

Photo emission electron microscopy

Outline

Introduction

Electron optics:

Resolution

Transmission

Instruments:

PEEM

PEEM + LEEM

SPELEEM

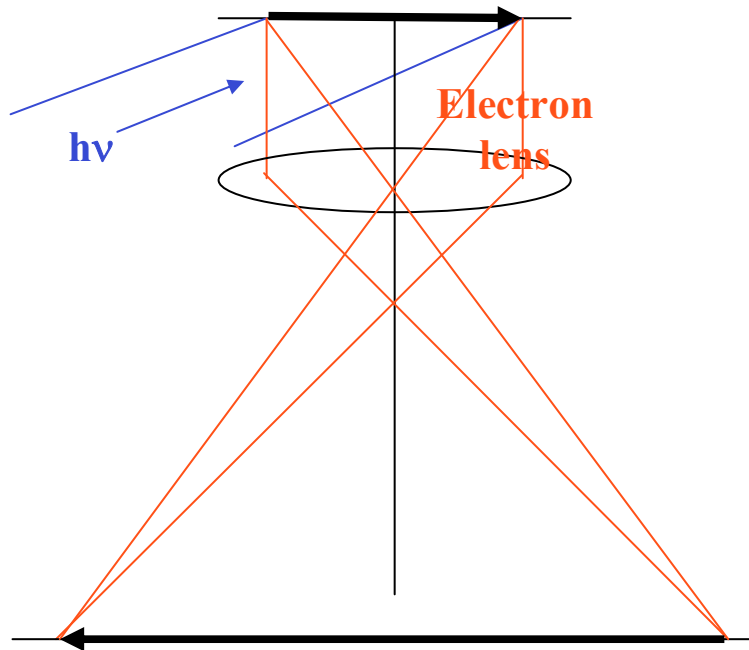
Methodic

Applications:

Magnetic imaging

Photo Emission Electron Microscopy (PEEM)

2 types

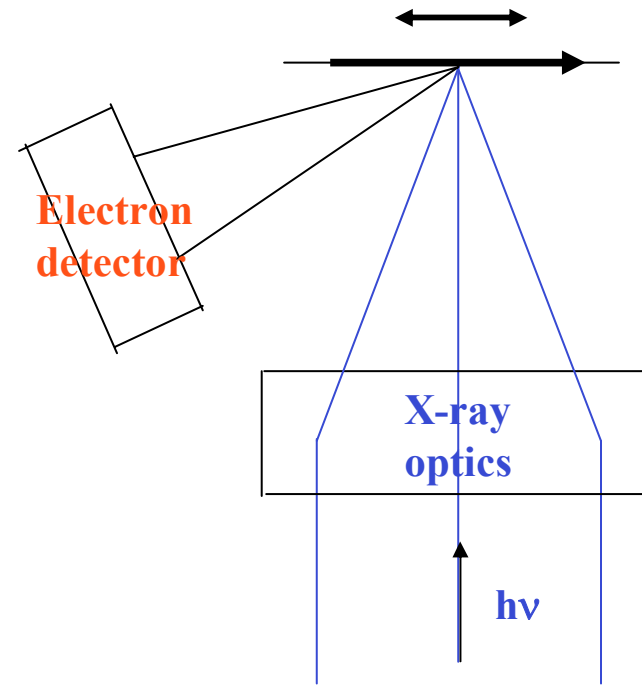


broad illumination

Full field

sample fixed

Bauer, Locatelli



focused illumination

Scanning

sample scanned

Kiskinova

PEEM

3 imaging modes

1 XPEEM

Photo electrons PE

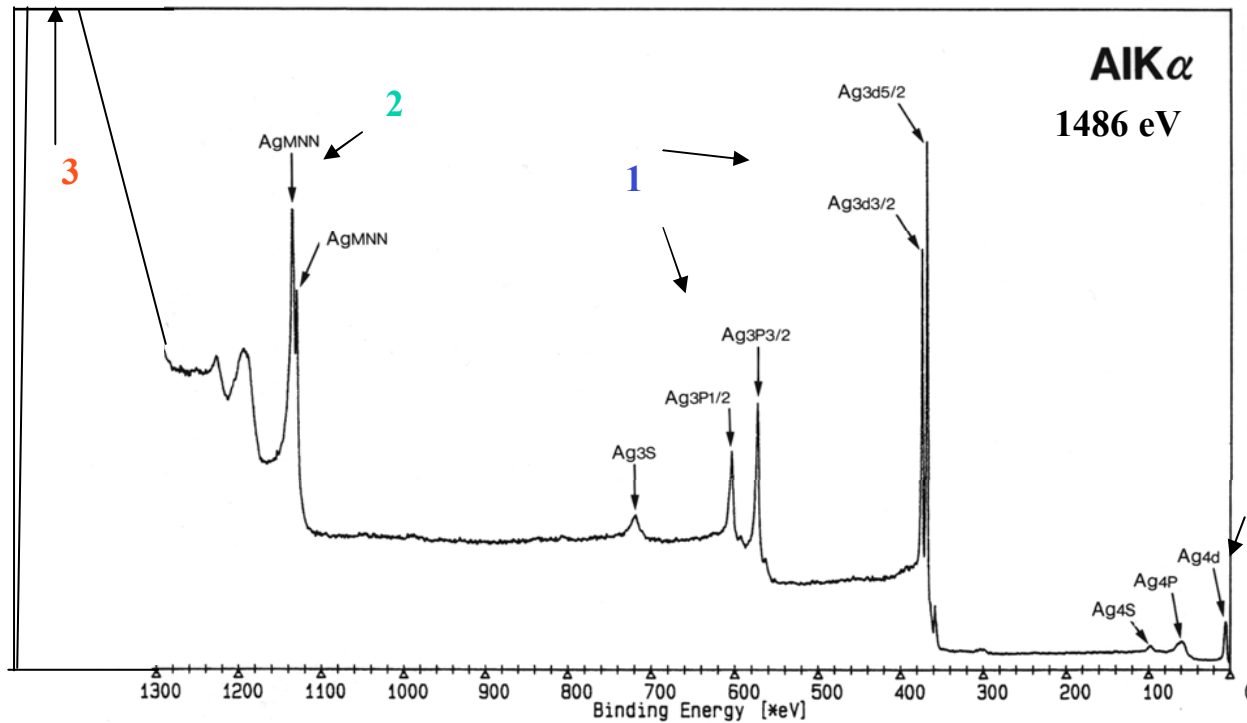
2 XAEEM

Auger electrons AE

} with energy filter

3 XSEEM

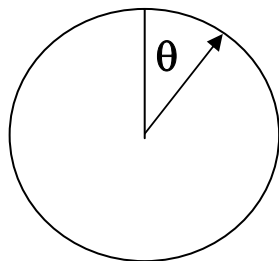
Secondary electrons SE



Angular distribution

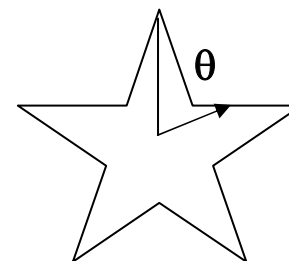
Internal

Amorphous, polycrystalline, SE



$$I_i(\theta) = \text{const.}$$

single crystalline, PE, AE

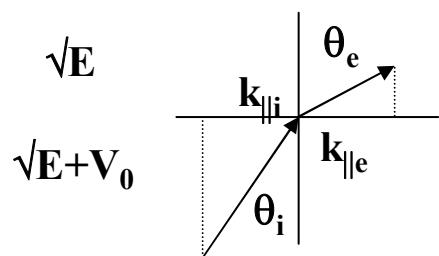


$I_i(\theta)$ due to diffraction

Internal (i) → External (e)

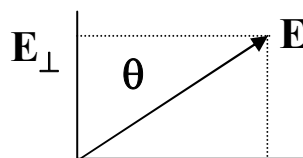
$n \sim$

Refraction



$k_{||}$ conservation $k_{||e} = k_{||i}$

For escape

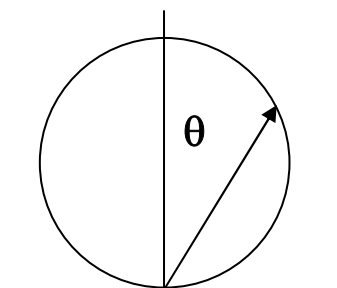


$$E_{\perp} = E \cos \theta$$

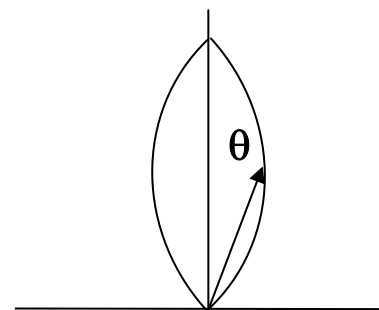
necessary:

$E_{\perp} > \Phi$ (work function)
 I (ionization energy)
 U (HOMO)

External



$$I_e(\theta) = \cos \theta$$



Electron optics

The cathode lens

In emission microscopy $\theta \equiv \alpha_0$ is large

Electron lenses can accept only small $\theta \equiv \alpha_0$ because of large chromatic and spherical aberrations

Solution of problem: accelerate electrons to high energy before lens



Immersion objective lens = cathode lens

$$n \sin \theta = \text{const}$$

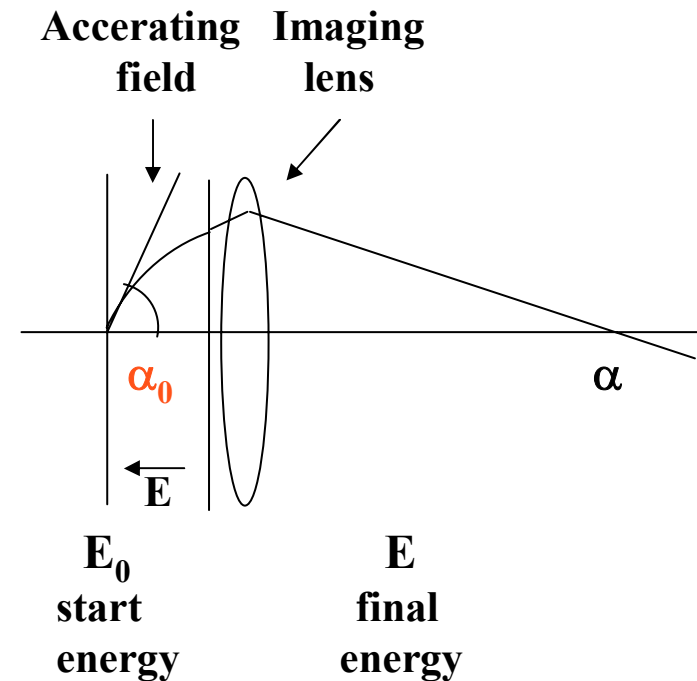
$$n \sim v \sim \sqrt{E}$$

$$\theta \rightarrow \alpha$$

$$\sin \alpha / \sin \alpha_0 = \sqrt{E_0 / E}$$

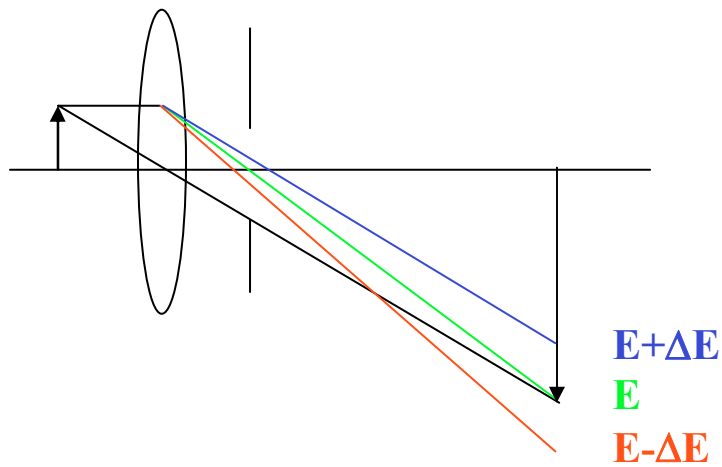
Example for $E = 20000$ eV:

| | | |
|------------------------------------|-------------|-------------|
| E_0 | 2 eV | 200 eV |
| α for $\alpha_0 = 45^\circ$ | 0.4° | 4.5° |



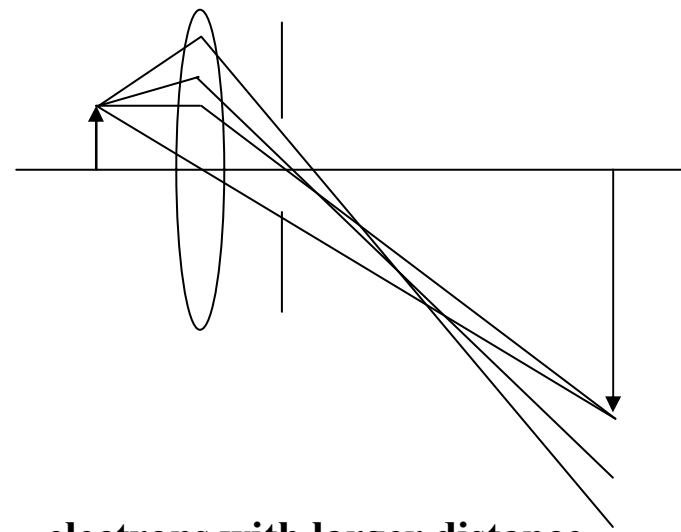
Aberrations

chromatic



**slower (faster) electrons
are more (less) deflected**

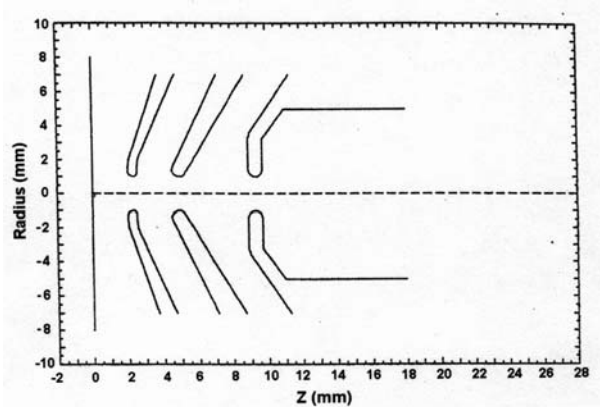
spherical



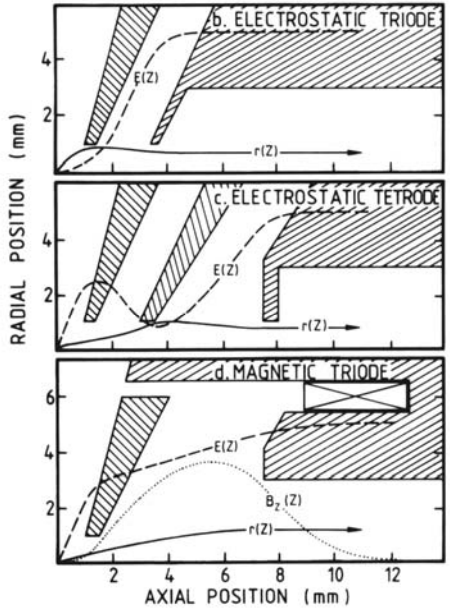
**electrons with larger distance
from axis are more deflected
(stronger field!)**

Cathode lens types

Electrostatic tetrode

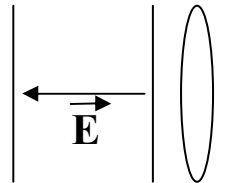


Lens comparison



←
Magnetic diode

Estimation of aberrations:
Separate lens into acceleration and imaging regions

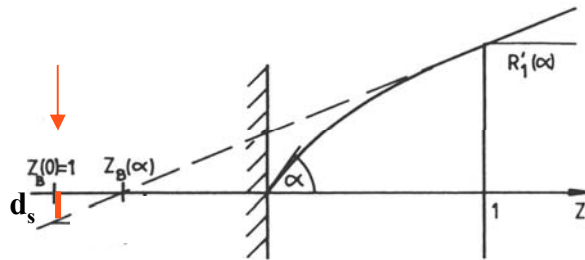


At low energies the aberrations of the accelerating region dominate

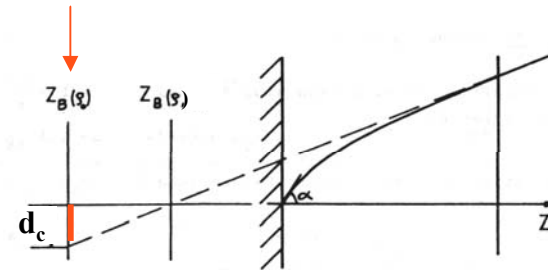
Aberrations of homogeneous acceleration field

$$\rho_0 = E_0/E \quad \varepsilon = \Delta E_0/E \quad \rho = \rho_0 + \varepsilon$$

Spherical aberration d_s



Chromatic aberration d_c



Analytical solution

Approximation: ρ_0 and $\varepsilon \ll 1/\cos^2 \alpha > 1$

Example: $E_0 = 100 \text{ eV}$, $\Delta E_0 = 1 \text{ eV}$, $E = 20000 \text{ eV}$

$$\varepsilon = \rho_0 / 100, \quad \rho_0 = 1/200$$

$$d_s \approx 2 \rho \sin \alpha (1 - \cos \alpha)$$

$$\approx \rho \alpha^3 \quad \text{for small } \alpha$$

$$d_c \approx 2 \rho \sin \alpha (\sqrt{\rho_0 / \rho} - 1)$$

$$\approx \varepsilon \sin \alpha \quad \text{for } \varepsilon \ll \rho_0 \approx \rho$$

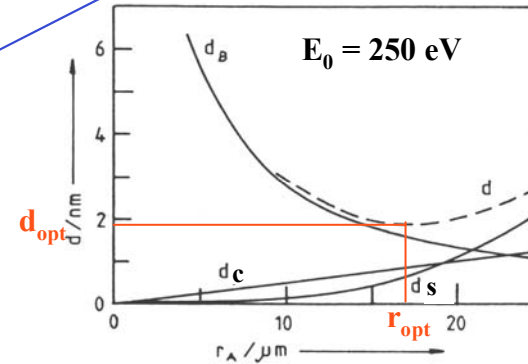
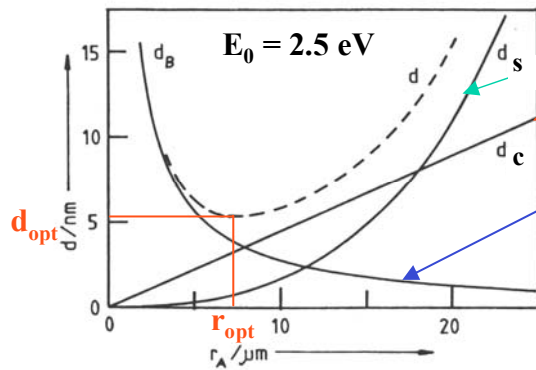
$$\approx \varepsilon \alpha \quad \text{for small } \alpha$$

α -dependent aberrations require α -limitation by angle-limiting aperture (“contrast aperture”) with radius r_A



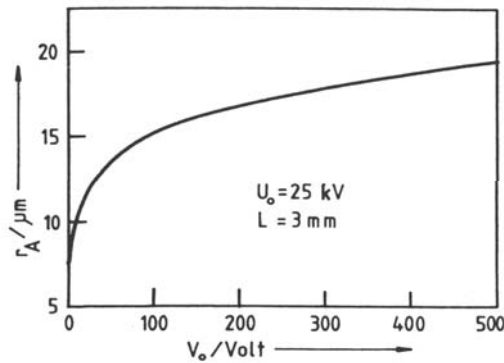
Diffraction by aperture: diffraction disc of confusion $d_B = 0.6 \lambda / r_A$

Approximate resolution $d = \sqrt{d_s^2 + d_c^2 + d_B^2}$

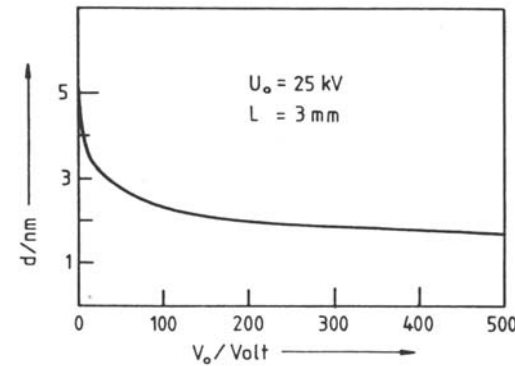


$L = 3 \text{ mm}$ $E = 25000 \text{ eV}$ $\Delta E_0 = 0.25 \text{ eV}$

Optimum aperture radius



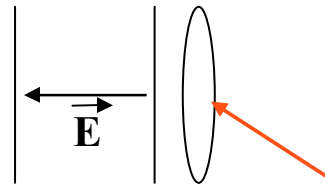
Optimum resolution



Note: small angle approximation $\sin \alpha \approx \alpha \sim r$

Complete lens

Combine acceleration and imaging regions



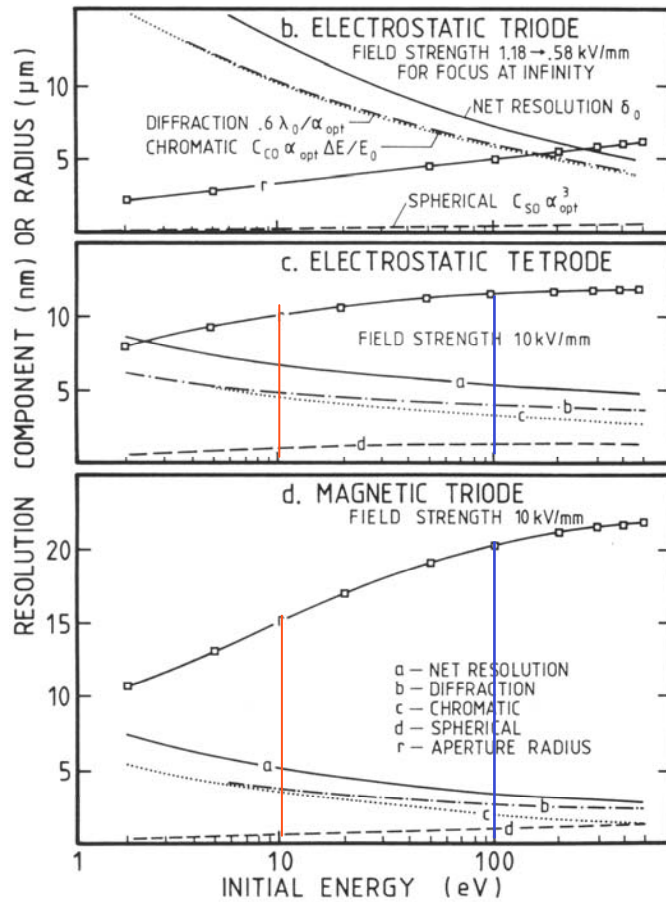
At low energies aberrations of accelerating region dominate

but

at high energies the spherical aberration of second part of lens becomes important

Resolution and optimum aperture of real lenses

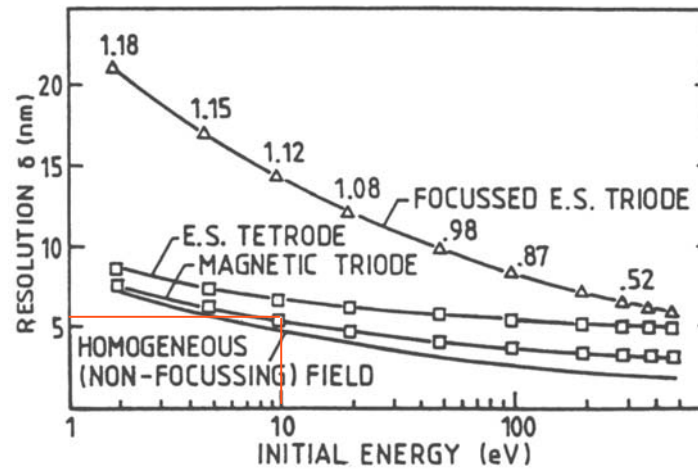
Optimum aperture r and resolution-limiting contributions



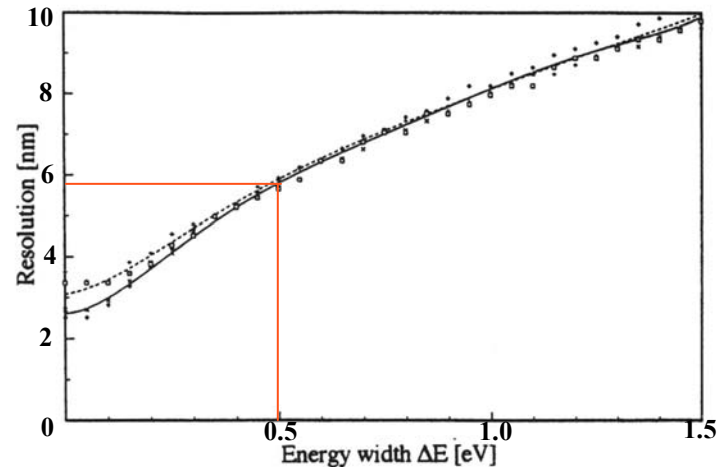
J. Chmelik et al, Optik 83 (1989)155

T. Müller, M.S. thesis, TU Clausthal 1995

Resolution with optimum aperture
E-dependence at fixed $\Delta E = 0.5$ eV, $U_0 = 20$ keV



ΔE -dependence at fixed $E = 10$ eV, $U_0 = 18$ kV magnetic triode



Transmission

limited by angle accepted by contrast aperture (r_A)

Axial distance (in back focal plane) of electron starting at angle α
 $r \approx f \sin \alpha \sqrt{E_0/E}$ (f focal length)



$$\sin \alpha \approx (r/f) \sqrt{E_0/E}$$

Examples for $f = 10$ mm, $E = 20000$ eV, $r_A = 10$ μ m

| | | |
|---------------|-------|--------|
| E_0 | 2 eV | 200 eV |
| $\sin \alpha$ | 0.2 | 0.02 |
| α | 11.5° | 1.15° |

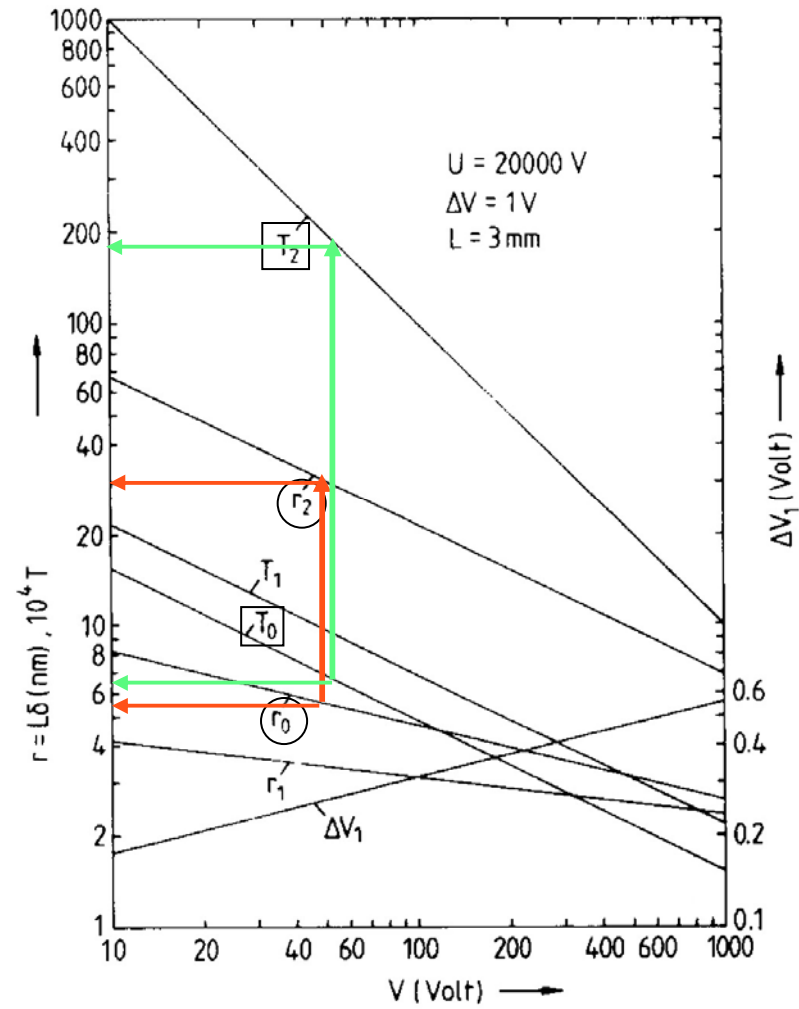
In emission microscopy (wide α range) optimum resolution condition
reduces transmission T , therefore

optimize T^n/d^2 instead of $1/d^2$

For $\cos \alpha$ distribution $T = \pi \sin^2 \alpha$

$$T^n/d^2 = \pi \sin^{2n}/d^2$$

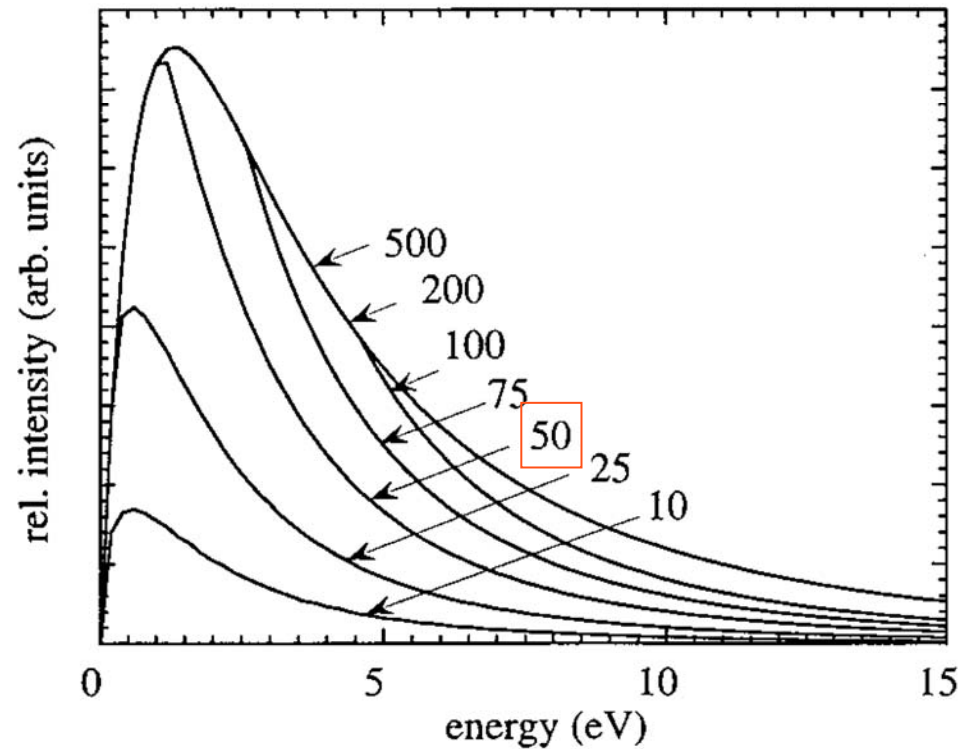
Transmission T_n , resolution r_n of homogeneous field



T_2
T and d
equally
weighted:

50 eV:
 $\Delta T \cong 30$
 $\Delta d \cong 5$

Influence of angle-limiting aperture on the energy distribution of secondary electrons

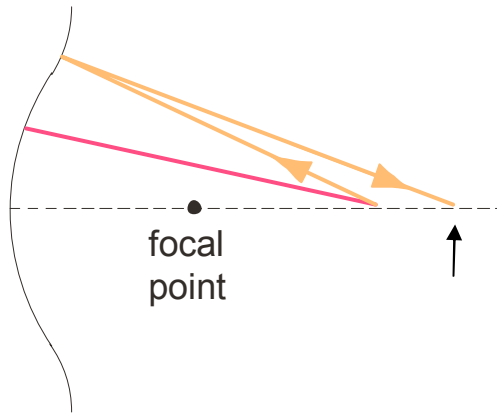
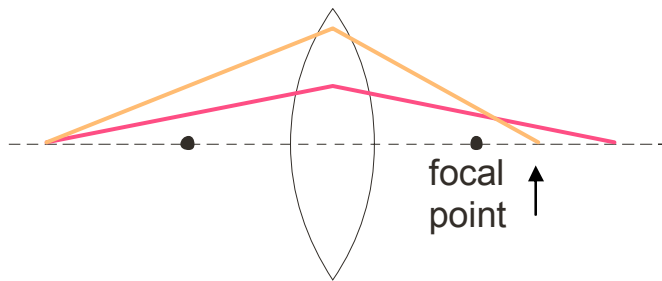


Work function $\Phi = 4 \text{ eV}$, accelerating voltage $V = 20 \text{ kV}$
Parameter: aperture diameter in μm , ALS PEEM

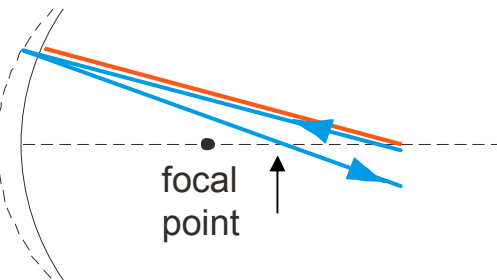
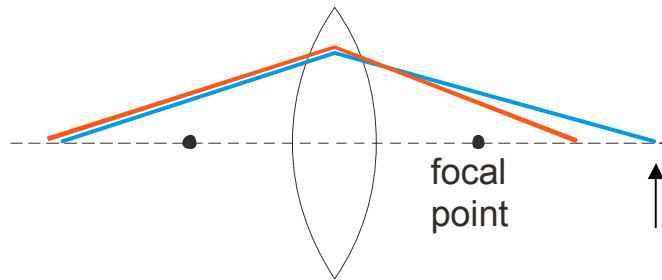
Aberration correction in electron optics

Round **convex** lenses

electrostatic mirror

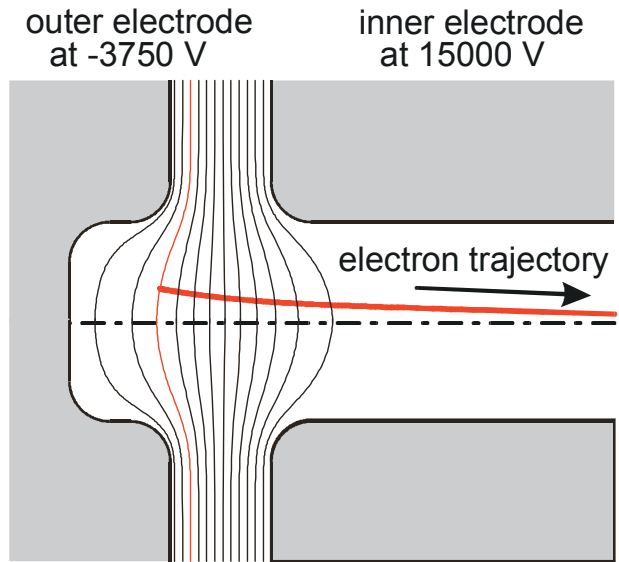


Spherical aberration



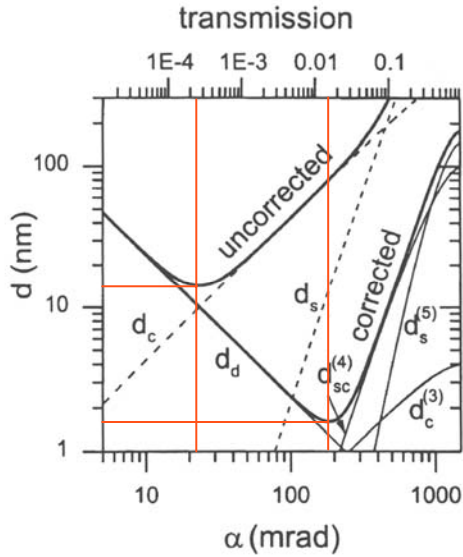
Chromatic aberration

Equipotential surfaces
in a diode mirror



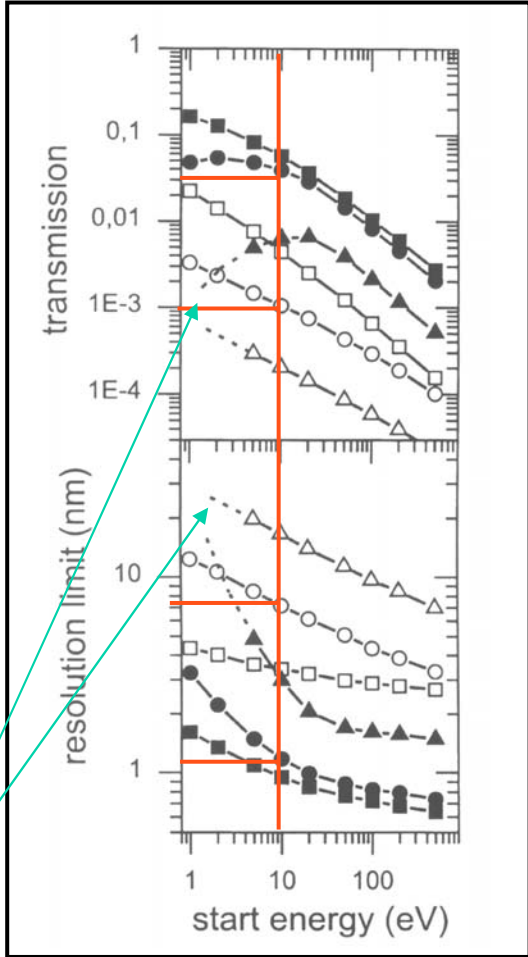
Resolution and transmission improvement with aberration correction

Example: SMART



$E_0 = 10 \text{ eV}$, $\Delta E = 2 \text{ eV}$, $F = 5 \text{ kV/mm}$

Calculations: D. Preikszas
From Th. Schmidt et al,
Surf. Rev. Lett. 9 (2002) 223

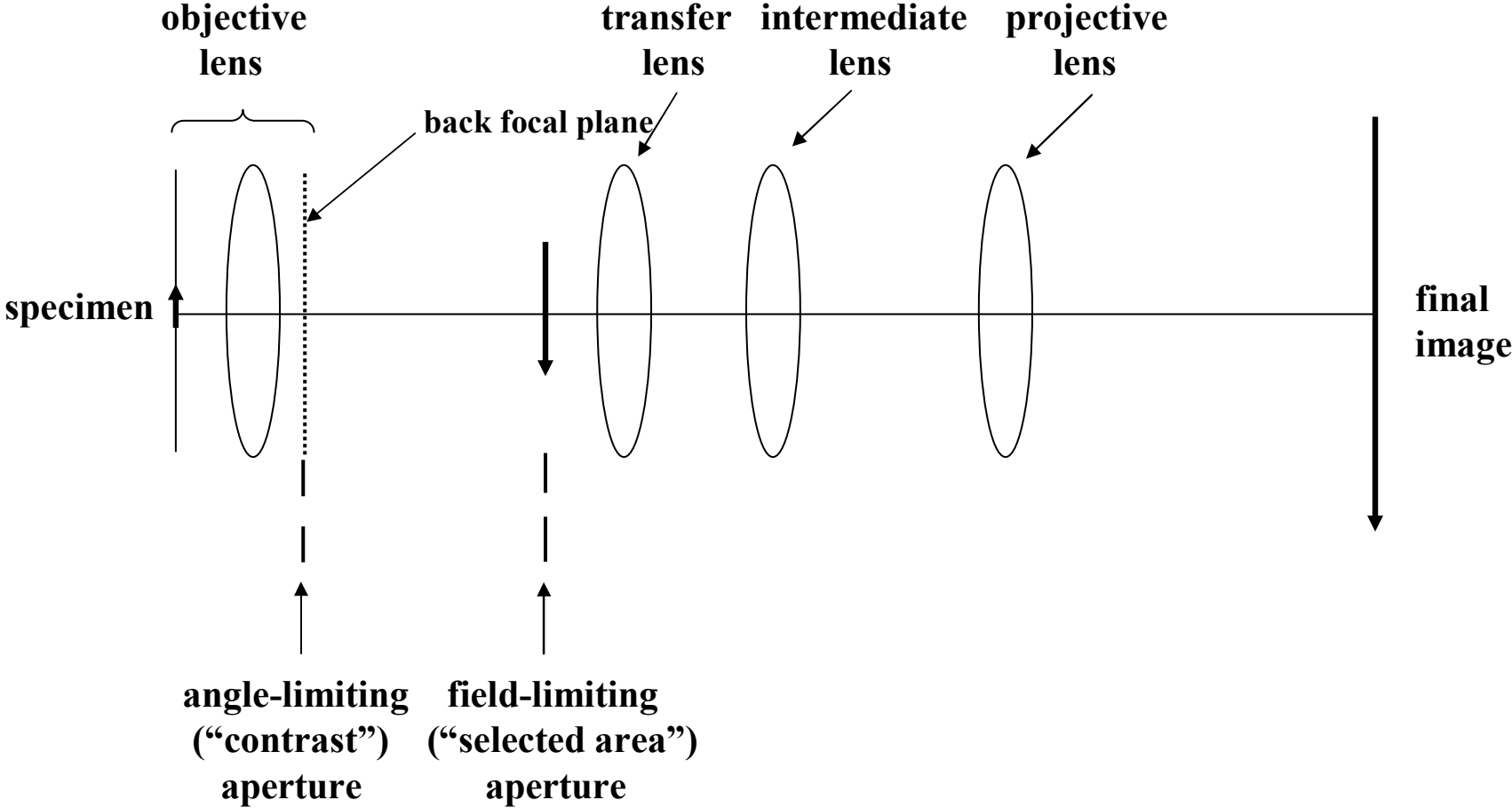


- | | | |
|-----------------|------------|------------|
| ΔE (eV) | | |
| 0.1 | □ | ■ |
| 1.0 | ○ | ● |
| 5.0 | △ | ▲ |
| | without | with |
| | correction | correction |

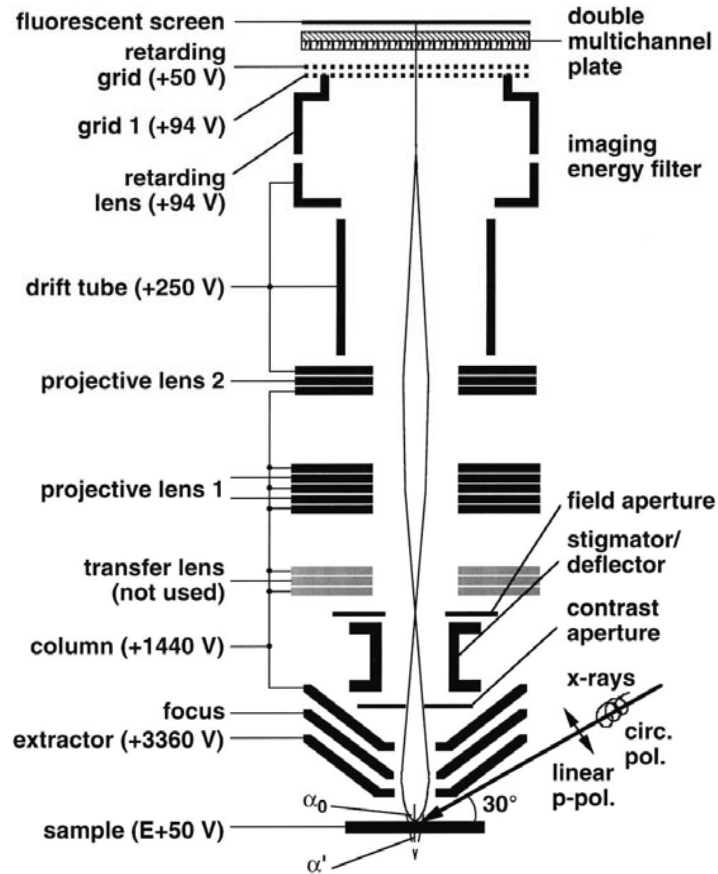
Energy filter needed for secondary electrons

Instruments

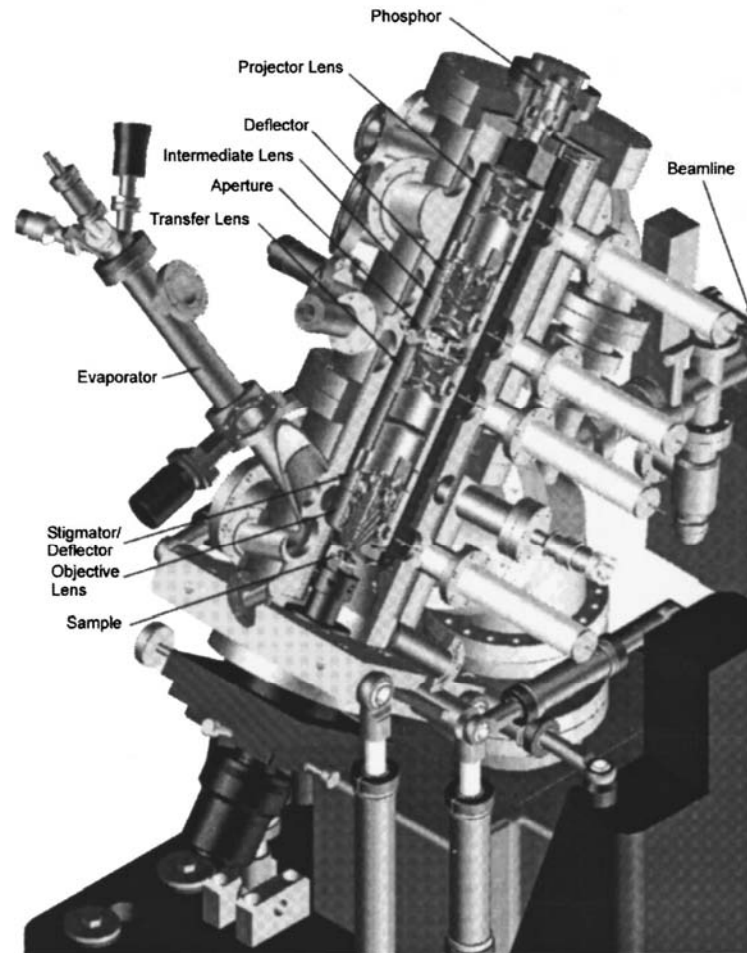
Basic PEEM schematic



Electrostatic PEEM examples

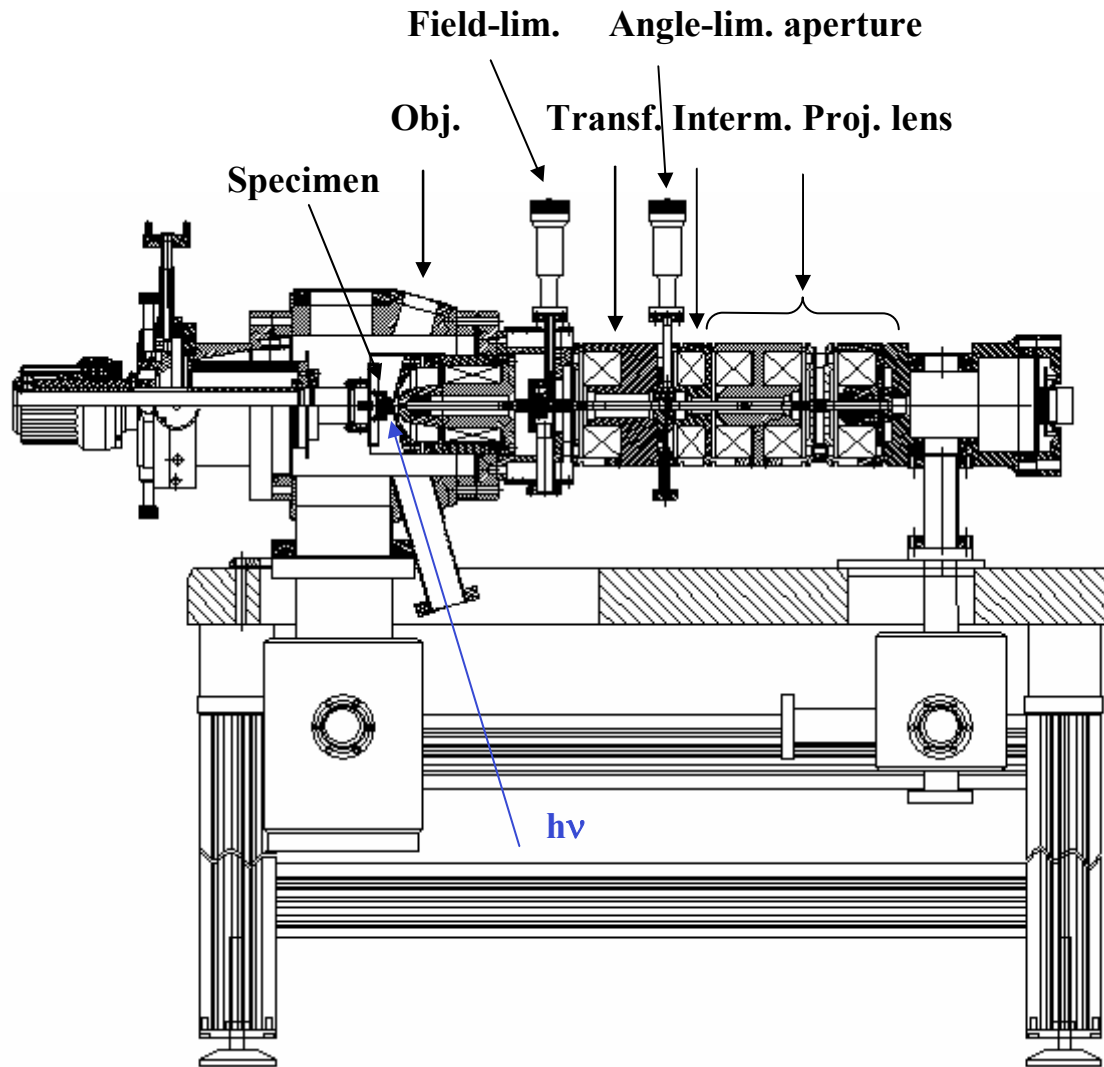


**Focus PEEM
with high pass filter**



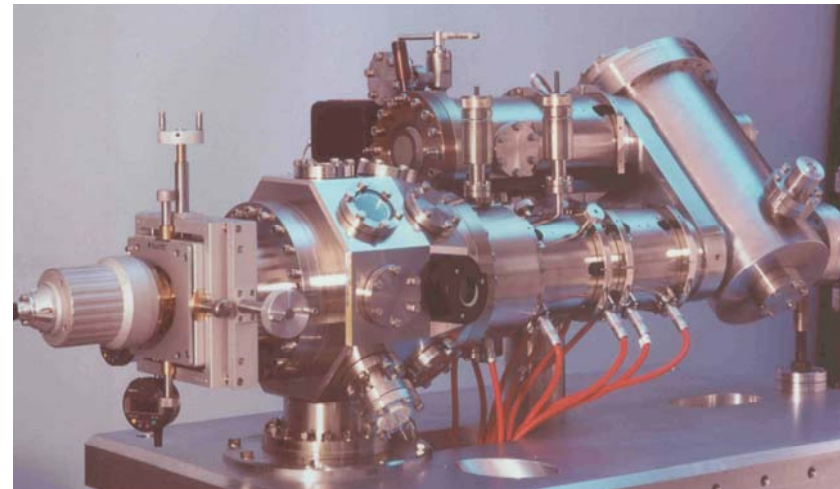
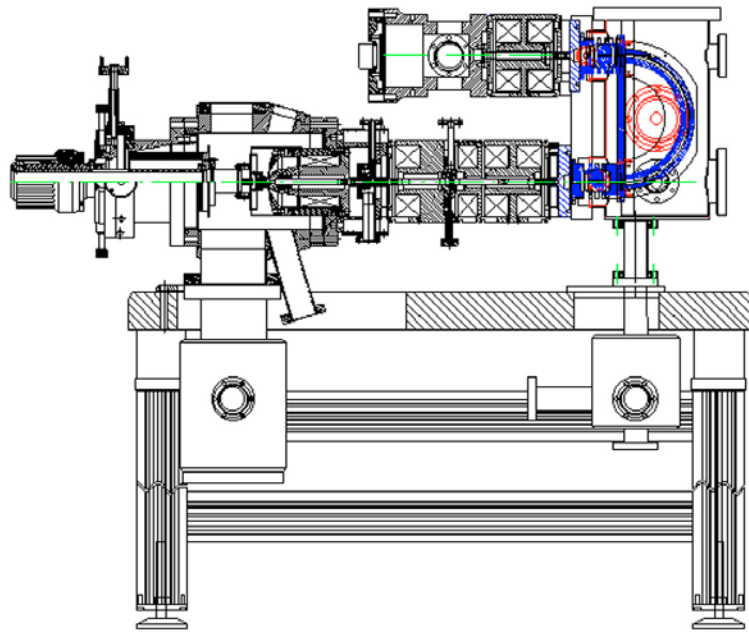
ALS PEEM II

Magnetic PEEM (ELMITEC)



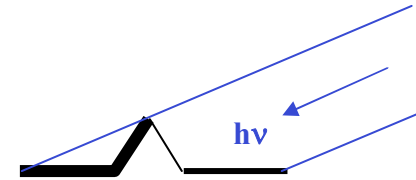
Spectroscopic PEEM with band pass filter

ELMITEC



Contrast mechanisms

- 1 Topographic contrast due to oblique illumination and field distortion
- 2 Work function contrast at low E_0 (escape probability!)
- 3 Chemical contrast due to inner shell ionization
- 4 Magnetic contrast via XMCD and XMLD

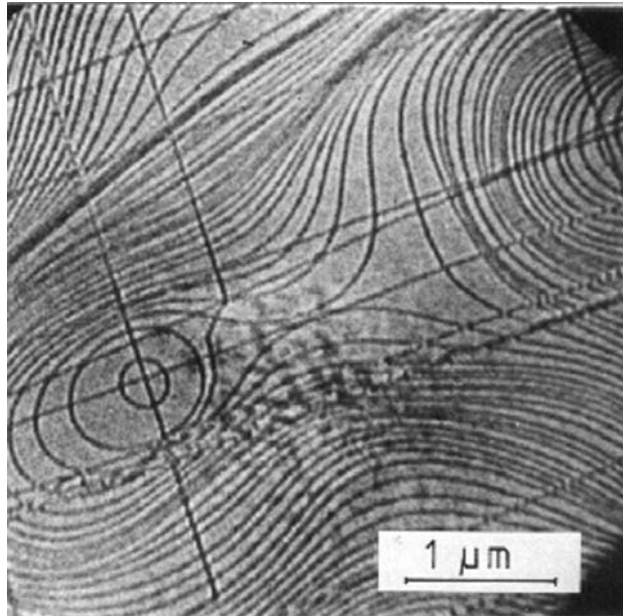


No structural contrast, therefore combination with
Low Energy Electron Microscopy
(LEEM)

The usefulness of LEEM

Properties not visible with PEEM, but with LEEM

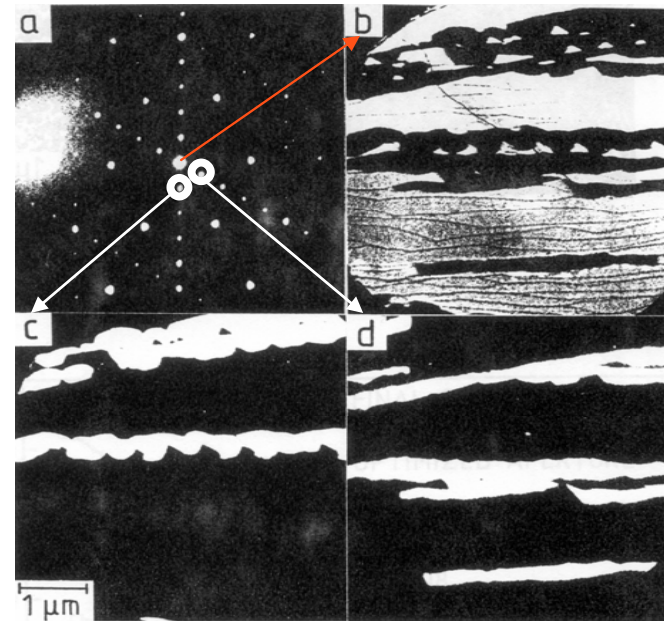
atomic steps



Mo(110)

Interference contrast

domain orientations



Au($\sqrt{3}\times\sqrt{3}$)-R30° + Au(5 × 2) on Si(111)

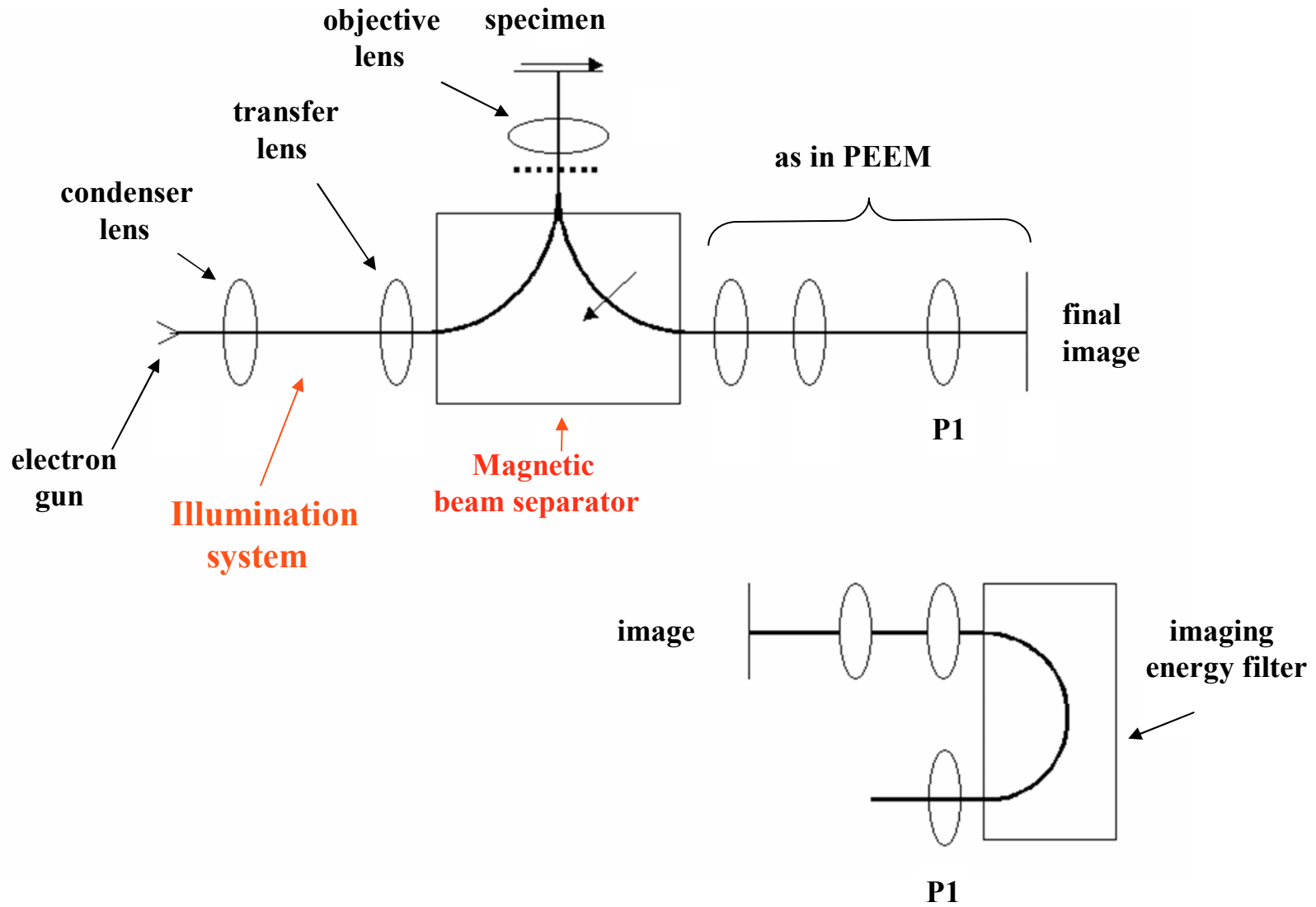
b **c,d**

Diffraction contrast

LEEM also much brighter and better resolution ⇒ use for focusing in XPEEM

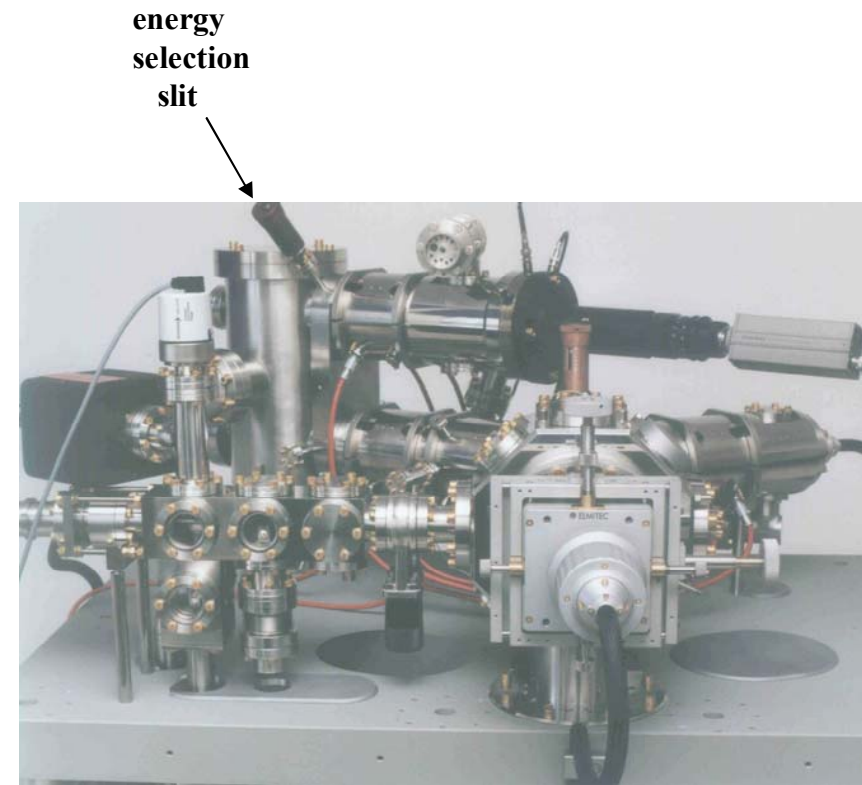
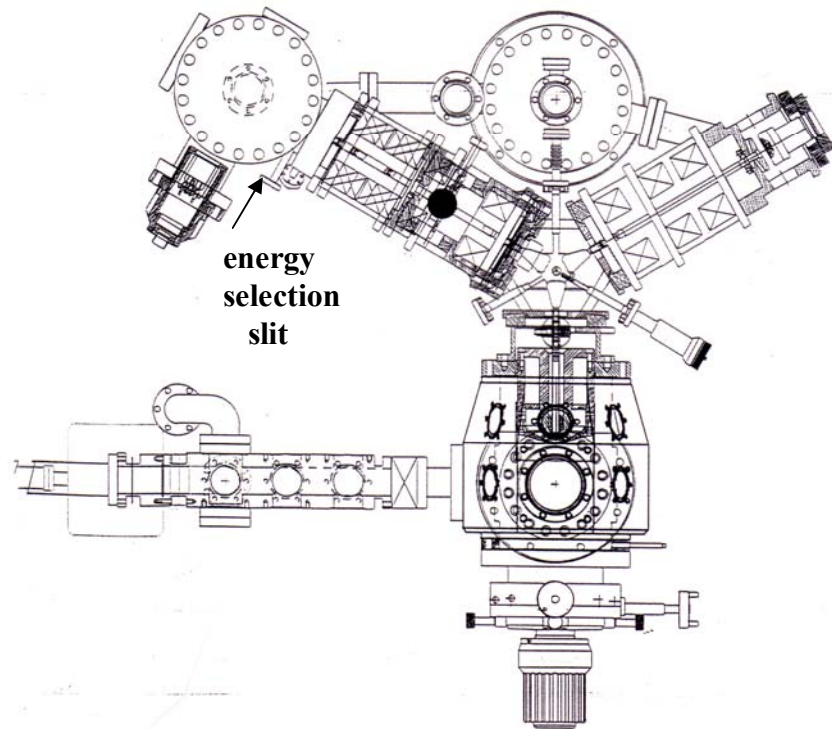
LEED much easier to interpret than PED ⇒ use for structure analysis

Basic LEEM schematic

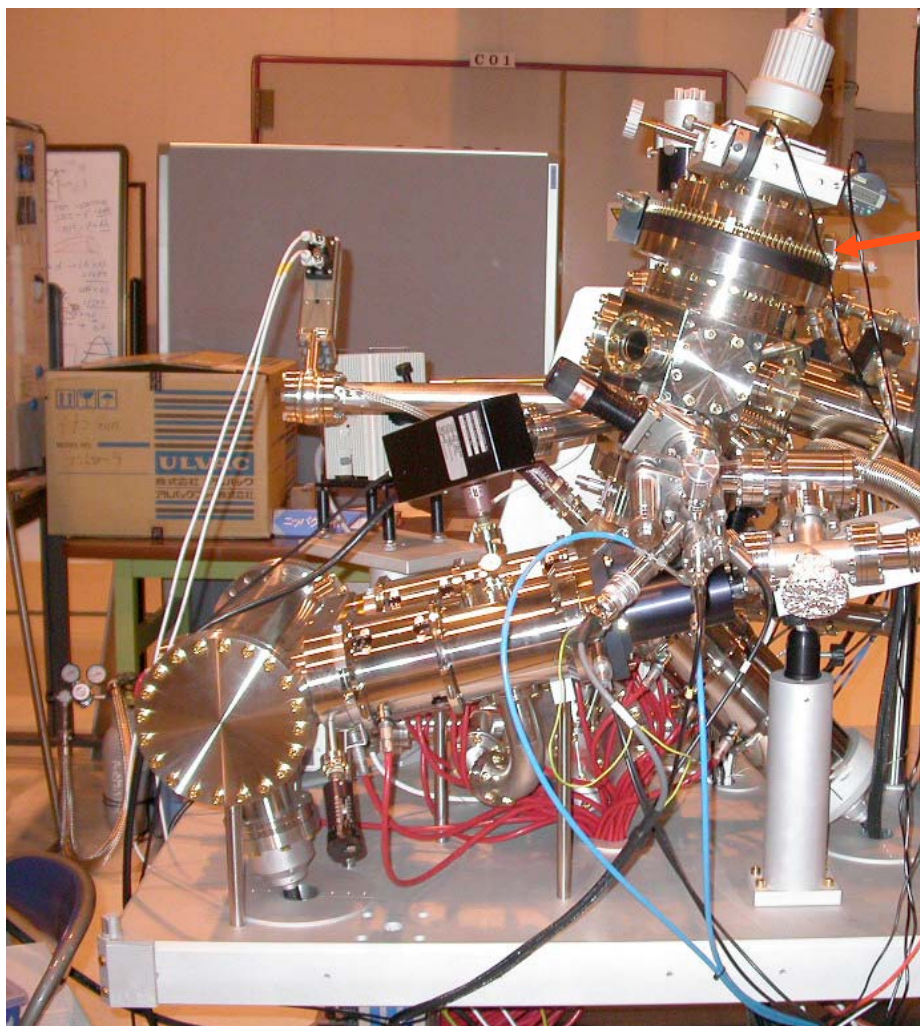


Spectroscopic Photo Emission and Low Energy Electron Microscope

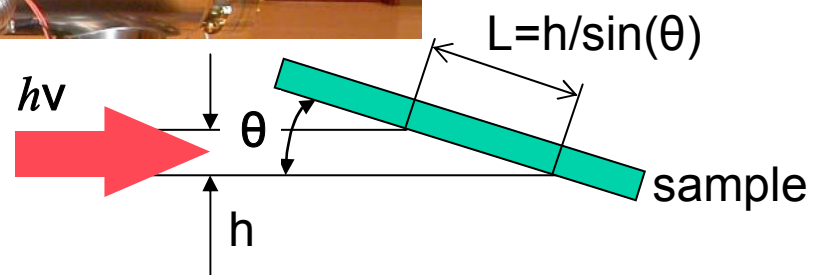
SPELEEM ELMITEC



SPELEEM side view



Differentially Pumped Rotary Platform



$$L \approx 3.6h \quad (\theta = 16^\circ)$$

h : vertical size of $h\nu$

θ : inclination angle

L : effective irradiation width

SPring8

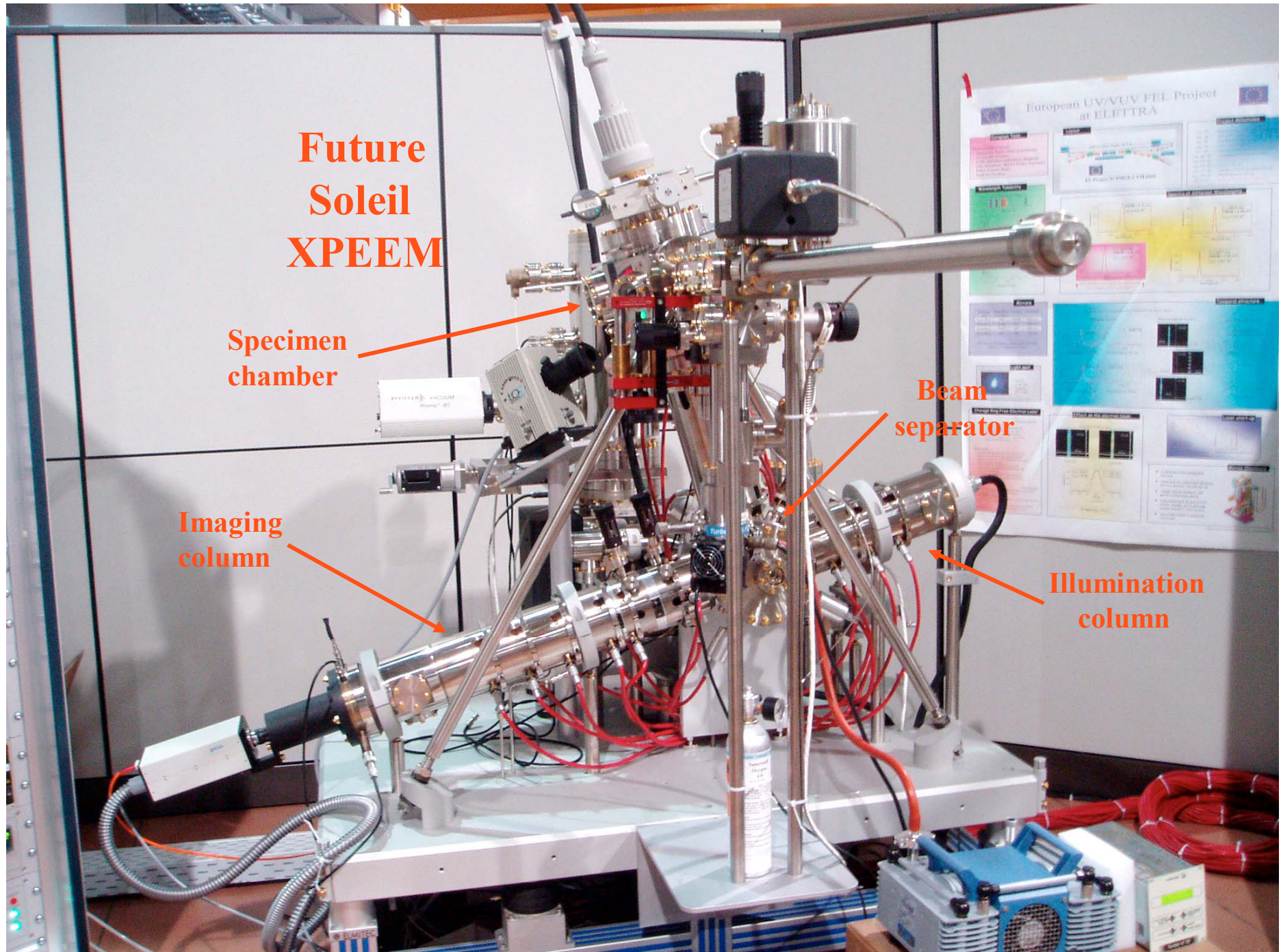
Future Soleil XPEEM

Specimen
chamber

Imaging
column

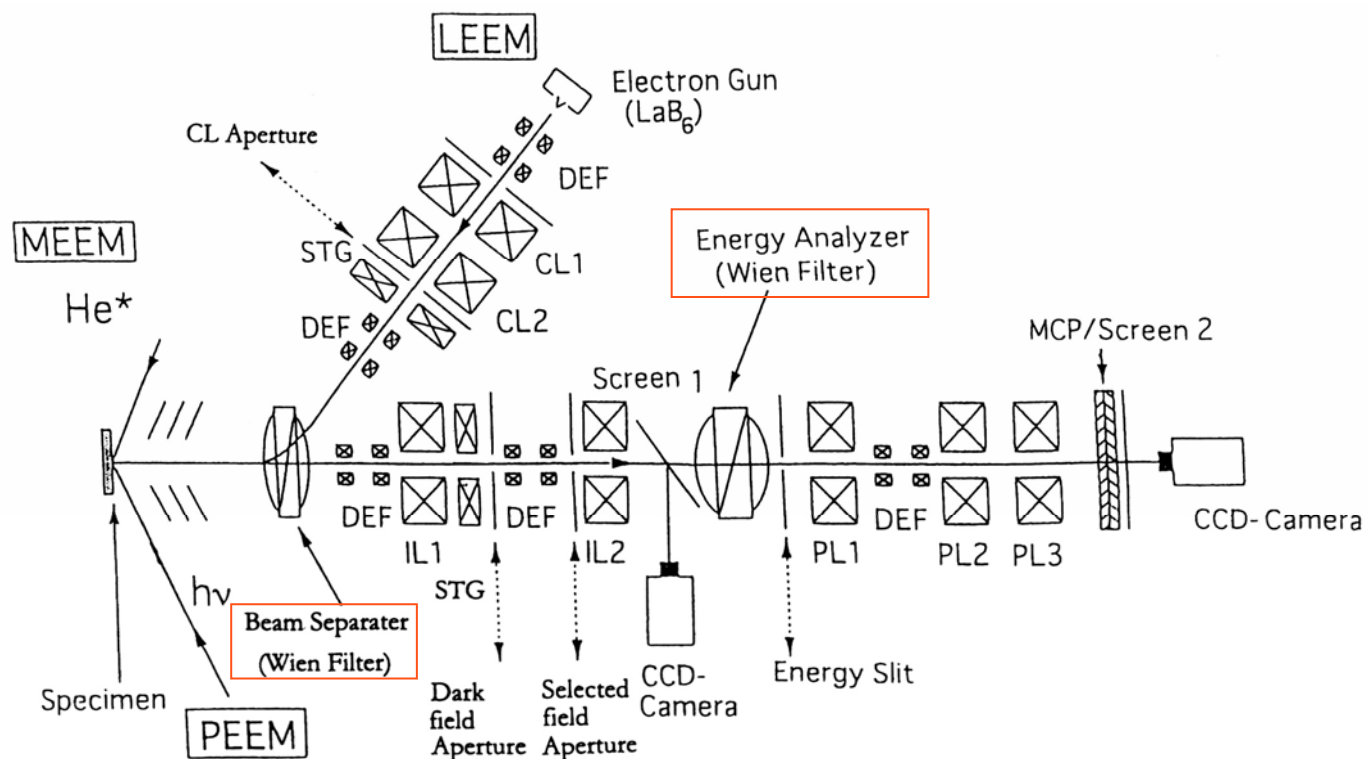
Beam
separator

Illumination
column



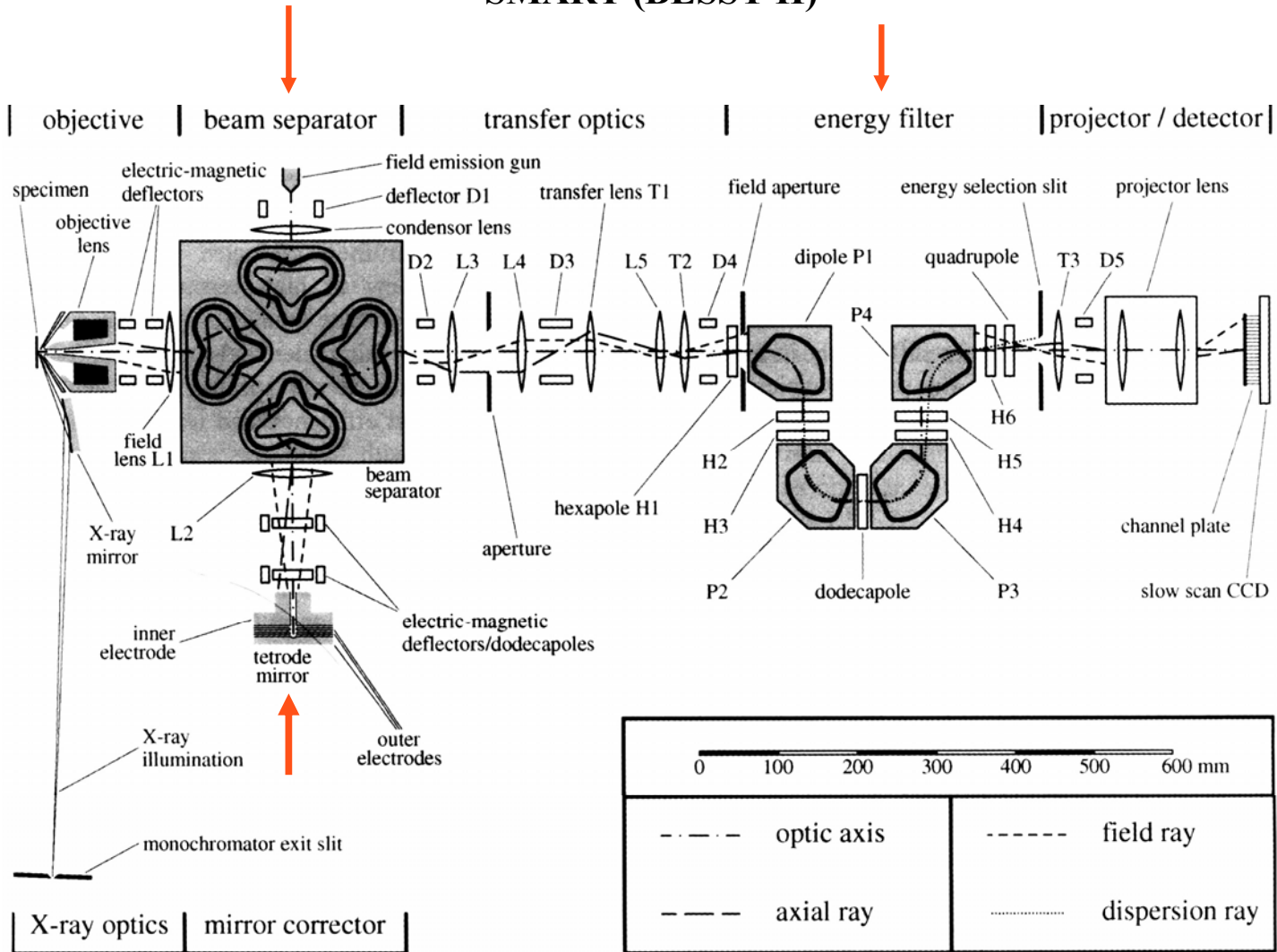
LEEM with energy filter

JEOL



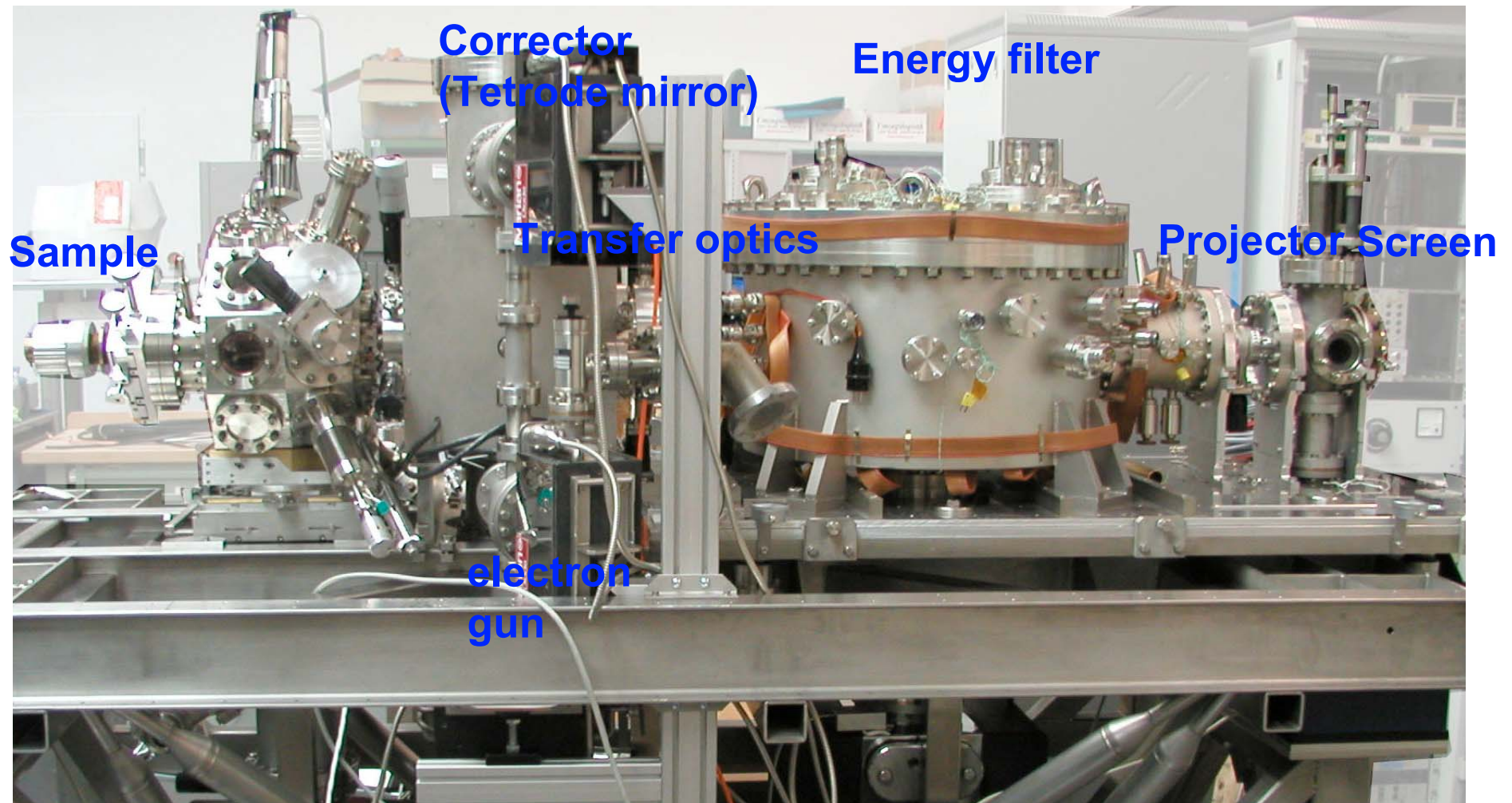
Aberration-corrected SPELEEM

SMART (BESSY II)



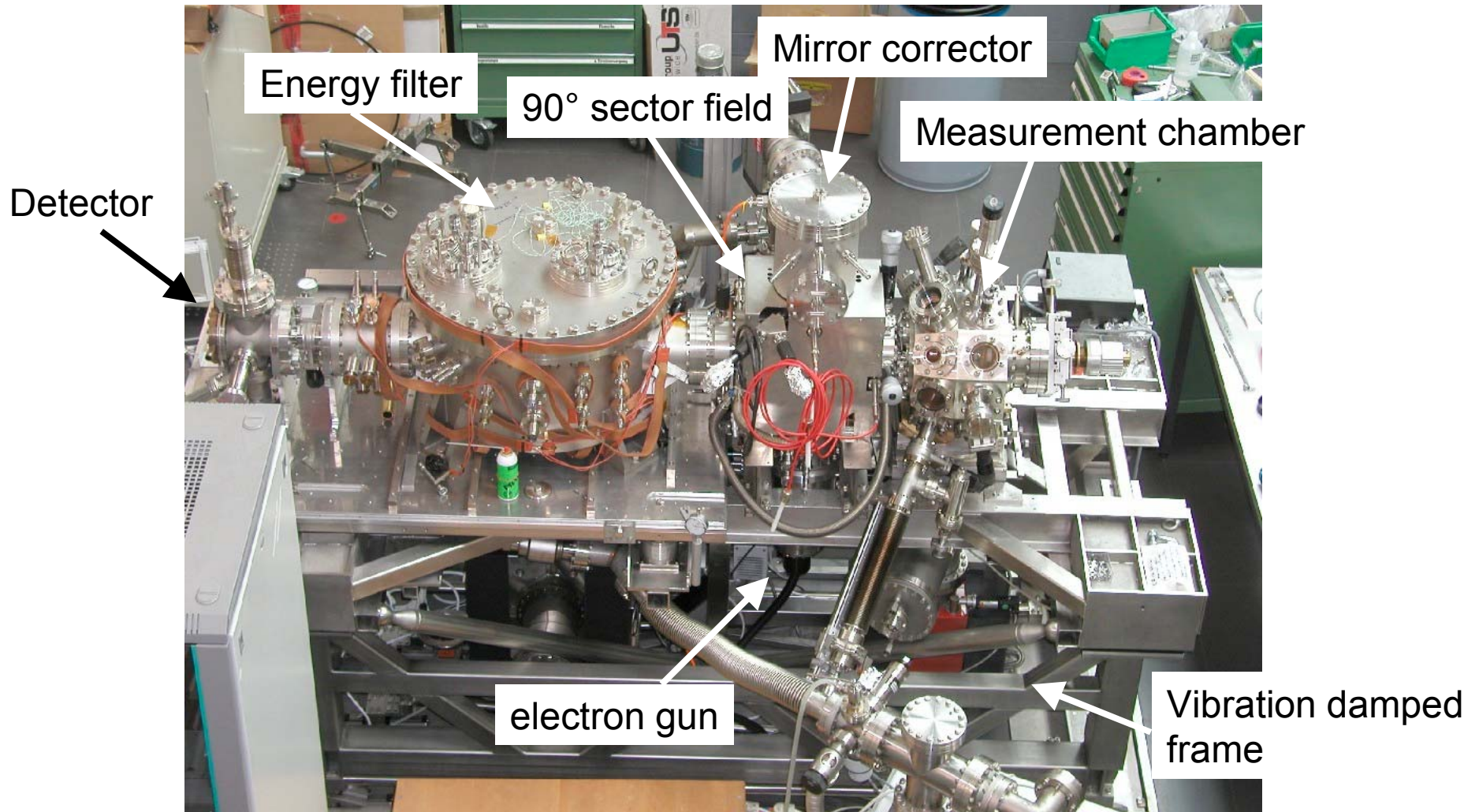
SMART side view

Aberration corrected PEEM/LEEM with energy filtering



Th. Schmidt April 2004

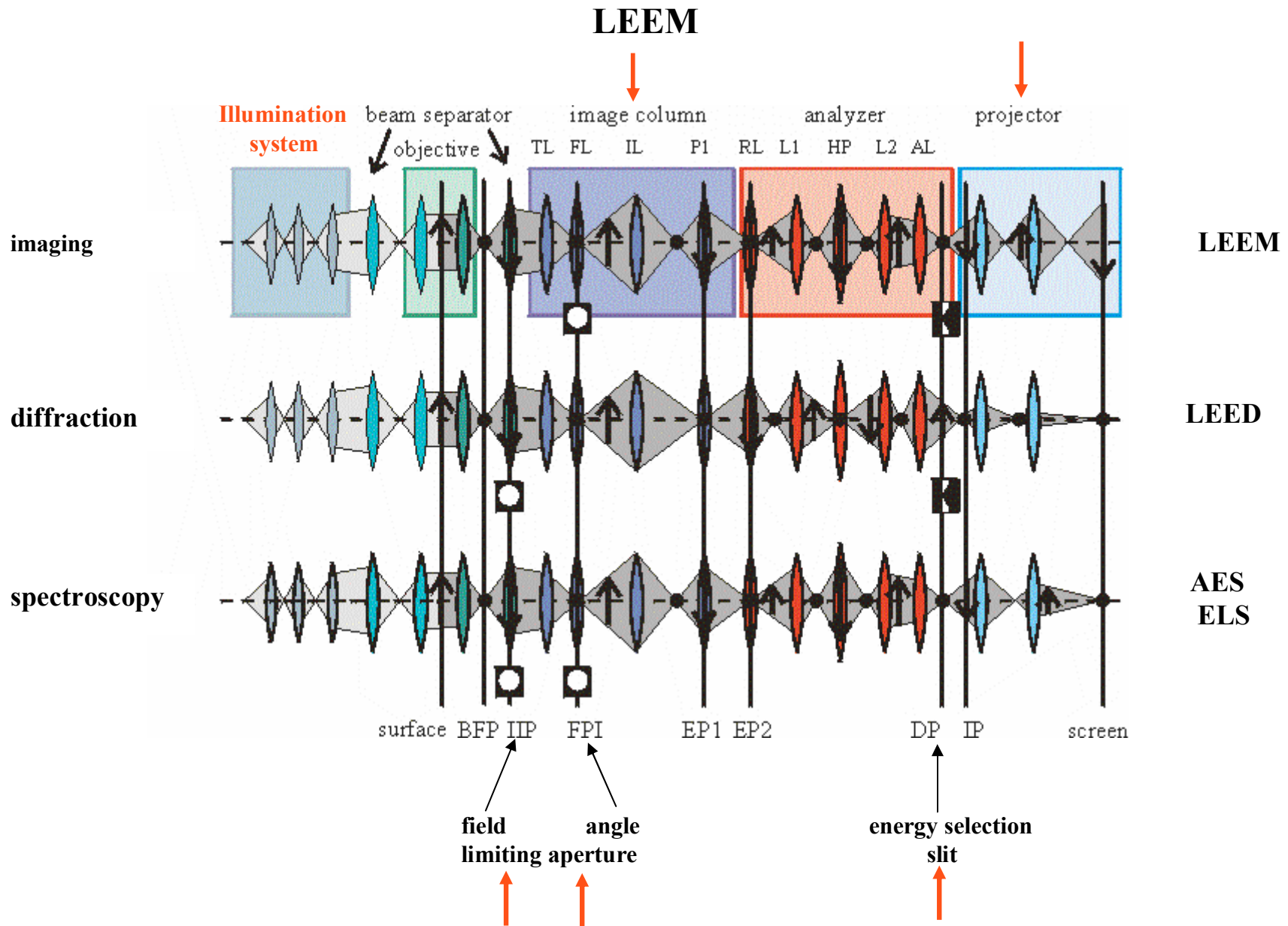
SMART top view



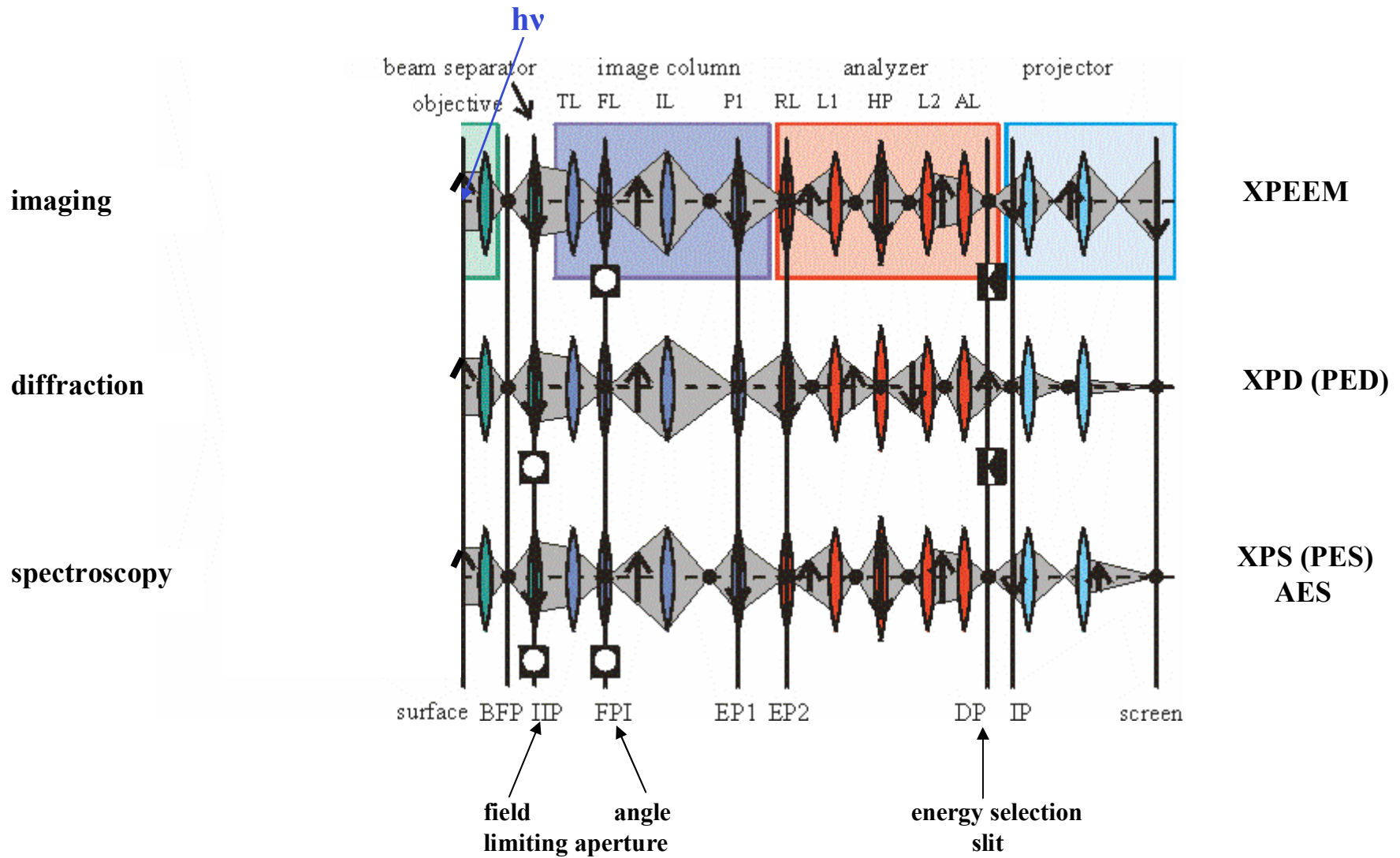
Th. Schmidt April 2004

Methodic

Operation modes of a SPELEEM



Operation modes of a SPELEEM with photons



PEEM practice

Ultrahigh vacuum (low 10^{-10} torr range)
but experiments up to 10^{-5} torr range possible

Surface cleaning: heating, sputtering or chemical reactions, e.g with oxygen for carbon removal

Choice of optimum photon energy:

Secondary electron imaging: $h\nu \approx E_i$ in:
XANES, NEXAFS, XMCD, XMLD

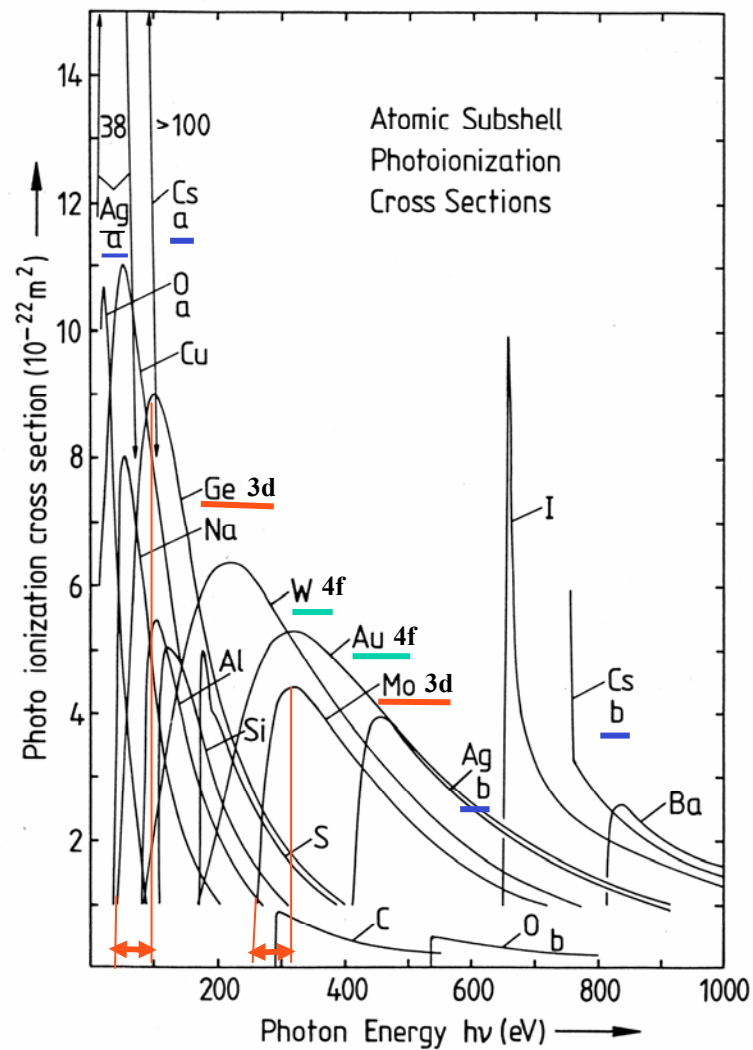
Photo electron imaging: $E_i + 50 \text{ eV} < h\nu < E_i + 200 \text{ eV}$

to minimize
secondary electron
background

to maximize
transmission
and
photo electron yield

Photo ionization cross sections

Photon energy selection



Binding energies (eV)

Ge 3d 29.8, 29.2

Mo 3d 231.1, 227.9

W 4f 33.6, 31.4

Au 4f 87.6, 84.0

Ag

a 4d ≈ 5

b 3d 374.0, 368.3

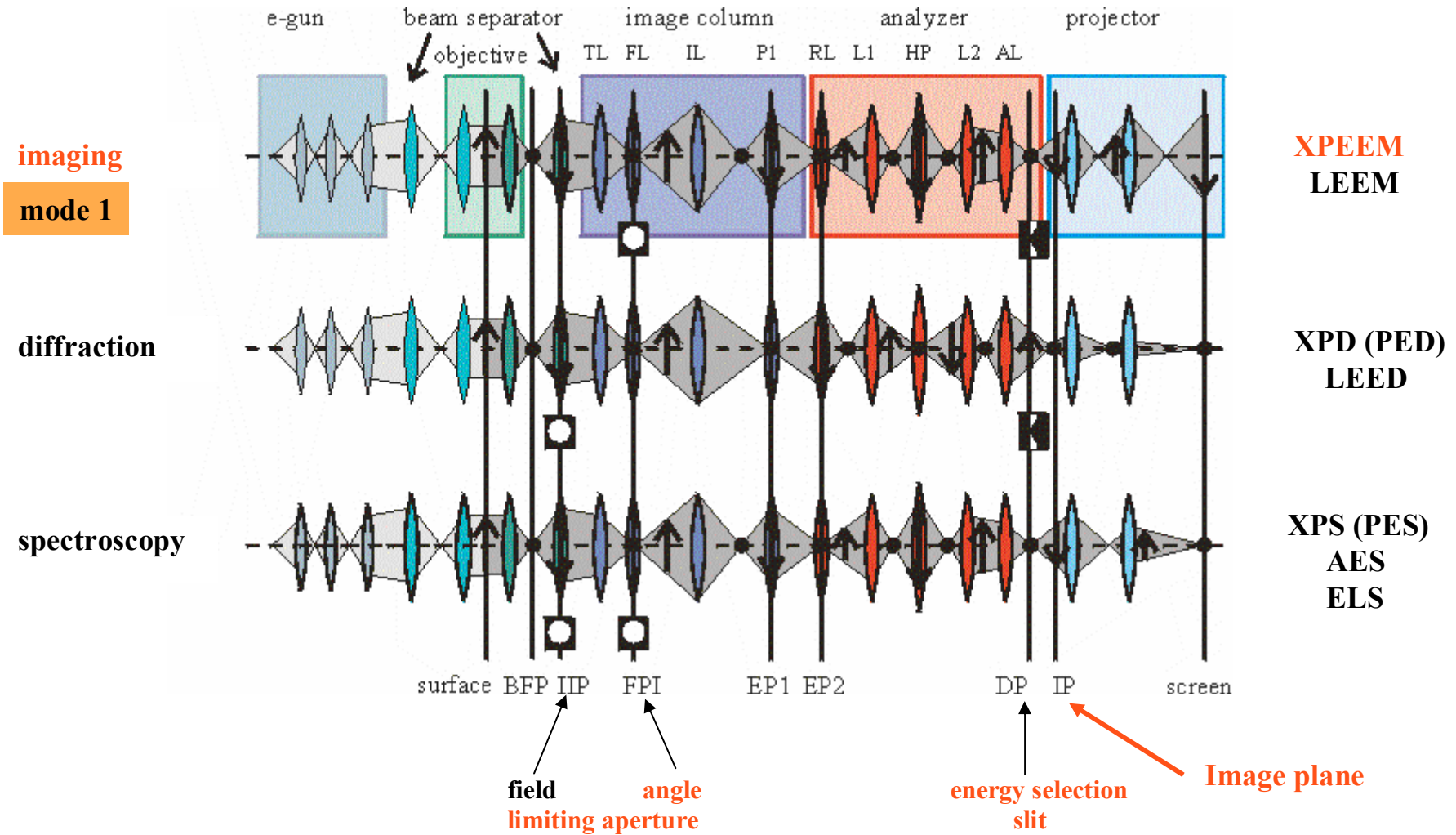
Cs

a 4d 79.8, 77.5

b 3d 740.5, 726.6

J.J. Yeh and I. Lindau,
Atomic Data 1985

Operation modes of a SPELEEM



Chemical imaging (mode 1)

secondary electrons
spatial resolution

$$\sigma_{\text{Ag}4d} \approx 5 \sigma_{\text{W}5d}$$

$$h\nu = 65 \text{ eV} \gg E_i < 10 \text{ eV}$$

$$\Delta E_F \leq 1 \text{ eV}$$

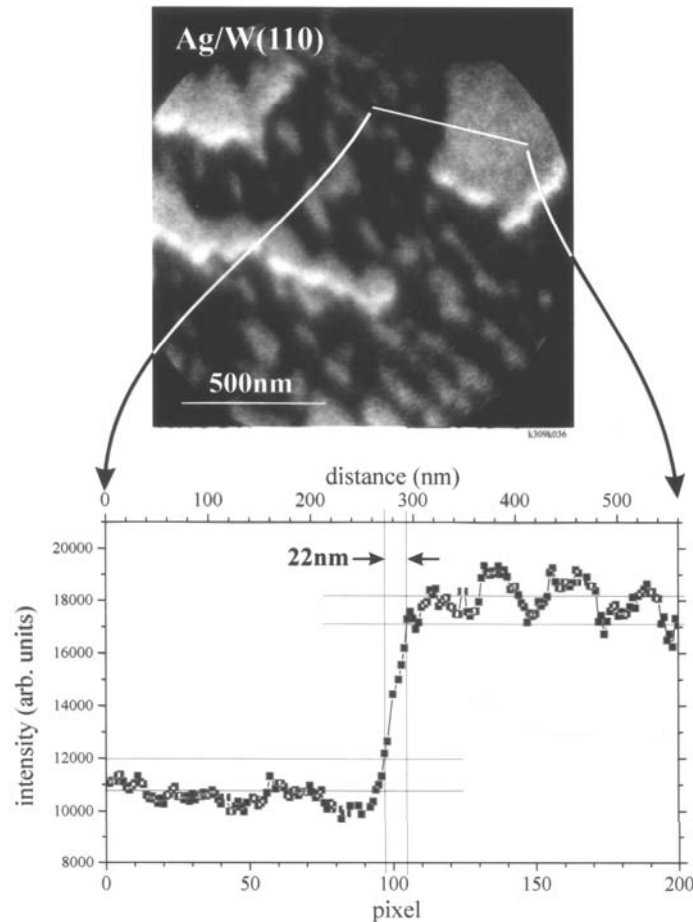
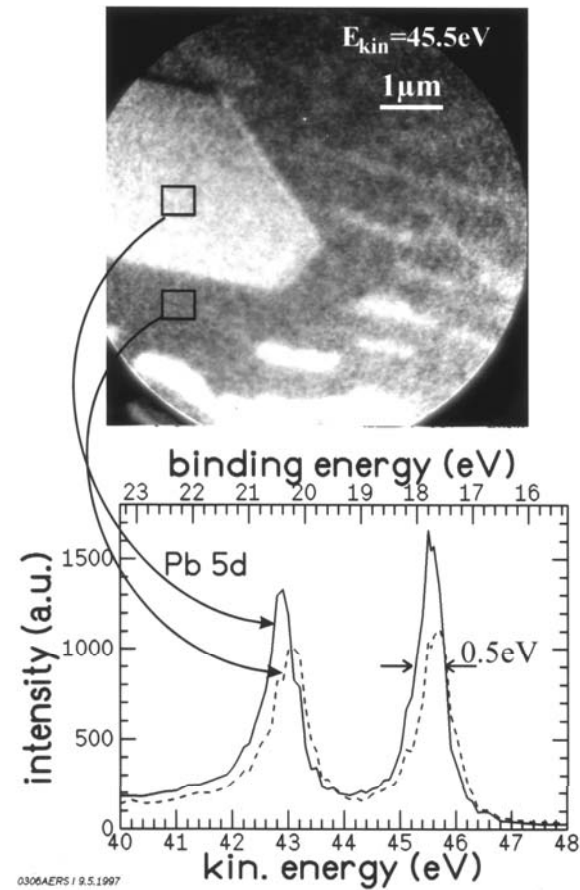


photo electrons
energy resolution

$$h\nu = 65 \text{ eV, images in } 0.2 \text{ eV steps}$$

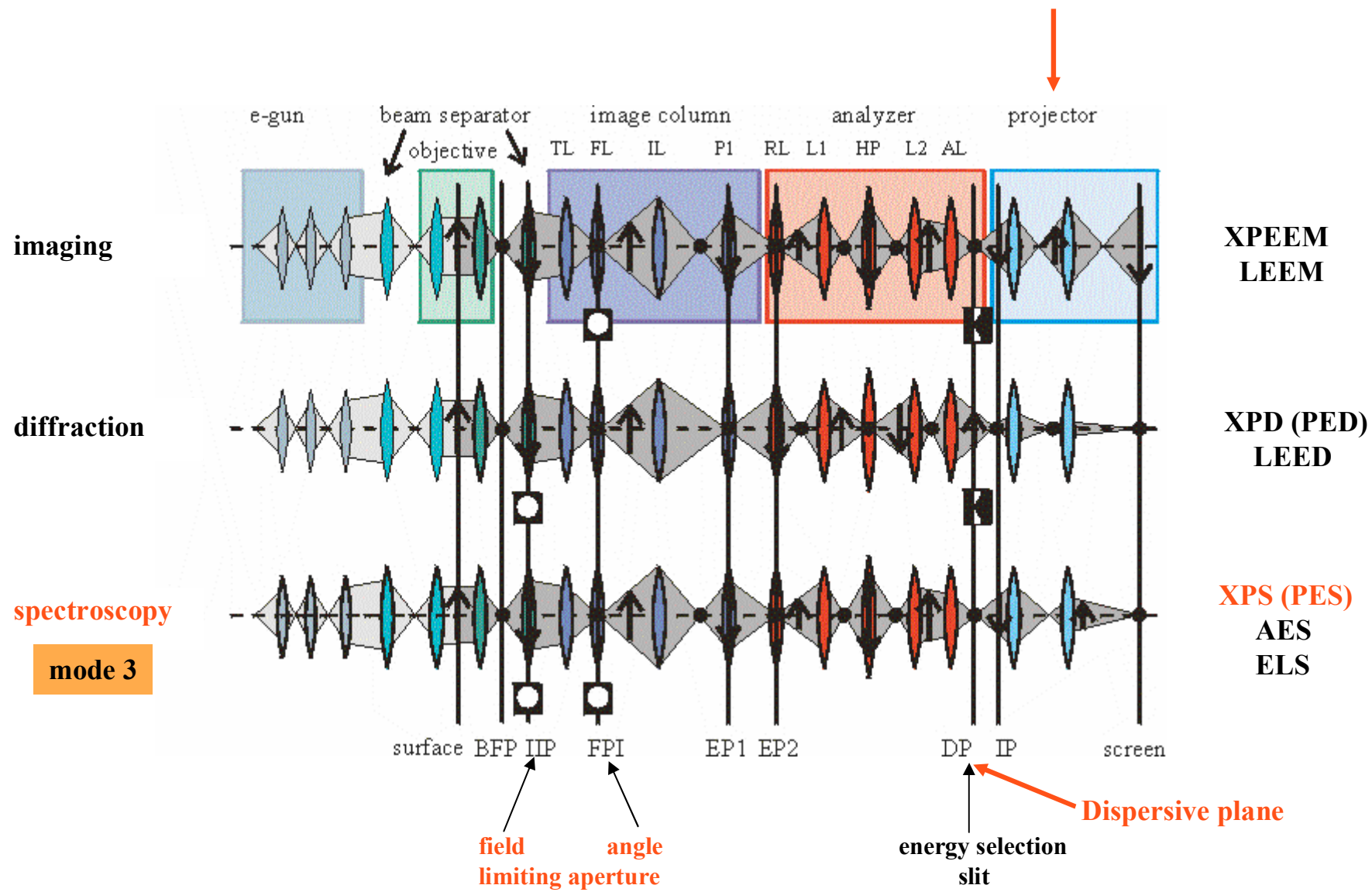
$$10\text{-}60 \text{ sec/image } 0.25 \mu\text{m}^2 \text{ areas}$$

$$\Delta E_F \leq 0.5 \text{ eV, } \Delta E_{\text{chem}} \approx 0.15 \text{ eV}$$



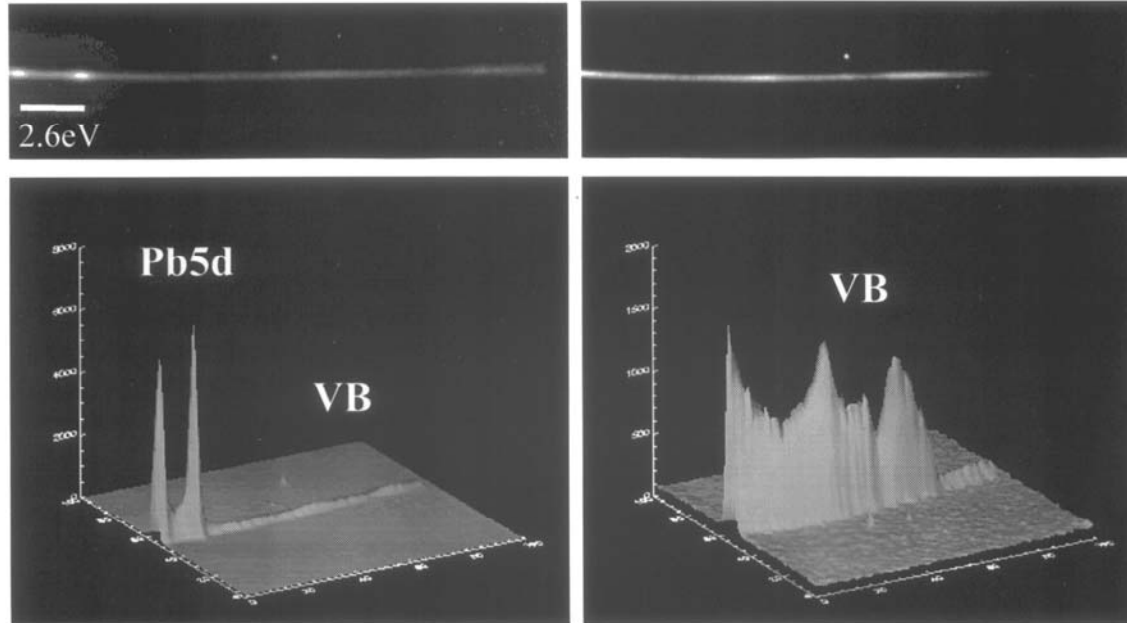
Pb/W(110)

Operation modes of a SPELEEM



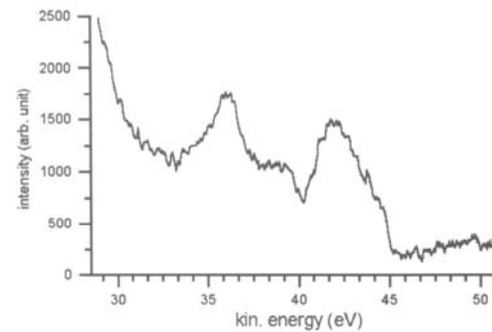
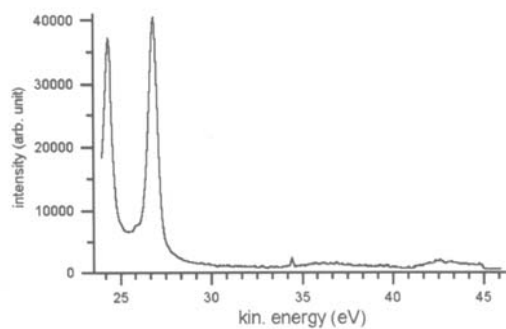
Fast local spectroscopy by imaging the dispersive plane (mode 3)

$\alpha = 8^\circ$ (contrast aperture), $0.8\mu\text{m}^2$ area (selected field aperture)
20 eV full dispersion, 60 sec
 $h\nu = 48\text{ eV}$



Dispersive plane

8 monolayers
Pb on Si(111)-
 $\sqrt{3}\times\sqrt{3}$ -Ag

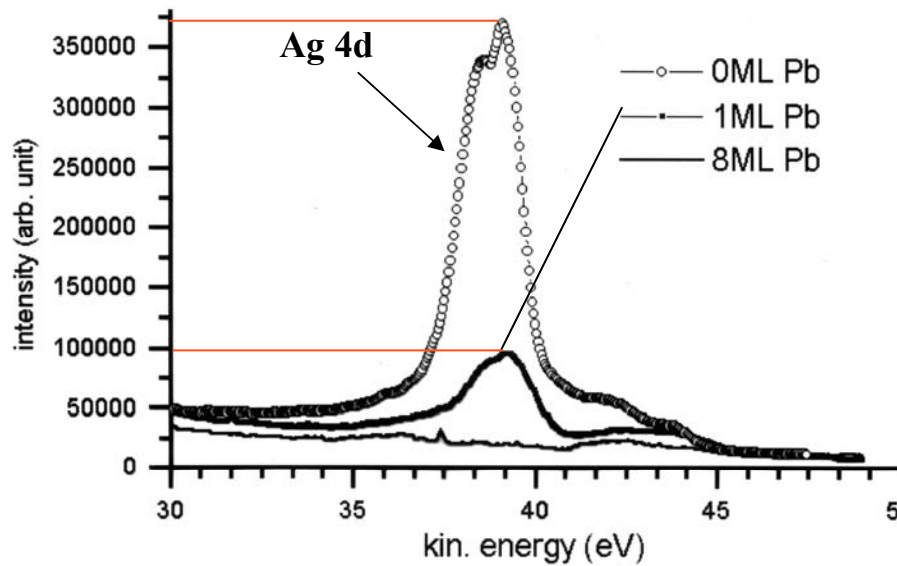


Surface sensitivity of photo electrons

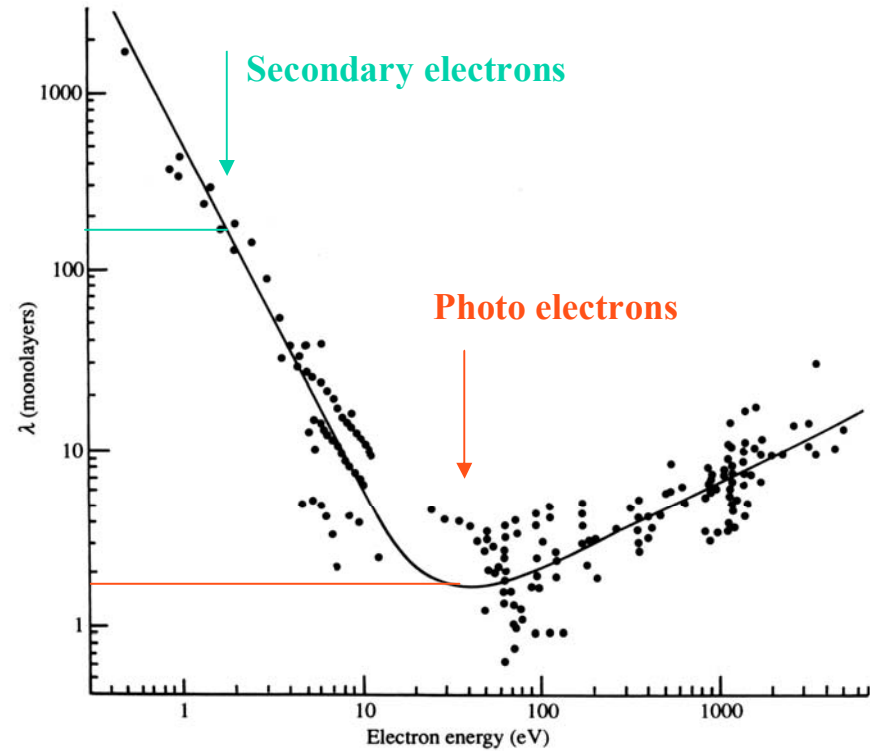
versus secondary electrons

valence band region

$h\nu = 48 \text{ eV}$

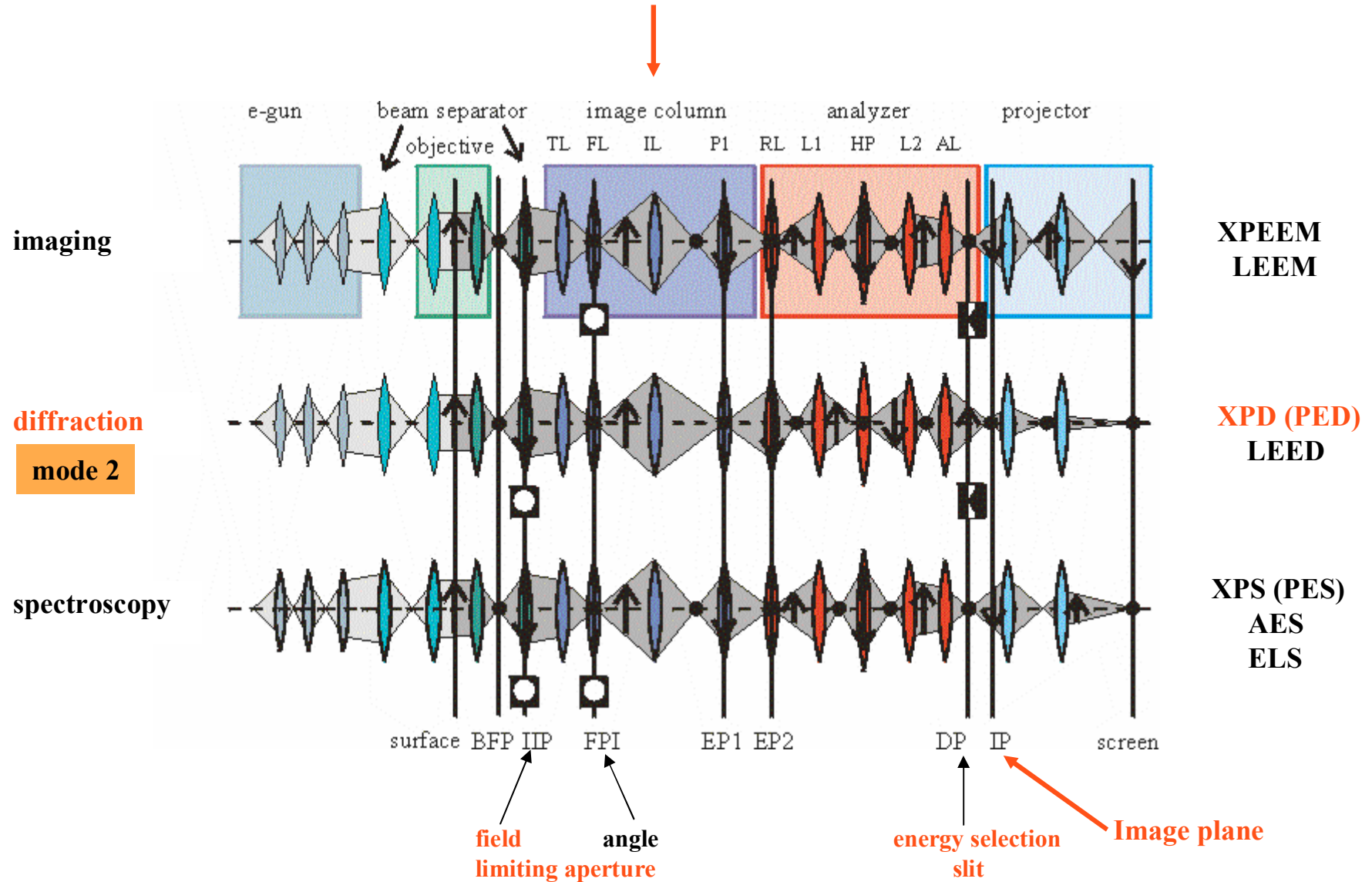


Pb on Si(111) – Ag ($\sqrt{3} \times \sqrt{3}$) – R30°
(1 monolayer Ag)



Inelastic mean free path
("universal curve")
determines sampling depth

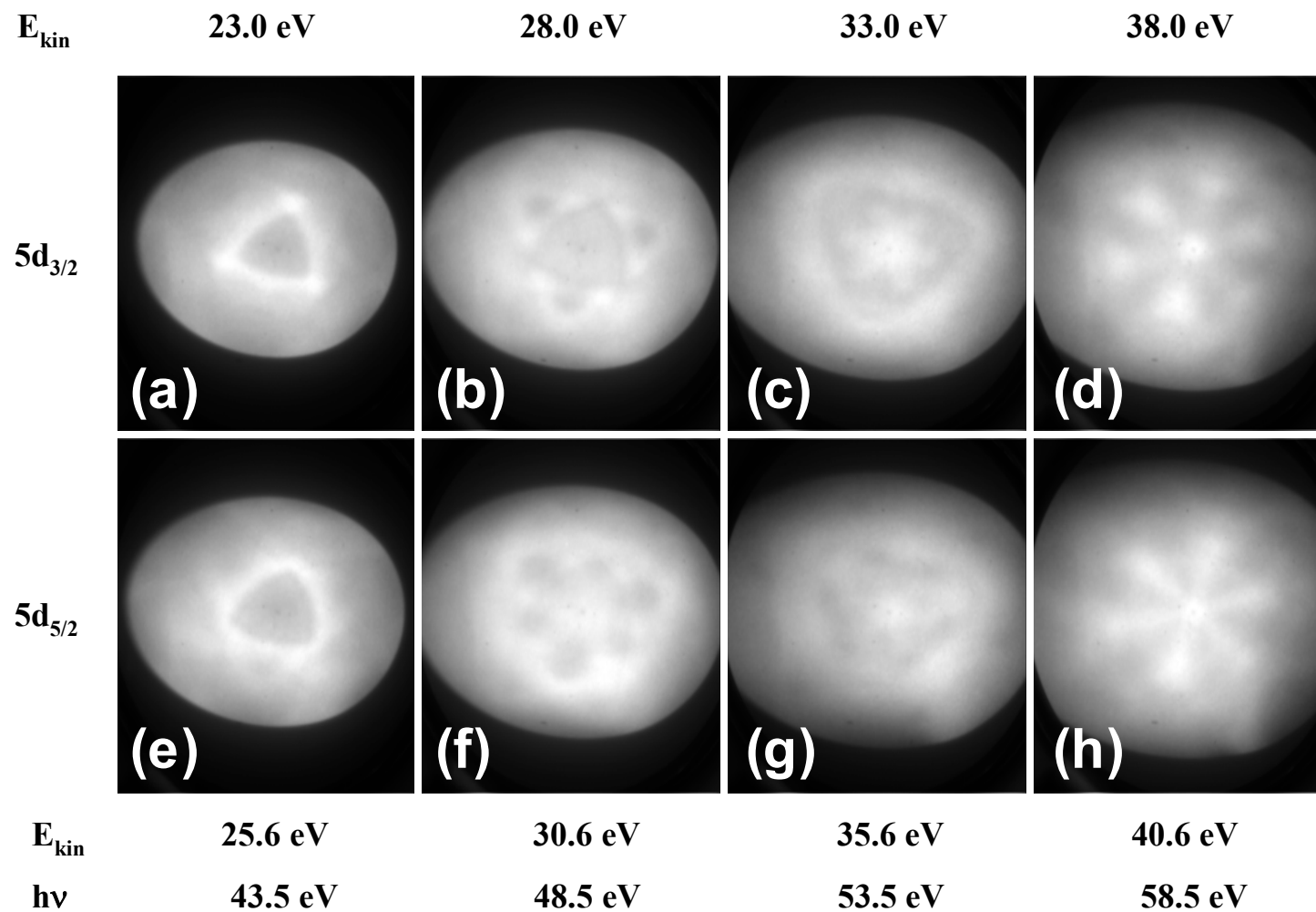
Operation modes of a SPELEEM



Local photo electron diffraction (mode 2)

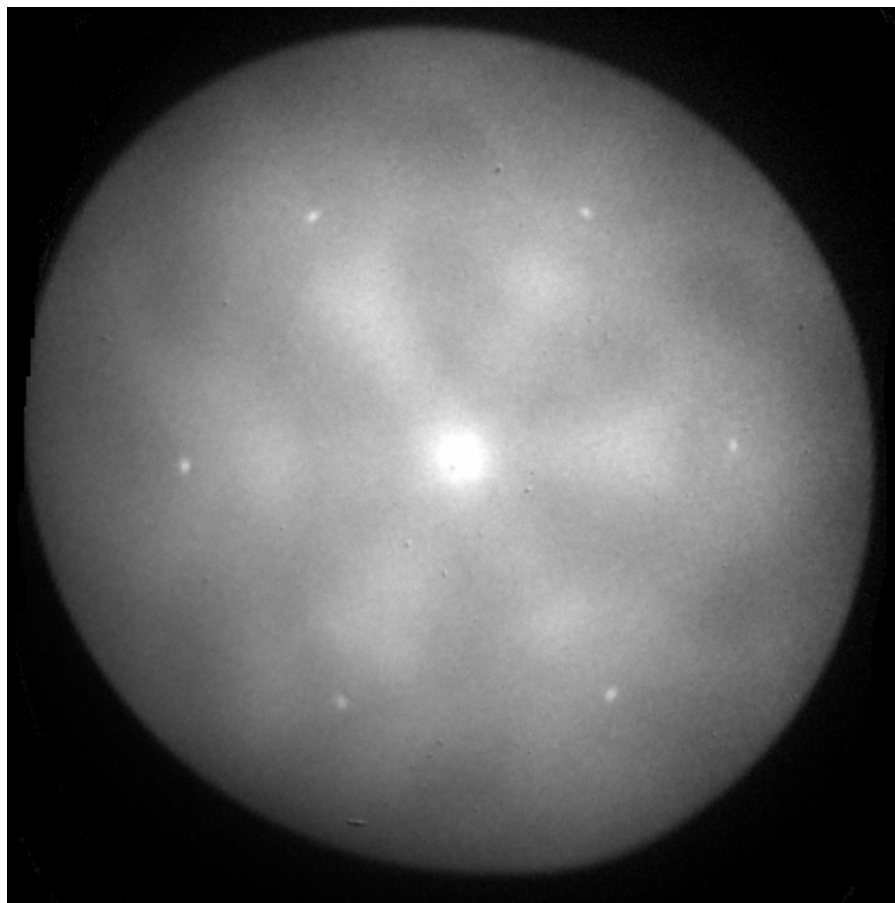
Pb 5d photo electrons

from $0.8\mu\text{m}^2$ area (selected field aperture)



Simultaneously acquired PED and LEED pattern

Pb 5d 38 eV

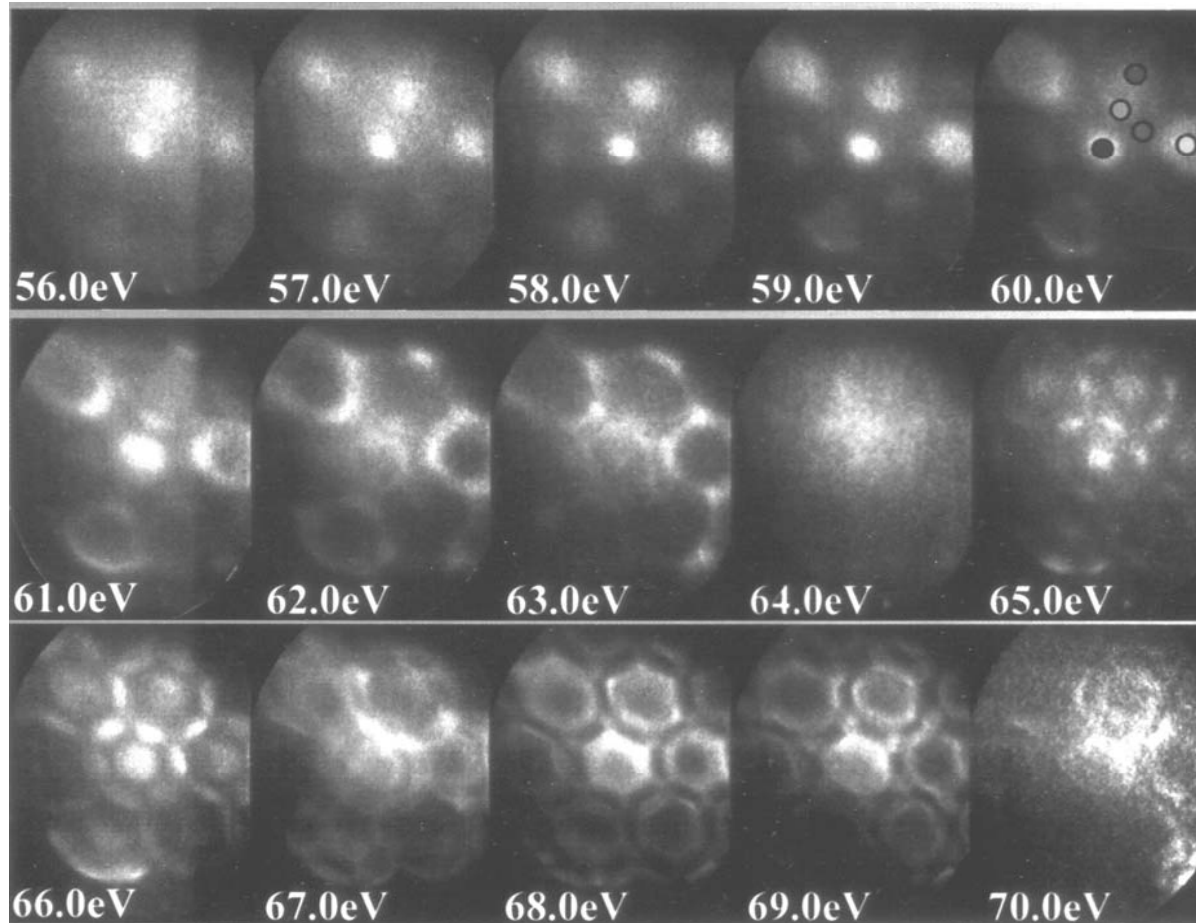


Local band structure analysis (mode 2)

Conduction band of Pb(111)

5 Pb monolayers on Si(111) – Au $\sqrt{3}\times\sqrt{3}$ – R30°

$h\nu = 73$ eV, $0.8 \mu\text{m}^2$ area (selected field aperture)



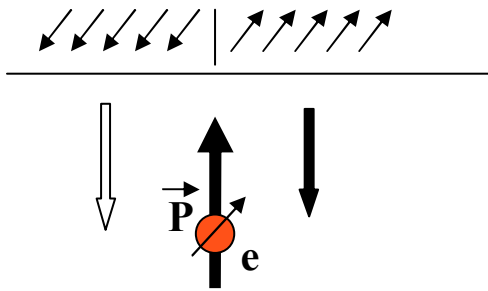
Parameter: E_{kin}

Magnetic imaging

XMCD, XMLD

Methods

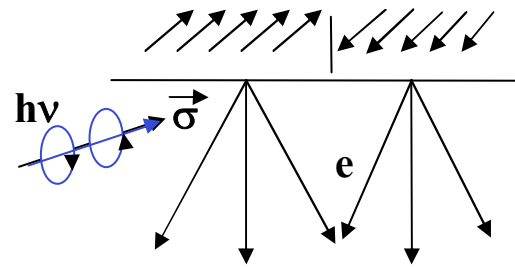
SPLEEM



$$E_i = E_r = 0 - 20 \text{ eV}$$

$$I = I_{0r} + c \vec{P} \cdot \vec{M}$$

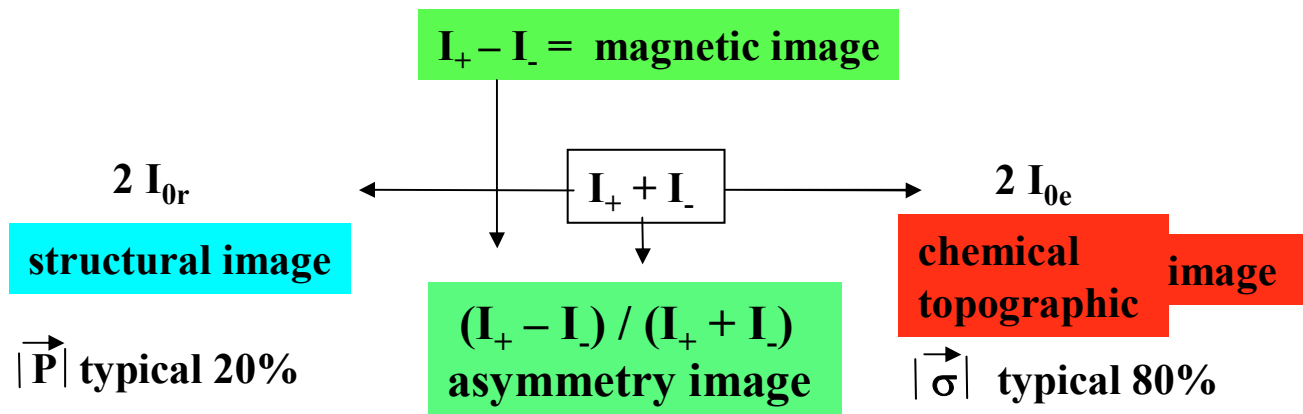
XMCDPEEM



$$E_e = 0 - 20 \text{ eV (SE)}$$

$$I = I_{0e} + c \vec{\sigma} \cdot \vec{M}$$

2 images with opposite \vec{P} or $\vec{\sigma}$: I_+ , I_-



Basic mechanisms

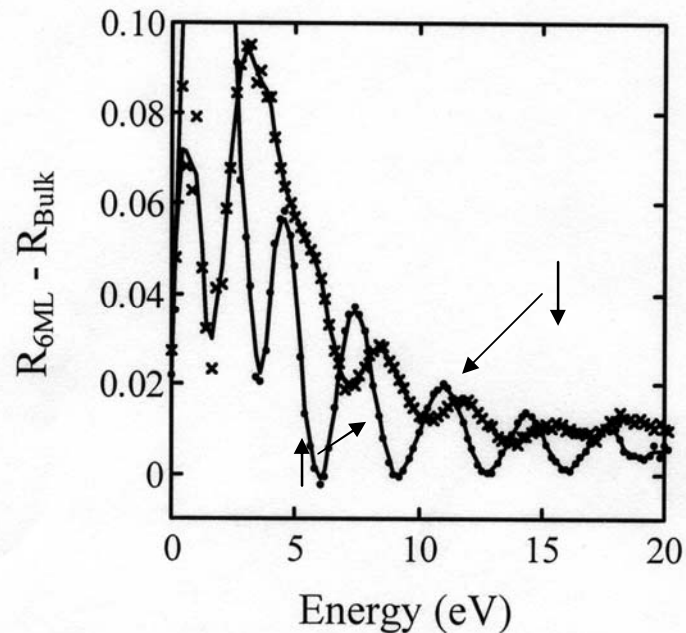
SPLEEM

Spin-dependent scattering cross-section
due to exchange interaction



Spin-dependent reflectivity

6 ML Fe on W(110)



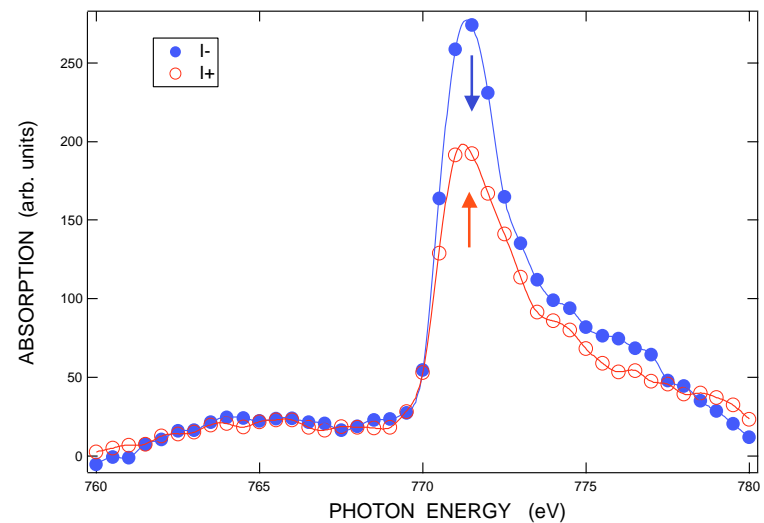
XMCDPEEM

Helicity-dependent transition probability
from 2p to unoccupied 3d ($\uparrow\downarrow$) states



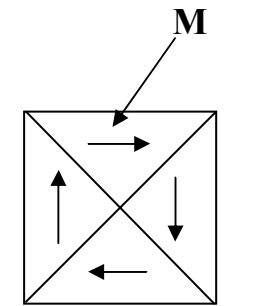
Spin-dependent secondary electron emission

Secondary electron yield around Co $2p_{3/2}$ edge

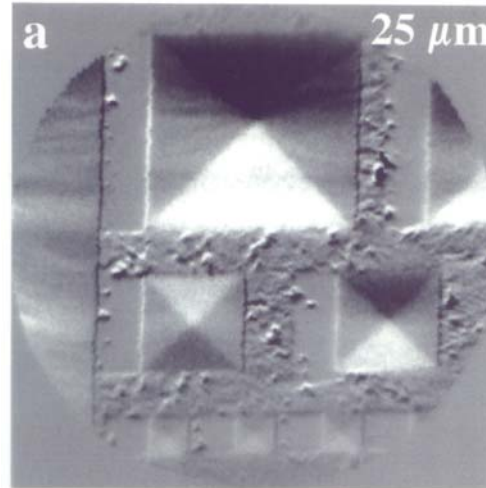
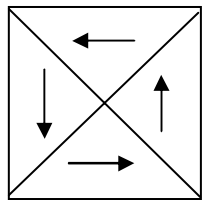


Closure domains in permalloy squares

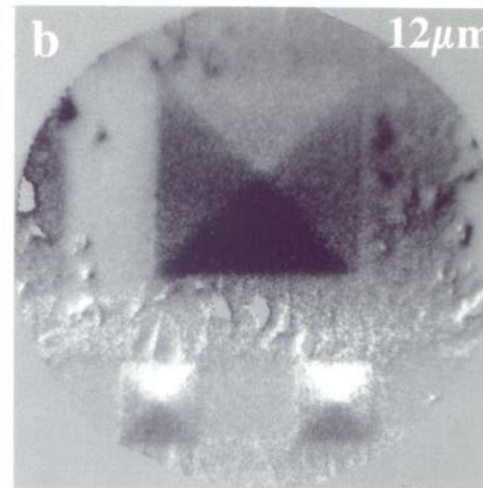
Ion beam milled from permalloy film



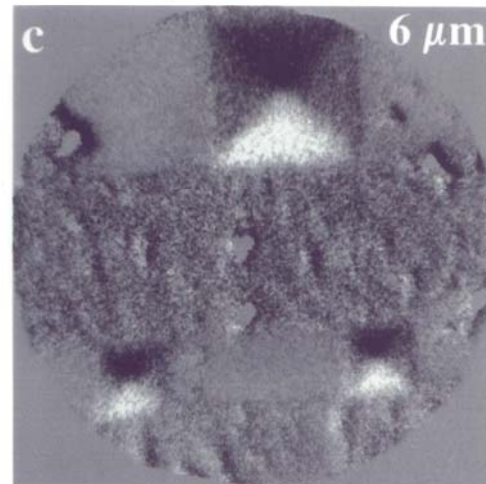
$h\nu$
→ σ
Fe L_3 edge



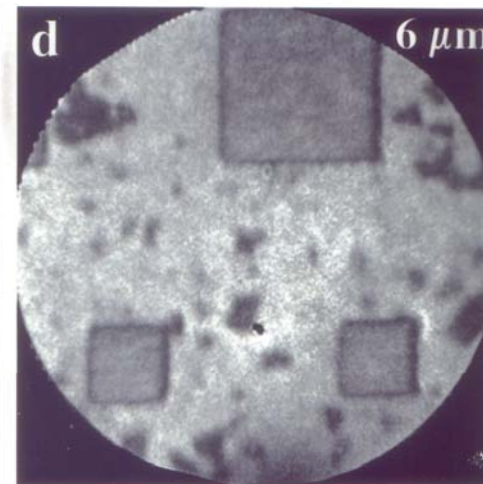
10 μm and 5 μm



5 μm and 2 μm



2 μm and 1 μm



2 μm and 1 μm (LEEM)

Conclusion

Full field XPEEM is one of the most important applications of the high brilliance of third generation synchrotron light, in particular when combined with a **band pass energy filter** and with **LEEM** because it allows a complete characterization of surfaces and thin films, presently on the 10 nm lateral **resolution** scale, with aberration correction in the future on the 1 nm scale (hopefully).

The main benefit of **aberration correction**, however, will be the strong increase of the **transmission** of the system which will reduce image acquisition time considerably. This will allow dynamical studies, which are presently limited to LEEM, also with XPEEM.

General references

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