



**The Abdus Salam
International Centre for Theoretical Physics**



1952-15

**School on Stochastic Geometry, the Stochastic Loewner Evolution,
and Non-Equilibrium Growth Processes**

7 - 18 July 2008

**From Critical to Off-Critical Interfaces
(Statistical interfaces and SLE)**

Denis BERNARD
*ENS, Lab. de Physique Theorique
F-75231 Paris Cedex 05
France*

FROM CRITICAL TO OFF-CRITICAL INTERFACES (07/2008)

I/ statistical interfaces and SLE

D. Bernard
with M. Bauer

1. Interfaces
2. statistical mechanics and curves
3. statistical martingales
4. SLE/CFT correspondence

II/ SLE/CFT and applications

1. Loewner eq, (algebraic version)
2. CFT martingales
3. Application 1: coupling to Gaussian field.
4. Application 2: $SLE(\kappa, \rho)$ and generalisation.
- (5. : N -SLE

III/ Off-critical SLE

1. (Warm-up) random walk / Girsanov theorem.
2. S.A.W as an example (of off-critical curves)
3. Off-critical drift and partition fct.
" SLE and field theory
4. Example 1: $[c=1+\text{mass}]$ $SLE_{4+\epsilon}$ (massive free field)
5. Example 2: off-critical LERW

I. Statistical interfaces and SLE

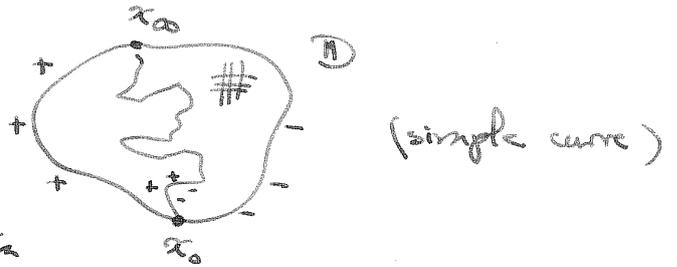
1/ What we describe? interfaces

Some of these have been described by M. Bower, let us me just recall few example to fix the ideas.

i) Ising model (spin clusters)

in domain \mathbb{D}
lattice $\Lambda \subset \mathbb{D}$

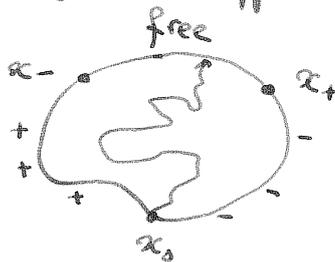
fix appropriate boundary cond
to impose that for each realization
there is an interface $\gamma_0 \rightsquigarrow \gamma_\infty$



at \bar{E}_0 , in the limit mesh $\rightarrow 0 \rightarrow$ SLE ($\kappa=3$) chordal.

we can play with different boundary condition

e.g

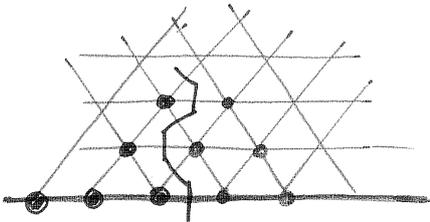


convergence towards
SLE dipolar
(but locally the same)

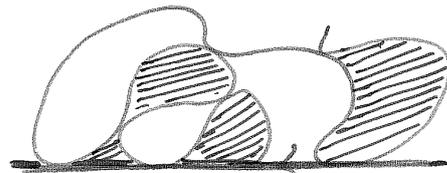
ii) Percolation

Eg on hexagonal
in the upper half plane

simple curve on the lattice
but not in the continuum limit
 \rightarrow double points ($\infty \neq$)



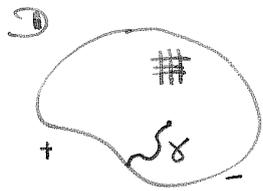
at criticality $p = p_c (=1/2)$



hull
which is \neq from cluster

2/ statistical mechanics and curves

c) stat. mech / Boltzmann weight \Rightarrow induces a measure on curve.



local degree of freedom \rightarrow configuration $c \in C'$
 Boltzmann weight w_c (for each configuration)

Probab of a config $P_c = \frac{1}{Z_D} w_c$

$$Z_D = \sum_{c \in C} w_c$$

\rightarrow measure on (portion of) curves.

fix a curve γ (view as a portion of an interface induced by the b.c.)

this selects a subset of configurations $C[\gamma]$ because 'spins' on the side of the curve are fixed

By Boltzmann rule $\frac{P[\gamma]}{Z_D} = \frac{Z_D[\gamma]}{Z_D}$, $Z_D[\gamma] = \sum_{c \in C[\gamma]} w_c$

Plc: By construction, this is a probability measure ($\sum P[\gamma] = 1$)
 because summing over all curves (say of a given length, call it T)
 reproduces the full sum: $\sum_{\gamma_T} Z_D[\gamma] = Z_D \iff \bigcup_{\gamma_T} C[\gamma] = C$

ii) domain Markov property



what to compare? $\left\{ \begin{array}{l} P_{D'}[\gamma|\gamma'] \text{ in } D', \text{ condit-d on } \gamma \\ P_{D|\gamma}[\gamma'] \text{ in } D, \gamma \end{array} \right.$

(both measure of identical support)

$$\frac{P_{D'}[\gamma|\gamma']}{P_{D'}[\gamma]} = \frac{Z_{D'}[\gamma\gamma']}{Z_{D'}[\gamma]} = \frac{Z_D[\gamma\gamma']}{Z_D} \cdot \frac{Z_D}{Z_D[\gamma]} = \frac{Z_D[\gamma\gamma']}{Z_D[\gamma]}$$

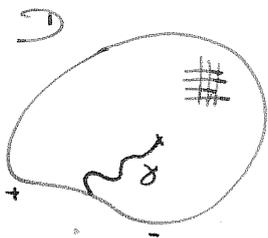
$$P_{D|\gamma}[\gamma'] = \frac{Z_{D|\gamma}[\gamma']}{Z_{D|\gamma}} \quad \text{but } \left\{ \begin{array}{l} Z_{D|\gamma}[\gamma'] = e^{E_\gamma} Z_D[\gamma\gamma'] \\ Z_{D|\gamma} = e^{E_\gamma} Z_D[\gamma] \end{array} \right. \quad \text{because}$$

iii) continuum limit

As explained in the previous week
domain Markov prop.
⊕ conf. transport } → SLE

3) Statistical martingale

The aim is to show that statistical correlation f^d are martingales for the process describing the growth of the curve.



O: observables are maps: $c \rightarrow O_c$
example
correlation f^d are expectation

$$\langle O \rangle_D \equiv \mathbb{E}_D [O] = \frac{1}{Z_D} \sum_c w_c O_c$$

consider now this expectation but conditioned on γ (of length τ , say)
(As for the Markov prop, it is easy to see that ~~conditioned~~ the conditioned expectation is the same as the expectation in the cut domain)

$$\langle O \rangle_{D|\gamma_\tau} \equiv \langle O \rangle_{D, \gamma_\tau} = \frac{1}{Z_{D|\gamma_\tau}} \sum_{c \in C|\gamma_\tau} w_c O_c$$

Since γ_τ is random (with law $P(\gamma_\tau)$), τ are random number
We have

$$\overset{\substack{\text{average} \\ \text{over curve } \gamma_\tau}}{\mathbb{E}[\langle O \rangle_{D|\gamma_\tau}]} = \langle O \rangle_D \quad \text{ie } \parallel \text{Reorganisation of the statistical sum}$$

(the proof is easy ...)

More, specifying γ_τ of a given length, partition the configurati. space
increasing the length amounts to consider finer and finer
partitions. \rightarrow we thus define a filtration (increase the knowledge)

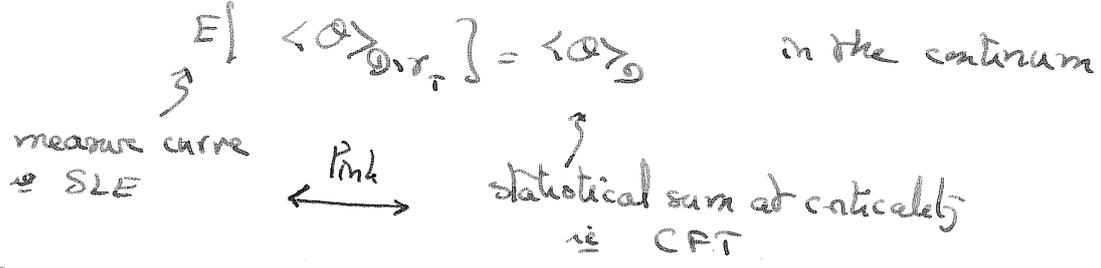
So we have

$$\langle O \rangle_{D, \gamma_t} = E[\langle O \rangle_{D, \gamma_t} | \gamma_t]$$

$$\Rightarrow \begin{cases} E[\langle O \rangle_{D, \gamma_s} | \gamma_t] = \langle O \rangle_{D, \gamma_t} & \text{for } s > t \\ \text{i.e. } t \mapsto \langle O \rangle_{D, \gamma_t} \text{ is a martingale} \end{cases}$$

4. SLE/CFT correspondence

Identification between CFT/SLE comes by taking the continuous limit of the previous formulae. eg



1) $\langle O \rangle_D$ as ratio of partition function
as ratio of CFT correlation f^{ch}



need to code for the bdy cond. and their change
 \Rightarrow insertion of bdy cond. changing operator $\psi(x_0)$ and $\psi(x_\infty)$

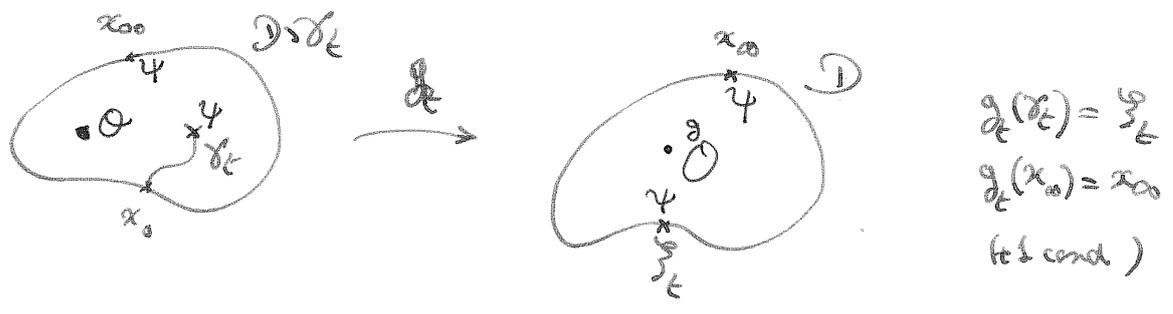
$$\langle O \rangle_D \xrightarrow{\text{in the continuum}} \frac{\langle \psi(x_0) O \psi(x_\infty) \rangle_D}{\langle \psi(x_0) \psi(x_\infty) \rangle_D} \quad (\text{here cft correlation})$$

\uparrow represents some kind of the continuum limit of the partition function.

(*) gives the SLE/CFT correspondence.

i.e. this ratio of cft correlator but in D, γ_t are martingales

The advantage is that (because of conformal invariance) the transformation from $\mathcal{D} \ni \mathcal{D}_t$ to \mathcal{D} is purely kinematical.



Then $\langle \Psi(x_0) \circ \Psi(x_t) \rangle_{\mathcal{D} \ni \mathcal{D}_t} = |g'_t(x_0)|^h |g'_t(x_t)|^h \langle \Psi(x_0) \circ \Psi(x_t) \rangle_{\mathcal{D}}$

with h the scaling dim of the operator Ψ .

\Rightarrow The process $t \rightarrow \frac{\langle \Psi(x_0) \circ \Psi(x_t) \rangle_{\mathcal{D}}}{\langle \Psi(x_0) \Psi(x_t) \rangle_{\mathcal{D}}}$ are (chordal) SLE... martingales

Pr: The Jacobian $|g'_t(x_t)|^h$ (or both) are infinite cancel

eq: $\circ \Psi$ = transform of operator Ψ by conformal map.

ii) CFT/SLE (chordal)

The martingale relation is true iff. we chose correctly the CFT and the operator Ψ .

CFT: with central charge $c = \frac{(6-k)(3k-8)}{2k} < 1$
 Ψ : with dimension $h = \frac{6-k}{2k}$
 and "degenerate" at level 2.

iii) Examples of CFT/SLE

$c=1$ Gaussian free field $\kappa=4$

$c=1/2$ Ising $\kappa=3$ open cluster
 $\kappa=16/3$ FK (high temp^o) cluster

$c=0$ $\kappa=6$ percolation

$\kappa=8/3$ S.A.W.