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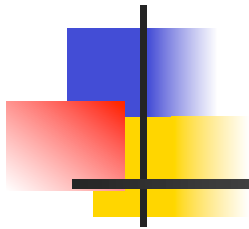
**School on Synchrotron and Free-Electron-Laser Sources and their  
Multidisciplinary Applications**

*26 April - 7 May, 2010*

**Terahertz Spectroscopy**

Stefano Lupi  
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Italy*

# Terahertz Spectroscopy



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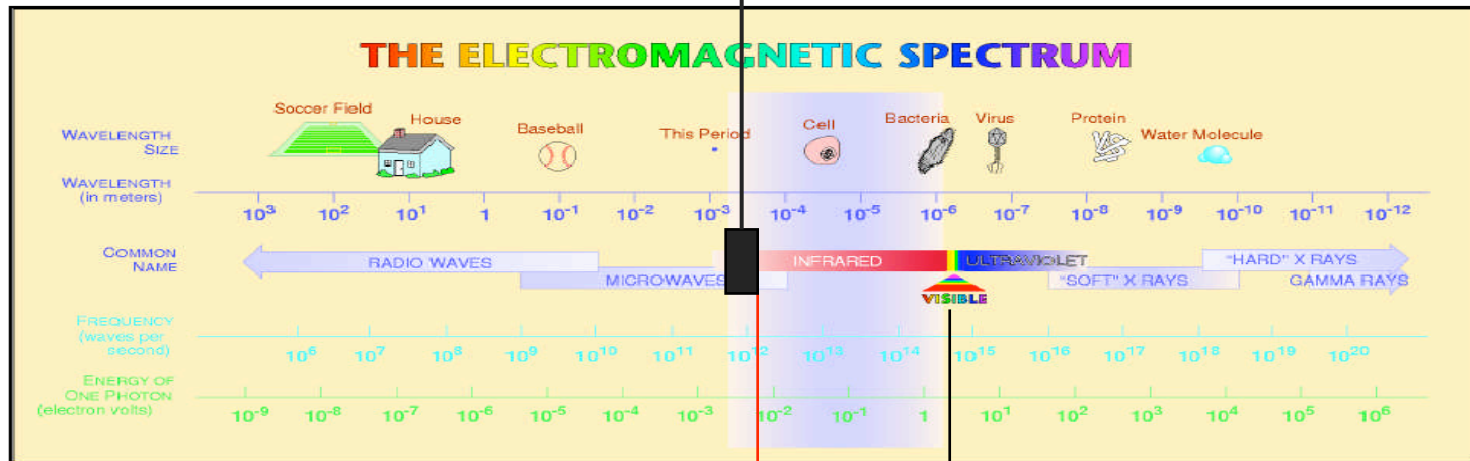
# Outline

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- THz gap
- Overview of existing and planned THz sources
- THz Science

# Electromagnetic Spectrum

The “**THz**gap”, Collective Excitations in Macromolecules and exotic electronic materials



IR Units:  $200 \text{ cm}^{-1} = 300 \text{ K} = 25 \text{ meV} = 50 \mu\text{m} = 7 \text{ THz}$

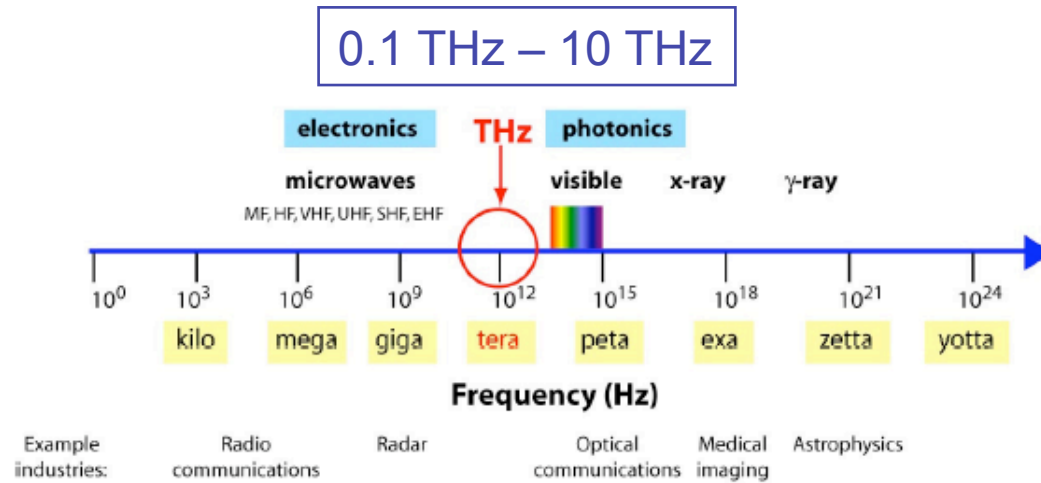
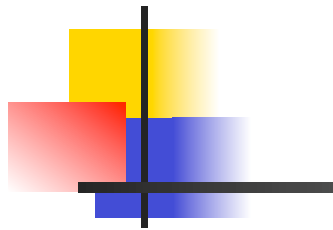
FIR MIR NIR

Phonons;  
Drude absorption;  
Gaps in superconductors;  
Molecular Rotations;

Molecular Vibrations  
**Fingerprints** for  
Chemistry, Biology,  
And Geology

Molecular Overtones and  
Combinations bands;  
Excitons;  
Gaps in semiconductors

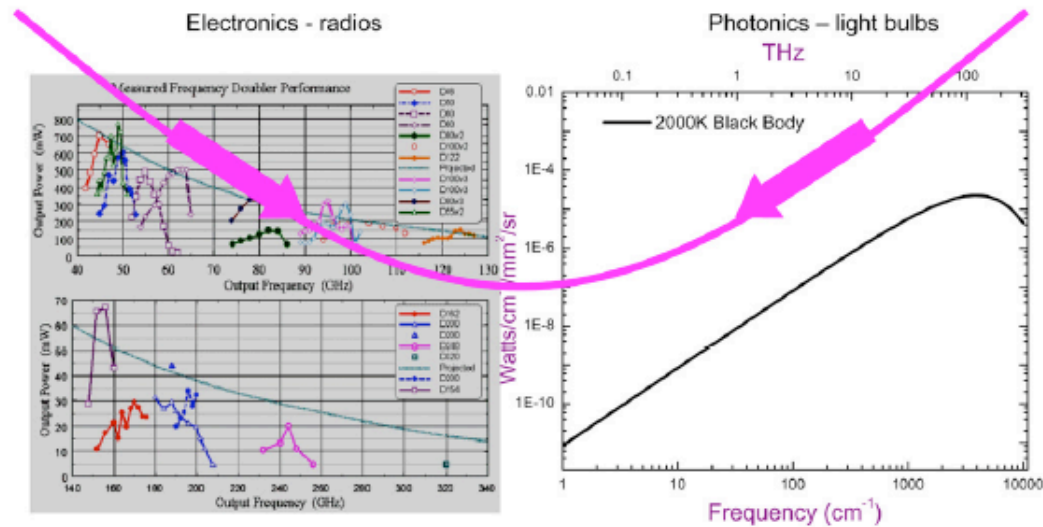
# The Terahertz gap



**1 THz ~ 1 ps ~ 300  $\mu$ m ~ 33  $\text{cm}^{-1}$  ~ 4.1 meV ~ 47.6 $^\circ$ K**

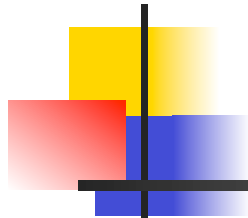
**Figure 1.** Schematic of the electromagnetic spectrum showing that THz light lies between electronics and photonics.

**No electronics, few microwaves generators**

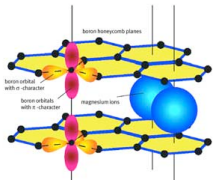


**Vanishing thermal power, few tunable and pulsed lasers.**

# THz Science



## Condensed Matter Physics



### Superconductivity

- Energy gap
- Symmetry of the order parameter
- Direct determination of the superfluid density
- Dynamics of Cooper pairs

### Low-dimensional effects

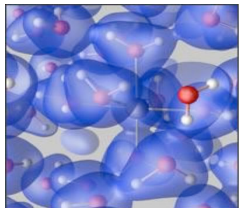
- Dimensionality crossover
- Non-Fermi liquid normal states
- Broken symmetry ground states

### Strongly correlated electrons

- Kondo problem
- Heavy fermions

### Magnetic Dynamics

## Physical and Analytical Chemistry



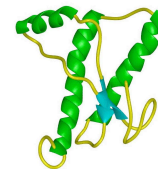
### Polar liquids

- Hydrogen bond
- Van der Waals interactions
- Acoustic-Optic phonon mixing in water

### Solutions

- Static and dynamic interactions between solvated ions and solvent

## Life Sciences



### Macromolecules conformation

- Secondary and tertiary structure
- Coherent dynamic development

### Metabolism

- Water distribution
- Ionic channel dynamics

### Imaging

- 3D tomography of dry tissues
- Near-field sub-wavelength spatial resolution

## New Technologies

### THz technologies

- Array THz detectors
- Metamaterials

### Medical diagnostic

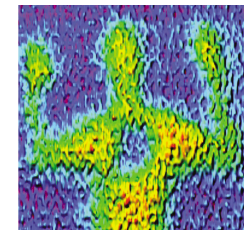
- Skin cancer detection

### Industrial production

- Material inspection
- Production line monitoring

### Defense industry/Homeland security

- Detection of explosives and biohazards





# Available THz sources

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## Terahertz Lasers

- Quantum Cascade Lasers (>2 THz),
- Si- and Ge-lasers (>1.5 THz),
- Far-Infrared FELs: FELIX, FELBE,
- Gas-based lasers (only at some given frequency),

## Sources based on electron acceleration

### Narrow Band

- Stanford (>3 THz),
- ENEA compact FEL (<0.1 THz),
- Far-IR undulator @ FLASH (Hamburg),
- Backward-wave oscillators,

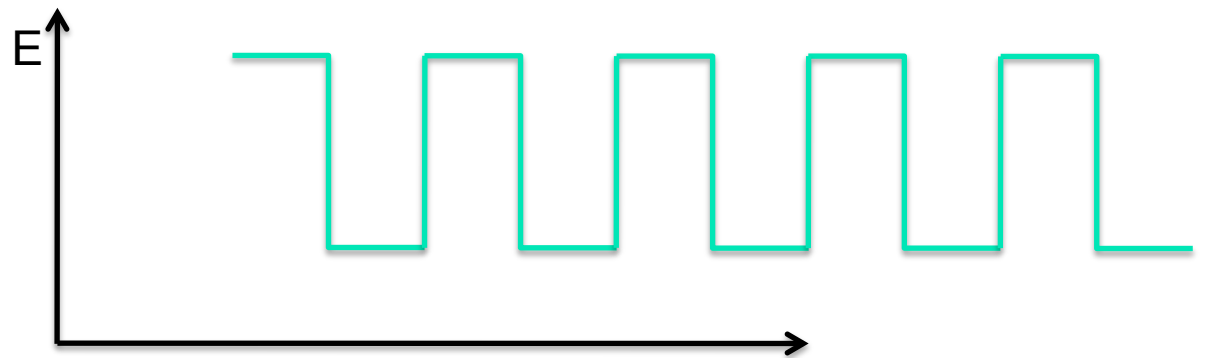
### Broad Band

- Laser Amplifiers: plasma and non linear crystals,
- Coherent Synchrotron Radiation @ III Generation Machines,
- LINAC-based Coherent Radiation Sources,

# Quantum Cascade Laser

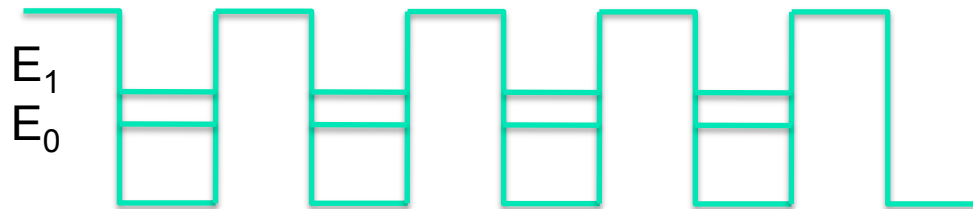
Federico Capasso, Bell Labs

Quantum Wells are periodic potential wells built by semiconductor multilayers



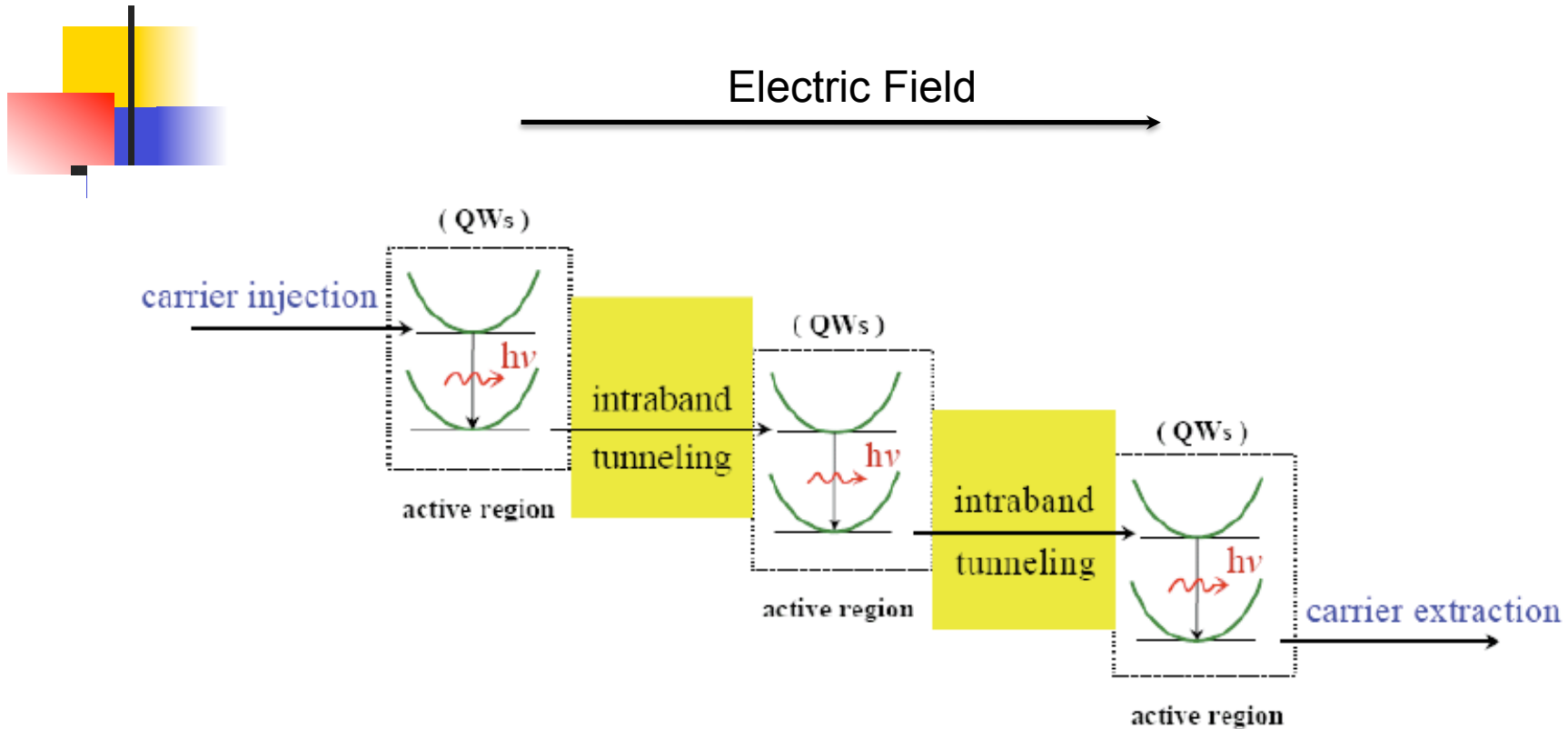
Grown direction of semiconductor multilayers

Through electrons or holes doping electronic states are created in the wells



The energy difference  $E_1 - E_0$  can be engineered from the THz to the mid-IR region

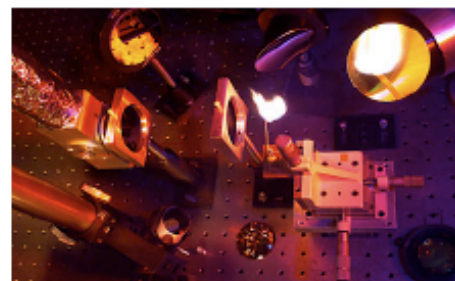




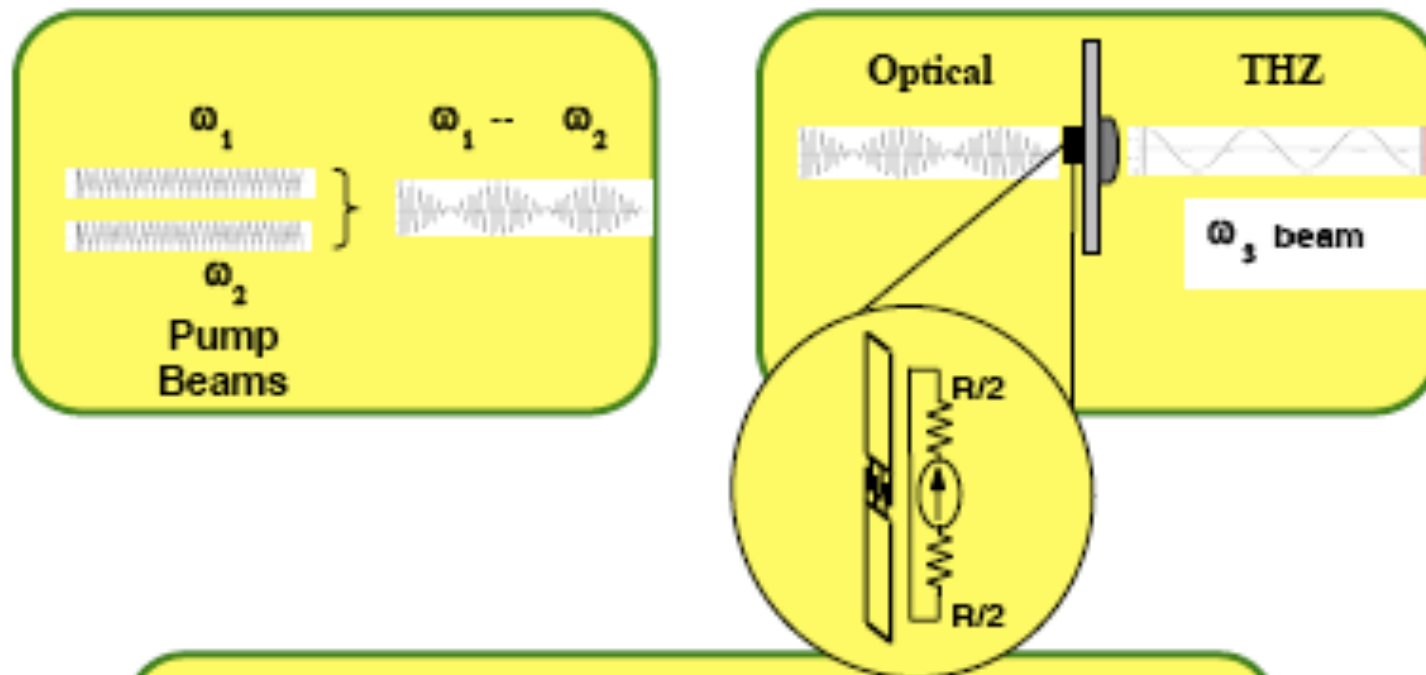
The relaxation from the excited state  $E_1$  to the ground state  $E_0$  produces in-phase photons which are amplified due to the optical cavity effect induced by the different semiconductors constituting the layers

# What makes the QC-laser special?

- Wavelength agility
  - layer thicknesses determine emission wavelength
- Demonstrated applications in mid/far-IR gas sensing
- High optical power  $\sim 1\text{W}$ , room-T operation
  - cascading re-uses electrons
- Ultra-fast carrier dynamics
  - no relaxation oscillations
- Pure TM-polarization – efficient in-plane light coupling
  - Micro-lasers
- Small linewidth enhancement factor
- Intrinsic “design potential”



## CW THz Generation by Photomixing or Heterodyning

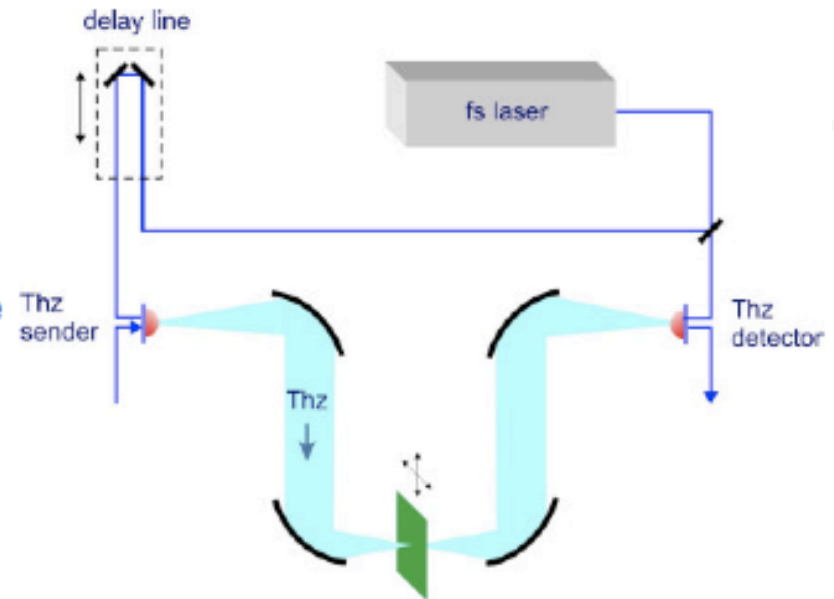


$$P_3 = \frac{R}{2} \eta_1 \lambda_1 \eta_2 \lambda_2 \left( \frac{e \hbar}{hc} \right)^2 \frac{P_1 P_2}{\left[ 1 + (\omega_3 \tau)^2 \right] \left[ 1 + (\omega_3 RC)^2 \right]}$$

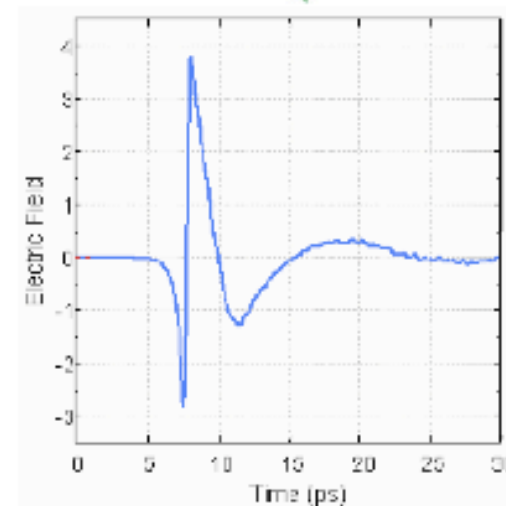
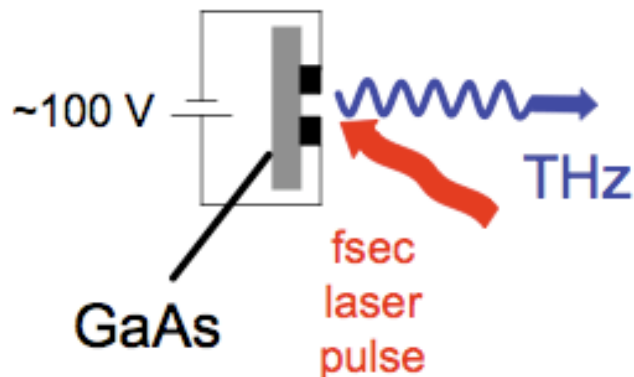
# Table-Top Laser Amplifier

## Time-domain and Scanning Techniques

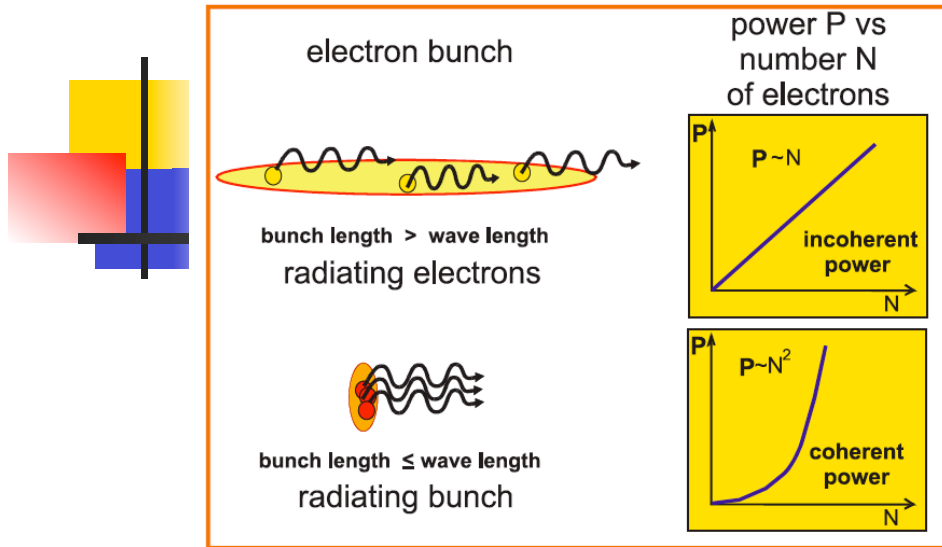
- Based on time domain spectroscopy (TDS)
- Uses recent advances in fs-laser technology.
- Measures intensity and phase of E-field vs time of a single cycle THz pulse
- Typically 0.2 – 2.5 THz, depending on fs laser



## Auston switch

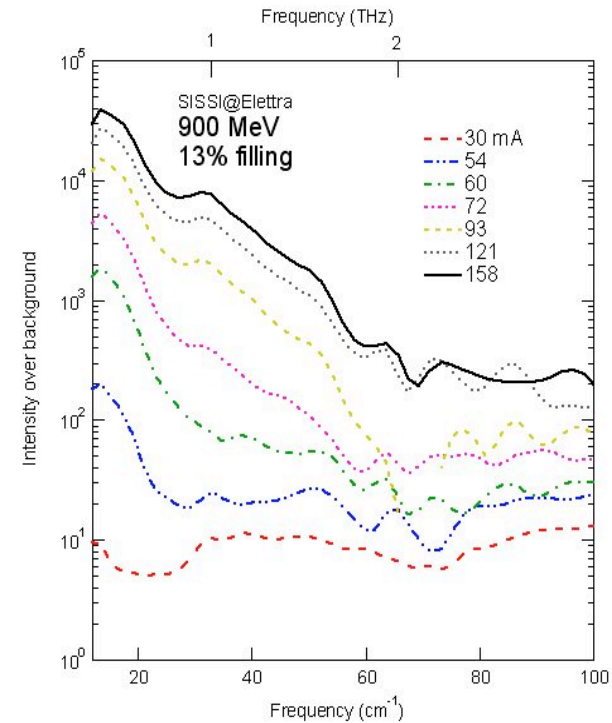
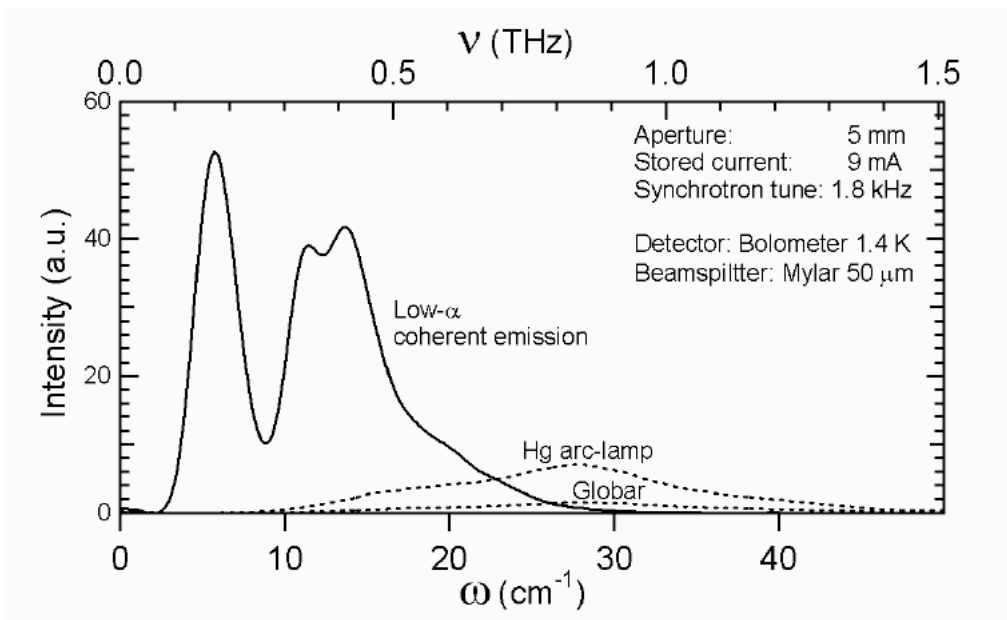


# Coherent Synchrotron Radiation (CSR)



$$I = I_{incoh} + I_{coh} = (N(1 - f_v) + N^2 f_v) I_{incoh}$$

$$f_v = \left| \int n(z) e^{i\pi \cos(\theta) z} dz \right|^2$$



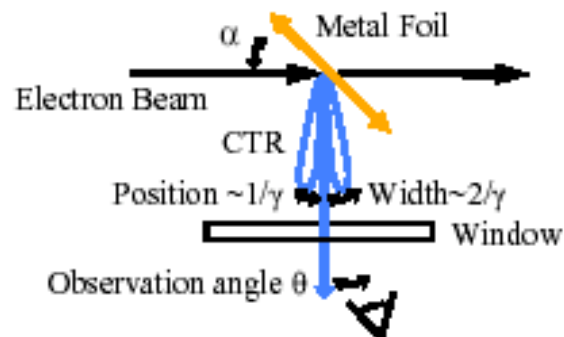
**low-e beam energy**

SISSI@Elettra: E. Karanzoulis, A. Perucchi, S.L et al, ,2007

**Low- $\alpha$  mode** IRIS@Bessy-II: U. Schade et al, PRL 2003

# Transition THz Radiation (CDR/CTR)

Transition Radiation occurs when an electron crosses the boundary between two different media



$N_e$  = number of electrons in single bunch  
 $\sigma$  = bunch-length

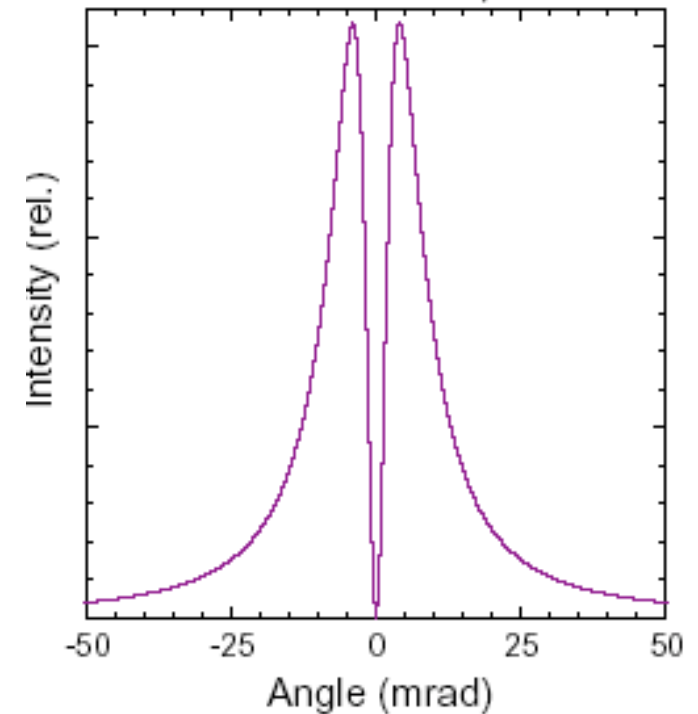
- Coherent Transition Radiation

$$\frac{dI_{CTR}}{d\omega d\Omega} = \frac{e^2 \beta^2}{\pi^2 c} \left( \frac{\sin \theta - \beta \cos \alpha}{(1 - \beta \sin \theta \cos \alpha)^2 - \beta^2 \cos^2 \theta} \right)^2 N_e \left[ 1 + N_e \exp\left[-\frac{\sigma^2 \omega^2}{c^2}\right] \right]$$

- Position, width of maxima
- Coherence

Intensity is 0 on axis and peaked at  $\Theta \sim 1/\gamma$   
 Polarization is radial

Far field distribution for  $\gamma = 200$

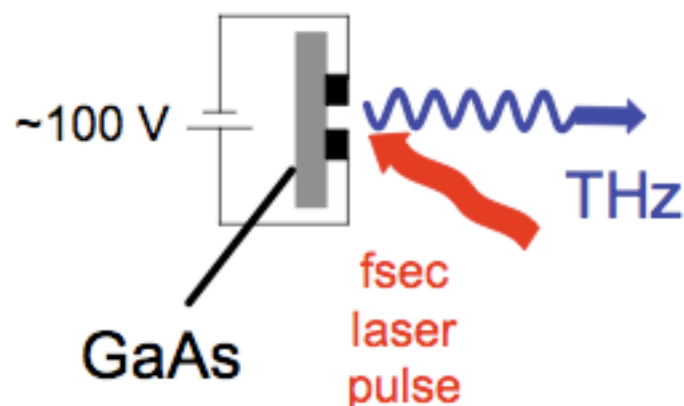


# Comparing Coherent THz Synchrotron and Conventional THz Sources

Larmor's Formula : Power =  $\frac{3e^2 a^2}{2c^3} \gamma^4$  (cgs units)

a=acceleration  
c=vel. of light  
 $\gamma$ =mass/rest mass

## Auston switch

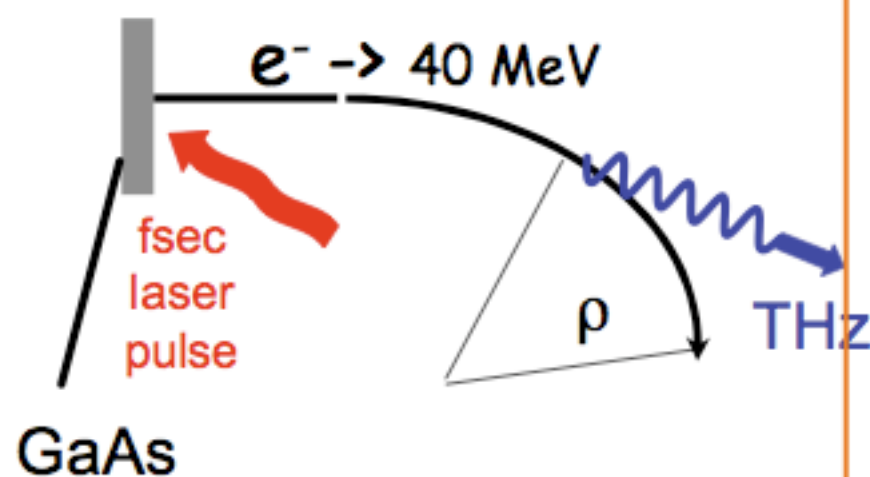


$$E = \frac{100V}{10^{-4}m} = 10^6 V/m$$

$$a = \frac{F}{m} = \frac{10^6 V \cdot e / m}{.5MeV / c^2} = \frac{10^6 (3 \times 10^8)^2}{0.5 \times 10^6}$$

$$\cong 10^{17} m/sec^2$$

## Synchrotron radiation



$$a = \frac{c^2}{?} = \frac{(3 \times 10^8)^2}{1} \cong 10^{17} m/sec^2$$

if ? = 1 m

## Comparing Coherent THz Synchrotron and Conventional THz Sources

Larmor's Formula :  $\text{Power} = \frac{3e^2 a^2}{2c^3} \gamma^4$  (cgs units)

$$\text{Power} \propto \gamma^4$$

### Synchrotron

### THz Antenna

$$\gamma = 1$$

To compare radiation in the THz region, ~40 MeV electrons will get the critical energy into the IR. So,

$$\gamma \approx 75$$

$$\gamma^4 \approx 10^7!$$

Relativistic electrons gain a huge factor in THz power.

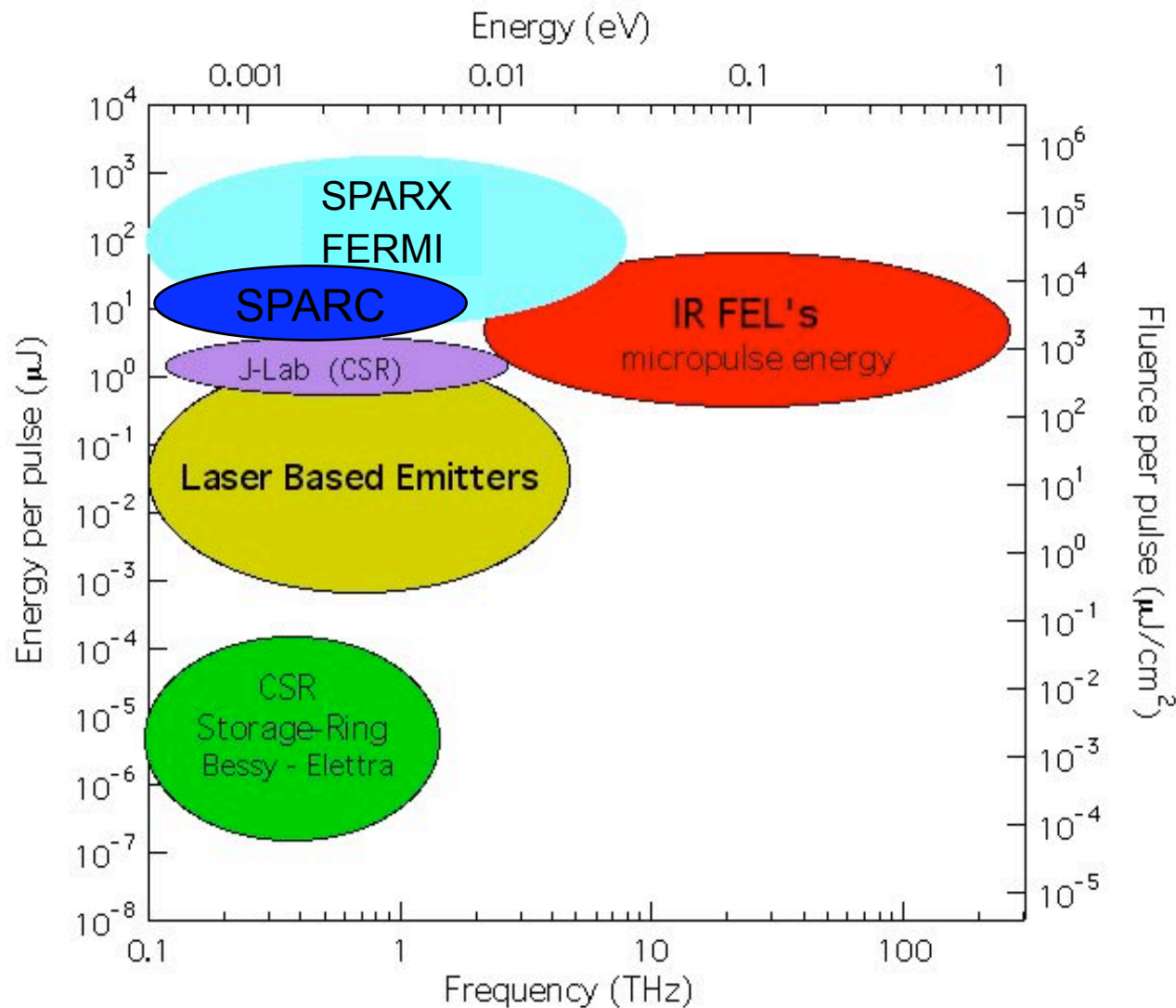
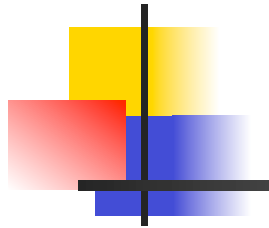


# Figures of Merit of THz Sources

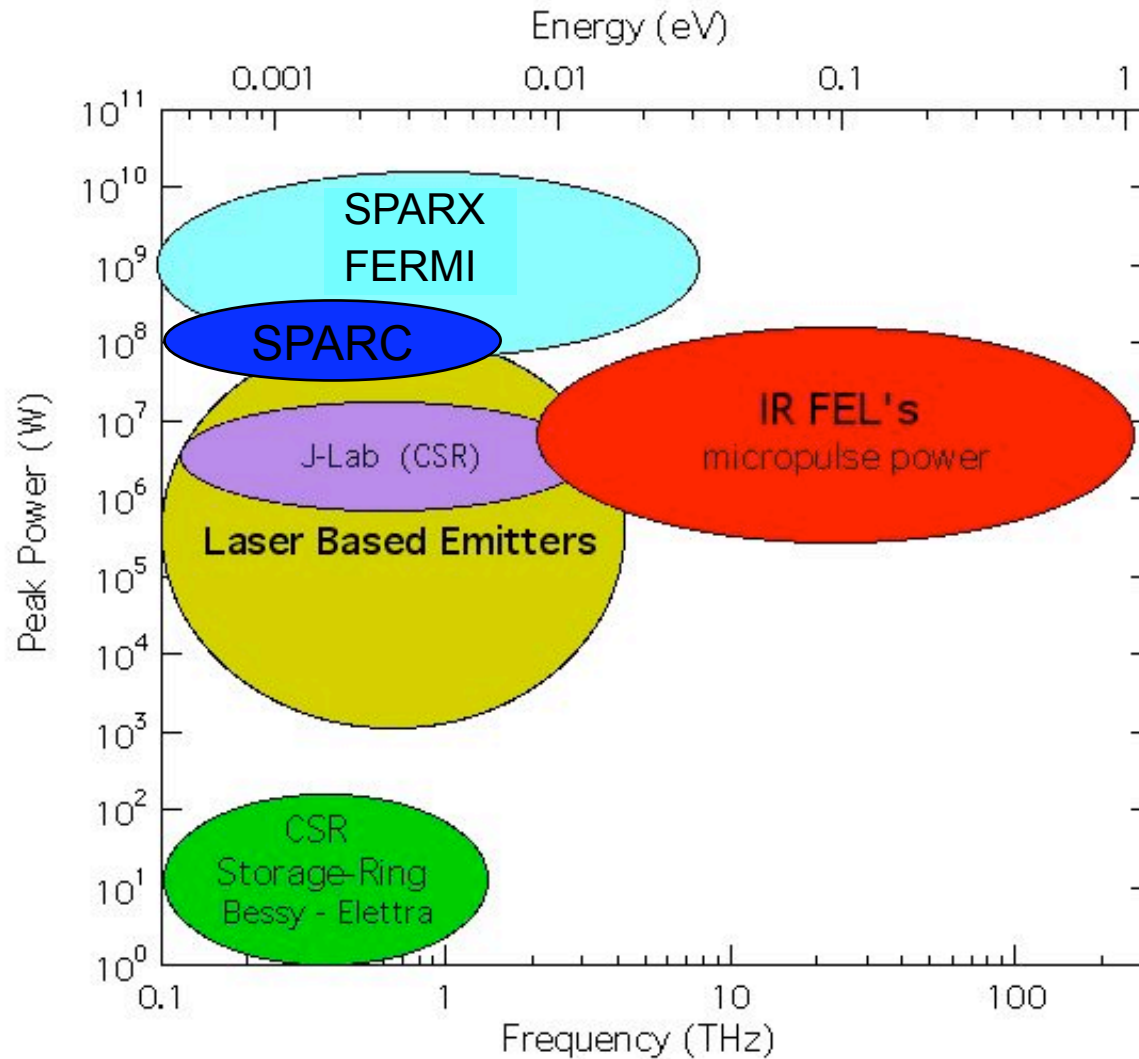
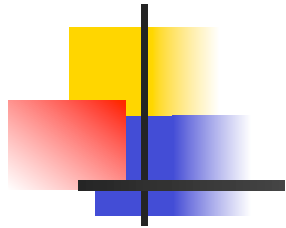
<b>Dedicated Broad Band Sources</b>	<b>Bandwidth</b>	<b>Pulse Width</b>	<b>Rep. Rate</b>	<b>Average Power</b>	<b>Pulse Energy</b>	<b>Peak E-Field</b>
Linac (BNL)	2 THz	300 fs	10 Hz	mW	100 mJ	1 MV/cm
ER-Linac (Jeferson Lab.)	2 THz	300 fs	75 MHz	100 W	$\mu$ J	10 KV/cm
Storage Ring (BessyII, Elettra)	1 THz	1 ps	500 MHz	< W	<nJ	< 10V/cm
Table-top Laser	5 THz	100 fs	1 KHz	nW to $\mu$ W	nJ to $\mu$ J	10 KV/cm
<b>Narrow Band Sources</b>	<b>Op. Range</b>	<b>Pulse Width</b>	<b>Rep. Rate</b>	<b>Average Power</b>	<b>Pulse Energy</b>	<b>Peak E-Field</b>
Stanford FEL	3.75 -20 THz	2-10 ps	11.8 MHz	<1 W	$\mu$ J	<250 kV/cm
UCSB FEL	0.88-4.8 THz	/	/	5-100 mW	/	<70 kV/cm
UCSB FEL	0.12-0.88 THz	/	/	5-100 mW	/	<20 kV/cm
FELIX	1.2-100 THz	6-100 cycles	1000/50/25 MHz	<1 W	1-50 $\mu$ J	<10 V/cm
ENEAL FEL CATS	0.09-0.15 THz	50 ps	3 GHz	<1 W	0.5 $\mu$ J	<10 V/cm
FELBE	1.5-100 THz	1-10 ps	13 MHz	50 mW@2THz	0.01-2 $\mu$ J	10 kV/cm
<b>Continous wave Sources</b>	<b>Op. Range</b>			<b>Average Power</b>		
Backward Wave Oscillator	0.1-1.2 THz			1 mW@1.2 THz		
Gas Lasers (CO2-pumped)	not cont. tunable			1 W@2.5 THz		
Quantum Cascade Lasers	2-10 THz			10 mW@3.5 THz		
<b>IV UV-X FEL</b>	<b>Bandwidth</b>	<b>Pulse Width</b>	<b>Rep. Rate</b>	<b>Average Power</b>	<b>Pulse Energy</b>	<b>Peak E-Field</b>
SPARX-FERMI-SPARC (long bunch)	1 THz	1 ps	0-50 Hz	1-5 mW	100 $\mu$ J	MV/cm
SPARX-FERMI (medium bunch)	2 THz	500 fs	0-50 Hz	2-7 mW	150 $\mu$ J	MV/cm
SPARX-FERMI (short bunch)	10 THz	50 fs	0-50 Hz	>10 mW	>nJ	10 MV/cm

Data from: Biedron et al 2007; Riemann 2007; A. Perucchi et al 2007

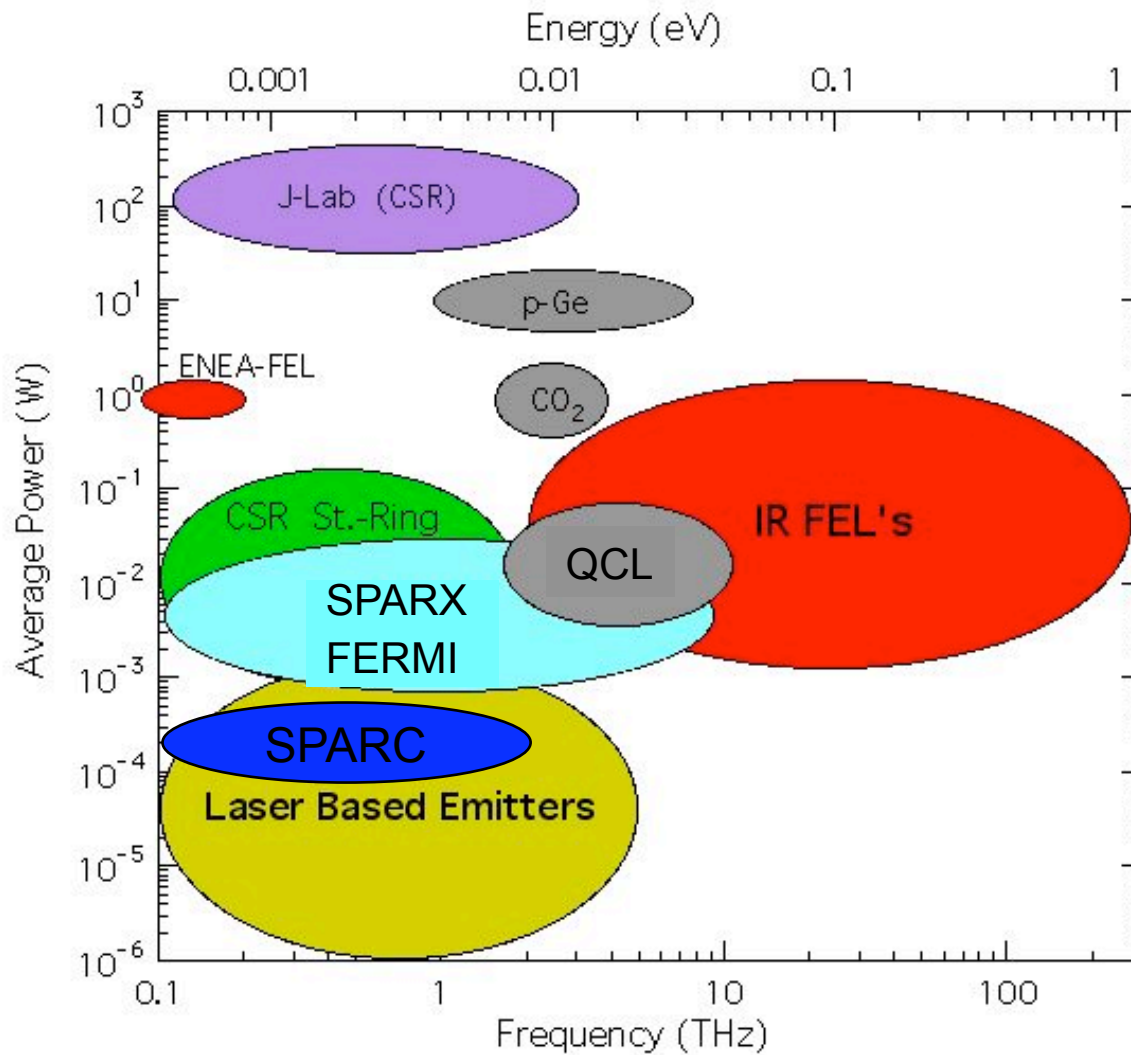
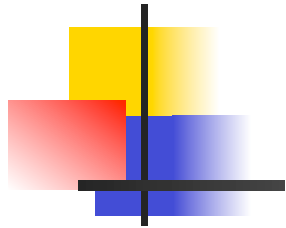
# Figures of Merit of THz sources: energy/pulse

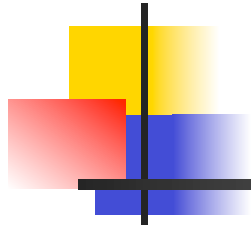


# Figures of Merit of THz sources: peak power



# Figures of Merit of THz sources: average power





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# THz Science

## Linear Spectroscopy



# Superconductivity today:

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## 1. Looking for new materials

- 1986: High- $T_c$  cuprates
- 1990:  $C_{60}$  fullerene
- 2000:  $MgB_2$
- 2003:  $Na_xCoO_2 \cdot 3H_2O$
- 2004: B-doped diamond

## 2. Understanding new materials



THz spectroscopy plays a role...



... because

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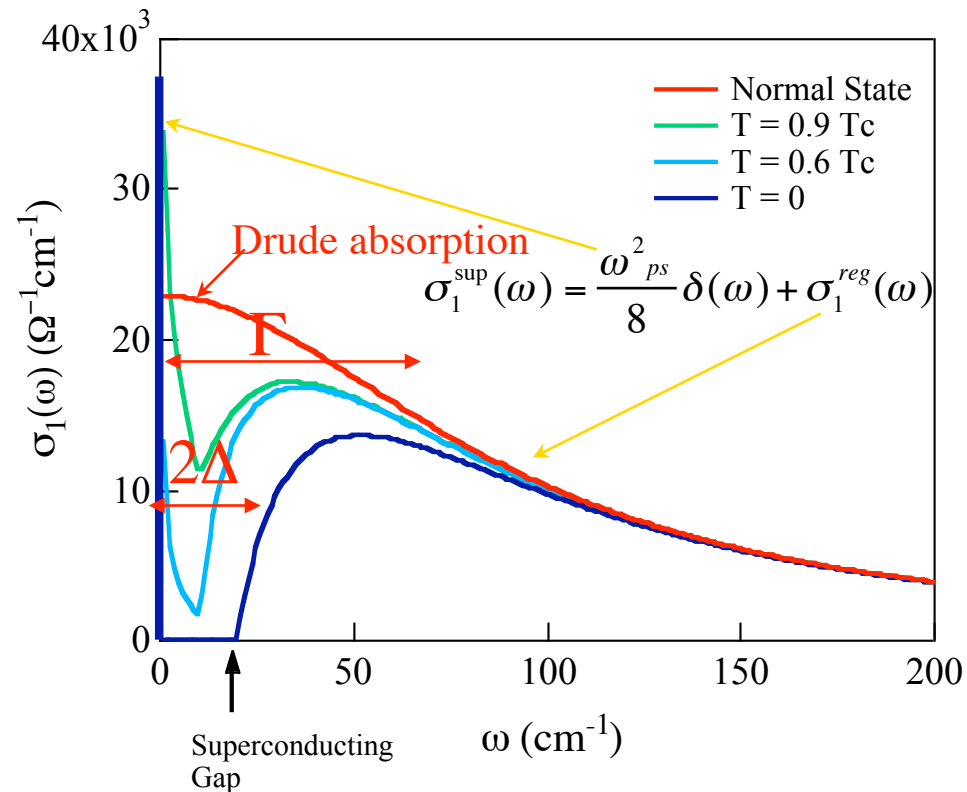
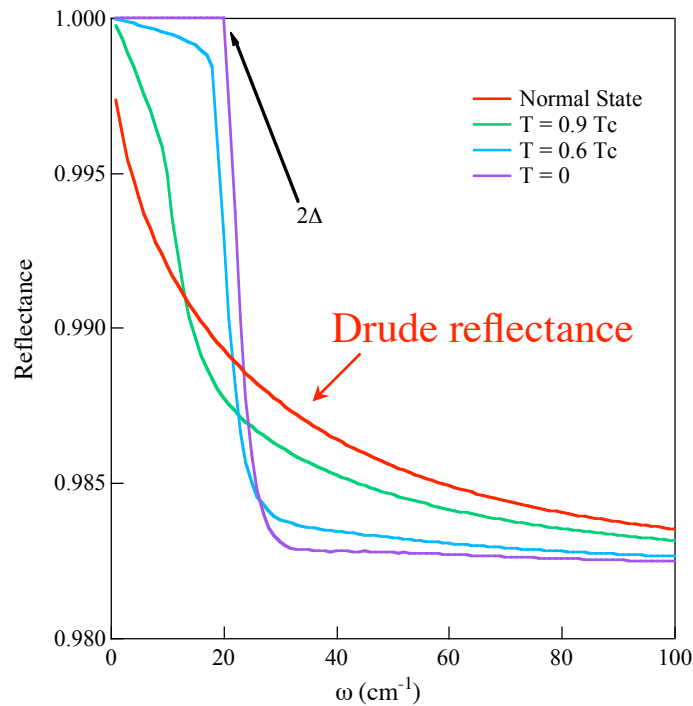
Superconductivity is ruled by **low-energy** electrodynamics:

- Superconducting gap : THz-range
- Spectral weight of condensate and penetration depth: THz-range
  - Mediators of pairing (phonons, etc.): Far-infrared
  - Free-carrier conductivity above  $T_c$ : Infrared

# Basic optics of Superconductors

Minimum excitation energy:  
Cooper-pair breaking  $2\Delta$

- Superconducting gap observed if:
- sample in the dirty-limit ( $2\Delta < \Gamma$ )
  - Cooper pairs in **s-wave** symmetry



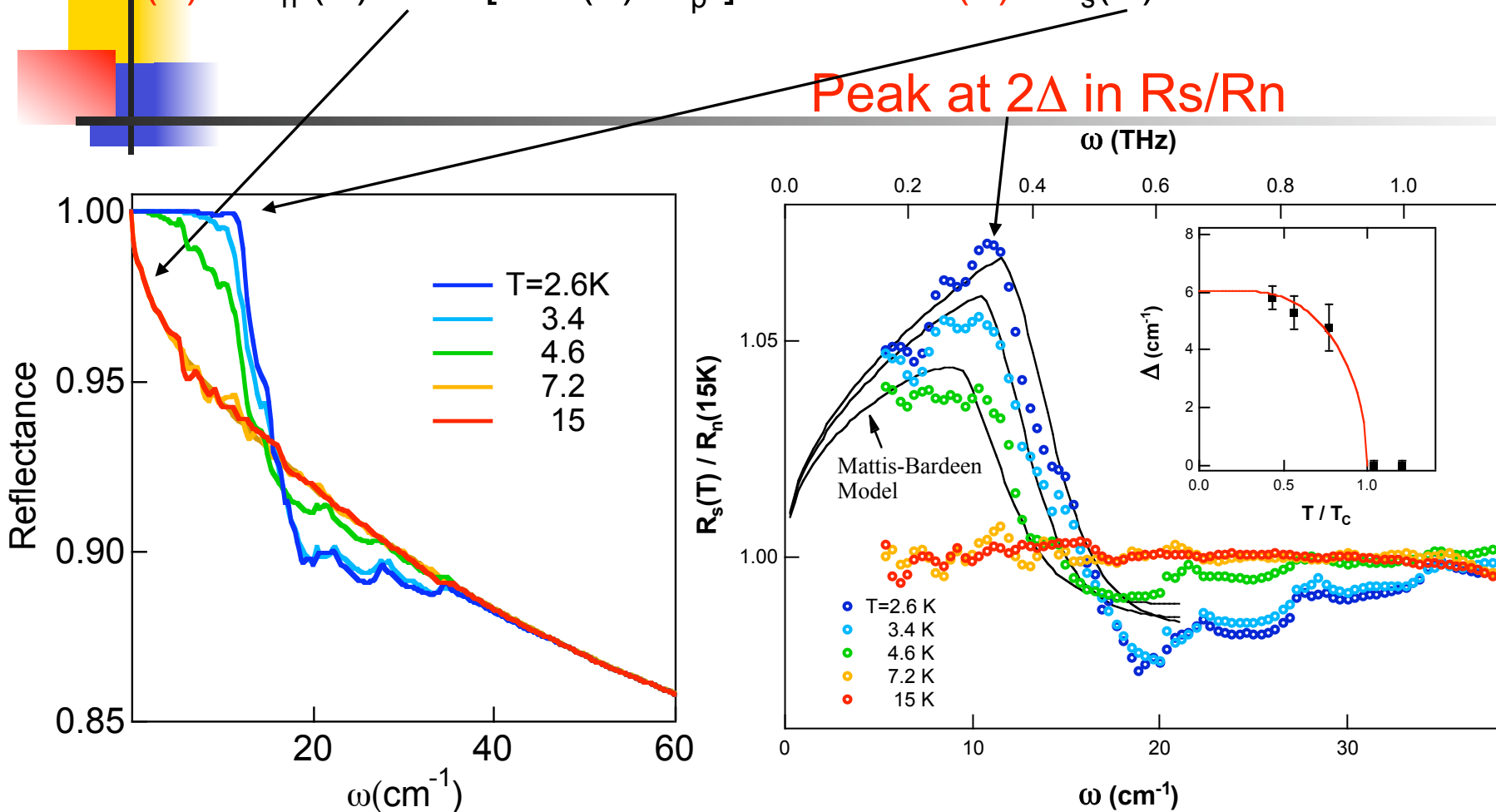
$$\int [\sigma_1(\omega, T > T_c) - \sigma_1(\omega, T < T_c)] d\omega = \omega_{ps}^2 / 8 = n_s e^2 / m^* \rightarrow \lambda = c / \omega_{ps}$$

**Ferrel-Glover-Tinkham Rule**



# Superconducting Diamond

$$\omega \leq \Gamma(T) : R_n(\omega) = 1 - [8\omega\Gamma(T)/\omega_p^2]^{1/2} \quad \omega \leq 2\Delta(T) : R_s(\omega) = 1$$



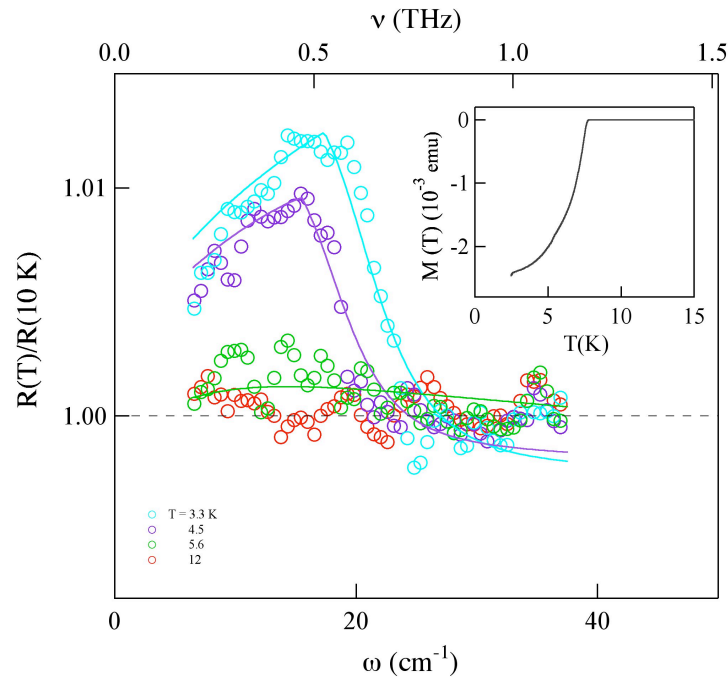
**s-wave Dirty-Limit Regime;  $2\Delta(2.6\text{ K})=12\pm 1\text{ cm}^{-1}$ ---->  $2\Delta/k_B T_c=3.2 \pm 0.5$**

# CaAlSi (anisotropic superconductor)

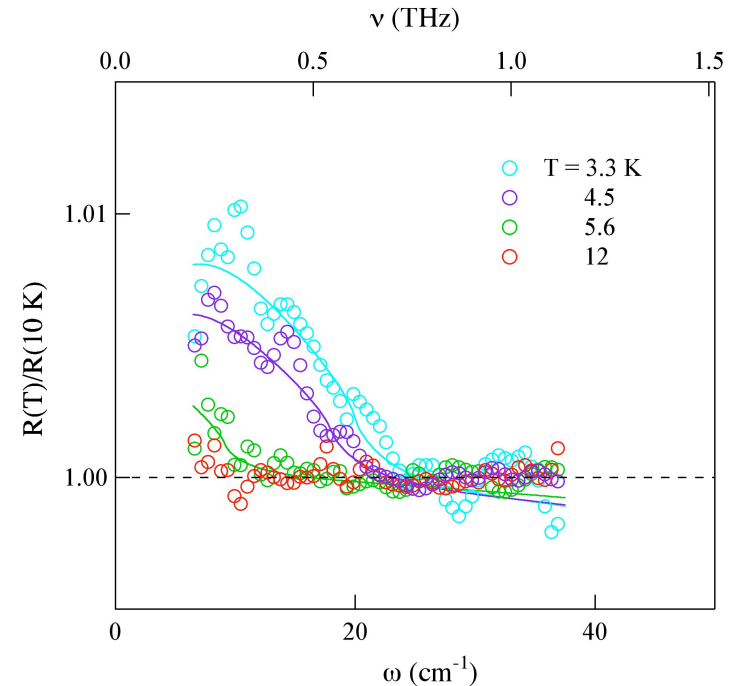
## polarized gap measurements



exagonal plane



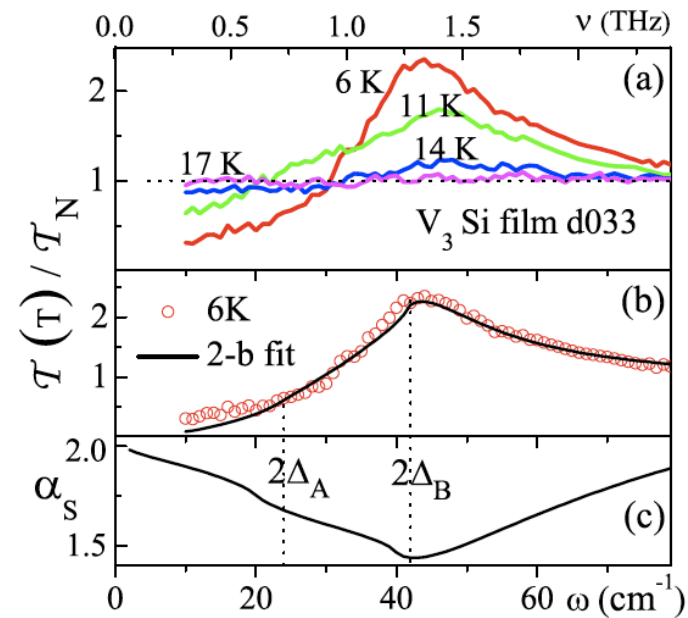
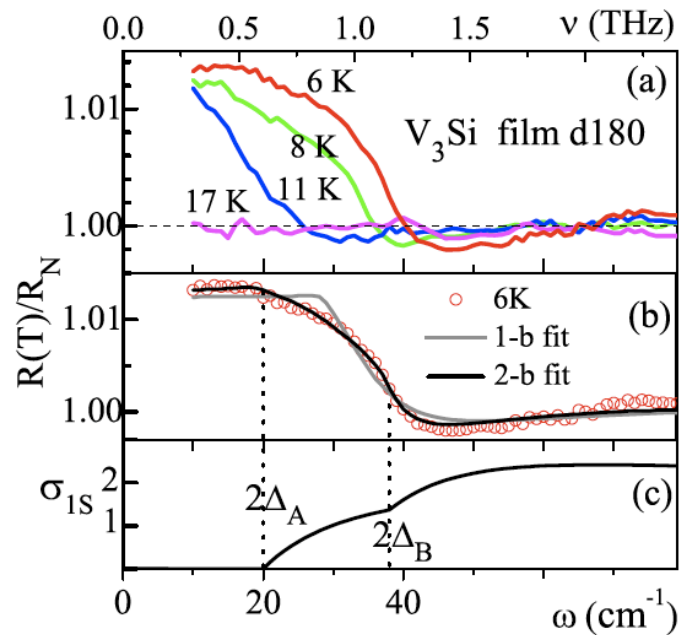
c-axis



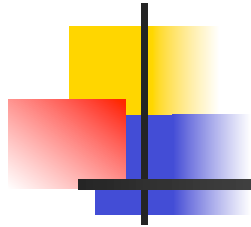
**s-wave Dirty-Limit Regime both along the exagonal plane and along the c-axis:**

$$\begin{aligned}
 2\Delta_{ab}(0 \text{ K}) &= 19 \pm 1 \text{ cm}^{-1} \\
 2\Delta_c(0 \text{ K}) &= 23 \pm 1 \text{ cm}^{-1}
 \end{aligned}
 \longrightarrow
 \begin{aligned}
 2\Delta_{ab}/k_B T_C &= 3.9 \pm 0.2 \\
 2\Delta_c/k_B T_C &= 4.5 \pm 0.2
 \end{aligned}$$

# Multigap Superconductivity: V<sub>3</sub>Si



**THz transmittance and reflectance demonstrate the presence of two BCS gaps  $2\Delta_1=0.71$  THz,  $2\Delta_2=1.27$  THz**

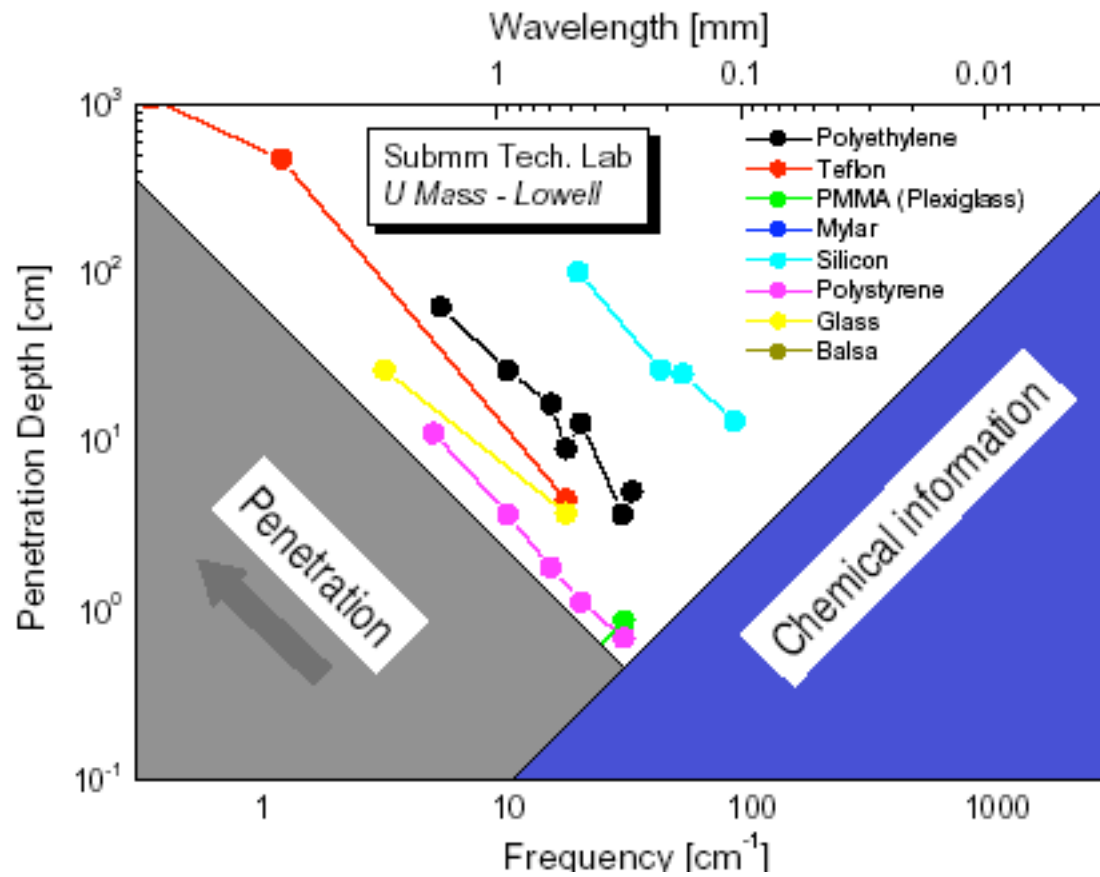
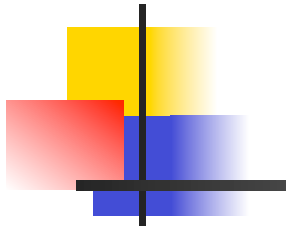


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# THz Science

Imaging

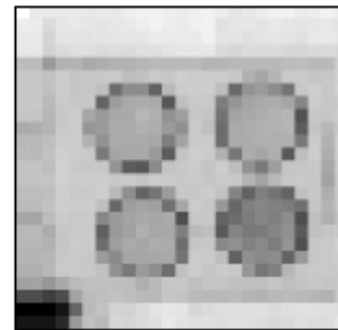
# Imaging vs Penetration



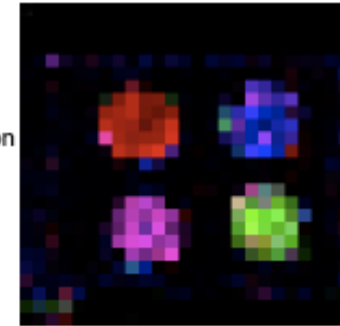
# Chemical Pharmaceutical recognition



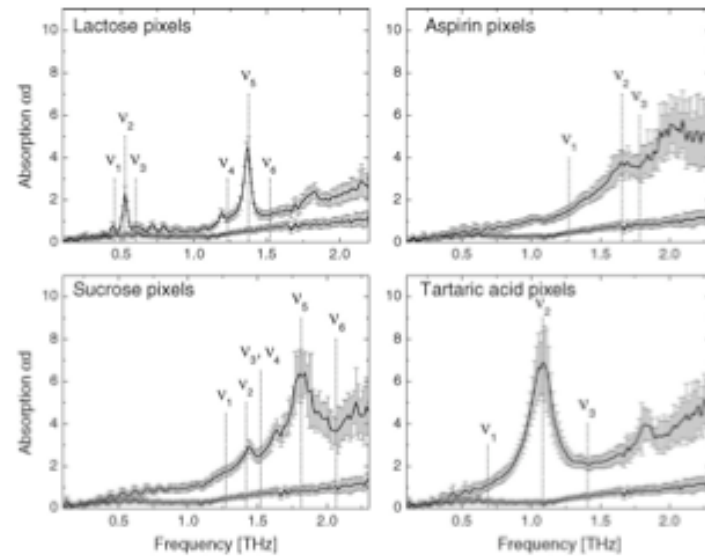
**Fig. 8.** Visible image of sample with four pellets containing different chemicals: (1) lactose, (2) aspirin, (3) sucrose, and (4) tartaric acid.



Chemical recognition



**Figure 9.2** The left panel shows a traditional THz transmission image of four pills containing different chemicals (lactose, aspirin, sucrose, and tartaric acid). The right image has color-coded according to the spectral information contained in the full THz trace recorded at each pixel, and is known as "functional imaging". Figure courtesy of Y. Watanabe et al., Yamagata University.



Background absorption

**Fig. 10.** Solid lines show the average absorption of lactose (top, left), aspirin (top, right), sucrose (down, left), and tartaric acid (down, right) in the sample. The lower curve in each panel shows the absorption of the packaging material. The error bars represent one standard deviation from the mean of typically 20-30 measurements. The indicated frequencies are used for chemical recognition. After [20].

# Imaging of Bio-materials, molecular in-vivo imaging of pathogenesis

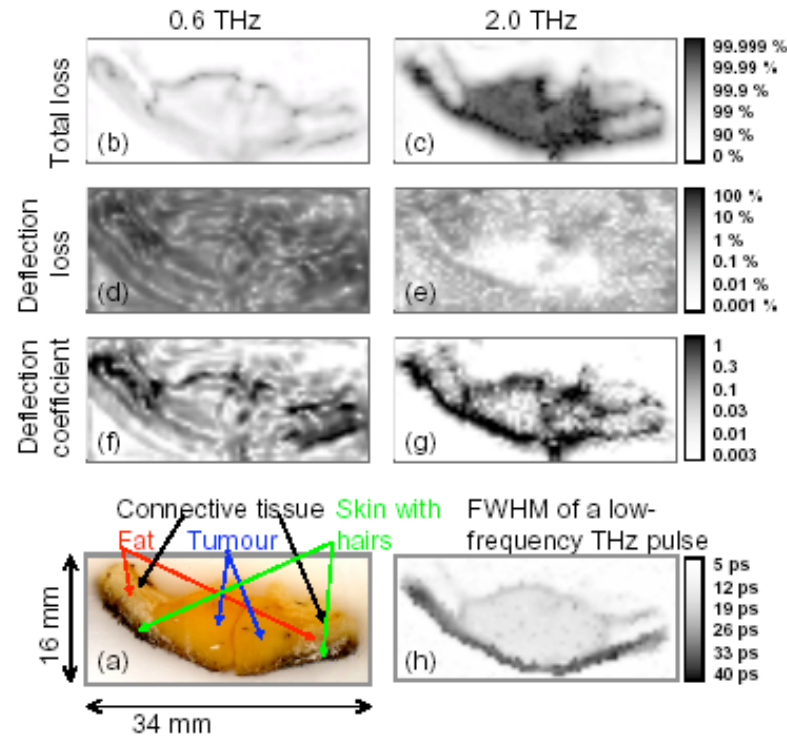
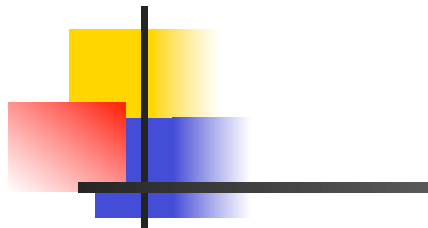


Fig. 3. (a) Optical image of the sample; (b) and (c): Total loss in transmission, (d) and (e): Loss induced by deflection; (f) and (g): Deflection coefficient. (h) Pulse duration (FWHM) of a low-frequency THz pulse. Click on Fig. 3(b,d,f) to see the data as a function of the frequency. (426 kB QuickTime movie)

T. Lffler et al, Optics Express  
9, 616 (2001)  
Ferguson et al, Nature  
Materials1, 26 (2002)  
X.-C Zhang Phys. Med. Biol.  
47, 3667 (2002)

THz do not subject a biological tissue to harmful radiation and may provide both imaging and spectroscopic information on biological materials.

Needs of:

1. High S/N ratio
2. High acquisition rate and resolution
3. THz database for biological tissues
4. High power to increase sensing and penetration
5. Near-field imaging to increase spatial resolution (up to now 1 microns resolution has been obtained)

*PROTOTYPE DESIGN OF T-RAY  
SCIENCE'S COMPACT SKIN CANCER  
DETECTION SYSTEM*





# Quality Assurance of Chocolate Products with Terahertz Imaging

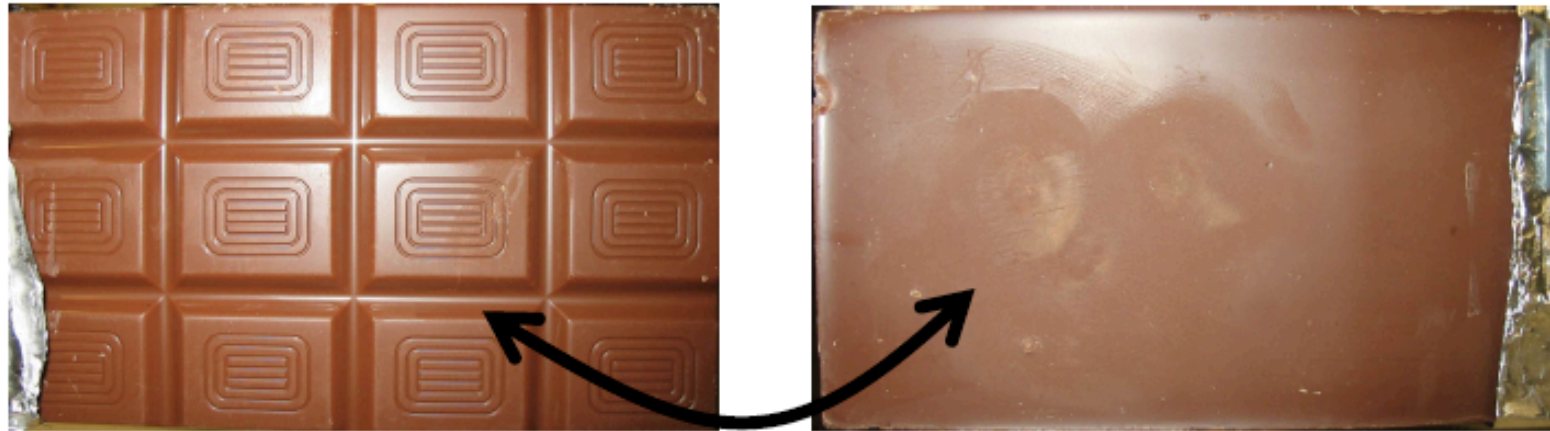
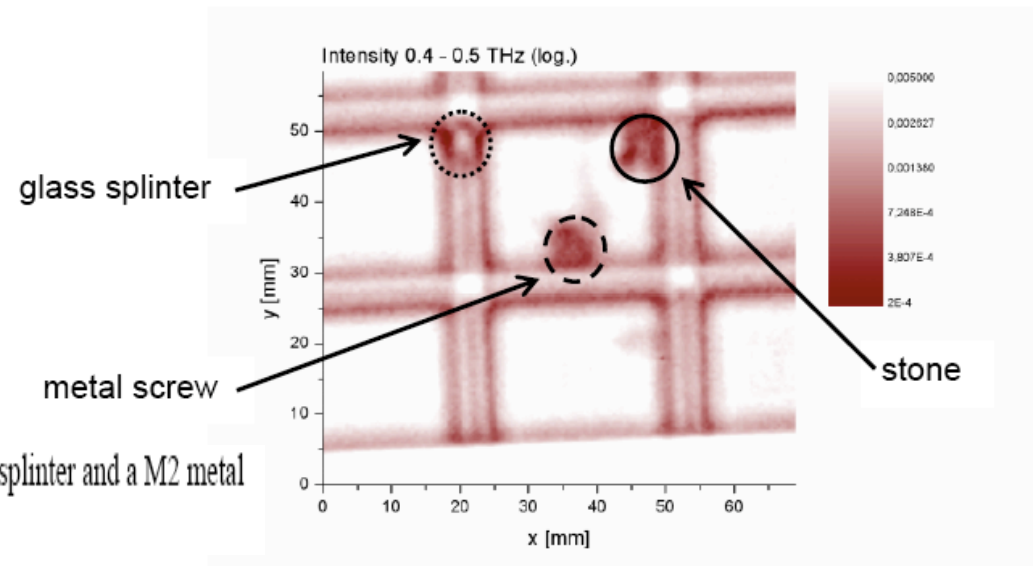
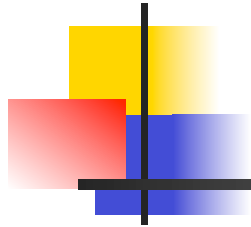


Figure 3. Front and back side of a chocolate bar after artificial contamination with a stone, a M2 metal screw and a glass splinter.



Figure 4. Photograph of different kind of contaminations like a small stone, a glass splinter and a M2 metal screw. The M2 screw may help for size reference.



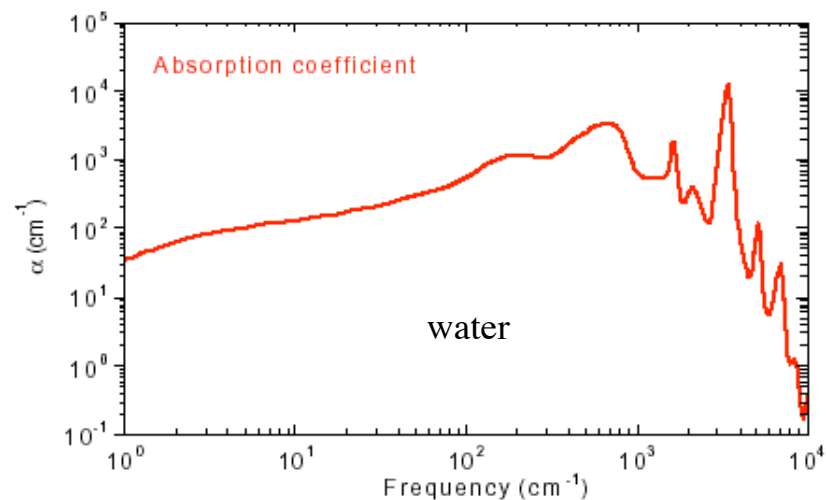
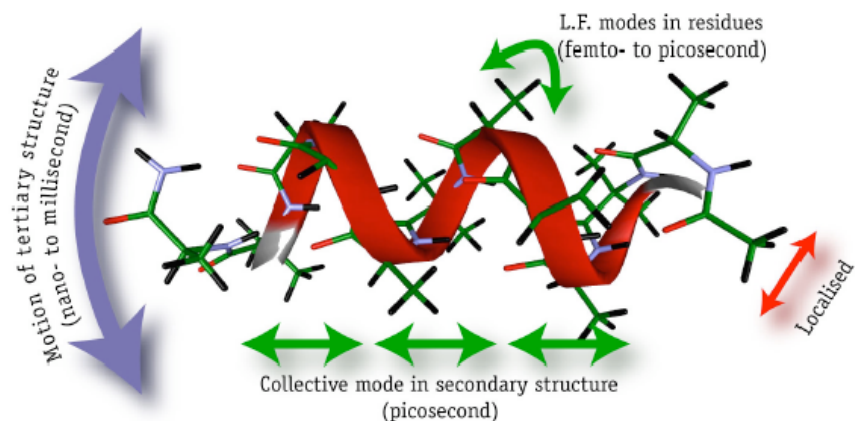


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# THz Science

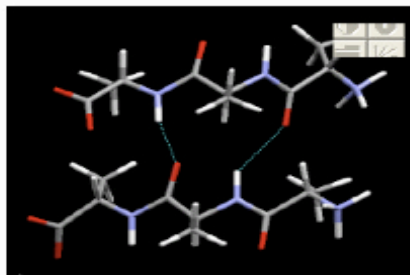
Linear Spectroscopy in Life Science

# Conformational Collective modes of macromolecules

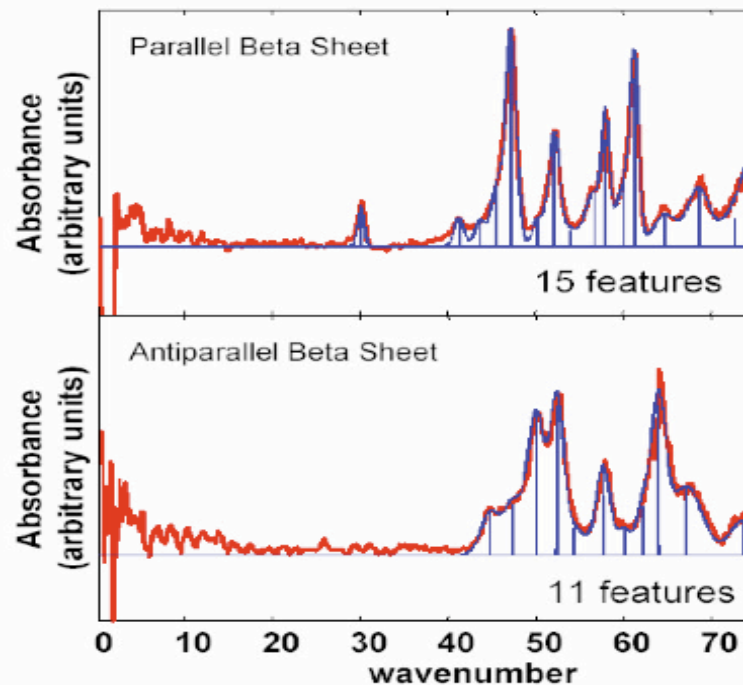
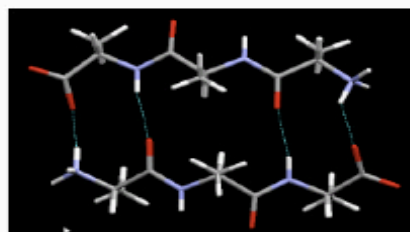


THz Spectra for parallel and antiparallel forms of trialanine show extreme sensitivity to the molecular environment.

Parallel Beta Sheet conformation



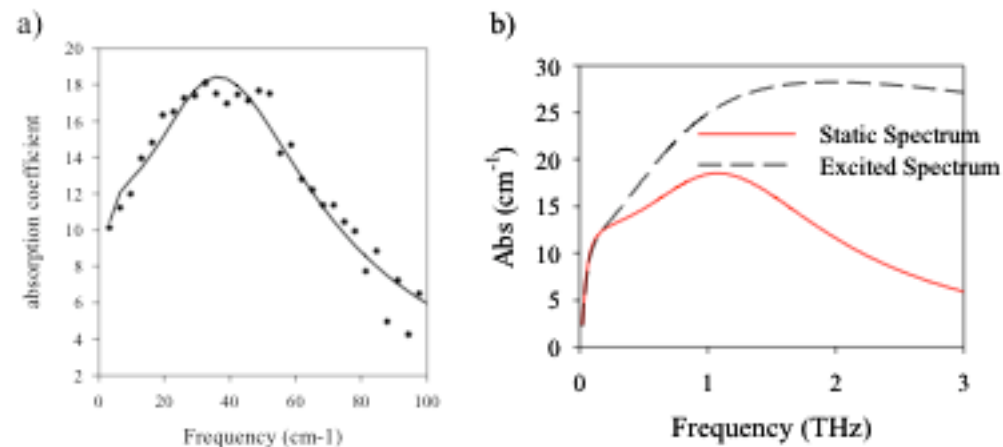
Antiparallel Beta Sheet conformation



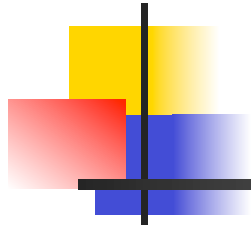
# Femto-Chemistry

THz is sensitive to interaction time between a molecule and the surrounding solvent on sub-ps scale

Standard high-frequency time-resolved experiments:  
excite in the UV-VIS a soluted molecule -----> probe in the THz the dynamical effects on solvent



**Figure 6.2.** Static spectrum of chloroform (part a). Absorption spectrum of collective solvent modes near photoexcited dye molecule (part b). Image courtesy of Charles Schmittenmaer, Yale University.



# THz Science

Pump-Probe Time Resolved  
Spectroscopy

# Pump-Probe Spectroscopy

**1) Optical Pump-Optical Probe Spectroscopy** Pump and Probe pulses (often at a single frequency) fall in the visible (near-IR)

- High energy excitation;
- Strong scattering effects;
- High energy dynamics;

**Extrinsic dynamics**

**2) Optical Pump-THz Probe Spectroscopy** Pump (single frequency) in the visible (near-IR) Probe in the far-IR and THz range;

- Similar inelastic effects in the Quasi-Particle decay like in (1) but investigation of the low-energy dynamics;

**3) THz Pump-THz Probe Spectroscopy** Affordable Pump pulses falling in the far-IR and THz;

- Possibility to resonate and/or selectionate several fundamental excitations;

**Intrinsic dynamics**

**4) THz Pump- IR+VIS SR Probe Spectroscopy**

- Affordable Pump pulses falling in the far-IR and THz;
- Possibility to resonate and/or select several fundamental excitations;

**Intrinsic strongly-coupled different energy scale dynamics**

VOLUME 85, NUMBER 14 PHYSICAL REVIEW LETTERS 2 OCTOBER 2000  
**Exploring the Dynamics of Superconductors by Time-Resolved Far-Infrared Spectroscopy**

G. L. Carr\*

R. P. S. M. Lobo, J. LaVeigne, D. H. Reitze, and D. B. Tanner

High Energy (Visible) Pump  
THz Probe

Many Cooper pairs are broken

→ Extrinsic Dynamics

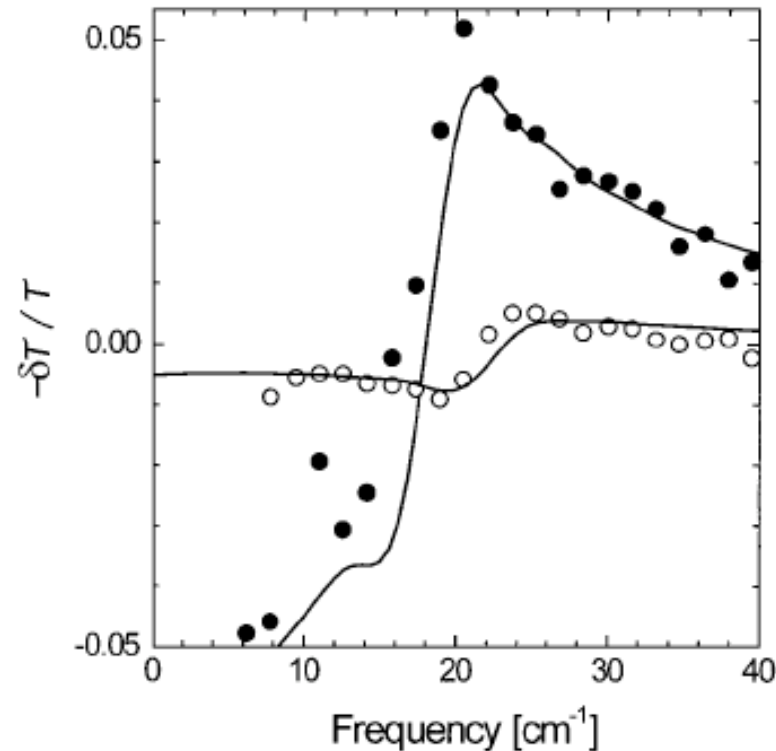



FIG. 3. Photoinduced spectral changes for two samples. The laser pulse energies were 1.8 nJ (solid circles) and 0.4 nJ (open circles). Also shown are fits (solid curves) assuming an  $\sim 3\%$  and  $0.6\%$  reduction in the energy gap, respectively. All other parameters were held fixed at the values determined from the temperature-dependent  $T_S/T_N$  fits.



# New paradigm: THz pump → THz probe

---

- ★ Use THz pulses as the pump (excitation) source in “all THz” pump probe experiments.
- ★ Study non-linear absorption in a variety of materials, including nanoparticles / quantum dots.
- ★ Induce coherent current excitations
- ★ Move atoms in their local potential wells
- ★ Orient spins in magnetic systems.



# High-power half-cycle THz pulses

A 100  $\mu\text{J}$ , half-cycle THz pulse, focused into a volume of 1  $\text{mm}^3$  or less.

- **E-field** =  $[2D_E/\epsilon_0]^{1/2} \sim 10^8 \text{ V/m}$  ( $\sim 1 \text{ MV/cm}$ ).
- $\Rightarrow$  Use large electric field to displace atoms in polar solids (structural phase transitions, soft modes, ferroelectricity, ...)
- **H-field** =  $E/c \sim 0.3 \text{ T}$
- $\Rightarrow$  Use transient magnetic field to create magnetic/spin excitations and follow dynamics on ps time scale (e.g., time-resolved MOKE).

Or, some other shape pulse?

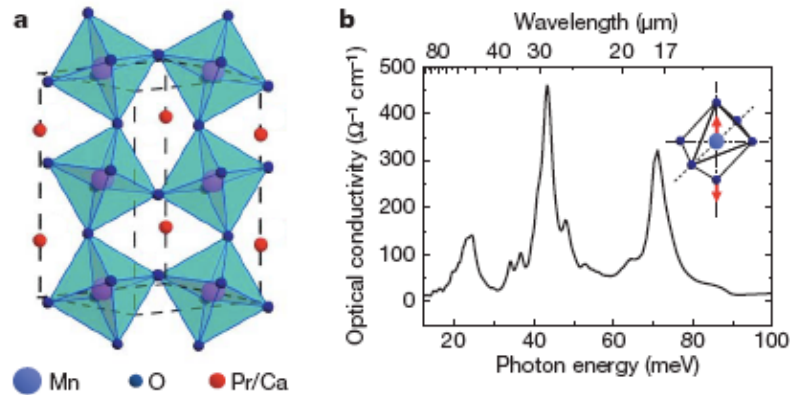
$$\frac{dI(\omega)}{d\omega} \underset{\text{multiparti de}}{=} [N + N(N-1)f(\omega)] \frac{dI(\omega)}{d\omega} \quad f(\omega) = \left| \int_{-\infty}^{\infty} e^{i\omega \hat{n} \cdot \vec{r}/c} S(r) dr \right|^2$$

$\Rightarrow$  shape electron bunch profile to control E-field shape (coll. W/J. Neuman, U. Md.)

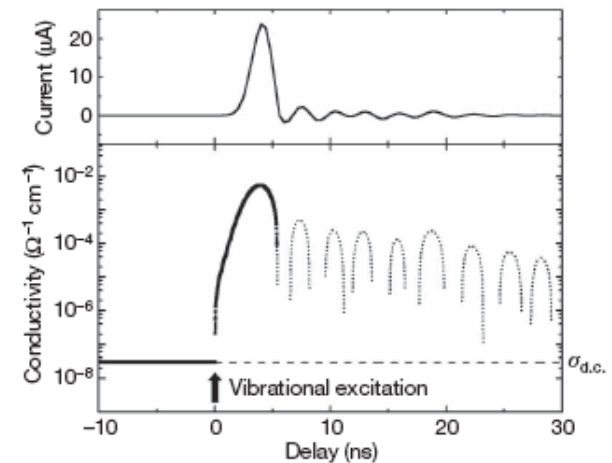
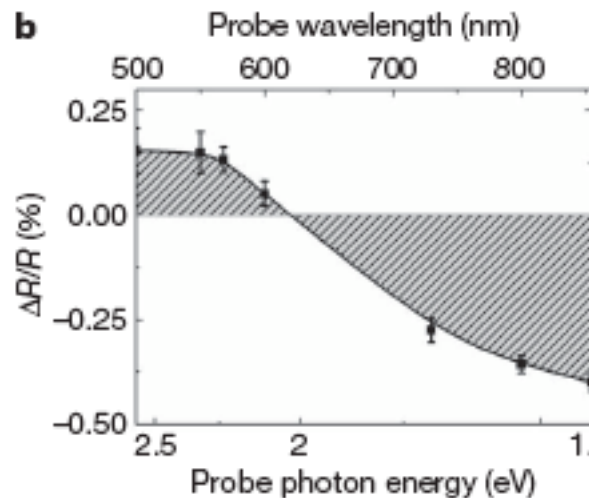
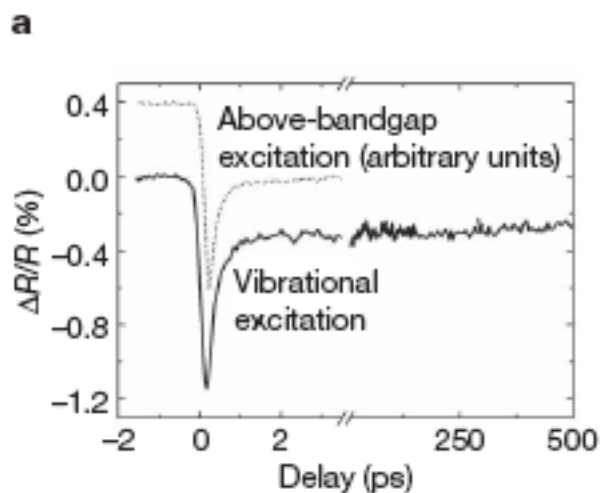
# Coherent Control of Electronic Properties of Manganites

Matteo Rini<sup>1</sup>, Ra'anan Tobey<sup>2</sup>, Nicky Dean<sup>2</sup>, Jiro Itatani<sup>1,3</sup>, Yasuhide Tomioka<sup>4</sup>, Yoshinori Tokura<sup>4,5</sup>, Robert W. Schoenlein<sup>1</sup> & Andrea Cavalleri<sup>2,6</sup>

NATURE|Vol 449|6 September 2007

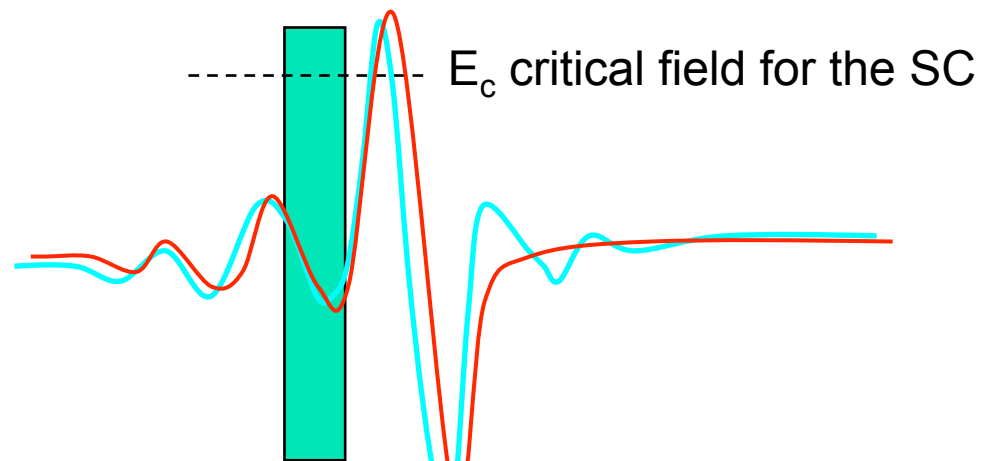


Several exotic materials show a strong interaction between low and high energy degrees of freedom  
 → Through the pump of low energy states one can coherently modify the high energy states

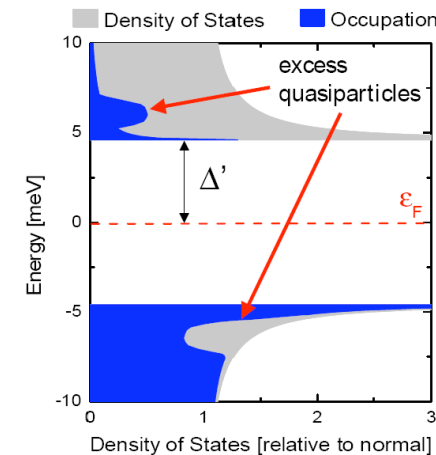


# New Experiment in Superconductivity

Breaking the superconducting state through an electric field exceeding the critical electric field  $E_c$



Superconducting film,  $T < T_c$

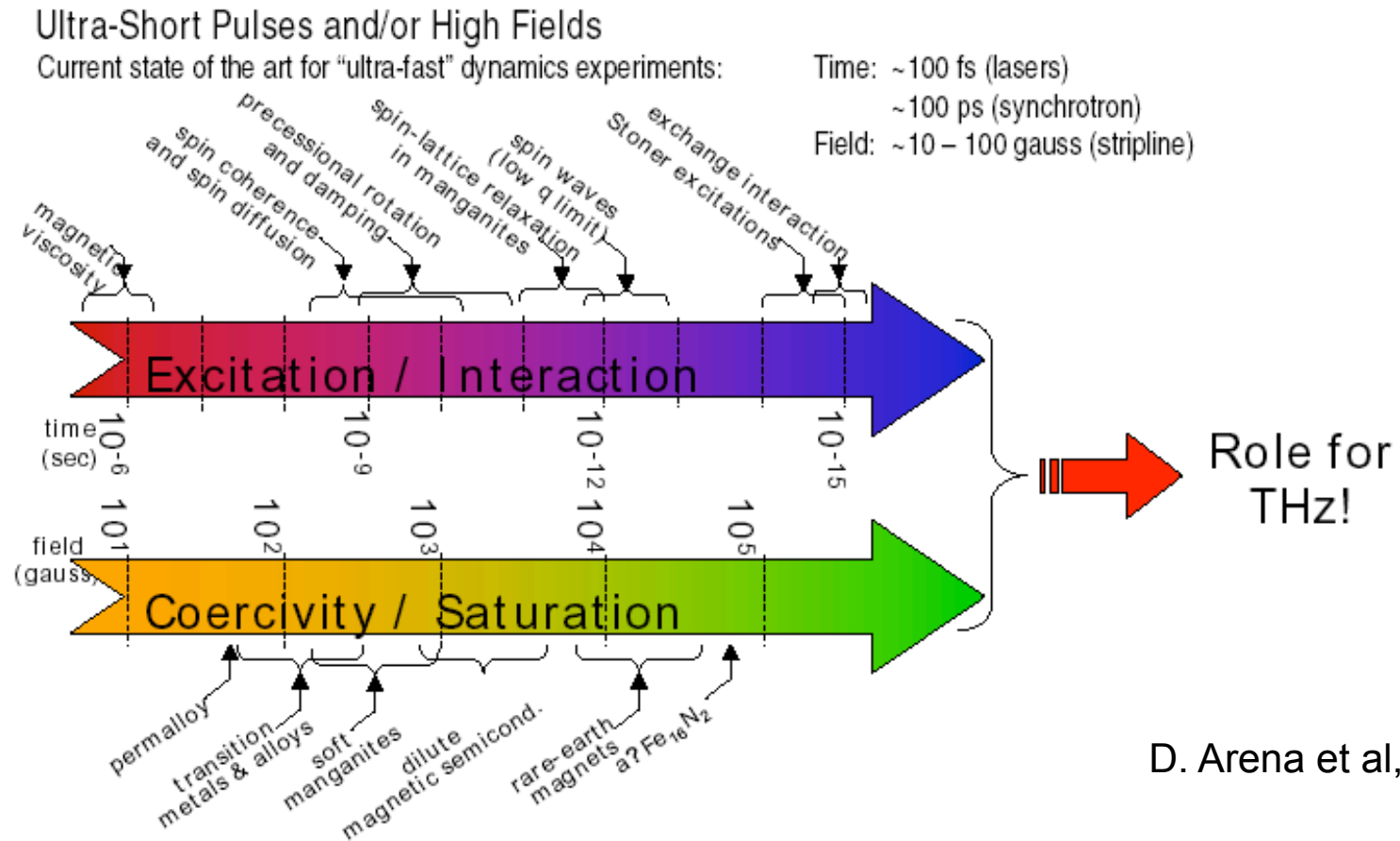


**Ultrashort, intense ( $E > 1$  MV/cm) THz pulses needed!**

# Magnetic Dynamics with THz pulses

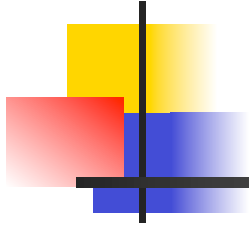
A sub-ps THz pulse associate to an E field of 1 MV/cm a B field of 0.5 T

## THz induced magnetic transitions



D. Arena et al, 2006

# Conclusion



- Relevant scientific cases for Linear and Pump-Probe THz science;
- Relevant technological applications for THz radiation;
- Lack of broad-band ultra-fast sources in the THz range;
- Most of those experiments could be performed through THz radiation produced in LINAC sources;

# Acknowledgments



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- Andrea Perucchi ([SISSI@ELETTRA](mailto:SISSI@ELETTRA))
- Michele Ortolani (CNR-IFN)
- E. Karanzoulis (ELETTRA)
- M. Veronese (ELETTRA)
- G. De Ninno (ELETTRA)
  
- P. Calvani (INFM and University La Sapienza)
- A. Nucara (INFM and University La Sapienza)
  
- G. Gallerano (ENEA)
- G. Williams (Jefferson Laboratory)
  
- The Elettra management