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Missing Partner Mechanism in SO(10) Grand Unification

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**Missing Partner Mechanism in
SO(10) Grand Unification**

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SUSY SU(5) GUT

Higgs sector \longrightarrow $5, \bar{5}, 24$

$$\langle 24 \rangle = V \times \text{diag}(1, 1, 1, -3/2, -3/2)$$

$$\text{SU}(5) \rightarrow \text{SU}(3)_c \times \text{SU}(2)_L \times \text{U}(1)_Y$$

$$5 = h_u(1, 2, 1/2) + T(3, 1, -1/3)$$

$$\bar{5} = h_d(1, 2, -1/2) + \bar{T}(\bar{3}, 1, 1/3)$$

$$W \supset M_5 5 \cdot \bar{5} + \lambda 5 \cdot \bar{5} \cdot 24$$

$$m_D = M_5 - (3/2)\lambda V \quad \text{and} \quad m_T = M_5 + \lambda V$$

The possible mechanisms

- **The Missing VEV Mechanism**
- **The Sliding-singlet Mechanism**
- **The pseudo–Goldstone mechanism**
- **The Missing Partner Mechanism**

The Missing VEV Mechanism SUSY SO(10) GUT

$$3 \times 16, 10, 10', 45, 54, \dots$$

Yukawa interaction : $16 \cdot 16 \cdot 10$

$$\langle 45 \rangle = \text{diag}(a, a, a, 0, 0) \otimes i\tau_2.$$

$$W \supset M_{10} 10' 10' + \lambda 10 \cdot 10' \cdot 45,$$

The mass matrix for color-triplets is

$$(\bar{\mathbf{3}} \ \bar{\mathbf{3}}) \begin{pmatrix} 0 & \lambda_1 a \\ -\lambda_1 a & M_{10} \end{pmatrix} \begin{pmatrix} \mathbf{3} \\ \mathbf{3} \end{pmatrix},$$

The mass matrix for doublets is

$$(\bar{\mathbf{2}} \ \bar{\mathbf{2}}) \begin{pmatrix} 0 & 0 \\ 0 & M_{10} \end{pmatrix} \begin{pmatrix} \mathbf{2} \\ \mathbf{2} \end{pmatrix},$$

Realization of Missing VEV Mechanism for SUSY
SU(N) GUT see: J. L. Chkareuli, I. G. Gogoladze
and A. B. Kobakhidze, Phys. Rev. Lett. 80 (1998)
912.

The Sliding-singlet Mechanism SUSY SU(5) GUT

$$W = \bar{5} \cdot (24 + 1) \cdot 5$$

$$\langle 24 \rangle = V \text{diag}\left(-\frac{2}{3}, -\frac{2}{3}, -\frac{2}{3}, 1, 1\right)$$

$$F_5 = 0 \Rightarrow (\langle 24 \rangle - \langle 1 \rangle) \cdot \langle \bar{5} \rangle = 0$$

Since the SM Higgs doublet develops VEV

$$\langle 1 \rangle = -V \Rightarrow \langle 24 \rangle + \langle 1 \rangle = V \text{diag}\left(-\frac{5}{3}, -\frac{5}{3}, -\frac{5}{3}, 0, 0\right)$$

Unfortunately the mechanism doesn't work in SU(5) case. To make it work one needs to go SU(6) theory.

The pseudo–Goldstone mechanism SUSY SU(6) GUT

35 dos't couple to $(6 + \bar{6})$

$$\Rightarrow \text{SU}(6)_{35} \otimes \text{SU}(6)_{6,\bar{6}}$$

$$\langle 35 \rangle = V \cdot \text{diag}(1, 1, 1, 1, -2, -2)$$

$$\text{SU}(6)_{35} \rightarrow \text{SU}(4) \times \text{SU}(2) \times \text{U}(1) :$$

$$35 - (15 + 3 + 1) = 16$$

$$\text{SU}(6)_{(6,\bar{6})} \rightarrow \text{SU}(5) :$$

$$35 - 24 = 11$$

$$\text{SU}(6) \rightarrow \text{SU}(3)_C \times \text{SU}(2)_L \times \text{U}(1)_Y :$$

$$35 - (8 + 3 + 1) = 23$$

$$16 + 11 - 23 = 4 = 2 + 2$$

The Missing Partner Mechanism SUSY SU(5) GUT

Higgs sector

$$5 + \bar{5}(h + \bar{h}), \quad 50 + \bar{50}(\psi + \bar{\psi}), \quad 75(\phi)$$

$$SU(5) \rightarrow SU(3)_c \times SU(2)_L \times U(1)_Y$$

$$5 = (3, 1)_{-2} + (1, 2)_3$$

$$50 = (1, 1)_{-12} + (3, 1)_{-2} + (3, 2)_{-7} + (6, 3)_{-2} \\ + (6, 1)_8 + (8, 2)_3$$

$$75 = (1, 1)_0 + (3, 1)_{10} + (3, 2)_{-5} \dots$$

$$W \supset \phi h \bar{\psi} + \phi \bar{h} \psi + M_\psi \bar{\psi} \psi \quad \text{after} \quad \langle \phi \rangle \sim V_\phi$$

$$\begin{pmatrix} \bar{\mathbf{T}}_h & \bar{\mathbf{T}}_\psi \end{pmatrix} \begin{pmatrix} 0 & \mathbf{V}_\phi \\ V_\phi & \mathbf{M}_\psi \end{pmatrix} \begin{pmatrix} \mathbf{T}_h \\ \mathbf{T}_\psi \end{pmatrix},$$

At this level the doublets from $5 + \bar{5}$ are massless since there are no doublets in the $50 + \bar{50}$.

SUSY SO(10) GUT

The lowest dimensional Higgs representations which has a missing doublet in SO(10) is the $126 + \overline{126}$

SU(5) \times U(1) decomposition

$$126 = 1_{-10} + \bar{5}_{-2} + 10_{-6} + \overline{15}_6 + 45_2 + \overline{50}_{-2}$$

$$16 \cdot 16 \cdot 10 \Rightarrow 10 \cdot 126 \cdot (210 + 1050)$$

$$210 = 1_0 + 24_0 + 75_0 + 5_{-8} + \bar{5}_8 \dots$$

$$120 = 5_2 + \bar{5}_{-2} + 10_{-6} + \overline{10}_6 + 45_2 + \overline{45}_{-2}$$

$$210 \cdot 126 \cdot (10 + 120) + 210 \cdot \overline{126} \cdot (10 + 120) \\ + M_\Delta 126 \cdot \overline{126}$$

$$\langle 126 \rangle = \langle \overline{126} \rangle = 0 \text{ and } \langle 210 \rangle \neq 0$$

$$126 + \overline{126} \supset 2 \times (2 + 2) + 3 \times (3 + \bar{3})$$

$$120 \supset 2 \times (2 + 2) + 2 \times (3 + \bar{3})$$

$$10 \supset 1 \times (2 + 2) + 1 \times (3 + \bar{3})$$

One needs to forbid the couplings

$$M_1 120 \cdot 120 + M_2 10 \cdot 10 + 120 \cdot 10 \cdot 210$$

Explicit SO(10) Models

It is much more convenient if the 126-plets involved in the DT splitting have no VEVs.

For the rank breaking we can use either a scalar $16 + \overline{16}$ -plets or another $126 + \overline{126}$ -plets.

$$(a) : C = 16 \quad \overline{C} = \overline{16}$$

$$(b) : C = 126 \quad \overline{C} = \overline{126}$$

The goal is to have all order DT splitting. We introduce additional $U(1)_A$ symmetry. $U(1)$ charges of the superfield are

$$H(10) : 1, \quad \Sigma(120) : 1, \quad \Delta(126) : -1,$$

$$\overline{\Delta}(\overline{126}) : -1, \quad \Phi(210) : 0, \quad 16_i : -1/2$$

$$C(16) : 3/2, \quad \overline{C}(\overline{16}) : -3/2$$

$$C(126) : 3, \quad \overline{C}(\overline{126}) : -3$$

$$X : 2$$

$$W_{DT} = \Phi \Delta (H + \Sigma) + \Phi \overline{\Delta} (H + \Sigma) + X \overline{\Delta} \Delta$$

The couplings important for the symmetry breaking are

$$W(\Phi, C) = \frac{\lambda}{3}\Phi^3 + \frac{M_\Phi}{2}\Phi^2 + \bar{C}C(M_C + \sigma\Phi)$$

$$\langle \Phi \rangle \simeq \langle C \rangle \simeq \bar{C} = M_{GUT}$$

Having $U(1)_A$ symmetry we can forbid $H^2, \Sigma^2, \Phi H \Sigma$ coupling and it also guarantees that $\langle \Delta \rangle = 0$. Because operators such as $\bar{\Delta}\Delta(H^2 + \Sigma^2)$ will destroy DT splitting.

Yukawa Sector

$$\sum_{k=0} \left(\frac{\Phi}{M}\right)^k 16_i 16_j 10 + \sum_{k=0} \left(\frac{\Phi}{M}\right)^k 16_i 16_j 120$$

only renormalizable couplings $16 \cdot 16 \cdot 10$ and $16 \cdot 16 \cdot 120$ do not give desirable fermion mass pattern.

Majorana masses for the right handed neutrinos

$$\frac{X^2}{M_*^3} 16_i 16_j \bar{C} \bar{C}$$

such couplings can be generated from renormalizable interactions.

Conclusion

we have proposed a new solution of the DT splitting problem within SO(10) GUT via a missing partner mechanism. For this mechanism to be realized through renormalizable superpotential couplings we have considered the scalar superfield content $10 + 120 + 126 + \overline{126} + 210$ and the SO(10) rank breaking states C, \bar{C} . For the latter, two possibilities (a): $C = 16, \bar{C} = \overline{16}$ and (b): $C = 126, \bar{C} = \overline{126}$ can be considered with equal success.

- Gauge coupling perturbativity above the GUT scale.
- Proton stability
- more detailed study with accurate calculation of the mass spectrum is needed