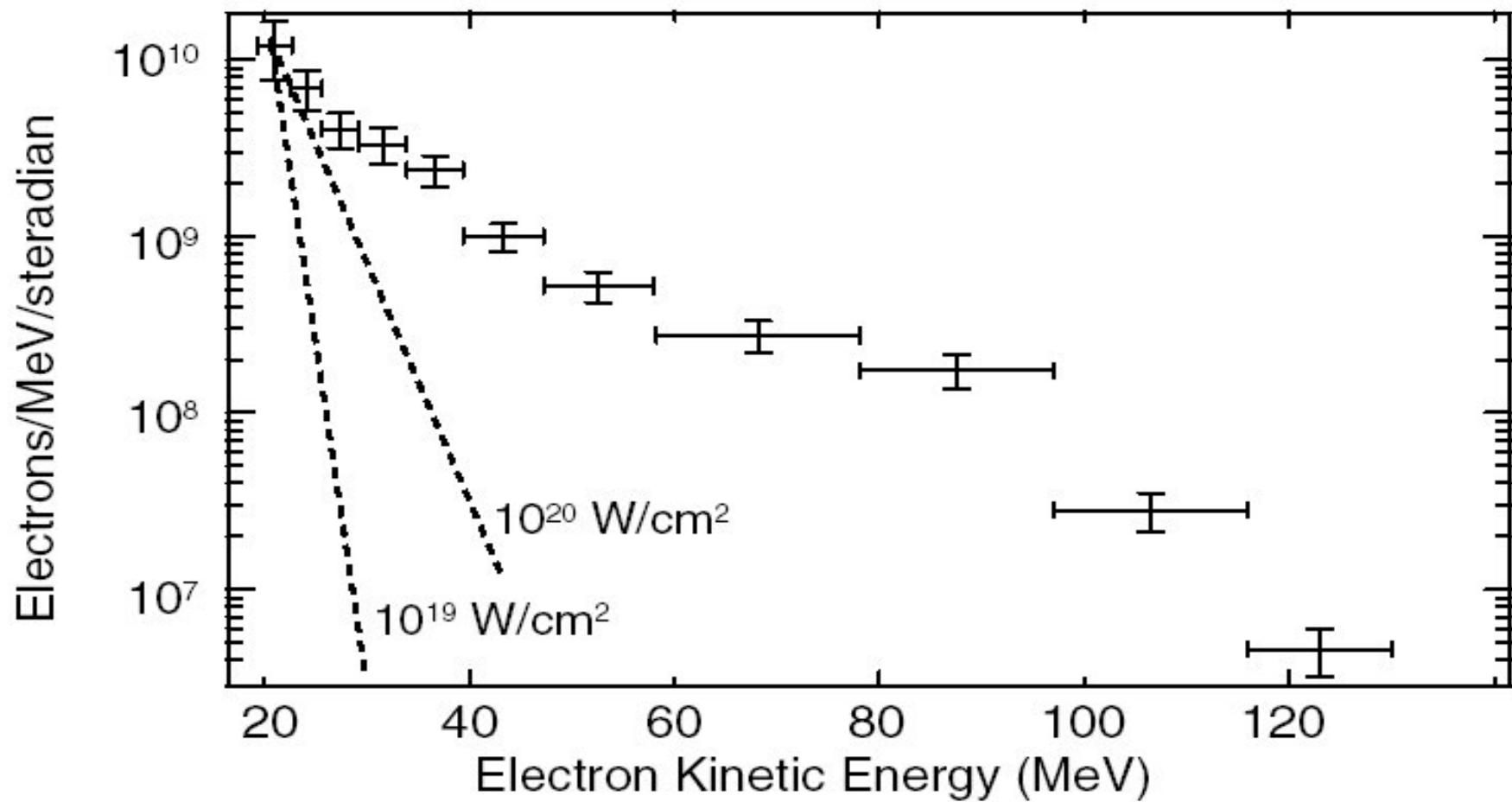
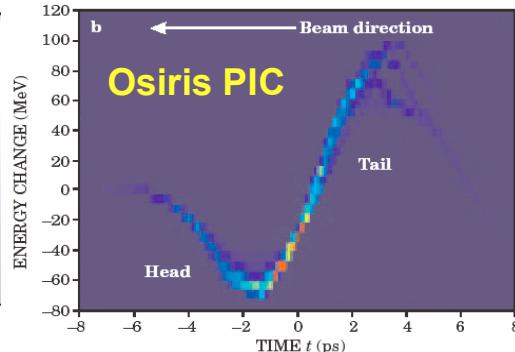
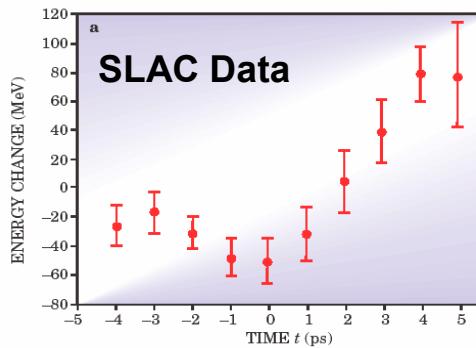


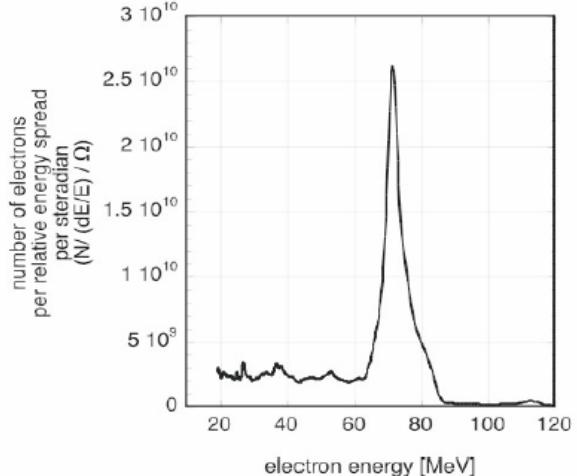
FIG. 3. A typical electron spectrum (unfolded) measured by an on-axis electron spectrometer. Ponderomotive scalings (Ref. [12]) at 10^{19} and 10^{20} W/cm^2 are also shown.

Plasma Accelerators

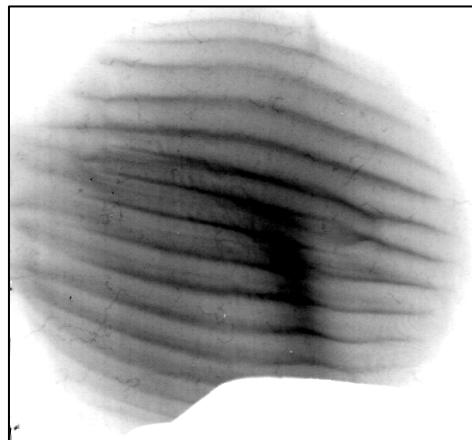
1. Positron acceleration/ 3-D Modeling



2. Monoenergetic electron beams

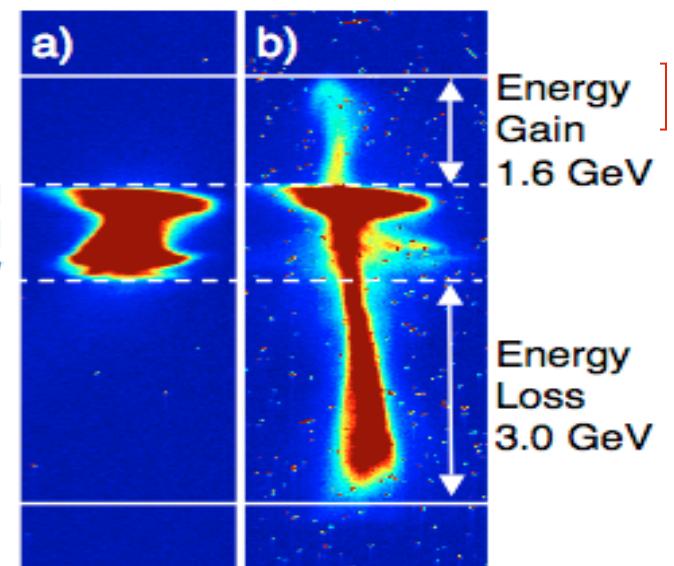


3. Multi-MeV proton beams

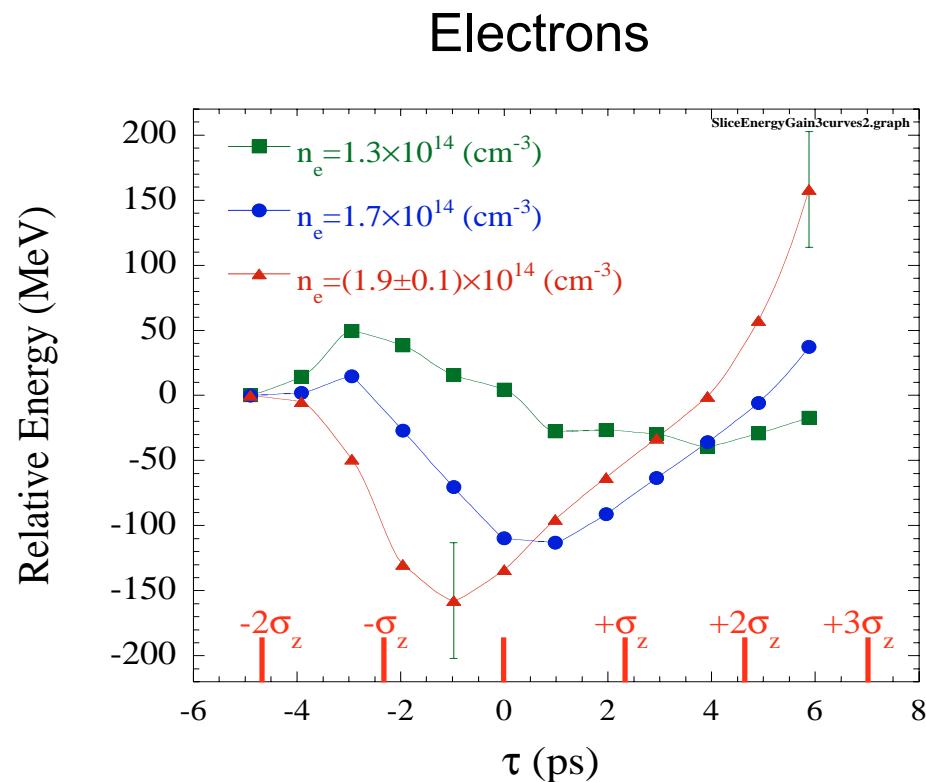


Four
Highlights!

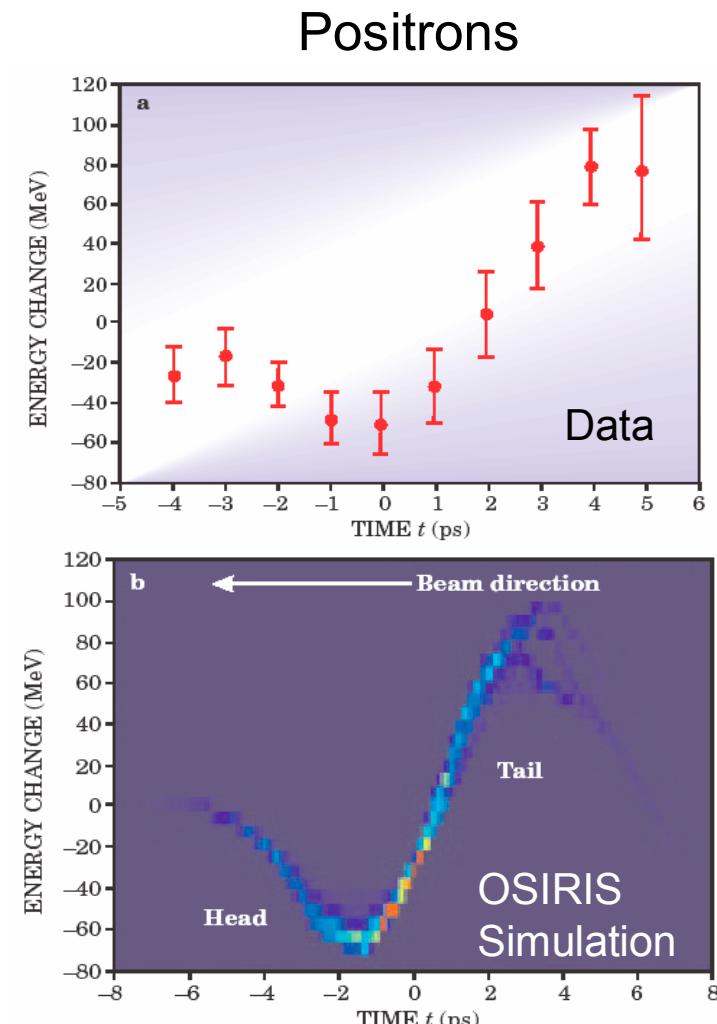
Incoming
Energy Spread
1.4 GeV



Acceleration Of Electrons & Positrons: E-162



- Some electrons gained 280 MeV (200 MeV/m)
- Now going for 2 GeV at a rate of 10,000 MeV/m this month at SLAC



- Loss ≈ 50 MeV
- Gain ≈ 75 MeV

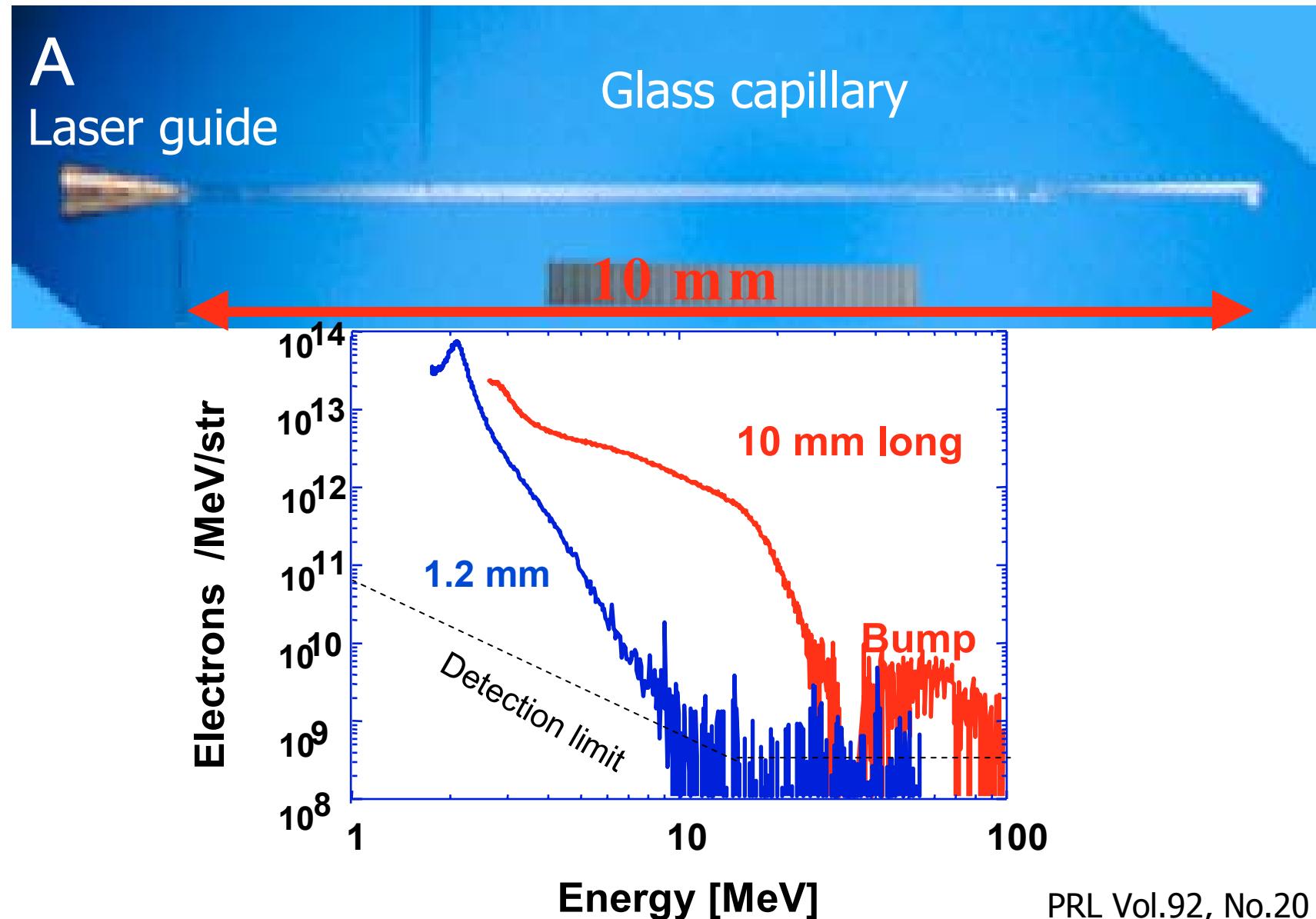
Summary of Experimental Results

Mechanism	Labs	Energy gain	Acc field	Acc length
BW	UCLA, LULI, Canada, ILE	1 to 30 MeV	1 GV/m	1 to 10 mm
Laser Wakefield	LULI	1.5 MeV	1 GV/m	2 mm
Plasma Wakefield	SLAC-UCLA- USC-Berkeley	80 - 150 MeV	70 MV/m	1.4 m
SM Wakefield	RAL, LULI, LOA	60 to 200 MeV	100 to 400 GV/m	1 mm

- ➡ Large accelerating gradients
- ➡ Agreement with theoretical predictions
- ➡ Broad spectrum due to inadequate injectors
- ➡ Improvements are necessary**

Intense Relativistic Beams in Plasmas: New Plasma Physics

- Wake generation/ particle acceleration
- Focusing
- Hosing
- “Collective Refraction”
- Radiation generation
- Ionization effects
- Compact accelerators
- Plasma lens/astro jets
- E-cloud instability/LHC
- Fast kickers
- Tunable light sources
- Beam prop. physics/X-ray lasers



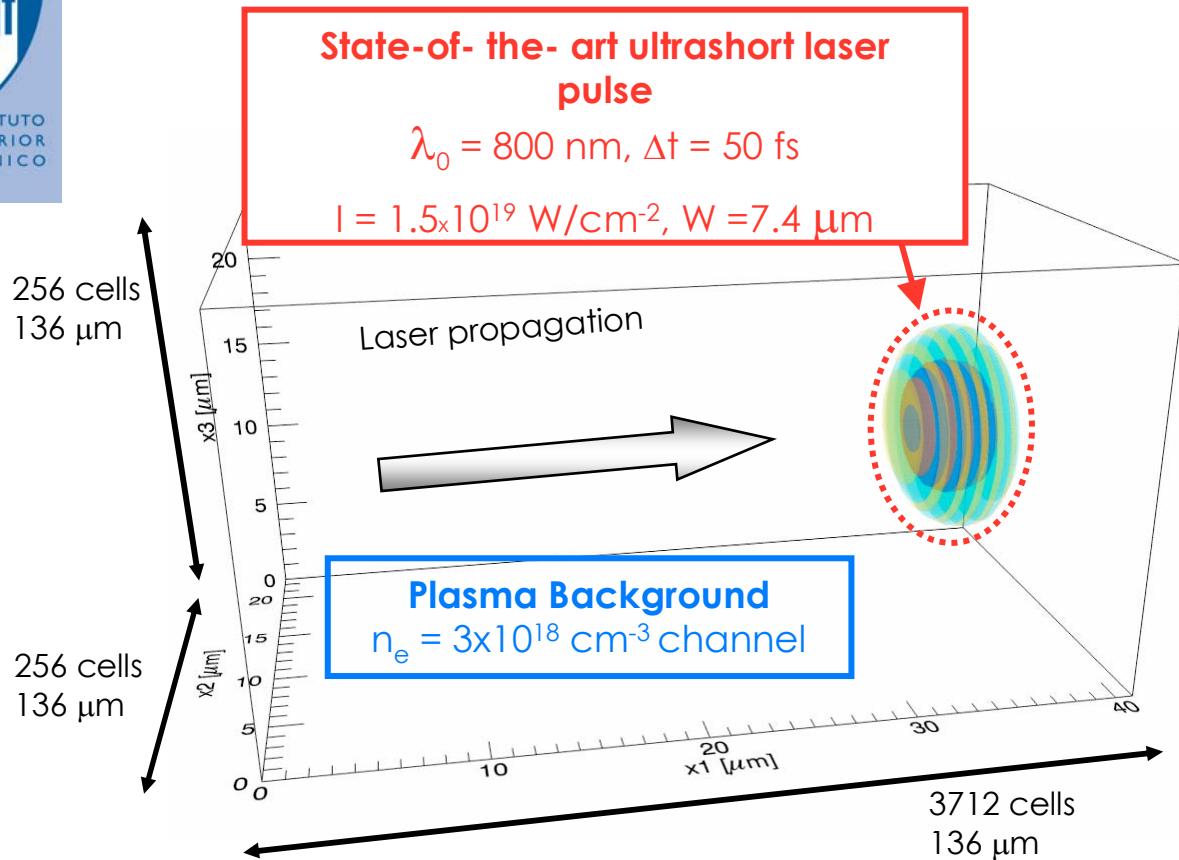
Electron Accelerators

- 1) Laser-induced Plasma Wakefield Accelerators
 - a) Plasma Beat-Wave Accelerator (PBWA)
 - b) Laser Plasma Wakefield Accelerator (LPWA)
 - c) Self-Modulated Laser Wakefield Accelerator (SMLWA)
- 2) Electron beam-induced Plasma Wakefield Accelerator (PWA)
- 3) Inverse Cherenkov Laser Accelerator
- 4) Surfatron Accelerators

Conclusions

- Laser Plasma Accelerators > 100 MeV
- Numerous applications for 100 MeV-1 GeV beams – Medicine, Light Sources, Industry
- Ultra-High energies can be achieved by using a plasma afterburner on existing facilities – energies can be boosted up to 100 GeV
- 1 GeV barrier was broken by SLAC e-beam Wakefield Experiment.

Full 3D LWFA Simulation



Simulation ran for 200,000 hours
 $(\sim 40 \text{ Rayleigh lengths})$



imulation Parameters

-Laser:

- $a_0 = 3$
- $W_0 = 9.25 \lambda = 7.4 \mu\text{m}$
- $\omega_l/\omega_p = 22.5$

-Particles

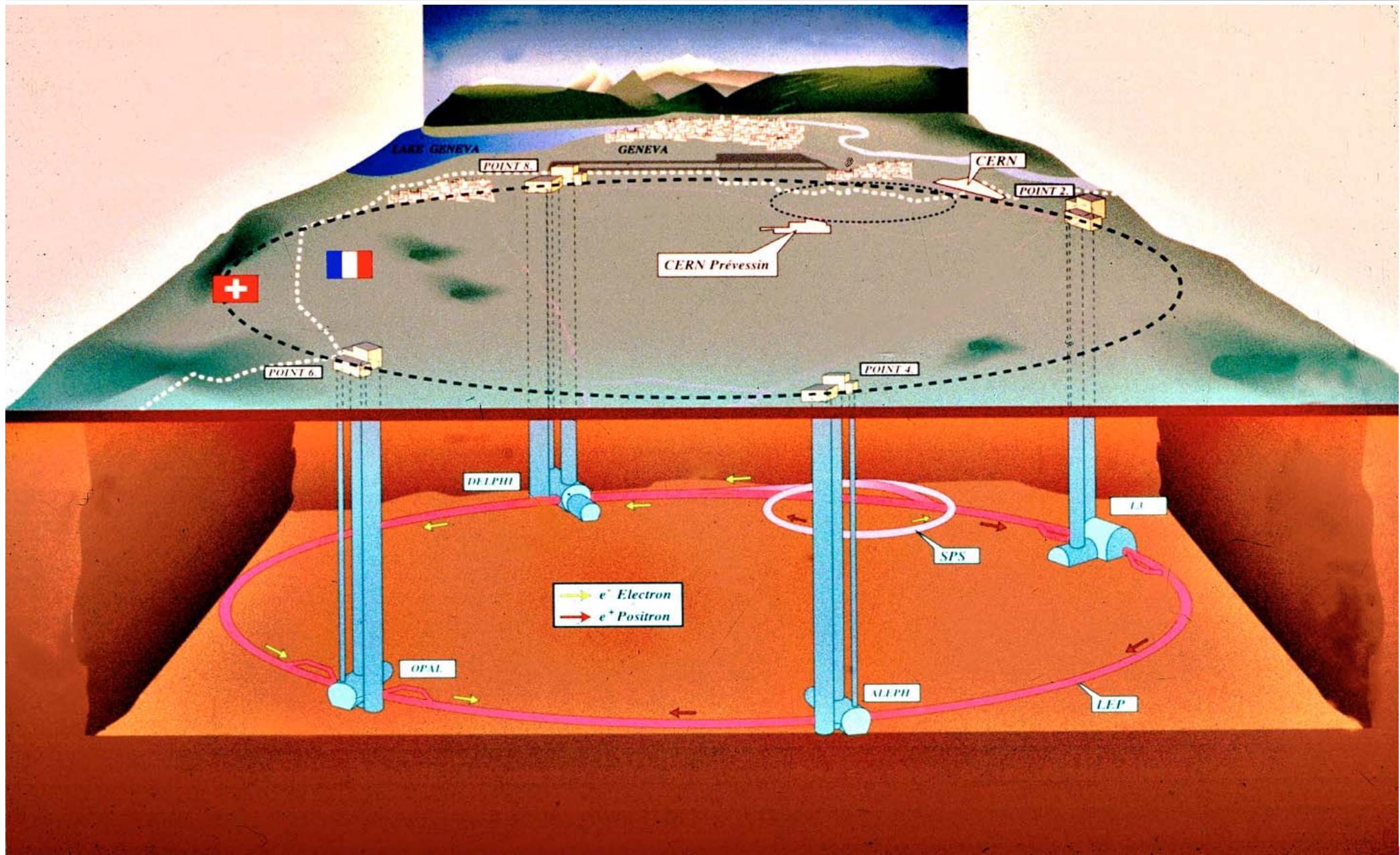
- 1x2x2 particles/cell
- 240 million total

-Channel length

- $L = .828 \text{ cm}$
- 300,000 timesteps

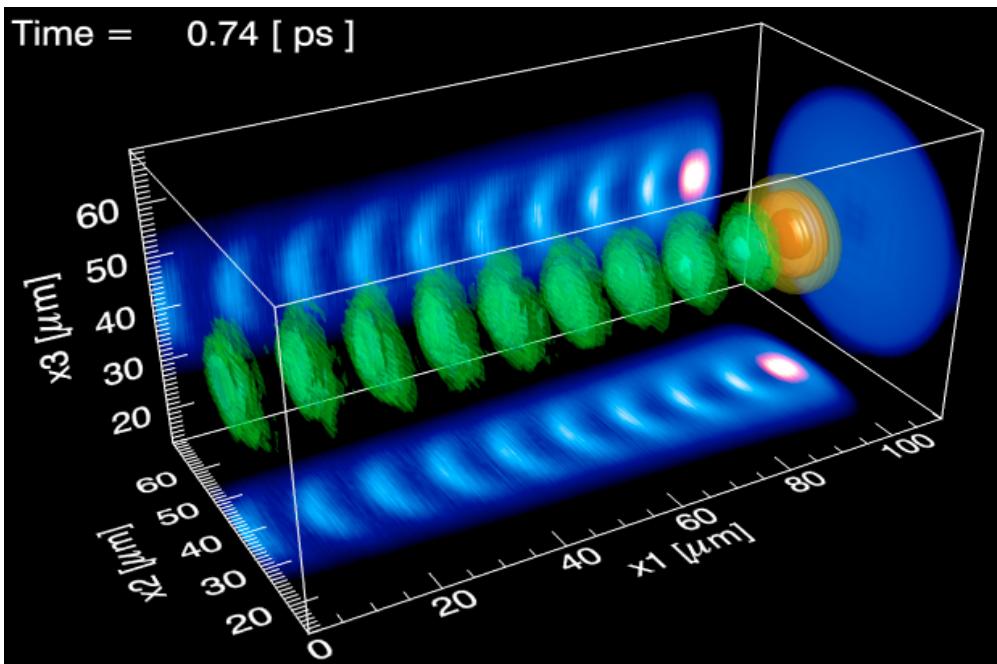
• The parameters are similar to those at LOA and LBNL

CERN – LEP schematic

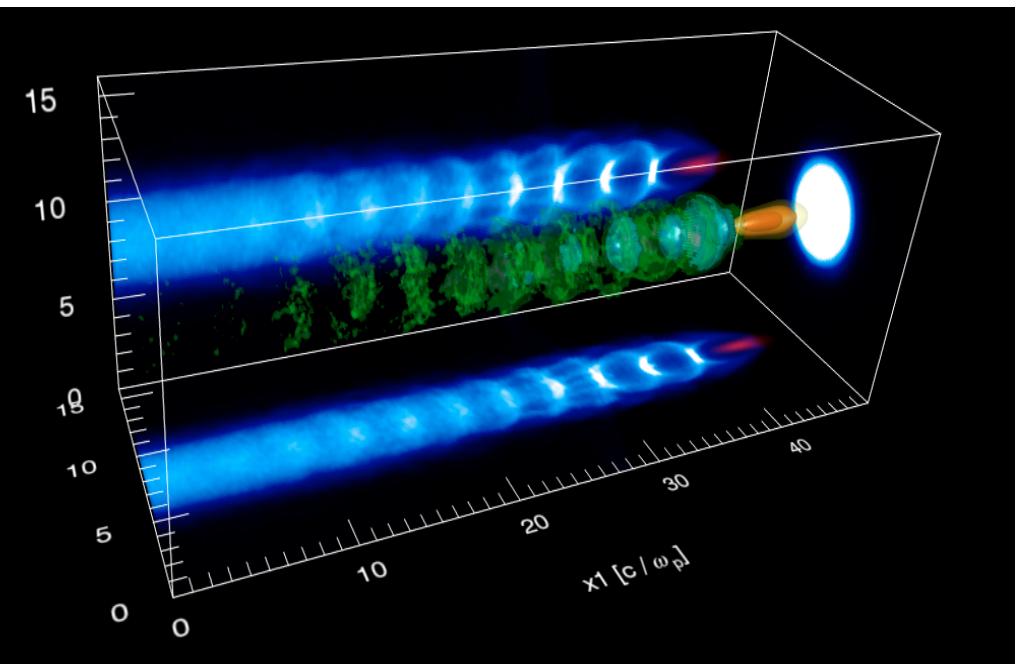


Laser & Electron Wakes

Nonlinear wakes are *similar* with laser or particle beam drivers:
3-D PIC OSIRIS Simulation
(self-ionized gas)



Laser Wake



Electron beam Wake

