



*The Abdus Salam  
International Centre for Theoretical Physics*



**2054-6**

## **Structure and Dynamics of Hydrogen-Bonded Systems**

*26 - 27 October 2009*

### **Concerted Proton Tunneling in Ice Ih?**

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# Concerted proton tunneling in Ice Ih?

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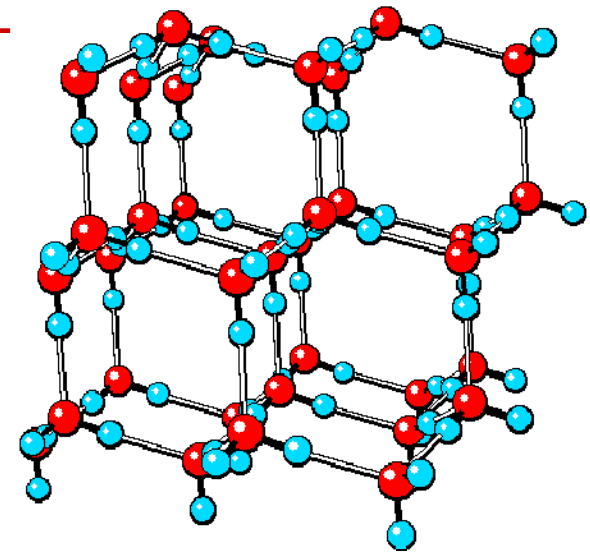


# Incoherent cross section of ice Ih

## ✓ Hexagonal ice: the prototype of *ice disorder*

Principal building units: buckled hexagons with O at corners and 2 H-sites in between (average occupation  $\frac{1}{2}$ )

Protons occupy randomly the two possible sites



## ✓ Incoherent neutron scattering

Measure the FT of the **spatial probability density function of a single proton** :

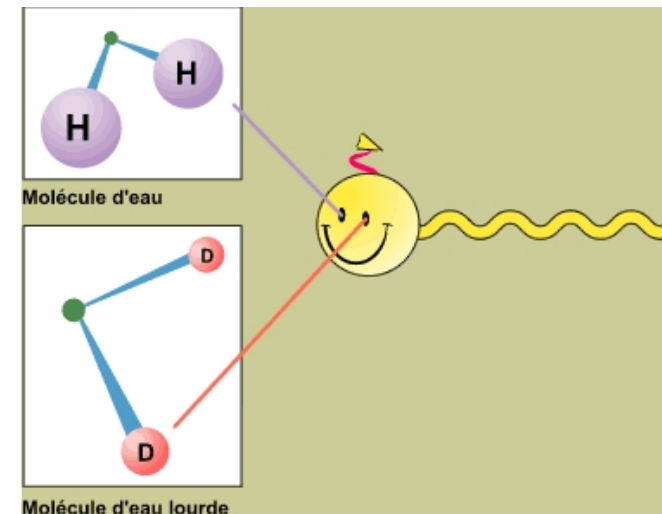
'elastic'  $\omega \sim 0$  long time scale configurations are probed

'quasi-elastic' as a function of time (ps-ns)

For a harmonic crystal  $\rightarrow$  Debye Waller factor

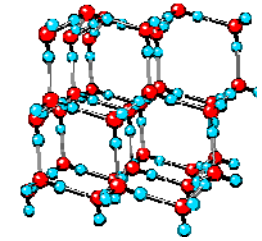
$$S(Q, \omega \sim 0) = \sigma_{\text{inc}} / 4\pi \cdot a_0 \exp(-Q^2 \langle \Delta x^2(T) \rangle)$$

High visibility of H:  $\sigma_{\text{inc}} = 81$  barns

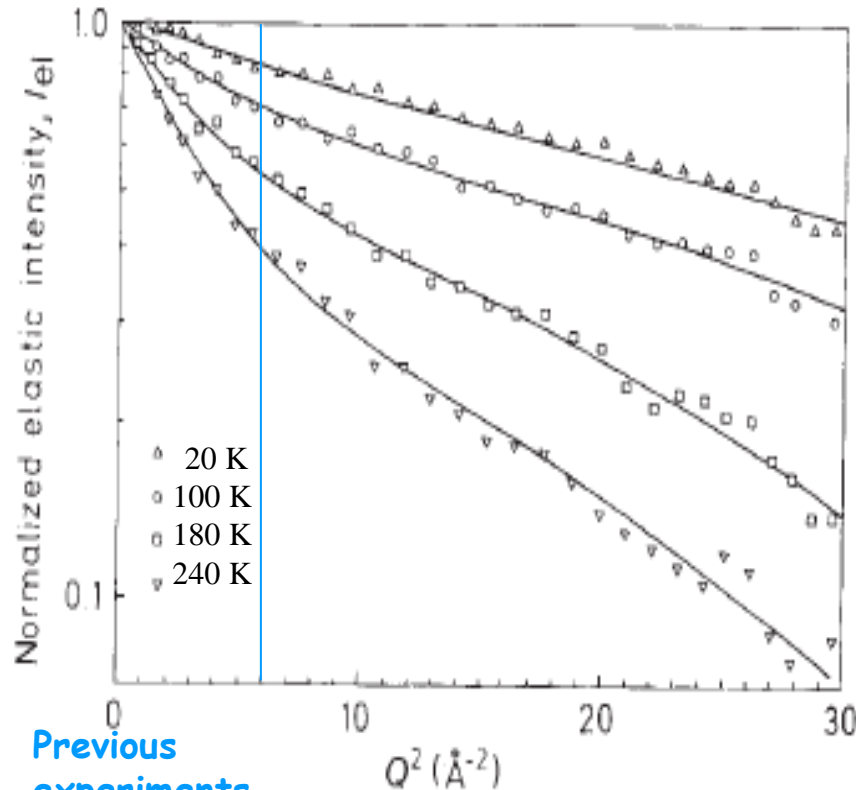


Wavelength  $\sim$  atomic distances

# Non harmonic behavior of ice Ih! [IN13-ILL]



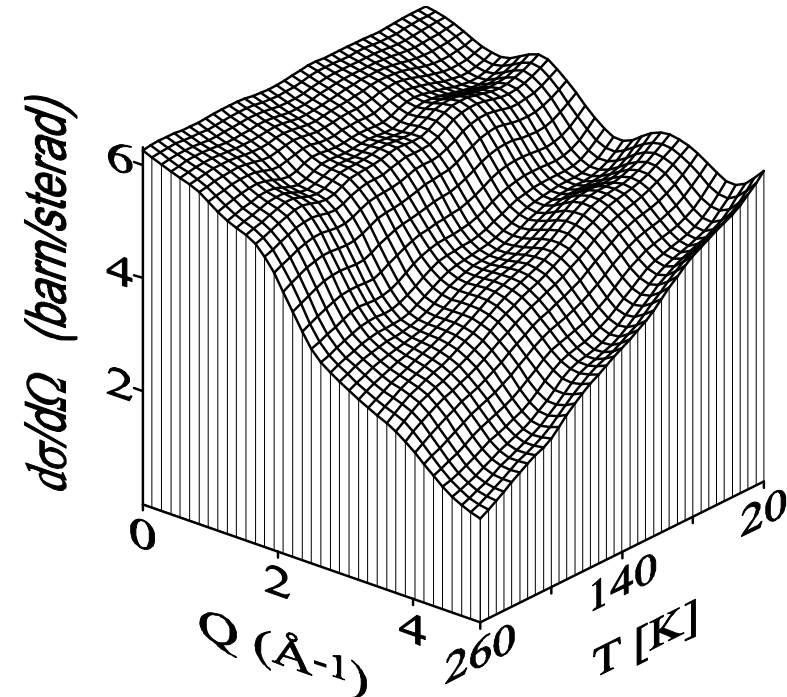
## Normalized elastic intensity



Previous  
experiments  
Q range

The elastic intensity can  
not be fitted by a Q- Gaussian function :  
**non-harmonic motion of protons!**

## Incoherent cross section



15 temperatures from 20 K to 260K

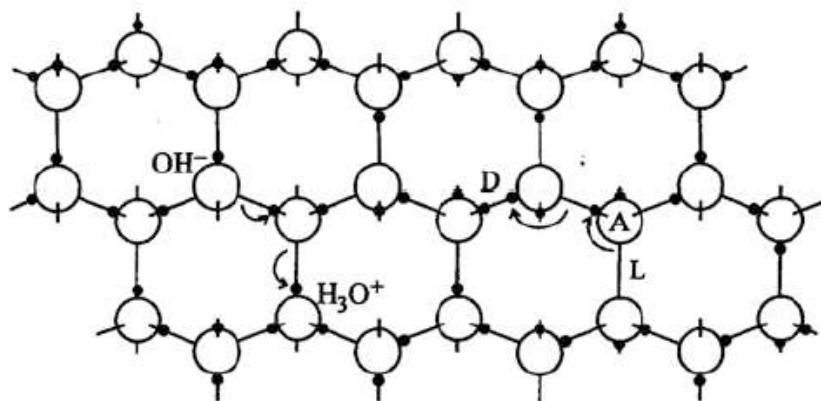
Oscillatory trend → at least one special  
distance in the single-proton probability  
density function

**Coherence effects**

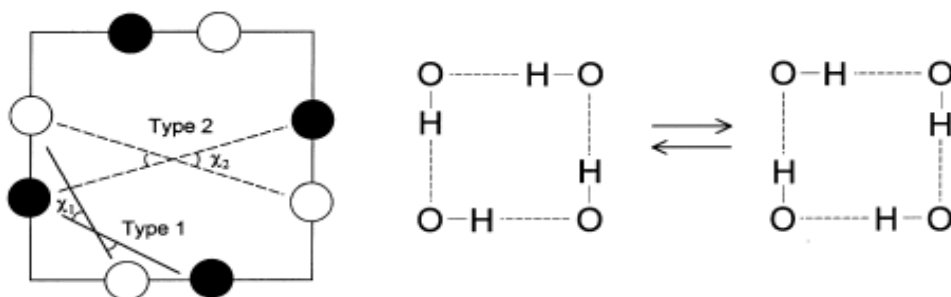
Number of protons involved ~%

L.E.B. et al, in publication

# How can a proton move in a cyclic network?



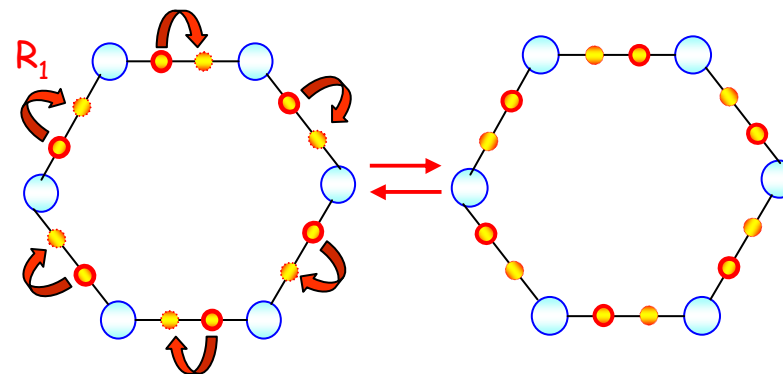
A single proton cannot jump from one site to the other without producing defects :  
High activation barrier  $\sim \rightarrow 10^{-6}$  events



**Saenger et al., Nature 296 (1992)**  
*Flip-Flop H bonding in a partially disordered Cyclodextrin*

**Brougham et al., Nature 397 (1999)**  
*Coordinated proton tunneling in cyclic array of 4 H- bonds (calix[n]arene)*

*Ordered hexagonal loops:* the H occupies the same site in the six bonds



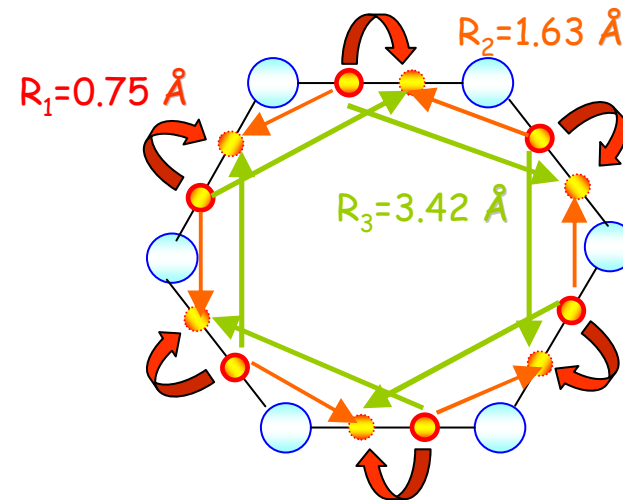
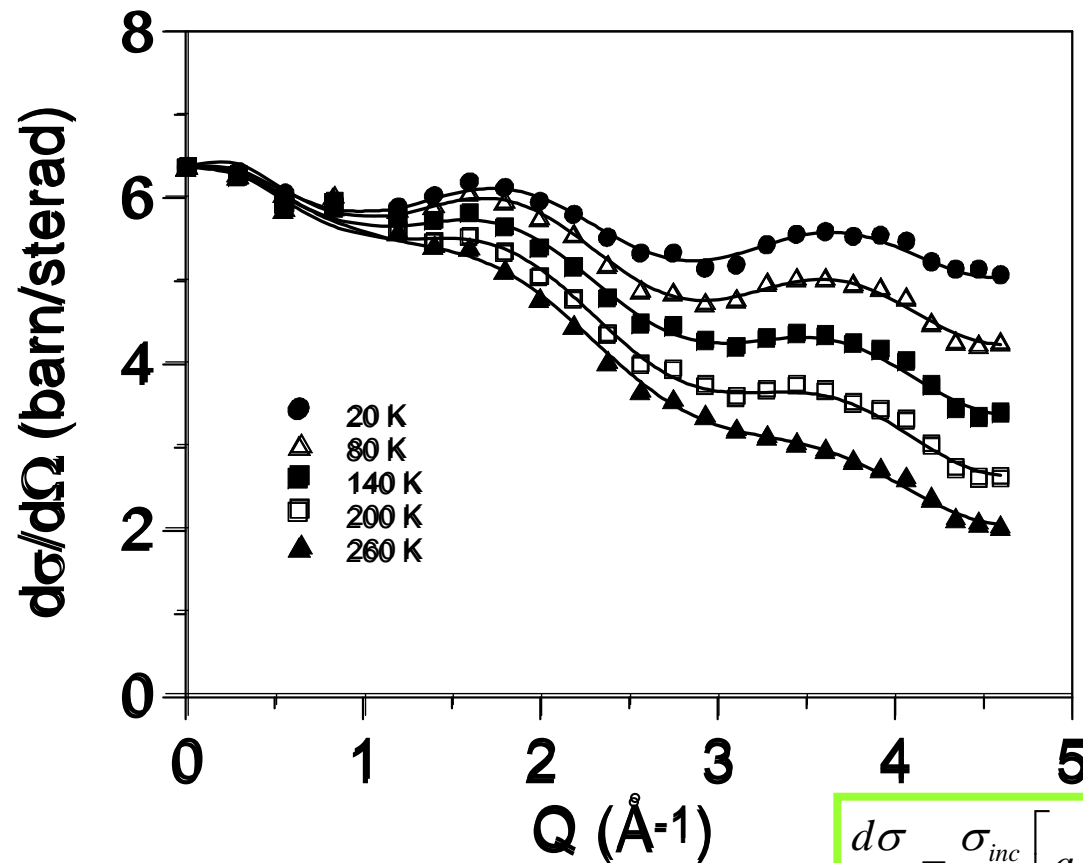
In the ordered loops 6 protons can move with no change of the total  $E_{\text{cris}}$

Coordinate motion highly favored:  
Most likely lower activation barrier

*«.. a jump of a H atom from one position to another in ordered loops causes all the connected H bonds to change in a cooperative concerted mechanism (domino effect) »*

L. Pauling, J Am. Chem. Soc. 57 (1935)

# Concerted proton jump model



$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{inc}}{4\pi} \left[ a_0 \exp[-\langle \Delta x^2 \rangle Q^2] + \sum_i^N a_i \cos(Q \cdot R_i) \right]$$

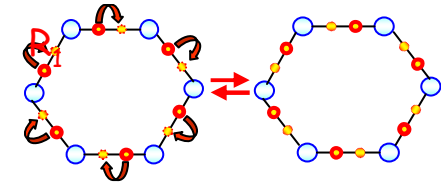
The best fit is obtained using two distances parameters only:

$R_1 = 0.75 \pm 0.03 \text{ \AA} \rightarrow$  H sites along one O-O side; weight factor 0.9; slightly T dependent

$R_{av} = 3.4 \pm 0.05 \text{ \AA} \rightarrow$  average of all others distances?; weight factor 0.1, T independent

L.E.B. et al, in publication

# First conclusions and new challenges:



✓ Non-harmonic motion of H in ice Ih, faster than our t-window (0.5 ns)

Time scale? T-behaviour of the associated time



Quasielastic Neutron Scattering experiment on a shorter time scale

✓ Coherence effects in the incoherent cross section on a main distance of 0.75 Å

Which kind of motion? Associated with H-disorder?



Comparison with a H-ordered ice form (Ice VIII) and with a different H-disordered form (Ice Ic)

✓ H involved ~%  $\rightarrow$  low energy barrier

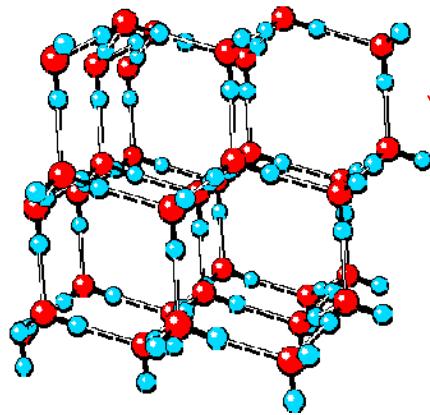
Concerted mechanism? Role of the ordered loops?



Partial deuteration to broke the loop symmetry

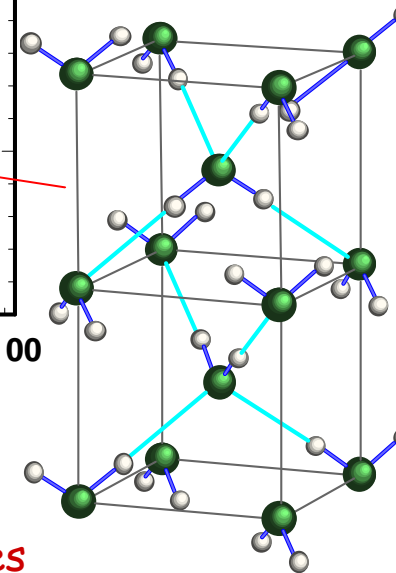
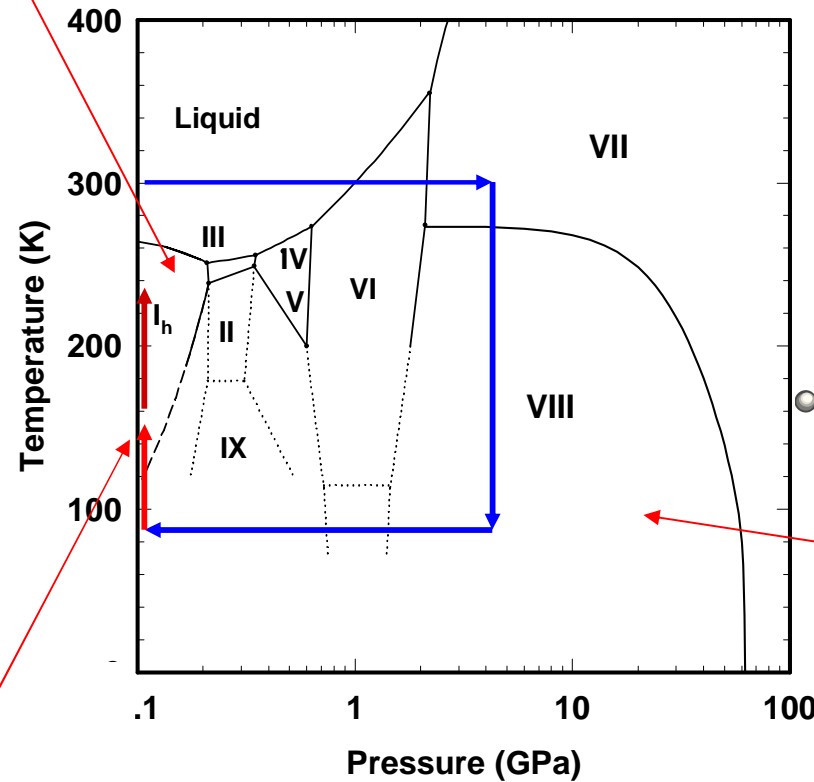
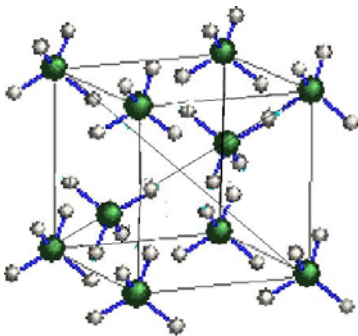


# Ice ordered and disordered structures:



At 190K  $\rightarrow$  Ice Ih  
H-disordered

At 145K  $\rightarrow$  Ice Ic  
metastable  
H-disordered

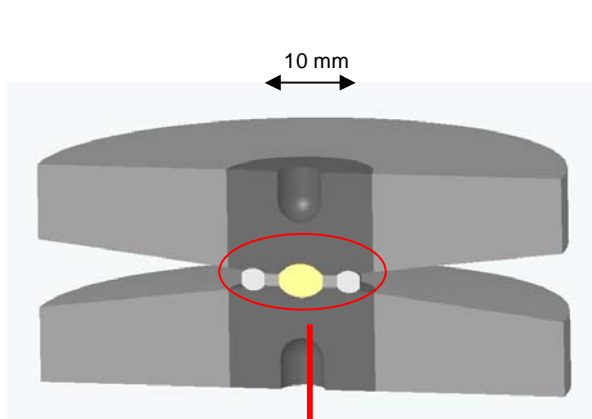


Ice VIII  
H- ordered

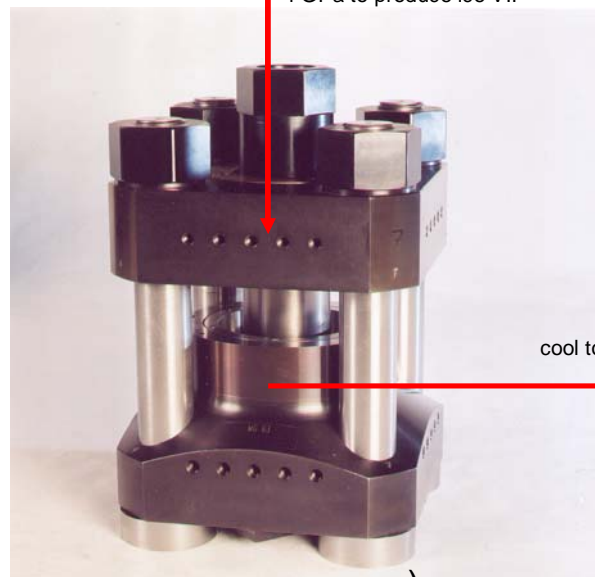
Mutual transformation of the 3 phases  
Minimize the possible spurious differences



# Making large quantities of H<sub>2</sub>O ice VIII

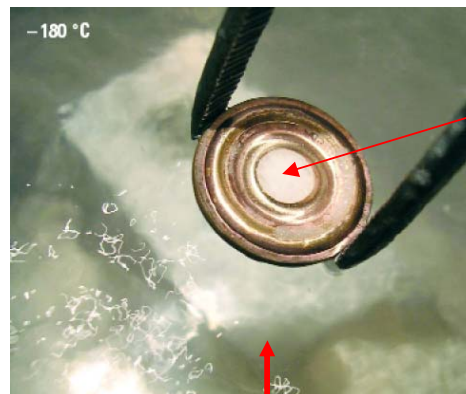


Increase load to  
4 GPa to produce ice VII



Paris-Edinburgh Press (50 kg)

cool to 77 K

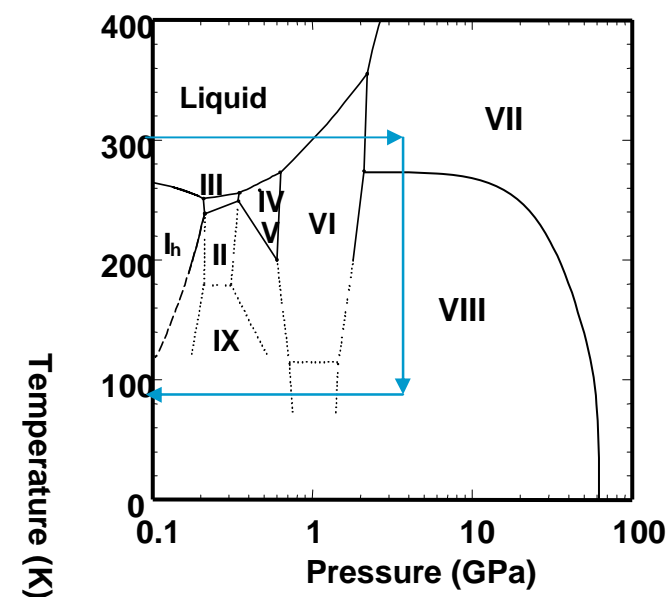
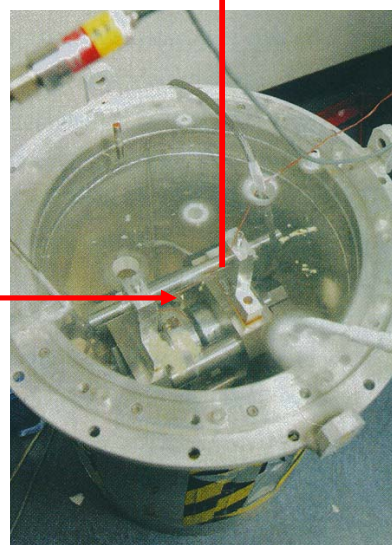


Ice VIII



40 mm<sup>3</sup>/loading  
1 loading/day  
25 loadings:  
1 cm<sup>3</sup>/month

Decrease load  
and recover

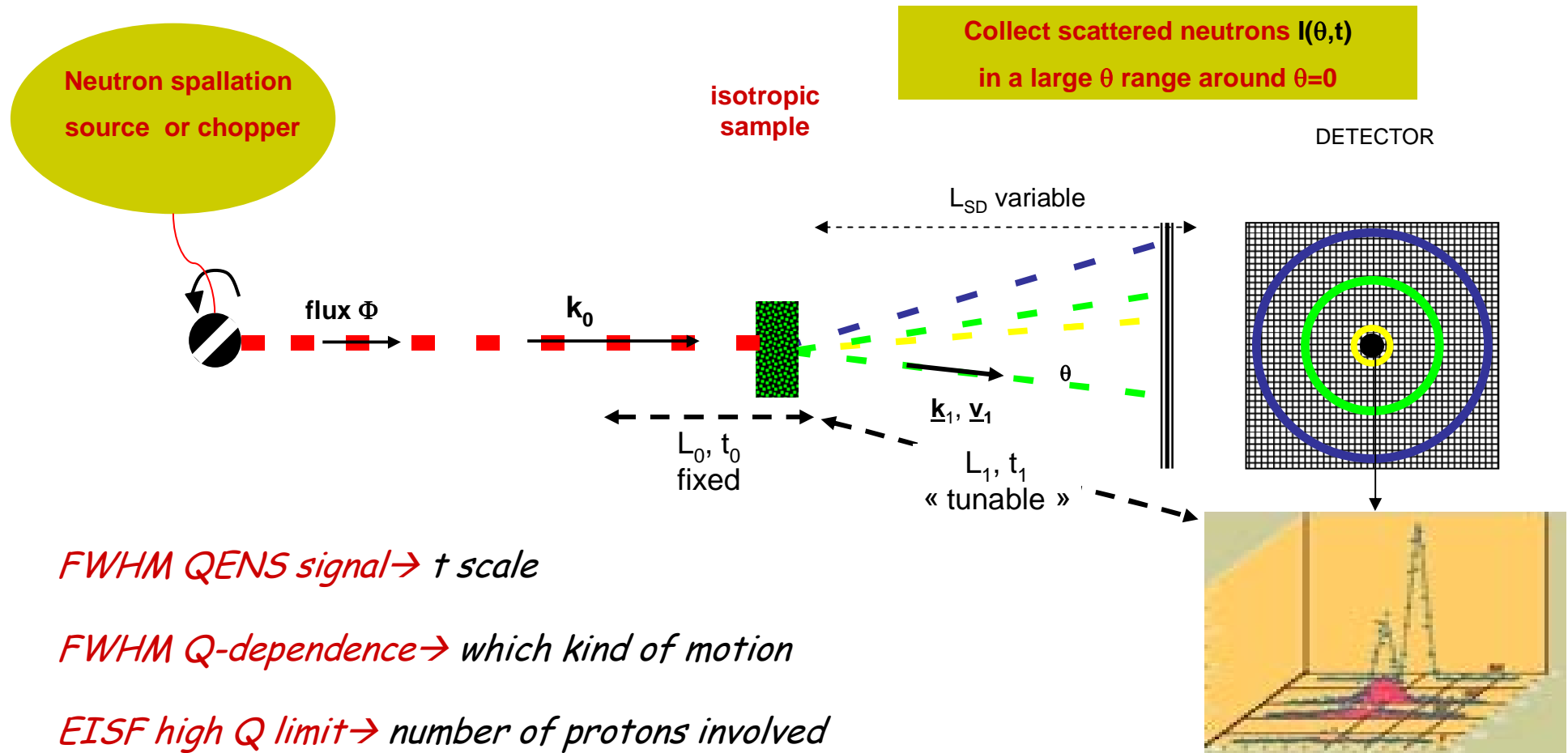


## Ice VIII→Ice Ih conversion



# Incoherent Quasi-Elastic Neutron Scattering [IRIS-ISIS]:

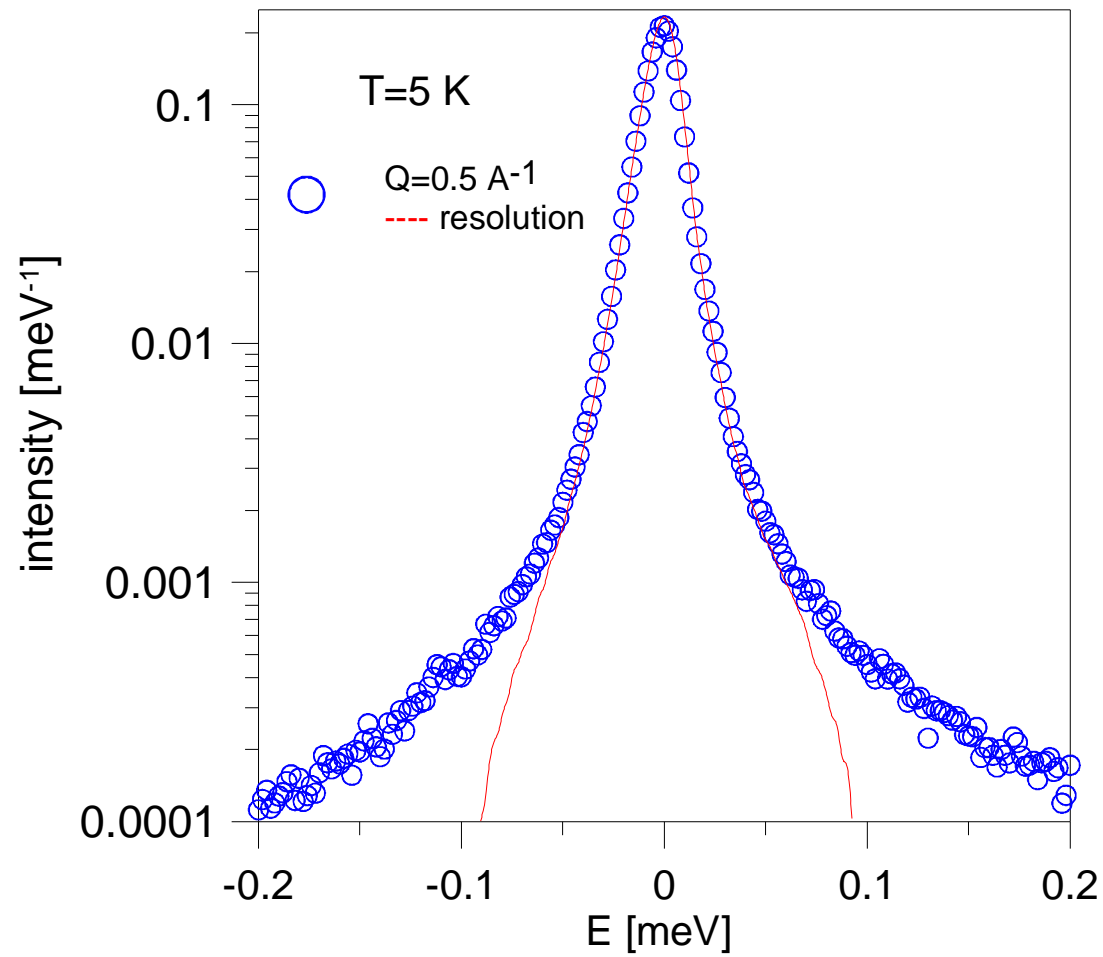
Probes motions of single proton:  $\Delta E = 15 \mu\text{eV} \rightarrow$  dynamics faster than 100 ps



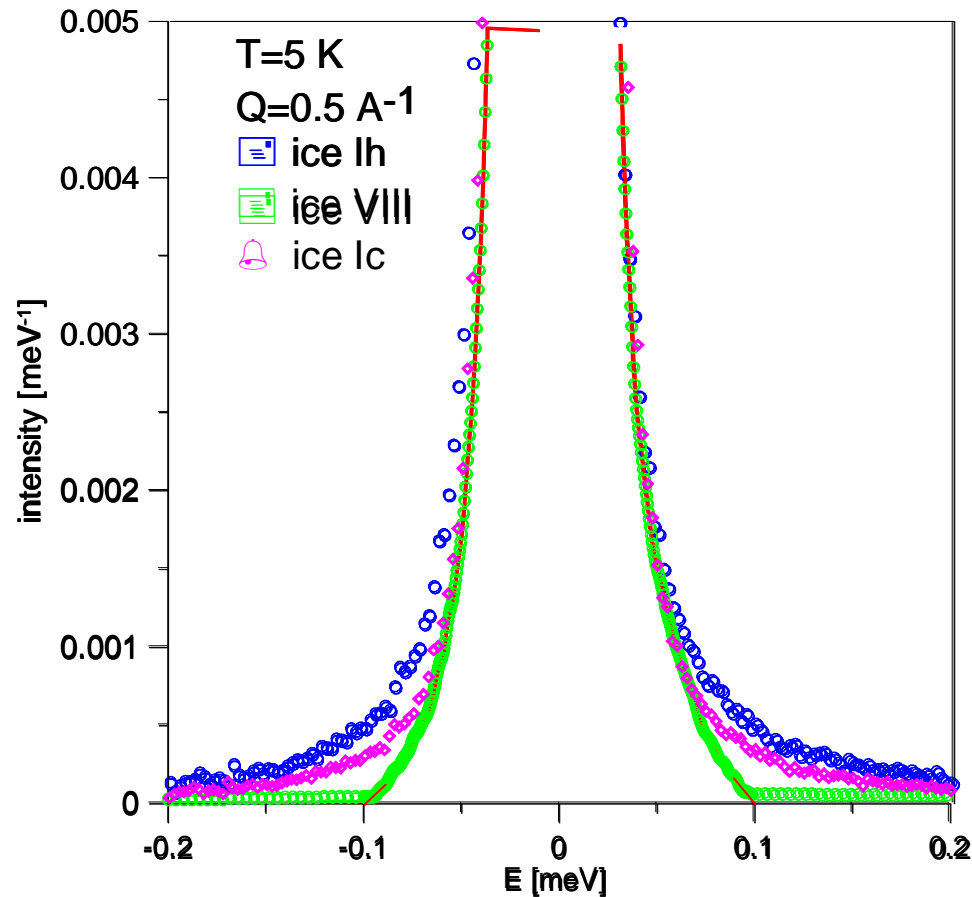
$$I(\theta, E) = \Phi N a(k_1) \frac{k_1}{k_0} \tilde{S}(QE)$$

# Quasi-elastic contribution in H-disordered ice down to 5K!

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# Quasi-elastic contribution in H-disordered ice down to 5K!



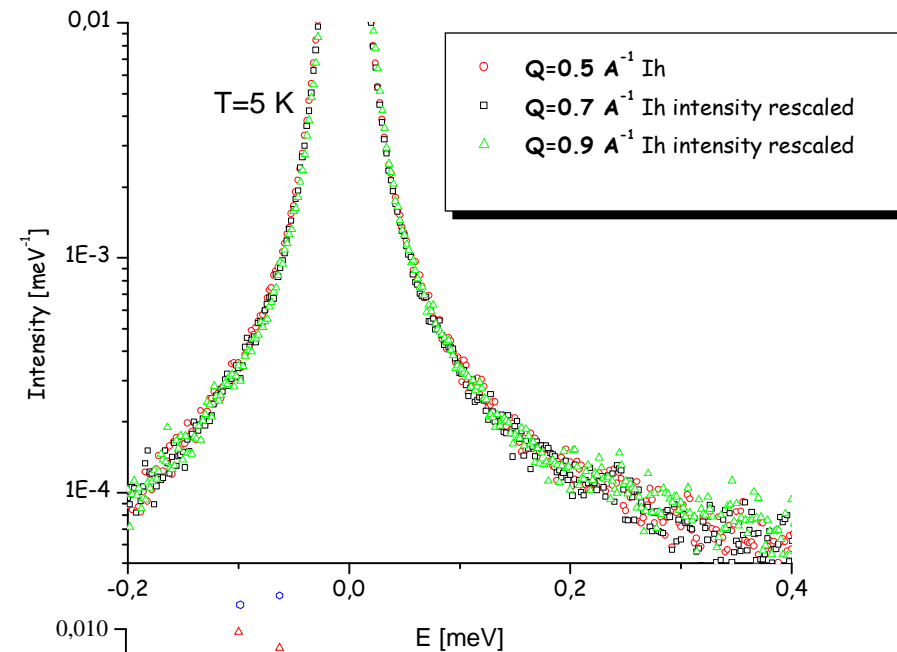
✓ Proton dynamics in ice Ih active at 5K

✓ The dynamics is absent in the H-ordered phase Ice VIII

✓ The dynamics is present in the other H-disordered phase Ice Ic

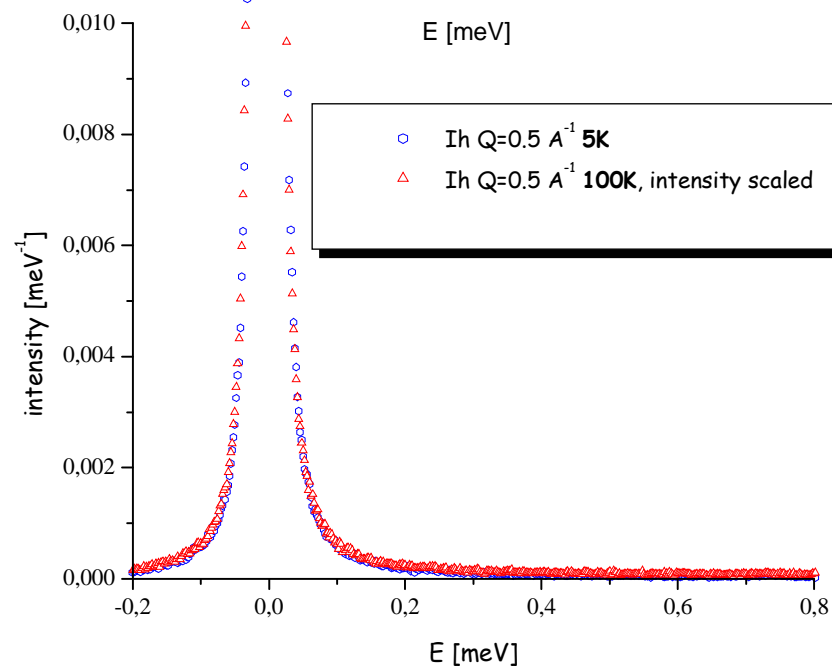
Dynamics connected to H-disordered structure!

# Temperature and wavevector characterization

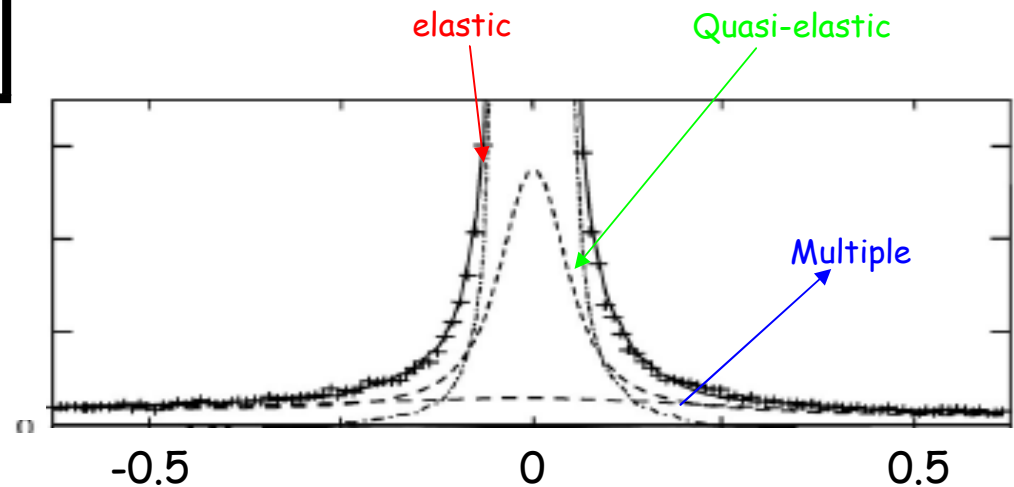


✓ T independent FWHM → non Arrhenius

✓ Q-independent FWHM : localized mode

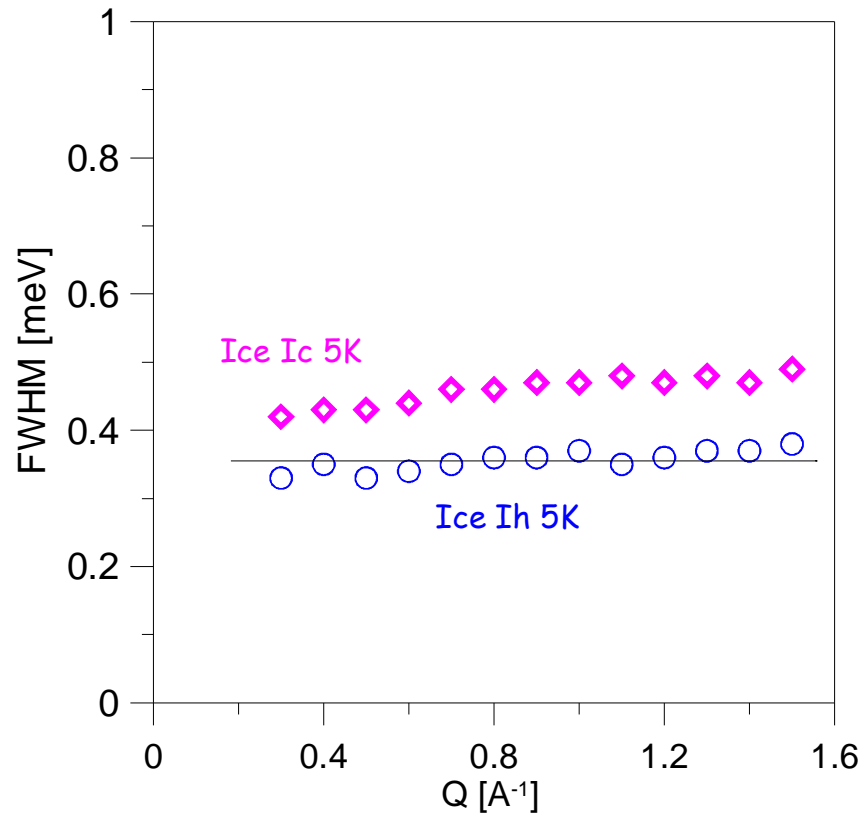


$$S(Q, \omega) = e^{-2/3 \langle \Delta x^2(T) \rangle Q^2} [S_0 \delta(\omega) + S_1 \tau_c / \pi (1 + \omega^2 \tau_c^2)]$$



## Time scales:

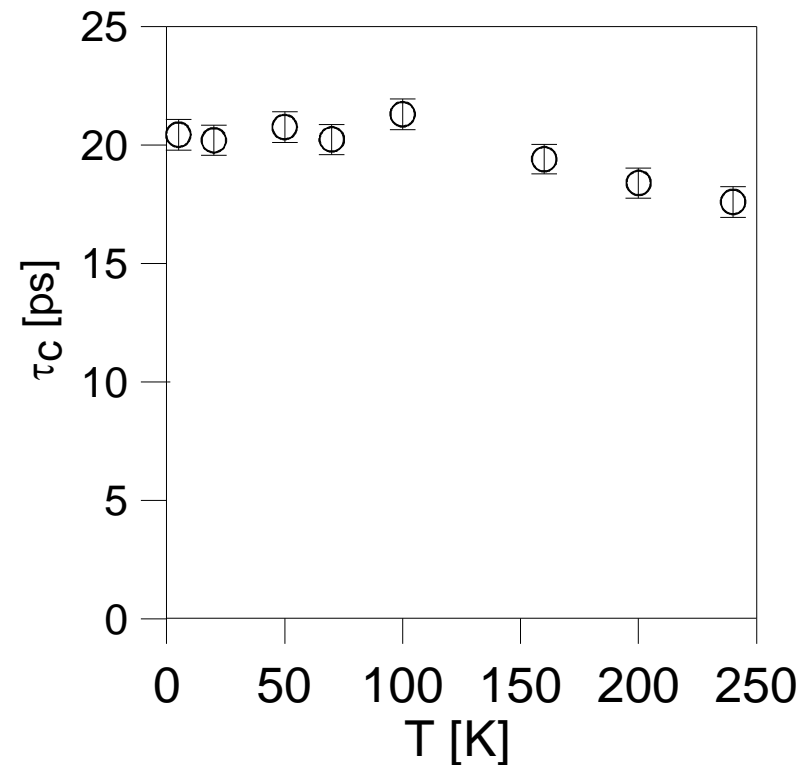
✓ Q independent FWHM



$\tau_c \sim 3.6 \pm 0.3$  ps in ice Ih  
 $\tau_c \sim 2.9 \pm 0.3$  ps in ice Ic

The inverse correlation time in the low T limit determines the hopping rate  $k_0$

✓ FWHM (T) → non Arrhenius



$k_0 \sim 2.7 \times 10^{11} \text{ s}^{-1}$  T independent  
 in 5-100 K range

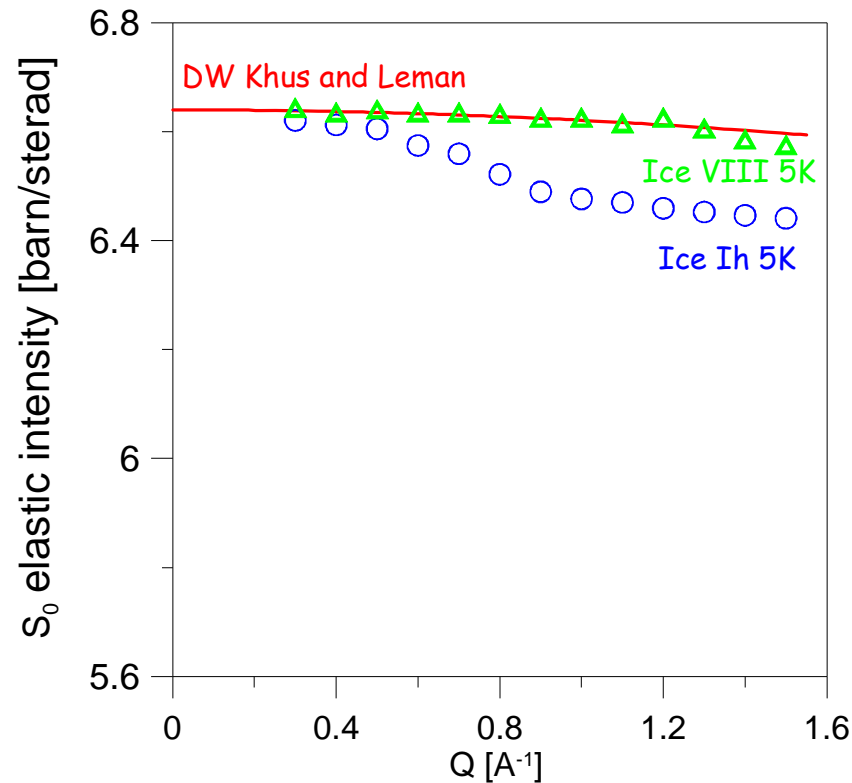
Excludes classical hopping or cage motion  
 of the proton and stepwise tunneling

L.E.B. et al, PRL 103, 165901 (2009)



# Length scales and number of protons involved:

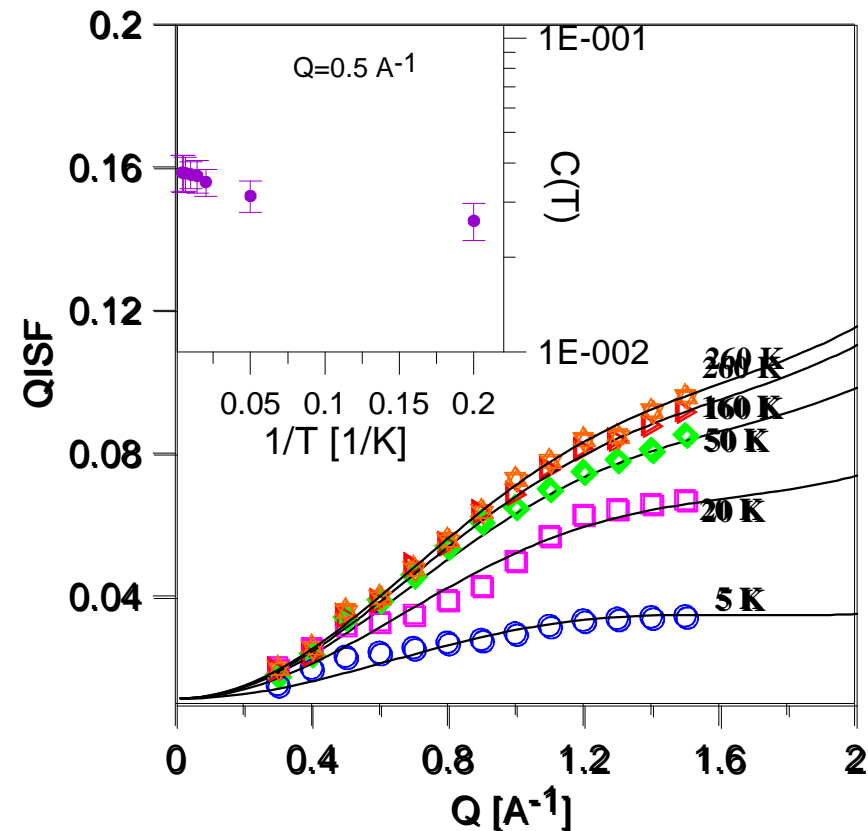
✓EISF anomalous decay



EISF: energy integrated elastic contribution

~3% H involved in ice Ih  
~2% H involved in ice Ic

✓QISF wavevector evolution

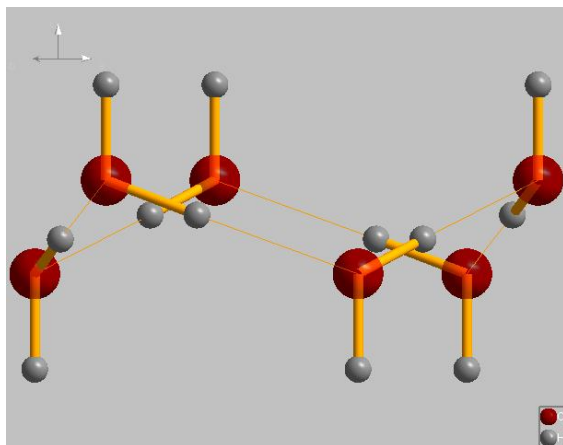
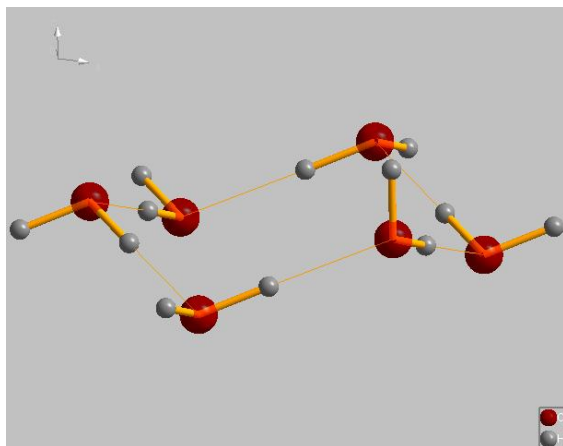
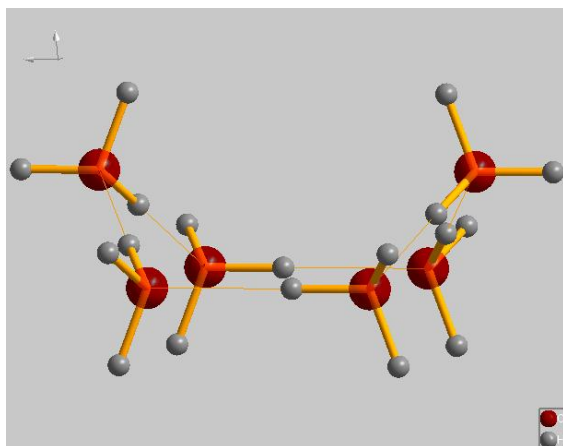


Double well model:  $2C(T)[1-\sin(Qd)/Qd]$

Low  $E_{\text{barrier}} \sim \text{meV}$

Distance  $d = 0.75 \pm 0.05 \text{ \AA}$

L.E.B. et al, PRL 103, 165901 (2009)

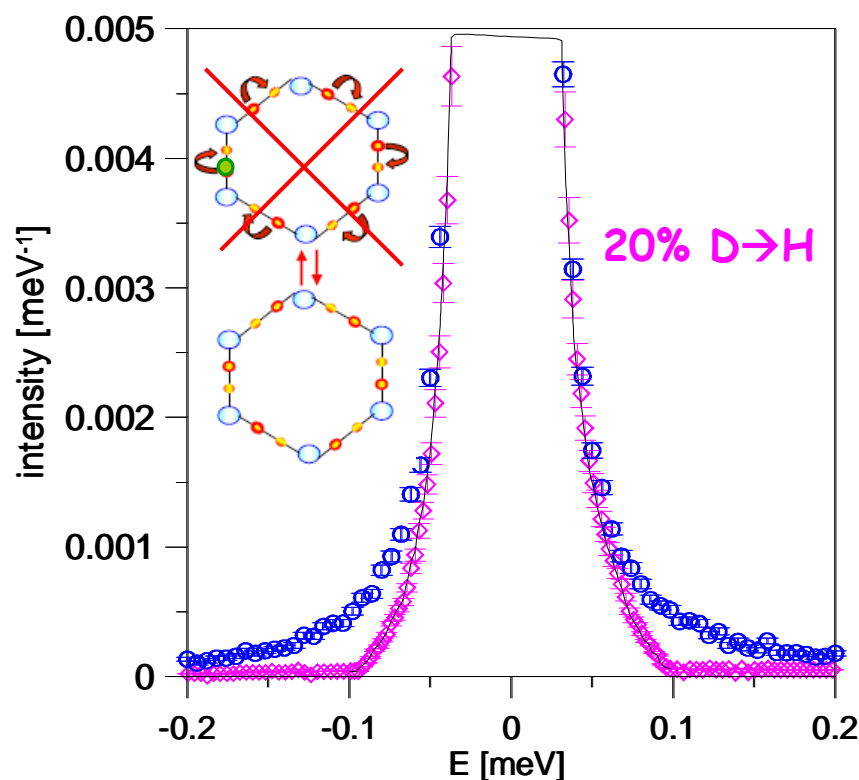


*The results obtained so far are inconsistent with any known sequential or stepwise motion of the protons, a mechanism that would have an E-barrier at least one order of magnitude higher:*

Ordered loops  $\rightarrow$   $1/32$  of total loops  $\sim$  number of H involved

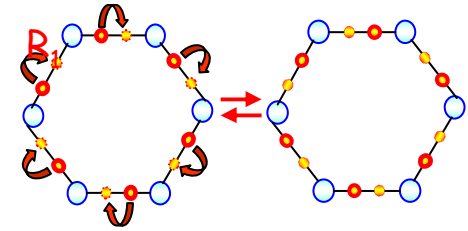
6 protons can move with no change of the total  $E_{\text{cryst}} \rightarrow$  lower E barrier

What happens if we break the symmetry?  $\rightarrow$  DEUTERATION!



# Conclusions II and future work:

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## Elastic experiment:

- ✓ Non-harmonic motion of H in ice Ih
- ✓ Main distance involved of  $0.75 \text{ \AA}$
- ✓ Number of protons involved:  $\sim 10\%$

## QENS experiment:

- ✓ time scale  $\rightarrow \sim 3 \text{ ps}$ ; localized motion  $\sim 0.75 \text{ \AA}$
- ✓ Non-Arrhenius time rate  $\rightarrow$  no classical hopping, neither cage motion
- ✓ It occurs in the proton disordered phases only  $\rightarrow$  linked to proton disorder
- ✓ High time rate ( $2.7 \cdot 10^{11} \text{ s}^{-1}$ ), low E-barrier  $\sim \text{meV}$   $\rightarrow$  quantum concerted proton tunnelling
- ✓ Disappears with partial deuteration  $\rightarrow$  most likely associated with ordered loops

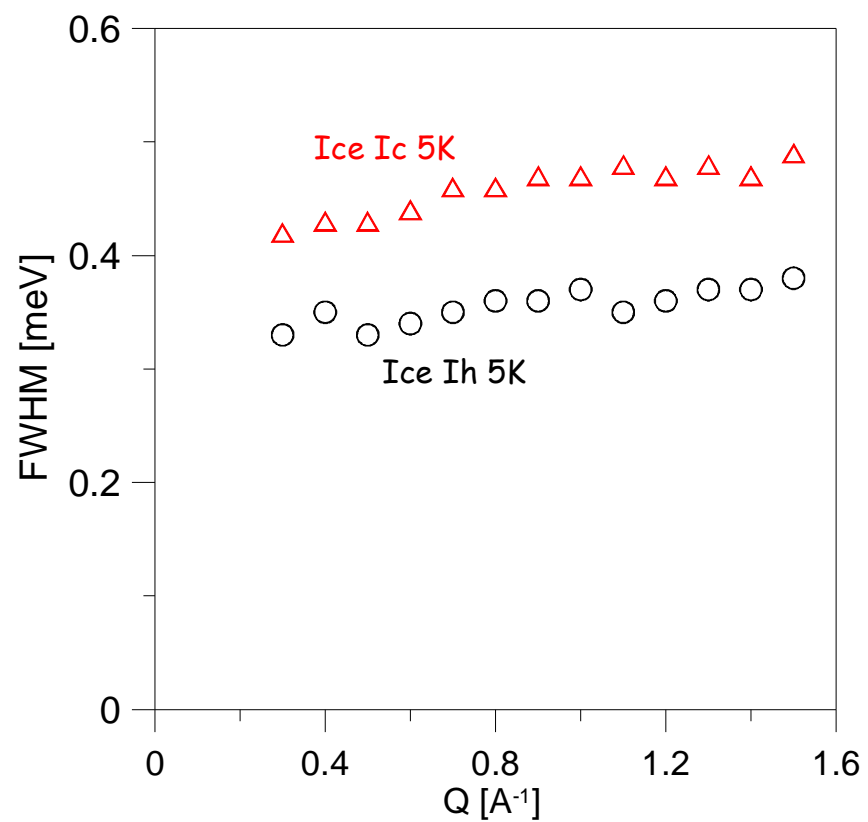
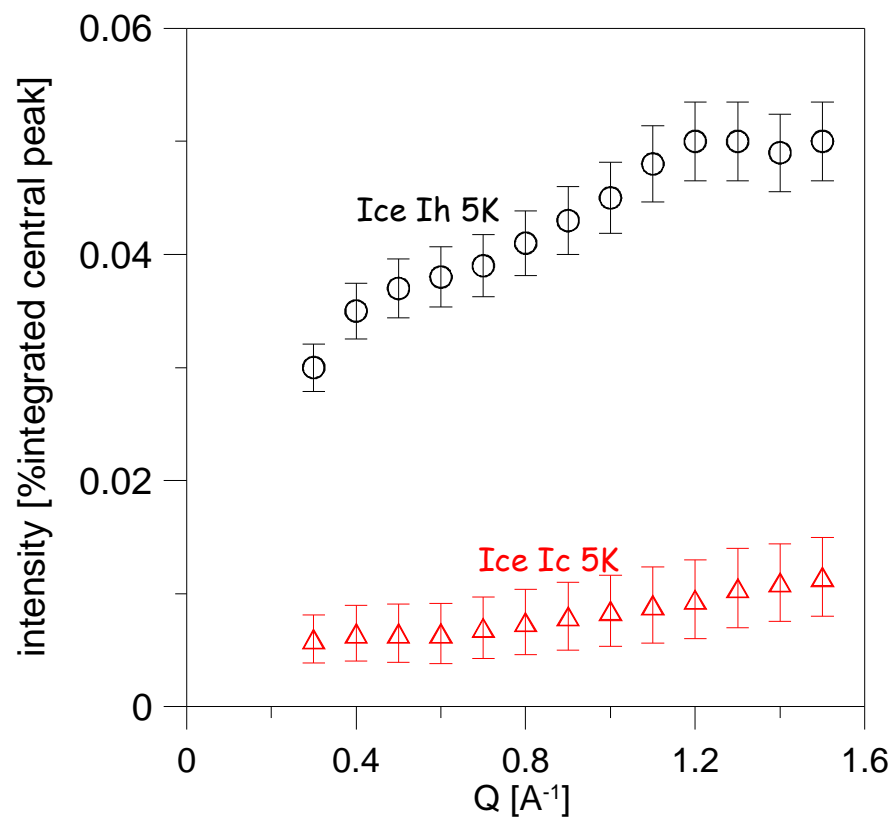
## .. Work in progress

- ✓ Detailed characterization at low temperatures (1-20 K)  $\rightarrow E_{\text{barrier}}$
- ✓ Higher resolution measurements  $\rightarrow$  check if the process is really stochastic
- ✓ Measurements on a single crystal  $\rightarrow$  better definition of the geometry

..we need full path integral calculations!!!

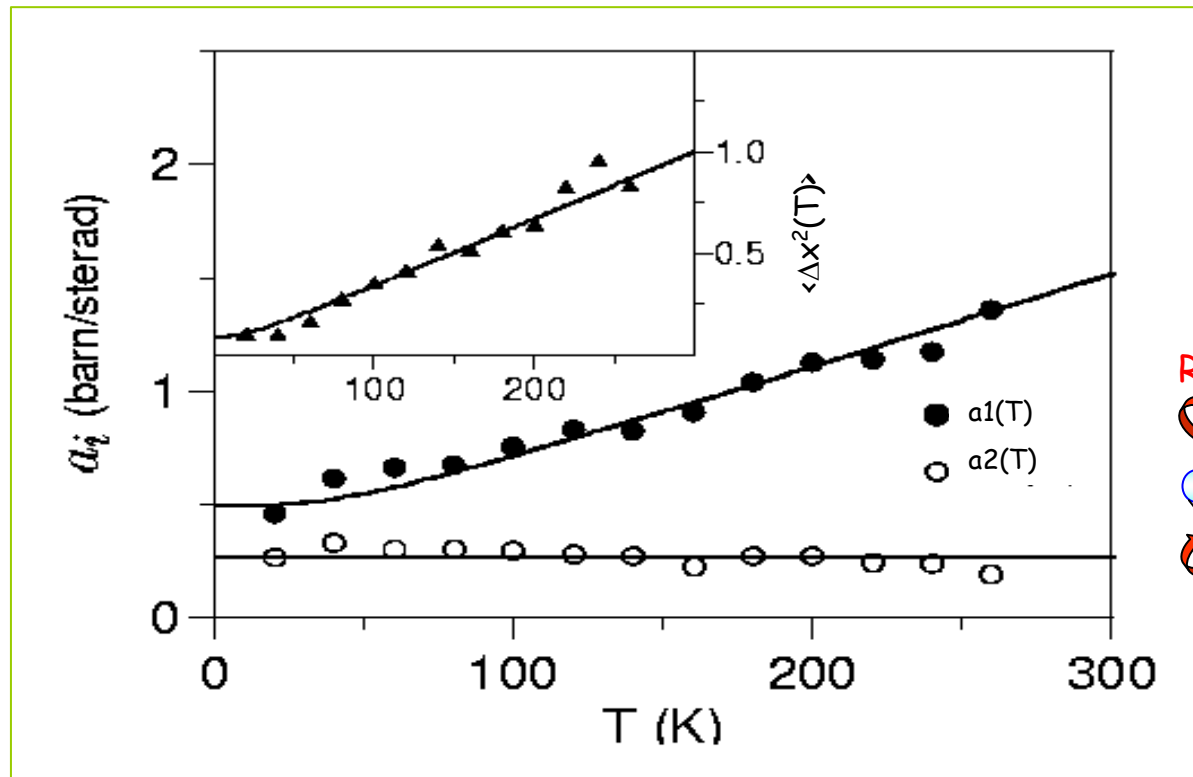
**Thank you!**

## What about ice Ic?

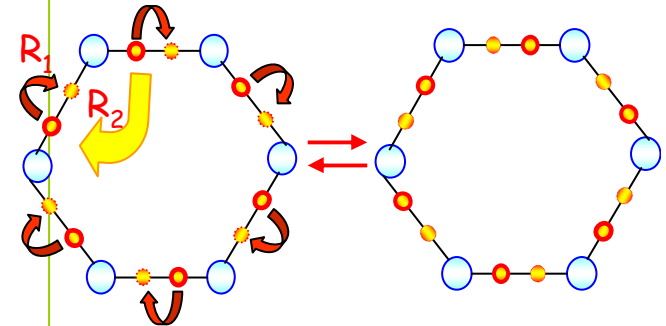


Two oscillators model:

$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{inc}}{4\pi} \left[ a_0 \exp[-\langle \Delta x^2(T) \rangle Q^2] + a_1 \cos(Q \cdot R_1) + a_2 \cos(Q \cdot R_2) \right]$$



$a_i$  proportional to  $p_a p_b$   
(equilibrium population  
of the two tautomeres)

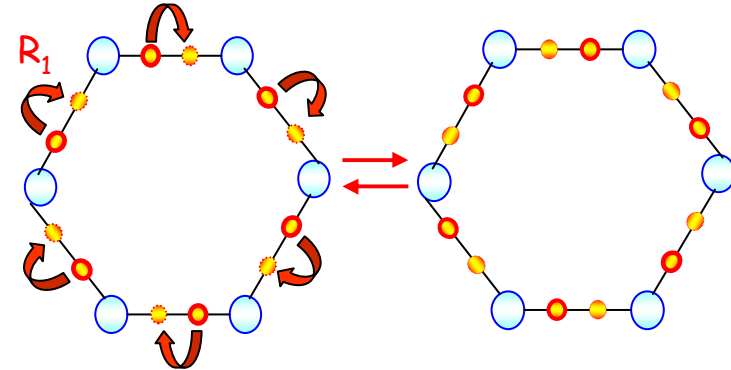
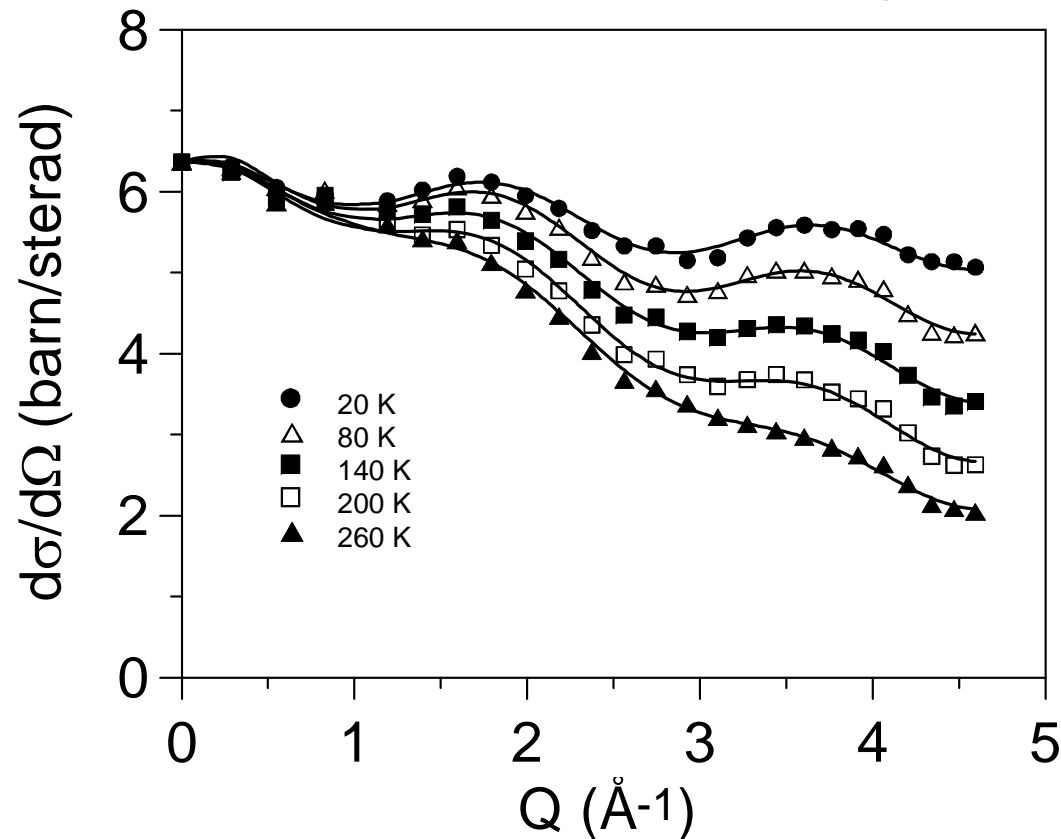


✓  $a_1$  increases with  $T$ , fitted by a Bose factor at fixed energy of  $10.3 \pm 0.5$  meV:  
phonon assisted process

✓  $a_2 \sim T$  independent: intrinsic feature of the ground-state H wavefunction

✓  $\langle \Delta x^2(T) \rangle$  in agreement with diffraction data (Kuks and Lehmann, J. Phys. C 48 (1987) )  
Fitted with 2 isolated oscillators at 10.6 and 4.3 meV

Fitting with a coordinated proton jump model

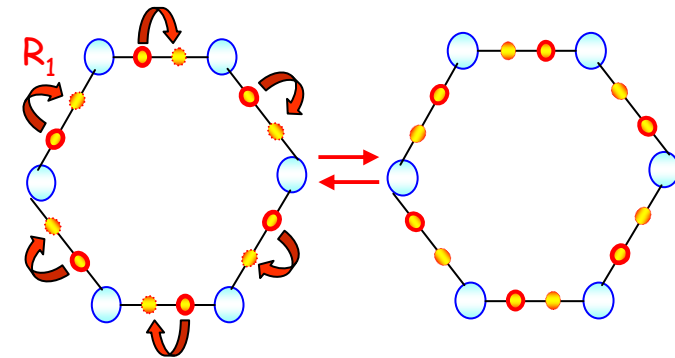
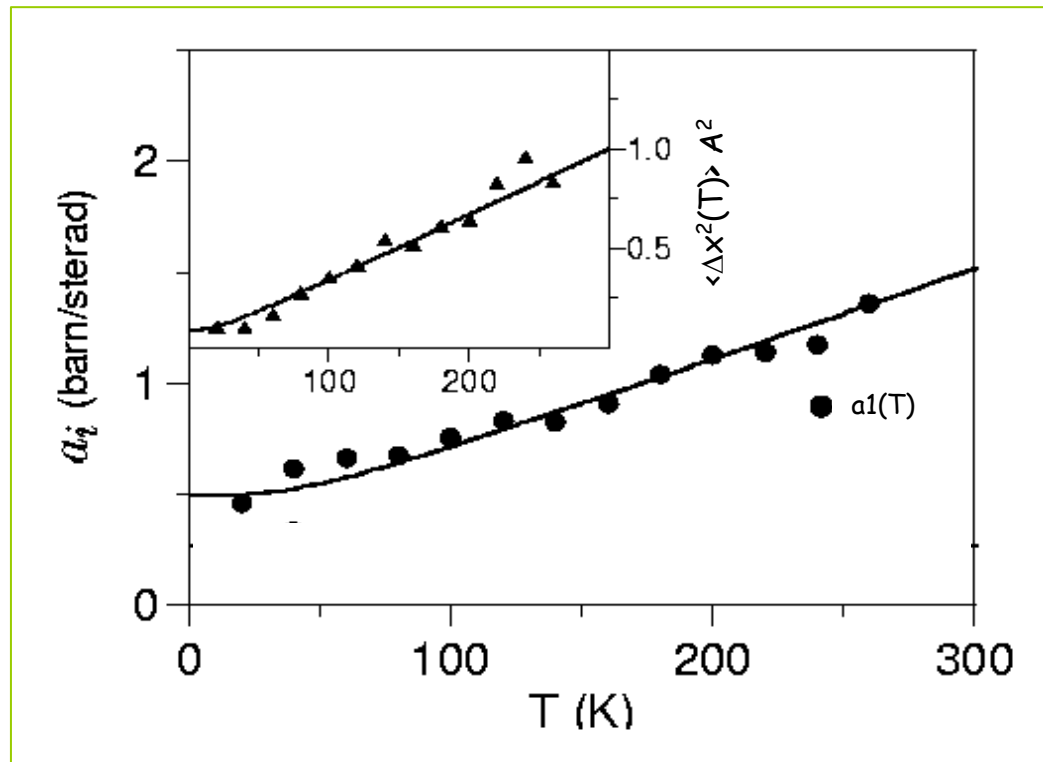


$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{inc}}{4\pi} e^{-\langle \Delta x^2 \rangle Q^2} [1 - 2p_1 p_2 (1 - \sin(Qd)/(Qd))]$$

The best fit is obtained with a distance of  $d=0.75 \pm 0.05 \text{ \AA}$  (H sites along one O-O bond  $R1=0.78 \text{ \AA}$  )

## Two oscillators model:

$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{inc}}{4\pi} e^{-\langle \Delta x^2 \rangle Q^2} [1 - 2p_1 p_2 (1 - \sin(Qd)/(Qd))]$$

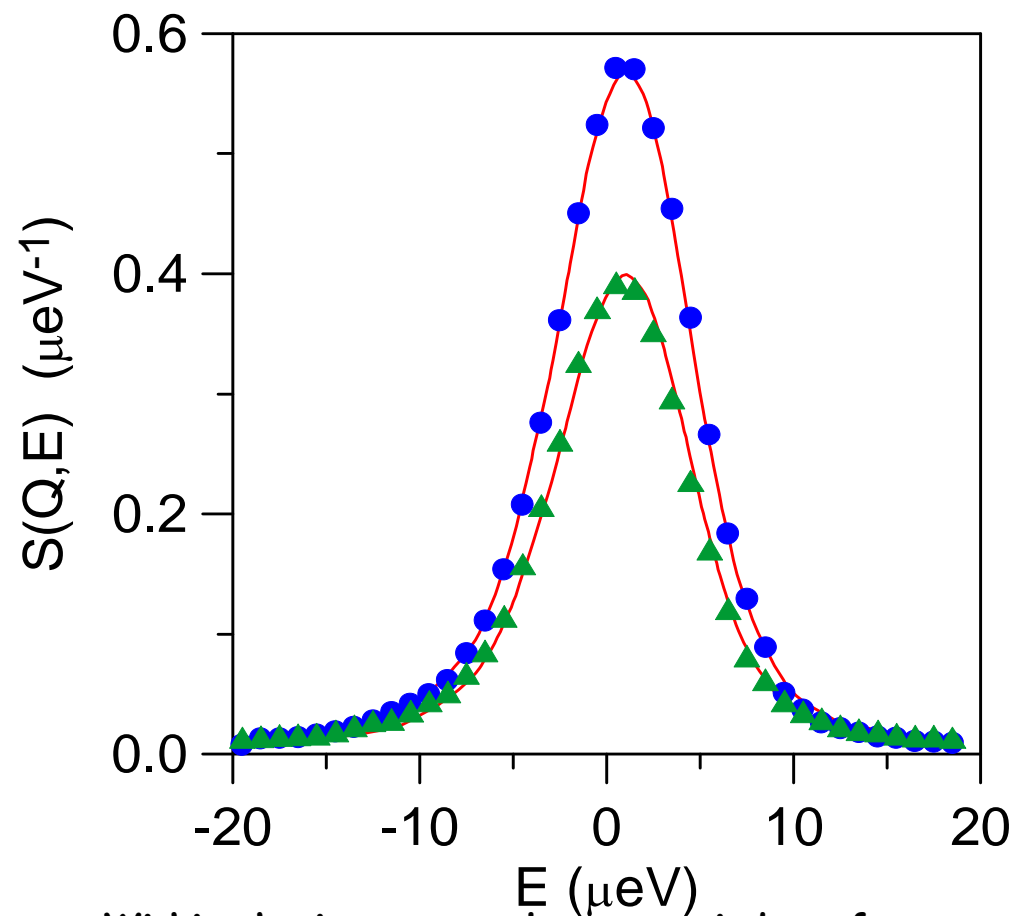


$p_1, p_2 \rightarrow$  equilibrium population of the two tautomeres

•  $a_i = 2p_1 p_2$  increases with  $T$ , fitted by a Bose factor at fixed energy of 10.3 meV: **phonon assisted process**

•  $\langle \Delta x^2(T) \rangle$  was fitted by a *two isolated oscillators* model: best fit energies  $\rightarrow$  4.3 meV and 10.6 meV in agreement with diffraction data (Khuss and Lemann,.....)





Within the instrumental energy window of  
IN13 (ILL) the scattered intensity is produced  
by an elastic process (resolution 9  $\mu\text{eV}$ )