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Turbulence spreading and nonlocal transport in magnetized plasmas

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Turbulence Spreading and Nonlocal Transport in Magnetized Plasmas

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JET-EFDA Contributors

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- Definition of turbulence spreading, TS →
- Motivation, how can it influence the turbulent transport
- 1D transport model accounting for turbulence spreading
- Comparison with a “simple” turbulence simulation: 2D interchange model
- Applying the TS transport model to describe perturbative transport experiment in JET
- Heat modulation and fast heat pulse propagation.



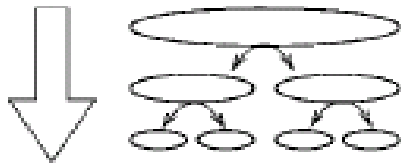


Turbulence Spreading

The turbulence itself is a transported quantity: may spread from unstable regions of generation into stable regions.

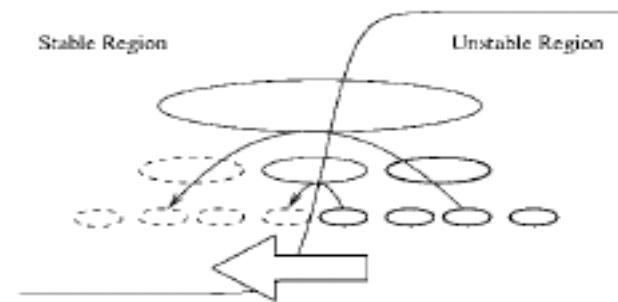
Plasma Turbulence: Garbet *et al*, NF **34**, 963 (1994); Gürçan, Diamond, Hahm, PoP, **13**, 052306 (2006); **14**, 055902 (2007)

Richardson Cascade

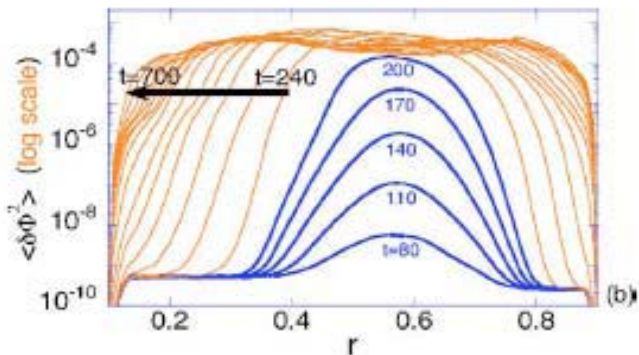


Spreading in k-space, from "unstable" k to "stable" k

Turbulence Spreading



Spreading in configuration space from unstable region into stable region



Gyrokinetic simulations

Wang *et al*. PoP **14**, 072306 (2007)

Turbulent spreading well-known in fluid turbulence: turbulent overshoot, penetration



Quasi-linear approach and mixing length theory [Kadomtsev 1965]: **Balance the linear growth of instability, γ , against D , the “*turbulent diffusion*”**

$$D \propto \frac{\gamma}{k^2}$$

k is a perpendicular wave number of the turbulence.

More refined versions look at turbulent spectra and include off-diagonal terms (f.x. Weiland transport model).

Successful approach, but drawbacks:

Local transport models do not account for:

- high transport at low gradients
- Up-gradient transport
- burstiness
- avalanches
- fast transport events

All observed effects!



- Turbulent spreading accounts for nonlocal effects
- Include naturally up-gradient transport
- May account for avalanche and bursty transport
- Non-diffusive effects: breaking direct relationship between gradients and fluxes: Ficks law

Note: Turbulent spreading may work in basically two ways:

1. The turbulence spreading takes the stability boundary with it implying enhanced turbulence and transport in the “new” region
2. The turbulence penetrates into a stable region where the modes are damped, implying different transport characteristics as, e.g., up-gradient transport.



**1D turbulent transport model of heat in the core (density profile fixed), accounting for spreading of turbulence into stable regions:
Describe pinch effects and profile stiffness.**

$$\left. \begin{aligned} \frac{\partial E}{\partial t} &= \frac{1}{r} \frac{\partial}{\partial r} r \left[D_0 E \frac{\partial}{\partial r} E \right] + \gamma E - (\gamma_0 + \beta E^2) E, \Rightarrow \text{Turbulent energy, growth and spreading } D_0 \text{ const.} \\ \frac{\partial T}{\partial t} &= -\frac{1}{r} \frac{\partial}{\partial r} r q + \chi_0 \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} T + S(r) \quad \Rightarrow \text{Temperature transport} \end{aligned} \right\}$$

Growth rate: $\gamma = \lambda [\kappa_T - \kappa_c]$; λ : free parameter; $\kappa_T \equiv |\partial_r T|/T$, κ_c : critical.

The heat flux: $q = \langle \tilde{T} v_r \rangle$, $E \approx \langle v_r^2 \rangle$. cross coherence $\xi \propto \gamma$ between \tilde{T} and v_r :

$$\Rightarrow q = \xi \sqrt{\langle \tilde{T}^2 \rangle \langle v_r^2 \rangle}, \tilde{T}/T = C\sqrt{E}, \text{ i.e., and } \langle \tilde{T}^2 \rangle = C^2 \langle E \rangle \langle T^2 \rangle,$$

$$q = C\gamma ET = C\lambda ET [\kappa_T - \kappa_c]. \quad q < 0 \text{ for } \kappa < \kappa_c: \text{ Up-gradient transport!}$$



Similar transport model applies to the transport of density:

$$\frac{\partial n}{\partial t} = -\frac{\partial}{\partial x}\Gamma + S(x, y) \quad \Gamma = \langle \tilde{n}v_r \rangle$$

Apply a simple 2D interchange model for testing the TS model:

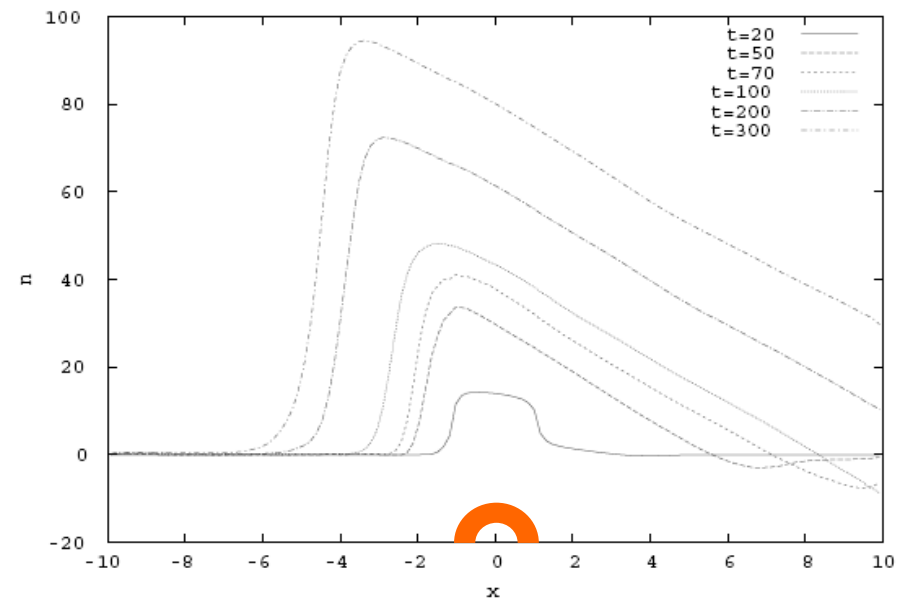
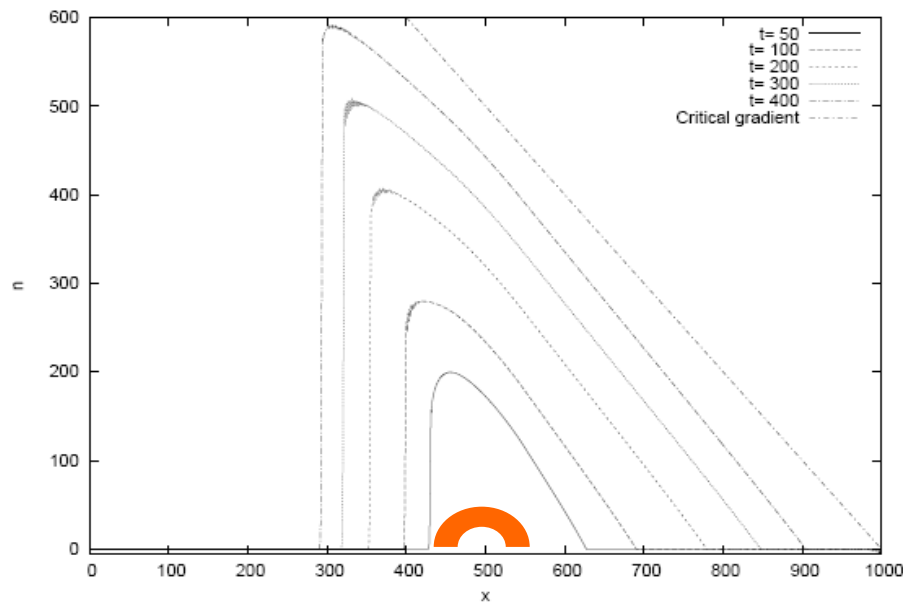
$$\partial_t \Omega + \vec{v} \cdot \nabla \Omega + \mathcal{K}(n) = \mu_\omega \nabla^2 \Omega$$

$$\partial_t n + \vec{v} \cdot \nabla n - \underline{n\mathcal{K}(\phi) + \mathcal{K}(n)} = \mu_n \nabla^2 n + S(n)$$

Ω is the vorticity ($\Omega = \nabla^2 \phi$) perpendicular to the (x, y) -plane

Contains **curvature** as drive, density source and exhibits threshold for instability.

Naulin et al Phys Plasma **12**, 122306 (2005)



Evolution of profiles in 2d simulation (left), and in TS model (right), location of source indicated

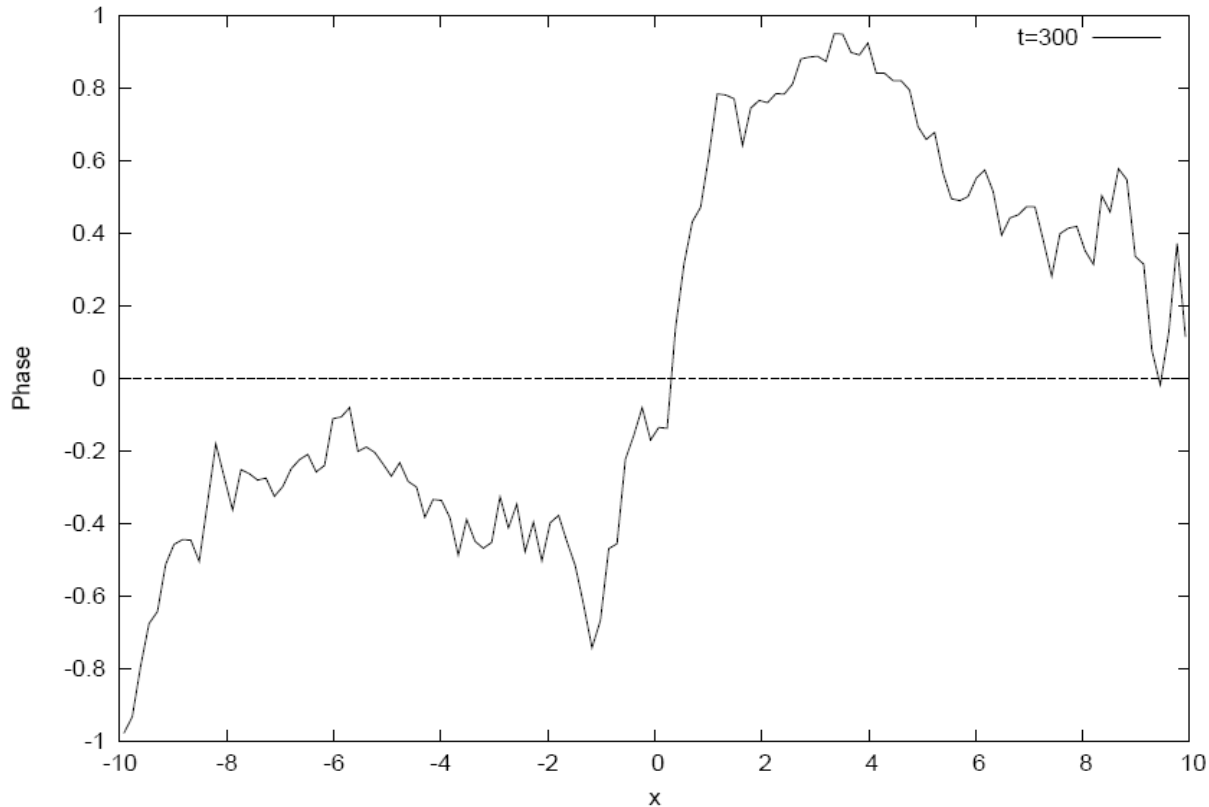
The profile evolves toward a critical profile “marginally stable”



Where does the energy come from?

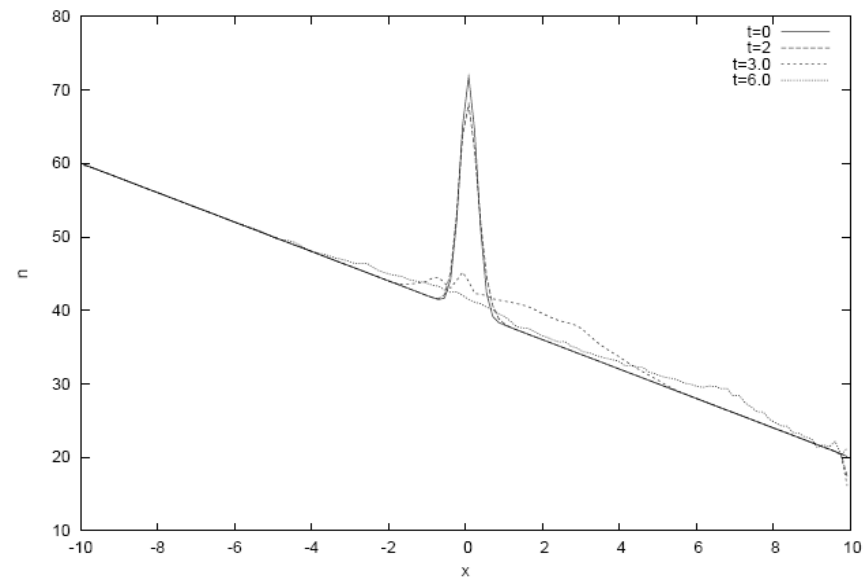
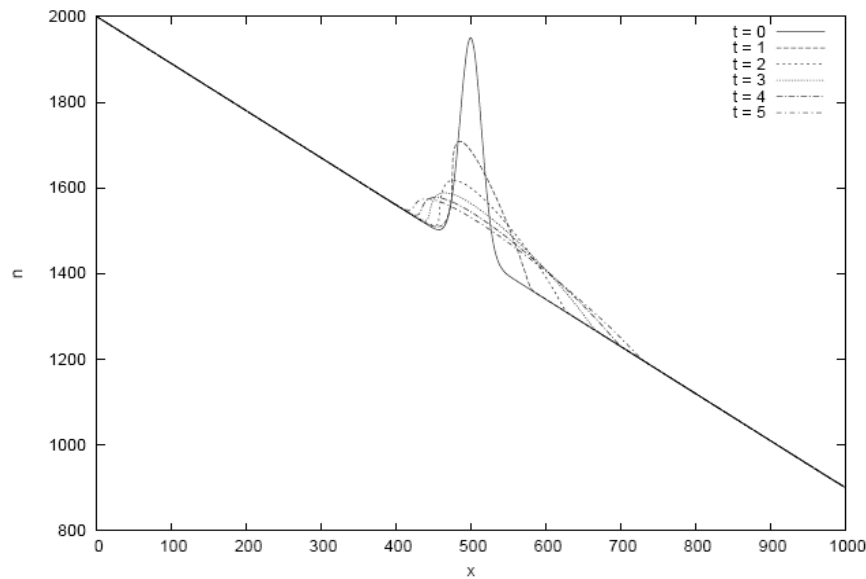
- Gradient is a drive for turbulence.
- Energetically the energy input into turbulence is connected with the flux.
- Turbulence is damped in stable regions.
- Turbulent energy is exhausted not in dissipative effects, but in reversal of transport direction.
- **Role of stable modes!!!!** (which are completely ignored in traditional models)

Stable modes show different phase relationship between turbulence and the transported quantity

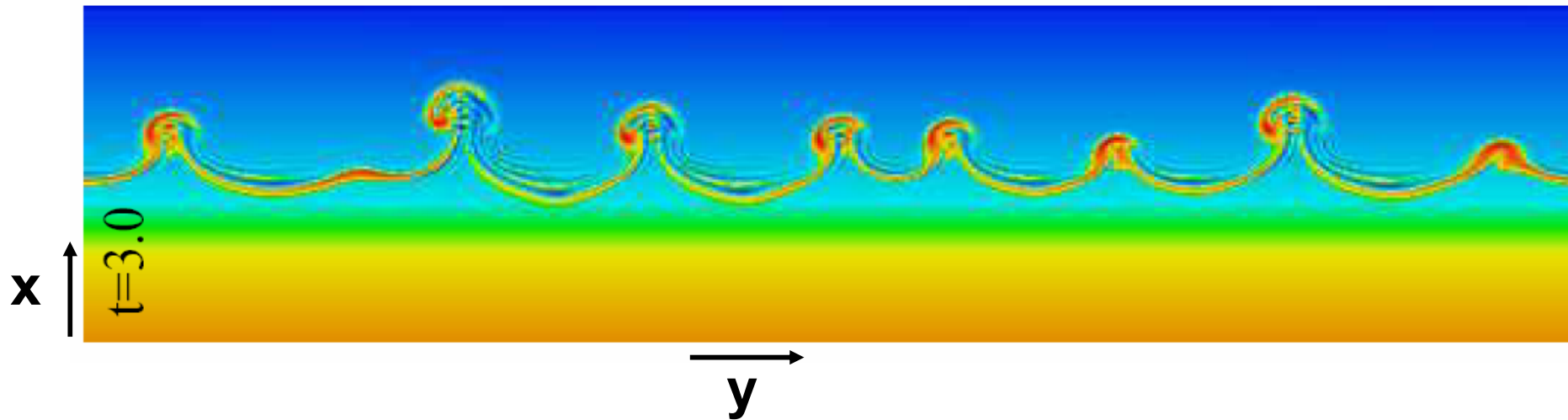


Average phase between density and radial velocity fluctuations

Note: sign of $\Gamma(= \langle \tilde{n}v_r \rangle)$ is determined by sign of phase between n and v_r



Density pulse initialized on a subcritical background gradient.
Model (left) and 2D simulation (right). Up-gradient transport.



Initialize pulse on a subcritical background gradient. Instability develops into mushroom like structures, clearly neither local nor diffusive transport...

However, the mean profile appears to be well described by the 1D transport model!



$$\frac{\partial E}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} r \left[D_0 E \frac{\partial}{\partial r} E \right] + \gamma E - (\gamma_0 + \beta E^2) E,$$
$$\frac{\partial T}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} r q + \chi_0 \frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} T + S(r)$$

$D_0 = 0$ Critical Gradient
Transport model.

$$\longrightarrow q_t = C \frac{\lambda^{3/2}}{\beta^{1/2}} T [\kappa_T - \kappa_c]^{3/2} H [\kappa_T - \kappa_c] - \chi_0 \partial_r T$$

The model then reduce to the standard **critical gradient model**, **CGM**, widely used in describing perturbative transport experiments

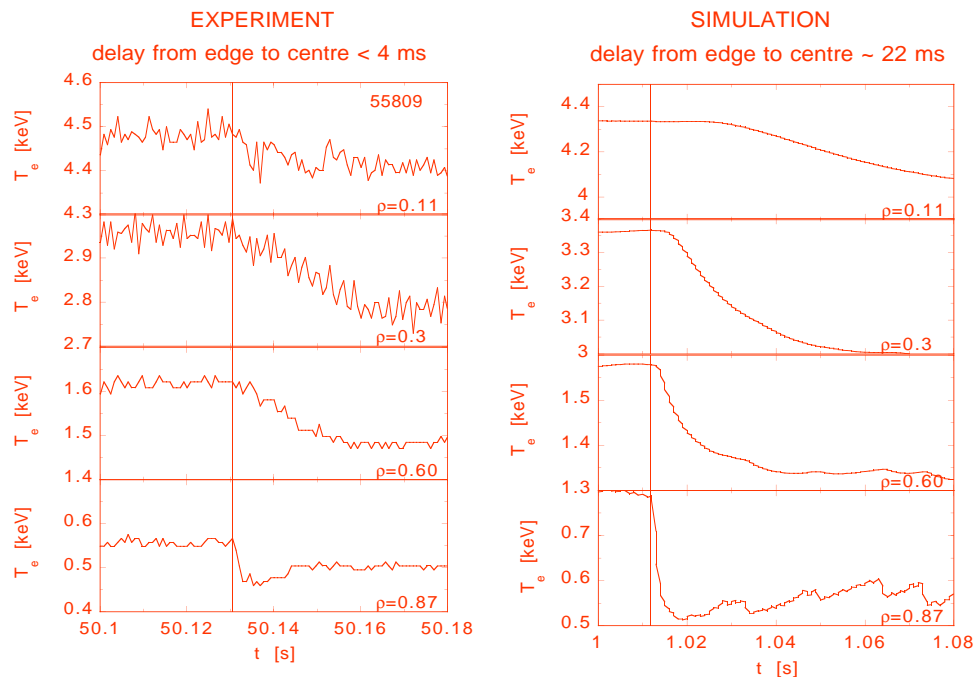
Imbeaux *et al.* PPCF **43**, 1503 (2001)

Mantica & Ryter, C.R. Physique **7**, 634 (2006)



Application to transient transport events in JET: Fast propagation of cold pulse

P. Mantica et al., (Proc. 19th IAEA Conf. Lyon, 2002) EX/P1-04, IAEA 2002

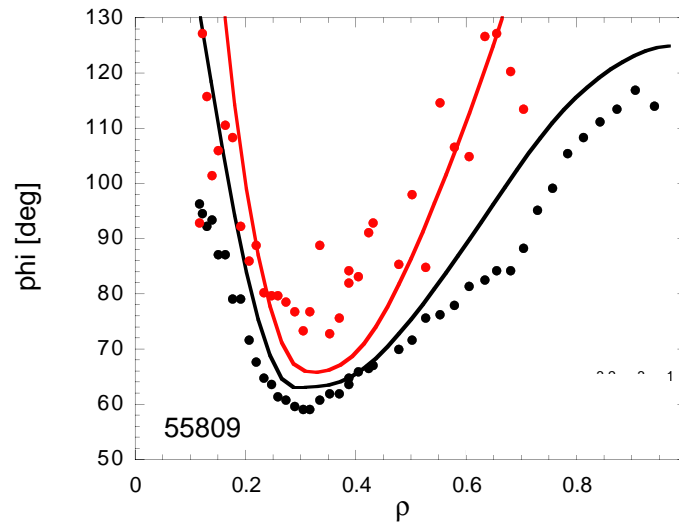
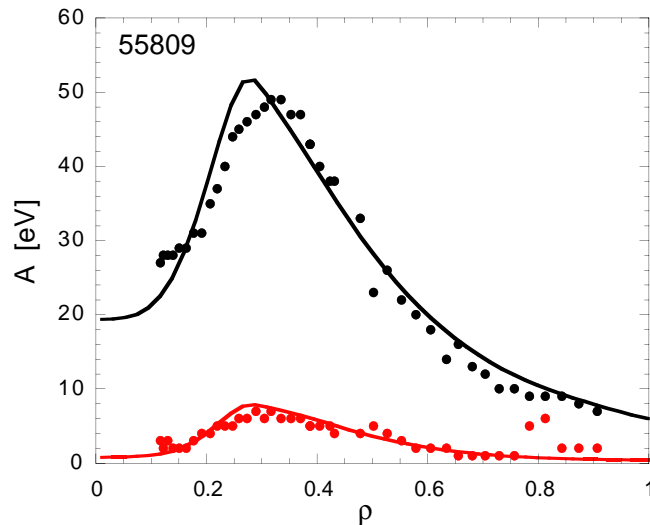


Cold pulse experiment, JET # 55809; CGM simulation with coefficients fitting heat modulation experiment, **too slow for cold pulse.**

Transient cold pulse, initiated by local cooling at the edge

The pulse propagates much faster than heat modulation, which is well described by a “standard” critical gradient model, CGM.

Challenge: explain both effects within the same model!



T_e modulation by off-axis ICH

Compared with CGM

For $\rho > 0.3$ unstable profile, for $\rho < 0.3$ stable profile

Compared with cold pulse propagation:

Surprising result:

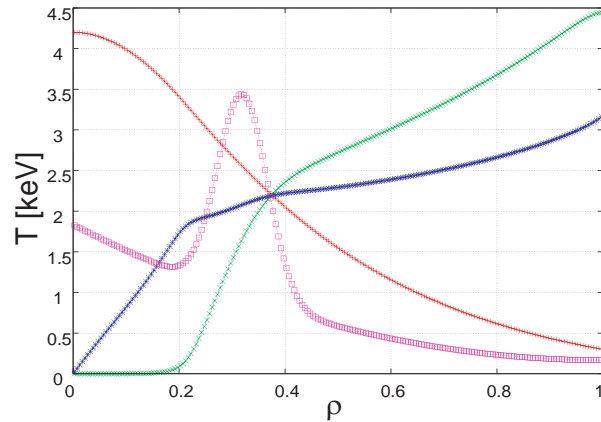
$\rho > 0.3$: the plasma is stiff and both perturbations propagate fast

$\rho < 0.3$: the plasma is below threshold and heat wave slows down and is damped **BUT cold pulse still travels fast!**

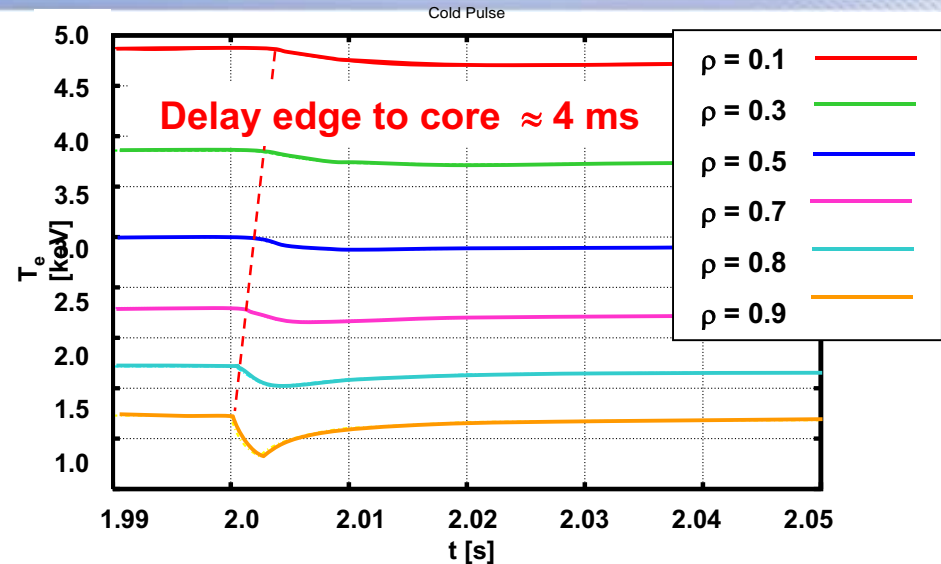
Incompatible with local transport models!



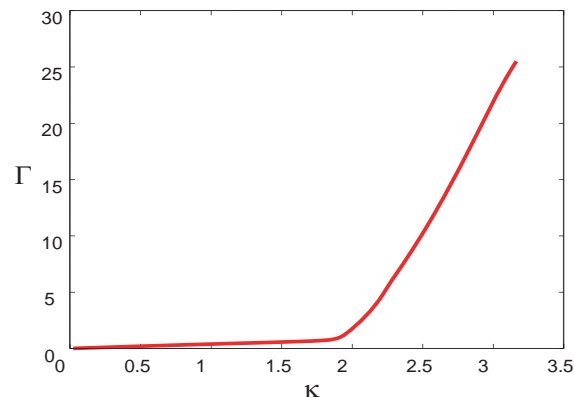
Perturbative Transport I



Profiles: T , E , κ_T , $S(r)$
off-axis ICH + NBI



Cold pulse propagation

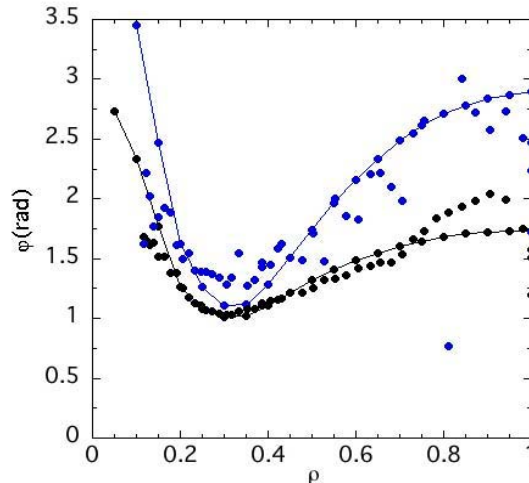
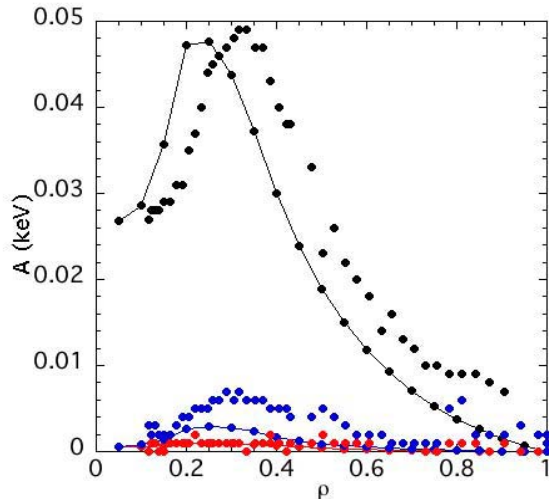


Stiffness: $\Gamma = q_t / T$
($q_t = q - \chi_0 \nabla T$) vs κ_T

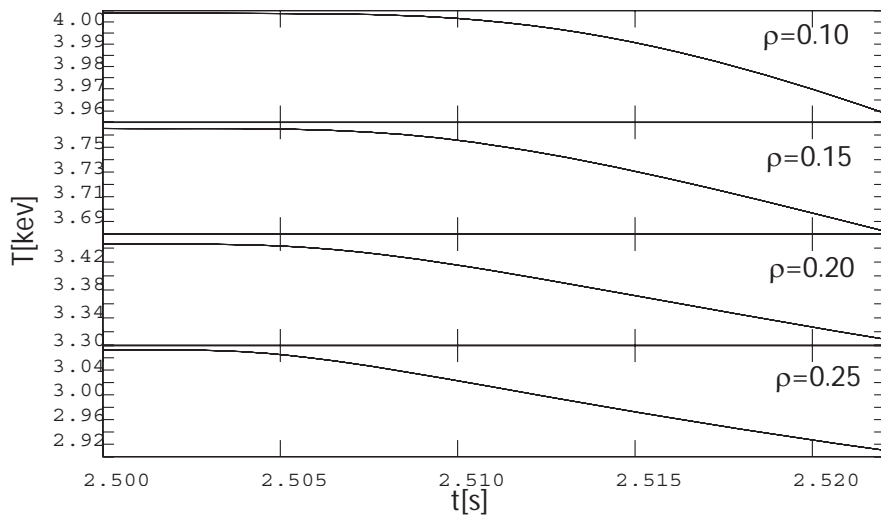
Fast cold pulse propagation easily modelled for a wide parameter range; even pulse reversal (increasing centre temperature response) have been modelled.

Challenge: model both heat modulation experiment and cold pulse propagation!

Rasmussen *et al.*, EPS 2006 Rome



The heat modulation comparing JETpulse 55809; numerical results: linespoints. Amplitude and phase: Fundamental (black), second (red) and third harmonic (blue)



Cold pulse evolution: < 18 ms

A step in the right direction!

But room for improvements:

Account for more than one transport channel

Investigate influence of ITB

Integrate TS model with real dispersion relation solver



- Turbulent spreading is important in many systems – in plasmas often observed, also in turbulence simulations, incl. gyro-kinetic simulations
- Turbulence spreading models allows for introduction of “non-locality” into transport models
- It accounts directly for pinch effects – up-gradient transport
- Results of the TS model compare well with turbulence simulations – at least for interchange model
- A first step to include more complicated transport processes into “simple” models