



*The Abdus Salam
International Centre for Theoretical Physics*



2035-1

Conference on Superconductor-Insulator Transitions

18 - 23 May 2009

Comparison of superconductor-insulator transitions in different materials

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142432 Chernogolovka, Russia

Superconductor – insulator transitions

in

different materials

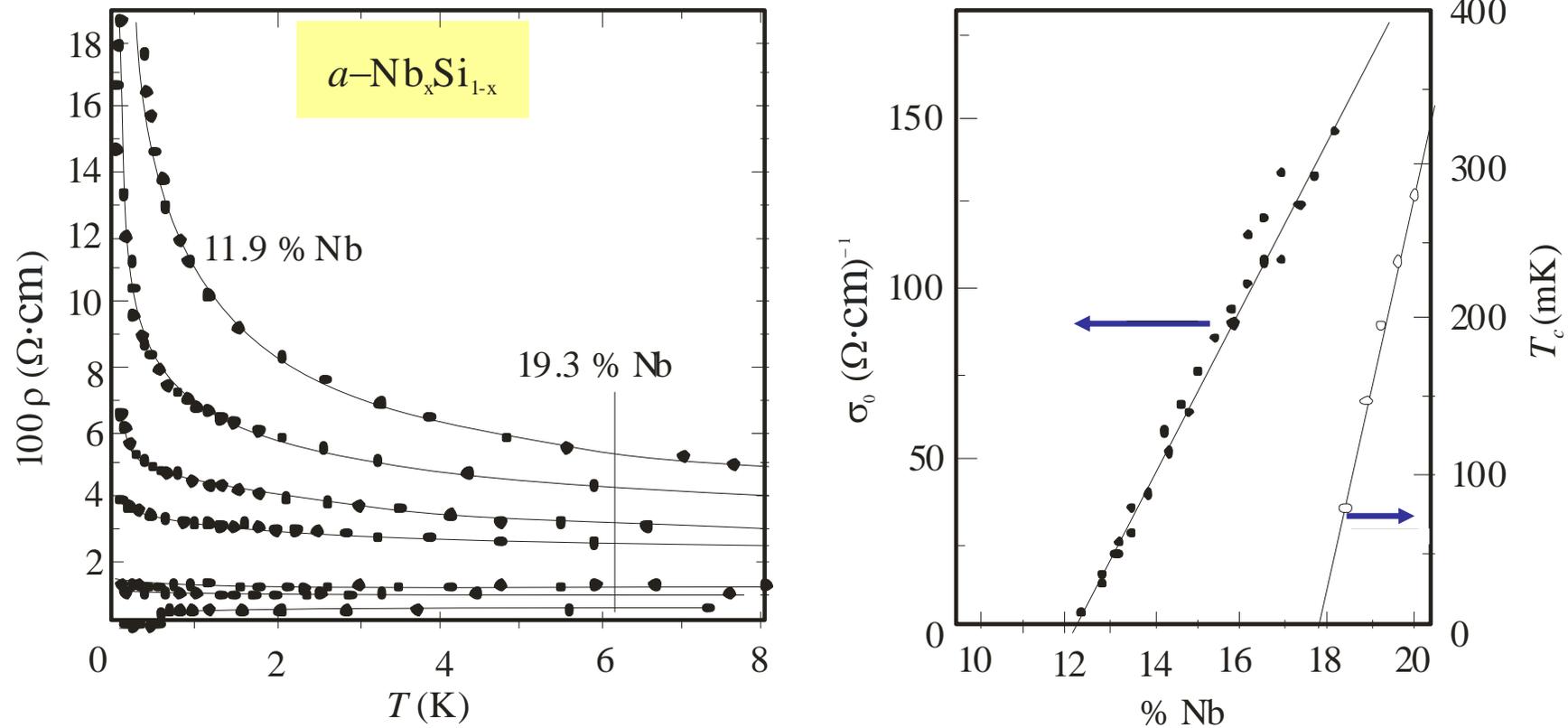
Trieste

May 2009

Bureaucratic classification

- 1. Ultrathin films* **Bi, Be, Pb, Al, Ga, Ta**
- 2. Materials with variable content* **InO, TiN, NbSi,
MoGe, MoSi**
- 3. High- T_c superconductors* **BiSrCaCuO, YBaCuO,
LaSrCuO, NdCeCuO**

*Split transition :
superconductor – normal metal – insulator*



D.J. Bishop, E.G. Spencer, and R.C. Dynes,
Solid St. Electron. **28**, 735 (1985)

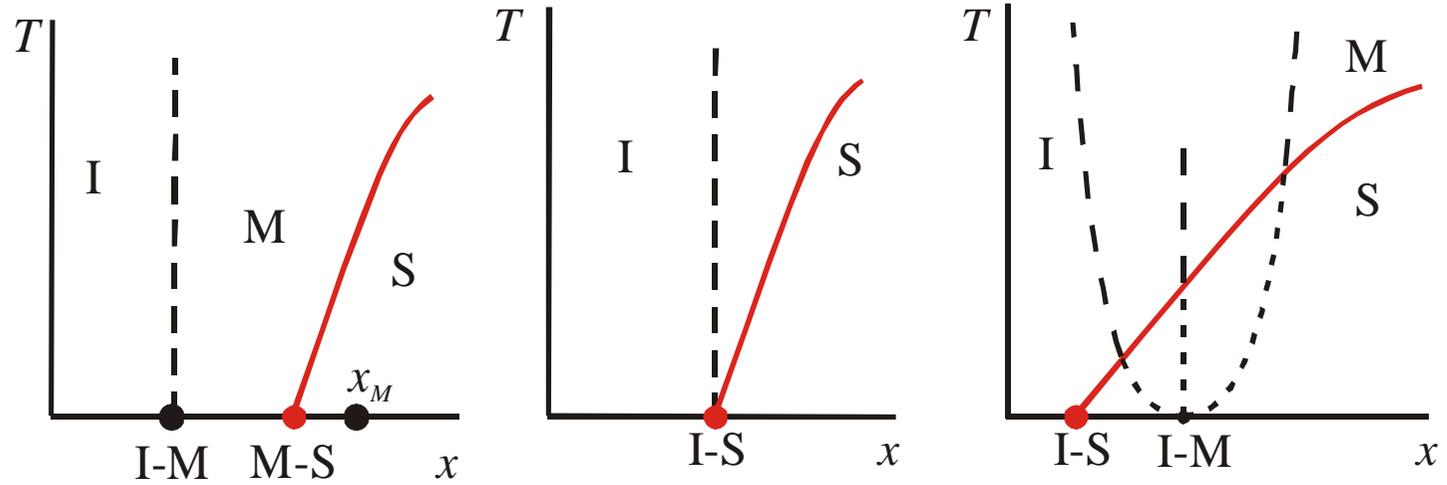
*Let quantum phase transition
from superconducting ground state to nonsuperconducting
at $T = 0$ in the disordered media
be tuned by some control parameter*

*It should be compared with
Anderson normal-metal – insulator transition*

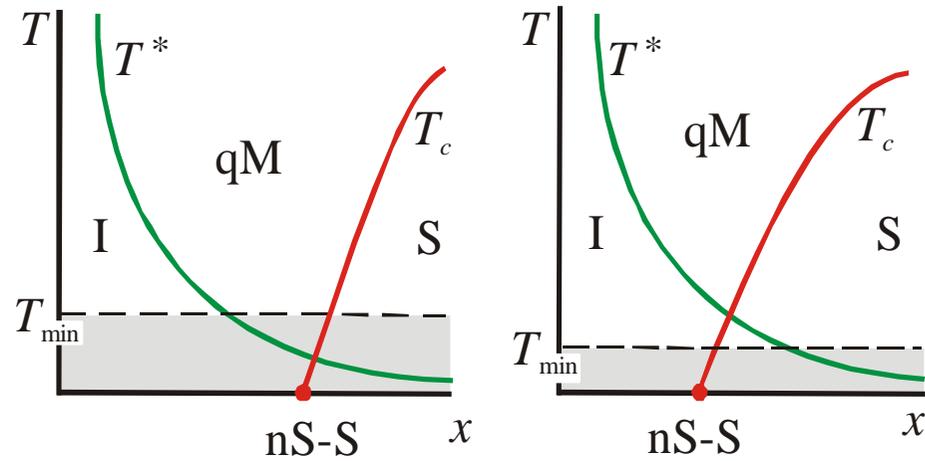
What happens first ?

Variants of the (x, T) – phase diagram

3D



2D



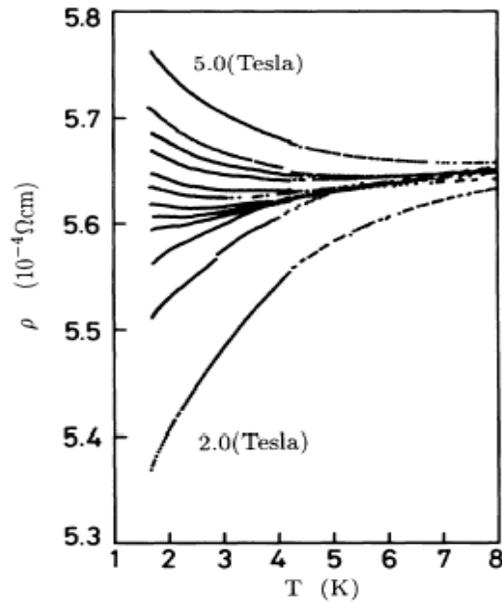
Bureaucratic classification

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semiBureaucratic classification

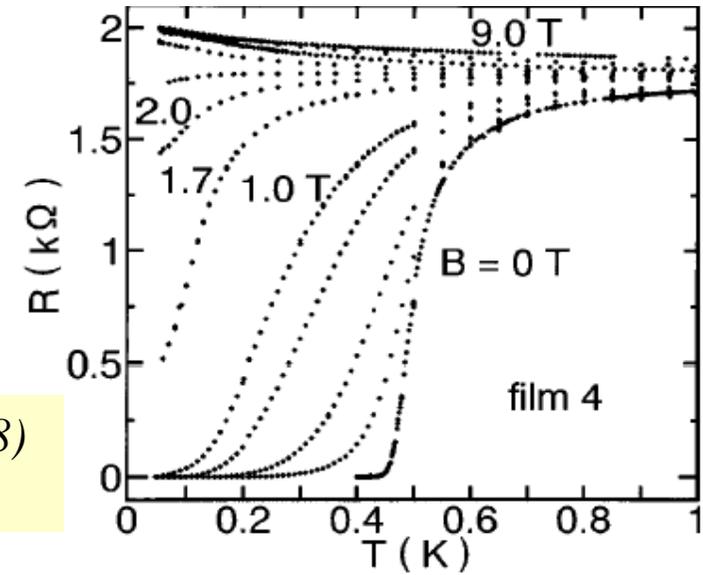
- 1. Ultrathin films* **Bi, Be, Pb, Al, Ga, Ta**
- 2. Materials with variable content* **InO, TiN, NbSi, MoGe, MoSi**
- 3. High- T_c superconductors* **BiSrCaCuO, YBaCuO, LaSrCuO, NdCeCuO**
- 4. Crossover from superconductor – normal-metal to superconductor – insulator transition*
Ta, NbSi, MoGe, MoSi, NdCeCuO

Transitions into “bad” metals



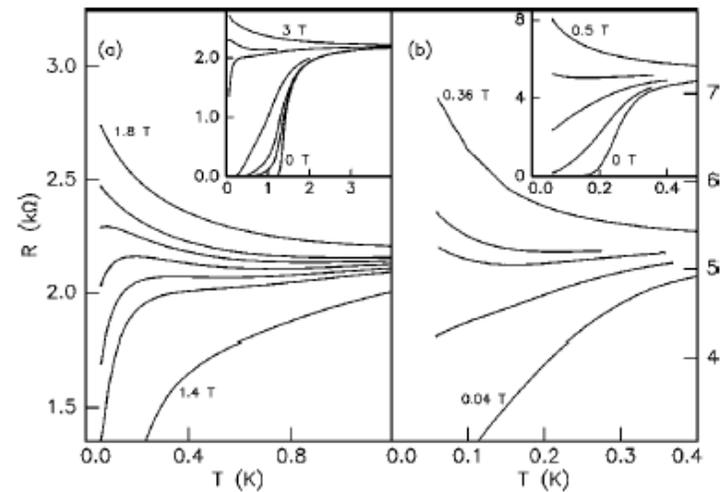
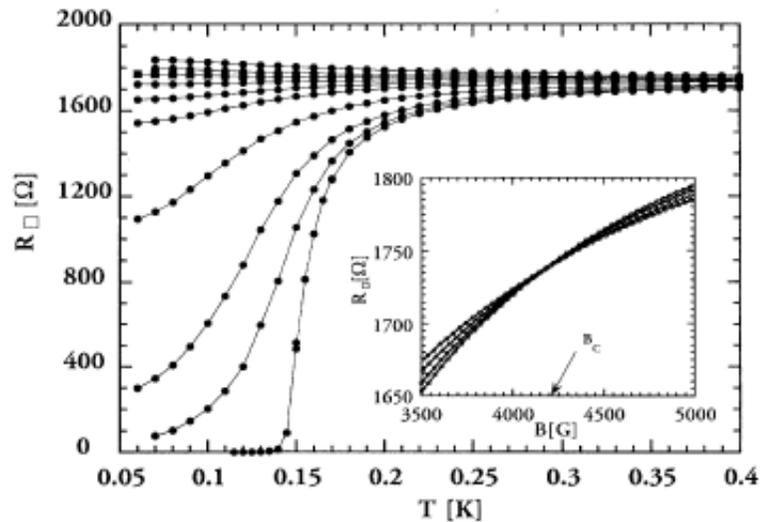
Tanda et al. (1992)
 $\text{Nd}_{2-x}\text{Ce}_x\text{CuO}_4$

S. Okuma et al. (1998)
 $\text{Mo}_x\text{Si}_{1-x}$

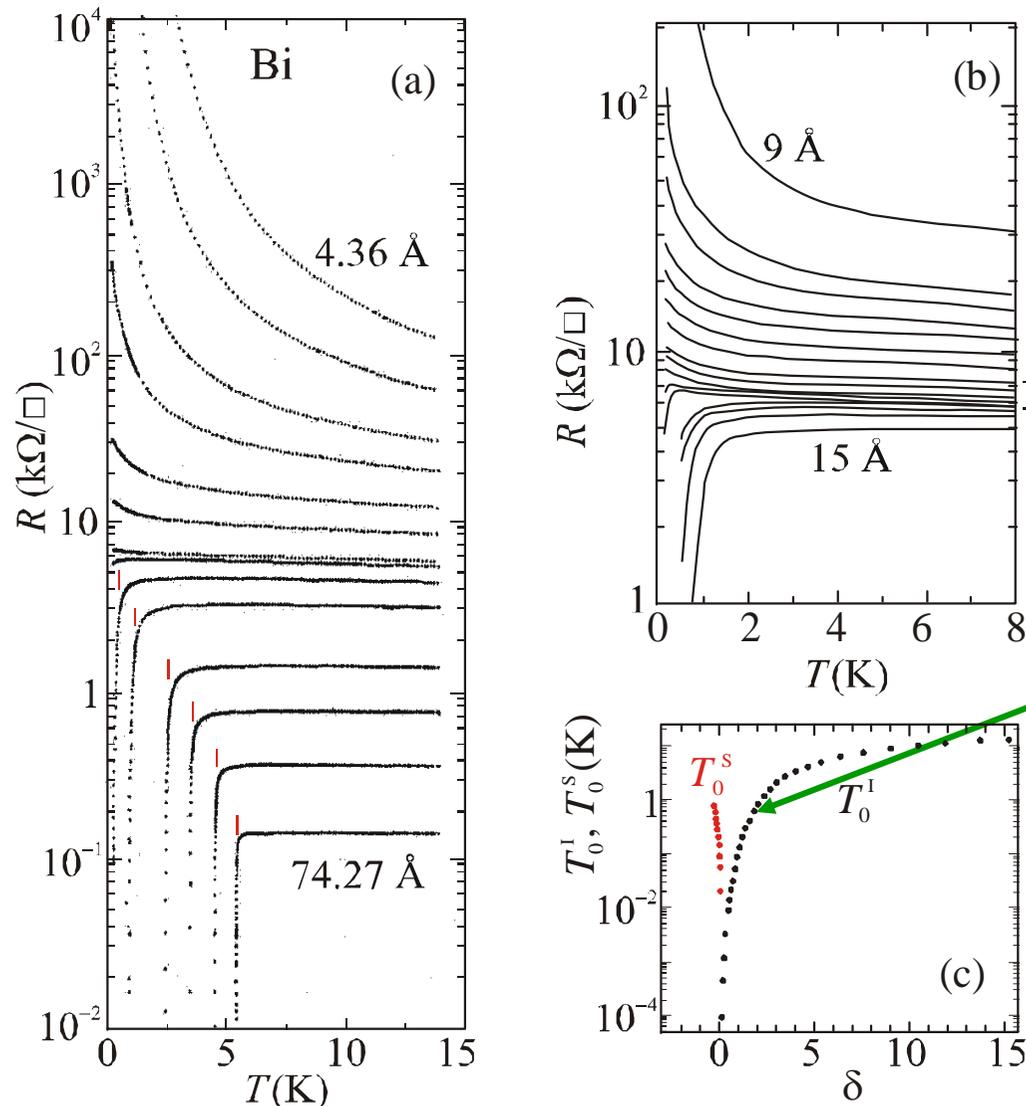


A. Yazdani and A. Kapitulnik (1995)
 $\alpha\text{-MoGe}$

J. A. Chervenak and J. M. Valles, Jr (2000)
 Bi/Sb



Ultra thin Bi films – transition induced by thickness changes



Experimental curves are analysed by the help of one of these three expressions

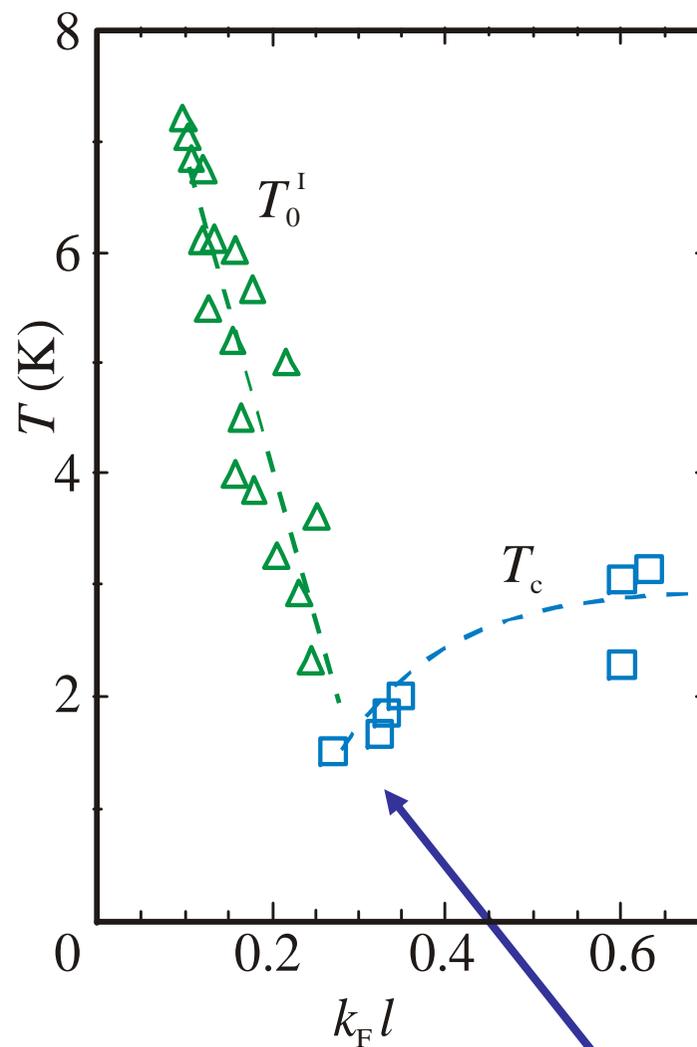
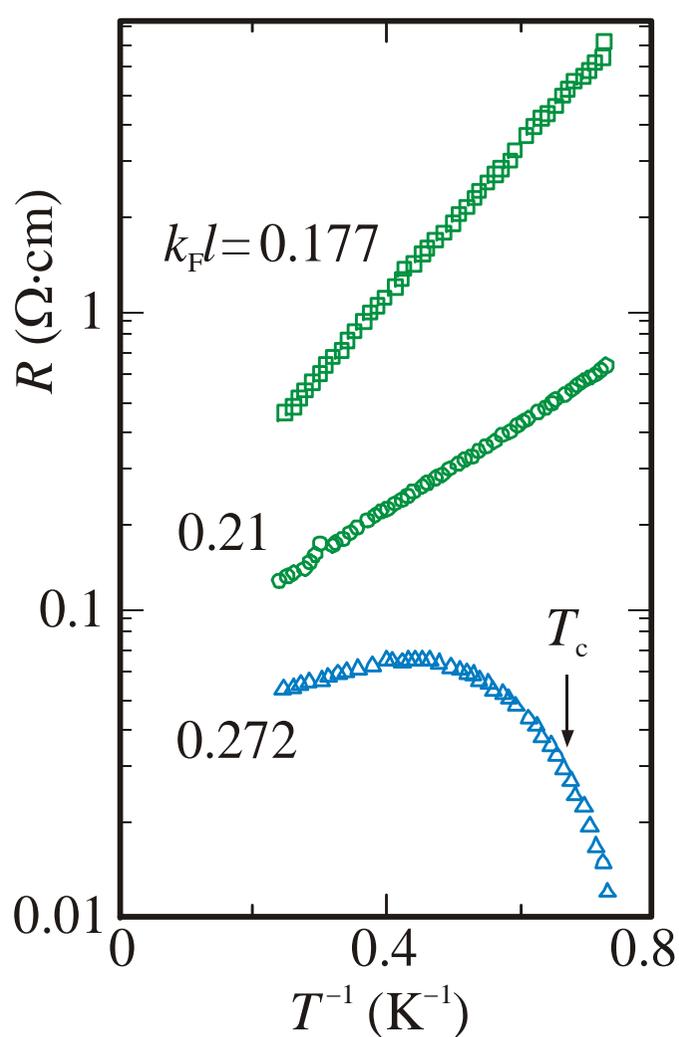
$$g = g_0 \exp(-T_0^I / T)$$

$$g = g_0 \exp[-(T_0^I / T)^{1/2}]$$

$$g = g_0 - \ln(T / T_0^I)$$

D.B. Haviland, Y. Liu, A.M. Goldman et al.,
 Phys. Rev. Lett. **62**, 2180 (1989); **81**, 5217 (1998);
 Phys. Rev. B **47**, 5931 (1993)

a – InO_x films



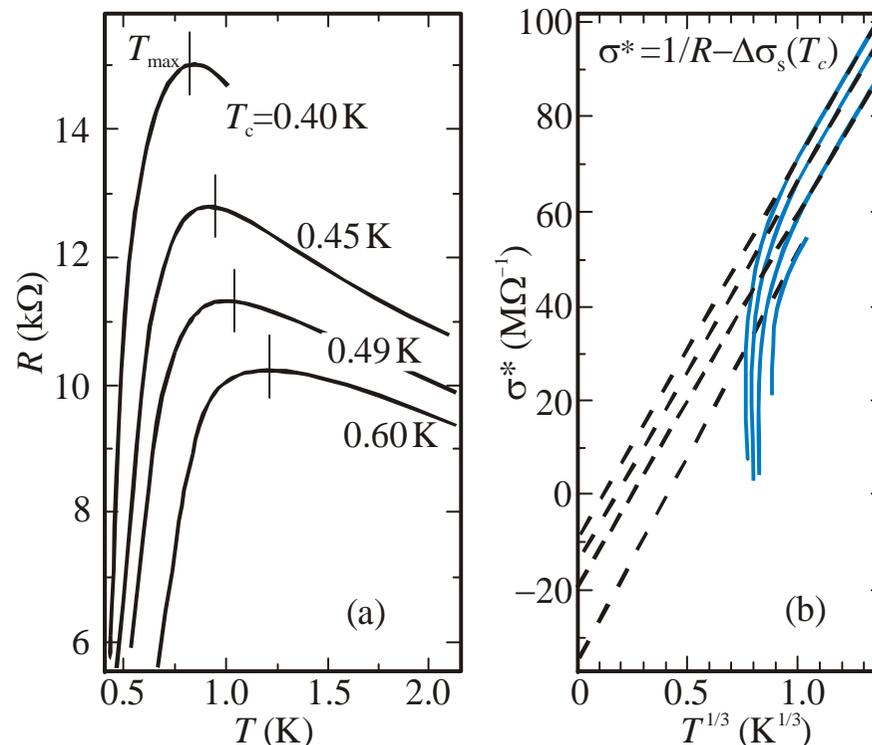
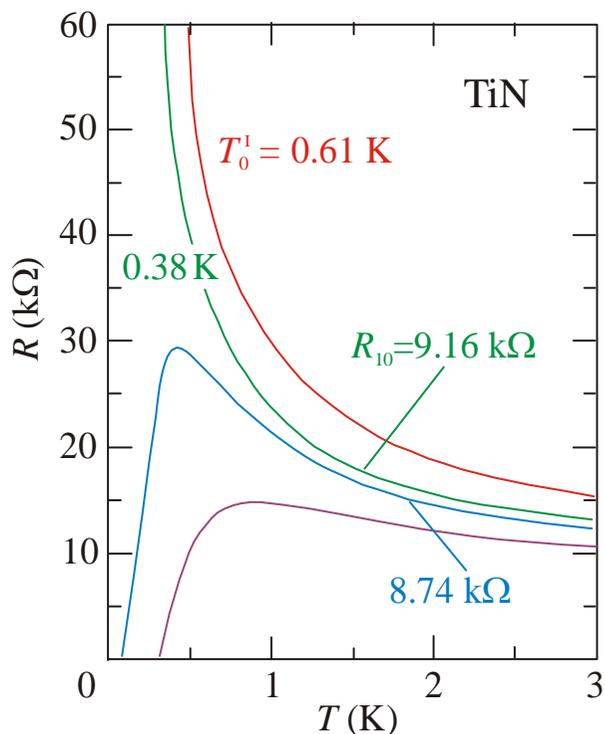
D. Shahar and Z. Ovadyahu,
Phys. Rev. B **46**, 10917 (1992)

Transition is not split

TiN films, thickness 5 nm

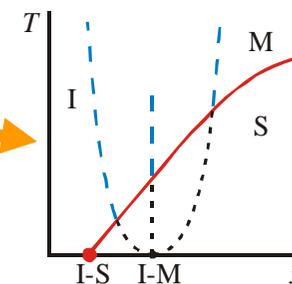
**T.I. Baturina, A.Yu. Mironov,
V.M. Vinokur, M.R. Baklanov,
and C. Strunk**

Phys. Rev. Lett. **99**, 257003 (2007)



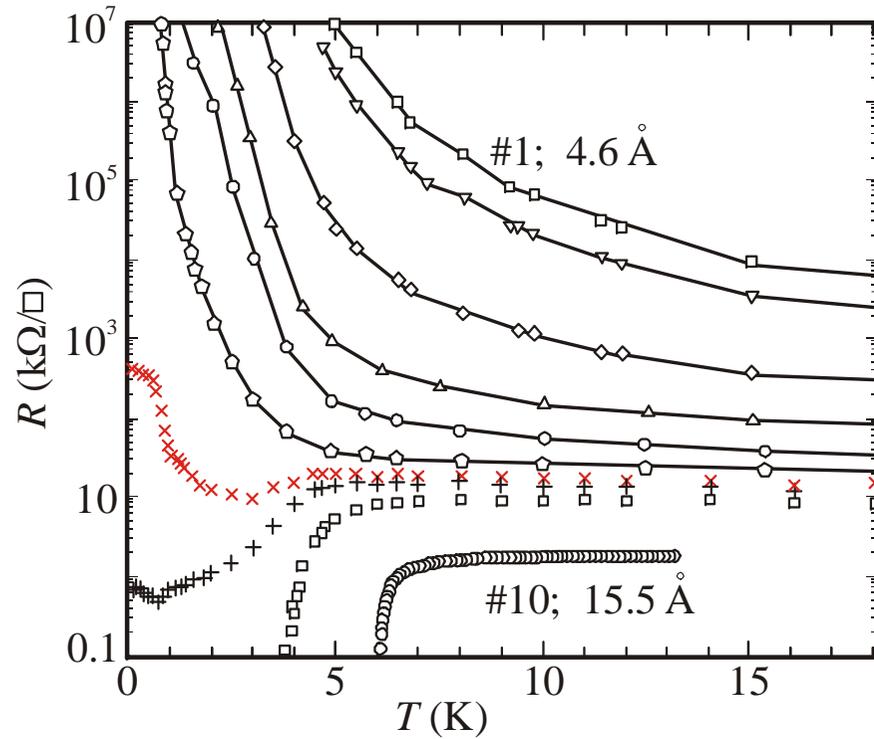
**T.I. Baturina, D.R. Islamov, J. Bentner,
C. Strunk, M.R. Baklanov, and A. Satta,**
JETP Lett. **79**, 337 (2004)

$$a + bT^{1/3}$$

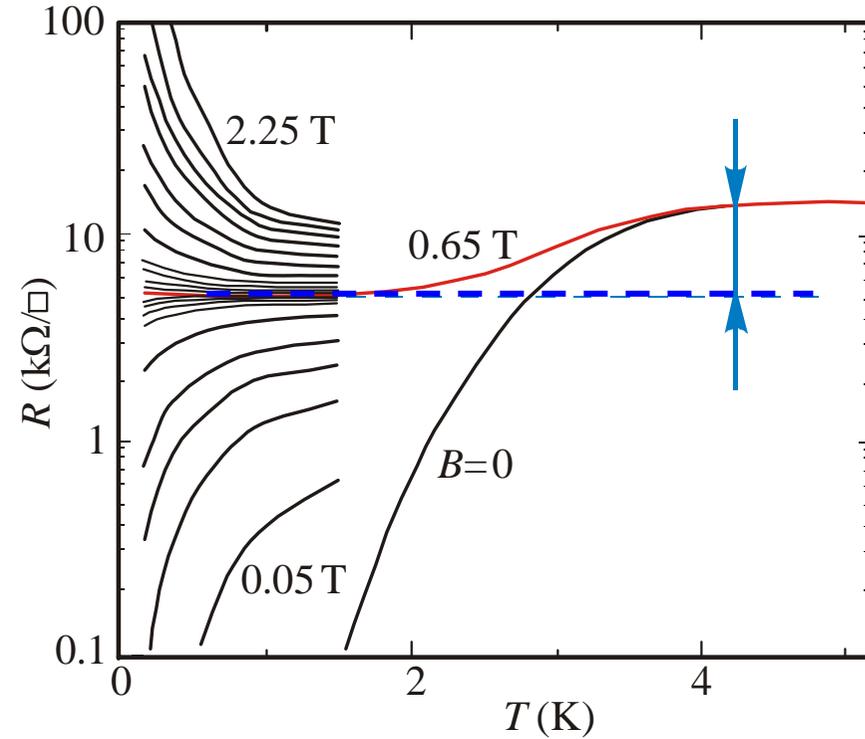


A.I. Larkin,
Ann. Phys. (Leipzig) **8**, 794 (1999)

Ultra thin Be films

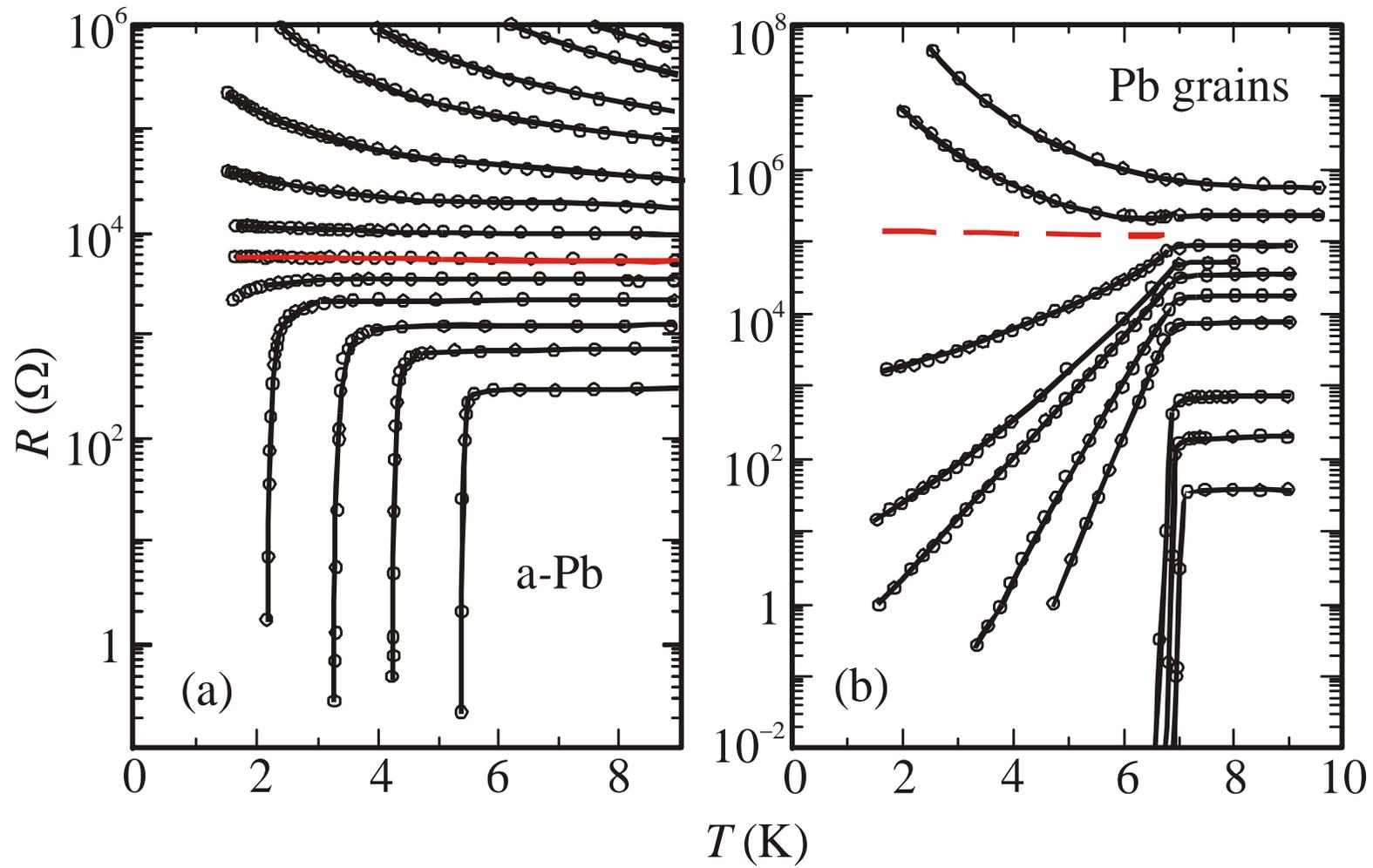


Transition induced by thickness changes



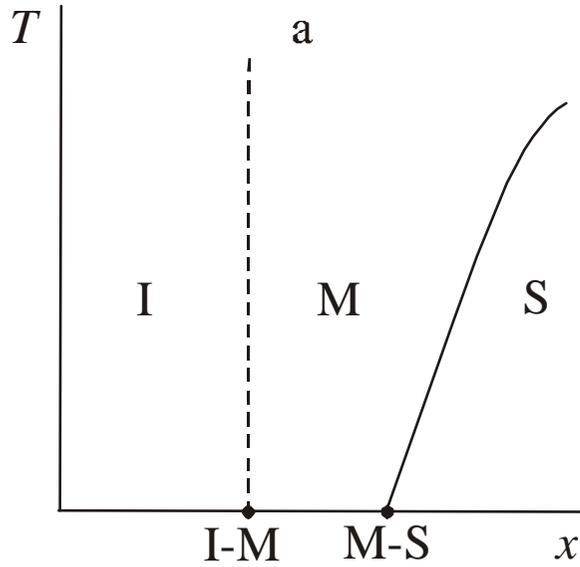
Transition induced by magnetic field

E. Bielejec and Wenhao Wu, Phys. Rev. Lett. **88**, 206802 (2002)



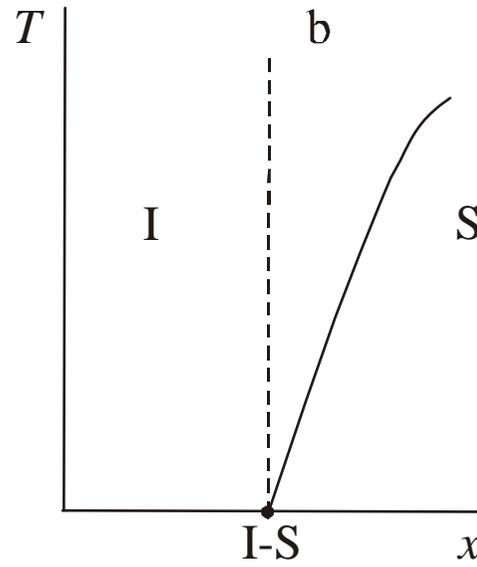
A. Frydman, Physica C **391**, 189 (2003)

Three different phase transition diagrams



split

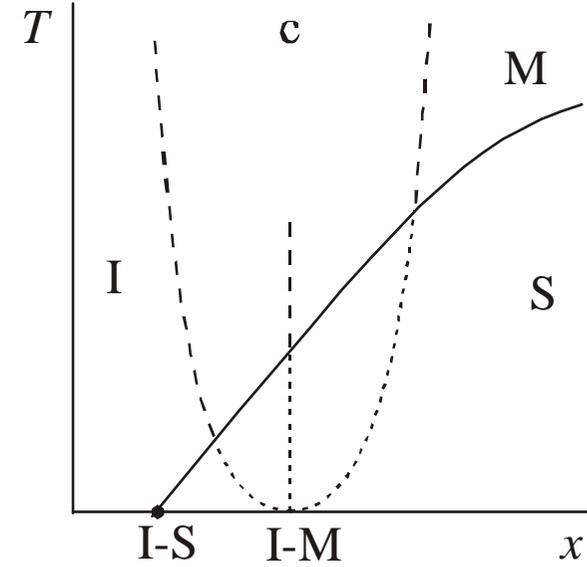
$\text{Nb}_x\text{Si}_{1-x}$
 NdCeCuO
Ta ??



single

$a\text{-Bi}$

InO_x
 $a\text{-Be}$

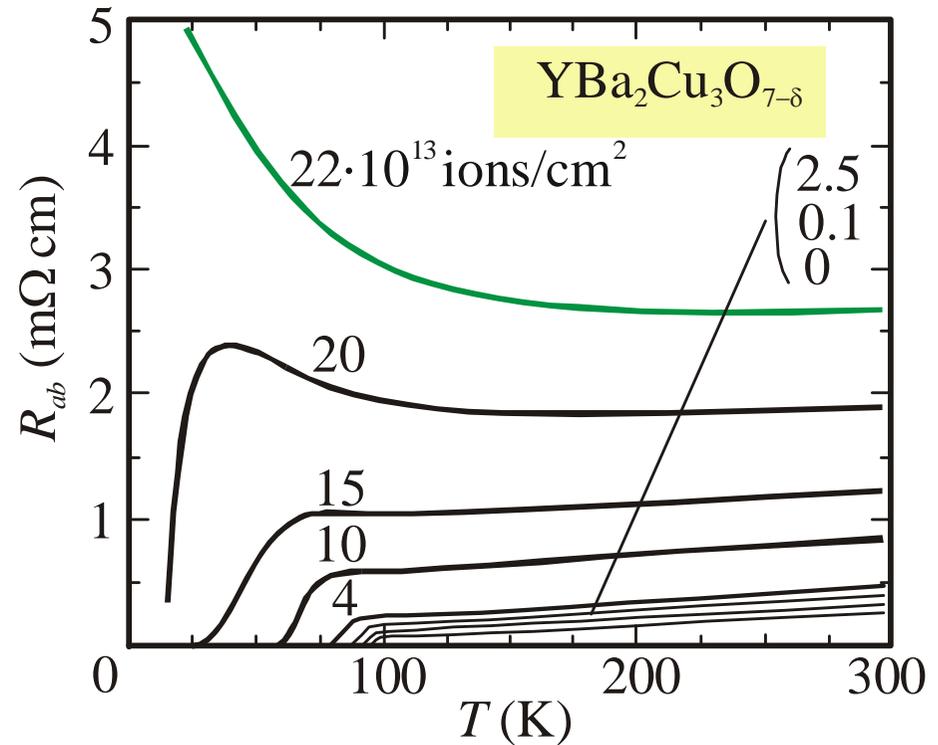


overlapped

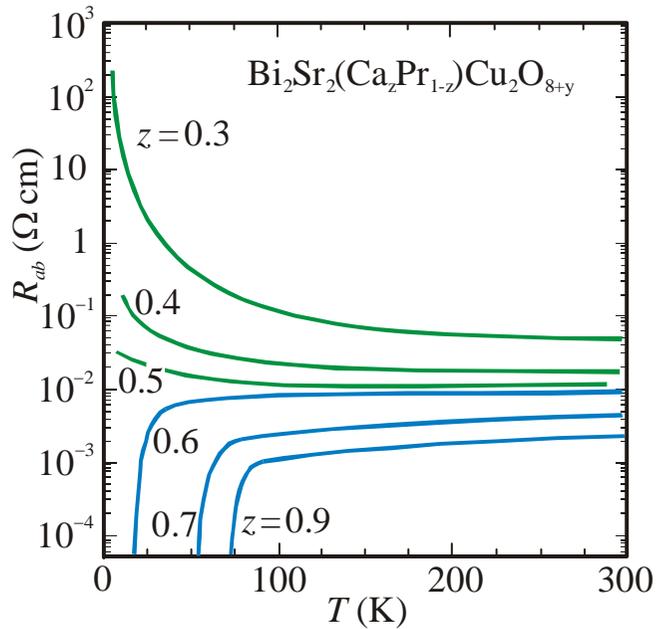
TiN_x

S – I transitions in High- T_c Superconductors

*High- T_c superconductivity
to arise
the complicate anisotropic
crystalline field is
necessary*



*J.M. Valles, Jr., A.E. White, K.T Short, R.C. Dynes,
J.P.Garno, A.F.J. Levi, M. Anzlowar, and K. Baldwin,
Phys. Rev. B **39**, 11599 (1989)*

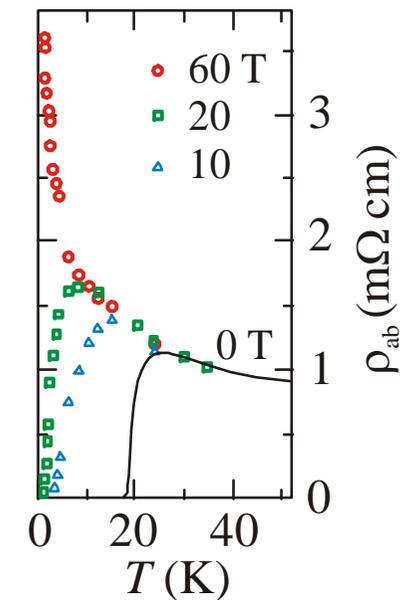
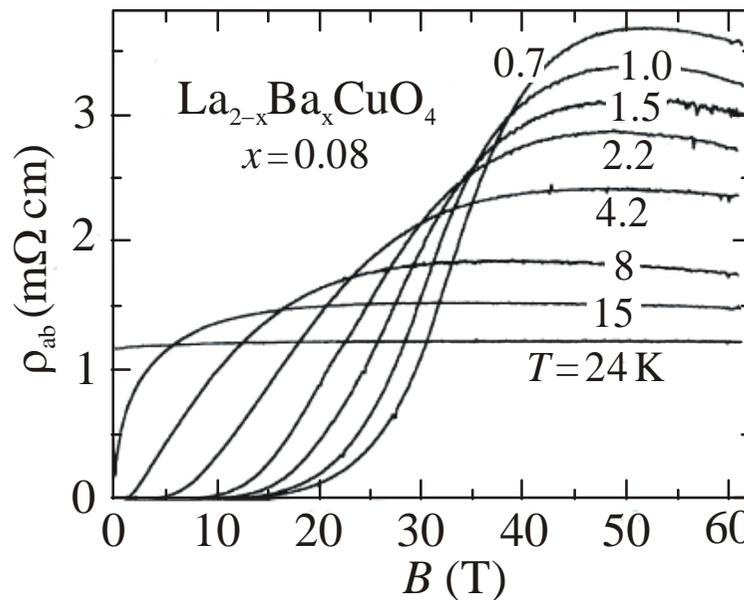


Transition induced by underdoping

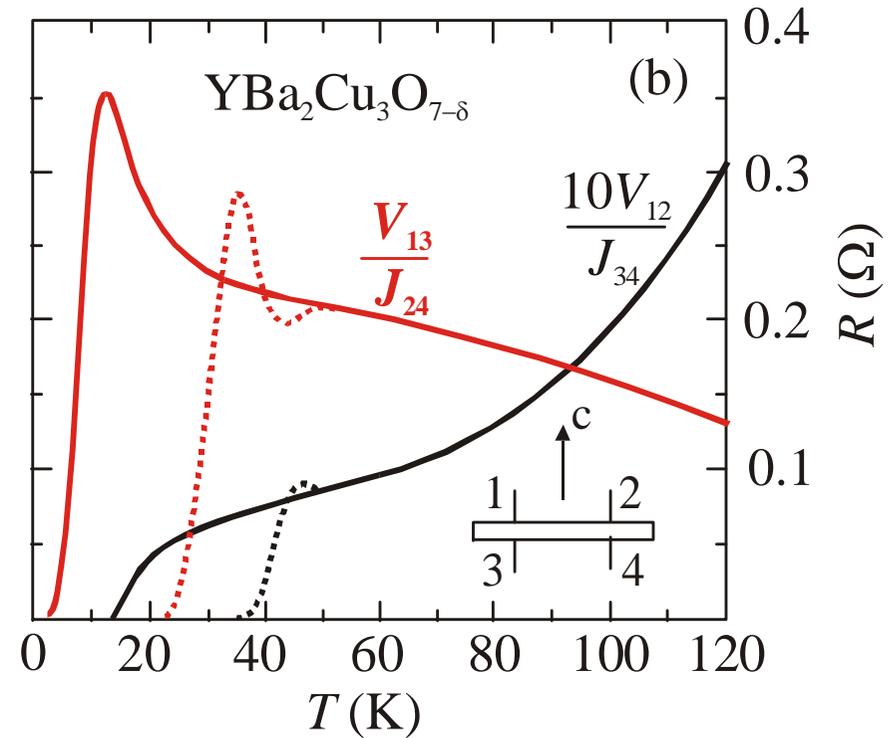
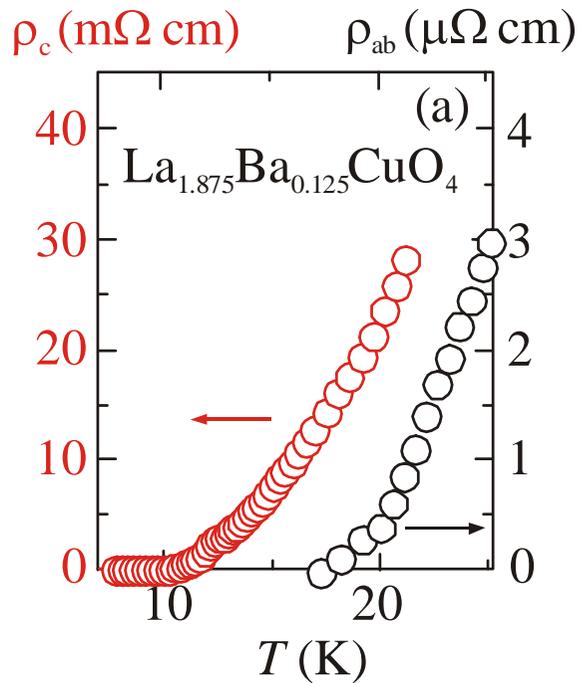
*B. Beschoten, S. Sadewasser,
G. Guntherodt, and C. Quitmann,
Phys. Rev. Lett. **77**, 1837 (1996)*

Field induced transition

*Y. Ando, G.S. Boebinger,
A. Passner, T. Kimura, and
K. Kishio,
Phys. Rev. Lett. **75**, 4662
(1995)*



Friedel effect



*Q. Li, M. Huecker, G.D. Gu,
A.M. Tsvelik, and J.M. Tranquada,*
Phys. Rev. Lett 99, 067001 (2007)

V.N. Zverev, D.B. Shovkun, I.G. Naumenko,
JETP Lett. **68**, 332 (1998)

Two possible scenarios

The order parameter

$$\Phi(\mathbf{r}) = \Delta \exp(i\varphi(\mathbf{r}))$$

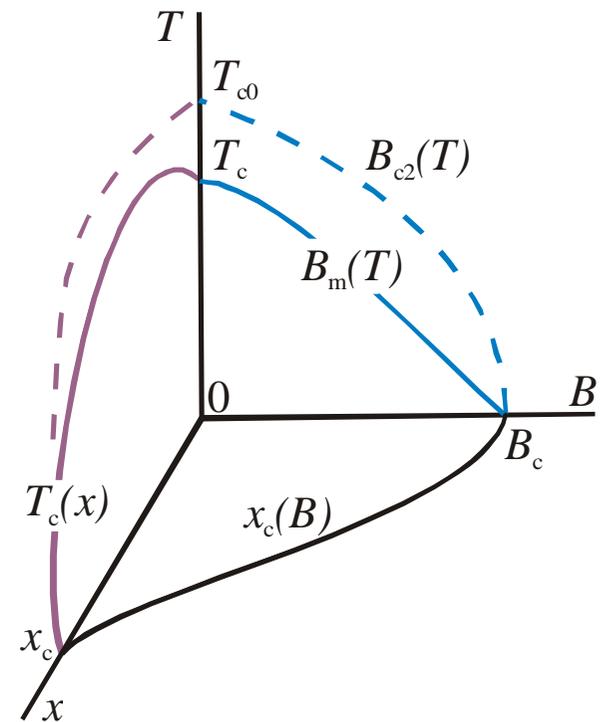
Fermionic scenario

$$\Delta \rightarrow 0$$

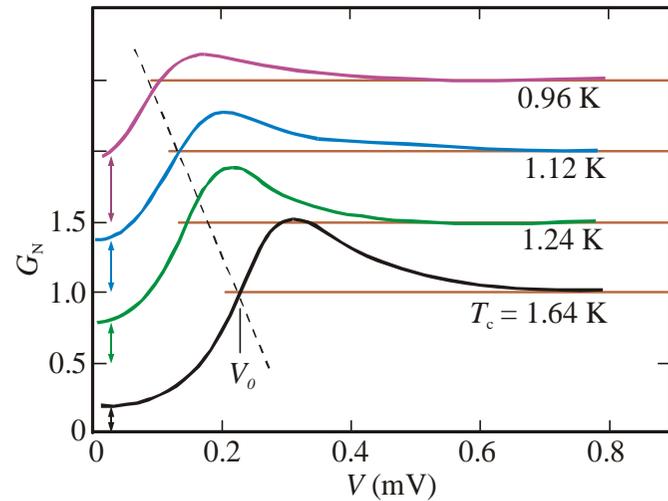
Bosonic scenario

Δ remains finite, but

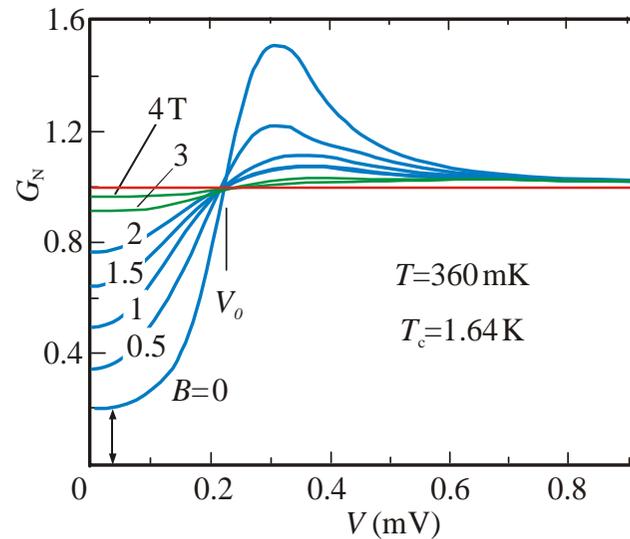
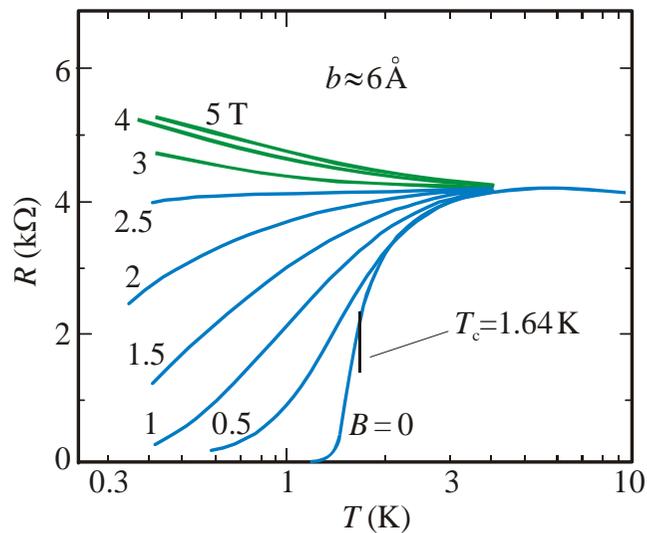
$$G(\mathbf{r}) = \langle \Phi(\mathbf{r})\Phi(0) \rangle \rightarrow 0 \text{ when } |\mathbf{r}| \rightarrow \infty$$



Ultrathin Bi films – tunnel measurements



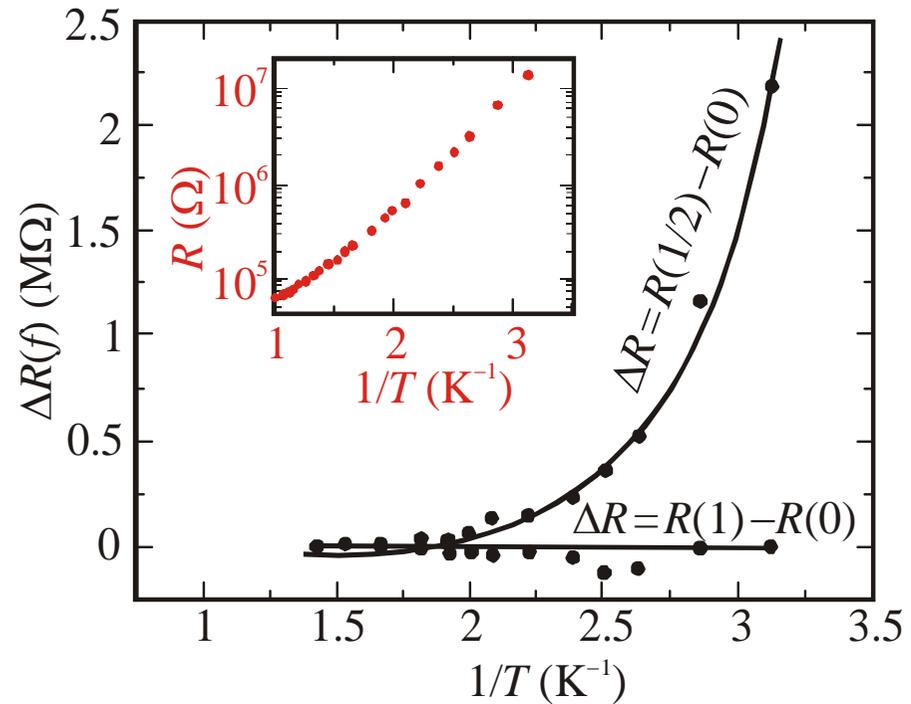
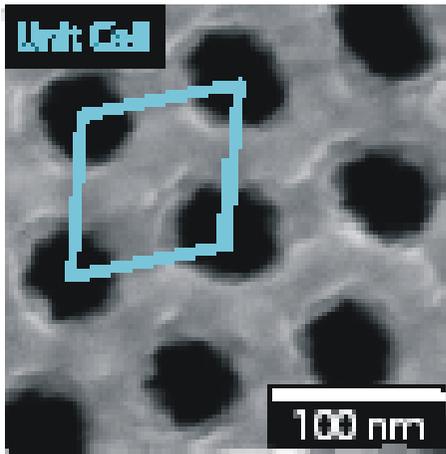
Four films of different thickness in zero magnetic field



Field dependence for one of these films

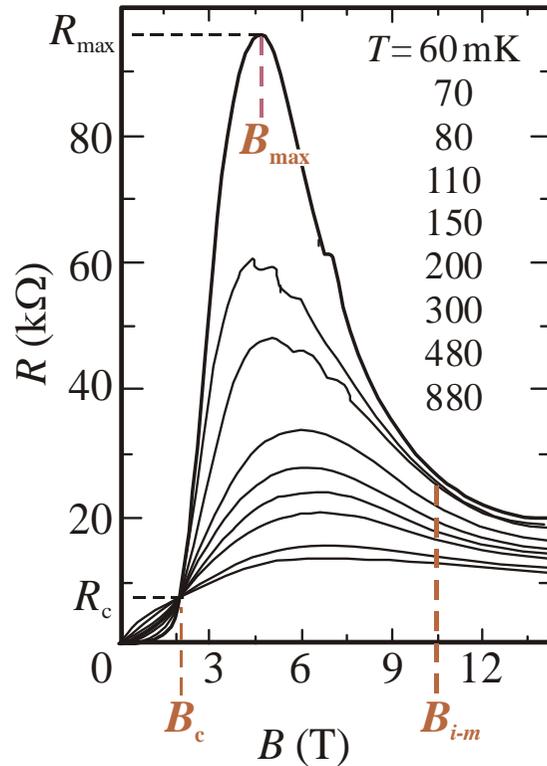
S.-Y. Hsu, J.A. Chervenak, and J.M. Valles, Jr., Phys. Rev. Lett. 75, 132 (1995)

Finite ξ of localized superconducting pairs

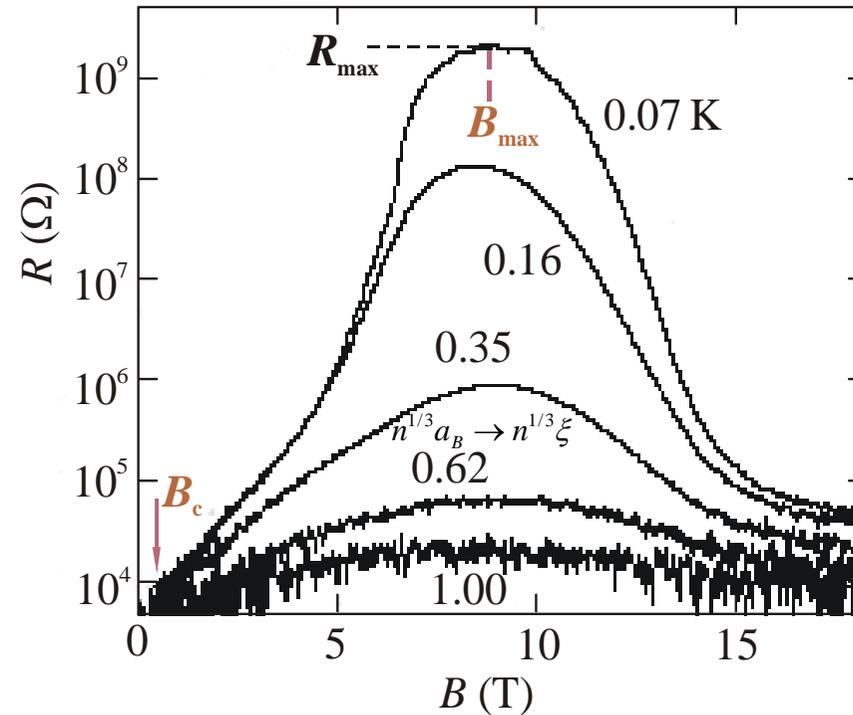


M.D. Steward Jr, A. Yin, J.M. Xu, and J.M. Valles Jr,
Science **318**, 1273 (2007)

Negative magnetoresistance in InO films



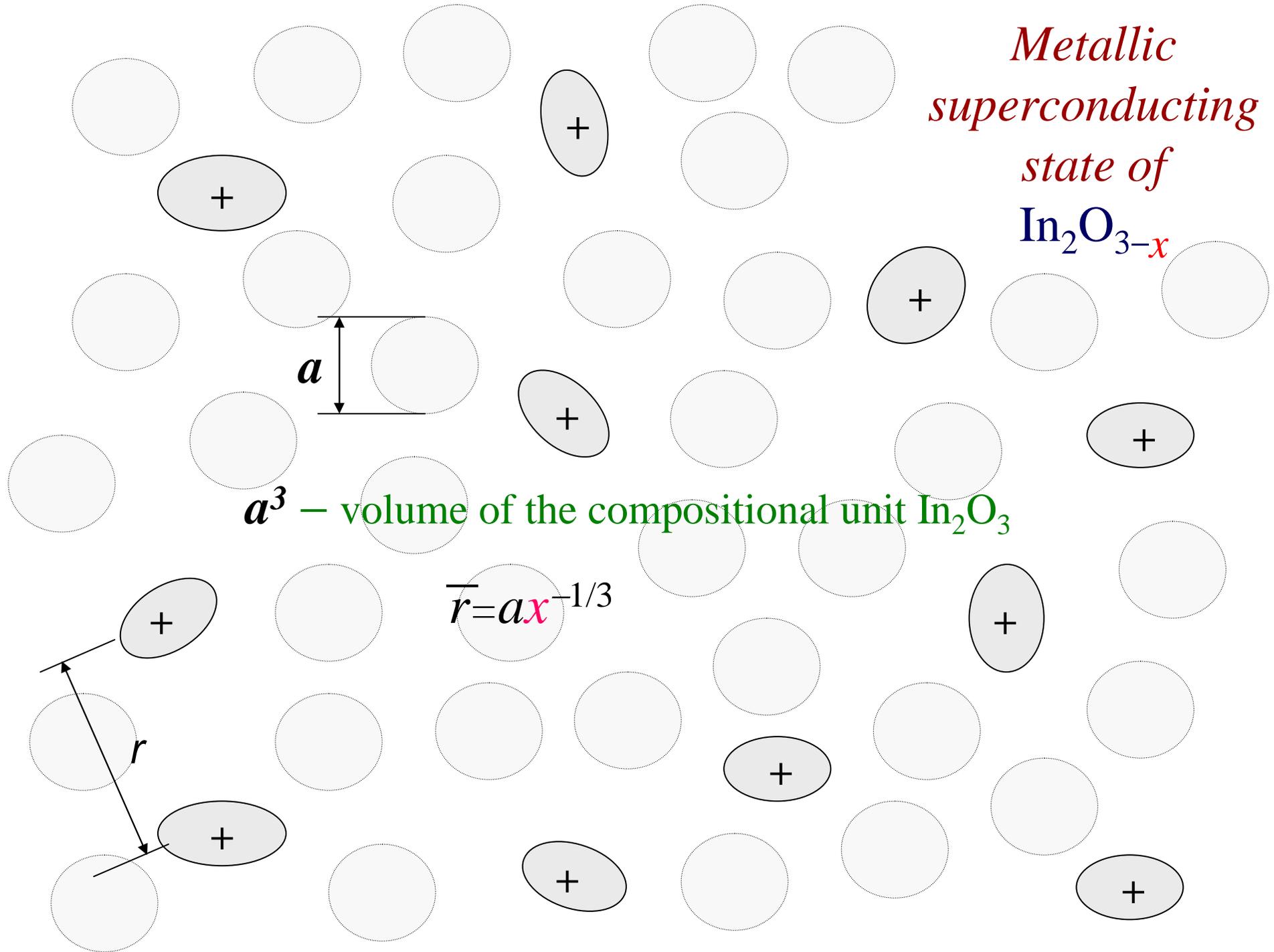
*V.F. Gantmakher, M.V. Golubkov,
V.T. Dolgoplov, A.A. Shashkin,
and G.E. Tsydynzhapov,
JETP Lett. **71**, 473 (2000)*



*G. Sambandamurthy, L.W. Engel,
A. Johansson, and D. Shahar,
Phys. Rev. Lett. **92**, 107005 (2004)*

$$B_{max} \gg B_c \quad n^{1/3} a_B \rightarrow n^{1/3} \xi \quad \xi = \xi(B - B_c, \varepsilon(B))$$

*Metallic
superconducting
state of
 $\text{In}_2\text{O}_{3-x}$*



a^3 – volume of the compositional unit In_2O_3

$$\bar{r} = ax^{-1/3}$$