# Entanglement Complexity of Systems of Self-Avoiding Walks in Lattice Tubes 

Conference on Knots and other Entanglements in Biopolymers: Topological and Geometrical Aspects of DNA, RNA and Protein Structures

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## Talk Outline

- Review of Lattice Models of Polymers
- Orlandini, Tesi and Whittington proposal for a "good" measure of entanglement complexity for dense polymer systems
- Results on what can be proved for this measure of the entanglement complexity for a system of self-avoiding walks in a tube (work with M. Atapour).


## Introduction to Polymers

Polymer: Large molecule made of repeated molecular units called monomers; if there is more than one type of monomer Copolymer


Fundamental Question of Interest: What properties of polymer solutions are primarily a result of the fact that a polymer is a very large molecule made up of repeated molecular units?

## Modelling Polymers in Solution

## ADVANTAGES of LATTICE MODELS:

excluded volume property is easily incorporated substantial conformational freedom available combinatorial analysis possible qualitative features of phase diagrams expected to be correct values of critical exponents expected to be exact

## MODELS of LINEAR and RING POLYMERS:



SELF-AVOIDING WALK


CIRCLE GRAPH


SELF-AVOIDING POLYGON

## Properties of Self-avoiding Polygons (SAPs) in $\mathbb{Z}^{d}$ <br> Standard Concatenation Argument



$$
\begin{gathered}
\begin{array}{|c}
p_{n} p_{m} \leq(d-1) p_{n+m} ; p_{n} \leq(2 d)^{n}
\end{array} \Rightarrow \lim _{n \rightarrow \infty} \frac{1}{2 n} \log p_{2 n} \equiv \log \mu_{d}=\kappa_{d} \\
\text { (Hammersley Proc.Camb.Phil.Soc. 58(1961), 235-8) } \\
p_{n}(\phi) p_{m}(\phi) \leq 2 p_{n+m}(\phi) ; p_{n}(\phi) \leq p_{n} \Rightarrow \lim _{n \rightarrow \infty} \frac{1}{2 n} \log p_{2 n}(\phi) \equiv \log \mu_{0}=\kappa_{0} \\
\text { (Sumners and Whittington JPA 21 } 1988), 1689-94) \\
\kappa_{o}<\kappa_{3} \Rightarrow \text { Prob. of Knotting }=1-\frac{p_{n}(\phi)}{p_{n}}=1-e^{-\left(\kappa-\kappa_{o}\right) n+o(n)}
\end{gathered}
$$

Key ingredient: Pattern theorem (Kesten, 1963) used to prove that "tight trefoil" pattern occurs at least once in all but exponentially few sufficiently long SAPs.


## Good Measures of Knot Complexity

Results extended in several directions ... Soteros, Sumners and Whittington (1992
MathProcCambPhilSoc 111 75)
Good Measures of Knot Complexity: Function $F: \mathcal{K} \rightarrow[0, \infty)$ s.t.
(i) $F(\phi)=0$
(ii) $\exists K \in \mathcal{K}$ s.t. $F(n K \# L) \geq n F(K)>0 \forall L \in \mathcal{K}$ (i.e. roughly additive w.r.t. knot product)

Then Sumners \& Whittington (1988) $\Rightarrow$
If $F$ is a good measure of knot complexity, let $K$ be a knot which satisfies part (ii) above. Then, there exists $n_{K}, A_{k}>0$ s.t. $\forall n>n_{K}$ all but exponentially few $n$-SAPs have $F$-complexity exceeding $A_{K} n$.

Examples of Good Measures: Crossing number, number of prime factors, genus, bridge number minus one, span of any non-trivial Laurent knot polynomial, log(order), unknotting number, minor index, braid index minus one.

# For Dense Polymer Solutions or Polymer Melts, What is a "Good" Measure of Entanglement Complexity? 



Electron micrograph of a $0.4 \mathrm{mg} / \mathrm{ml}$ actin solution polymerized in vitro. Bar $=1 \mu \mathrm{~m}$. (Gotter et al arXiv:cond-mat/9611097 v1 20 Nov 1995.)

Orlandini, Tesi and Whittington 2000 - characterize entanglement by "linking number" for a tube or cube from solution.

## Orlandini et al (2000) Measure for Entanglement Complexity in Dense Systems

Proposal: Take a random cube (or tube) from $n$-edge system composed of $c$ chains. Assume endpts are in the cube boundary and no edges in the boundary. For each pair $(i, j)$ of chains, join up the two ends of each chain by an arc outside the cube and compute linking number $L k(i, j)=\frac{1}{2} \sum_{m} \sigma_{m}$.
Entanglement Measure: $E C=\sum_{i=1}^{k-1} \sum_{j>i}|L k(i, j)|$


Regular Projections of Two Chains


Density dependence of $\langle E C\rangle$ for cube size 10

## What Can Be Proved About This Entanglement Measure? SSAWs in $\infty \times N \times M$ Tube



Define: A System of Self-avoiding Walks (SSAW) of size $n$, span $s$ and with $c$ components in an ( $N, M$ )-tube is a finite subgraph of the tube with $n$-edges and $c$ connected components s.t.
(i) each connected component is an undirected SAW (USAW) with its endpts in a tube wall (e.g. $y=0$, $y=N, z=0$ or $z=M$ ) and no degree two vertices in any wall of the tube
(ii) for each integer $m \in[0, s]$, there is at least one vertex of the SSAW in the plane $x=m$ and no vertices in $x=m, \forall m \notin[0, s]$.
$\boldsymbol{q}_{\boldsymbol{s}}(\boldsymbol{N}, \boldsymbol{M} ; \boldsymbol{n}, \boldsymbol{c})$ - \# of $\boldsymbol{n}$-edge, span $\boldsymbol{s}, \boldsymbol{c}$-component SSAWs in ( $\boldsymbol{N}, \boldsymbol{M}$ )-tube (up to $\boldsymbol{x}$-translation) Note: $c=\frac{1}{2} \times(\#$ of degree one vertices in SSAW)

Concatenation of SSAWs:

$G_{1}$


$G$
$\mathrm{SSAW} \boldsymbol{G}(\boldsymbol{s}=\mathbf{7}, \boldsymbol{n}=\mathbf{3 1}, \boldsymbol{c}=\mathbf{9})$,
the concatenation of $\boldsymbol{G}_{\mathbf{1}}$ and $\boldsymbol{G}_{\mathbf{2}}$.

Concatenation gives
$\boldsymbol{q}_{s_{1}}\left(N, M ; n_{1}, c_{1}\right) q_{s_{2}}\left(N, M ; n_{2}, c_{2}\right) \leq q_{s_{1}+s_{2}+1}\left(N, M ; n_{1}+n_{2}, c_{1}+c_{2}\right)$
$\Rightarrow$ (via standard arguments, c.f. Janse van Rensburg 2000)
the existence of the limiting free energy:

$$
F(N, M ; x, y)=\lim _{s \rightarrow \infty} \frac{1}{s} \log Z_{s}(N, M ; x, y)
$$

where

$$
Z_{s}(N, M ; x, y)=\sum_{n, c} q_{s}(N, M ; n, c) x^{n} y^{c}
$$

$F(N, M ; x, y)=\lim _{s \rightarrow \infty} \frac{1}{s} \log Z_{s}(N, M ; x, y) ; \quad Z_{s}(N, M ; x, y)=\sum_{n, c} q_{s}(N, M ; n, c) x^{n} y^{c}$.
Plus $\boldsymbol{F}(\boldsymbol{N}, \boldsymbol{M} ; \boldsymbol{x}, \boldsymbol{y})$ is
(i) a convex function of $\log \boldsymbol{x}($ for fixed $\boldsymbol{y})$ and $\operatorname{of} \log \boldsymbol{y}($ for fixed $\boldsymbol{x})$;
(ii) its right and left derivatives in $\boldsymbol{x}$ (for fixed $\boldsymbol{y}$ ) and in $\boldsymbol{y}$ (for fixed $\boldsymbol{x}$ ) exist everywhere in $(\mathbf{0}, \boldsymbol{\infty})$;
(iii) it is differentiable almost everywhere, and when the derivative exists the order of the limit and derivative can be interchanged.
$\Rightarrow$
For r.v. $\boldsymbol{X}$ with state space span $s$ SSAWs and $\mathbb{P}(X=G)=\frac{x^{n(G)} y^{c(G)}}{Z_{s}(N, M ; \boldsymbol{x}, \boldsymbol{y})}$
$\lim _{s \rightarrow \infty} \frac{\mathbb{E}[\boldsymbol{n}(\boldsymbol{X})]}{\boldsymbol{s}}$ exists a.e. and is non-decreasing in $\boldsymbol{x}$, i.e.
For fixed $(\boldsymbol{x}, \boldsymbol{y})$, the avg. edge-density $\frac{\mathbb{E}[\boldsymbol{n}(\boldsymbol{X})]}{\boldsymbol{s N M}}$ goes to $\rho_{\boldsymbol{e}}(\boldsymbol{x}, \boldsymbol{y})$ as $\boldsymbol{s} \rightarrow \infty$, and $\uparrow$ as $\boldsymbol{x} \uparrow$.
$\lim _{s \rightarrow \infty} \frac{\mathbb{E}[c(\boldsymbol{X})]}{s}$ exists a.e. and is non-decreasing in $\boldsymbol{y}$, i.e.
For fixed $(\boldsymbol{x}, \boldsymbol{y})$, the avg. walk-density goes to $\rho_{\boldsymbol{w}}(\boldsymbol{x}, \boldsymbol{y})$ as $\boldsymbol{s} \rightarrow \infty$, and $\uparrow$ as $\boldsymbol{y} \uparrow$.

## Entanglement Complexity of SSAWs in $\infty \times N \times M$ Tube



A two-component link $K_{12}$ is associated to a pair of USAWs $\left(\bar{w}_{1}, \bar{w}_{2}\right)$.
The following result has been proved (c.f. Atapour, PhD thesis):
Corollary 1 Given an n-edge SSAW $G$, let $w_{1}, \ldots w_{c}$ be the sequence of USAWs in $G$. Let $\bar{G}$ be the SSAW associated to $G$ as prescribed in Lemma 6.2.1. For any $1 \leq i<j \leq k$, there exists a two-component polygonal link $K_{i j}=\left(K_{i}, K_{j}\right)$ corresponding to the pair of USAWs $\left(w_{i}, w_{j}\right)$ such that when it is projected into the xy-plane those edges of $K_{i j}$ which lie outside $B_{G}$ will create no crossing with the edges inside $B_{G}$ and at most one crossing, involving both polygons $K_{i}$ and $K_{j}$, amongst themselves.

The Entanglement Complexity, $E C(G)$, of any SSAW $G$ is now defined as follows:

$$
\begin{equation*}
E C(G)=\sum_{i=1}^{c-1} \sum_{j=i+1}^{c}\left|L k\left(K_{i j}\right)\right| \tag{1}
\end{equation*}
$$

where $L k\left(K_{i j}\right)$ is the linking number of $K_{i j}$.

Upper Bound on EC of SSAWs: Since one can bound the number of crossings associated with an edge of the SSAW by a constant (dependent only on $N$ and $M$ ), the $E C$-complexity can be bounded above by a linear function of $n$ or $s$ :
(i) Given a pair of SAWs, we can extend their endpts (by adding at most a( $N, M$ ) edges) so that the new endpoints are "away" from the tube.
(ii) For any $\epsilon>0$, there exists a "push-off" of the walk pair which has a regular projection in the $(x, y)$-plane such that the $(x, y)$ coordinates of walk vertices are within $\epsilon$ of their values in the regular projection. The endpoints will be on the boundary of the projection.
$\Rightarrow E C=\sum_{i=1}^{c-1} \sum_{j>i}|L k(i, j)| \leq \sum_{i=1}^{c} \sum_{m=1}^{n_{i}} C(i, m)+\sum_{i=1}^{c} e x(i)$
$n_{i}:=$ number of edges in the $i$ th SAW

$C(i, m):=$ the number of edges crossed by the $m$ th edge of the $i$ th SAW $\leq$ the number of edges $E$ in a $3 \times(N+2) \times M$ slice of the lattice
$e x(i):=$ the number of walk pairs involving the $i$ th walk which could result in an extra crossing $\leq$ the number of walks (other than the $i$ th) which have vertices in a slice of the tube determined by the location of the endpts of the $i$ th walk $\leq$ number of vertices in that tube slice $\leq(N+2)(M+1) n_{i}$

$$
\Rightarrow E C=\sum_{i=1}^{c-1} \sum_{j>i}|L k(i, j)| \leq d(N, M) n \leq d d(N, M) s
$$

## A Lower Bound for EC of SSAWs in Tubes

$q_{s}(N, M ; n, c ;<m, P)$ - \# of $n$-edge, span $s, c$-component SSAWs in ( $N, M$ )-tube (up to $x$-translation) which contain less than $m$ translates of pattern $P$
A Pattern Theorem would allow us to compare:

$$
\begin{aligned}
& F(N, M ; x, y ; \epsilon, P)=\lim _{s \rightarrow \infty} \frac{1}{s} \log Z_{s}(N, M ; x, y ; \epsilon, P) \\
& \quad Z_{s}(N, M ; x, y ; \epsilon, P)=\sum_{n, c} q_{s}(N, M ; n, c ;<\epsilon s, P) x^{n} y^{c}
\end{aligned}
$$

to

$$
\begin{gathered}
F(N, M ; x, y)=\lim _{s \rightarrow \infty} \frac{1}{s} \log Z_{s}(N, M ; x, y) \\
Z_{s}(N, M ; x, y)=\sum_{n, c} q_{s}(N, M ; n, c) x^{n} y^{c}
\end{gathered}
$$

If $F(N, M ; x, y ; \epsilon, P)<F(N, M ; x, y)$ then
$\Rightarrow$ All but exponentially few sufficiently wide span

$s$ "weighted" SSAWs contain at least $\epsilon s$ translates of $P$.
Since $P$ cannot occur
more than $\epsilon s$ times in an SSAW with $E C \leq \epsilon s \Rightarrow$
(i) All but exponentially
few sufficiently wide SSAWs have entanglement measure $E C>\epsilon s$, i.e. $E C$ grows at least linearly in the span.
(ii) The probability that a span
$s$ SSAW has $E C$ greater than $\epsilon s$ goes to 1 as $s \rightarrow \infty$
HENCE A PATTERN THEOREM IS NEEDED - TRANSFER MATRIX WILL BE USED

Transfer Matrix for SSAWs: Given $k \geq 2$, an SSAW is a sequence of overlapping $k$-configs.

$\mathcal{S}_{k}$ : set of all possible SSAW $k$-configs
$\mathcal{S}_{k, o}$ : subset of $\mathcal{S}_{k}$ that are start $k$-configs
$\mathcal{S}_{k, f}$ : subset of $\mathcal{S}_{k}$ that are end $k$-configs
$\mathcal{G}_{k}=(V, A):$
graph with vertex set $V=\mathcal{S}_{k}=\left\{P_{1}, P_{2}, \ldots\right\}$
and $\operatorname{arc}\left(P_{i}, P_{j}\right) \in A$ iff $P_{i} P_{j}$ is in $\mathcal{S}_{k+1}$.
Any walk on $\mathcal{G}_{k}$ starting
in $\mathcal{S}_{k, o}$ and ending in $\mathcal{S}_{k, f}$ corresponds to an SSAW.


Transfer Matrix for SSAWS: Given $k \geq 2$ and $\mathcal{G}_{k}$
$\mathcal{G}_{k}=(V, A):$
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Any walk on $\mathcal{G}_{k}$ starting
in $\mathcal{S}_{k, o}$ and ending in $\mathcal{S}_{k, f}$ corresponds to an SSAW.
Simplest Transfer Matrix: $T=\left(T_{i, j}\right)$
$T_{i, j}=\left\{\begin{array}{lr}1, & \left(P_{i}, P_{j}\right) \in A \\ 0, & \text { otherwise }\end{array}\right.$

$\Rightarrow$
$\left(T^{r}\right)_{i, j}=\#$ of $(k+r)$-configs starting with $P_{i} \&$ ending with $P_{j}$.
If $P_{i} \in \mathcal{S}_{k, o}$ and $P_{j} \in \mathcal{S}_{k, f}$, then they are SSAWs.

## More General Transfer Matrix for SSAWS:

Recall: $X$ a r.v. with state space all span $s$ SSAWs.
$\mathbb{P}(X=G)=\frac{x^{n(G)} y^{c(G)}}{Z_{s}(N, M ; x, y)}$
$n(G)$ - \# of edges of $G ; c(G)$ - \# of walks in $G$ (i.e. \# of connected components)
Grand Canonial Partition function:

$$
Q(x, y, z)=\sum_{s, n, c} q_{s}(N, M ; n, c) x^{n} y^{c} z^{s}=\sum_{s} Z_{s}(N, M ; x, y) z^{s} .
$$

$-\log R_{Q}(x, y)=$
$F(N, M ; x, y)=\lim _{s \rightarrow \infty} s^{-1} \log Z_{s}(N, M ; x, y)$
Given $k \geq 2$ and $\mathcal{G}_{k}$
General Transfer Matrix: $T=\left(T_{i, j}\right)$
$T_{i, j}=\left\{\begin{array}{lr}z^{s\left(P_{i}, P_{j}\right)} x^{n\left(P_{i}, P_{j}\right)} y^{c\left(P_{i}, P_{j}\right)}, & \left(P_{i}, P_{j}\right) \in A \\ 0, & \text { otherwise }\end{array}\right.$
$\Rightarrow$

$Q(x, y, z)=\sum_{r=0}^{\infty} \sum_{\left\{i: P_{i} \in \mathcal{S}_{k, o}\right\}} \sum_{\left\{j: P_{j} \in \mathcal{S}_{k, f}\right\}}\left(T^{r}\right)_{i, j}=\frac{f(x, y, z)}{\operatorname{det}(I-T)}$.
where $f(x, y, z)$ is analytic.

## Consequences of the Transfer Matrix Representation

$q_{s}(N, M ; n, c ; \bar{P})$ - \# of $n$-edge, span $s, c$-component SSAWs in ( $N, M$ )-tube (up to $x$-translation) which DO NOT contain pattern $P$
$Z_{s}(N, M ; x, y ; \bar{P})=\sum_{n, c} q_{s}(N, M ; n, c ; \bar{P}) x^{n} y^{c}$.
The transfer matrix and Perron-Frobenius Theory $\Rightarrow$ $F(N, M ; x, y ; \bar{P})=\lim _{s \rightarrow \infty} \frac{1}{s} \log Z_{s}(N, M ; x, y ; \bar{P})$
is strictly less than
$F(N, M ; x, y)=\lim _{s \rightarrow \infty} \frac{1}{s} \log Z_{s}(N, M ; x, y)$


Moreover, for any non-negative integer valued additive functional $\psi$ defined for SSAWs, there exists $\gamma_{\psi, x, y}>0$ (can be determined from the eigenvalues and eigenvectors of $T$ ) such that as $s \rightarrow \infty$

$$
\mathbb{E}[\psi(X)]=\left(\gamma_{\psi, x, y}\right) s+O(1)
$$

## HENCE:

For the edge count: $\quad \mathbb{E}[n(X)]=\left(\gamma_{e, x, y}\right) s+O(1)$
For the walk count: $\quad \mathbb{E}[c(X)]=\left(\gamma_{w, x, y}\right) s+O(1)$
For $n_{P}(X)$ - \# of translates of $P$ occurring in $X: \quad \mathbb{E}\left[n_{P}(X)\right]=\left(\gamma_{P, x, y}\right) s+O(1)$
For $n_{\times}(X)$ - \# of crossings in a regular projection of $X: \quad \mathbb{E}\left[n_{\times}(X)\right]=\left(\gamma_{\times, x, y}\right) s+O(1)$

For the Entanglement Complexity Measure EC:

$$
\left(\gamma_{P, x, y}\right) s+O(1) \leq \mathbb{E}[E C(X)] \leq\left(\gamma_{\times, x, y}\right) s+d^{\prime}(N, M) s+O(1)
$$

EC IS A GOOD MEASURE OF ENTANGLEMENT COMPLEXITY

## Future Work

For the Entanglement Complexity Measure $E C$ :

$$
\left(\gamma_{P, x, y}\right) s+O(1) \leq \mathbb{E}[E C(X)] \leq\left(\gamma_{\times, x, y}\right) s+d^{\prime}(N, M) s+O(1)
$$

How do the $\gamma$ 's change with $(x, y)$ ? i.e. How does avg. $E C(X)$ change as the edge-density and the walk-density change?

For fixed limiting edge-density:

$$
f(N, M ; \alpha) \equiv \lim _{s \rightarrow \infty} s^{-1} \log q_{s}(N, M ;\lfloor\alpha s\rfloor)
$$

where $q_{s}(N, M ;\lfloor\alpha s\rfloor)$ is the \# of span $s,\lfloor\alpha s\rfloor$-edge SSAWs (up-to $x$-translation). What do these results mean for $E C$ ?

What happens as $N$ and $M$ change?
What about other measures of Entanglement Complexity?

## Stretching a Polygon in an ( $N, M$ )-Tube

$p_{n}(N, M ; s)$ : \# of $\boldsymbol{n}$-edge, span $s$ SAPs in ( $N, M$ )-tube
$\hat{Z}_{n}(N, M ; x)=\sum_{s} p_{n}(N, M ; s) e^{f s}$
where $x=e^{f}$
Standard concatenation argument for SAPs in a tube $\Rightarrow$
$\hat{F}(N, M ; x)=\lim _{n \rightarrow \infty} n^{-1} \log \hat{Z}_{n}(N, M ; x)$ exists.

For r.v. $\boldsymbol{X}$ with state space $\boldsymbol{n}$-edge SAPs,
$\mathbb{P}(X=\omega)=\frac{x^{s}}{\hat{Z}_{n}(N, M ; x)}=\frac{e^{f s}}{\hat{Z}_{n}\left(N, M ; e^{f}\right)}$
Transfer matrix arguments $\Rightarrow$
$\mathbb{E}[s(X)]=\frac{1}{\rho_{e}(f)} n+O(1)$
and is non-decreasing in $\boldsymbol{f}$.
 fixed force $f>0$; note that the tube is rotated 90 degrees counterclockwise.

## Knot Complexity of Stretched Polygons in a Tube

$\boldsymbol{p}_{\boldsymbol{n}}(\boldsymbol{N}, \boldsymbol{M} ; \boldsymbol{s} ; \overline{\boldsymbol{P}})$ : \# of $\boldsymbol{n}$-edge, span $\boldsymbol{s}$ SAPs in $(\boldsymbol{N}, \boldsymbol{M})$-tube which do NOT contain pattern $\boldsymbol{P}$ $\hat{Z}_{n}(N, M ; x ; \bar{P})=\sum_{s} p_{n}(N, M ; s ; \bar{P}) e^{f s}$, where $x=e^{f}$

A concatenation argument for SAPs in a tube $\Rightarrow$
$\hat{F}(N, M ; x ; \bar{P})=\lim _{n \rightarrow \infty} n^{-1} \log \hat{Z}_{n}(N, M ; x ; \bar{P})$ exists.

Transfer-matrix arguments $\Rightarrow$
$\hat{\boldsymbol{F}}(\boldsymbol{N}, M ; \boldsymbol{x} ; \overline{\boldsymbol{P}})<\hat{\boldsymbol{F}}(\boldsymbol{N}, M ; \boldsymbol{x})$
$\mathbb{E}\left[n_{P}(X)\right]=\left(\gamma_{P, f}\right) n+O(1)$
For any $\boldsymbol{f}$, all sufficiently long
$\boldsymbol{n}$-edge SAPs under the influence of the force $\boldsymbol{f}$ contain at least $\gamma_{P, f} \boldsymbol{n}$ trefoils in their knot decomposition.


Part of a tight
trefoil pattern for
the tube.

