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Workshop on the original of P, CP and T Violation

2 - 5 July 2008

Experimental searches for mirror matter

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Experimental searches for mirror matter

S.N. Gninenko INR Moscow

Workshop on origin of P, CP and T violation July 2-5, 2008 Trieste, Italy

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- Experimental searches for mirror matter -

mirror matter model

□ searches at low energies

- Ps Ps` oscillations (focus of the talk)
- n-n` transitions (Berezhani, this workshop)

□ searches at LHC

- H-H`mixing (Chacko, this workshop)

□ Summary

P violation: two classes of models

Great mistery: why only left-handed fermions feel week interactions?

Wu et al.'56: decays of polarized ${}^{60}Co > {}^{60}Ni e - v$.

Left-right symmetric models-

Parity restoration at high energy scale > 1 TeV, New heavy right handed W_R LHC can probe W_R mass up to 3 TeV, (N.V.Krasnikov, this Workshop) however, already high limits on M(W_R): > 2.5 TeV from $\Delta M_{K/B}$ and > 4 TeV from K0-decays (R.Mohapatra,X.Ji, this Workshop)

□ Mirror matter model -

So far no data confronting the model Effects of parity restoration can be directly observed at low energies in a table top experiment (a'la Wu)

mirror matter (mm) model

\Box old idea: P -> CP -> CPA symmetry

Nature is intrinsically L-R symmetric with L-R particle

- properties exchanged: V-A->V+A
- New CPA-sector must be hidden; connected to our world by gravity

Kobzarev, Okun, Pomeranchuk'65; Lee, Yang'56; Pavsic'74;

Blinnikov, Khlopov'83.

modern mm model -

based on minimal symmetry:

 $(SU(3)_C \times SU(2)_L \times U(1)_Y) \times (SU(3)_{CM} \times SU(2)_R \times U(1)_{YM})$ SM fermions and gauge bosons are accompanied by identical mirror partners

- mm must have different cosmology,
- is a good candidate for dark matter
- can be linked to string theory, extra dimen., ...

Foot, Volkas'91 Berezhiani, Mohapatra'95, Berezhiani et al.'00-08, Akhmedov, Senjanovic'92,

For review see: L.Okun hep-ph/0606202

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ordinary -mirror particle interaction

Conservation laws for ordinary matter and mirror matter prevent particles with colour and charge from interacting between two sectors.



But, interaction between colourless neutral particles is allowed: ordinary and hidden sectors can communicate trough mixing of H-H^{*}, kinetic mixing of photons, mass mixing between neutrinos, neutrons, etc...(Z.Berezhiani, this Workshop)



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possible mm effects:

Higgs mixing

Ignatiev, Volkas '01; Barbieri et al. '05, Wilczek' 07, Li et al. '07.....

Ps-Ps` oscillations

Glashow '86, SG '95, Foot, SG '01; Atoyan et al. '89, Mitsui et al. '95, Badertsher et al. 07

\Box n-n` mixing

Berezhiani, Bento '05; Pokotilovski '06, Ben et al.(PSI) '07, Serebrov et al. (PNPI) '07, Mohapatra et al. '05.

dark matter

DAMA '05; DAMA/LIBRA '07, Foot '01-07; Ignatiev,Volkas'03, Mitra'03-06,...

\Box v-v` mixing

Berezhiani, Mohapatra '95, Foot, Volkas '00; Mohapatra, Nasri '05

cosmology

Blinnikov, Khlopov'82,83, Khlopov'91,00, Berezhiani'95-08, Ciarcelluti'03-05,....

Image: Image:

Holdom '85; Ignatiev '91; Gninenko et al.'07....

□ anomalous events,

Foot, Silagadze'01-05, Foot,Mitra'02-03,...

γ - γ ` kinetic mixing



 \Box new mass eigenstates: oPs± = (oPs±oPs`)/ $\sqrt{2}$

u energy splitting: $\Delta E = 2\epsilon f$, f=8.4x10⁴ MHz from oPs-pPs splitting

□ oscillation probability: $P(oPs-oPs)(t)=sin^2(2\pi\epsilon ft)$

Holdom'86, Glashow '86

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Experimental signatures for oPs-oPs' oscillations

- modification of oPs decay curve very difficult to measure, high statistics required
- OPs`-> 3γ' -> invisible decay more convenient: few events need to be observed

Branchnig ratio in vacuum:

$$Br(oPs \rightarrow invisible) = \frac{2(2\pi\varepsilon f)^2}{\Gamma_{SM}^2 + 4(2\pi\varepsilon f)^2}; t \gg \frac{1}{\Gamma_{SM}}$$

in a target (in presence of collisions):

$$Br(oPs \rightarrow invisible) \simeq \frac{2(2\pi \varepsilon f)^{2}}{\Gamma_{SM}}; t >> \frac{1}{\Gamma_{COLL}};$$

SG'95, Foot and SG '00 Suppression factor

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Positronium decays in the SM

□ parapositronium g.s.: L=0, S=0, pPs-> 2, 4, ... γ , τ ~10⁻¹⁰ s

- □ orthopositronium g.s.: L=0, S=1, oPs-> 3, 5, ... γ , τ ~10⁻⁷ s
- OPs-->2γ forbidden
- Br (pPs-->4γ)~ Br(oPs-->5γ)~10⁻⁶

 \Box e⁺ + e⁻ + M -> γ + M or e⁺ + e⁻ + M -> M* are small

□ Br(oPs--> invisible)<~10⁻¹⁸ is extremely small

Only 𝔅, 𝔅, 𝔅, 2, 3, ¼, ..gamma in the final state. for Br(oPs--> anything)>~10⁻⁵ 9

o-Ps decay rate puzzle (1982-2002) History of oPs decay rate measurements



Discrepancy $\Gamma_{exp} > \Gamma_{SM}$: -unknow contribution at the level Br(oPs->X)~10⁻³ or

- experimental problems

Tokyo measurements (oPs formation in target) and Ann Arbor experiment (oPs formation in vacuum,) agree to each other and also agree with QED predictions



Experimental setup

spectrum from Ge detector

oPs formation and decay in a 0.1 g/cm³ dense target ~10⁴ collisions/lifetime results in pick off rate: e+e- + (e-M)-> e- + (e+e- -> 2γ)+ M ~10⁻² Γ_{oPs} and suppression of oPs-oPs' transitions

Asai et al.'08

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Discrepancy triggers searches for new physics.

all exotic modes oPs-> wrong number of photons : 0,1, 2, 4.... are excluded at the level $Br < ~10^{-5}$. However, measurements are performed whith oPs formation in a target, i.e. oscillations oPs-oPs' are suppressed and cannot contribute through oPs-> invisible, hence no constraint on oPs->invisible through mirror oscillations.

Experiments that agree with QED predictions,

-Tokyo measurements are not sensitive to mirror effect: oscillations oPsoPs' are suppressed and cannot contribute to the decay rate - Ann Arbor vacuum experiment: small cavity size is used and oPs collisions with the cavity walls may dump oscillations.

-In addition.....

more interesting observation:

- two Ann Arbor vacuum experiments with big and small cavities disagree:
 oPs decay rate for big cavity measurements is higher ~10⁻³
- if "acceleration" of oPs decay is due to a SM process, e.g. due to oPs positron annihilation with wrong electron (pick off process) or due to oPs-pPs mixing,

then

- 2γ (511 keV line) must be seen in the final state in the big cavity experiment in addition to the main spectrum from oPs-> 3 γ !

Two Ann Arbor experiments on oPs decay rate in vacuum

 $arGamma_{ ext{exp}} > arGamma_{ ext{SM}}$

$$\Gamma_{\rm exp} = \Gamma_{\rm SM}$$



in the SM oPs can decay faster only due to additional SM 2γ annihilation process, i.e.

- pick-off annihilation, or
- oPs-pPs mixing (external fields)

Search for 511 keV line.

 $arGamma_{ ext{exp}} > arGamma_{ ext{SM}}$





check presence of 2γ (511 keV line) with intensity ~10⁻³ of the total rate

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No 511 keV line - hint for mirror matter ?

VOLUME 66, NUMBER 10 PHYSICAL REVIEW LETTERS

11 MARCH 1991

Direct Search for Two-Photon Decay Modes of Orthopositronium

D. W. Gidley, J. S. Nico, and M. Skalsey Department of Physics, University of Michigan, Ann Arbor, Michigan 48109 (Received 13 December 1990)

A direct search for γ rays from the forbidden decay o-Ps \rightarrow 2y has been performed using the same apparatus that recently measured an anomalously high o-Ps decay rate. Using a high-resolution Ge γ -ray detector, a 233-ppm limit is set on the branching ratio to a pair of 511-keV γ rays, and a 200-ppm limit is set on the branching ratio to a pair of 512-keV. Hence, these 2 γ modes cannot be responsible for the o-Ps decay-rate discrepancy between theory and experiment of 1400 \pm 230 ppm.

PACS numbers: 36.10.Dr, 11.30.-j, 12.20.Fv



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What size ε , if oPs escapes to mirror world?

 $\Delta\Gamma(oPs \rightarrow invisible) = \Gamma_{EXP} - \Gamma_{SM}$ $\Delta\Gamma(oPs \rightarrow invisible) \simeq 10^{-3}$ $\Delta\Gamma(oPs \rightarrow invisible) = \frac{2(2\pi\varepsilon f)^2}{\Gamma_{ov}\Gamma_{ov}}$ $\Gamma_{COLL} \simeq 3 \Gamma_{SM}$ $\varepsilon \simeq (5 \pm 1) \times 10^{-7}$

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Existing limits on $\boldsymbol{\epsilon}$

□ $\epsilon < \sim 10^{-5}$ from milli-charged particle searches SLAC'01

□ ε< 10⁻⁸ Tokyo group claim. However, suppression due to collisions are not considered.
 ε< ~10⁻⁶ after correction. → Br(oPs->inv)<~ 10⁻²

Mitsui et al. '93; SG '95;

 $\Box \ \varepsilon < 3x10^{-8} \text{ from BBN} \qquad \Longrightarrow \ Br(oPs->inv) <~ 10^{-6}$

Glashow, Carlson'87;

□ $\epsilon \sim 4x10^{-9}$ from DAMA, DAMA/LIBRA observation of modulation \implies Br(oPs->inv)~ 10^{-7}-10^{-8} Foot'01-08;

Program of experiments

- search for oPs escape into extra dimensions
 sensitivity in Br~ 10⁻⁸ 10⁻⁹
- 2. measurements of free gravity fall of antihydrigen (and positronium), very cold < 10 mK Rydberg oPs formation is required</p>
 - tests of mirror matter effect in oPs decays well fit into this program

Moscow experiment on o-Ps->invisible

"A search for photonless annihilation of orthopositronium," Atoyan, SG, Razin, Ryabov, Phys. Lett. B 220, 317 (1989). LEP Z->*invisible*

- Radioisotope Ps source (²²Na)
- Generate trigger on ²²Na decay (positron + 1.2 MeV photon)
- Detect energy of all e+ annihilations
- Subtract p-Ps -> 2γ events in Nal spectrum
- "Difference of two large numbers" problem
- Statistics, background limited





Fig. 1. Schematic view of the set-up: (1): Nal calorimeter; (2): Nal counter (3): target; (4): proportional counter; (5): the positron source ²²Na.

$$\frac{\Gamma(O - Ps \rightarrow nothing)}{\Gamma(O - Ps \rightarrow 3\gamma)} < 5.8 \times 10^{-4}$$

Cannot be resnonsible for o-Ps decay rate anomaly.

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Tokyo experiment on o-Ps -> invisible

- Radioisotope source: 22Na
- Composite trigger: (e+) & (1275 keV)
- 840 kg calorimeter mass
- "High resolution" 1275 keV trigger
- Statistics, PMT noise limited

T. Mitsui et al., PRL 70, 2265 (1993).

$$\frac{\Gamma(O - Ps \rightarrow \text{nothing})}{\Gamma(O - Ps \rightarrow 3\gamma)} < 2.8 \times 10^{-6}$$



FIG. 1. Schematic of the experimental setup.

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ETH-INR-LAPP experiment on o-Ps -> invisible



oPs target



Photograph of the calorimeter (assembling phase)



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Picture of the lab





Signal selection

Charged particle VETO



Rejection of shake-off electrons

The atomic shell electrons ejected in the EC process are a source of background. The ejection probability decreases strongly with the energy of the ejected electrons thus a cut on the energy deposited in the fiber can suppress this background to the required level of 10⁻⁰.

The probability for the atomic shell electrons to be ejected in the EC process was measured as a function of the energy deposited in the fiber.



Fiber energy spectroscopy

The energy is read with the FBGO through the aerogel and a hole in the wrapping using the BGO as a light guide.



Data selection



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Background estimation

The table summarizes the expected background for the experiment estimated from the simulation and the measurement of the shake-off probability

BACKGROUND SOURCE	EXPECTED LEVEL	
Hermiticity Dead Material Resolution	< 10 ⁻⁹	
Absorption in trigger Energy window	< 10 ⁻⁹	
MS positron 546 keV	< 10 ⁻⁹	Most dangerous
MS positron 1.83 MeV	< 10 ⁻⁹	background
Compton EC photon	< 10 ⁻⁹	1
Accidental noise and EC photon	3.2×10^{-11}	1
Source contamination and EC photon	< 1.6 × 10 ⁻¹⁰	
Shake-off electrons in EC process	$\simeq 10^{-8}$ (for 140 keV threshold)	
Physical backgrounds	10-10	
Total	≃ 10 ⁻⁸	

Results A.Baderscher et al., PRD(2007)

DATA	Air	Nitrogen	Combined
Fiber triggers	0.6×10^{10}	0.79×10^{10}	1.39×10^{10}
Selected events	0.61×10^{8}	0.8×10^{8}	1.41×10^{8}
o-Ps fraction	3.41 %	5.29 %	4.48 %
Number of o-Ps	2.08×10^{6}	4.23×10^{6}	6.31×10^{6}

After the selection cut one can perform the sum of the

total energy in the calorimeter

 $E_{tot} = \sum_{i}^{all} E_i - E_{TBGO}$

Data taking period: 4.5 months 1.39x10¹⁰ triggers



Since no event is observed in the signal region, this result provides an upper limit on the o-Ps -> invisible

 $Br(o - Ps \rightarrow invisible) = 2.3/(N_{o-Ps} \cdot \epsilon) \le 4.2 \times 10^{-7}$

$$Br(p - Ps \to invisible) = 2.3/(N_{p-Ps} \cdot \epsilon) \le 4.3 \times 10^{-7} (90\% \text{ C.L.})$$
$$Br(e^+e^- \to invisible) = 2.3/(N_{e^+e^-} \cdot \epsilon) \le 2.1 \times 10^{-8} (90\% \text{ C.L.})$$

factor ~7 better than Tokyo result

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New limit on $\boldsymbol{\epsilon}$

α ε< (1.6 - 3)x10⁻⁷

Main uncertainty from collision rate of oPs in the aerogel target: $\varepsilon \sim (5 \pm 1) \times 10^{-7}$ cannot be reliably excluded. still could be Br(oPs->inv)~ 10⁻³

Badertsher et al.'07;

□ vacuum experiment is needed

How to search for hidden world with oPs?



Cortesy New Scientists, 2004.

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Search for mirror matter via o-Ps->invisible decays in vacuum

"An apparatus to search for mirror dark matter via the invisible decay of orthopositronium in vacuum," SG'03, Badertscher *et al.'03*





cold oPs formation in vacuum: minimum collisions, enough time for oscillations, minimize leak through entrance window

- □ e+ tagging system: timing coincidence of e+ bunch and MCP signal from secondary e-'s emission, inefficiency < 1.e-8
- pulsed slow positron beam: high efficiency & compression factor
- hermetic calorimeter in magnetic field
- very thin vacuum pipe
- □ expected sensitivity in Br(oPs->inv) ~ 10⁻⁸

cold Ps formation target





Ps formation in porous Si films to minimize the leak.



Figure 4 Typical Ps lifetime spectra for a film with open porous network (black) and after capping (red).



FIG. 2. The energy distribution of *o*-Ps emitted from a typical porous silica film. Note that higher energy positrons implanted more deeply (solid circles) in the film produce more thermalized *o*-Ps with fewer events in the epithermal tail.

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Results on positron pulse width

Alberola et al., Nucl. Instr. Method A 560 (2006) 224-232



positron tagging system

expected inefficency < 10^{-8} , Δt =start-stop~ 10^{-9} sec



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Cross-check of oPs disappearance



 $\Delta n = n_{p1}(E < E_{thr}) - n_{p2}(E < E_{thr})$

no 0-peak if "bad" vacuum

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Search for n-n` oscillations

Given the set of the

□ the small mass n-n`mixing could have direct consequences for the propagation of ultra-high energy cosmic rays at cosmological distances (GZK cut off) Oscillations could be fast $\tau \sim 1$ sec

Berezhiani, Bento '05;

- experimental technique: n decays in a trap:
 - no magnetic field
 - as less as posible collisions

Pokotilovski'06

GZK cut off

Explanation with fast $(T_{osc} \ll T_n)$ neutron – mirror neutron oscillations



Proposal of PNPI for $n \rightarrow n'$ **experiment**

Experiment for $n \rightarrow n'$ oscillations can be realized using PNPI EDM spectrometer



 $\tau_{\text{storage H=0}}$ (without magnetic field)

 $\tau_{\text{storage H}\neq 0}$ (with magnetic field)

 $\Delta \tau = \tau_{\text{storage H}\neq 0} - \tau_{\text{storage H}=0} = (\neq)0$ It is the main question

A.Serebrov Talk at B-L Workshop'07

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Scheme of experimental installation to search for n-n' transitions



1: UCN input guide; 2: UCN storage chamber; 3: magnetic shielding; 4: solenoid; 5-6: UCN detectors; 7-9: valves; 10: Cs-magnetometers, 11: monitor detector, 12: entrance valve.



Count rates of UCN detectors (5 and 6) in log scale during measurements. The filling time is 100 s. Holding times were $t_1 = 50$ s and $t_2 = 470$ s. Empting time is 150 s. The time of background measurement is 150 s. The region 3 in these plots is used to deduce the numbers N_1 and N_2 required for the determination of the storage time, respectively the ratio *R* (after background subtraction). This picture was obtained after 130 cycles. $\tau_{\text{fill}} = 35$ s, $\tau_{\text{emp}} = 30$ s after holding time 50 s, $\tau_{\text{emp}} = 38$ s after holding time 470 s. **17**

A.Serebrov Talk at B-L Workshop'07

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Present experimental status of n-n` mixing

U. Shmidt et al'07 (UFRM II, Munchen) $\tau_{osc} > 2 s$

□ G. Ban et al.'07 (ILL, Grenoble) $\tau_{osc} > 103 \text{ s}$

A.Serebrov et al. (ILL, Grenoble)

$$\tau_{\rm osc}$$
 > 414 s

Further improvement $\tau_{osc} > 10^4$ sec

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Higgs:

<u>Higgs Boson</u>

Ordinary matter Higgs: $\Phi_0 \sim (1,2,1)(1,1,0)$

Mirror matter Higgs: $\Phi \sim (1,1,0)(1,2,1)$

The Z₂ symmetry implies an operator which transforms the two kinds of matter between each other, $\Phi_0 \leftrightarrow \Phi$ `

interaction: $L_{H-H} = \eta \phi^+ \phi \phi^{+} \phi^{+}$

two mass eigenstates are created by mixture of our Higgs field and the mirror field

$$\Phi_{\pm} = \frac{1}{\sqrt{2}} (\Phi_{O} \pm \Phi)$$

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Search for H-H`mixing

overall Higgs production rate now is devided between two channels instead of one channel with significance S (=Siganl/Noise) for the same running time one will get significance S/2

Higgs signal at LHC might be much weaker.

Foot, Lew, Volkas'Ignatiev, Volkas '95; Barbieri et al. '05; Wilczek '06, Zhu et al. '07

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SUMMARY

what is the origin of parity violation?
 50 years old question

experiments results

- on H-H` at LHC
- on n-n` and Ps-Ps` at low energies might give an answer to this question in the near future