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**Pressure induced structural transformation in
nano-particles and carbon nano-tubes**

Xin-Gao GONG
Fudan University
Department of Physics
220 Handan Road
200433 Shanghai
CHINA

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Xin-Gao Gong

*Department of Physics,
Fudan University, Shanghai-200433, China*

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Supported by:

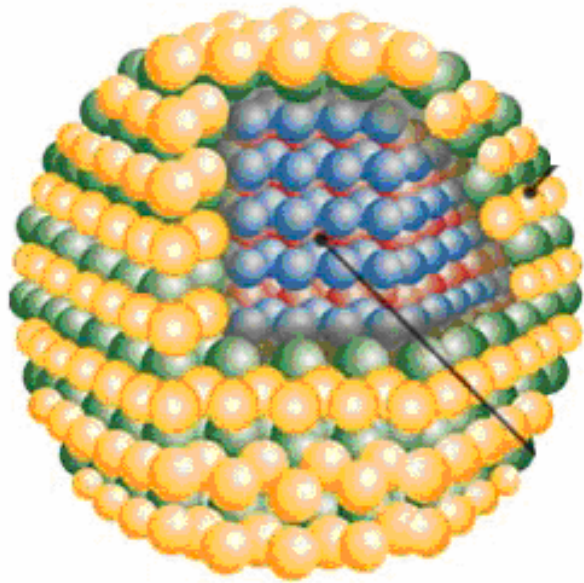
National Natural Science Foundation

National Program for Basic Research

Outline:

- **Introduction**
- **Constant-pressure MD method for finite system**
- **Some applications:**
 - (1) Melting of nano-particles*
 - (2) Solid-solid transition in nano-particles*
 - (3) Pressure-induced hard-soft transition of CNTs*
- **Summary**

Why finite systems under pressure?



- Nanocrystal has much larger surface to volume ratio, influence elastic and thermodynamic properties

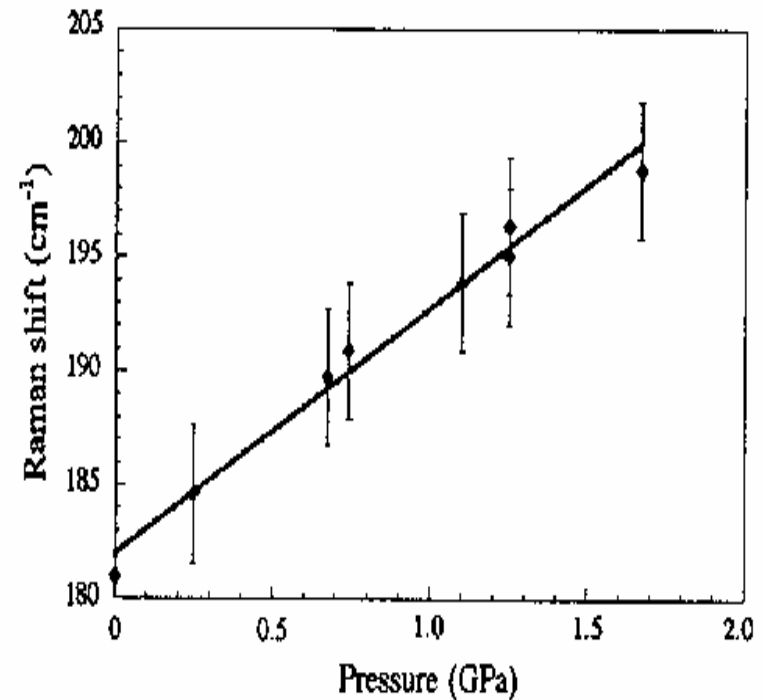
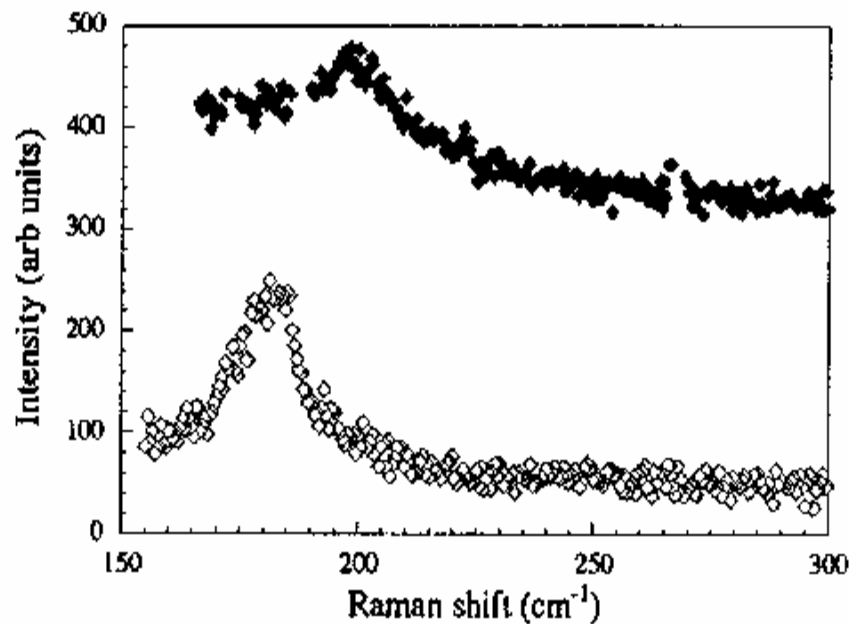
How different the small is?

- Helping to understand the micro-mechanism of bulk material

Carbon Nano-tube under the pressure:

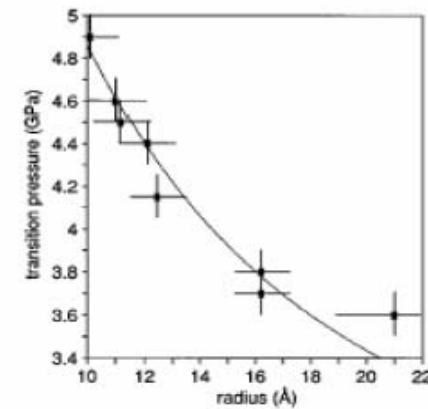
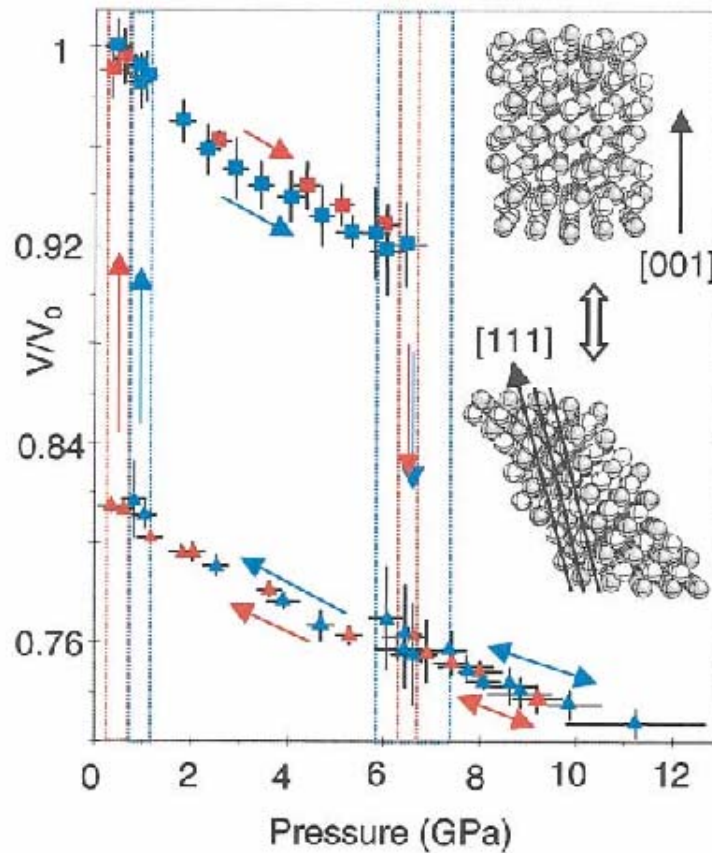
Exp. Observation

Breathing mode disappeared at $P=1.7$ GPa



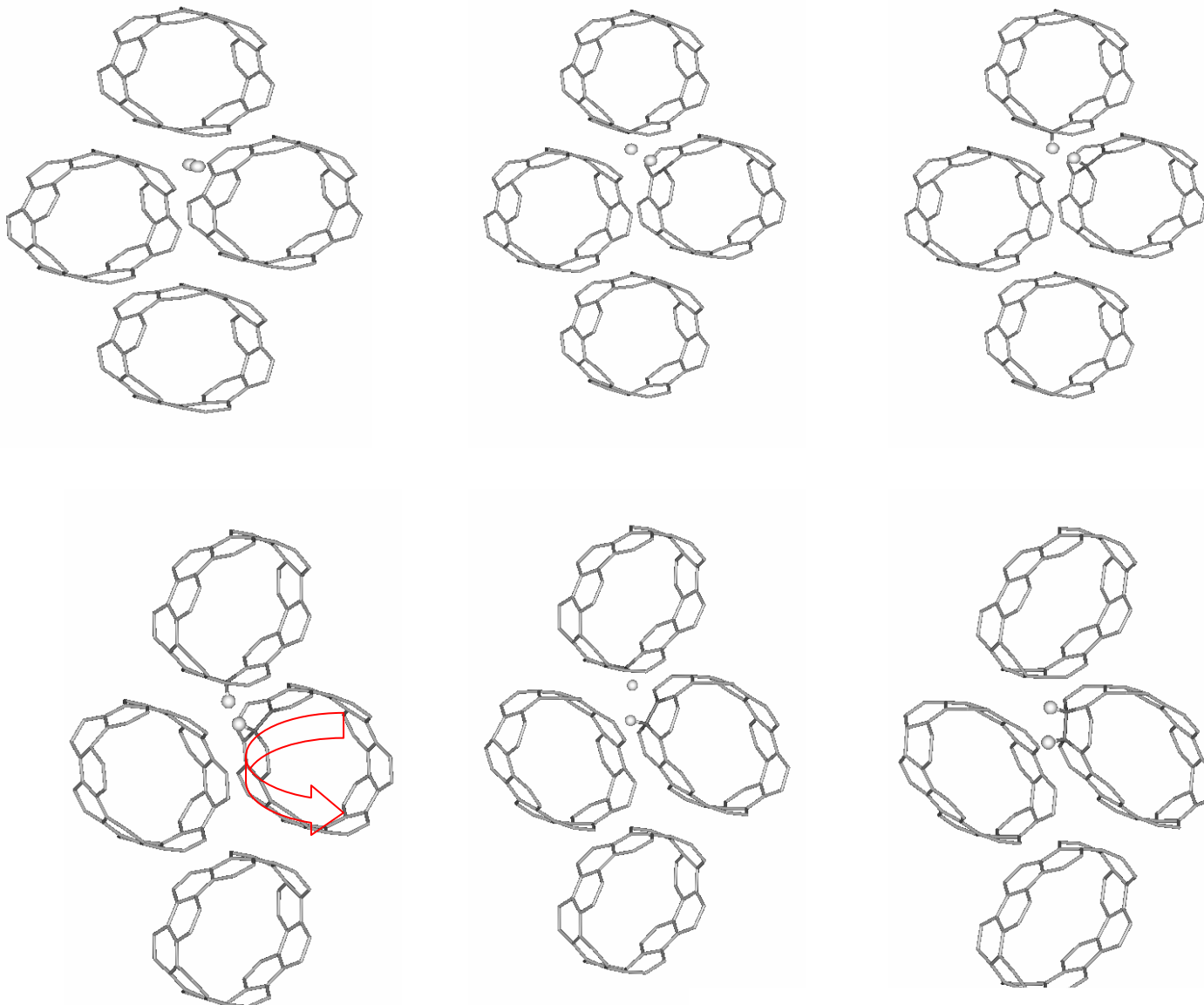
Peters et al. PRB, 61(5939)

Structural transformation of CdSe



Wurtzite to rock salt transformation pressure vs. size of CdSe nanocrystals (Alivisatos group).

Pressure induced H_2 dissociation in tube bundles:



Chan, Chen, Gong and Liu, PRL, 8720(2001)

- **Conventional constant-pressure MD:**

Parrinello-Rahman

$$H = \sum_i^N \frac{1}{2} m_i \dot{\vec{R}}_i^2 + \Phi(\{\vec{R}_i\}) + \frac{1}{2} \mu \dot{\Omega}^2 + P\Omega$$

Where volume: $\Omega = \vec{a} \times \vec{b} \cdot \vec{c}$

Finite system without any periodicity:???

We propose:

JPCM 14, L487(2002)

$$L = \sum_i^N \frac{1}{2} m_i \dot{\vec{R}}_i^2 - \Phi(\{\vec{R}_i\}) - P_{ext} \Omega(\{\vec{R}_i\})$$

Where volume is the whole volume of the finite system.

If we write the volume of the system into in following form

$$\Omega = \Omega(R_{ij}^3)$$

One can easily prove that:

$$P_{ext} = P_{int} = \frac{1}{3\Omega} \left(\sum_i^N m_i \dot{\vec{R}}_i^2 - \sum_i^N \vec{R}_i \cdot \nabla \Phi \right)$$

Motion of Equations:

$$\frac{d^2 R_i}{dt^2} = -\nabla_i \Phi - P_{ext} \nabla_i \Omega$$

Main Features:

Computation overhead is small!

Parameter free!

Easy to code, either in DFT based code or MD!

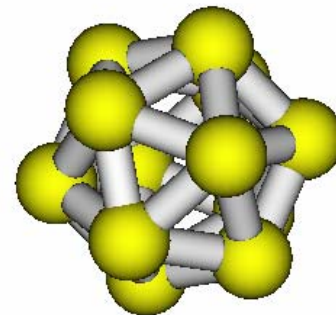
How to calculate volume of a finite system:

$$\sum_i^N \vec{R}_i \cdot \nabla \Omega = 3\Omega \quad \Omega = \Omega(R_{ij}^3)$$

A simplest scheme: atomic sphere to approximate WS cell
(not too bad for the metal system)

$$\Omega = \sum_i \gamma_i \frac{4\pi}{3} \sum_{j \neq i} \left(\frac{r_{ij}}{2}\right)^3 \quad \begin{array}{l} \gamma_i \text{ constant,} \\ r_{ij} \text{ interatomic distance} \end{array}$$

A good approximation: the volume of the polyhedron



Recent Progress by M. Cococcioni et al. : PRL(2005)

The volume is defined from the charge density:

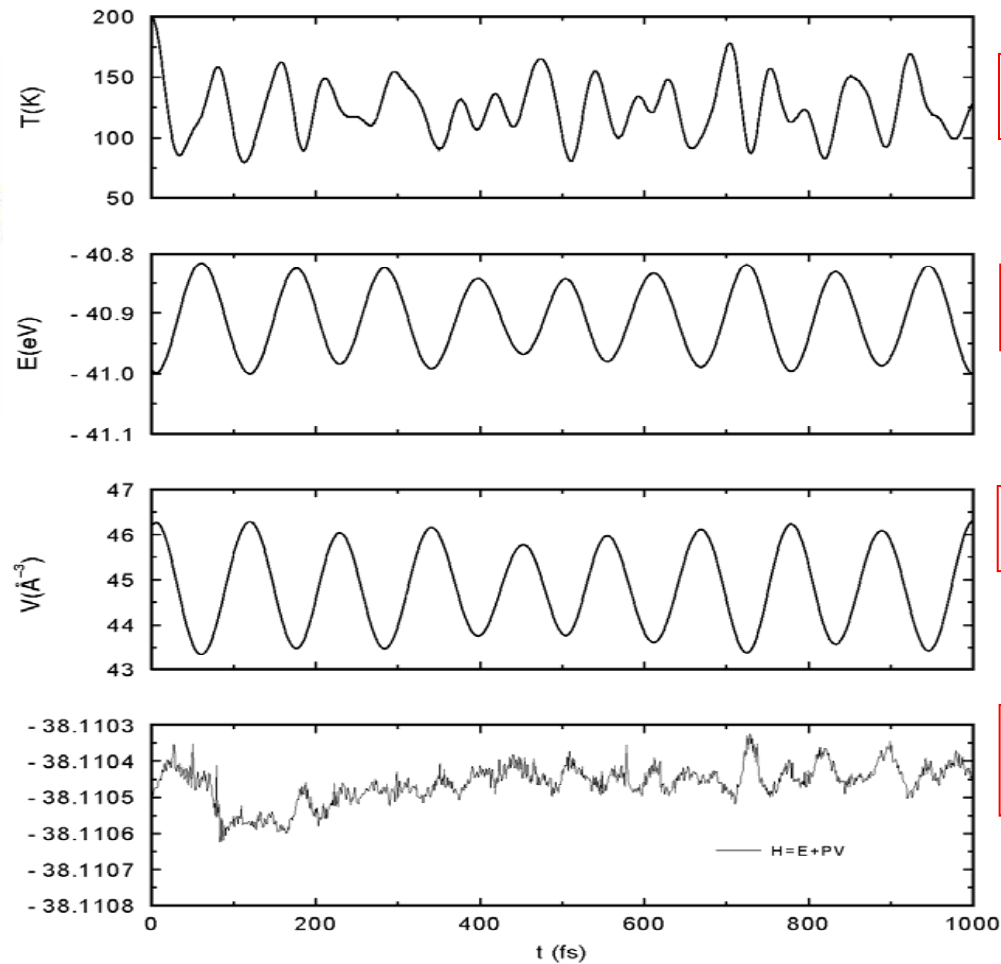
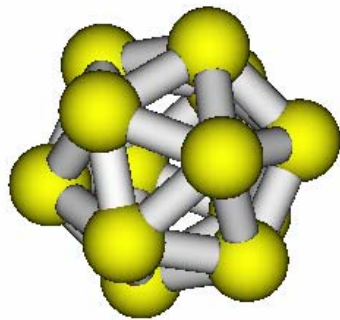
$$V_q = \int d\mathbf{r} \vartheta(\rho(\mathbf{r}) - \alpha),$$

If the step function is approximated by a gaussian:

$$\Phi_V(\mathbf{r}) = P \frac{\delta V_q}{\delta \rho} \Big|_{\rho=\rho(\mathbf{r})} = \frac{P}{\sigma \sqrt{2\pi}} e^{-(\rho(\mathbf{r})-\alpha)/2\sigma^2}.$$

Constant Pressure in Ab-initio MD

Al_{13} cluster at 10 GPa



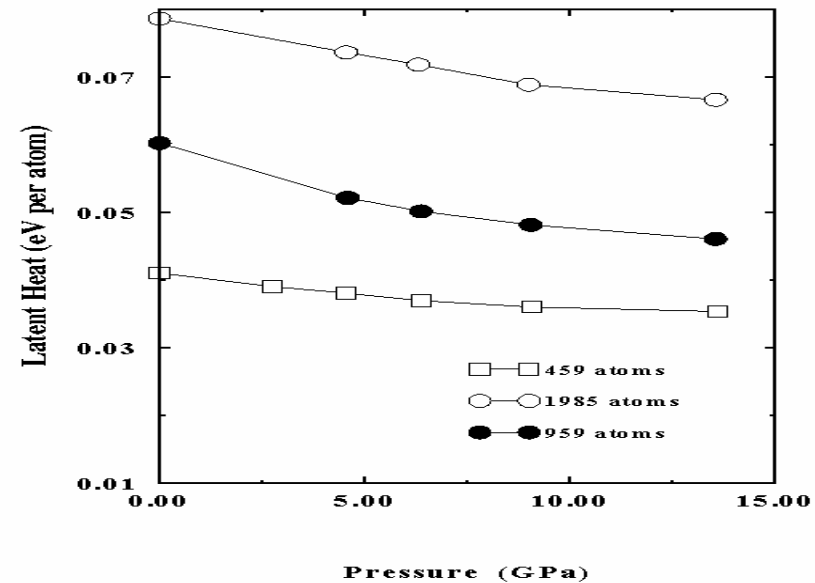
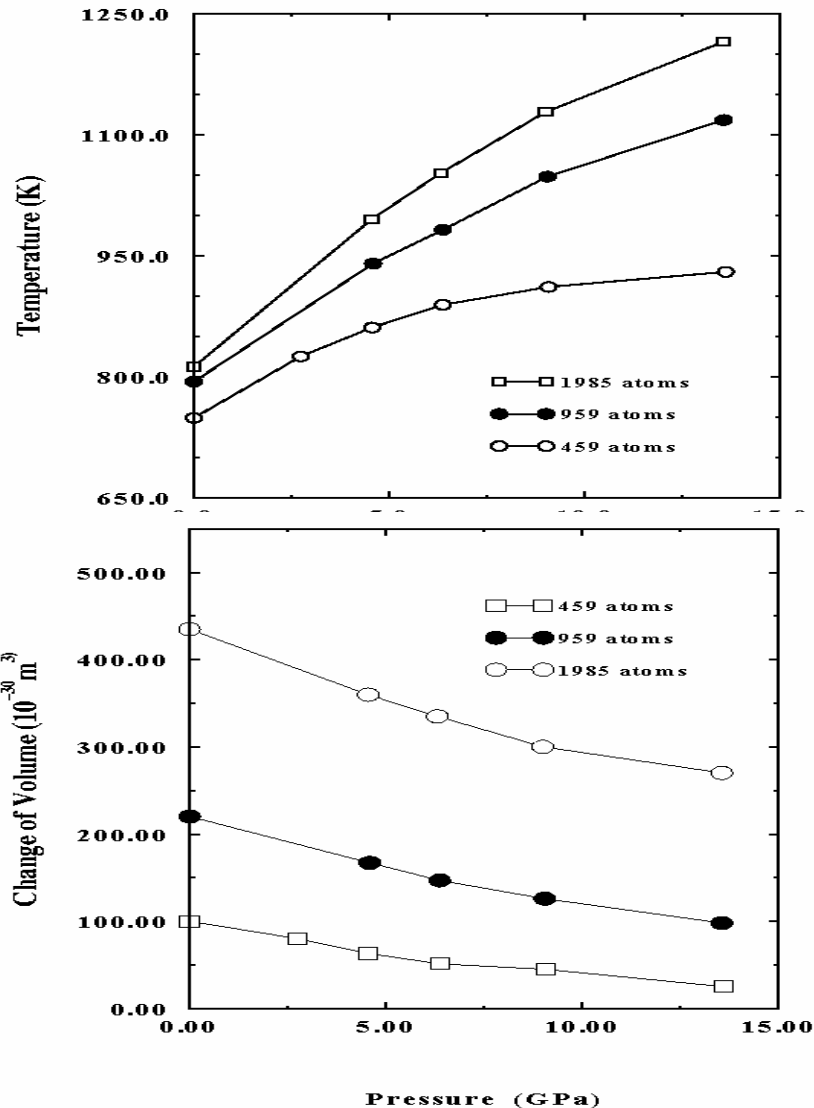
Constant Temperature

Total energy fluctuation

Volume fluctuation

Constant Enthalpy

Melting of small Ni particles under pressure: (classical MD with Sutton-Chen potentials)

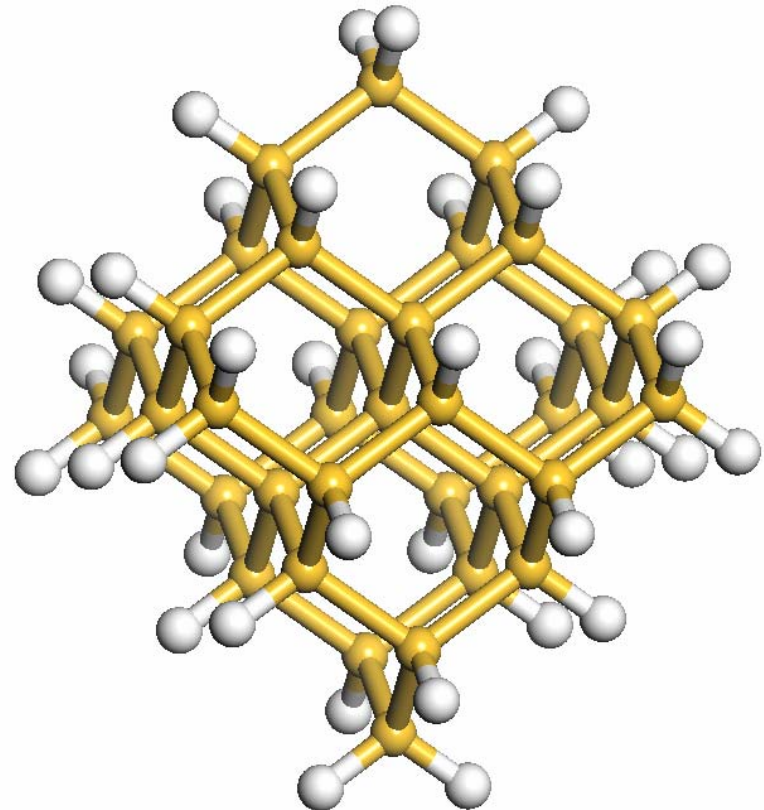


$$\frac{dp}{dT} = \frac{l}{T(V_2 - V_1)}$$

Ni particles: Small is the same!

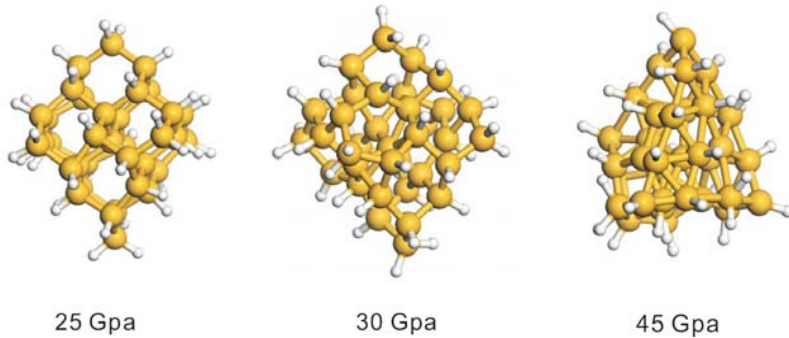
Si cluster under pressure

- Si_{35} saturated with 36 hydrogen atoms
 - Si_{35} structure is cut from bulk silicon.
 - Constant pressure is exposed on Si atoms.



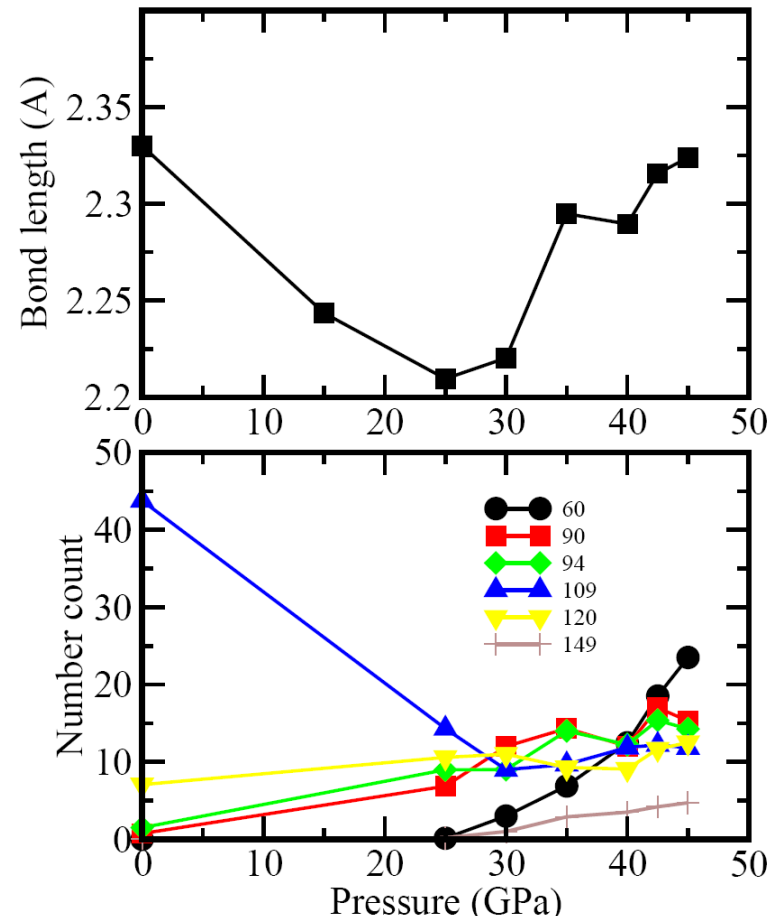
Si₃₅ cluster under pressure

Structure changes at 30 GPa



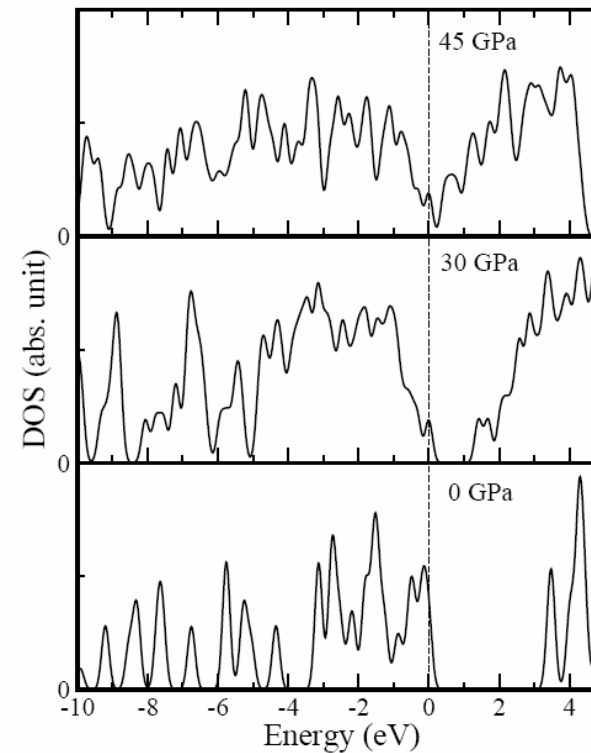
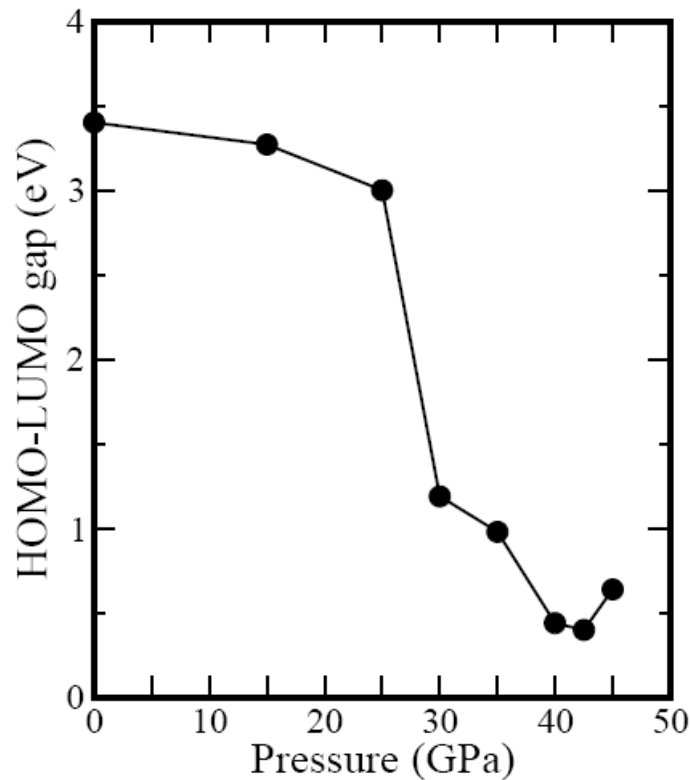
Bond length and angle

- Bond length increase beyond P_c
- 60 degree angle which refers to hexagonal structure appears.



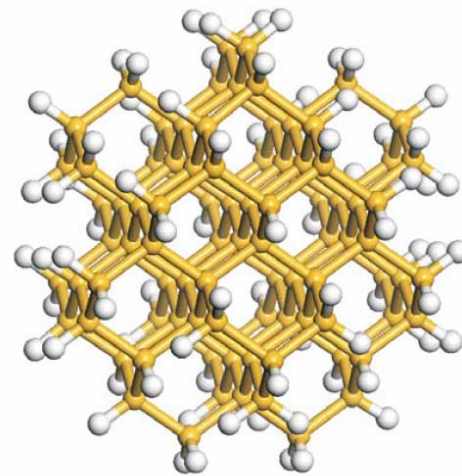
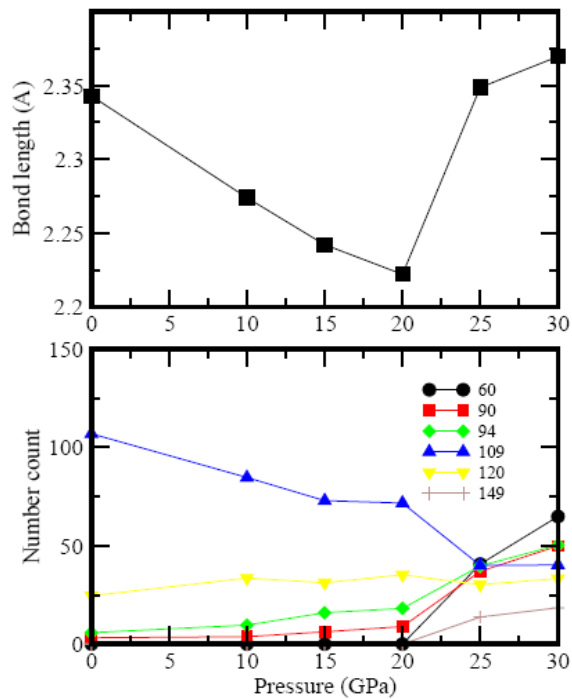
Si₃₅ cluster under pressure

- Metallization driven by pressure
—HOMO-LUMO gap decreases

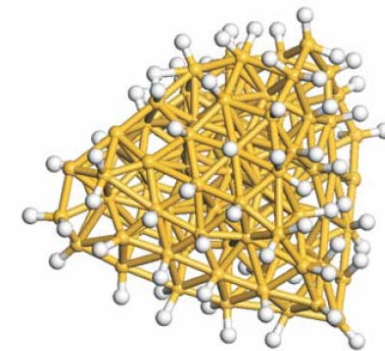


Si₈₇ cluster under pressure

Si₈₇ has a lower P_c



0 Gpa



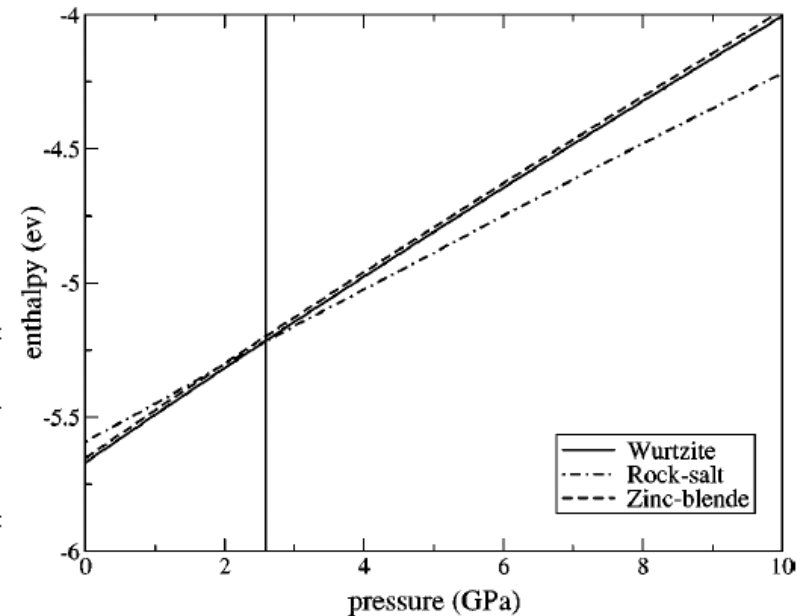
35 Gpa

Pressure Induced Structure transformation of large CdSe particles

Classical Potential:

$$V_{ij} = \frac{q_i q_j}{r_{ij}} + 4\epsilon_{ij} \left\{ \left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right\}$$

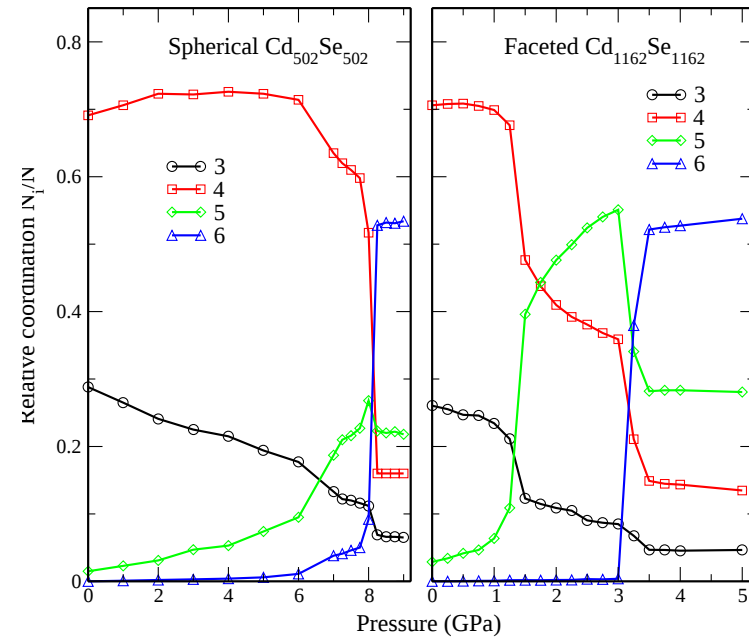
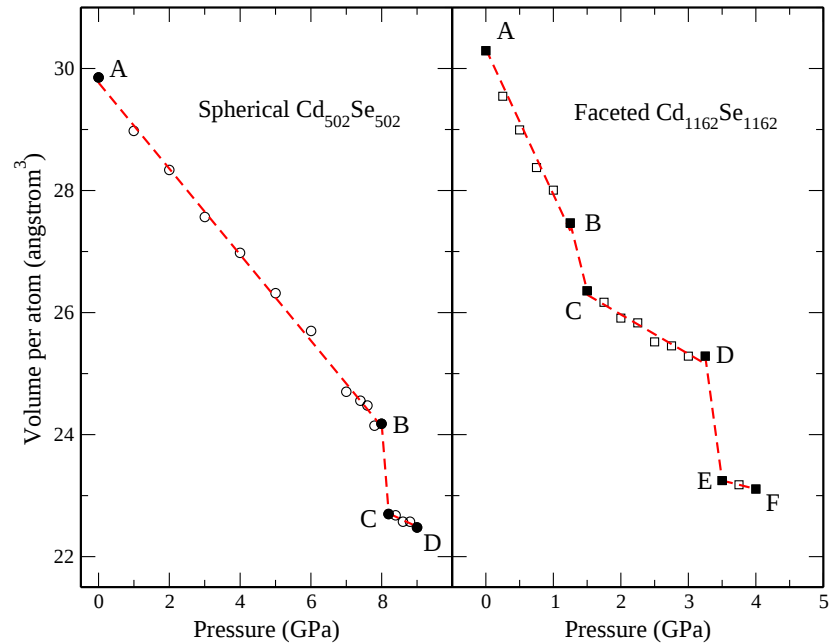
	q	σ (Å)	ϵ (K)
Cd	1.18	1.98	16.8
Se	-1.18	5.24	14.9



The two-body potential consists of long-range Coulomb part and a short-range part. The phase transformation of bulk CdSe from wurtzite structure to rocksalt structure occurs at about **2.5** GPa, in agreement with experimental measurements.

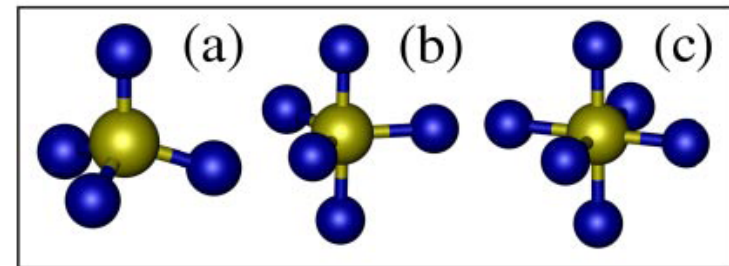
E. Rabani, J. Chem. Phys. **116**, 258 (2002)

Transformation path

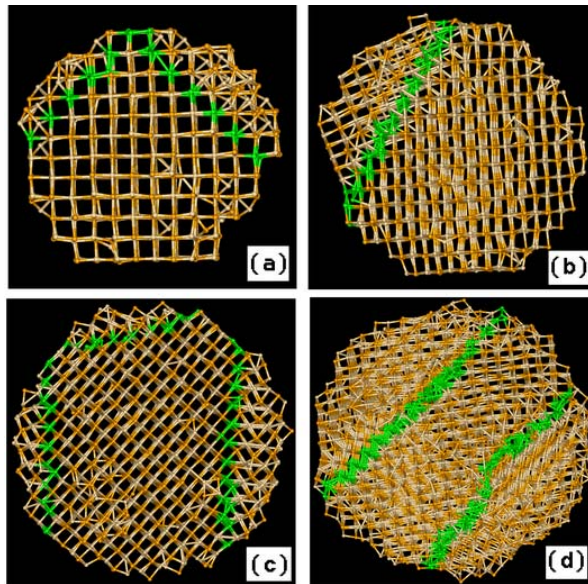


Spherical: Wurtzite $\rightarrow$ Rocksalt

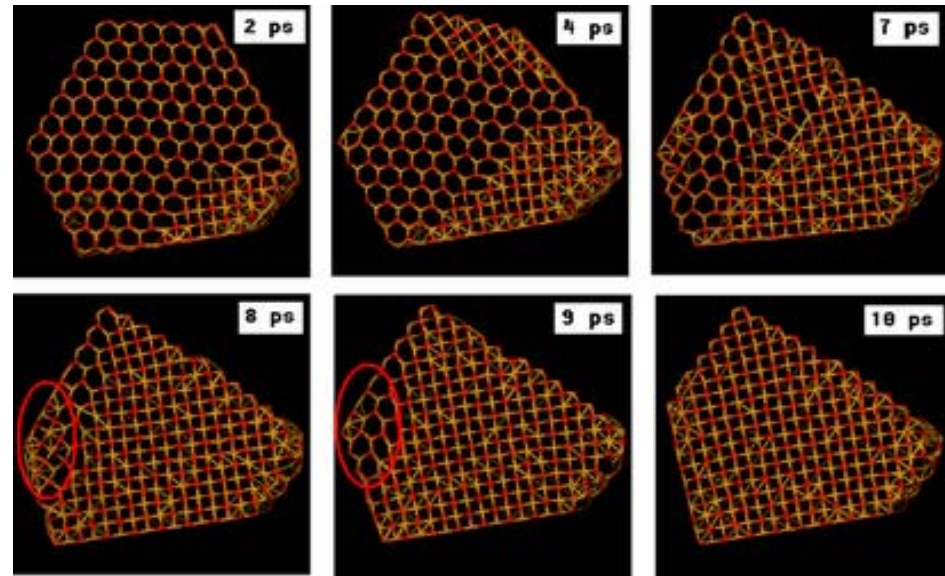
Faceted: Wurtzite $\rightarrow$ h-MgO (five coordinated) $\rightarrow$ Rocksalt



Domains after transformation

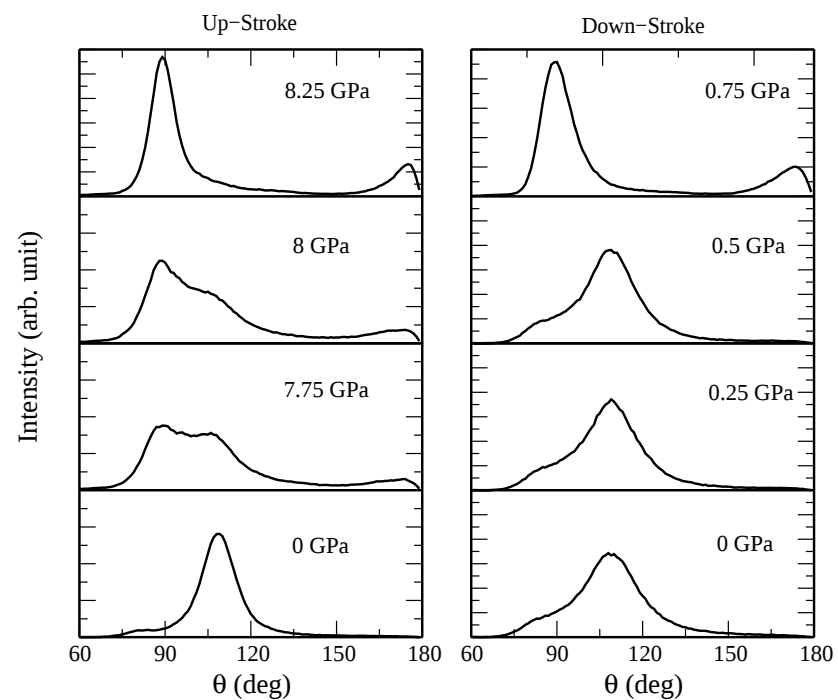


Strain domains after the structural transformation in spherical nanocrystals. The grain boundary is shown as green atoms



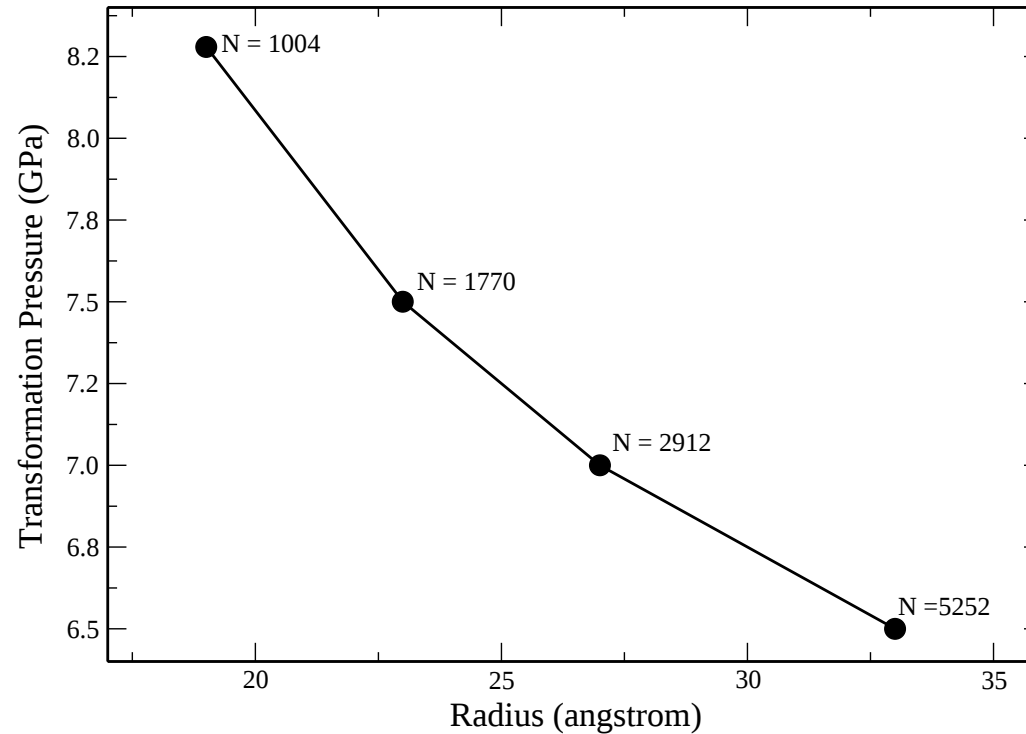
The snapshots of the MD simulation for the faceted $\text{Cd}_{1162}\text{Se}_{1162}$ nanocrystal transforming from five-fold coordination structure to rocksalt structure.

The hysteresis behavior

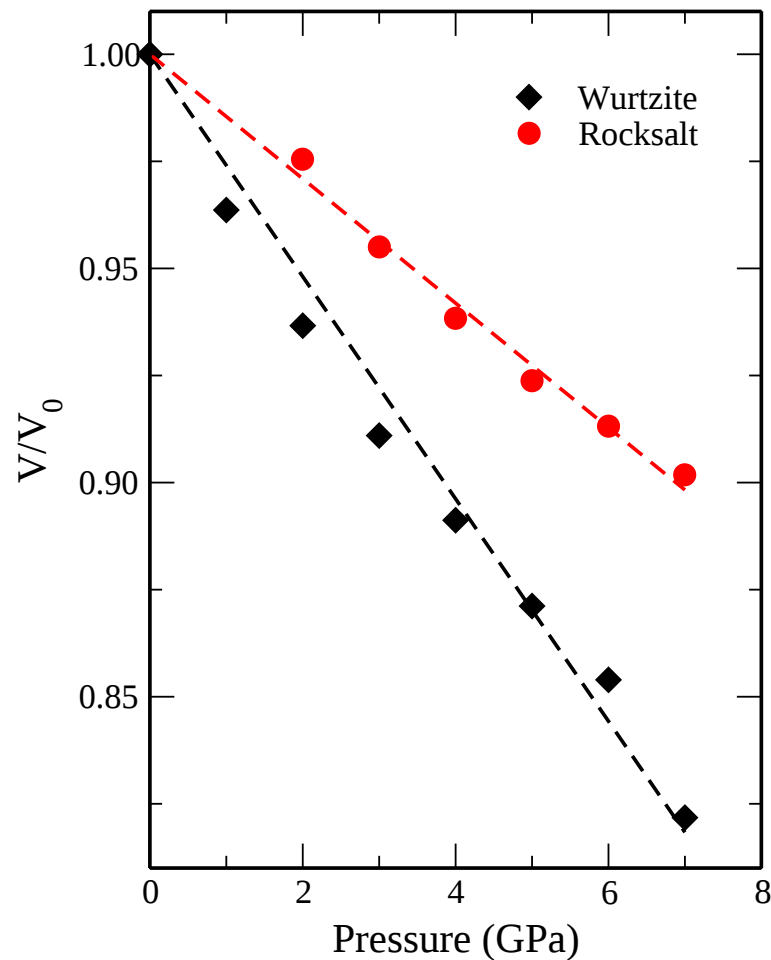


Averaged bond angle distribution of spherical $Cd_{502}Se_{502}$ nanocrystals.

Size effect on the transformation pressure



Variation of the transformation pressure with radius for nanocrystals at 300 K. With increasing nanocrystal size, the transformation pressure approaches the bulk value from above.



Bulk modulus

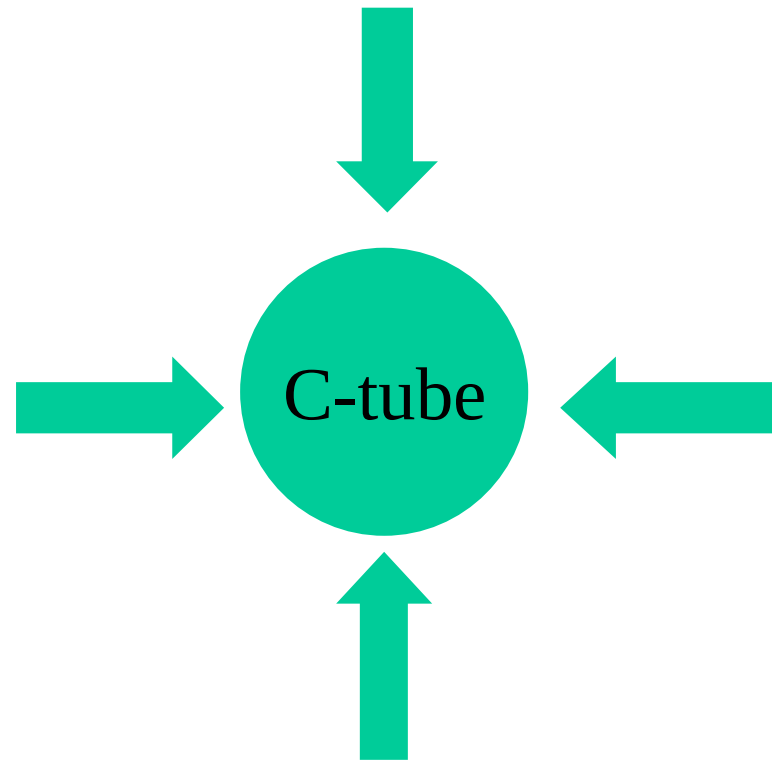
Physical constants	Wurtzite	Rock salt
B_0 (GPa)	37 ± 5	74 ± 2
B'_0	11 ± 3	
V_0 (m ³)	5.62×10^{-29}	4.36×10^{-29}
c_1 (N/m)	0.34	0.63
c_2 (N/m) (Å ²)	84	83

Radius 21 angstrom

S. H. Tolbert and A. P. Alivisatos,
J. Chem. Phys. **102**, 4642 (1995).

Volume as a functions of external pressure for spherical $Cd_{502}Se_{502}$ nanocrystal. The data can be fit with a linear volume compressibility of $B_0 = 38.5$ GPa for wurtzite and $B_0 = 68.8$ GPa for rocksalt

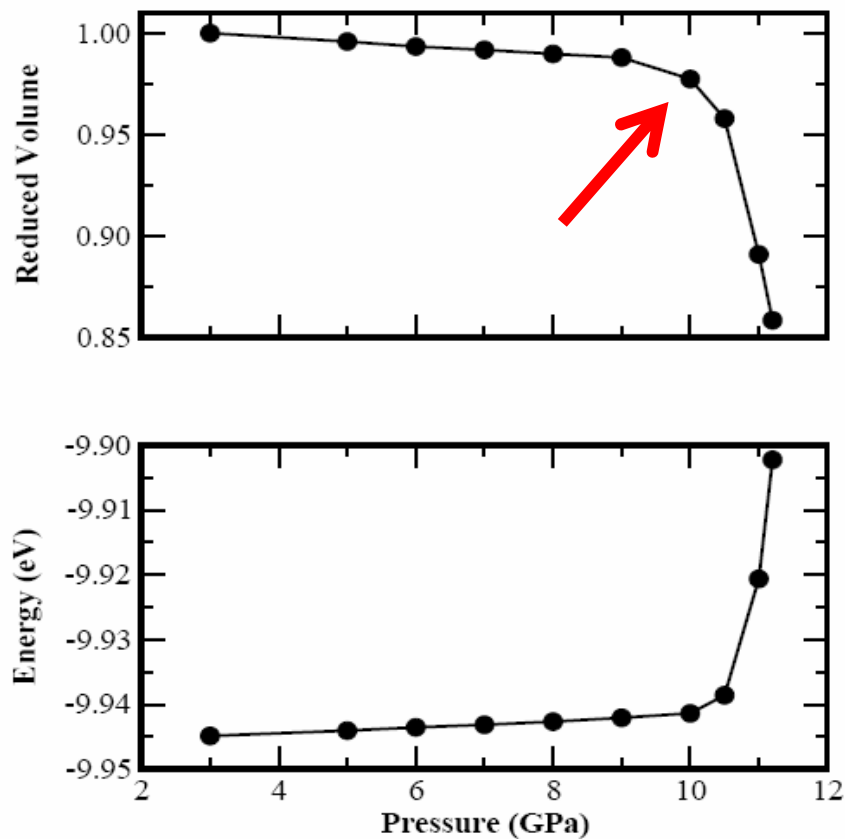
Schematic show of carbon tube under hydrostatic pressure



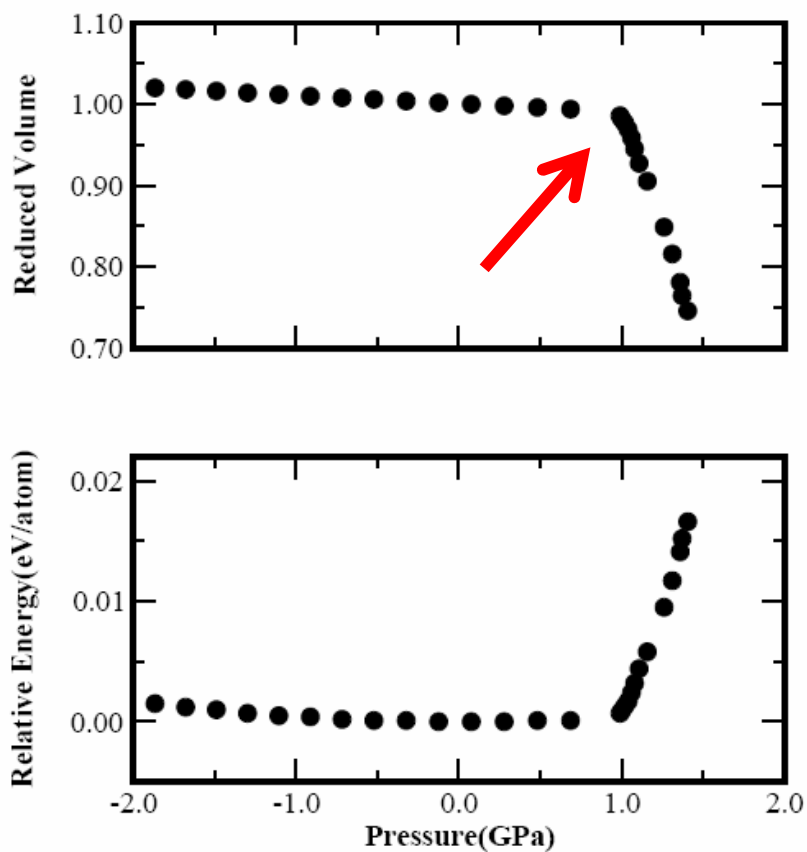
Periodic boundary condition in axial direction

PRB70, 165417(2004)

Pressure Induced hard-soft transition in CNTBs

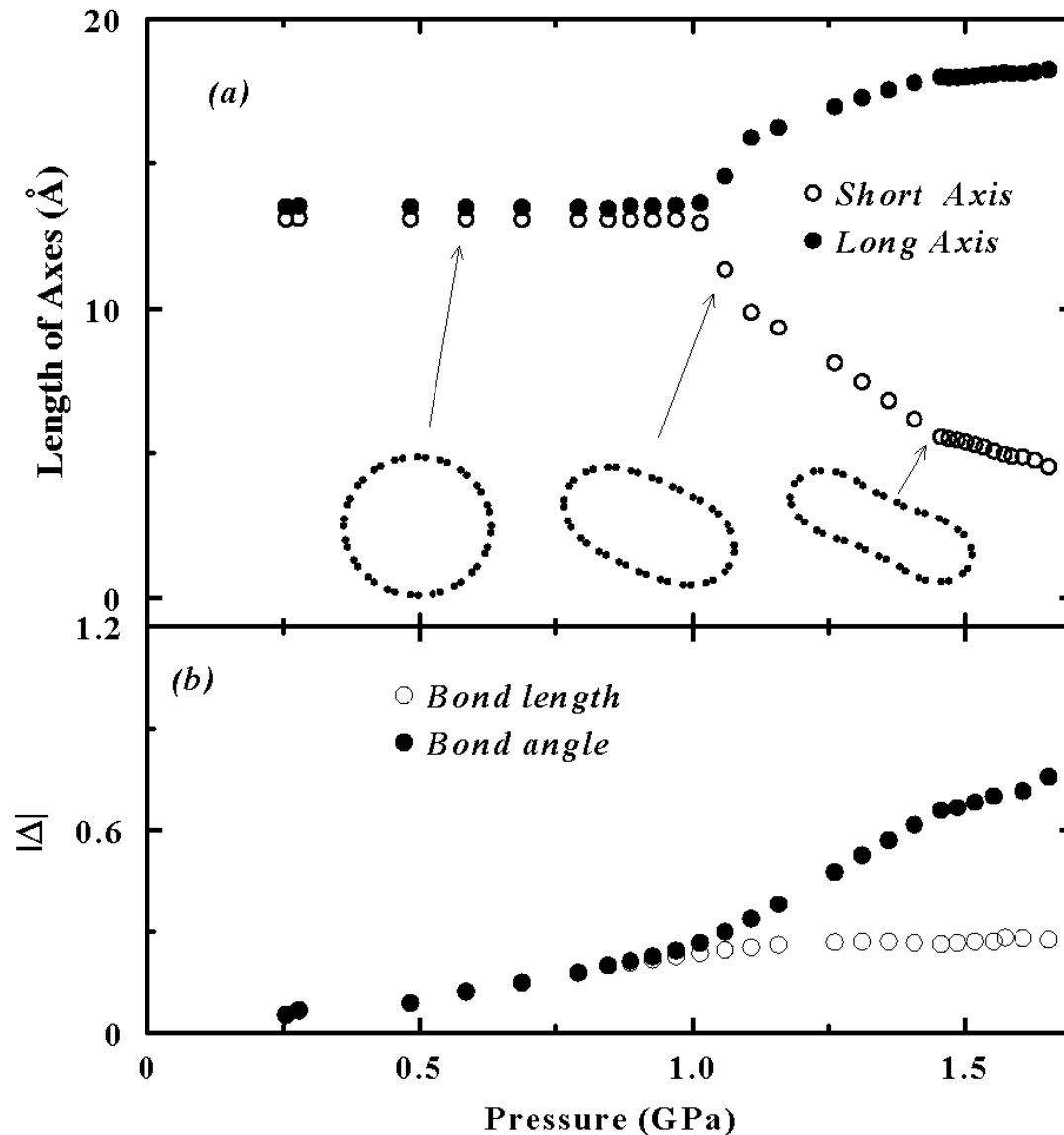


Ab-initio MD, (6, 6)



MD, (10,10)

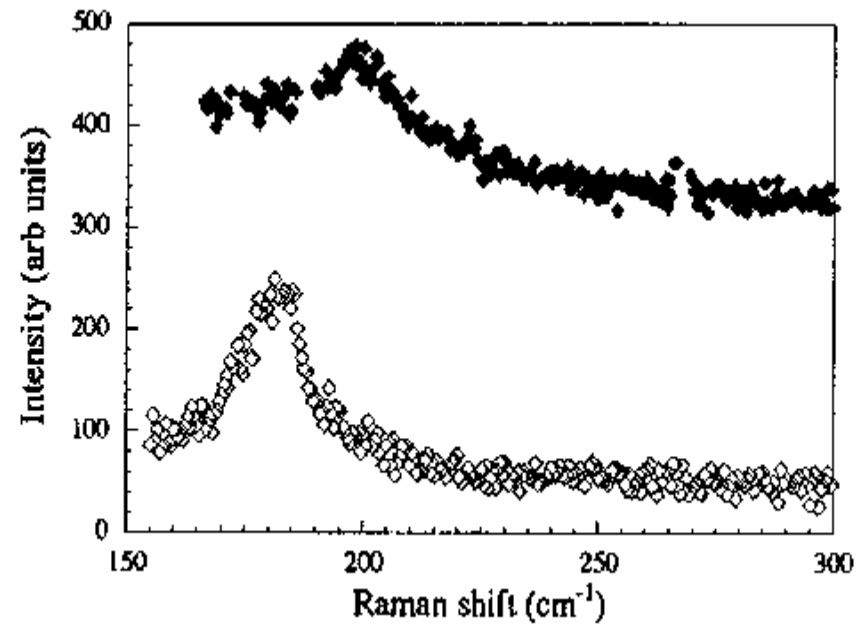
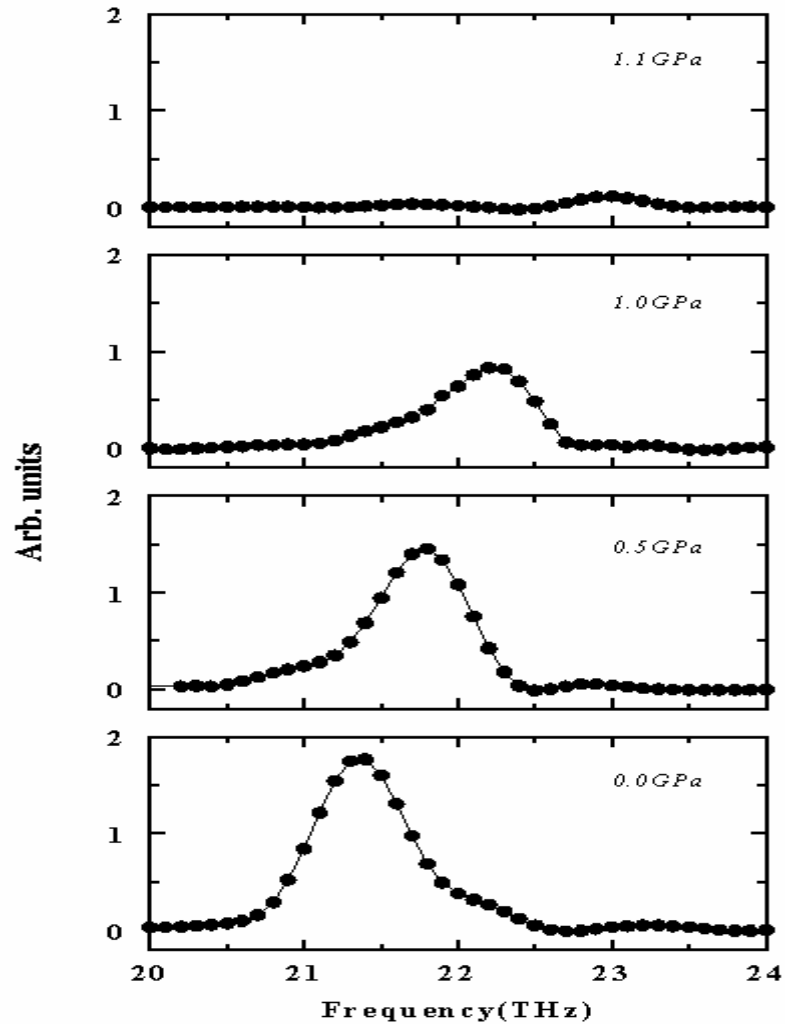
Shape Changes under pressure



Shape change with hard-soft transition

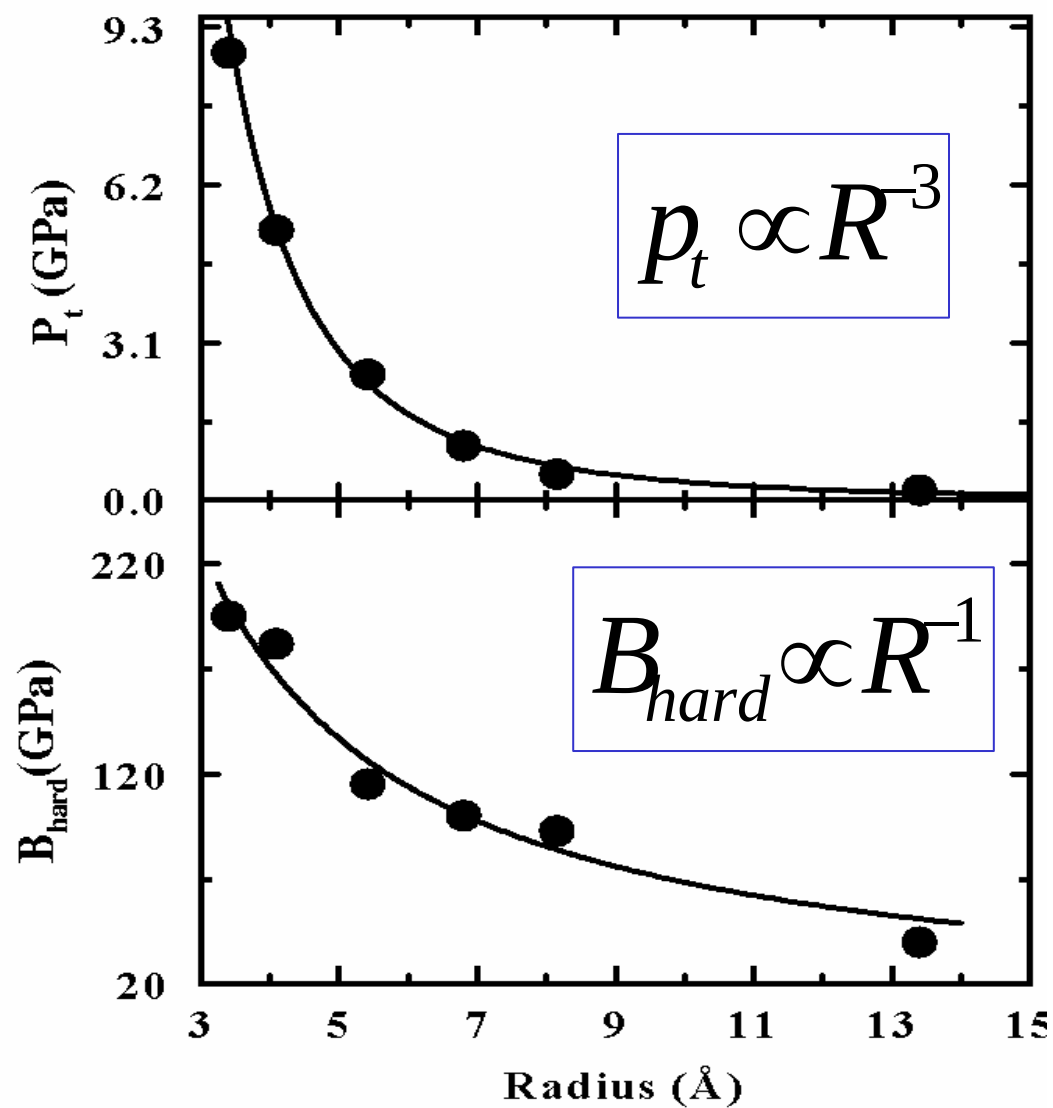
Bond angle responsible to the pressure after the change

Softening the breath mode of the Carbon tube:



Peters et al. PRB, 61(5939)

Size dependence of transition pressure



Elastic model for tube under pressure:

- Elastic energy of the tube per unit length:

$$E = \frac{D}{2} \oint \frac{1}{\rho^2} dl + \frac{C}{2} \oint \left(\frac{\oint dl - L_0}{L_0} \right)^2 dl + PA,$$

where $D = Yh^3/12(1-\nu^2)$ and $C = Yh/(1-\nu^2)$

Y: Youngs modulus

ν : Poisson ratio

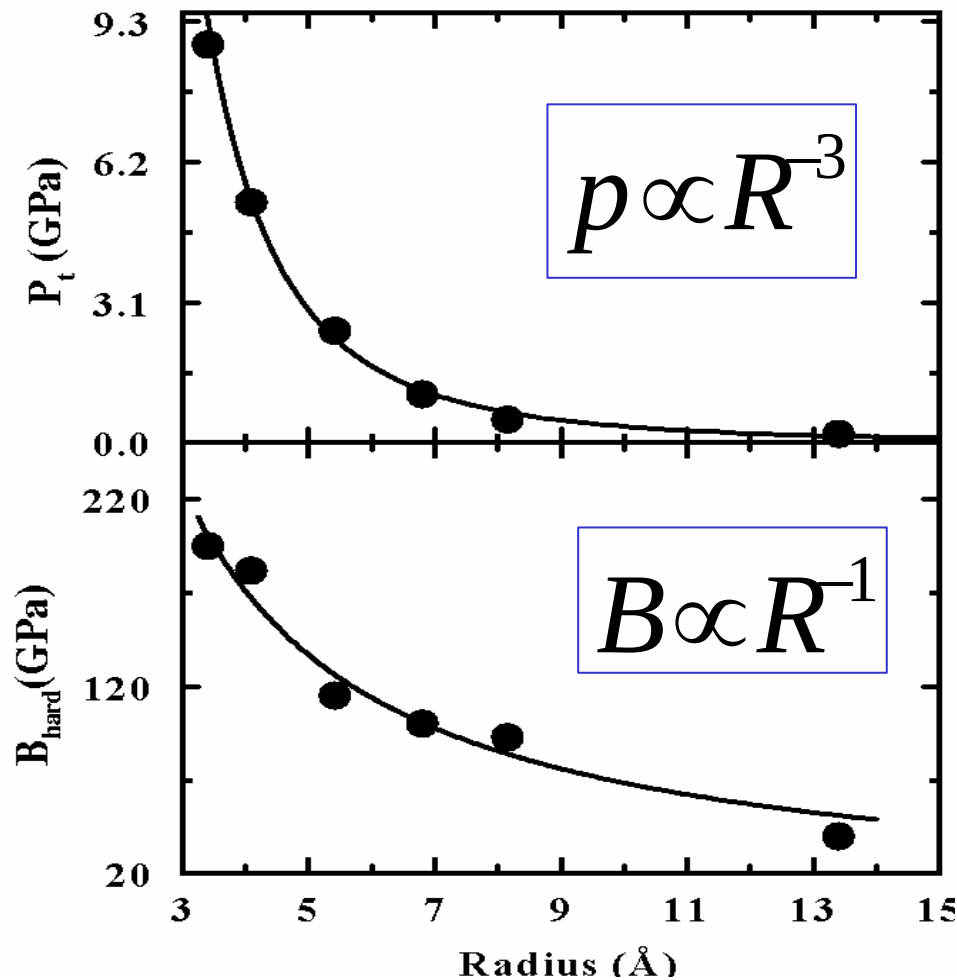
h: thickness of tube

A: area of tube cross section

L_0 : perimeter of thue cross section

Size dependence of transition pressure and bulk modulus: Simulation and modeling

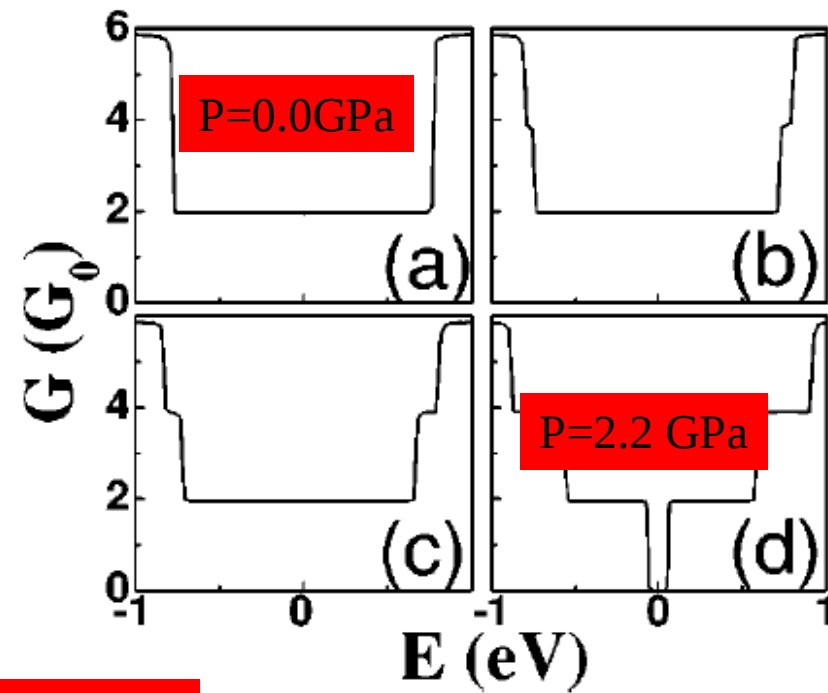
$$P_t = \frac{3D}{R_0^3(1 + \epsilon_c)^3} \approx \frac{3D}{R_0^3}, \quad B_h = \frac{C}{2R_0}.$$



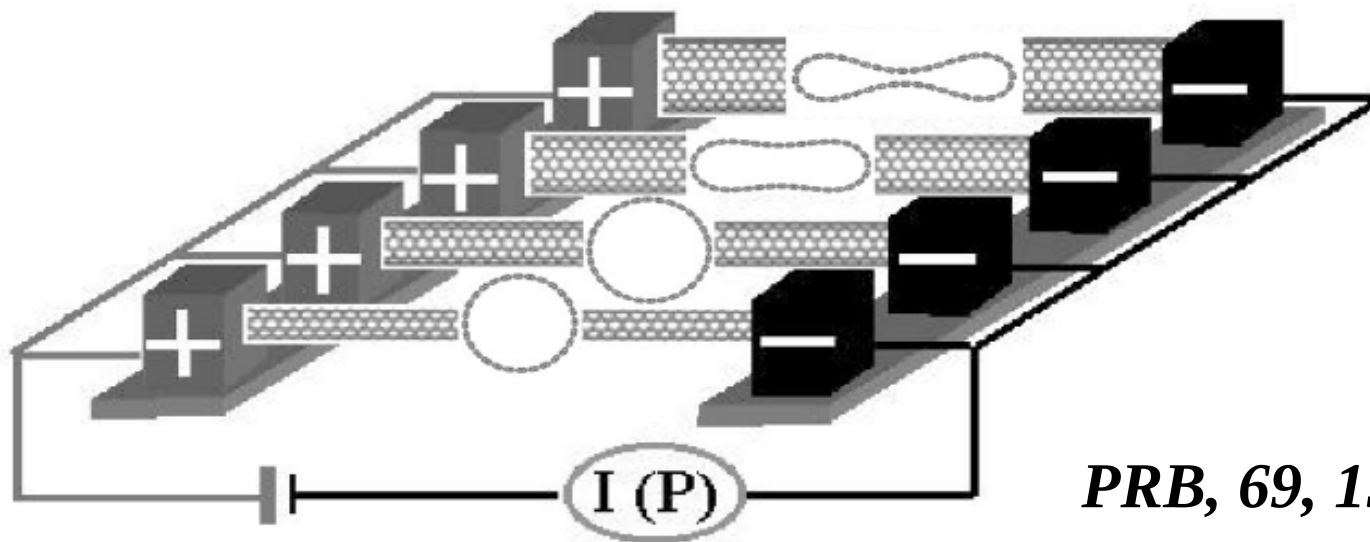
$$D = 0.76 \text{ eV}$$

In agreement with
available data

Conductivity of CNT Under Pressure:

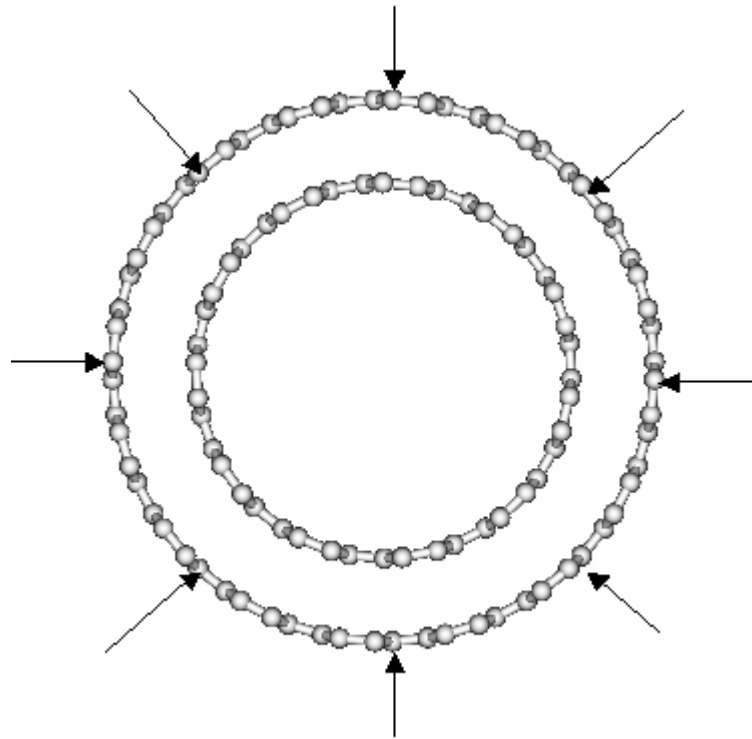


Conceptional Pressure Sensor:



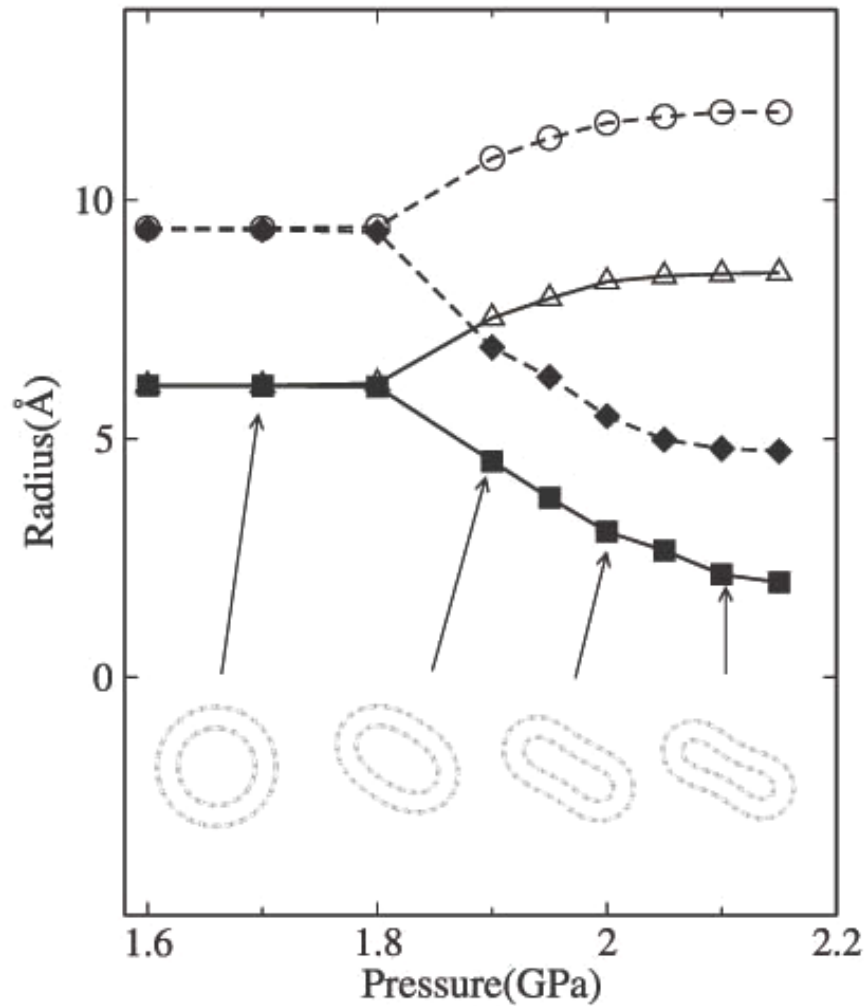
PRB, 69, 153406

Double Wall Carbon nano-tube Under Pressure:

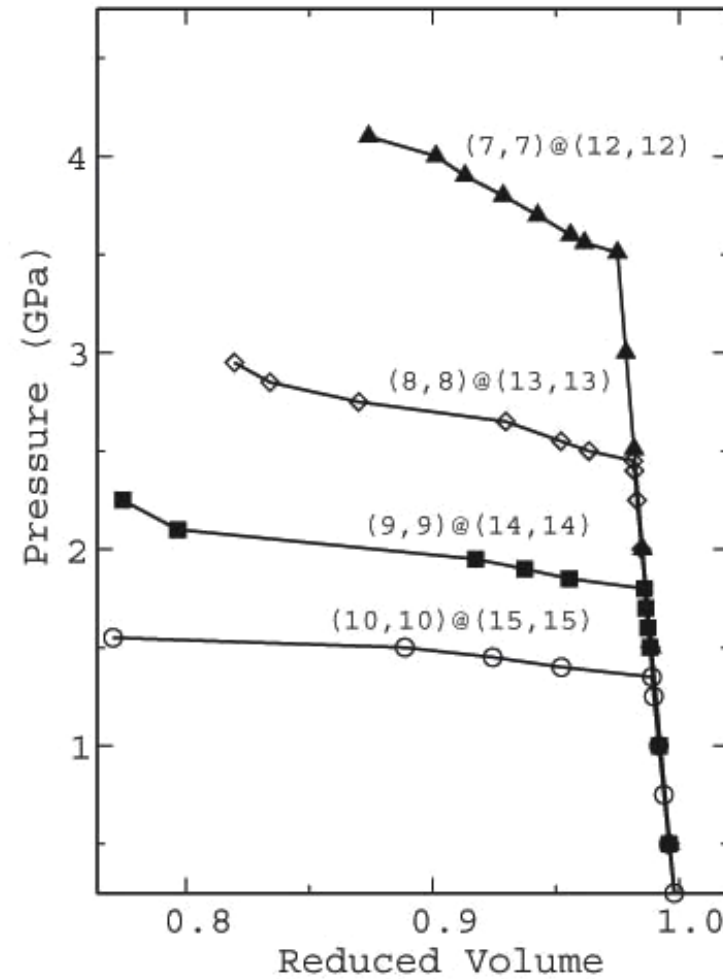


outer tube: uniform pressure
inner tube: response to the
VdW potential

YE et al. PRB72, 035454(2005)

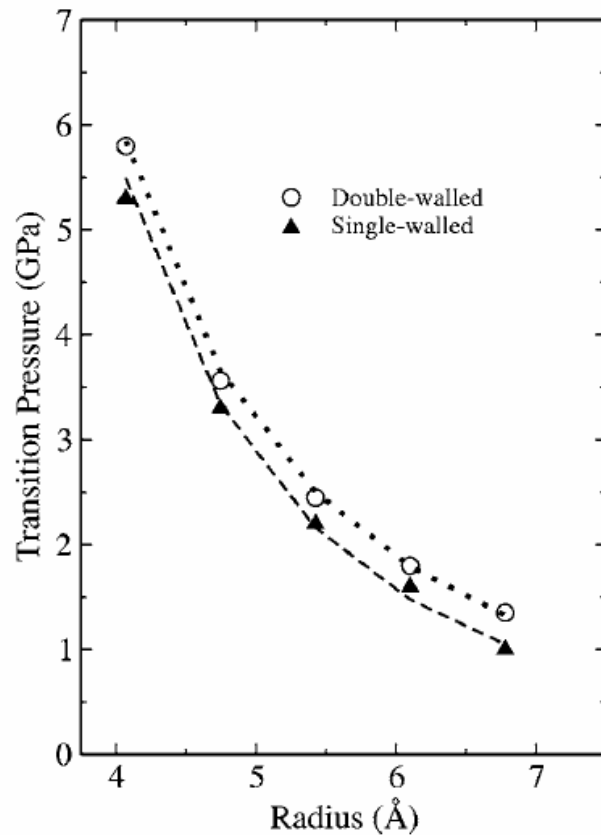


The long radius, short radius of tube as a function of pressure for (9,9)@(14,14) nanotube. At the same pressure, outertube and innertube collapse from circle to oval

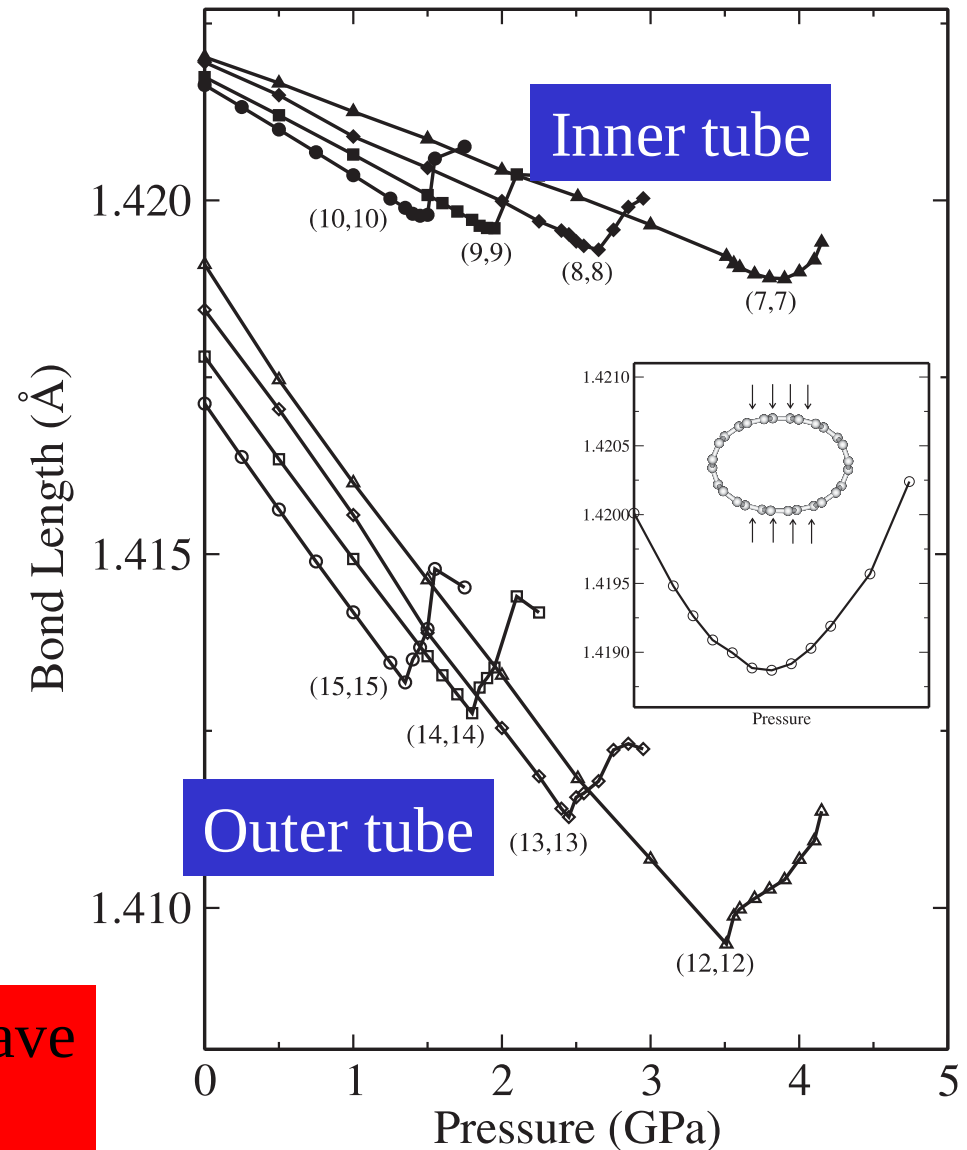


Pressure as a function of the reduced volume for DWCNTs at 300 K. The discontinuity of the slopes indicates hard-soft transition under hydrostatic pressure

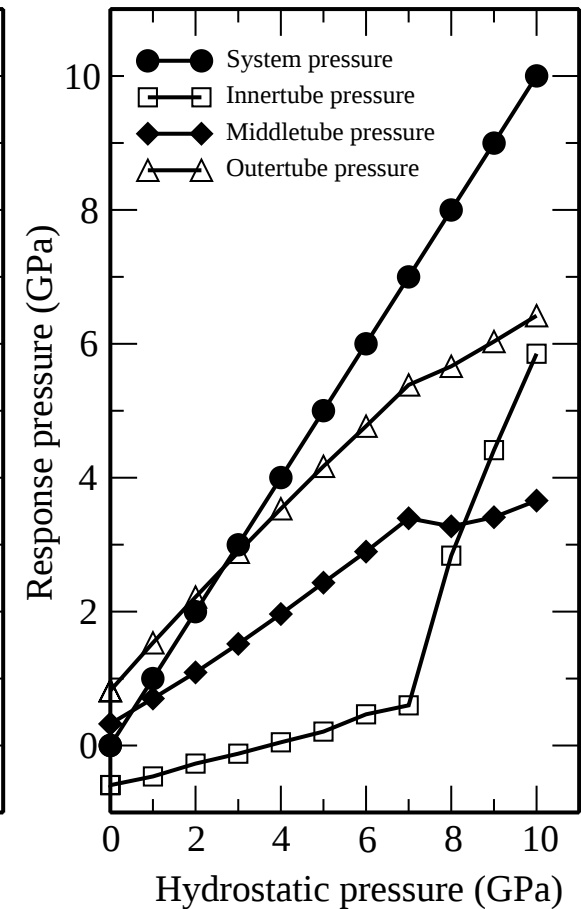
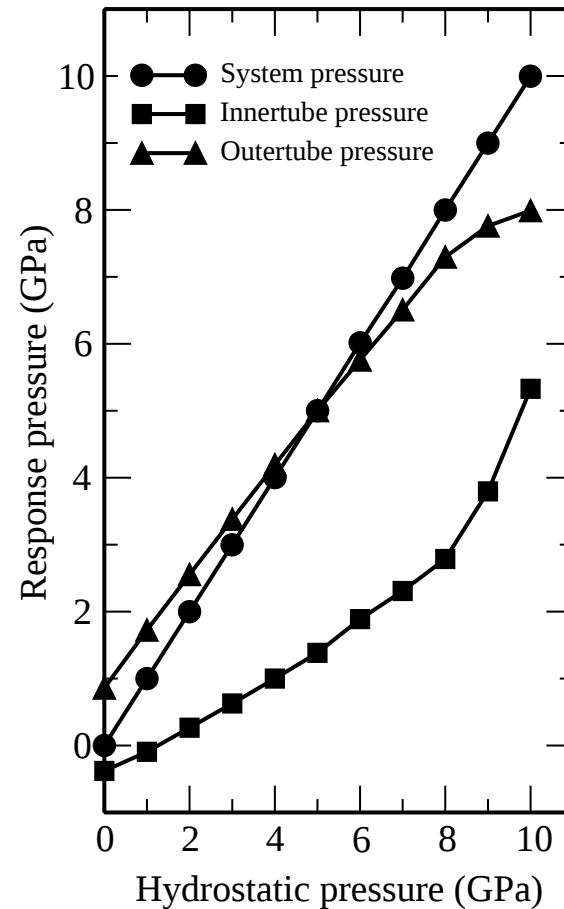
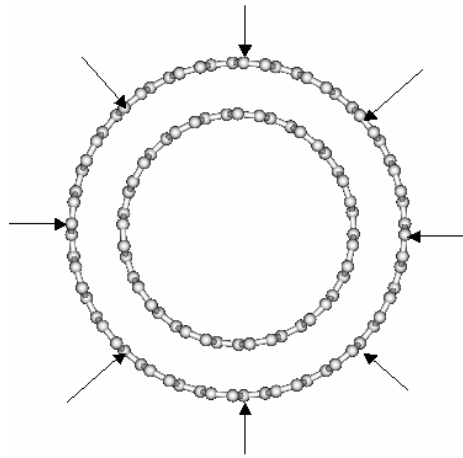
Transition pressure for Double-walled tubes



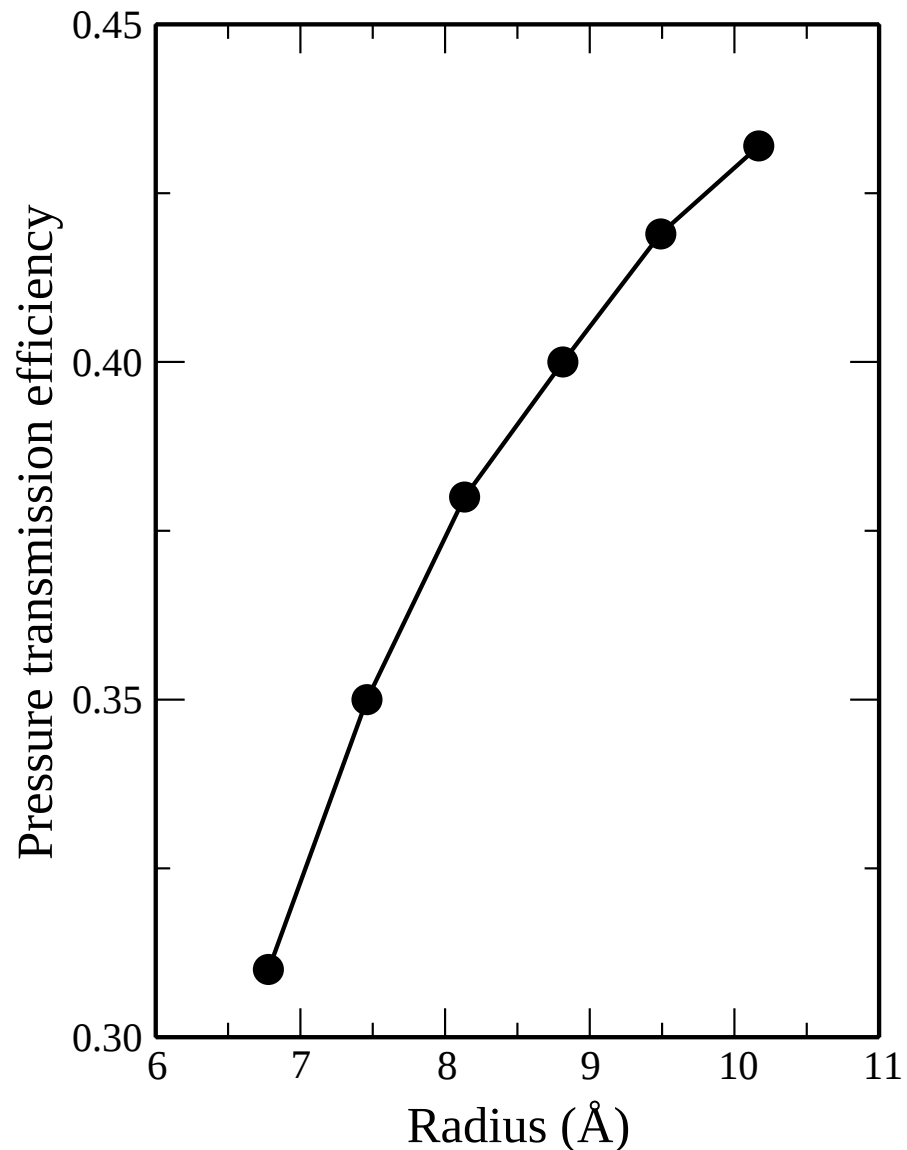
Inner and outer tubes have different behavior



Pressure transfer Efficiency from outer Inner tuber



Response pressure as a function of the external hydrostatic pressure for (5,5)@(10,10) DWCNT(left) and (5,5)@(10,10)@(15,15) TWCNT(right). The fitted relation of leftpanel relation is $y=0.31x-0.36$, we define x as the pressure transmission efficiency

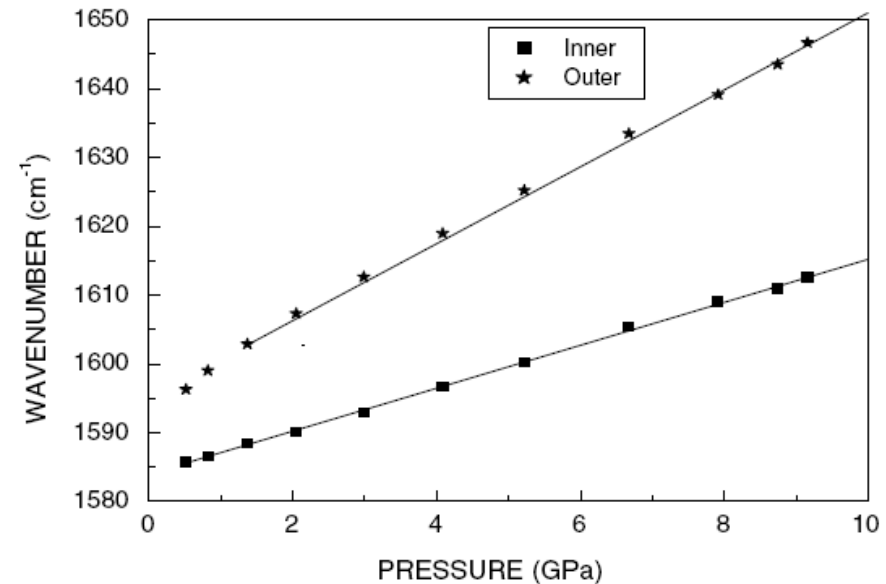
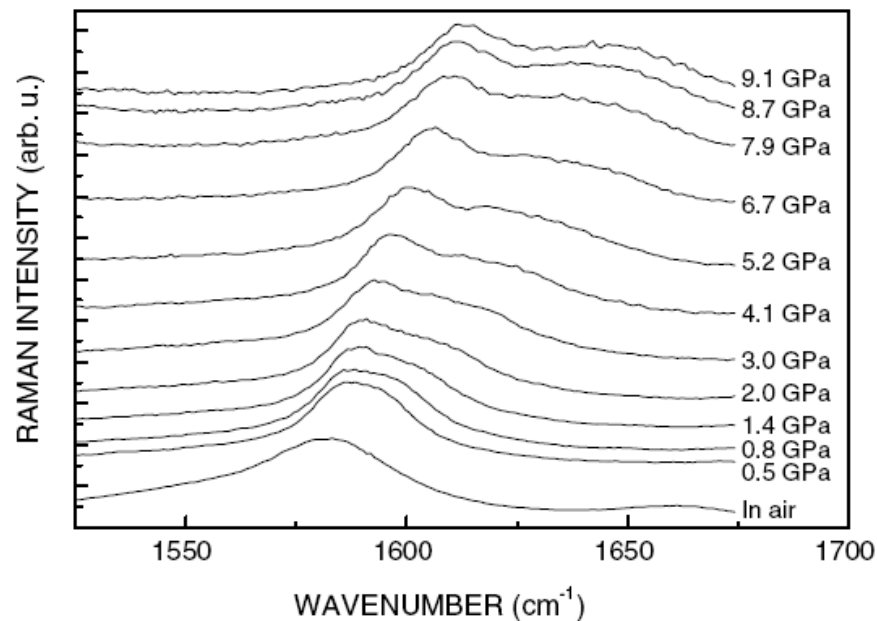


The transmission efficiency increases as tube radius increases

The pressure transmission efficiency of DWCNT as a function of the outertube radius.

Ye et al. PRB (2007)

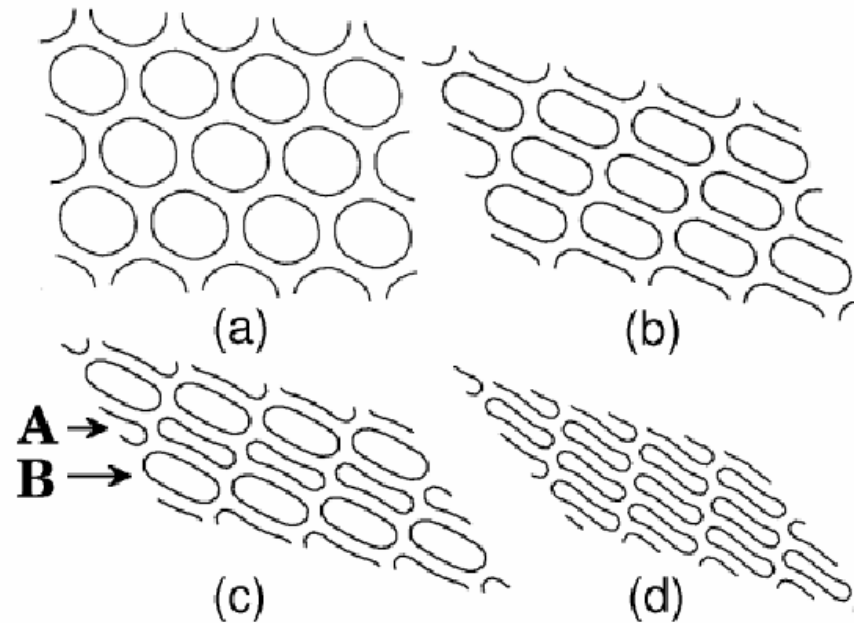
Raman spectra of double-walled carbon nanotube in the region of optical mode at different hydrostatic pressure.



The radius of the inner tube is 5 angstrom, the pressure coefficient of the inner tube is 3.11 cm⁻¹/GPa, the outer tube is 5.59 cm⁻¹/GPa. The pressure coefficient of the inner tube is 45% smaller than the outer tube's

P. Puech, H. Hubel, D. J. Dunstan, R. R. Bacsa, C. Laurent, and W. S. Bacsa, Phys. Rev. Lett. **93**, 095506 (2004).

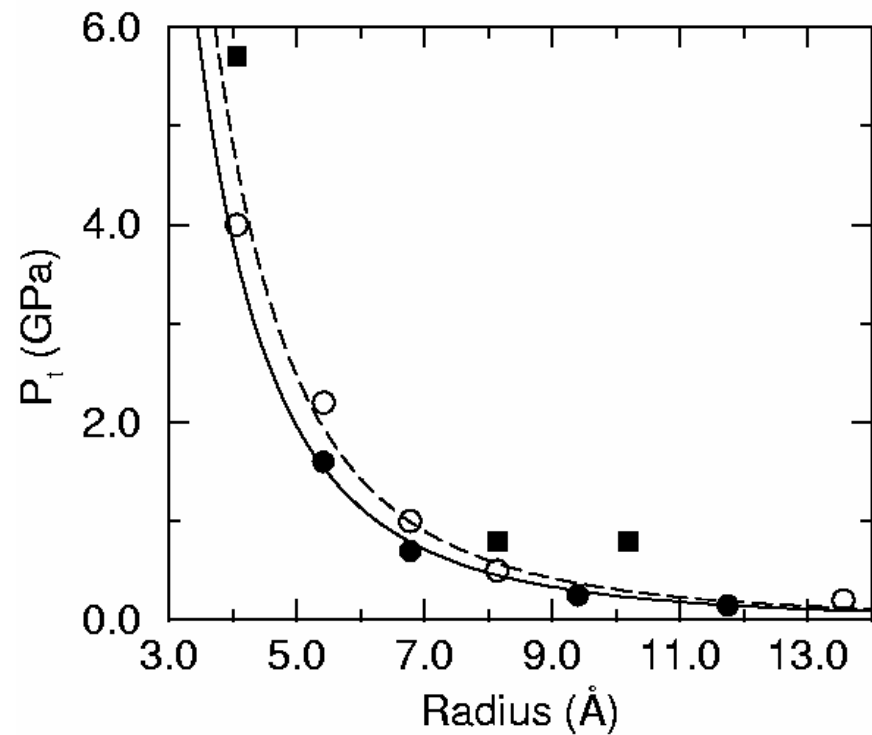
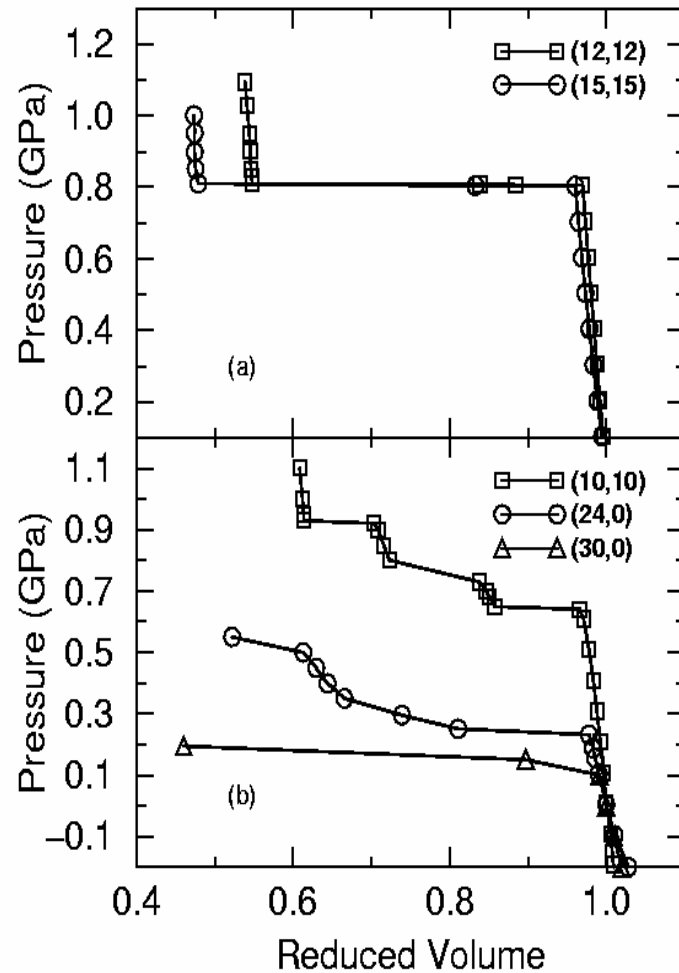
Tube lattice Under pressure:



	Energy/atom (eV)			PV /atom	enthalpy/atom
	E	E_1	E_2	(eV)	(eV)
A	-7.366	-7.330	-7.318	0.012	-7.354
B	-7.378	-7.342	-7.341	0.021	-7.357
Δ	0.012	0.012	0.023	0.009	0.003

Zhang, Liu and Gong, PRB70, 035422(2004), PRL93, 149601(2004)

Transition pressure: isolated tube and lattice



Symmetry and Radius Dependent

Summary:

- A new algorithm for constant pressure MD for finite system, which provides a window to study the properties of finite-system.
- Applications:
Si cluster has a higher transition pressure than bulk. Larger clusters, lower P_c . Pressure driven Si cluster to simple hexagonal-like structure, not β -tin!

A hard-soft transition of carbon nanotube is identified.

.....

Thank You!

Martonark, Molteni and Parrinello Proposed(PRL, 2001):

$$L = \frac{1}{2} \sum_i \mu \int d\mathbf{r} |\dot{\psi}_i(\mathbf{r})|^2 + \frac{1}{2} \sum_I M_I \dot{\mathbf{R}}_I^2 - E[\{\psi_i\}, \{\mathbf{R}_I\}] + \sum_{i,j} \Lambda_{ij} (\langle \psi_i | \psi_j \rangle - \delta_{ij})$$

$$+ \frac{1}{2} \sum_I m \dot{\mathbf{X}}_I^2 - \sum_{I,J} V_{C-L}(|\mathbf{R}_I - \mathbf{X}_J|) - \sum_{I < J} V_{L-L}(|\mathbf{X}_I - \mathbf{X}_J|).$$

N_L : No. of atom for pressure transfer medium.

Large enough !

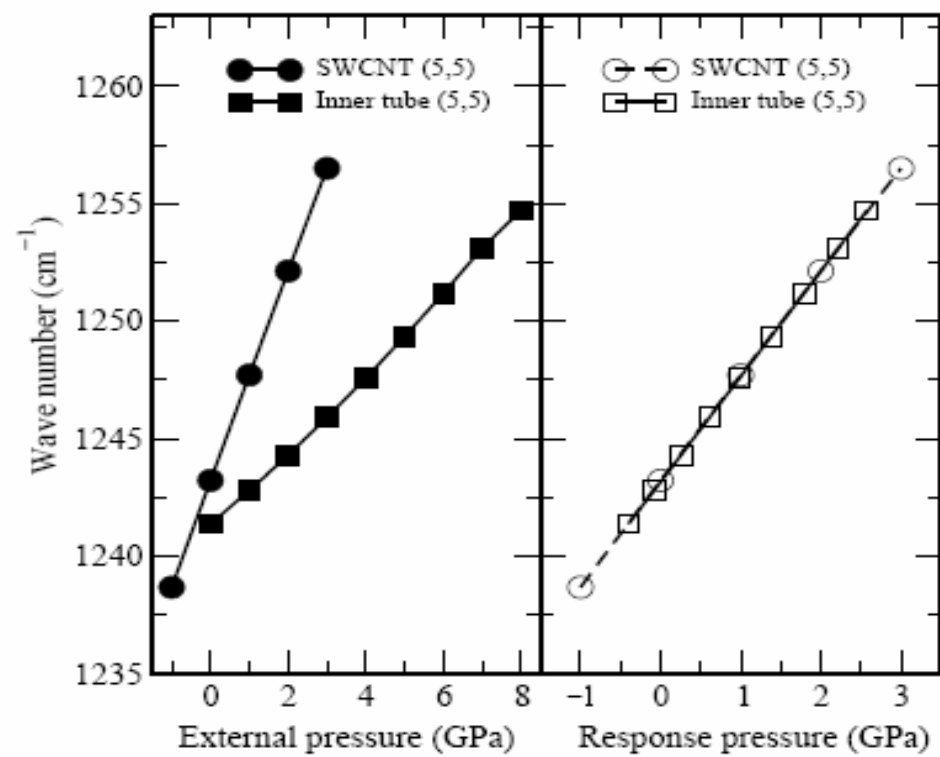
V_{C-L} : Interaction between cluster and medium.

Accurate !

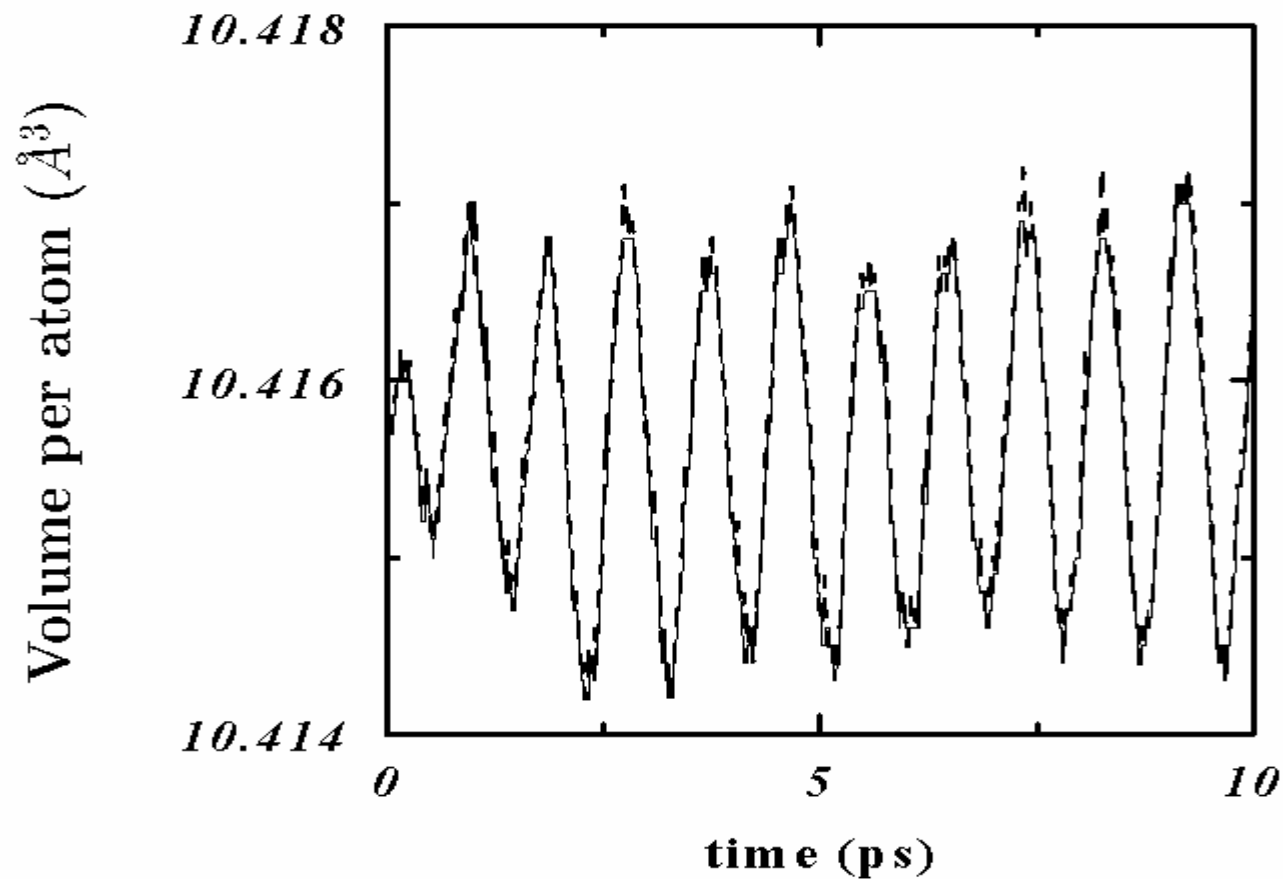
V_{L-L} : Interaction for the atoms in the medium.

Good !

simple, but large computational overhead !

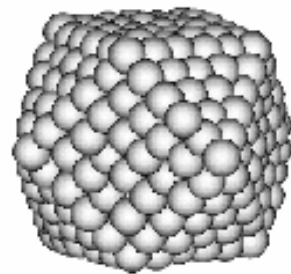


Validity of the volume decomposition:

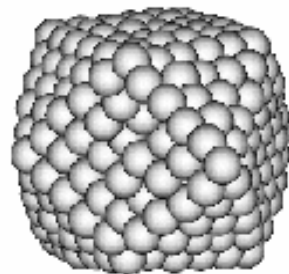


Comparison between the “exact” and atomic volume:
The agreement is good!

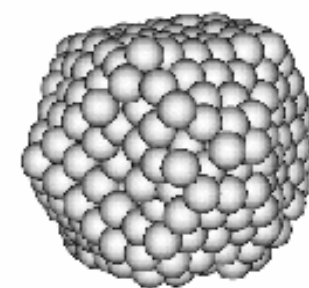
Temperature Dependence: Structure of Ni₅₆₁ without Pressure



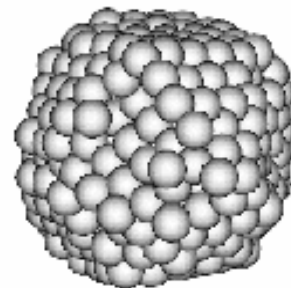
493 K



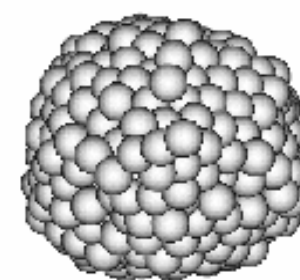
603 K



663 K

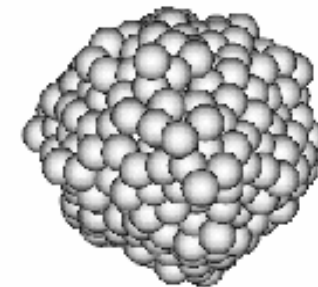


747 K



791 K

↑
 $\sim T_m$

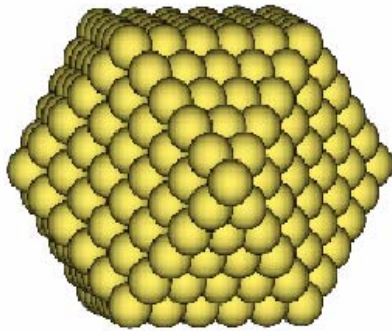


949 K

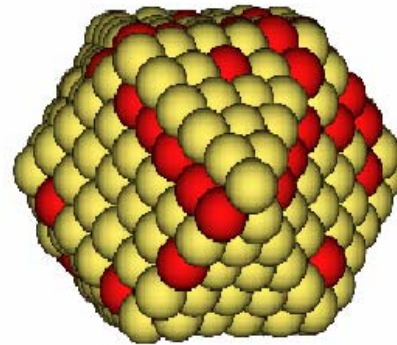
No structure transition before melting!

Temperature dependence:
structure Ni_{561} under pressure 10GPa

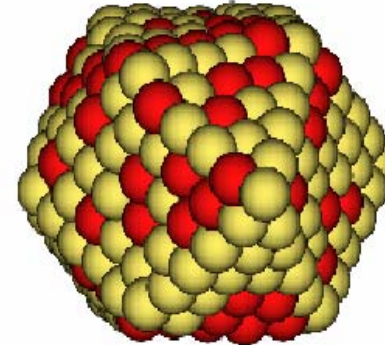
FCC-like structure to Icosahedral structure transition



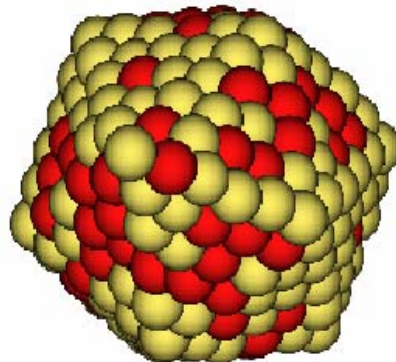
0 K



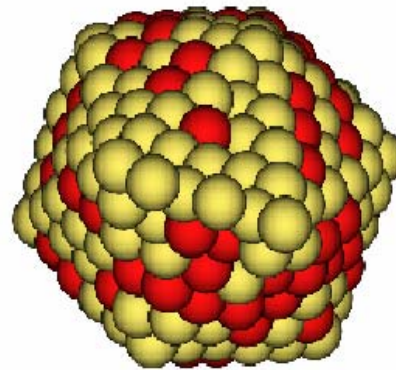
513 K



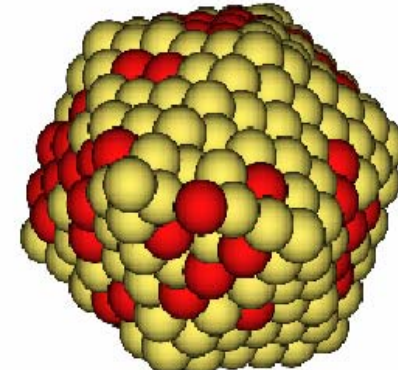
559 K



592 K

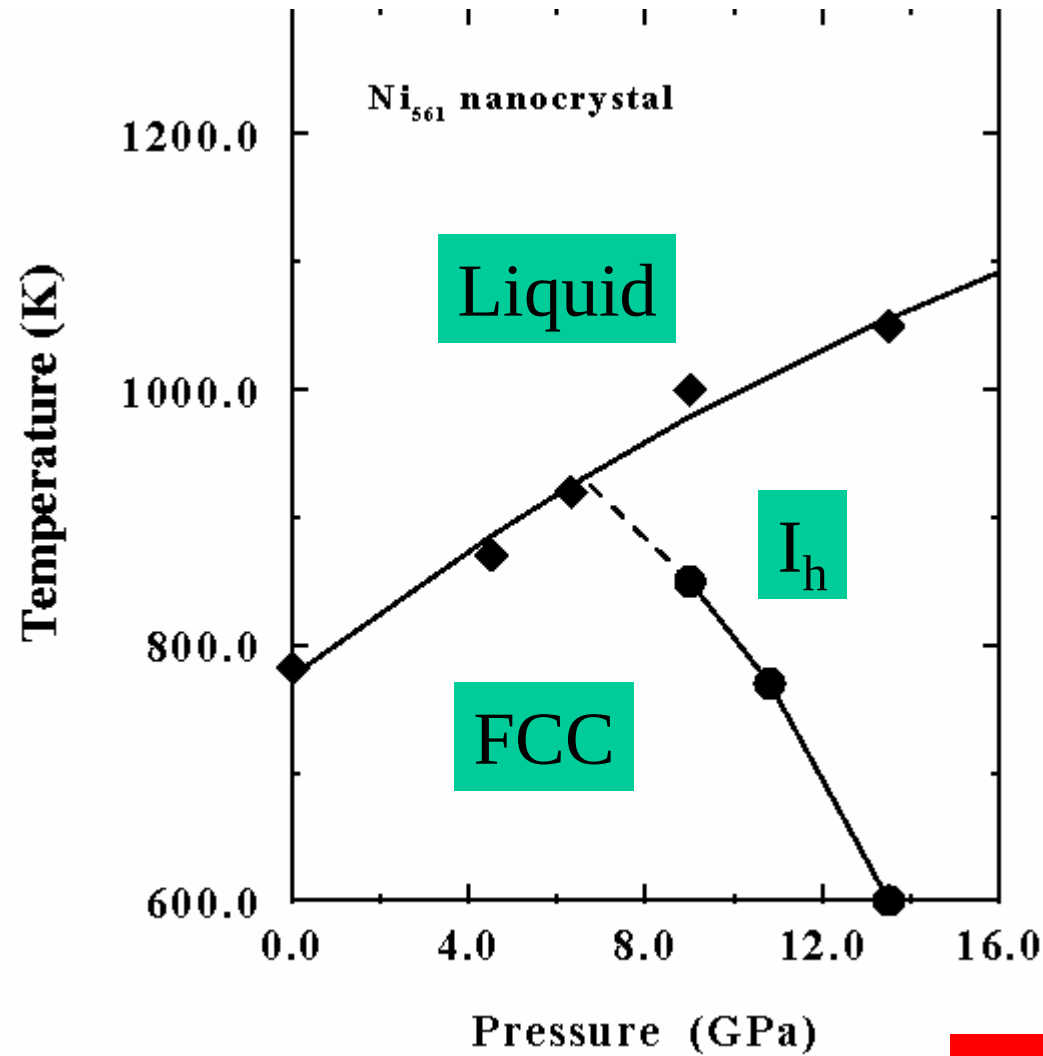


603 K

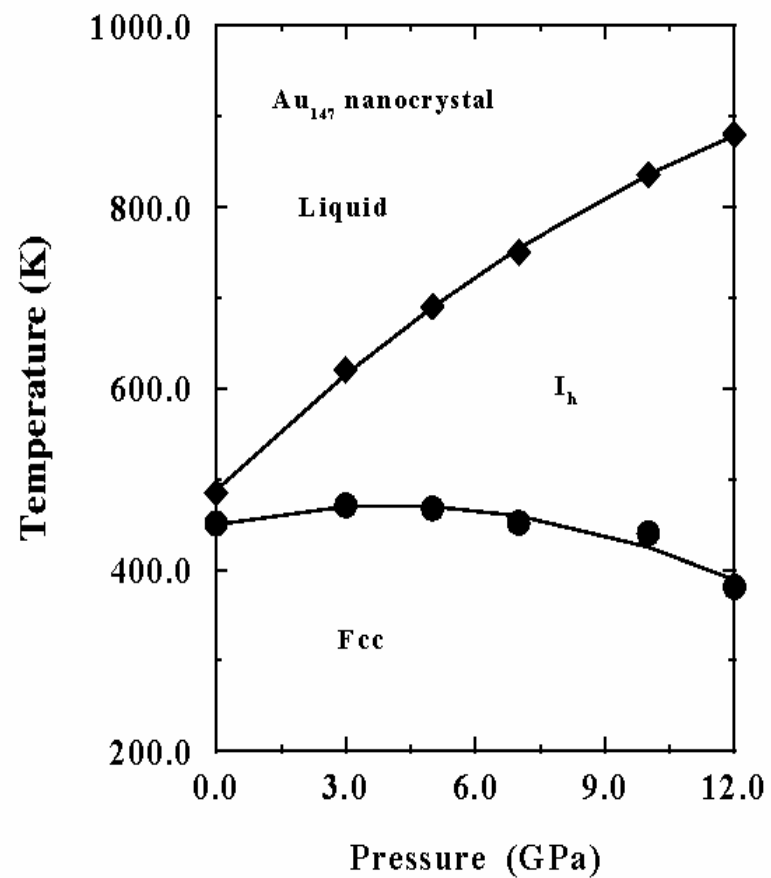
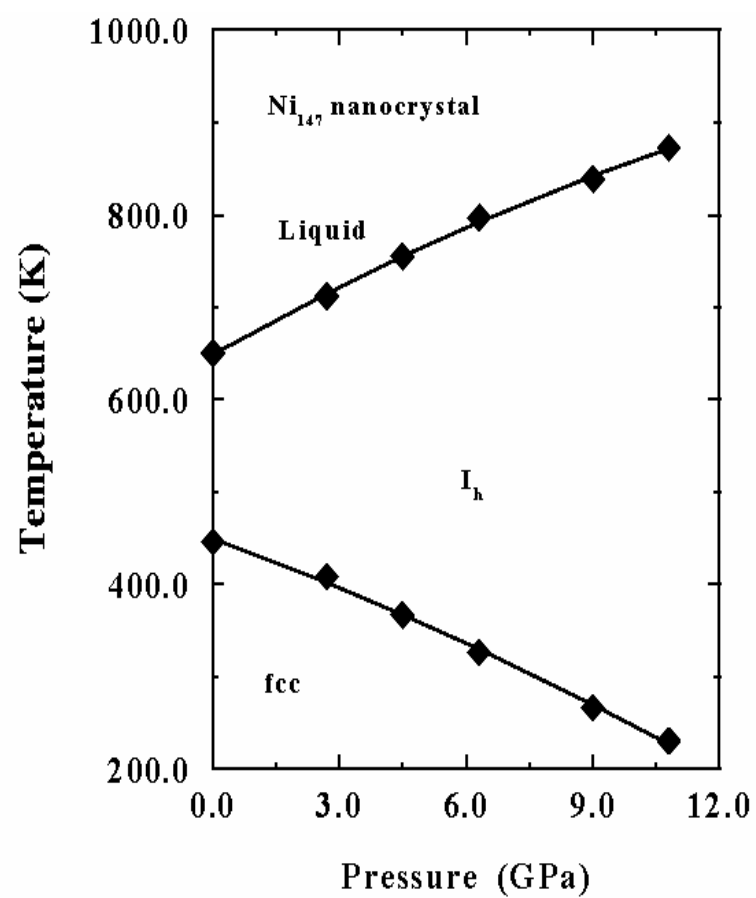


636 K

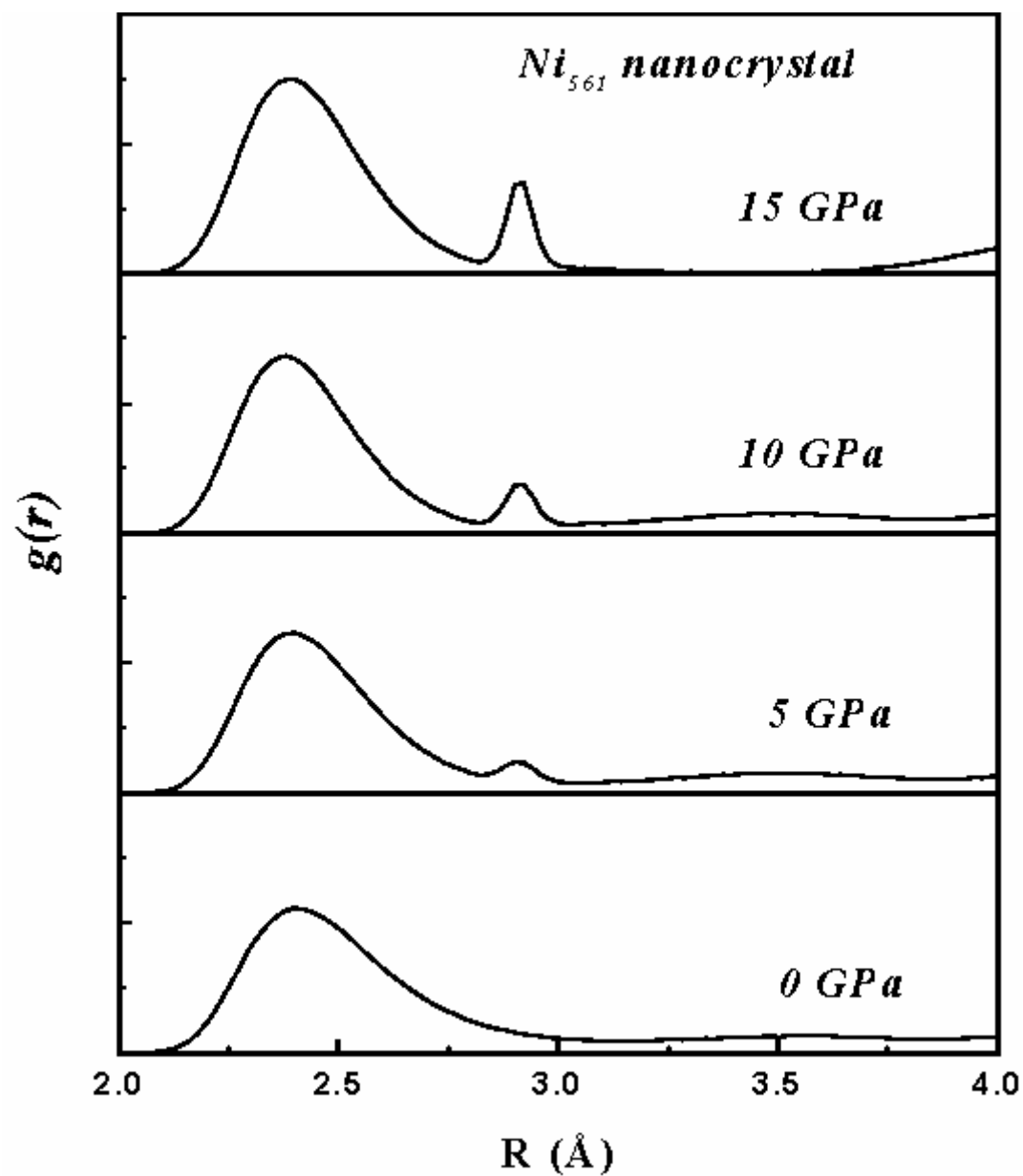
“Phase Diagram of nano-particle Ni”



Small is different!



Under pressure: bcc-Ni locally appeared.



The critical transition pressure is defined by the conditions $\frac{d^2 E}{d\omega^2}|_{\omega=1}=0$ and $\frac{dE}{d\epsilon}|_{\omega=1}=0$, which give rise to

$$P_t = \frac{3D}{R_0^3(1 + \epsilon_c)^3} \approx \frac{3D}{R_0^3},$$

$$\epsilon_c = -\frac{5D}{2CR_0^2}.$$

The ratio of bulk modulus in hard and soft phase:

$$\frac{B_s}{B_h} = \frac{19}{12} \left(\frac{h}{R_0} \right)^2.$$

$$\frac{B_s}{B_h} \propto \left(\frac{h}{R_0} \right)^2 \approx 0.01$$

- Pressure:
powerful tool to explore the meta-stability,
phase transitions of solids,
new structures, new physics, new properties...
- size effect, surface effect:
more rich physics in meta-stability!