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Workshop on Supersolid 2008

18 - 22 August 2008

The structure and dynamics of hcp He-4 below 1 K

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# Structure and dynamics of hcp <sup>4</sup>He below 1 K

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#### ICTP Trieste – Supersolid 2008



# Structure and dynamics of hcp <sup>4</sup>He below 1 K

#### Structure

Joost van Duijn (Madrid)

Oleg Kirichek

Chris Frost



Richard Down

#### Both

John Goodkind Sunil Sinha





John Copley Yiming Qiu Ross Erwin

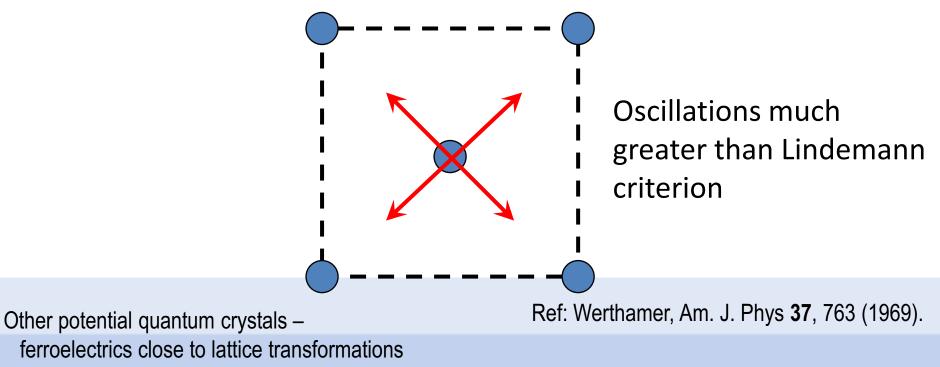


Collin Broholm



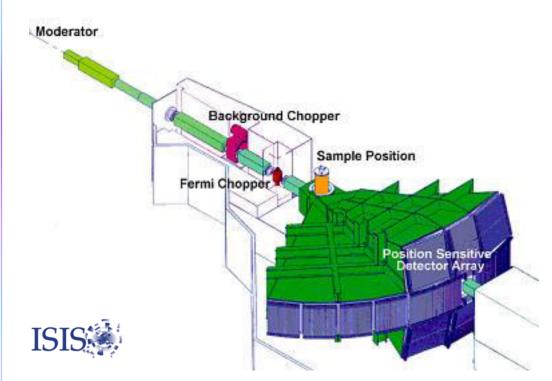
## A quantum crystal

- Helium only solidifies under pressure
  - light atom; no overall charge
  - tightly bound electronic shells  $\rightarrow$  small vdW forces
  - large zero-point fluctuations



### Are there any structural changes?

#### **MAPS time-of-flight spectrometer**



#### **Characteristics**

Incident white beam of neutrons Use time-of-flight detection to extract *d*-spacings. No need to move

sample once grown

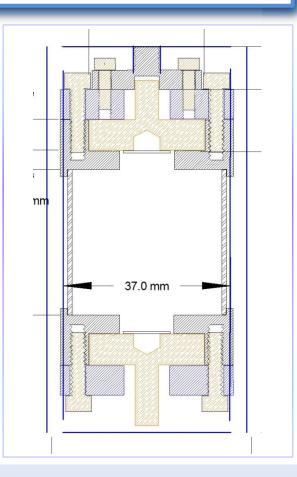
#### **Sample chamber**

Stainless steel walls.

Fill line at top; cold spot at bottom.

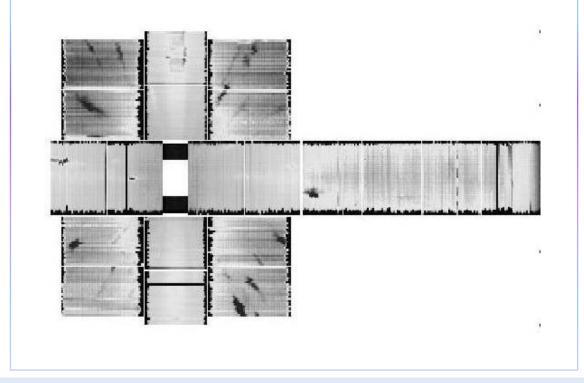
Quartz transducers covering vertical path.

Sample grown *in situ* starting with a pressure of 45 bar at ~ 2 K using the blocked capillary technique.



# Finding Bragg peaks

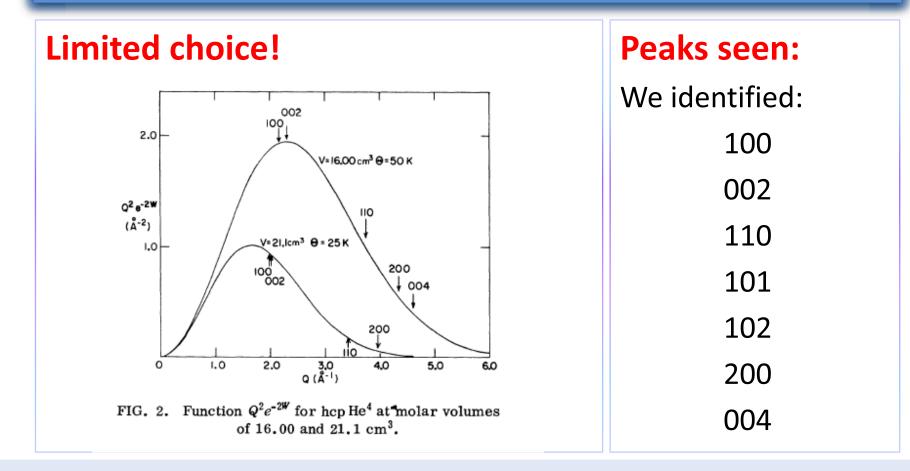
#### **MAPS detector banks**



#### **Characteristics**

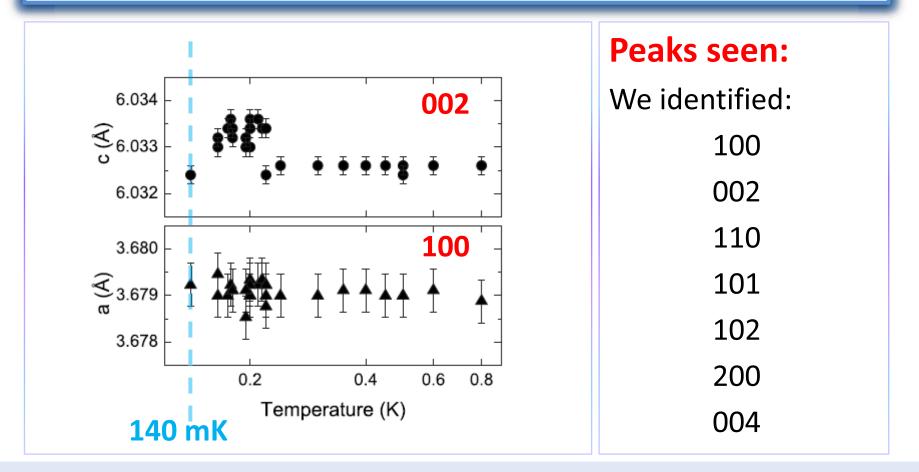
Incident white beam of neutrons Use time-of-flight detection to extract *d*-spacings. No need to move sample once grown.

## Finding Bragg peaks



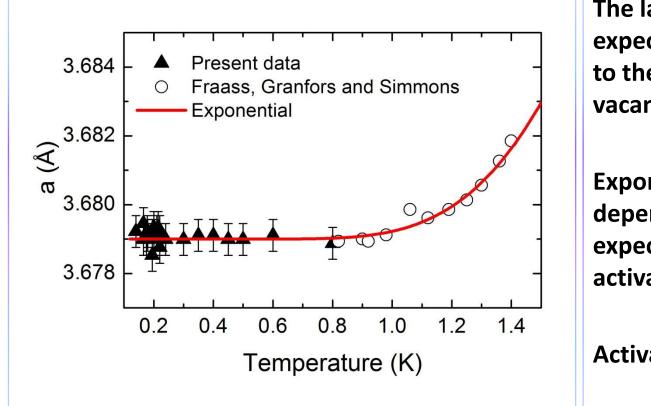
Ref: Reese et al., Phys. Rev. A 3, 1688 (1971).

#### Lattice parameters



Ref: Blackburn et al., Phys. Rev. B 76, 024523 (2007).

## Vacancy population



The lattice parameter is expected to shift with to the number of vacancies.

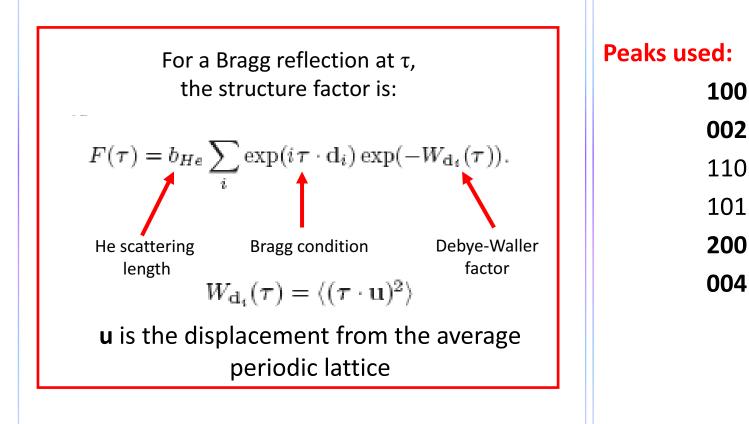
Exponential dependence as expected for thermal activation.

Activation energy 8.6 K

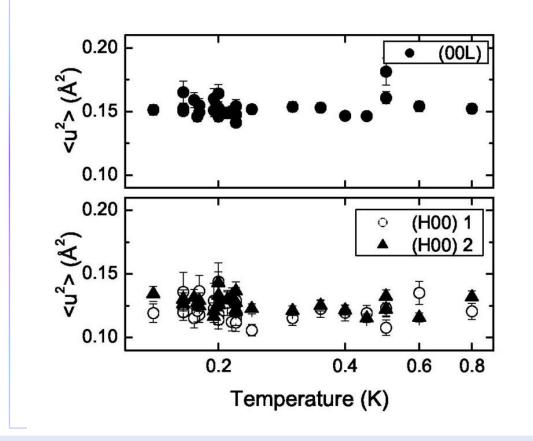
Caveat: mass density assumed to remain constant as function of temperature.

Ref: Fraass, Granfors and Simmons, Phys. Rev. B 39, 124 (1989).

### Mean square atomic displacement



### Mean square atomic displacement



The mean square displacement is anisotropic, as expected for an hcp crystal. It does not change with temperature;

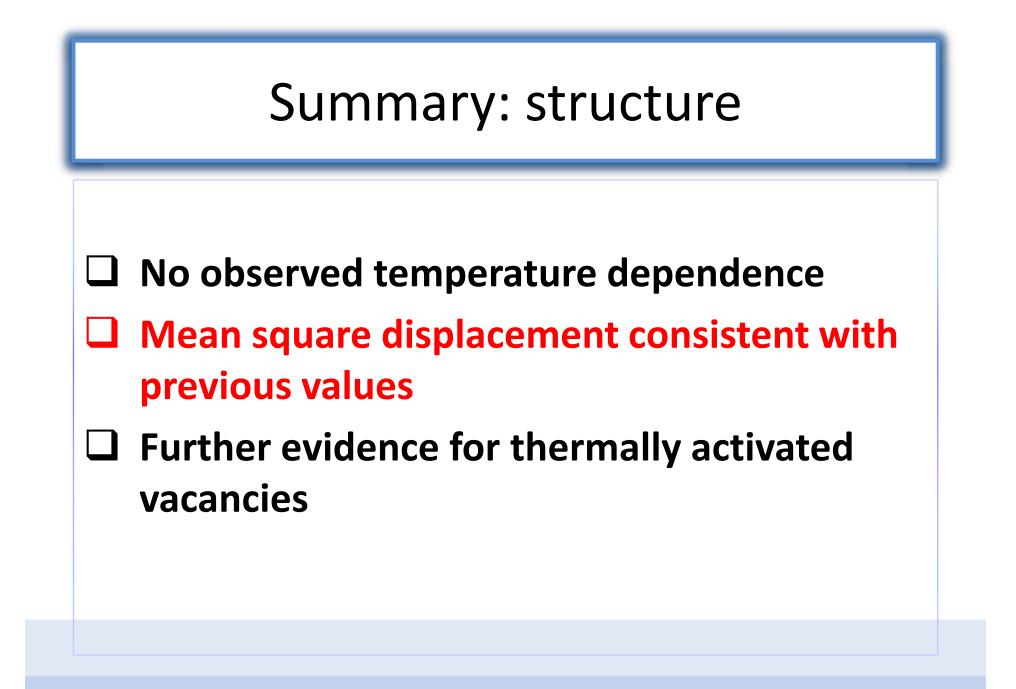
we are looking at zero-point motion.

Ref: Blackburn et al., Phys. Rev. B 76, 024523 (2007).

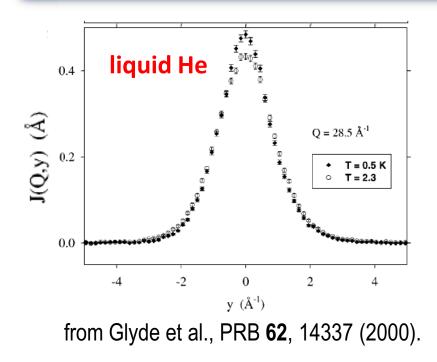
### Mean square atomic displacement

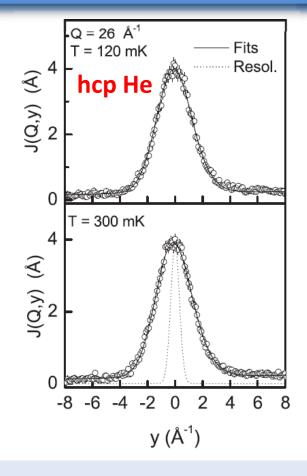
TABLE II. Published values for  $\langle u^2 \rangle$  in solid <sup>4</sup>He obtained by x-ray and neutron diffraction. In all cases, the harmonic approximation is assumed and for the mixed peak data, no distinction between in-plane and out-of-plane displacements was made.

Molar volume (cm <sup>3</sup> )	Temperature (K)	$\langle u^2 \rangle$ (Å <sup>2</sup> )	Peak type	Reference
11.01	15	0.0593(1)	mixed	Venkataraman and Simmons, PRB 68 (2003) – x-ray
12.06	5.8	0.0466(3)	mixed	Stassis, Khatamian and Kline, SSC 25 (1978) – neutron
12.12	14.8	0.0563(14)	mixed	Arms, thesis (1999) – x-ray
12.13	14.8	0.0513(10)	mixed	Arms, thesis (1999) – x-ray
15.72	5.8	0.0861(9)	mixed	Stassis, Khatamian and Kline, SSC 25 (1978) – neutron
20.9	0.7	0.1537(7)	(00L)	Burns and Isaacs, PRB 55 (1997) – x-ray
21.3	<1	0.150(1)	(00L)	This work (neutron)
21.3	<1	0.122(1)	(H00)	This work (neutron)



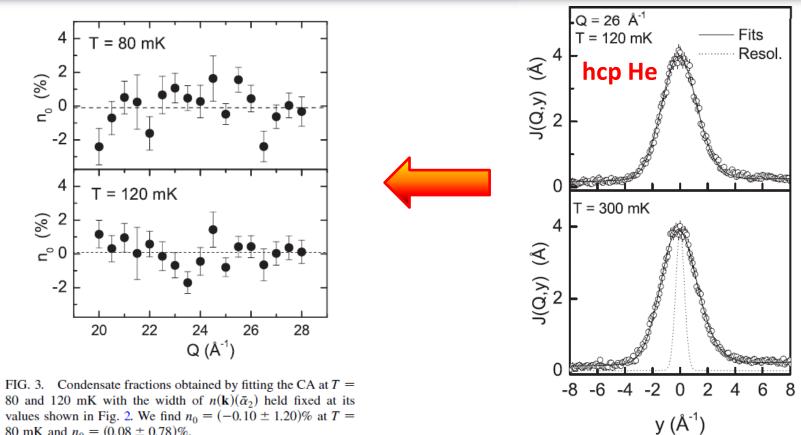
#### Momentum distribution – signs of $\rho_s$





from Diallo et al., Phys. Rev. Lett. 98, 205301 (2007).

#### Momentum distribution – signs of $\rho_s$

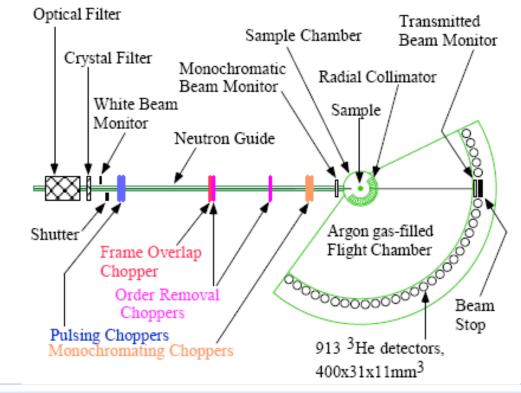


80 mK and  $n_0 = (0.08 \pm 0.78)\%$ .

Refs: Adams et al., Phys. Rev. Lett. 98, 085301 (2007), Diallo et al., Phys. Rev. Lett. 98, 205301 (2008).

#### What about the excitation spectrum?

#### **DCS time-of-flight spectrometer**



#### **Characteristics**

Monochromatic beam.

Use time-of-flight detection to extract energy transfer information.

Sample needs limited rotation to get full coverage.



#### Sample chamber

Stainless steel walls.

Fill line at top; cold spot at bottom.

Quartz transducers covering vertical path, and X-cut transducers horizontally at top.

Sample grown *in situ* along the melting curve, at ≈ 1.2 K.



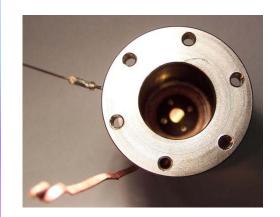
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Stainless steel walls.

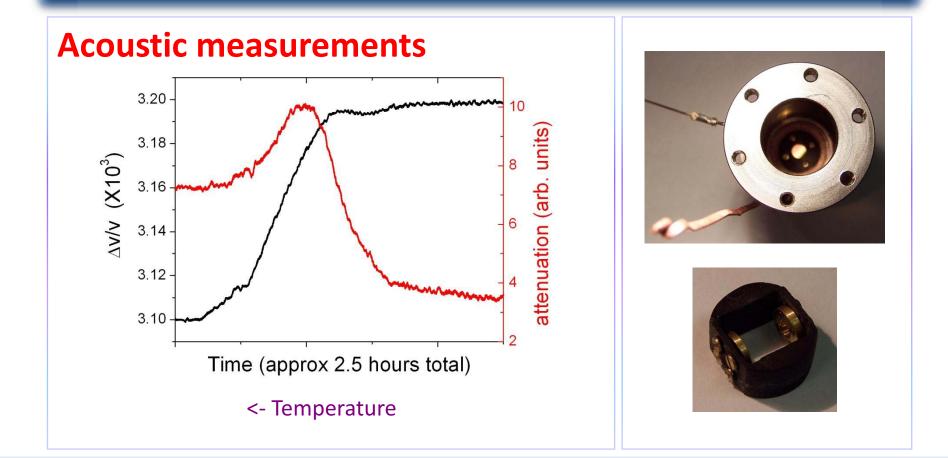
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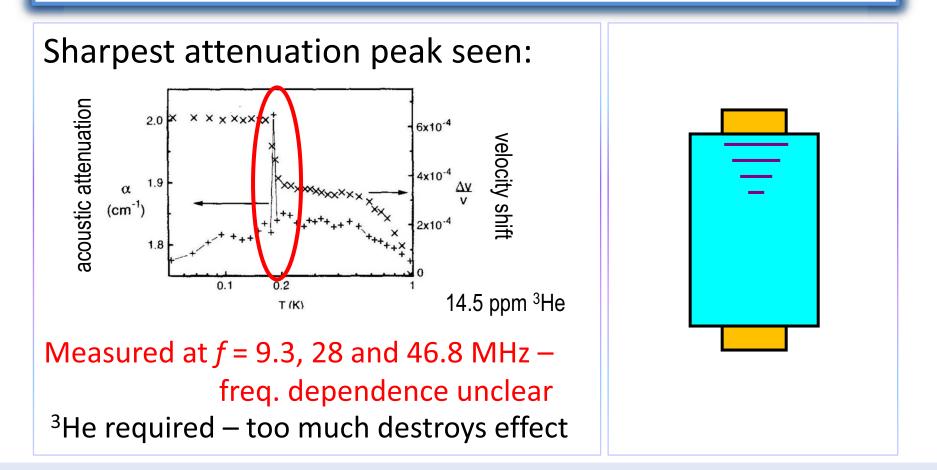






Definitely got below 200 mK; thermometry seemed correct. Closest thermometer was on mixing chamber.

## Ultrasound measurements



Ref: Ho, Bindloss and Goodkind, J. Low. Temp Phys. 109, 409 (1997).

### Ultrasound measurements

Non phonon thermal excitation seen.

The excitation has a speed of 226 ms<sup>-1</sup>.

Indication that for T > 200 mK these excitations are gapped.

This gap disappears below 200 mK.

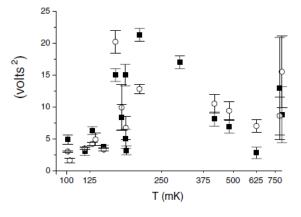
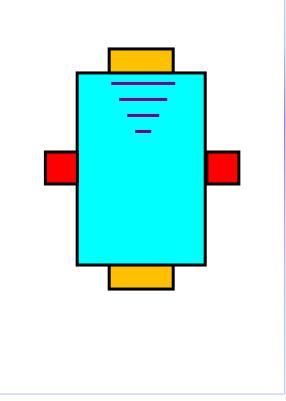


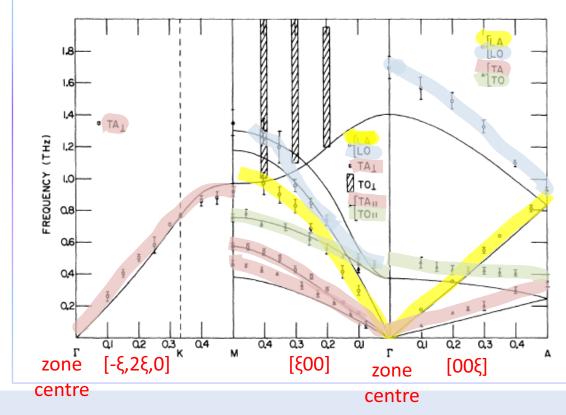
FIG. 5. Values of  $\phi$  from fits of an exponential function to the shifts of  $\alpha$  and  $\nu$  vs  $V_{hp}^2$ .  $\blacksquare$  from attenuation;  $\circ$  from velocity.

Ref: Goodkind, Phys. Rev. Lett. 89, 095301 (2002).



#### What has previously been measured?

#### **Summary of measured spectra**

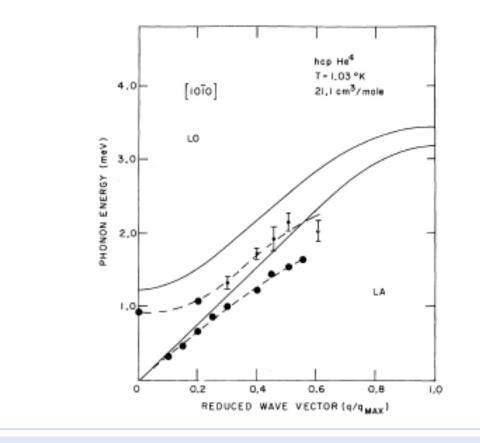


#### References

Molar volume 16cm<sup>3</sup> Reese et al., PRA **3**, 1688 (1971).

Molar volume 21.1cm<sup>3</sup> Minkiewicz et al., PR **174**, 267 (1968).

#### Phonon spectrum in the basal plane

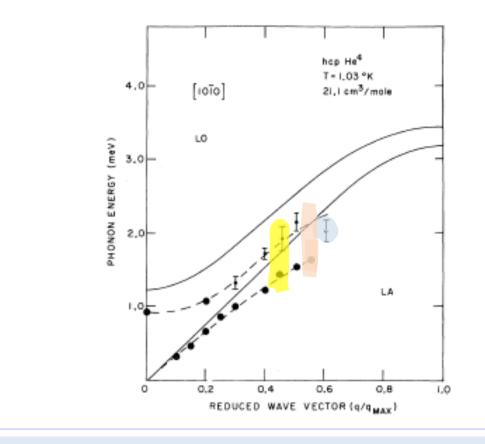


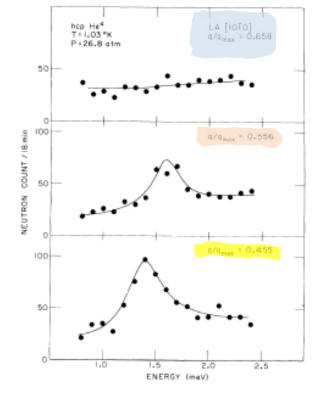
#### References

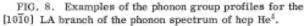
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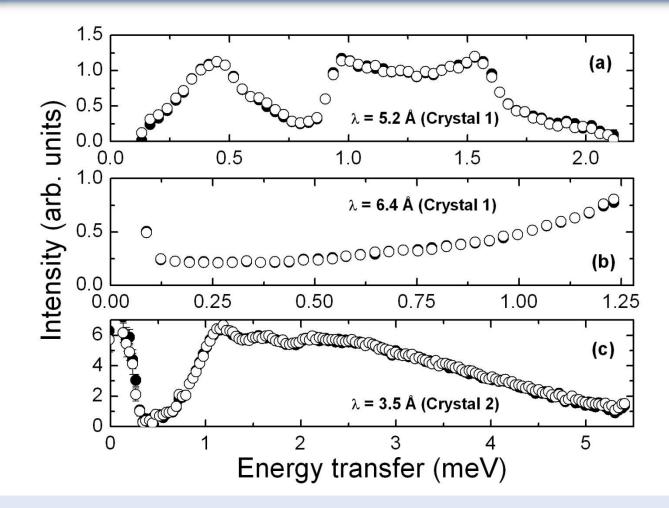
Molar volume 21.1cm<sup>3</sup> Minkiewicz et al., PR **174**, 267 (1968).

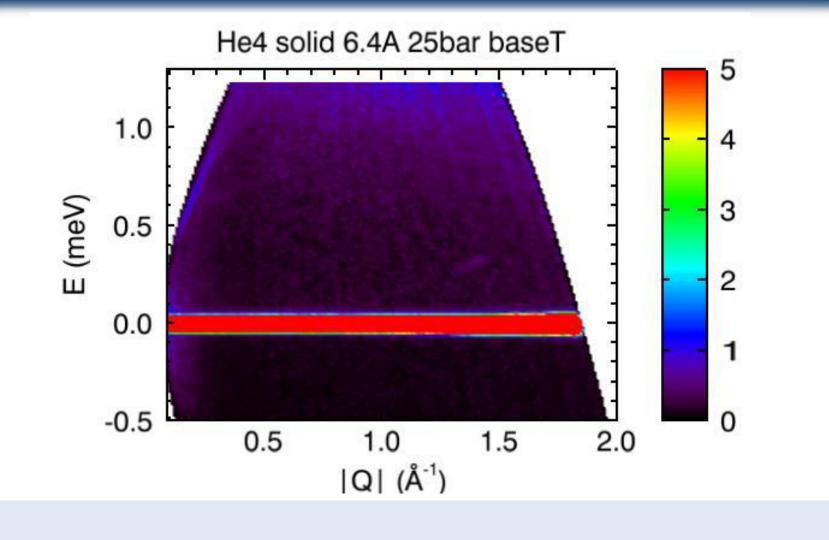
#### Phonon spectrum in the basal plane

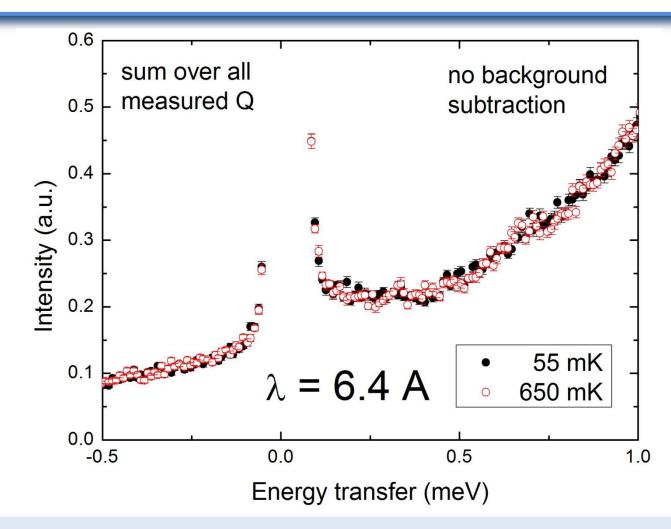


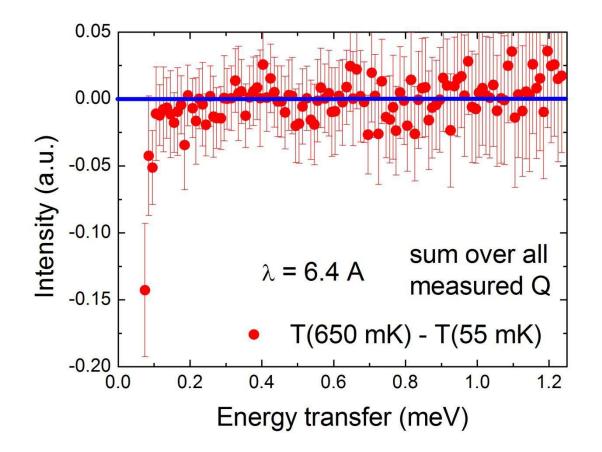


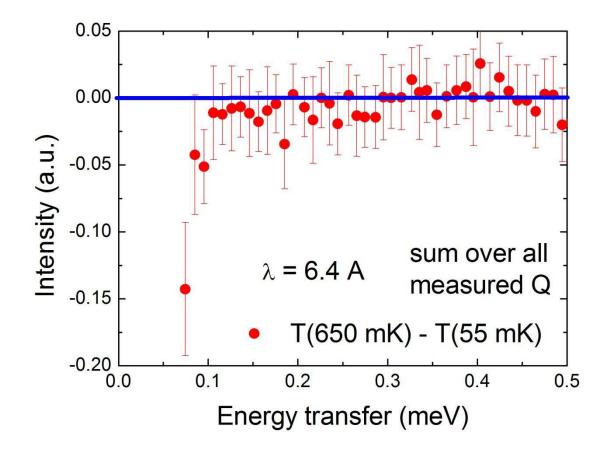




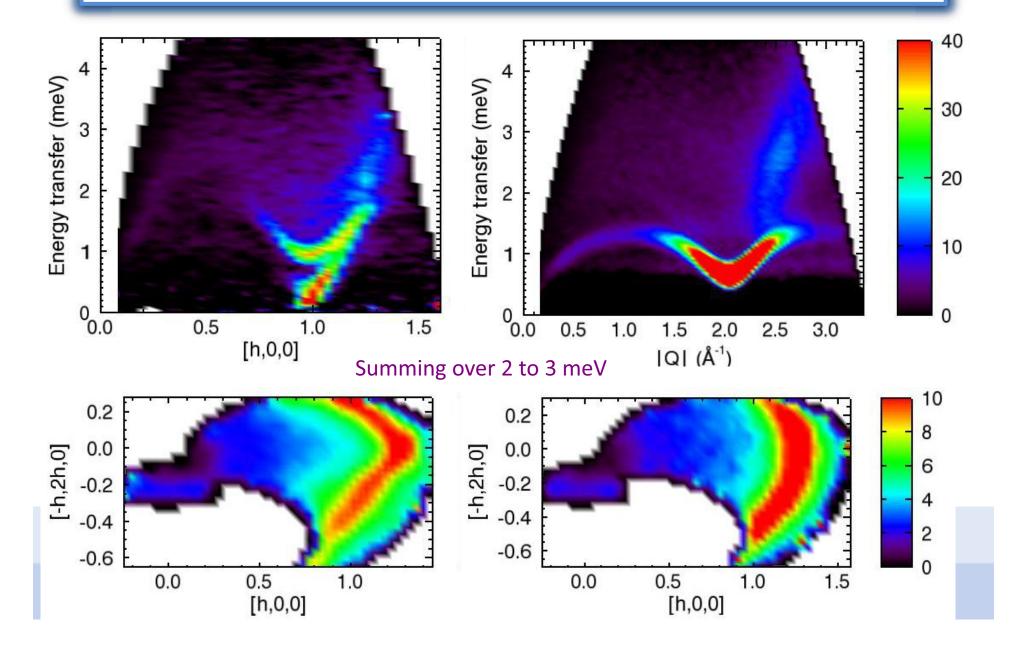




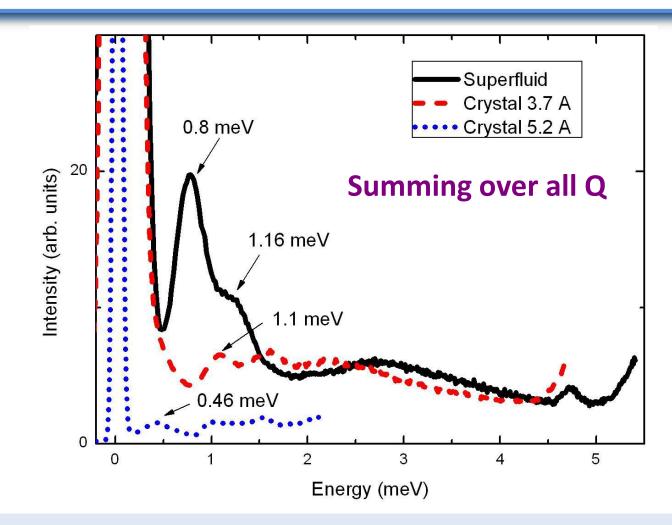


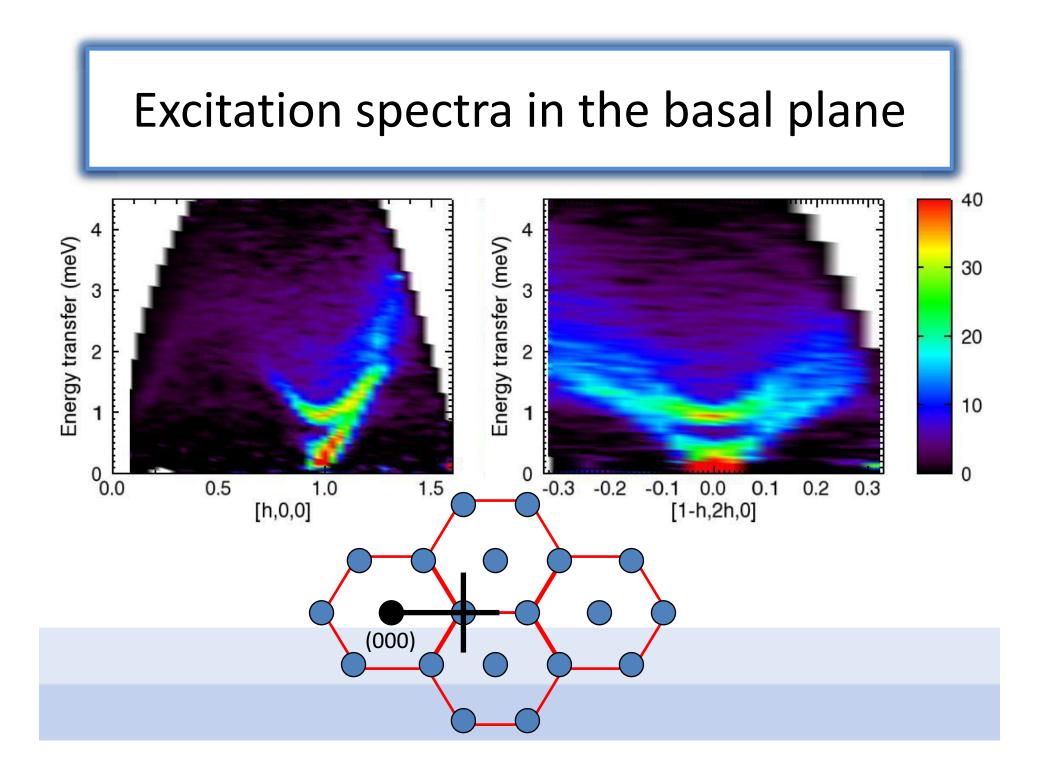


## Comparison with the superfluid

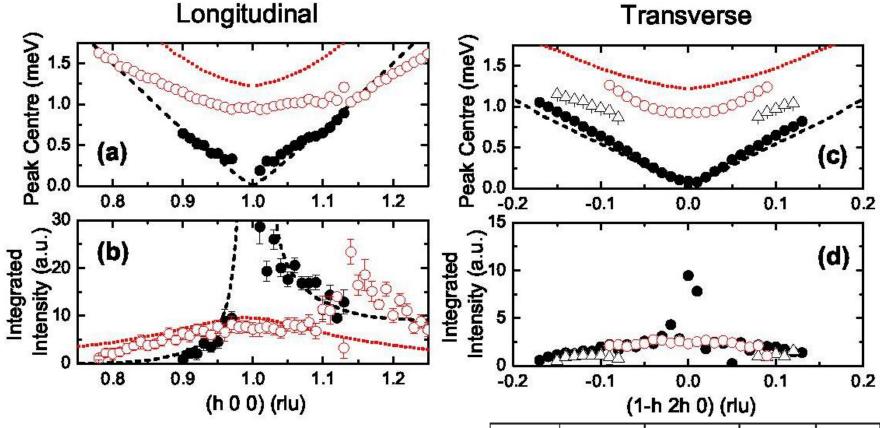


# Any superfluid in the sample?





# Sharp, resolution-limited modes

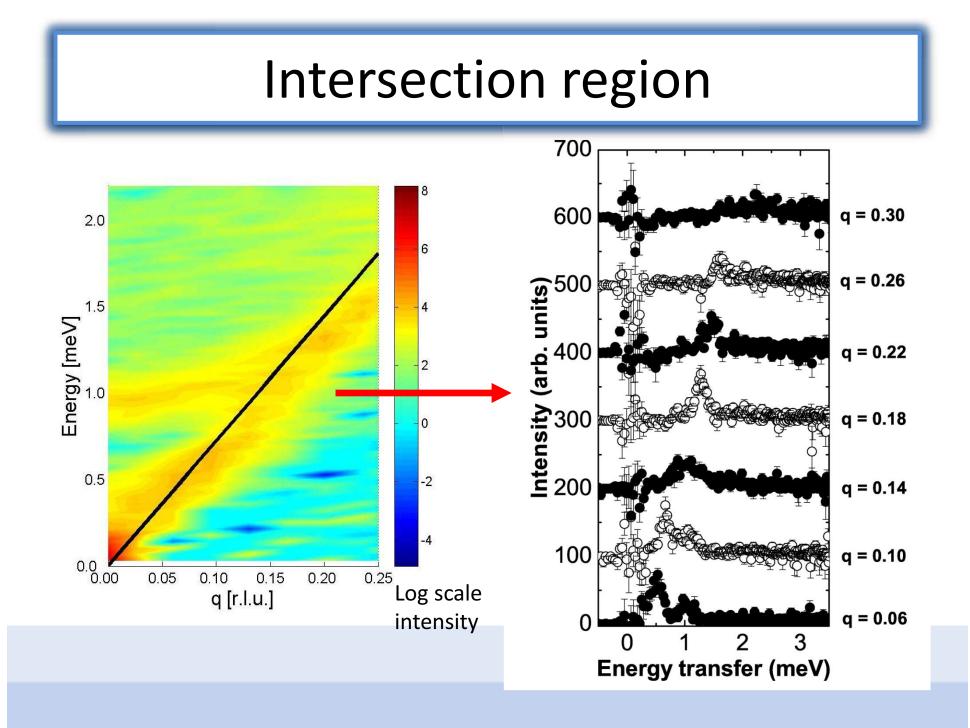


Lines are from calculations using a selfconsistent phonon theory

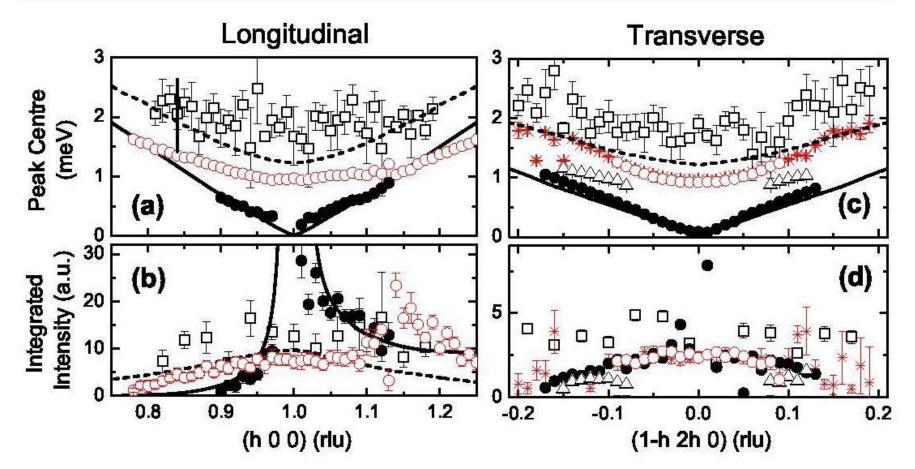
Ref: Gillis, Koehler and Werthamer, Phys. Rev. **175**, 1110 (1968).

Material	Mode	Effective mass	Energy gap	$Q_{\min \text{imum}}$
		$\rm m_{He}$	$\mathrm{meV}$	$Å^{-1}$
Crystal 1	Longitudinal	0.107(3)	0.951(4)	1.964(6)
Crystal 2	Longitudinal	0.114(5)	1.002(3)	1.970(6)
Crystal 1	Transverse	0.144(2)	0.916(2)	1.970(1)
Superfluid	Roton	0.094(2)	0.680(1)	2.054(1)

Table I: Characteristics of the vacancy mode in the solid and the roton mode in the superfluid at 24 bar.

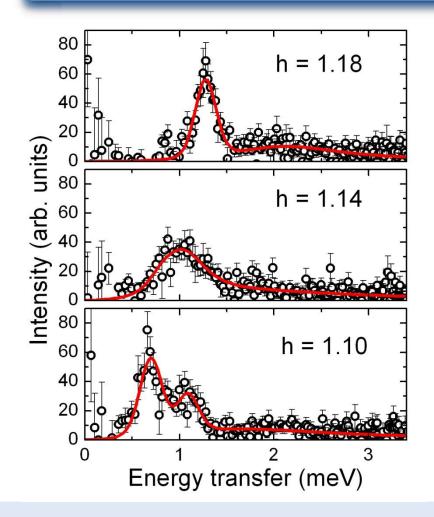


## Above the sharp mode





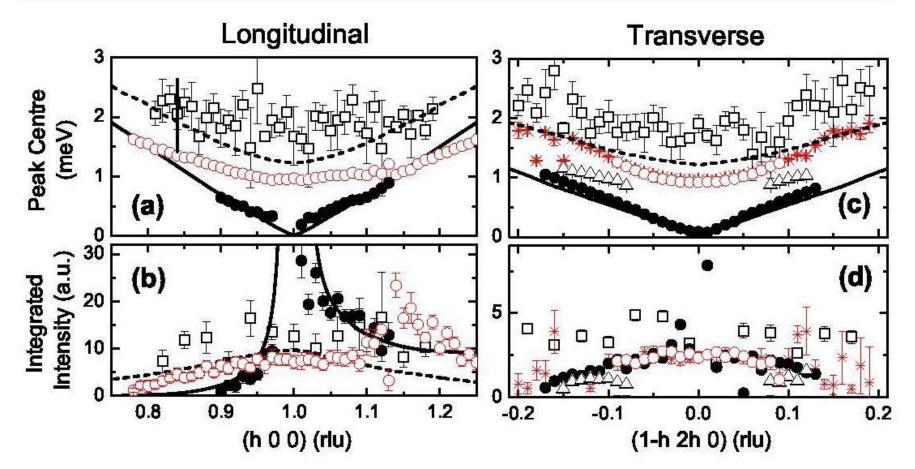
# So where is the optic phonon?



Fitting using damped harmonic oscillators for all of poles, with positions and relative intensities of the acoustic and optic modes fixed using the GKW calculations.

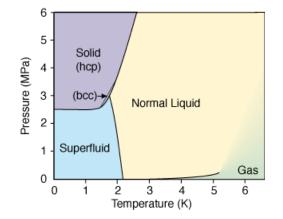
Convoluted with instrumental resolution.

## Above the sharp mode

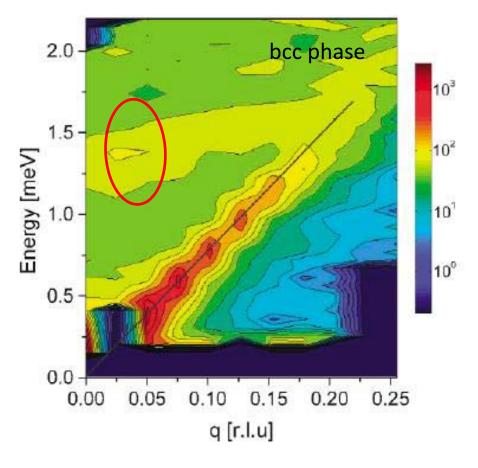




## Comparison with bcc phase

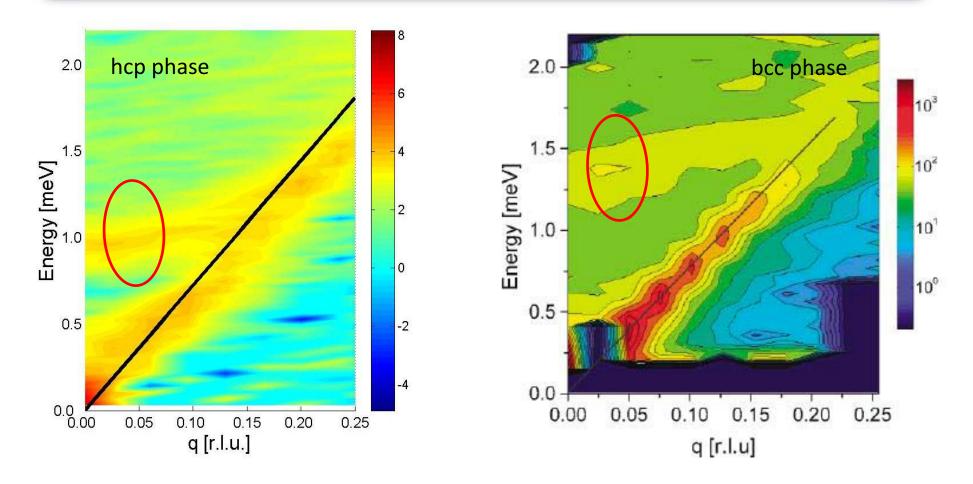


The bcc phase has one atom per unit cell therefore there is no optic phonon



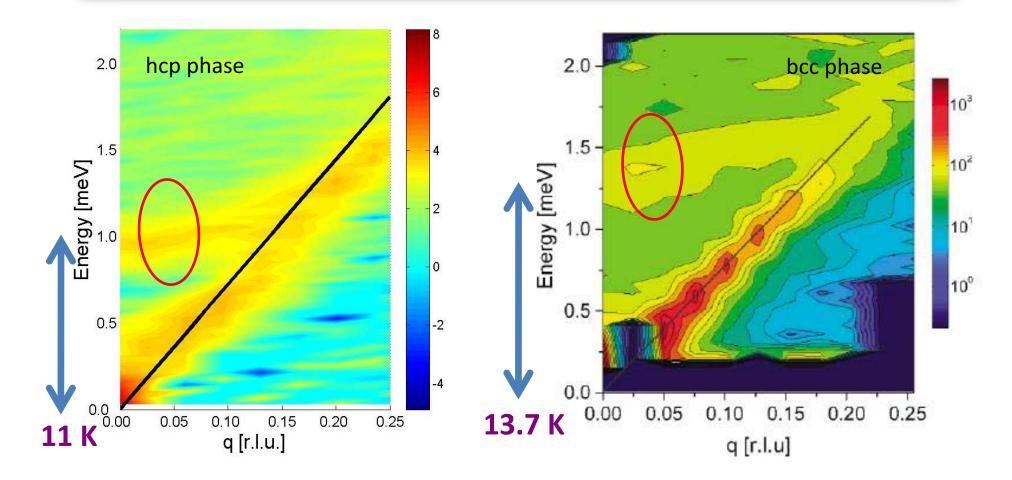
from: Markovich et al., Phys. Rev. Lett. **88**, 195301 (2002).

## Comparison with bcc phase



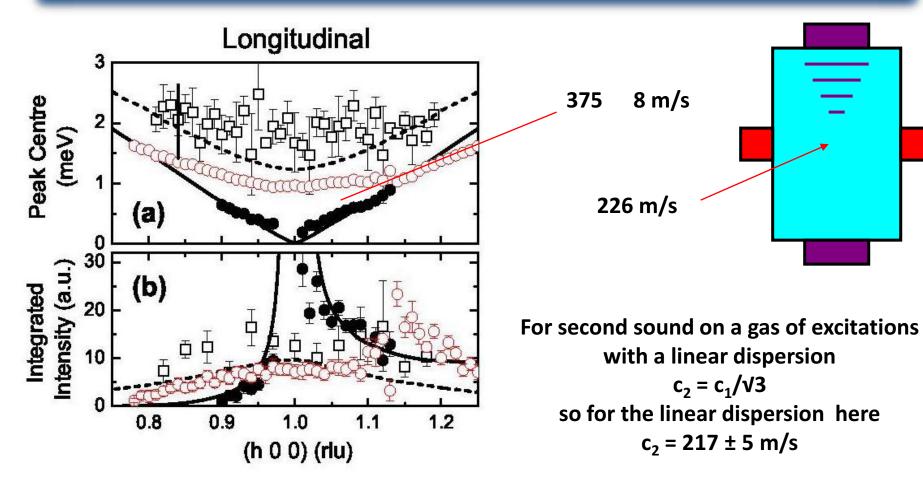
from: Markovich et al., Phys. Rev. Lett. **88**, 195301 (2002).

## Comparison with bcc phase



from: Markovich et al., Phys. Rev. Lett. **88**, 195301 (2002).

## Relation to bolometer expt?

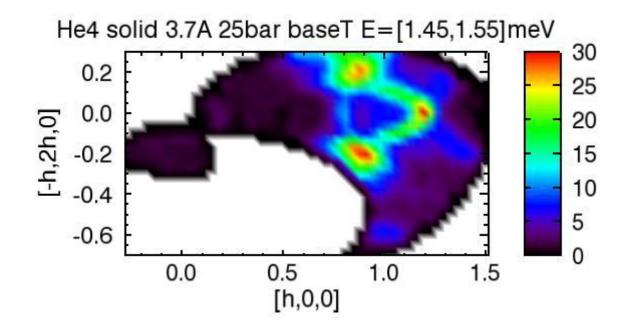




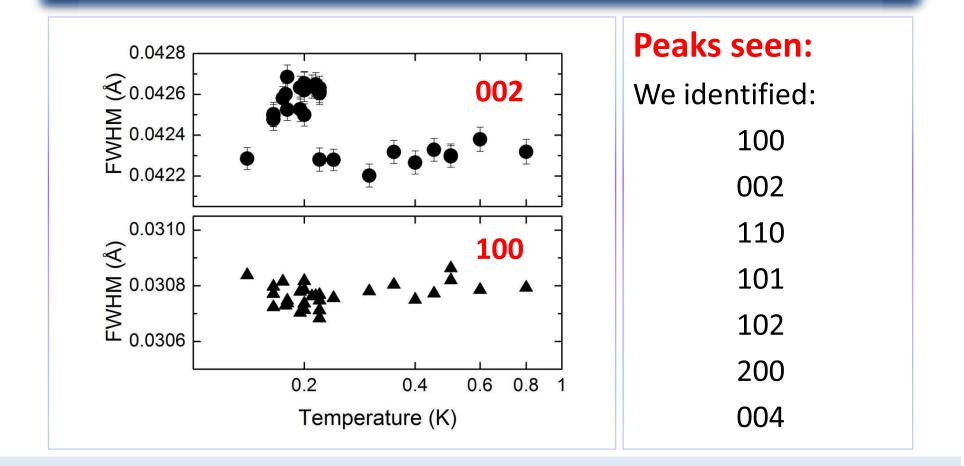
 $c_2 = c_1 / \sqrt{3}$ 

## Summary

- Observation of a possible delocalized vacancy excitation
- No temperature dependence observed

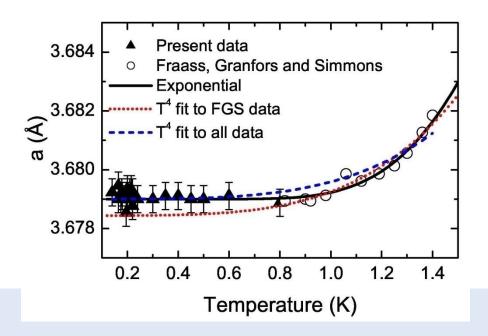


# Bragg peak width



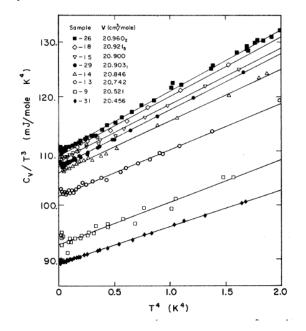
# Vacancies in helium crystals

There is a lot of evidence for the presence of thermally activated vacancies above 1 K. Anderson, Brinkman and Huse proposed a model of (delocalized) zero-point vacancies forming a condensate at low temperatures.



Ref: Blackburn et al., Phys. Rev. B 76, 024523 (2007).

The heat capacity of solid helium shows a roughly T<sup>7</sup> dependence, as yet unexplained.

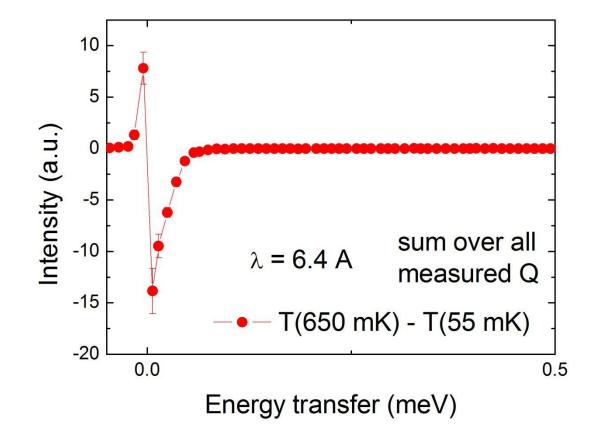


from: Gardner, Hoffer and Phillips, Phys. Rev. A **7**, 1029 (1973).

ABH used this to calculate a vacancy concentration with a T<sup>4</sup> dependence.

Ref: Anderson, Brinkman and Huse, Science **310**, 1164 (2005).

## Temperature dependence - extra



### A new mode?

