



The Abdus Salam  
International Centre for Theoretical Physics



SMR.1738 - 9

WINTER COLLEGE  
on  
QUANTUM AND CLASSICAL ASPECTS  
of  
INFORMATION OPTICS

*30 January - 10 February 2006*

**Spin of Orbital Angular Momentum**

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*Trieste Feb 2006*

# Spin cf. Orbital Angular Momentum

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*Miles Padgett*



**UNIVERSITY**  
*of*  
**GLASGOW**



# Transfer of AM to micro-objects

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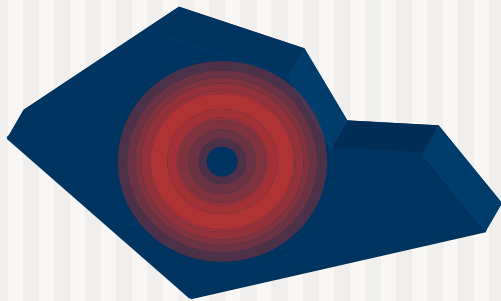
He *et al.* Phys Rev. Lett. 1995

Simpson *et al.* Opt. Lett. 1997

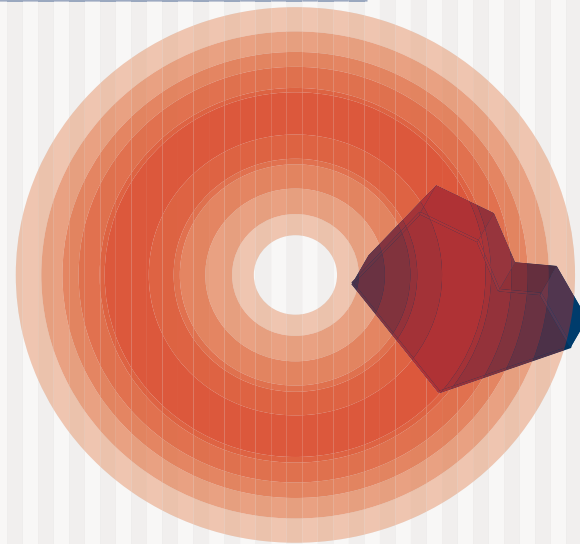
O'Neil *et al.* Phys. Rev. Lett. 2002

# Angular momentum interactions with particles

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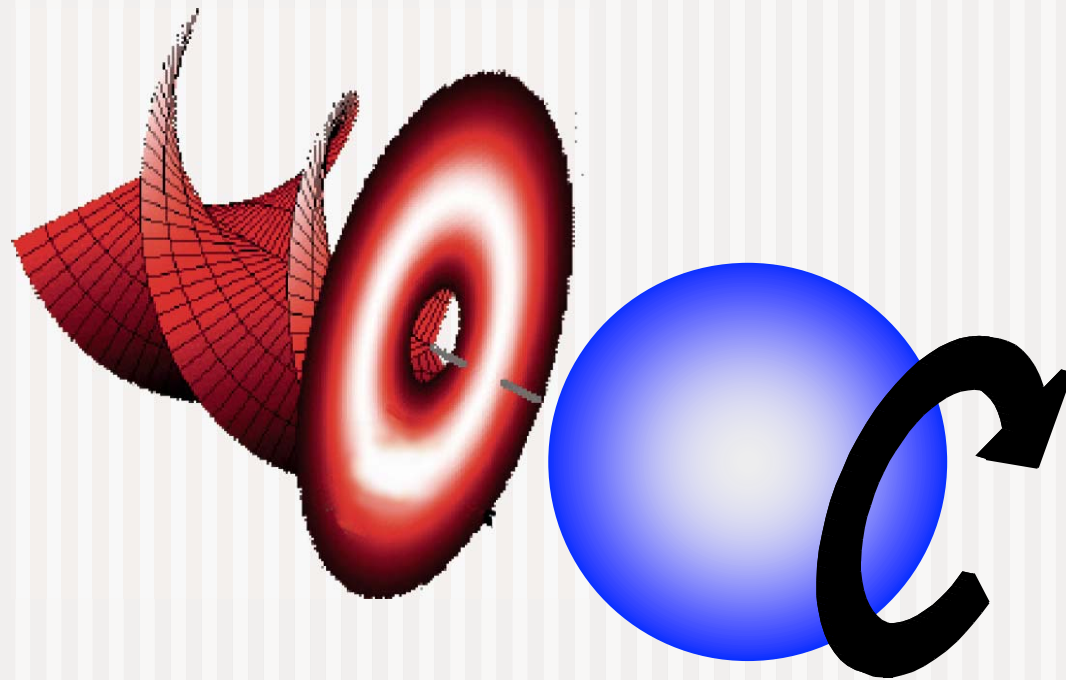
- Object larger than beam
  - Spin AM = Orbital AM (for absorption)



- Beam larger than object
  - Spin AM  $\neq$  Orbital AM

# On-axis Spin and Orbital transfer

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**SAM &or OAM**

Particle spins on beam axis

# OAM / SAM transfer to particle held in optical tweezers

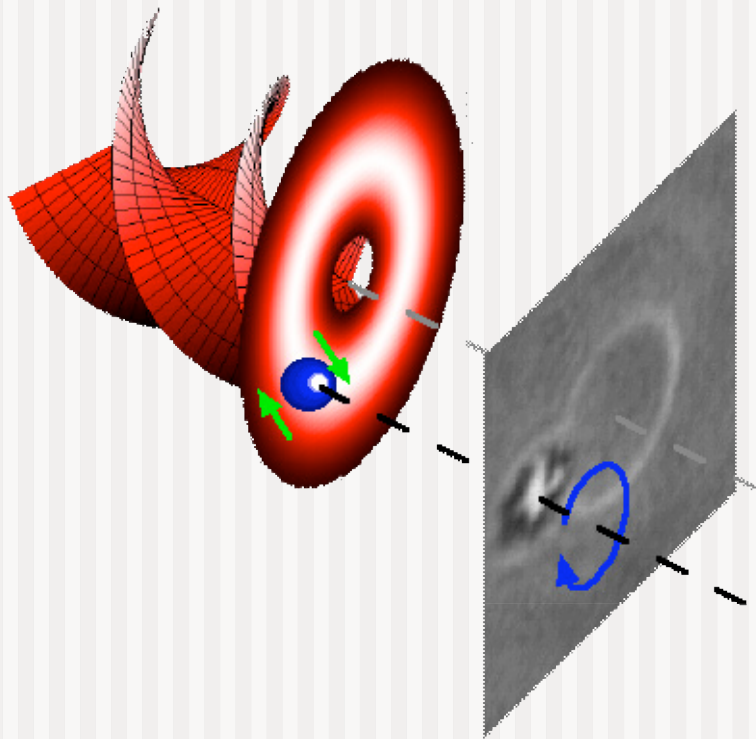
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$$\text{OAM } (\hbar) \text{ +/- SAM } (\hbar)$$

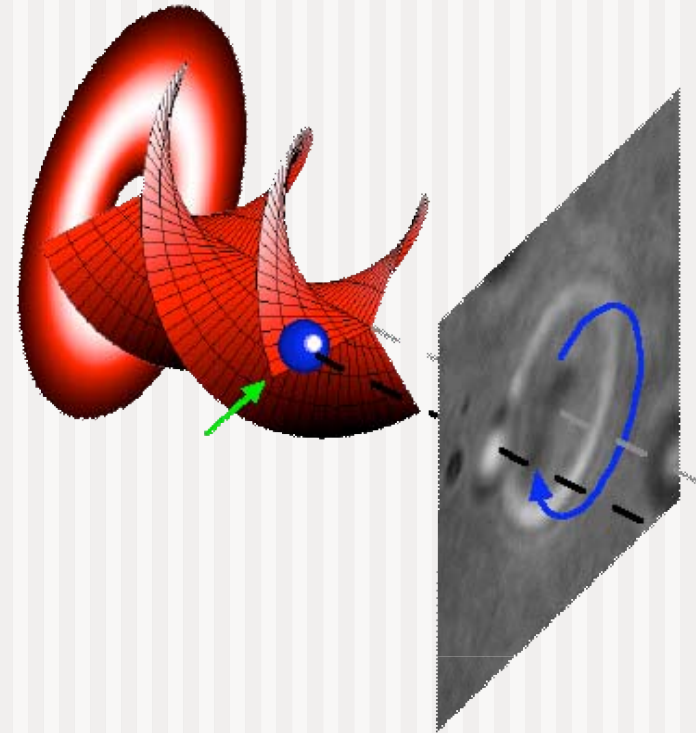
Particle spins and stops

# Off-axis Spin and Orbital transfer



**SAM**

Particle spins on its own axis



**OAM**

Particle orbits the beam axis

# OAM / SAM transfer to particle held in optical tweezers

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**SAM**

Particle spins on its own axis



**OAM**

Particle orbits the beam axis



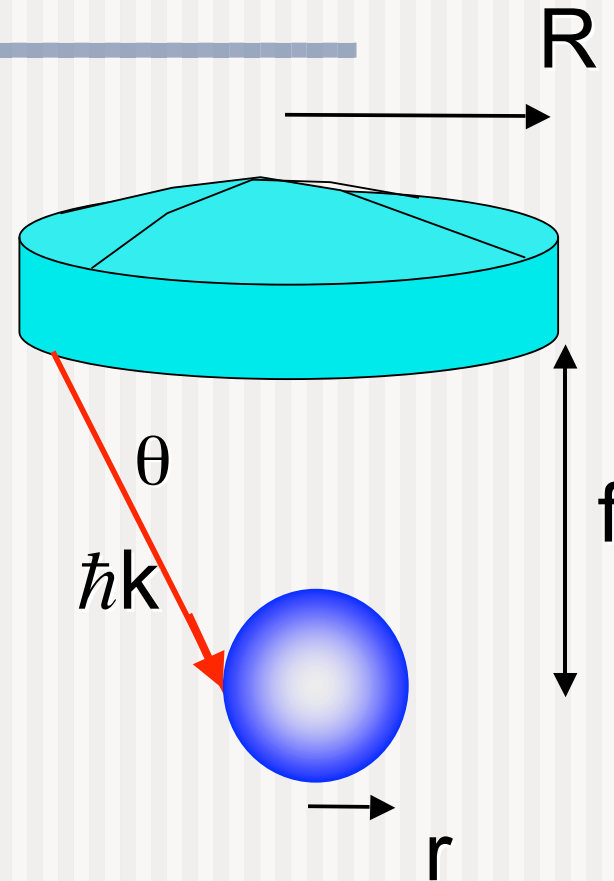
# Ray-optics to model OAM

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Courtial and Padgett Opt. Commun.  
2000

# Transfer of angular momentum

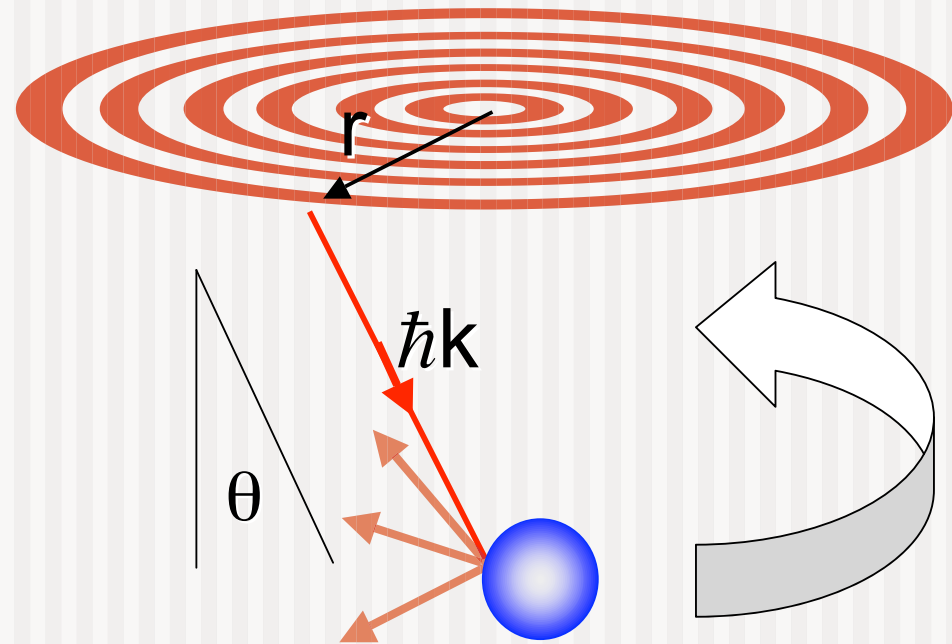
- Angular momentum arises from skew rays
  - $\theta = \ell / kr$
- The skew angle sets the azimuthal component to the momentum density
  - $p_\phi = \hbar k \theta = \ell \hbar / r$
  - $L = p_\phi r = \ell \hbar$
  - $L_{\max} = \hbar k r (R/f)$



Ray-Optics gives the right answer

# Transfer of orbital AM (e.g. from Bessel beam)

- Local intensity (Bessel Beam)
  - $I \propto 1/r$
- Angular momentum arises from skew rays
  - $\theta = \ell / kr$
  - i.e.  $\theta \propto 1/r$
- Circumference of ring
  - $\propto r$
- (orbital) rotation rate
  - $\Omega \propto 1/r^3$



Ray-Optics gives the right answer

# Rotational frequency shifts

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Garetz and Arnold Opt. Commun. 1979

Courtial *et al.* Phys. Rev. Lett. 1998

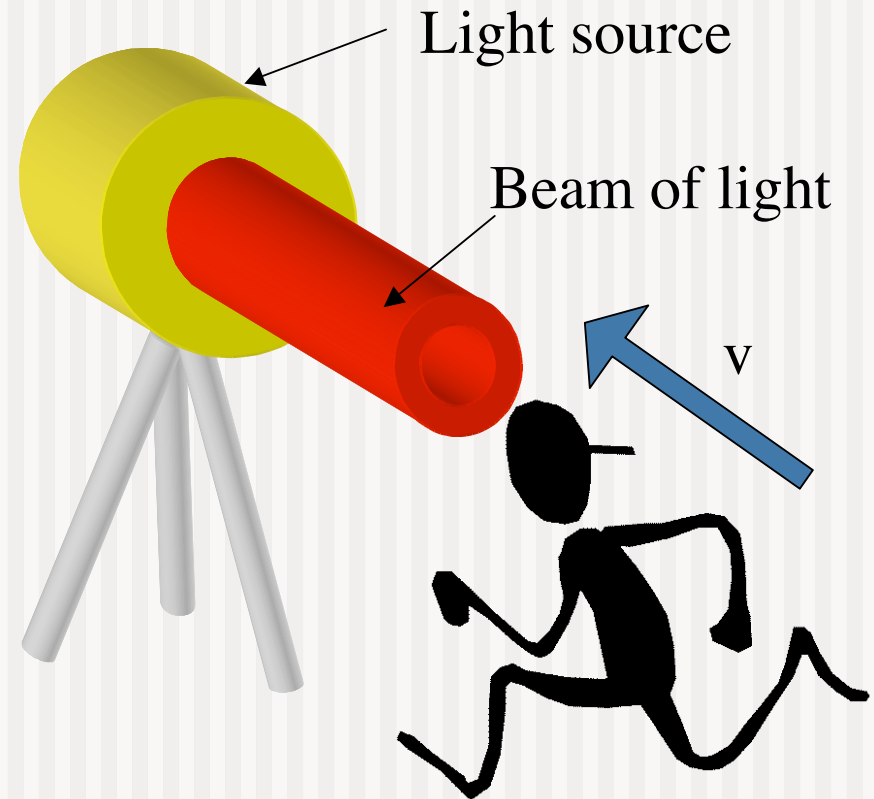
# The linear Doppler shift

- Light source moves towards or away from detector giving Doppler shift

- $\Delta\omega = \omega_0 \times v/c$

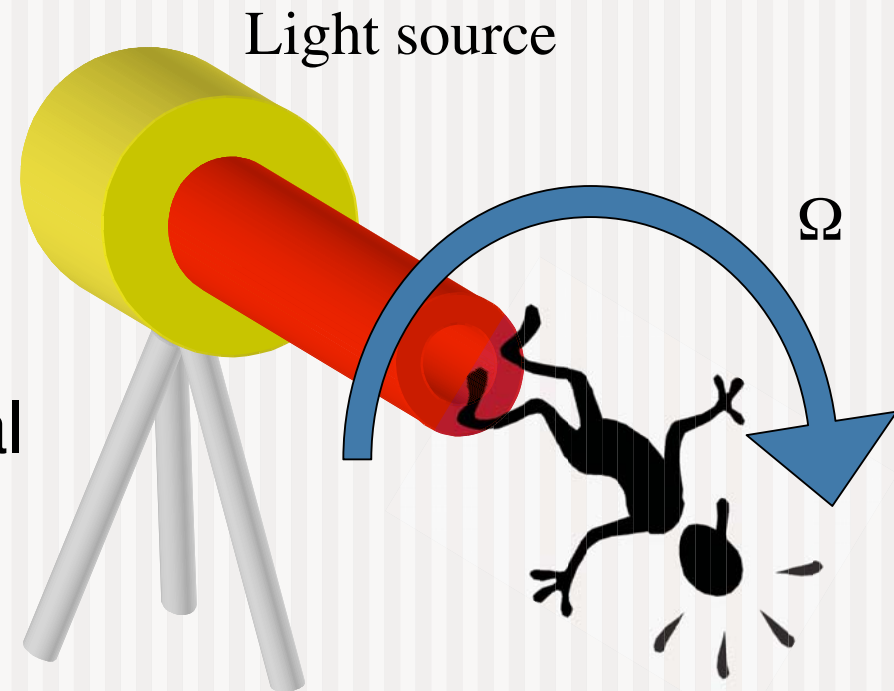
- Re-express in terms of linear momentum per photon,  $p$

- $\Delta\omega = v \times p/\hbar$



# The annular Doppler shift

- Light source rotates with respect to detector giving Doppler shift
  - $\Delta\omega = \Omega \times (\ell + \sigma)$
- Also called rotational frequency shift



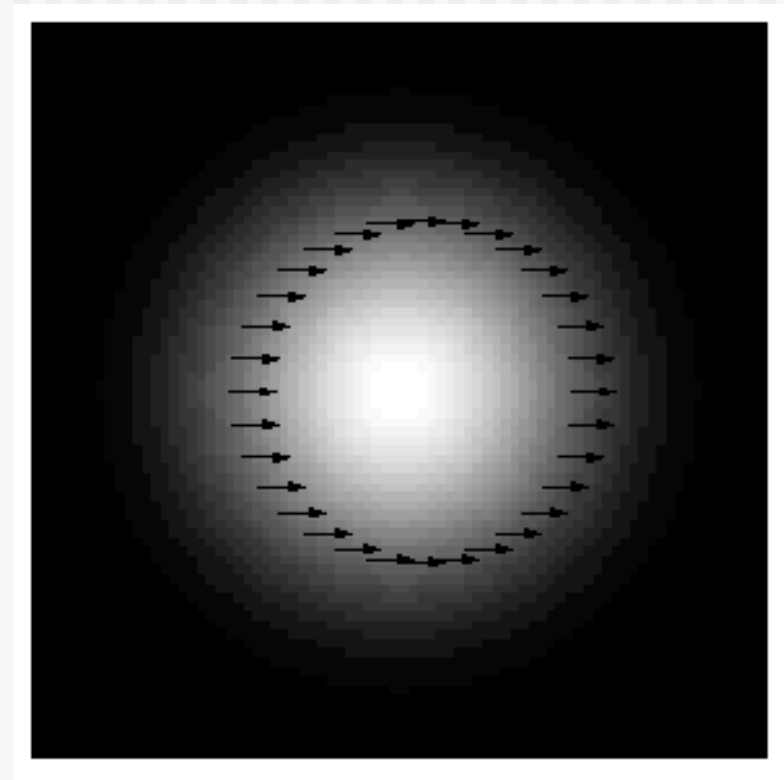
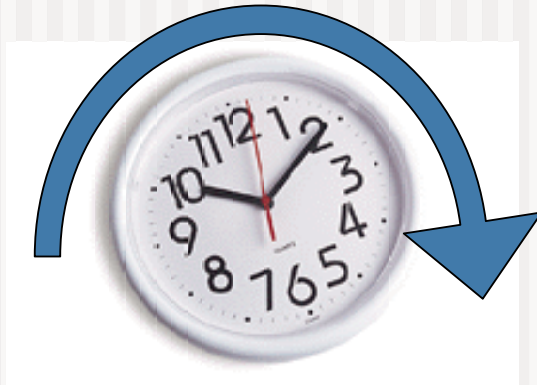
# Annular Doppler shift for circularly polarised light

- Additional rotation of polarisation (at  $\omega$ ) shifts frequency

$$\Delta\omega = \Omega$$

$$= \sigma\Omega \quad (\sigma=\pm 1)$$

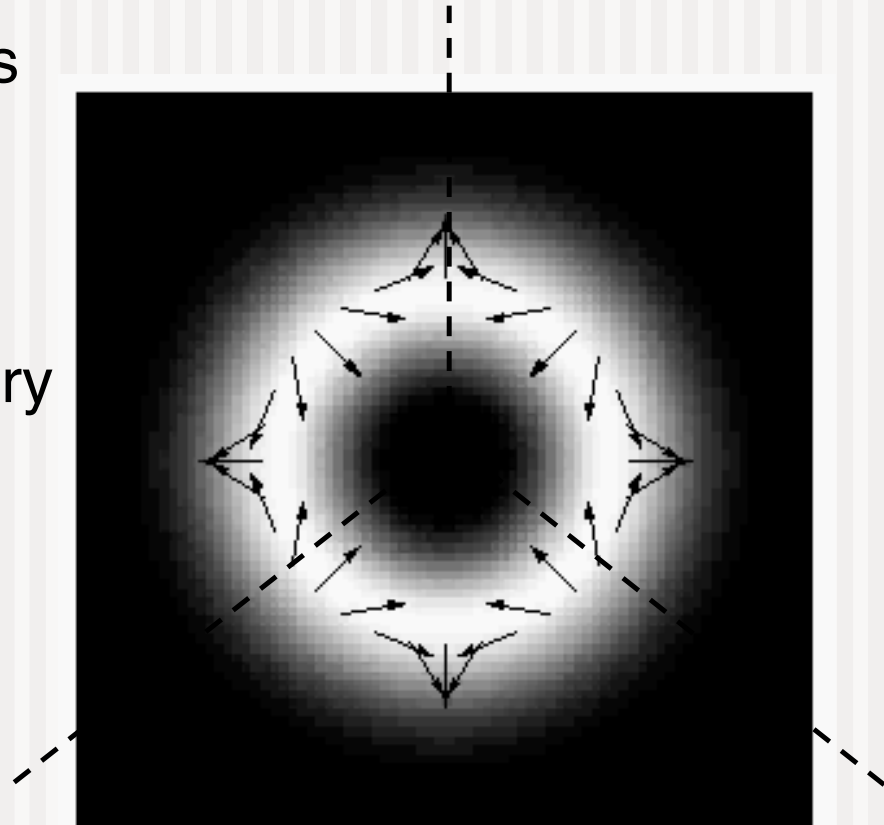
- c.f. time speeds up if you rotate a clock!



# Annular Doppler for helically phased circ. polarised light -1

- Such a beam contains both SAM and OAM
- Example 1  
 $l = 3, \sigma = +1$
- Four fold rot. Symmetry
- Rotate beam at  $\Omega$

$$\begin{aligned}\Delta\omega &= (l+\sigma)\Omega \\ &= J\Omega \\ &= 4\Omega\end{aligned}$$





# Rot. Doppler for helically phased, circ. polarised light -2

- The SAM and OAM add or subtract

- Example 2

$$l = -3, \sigma = +1$$

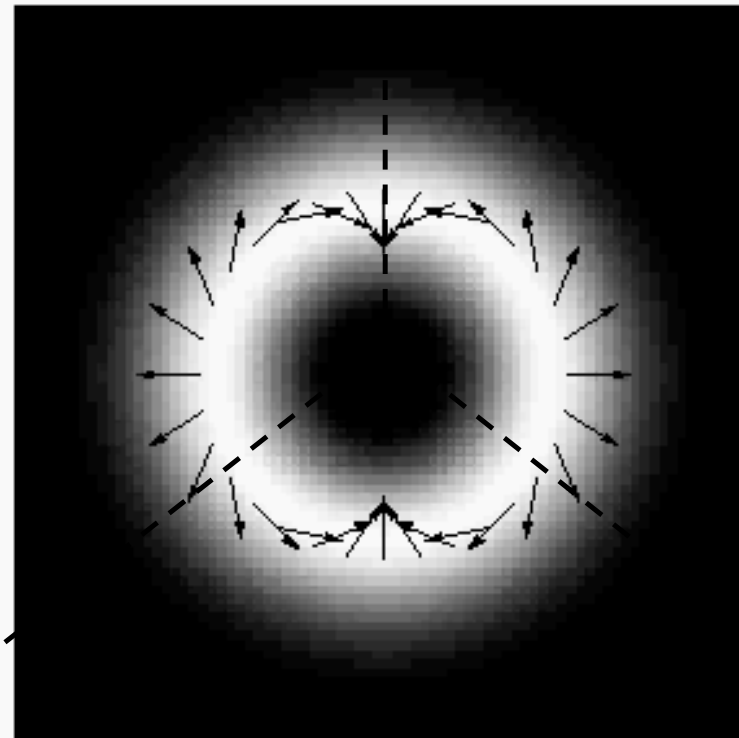
- Two fold rot. Symmetry

- Rotate beam at  $\Omega$

$$\Delta\omega = (l + \sigma)\Omega$$

$$= J\Omega$$

$$= 2\Omega$$



# Ray-optics to model OAM

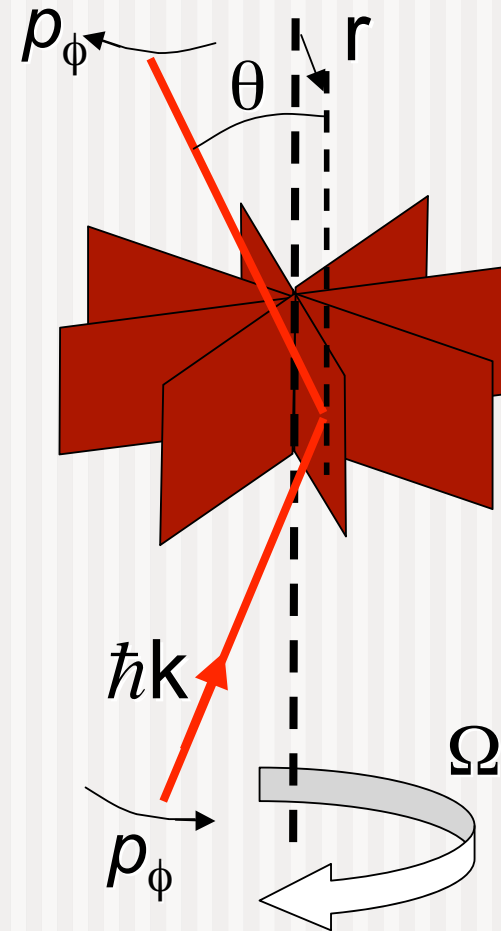
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Padgett, J Opt A 2004

# Rotationally induced frequency shifts

- Waveplate reverses  $p_\phi$
- Exerts force
  - $F = 2\hbar k \sin\theta$
- Skew angle of ray
  - $\theta = \ell / kr$
- Work done (per photon)
  - $W = Fv = 2 \ell \hbar \Omega$
- Frequency shift
  - $\Delta\omega = 2\ell\Omega$

Ray-Optics/Photon Pressure gives the right answer



# Non-linear optics

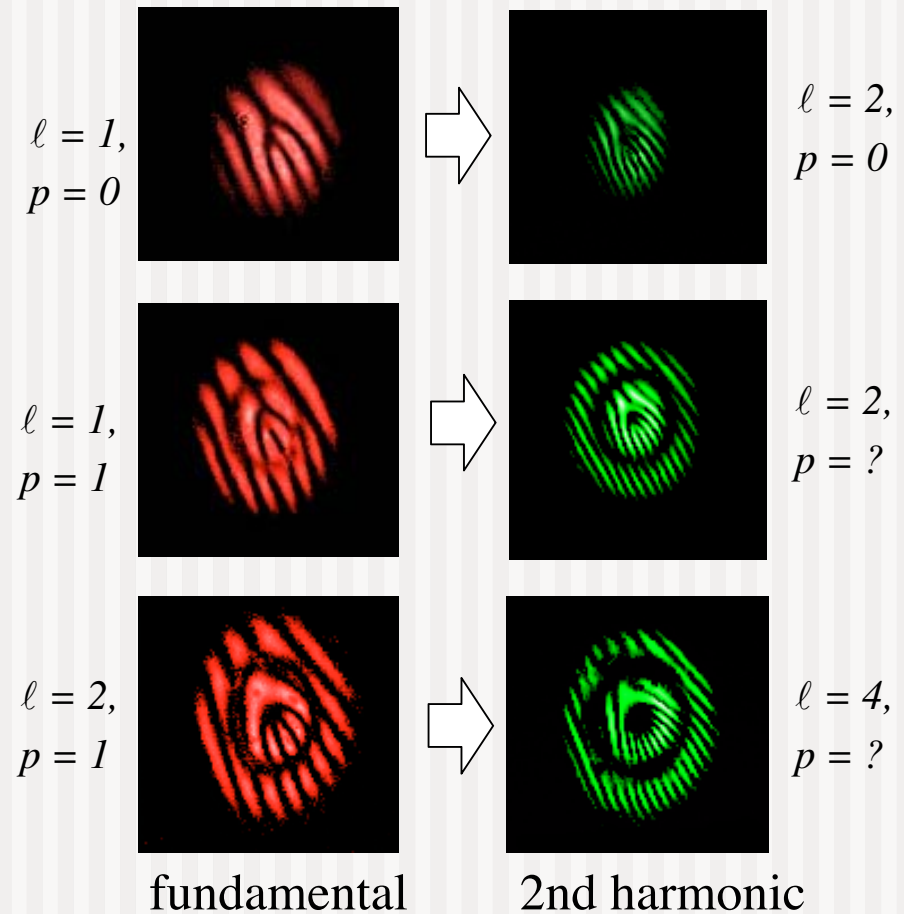
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Courtial *et al.* Phys. Rev. A 1997

Mair *et al.* Nature 2001

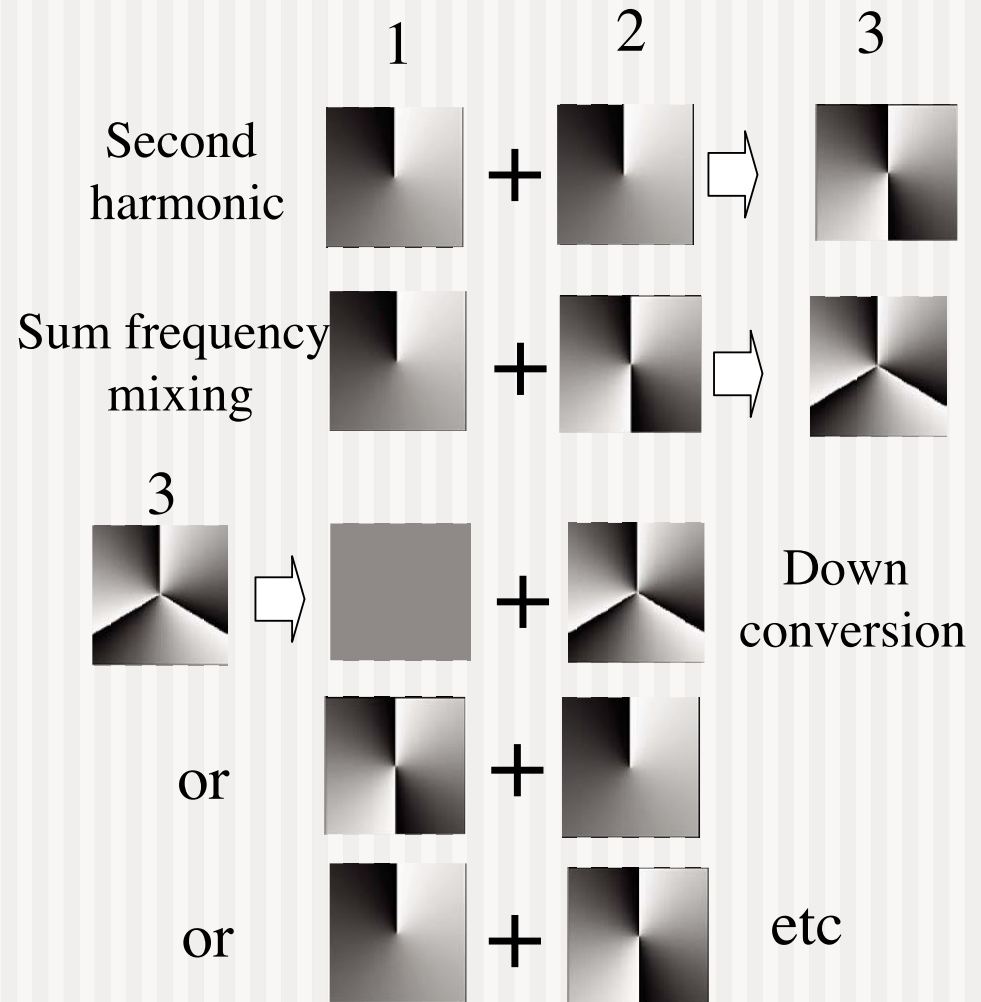
# OAM conserved in SHG

- OAM conserved in the light beam
- c.f. SAM in which OAM is not conserved
- But, down conversion is more complicated!



# OAM in three-wave interactions

- Fixed phase relationship between three fields
  - $\psi_1 + \psi_2 + \psi_3 = \pm\pi/2$
- Azimuthal phase terms are linked to each other, giving
  - $l_1 + l_2 = l_3$



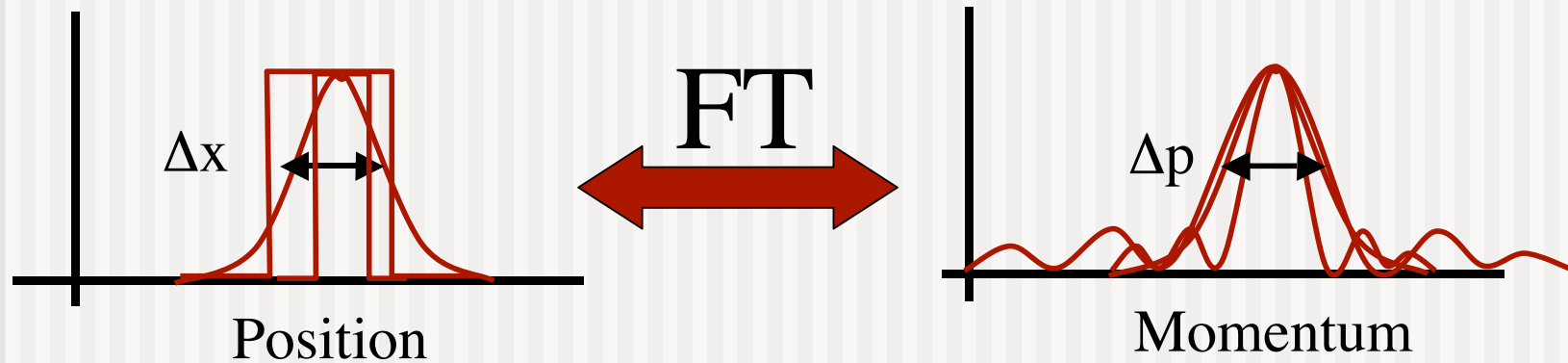
# The (angular) uncertainty principle

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Franke-Arnold *et al.* NJP 2004

EPSRC 2004-2006

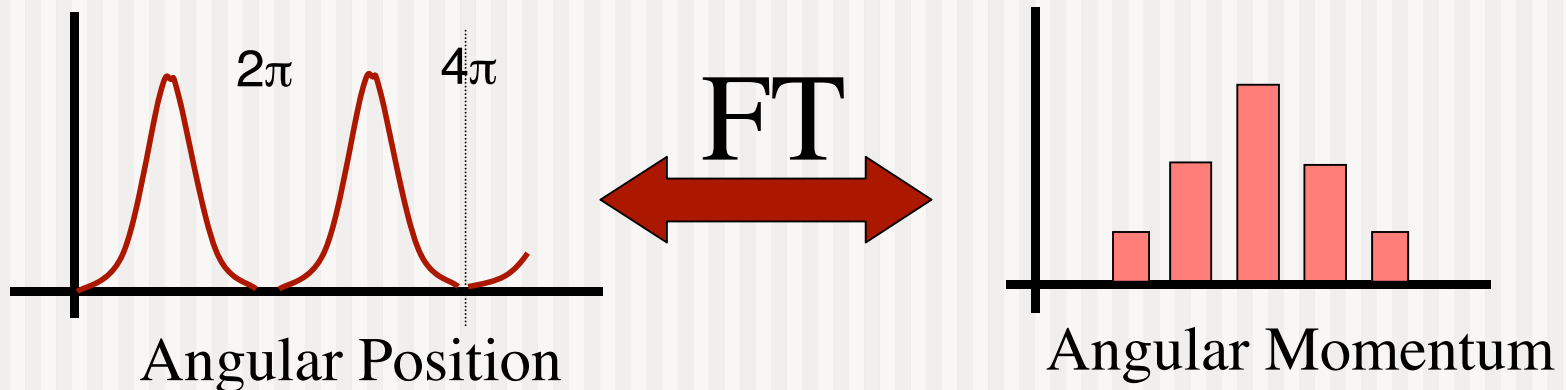
# Uncertainty relationships



- Heisenberg's Uncertainty principle
  - $\Delta x \Delta p \geq \hbar/2$
- For Gaussian distribution
  - $\Delta x \Delta p = \hbar/2$  (Gaussian gives minimum uncertainty state product)
- What about angular momentum?



# Uncertainty in Angular Momentum



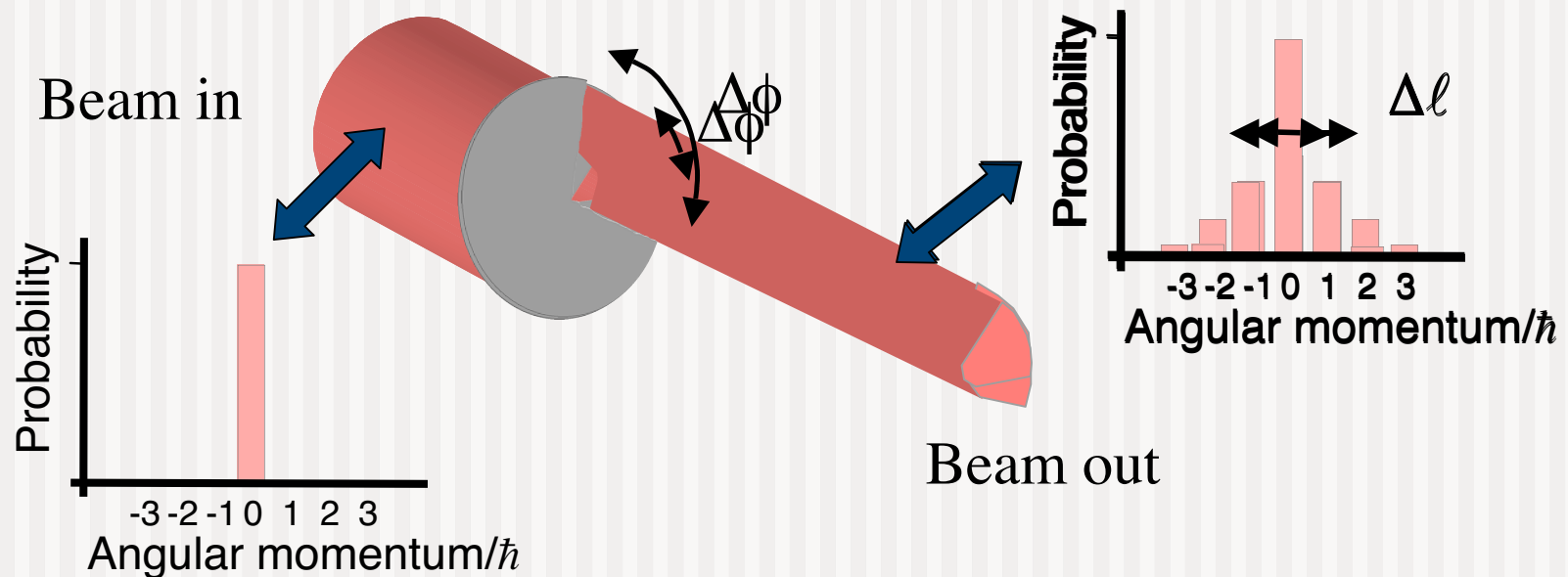
- Angular position repeats modulo  $2\pi$ ,
  - FT of repeating position gives discrete angular momentum values
  - $\Delta\phi\Delta L \geq \hbar/2$  ?
- For no restriction on  $\phi$ ,  $\Delta\phi$  still finite, but L can be measured exactly
  - With no restriction,  $\Delta\phi = \pi/\sqrt{3}$ , but  $\Delta L = 0$ , i.e.  $\Delta\phi\Delta L = 0$
- What are the minimum uncertainty states?

# Angular position

# Angular momentum

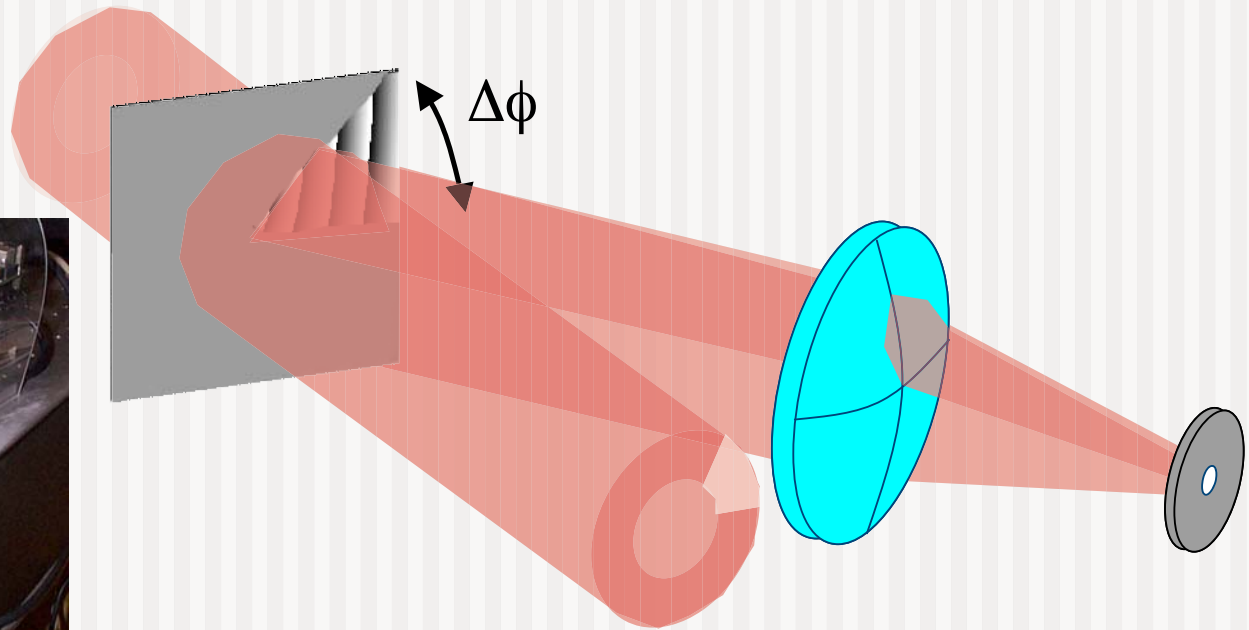
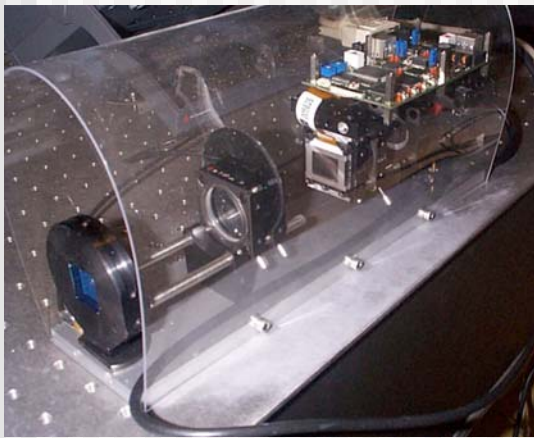


- A “cake-slice” aperture placed in a light beam restricts the angular position ( $\approx$  of the photon)



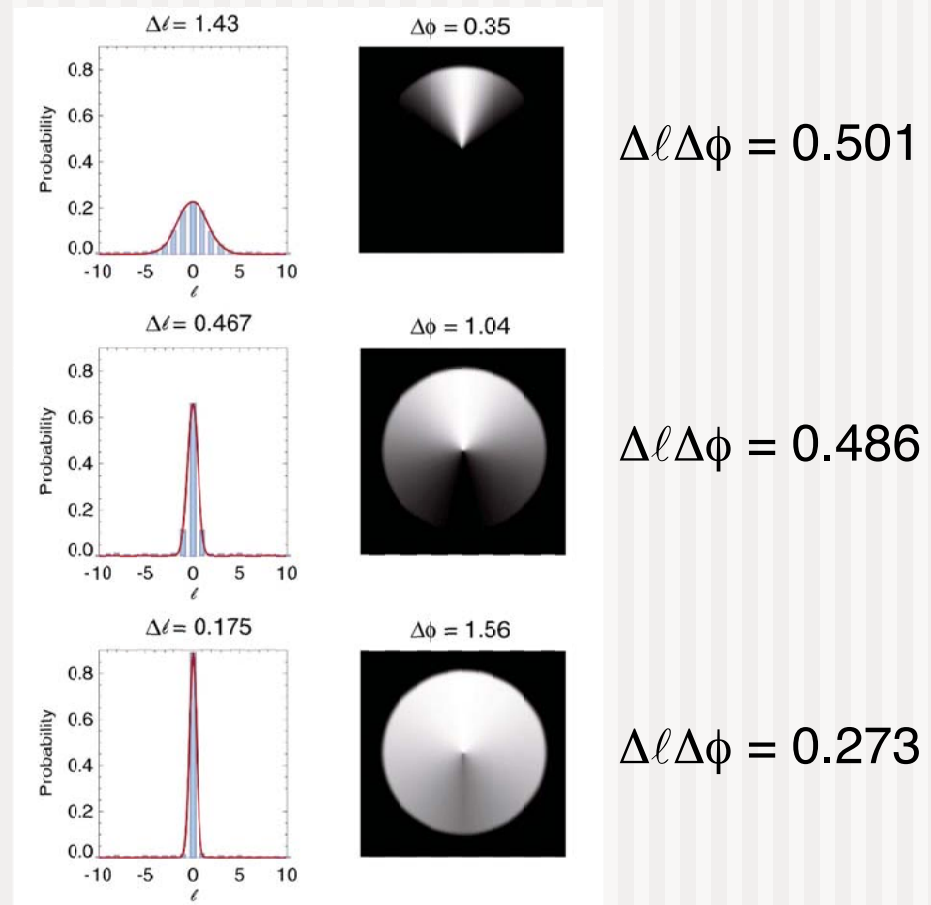
# Doing the experiment

- SLM used to measure  $\ell$
- Same SLM used to impose aperture
- Transmission through aperture gives  $\ell$  components



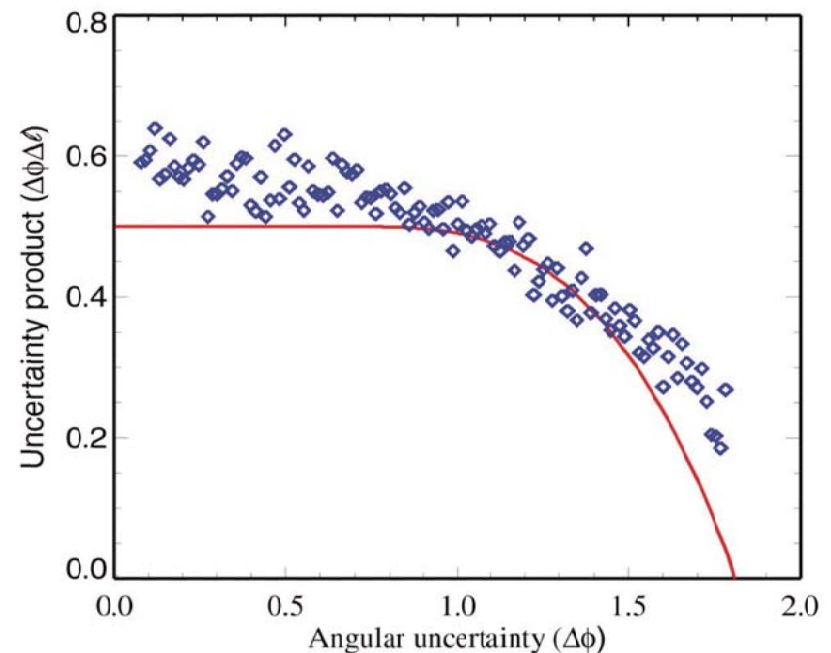
# The uncertainty relationship for angular momentum

- For small “slices” we find an uncertainty relationship
  - $\Delta L \Delta \phi \geq \hbar/2$
  - i.e.  $\Delta \ell \Delta \phi \geq 1/2$
- For no aperture
  - $\Delta L = 0$
- More complicated for large “slices” .....



# The minimum uncertainty states

- More complicated for large “slices”.....
  - $\Delta\ell\Delta\phi \geq 1/2(1-2\pi(P_{\text{edge}}))?$
- Minimum uncertainty position state
  - Gaussian symmetrically truncated by  $-\pi$  to  $+\pi$ ?
- Results inconclusive?



# The OAM communicator

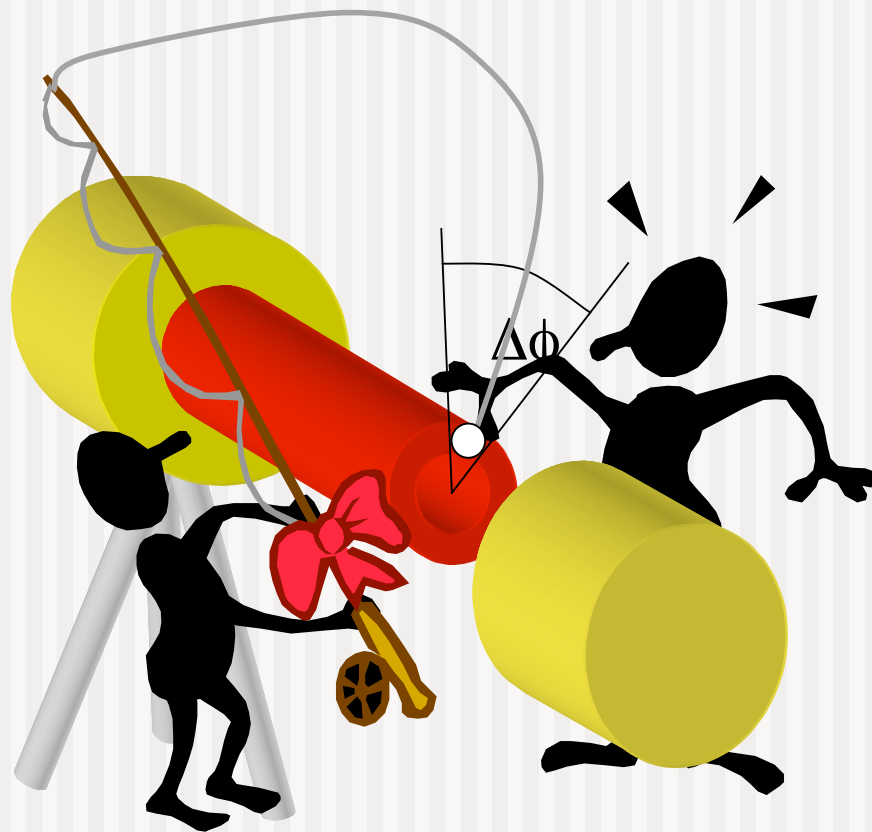
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Gibson *et al.*  
Opt Express **12**,  
5448 (2004)



# Measuring OAM gives secure(ish) communication

- Use OAM (i.e.  $l$ -index) to encode data on a light beam
  - Each photon can take ANY value of  $l$  hence increases data capacity
- $\Delta l \Delta \phi \geq 1/2$  gives security
  - Can't measure  $l$  from only part of beam
  - e.g. Can't measure  $l$  from scattered light
  - e.g. Can't measure  $l$  from side lobe



# Free-space comms

supported by  **Scottish Enterprise**  
proof of concept fund



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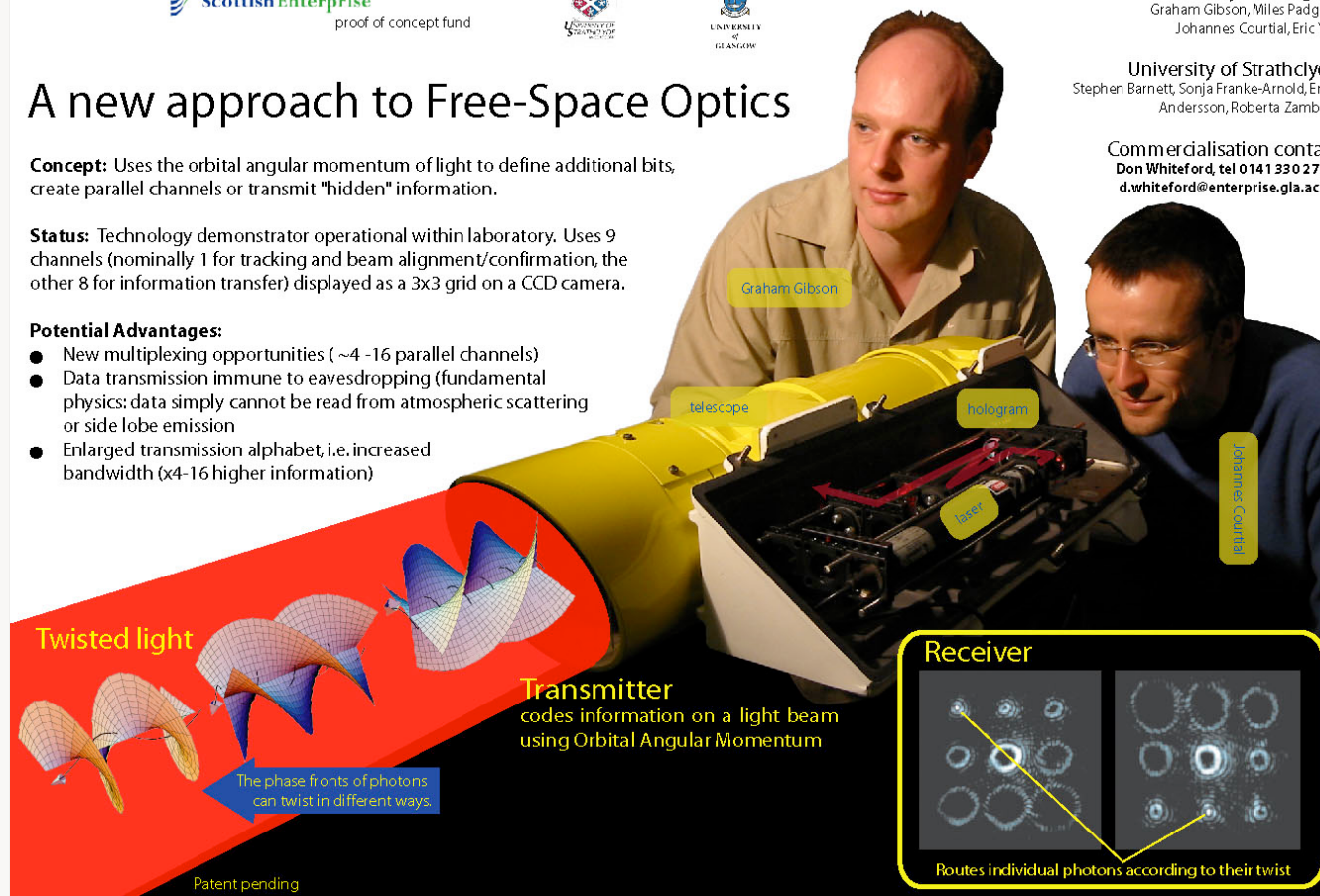
## A new approach to Free-Space Optics

**Concept:** Uses the orbital angular momentum of light to define additional bits, create parallel channels or transmit "hidden" information.

**Status:** Technology demonstrator operational within laboratory. Uses 9 channels (nominally 1 for tracking and beam alignment/confirmation, the other 8 for information transfer) displayed as a 3x3 grid on a CCD camera.

### Potential Advantages:

- New multiplexing opportunities (~4-16 parallel channels)
- Data transmission immune to eavesdropping (fundamental physics: data simply cannot be read from atmospheric scattering or side lobe emission)
- Enlarged transmission alphabet, i.e. increased bandwidth (x4-16 higher information)



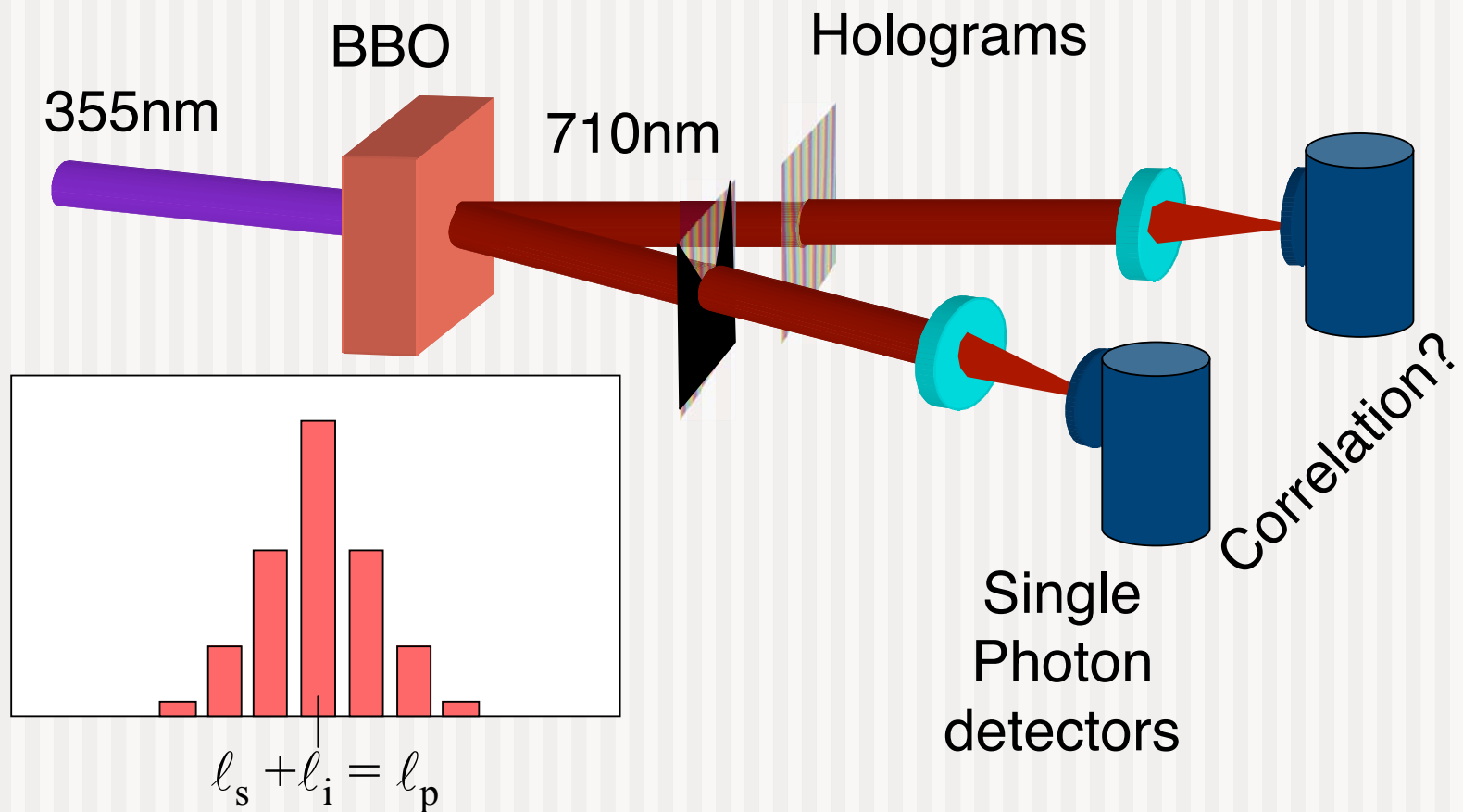


# Entanglement of angular position and angular momentum

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EPSRC 2004-2006

# Measuring angle and angular momentum



# The (angular) momentum paradox

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Brevik Phys Rep 1972

Loudon Phys Rev. A 2003

Padgett *et al.* J. Mod Opt 2003

Mansuripur Opt. Exp 2005

EPSRC 2005-2008

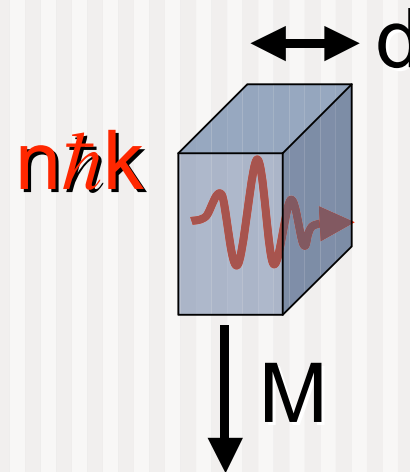
# The momentum of light in a dielectric

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- What is the momentum of light (a photon) inside a dielectric (refractive index  $n$ )?
  - $p = n\hbar k$  (Minkowski) {equiv.  $\mathbf{p} = \mathbf{D} \times \mathbf{B}$ }
  - $p = \hbar k/n$  (Abraham) {equiv.  $\mathbf{p} = \mathbf{E} \times \mathbf{H}/c^2$ }
- What is the angular momentum of light (a photon) inside a dielectric (refractive index  $n$ )?
  - $L = (l + \sigma) \hbar$  (Minkowski)
  - $L = (l + \sigma) \hbar/n^2$  (Abraham)

# Minkowski

- A short, single-photon pulse traverses a block
- During transit of block, momentum of photon is increased
- Block moves towards source



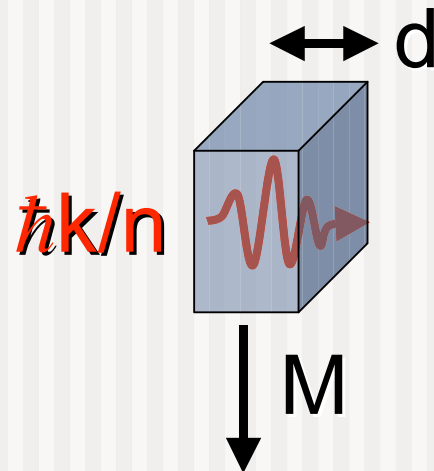
$$p_{block} = (1-n)\hbar k$$

block moves

$$\Delta z = (n-n^2)\hbar k d/cM$$

# Abraham

- A short, single-photon pulse traverses a block
- During transit of block, momentum of photon is decreased
- Block moves away from source



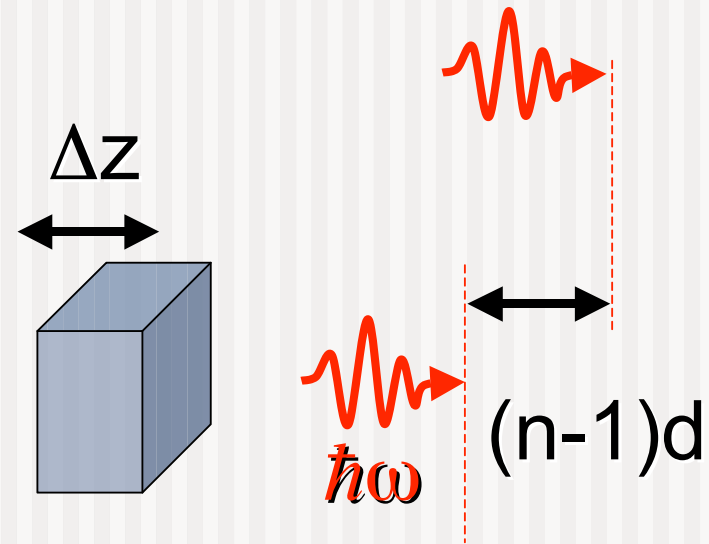
$$p_{block} = (1 - 1/n)\hbar k$$

block moves

$$\Delta z = (n-1)\hbar k d/cM$$

# Einstein Box $\rightarrow$ Abraham

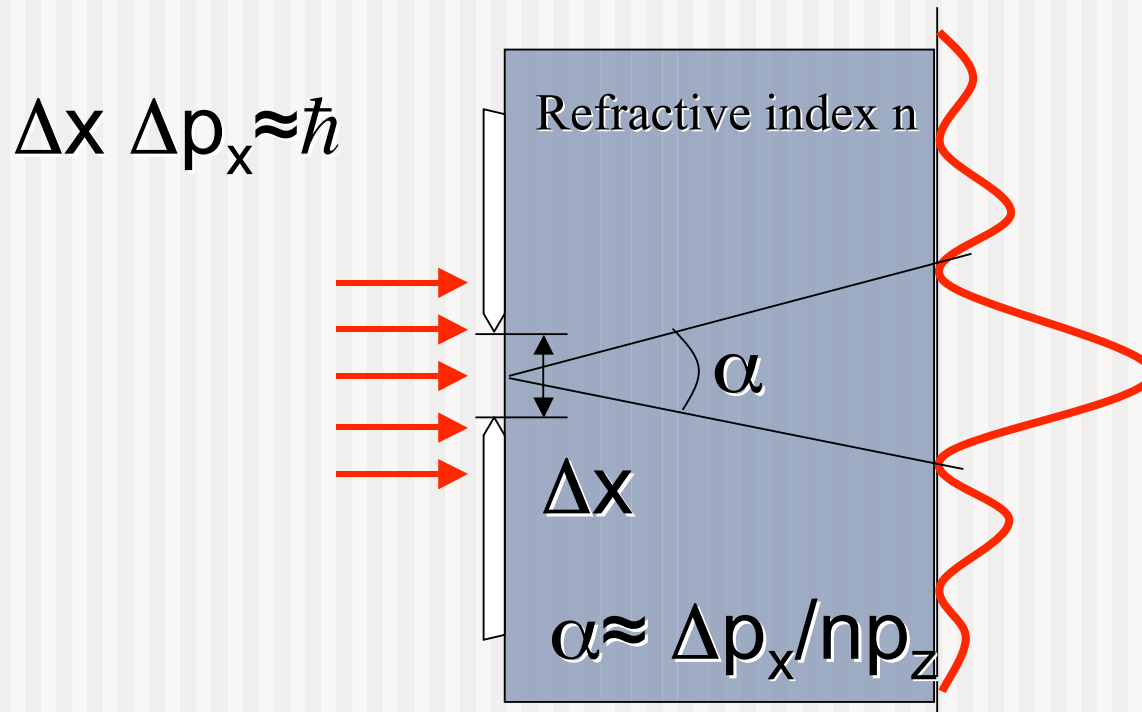
- Alternative argument based only on displacement of the centre of mass-energy
- Delay of photon energy equated to energy displacement of block
- Agrees with Abraham formulation



$$\Delta z M c^2 = (n-1) d \hbar \omega$$

$$\Delta z = (n-1) \hbar k d / c M$$

# Diffraction -> Minkowski

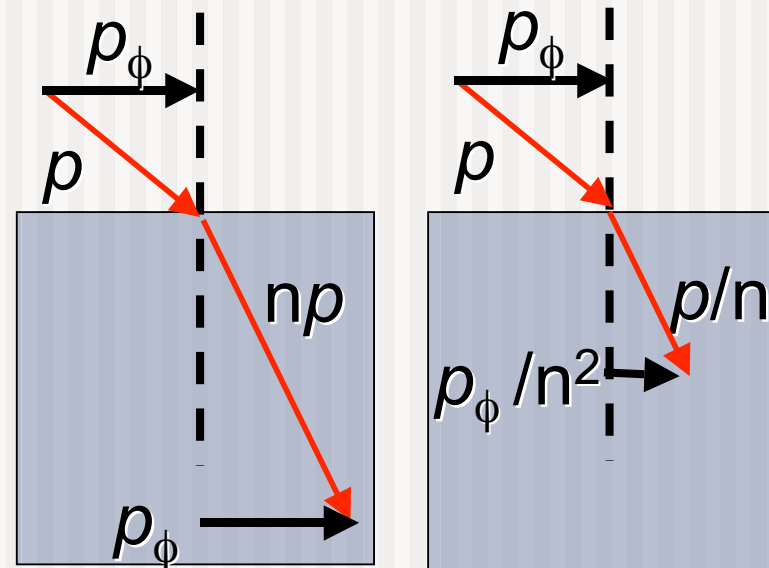


- Diffraction arises from Uncertainty Principle
- Shrinking pattern implies Minkowski formulation



# Angular momentum

- Angular momentum arises from  $\phi$  comp. of  $p$
- At interface,  $\phi$  comp. reduced by  $1/n$  (Snell's Law)
- True for both Abraham and Minkowski
- Implications re AM
  - Minkowski  $L = \ell \hbar$
  - Abraham  $L = \ell \hbar / n^2$



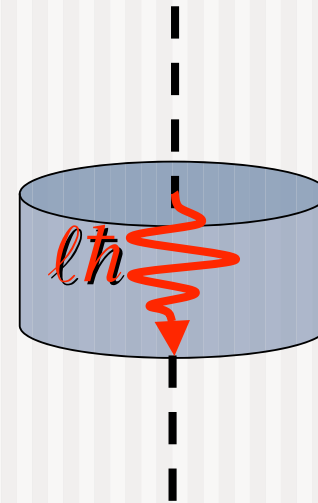
Minkowski

Abraham

# Minkowski equiv for AM

- Single photon pulse carrying angular momentum
- During transit of block, angular momentum is unchanged
- Block does not rotate

$$\Omega_{block} = 0$$

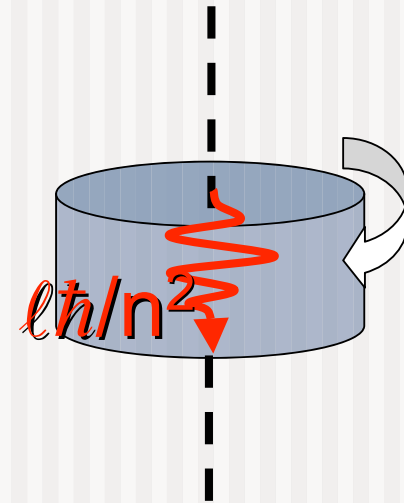


$$\Delta\phi_{block} = 0$$

# Abraham equiv for AM

- Single photon pulse carrying angular momentum
- During transit of block, angular momentum is changed
- Block does rotate

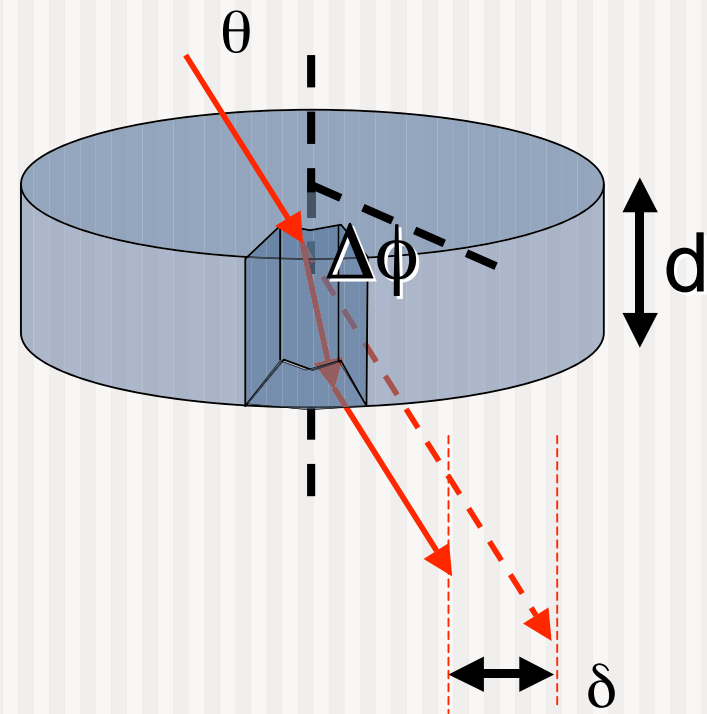
$$\Omega_{block} = (1 - 1/n^2) \ell \hbar / I$$



$$\Delta\phi_{block} = (n - 1/n) \ell \hbar d / cI$$

# Einstein Box for Angular Momentum

- Equate the lateral delay of the photon energy to the mass energy displacement of the disc element
- Sub. in for skew angle ( $\theta = \ell / kr$ ) and integrate over disc
  - $\Delta\phi_{block} = (n-1/n)\ell\hbar d/cI$   
( $I =$  moment of inertia)
- The Abraham result!
- Also true for spin AM?
  - $\Delta\phi_{block} = (n-1/n)(\ell + \sigma)\hbar d/cI$



$$\Delta\phi_{block} r M c^2 = \delta \hbar \omega = \hbar \omega (n-1/n) d \theta$$

# The Mechanical rotation (Faraday) Effects

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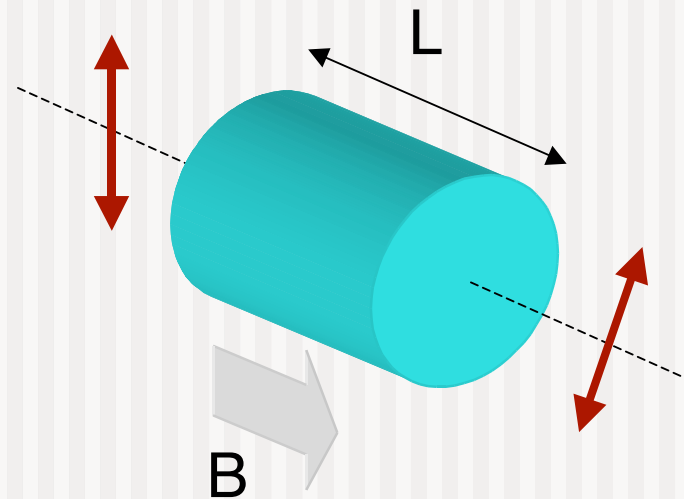
Jones *et al.* J. Proc. Roy Soc. A 1976

Nienhuis *et al.* Phys Rev A 1992

SUPA 2005-2006

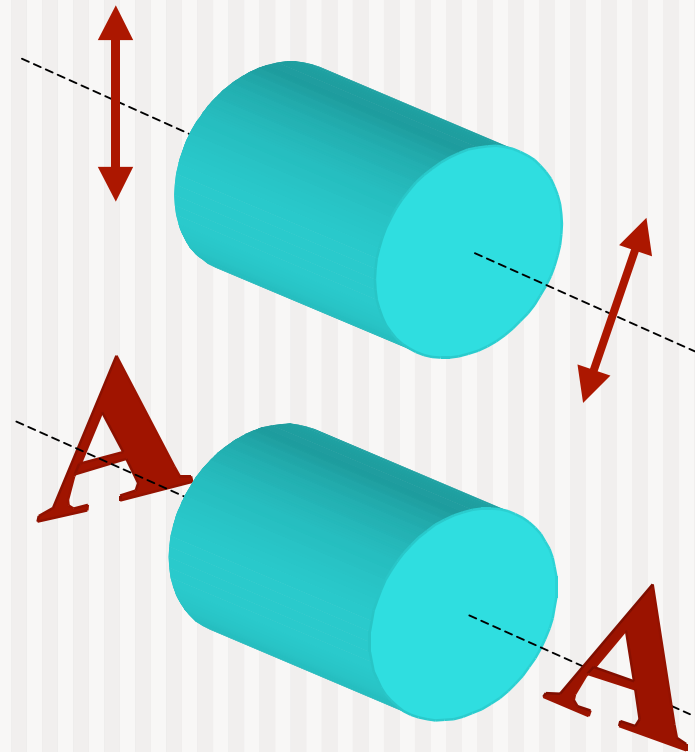
# Magnetic Faraday effect

- Rotation of plane polarised light
  - $\Delta\theta = BLV$ 
    - $V$  Verdet constant
- OR treat as phase delay of circularly polarised light
  - $\Delta\phi = \sigma BLV$
- Are SAM and OAM equivalent?



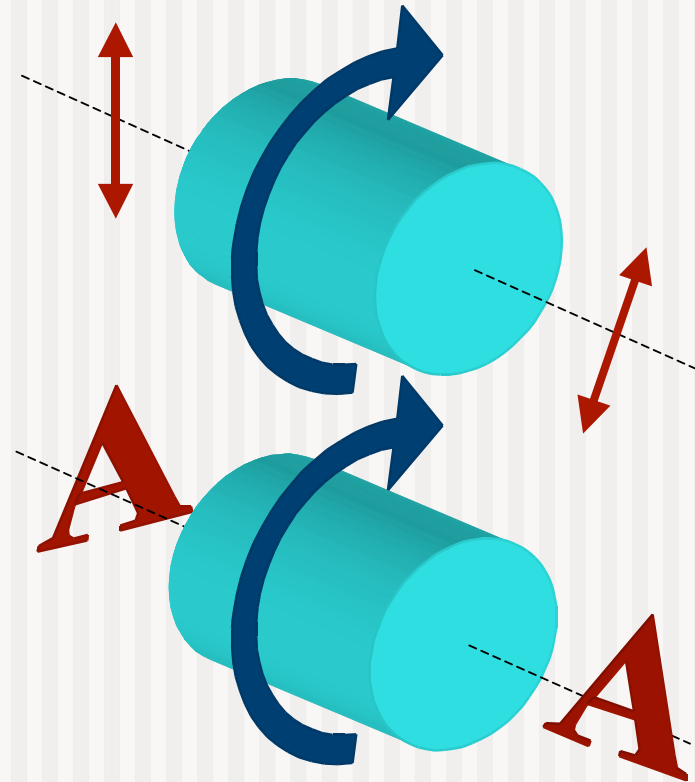
# Rotations of Polarisation and Image

- SAM -> Polarisation rotation
- OAM-> Image rotation
- Look through a Faraday isolator ( $\Delta\theta\approx 45^\circ$ ), is the “world” rotated - NO
  - SAM and OAM are not equivalent in the Magnetic Faraday effect



# Mechanical Faraday effect

- Photon drag, gives Polarisation rotation
  - $\Delta\theta = \Omega(n-1/n)L/c$
- Phase delay equiv.
  - $\Delta\phi = \sigma\Omega(n-1/n)L/c$
- Does photon drag give image rotation?
  - $\Delta\theta = \Omega(n-1/n)L/c$





# Photon drag

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- Transverse photon drag
  - $u_x = (1 - 1/n^2)v$
- For transit time  $Ln/c$ , gives displacement
  - $\Delta x = L(1 - 1/n^2)nv/c$
- In cylindrical frame
  - $\Delta\theta r = L(1 - 1/n^2)nr\Omega/c$
  - $\Delta\theta = L(n - 1/n)\Omega/c$
- OAM equiv. SAM?

