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Signaling the Arrival of the LHC Era

8 - 13 December 2008

Searches for Higgs Bosons at the LHC

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Searches for Higgs Bosons at the LHC

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• Searching for Higgs Bosons at ATLAS & CMS - I

- Strategy to new physics
- Before collisions

• Test beams

- Commissioning with cosmic rays
- Alignment strategies and pre-alignments
- Full simulations of realistic "as-built" and "as-installed" detector
- Triggers for initial luminosities

– In the early going ...

- Luminosity measurements
- Sanity checks $J/\Psi(Y) \rightarrow \mu\mu$ and $Z \rightarrow \mu\mu/ee \dots$
- Minimum bias and underlying events
- Constraining PDF with early data
- Inclusive jet cross section
- Commissioning with top events
- Jet energy scale and b-tagging efficiency
- First discoveries

• Searching for Higgs Bosons at ATLAS & CMS - II

- Standard Model Higgs
- Higgs in Extended Models
- Higgs property measurements

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~ already introduced by previous speakers

Outline

ATLAS and CMS ...





| Parameter | ATLAS | CMS |
|--|-------------------|------------------|
| Total weight (tons) | 7,000 | 12,500 |
| Overall diameter (m) | 22 | 15 |
| Overall length (m) | 46 | 20 |
| Magnetic field for tracking (T) | 2 | 4 |
| Solid angle for lepton ID or tracking $(\Delta \phi \times \Delta \eta)$ | $2\pi \times 5.0$ | $2\pi 	imes 5.0$ |
| Solid angle for energy measurements $(\Delta \phi \times \Delta \eta)$ | $2\pi 	imes 9.6$ | $2\pi 	imes 9.6$ |
| Total cost (MCHF) | 550 | 550 |

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Strategy to new physics discovery ...

- Before data taking:
 - Quality controls of detector construction to meet physics requirements
 - Test beam: several years of activities culminating in the combined test beam of 2004/2007 to understand and calibrate sub-detectors, and to validate/tune software tools, e.g., Geant4 simulation
 - Full simulations of realistic "as-built" and "as-installed" detector (misalignments, material non-uniformity, dead channels): test and validate calibration and alignment strategies
 - Some aspects of commissioning with cosmics are being addressed now or have already been addressed:
 Dre clignment and collibration
 - Pre-alignment and calibration
 - Initial detector shake-down
 - Data processing at the Tier 0 (CERN), distributed to Tier 1's and some Tier 2's. Analysis at the Tier 2's.
- With first data
 - Commission and calibrate detector and trigger in situ with minimum bias, Z→ll, etc
 - Rediscover SM at $\sqrt{s} = 10$ TeV (minimum bias, W, Z, tt, QCD jets, etc)
 - Validate and tune tools (MC generators)
 - Measure main backgrounds to new physics: W+jets, Z+jets, tt+jets,

OCD multijets, ... Kétévi A. Assamagan

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almost here

Luminosity - why measure it?

Measure the cross sections for "Standard " processes

Jet production

.....

Top pair production Theoretically known to ~ 10 %

- New physics manifesting in deviation of $\sigma \times BR$ relative to the Standard Model predictions. Precision measurement becomes more important if new physics not directly seen. (characteristic scale too high!)
 - Important precision measurements
 - = Higgs production $\sigma \times BR$
 - tanβ measurement for MSSM Higgs
 - **.....**

Significant uncertainty from luminosity unless measuring ratio of observables where luminosity cancels out in the ratio



Relative precision on the measurement of σ_{H} ×BR for various channels, as function of m_{H_2} at $|Ldt = 300 \text{ fb}^{-1}$. The dominant uncertainty is from Luminosity: 10% (open symbols), 5% (solid symbols).

(ATLAS Physics TDR, May 1999)

Luminosity measurements ...

CMS/TOTEM and ATLAS forward detectors for forward physics, heavy ion, ... and luminosity measurements



Higgs Searches: present limits



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Higgs Production at the LHC



SM Higgs Decays at the LHC





Many channels explored! All the mass range is covered!



Light Higgs Boson ...



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In the early going ...



Expected data samples with only 100 pb⁻¹ ~ 38.5 days at 10³² cm⁻² s⁻¹ with 30% eff.

| Channels (<u>examples</u>) | Events to tape for 100 pb ⁻¹ (ATLAS) | Total statistics from LEP and Tevatron |
|--|--|--|
| $W \rightarrow \mu v$ $Z \rightarrow \mu \mu$ $tt \rightarrow W b W b \rightarrow \mu v + X$ QCD jets $p_T > 1 \text{ TeV}$ $\tilde{g}\tilde{g}$ m = 1 TeV | ~ 10 ⁶ ~ 10 ⁵ ~ 10 ⁴ > 10 ³ ~ 50 | ~ 10 ⁴ LEP, ~ 10 ⁶⁻⁷ Tevatron ~ 10 ⁶ LEP, ~ 10 ⁵⁻⁶ Tevatron ~ 10 ³⁻⁴ Tevatron |
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$J/\Psi(Y) \rightarrow \mu\mu$ and $Z \rightarrow \mu\mu$...



EM Calorimeter uniformity...

- Use Z→ee to correct for residual long range nonuniformity due to
 - Module-to-module variation
 - Temperature effects



Minimum bias/Underlying Event ...

Minimum bias events

- Inelastic hadron-hadron events selected with an experiment's "minimum bias trigger".
- Usually associated with inelastic nonsingle-diffractive events (NSD) (e.g. UA5, E735, CDF ... ATLAS/CMS?)



- Need minimum bias data to:
 - Study general characteristics of protonproton interactions
 - Investigate multi-parton interactions and the structure of the proton etc.
 - Understand the underlying event: impact on physics analyses?

The underlying event (UE)

The "soft part" associated with hard scatters



- In parton-parton scattering, the UE is usually defined to be everything except the two outgoing hard scattered jets:
 - Beam-beam remnants.
 - D Additional parton-parton interactions.
 - ISR + FSR
- Can we use "minimum bias" data to model the "underlying event"?
 - At least for the beam-beam remnant and multiple interactions?





Addition theoretical errors from PDF uncertainties ~15% at 1 TeV

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Constraining PDF with early data ...

Using $W \rightarrow I_V$ angular distribution





Constraining PDF with early data ...

- What will be the effect of including LHC data in global PDF Fits?
 - How much can we reduce the error?
- Take 10⁶ W \rightarrow ev events generated with CTEQ6.1 PDF and ATLFAST detector simulation to probe the low-x gluon PDF at Q²=m_W²: W⁺ \rightarrow e⁺v rapidity spectrum sensitive to gluon shape parameter λ
 - The statistics corresponds to 150 pb⁻¹
 - Introduce 4% systematic errors by hand
- This "data" included in the global fit



The central value of the ZEUS-PDF prediction shifts and uncertainty is reduced; Error on low-x gluon shape parameter $\lambda [xg(x) \sim x^{-\lambda}]$ reduced by 35%

The systematics on electron acceptance versus η , will be controlled to a few percent using Z→ee (~3 10⁴ events at 100 pb⁻¹)

Jet Energy Scale / b-tagging Efficiency...

- Calibration of light jet energy scale, to complement γ+jets sample
 - Remove $|m_{jj}-m_W| < 10 \text{GeV/c}^2$ (bias requirement) and look at m_{jj} for all the 3jet combinations in m_t mass window

Jet E-scale: ■ initially to ~ 10% from test beam + simulation (Geant4 reproduces test-beam pion response of hadronic end-cap calorimeter to ~2%) ■ 'Then eventually from data (γ/Z+jets, W→jj in tt events) + simulation →1%

- Calibration of b-tagging efficiency
 - Select 3jets from the hadronic top
 - Perform a fit using resolution, m_t , and m_W as constraints
 - Measure the b-tag efficiency for the 4th jets a function of E_T and eta



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top reconstruction & mass measurement ...

• Several methods for early top mass measurement - for example, kinematic fit to reconstruct hadronic top in lepton+jet sample. Implications for Higgs:

− ttH (H→bb); H⁺ → tb; t → bH⁺



Excess number of events in tails...



Deviation from the SM could be seen quickly: but need time to demonstrate that it is SUSY



With only 100 pb⁻¹, deviation from the SM up to ~ 1 TeV. But understanding backgrounds would require 1 fb⁻¹: Implications for Higgs

- Higgs production from SUSY cascade decays
- Higgs decays to SUSY particles (and invisibles)

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Narrow Resonance \rightarrow 11 at \sim 1 TeV...



• With 100 pb⁻¹, signal large enough for discovery up to m>1 TeV

• Signal is a narrow resonance on top of a small Drell-Yan background

Ultimate calorimeter performance not needed

Discovery likely in the e+e- channel but μ + μ - channel needed for couplings, asymmetry, etc

Di-jet Resonances ...

• Physics interest is in the high mass tail:

- Sensitivity to excited quarks, W', Z', etc.
- Limits from CDF and D0 are in the range < 0.78 TeV
- With few pb⁻¹ at 14 TeV we can extend the range
- Crucial experimental parameter is the energy resolution in measuring jet energy (They are narrow resonances)



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Summary

- With first data, up to $100 \text{ pb}^{-1}(2009)$
 - Detector and Trigger commissioning and calibration in situ
 - Simulation/reconstruction software tuning
 - "Re-discover" the SM: W, Z, top, jets
 - Help constrain PDF uncertainties
 - Could discovery some new physics
 - \sim TeV scale resonance X \rightarrow ll
 - Narrow resonance in di-jet mass tail
 - Hint for SUSY
 - Hint for a light neutral scalar
- With more data ($\geq 1 \text{ fb}^{-1}$: 2010)
 - Discover a TeV scale SUSY
 - Discover at least one Higgs boson
 - Understand deviation, excess as signal for new physics

→ Searches for Higgs Bosons at the LHC - ||

Outline

- Introduction
- Standard Model Higgs
- MSSM Higgs
- Higgs Property measurements

Some of the questions ...

- What is the origin of mass?
 - How is the Electroweak symmetry broken?
 - Do Higgs bosons exist?
- What is the underlying theory of unification?
 - Are the interactions unified at some larger energy scale?
 - How is gravity incorporated?
 - Is there supersymmetry?
- Flavor
 - Why are there 3 flavor families?
 - Neutrino masses and mixing?
 - The origin of CP violation

Answers to some of these questions are expected at the TeV mass scale, i.e. at the LHC

Improvement in Higgs Studies at the LHC

- •Many studies have meanwhile been performed using detailed GEANT simulations of the detectors
 - Physics Performance Technical Design Report from the CMS collaboration
 - ATLAS CSC (Computing System Challenge) notes in preparation, to be released towards the end of 2008
- New (N)NLO Monte Carlos (also for backgrounds)
 - MCFM Monte Carlo, J. Campbell and K. Ellis, http://mcfm.fnal.gov
 - MC@NLO Monte Carlo, S.Frixione and B. Webber, www.eb.phy.cam.ar.uk/theory/
 - T. Figy, C. Oleari and D. Zeppenfeld, Phys. Rev. D68, 073005 (2003)
 - E.L.Berger and J. Campbell, Phys. Rev. D70, 073011 (2004)
 - C. Anastasiou, K. Melnikov and F. Petriello, hep-ph/0409088 and hep-ph/0501130

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Physics Performances Physics Technical Design Report Vol II

CMS: CERN / LHCC 2006-021

- New approaches to match parton showers and matrix elements
 - ALPGEN Monte Carlo + MLM matching, M. Mangano et al.
 - SHERPA Monte Carlo, F. Krauss et al.

÷...

Tevatron data are extremely valuable for validation, work has started

- More detailed, better understood reconstruction methods (partially based on test beam results,...)
- Further studies of new Higgs boson scenarios

(Various MSSM benchmark scenarios, CP-violating scenarios, Invisible Higgs boson decays,....)Kétévi A. AssamaganICTP Trieste, 09-10.12.0829

ATLAS: CERN-OPEN 2008-020 (to appear soon)

Standard Model Higgs

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Higgs → gamma gamma



1000

Higgs → gamma gamma backgrounds

| Within a mass window M_H +/- 1.4 σ GeV: | | |
|--|--------------------|--|
| Background Process | Cross-section (fb) | |
| γγ | 562 | |
| Reducible γj | 318 | |
| Reducible jj | 49 | |
| Drell Yan | 18 | |

- Background is evaluated with *NLO* simulations.
- It will be measured from data sidebands.



Strategy for jet rejection:

- Longitudinal segmentation of the calorimeter.
- Fine segmentation of the first layer (η -strips) => good π^0 rejection.
- Isolation of the electromagnetic cluster.
- Isolation based on tracks reconstructed by the inner detector.



Higgs \rightarrow gamma gamma reconstruction

PRIMARY VERTEX If the vertex is unknown, add 1.4 GeV to the mass resolution. Combine calorimeter and tracker informations!



- Calorimeter \rightarrow vertex position accuracy of *19 mm* - Combining with the tracker information $\rightarrow \sim 0.1 \text{ mm}$

Calorimeter information is useful in case of pile-up or events with low tracks multiplicity.

0.14

CONVERSIONS

- ~50% of the events with at least one converted γ !
- ad hoc energy calibration required in late conversions;
- conversion vertex used in computation of the direction;

- used for gamma-jet background estimation.





Higgs \rightarrow gamma gamma reconstruction

9000 Signal 8000 Irreducible bkg 7000 6000 Reducible bkg 5000 4000 3000 2000 1000 ATLAS preliminary 115 120 125 130 135 140 145 110 145 150 M_{ov} (GeV) H+1 jet 600 Signal 🗏 Irreducible bkg 500 Reducible bkg 400 300 200 100 ATLAS preliminary THE PROPERTY OF THE PROPERTY O 115 120 125 110 130 135 140 145 150 M_{rr} (GeV) 25 H+2jet (VBF) Signal Irreducible bkg 20 Reducible bkg

> 145 150

M₇₇ (GeV)

140

135

Events/GeV

Events/GeV

Events/GeV

ATLAS preliminary

115 120 125 130

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fio

From associated production, high levels of purity for signal achieved thanks to additional requirements (but less events....)



Higgs \rightarrow gamma gamma significance Events/GeV 200 400 Higgs M,=140 GeV (x10) E Higgs M,=130 GeV (x10) New elements of the analyses: CMS Higgs M_L=120 GeV (x10) 🛅 Higgs M, =115 GeV (x10) ee Drell Yan iets pmar > 50 GeV y+jets (1 prompt y + 1 fake) - NLO calculations available y+jets (2 prompt y) 300 🔛 γγ box (Binoth et al., DIPHOX, RESBOS) 🖂 yy born 200 - Realistic detector material 100 - More realistic K factors (for signal and background) - Split signal sample acc. to resolution functions 80 160 170 180 100 110 120 130 140 150 M_{yy} (GeV) Signal Significance Discovery Significance (No) Int L = 30 fb⁻¹ Combined fit based with M ed, fit based with M_floate ATLAS 10 usive, fit based with M., fixed 6 CMS $I = 10 \, \text{fb}$ 5 ЗE ← Cut-Based Analysis (with syst. err.) 2 - Cut-Based Analysis (no syst. err.) 1Ē Optimized Analysis (with syst. err.) Preliminary Optimized Analysis (no syst. err.) 120 135 125 130 140 110 115 120 125 130 135 140 145 150 155 Higgs boson mass [GeV] M_u (GeV) Comparable results for ATLAS and CMS • Improvements possible by using more exclusive $\gamma\gamma$ + jet topologies

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events.

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Selection efficiency as a function of the Higgs mass, for each of the three decay channels.



Mass resolution as a function of the Higgs mass. Open circles denote the resolution Z mass constraint improves resolution.

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Higgs $\rightarrow ZZ^* \rightarrow 41$

- Main backgrounds: ZZ (irreducible), tt, Zbb (reducible)
- Main experimental tools for background suppression:
 - lepton isolation in the tracker and in the calorimeter
 - impact parameter

Updated ATLAS and CMS studies:

- ZZ background: NLO K factor used
- background from side bands





ZZ background estimation

What is the expected number of ZZ background events under the signal region?

■ <u>direct simulation ZZ→41 process :</u>

%large uncertainties theoretical knowledge, MC modelling detector acceptance
%statistical precision "only" limited by Monte Carlo statistics

<u>control samples :</u>

& full cancellation pp luminosity uncertainty & reduced sensitivity to theoretical and experimental uncertainties

normalization to single Z data

$$R = \frac{(\sigma_{ZZ \to 41} \epsilon_{41} \int Ldt)}{(\sigma_{Z \to 21} \epsilon_{21} \int Ldt)}$$

 $\frac{1}{2}$ theoretical uncertainty : 2–8%

 $\frac{1}{2}$ statistical precision : high thanks to large statistics of $Z \rightarrow l+l$ - data

normalization to sidebands

$$N_{estimated \ bkg}^{signal \ region} = \alpha_{MC} N_{observed \ bkg}^{sidebands}$$
$$\alpha_{MC} = \frac{N_{expected \ bkg}^{signal \ region}}{N_{expected \ bkg}^{sidebands}}$$

\$<u>theoretical uncertainty:</u> 0.5-4%

statistical precision : available statistics in sidebands limiting factor?





Vector Boson Fusion qq H

Motivation: Increase discovery potential at low mass Improve and extend measurement of Higgs boson parameters (couplings to bosons, fermions)

> Established (low mass region) by D. Zeppenfeld et al. (1997/98) Earlier studies: R.Kleiss W.J.Stirling, Phys. Lett. 200 (1988) 193; Dokshitzer, Khoze, Troyan, Sov.J. Nucl. Phys. 46 (1987) 712; Dokshitzer, Khoze, Sjöstrand, Phys.Lett., B274 (1992) 116.





Rapidity distribution of jets in tt and



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4

η

2



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$H \rightarrow WW \rightarrow \ell \nu \ell \nu$



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Vector Boson Fusion $qq H(\rightarrow \tau\tau)$

- High BR in the low mass region.
- 3 channels: ll, lh, hh (65% of τ gives hadrons)

SELECTION

• **Trigger**: isolated electrons (μ) with $p_T > 22 (20) GeV/c$ ($\varepsilon - 10\%$)

 $\tau + E_T^{miss}$ ($\epsilon \sim 3.7\%$) for the *hh* channel

Isolation cut

- Likelihood exploiting the shower shape and the track quality to separate τ and jet.
- b-jet veto to kill tt(+jets)→lvb lvb (+jets) (background for the ll channel)
- select highest E_T jets in opposite hemispheres
- Central jet veto

BACKGROUNDS

- $Z \rightarrow \tau \tau + jets$
- $W \rightarrow \tau v + jets$
- tt+jets
- QCD multi-jets for the hh channel



MAIN ISSUES:

- Discrepancies in Monte Carlo generator → impact on veto efficiency
- Pileup \rightarrow impact on E_T^{miss} and jet veto
- Estimation of QCD multi-jet \rightarrow no sensitivity yet on *hh* channel

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<u>Vector Boson Fusion qq H($\rightarrow \tau \tau$)</u>

 $\begin{array}{rcl} qq \ H & \rightarrow \ qq \ \tau \ \tau \\ & \rightarrow \ qq \ \ell \nu \nu \ \ell \nu \nu \\ & \rightarrow \ qq \ \ell \nu \nu \ h \nu \end{array}$

ATLAS

VBF H(120)→ττ→lh

 $\sqrt{s} = 14 \text{ TeV}, 30 \text{ fb}^{1}$

≥14

Events / (5

6

Experimental challenges:

- In-time pileup, out-of-time pileup, underlying event.
- test simulations & use vertexing for the jet
- calorimeter timing
- early data underlying event measurement
- Identification of hadronic τ
- Good E_T^{miss} resolution (since there are neutrinos...)
- \bullet Knowledge of the $Z \to \tau \tau$ background shape in the

p_{T.miss}

V e/u

high mass region: use data $Z \rightarrow \mu\mu$ to emulate it!

Higgs mass reconstructed using the angle

between the two τ and the **collinear**

approximation:

 $m_{\tau\tau} = m_{ll} / sqrt(X_1 X_2)$ with $X_i = P_T(l_i) / P_T(\tau_i)$







ttH → ttbb

Complex final states: $H \rightarrow bb, t \rightarrow bjj, t \rightarrow bk$

 $t \rightarrow b\ell v, t \rightarrow b\ell v$

 $t \rightarrow bjj, t \rightarrow bjj$

Main backgrounds:

- combinatorial background from signal (4b in final state)
- ttjj, ttbb, ttZ,...
- Wjjjjjj, WWbbjj, etc. (excellent b-tag performance required)



 Updated CMS study (2006): ALPGEN matrix element calculations for backgrounds → larger backgrounds (ttjj dominant), experimental + theoretical uncertainties, e.g. ttbb, exp. norm. difficult.....





- Full mass range can already be covered after a few years at low luminosity
- Similar performance in ATLAS
- Several channels available over a large range of masses Vector boson fusion channels play an important role at low mass !

Important changes w.r.t. previous studies:

- $H \rightarrow \gamma \gamma$ sensitivity of ATLAS and CMS comparable
- ttH \rightarrow tt bb disappeared in both ATLAS and CMS studies

The Extended Higgs Sector

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MSSM Higgs

- One doublet of Higgs pseudo-scalar fields is replaced with two:
 - One couples to up-fermions and has $vev=v_u$
 - One couples to down fermions and has vev=v_d
- 2X4-3=5 physical scalar fields/particles: *h*, *H*, *A*, *H*[±]
- •Properties at tree level:
- fully defined by 2 free parameters: m_A , $tan\beta = v_u/v_d$
- CP-odd A:
- never couples to Z or W;
- decays to bb, $\tau\tau$ (and additionally tt for small $tan\beta$).
- CP-even h and H:

SM-like near their mass limits vs m_A ; at large $tan\beta$ enhanced coupling with down fermions, suppressed couplings to W and Z.

- H^{\pm} "strongly" couples to *tb* and τv
- All Higgs bosons are narrow (Γ<10GeV)

We choose the benchmark scenario m_h^{max} corresponding to maximal theoretically allowed region for m_h .



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Discovery potential for SUSY Higgs bosons



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$A/H \rightarrow \mu\mu$



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$H \rightarrow \tau \tau (\rightarrow 11)$

SELECTION.

- •Trigger: isolated $\mu(e)$ with $p_T > 20$ (25)GeV || two isolated $e \parallel$ or one e & one μ
- b-tagging on at least one jet to suppress light jets
- Cuts on missing E_T , b momentum, lepton momentum, number
- of jets (<3) to reject Z and tt backgrounds
- Collinear approximation



• Studies ongoing on hadronic τ decay mode





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Charged Higgs









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Invisible Higgs decays ?

Possible searches: tt H $\rightarrow \ell v b qqb + P_T^{miss}$ $Z H \rightarrow \ell \ell + P_T^{miss}$ $qq H \rightarrow qq + P_T^{miss}$ - J.F. Gunion, Phys. Rev. Lett. 72 (1994)

- D. Choudhury and D.P. Roy, Phys. Lett. B322 (1994)

- O. Eboli and D. Zeppenfeld, Phys. Lett. B495 (2000)



All three channels have been studied:

key signature: excess of events above SM backgrounds with large P_T^{miss} (> 100 GeV/c)



Problems / ongoing work:

- ttH and ZH channels have low rates
- More difficult trigger situation for qqH
- backgrounds need to be precisely known (partially normalization using ref. channels possible)
- non SM scenarios are being studied at present first example: SUSY scenario

Invisible Higgs decays ?



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Higgs decays via SUSY particles

If SUSY exists : search for $H/A \rightarrow \chi^0_2 \chi^0_2 \rightarrow \ell \ell \chi^0_1 \ell \ell \chi^0_1$



CMS: special choice in MSSM (no scan)

 $M_1 = 60 \text{ GeV/c}^2$ $M_2 = 110 \text{ GeV/c}^2$ $\mu = -500 \text{ GeV/c}^2$

 $H/A \rightarrow \chi^0_2 \chi^0_2 \rightarrow \ell \ell \chi^0_1 \ \ell \ell \chi^0_1 \qquad gb \rightarrow tH^+, \ H^{\pm} \rightarrow \chi_{2,3}^0 \ \chi_{1,2}^{\pm} \rightarrow 3\ell + E_T^{miss}$



ATLAS: special choice in MSSM (no scan) $M_1 = 60 \text{ GeV/c}^2$ $M_2 = 210 \text{ GeV/c}^2$ $\mu = 135 \text{ GeV/c}^2$ $m(s - \ell_R) = 110 \text{ GeV/c}^2$ $m(s - \tau_R) = 210 \text{ GeV/c}^2$

- Exclusions depend on MSSM parameters (slepton masses, m)
- More systematic studies are needed (initiated by A. Djouadi et al., also started in ATLAS)

Updated MSSM scan for different benchmark scenarios

Benchmark scenarios as defined by M.Carena et al. (h mainly affected)



ATLAS preliminary, 30 fb^{-1,} 5o discovery

MHMAX scenario $(M_{SUSY} = 1 \text{ TeV/c}^2)$ maximal theoretically allowed region for m_h

Nomixing scenario $(M_{SUSY} = 2 \text{ TeV/c}^2)$ (1TeV almost excl. by LEP) small $m_h \rightarrow difficult$ for LHC

Gluophobic scenario ($M_{SUSY} = 350 \text{ GeV/c}^2$) coupling to gluons suppressed (cancellation of top + stop loops) small rate for g g \rightarrow H, H \rightarrow $\gamma\gamma$ and Z \rightarrow 4 ℓ

Small α **scenario** (M_{SUSY} = 800 GeV/c²) coupling to b (and t) suppressed (cancellation of sbottom, gluino loops) for large tan β and M_A 100 to 500 GeV/c²

Higgs search at the LHC in CP-violating scenarios

- Tree level CP of Higgs Potential in MSSM may be violated sizably at higher orders by loop effects involving CP-violating interactions of Higgs to stop and sbottom



- CP eigenstates h, A, H mix to mass eigenstates H₁, H₂, H₃



- Effect maximized in a defined benchmark scenario (CPX)
 (M. Carena et al., Phys.Lett. B 495 155 (2000))
 arg(A_t) = arg(A_b) = arg(M_{gluino}) = 90°
- No lower mass limit for H₁
 from LEP !
 (decoupling from the Z)

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details depend on m<sub>top</sub> and on
theory model
(FeynHiggs vs. CPsuperH)
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MSSM discovery potential for the CPX scenario



- Large fraction of the parameter range can be covered, however, small hole at (intermediate tan β , low m_{H+}) corresponding to low m_{H+}
- More studies needed, e.g. investigate lower H₁ masses,

additional decay channels:

 $tt \to Wb \ H^+b \to \ell_V b \ WH_1 b, \ H_1 \to bb$

Higgs Property Measurements

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Higgs Property Measurements

- If a Higgs-like particle is discovered, its properties must be measured to establish its nature:
 - Mass
 - Couplings to fermions and bosons
 - Spin and CP
 - Higgs self coupling

Higgs mass measurements

• Higgs boson mass can be measured with a precision of 0.1% over a large mass range $(130 - \sim 450 \text{ GeV/c}^2)$ ($\gamma\gamma$ and ZZ $\rightarrow 4\ell$ resonances, el.magn. calo. scale uncertainty assumed to be $\pm 0.1\%$)



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Higgs couplings measurements

- Different quantities can be measured depending on assumptions:
 - CP-even, Spin 0 implicit in H -> WW event selection. Measure only the rate: σ^*BR
 - Assuming only one Higgs boson, no degenerate Higgs pair, extract relative branching ratios
 - Assuming Standard Model couplings, no extra particles, or strong couplings to light fermions: extract ratios of Higgs couplings
 - Assuming that the sum of the visible branching ratios is the same as in the SM: extract absolute couplings

Higgs couplings measurements

Global likelihood-fit (at each possible Higgs boson mass) Input: measured rates, separated for the various production modes Output: Higgs boson couplings, normalized to the WW-coupling.

Channels used in coupling fit:

| Production | | Decay | Mass range |
|------------------------|----------------------|---|-----------------------------------|
| | GF: | $H \to ZZ^{(*)} \to 4l$ | 110 GeV - 200 GeV |
| | Gluon Fusion | $H 	o WW^{(*)} 	o l u l u$ | 110 GeV - 200 GeV |
| g | $(gg \rightarrow H)$ | $H ightarrow \gamma \gamma$ | $110~{\rm GeV}$ - $150~{\rm GeV}$ |
| | WBF: | $H \to ZZ^{(*)} \to 4l$ | 110 GeV - 200 GeV |
| W, Z, H | Weak Boson | $H \to WW^{(*)} \to l\nu l\nu$ | 110 GeV - 190 GeV |
| W, Z | Fusion | $H \to \tau \tau \to l \nu \nu l \nu \nu$ | $110~{\rm GeV}$ - $150~{\rm GeV}$ |
| q | (qq H) | $H \to \tau \tau \to l \nu \nu \mathrm{had} \nu$ | $110~{\rm GeV}$ - $150~{\rm GeV}$ |
| | | $H \rightarrow \gamma \gamma$ | $110~{\rm GeV}$ - $150~{\rm GeV}$ |
| eeee t | $t\bar{t}H$ | $H \to WW^{(*)} \to l\nu l\nu (l\nu)$ | $120~{\rm GeV}$ - $200~{\rm GeV}$ |
| $t \rightarrow -H_{-}$ | | $H ightarrow b \overline{b}$ | 110 GeV - 140 GeV |
| occoc t | | $H \to \tau \tau$ (not included) | $110~{\rm GeV}$ - $150~{\rm GeV}$ |
| | | $H \rightarrow \gamma \gamma$ | 110 GeV - 120 GeV |
| | WH | $H \to WW^{(*)} \to l\nu l\nu (l\nu)$ | 150 GeV - 190 GeV |
| | | $H ightarrow \gamma \gamma$ | 110 GeV - 120 GeV |
| Ň | ZH | $H ightarrow \gamma \gamma$ | 110 GeV - 120 GeV |
| Kétévi A. Assamagan | | ICTP Trieste, 09-10.12.08 | 65 |

Rate Measurements



- ► Left: for M_H<150 GeV, rates can be measured with an accuracy typically between 30-100% with 30 fb⁻¹ of integrated luminosity
 - This assumes a 5% uncertainty on the luminosity measurement.
- ▶ Right: for channels that can be seen in the mass range $110 < M_H < 190$ GeV, error is smaller

Ratios of partial widths Assume only one Higgs الالجاح

Total width cannot be measured due to detector resolution.

►H→WW used as normalization since it is measured with the smallest error

Right: with 30 fb⁻¹ of integrated luminosity, ratios of partial widths can be measured with an accuracy better than 60%



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Ratios of Couplings

Assume no unobservable decay modes, etc.

Cross-sections and branching ratios are proportional to the squares of couplings

Proportionality constants, along with systematic errors, are computed from theory

► With 30 fb⁻¹ of integrated luminosity, Ratios of Higgs couplings can be measured with an accuracy ranging from ~20% to ~100%



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Absolute couplings

Additional assumption: sum of visible branching ratios (H \rightarrow WW, ZZ, plus $\gamma\gamma$, $\tau\tau$, and bb, depending on mass) is the same as in Standard Model, with an error corresponding to the sum of undetected branching ratios

Couplings to Z, W, and τ can be measured with a precision better than 100%



Spin/CP measurements

The SM Higgs is CP even scalar (spin 0). Spin information in its decay products. The Z bosons from Higgs decay are polarized. Polarizations depend on Higgs mass. For m_H >300 GeV, the Z's mostly longitudinally polarized

In the center-of-mass frame of the Z, longitudinal and transverse polarized Z bosons lead to different angular distributions of the decaying leptons:



Spin/CP measurements

The decay planes of the 2 Z's are expected to be correlated, due largely to the transversely polarized Z's. For m_H >300 GeV, correlation would "disappears" as the Z's becomes longitudinally polarized



Aspects of CP-violation in HZZ coupling, R. Godbole, et al, hep-ph/07080458

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Spin/CP measurements

<u>Results:</u>

[*C.P. Buszello et al., Eur. Phys. J. C 32, 209 (2004)*]

• For Higgs masses larger than about 230 GeV/c², a spin 1 hypothesis can be ruled out with 100 fb⁻¹.

 A spin-CP hypothesis of 0⁻ can be ruled out with less than 100 fb-1.

Analysis of H and Z may exclude Higgs-like particles with unusual spin/CP properties. Not for early data However.



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Spin/CP measurements



 Angle between struck quarks in VBF events are sensitive to tensor structure of HVV vertex
 Typically: P_T(j_{1,2})>20-40 GeV, Δη_{ij}>3-4.5; M_{ij}>500-1000 GeV
 H→ττ for M_H~120 GeV, H→WW for M_H~160 GeV





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0.02

0.01

0.5

1.5

Eur. Phys. J. C51:385-414

2

 $\Delta \phi_{ii}$

2.5

Excluding non-SM Higgs



In some regions of the MSSM parameter space, only one light Higgs is visible (left)
Try to exclude MSSM using a 2 analysis of coupling fits (right)

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Higgs Self coupling

To establish the Higgs mechanism the Higgs boson self-coupling has to be measured:

Cross sections for HH production:



 ϕ_1 (ϕ_1, ϕ_2 .) ϕ_2 ϕ_1 Kreis der Minima



small signal cross sections, large backgrounds from tt, WW, WZ, WWW, tttt, Wtt,...

⇒ no significant measurement possible at the LHC need Super LHC $L = 10^{35} \text{ cm}^{-2} \text{ sec}^{-1}$, 6000 fb⁻¹

Most sensitive channel: $gg \rightarrow HH \rightarrow WW WW \rightarrow \ell_{V} jj \ell_{V} jj$

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BACKUP

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Measurement of Higgs-Boson Coupling Ratios

assumptions: only SM particles couple to Higgs boson,

no large couplings of light fermions

Global fit(ATLAS study)(all channels at a given mass point)

Production cross sections

$$\sigma_{ggH} = \alpha_{ggH} \bullet g_t^2$$

$$\sigma_{VBF} = \alpha_{WF} \bullet g_w^2 + \alpha_{ZF} \bullet g_Z^2$$

$$\sigma_{ttH} = \alpha_{ttH} \bullet g_t^2$$

$$\sigma_{WH} = \alpha_{WH} \bullet g_W^2$$

$$\sigma_{ZH} = \alpha_{ZH} \bullet g_Z^2$$

b loop neglected for now in ggH

Fit parameters:

$$\frac{g_Z^2}{g_W^2} \ \frac{g_\tau^2}{g_W^2} \ \frac{g_b^2}{g_W^2} \ \frac{g_t^2}{g_W^2} \ \frac{g_t^2}{g_W^2} \ \frac{g_w^2}{\sqrt{\Gamma_H}}$$

 $\alpha\,$ from theory with assumed uncertainty $\,\Delta\alpha\,$

$$\Delta \alpha_{ggH} = 20\%$$

$$\Delta \alpha_{WF} = \alpha_{ZF} = 4\%$$

$$\Delta \alpha_{ttH} = 15\%$$

$$\Delta \alpha_{WH} = \Delta \alpha_{ZH} = 7\%$$

Ratio of Higgs-Boson Couplings

Branching ratios

$$BR(H \rightarrow WW) = \beta_{W} \frac{g_{W}^{2}}{\Gamma_{H}}$$

$$BR(H \rightarrow ZZ) = \beta_{Z} \frac{g_{Z}^{2}}{\Gamma_{H}}$$

$$BR(H \rightarrow \gamma\gamma) = \frac{\left(\beta_{\gamma(W)}g_{W} - \beta_{\gamma(t)}g_{t}\right)^{2}}{\Gamma_{H}} \Delta\beta = 1\%$$

$$BR(H \rightarrow \tau\tau) = \beta_{\tau} \frac{g_{\tau}^{2}}{\Gamma_{H}}$$

$$BR(H \rightarrow bb) = \beta_{b} \frac{g_{b}^{2}}{\Gamma_{H}}$$
Rate as function of x_i, e.g.

$$(\sigma \bullet BR)_{ggH,H\to ZZ} =$$

$$\alpha_{ggH} \frac{g_t^2}{g_W^2} \frac{g_W^2}{\sqrt{\Gamma_H}} \beta_Z \frac{g_Z^2}{g_W^2} \frac{g_W^2}{\sqrt{\Gamma_H}}$$

