



**The Abdus Salam
International Centre for Theoretical Physics**



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Spring School on Superstring Theory and Related Topics

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Superstring multiloop amplitudes and non-renormalization theorems Lecture I - Motivation and Summary of Results -

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Lecture Plan: Multiloop superstring amp's and non-renorm. thms

- I. Motivation and Summary of results
- II. Basics of Pure Spinor Formalism
- III. Loop amplitude prescription
- IV. Derivation of results

I. Motivation for studying superstring amp's

A. Duality conjectures

Type IIB superstrings have $SL(2, \mathbb{Z})$ duality

$$\Omega = a + i e^{-\varphi} \rightarrow \frac{A\Omega + B}{C\Omega + D} \quad AD - BC = 1$$

\uparrow \uparrow
 Ω_1 Ω_2

(B=C=-1 \Rightarrow $\varphi \rightarrow -\varphi$)

Type IIB effective action: $\int d^{10}x \sqrt{-g} \Omega_2^{-1/2} Z_{3/2}(\Omega, \bar{\Omega}) R^4 + \dots$

from string gauge

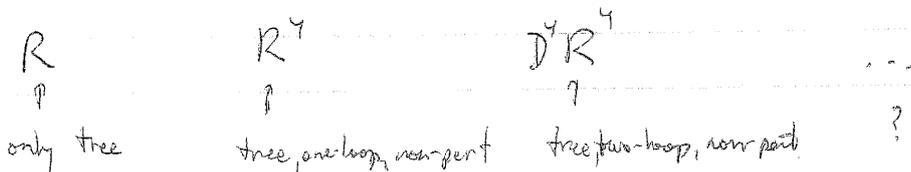
$$Z_{3/2} = \sum_{(m,n) \neq (0,0)} \frac{\Omega_2^s}{|m+n\Omega|^{2s}} = c_{1,s} \Omega_2^s + c_{2,s} \Omega_2^{1-s} + \sum_{k \neq 0} d_{k,s} e^{-2\pi i(k\Omega_2 - ik\Omega_1)}$$

$\dots = \int d^{10}x \sqrt{-g} \Omega_2^{-1/2} Z_{5/2}(\Omega, \bar{\Omega}) D^4 R^4 + \text{more complicated } SL(2, \mathbb{Z})\text{-inv. terms.}$

(Einstein gauge $\int d^{10}x \sqrt{-g} R \rightarrow \int d^{10}x \sqrt{-g} \Omega_2^2 R$ String gauge.)

$g \rightarrow g e^{-\varphi/2} = g e^{-\frac{\varphi}{4}} e^{-\frac{\varphi}{4}}$

Green, Russo, Vanhove 0610299



(Also have Heterotic-Type I duality \Rightarrow structure of $\text{Tr}(F^N)$ from Born-Infeld in Type I)

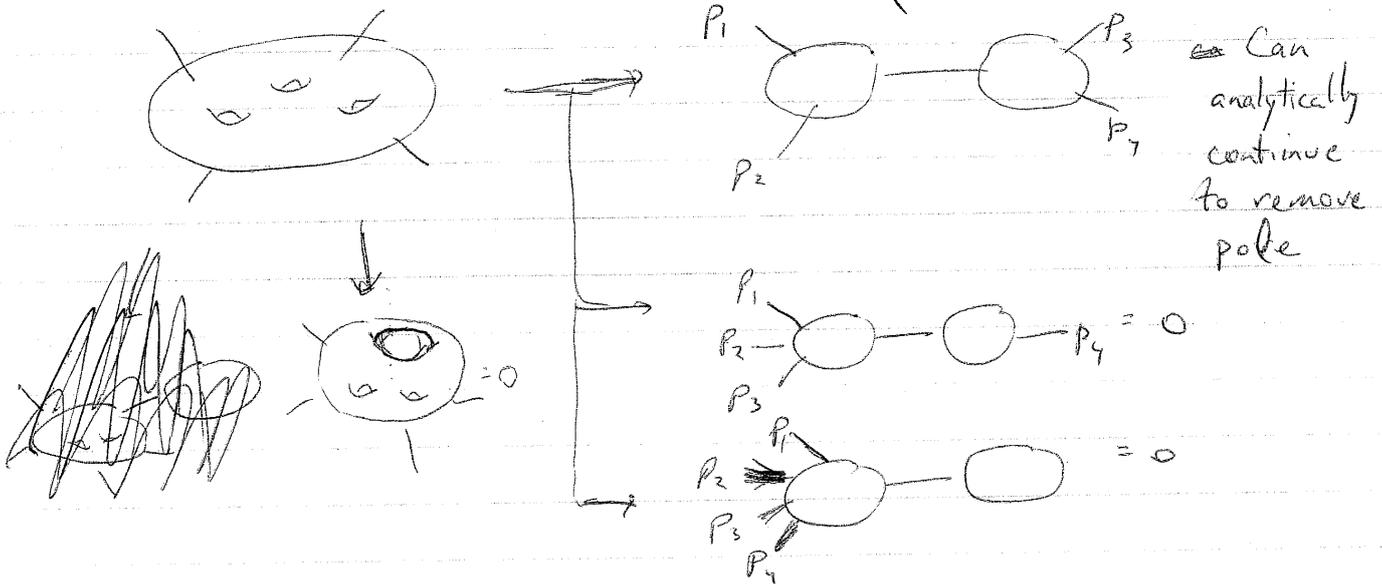
B. Finiteness conjectures

1. No unphysical poles in superstring amplitudes

\Rightarrow 0-pt, 1-pt amp's get no ~~tree~~ contributions (no cosm. constant and no tadpole)

2-pt massless amp. gets no loop contributions

(no mass shift \neq of massless particles)



2. 3-pt massless amp. gets no loop contributions

\Rightarrow coupling constant not renormalized

3. Finiteness of $N=4$ $d=8$ sugra

Dimensional analysis implies that $\int d^d x \sqrt{g} e^{(2h-2)\varphi} D^k R^4$ ^{loop}

$$\text{dim as } \Lambda^{(d-2)(h-1) - k - 8 + d} = \Lambda^{(d-2)h - 6 - k}$$

If $k \geq 2h$, no divergence when $\Lambda \rightarrow \infty$ if $d=4$.

Exp. evidence: $h=2 \Rightarrow k \geq 4$, $h=3 \Rightarrow k \geq 6$, $h=4 \Rightarrow k \geq 8$

(3)

4. Finiteness of ~~$N=4$~~ $N=4$ super-YM above $d=4$

$$\int dx e^{-\frac{1}{\Lambda^4} \text{Tr}(F^2)} + \int dx e^{(h-1)\frac{1}{\Lambda^4} D^k \text{Tr}(F^4)} + \int dx e^{(h-1)\frac{1}{\Lambda^4} D^{2k} (\text{Tr}(F^2))^2}$$

~~diverges~~ \downarrow ~~term~~
 diverges like $\Lambda^{(h-1)(d-4) - 8 - k + d} = \Lambda^{(d-4)h - 4 - k}$

Corrections proportional to $\left(\frac{1}{\Lambda^4} F^4\right)$ and $\left(\frac{1}{\Lambda^4} R^4\right)$ because of maximal susy



$\begin{matrix} a & b & c & d \\ \text{---} & \text{---} & \text{---} & \text{---} \\ A & & A & \end{matrix} \Rightarrow$ polarization dependence is fixed by A
 \downarrow
 massive has $2^7 \times 2^4 = 2^{11}$ comp's (or $2^6 \times 2^5 = 2^{11}$ comp's)
 (Albay + Maldacena)

C. Superstring methods

1. RNS

Multiloop amplitudes are complicated since need to sum over spin structures \Rightarrow spacetime susy is not manifest.

Explicit computations up to 4-point NS-NS scattering at 2-loops (d'Hoker, Phong; Zhu, et al.) \Rightarrow finiteness up to 2-loops

Difficult to prove non-renorm. theorems beyond 2-loops.

2. GS

Amplitudes can be computed in light-cone gauge. No sum over spin structures, but need "interaction pt. operators" and contact terms. Computations only up to 4-pt. one-loop (Green-Schwarz),

3. Pure Spinor formalism

In ~~flat background~~ Can compute loop amplitudes with manifest $d=10$ Poincare inv. and spacetime susy.

Explicit computations up to ~~5pt~~ 4-pt two-loop with NS or R external states.

Normalization of two-loop agrees with duality conjecture (Mafra, Gomez) (cannot be checked directly in RNS)

Can prove various non-renorm. theorem to all loops

D. Results

Can prove (assuming factorization) that no ~~loop~~ contribution at any loop to 0-pt, 1-pt, ~~2pt~~ 2-pt, 3-pt massless amplitudes

Can prove that first term which can contribute at h-loops is

$h=1$	2	3	4	5	6	...
R^4	$D^4 R^4$	$D^6 R^4$	$D^8 R^4$	$D^8 R^4$	$D^8 R^4$...
$\text{Tr}(F^4)$	$D^2 \text{Tr}(F^4)$	$D^2 \text{Tr}(F^4)$	$D^2 \text{Tr}(F^4)$...		

$$(\text{Tr}(F^2))^2 \quad D^2(\text{Tr}(F^2)^2) \quad D^4(\text{Tr}(F^2)^2) \quad D^4(\text{Tr}(F^2)^2) \quad D^4(\text{Tr}(F^2)^2) \quad \dots$$

\Rightarrow Divergence at h-loops when $d \geq d_c$ where $d_c =$

	$h=1$	2	3	4	5	6	7	Suggests $N=4$ $d=8$ Sugra diverges at 7 loops
Sugra	8	7	6	$1/2$	$24/5$	$13/3$	4	
$\text{Tr}(F^4)$	8	7	6	$1/2$	$24/5$	5	$34/7$	
$(\text{Tr}(F^2))^2$	8	7	$20/3$	6	$28/5$	$16/3$	$36/7$	

D-terms vs. F-terms