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Collective excitations in Bose and Fermi gases

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Collective oscillations of superfluid atomic Bose and Fermi gases

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CNR-INFM





- When T tends to 0 a macroscopic fraction of **bosons** occupies a single particle state (**BEC**)
- Wave function of macroscopically occupied single particle state defines order parameter
- Actual form of order parameter depends on two-body interaction (scattering length) and is described by Gross-Pitaevskii equation

Fermions

- In the absence of interactions the physics of **fermions** deeply differs from the one of bosons (consequence of **Pauli** principle)
- Interactions can change the scenario in a drastic way:
 - pairs of atoms can form a bound state (molecule) and give rise to **BEC**
 - pairing can affect the many-body physics also in the absence of two-body molecular formation (many-body or Cooper pairing) giving rise to BCS superfluidity.
 - In most cases the physics of interacting fermions is characterized by s-wave scattering length a





Interactions provide a continuous link between the physics of BEC and BCS superfluids



HYDRODYNAMIC THEORY OF SUPERFLUIDS

Basic assumptions:

- **Irrotationality** constraint (follows from the phase of order parameter)

- Conservation laws

(equation of continuity, equation for the current)

Basic ingredient:

- Equation of state

Consequence of Galilean invariance: from the equation for the field operator to the hydrodynamic equations of superfluids

Heisenberg equation for the field operator in uniform systems (Bose field)

$$i\hbar\frac{\partial}{\partial t}\hat{\Psi}(r,t) = [\hat{\Psi}(r,t),H] = [-\frac{\hbar^2}{2m}\nabla^2 + 2\int dr'\hat{\Psi}^+(r',t)V(r-r')\hat{\Psi}(r',t)]\hat{\Psi}(r,t)$$

(similar equation for Fermi field operator)

If $\hat{\Psi}(r,t)$ is solution, $\hat{\Psi}(r-vt,t) \exp[\frac{i}{\hbar}(mrv-\frac{1}{2}mv^2t)]$ is also solution (Galilean transformation with velocity v)

Order parameter
$$\langle \hat{\Psi} \rangle$$
 ($\langle \hat{\Psi} \hat{\Psi} \rangle$ in Fermi case)
acquires phase $S(r,t) = [mrv - (\frac{1}{2}mv^2 + \mu)t]/\hbar$ $m \rightarrow 2m$
in Fermi case





Relationship for **superfluid velocity** and **equation for the phase** are expected to hold also if order parameter varies **slowly** in space and time as well as in the presence of a smooth external potential. $\mu \rightarrow \mu + V_{ext}(r)$

HYDRODYNAMIC EQUATIONS **AT ZERO TEMPERATURE**

$$\frac{\partial}{\partial t}n + \nabla(vn) = 0$$

$$m\frac{\partial}{\partial t}v + \nabla(\frac{1}{2}mv^{2} + \mu(n) + V_{ext}) = 0$$

$$\int_{irrotationality}$$

Hydrodynamic equations of superfluids **Closed** equations for density and superfluid velocity field

KEY FEATURES OF HD EQUATIONS OF SUPERFLUIDS

- Have classical form (do not depend on Planck constant)
- Velocity field is irrotational
- Are equations for the total density (not for the condensate density)
- Should be distinguished from rotational hydrodynamics.
- Applicable to low energy, macroscopic, phenomena
- Hold for both **Bose** and **Fermi** superfluids
- Depend on equation of state $\mu(n)$ (sensitive to quantum correlations, statistics, dimensionality, ...)
- Equilibrium solutions (v=0) consistent with LDA $\mu(n) + V_{ext}(r) = \mu_0$

What do we mean by **macroscopic**, **low energy** phenomena?



WHAT ARE THE HYDRODYNAMIC EQUATIONS USEFUL FOR ?

They provide quantitative predictions for

- Expansion of the gas following sudden release of the trap
- Collective oscillations excited by modulating harmonic trap

Quantities of highest interest from both **theoretical** and **experimental** point of view

- Expansion provides information on release energy, sensitive to anisotropy
- Collective frequencies are measurable with highest precision and can provide accurate test of equation of state

Collective oscillations in trapped gases

Collective oscillations: unique tool to explore consequence of **superfluidity** and test the **equation of state** of interacting quantum gases (both **Bose** and **Fermi**)

> **Experimental data** for collective frequencies are available with **high precision** (recent application to the study of the Casimir force, Obrecht et al. 2007))

Propagation of sound in trapped gases

In **uniform** medium HD theory gives sound wave solution

$$\delta n \propto e^{i(qz-\omega t)}$$
 with $\omega = cq$; $mc^2 = \partial \mu / \partial n$

In **trapped** gases sound waves can propagate if wave length is smaller than axial size of the condensate. Condition is easily satisfied in elongated condensates.



Propagation of sound in elongated traps



For BEC gas (
$$\mu \propto n$$
)
 $c_{1D} = c_{bulk} / \sqrt{2}$ (Zaremba, 1998)
For unitary Fermi gas ($\mu \propto n^{2/3}$)
 $c_{1D} = c_{bulk} \sqrt{3/5}$ (Capuzzi et al, 2006)

Bosons

Sound wave packets propagating in a BEC (Mit 97)

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Sound wave packets propagating in an Interacting Fermi gas (Duke, 2006) behavior along the crossover



Difference bewteen BCS and QMC reflects: -at unitarity: different value of β in eq. of state $\mu \propto (1+\beta)n^{2/3}$ -On BEC side different molecule-molecule scattering length

Collective oscillations in harmonic trap

When wavelength is of the order of the size of the atomic cloud sound is no longer a useful concept. Solve linearized 3D HD equations

$$\frac{\partial^2}{\partial t^2} \delta n = \nabla (n_0 \nabla (\frac{\partial \mu}{\partial n} \delta n))$$

 $n_0(\vec{r})$ is non uniform equilibrium Thomas Fermi profile

$$m\frac{\partial}{\partial t}v = -\nabla(\frac{\partial\mu}{\partial n}\delta n)$$

Solutions of HD equations in harmonic trap predict both **surface** and **compression** modes (first investigated in dilute BEC gases (Stringari 96)

Surface modes





Fermions



- Experiments on collective oscillations show that on the BCS side of the resonance superfluidity is broken for relatively small values of $1/k_F|a|$ (where gap is of the order of radial oscillator frequency)

- Deeper in BCS regime frequency takes collisionless value
- Damping is minimum near resonance

Compression modes





Equation of state along BCS-BEC crossover

- Fixed Node Diffusion MC (Astrakharchick et al., 2004)
- Comparison with mean field BCS theory (- -



Radial breathing mode at Innsbruck (Altmeyer et al., 2007)

Fermions

MC equation of state (Astrakharchick et al., 2005)



Main conclusions concerning the m=0 radial compression mode in superfluid Fermi gases

- Accurate confirmation of the universal HD value $\sqrt{10/3\omega_{\perp}}$ predicted at unitarity.
- Accurate confirmation of **QMC equation of state** on the BEC side of the resonance.

- First evidence for Lee Huang Yang effect (enhancement of frequency with respect to BEC value (role of quantum fluctuations)

Landau's critical velocity (beyond HD)

While in BEC gas **sound velocity** provides critical velocity, in a Fermi BCS superfluid critical velocity is fixed by pair breaking mechanisms (role of the **gap**)

Landau's critical velocity **Dispersion law of** $v_{cr} = \min_p \frac{\varepsilon(p)}{\varepsilon(p)}$ elementary excitations - Landau's criterion for superfluidity (metastability): fluid moving with velocity **smaller** than critical velocity cannot decay (persistent current) - Ideal Bose gas and ideal Fermi gas one has $v_{cr} = 0$ - In interacting Fermi gas one predicts two limiting cases: **BCS** (role of the gap) **BEC** (Bogoliubov dispersion) $v_{cr} = c \propto \sqrt{a}$ (sound velocity) $v_{cr} = \Delta / p_F \propto \exp(\pi / 2k_F a)$

Dispersion law of density excitations along BCS-BEC crossover



(Combescot, Kagan and Stringari 2006)

Landau's critical velocity



Landau's critical velocity is highest near unitarity !!

Some conclusions



- high Landau's critical velocity

Recent reviews on ultracold Bose and Fermi gases

- Ultracold atomic Fermi gases Proceedings of 2006 Varenna Summer School W. Ketterle, M. Inguscio, and Ch. Salomon (in press)

- Many-body physics with ultracold atomic gases I. Bloch, J. Dalibard and W. Zwerger cond-mat/0704.3011 (submitted to Rev.Mod.Phys.)

Theory of Ultracold Fermi gases
 S. Giorgini, L. Pitaevskii and S. Stringari
 cond-mat/0706.3360 (submitted to Rev.Mod.Phys.)