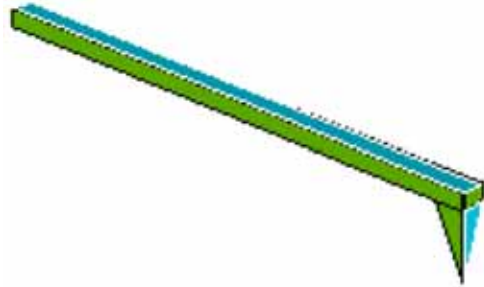


# Electron counting with quantum dots



Klaus Ensslin



Solid State Physics



**Zürich**

with

S. Gustavsson

I. Shorubalko

R. Leturcq

T. Ihn

A. C. Gossard

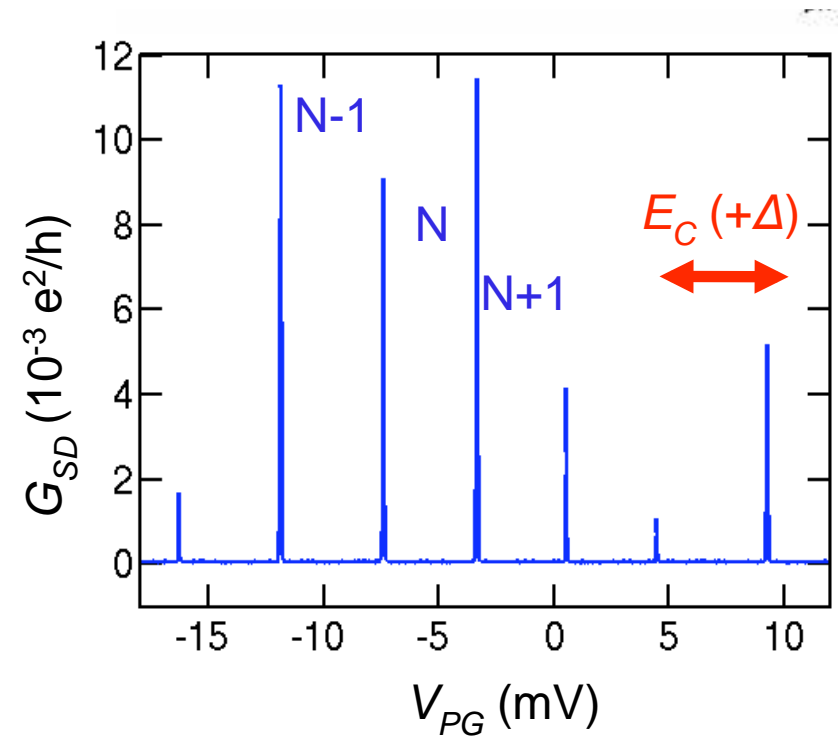
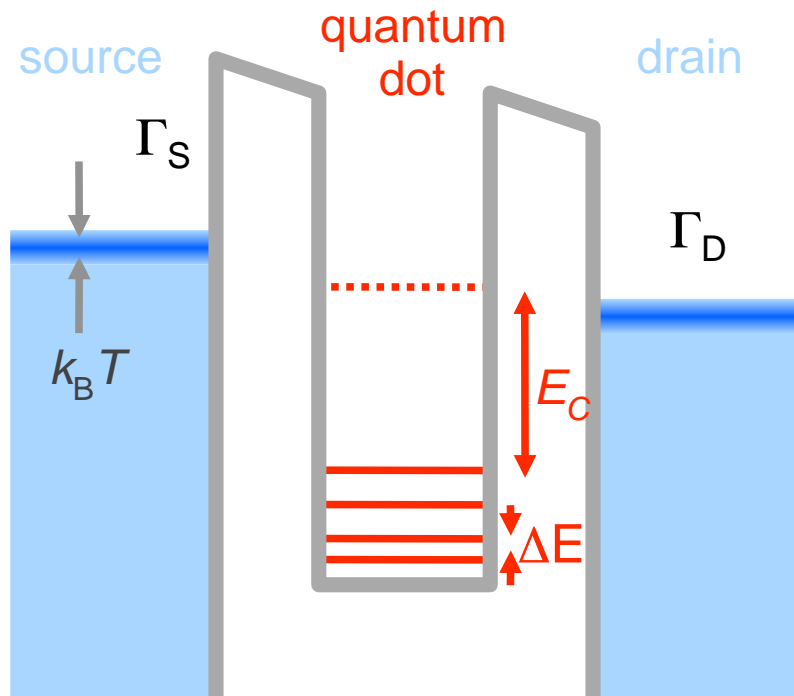
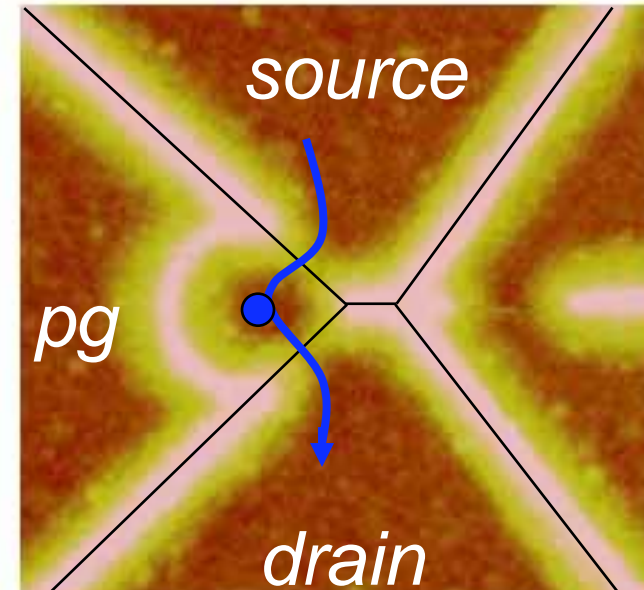


Time-resolved charge detection

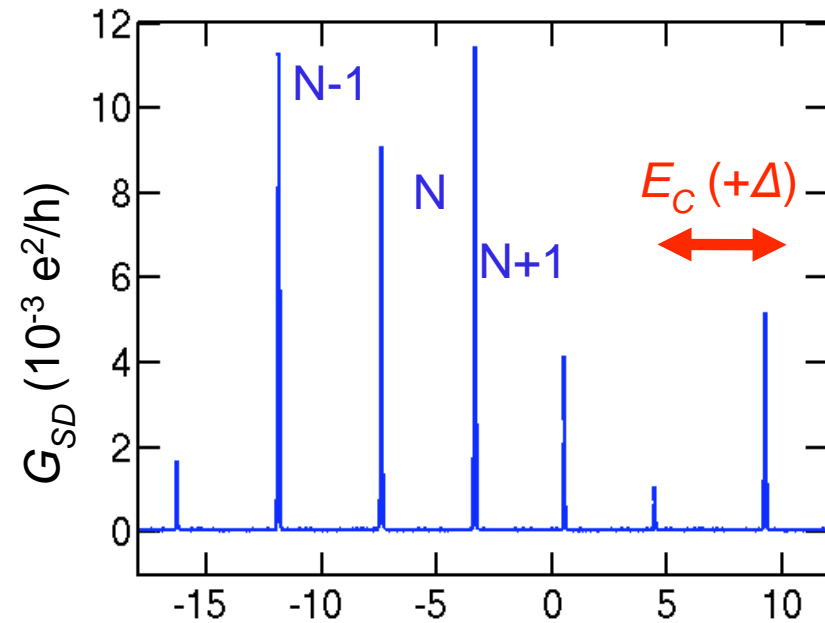
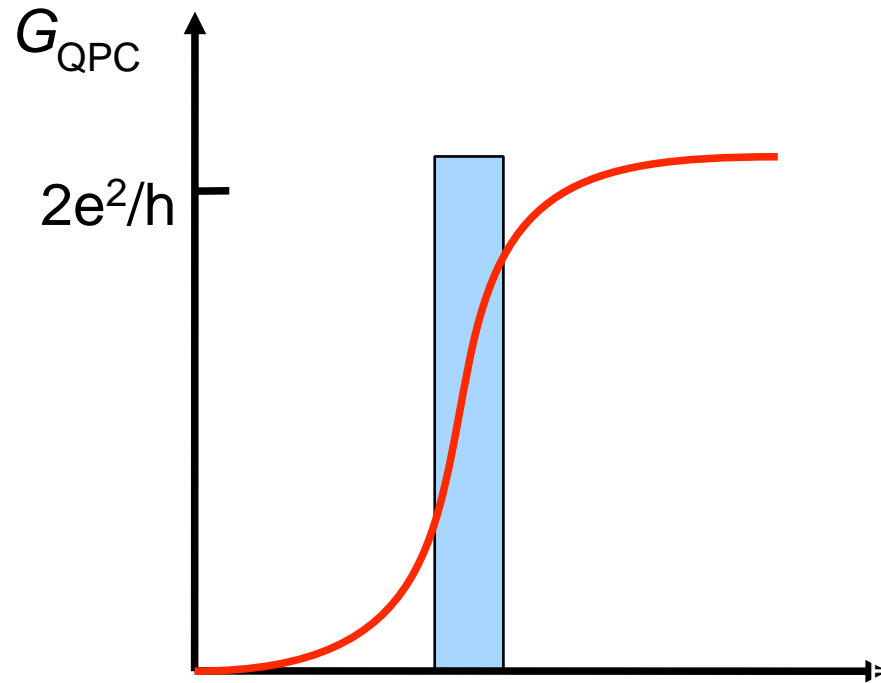
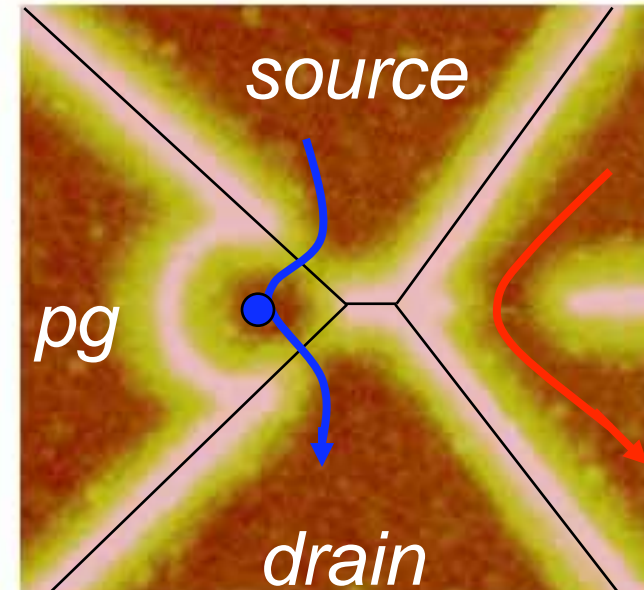
Single photon detection

Time-resolved single electron interference

# Spectroscopy of electronic states

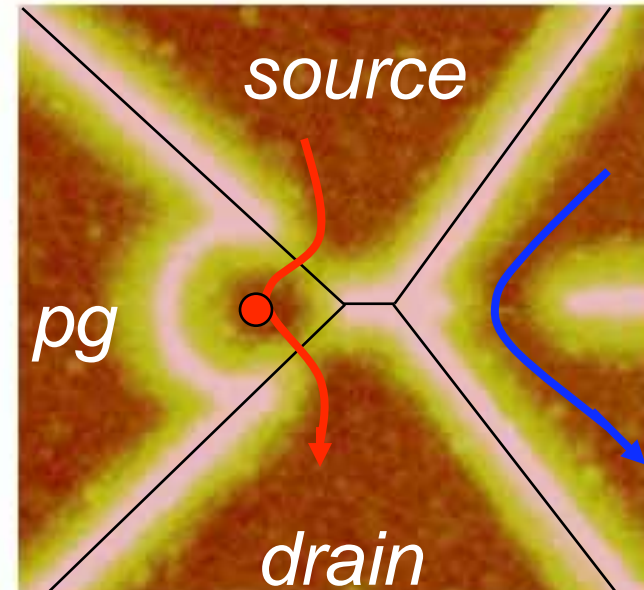
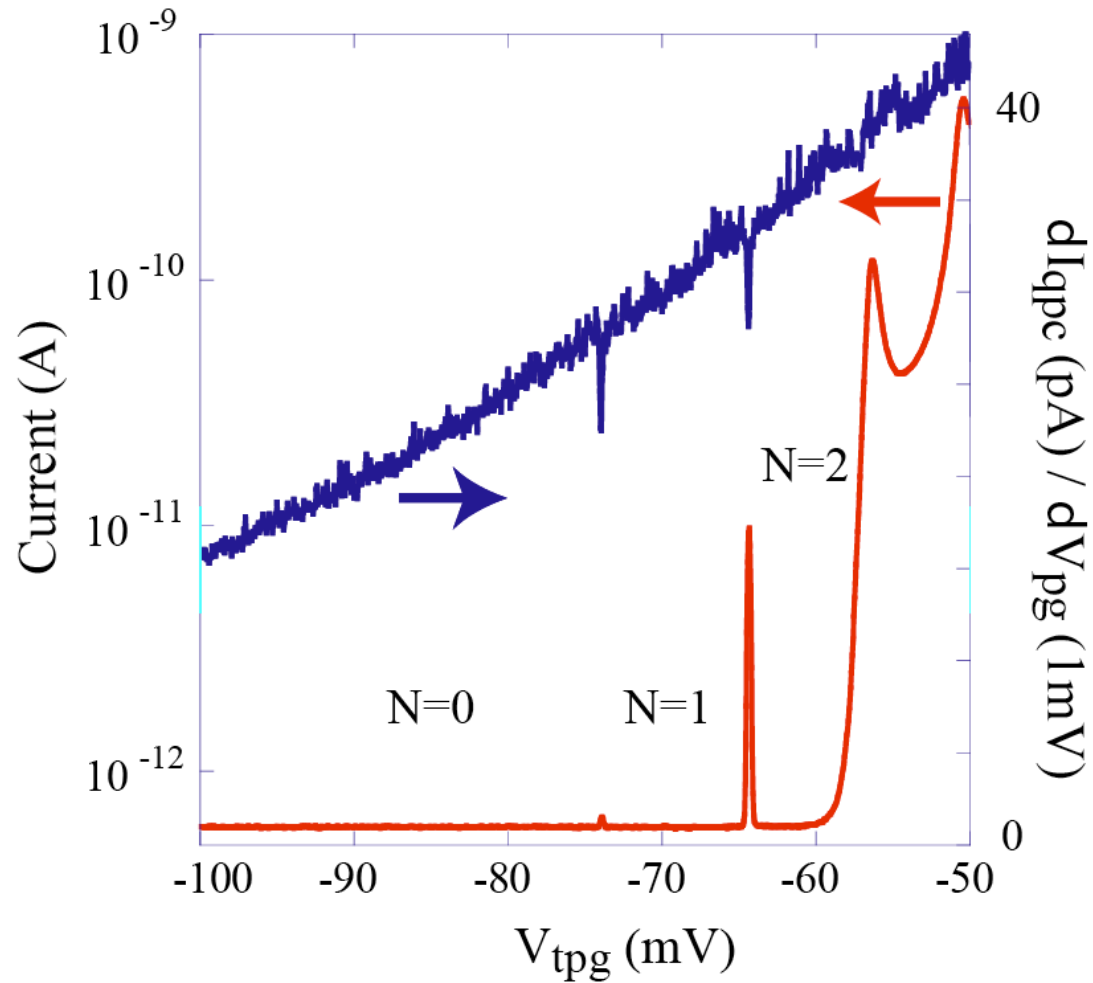


# Quantum point contact as a charge detector



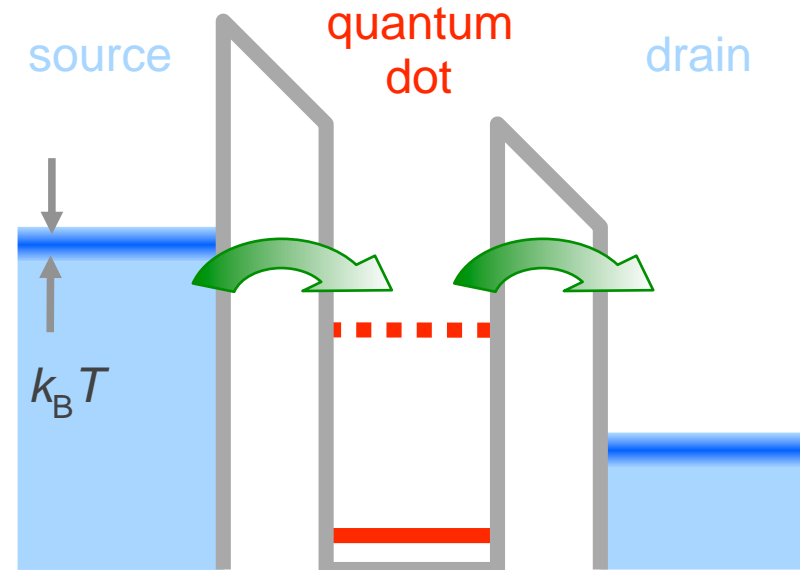
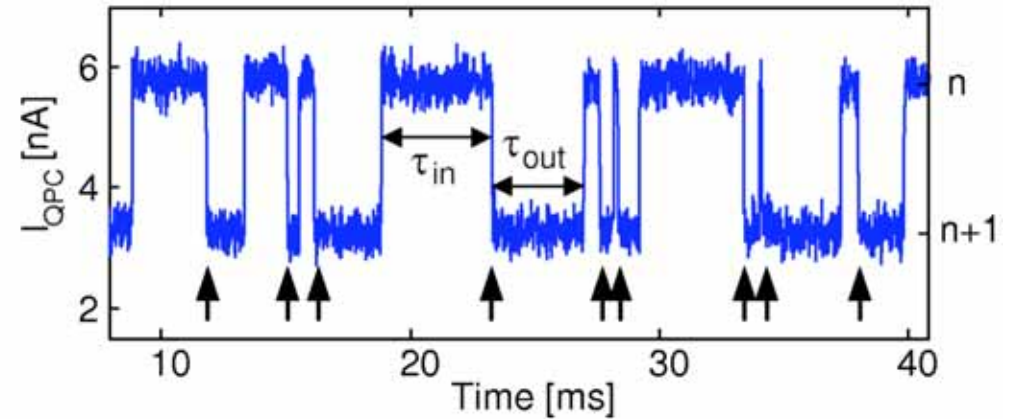
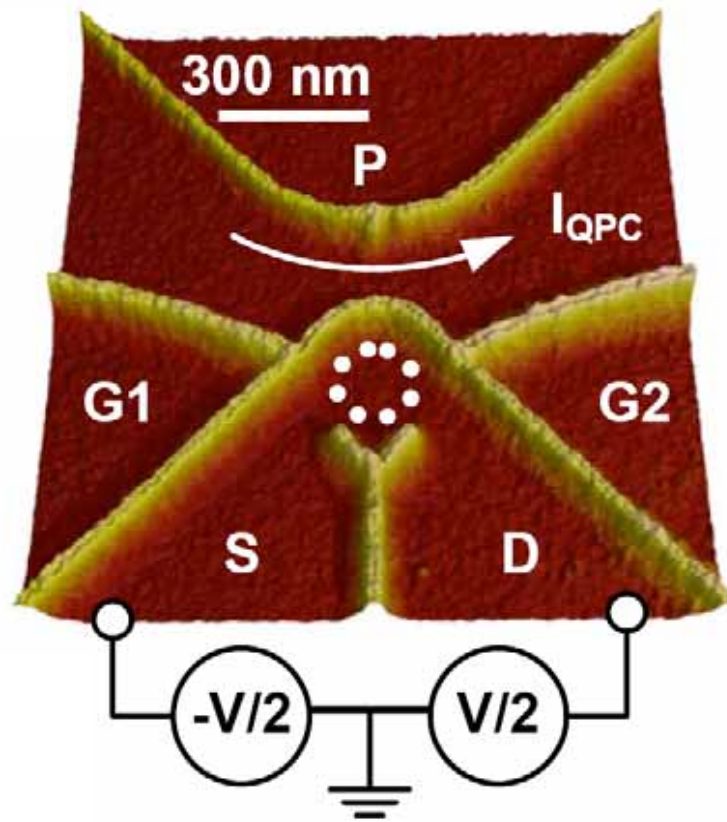
M. Field et al., Phys. Rev. Lett. 70, 1311 (1993)  $V_{PG}$  (mV)

# A few electron quantum dot



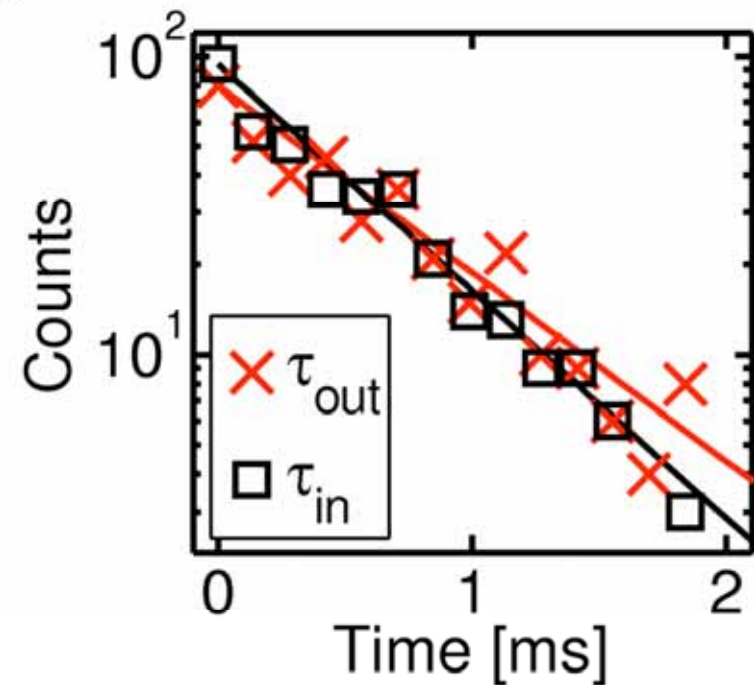
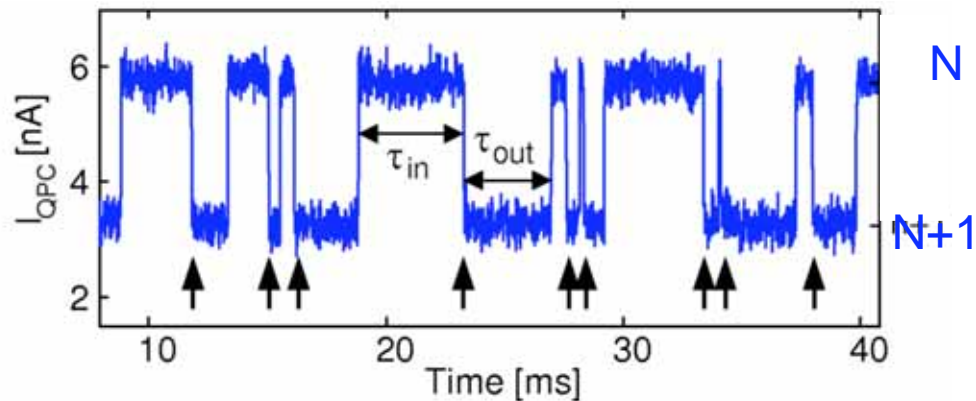
M. Sigrist

# Detection of single electron transport

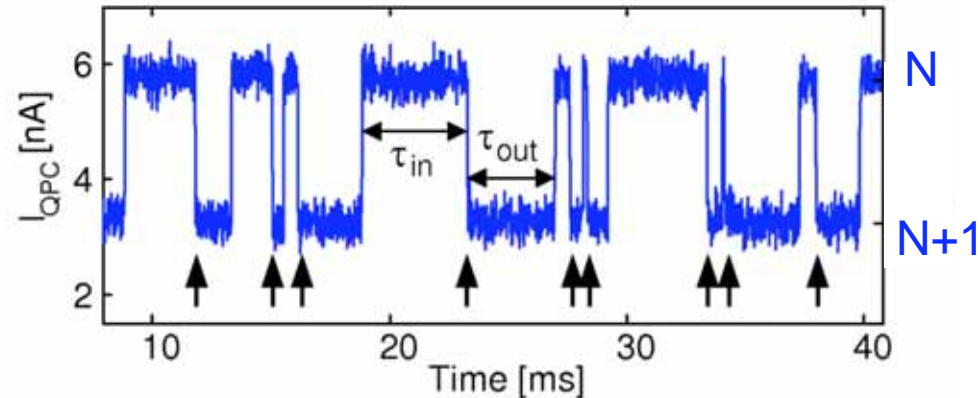


# Determination of the individual tunneling rates

- Exponential distribution of waiting times for independent events
- $\Gamma_S = \langle \tau_{in} \rangle$ ,  $\Gamma_D = \langle \tau_{out} \rangle$



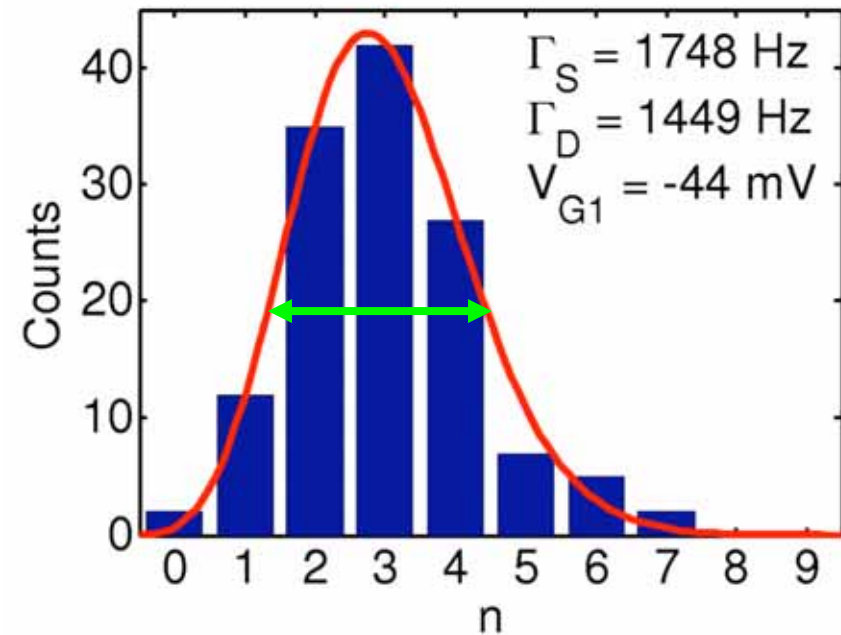
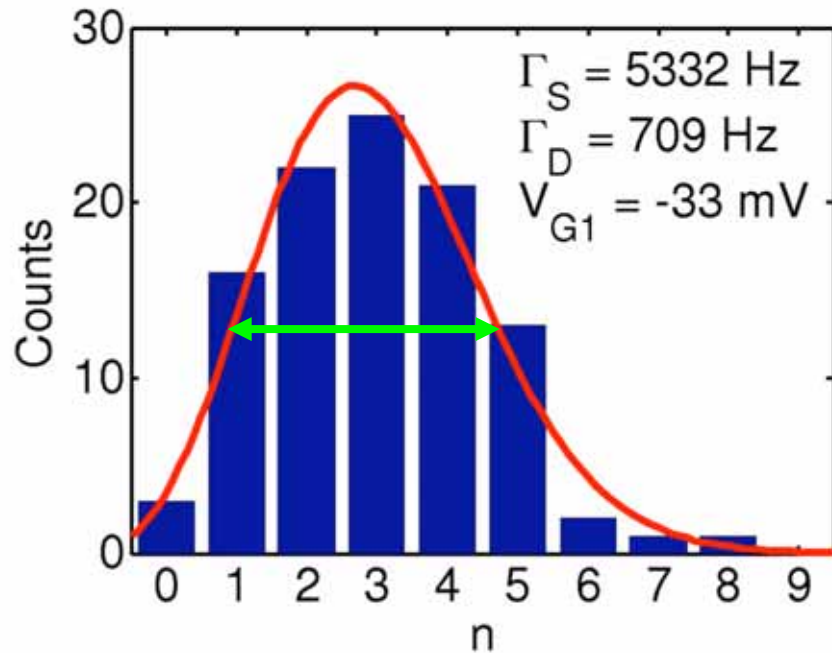
# Measuring the current by counting electrons



- Count number  $n$  of electrons entering the dot within a time  $t_0$ :  $I = e\langle n \rangle / t_0$
- Max. current = few fA (bandwidth = 30 kHz)
- BUT no absolute limitation for low current and noise measurements
  - here:  $I \approx$  few aA,  $S_I \approx 10^{-35} \text{ A}^2/\text{Hz}$



# Histogram of current fluctuations



- Poisson distribution for asymmetric coupling

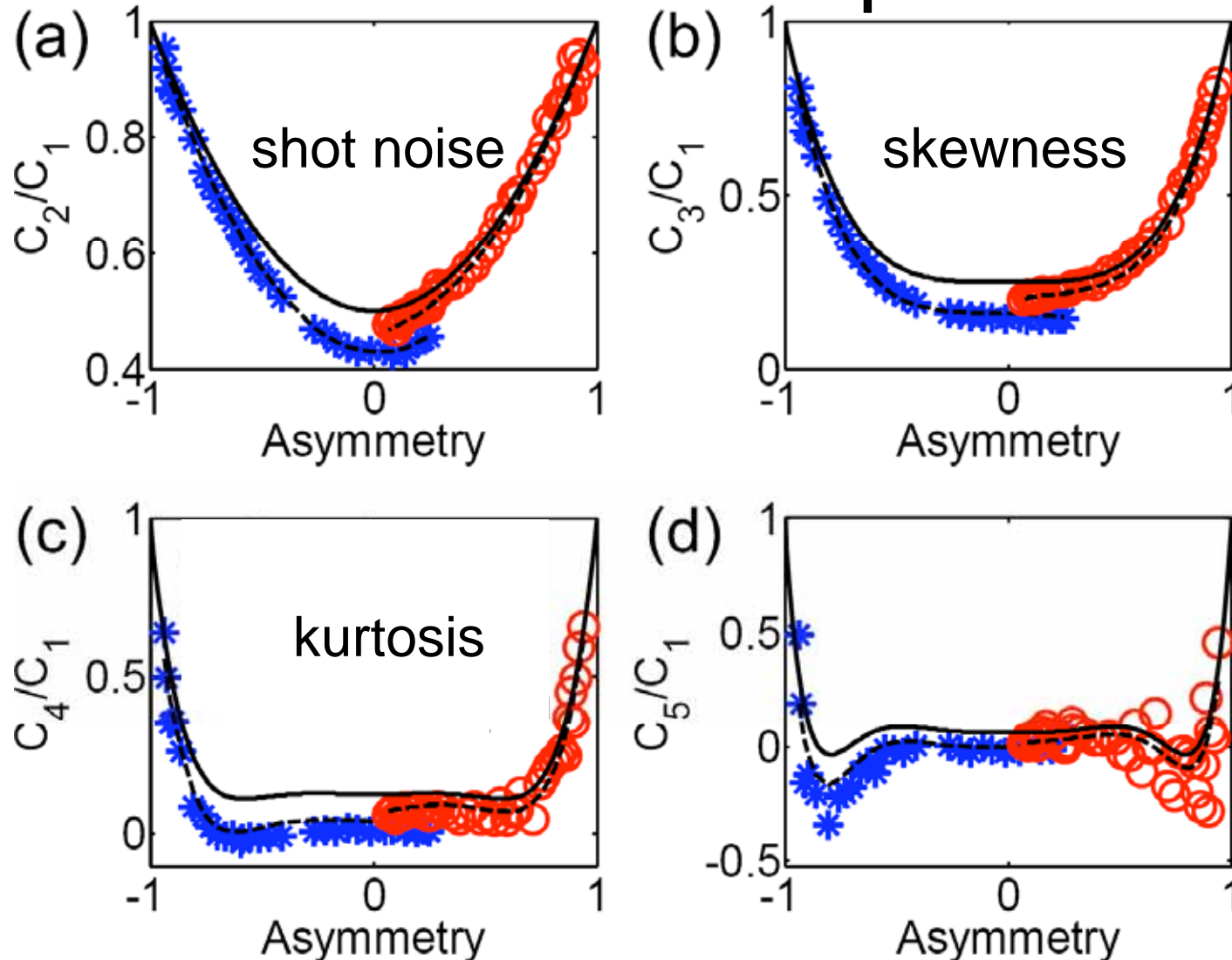
- Sub-Poisson distribution for symmetric coupling

Theory: Hershfield *et al.*, PRB **47**, 1967 (1993)  
Bagrets & Nazarov, PRB **67**, 085316 (2003)

Expt: Gustavsson *et al.*, PRL **96**, 076605 (2006)

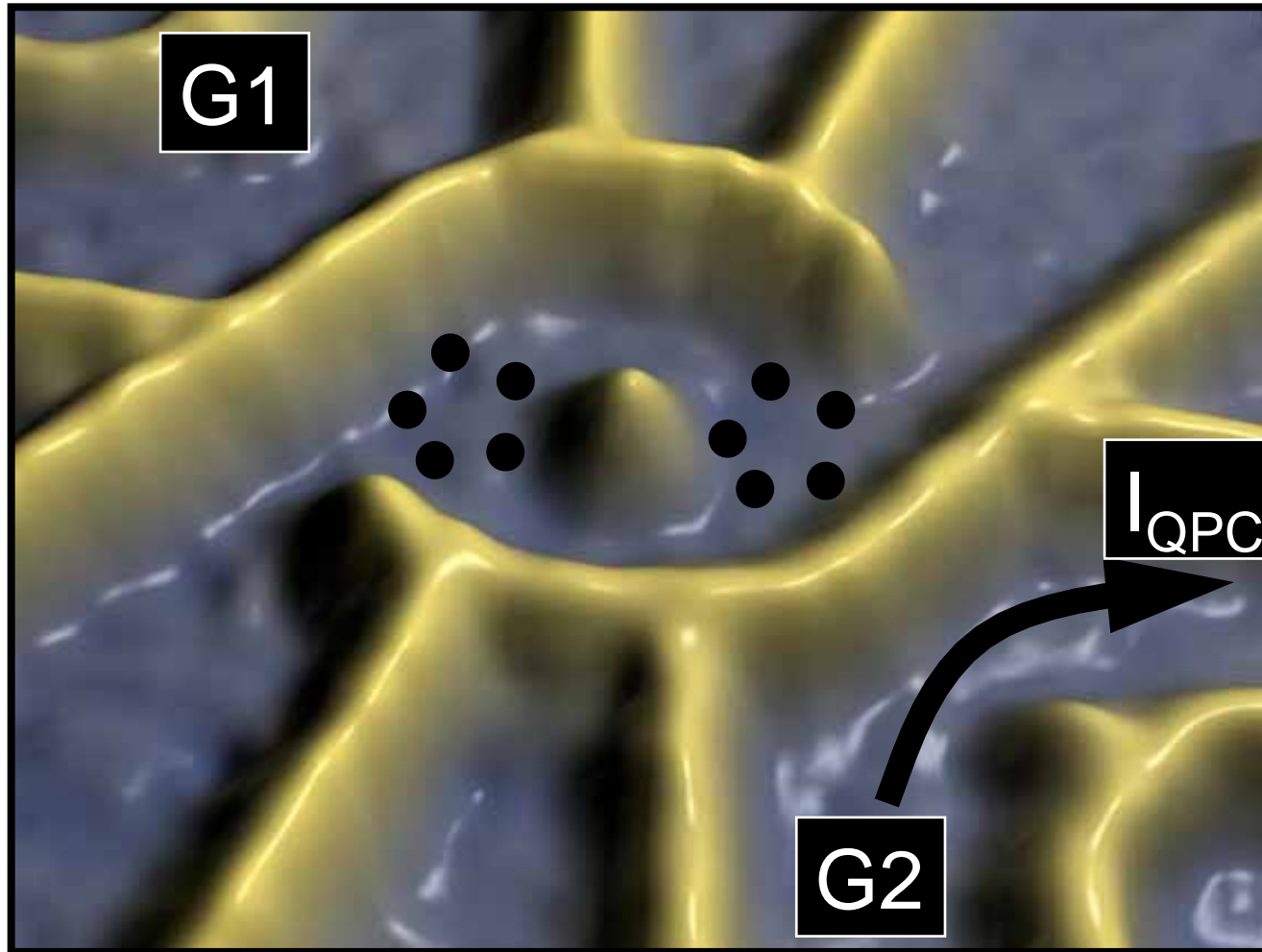


# Higher order correlations of electron transport

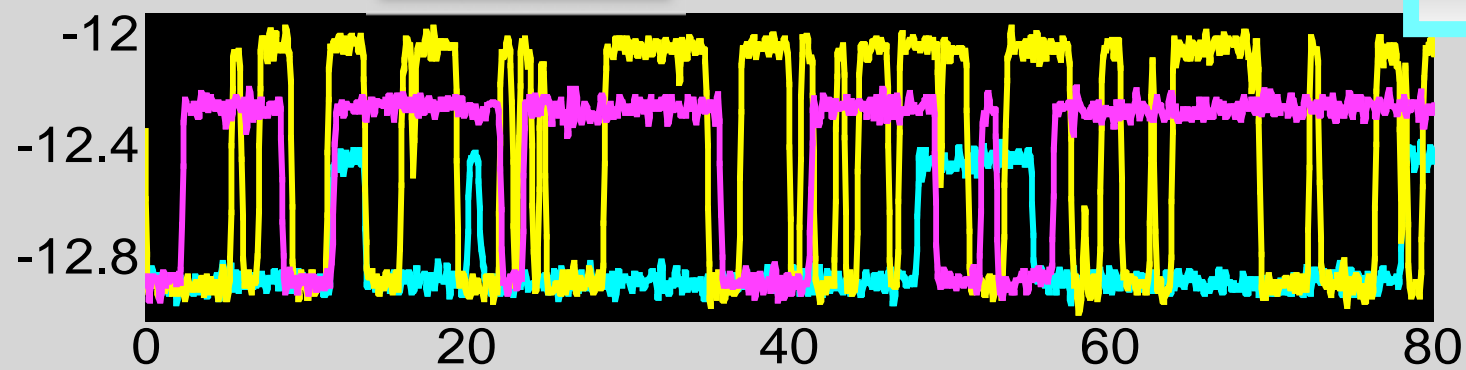
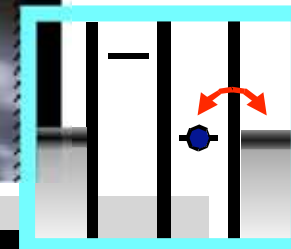
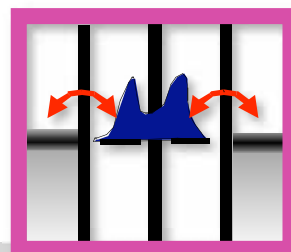
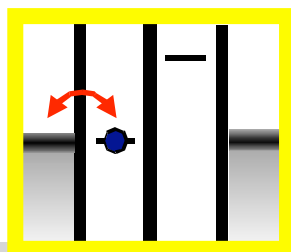
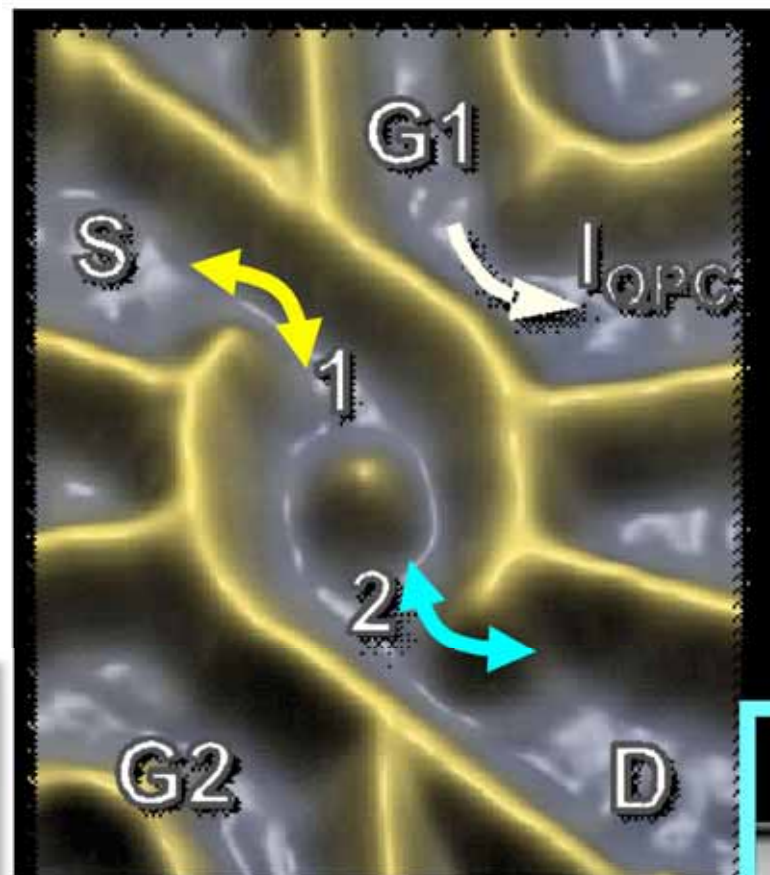
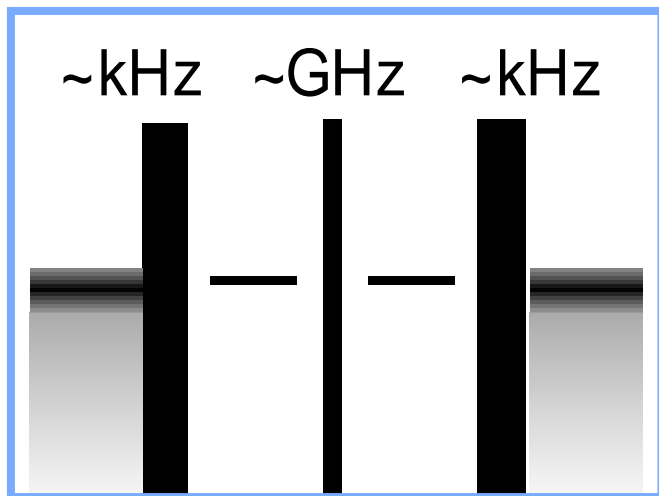


Gustavsson et al, PRB 75, 075314 (2007)

# Double quantum dot in a ring



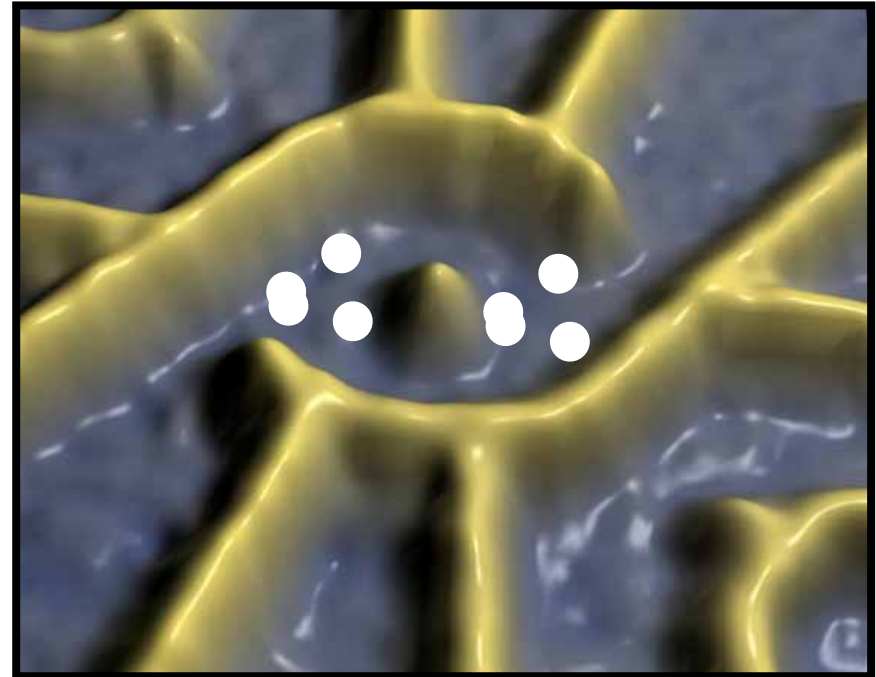
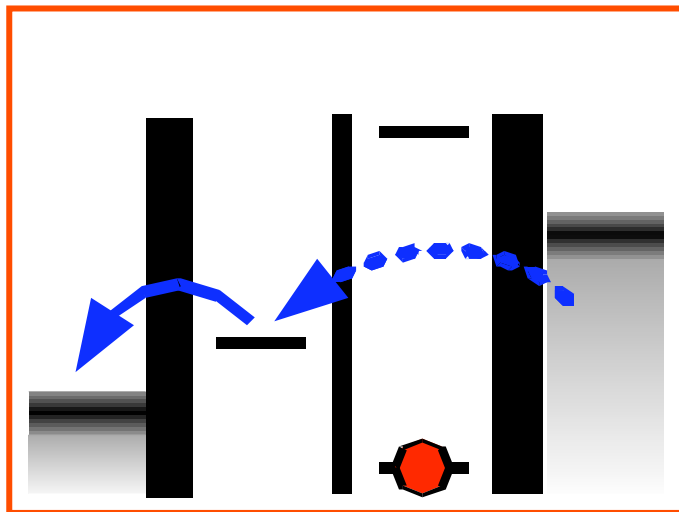
see also: electron counting in double dots: Fujisawa *et al.*, Science **312**, 1634 (2006)



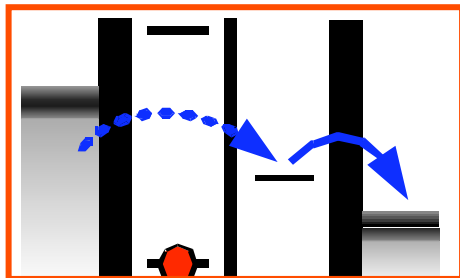
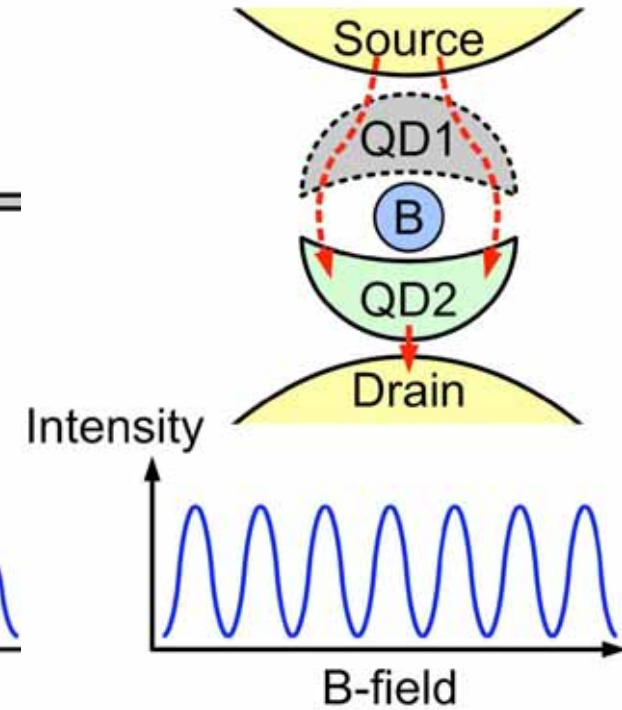
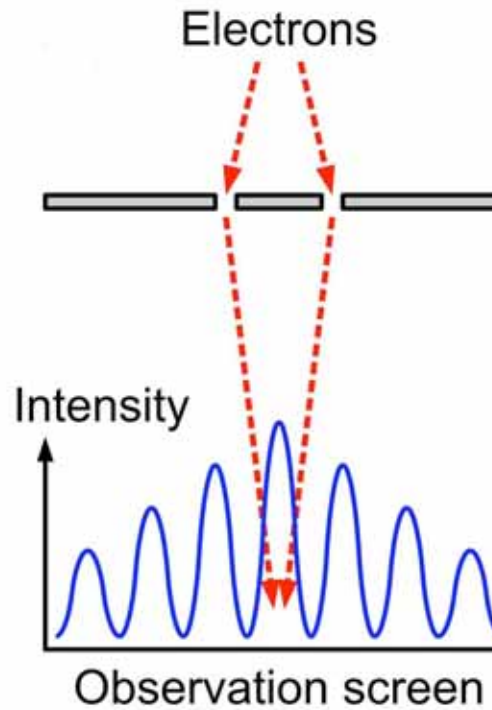
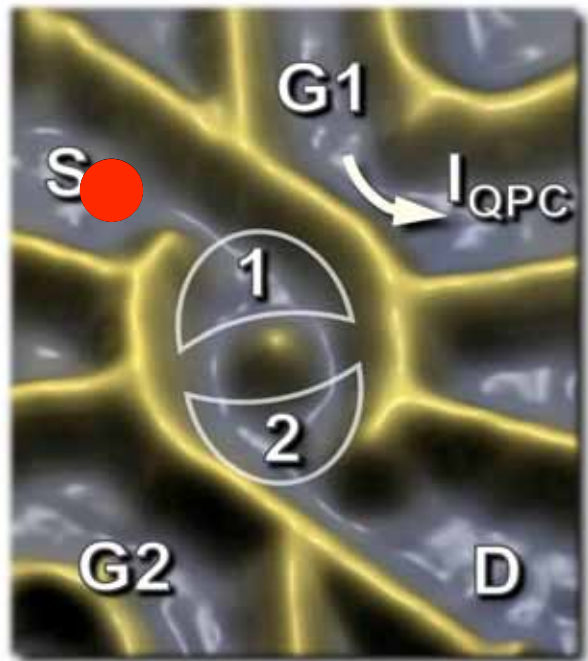
# Aharonov-Bohm with cotunneling

## Co-tunneling

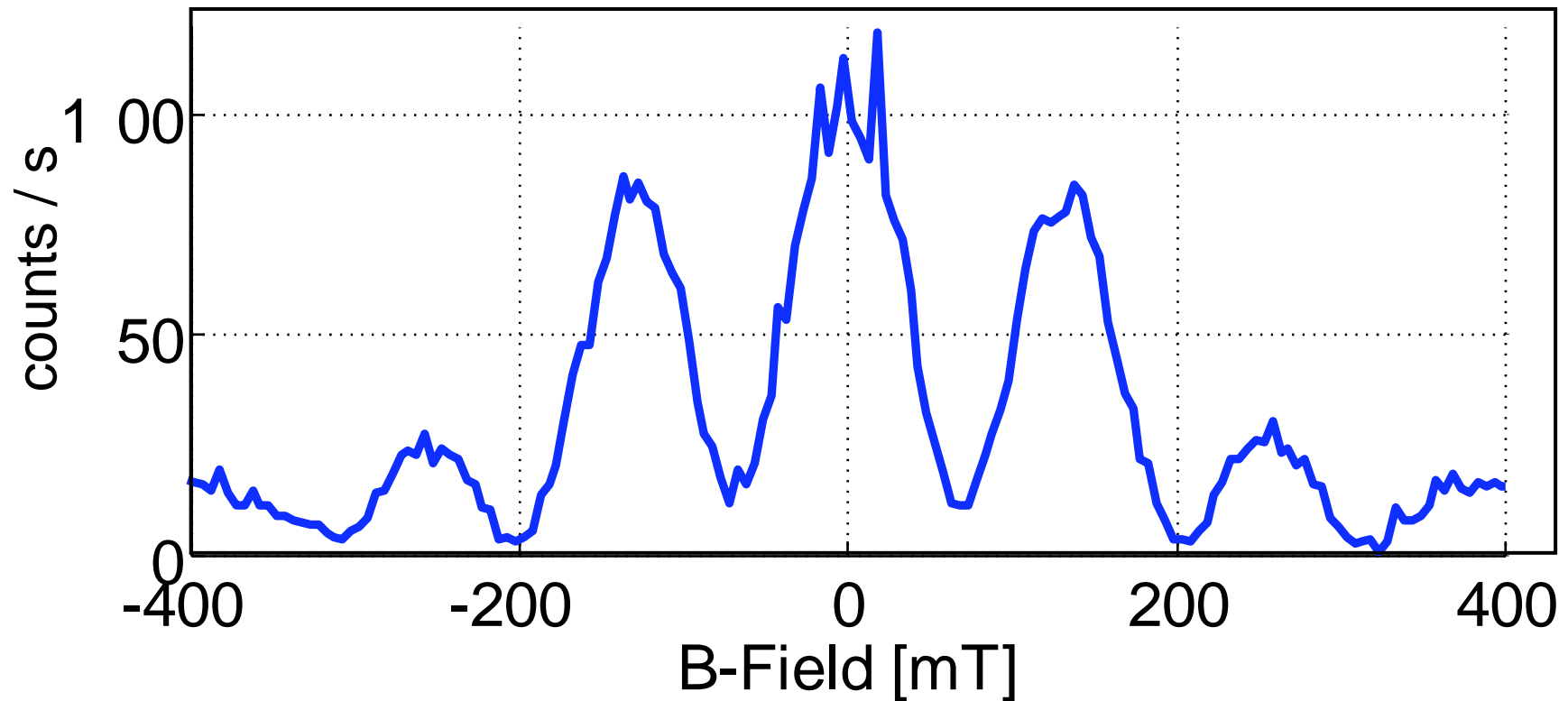
- Electrons are injected from the right lead
- They pass through either the upper or lower arm
- The interference takes place in the left QD



# Double slit experiment <-> Aharonov Bohm



# Aharonov-Bohm oscillations



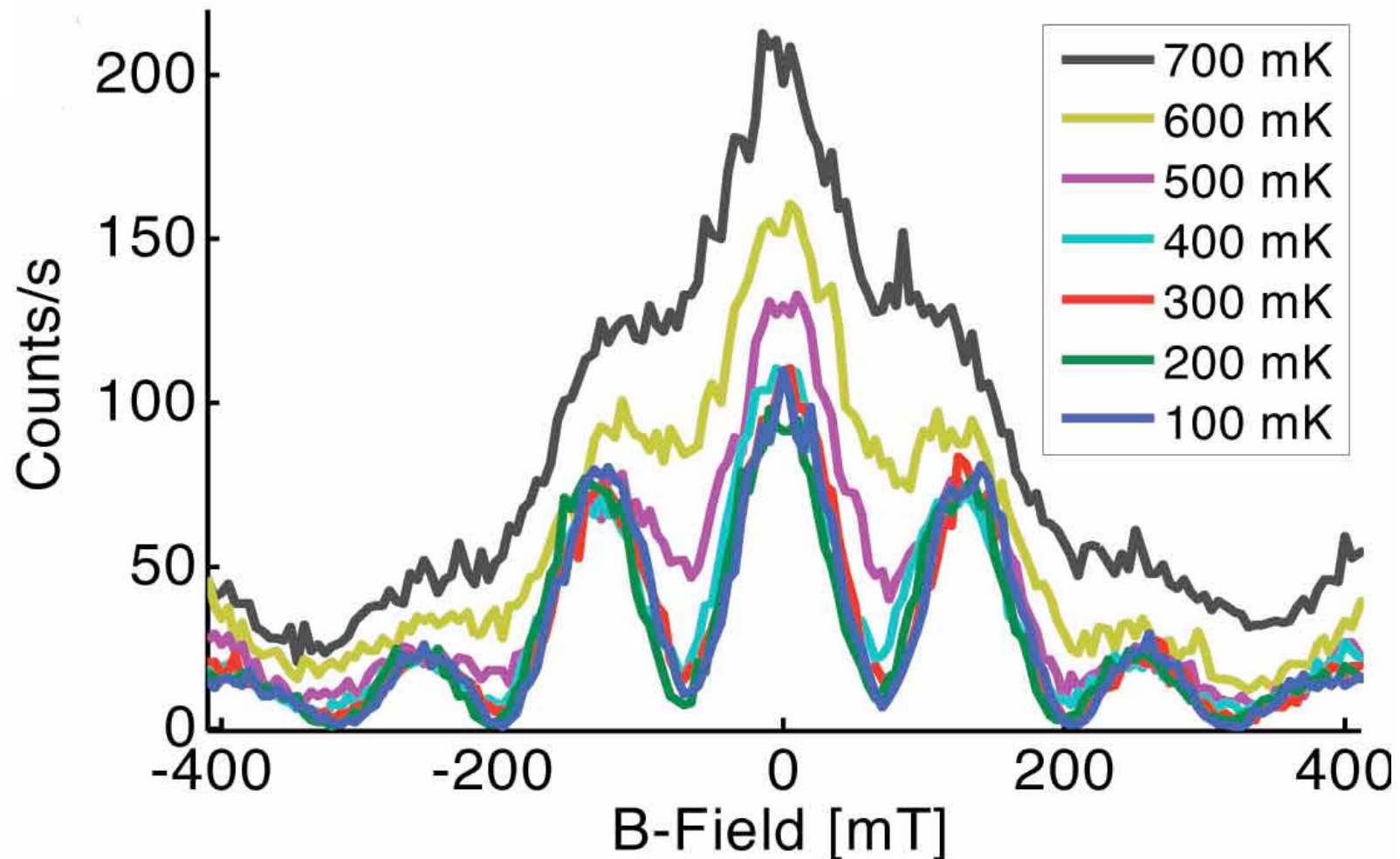
huge visibility! >90%

little decoherence - > due to long dwell time in the collecting dot?

requires the couplings of upper and lower arm to be well symmetrized



# Temperature dependence

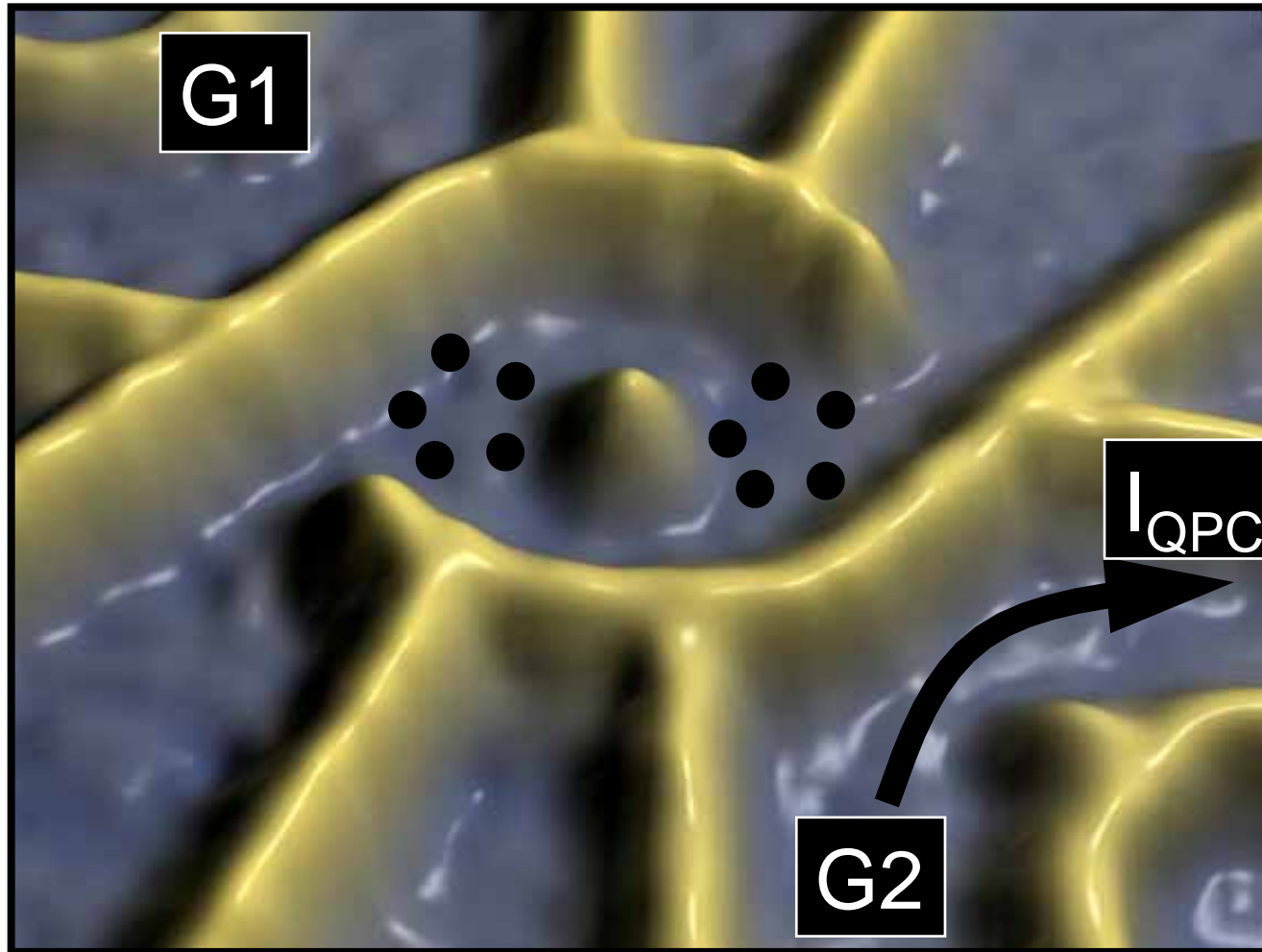


AB amplitude stable below  $T=400\text{mK}$

Destruction most likely due to thermal broadening

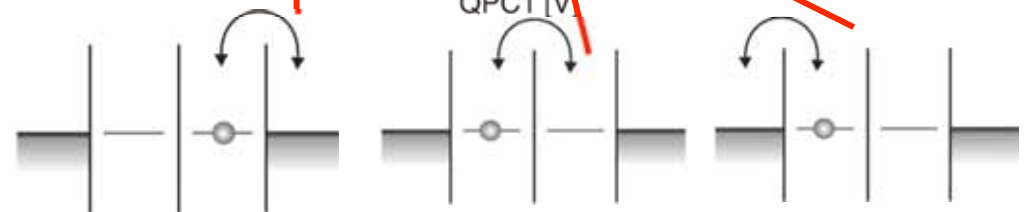
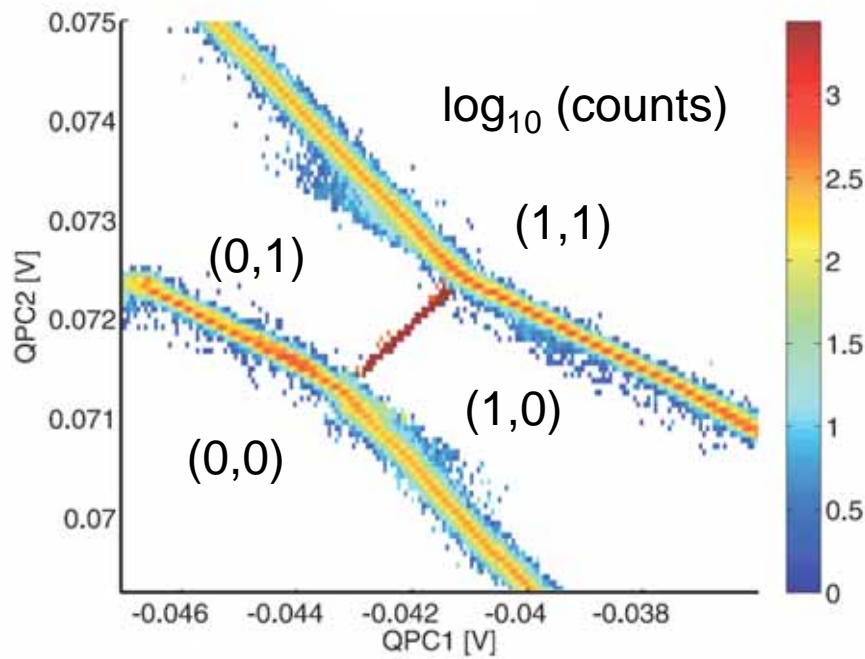
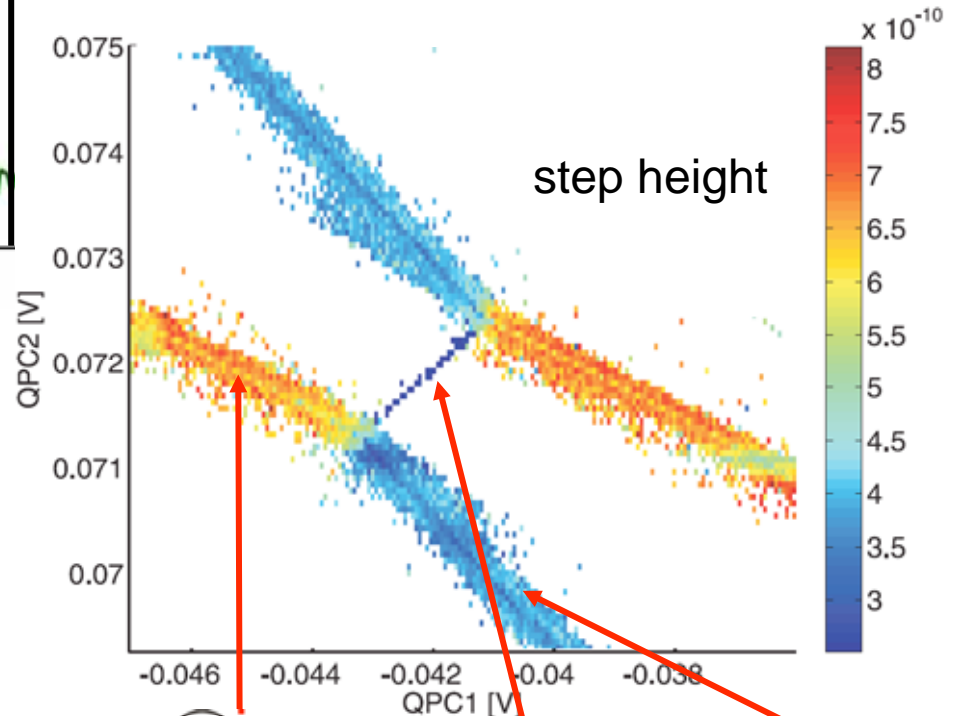
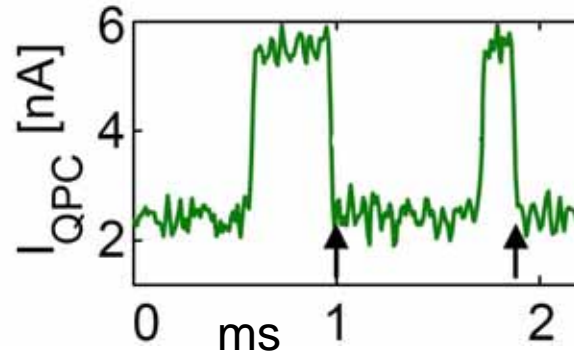
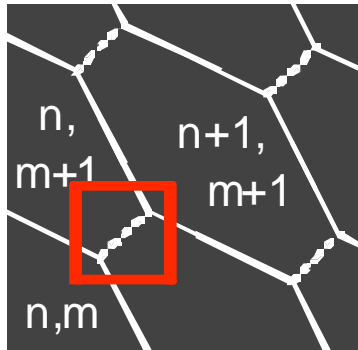


# Double quantum dot in a ring



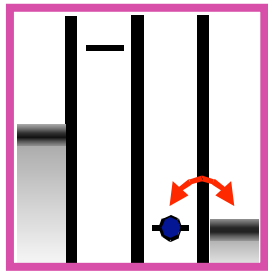
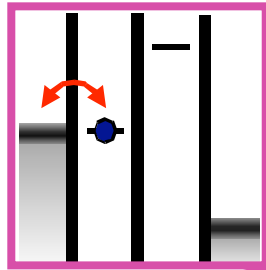
characterization of tunnel rates

# Charge stability from counting

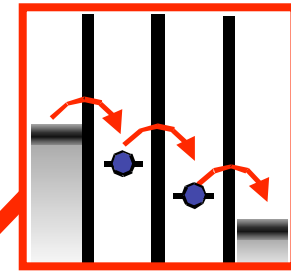
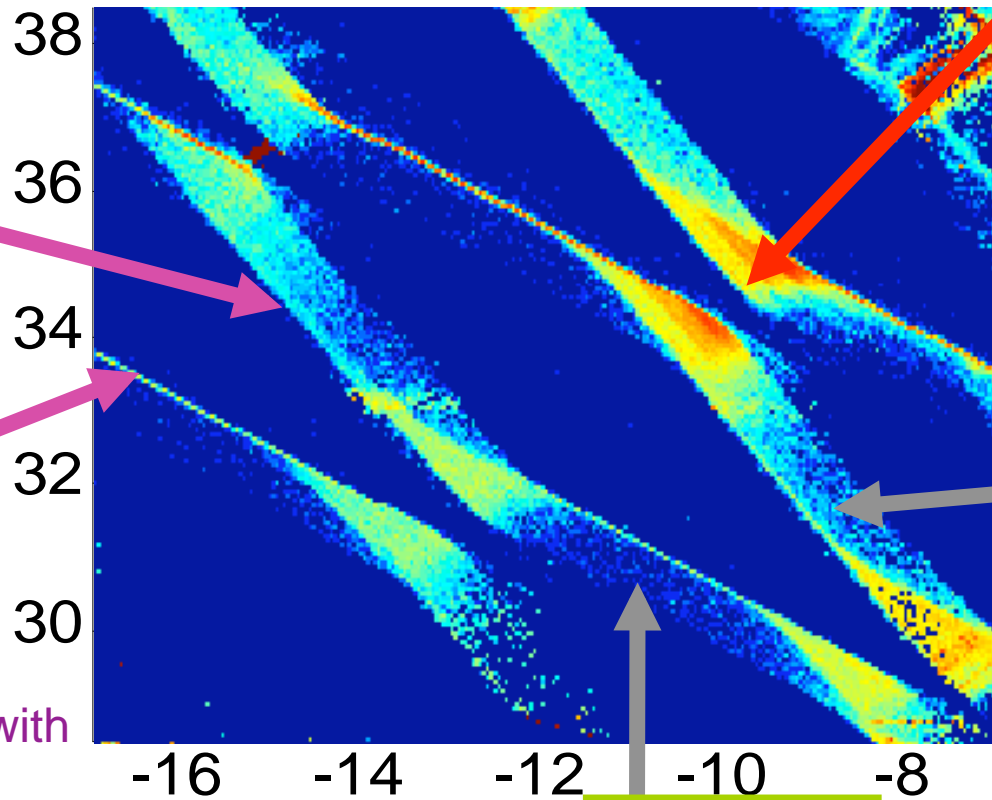


# High bias regime

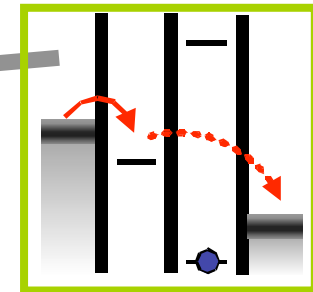
Counts/s



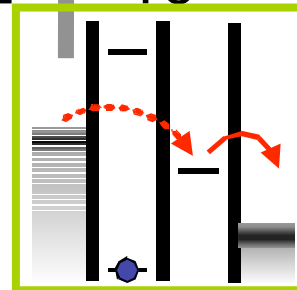
dot states aligned with source or drain



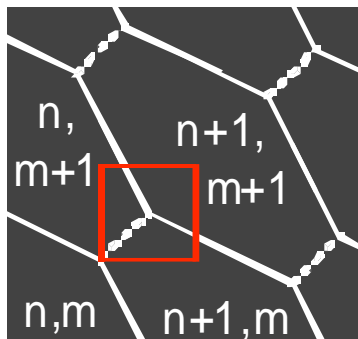
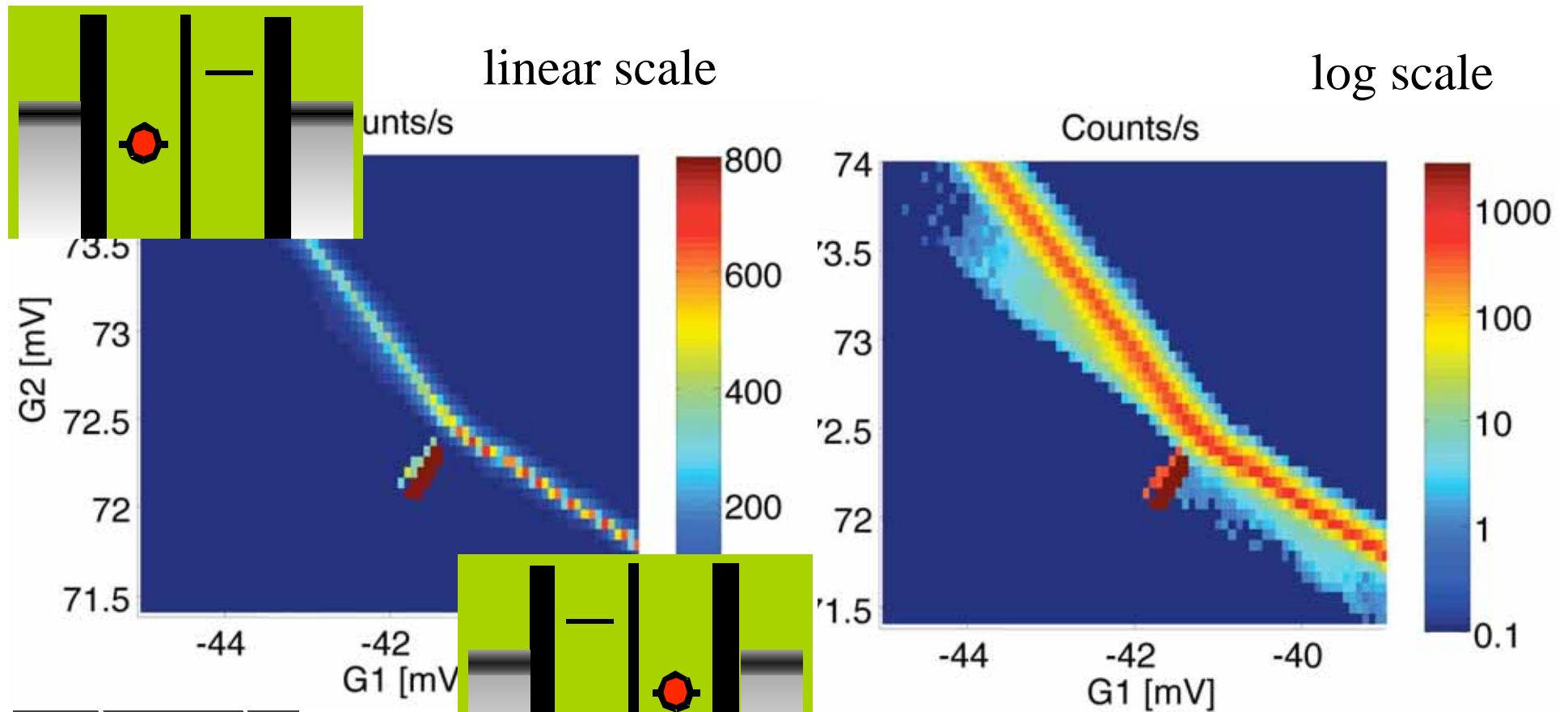
sequential tunneling



cotunneling

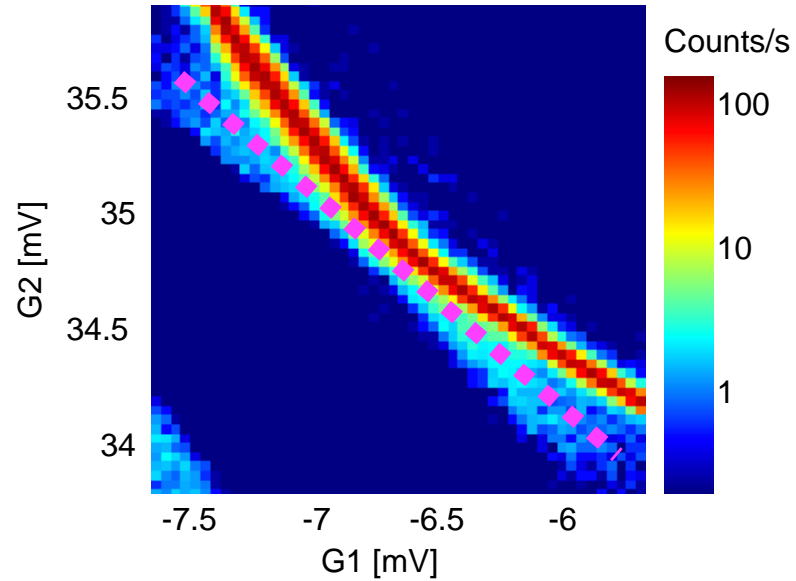


# Triangles at zero bias across dot

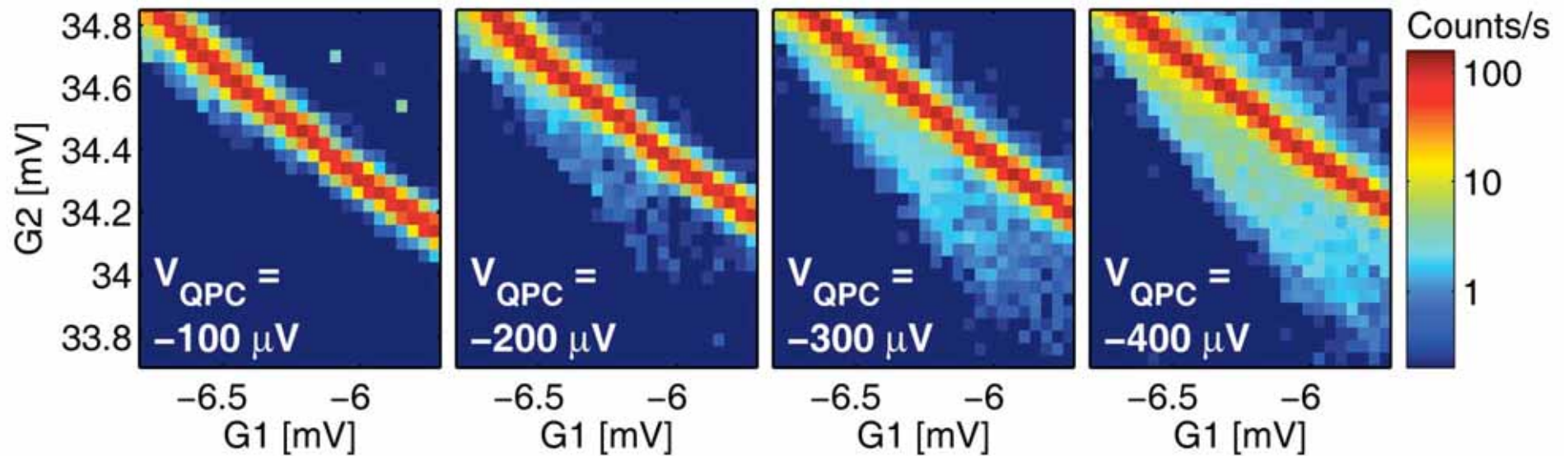


$$V_{\text{QPC}} = 300 \mu\text{V}$$

# Different biases across the QPC

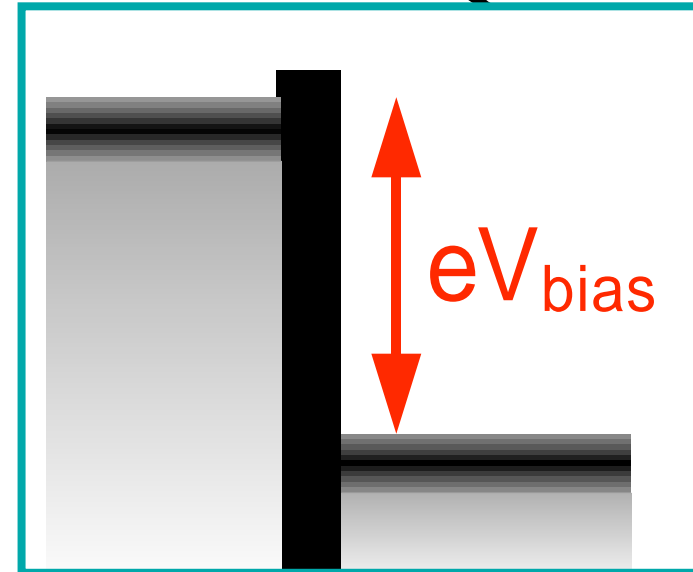


The triangles grow with increasing bias



# Microwave emission of a QPC

- Voltage biased tunnel junction
- Emission spectrum
  - Linear increase with bias
  - Cut-off at  $f=eV_{\text{bias}}/h$



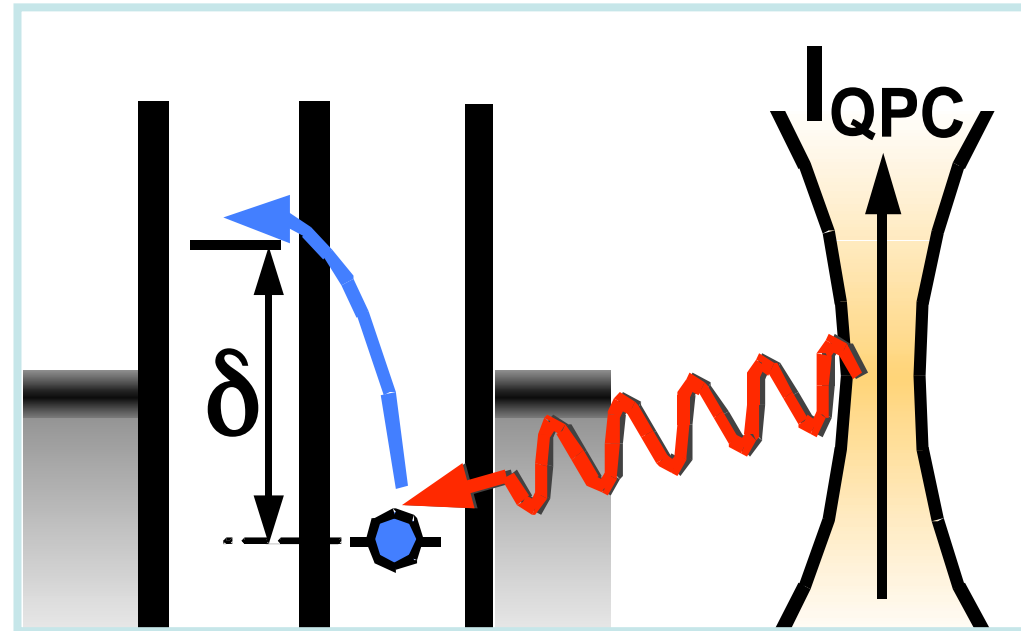
$$S_I(\omega) = \frac{4e^2}{h} T(1-T) \frac{eV - \hbar\omega}{1 - e^{-(eV - \hbar\omega)/k_B T}}$$

spectral noise density for the emission side ( $\omega > 0$ )

R. Aguado and L. Kouwenhoven,  
PRL **84**, 1986 (2000)

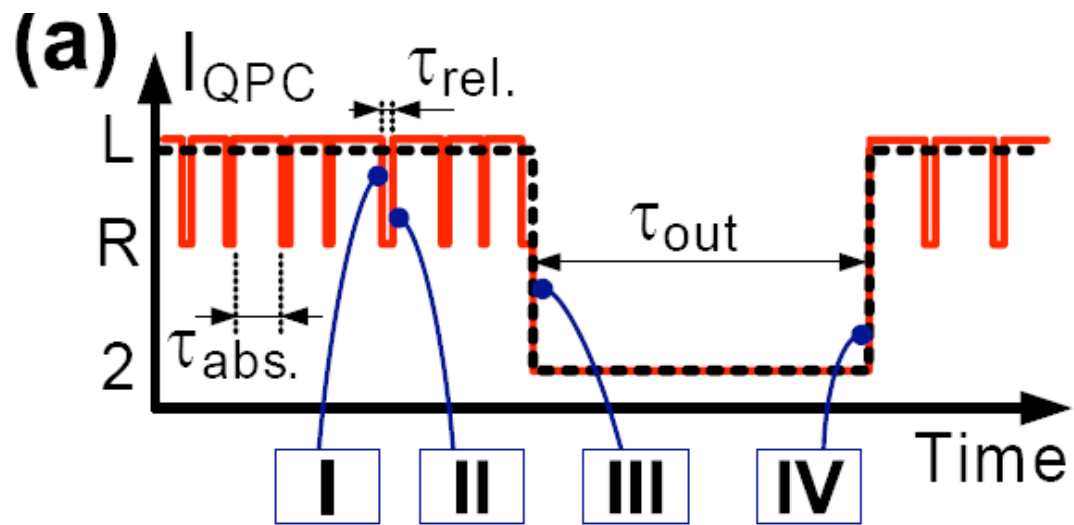
# Tunable noise detector

- The detuning of the quantum dots acts as a selective frequency filter
- The detuning is easily changed with gate voltages

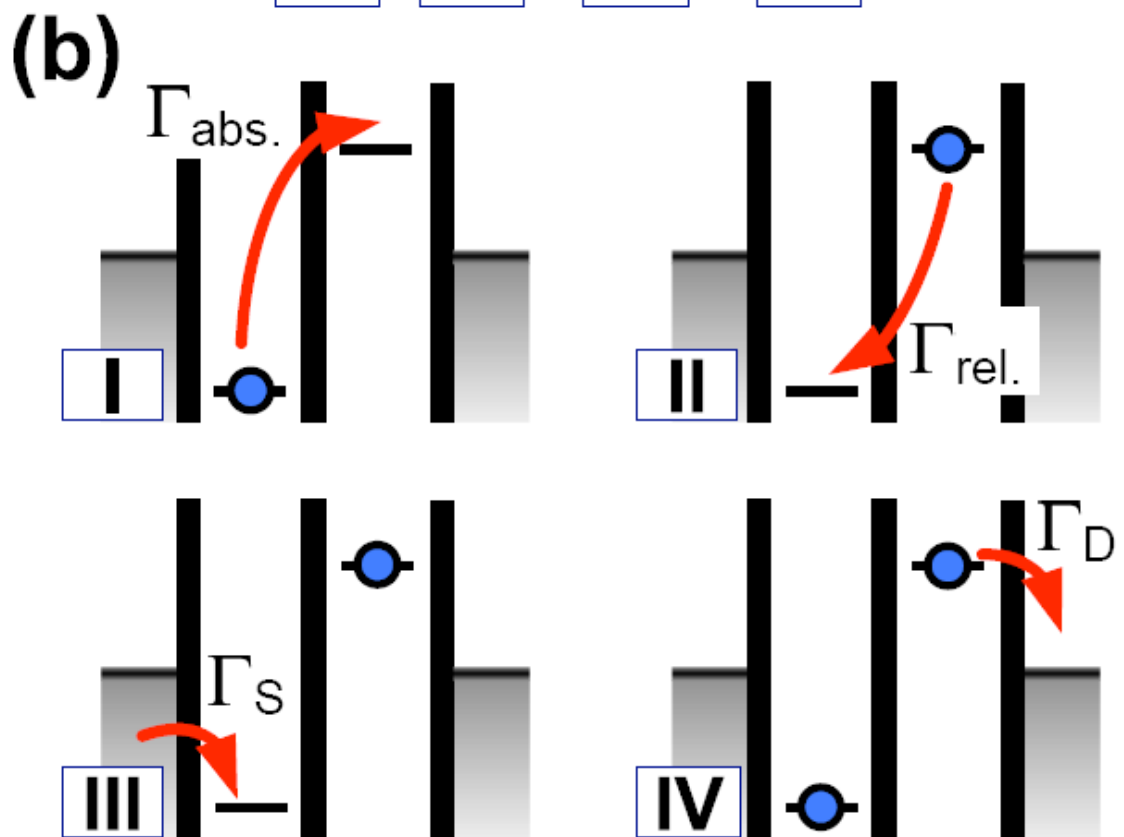


R. Aguado and L. Kouwenhoven,  
PRL **84**, 1986 (2000)





detector signal



absorption process

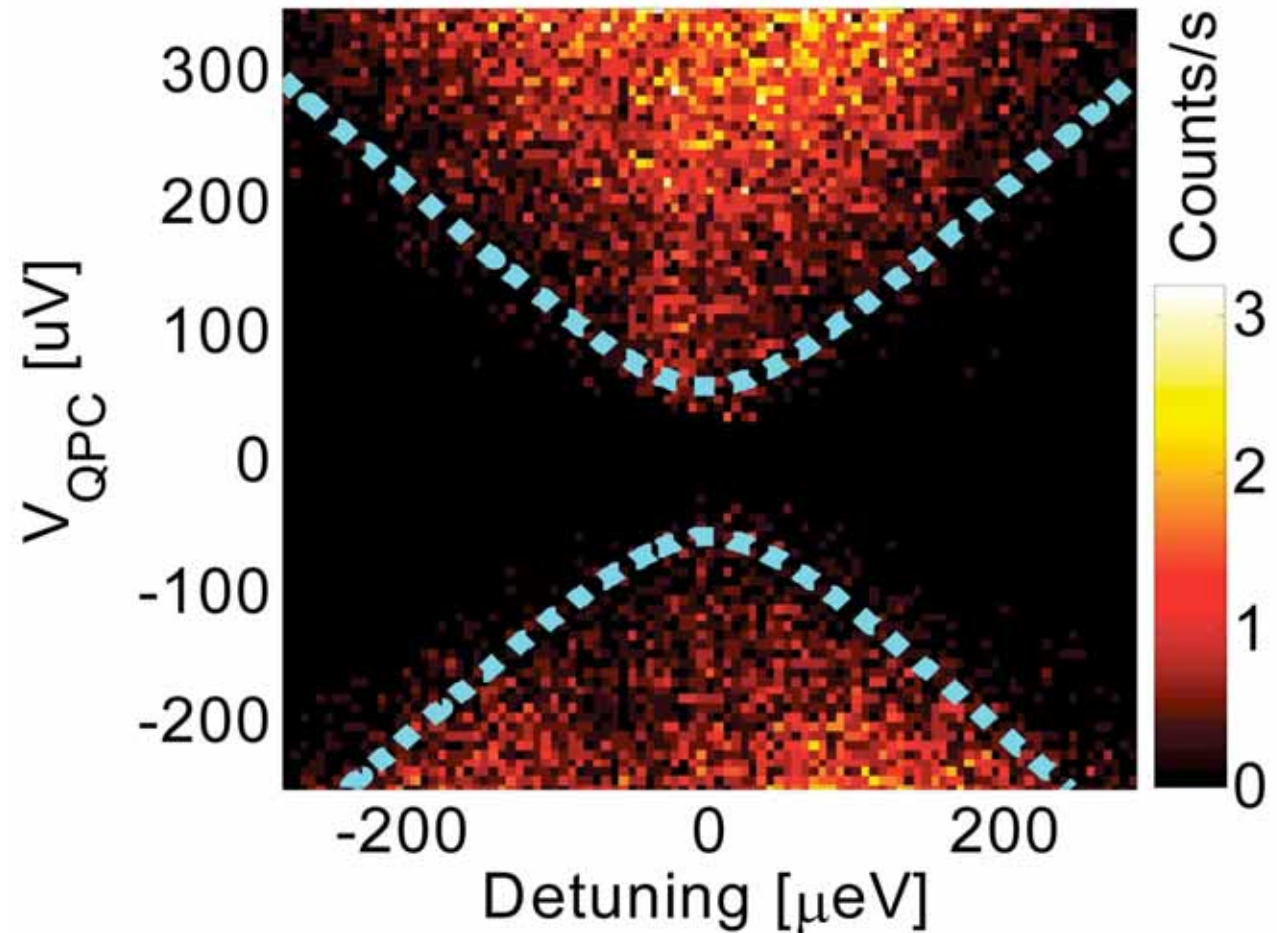
# Double dot detuning vs. QPC bias

- Level separation of the DQD dashed line:

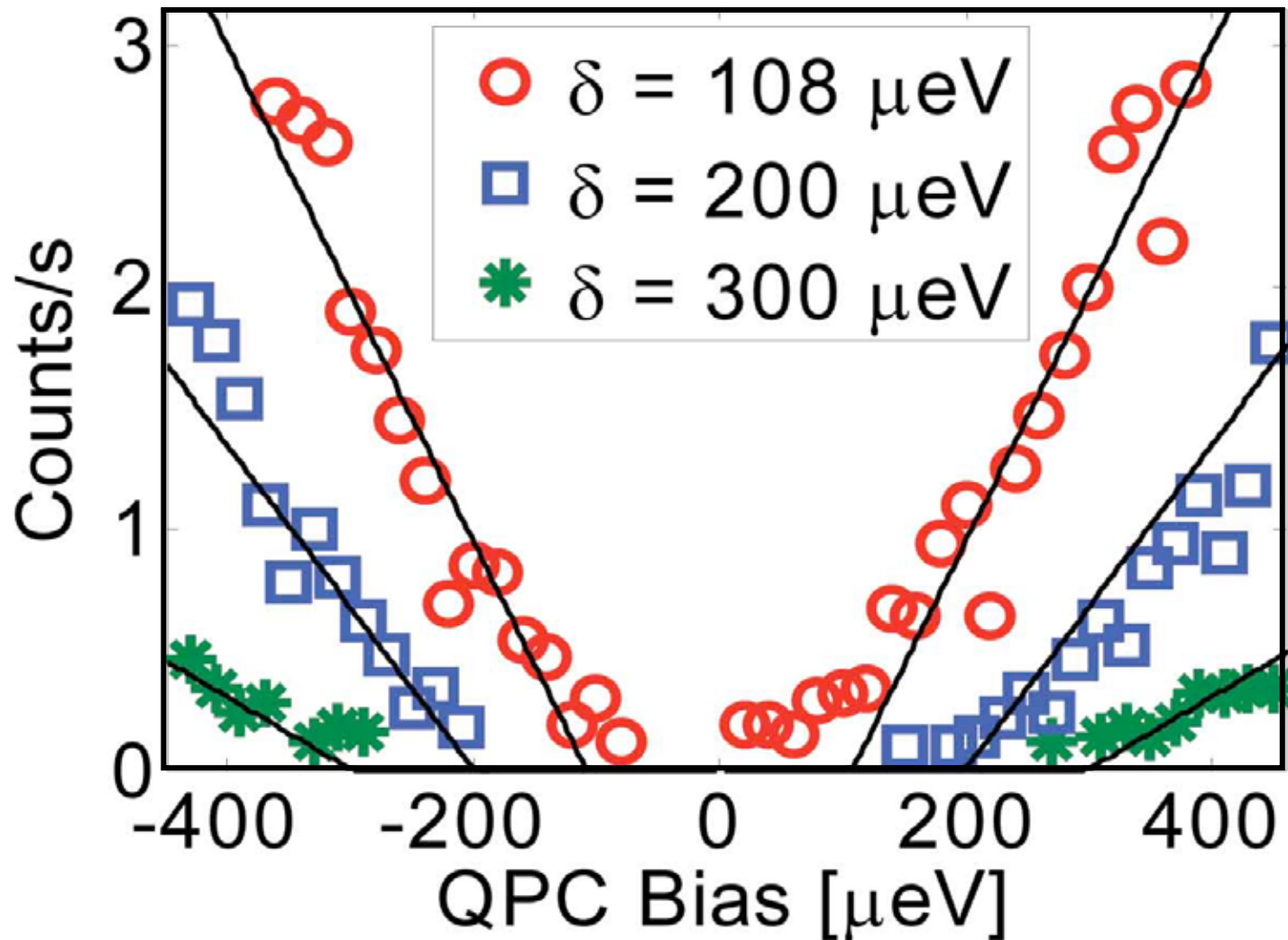
$$\Omega = \sqrt{4t^2 + \delta^2}$$

- No counts in the region with  $eV_{\text{QPC}} < \Omega$ !

$t$  : tunnel coupling  
 $\delta$  : detuning



# Bias dependence of the count rate

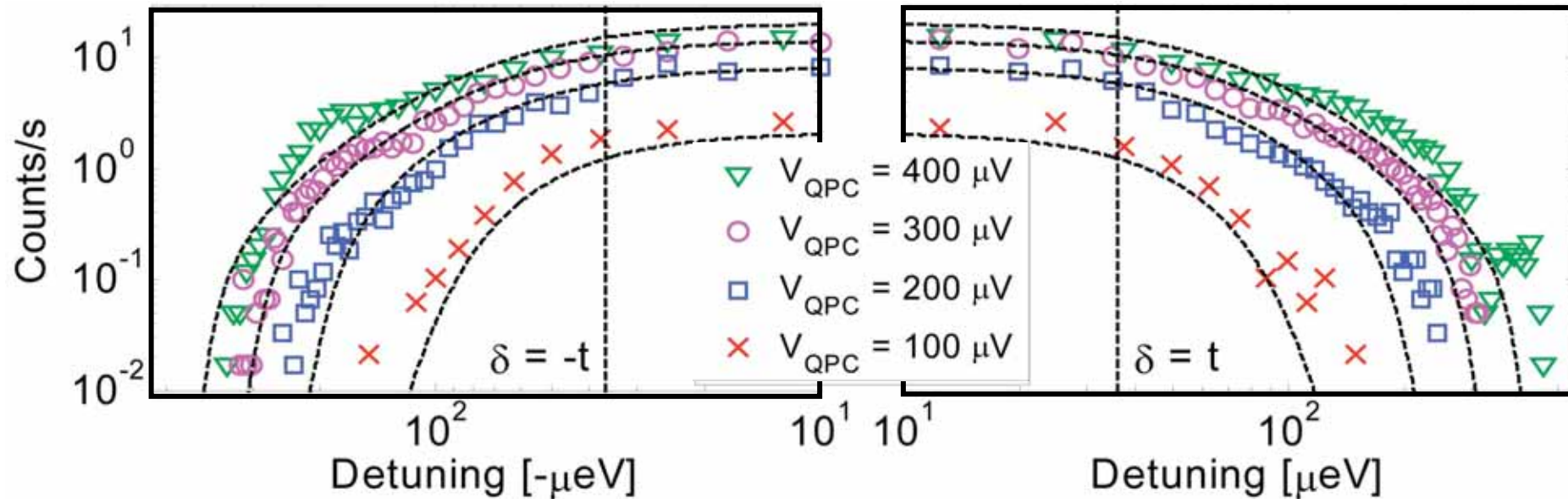


- Linear increase of absorption rate as soon as  $eV_{\text{QPC}} > \delta$

# Measuring the spectrum

absorption rate of the DQD in the presence of the QPC:

$$\Gamma_{abs} = \frac{4\pi e^2 \kappa^2 t^2 Z_l^2}{h^2} \frac{S_l(\Omega/\hbar)}{\Omega^2}$$



$\kappa$ : capacitive lever arm of QPC on DQD

$Z_l$ : zero frequency impedance of leads connecting QPC to voltage source

Clear cut-off at  $\delta = eV_{QPC}$

Gustavsson et al., PRL **99**, 206804 (2007)

# Single photon detection by a quantum dot

**quantum optics**

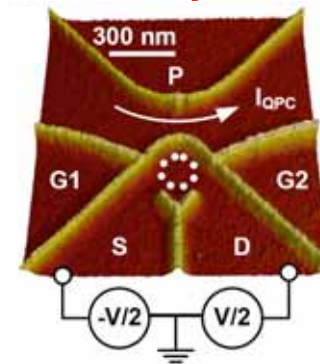
wave length  
of photon:  
500 nm



size of atom:  
1 nm

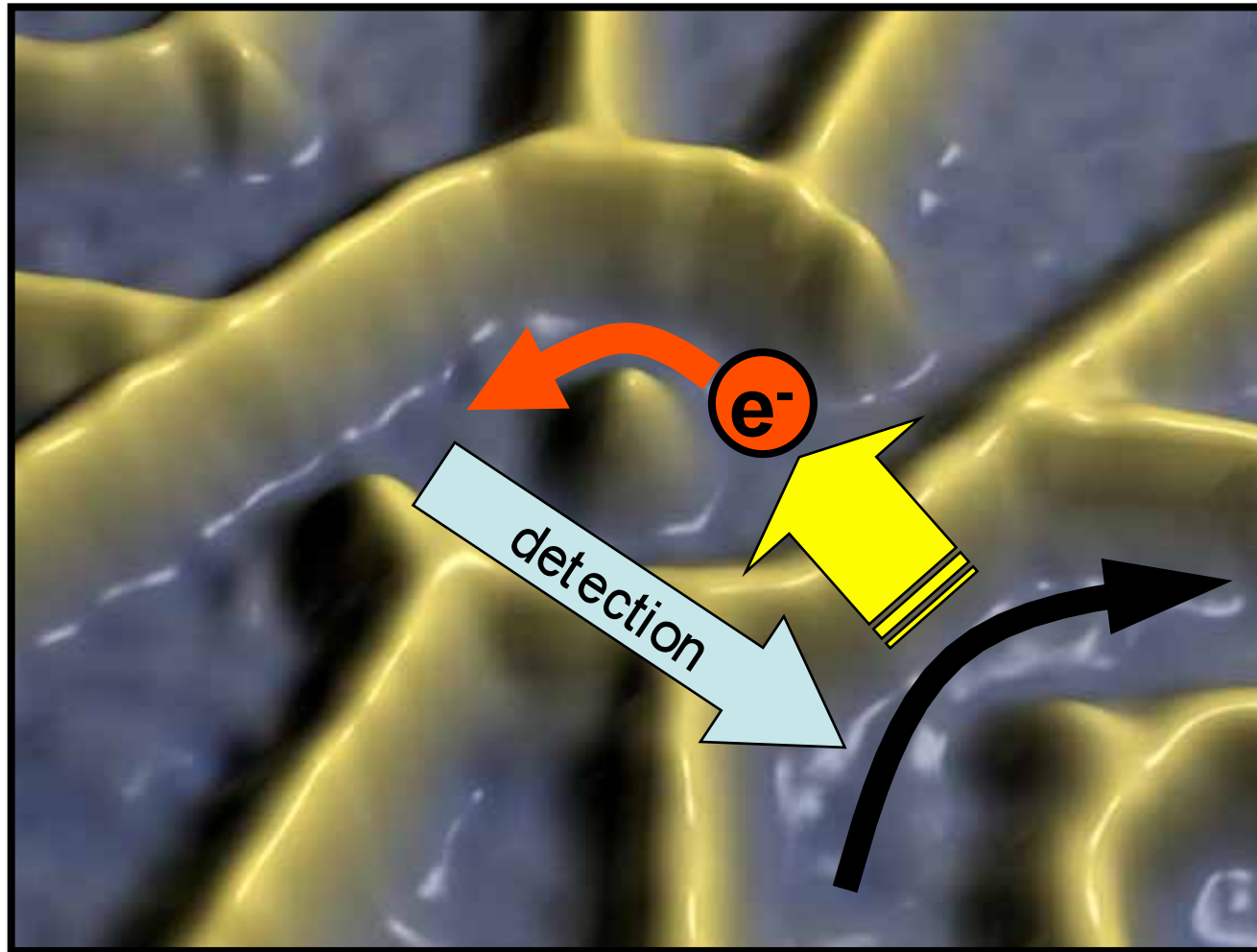
**semiconductor  
nanostructures**

wave length  
of photon:  
10 μm



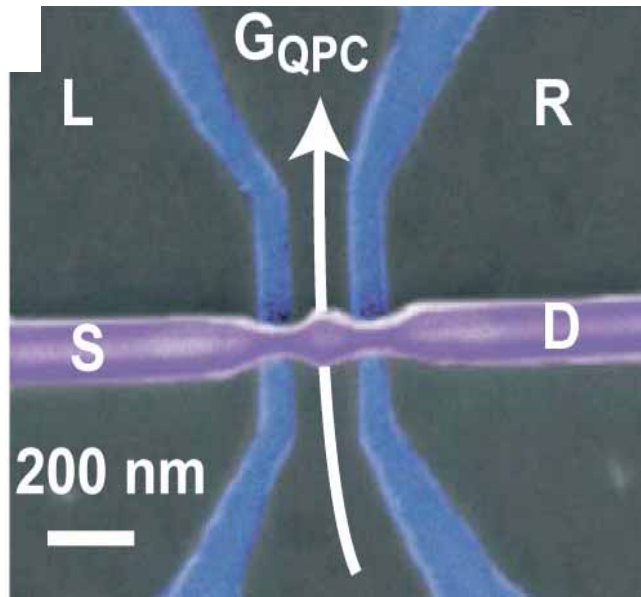
size of  
quantum dot:  
100 nm

# Single-photon, single-electron detection



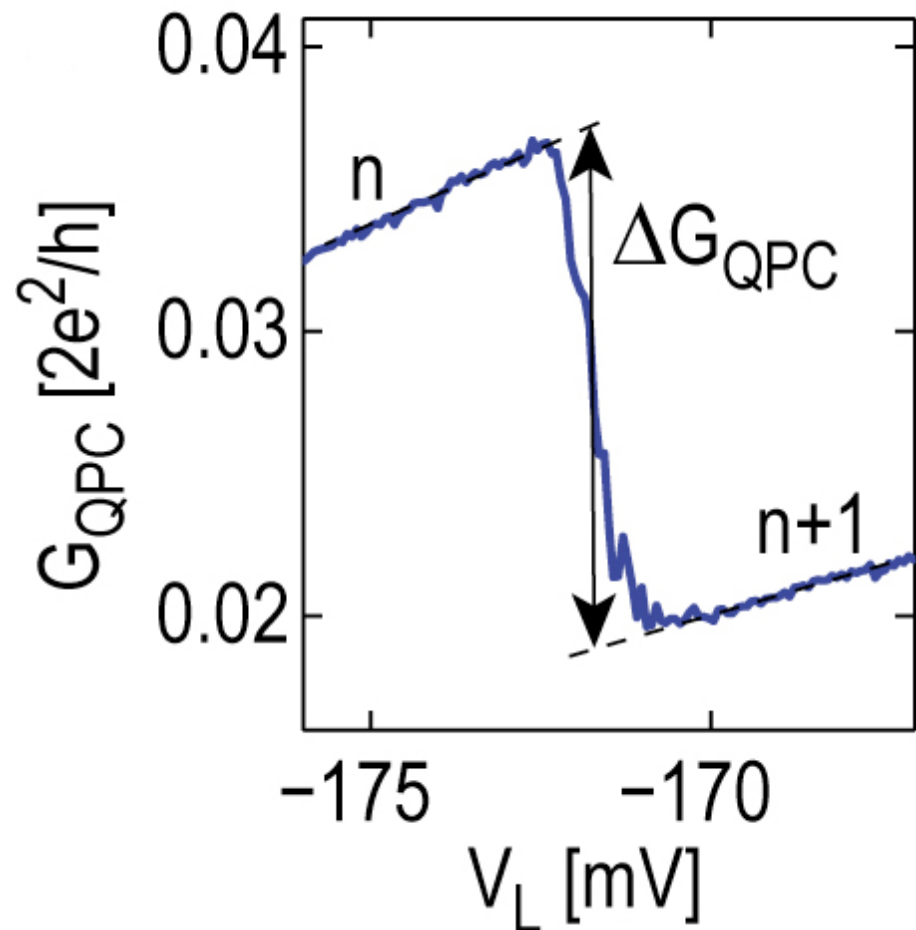


# Towards THZ photons – InAs nanowire dots



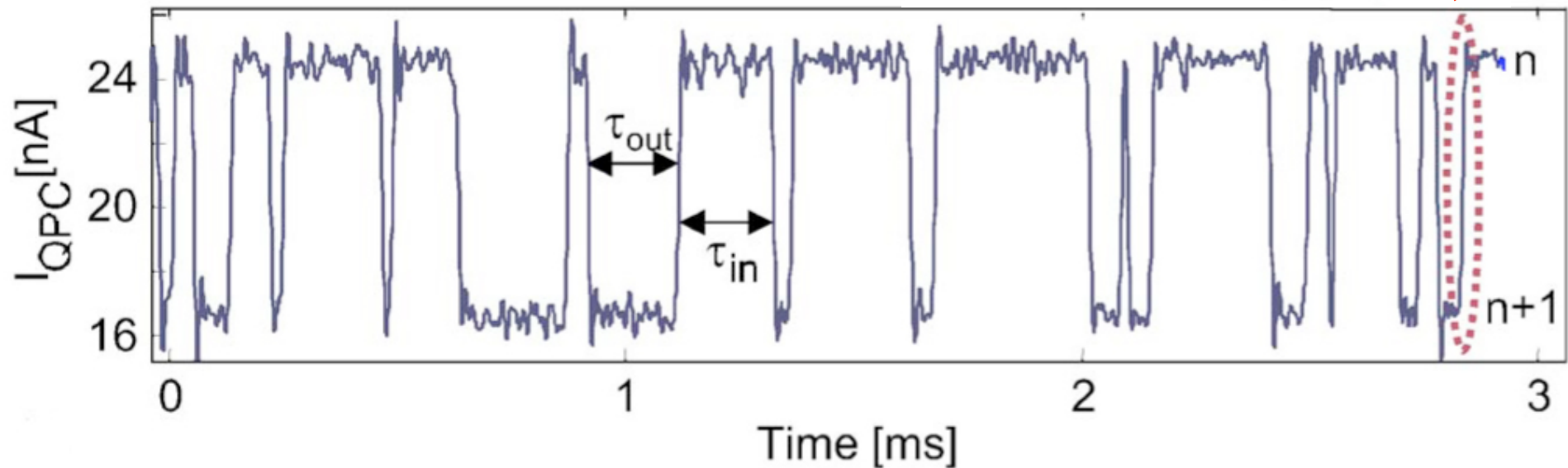
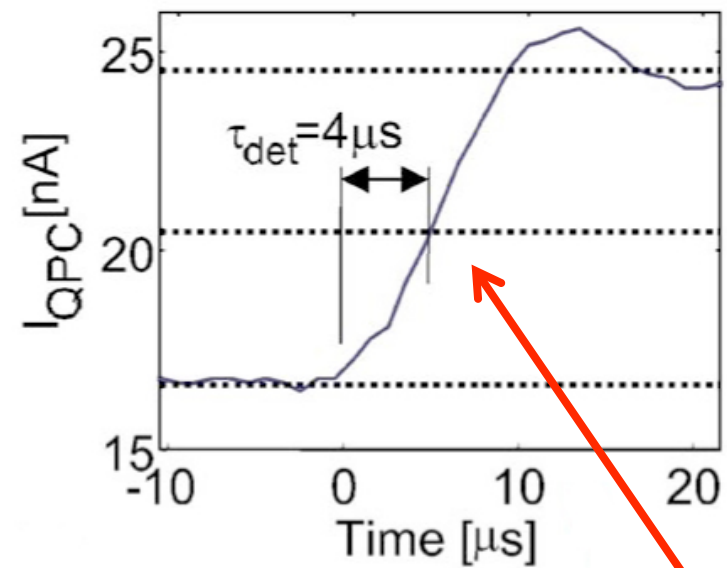
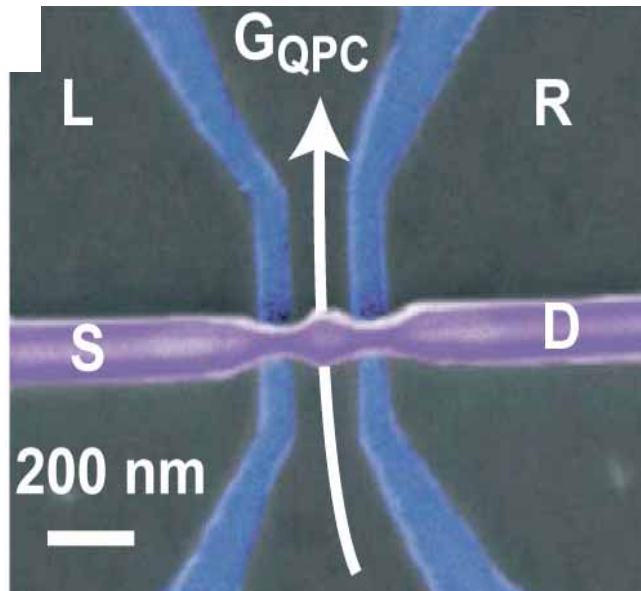
Strong coupling between dot (InAs) and detector (GaAs 2DEG)

up to 50%  
detector signal





# Time-resolved charge detection in InAs nanowire dots



Simon Gustavsson



# Thank you

Renaud Leturcq



Ivan Shorubalko



Thomas Ihn



## Plans:

- time resolution
- correlation experiments
- spin blockade
- graphene