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Grazing Incidence Small Angle X-ray Scattering (GISAXS)

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Introduction:

- Scattering with incidence at high angle -> grazing angle: what changes ?
- When is grazing incidence scattering useful ?
- Other grazing incidence techniques:



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The SAXS technique



Bragg's Law $\lambda = 2d \sin\theta$

Scattering Vector q



$$|\mathbf{q}|=2 \text{ k sin}\theta$$

 $\mathbf{q} = 2\pi / d$ inverse distance

Method of Small Angle X-Ray Scattering SAXS



SAXS technique

$$I(q) = \frac{1}{V} \left| \int \rho(r) e^{iq \cdot r} d^{3}r \right|^{2} \qquad \rho(r) = \text{electron density}$$

$$\rho_{1(2)} = \text{electron density} \qquad I(q) = \frac{N}{V} (\rho_{2} - \rho_{1})^{2} \left| F(q) \right|^{2} S(q)$$
single particle inter particle

Principle of Babinet : pores = particles

$$\int_{V_1} e^{i q \cdot r} d^3 r$$

SAXS

scattering

Grazing Incidence Small Angle X-Ray Scattering (GISAXS)



From: Lazzari R., IsGISAXS: a program for grazing-incidence small-angle X-ray scattering analysis of supported islands , J. Appl. Cryst. 35, (2002) 406

GISAXS: refraction effects of X-rays

α_i is small ->effects of refraction at the surface needs to be considered:

Refractive index n of matter for x-rays is:
with:

$$\delta = \frac{1}{2\pi} \times \frac{e^2}{mc^2} \times \frac{N_a \sum_i (Z_i - f_i')}{\sum_i A_i} \times \rho \lambda^2 \qquad \beta = \frac{1}{2\pi} \times \frac{e^2}{mc^2} \times \frac{N_a \sum_i f_i''}{\sum_i A_i} \times \rho \lambda^2 = \mu \frac{\lambda}{4\pi}$$

$$= 2.701 \times 10^{-6} x \frac{\sum_i (Z_i - f_i') \text{ [el. units]}}{\sum_i A_i \text{ [g / mol]}} \times \rho \text{ [g/ cm^3] } \lambda^2 \text{ [}A^2 \text{]}$$

summation over all atomic species i:

 $(Z_i - f_i)$ scattering factor

- anomalous dispersion factor
- A_i atomic weight

f,"

- N_a Avogadro's number
- ρ density
- λ wave length
- μ photo electric absorption coefficient

GISAXS: refraction effects of X-rays

Refractive index $n = 1 - \delta + i\beta$

critical angle of total reflection $\alpha_c \propto \sqrt{2\delta}$



Note:

- Real part of $n = 1 \delta + i\beta$ is always slightly < 1
- Because of refraction, a beam passing from vacuum into a sample comes closer to the sample surface
- For $\alpha_i < \alpha_c$, the beam is totally reflected, and only an evanescent wave, which decays over some nm, is present below the surface
- For $\alpha_i > \alpha_c$ the transmitted wave propagates

Scattering Depth

The perpendicular components of the incident and emergent wave vectors are modified upon crossing the surface, and become complex due to refraction and absorption: 2π

$$k'_i,_f = \frac{2\pi}{\lambda}(A_i,_f - iB_i,_f)$$

For small incidence and emergence angles:

$$A_{i},_{f} = \frac{1}{\sqrt{2}} (\sqrt{(\alpha_{i},_{f}^{2} - \alpha_{c}^{2})^{2} + 4\beta^{2}} + \alpha_{i},_{f}^{2} - \alpha_{c}^{2})^{0.5}$$
$$B_{i},_{f} = \frac{1}{\sqrt{2}} (\sqrt{(\alpha_{i},_{f}^{2} - \alpha_{c}^{2})^{2} + 4\beta^{2}} + \alpha_{c}^{2} - \alpha_{i},_{f}^{2})^{0.5}$$

-> the perpendicular momentum transfer $Q'_{\perp} = k'_{f\perp} - k'_{i\perp}$ inside the sample becomes complex

The scattering depth
$$\Lambda = \frac{1}{\text{Im}(Q'_{\perp})} = \frac{\lambda}{4\pi(B_i + B_f)}$$
 is thus strongly affected
by refraction when α_i and / or α_f are close to α_c



G. Renaud/Surface Science Reports 32 (1998) 1-90

Example: Quantum Dots studied by GISAXS

II-VI semi-conductors Nanostructures: what for ?

quantum confinement of excitons, increase of the energy gap

- variable length optical absorbers
- high-speed non-linear optical switches
- photoluminiscence with variabel wavelength by changing particle size (-> quantum dot laser)
- formation of nanoparticles in dielectrics and nonabsorbing materials

BUT: working devices will depend on a precise control of size and density of an ensemble of Quantum Dots (QD)!



Fig. 5. Calculated and measured bandgap shift for CdS nanocrystals. The solid line is calculated from Eq. (2). The points are measurements of the bandgap shift from Eq. (1) for the samples in Fig. 4.



C.W. White et al. | Nucl. Instr. and Meth. in Phys. Res. B 141 (1998) 228

Nanocrystals in Substrates



GISAXS: what information do we get ?



 $\Lambda = 2\pi/q_m$, $\Lambda = average nanocrystals distance - obtained from interference peak, <math>q_{max}$

R_g = radius of gyration ⇒ D =<u>average cluster diameter</u> obtained from slope of linear region of I(q) versus q²

2D GISAXS pattern





U.V. Desnica, P. Dubcek, I.D. Desnica-Frankovic, M. Buljan, S. Bernstorff and C.W. White, J.Appl.Cryst. 36, 443-446 (2003)

Synthesis and growth of CdS Quantum Dots in SiO₂

GISAXS-pattern

Intensity profiles



U.V. Desnica, P. Dubcek, I.D. Desnica-Frankovic, M. Buljan, S. Bernstorff and C.W. White, J.Appl.Cryst. 36, 443-446 (2003)

GISAXS data evaluation

Scattering from embedded particles: Local Monodisperse Approximation - LMA -system is approximated by monodisperse subsystems, weighted by the size distribution -positions of the particles are completely correlated with their sizes

$$I(q) \propto \left|T(\alpha_i)\right|^2 \left|T(\alpha_f)\right|^2 \int_0^\infty P(q,D) S(q,D_{hs},\eta_{hs}) N(D,w) dD$$

 $T(\alpha_i)$ and $T(\alpha_f)$ - Fresnel transmission coefficients for angle of incidence α_i and exit α_f P(q,D) - form factor of a homogeneous sphere of diameter D, (D=2R); $S(q,D_{hs},\eta_{hs})$ - structure factor (η_{hs} - volume fraction and diameter of the hard spheres) N(w,D) - Gaussian size distribution function

Surface Scattering: Distorted Wave Born Approximation - DWBA

$$\mathbf{I}_{\text{total}} = \mathbf{I}_{\text{LMA}} + \mathbf{I}_{\text{surface}}$$

Distorted Wave Born Approximation: Vineyard 1982, Shinha et al. 1988 M. Rauscher, T. Salditt and H. Spohn, Phys. Rev. B 52, 16855 (1995) M. Rauscher et al., J. Appl. Phys. 86 (12), 6763 (1999)



Quantitative analysis of GISAXS



Synthesis and growth of CdS Quantum Dots in SiO₂





X-ray Diffraction

•Broad signal from amorphous SiO₂ substrate

•Superimposed: sharp peaks identified as CdS (hexagonal)

•Determination of average nanocrystal size (from the broadening of lines)

Bimodal size distribution of particles: GISAXS + TEM



Cross-sectional **TEM** image:

depth profiling by small increase of the incidence angle (0.02° step) ↓ closely spaced population of large particles with narrow size distribution R_G≈ 50 -60 nmnm shoulder - second group of smaller particles of broader size distribution R_G≈ 5-7 nm (CdSe in SiO₂, +0% Cd, 800 °C in Ar+H₂)



Size distribution for **ZnSe nanocrystals in SiO**₂ after annealing at 1000 °C, 30s, Ar. Sample was implanted with 33% surplus Cd. Among large particles positioned at approx. 100nm, there is also a band of small particles with diameters below 10 nm

U.V. Desnica, M. Buljan, I.D. Desnica-Frankovic, P. Dubcek, S. Bernstorff, M. Ivanda et al, NIM B 216, 407 (2004)

Complementary Methods



GISAXS study of Ge islands on a Si(111) substrate

Sample prepared by molecular beam epitaxy at substrate temperature 530°C



Scattering geometry for grazing incidence small angle scattering (GISAXS). The PSD is placed parallel to the sample surface is collects the small angle signal as a function of $q_{\parallel} = 2\pi/\lambda \sin(2\theta)$ at constant angles of incidence and exit, $\alpha_i \neq \alpha_f$. The axial symmetry the quantum dots are studied by rotating the azimuthal angle ω .

T.H. Metzger et al., J. Phys D: Appl. Phys 32 (1999) A202-A207

GISAXS study of Ge islands on a Si(111) substrate



symmetry. The full curve is a fit to the experimental data (0000) using the structure factor of a pyramid

3-fold symmetry of scattering pattern

means

3-fold symmetry of triangular pyramids

- Relation of
 w and known cleavage edge direction of the sample -> sides of
 pyramids are aligned along all possible <110> directions in the (111) oriented
 surface plane
- Fit of structure factor to experimental data -> dimensions of the pyramids



nm AFM (atomic force microscope) -> islands have pyramidal shape with well defined {113} facets; some exhibit also a small (111) terrace on a truncated top

T.H. Metzger et al., J. Phys D: Appl. Phys 32 (1999) A202-A207

Towards Quantum Dot Laser: the samples



Z-contrast cross section TEM for 10-fold stack sample (QD layers: 0.6 nm, spacer: 4.2 nm T_{growth} = 280 °C

Th. Schmidt et al., Physical Review B72, 195334 (2005)

Samples grown by migration enhanced molecular beam epitaxy with:

N = 3, 5 or 10

CdSe layers thickness: 1.8 monolayers (i.e. 0.5-0.6nm)

d_{spacer} = fix value in the range 2.0 nm - 8.0 nm Substrate temperature: 230, 280 or 310 °C

GISAXS reveals CdSe/ZnSSe Quantum Dot ordering



Th. Schmidt et al., Physical Review B72, 195334 (2005) Th. Schmidt, T. Clausen, J. Falta, G. Alexe, D. Hommel, S. Bernstorff, *Elettra Highlight 2004, p. 63*

Grazing Incidence Diffraction (GID)

Scattering geometry: defined by directions of incident beam and detector position

To reach diffraction condition: rotate sample about its surface normal so that netplanes make angle Θ with respect to both incident and scattered beam -> probe long range periodicity parallel to sample surface

To see <u>reflected beam</u>: $\alpha_i = \alpha_f$, $2 \Theta = 0$



Grazing Incidence Diffraction (GID)



Scattering geometry for grazing incidence diffraction to study the crystalline properties of quantum dots. Three-dimensional reciprocal space maps are recorded in this geometry by collecting the scattered intensity as a function of the exit angle α_f , the scattering angle 2θ and the sample angle θ .



Important directions in reciprocal space close to the (220) surface Bragg reflections of GaAs and InAs. By adjusting q_r the lattice relaxation in the dots is measured, along q_e information on the lateral form factor of the dots is obtained, the q_c -dependence determines the maximum depth and reveals the form factor of the dots in the growth direction. InAs islands grown by molecular beam epitaxy on GaAs(100) at 530°C substrate temperature

GID -> crystalline properties of coherent quantum dots

Analysis of reciprocal space maps -> interdependence of shape and elastic strain within the quantum dots



'Isostrain-scattering' between the (220) surface Bragg reflections of GaAs and InAs. The intensity distribution in angular scans is shown for different radial positions (q_r) as indicated in the inset. The angular dependence (q_a) of the scattering signal is characteristic for the lateral shape of the InAs dots. The central maximum broadens with increasing relaxation $|q_r|$



Form factor-induced scattering signal as a function of q_u for ring-shaped InAs islands overgrown by GaAs as a function of the relaxation q_r . The corresponding shapes of the islands are indicated schematically on the right-hand side.

T.H. Metzger et. al., J. Phys. D.: Appl. Phys. 32 (1999) A202-A207

X-ray Reflectivity measurements

Determine

- Surface / interface roughness
- Layer thickness



Why layer thickness ?

Because function of devices depends on thickness:

- IC (semiconductors)
- Solar cells (photovoltaic materials)
- Displays (liquid crystals)
- Data storage (magnetic multilayer)
- Biomaterial (proteins, lipids -> biochip)
- LED (electroluminescence)
- Coatings and platings (optical materials)

X-ray Reflectivity measurements

Thin films on substrates:



high correlation between the film and substrate surface



→ interference patterns in the specular plane
→ precise film thickness determination

concentric rings around the direct beam direction correspond to inner structure of the film





Sample rotation angle

Rough surfaces -> diffuse scattering -> lateral features of the roughness: height-height correlations

Basic experimental considerations

<u>High photon flux</u> is needed for GI-scattering: Nr. of atoms in interfaces / surfaces is << as in bulk ! -> use rotating anodes or synchrotron sources !

Critical angles for e.g. 8 kV photons: .

organic filmsac = $0.1^{\circ} <-> 1.7$ mradsubstratesac = $0.2^{\circ} <-> 3.4$ mrad

Sample size (length L x width W): accepted beam height H = L x aL= 20 mm, $a = 0.2^{\circ} \rightarrow H = 70 mm$ W > beam width of ca. 5 mm

Sample surrounding

avoid photons scattered from air -> work in vacuum, He-atmosphere, or use very short airpaths

avoid photons scattered from slits, windows, ...

--> • very flat substrates needed (e.g. silicon wafers, float glass)

• very flat films (spin coating, dip coating, implanting, ...) needed



GISAXS is an ideal tool to study:

- non-destructive
- no special sample preparation / samples as prepared (implanted, spin / dip coating, ...)
- in ambient conditions, or under vacuum
- mesoscopic scale: from molecular size to ca 150 nm
- extremely sensitive to lateral and normal structure
- accurate statistical average
- compatible with in-situ experiments (parameters can be temperature, shear, tear, chemical mixing, ...)
- Low-contrast systems possible
- very weakly-scattering samples possible (probe shape and 2D organization of interfaces and QD, of biological molecules deposited on surfaces - > bio-chips, function of selected bio-molecules)

ESSENTIAL

- source providing high flux / brilliance
- 2D detector with high dynamic range, and very low background / noise

Thank you for your attention !





If you now want to perform your own SAXS or GISAXS experiment with SR:

Info: www.elettra.trieste.it www.ibn.oeaw.at/beamline/

you are also welcome to e-mail me at

bernstorff@elettra.trieste.it



Some useful literature ...

SAXS:

- · A. Guinier and G. Fournet, "Small-Angle Scattering of X-rays", Wiley, New York, 1955
- · O. Glatter and O. Kratky, editors of "Small-Angle X-Ray Scattering", Academic Press, New York 1982
- L. Feigin and D. Svegun, "Structure Analysis by Small-Angle X-Ray and Neutron Scattering, Plenum Press, New York, 1987
- O. Kratky and P. Laggner, "X-Ray Small-Angle Scattering", Encyclopedia of Physical Science and Technology, Vol 17, pp. 727-781, Academic Pres, Orlando, 1992
- H. Brumberger, editor of "Modern Aspects of Small-Angle Scattering", NATO Advanced Series, Kluwer, Dordrecht, 1993
- · Manfred's homepage http://scattering.tripod.com/

GISAXS:

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- Till Hartmut Metzger, Tobias Urs Schuelli and Martin Schmidbauer, "X-ray Methods for Strain and Composition Analysis in Self-Organized Semiconductor Nanostructures", lab-neel.grenoble.cnrs.fr/pageperso/fruche/crphys/contrib/metzger.pdf
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- Salditt T., Metzger T.H. and Peisl J., Kinetic Roughness of Amorphous Multilavers Studied by Diffuse X-Ray Scattering, Phys. Rev. Lett. 73, (1994) 2228.
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- M. Rauscher et al., J. Appl. Phys. 86 (12), 6763 (1999)

Infos on Austrian SAXS beamline at ELETTRA: http://www.ibn.oeaw.ac.at/beamline/