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Wave breaking of electrostatic waves in warm plasma

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Motivation

Wave breaking is a poorly understood phenomenon in plasma physics... but why?

- Too many conflicting definitions
- Too many models
- No effort to understand consequences of one's own model
- No effort to compare different models



Definition of wave breaking

“Wave breaking occurs when the forward fluid or plasma motion is larger than the wave phase speed”

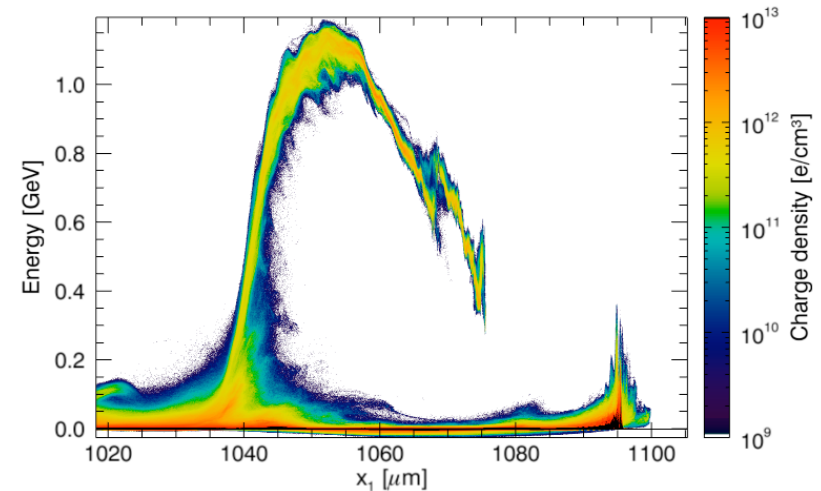


Image: L.O. Silva



Wave breaking in cold plasma

Wave breaking for an oscillation in cold plasma with wave number k and amplitude A , according to

Dawson:

- Neighbouring plasma fluid elements “collide” during oscillations
- $k*A = 1$
- $v = \omega*A = (\omega/k)*k*A = v_{\phi}$
- Plasma fluid elements **overtake** (and get **trapped** in) the wave

J.M. Dawson, Phys. Rev. **113**, 383 (1959)

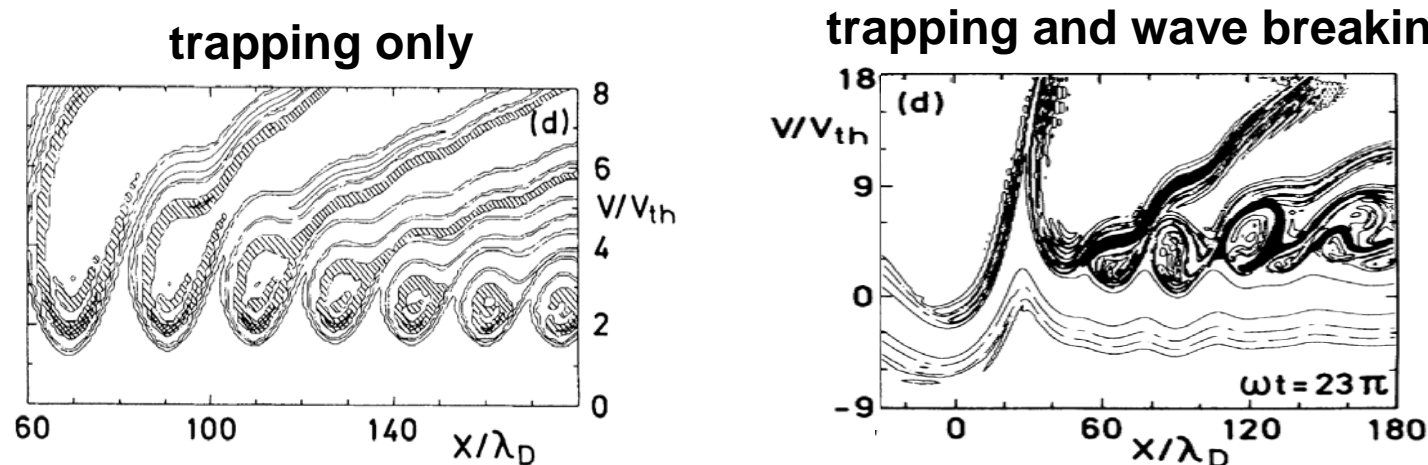


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Wave breaking in warm plasma

There is always **some** trapping in a **warm** plasma

Wave breaking implies **heavy** particle trapping



Definition: a wave breaks when it traps particles at the **electron sound speed**: $s_0^2 = 3k_B T/m_e$

A. Bergmann and P. Mulser,
Phys. Rev. E **47**, 3585 (1993)

Wave breaking in practice

Most theoretical papers use the following approach:

- Derive **warm-fluid model** from Vlasov equation
- Push model until it **breaks down**
- **Confuse** model breakdown with wave breaking
- If model **does not break down**, wave breaking **does not exist ???**



Problem with standard approach

- **Breakdown** is a mathematical property of a model
- One must **prove separately** (e.g. by relating breakdown to particle trapping) that this corresponds to a physical phenomenon like wave breaking
- **Few authors bother to do this**
- Result: different **models**, different **breakdown limits**, different “**wave breaking limits**”



Case 1: quasi-static waves

Wave in a homogeneous plasma, only depends on coordinate $\zeta = z - v_\phi * t$

Cold, non-relativistic: $E_{WB} = v_\phi$

J.M. Dawson, Phys. Rev. **113**, 383 (1959)

Cold, relativistic: $E_{WB} = \sqrt{2(\gamma_\phi - 1)}$

Akhiezer and Polovin, Sov. Phys. JETP **3**, 696 (1956)

Warm, non-relativistic:

$$E_{WB} = v_\phi \left[1 - \frac{8}{3} \left(\beta / v_\phi^2 \right)^{1/4} \right]^{1/2}, \quad \beta = 3k_B T / (mc^2)$$

T.P. Coffey, Phys. Fluids **14**, 1402 (1971)

Warm, relativistic: **needs more attention**



Case 1: warm, (ultra-)relativistic

Many models for this case, most are **invalid**.

Relativistic waterbag: $E_{WB} \propto \sqrt{\ln(\gamma_\phi^2 \beta)} / \beta^{1/4}$

- **Proper** relativistic plasma pressure
- **Good correspondence** between model breakdown and particle trapping
- Particle trapping handled **properly**
 - Use of **warm-plasma** potential to study trapping in **warm plasma**
 - **No separate plasma pressure term** in particle Hamiltonian, as it should be
- $E_{WB} \rightarrow \infty$ for $\gamma_\phi \rightarrow \infty$, as it should be

Katsouleas and Mori, PRL **61**, 90 (1988)

Trines and Norreys, Phys. Plasmas **13**, 123102 (2006)

Case 1: alternative models

Warm-plasma approximation: $E_{WB} \propto 1/\beta^{1/4}$
misbehaves for $\gamma_\phi \rightarrow \infty$ (E_{WB} **fails** to diverge)

J.B. Rosenzweig, Phys. Rev. A **38**, 3634 (1988)

Z.M. Sheng and J. Meyer-ter-Vehn, Phys. Plasmas **4**, 493 (1997)

Schroeder, Esarey and Shadwick, PRE **72**, 055401(R) (2005)

Schroeder *et al.*, Phys. Plasmas **13**, 033103 (2006)

Esarey *et al.*, Phys. Plasmas **14**, 056707 (2006)

Three-fluid model: $E_{WB} \propto \sqrt{\gamma_\phi/3}$
behaves too much like **cold** plasma

J.B. Rosenzweig, Phys. Rev. A **40**, 5249 (1989)

Method of characteristics: $E_{WB} = \infty$
wave **never** breaks in this model

Aleshin *et al.*, Plasma Phys. Rep. **19**, 523 (1993)

A. Khachatryan, Phys. Plasmas **5**, 112 (1998)



Case 2: non-quasi-static waves

Non-quasi-static plasma waves exhibit a spectrum of “weird” phenomena, such as:

- **position**-dependent frequency
- **amplitude**-dependent frequency
- **secular** behaviour
- wave number **advection** in thermal plasma
- **curbing** of wave breaking by thermal effects (unheard of for quasi-static waves)



Non-constant frequency

Relativistic cold-plasma oscillations according to Polovin:

$$\frac{\partial^2 p}{\partial \tau^2} + \frac{\omega_p^2(\underline{x})}{\sqrt{1+p^2}} p = 0$$

Position-dependent frequency can result from:

- non-constant **density**
- position-dependent **amplitude**

R.V. Polovin, Sov. Phys. JETP 4, 290 (1957)



Secular behaviour

Consider the following plasma oscillation:

$$\delta = A \sin[\psi(\underline{x}, \tau)], \quad k \equiv \partial\psi / \partial \underline{x}, \quad \omega \equiv -\partial\psi / \partial \tau$$

It follows (Whitham): $\frac{\partial k}{\partial \tau} + \frac{\partial \omega}{\partial \underline{x}} = 0$

If ω depends on \underline{x} , then k depends on τ !

Example: k will grow **linearly** on a downward **density ramp**

Significant consequences for **wave breaking**

J.M. Dawson, Phys. Rev. **113**, 383 (1959)

J.F. Drake *et al.*, Phys. Rev. Lett. **36**, 196 (1976)

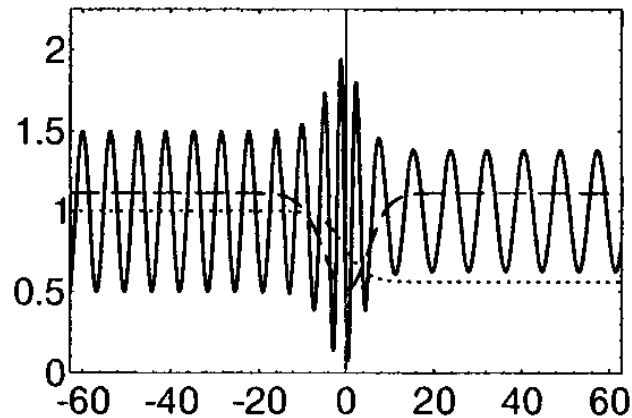


Secular behaviour II

Wave breaking condition (Dawson, Coffey):

$$kA = 1 - \left(\beta/v_{\phi}^2\right)^{1/4}$$

If k **grows**, then v_{ϕ} will **decrease** (Bulanov) and wave breaking amplitude will decrease



Dotted: plasma density
Solid: plasma oscillations
Dashed: local phase speed

Bulanov *et al.*, PRE 58, R5257 (1998)



Wave number advection

Add **thermal effects** (Bohm-Gross):

$$\omega^2 = \omega_p^2 + 3v_T^2 k^2$$

This leads to:

$$\frac{\partial k}{\partial \tau} + 3v_T^2 \frac{k}{\omega} \frac{\partial k}{\partial x} = -\frac{\omega_p}{\omega} \frac{\partial \omega_p}{\partial x}$$

Thermal effects make k **advect** away from the “hot spot” !

Which will win, secular behaviour or thermal effects?

R. Trines, PRE **79**, 056406 (2009)



Curbing of wave breaking

Wave on density ramp, total density drop Δn :

- In **cold** plasma, k will **grow indefinitely** and wave will **always** break
- In **warm** plasma, growth of k is **curbed** by thermal effects
- For small amplitude, small Δn and **large plasma temperature**, wave will **not break**
- Compare with quasi-static waves, where thermal effects facilitate wave breaking

R. Trines, PRE **79**, 056406 (2009)



Case 3: driven waves (resonance absorption)

- **Related** to previous case (density ramp), but waves are **driven** by long laser beam
- Found **4** different models for **warm** plasma:
 - Ginzburg, Propagation of EM waves in plasma (1960)
 - Kruer, Phys. Fluids 22, 1111 (1979)
 - Bezzerides and Gitomer, Phys. Fluids 26, 1359 (1983)
 - Bergmann and Mulser, PRE 47, 3585 (1993)
- Will have to pull all these together to solve this problem: not easy



Conclusions

- **Wave breaking**: an important phenomenon, often **misunderstood** because of **bad definitions**
- With **good** definitions and **careful** analysis, one can make significant headway
- **Thermal** effects and **secular** behaviour together lead to fascinating behaviour
- Resonance absorption will be next
- Always some other aspect to explore, e.g. using **more than 1** dimension!

