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Dark Matter
Lecture 2: Principles of direct WIMP detection

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Lecture 2: Principles of direct WIMP detection

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Content

- Direct detection of WIMPs

  expected rates in a terrestrial detector

  kinematics of elastic WIMP-nucleus scattering

  differential rates

  corrections I: movement of the Earth

  corrections II: form factors

  cross sections for scattering on nucleons

  - spin independent

  - spin dependent

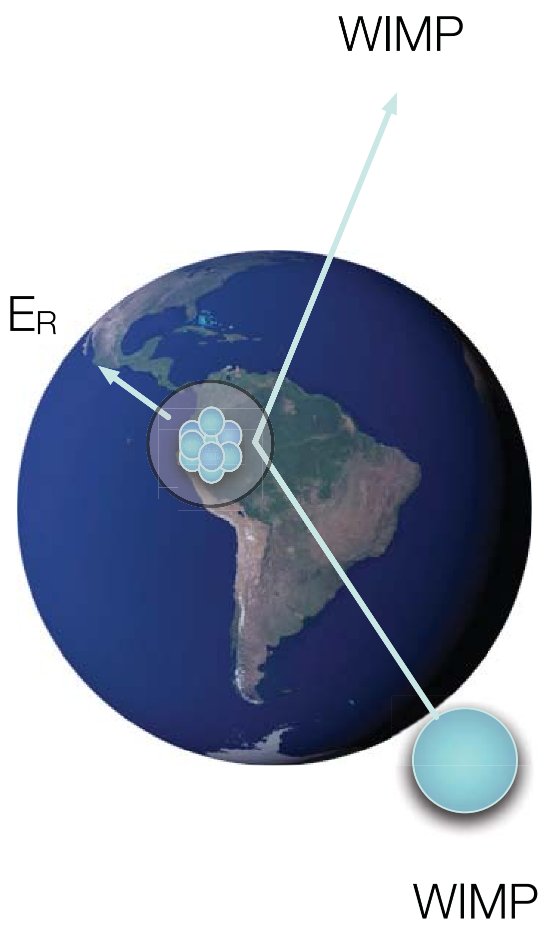
- Expected WIMP signal and backgrounds

  quenching factors and background discrimination

  main background sources in direct detection experiments

  detector strategies

Direct Detection of WIMPs



- **Elastic collision with atomic nuclei**
- The recoil energy of the nucleus is:

$$E_R = \frac{|\vec{q}|^2}{2m_N} = \frac{\mu^2 v^2}{m_N}(1-\cos\theta)$$

- q = momentum transfer $|\vec{q}|^2 = 2\mu^2 v^2 (1-\cos\theta)$
- μ = reduced mass (m_N = nucleus mass; m_χ = WIMP mass)

$$\mu = \frac{m_\chi m_N}{m_\chi + m_N}$$

- v = mean WIMP-velocity relative to the target
- θ = scattering angle in the center of mass system

Expected Rates in a Detector

- For now **strongly simplified**:

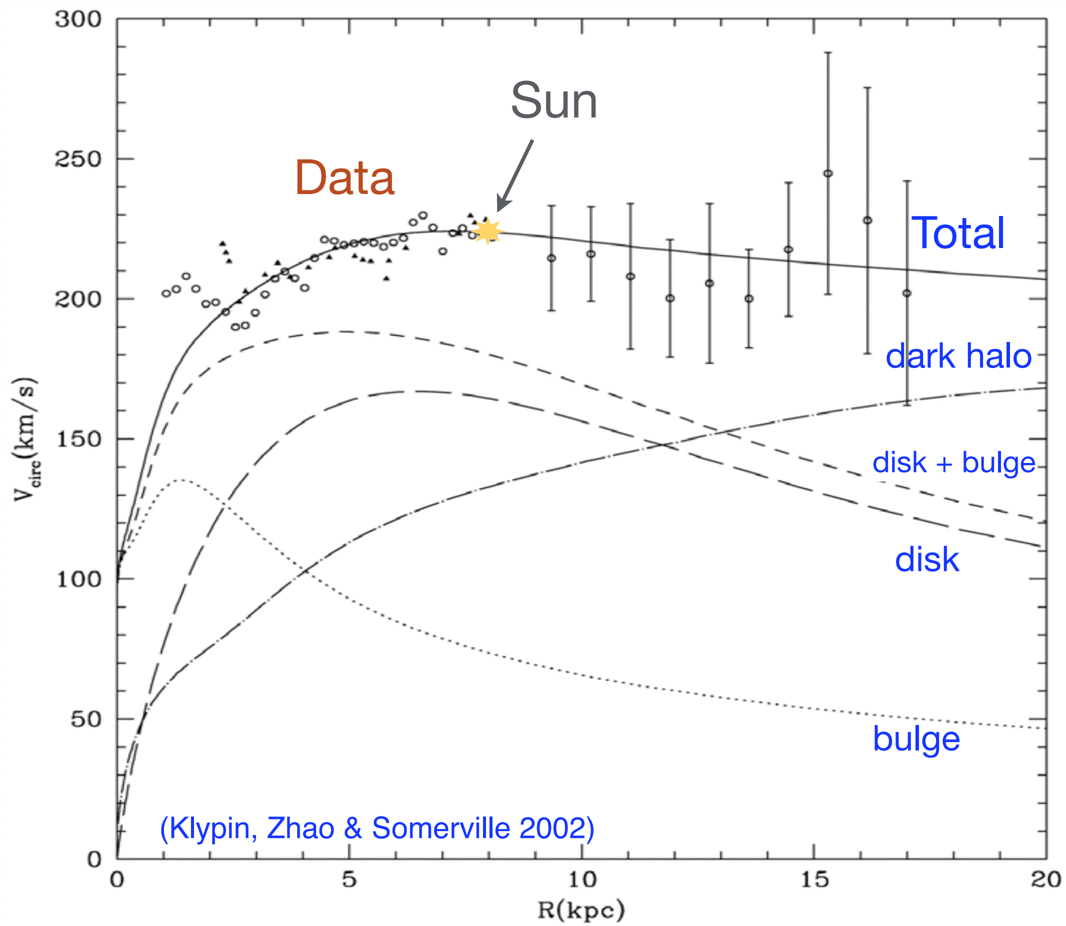
Astrophysics

$$R \propto N \frac{\rho_\chi}{m_\chi} \sigma_{\chi N} \cdot \langle v \rangle$$

Particle physics

- N = number of target nuclei in a detector
- ρ_χ = local density of the dark matter in the Milky Way
- <v> = mean WIMP velocity relative to the target
- m_χ = WIMP-mass
- $\sigma_{\chi N}$ =cross section for elastic scattering

Local Density of WIMPs in the Milky Way



Particle data group:

$$\rho_{halo} = 0.1 - 0.7\ GeVcm^{-3}$$

$$\rho_{disk} = 2 - 7\ GeVcm^{-3}$$

'Standard' value:

$$\rho_{\chi} \simeq 0.3\ GeVcm^{-3}$$

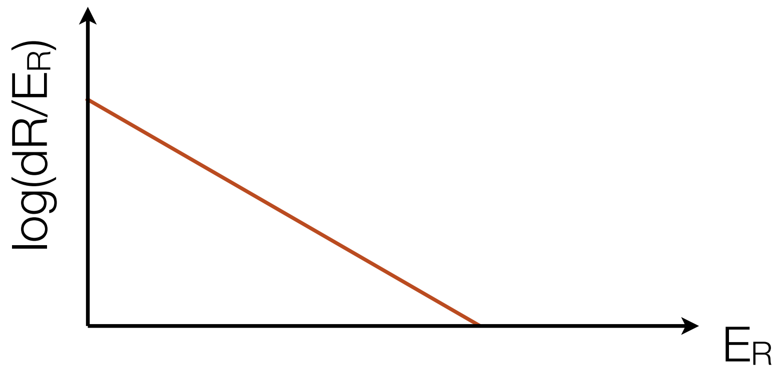
$$\rho_{\chi} \simeq 3000\ \mathrm{WIMPs} \cdot m^{-3}$$

$$(M_{WIMP} = 100\ \mathrm{GeV})$$

Expected Rates in a Detector

- **The differential rate (still strongly simplified) is:**

$$\frac{dR}{dE_R} = \frac{R_0}{E_0 r} e^{-\frac{E_R}{E_0 r}}$$



- **R =** event rate per unit mass
- **E_R =** nuclear recoil energy
- **R_0 =** total event rate
- **E_0 =** most probable energy of WIMPs

  (Maxwell-Boltzmann distribution)
- **r =** kinematic factor $r = \frac{4 m_\chi m_N}{(m_\chi + m_N)^2}$

$$\int_0^\infty \frac{dR}{dE_R} dE_R = R_0$$

$$\langle E_R \rangle = \int_0^\infty E_R \frac{dR}{dE_R} dE_R = E_0 r$$

Some Typical Numbers

- **We assume that the WIMP mass and the nucleus mass are identical:**

$$m_\chi = m_N = 100\ GeV \cdot c^{-2}$$

$$\Rightarrow r = \frac{4 m_\chi m_N}{(m_\chi + m_N)^2} = 1$$

kinematic factor

$$\mathrm{v} \sim 220\ \mathrm{km\ s^{-1}} = 0.75 \times 10^{-3} c$$

mean WIMP velocity relative to target
(halo is stationary, Sun moves through halo)

$$\langle E_R \rangle = E_0 = \frac{1}{2} m_\chi \mathrm{v}^2$$

$$\langle E_R \rangle = \frac{1}{2} 100 \frac{GeV}{c^2} (0.75 \times 10^{-3} c)^2$$

$$\langle E_R \rangle \approx 30\ \mathrm{keV}$$

mean recoil energy deposited in a detector

Expected Rates in a Detector

- We have to take into account following facts:

➡ The WIMPs will have a velocity distribution f(v)

➡ the detector is on Earth, which moves around the Sun, which moves around the galactic center

➡ the cross section depends on whether the interaction is spin-independent, or spin-dependent

➡ the WIMPs scatter on nuclei, which have a finite size; we have to consider form-factor corrections < 1

➡ the nuclear recoil energy is not necessarily the *observed* energy, since in general the detection efficiency is < 1

➡ detectors have a certain energy resolution and energy threshold

$$\Rightarrow \frac{dR}{dE_R} = R_0 S(E_R) F^2(E_R) I$$

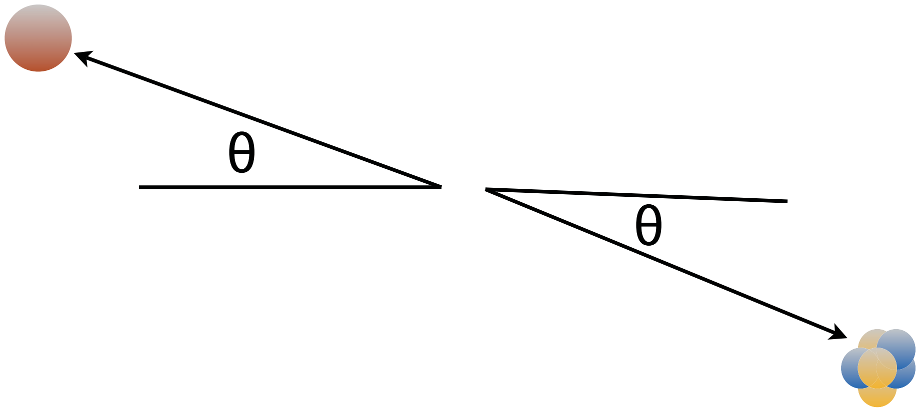
spectral function (masses and kinematics)

form factor correction

type of interaction

Kinematics

- WIMPs with **velocity v and kinetic energy** $E_i = \frac{1}{2} m_\chi v^2$ which are scattered under an angle θ in the center of mass system, will yield a **recoil energy E_R in the laboratory system:**



$$E_R = E_i r \frac{(1-\cos\theta)}{2}$$

$$r = \frac{4\mu^2}{m_\chi m_N} = \frac{4 m_\chi m_N}{(m_\chi + m_N)^2}$$

reduced mass $$\mu = \frac{m_\chi m_N}{m_\chi + m_N}$$

Kinematics

- **Assumption**: the scattering is isotropic => uniform in cosθ
- An incoming WIMP with energy E_i will deliver a recoil energy:

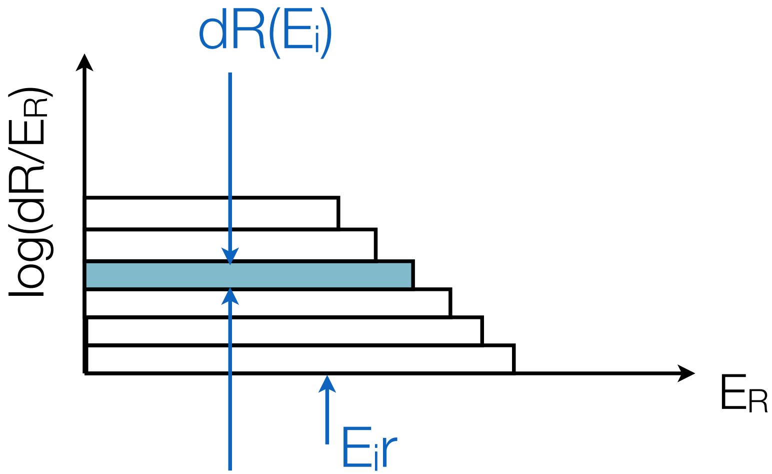
$$0 \leq E_R \leq E_i r$$

- We had looked at the case with r = 1 (equal masses), a stationary target and θ=180° (head-on collision)

$$E_R = E_i$$

- **How does the overall spectrum look like?** We will sample the incident spectrum. In each interval $E_i \to E_i + dE_i$ we will have a contribution to the spectrum in $E_R \to E_R + dE_R$ at rate $dR(E_i)$ of

$$d\left(\frac{dR}{dE_R}(E_R)\right) = \frac{dR(E_i)}{E_i r}$$



Kinematics

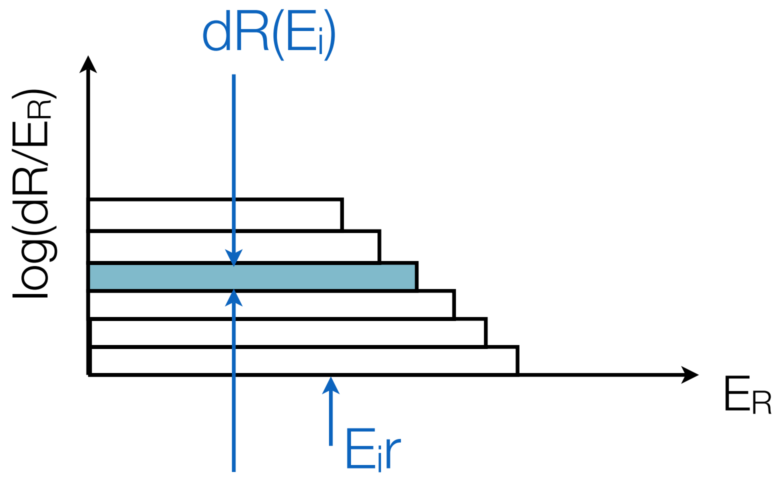
- **We have to integrate over all incoming WIMP energies:**

$$\frac{dR}{dE_R}(E_R) = \int_{E_{\min}}^{E_{\max}} \frac{dR(E_i)}{E_i r}$$

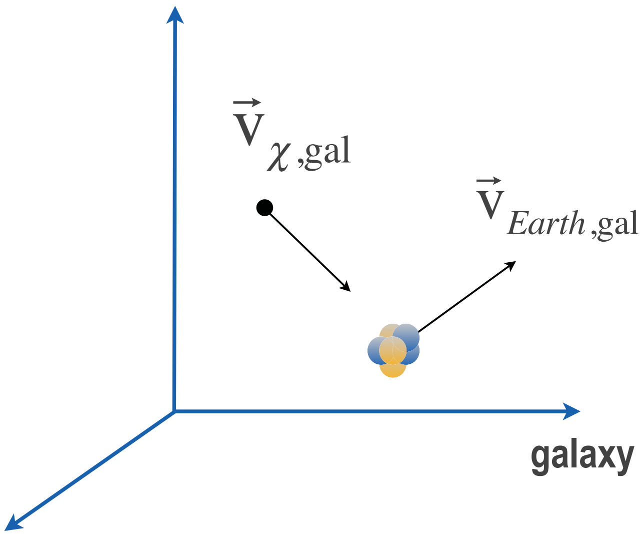
- **For E_{max}:** we will use either ∞ or v_{esc}
- **For E_{min}:** to deposit a recoil energy E_R, we need an incident WIMP energy:

$$E_i \geq \frac{E_R}{r} \equiv E_{\min}$$

- We will now determine the differential rate.



Coordinate System



collision kinematics

$\vec{v} = \vec{v}_{\chi,\mathrm{Earth}} =$ WIMP velocity in the target (Earth) frame

$\vec{v}_{Earth,\mathrm{gal}} =$ Earth velocity in the Galaxy frame

$\vec{v}_{\chi,\mathrm{gal}} =$ WIMP velocity in the Galaxy frame

galaxy dynamics

$$\vec{v}_{\chi,\mathrm{gal}} = \vec{v}_{\chi,\mathrm{Earth}} + \vec{v}_{Earth,\mathrm{gal}}$$

$$\vec{v}_{\chi,\mathrm{gal}} = \vec{v} + \vec{v}_{\mathrm{E}}$$

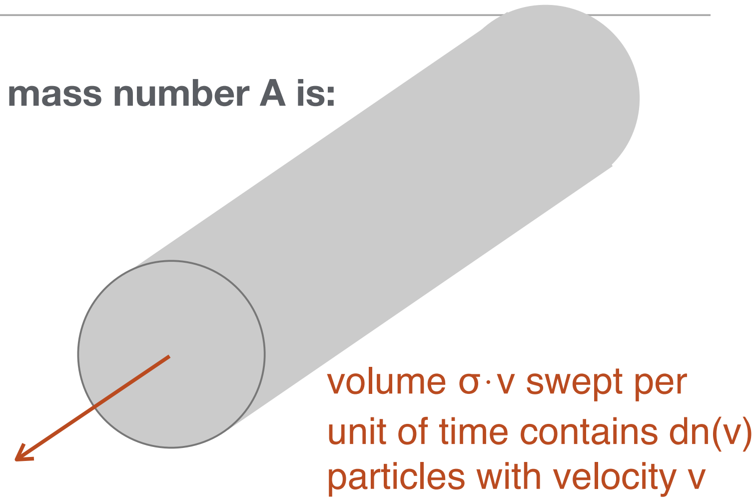
$$f(\vec{v}, \vec{v}_{\mathrm{E}}) = e^{-\frac{(\vec{v}+\vec{v}_{\mathrm{E}})^2}{v_0^2}}$$

Maxwell-Boltzmann velocity distribution

Differential Rate

- **The event rate per unit mass in a detector with nuclear mass number A is:**

$$dR = \frac{N_A}{A} \sigma \; v \; dn$$



➡ N_A =6.022×10^{26} kg^{-1} Avogadro number

➡ σ = cross section for the scattering on the nucleus

- **The differential particle density *dn*** is taken as a function of the **velocity v:**

$$dn = \frac{n_0}{k} f(\vec{v}, \vec{v}_E) d^3\vec{v}$$

- with the mean WIMP number density $n_0 = \frac{\rho_\chi}{m_\chi}$

➡ v = velocity relative to the target (which is on Earth)

➡ v_E = Earth velocity (and thus target velocity) relative to the dark matter distribution

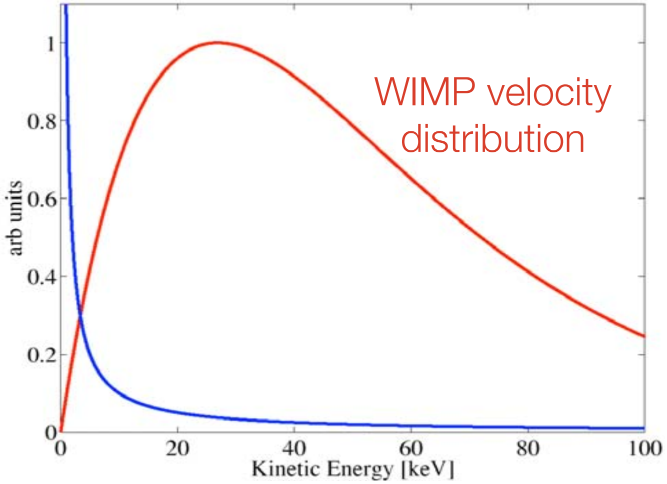
Differential Rate

- k is a normalization constant, so that:

$$\int_0^{v_{esc}} dn \equiv n_0$$

- where v_{esc} = local galactic escape velocity (≈ 544 km/s)
- this means:

$$k = \int f(\vec{\mathrm{v}},\vec{\mathrm{v}}_E)d^3\vec{\mathrm{v}}$$



$$k = \int_0^{2\pi} d\phi \int_{-1}^{+1} d(\cos\theta) \int_0^{\mathrm{v}_{esc}} f(\vec{\mathrm{v}},\vec{\mathrm{v}}_E)\mathrm{v}^2 d\mathrm{v}$$

- We assume a Maxwell-Boltzmann WIMP velocity distribution with respect to the galactic frame:

$$f(\vec{\mathrm{v}},\vec{\mathrm{v}}_E) = e^{-\frac{(\vec{\mathrm{v}}+\vec{\mathrm{v}}_E)^2}{\mathrm{v}_0^2}}$$

$\vec{\mathrm{v}}+\vec{\mathrm{v}}_E$ WIMP velocity in the galaxy frame

$\mathrm{v}_0 \approx 220 \text{ km s}^{-1}$

Differential Rate

- **We first look at the simplified case of $v_E = 0$ and $v_{esc} = \infty$**
- For this case, we have:

$$k = k_0 = \int_0^{2\pi} d\phi \int_{-1}^{+1} d(\cos\theta) \int_0^{\infty} e^{-\frac{(\vec{v}+0)^2}{v_0^2}} v^2 dv = 4\pi \int_0^{\infty} e^{-\frac{(\vec{v})^2}{v_0^2}} v^2 dv = (\pi v_0^2)^{3/2}$$

- and thus

$$dR = R_0 \frac{1}{2\pi v_0^4} v f(v,0) d^3v$$

- with the **total rate R_0** per unit mass (**$v_E = 0$ and $v_{esc} = \infty$**) being defined as:

$$R_0 = \frac{2}{\sqrt{\pi}} \frac{N_A}{A} \frac{\rho_\chi}{m_\chi} \sigma_0 \ v_0$$

Differential Rate

- For a Maxwellian-Boltzmann distribution

$$f(\mathrm{v},0) = e^{-\frac{\mathrm{v}^2}{\mathrm{v}_0^2}}$$

- isotropic: $d^3\mathrm{v} \rightarrow 4\pi \mathrm{v}^2 d\mathrm{v}$

- and with the incident, and most probable energy of WIMPs:

$$\mathrm{E_i} = \frac{1}{2} m_\chi \mathrm{v}^2 \quad \text{and} \quad E_0 = \frac{1}{2} m_\chi \mathrm{v}_0^2$$

- **we obtain for the differential rate:**

$$\frac{dR}{dE_R}(E_R) = \int_{E_R/r}^{\infty} \frac{dR(E_i)}{E_i r} = \frac{R_0}{r\left(\frac{1}{2} m_\chi \mathrm{v}_0^2\right)^2} \int_{\mathrm{v}_{\min}}^{\infty} e^{-\frac{(\vec{\mathrm{v}})^2}{\mathrm{v}_0^2}} \mathrm{v} d\mathrm{v} = \frac{R_0}{E_0 r} e^{-\frac{E_R}{E_0 r}}$$

$$\mathrm{v}_{\min} = \sqrt{\frac{2E_R}{r \cdot m_\chi}}$$

1. Correction: galactic escape velocity v_{esc}

- **For a finite escape velocity v_{esc} (and still $v_E = 0$)**

$$\left|\vec{v}+\vec{v}_E\right| = v_{esc}$$

WIMP velocity in the galaxy frame

- we obtain for the differential rate:

$$\frac{dR}{dE_R} = \frac{k_0}{k_1(v_{esc},0)} \frac{R_0}{E_0 r}\left(e^{-\frac{E_R}{E_0 r}} - e^{-\frac{(v_{esc})^2}{v_0^2}} \right)$$

- **Example:** if we use the value v_{esc} ~ 600 km/s, and v_0 = 220 km/s, we obtain:

$$\frac{k_0}{k_1} = 0.9965 \qquad \frac{R(0,v_{esc})}{R_0} = 0.9948$$

2. Correction: velocity of the Earth v_E

- Clearly the Earth is moving, thus $v_E \neq 0$, and $v_E \sim v_0 \approx 220$ km/s
- **A complete calculation yields (see Appendix in Ref [5]):**

$$\frac{dR}{dE_R} = \frac{k_0}{k_1}\frac{R_0}{E_0 r}\left\{\frac{\sqrt{\pi}v_0}{4v_E}\left[erf\left(\frac{v_{min}+v_E}{v_0}\right) - erf\left(\frac{v_{min}-v_E}{v_0}\right)\right] - e^{-\frac{(v_{esc})^2}{v_0^2}}\right\}$$

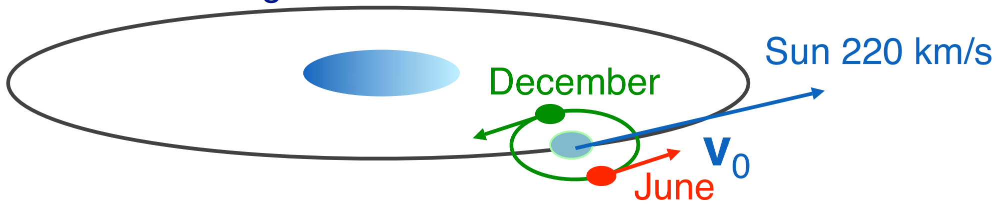
- with the error function being defines as: $erf(x) = \frac{2}{\sqrt{\pi}}\int_0^x e^{-t^2}\,dt$
- and: $v_{min} = v_0\sqrt{\frac{E_R}{E_0 r}}$

$$k_1 = k_0\left[erf\left(\frac{v_{esc}}{v_0}\right) - \frac{2}{\sqrt{\pi}}\frac{v_{esc}}{v_0} - e^{-\frac{(v_{esc})^2}{v_0^2}}\right]$$

Signal Modulation: Annual Effect

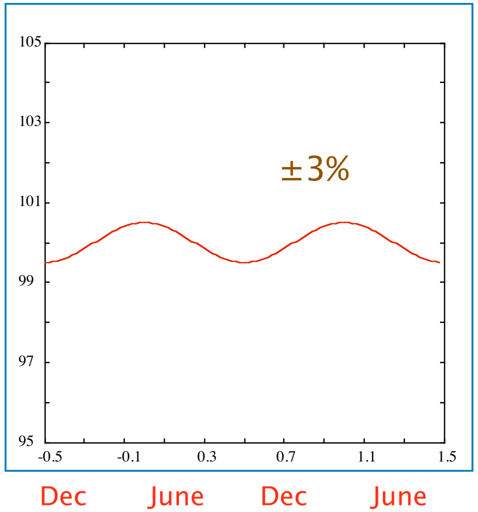
- The velocity of the Earth varies over the year as the Earth moves around the Sun, and can be written as [in km/s]:

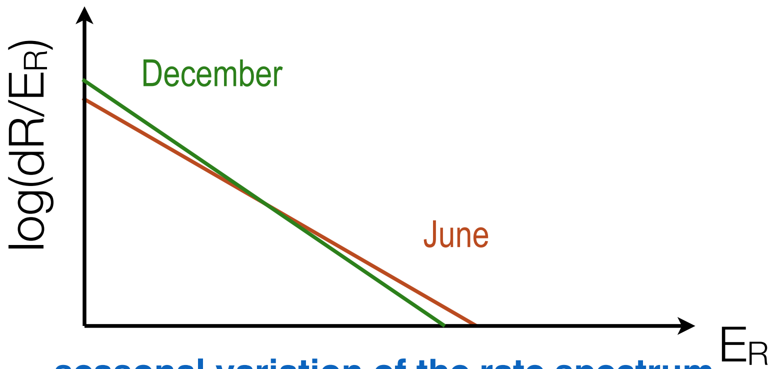
$$v_E(t) = v_0\left[1.05 + 0.07\cos\frac{2\pi(t-t_p)}{1yr}\right]$$



- t = days since January 1st
- t_p = 2. June (152.5 d) ± 1.3 d; 1 yr = 362.25 d

- the velocity modulation gives rise to a ~ 3% modulation in the rate $\frac{d}{dv_E}\left(\frac{R}{R_0}\right) \sim \frac{1}{2v_E}\frac{R}{R_0}$ (for $v_E \sim v_0$)

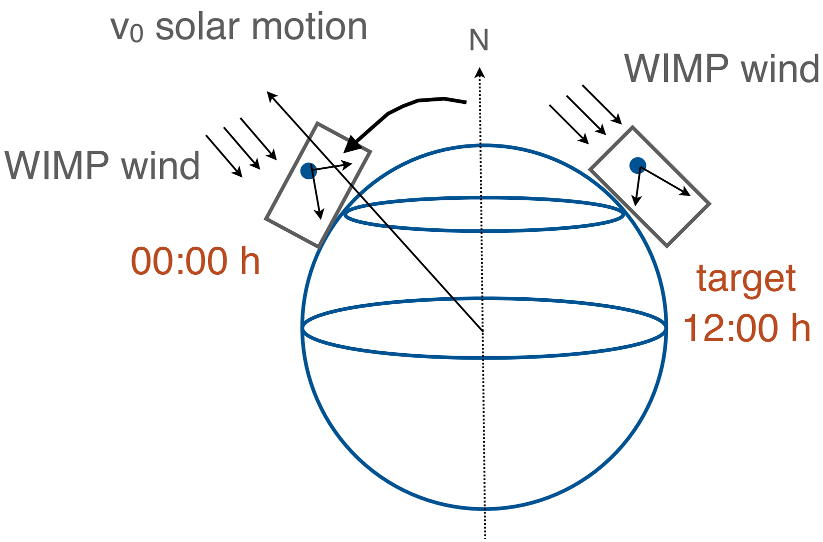


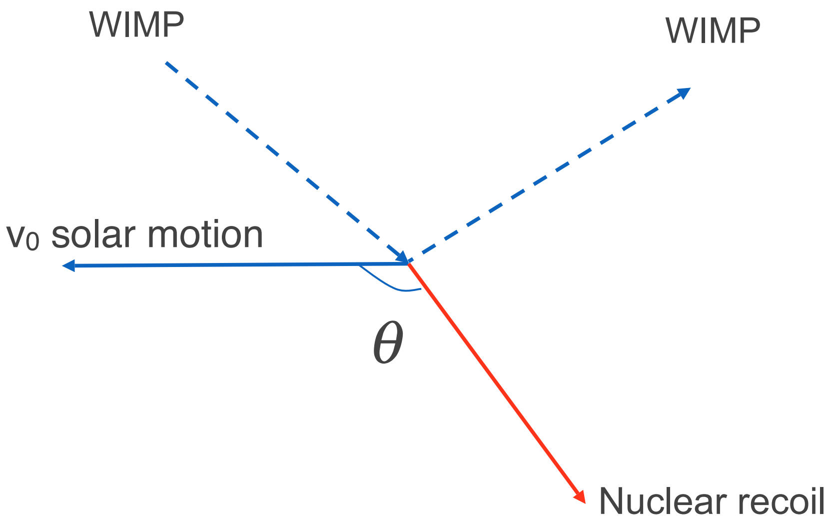


seasonal variation of the rate spectrum

Signal Modulation: Recoil Direction

The mean recoil direction rotates over one sidereal day





The differential angular spectrum is given by:

$$\frac{d^2R}{dE_R d(\cos\theta)} = \frac{1}{2}\frac{R_0}{E_0 r} e^{-\frac{(v_E \cos\theta - v_{min})^2}{v_0^2}}$$

=> asymmetry: more events in forward than in backward direction
- few 10s-100s events need to distinguish between halo models (depending on the E_{th} of the detector)

Additional "Corrections" to the Differential Rate

$$\Rightarrow \frac{dR}{dE_R} = R_0 S(E_R) F^2(E_R) I$$

- so far we have discussed the **spectral function $S(E_R)$**
- it contains the kinematic of the scattering, and the time dependence of the signal
- **we now discuss**
  - ➡ $F^2(E_R)$: form factor corrections, with $E_R = q^2/2m_X$
  - ➡ I: type of interaction
- in general, for NR particles (v << c) the scalar and axial interactions dominate; we will thus consider spin-independent and spin-dependent couplings

Nuclear form factor and spin-independent couplings

- Scattering amplitude: Born approximation $\vec{q} = \hbar\,(\vec{k}' - \vec{k})$
- Spin-independent scattering is coherent $\lambda = \hbar/q \sim$ few fm

$$M(\vec{q}) = f_n\, A \underbrace{\int d^3x\, \rho(\vec{x})\, e^{i\,\vec{q}\cdot\vec{x}}}_{F(\vec{q})} \quad \Rightarrow \quad \sigma \propto |M|^2 \propto A^2$$

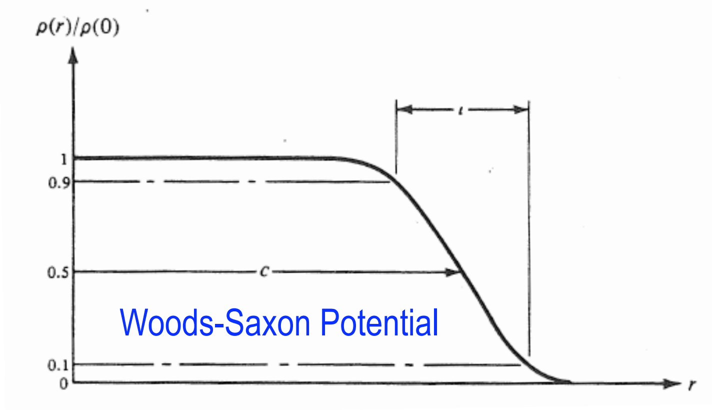
mass number

fundamental couplings to nucleons

Fourier-transform of the density of scattering centers

$$F(qr_n) = \underbrace{\frac{3[\sin(qr_n) - qr_n\cos(qr_n)]}{(qr_n)^3}}_{j_1(qr_n)}\, e^{-(qs)^2/2}$$

"Helm" form factor

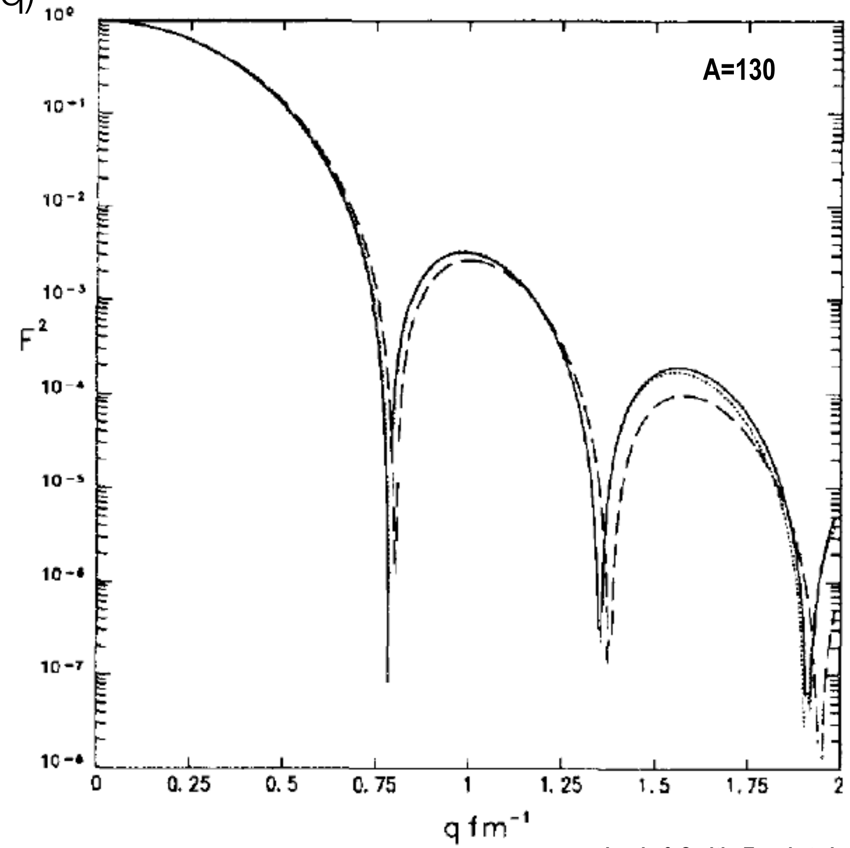


- with r_n = nuclear radius, $r_n \approx 1.2\, A^{1/3}$ fm, s = 1fm (skin thickness)

Nuclear form factor and spin-independent couplings

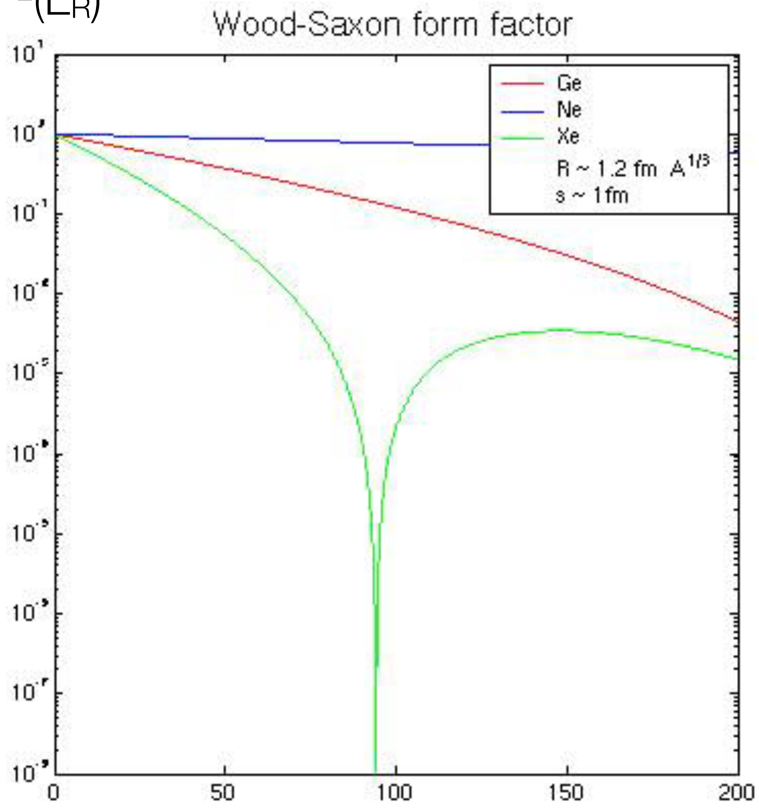
- Loss of coherence as larger momentum transfers probes smaller scales:

$F^2(q)$



Lewin & Smith; Engel et al

$F^2(E_R)$



E_R

Spin independent cross section

- The differential cross section can be written as:

$$\frac{d\sigma(q)}{dq^2} = \frac{\sigma_0 F^2(q)}{4\mu^2 \mathrm{v}^2}$$

relative velocity in center-of-mass frame

- where σ_0 = total cross section for F(q) = 1

- From Fermi's Golden Rule it follows:

$$\frac{d\sigma(q)}{dq^2} = \frac{1}{\pi \mathrm{v}^2}|M|^2 = \frac{1}{\pi \mathrm{v}^2} f_n^2 A^2 F^2(q)$$

- We can then identify the **total cross section σ_0 for F(q)=1**:

$$\sigma_0 = \frac{4\mu^2}{\pi} f_n^2 A^2 = \underbrace{\frac{4}{\pi} m_n^2 f_n^2}_{\sigma_n} \frac{\mu^2}{m_n^2} A^2$$

cross section for scattering off nucleus

cross section for scattering on nucleons

dependence on particle physics model for WIMP

Spin independent cross section and differential rate

- Putting now everything together:

$$\frac{d\sigma(q)}{dq^2} = \frac{1}{4m_n^2 v^2}\sigma_n A^2 F^2(q)$$

$$\frac{dR}{dE_R} = \frac{R_0}{E_0 r} e^{-\frac{E_R}{E_0 r}} F^2(q)$$

particle physics

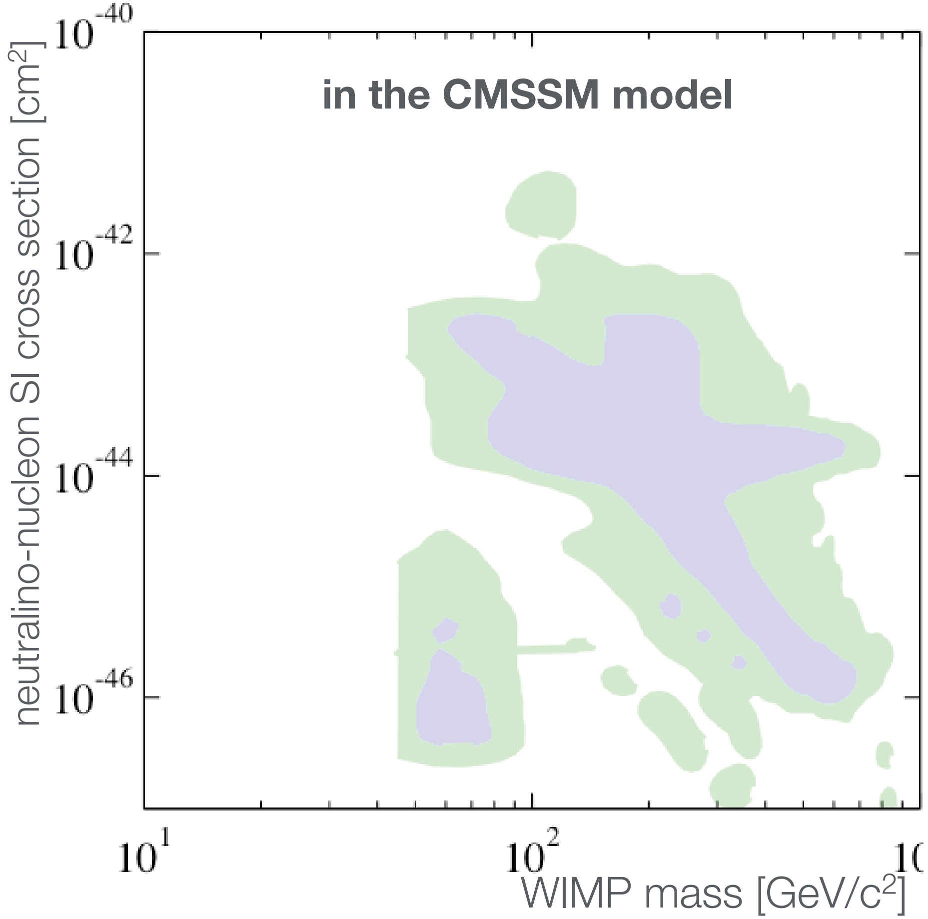
dark matter halo

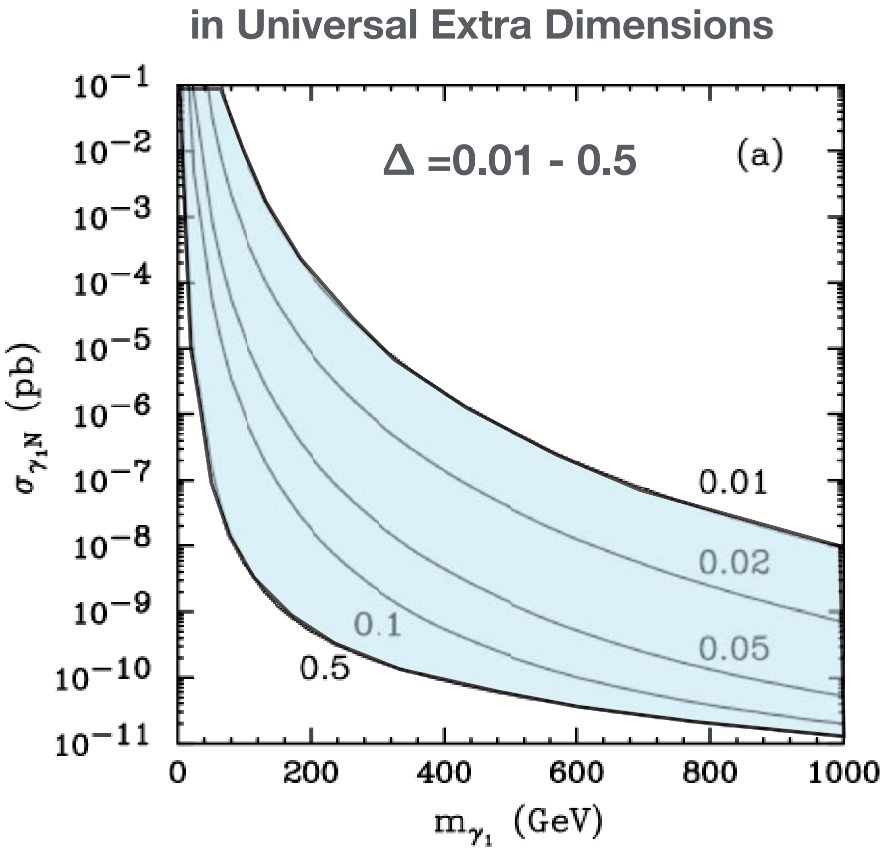
$$R_0 = \frac{2}{\sqrt{\pi}}\frac{N_A}{A}\frac{\rho_\chi}{m_\chi}\sigma_0 \, v_0 \qquad \sigma_0 = \sigma_n \frac{A^2}{m_n^2}\left(\frac{m_\chi m_N}{m_\chi + m_N}\right)^2$$

detector

WIMP Mass and SI Cross Section

- Predictions from supersymmetry (left) and UED (right) [10^{-8} pb = 10^{-44} cm^2]:

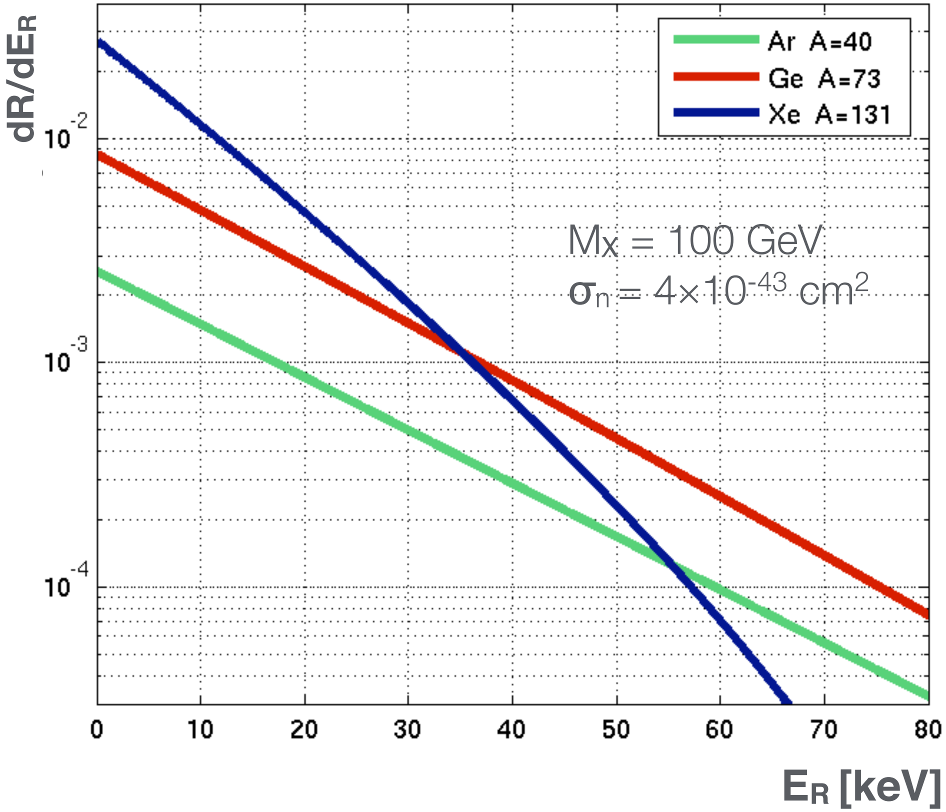




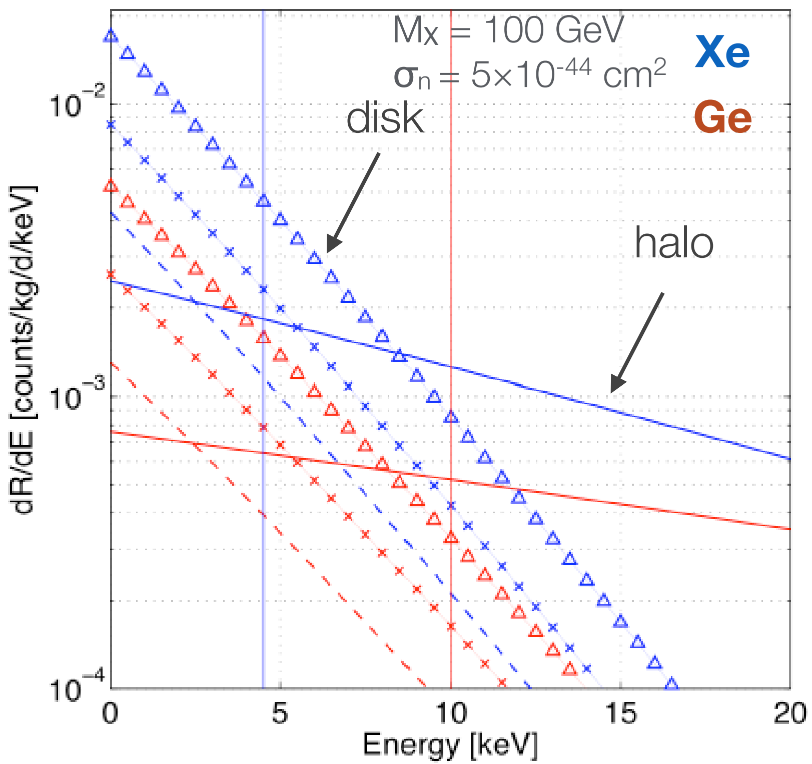
Spin independent cross section and differential rate

- Expected rates for different detector materials

from the dark matter halo



from the dark matter thick disk (and halo)



Nuclear form factor and spin dependent couplings

- **For spin dependent couplings the scattering amplitude is dominated by the unpaired nucleon:** the coupling is to the total nuclear spin J (paired nucleons ↑↓ tend to cancel):

$$\frac{d\sigma(q)}{dq^2} = \frac{8}{\pi v^2}\Lambda^2 G_F^2 J(J+1)F^2(q)$$

- with: G_F = Fermi constant, J = nuclear spin, $F^2(q)$ = form factor for spin dependent interactions

- and $$\Lambda = \frac{1}{J}\left[a_p\langle S_p\rangle + a_n\langle S_n\rangle\right]$$

- **a_p, a_n: effective coupling of the WIMPs to protons and neutrons, typically α/m_W^2**

- and the expectation values of the proton and neutron spins in the nucleus

$$\langle S_{p,n}\rangle = \langle N|S_{p,n}|N\rangle$$

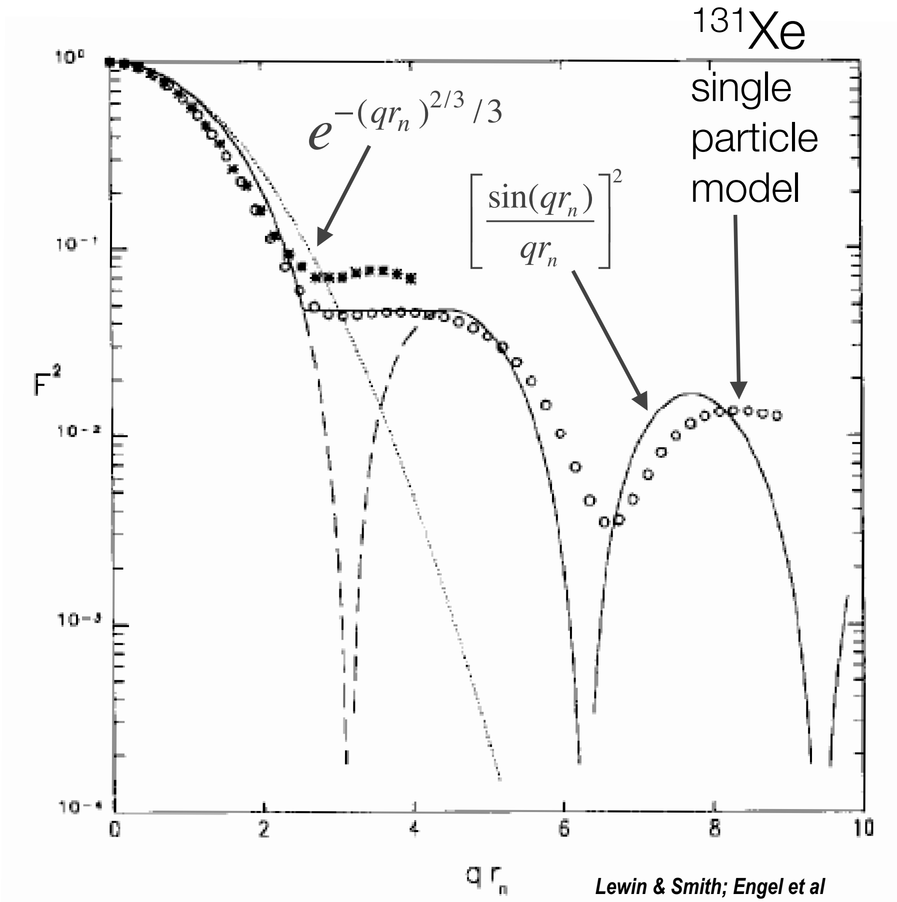
measure the amount of spin carried by the p- and n-groups inside the nucleus

Nuclear form factor and spin dependent couplings

- **Form factor example: simplified, based on model with valence nucleons in a thin shell:**

$$F(qr_n) = j_0(qr_n) = \frac{\sin(qr_n)}{qr_n}$$

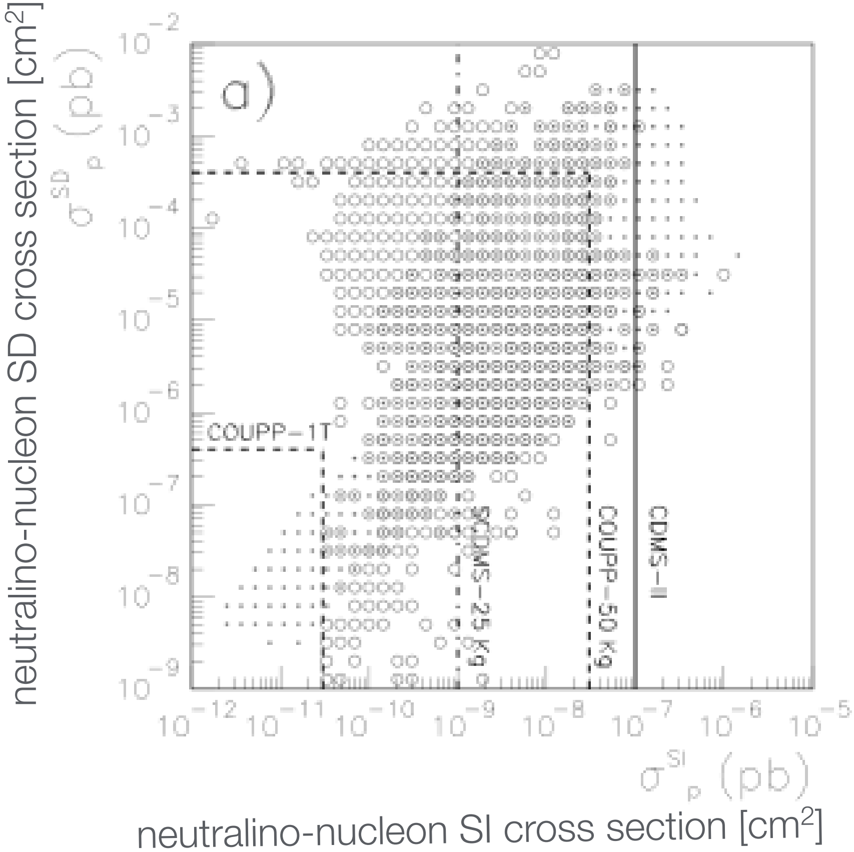
- **Better**: detailed calculations based on realistic nuclear models
  -for instance, the conventional nuclear shell model using reasonable nuclear Hamiltonians
  - cross check by agreement of predicted versus measured magnetic moment of the nucleus (since the matrix element for χN scattering is similar to the magnetic moment operator)



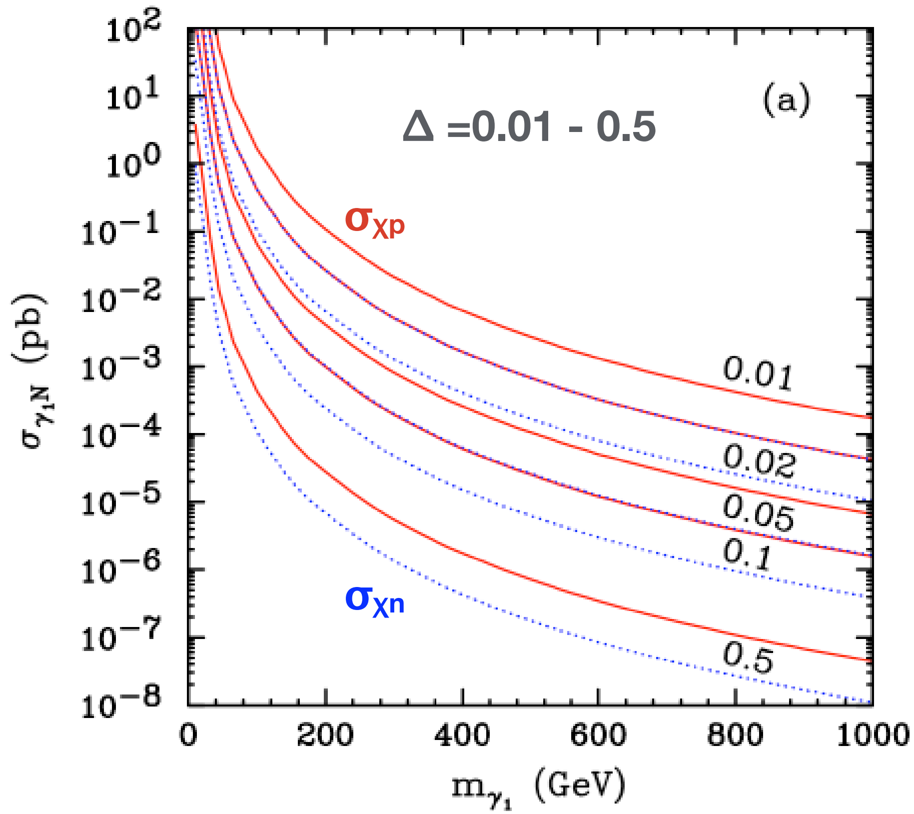
WIMP Mass and SD Cross Section

- Predictions from supersymmetry (left) and UED (right) [10^{-8} pb = 10^{-44} cm^2]:

in the CMSSM model

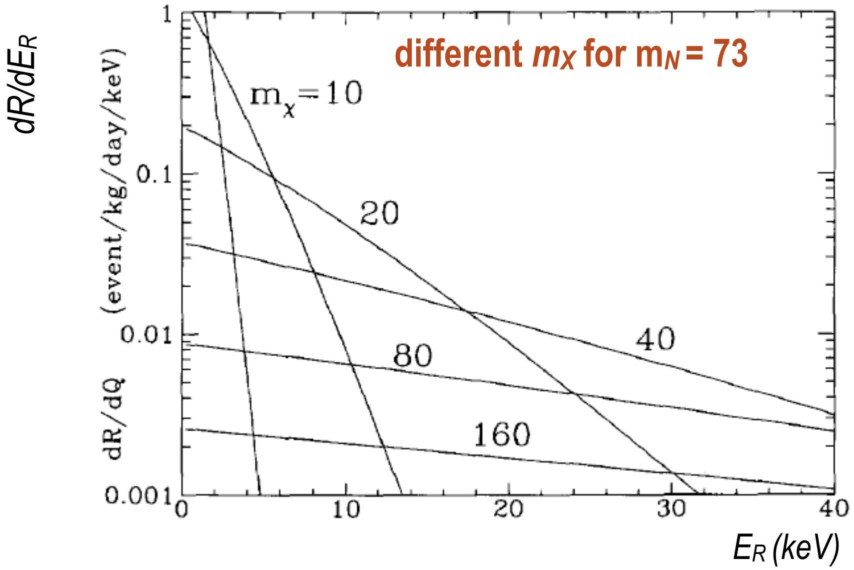


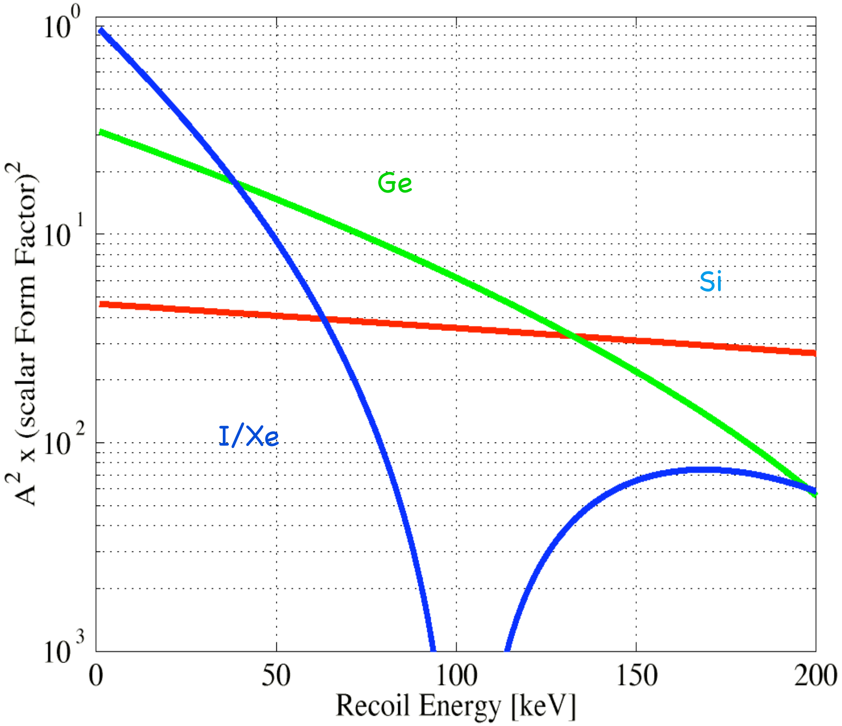
in Universal Extra Dimensions



Summary: Signal Characteristics of a WIMP

- A^2 - dependence of rates
- coherence loss (for $q \sim \mu v \sim 1/r_n \sim 200$ MeV)
- relative rates, for instance in Ge/Si, Ar/Xe,...
- dependance on WIMP mass
- time dependence of the signal (annual, diurnal)

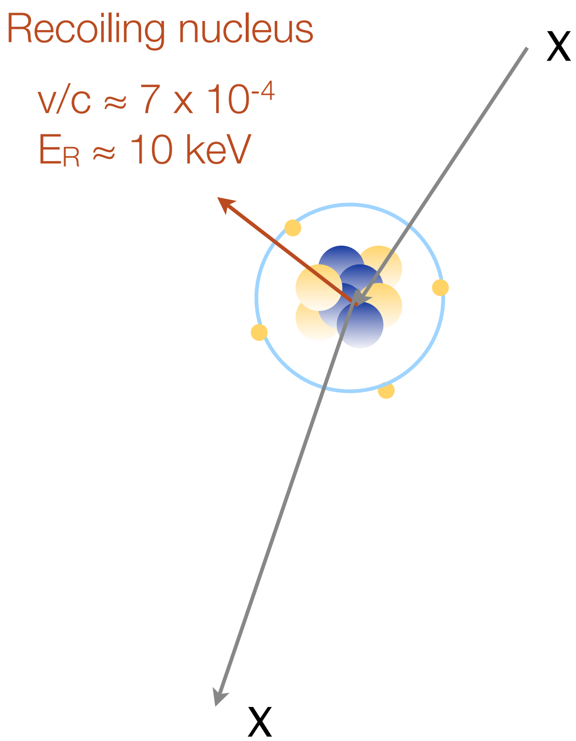


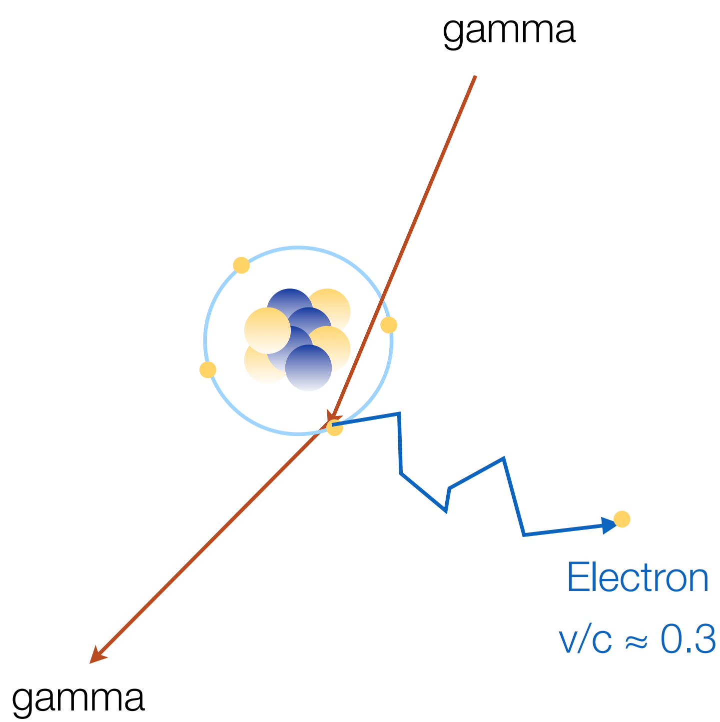


Detection of WIMPs: Signal and Backgrounds

Signal (WIMPs)

Background (gamma-, beta-radiation)





Quenching Factor and Discrimination

- WIMPs (and neutrons) scatter off nuclei
- Most background noise sources (gammas, electrons) scatter off electrons
- Detectors have a different response to nuclear recoils than to electron recoils
- **Quenching factor (QF)** = describes the difference in the amount of visible energy in a detector for these two classes of events
  - ➡ keVee = measured signal from an electron recoil
  - ➡ keVr = measured signal from a nuclear recoil

- **For nuclear recoil events:**

$$E_{visible}(keVee) = QF \times E_{recoil}(keVr)$$

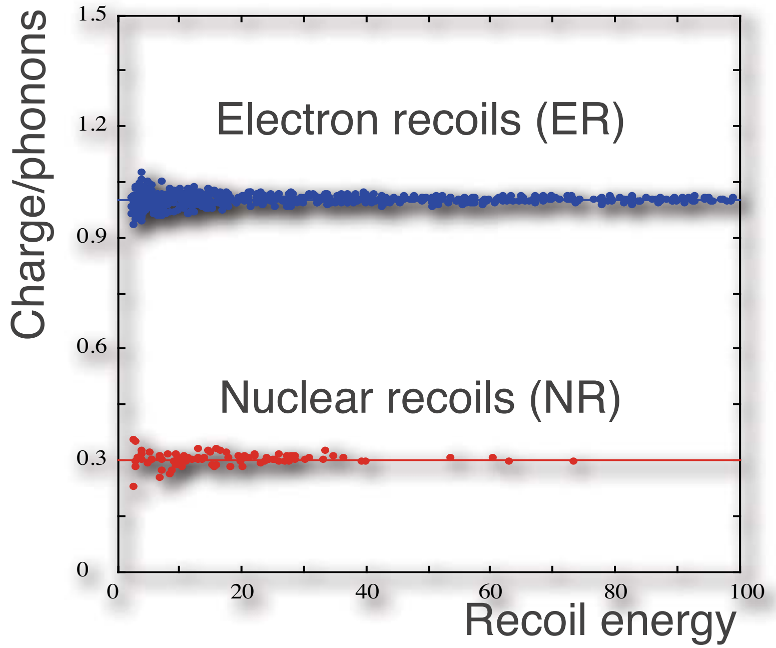
- The two energy scales are calibrated with gamma (^{57}Co, ^{133}Ba, ^{137}Cs, ^{60}Co, etc) and neutron (AmBe, ^{252}Cf, n-generator, etc) sources

Quenching Factor and Discrimination

- The quenching factor allows to distinguish between electron and nuclear recoils if two simultaneous detection mechanisms are used

- **Example**:
- charge and phonons in Ge
- $E_{visible}$ ~ 1/3 E_{recoil} for nuclear recoils
  - ➡ QF ~ 30% in Ge

- ER = background
- NR = WIMPs (or neutron backgrounds)



Backgrounds in Dark Matter Detectors

- Radioactivity of surroundings
- Radioactivity of detector and shield materials
- Cosmic rays and secondary reactions

- Remember: activity of a source $$A = \frac{dN}{dt} = -\lambda N$$
- **Do you know?**

N = number of radioactive nuclei
λ = decay constant, $T_{1/2} = \ln 2/\lambda = \ln 2\ \tau$
[A] = Bq = 1 decay/s (1Ci = 3.7×10^{10} decays/s = A [1g pure ^{226}Ra])

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1. how much radioactivity (in Bq) is in your body? where from?

2. how many radon atoms escape per 1 m^2 of ground, per s?

3. how many plutonium atoms you find in 1 kg of soil?

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1. how much radioactivity (in Bq) is in your body? where from?

1. 4000 Bq from ^{14}C, 4000 Bq from ^{40}K (e^- + 400 1.4 MeV γ + 8000 ν_e)

2. how many radon atoms escape per 1 m^2 of ground, per s?

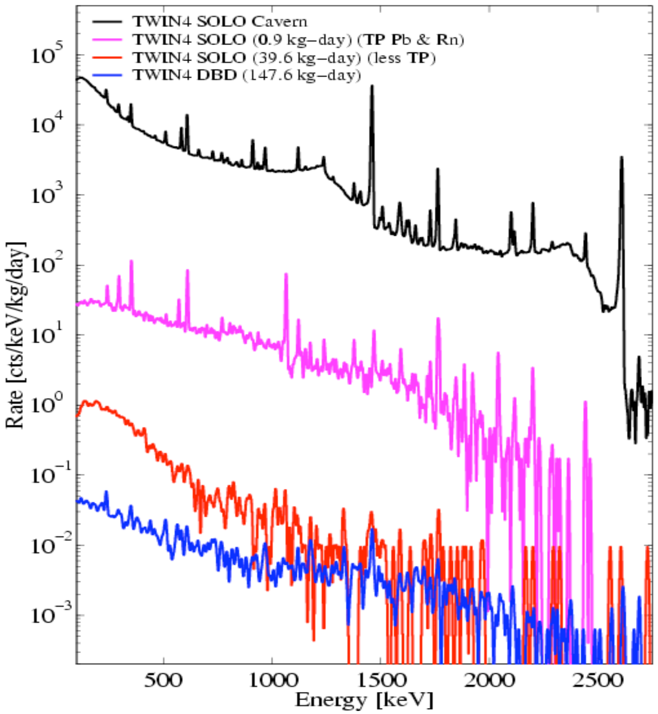
2. 7000 atoms/m^2 s

3. how many plutonium atoms you find in 1 kg of soil?

3. 10 millions (transmutation of ^{238}U by fast CR neutrons), soil: 1 - 3 mg U per kg

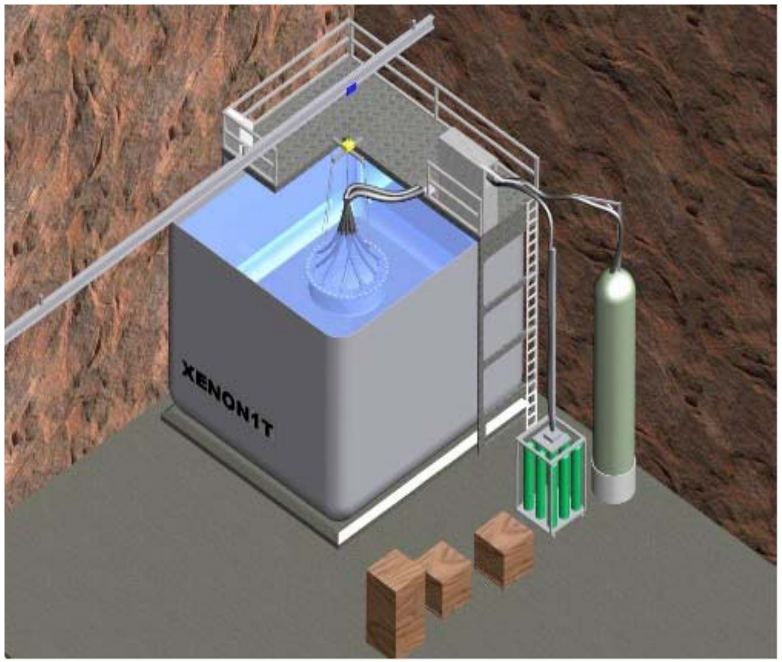
Backgrounds in Dark Matter Detectors

- **External, natural radioactivity:** ^{238}U, ^{238}Th, ^{40}K decays in rock and concrete walls of the laboratory => mostly gammas and neutrons from (α,n) and fission reactions
- **Radon decays in air:**
  - ➡ **passive shields:** Pb against the gammas, polyethylene/water against neutrons
  - ➡ **active shields:** large water Cerenkov detectors or scintillators for gammas and neutrons



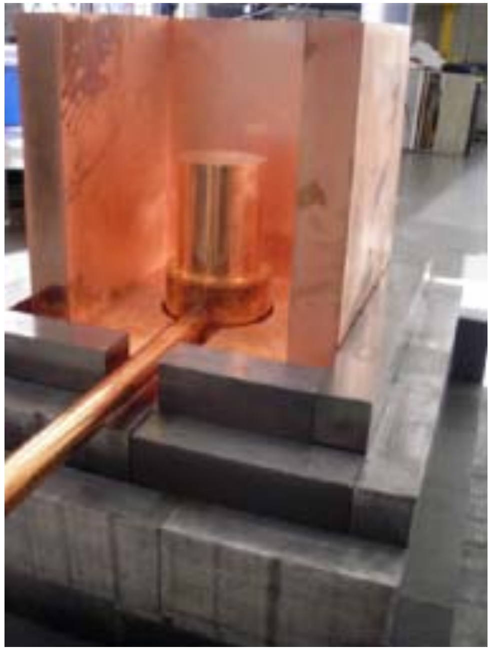
Ge detector underground, no shield

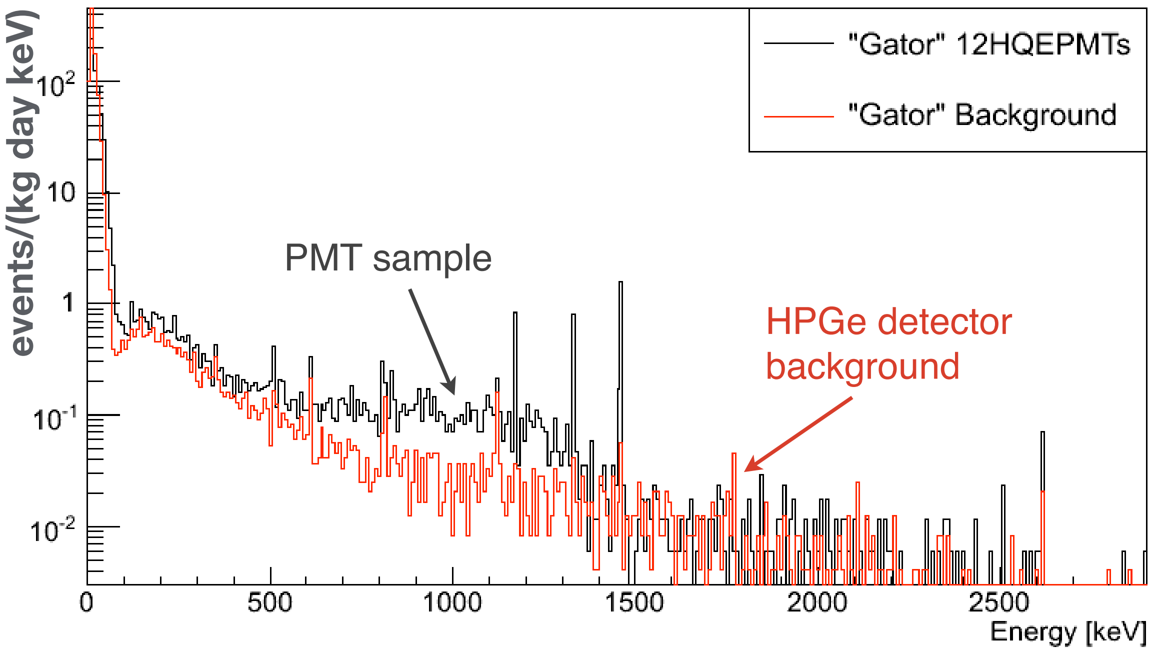
Ge detector underground, Pb shield and purge for Rn



Backgrounds in Dark Matter Detectors

- **Internal radioactivities:** ^{238}U, ^{238}Th, ^{40}K, ^{137}Cs, ^{60}Co, ^{39}Ar, ^{85}Kr, ... decays in the detector materials, target medium and shields
- Ultra-pure Ge spectrometers (as well as other methods) are used to screen the materials before using them in a detector, down to parts-per-billion (ppb) (or lower) levels

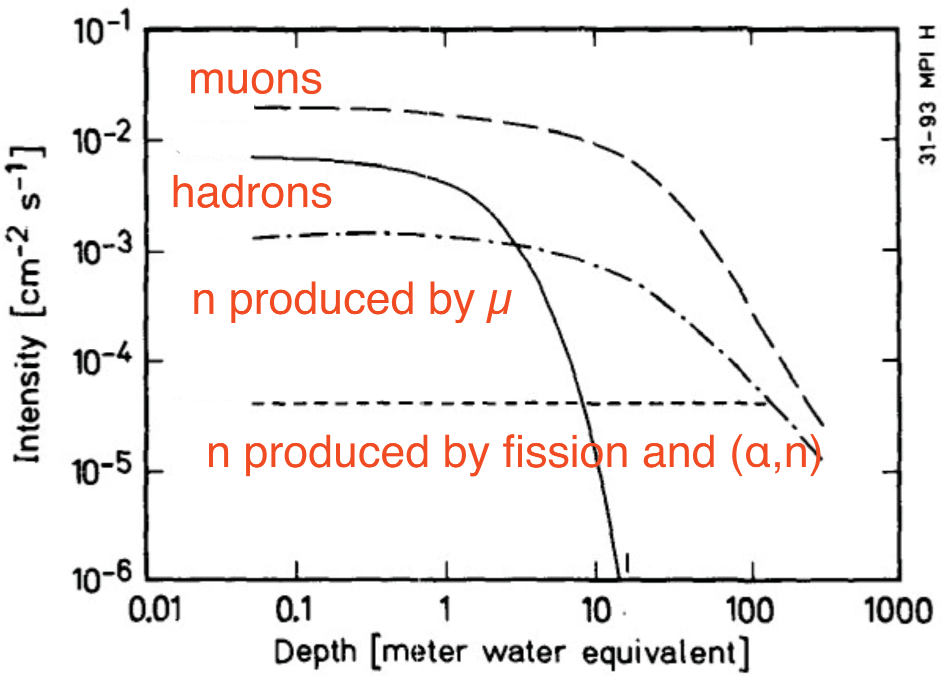




Backgrounds in Dark Matter Detectors

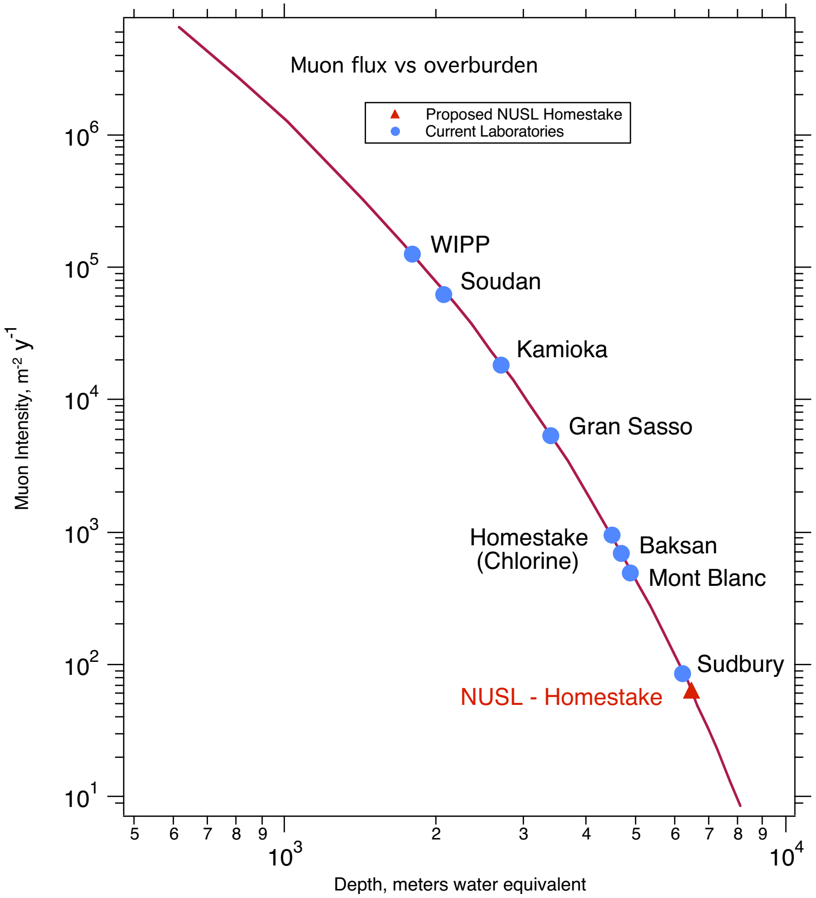
- Cosmic rays and secondary/tertiary particles: go underground!
- Hadronic component (n, p): reduced by few meter water equivalent (mwe)

Flux of cosmic ray secondaries and
tertiary-produced neutrons in a typical
Pb shield vs shielding depth
Gerd Heusser, 1995



Backgrounds in Dark Matter Detectors

- **Most problematic:** muons and muon induced neutrons
  ➡ **go deep underground**, several laboratories, worldwide



Site (multiple levels given in ft)	Relative muon flux	Relative neutron flux T > 10 MeV
WIPP (2130 ft) (1500 mwe)	× 65	× 45
Soudan (2070 mwe)	× 30	× 25
Kamioke	× 12	× 11
Boulby	× 4	× 4
Gran Sasso (3700 mwe)		
Frejus (4000 mwe)	× 1	× 1
Homestake (4860 ft)		
Mont Blanc	$\times 6^{-1}$	$\times 6^{-1}$
Sudbury	$\times 25^{-1}$	$\times 25^{-1}$
Homestake (8200 ft)	$\times 50^{-1}$	$\times 50^{-1}$

compiled by: R. Gaitskell

Backgrounds in Dark Matter Detectors

- **Activation of detector and other materials during production and transportation at the Earth's surface. A precise calculation requires:**
  - ➡ cosmic ray spectrum (varies with geomagnetic latitude)
  - ➡ cross section for the production of isotopes (only few are directly measured)
- production is dominated by (n,x) reactions (95%) and (p,x) reactions (5%)

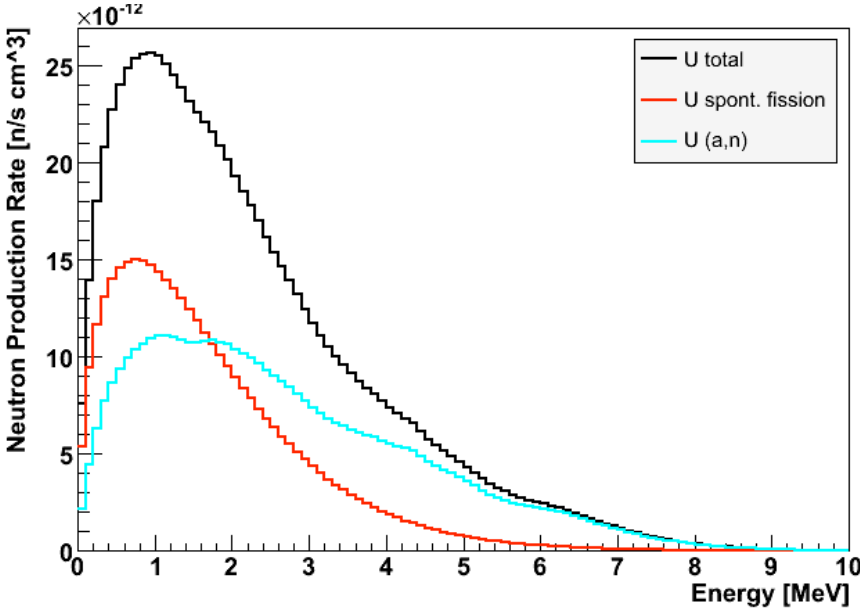
production in Ge after 30d exposure at the Earth's surface and 1 yr storage below ground

Isotope	**Decay**	**Half life**	**Energy in Ge [keV]**	**Activity [μBq/kg]**
^{3}H	β^-	12.33 yr	$E_{max(\beta-)}$=18.6	2
^{49}V	EC	330 d	$E_{K(Ti)}$ = 5	1.6
^{54}Mn	EC, β^+	312 d	$E_{K(Cr)}$ = 5.4, E_γ=841	0.95
^{55}Fe	EC	2.7 yr	$E_{K(Mn)}$ = 6	0.66
^{57}Co	EC	272 d	$E_{K(Fe)}$=6.4, E_γ=128	1.3
^{60}Co	β^-	5.3 yr	$E_{max(\beta-)}$=318, E_γ=1173,1333	0.2
^{63}Ni	β^-	100 yr	$E_{max(\beta-)}$=67	0.009
^{65}Zn	EC, β^+	244 d	$E^{K(Cu)}$ = 9, E_γ=1125	9.2
^{68}Ge	EC	271 d	$E_{K(Ga)}$ = 10.4	172

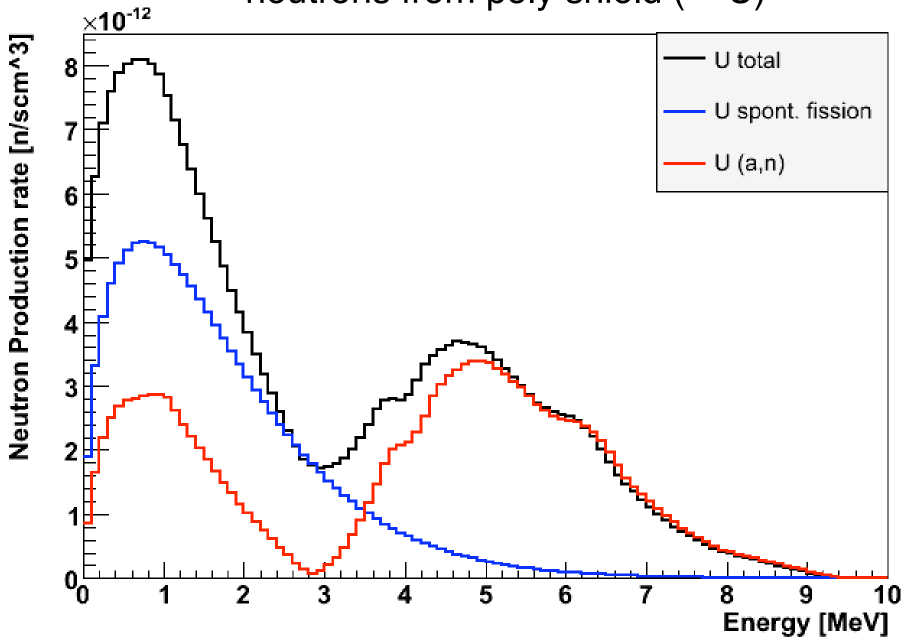
Neutron Backgrounds

- **MeV neutrons can mimic WIMPs** by elastically scattering from the target nuclei
- the rates of neutrons from detector materials and rock are calculated taking into account the exact material composition, the α energies and cross sections for (α,n) and fission reactions and the measured U/Th contents

neutrons from rock (^{238}U)



neutrons from poly shield (^{238}U)



E. Tziaferi

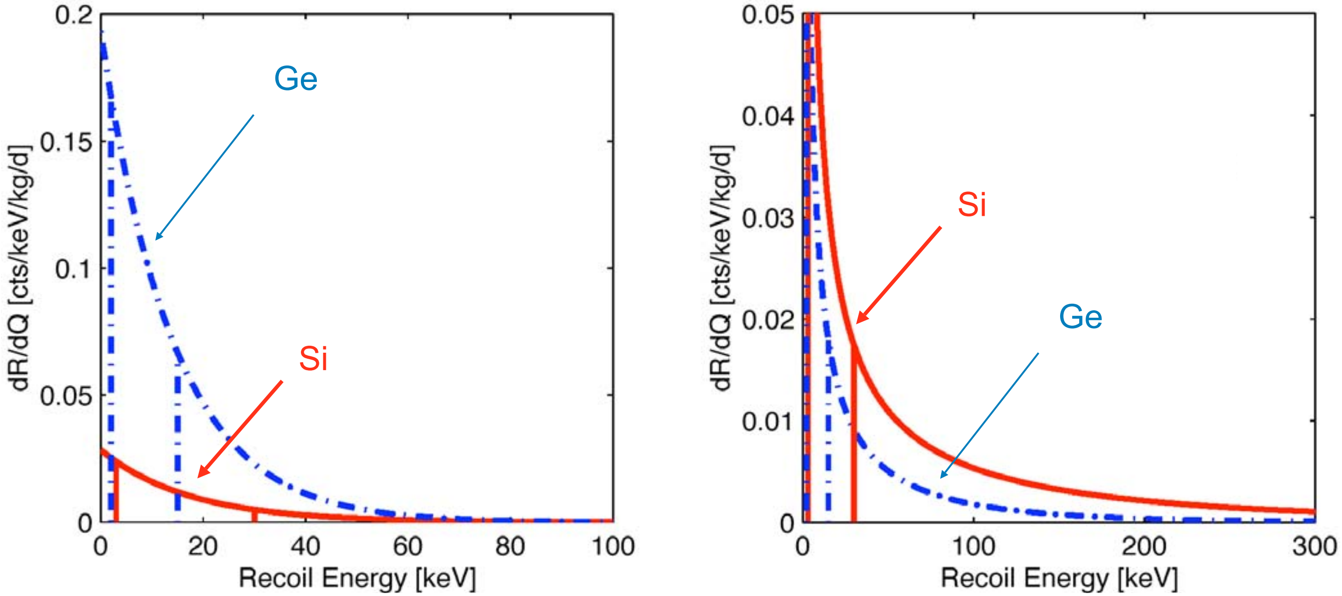
Neutrons: how can we distinguish them from WIMPs?

➡ mean free path of few cm (neutrons) versus 10^{10} m (WIMP)

➡ material dependence of differential recoil spectrum

➡ time dependence of WIMP signal (if neutron background is measured to be constant in time)

WIMP signal ~ A^2



WIMPs, M_X = 40 GeV **neutrons**

Detector strategies

Aggressively reduce the absolute background	Background reduction by pulse shape analysis and/or self-shielding	Background rejection based on simultaneous detection of two signals	Other detector strategies
State of the art: (primary goal is $0\nu\beta\beta$ decay): Heidelberg-Moscow HDMS IGEX Near future projects: GERDA MAJORANA	Large mass, simple detectors: NaI (DAMA, LIBRA, ANAIS, NAIAD) CsI (KIMS) Large liquid noble gas detectors: XMASS, CLEAN, DEAP	**Charge/phonon** (CDMS, EDELWEISS, SuperCDMS, EURECA) **Light/phonon** (CRESST, ROSEBUD, EURECA) **Charge/light** (XENON, ZEPLIN, LUX, ArDM, WARP)	Large bubble chambers - insensitive to electromagnetic background (COUPP, PICASSO) Low-pressure gas detectors, sensitive to the direction of the nuclear recoil (DRIFT)

In addition:
reject multiple scattered events and events close to detector boundaries
look for an annual and a diurnal modulation in the event rate

End