



**The Abdus Salam
International Centre for Theoretical Physics**



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**Summer School on Novel Quantum Phases and Non-Equilibrium
Phenomena in Cold Atomic Gases**

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From few-body to many-body physics in cold Fermi gases - Part II

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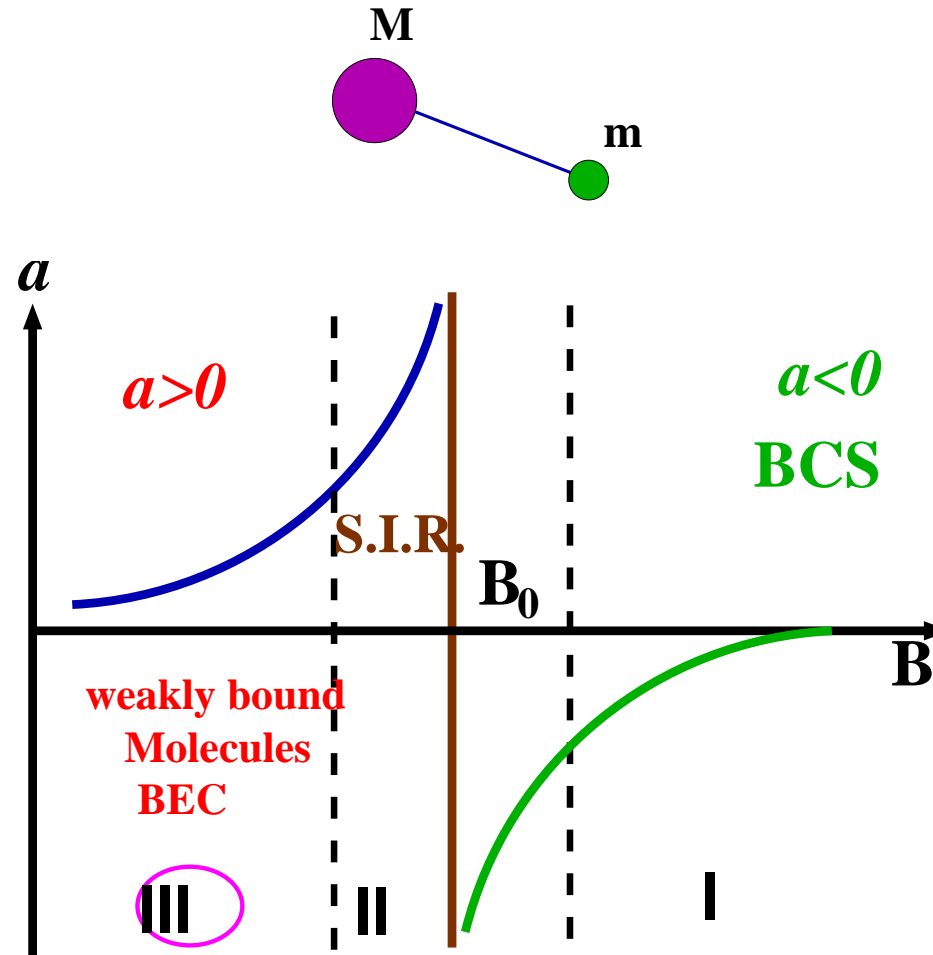
Lecture 2. Molecular regimes in Fermi-Fermi mixtures

Contents

- Introduction.
- Molecular regime. Interaction between molecules
- Collisional relaxation
- Pauli principle and effect of the mass ratio
- Crystalline phase and quantum transitions
- Conclusions

Mixtures of Fermi gases

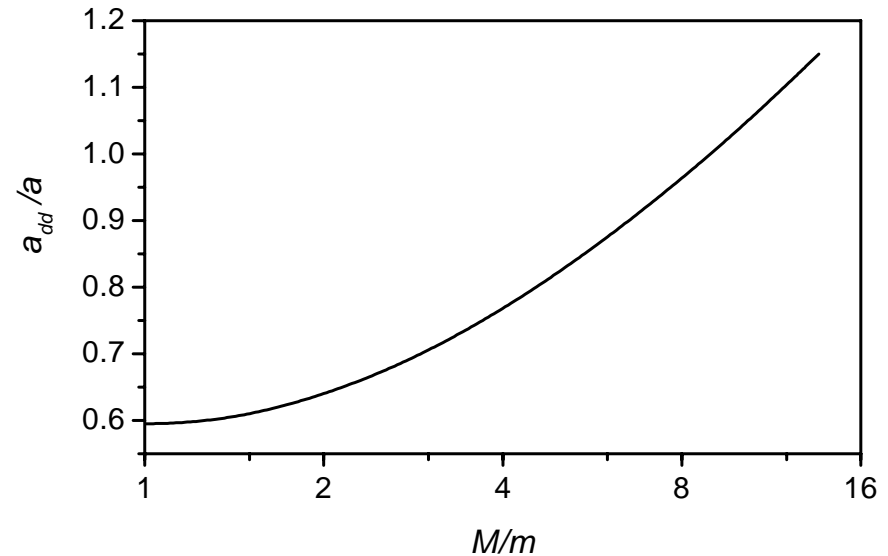
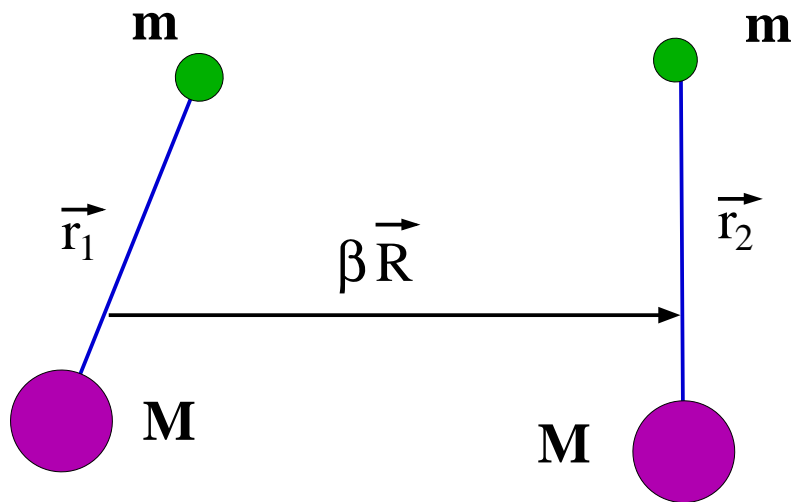
Molecules of different fermionic atoms ${}^6\text{Li}{}^{40}\text{K}$ ${}^6\text{Li}{}^{87}\text{Sr}$



What happens with collisional stability and molecular BEC?
Is there something else interesting ?

Molecule-molecule interaction

Interaction between the molecules (a_{dd}) Petrov et al 2005



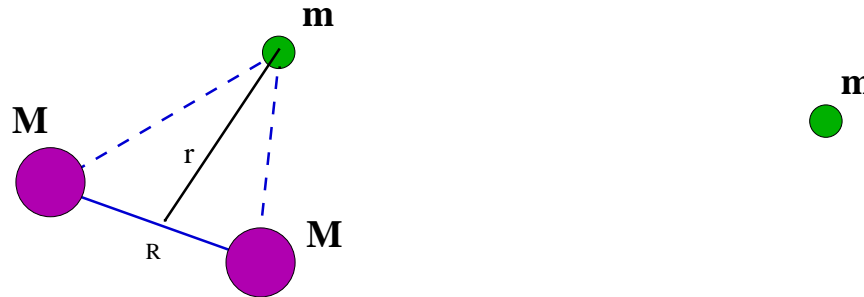
$$\beta = \sqrt{2Mm}/(M + m)$$

Nothing dramatic happens, but why I stop at $M/m = 13.6$?

Collisional relaxation. Mediated potential

Exact solution for the dependence on a and M/m

$M \gg m \rightarrow$ Born-Oppenheimer picture



$r \ll a \rightarrow$ One bound state of a light atom with two fixed heavy ones

$$\psi(\vec{r}) \propto \left(\frac{\exp(-\lambda|\vec{r}-\vec{R}/2|/R)}{|\vec{r}-\vec{R}/2|} + \frac{\exp(-\lambda|\vec{r}+\vec{R}/2|/R)}{|\vec{r}+\vec{R}/2|} \right)$$

Bethe-Peierls boundary condition

$$|\vec{r} \pm \vec{R}/2| \rightarrow 0 \quad \psi \propto (1 - a/|\vec{r} \pm \vec{R}/2|)$$

$$\lambda \approx 0.567 \quad \varepsilon(R) = U(R) = -\hbar^2 \lambda^2 / 2mR^2$$

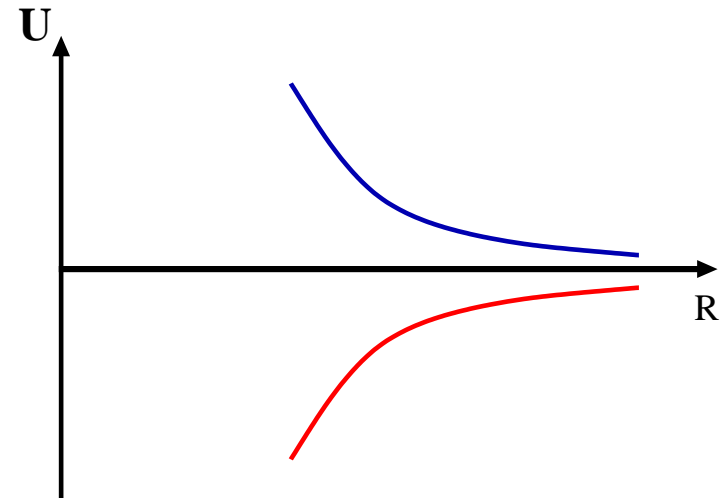
Mediated attractive potential $U(R) \approx -0.16\hbar^2/mR^2$

Effective potential

Pauli principle \Rightarrow Centrifugal potential $U_c = 2\hbar^2 / MR^2$

Mediated attraction competes
with Pauli principle

$$\begin{aligned} U_{eff}(R) &= U(R) + U_c(R) \\ &= -0.16\hbar^2 / mR^2 + 2\hbar^2 / MR^2 \end{aligned}$$

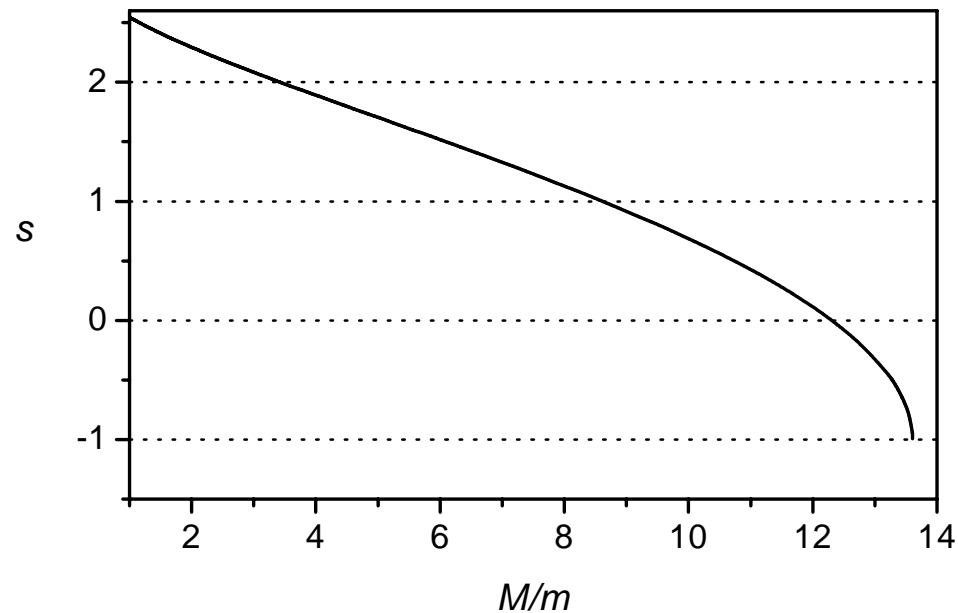


$M/m > 13.6 \rightarrow$ **fall into center** short-range physics

Many nodes of the wavefunction. Many (trimer) bound states

$$M/m < 13.6 \rightarrow \alpha_{rel} \sim a^{-s}$$

Relaxation rate



$$\alpha_{rel} \sim a^{-s}$$

$M/m > 12.33 \rightarrow U_{eff} > 0, \rightarrow s > 0$

$M/m > 12.33 \rightarrow U_{eff} < 0, s < 0 \rightarrow \alpha_{rel}$ increases with **a**

$M/m = 13.6 \rightarrow s = -1 \rightarrow \alpha_{rel} \sim a$

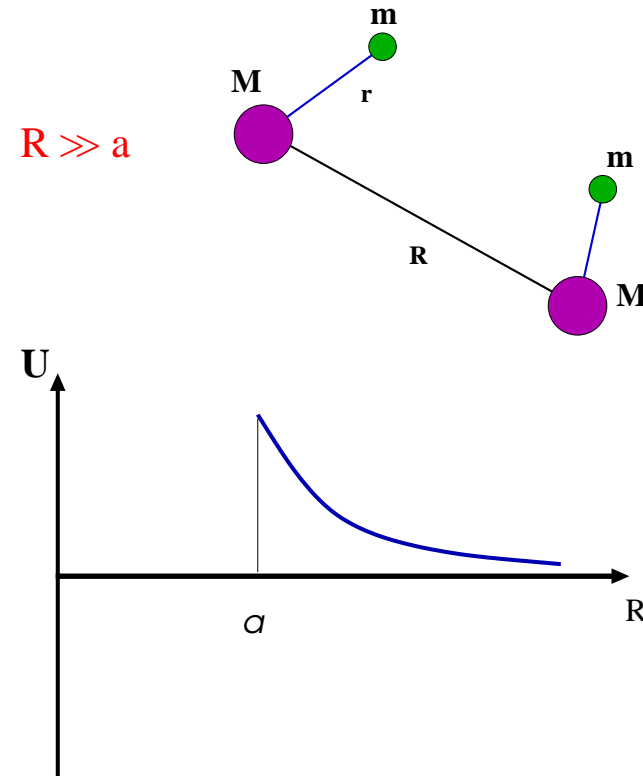
$M/m > 13.6 \rightarrow$ **fall into center** **short-range physics**

Long-range intermolecular repulsion

Molecules of heavy and light fermions **Born-Oppenheimer picture**

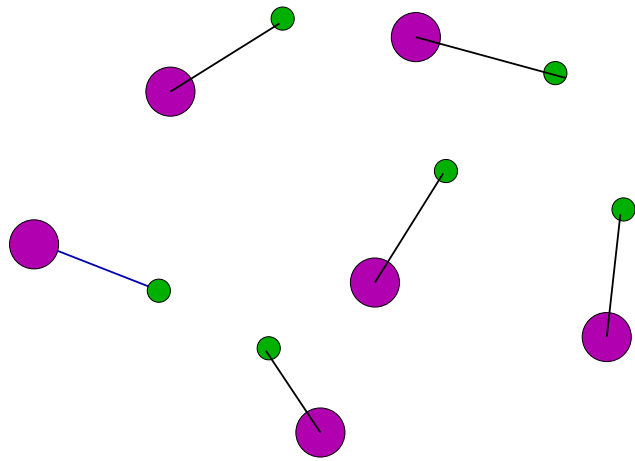
$$U(R) = 2 \left(\frac{\hbar^2}{maR} \right) \exp(-2R/a)$$

$$P \sim \exp \left(-0.9 \sqrt{\frac{M}{m}} \right)$$



$M \gg \gg m \rightarrow$ Collisional stability independent of a

Many-body system of molecules



Petrov, Astrakharchik,
Papoular, Salomon, GS

No interaction between light fermions

Born-Oppenheimer approach N lowest single-particle states for a light atom
Zero-range appr. for light-heavy interaction. Large inter-heavy distances \Rightarrow
Narrow band of N light-atom states, by $\sim \epsilon_0$ below the continuum

$$\text{Total energy } E = -N\epsilon_0 + (1/2) \sum_{i,j} U(R_{ij})$$

$$\epsilon_0 = \hbar^2 \kappa_0^2 / 2m \Rightarrow \text{molecular binding energy, } \kappa_0^{-1} \rightarrow \text{molecular size}$$

$$U_{3D}(R) = 4\epsilon_0 [1 - 2(\kappa_0 R)^{-1}] \exp(-2\kappa_0 R); \quad (1/\kappa_0 R) \exp(-\kappa_0 R) \ll 1$$

$$U_{2D}(R) = 4\epsilon_0 [\kappa_0 R K_0(\kappa_0 R) K_1(\kappa_0 R) - K_0^2(\kappa_0 R)]; \quad K_0(\kappa_0 R) \ll 1$$

$$R \approx 2/\kappa_0 \text{ or larger}$$

Phase diagram

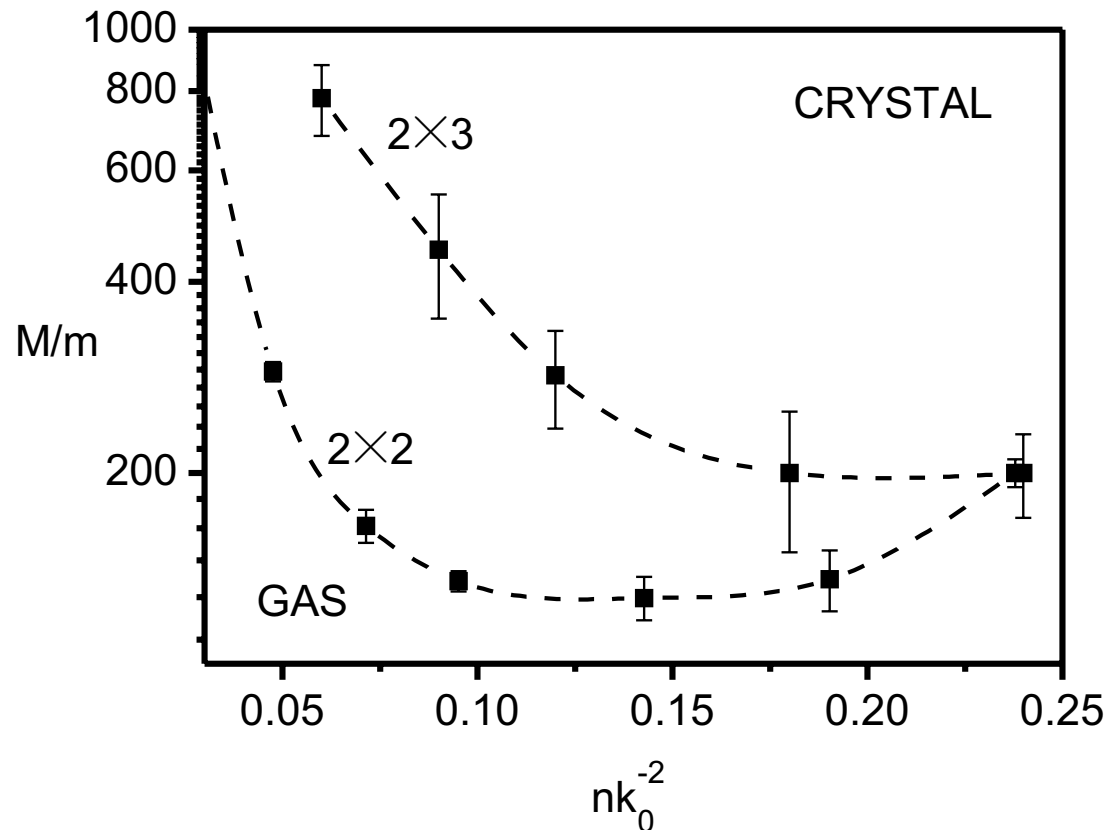
2D motion of heavy atoms

$$H = -(\hbar^2/2M) \sum_i \Delta_{R_i} + (1/2) \sum_{i,j} U(R_{i,j})$$

$(M/m) > (M/m)_c \rightarrow$ **crystalline phase**

2D motion of light atoms $\Rightarrow (M/m)_c = 120$ triangular lattice

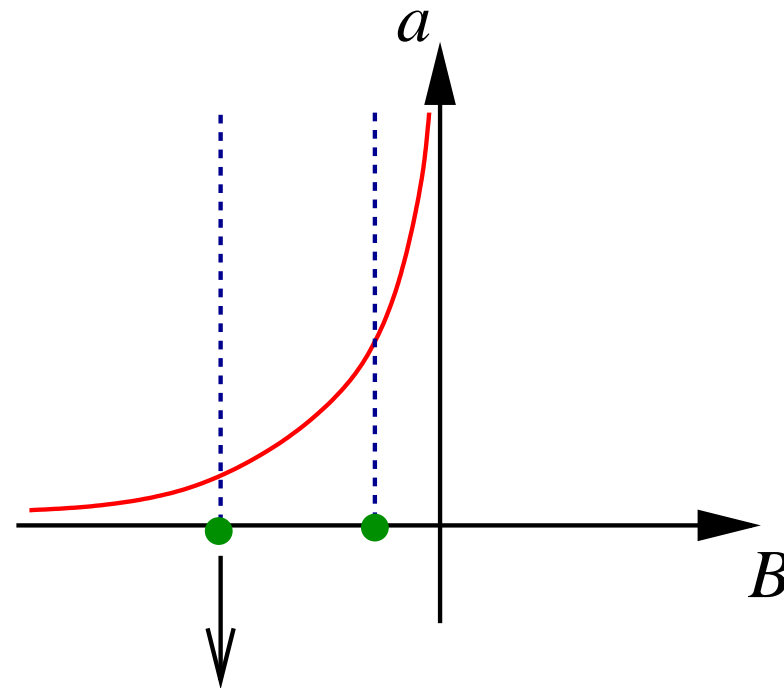
3D motion of light atoms $\Rightarrow (M/m)_c = 200$ triangular lattice



Quantum transitions

$$\frac{M}{m} > \left(\frac{M}{m}\right)_c \quad \text{and } n \text{ fixed}$$

Increase a



depends on $\frac{M}{m}$ but always $na^3 \ll 1$

first-order transition

Realization of the crystalline phase

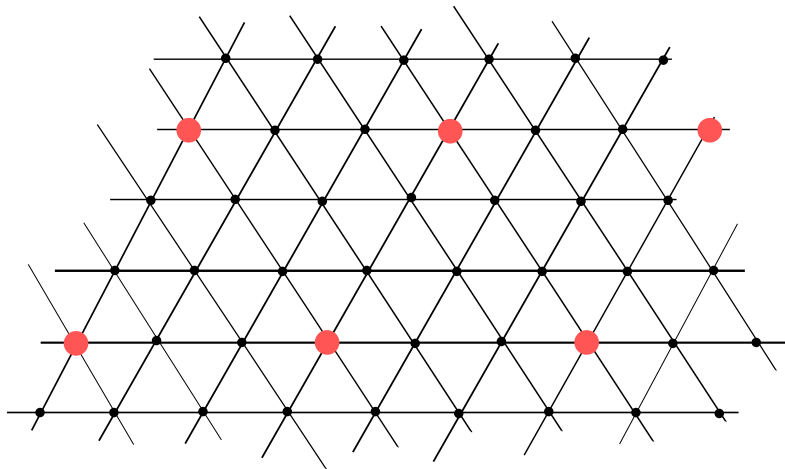
$$\frac{M}{m} \approx 200 \quad \text{or} \quad \frac{M}{m} \approx 100 \quad \rightarrow \text{no gas phase possible}$$

How to obtain the crystalline phase?

Optical lattice for heavy fermions

Small filling factor \Rightarrow Increase of M/m

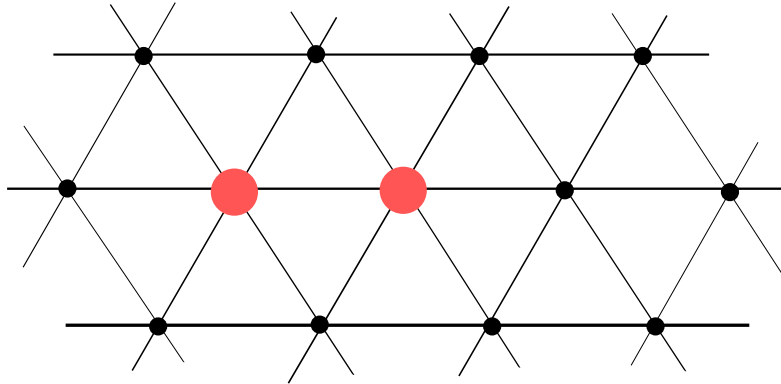
Increase of M by a factor of 20 or more is possible



Formation of a superlattice

Stability of the crystalline phase

Relaxation into deep bound states



heavy atoms in neighboring sites $\Rightarrow P_1 \sim nL^2 \exp(-\sqrt{M_*/m})$

jump to the same site $\Rightarrow P_2 \sim (t/U_0)^2$; $t = \hbar^2/M_*L^2$, $U_0 \sim \hbar\omega_l = \hbar^2/ML^2$

undergo relaxation process $\Rightarrow \tau_0^{-1} \sim (\hbar/ML)(1/l^3)$ at worst

Relaxation rate $\tau^{-1} \sim P_1P_2\tau_0^{-1} \sim nL^2(M/M_*)^2(l/L)^2(\hbar/M) \exp(-\sqrt{M_*/m})$

τ exceeds 10s even for $n \sim 10^9 \text{ cm}^{-2}$

Formation of trimer states (2 heavy and 1 light atom)

4-body problem in a lattice $\Rightarrow \tau$ can range from 0.1 to 100s for $n \sim 10^9 \text{ cm}^{-2}$

Conclusions

- Remarkable physics of weakly bound molecules in cold Fermi gases
- Novel physics of molecular collisional stability in mixtures of Fermi gases
- Possibilities to create new macroscopic quantum systems