



*The Abdus Salam
International Centre for Theoretical Physics*



2040-8

Workshop: Eternal Inflation

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The Lifetime of Stringy dS

M. Lippert

*Israel Institute of Technology
Israel*

B. Freivogel

*University of California At Berkeley
U.S.A.*

The Lifetime of Stringy dS

with Ben Freivogel

arXiv:0807.1104

What is the maximum lifetime of dS?

No eternal dS

Observables (Banks & Fischler, Witten, etc.)

Boltzmann Brains (DKS, etc.)

No string theory examples

Poincaré Recurrence Time: $t_r \sim e^{\mathbf{S}_{dS}} = e^{\frac{24\pi^2}{G_4^2 V_{dS}}}$

Upper bound: $t_d \leq t_r \sim e^{M_p^2/H^2}$ (Susskind, Giddings, etc)

Can we do better?

Boltzmann Brains

What are BBs?

observers formed by fluctuations
violation of 2nd law
not “ordinary” observers

Estimate BB entropy S_{BB} (rough)

$S_{BB} \sim$ minimum S for observer
estimate: # of dof $\sim$ # of particles

$S_{BB} > 10^{15}$ bacterium

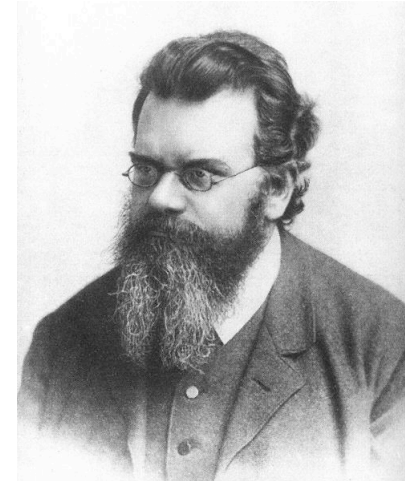
$S_{BB} < 10^{50}$ earth

Time to nucleate a BB

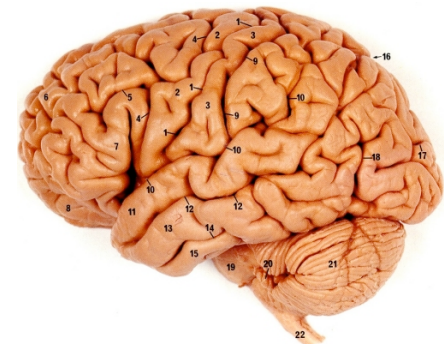
$$t_{BB} \gtrsim H^{-1} e^{S_{BB}}$$

$$t_{BB} \approx H^{-1} e^{10^{40 \pm 20}}$$

Imprecise, but $t_{BB} \ll t_r$



Boltzmann



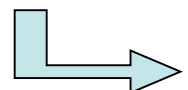
Brain

Boltzmann Brain Bound

How many BB in a causal patch? $N_{BB} \sim \frac{t_d}{t_{BB}}$

BB problem: don't want to be swamped by BBs
precise statement depends on measure

Want $t_d < t_{BB}$ to avoid conflict with observation

 nontrivial bound because $t_{BB} \ll t_r$

Is BB bound obeyed in string landscape?

Naively, easy to make $t_d \approx t_r$

→ Depends on potential, which can be anything

In string theory

potential constrained

illustrative example: KKLT

SUSY AdS → uplifted to dS
small ~~SUSY~~ (fined tuned)

CDGKL estimate $t_d \sim e^{M_p^2/m_{3/2}^2}$

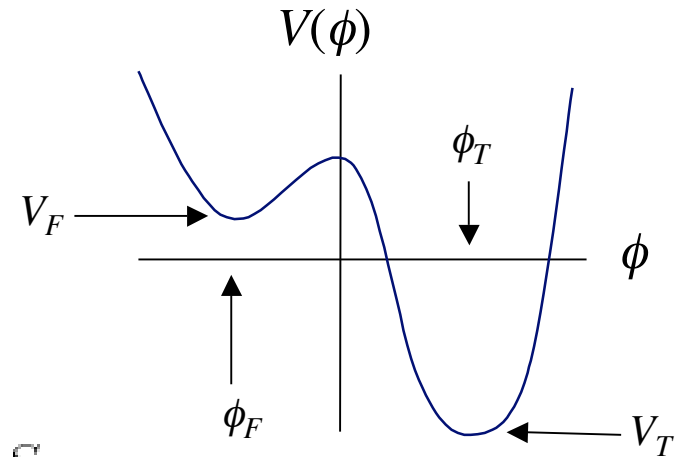
We find: this is an overestimate

False Vacuum Decay

CdL Formalism

- Scalar coupled to gravity
- Euclidean instanton for tunneling from ϕ_F to ϕ_T
- Decay rate: $t_d = \Gamma^{-1} \sim e^B$

where $B = S_{CdL} - S_F$



Lorentzian dynamics
expanding bubble of true vacuum

$$\left\{ \begin{array}{l} V_T > 0 \rightarrow \text{dS} \\ V_T = 0 \rightarrow \text{open FRW} \\ V_T < 0 \rightarrow \text{big crunch} \end{array} \right.$$

Thin-wall Limit

Transition from V_F to $V_T \approx$ Domain Wall

$$B = \underbrace{B_0(\tau, \delta V)}_{\downarrow} \underbrace{r(x, y)}_{\rightarrow \text{GR correction}}$$

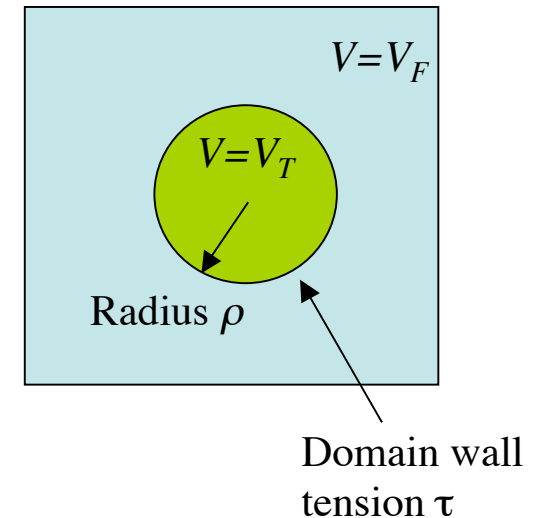
$$B_0 = \frac{27\pi^2\tau^4}{2\delta V^3} \quad \text{QFT result}$$

$$\text{where } x = \frac{\tau^2}{\tau_c^2} \quad y = \frac{V_F + V_T}{\delta V}$$

$$\text{Critical Tension } \tau_c = \sqrt{\frac{4\delta V}{3G_4}}$$

Bubble radius

$$\rho = \frac{\rho_0}{\sqrt{1 + 2xy + x^2}} \quad \text{where } \rho_0 = \frac{3\tau}{\delta V} \quad \text{QFT result}$$



Metastable dS decay

False vacuum - dS

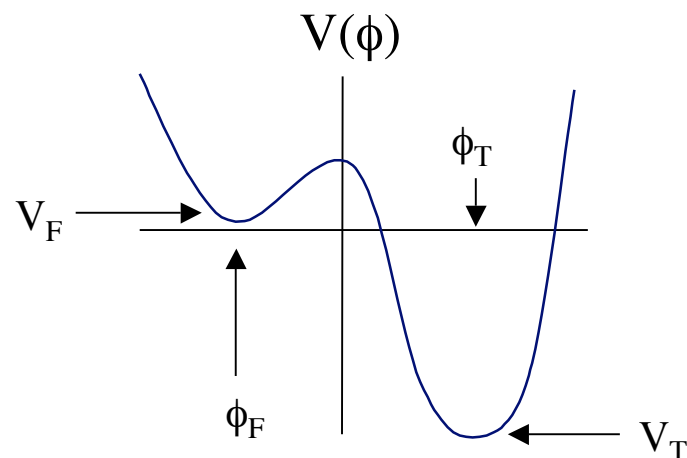
$$V_{dS} = V_F \gtrsim 0$$

True vacuum - AdS

$$V_{AdS} = V_T < 0$$

$$|V_{AdS}| \gg V_{dS}$$

$$E_{\cancel{SUSY}} = m_{3/2} \sim 1/R_{AdS} \sim \sqrt{G_4 V_{AdS}}$$



Decay Time bounded

$$B = S_{CdL} - S_F \leq -S_F = \frac{24\pi^2}{G_4^2 V_{dS}} \quad \longrightarrow \quad t_d \leq t_r$$

Limits

Subcritical Tension $x \ll 1$

gravity unimportant, use QFT results

$$B = B_0 \propto \frac{\tau^4}{\delta V^3}$$

$$\rho \sim \sqrt{x} R_{AdS} \ll R_{AdS} \ll R_{dS}$$

Critical Tension $x = 1$

gravity important

$$\rho = R_{dS}$$

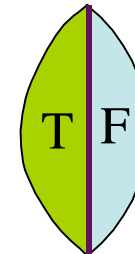
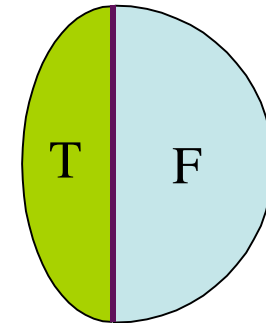
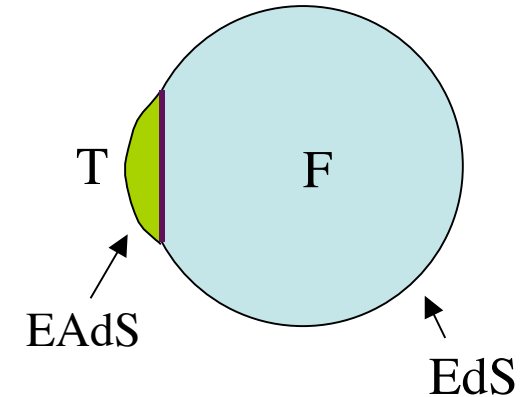
Supercritical Tension $x \gg 1$

gravity important

$$B_{KPV}^0 \approx \frac{24\pi^2}{G_4^2 V_{dS}} = S_{dS}$$

$$\rho \sim \frac{R_{AdS}}{\sqrt{x}}$$

Instanton



KKLT de Sitter Vacua

Flux Compactification

F theory compactified on elliptically fibered CY 4-fold

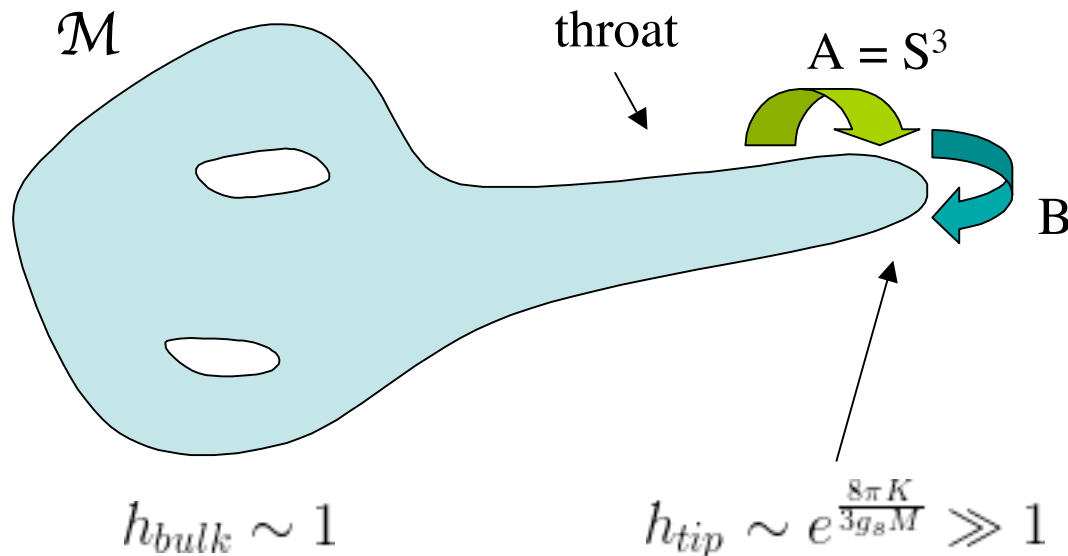


IIB compactified on 6d Orientifold $\mathcal{M}$ w/ D7's

Turn on ISD flux $G_3 = F_3 - \tau H_3$ on 3-cycles of $\mathcal{M}$

$$ds^2 = h^{-1/2} g_{\mu\nu}(x) dx^\mu ds^\nu + h^{1/2} e^{2u} dy^m dy^n$$

$\mathcal{M}$ near Conifold point $\sim$ Warped Deformed Conifold (KS throat)

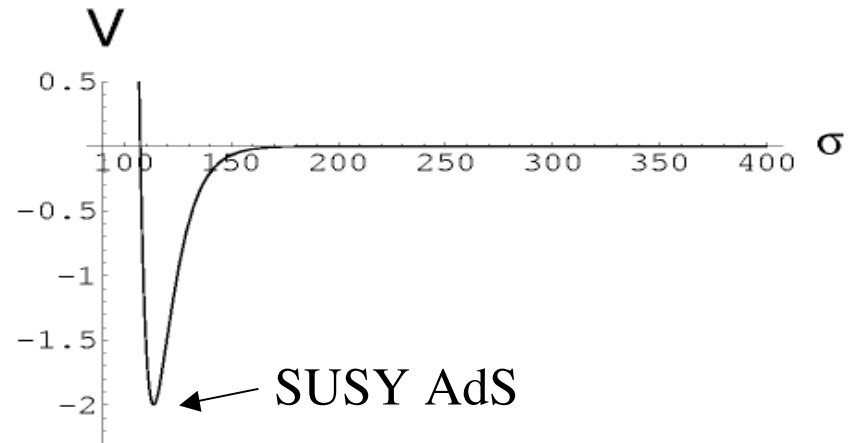


Fluxes

$$M \sim \int_A F_3$$

$$K \sim - \int_B H_3$$

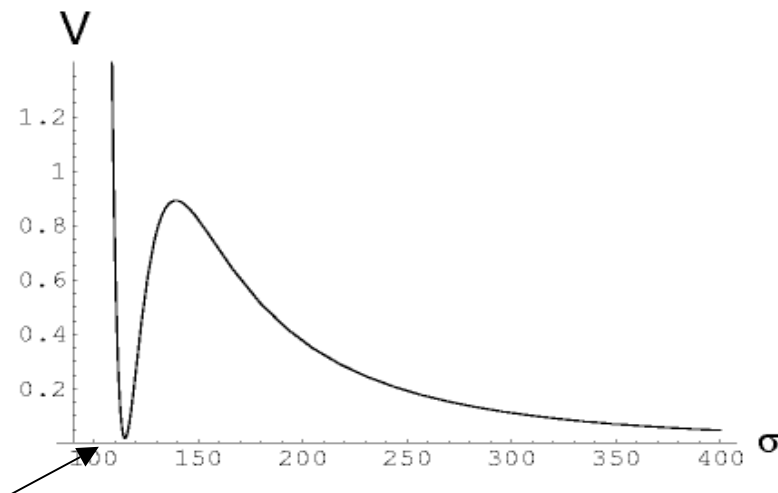
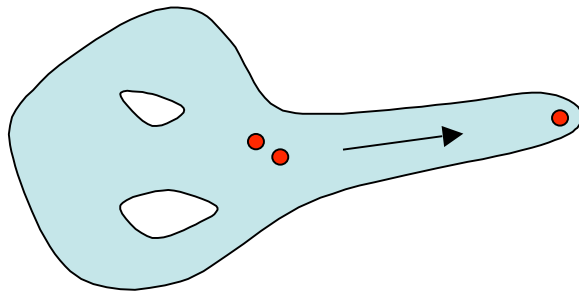
Non-pert. effects generate
potential for $\sigma = e^{4u}$
(other moduli fixed)



Add $N_{\overline{D3}} \overline{D3}$ s

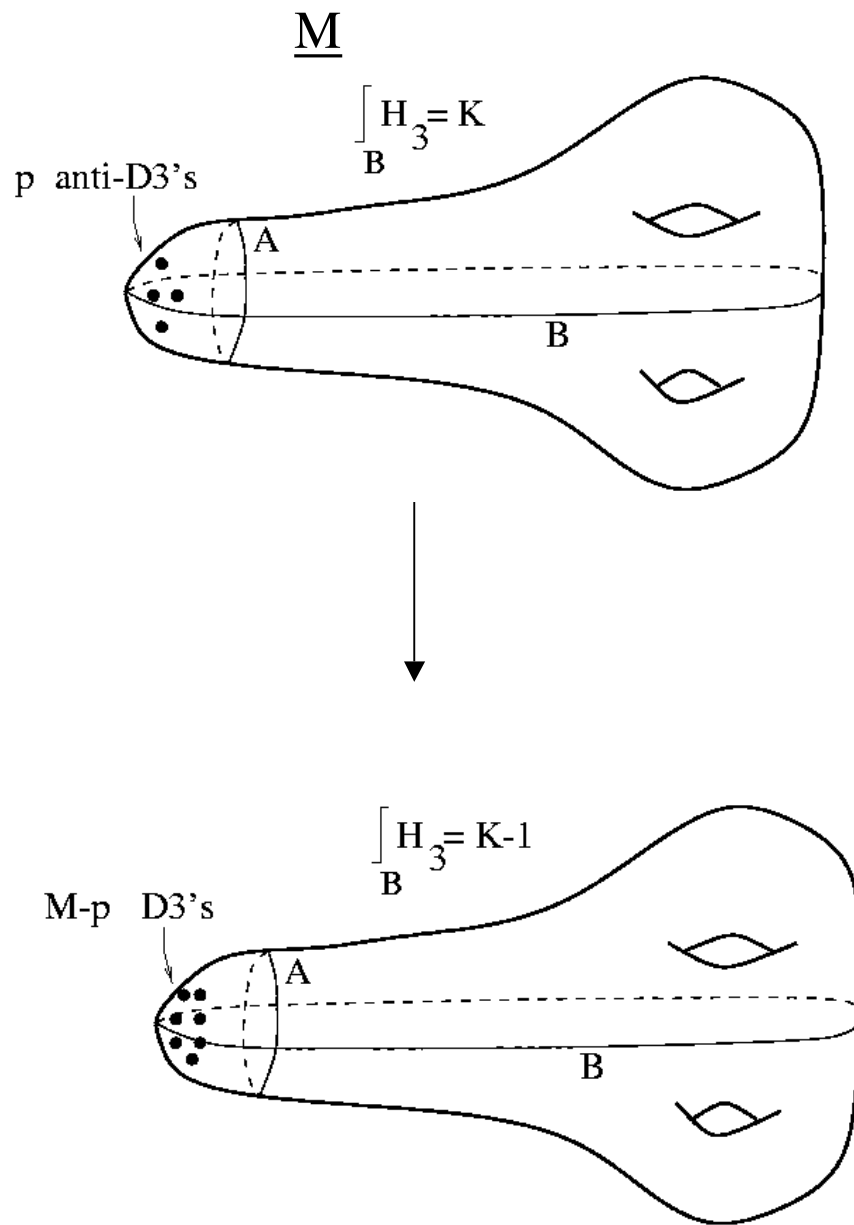
~~SUSY~~

migrate to tip
contribute δV

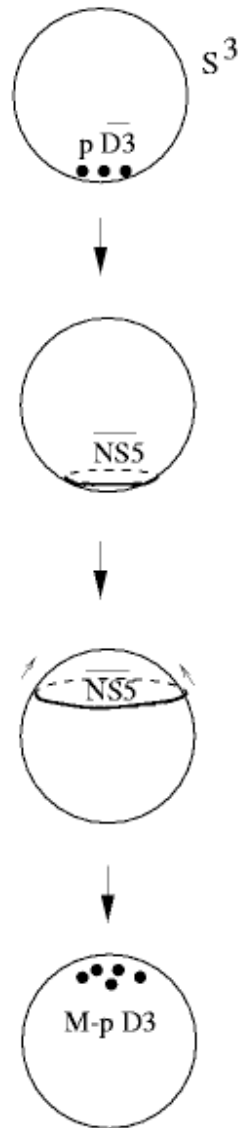


Metastable dS

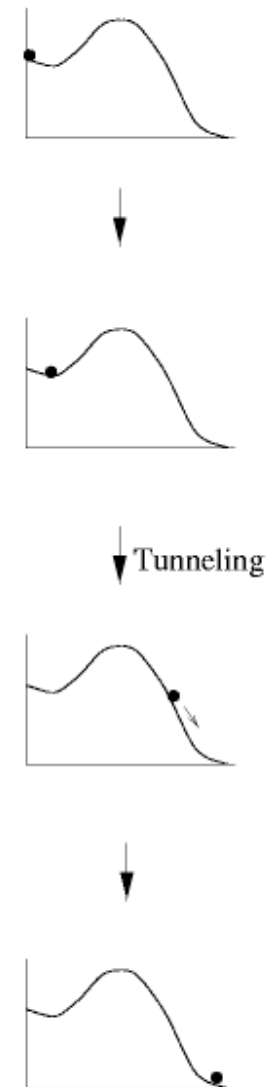
Brane/Flux Annihilation (KPV)



A-cycle



Potential



KPV decay - QFT approx

Instanton (thin-wall) = NS5 wrapped on S^3 , bubble in spacetime

$$\begin{array}{lcl}
 \text{4d tension } \tau = \tau_{NS5} V_{S^3} h_{tip}^{-3/4} = \frac{M^{3/2}}{16\pi^3 g_s^{1/2} l_s^3} h_{tip}^{-3/4} & & \\
 \text{Energy difference } \delta V \approx |V_{AdS}| = \frac{2N_{\overline{D3}}}{8\pi^3 g_s l_s^4} h_{tip}^{-1} & & \\
 & \searrow & \\
 & B_{KPV}^0 = 3 \cdot 10^{-3} \frac{g_s M^6}{(N_{\overline{D3}})^3} &
 \end{array}$$

Why doesn't this depend on h_{tip} ?

everything localized at tip

measure relative to bulk $\longrightarrow$ h_{tip} factors

use proper quantities $\longrightarrow$ independent of h_{tip}

How big can this be?

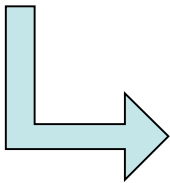
$$B_{KPV}^0 = 3 \cdot 10^{-3} \frac{g_s M^6}{(N_{\overline{D}3})^3}$$

choose $N_{\overline{D}3} = 1$

tadpole $MK < \frac{\chi_{CY_4}}{24}$

KS throat $K > g_s M > 1$

$$\left. \vphantom{\begin{matrix} MK < \frac{\chi_{CY_4}}{24} \\ K > g_s M > 1 \end{matrix}} \right\} M < \frac{\chi_{CY_4}}{24}$$


$$B_{KPV}^0 < 4 \cdot 10^{-10} \chi_{CY_4}^5$$

Largest known $\chi_{CY_4} \sim 10^6$

Assume this is a bound $\longrightarrow t_d < H^{-1} e^{10^{22}}$

Gravity Corrections

So far ignored gravity

Ok if $x \ll 1 \iff \rho \ll R_{dS}, R_{AdS}$

$$x = \frac{3\pi^4 g_s^2 M^3}{2e^{6u} N_{\overline{D}3}} h_{tip}^{-1/2}$$

Can x ever be > 1 ?

$$\left. \begin{array}{ll} \text{CY} > \text{throat} & e^{4u} > 3\pi^3 g_s M K \\ \text{KS throat} & K > g_s M > 1 \end{array} \right\} x < 10^4 h_{tip}^{-1/2}$$

$$h_{tip} \left\{ \begin{array}{ll} > 10^8 & \text{long throat} \implies \text{QFT result ok} \\ < 10^8 & \text{short throat} \implies t_{KPV} \approx t_r ! \end{array} \right.$$

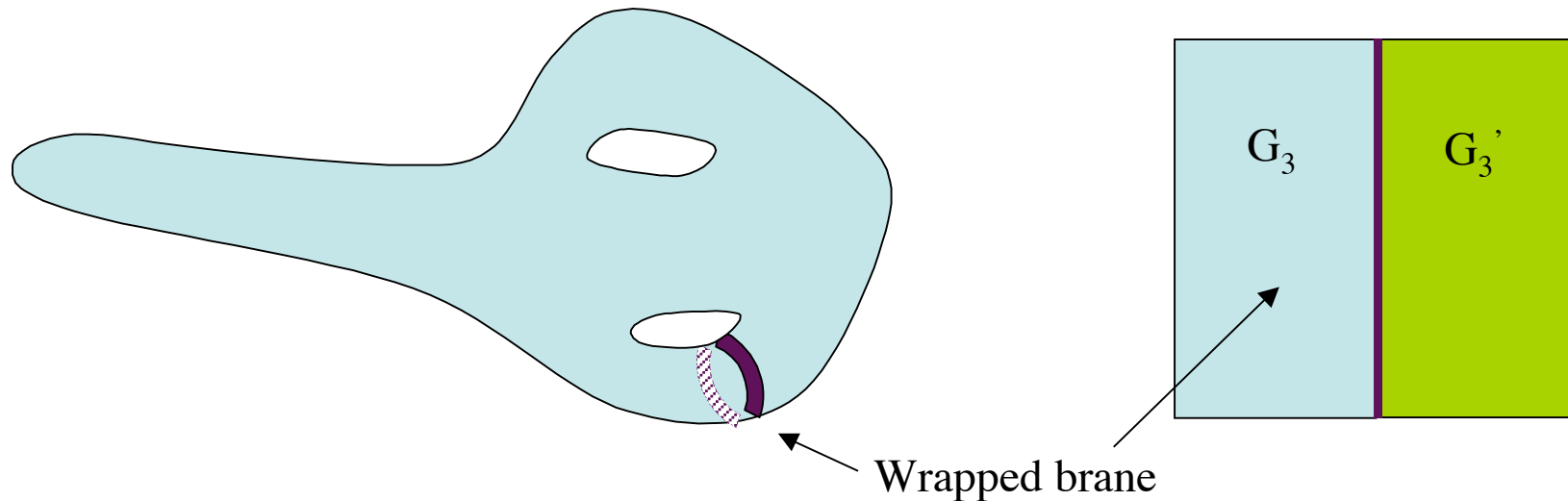
Bulk Flux Decay

KPV not the only decay channel

Before uplifting to dS $\Rightarrow$ SUSY AdS

wrapped branes change flux vacua

BPS domain walls



~~SUSY~~ dS

near-BPS walls $\Rightarrow$ instantons

estimate action

$$B_{CDGKL} = \frac{6\pi^2}{G_4^2 V_{AdS}}$$

(CDGKL)

KPV vs flux decay

$$B_{CDGKL} \sim \frac{1}{G_4^2 V_{AdS}} \sim R_{AdS}^4 V_{AdS} \quad \text{use } R_{AdS} = G_4 V_{AdS}$$

$$\sim \underbrace{\frac{R_{AdS}^4}{\rho_0^4}}_{x^{-2}} \underbrace{V_{AdS} \rho_0^4}_{B_{KPV}^0}$$

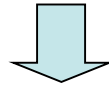
For long throat $x < 1 \Rightarrow B = B_{KPV}^0$ KPV
 For short throat $x > 1 \Rightarrow B = B_{KPV}^0 / x^2$ Bulk flux



$$B < B_{KPV}^0 \quad \forall x$$

Bound on KKLT lifetime

$$B_{KPV}^0 < 4 \cdot 10^{-10} \chi_{CY_4}^5$$



$$t_d < H^{-1} e^{4 \cdot 10^{-10} \chi_{CY_4}^5} < H^{-1} e^{10^{22}}$$

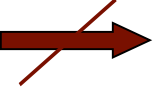
Compare with BB time

$$t_{BB} \approx H^{-1} e^{10^{40 \pm 20}}$$

What's not to like?

- Not airtight bound
e.g. CDGKL estimate
- Depends sensitively on bound on χ_{CY_4}
- Only one tiny corner of landscape
Is this a general feature of string landscape?

Slogans

Nearly SUSY  Nearly stable

Here, t_d independent of $m_{3/2}$ (for small $m_{3/2}$)

String vacua make V_{dS} small by accident

need ingredients at scales other than V_{dS}

t_d set by these scales

Conjecture: $t_d < t_{BB}$ for all string vacua

long-lived string vacua hard to make

HW: find counter-example