



The Abdus Salam  
International Centre for Theoretical Physics



SMR 1698/3

## WORKSHOP ON PLASMA PHYSICS

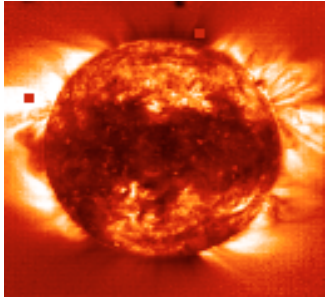
7 - 11 March 2005

# Measurements of Plasma Opacity using X-Ray Lasers

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These are preliminary lecture notes, intended only for distribution to participants.



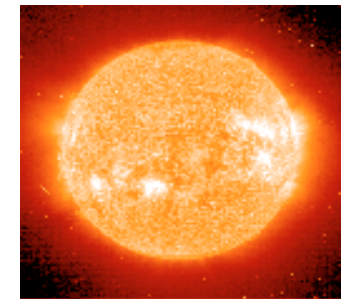
FeX11/FeX1 ratio

# Using lasers to measure the solar opacity

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University of York



He II emission

# Acknowledgements

- I would like to acknowledge members of my research group: Matthew Edwards, Pritesh Mistry, David Whittaker and the team at the Prague Asterix Laser headed by Bedrich Rus.
- Funding from EU Access to Large Scale Facilities and EPSRC.

# Outline

- Explanation of opacity in plasmas.
- Previous laboratory methods to measure opacity.
- Why opacity measurements are needed-relevant to solar plasmas?
- Opacity measurements using x-ray lasers.
- Conclusions.

# The equation of radiative transfer

- For a beam of light of single frequency, the change in light intensity  $dI_\nu$  in passing through a slab of matter of thickness  $ds$  is

$$- \frac{dI_\nu}{ds} = -k_\nu I_\nu + j_\nu$$

- where  $j_\nu$  is the emissivity (and scatter into the beam) and  $k_\nu$  is the total opacity at the particular frequency.

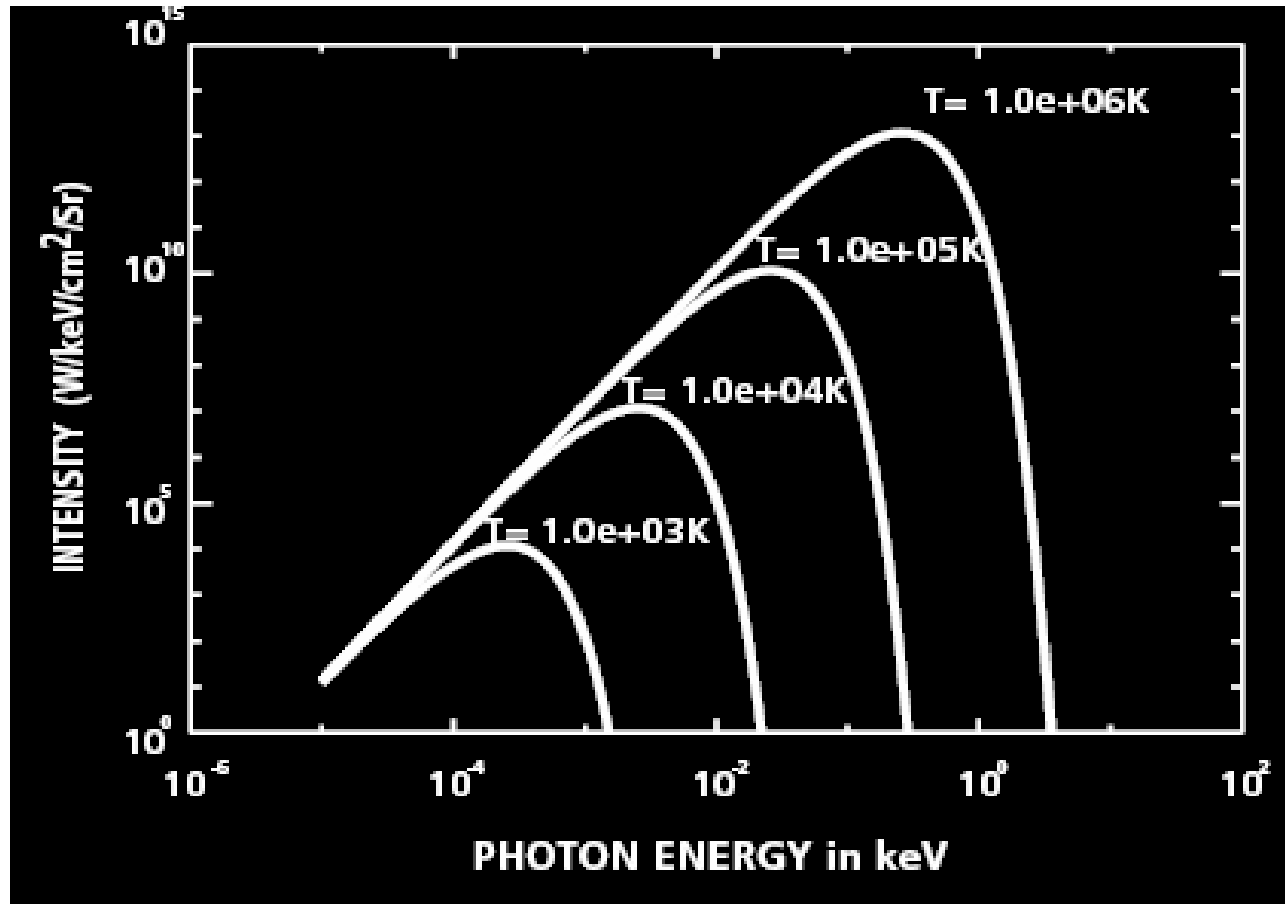
- In modelling solar and stellar energy flow (and other thermal plasmas), a frequency averaged measure of opacity  $k$  is used known as the Rosseland mean opacity so that radiation outflow can be modelled as a diffusion process

$$\frac{1}{k} = \frac{\int_0^\infty \frac{1}{k_\nu} \frac{dB_\nu}{dT} d\nu}{\int_0^\infty \frac{dB_\nu}{dT} d\nu} = \int_0^\infty \frac{1}{k_\nu} f(u) du$$

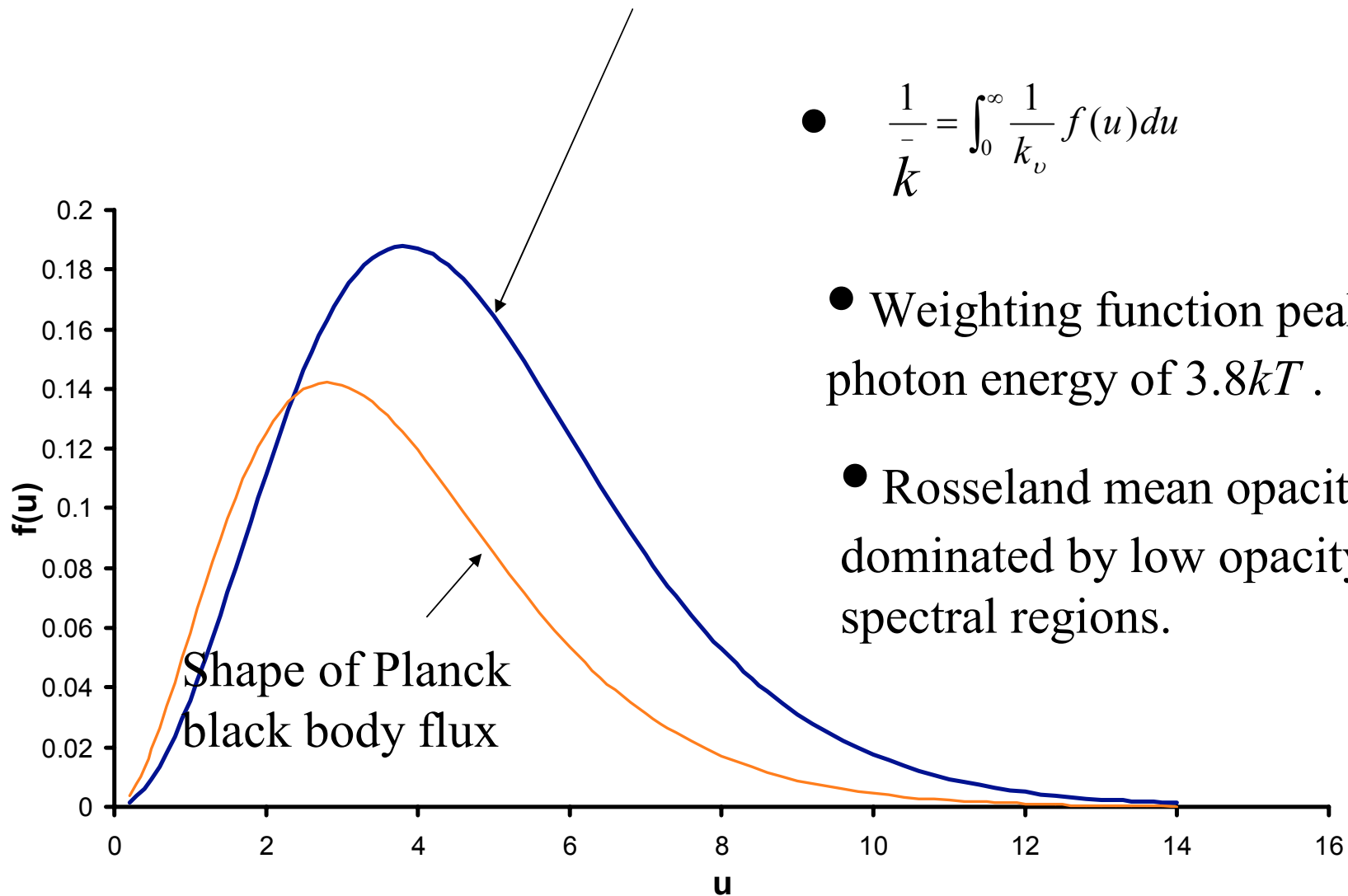
- where  $B_\nu$  is the black-body radiation distribution,  $f(u)$  is a weighting function and  $u = h\nu/kT$ .

# Black body radiation

- The radiation field when radiation is in equilibrium with matter.
- The peak of intensity is at photon energy  $2.8 kT$
- 11605 degrees Kelvin = 1 eV.



# The Rosseland mean opacity weighting

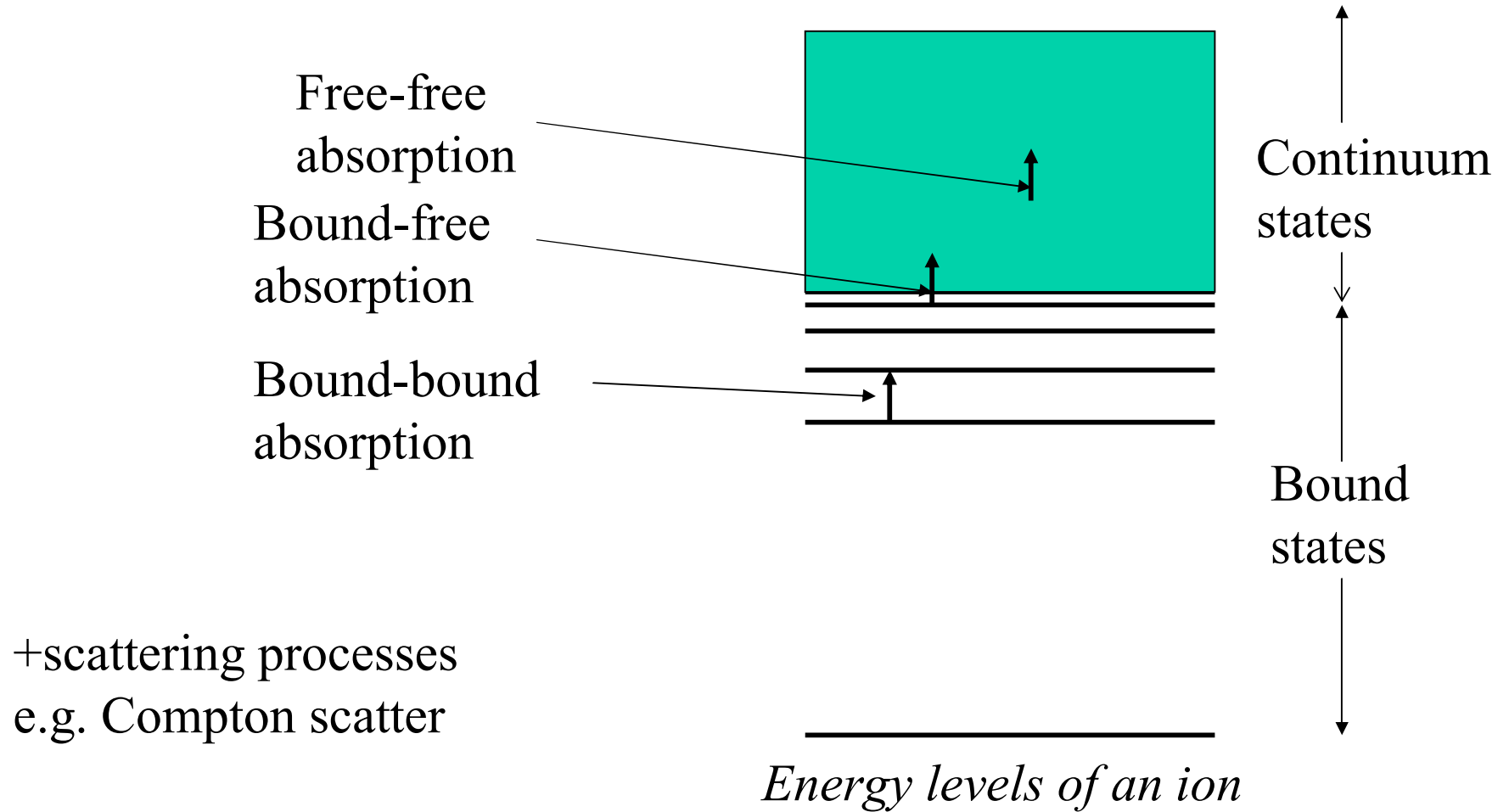


- $\frac{1}{\bar{k}} = \int_0^{\infty} \frac{1}{k_{\nu}} f(u) du$

- Weighting function peaks at photon energy of  $3.8kT$ .

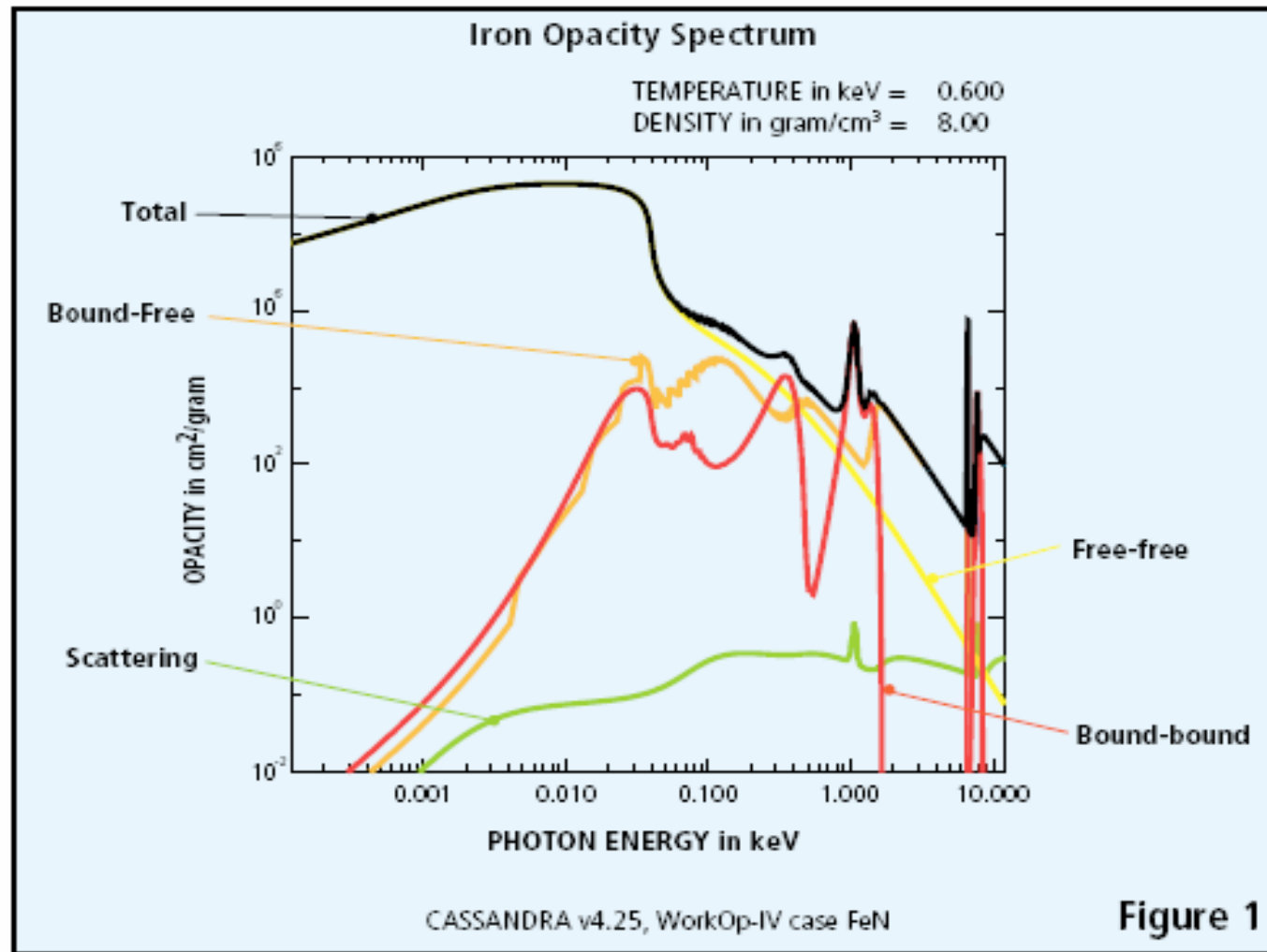
- Rosseland mean opacity is dominated by low opacity  $k_{\nu}$  spectral regions.

# Important photon absorption processes





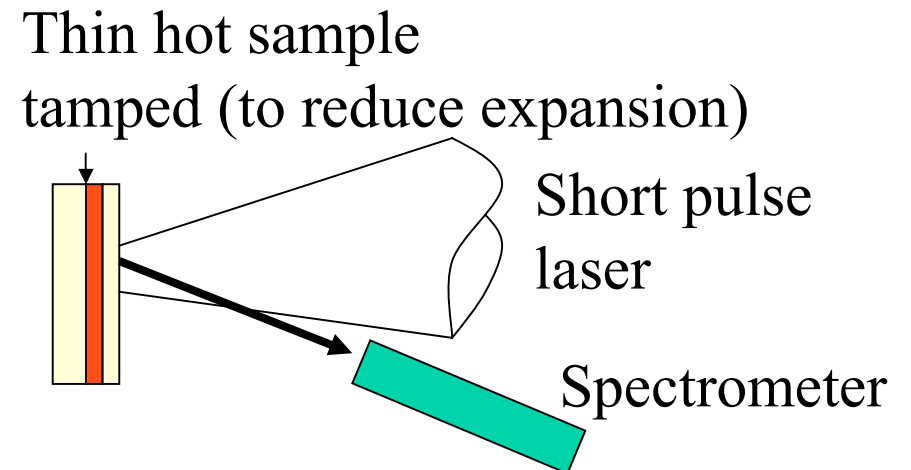
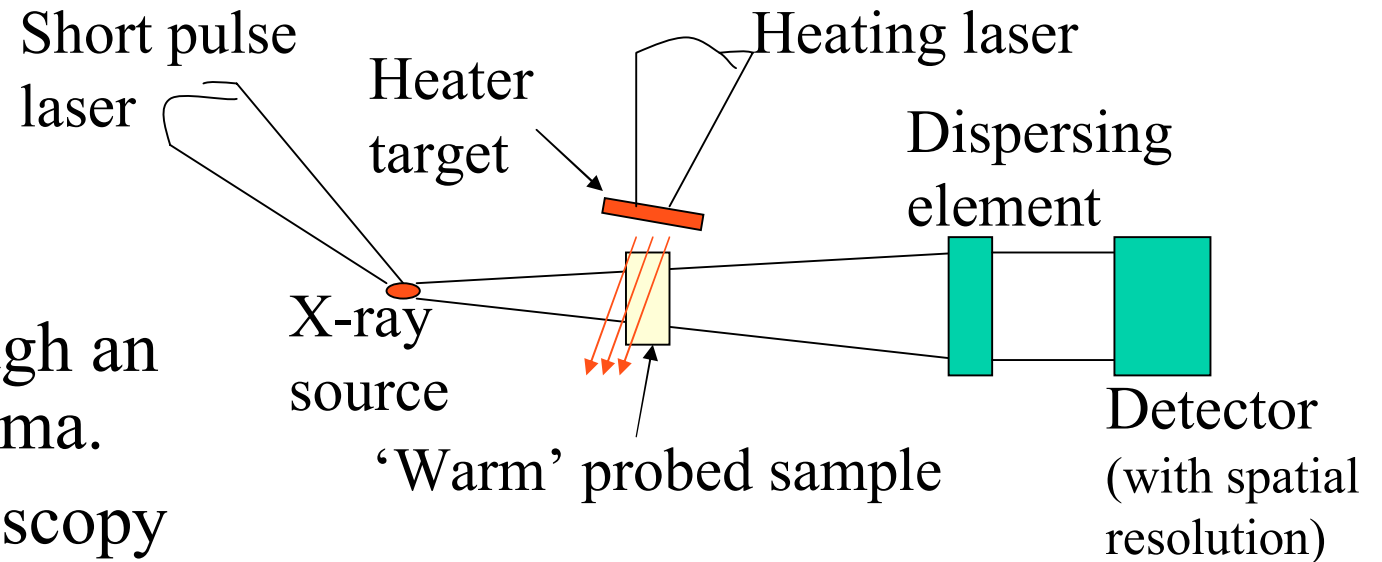
# Processes contributing to opacity



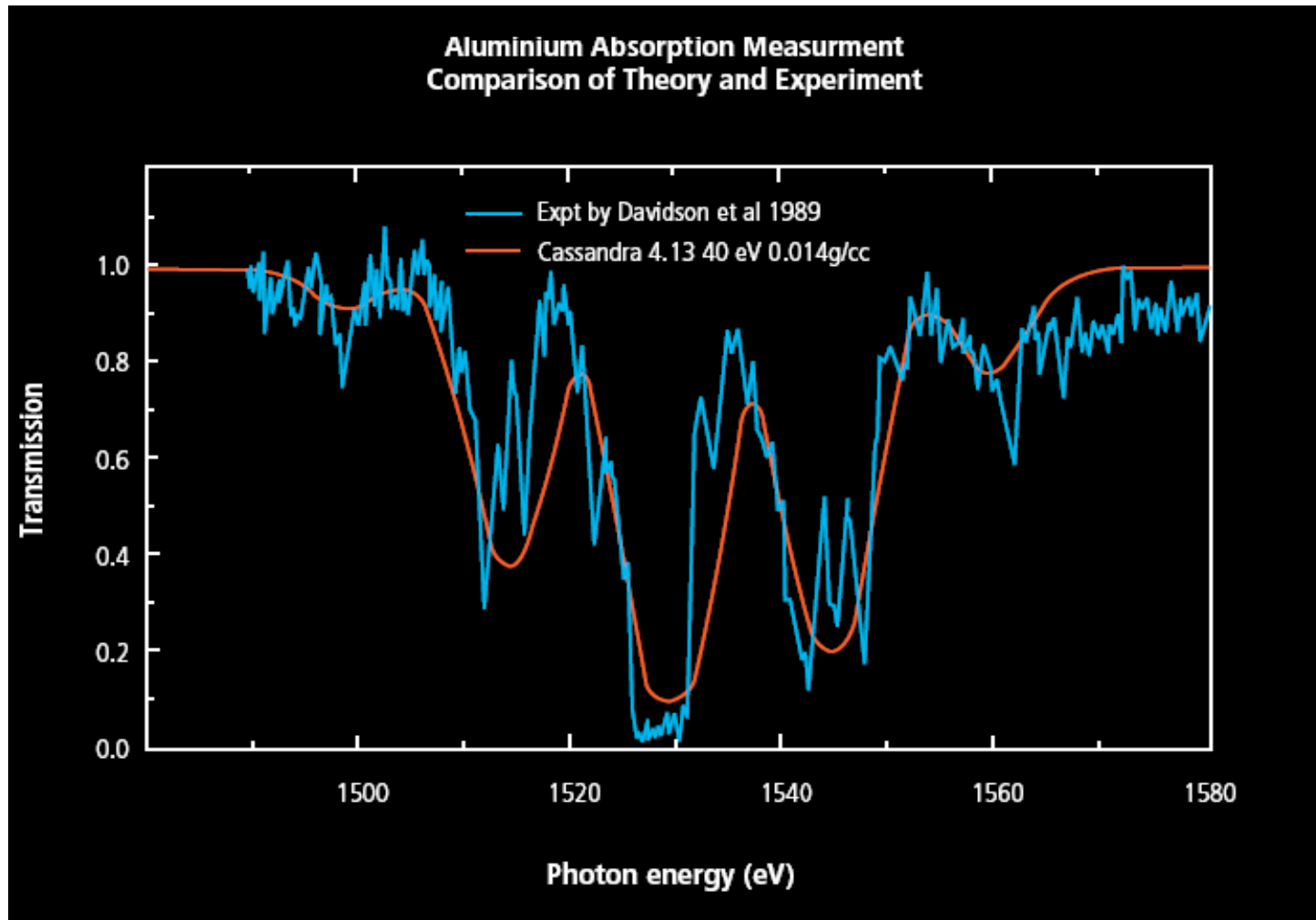
Taken from AWE website

# Previous measurements of plasma opacity in the laboratory

- Shine a bright backlighter through an x-ray heated plasma.
- Emission spectroscopy can be used:
  - Assume plasma is in LTE
  - Calculate opacity from emissivity of a thin sample.



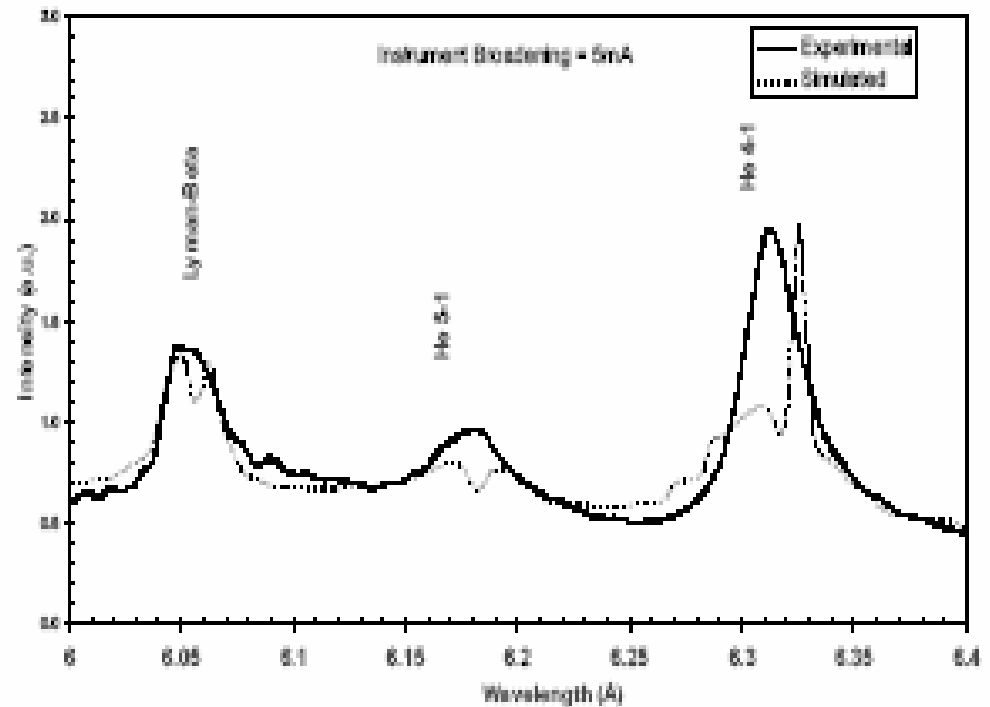
# Sample x-ray backlighter measurement of opacity



Plasma  
Conditions:  
4.13 eV and  
0.014 g cm<sup>-3</sup>.

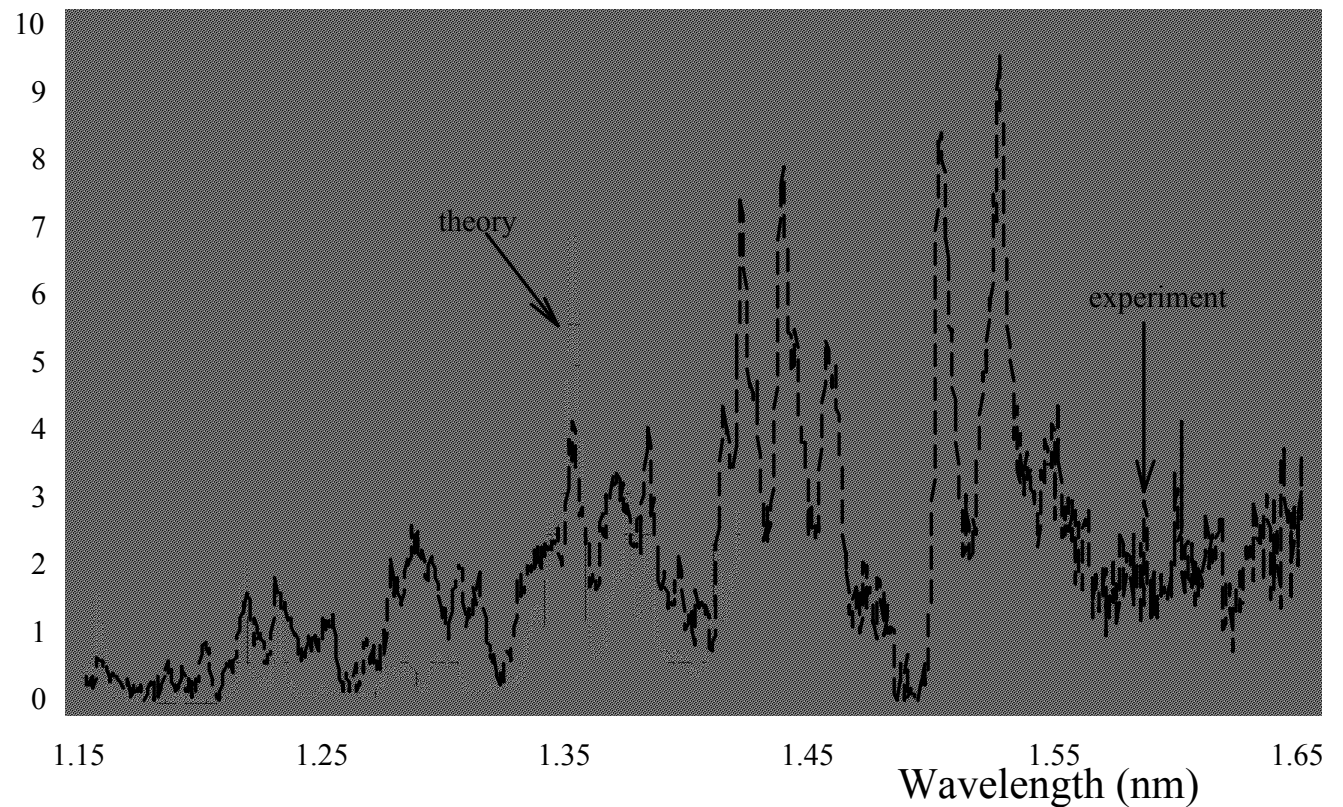
# Opacity from thin heated targets

Plasma conditions deduced by fitting simulation of emission for a buried Al layer in CH directly heated by a laser: electron density  $6.9 \times 10^{22}/\text{cm}^3$  ( $\sim 0.3 \text{ gcm}^{-3}$ ) electron temperature 400 eV.



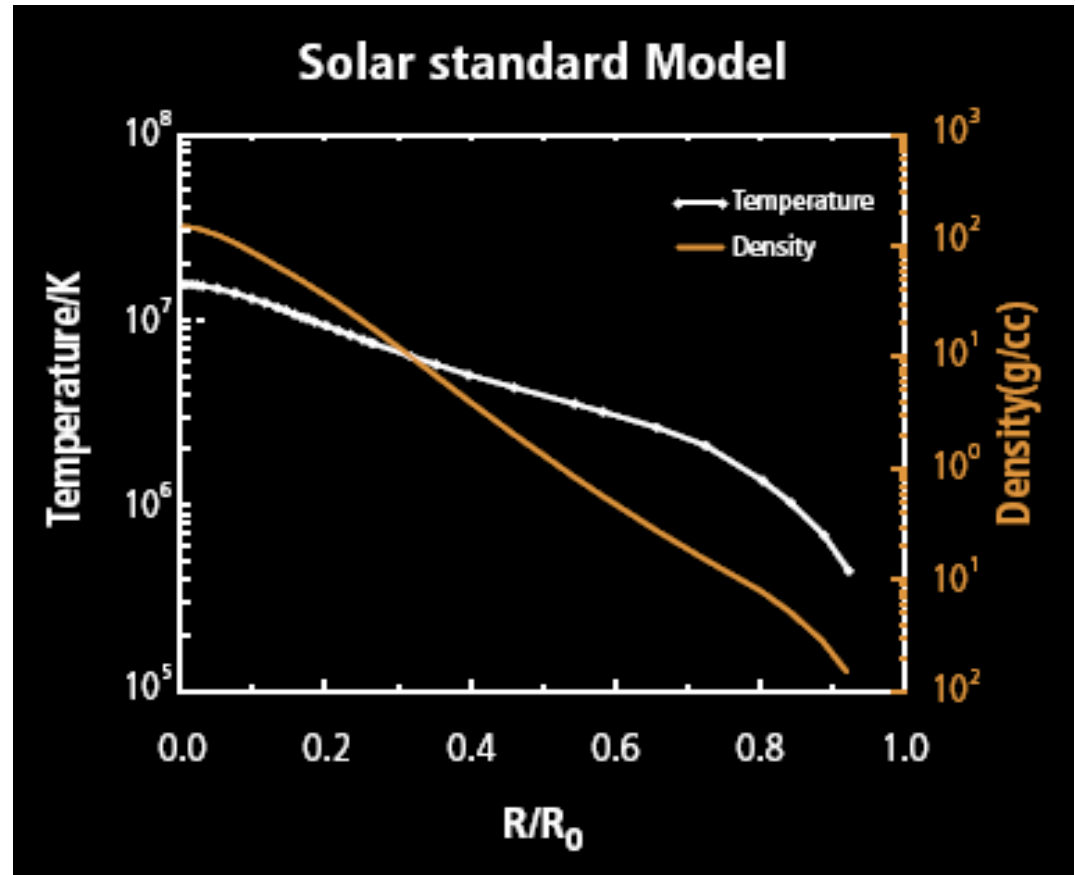
# Opacity deduction from thin heated targets

Plasma conditions deduced by fitting simulation to experimental Fe spectrum for Fe layer buried in CH: electron density  $10^{24} \text{ cm}^{-3}$  ( $\sim 3 \text{ gcm}^{-3}$ ), temperature 150 eV.

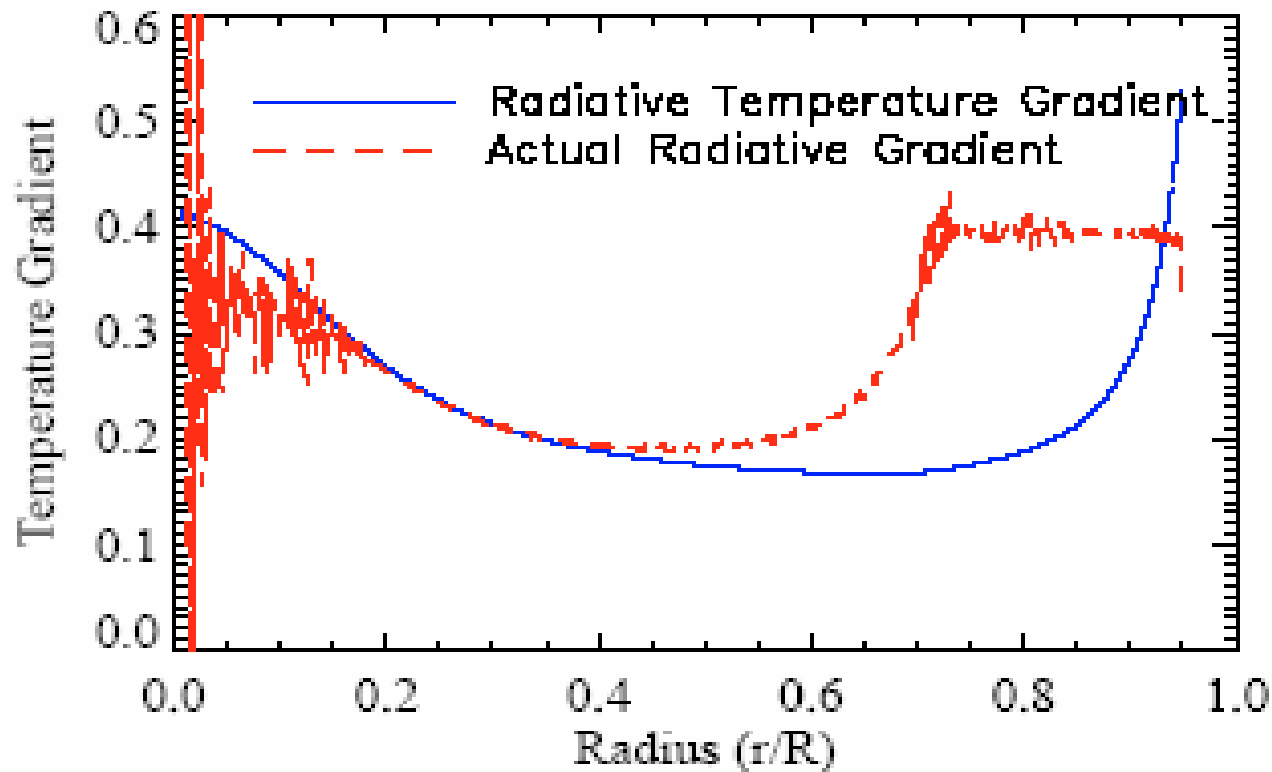


# Are laboratory opacity measurement relevant to the sun?

- Backlighter measurements with x-ray heating of targets have measured the opacity of plasmas around  $10^{-2} \text{ gcm}^{-3}$  and up to 40 eV: corresponding to the outer regions of the sun.
- Direct heating of thin layers and deduction from emissivity (assuming LTE) has measured opacities up to 500 eV and density close to solid ( $0.1 - 10 \text{ gcm}^{-3}$ ):  $\Rightarrow R/R_0 \sim 0.5$



# Opacity dominates solar energy outflow at the solar centre and is important in the outer solar region



- Takes  $\sim$  million years for light to travel from the Solar centre to the photosphere.
- Mean free path for light absorption at the Solar centre is 1 mm.

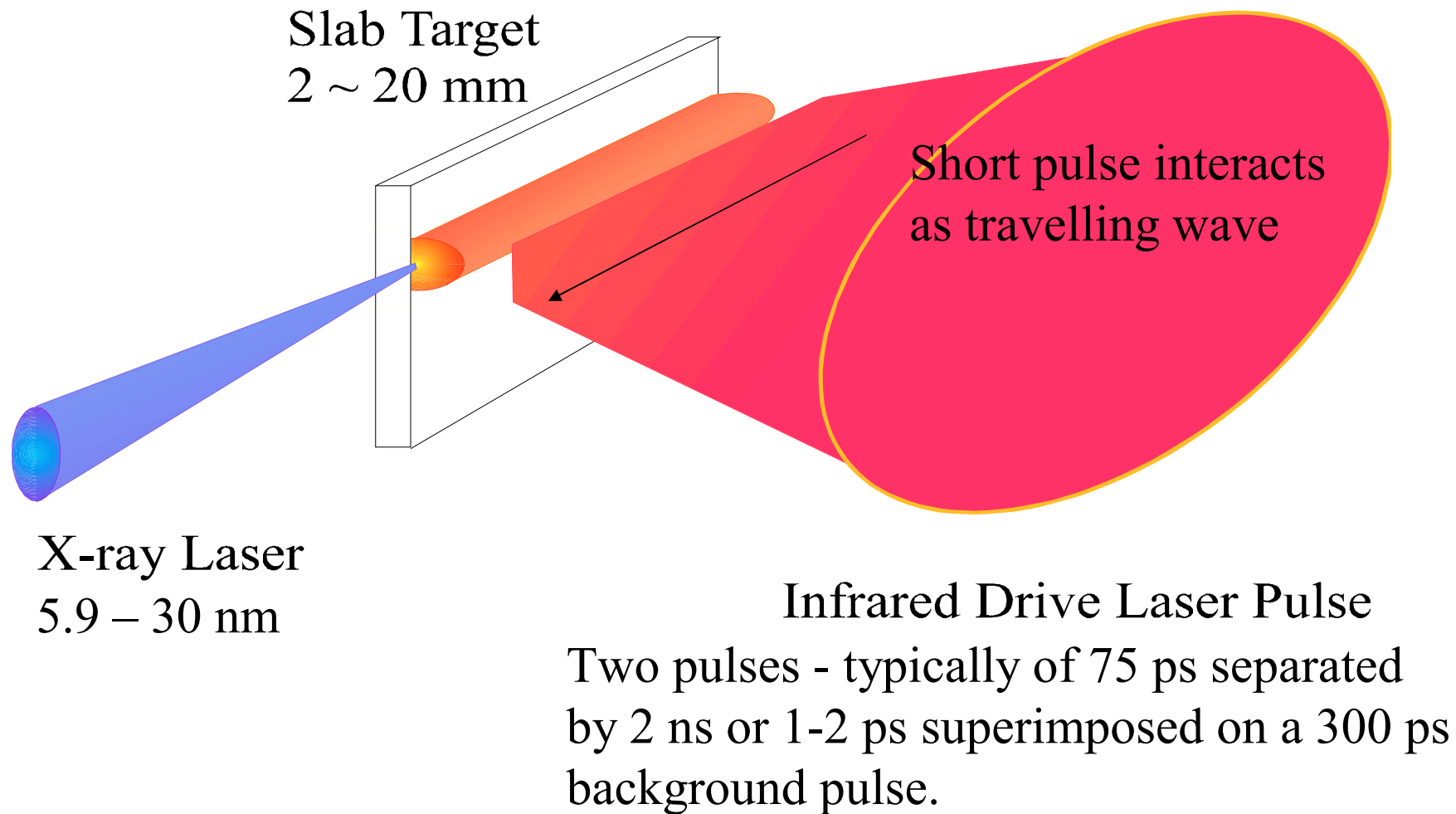
# X-ray lasers are much brighter than conventional backlighters

- X-ray lasers produce pulses of 60 – 200 eV photon energy over durations of 3-100 ps.
- Saturated irradiances are  $\sim 3 \times 10^{10} \text{ Wcm}^{-2}$ .
- Photons are produced in a beam with divergence 5 – 20 mrad.
- Spectral bandwidth is very narrow ( $\nu/\Delta\nu > 10^4$ ).

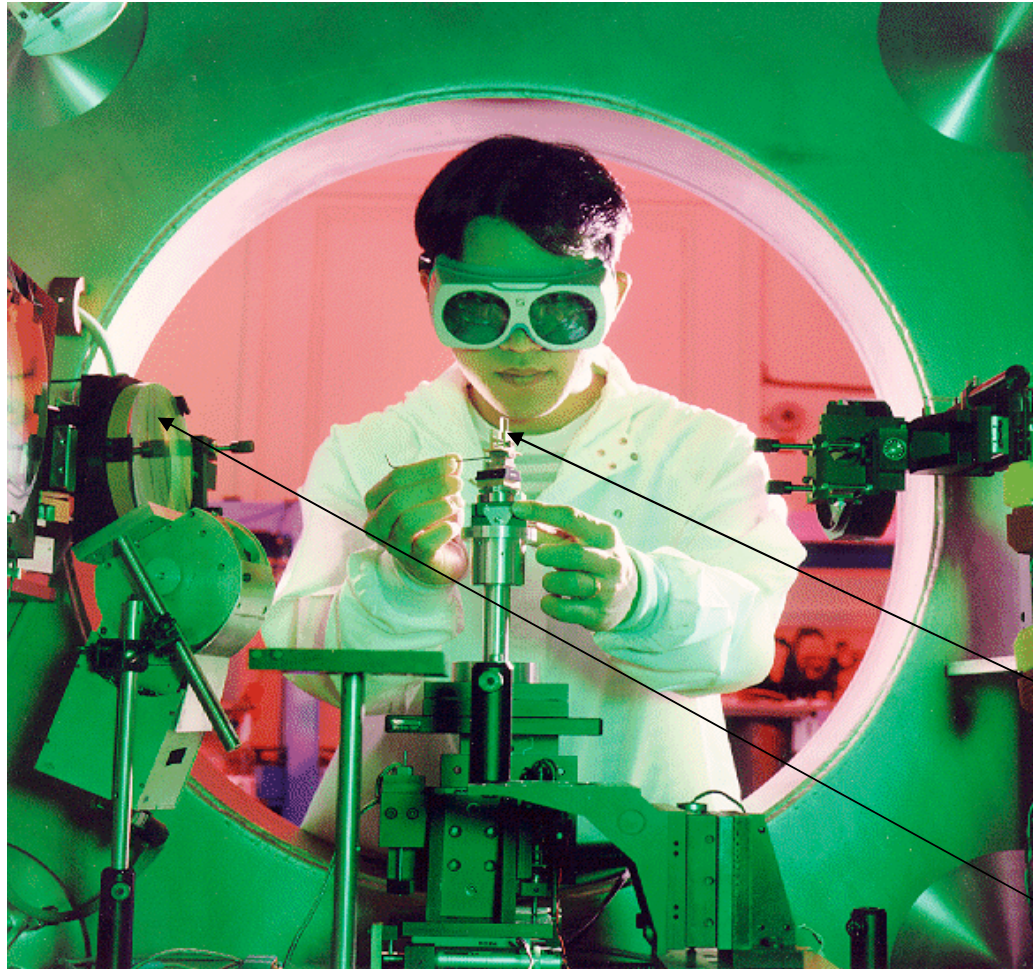
$\Rightarrow$  Beam-like nature and narrow spectral bandwidth should enable x-ray lasers to ‘out-shine’ emission from solar condition plasmas (solar surface flux is  $\sim 6 \times 10^{17} \text{ Wcm}^{-2}$ ).



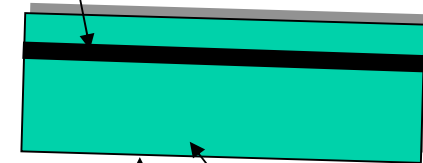
# X-ray lasers



# Adjusting an x-ray laser target



Stripe coated  
on slide



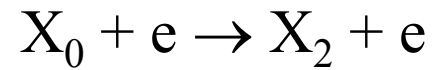
Glass slide

Target

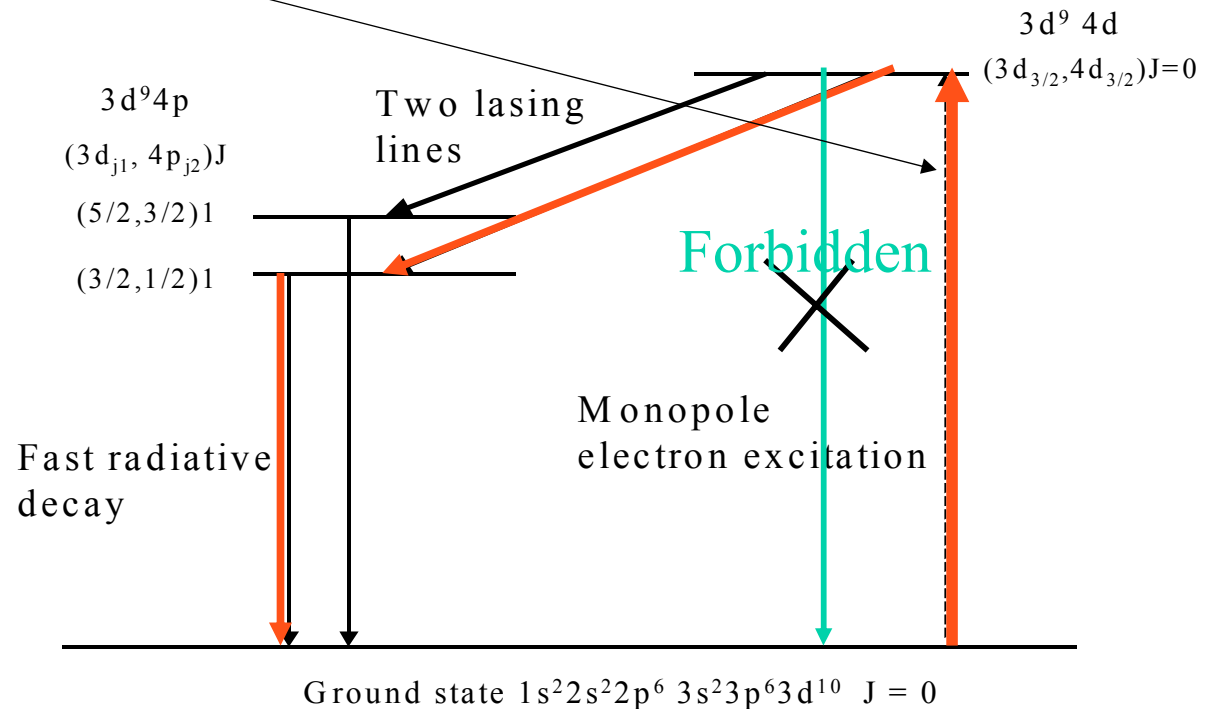
Mirror to focus pump laser

# Electron collisional pumping in laser-produced plasma creates x-ray lasing

Monopole excitation

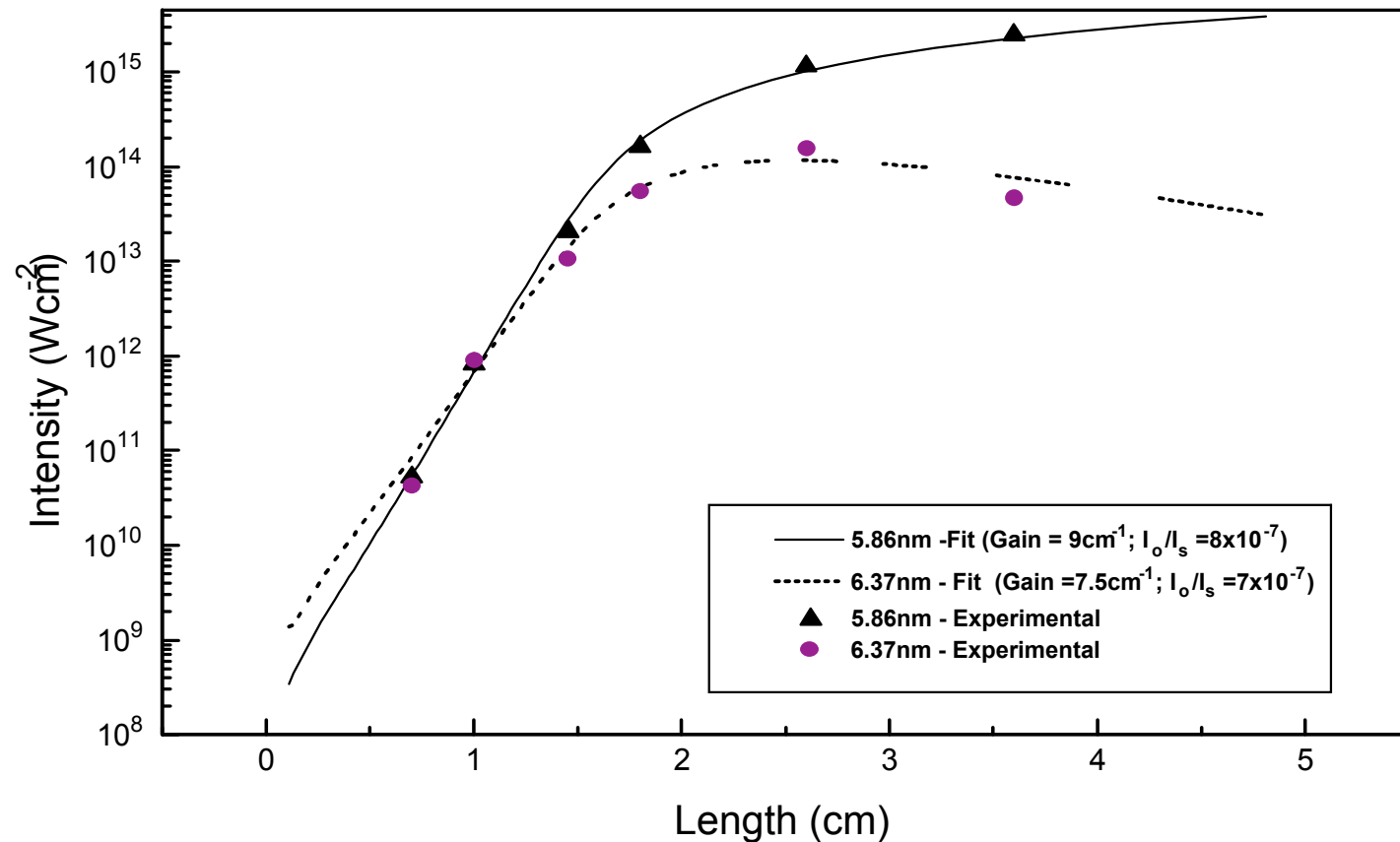


Energy levels important for lasing in Ni-like ions



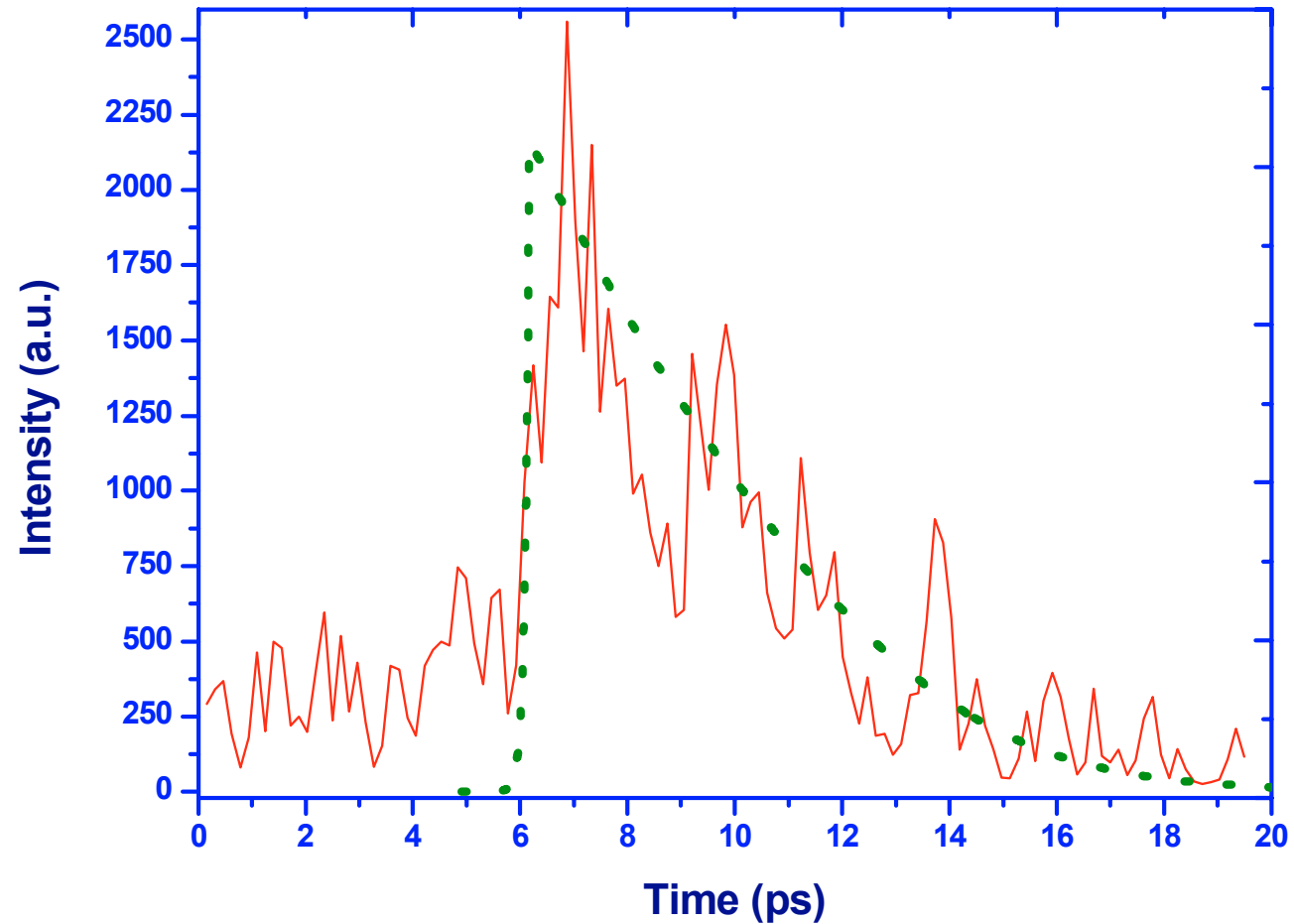
# Shortest wavelength saturated laser is at 5.9 nm (210 eV)

X-ray laser  
output of  
Ni-like Dy

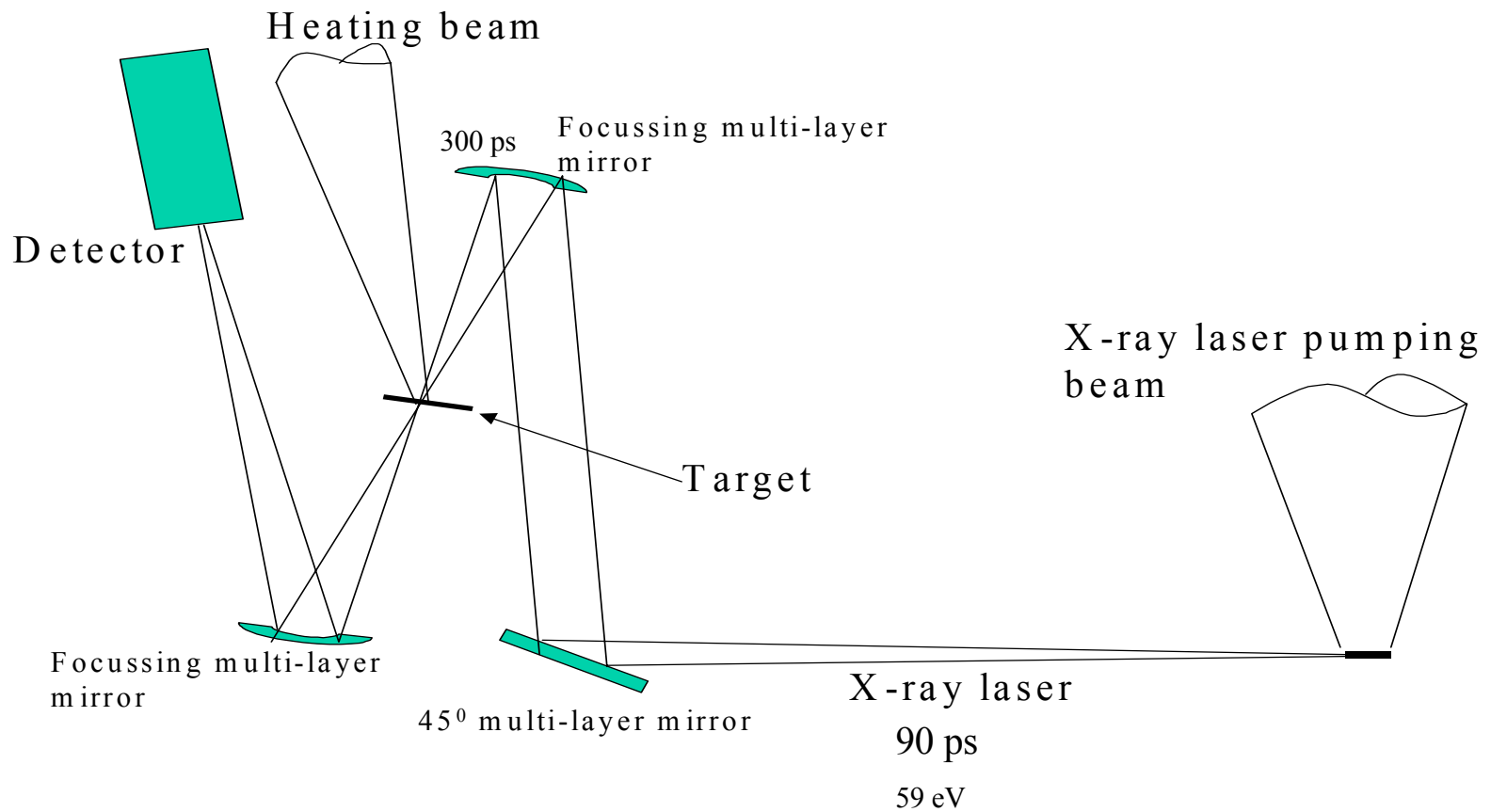


# X-ray laser pulse duration as short as 3 ps have been measured

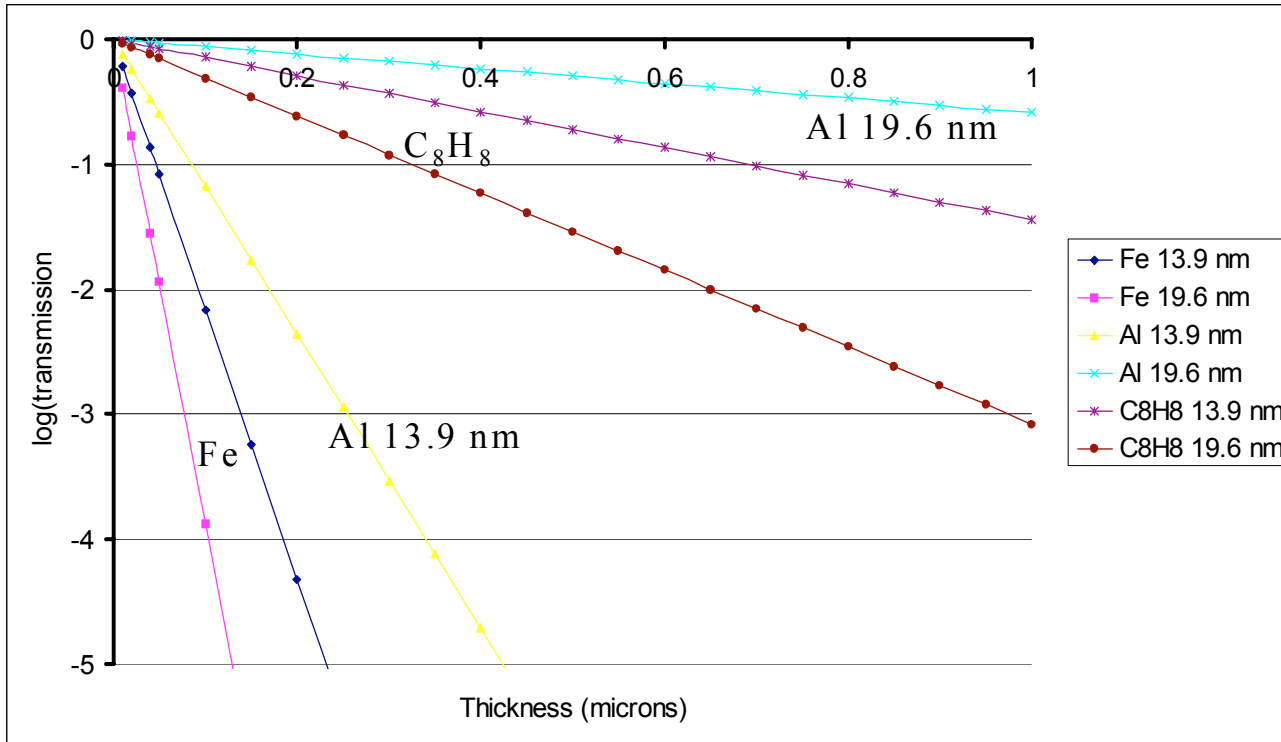
Ni-like Ag  
output at  
13.9 nm.



# An experiment to use x-ray laser output as a backlighter to measure plasma opacity



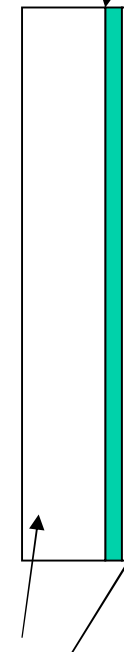
# Targets



Transmission of solid,  
room temperature material

Fe  
0.05 micron  
with some Mg  
as a signature

>>

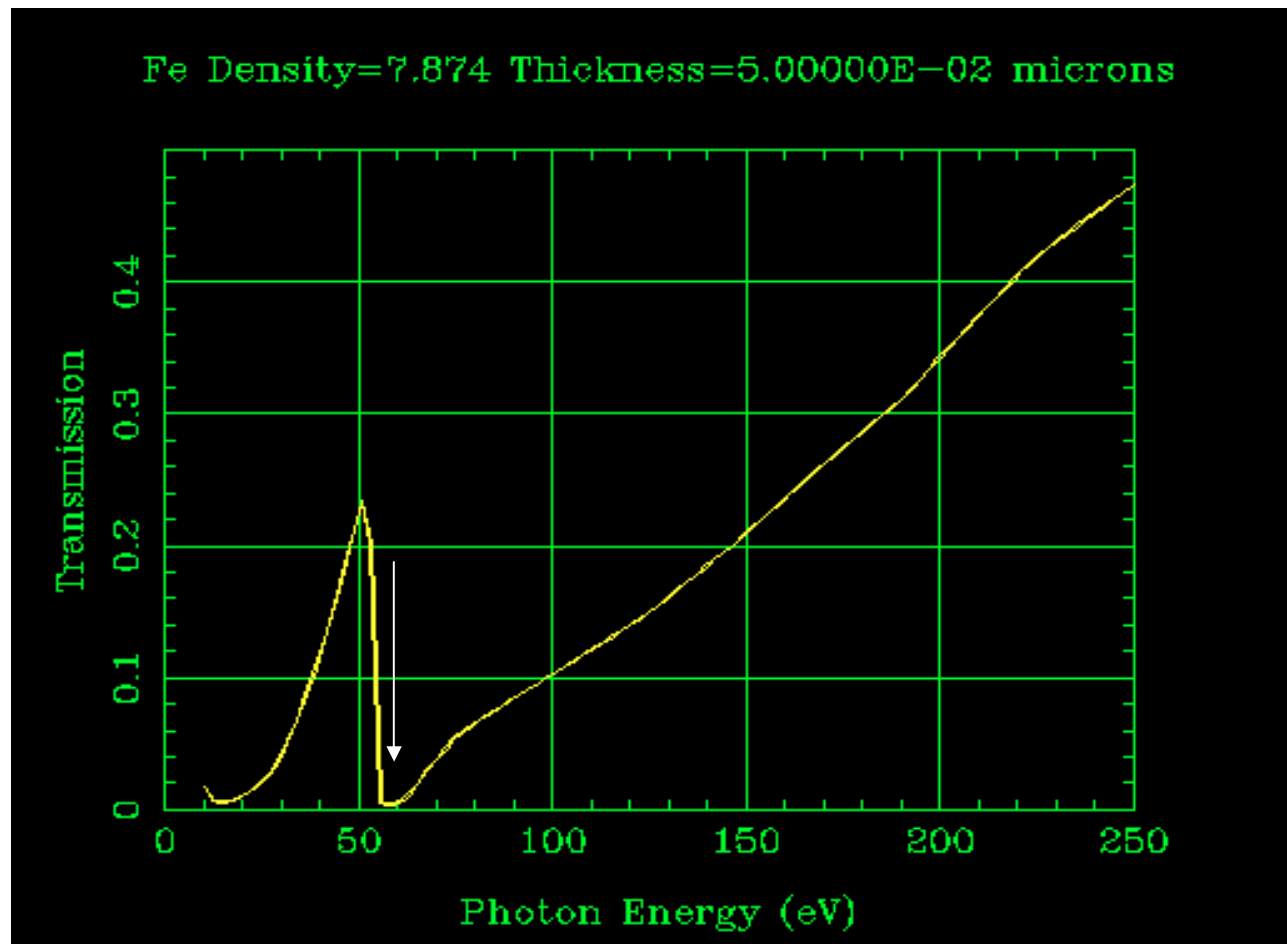


Al or CH

# The opacity of solid iron at EUV photon energies

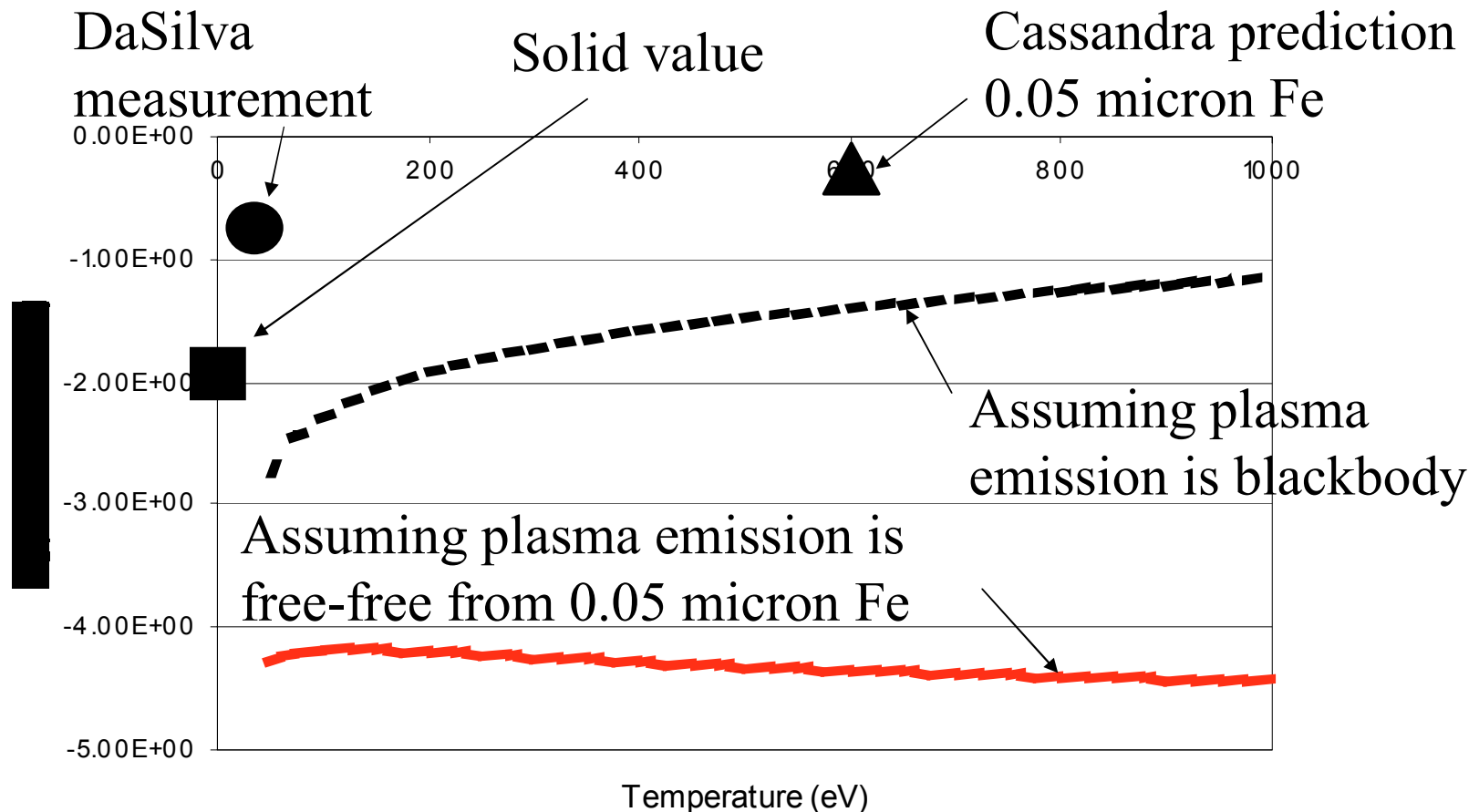
Experiment at 21.2nm has a high opacity in cold Fe. Opacity rapidly decreases with ionisation.

Arrow shows photon energy for 21.2 nm.





# Signal to noise for Fe target transmission measurement

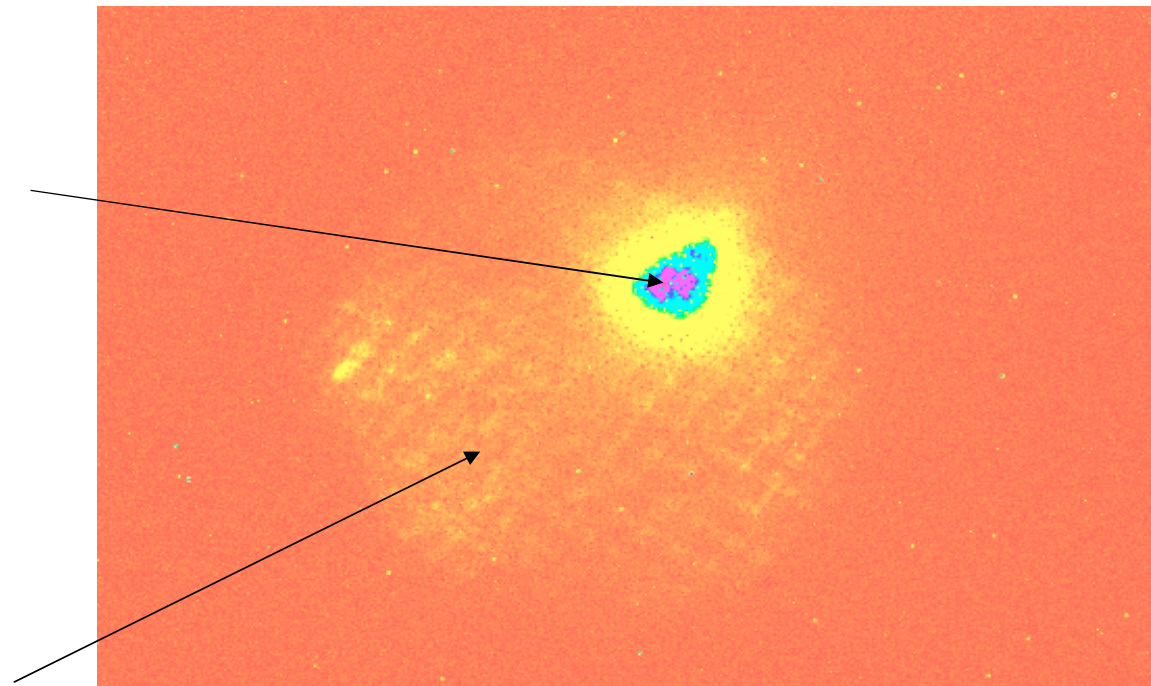


Transmission minimum producing 10X x-ray laser flux compared to plasma background (as labelled).

# Transmission of 21.2 nm x-ray laser through 0.8 $\mu\text{m}$ Al heated by laser

Transmission of X-ray laser through heated and ionised material

X-ray laser shining through unheated Al

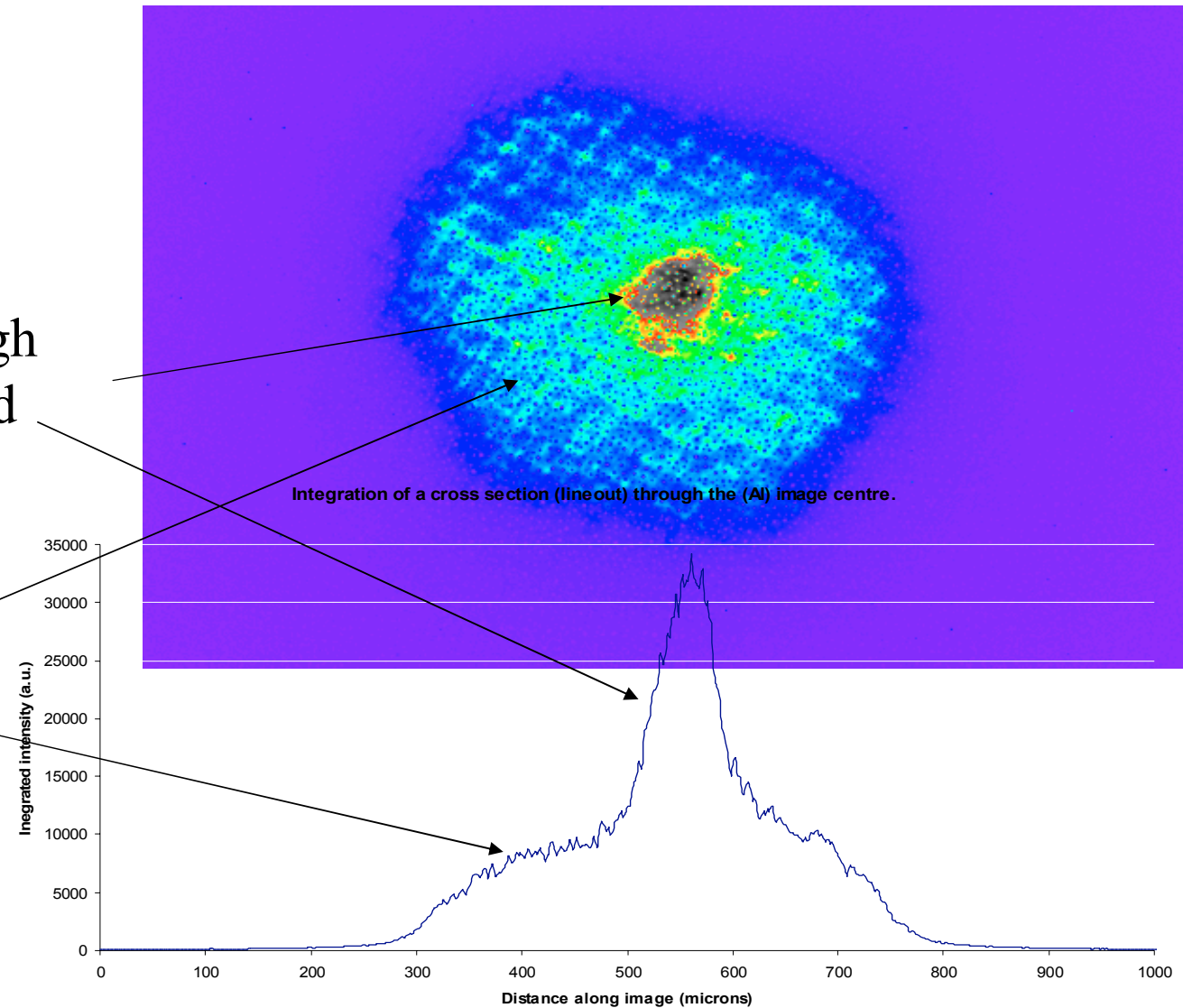


# Another probed Al plasma

Transmission of  
X-ray laser through  
heated and ionised  
material

X-ray laser shining  
through unheated Al

⇒44% transmission  
through ionised Al  
while unheated Al  
transmits 12%

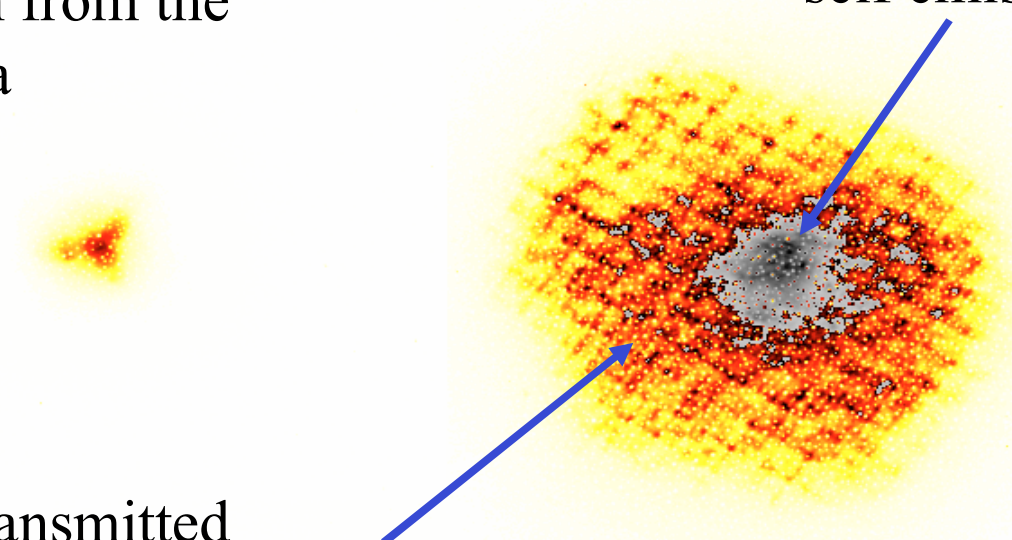


# Probed Al with a measure of self-emission

Self-emission from the hot Al plasma

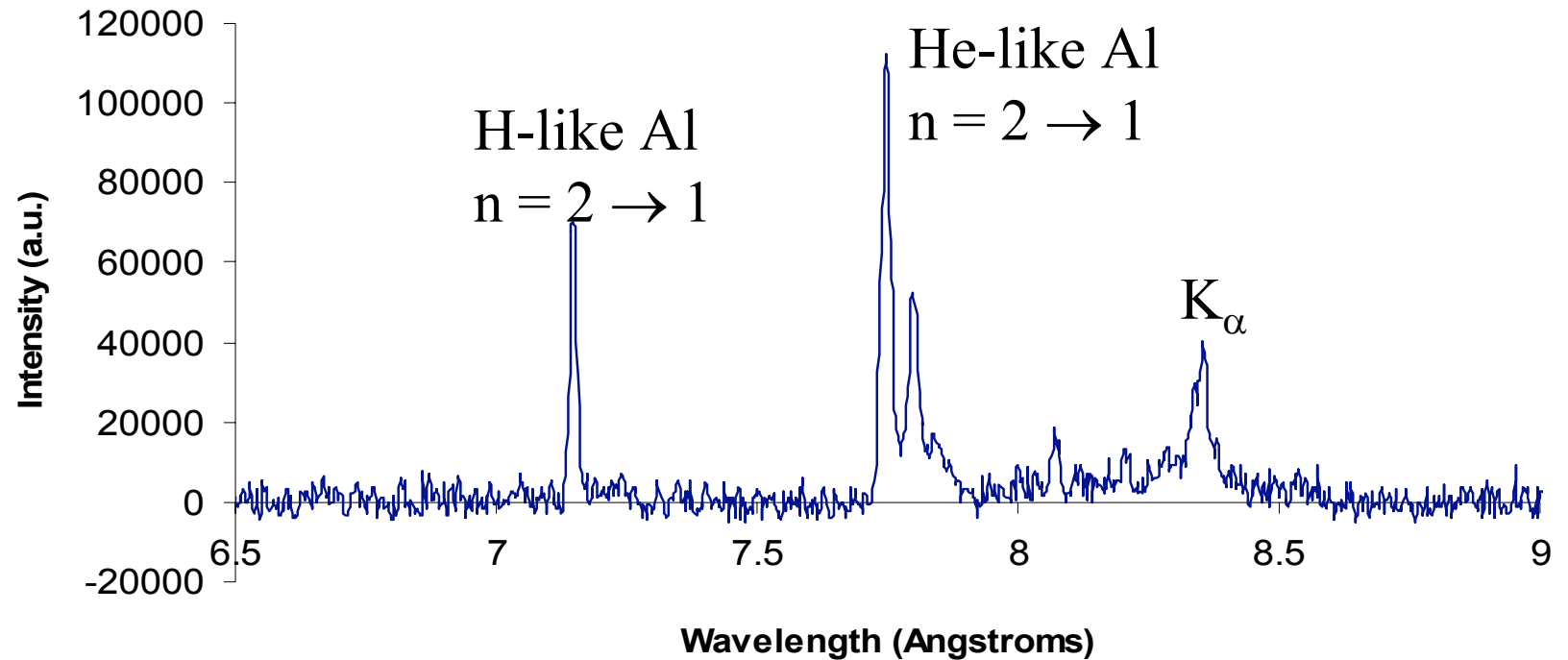
Transmitted X-ray laser beam through heated plasma + self emission

Transmitted X-ray laser beam through unheated Al



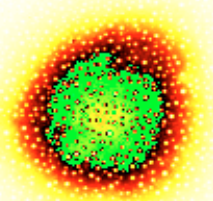
# Aluminium is well ionised

Aluminium spectrum produced by a 10J, 130 micron spot on a 0.8 micron Al foil

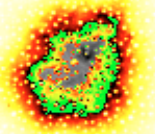


# Opacity of thin targets of Fe

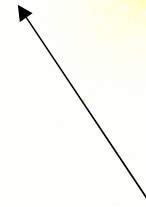
Self emission  
of Fe



Transmitted x-ray laser  
through heated Fe  
+self emission

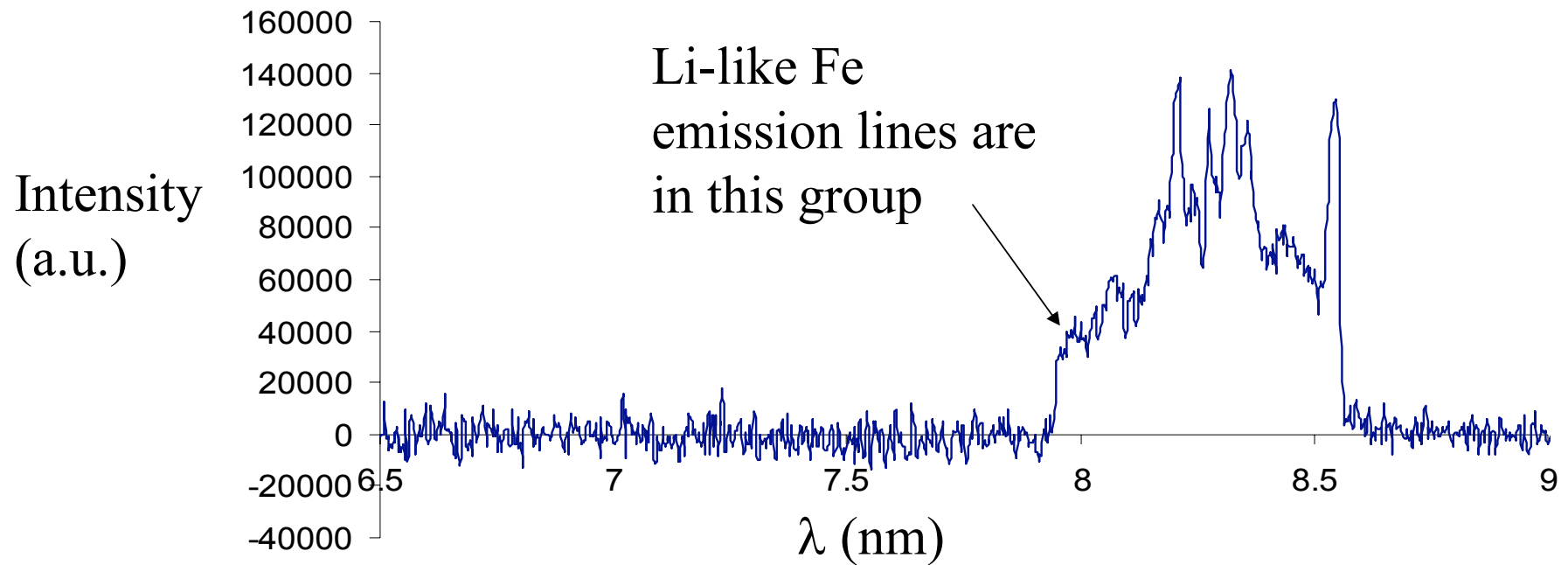


X-ray laser flux through  
cold Fe is small

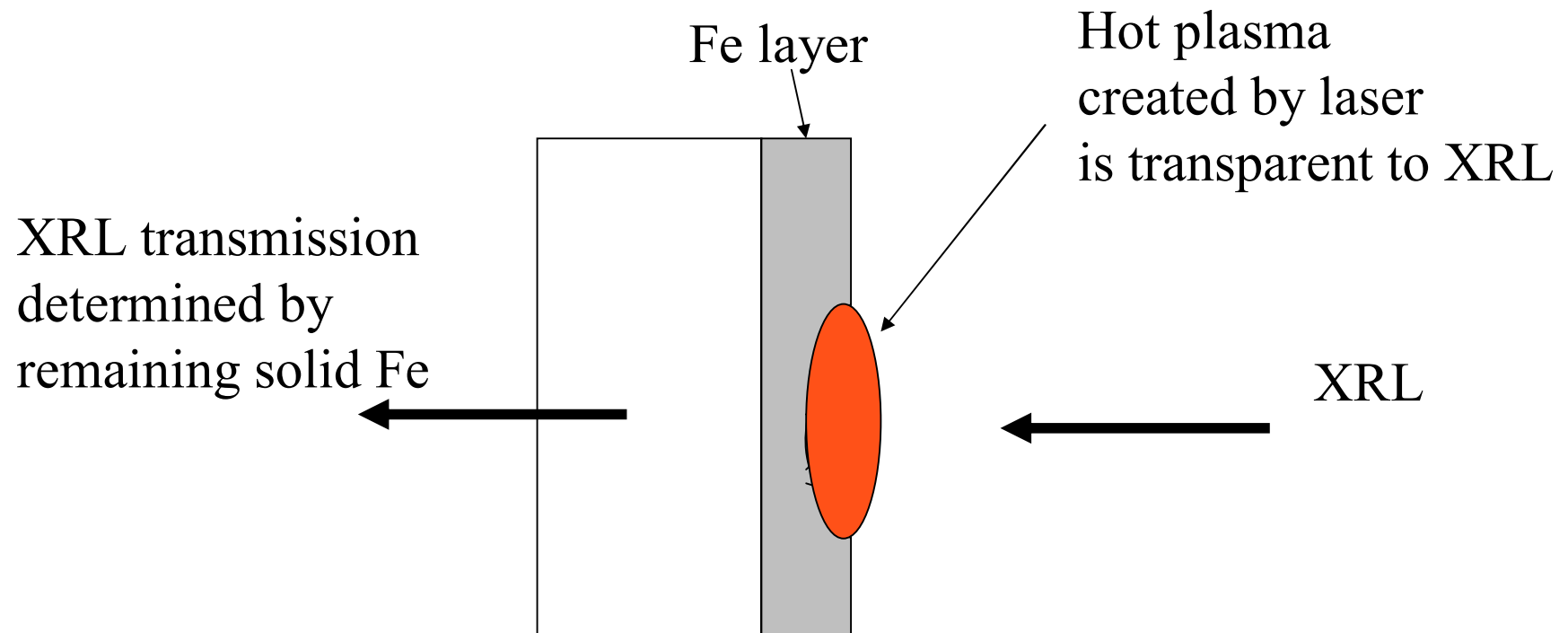


# Thin iron targets are also well ionised

Spectrum produced by a 10J, 130 micron spot on a 0.75 micron Al foil coated with 0.05 microns of Fe



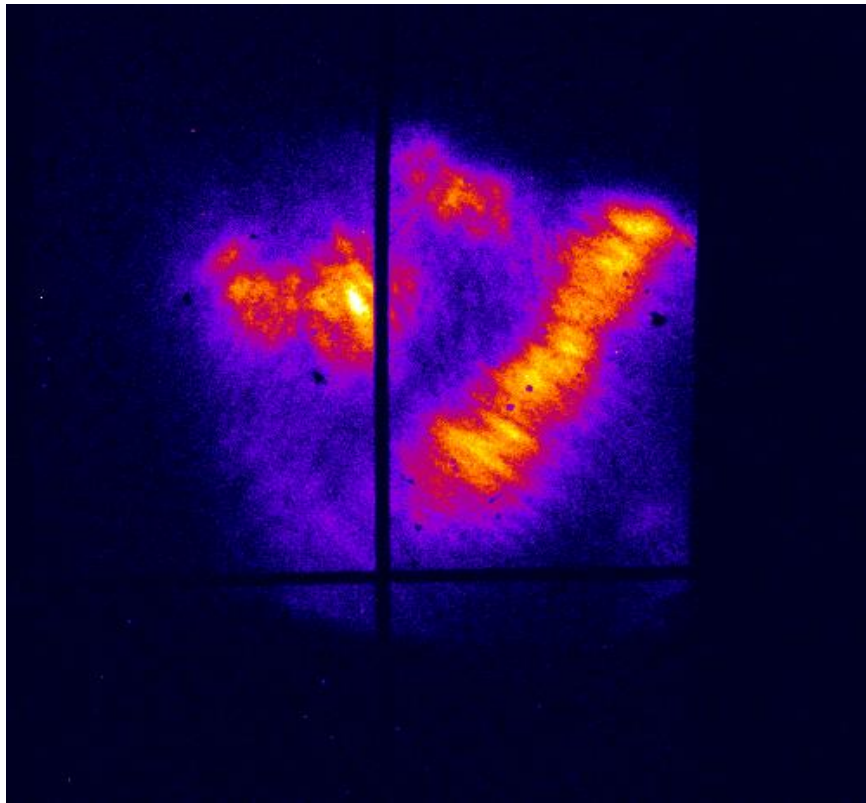
# Schematic of physics of transmission with thin non-tamped foils.



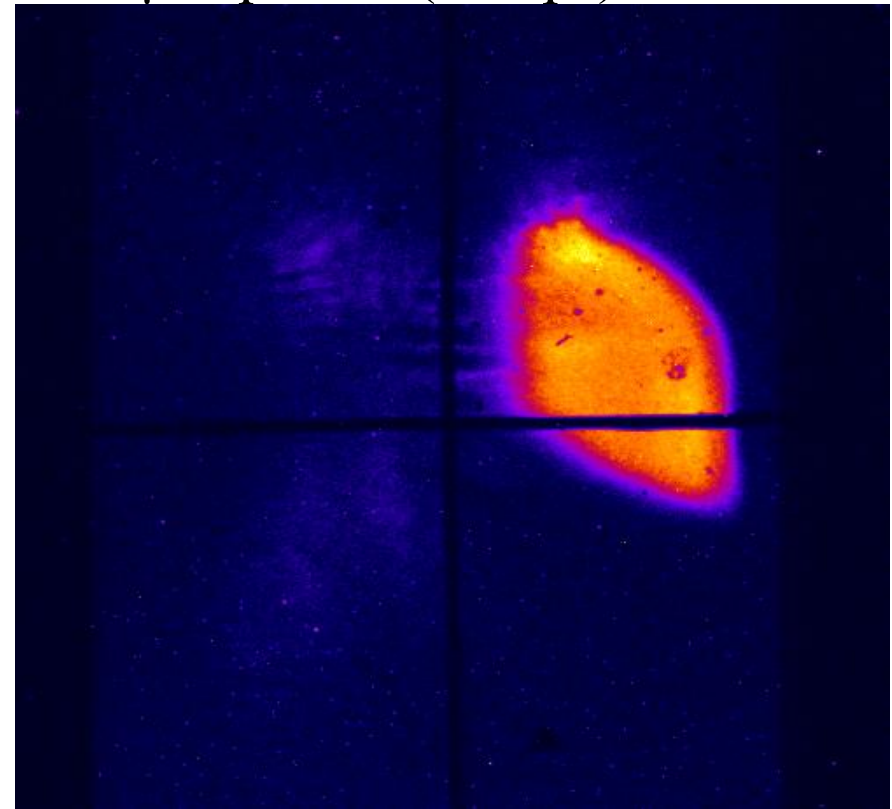


# Iron transmission over the laser heated area

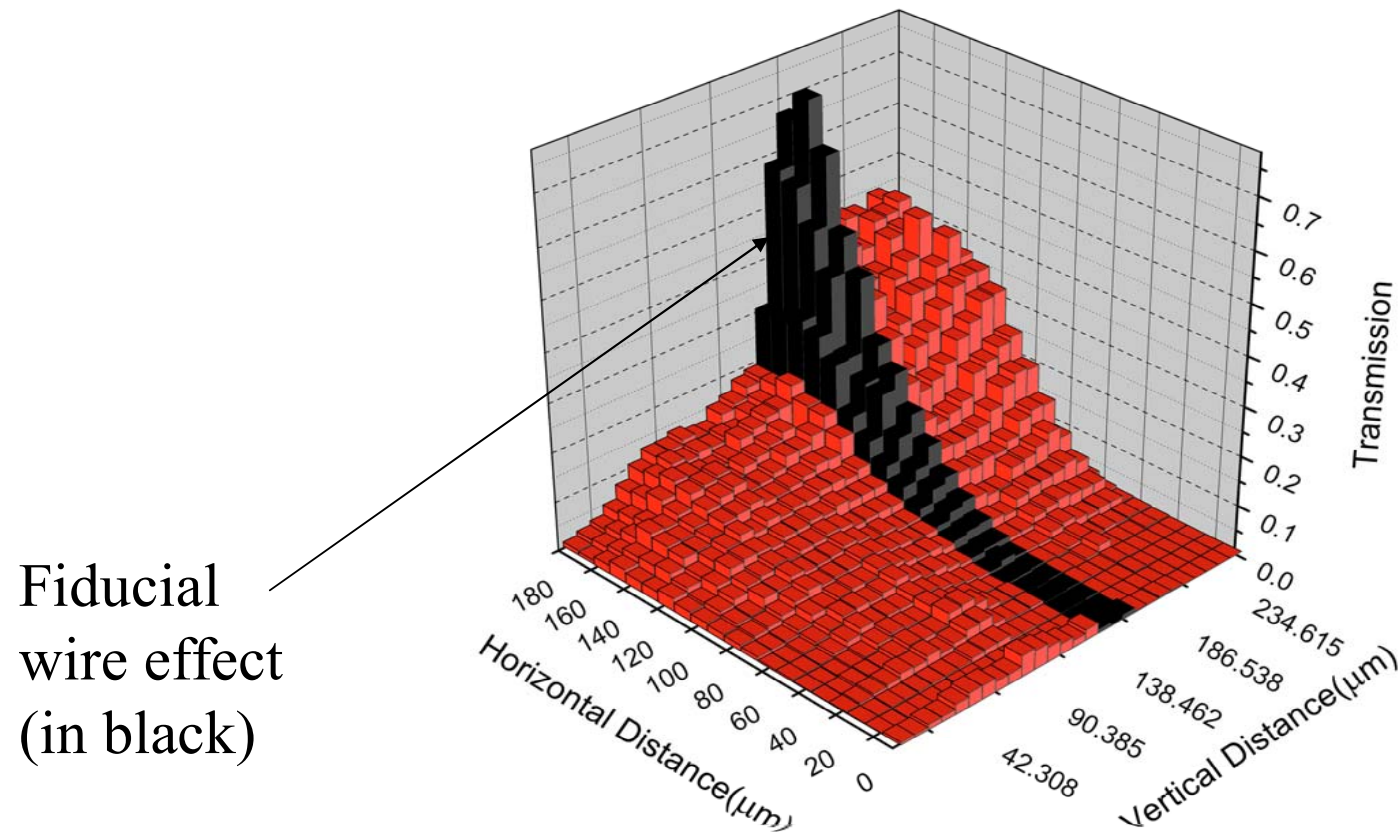
XRL beam profile



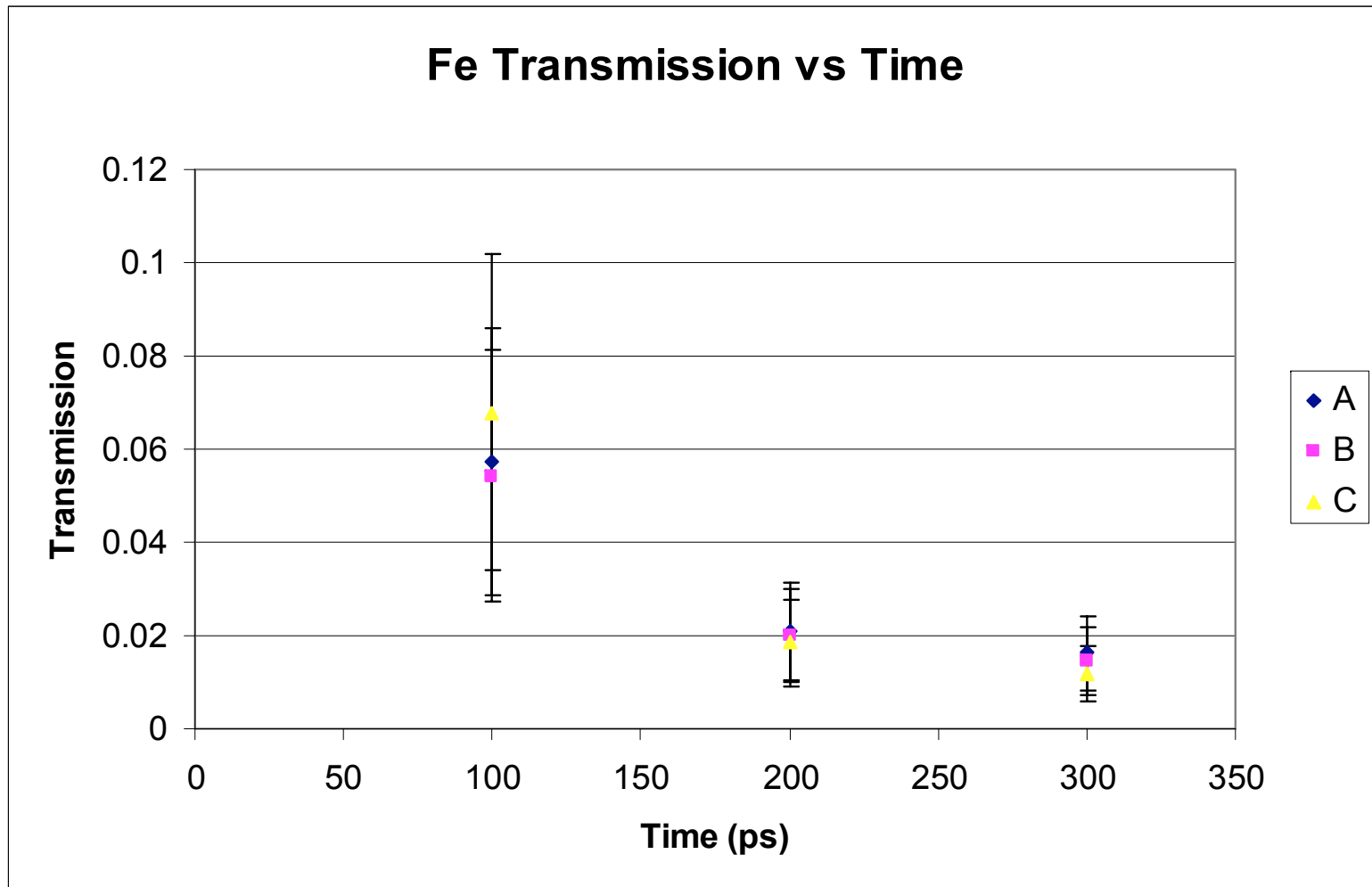
Transmission of XRL thru heated 0.05  $\mu\text{m}$  Fe under 0.05  $\mu\text{m}$  plastic (100 ps)



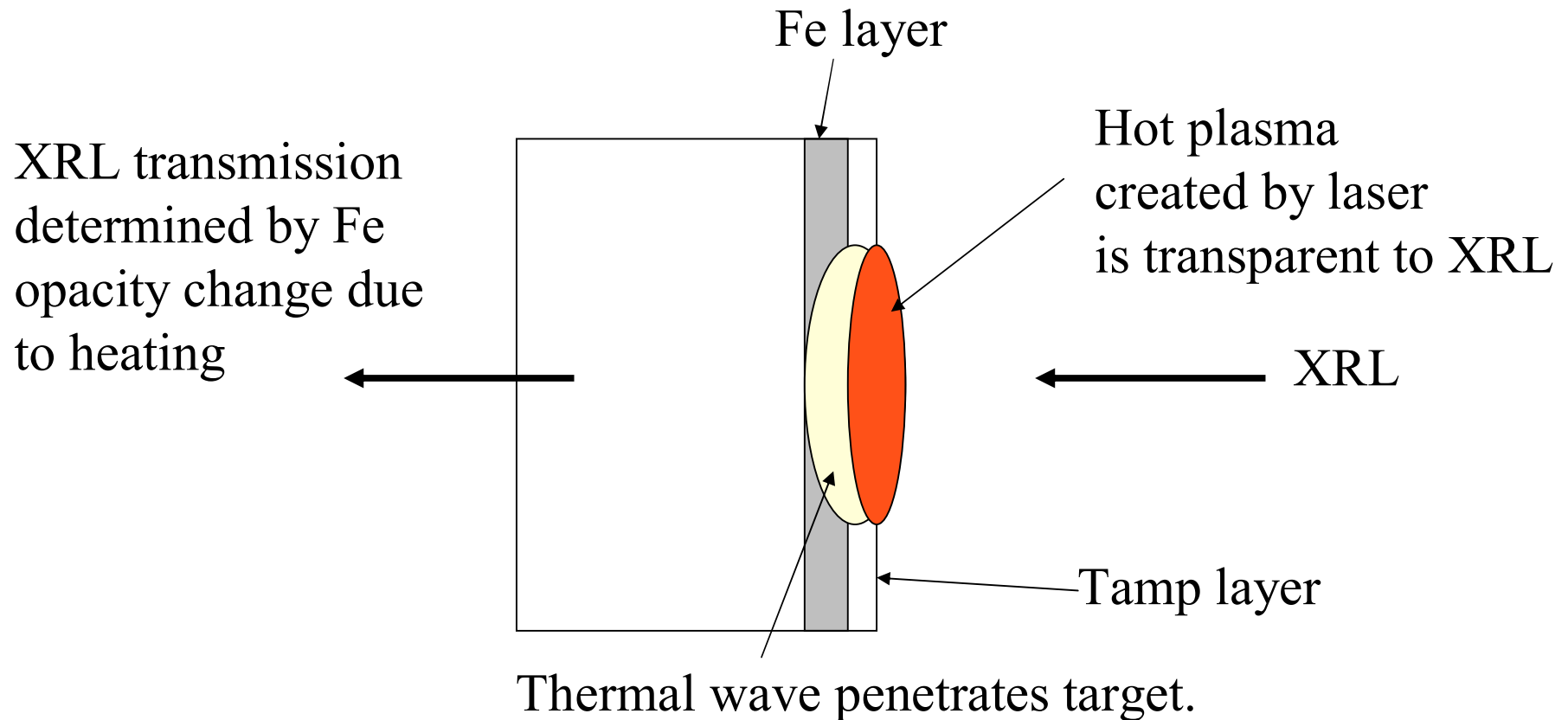
# Transmission can be deduced taking account of XRL variation



# Iron transmission when buried under plastic



# Schematic of physics with tamped Fe targets



# Conclusions

- Plasma opacity is important in astrophysical plasmas and can be measured in the laboratory using lasers to heat the plasma samples and to produce ‘backlighters’: opacities for temperatures up to 40 eV for densities  $\sim 10^{-2} \text{ gcm}^{-3}$  have been previously measured.
- An ‘X-ray laser’ backlighter at 60 eV photon energy has been demonstrated to be able to measure opacity in plasmas heated to 100’s eV at near solid density (1 – 8  $\text{gcm}^{-3}$ ).