



The Abdu Salam
International Centre for Theoretical Physics

United Nations
Educational, Scientific
and Cultural Organization

International Atomic
Energy Agency



SMR.1675 - 9

**Workshop on
Noise and Instabilities in Quantum Mechanics**

3 - 7 October 2005

Quantum noise and instabilities in cavity QED

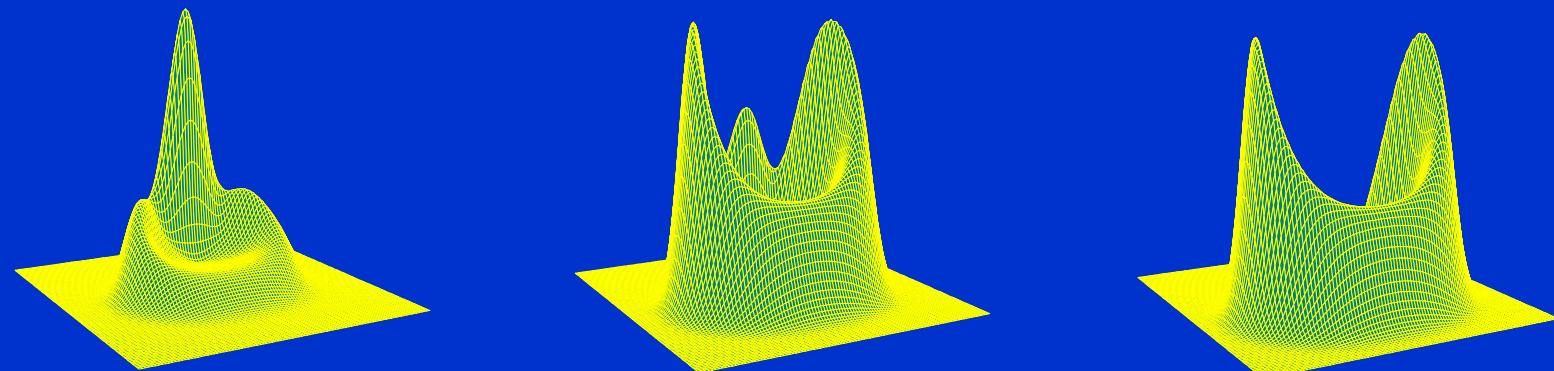
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Auckland
NEW ZEALAND

These are preliminary lecture notes, intended only for distribution to participants

Quantum Noise and Instabilities in Cavity QED

H. J. Carmichael

University of Auckland

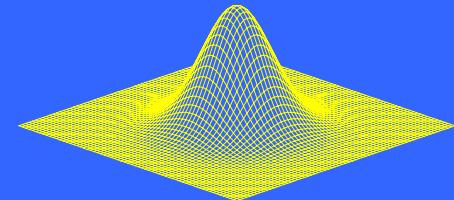


Support by the Marsden Fund of RSNZ



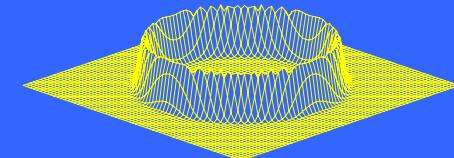
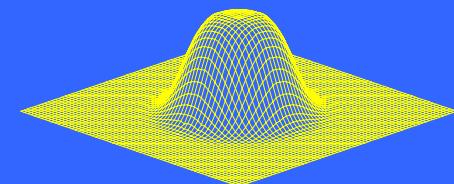
Background:

- quantum noise in the laser
- classical and non-classical fields
- squeezed light
- strong coupling/system size



Cavity QED:

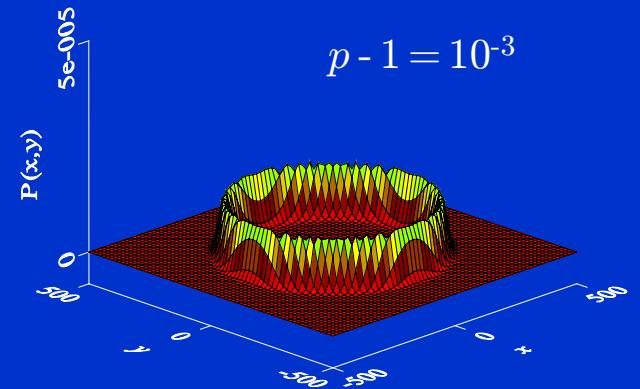
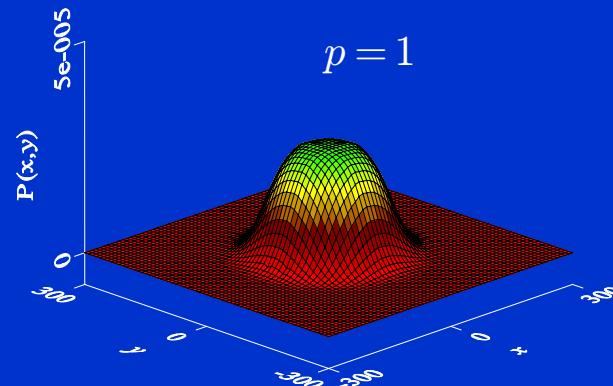
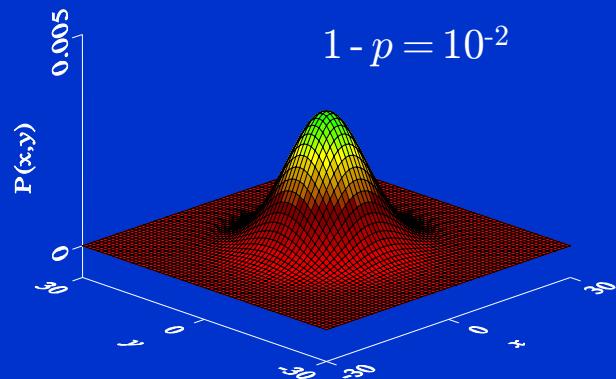
- optical bistability
- non-classical intensity noise
- squeezing/amplitude noise
- spontaneous dressed-state polarization



single-mode laser

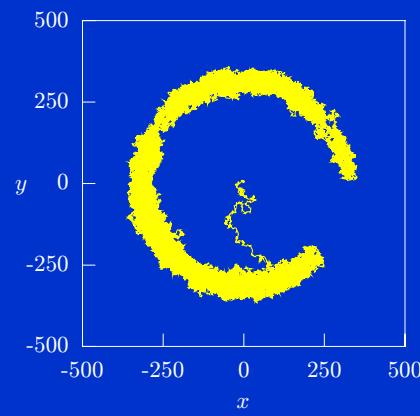
$$d\bar{\alpha} = -\bar{\alpha}(1-p + p|\bar{\alpha}|^2)dt + \sqrt{\frac{n_{\text{spon}}}{n_{\text{sat}}}}(dW_1 + i dW_2)$$

quantum noise

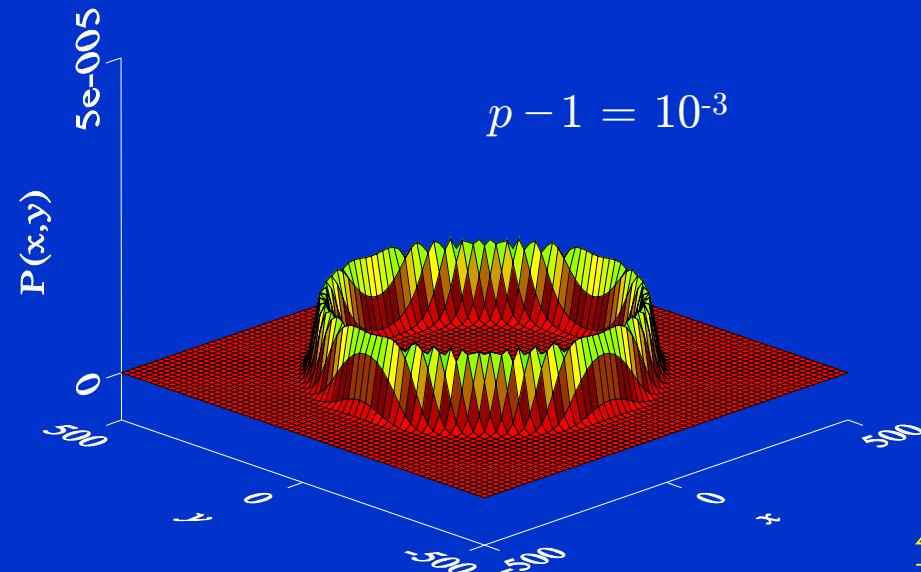


first-order coherence: laser linewidth

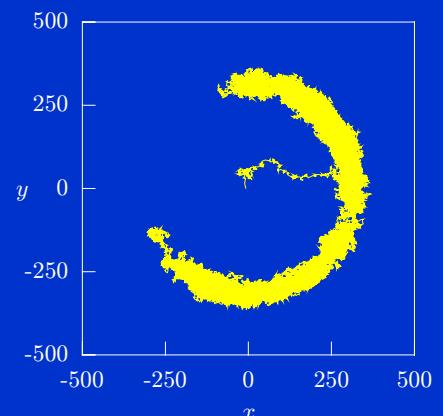
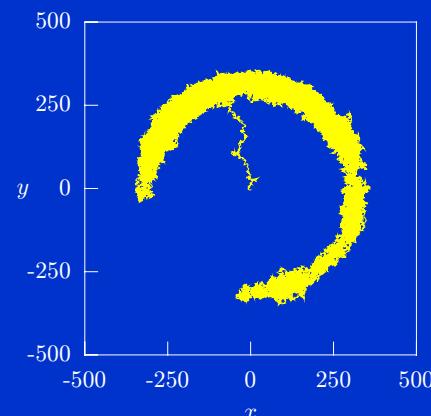
time = $10^6 \kappa^{-1}$



$$p-1 = 10^{-3}$$



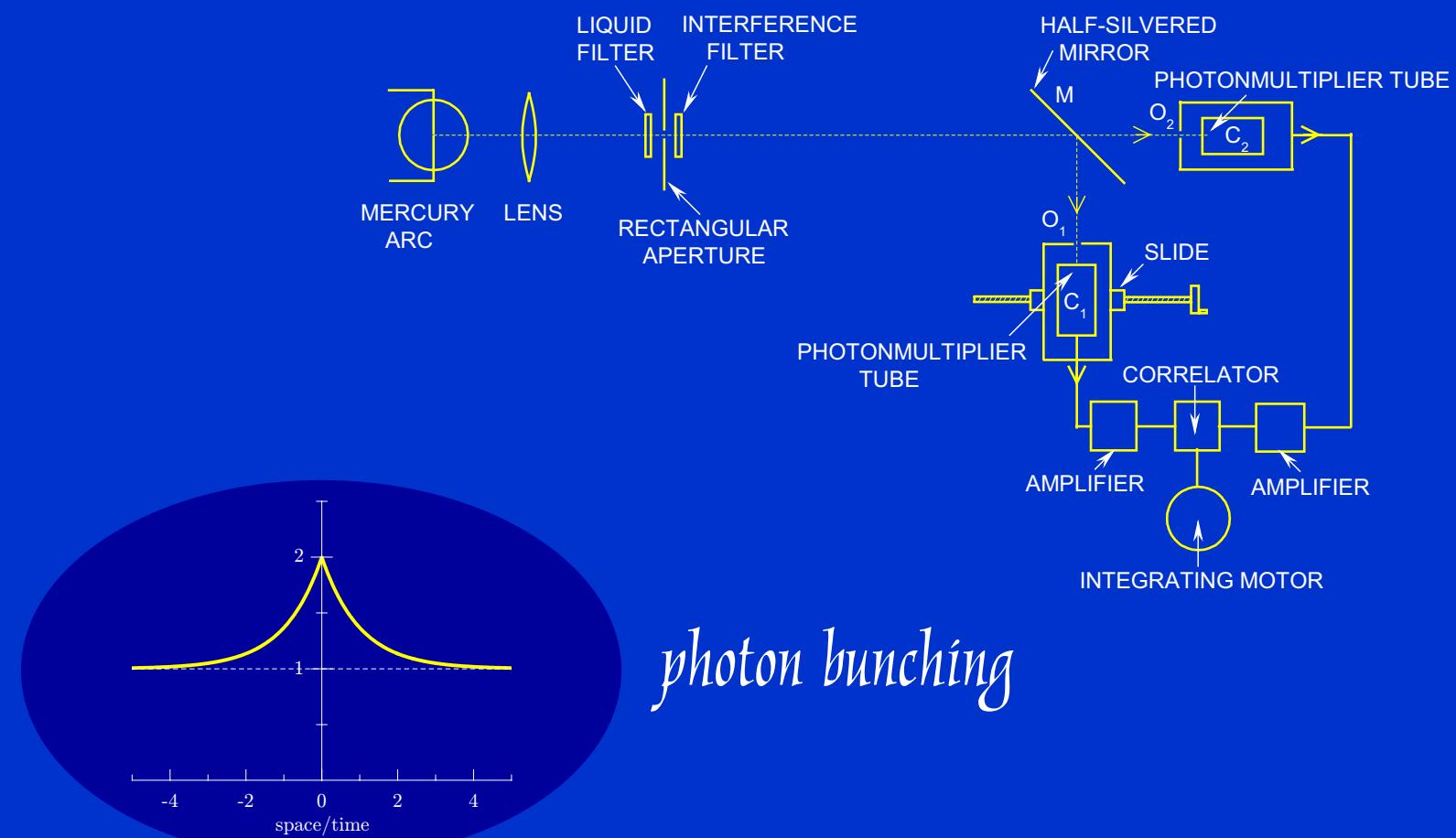
$$\frac{\Delta\omega}{\kappa} = \frac{n_{\text{spon}}}{2n_{\text{sat}}(p-1)}$$



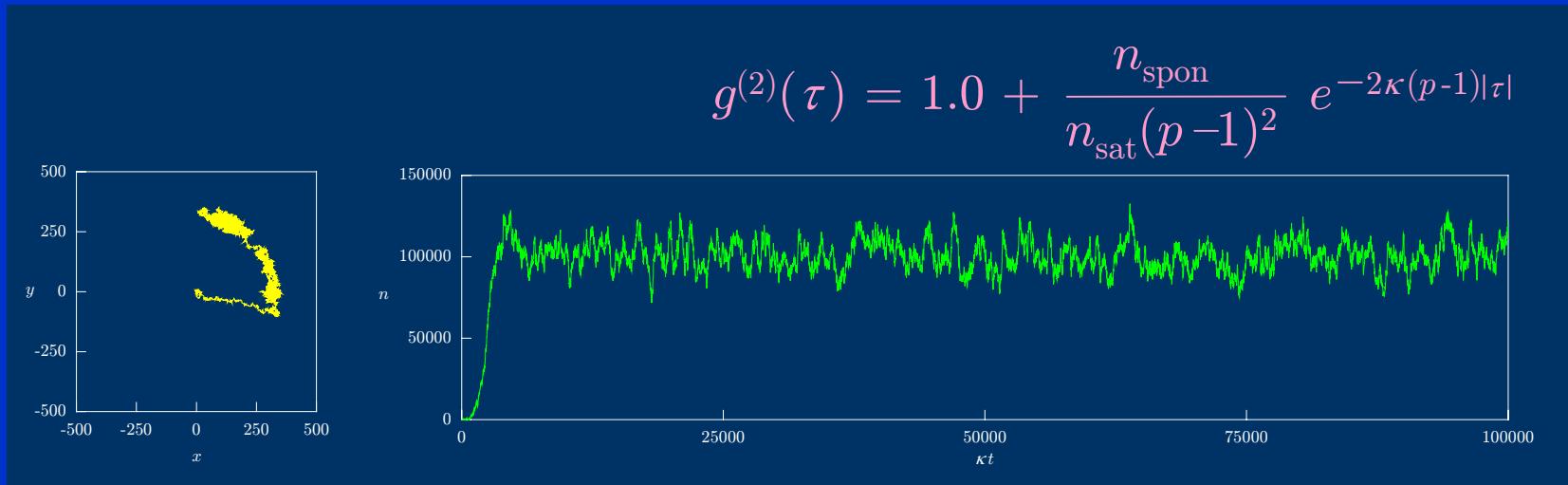
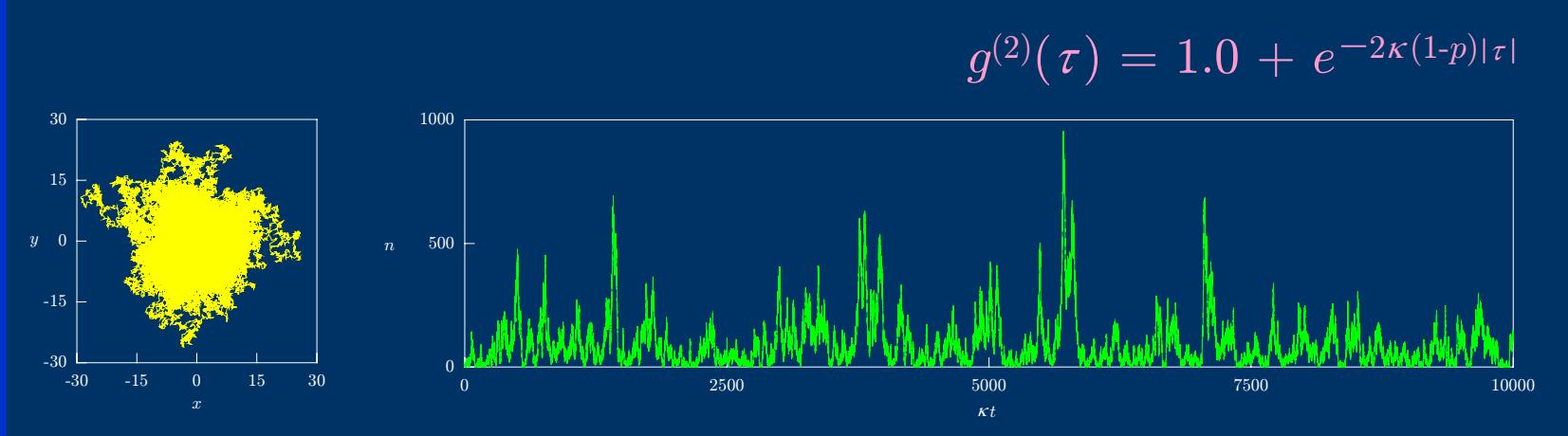
Hanbury Brown and Twiss light intensity interferometer

R. Hanbury Brown and R.Q. Twiss, Nature 177, 27 (1956)

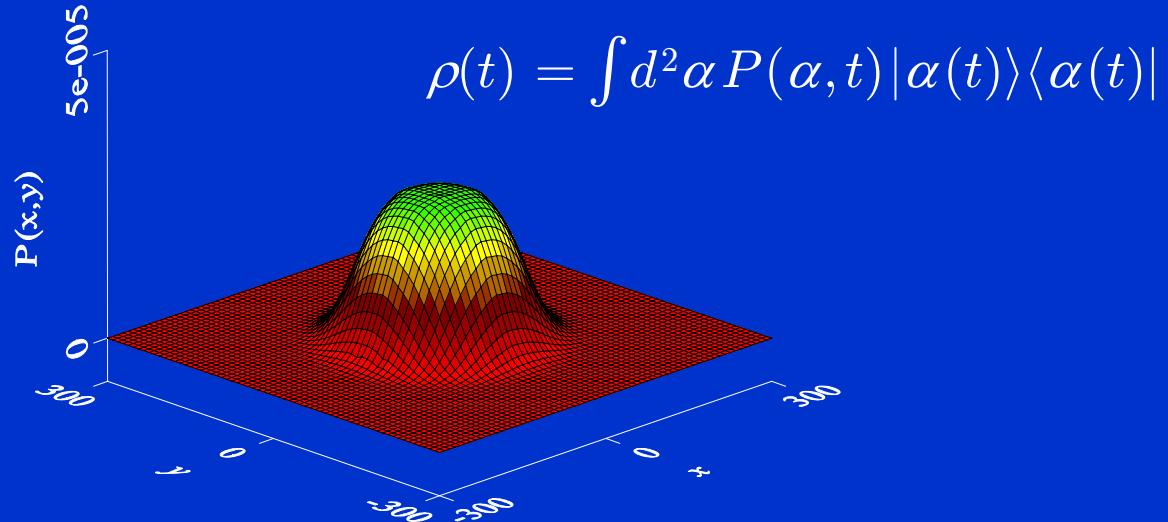
R.Q. Twiss, A.G. Little, and R. Hanbury Brown, Nature 180, 324 (1957)



second-order coherence

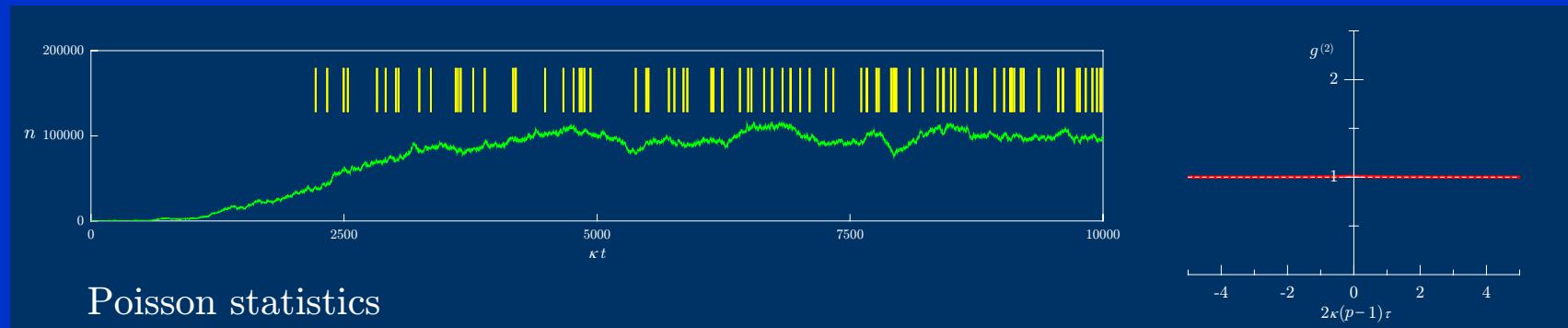
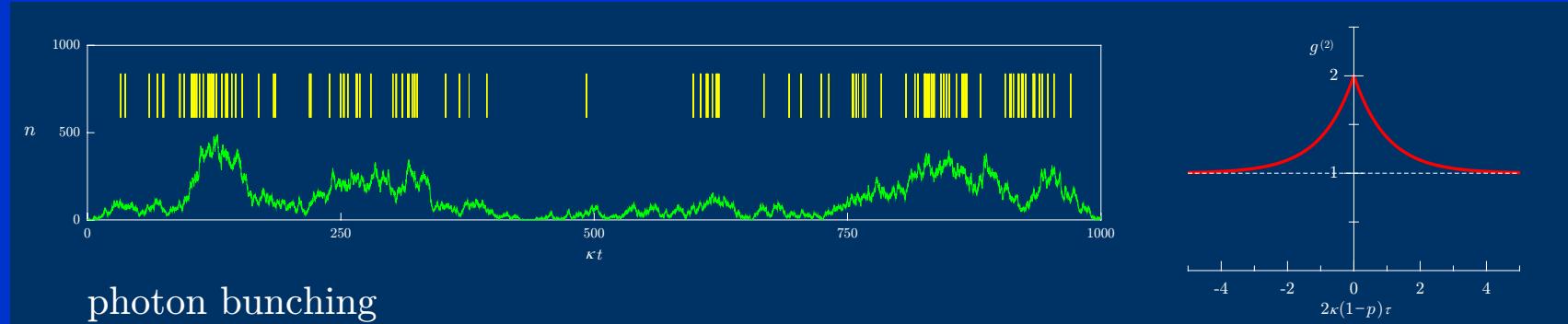


classical and non-classical fields



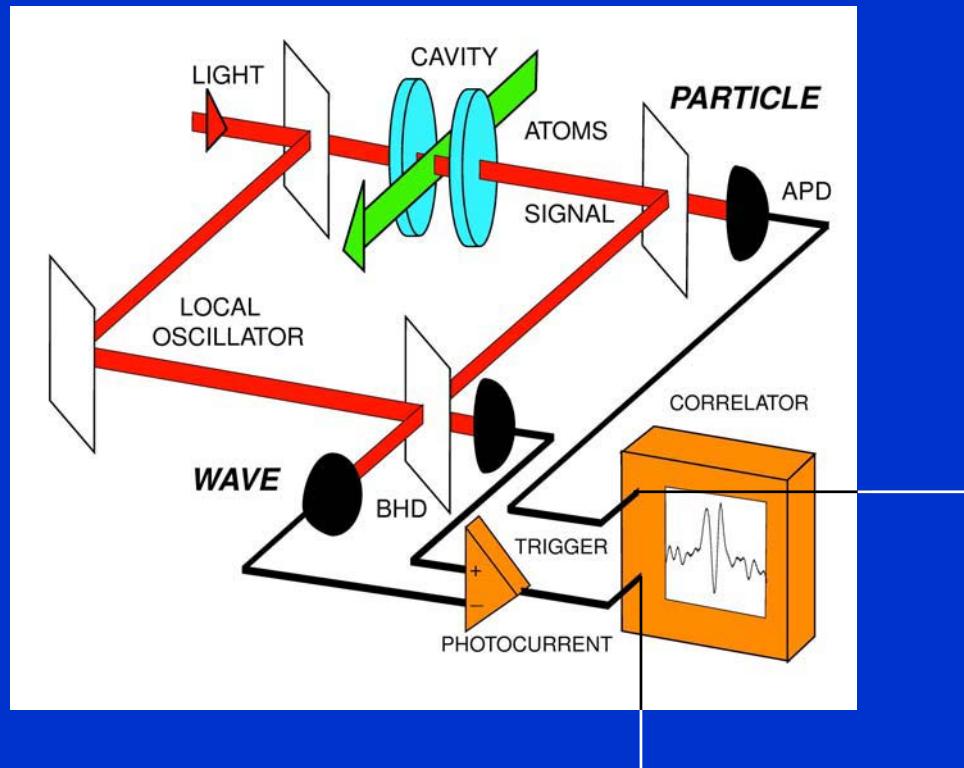
$$g^{(2)}(\tau) = \frac{\langle a^\dagger(t)a^\dagger(t+\tau)a(t+\tau)a(t) \rangle}{\langle a^\dagger(t)a(t) \rangle^2} = \frac{\langle |\alpha(t)|^2|\alpha(t+\tau)|^2 \rangle}{\langle |\alpha(t)|^2 \rangle^2}$$

└ quantum average └ time average

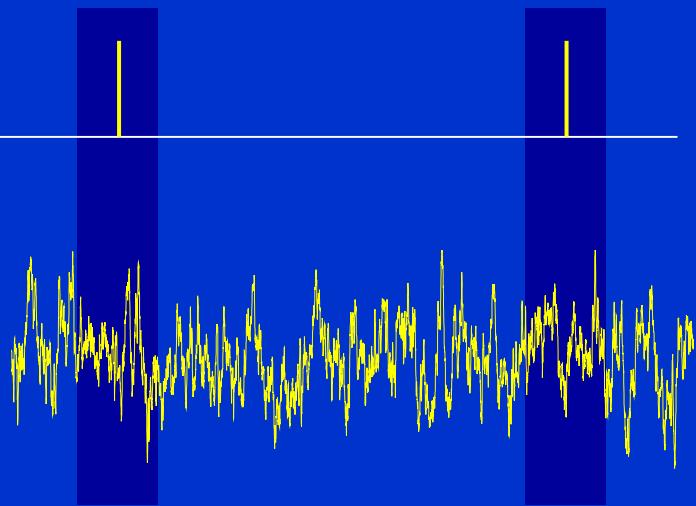


conditional homodyne detection

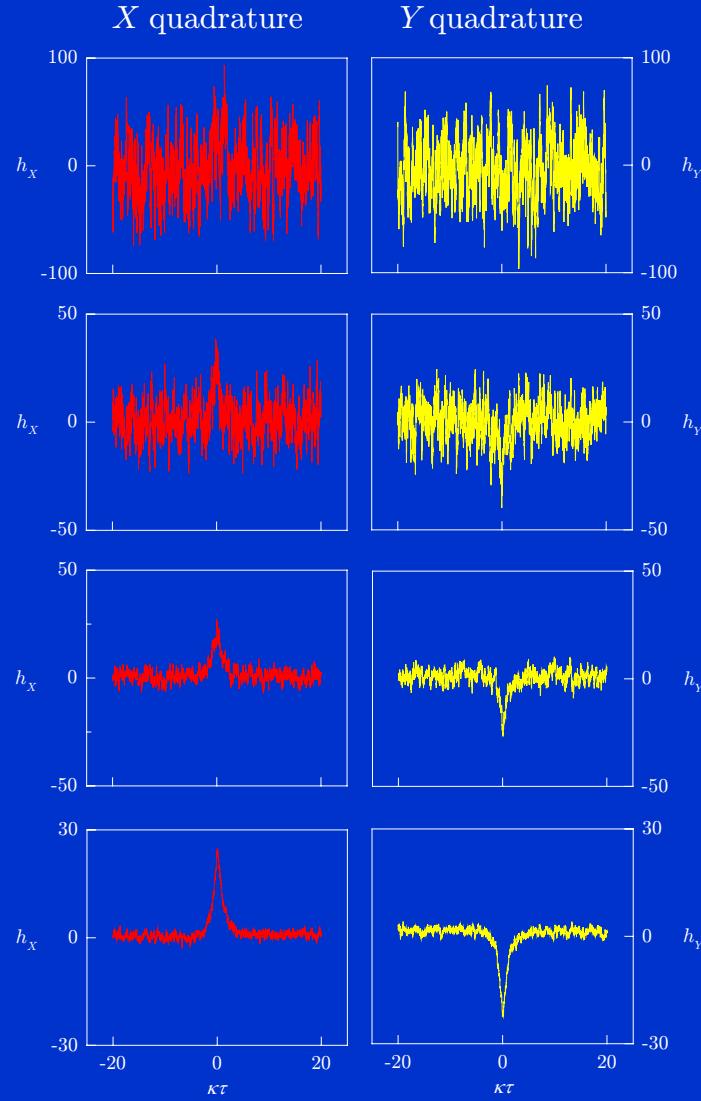
G.T. Foster, L.A. Orozco, H.M. Castro-Beltran, and H.J. Carmichael,
Phys. Rev. Lett. 85, 3149 (2000)



"start" 2



squeezed light



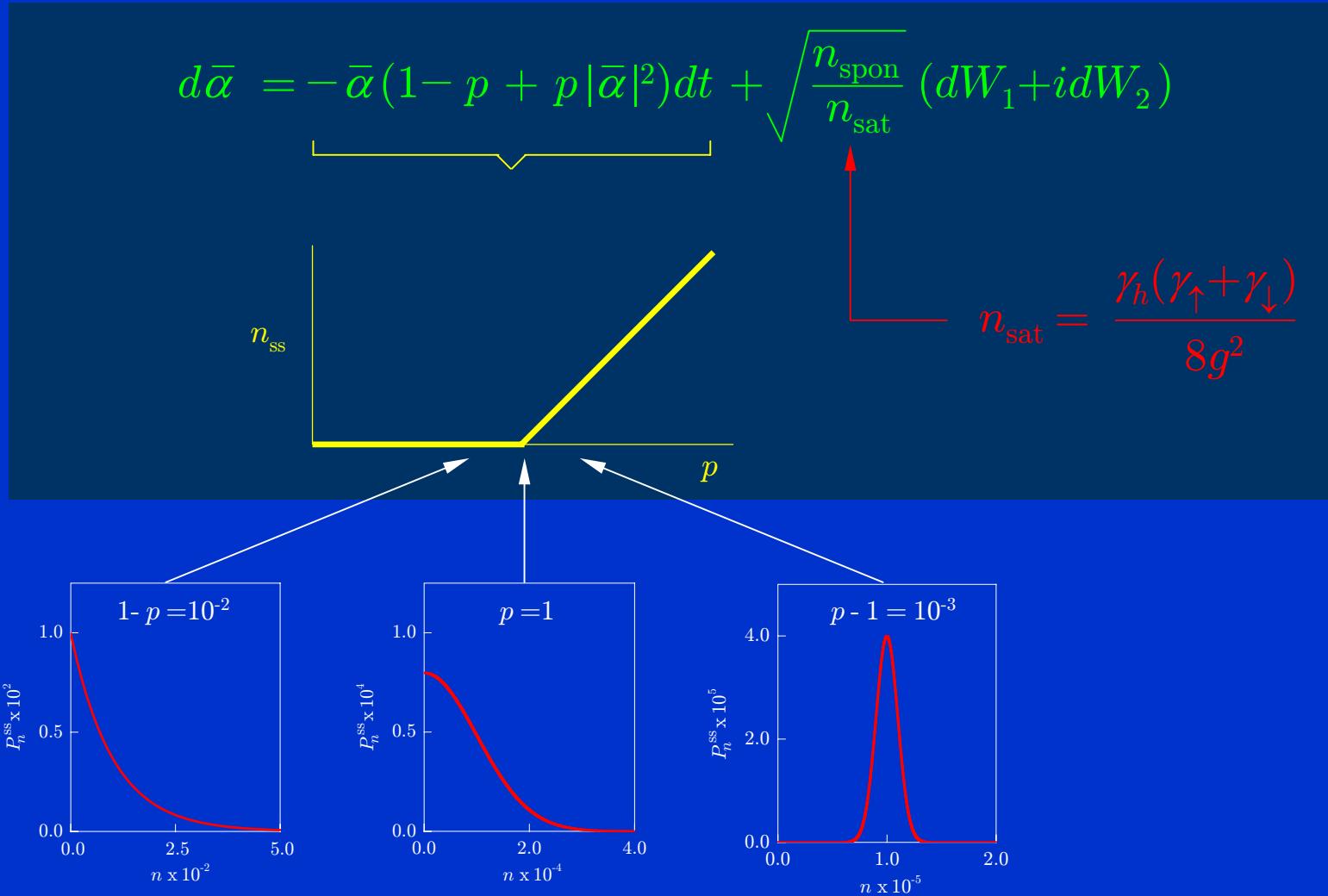
$$h_X(\tau)-1 = \frac{\langle : \hat{X}(t) \hat{X}(t+\tau) : \rangle}{\langle \hat{a}^\dagger(t) \hat{a}(t) \rangle} = \frac{\langle X(t) X(t+\tau) \rangle}{\langle |\alpha(t)|^2 \rangle}$$

$\frac{a+a^\dagger}{2}$ $\frac{\alpha+\alpha^*}{2}$

$$h_Y(\tau)-1 = \frac{\langle : \hat{Y}(t) \hat{Y}(t+\tau) : \rangle}{\langle \hat{a}^\dagger(t) \hat{a}(t) \rangle} = \frac{\langle Y(t) Y(t+\tau) \rangle}{\langle |\alpha(t)|^2 \rangle}$$

$\frac{a-a^\dagger}{2i}$ $\frac{\alpha-\alpha^*}{2i}$

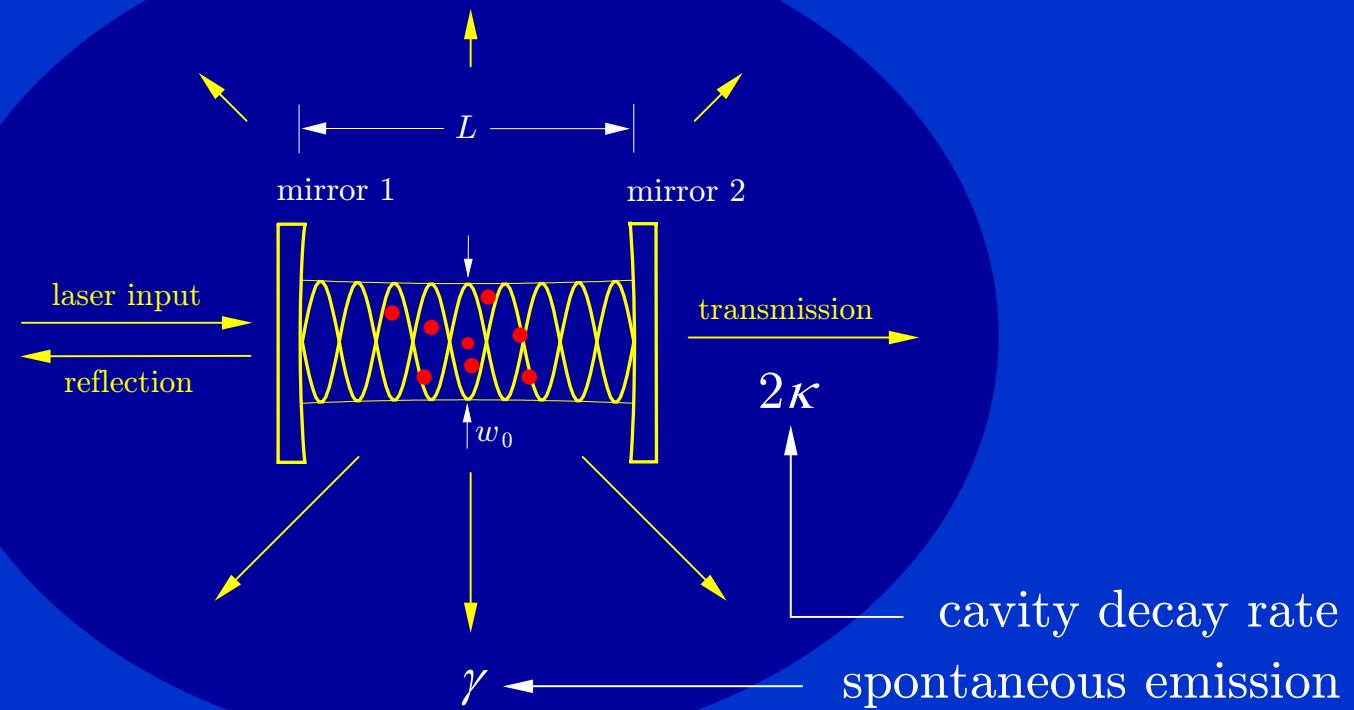
strong coupling and system size



cavity QED

dipole coupling strength:

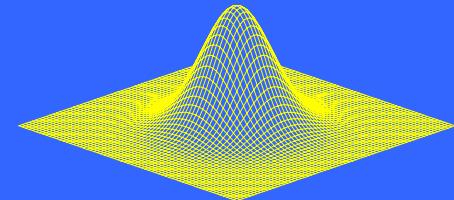
$$g = \sqrt{\frac{\omega_C d^2}{2\hbar\epsilon_0 V}}$$





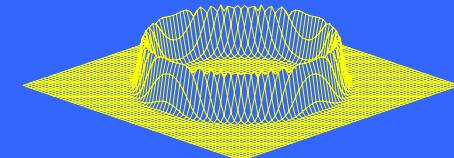
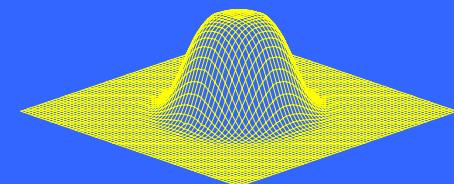
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Cavity QED:

- optical bistability
- non-classical intensity noise
- squeezing/amplitude noise
- spontaneous dressed-state polarization



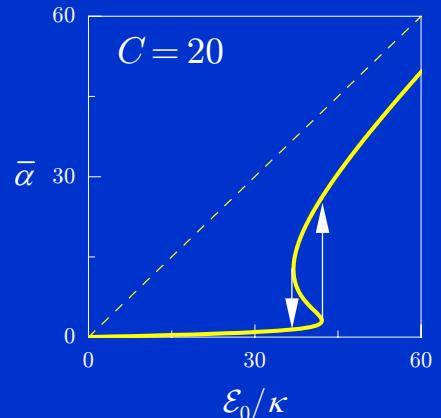
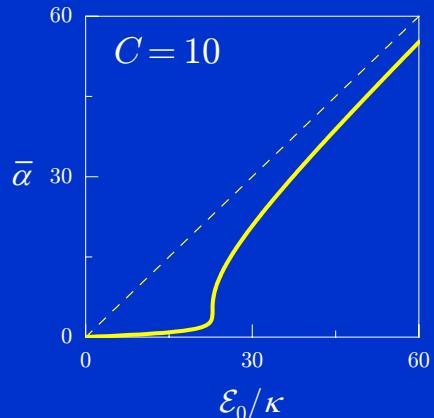
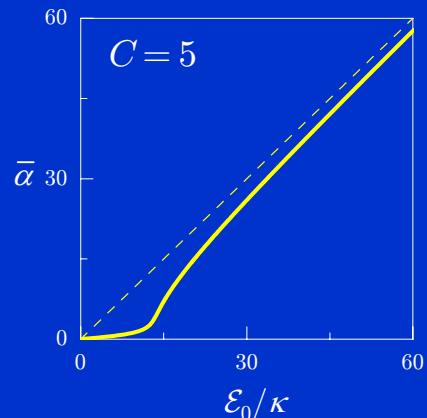
optical bistability

$$\dot{\bar{\alpha}} = -\kappa \bar{\alpha} + \sum_j g(\mathbf{r}_j) \bar{\beta}_j + \bar{\mathcal{E}}_0$$

$$\dot{\bar{\beta}}_j = -\frac{\gamma}{2} \bar{\beta}_j + g(\mathbf{r}_j) \bar{\alpha} \bar{m}_j$$

$$\dot{\bar{m}}_j = -\gamma (\bar{m}_j + 1) - 2g(\mathbf{r}_j) (\bar{\beta}_j^* \bar{\alpha} + \text{c.c.})$$

$$\begin{aligned} & \downarrow \text{input field} \\ \mathcal{E}_0/\kappa &= f(\bar{\alpha}; C) \\ & \uparrow \\ 2C &= \frac{\alpha_0 l}{1-R} \end{aligned}$$



quantum trajectories

non-unitary Schrödinger evolution:

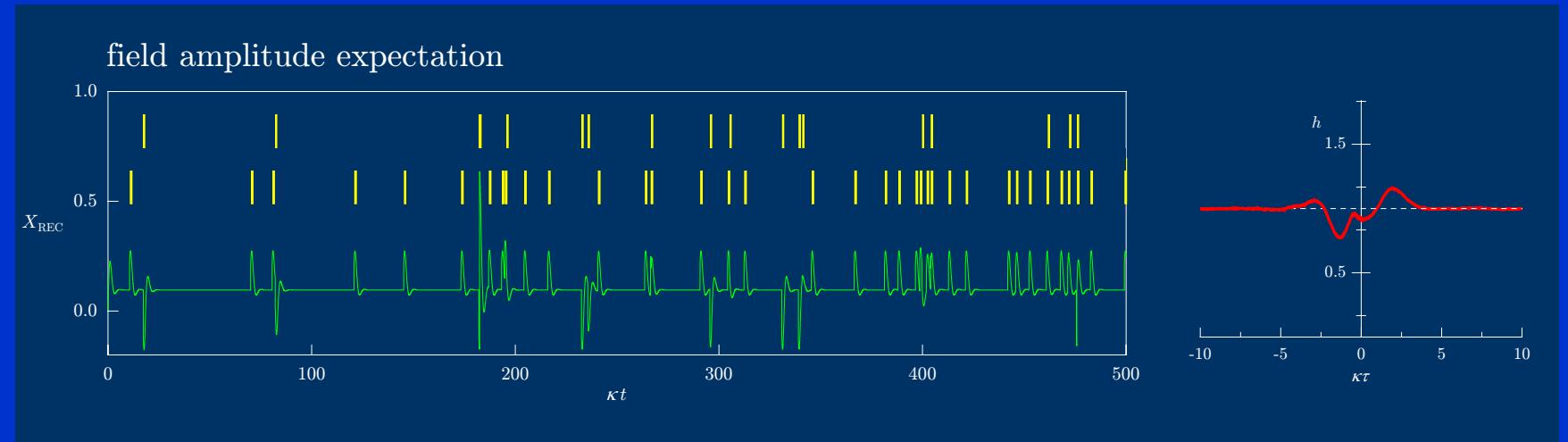
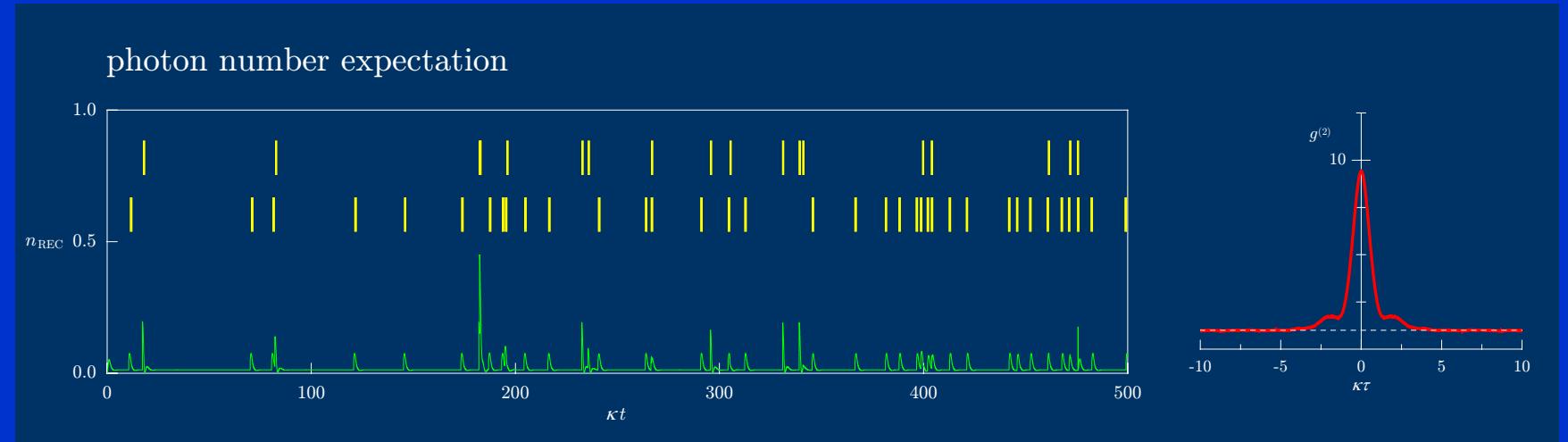
$$\frac{d|\bar{\psi}_{\text{REC}}\rangle}{dt} = \frac{1}{i\hbar} \hat{H}_B(t) |\bar{\psi}_{\text{REC}}\rangle$$

$$\frac{1}{i\hbar} \hat{H}_B(t) = \underbrace{\mathcal{E}(\hat{a}^\dagger - \hat{a})}_{\text{laser input}} + \underbrace{\sum_j g(r_j(t))(\hat{a}^\dagger \hat{\sigma}_{j-} - \hat{a} \hat{\sigma}_{j+})}_{\text{dipole interaction}} - \kappa \hat{a}^\dagger \hat{a} - \frac{\gamma}{2} \sum_j \hat{\sigma}_{j+} \hat{\sigma}_{j-}$$

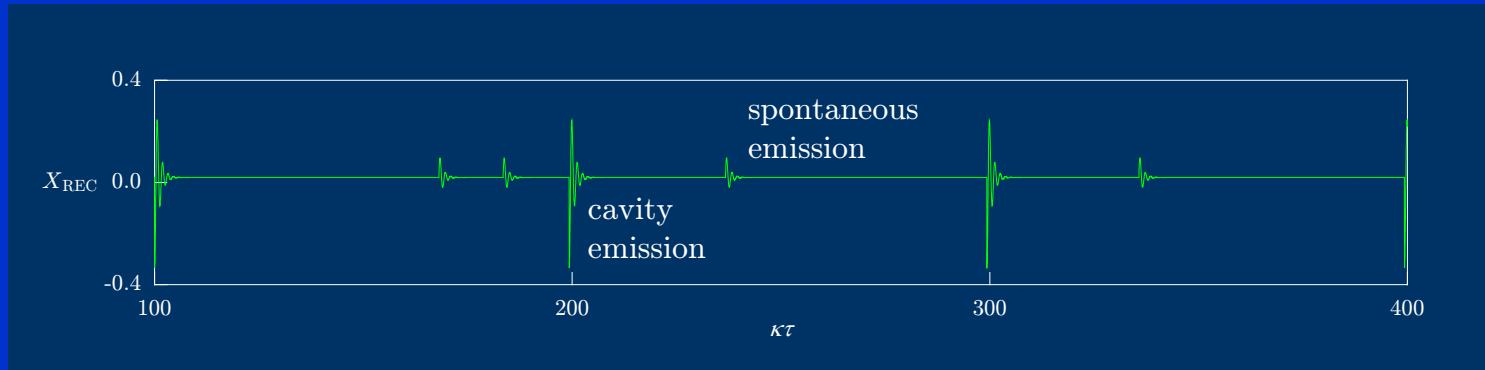
↑ spontaneous emission
cavity loss

quantum jumps:

$$|\bar{\psi}_{\text{REC}}\rangle \longrightarrow \begin{cases} \hat{a} |\bar{\psi}_{\text{REC}}\rangle & \text{at rate } 2\kappa \langle (\hat{a}^\dagger \hat{a})(t) \rangle_{\text{REC}} \\ \hat{\sigma}_{j-} |\bar{\psi}_{\text{REC}}\rangle & \text{at rate } \gamma \langle (\hat{\sigma}_{j+} \hat{\sigma}_{j-})(t) \rangle_{\text{REC}} \end{cases}$$



weak-excitation limit



$$g_{\{r_j\}}^{(2)}(\tau) = \left[h_{\{r_j\}}(\tau) \right]^2$$

$$h_{\{r_j\}}(\tau) = 1.0 - A_{\{r_j\}} e^{-\frac{1}{2}(\kappa+\gamma/2)|\tau|} \left[\cos(\Omega_{\{r_j\}} \tau) + \frac{\frac{1}{2}(\kappa+\gamma/2)}{\Omega_{\{r_j\}}} \sin(\Omega_{\{r_j\}} \tau) \right]$$

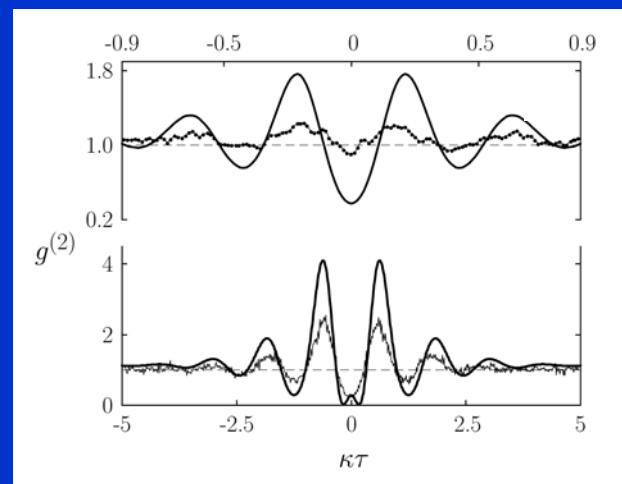
\uparrow

$$\Omega_{\{r_j\}} = \sqrt{\sum_j g^2(r_j) - \frac{1}{4}(\kappa-\gamma/2)^2}$$

experiments 1

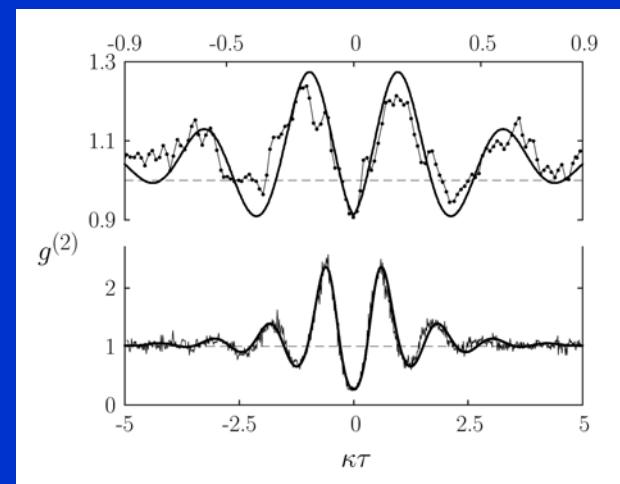
G. Rempe, R.J. Thompson, R.J. Brecha, W.D. Lee, and H.J. Kimble,
Phys. Rev. Lett. **67**, 1727 (1991)

G.T. Foster, S.L. Mielke, and L.A. Orozco, Phys. Rev. A **61**, 053821 (2000)



Rempe *et al.*

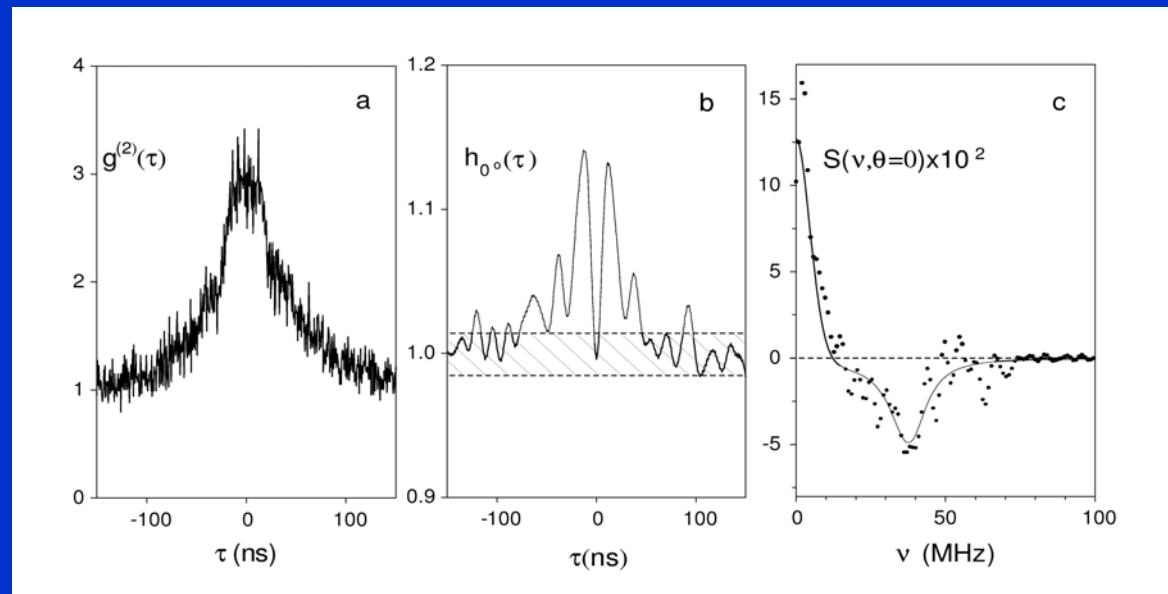
Foster *et al.*



non-classical intensity noise

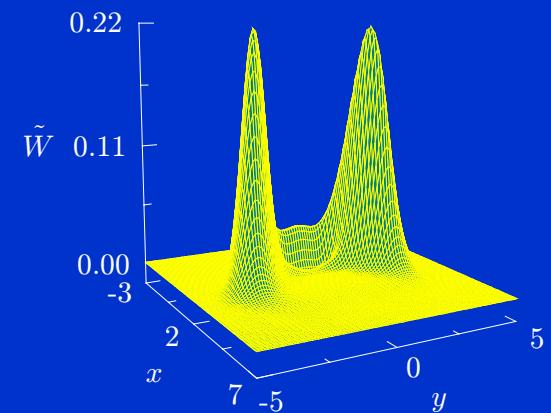
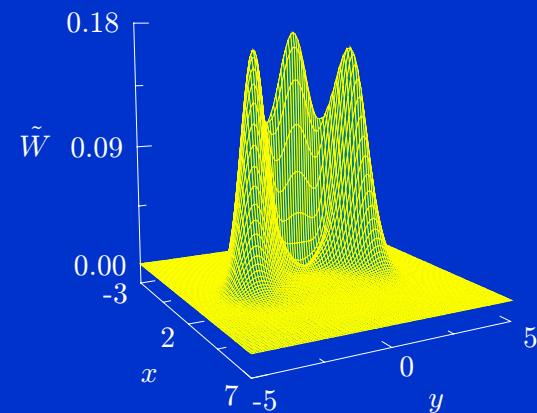
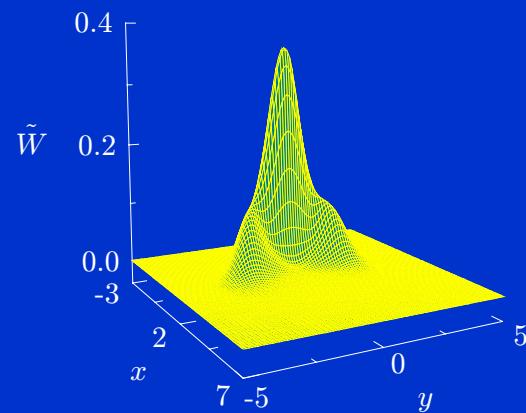
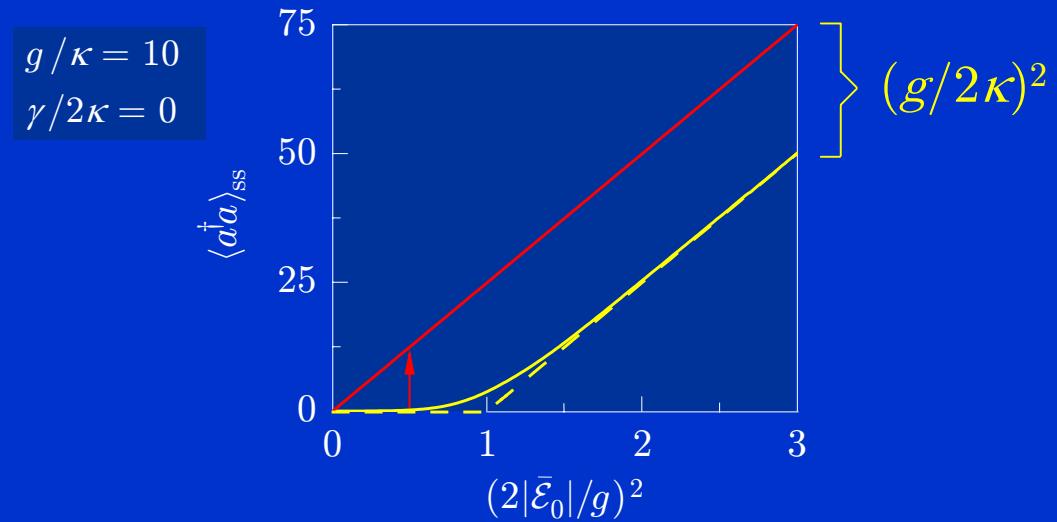
experiments 2

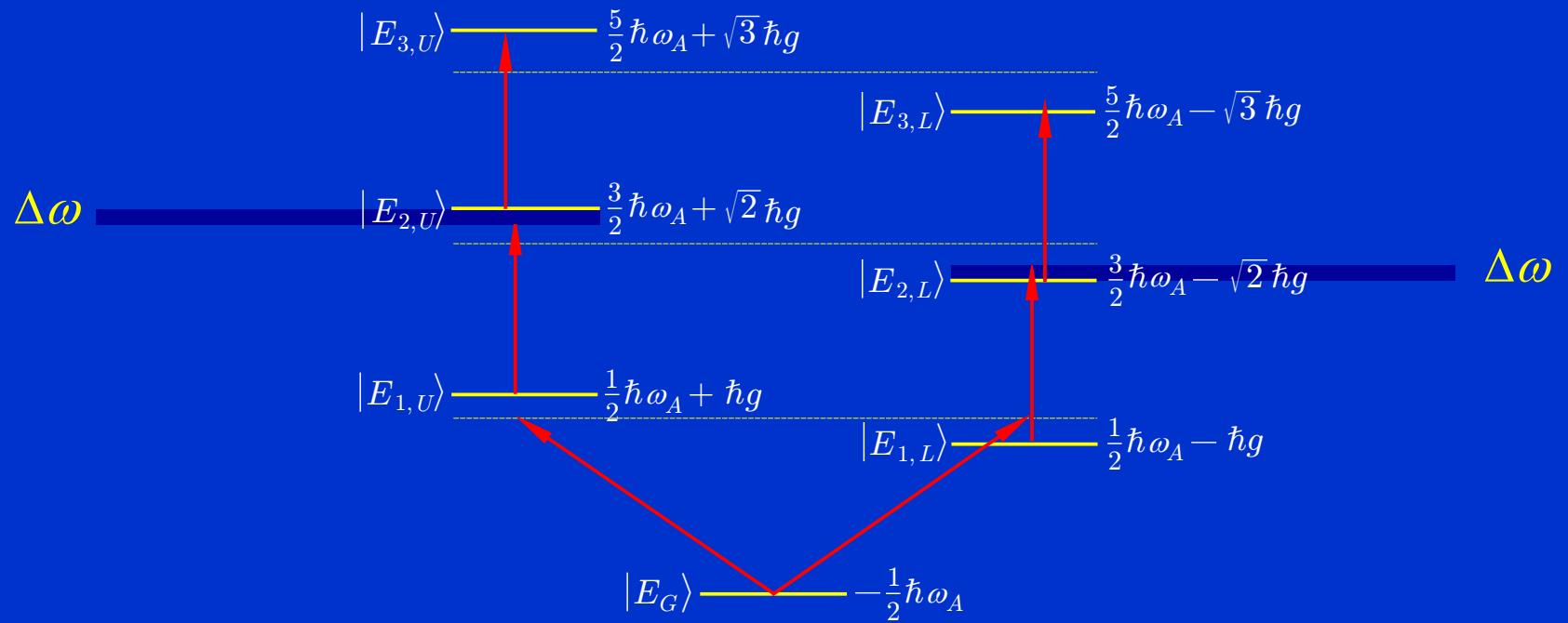
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Phys. Rev. Lett 85, 3149 (2000)



squeezing/amplitude noise

spontaneous dressed-state polarization





$$\Delta\omega = \pm(\sqrt{n+1} - \sqrt{n})g \approx \pm \frac{g}{2\sqrt{n}}$$

$$\bar{n} \left[1 - \frac{1}{1 + (\Delta\omega/\kappa)^2} \right] \approx \bar{n}(\Delta\omega/\kappa)^2 = (g/2\kappa)^2$$

spontaneous emission

$$\begin{aligned}g/\kappa &= 10 \\ \gamma/2\kappa &= 1\end{aligned}$$

