Surface and thin film structural determinations from combined Photoelectron Diffraction and Grazing Incidence X-ray Diffraction experiments

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e Dipartimento di Fisica Università di Trieste Surface and thin films crystallography.

Experimental techniques: Low Energy Electron Diffraction (LEED) High Energy Electron Diffraction (RHEED) He Atom Scattering (HAS) PhotoElectron Diffraction (PED) Grazing Incidence X-ray Diffraction (GIXD)

Surface sensitivity: short mean free path (LEED, PED, HAS) or grazing incidence (RHEED e GIXD).

## LEED: cheap experimental apparatus but dynamical calculations

GIXD: Kinematic scattering calculations  $\Rightarrow$  large structural parameter space **but** synchrotron radiation, high experimental apparatus complexity.

For thin films: GIXD variable penetration depth study of buried interfaces.

#### In plane X- Ray Diffraction measurements

In plane measurements: grazing incidence, grazing emission.

In plane structural parameters determination.



Rod scan: Grazing incidence Constant Q transfer variable exit angle  $Q_z$  variation Out of plane structural determination





## **ALOISA** beamline



.

## **Experimental apparatus**

# ALOISA

High Resolution Photoemission Spectroscopy X-Ray Photoelectron Diffraction Grazing Incidence X-Ray Diffraction X-Ray Reflectivity Absorption Spectroscopy



Fe magnetic properties depend on its lattice cell



Suitable choice of the substrate allows to grow thin Fe films with different structures

## Fe/Cu(100)

Structure: 1-4 ML pseudomorphic FM fcc, Fe on hollow 4-10 MLtopmost bilayer FM fcc; inner layers AF fcc > 10 ML FM bcc (110), Fe on bridge

**Morphology:** 1-10 MLlayer-by-layer growth > 10 MLVolmer-Weber growth

Magnetism: < 2 MLno magnetic signal 2-10 ML⊥ magnetization > 10 ML// magnetization

10 ML  $\rightarrow$  Martensitic fcc to bcc transition  $\rightarrow$  Spin Re-Orientation transition

## $Fe/Cu_3Au(100)$

Wuttig (LEED I-V, MOKE) Prep. 2 ML pre-deposited at LT Structure: > 3.5 MLbct > 5-6 MLbcc Morphology: ---Magnetism: 1.1 ML  $\perp$  magnetiz. onset 2.3 ML $\perp$  to // SRO transition > 6 MLeasy axis <110>

Kirschner (STM, MOKE) <u>Structure</u>: <3.5 MLpseudomorphic fct 3.5-6 MLfct + bct > 10-12 MLbcc <u>Morphology</u>: 1.1 ML (three levels populated) > 2.5 ML islands elongated along <110> <u>Magnetism</u>: 2.1 ML⊥ magnetiz. onset 3.3 ML⊥ to // SRO transition

No correlation between SRO and structural modifications

SRO transition  $\Leftrightarrow$  bct phase

#### Low Temperature Growth

#### **Wuttig Kirschner**

anticipated  $\perp$  onset 0.9 ML anticipated  $\perp$  onset: 1.1 ML delayed SRO: 3.3 ML delayed SRO: 5.5 ML anticipated bcc delayed bct: 5.5 ML (films annealed after dep.)

## Fe/Cu<sub>3</sub>Au(100) GROWTH

## X-ray reflectivity during deposition



X-ray reflectivity during Fe deposition X-ray specular intensity at fixed energy and angle

Accurate thickness determination Interfaces roughness determination Morphology

Fit to a bcc model deposition only works beyond 10 ML

#### In plane Surface X-ray Diffraction



Strain is gradually released with increasing film thickness

## Electron Diffraction from Fe LMM Auger (698 eV)



Out of plane structural determination

Polar scans taken at fixed incidence angle by rotating the analyzer in the scattering plane (TM polarization)

Gradual shift of the forward focussing peak from fcc to bcc

Data comparison to multiple scattering calculations

Isotropic background calculated for given instrumental conditions (sample illumination, escape depth, emission matrix element)

Input value for lateral lattice spacing taken from XRD data

#### Linear combination of three phases: pseudomorphic fcc, fct, bct



#### Fcc structures are obtained by straining the bcc structure Expansion ∝ 87 bcc Fe: Ñ B c/a=1 a=2.87 Å ∝ Contraction c=3.74 Contraction fcc Fe 1.5 c/a=√2 This work ○ B. Schirmer, PRB <u>58</u> (1998) 4984 △ 37 ML, P. Luches, SS <u>471</u> (2000) 32 □ 250 ML, R. Rochow PRB <u>41</u> (1990) 3426 a=2.65 Å AF fcc FM fcc 1.4 3Å pseudomorphic phase **Epitaxial plane** 1.3 bcc 10 Å $\gamma$ -phase c/a 1.2 • Three Fe phases are seen to coexist between 10 and 20 Å 3.3 ML 4.8 ML • a=2.74 Å is the maximum strain of the 1.1 18 ML 53 ML $\gamma$ -fct phase 20 Å • At 25 Å, the $\alpha$ phase becomes the 1.0 $\alpha$ -phase 30 Å 36 Å dominant one and it evolves towards the Fe bcc value. 0.9 ⊾ 2.5 2.6 2.7 2.8 2.9

3.0

a (Å)

Fe films 6÷13 Å thick deposited at 150 K or 300 K, annealed at 400 K for T~10 min and cooled to room temperature (RT);





#### **TEMPERATURE DRIVEN REVERSIBLE BREAKDOWN OF PSEUDOMORPHISM**

#### In plane XRD

Fe film of 7.5 Å deposited at 150 K and annealed at 425 K

Radial scan along [100]

Pseudomorphic Fe contributes to  $Cu_3Au$  bulk peak. Rod scans clearly indicate the presence of pseudomorphic layers.

#### □ peak reversibly disappears



#### Auger Electron Diffraction

Fe Auger LMM (KE 698 eV) electron peak High temperature phase  $\Rightarrow$  better defined angular structures Films grown on bulk substrate adopt substrate structure and lattice constant (**PSEUDOMORPHISM**) in early stages of growth and gradually relax towards unstrained structure with increasing thickness

Elastic energy is accumulated during pseudomorphic growth



Stress is released by the formation of edge dislocations

a)  $\Theta = 3.6 \text{ ML}$ 



c)  $\Theta = 11 \text{ ML}$ 



B. Müller et al. PRL 76 (1996) 2358



Cu/Ru(0001)

C. Gunther et al. PRL 74 (1995) 754

•The  $\gamma$  phase and the dislocation network disappear upon annealing at 400 K;

•Upon cooling back to RT both the  $\gamma$  and the dislocation network REVERSIBLY reappear  $\rightarrow$  **PHASE TRANSITION**;

- •The phase transition is observable within the 6÷13 Å Fe thickness range;
- •The transition temperature increases from T=345 K for 6 Å Fe to T=380 K for 11 Å Fe;

•No correlation with magnetic properties of the film is observed;

## Modelling the transition and its driving force

Phase coexistence

0.00



#### Gradual enlargement of pseudomorphic patches



2.70 a [Å] 2.80

2.60

## Elastic energy must play a dominant role

Temperature dependent misfit:

•Linear thermal expansion coefficient  $\alpha(Cu_3Au)=1.9\cdot10^{-6} K^{-1}$ 

•Linear thermal expansion coefficient  $\alpha(\mbox{Fe}) {=} 1.1 {\,}^{1} 10^{-f}$  K  $^{-1}$ 



Misfit decreases with temperature!! Critical thickness for pseudomorphism breakdown increases

## Phase stacking or phase mixing ? XRD vs Tilt In plane X-ray Diffraction





## Low temperature Fe bi-layer islands





Deposition at 140 K  $\Rightarrow$  Flat bi-layer islands are formed Annealing to 400 K  $\Rightarrow$  Third layer is formed Fit to multiple scattering calculations, assuming the lateral lattice spacing to be equal to the substrate one

## Low temperature Fe bi-layer islands











## Check for Au segregation





# The pseudomorphic to bulk fcc phase transition of thin Ni films on Pd(100)

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## In-situ X-ray Diffraction and Photoemission experiments at the INFM-ALOISA beamline (Trieste)

• In-plane Surface X-ray Diffraction (XRD) to measure the formation of the relaxed **bulk-like phase**.

- Out-of-Plane XRD (ROD scans) to measure the layered structure evolution of the **pseudomorphic phase** during the phase transition.
- X-ray Photoelectron Diffraction (PED) and Spectroscopy (XPS) of the core level shifts and Valence Band to follow the **intermixing at the interface**.

• X-ray Reflectivity (XRR) to measure the perpendicular structure of the **bulk-like phase**.



- Both Pd and Ni have a Face Centered Cubic (fcc) structure.
- The **Pd** lattice parameter is **10.5 % larger** than the **Ni** one.
- Ni grows pseudomorphic to Pd(100) with a tetragonally strained (fct) structure.







The vertical spacing of the Ni fct structure has been measured to be 1.52 Å.

G.A. Rizzi et al., Surf. Sci. **522** (2003) 1



## **In-plane X-ray Diffraction radial scan**



• Room Temperature growth

• Ni grows pseudomorphic to Pd(100) up to a 12 Å coverage.

• The appearance of a **shoulder** in the radial scan **at 16** Å coverage witnesses the formation of a relaxed **bulk-like Ni phase**.

• The new relaxed phase evolves towards the Ni bulk phase as the film thickness increases.

## Out-of-plane XRD, (-2 0 L) ROD scan

• Increase of the number of oscillations from 8 to 12 Å due to increase of the Ni pseudomorphic layers from 6 to 10.

• Decrease of the oscillation amplitude due to the decrease of the sharpness of the interfaces and coverage of the pseudomorphic film.

• No shift of the minima position from 12 to 18 Å ↓

no change of the number of Ni
pseudomorphic layers → limit
thickness for pseudomorphic growth

Lateral growth of the bulk-like phase domains.





## Structural analysis of pseudomorphic Ni films

The limit thickness of stability of the Ni pseudomorphic phase is 10 layers.





## X-ray Valence Band spectroscopy at 2.5 Å Ni coverage

• The difference spectrum between the 2.5 Å Ni/Pd VB and the clean Pd one displays a peak at B.E.=5 eV (this satellite peak shifts towards the Ni bulk value of 6 eV as the thickness increases).

• The shift of the Ni VB satellite peak indicates the formation of a **NiPd alloy** (in agreement with the observed Pd 3d core level shift).





X-ray Photoelectron Diffraction of Ni 2p core level

- Enhanced pseudomorphism of the 2.5 Å film, probably due to intermixing (cfr. RODs, XPS, VB).
- At 2.5 Å, a weak peak along the surface normal (due to 3rd layer atoms) indicates both roughness and intermixing.
- Due to elasticity laws, a vertical compression of the pseudomorphic phase sets in at 5 Å.
- The strain of the pseudomorphic Ni phase is released at a critical coverage of 16 Å (cfr. XRD radial scans).





## X-ray Photoemission of Pd 3d core level



• Pd  $3d_{5/2}$  from the clean substrate displays a bulk component at B.E.=335.2 eV and a surface component at  $\Delta E$ =-0.35 eV.

- The surface component is replaced by **a new component**  $(\Delta E=+0.65 \text{ eV})$  upon Ni deposition, which **becomes dominant at a 16 Å coverage.**
- The new component is **assigned to Pd/Ni alloying** from comparison with Valence Band spectroscopy.

• The Pd 3d signal disappears beyond a coverage of 22 Å.



## • The intensity of the Pd 3d bulk component follows a regular exponential decay as a function of the overlayer Ni coverage $\rightarrow$ **no** dewetting beyond the transition.

- A different decay law is observed for the alloying component.
- Normalization to the bulk exponential decay highlights the behavior of the interface component.
- Alloying suddenly increases at the critical coverage of 16 Å, where the film transforms into the bulk-like phase.
- The Pd 3d signal fully disappears below a 30 Å Ni film, i.e. the alloying component is not due to segregated Pd, but to intermixing at the Ni/Pd interface.



## X-ray Reflectivity of the bulk-like Ni phase



• Perpendicular momentum scan of the specular reflectivity by photon energy scan at fixed scattering geometry.

- Fitting to a Ni fcc bulk model (layer spacing of 1.76 Å).
- The fit quality is lower for the 20 Å film due to large residual strain of the Ni bulk-like phase.

• The relative surface roughness remains about 20 % of the film thickness.

## $\Downarrow$

**NO dewetting**  $(2D \rightarrow 3D)$ 



## Ni film evolution

• 1 - 3 Å : fully pseudomorphic Ni film, partially intermixed with Pd.

• 4 - 12 Å : tetragonally strained pseudomorphic structure, the Ni lattice cell is vertically compressed according to elasticity laws. The Ni film is fully wetting the substrate. No change of intermixing at the interface.

• 12 Å : the Ni film achieves the limit thickness of stability of the pseudomorphic phase, corresponding to 10 layers (some of them partially intermixed or partially filled).

• 15 - 18 Å : critical coverage for the onset of the structural transition. Coexistence of pseudomorphic fct and bulk-like fcc phases. Increase of Ni/Pd intermixing together with the appearance of a LEED Moirè pattern [decoration spots with (10x10) periodicity, in agreement with the 10% lattice mismatch].

• > 18 Å : the bulk-like phase evolves towards the bulk fcc structure by a gradual release of the reidual strain. The bulk-like phase is still fully wetting the substrate.



## Conclusions

- The fcc domains grow by lateral propagation of the domain walls while the residual fct domains preserve their thickness and layered structure.
- The fct  $\rightarrow$  fcc phase transitions is driven by defects (domain walls) propagating from the interface to the surface of the Ni film.
- LEED patterns suggest the domain walls to be shear planes in the <110> direction.
- The limited roughness across the critical coverage of 16 Å excludes a dewetting (2D  $\rightarrow$  3D) transition .
- The fct  $\rightarrow$  fcc phase transition is of the martensitic type (1st order, non-diffusive transition).

## Structure of Methylthiolates on the Au(111) Surface

Self Assembled Monolayers (SAMs) of thiols on gold are prototypical metal-organic junctions of ubiquitous application in nanoscience

- Molecular electronic devices
- Surface Coatings
- Nanolithography
- Biosensors



The metal-organic interaction is not well understood and even the structure of the simplest of these materials  $(Au(111)SCH_3)$  is under dispute.

- Theory (structural optimization, static structures)

#### bridge site

- Experiments - Photoelectron Diffraction, NIXSW (Kondoh, PRL,2003, Roper, Chem.Phys.Lett.,2004)





on-top site

## Photoelectron diffraction

 Normal-emission, energy-resolved data exhibit oscillations with a single period, which are characteristic of a 180 degrees scattering from a near neighbor.



MSCD (http://www.sitp.lbl.gov/mscdpack/mscdpack.html)

• Angle-resolved PED (k=4.6)





- reflectivity results: pronounced drop in intensities for large q<sub>z</sub> not accounted for by a flat, defect-free Au surface
- disordered structure: 1° layer Au atoms modeled by multiple sites with fractional occupancies
- variable occupancies associated with adatoms and vacancies (adatom population 0.3)







## Molecular Dynamics

- The formation of on-top sites (as signaled by the breaking of a Au-S bond) is dynamic.
- "Pull-out" mechanism:



## L-Cysteine adsorption on Au(110)

Simple aminoacid, model system for adsorption to compare with simpler S bound organic molecules (alkenthiols).

Can be used to bind other more complex organic molecules li proteins, DNA

Complex system for structural determination, molecule with MANY degrees of freedom

Possible formation of intermolecular bonds.

Combination of spectroscopic and structural techniques

 $Au(110) \Rightarrow$  anisotropy in the molecular self assembly process.

Molecular beam source for UHV deposition

Spectroscopic studies on L-Cysteine from solution and UHV evaporation on Au(111) O.Cavalleri et al Phys. Chem. Chem. Phys. 6 2004 4042, J. Phys. Cond. Matter 16 2004, S2477, G. Gonella J. Phys. Chem. B 109 (2005) 18003

Cu-phtalocyanine, pentacene induce restructuring 1x2 Au(110) reconstruction







Recent STM and DF theoretical results on L-Cysteine on Au(110)

Chiral recognition of Cysteine dimers at low coverages

Kühnle et al Nature 415 (2002)891



FIG. 1. Large-scale STM image showing coexisting molec lar structures of D cysteine on a Au(110)-(1 × 2) surface (425



FIG. 3 (color). Most stable cysteine double-row structure, as

Formation of molecular rows at higher coverages, restructuring of substrate, long range order?

Kühnle et al. PRL 93 (2004) 86101

### In plane X-ray diffraction: domains size





Strongly asymmetric domain size, consistent with STM picture

## The fitting model

## × 4 reconstructed surface

+

## × 2 missing row clean surface



Anti-symmetrical displacements in the unit cell

Two laying down molecules per unit c



\_\_\_\_\_

 → 20%

#### **Out of plane XRD measurements**



more than 5 Au layers involved

poor sensitivity to the molecules



## **NEXAFS** analysis

NEXAFS spectra calculation by Plashkevych et al. Chemical Physics 232 (1998) 49

Calculation for isolated molecule

C1 and C2 contribution can not be resolved for isolated molecules.

C3  $\Pi^*$  contribution

Experimental spectra on L-cysteine powder display similar behavior

Cysteine chemisorption  $\Rightarrow$  C2 atom interacts strongly with substrate  $\Rightarrow$  C2 contribution shifts strongly.

Our high resolution spectra  $\Rightarrow$  Tentative peak association

P1  $\Rightarrow$  C2  $\sigma^*$  state strong shift due to substrate interaction P2 FL step? P3 ?

 $\Pi^*$  mostly localized on COO group.







## NEXAFS C1s: azimuthal dependence











S-C peak evidence, S-C bond parallel to the surface

**COO** plane parallel to the surface

**Other S-C and C-N related features** 

## The copper-phthalocyanine (CuPC) molecule

Electronic structure: delocalized  $\pi$ Gap in the visible energy region

Conduction in molecular solid is anisotropic Self-assembling







#### Au(110) cell structure: out of plane X-ray diffraction

**@**.4

0.8

L

11.22



According to theory most energetically favourable 1x3 reconstruction deep MR

Deep 1x3 reconstruction proposed for perylene and T6 on Au(110)

## X-ray diffraction: ×3 phase rodscans





## **Molecular orientation**

The  $2p \rightarrow 3dx^2 - y^2$  transition is strongly polarized in the molecular plane. The corresponding white lines exhibit a strong dichroism.



Carniato et al. PRB 63 85105

## **Molecular orientation**





I = S(1-|sinγ sinα+ cosγ sinθcosα|<sup>2</sup>)



 $\alpha$ =7°, incidence angle  $\theta$  polar angle  $\gamma$  molecules tilt angle

The molecule is 32 deg tilted with respect to the surface

## The ×5 phase



≻Shallow reconstructed ×5 cell

 $\succ$ Composed of  $\times$ 3 and  $\times$ 2 cells

▶Buckling of the third layer







## Cu-phthalocyanine/Au(110)





## CHAIN-CHAIN REPULSION DRIVES Cu-Pc ARRAY PHASES

## Conclusions

Weaker interaction of CuPc on Au respect to Al substrate from spectroscopy (XPS, NEXAFS)

#### but

CuPc (1x2) MR reconstruction lifted; substrate reconstruct x5 and x3 (x5 composed of a  $x^2 + x^3$  units)

x3 unit is a "shallow" MR reconstruction

x3 asymmetric cell (asymmetric atomic relaxations) from GIXD

Consistent with tilted geometry of the molecule from NEXAFS

CuPc molecules aligned along [1-10] direction long range order implies interaction between rows At RT no long range order of the molecule for the x3 saturation phase observed from HAS.