



The Abdus Salam
International Centre for Theoretical Physics



2145-34

Spring College on Computational Nanoscience

17 - 28 May 2010

Electronic-structure and Transport in Nanostructures based on Wannier Function

Nicola MARZARI

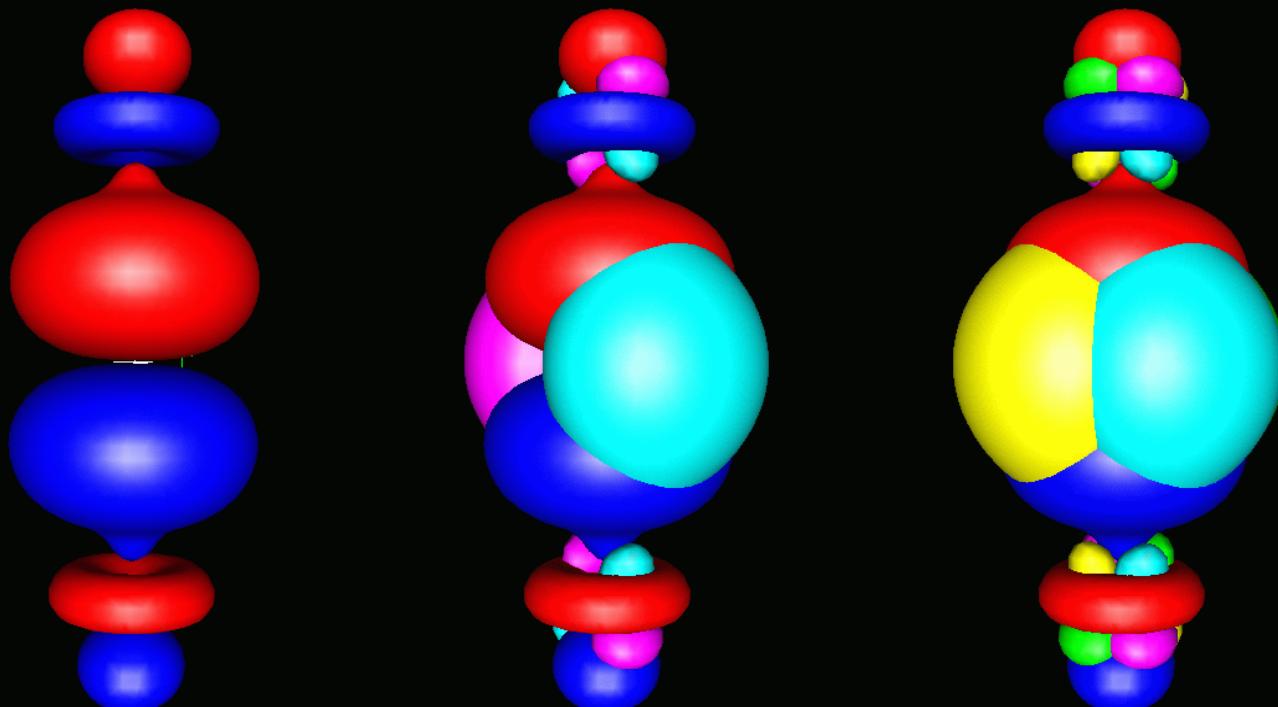
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Oxford, U.K. and Dept. Materials Science & Engineering, MIT, Cambridge
U.S.A.*

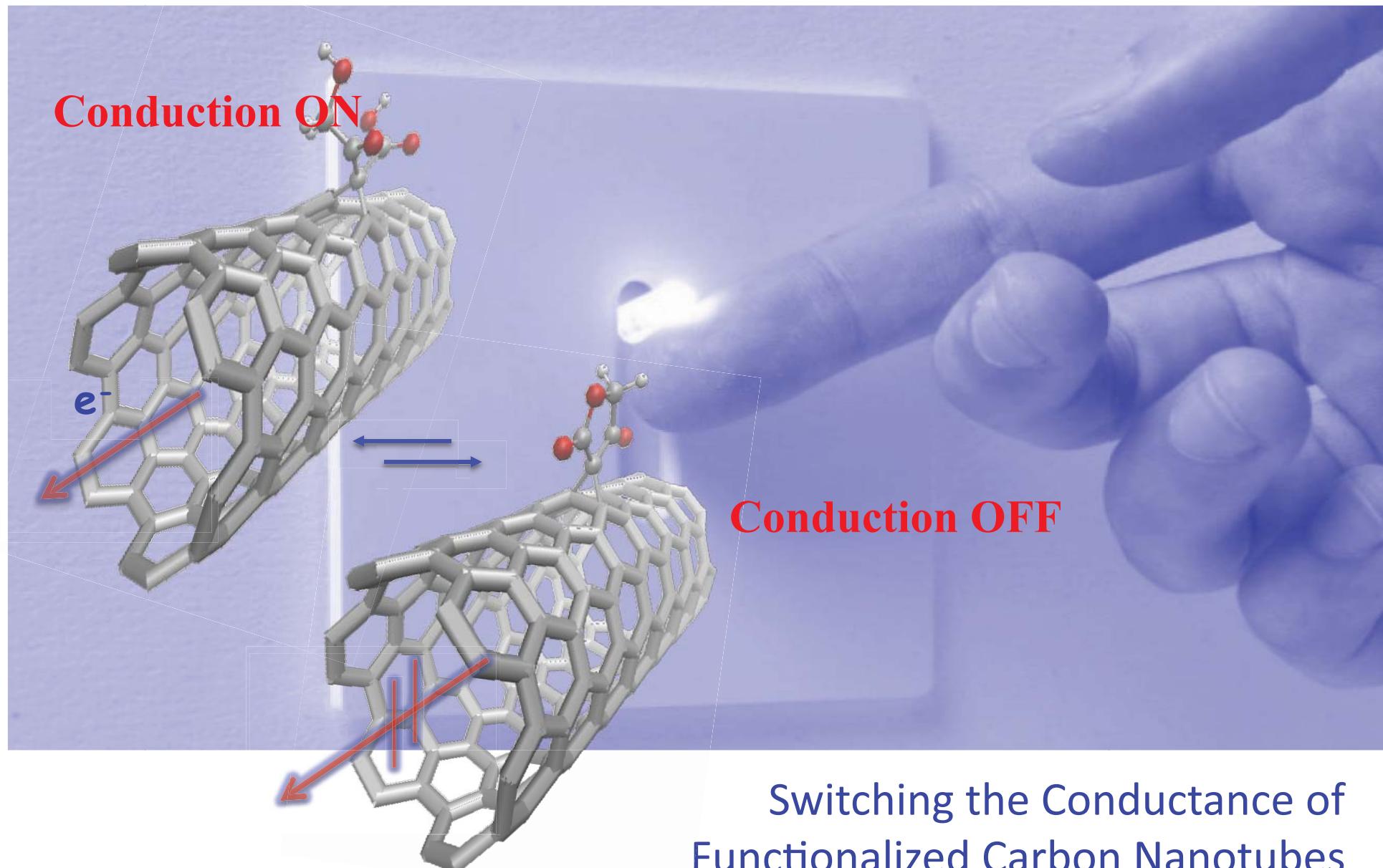
Building nanostructures Bloch by Bloch

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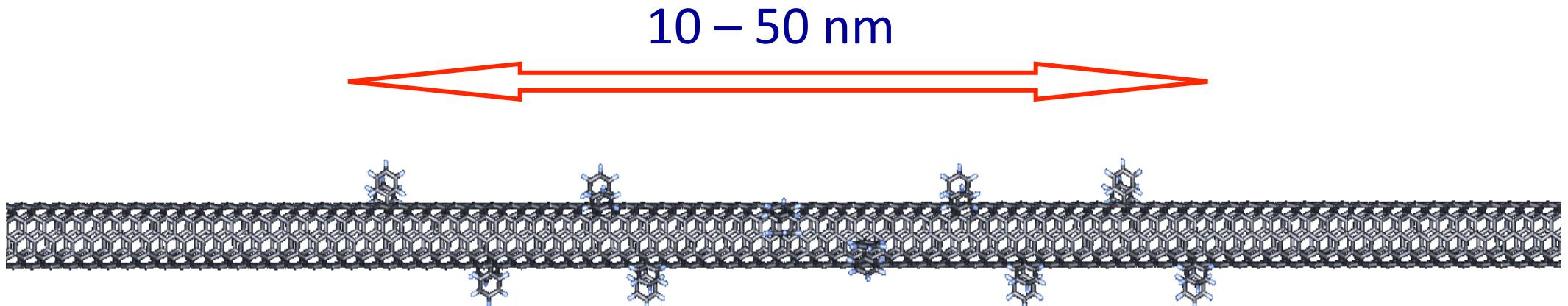
Department of Materials Science and Engineering, MIT





Switching the Conductance of
Functionalized Carbon Nanotubes

Electronic-structure and quantum conductance



- Electronic-structure of nanostructures with thousands of atoms from maximally-localized Wannier functions

Bloch Theorem

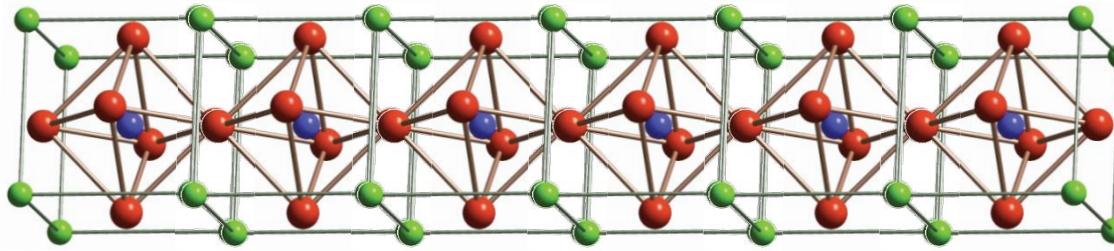
The one-particle effective Hamiltonian \hat{H} in a periodic lattice commutes with the lattice-translation operator \hat{T}_R , allowing us to choose the common eigenstates according to the prescriptions of Bloch theorem:

$$[\hat{H}, \hat{T}_R] = 0 \Rightarrow \Psi_{n\mathbf{k}}(\mathbf{r}) = u_{n\mathbf{k}}(\mathbf{r}) e^{i\mathbf{k}\cdot\mathbf{r}}$$

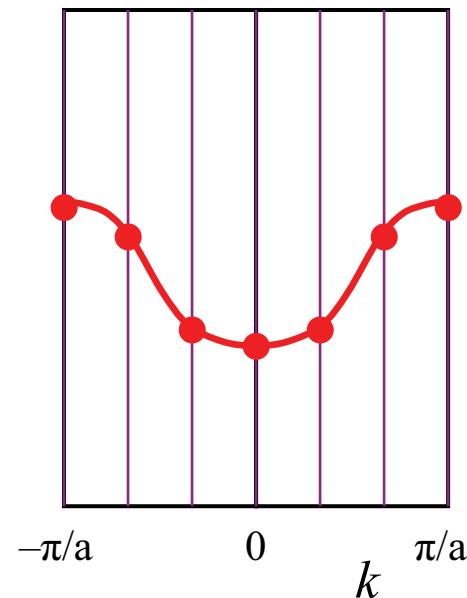
- n, k are the quantum numbers (band index and crystal momentum), u is periodic

Bloch Theorem

Crystal in real space:

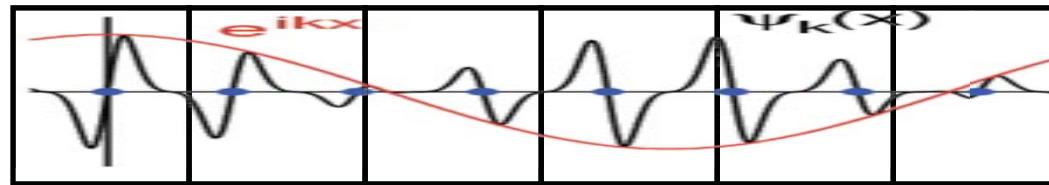


Brillouin zone in reciprocal space:

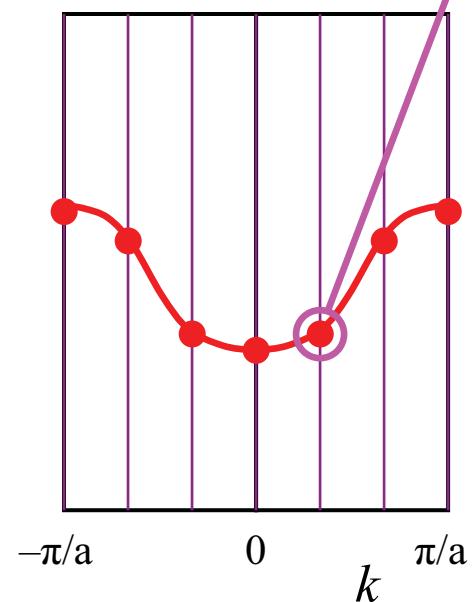


Bloch Theorem

Crystal in real space:

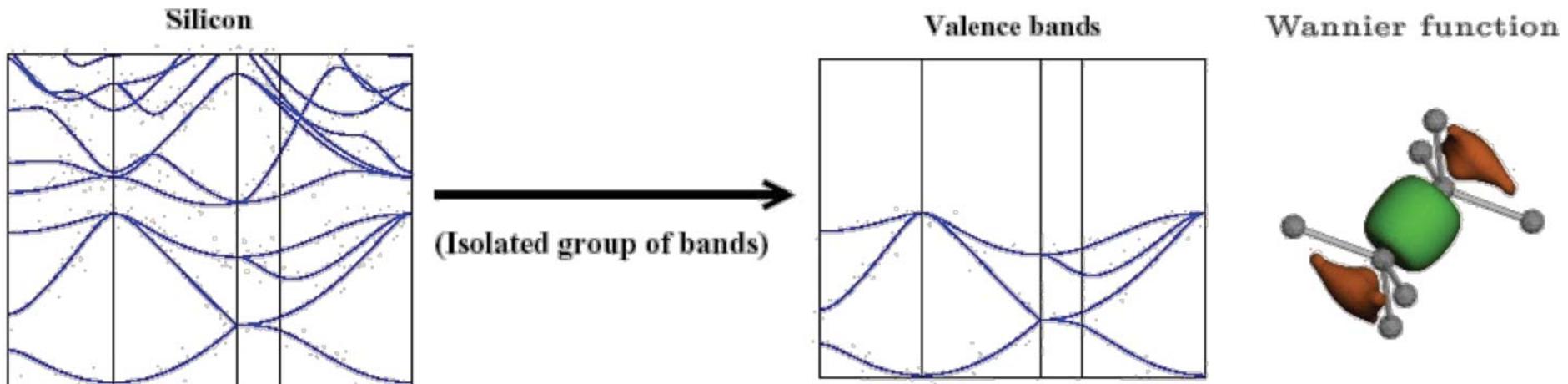


Brillouin zone in reciprocal space:



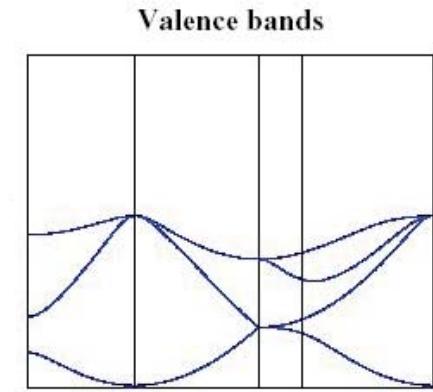
Generalized Wannier Functions for Composite Bands

- $\{|\mathbf{R}n\rangle\}$ span the same space as $\{|\Psi_{n\mathbf{k}}\rangle\}$
- $|\mathbf{R}n\rangle = w_n(\mathbf{r} - \mathbf{R})$ (translational images)
- $\langle \mathbf{R}n | \mathbf{R}'m \rangle = \delta_{n,m} \delta_{\mathbf{R},\mathbf{R}'}$
- “maximally” localized



From Bloch Orbitals to Wannier Functions

$$|\mathbf{R}n\rangle = \int_{BZ} \Psi_{n\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$



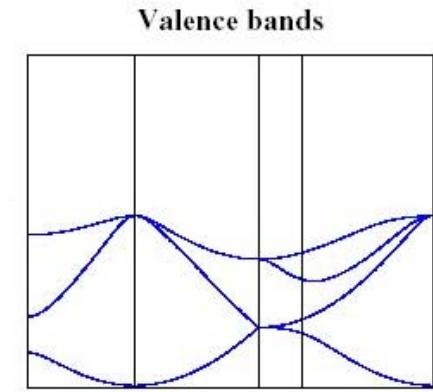
Gauge freedoms

- Arbitrary phase factor for every $n\mathbf{k}$ (Schrödinger)

$$|\mathbf{R}n\rangle = \int_{BZ} \left[e^{i\phi_n(\mathbf{k})} \psi_{n\mathbf{k}}(\mathbf{r}) \right] e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$

From Bloch Orbitals to Wannier Functions

$$|\mathbf{R}n\rangle = \int_{BZ} \Psi_{n\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$



Gauge freedoms

- Arbitrary phase factor for every $n\mathbf{k}$ (Schrödinger)
- Arbitrary unitary rotations $U_{mn}^{(\mathbf{k})}$ for every \mathbf{k} (DFT)

$$|\mathbf{R}n\rangle = \int_{BZ} \sum_m U_{mn}^{(\mathbf{k})} \Psi_{m\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$

The Localization Functional (Foster-Boys)

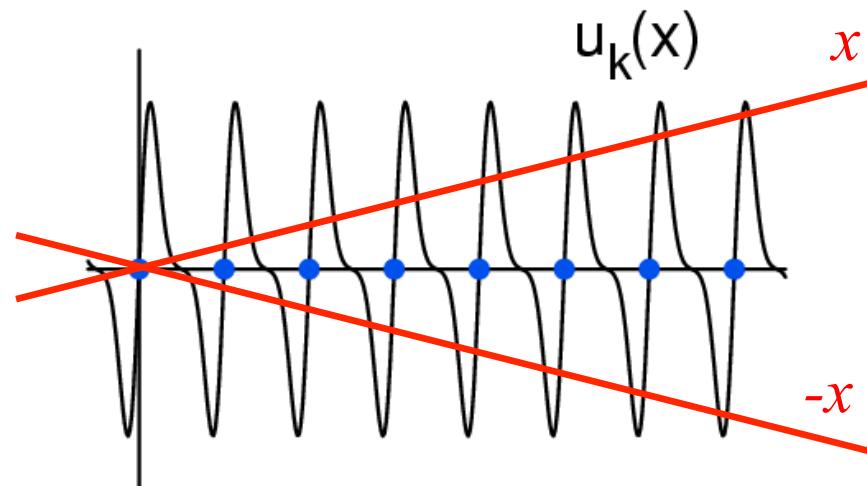
$$\Omega = \sum_n [\langle \mathbf{0}n | r^2 | \mathbf{0}n \rangle - \langle \mathbf{0}n | \mathbf{r} | \mathbf{0}n \rangle^2]$$

For a given set of Bloch orbitals, our goal is to minimize Ω with respect all the sets of unitary transformations $U_{mn}^{(k)}$

$$|\mathbf{R}n\rangle = \int_{BZ} \sum_m U_{mn}^{(k)} \Psi_{m\mathbf{k}}(\mathbf{r}) e^{-i\mathbf{k}\cdot\mathbf{R}} d\mathbf{k}$$

Position operator is ill defined !

$$\langle \psi_k | x | \psi_k \rangle = \int_{-\infty}^{\infty} x |u_k(x)|^2 dx$$



Blount identities

Position operator \Rightarrow Gradient

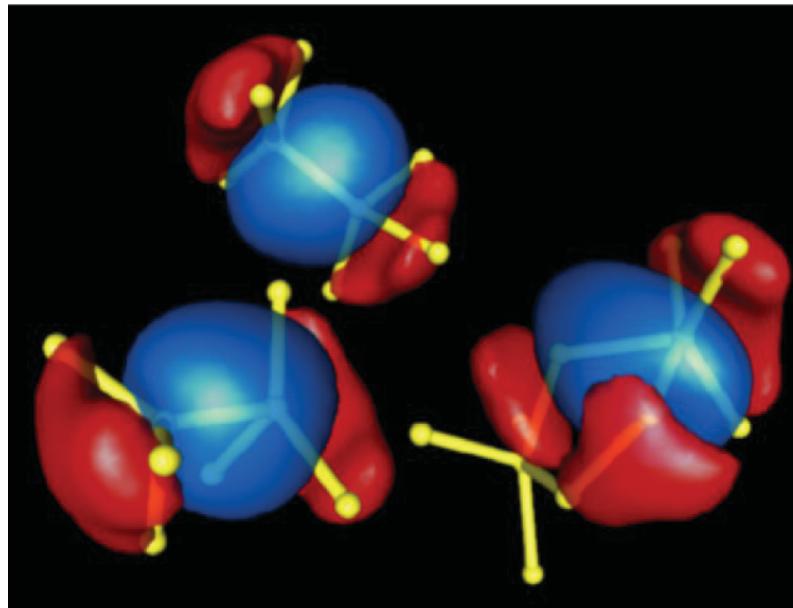
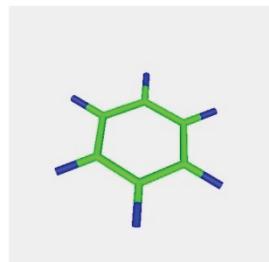
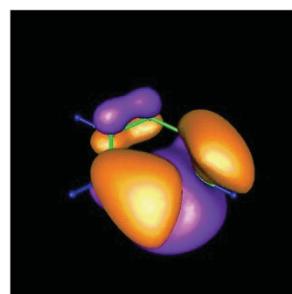
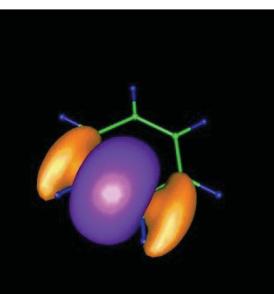
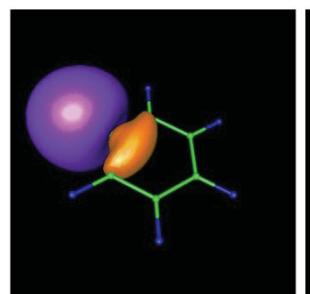
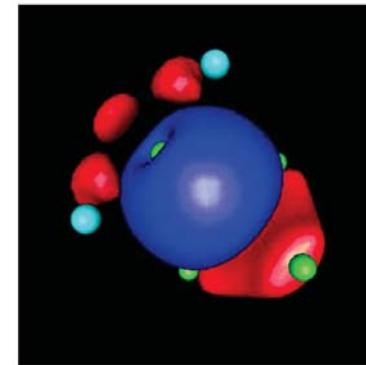
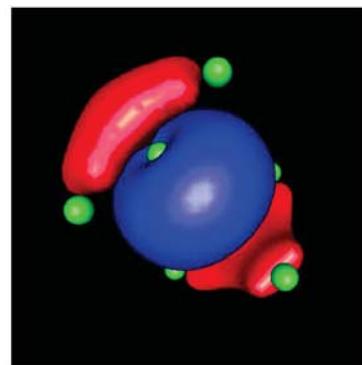
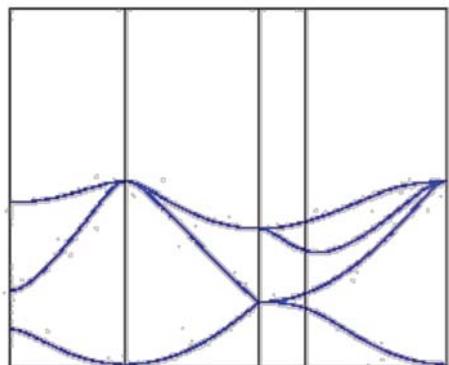
$$\langle \mathbf{0}n | \mathbf{r}^l | \mathbf{0}m \rangle = \frac{V}{(2\pi)^3} \int \langle u_{n\mathbf{k}} | \left(i \frac{\partial}{\partial \mathbf{k}} \right)^l | u_{m\mathbf{k}} \rangle d\mathbf{k}$$

We can then express positions and spreads as a function of
the phase relations between neighboring Bloch orbitals

$$M_{mn}^{(\mathbf{k}, \mathbf{b})} = \langle u_{m\mathbf{k}} | u_{n,\mathbf{k}+\mathbf{b}} \rangle$$

Silicon, GaAs, Amorphous Silicon, Benzene

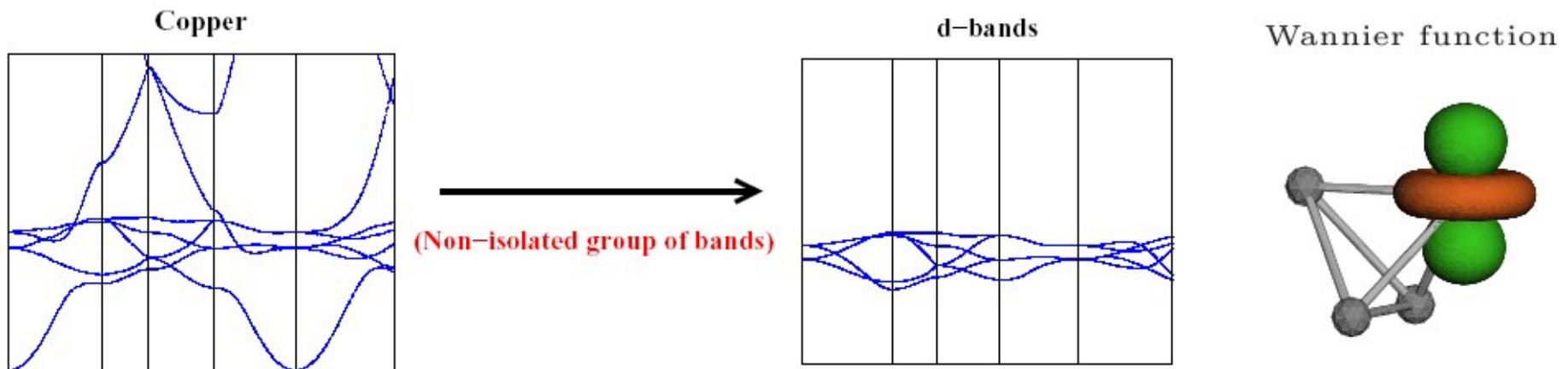
Valence bands



M. Fornari, N. Marzari, M. Peressi, and A. Baldereschi, Comp. Mater. Science 20, 337 (2001)

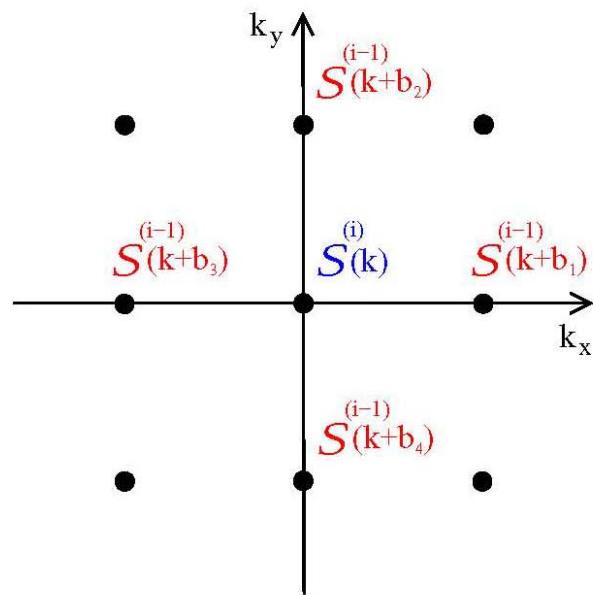
Disentanglement of Attached Bands

- Maximally-localized Wannier-like functions for conduction subspace
- Extract differentiable manifold with **optimal smoothness**



I. Souza, N. Marzari and D. Vanderbilt, Phys. Rev. B 65, 035109 (2002)

Iterative Minimization of Ω_I



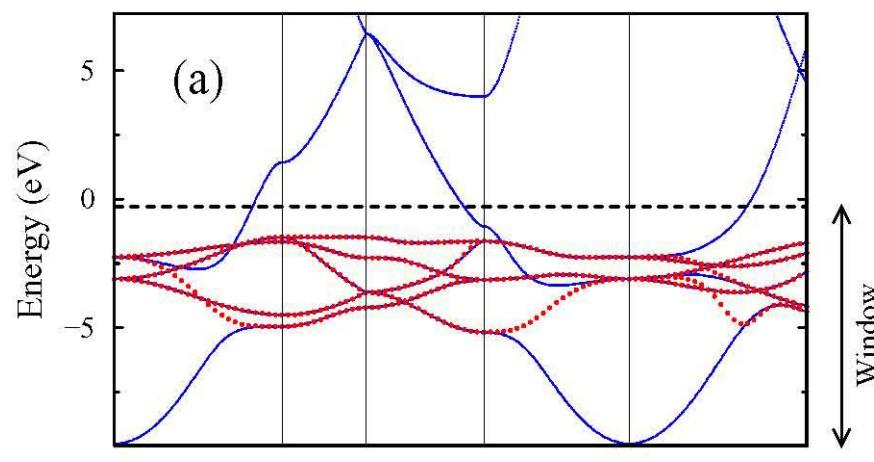
Minimize $\frac{\text{degree of mismatch}}{\text{between } \mathcal{S}^{(i)}(\mathbf{k}) \text{ and } \mathcal{S}^{(i-1)}(\mathbf{k} + \mathbf{b})}$, i.e.,
 maximize overlap $\sum_{\mathbf{b}} \sum_{m=1}^N \left| \langle u_{n\mathbf{k}}^{(i)} | u_{m,\mathbf{k}+\mathbf{b}}^{(i-1)} \rangle \right|^2$

- 1st iteration: Choose trial subspace at each \mathbf{k} (e.g. projected orbitals)
- i^{th} iteration: At each \mathbf{k} pick the N highest eigenvectors of

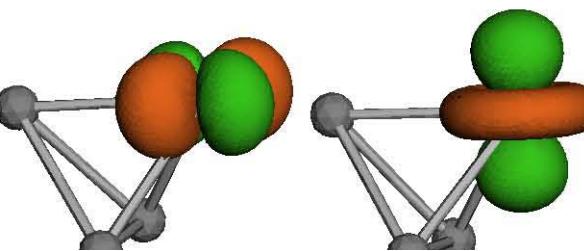
$$\left[\sum_{\mathbf{b}} \hat{P}_{\mathbf{k}+\mathbf{b}}^{(i-1)} \right] |u_{n\mathbf{k}}^{(i)}\rangle = \lambda_{n\mathbf{k}}^{(i)} |u_{n\mathbf{k}}^{(i)}\rangle \quad \hat{P}_{\mathbf{k}+\mathbf{b}}^{(i-1)} : \text{Projector onto } \mathcal{S}^{(i-1)}(\mathbf{k} + \mathbf{b})$$

- Repeat until self-consistency (when spaces $\mathcal{S}(\mathbf{k})$ stabilize)

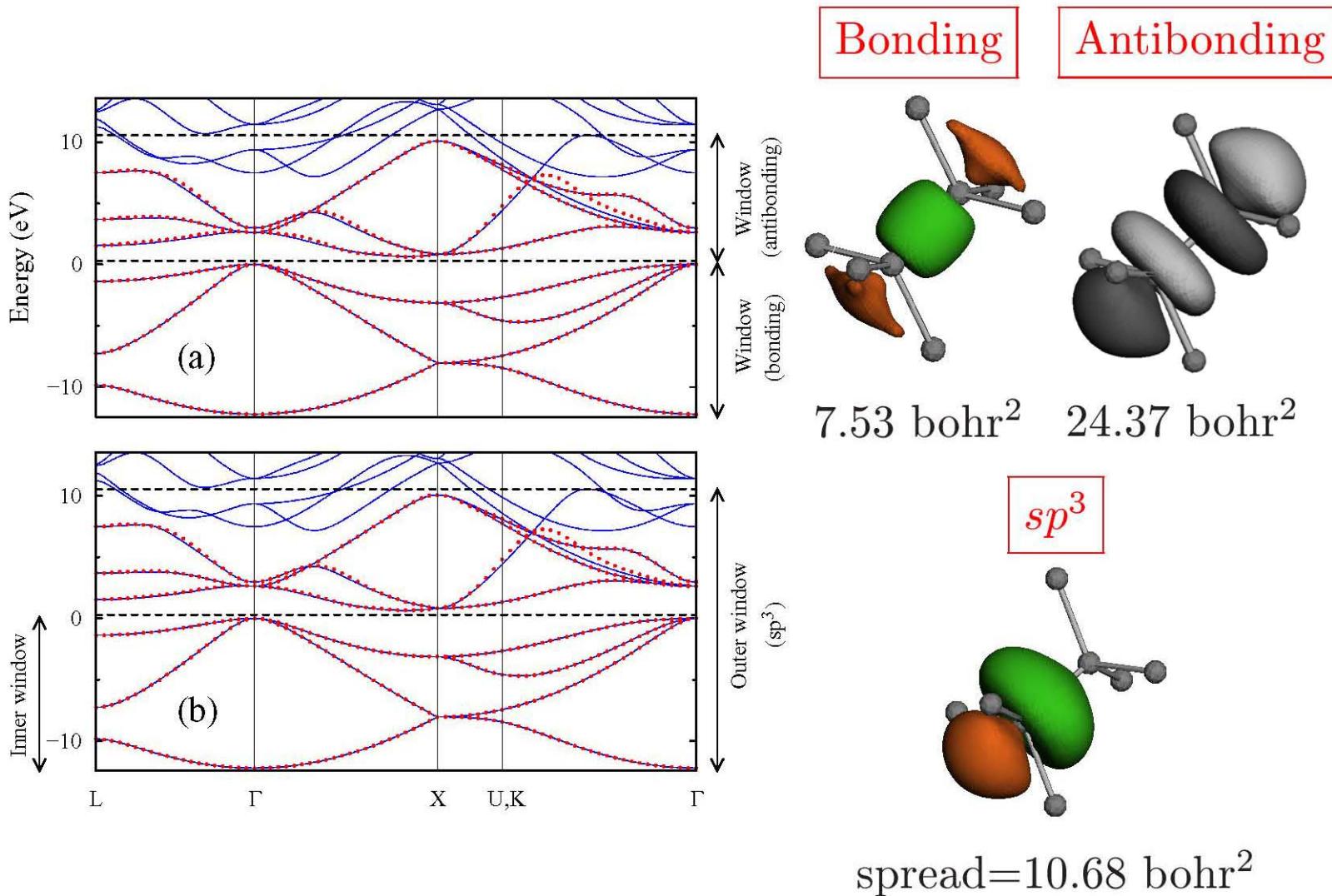
d Bands of Copper



The e_g *d* WFs of panel (b)

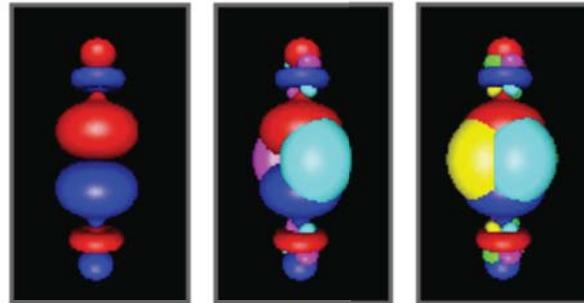


Silicon: Bonding and Antibonding Orbitals



Welcome!

This is the home of maximally-localised Wannier functions (MLWFs) and Wannier90, the computer program that calculates them. Wannier90 is released under the [GNU General Public License](#).

**Latest News****15 JAN 2010**

Wannier90 (v1.2) is now available for download [here](#).

**Wannier90 seamlessly interfaced to
Quantum-ESPRESSO, Abinit, Siesta, FLEUR ...**

See [here](#) for our news archive.

Please cite

Wannier90: A Tool for Obtaining Maximally-Localised Wannier Functions
A. A. Mostofi, J. R. Yates, Y.-S. Lee, I. Souza, D. Vanderbilt and N. Marzari
Comput. Phys. Commun. **178**, 685 (2008) [[ONLINE JOURNAL](#)]

in all publications resulting from your use of Wannier90.



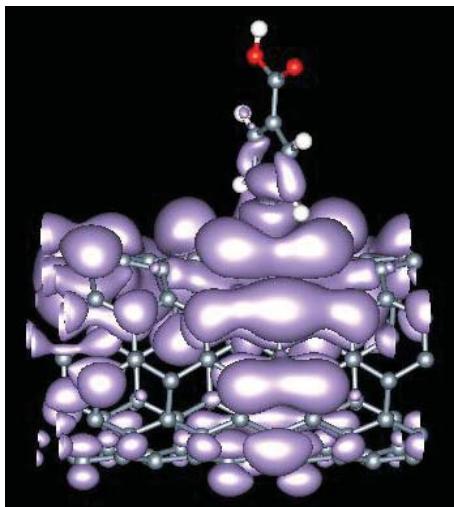
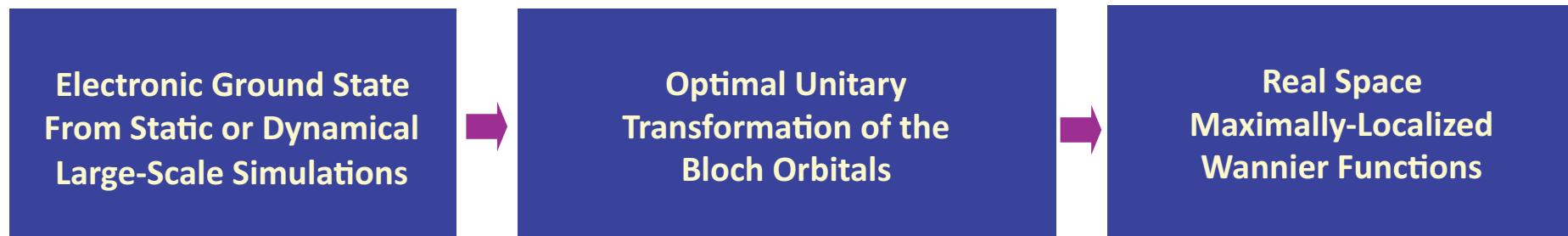
**Imperial College
London**



THE STATE UNIVERSITY OF NEW JERSEY
RUTGERS

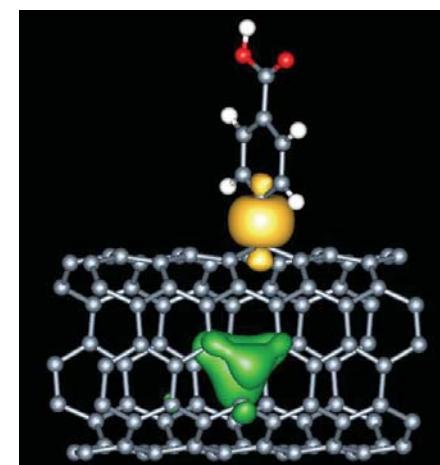


Electronic Structure of Large Nanostructures



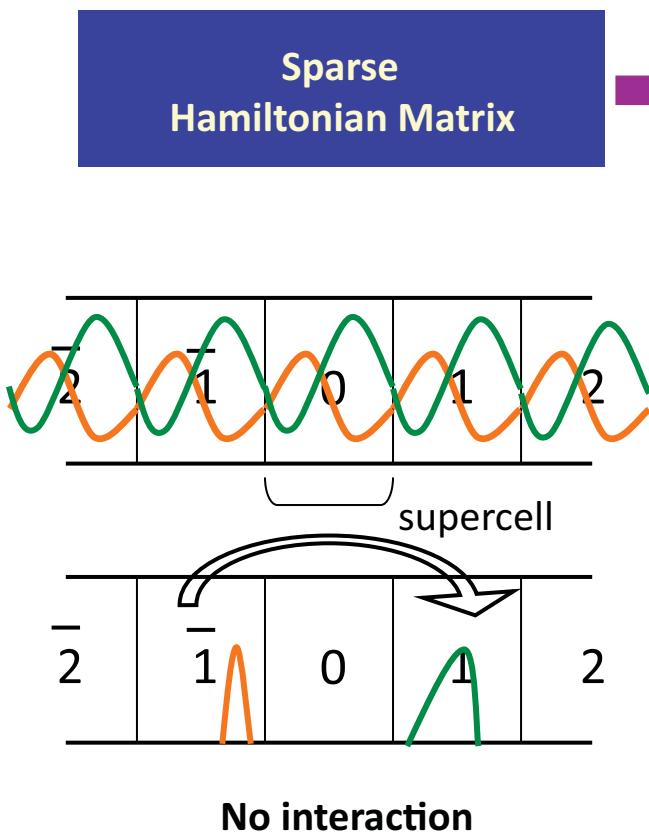
Minimization of the spread functional

$$\Omega = \sum_n [\langle r^2 \rangle_n - \langle \mathbf{r} \rangle_n^2]$$



N. Marzari and D. Vanderbilt, Phys. Rev. B 56, 12847 (1997)

Electronic Structure of Large Nanostructures



Green's Function
Transmission Function

Ballistic Conductance
Density of States

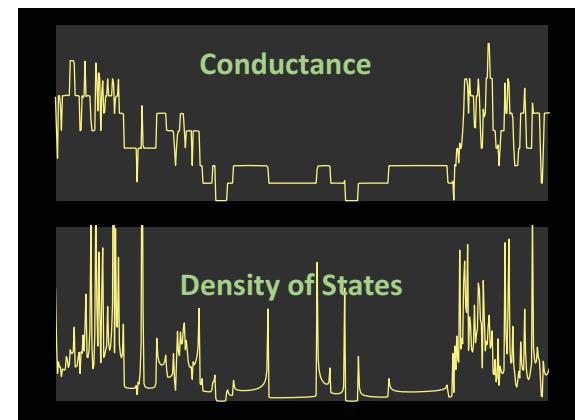
$$G(E) = \frac{2e^2}{h} T(E)$$

$$N(E) = -(1/\pi) \text{Im}[\text{Tr} G_C^r(E)]$$

$$H_{00}, H_{01} \Rightarrow G_C^r(E)$$

$$T(E) = \text{Tr}(\Gamma_L G_C^r \Gamma_R G_C^a)$$

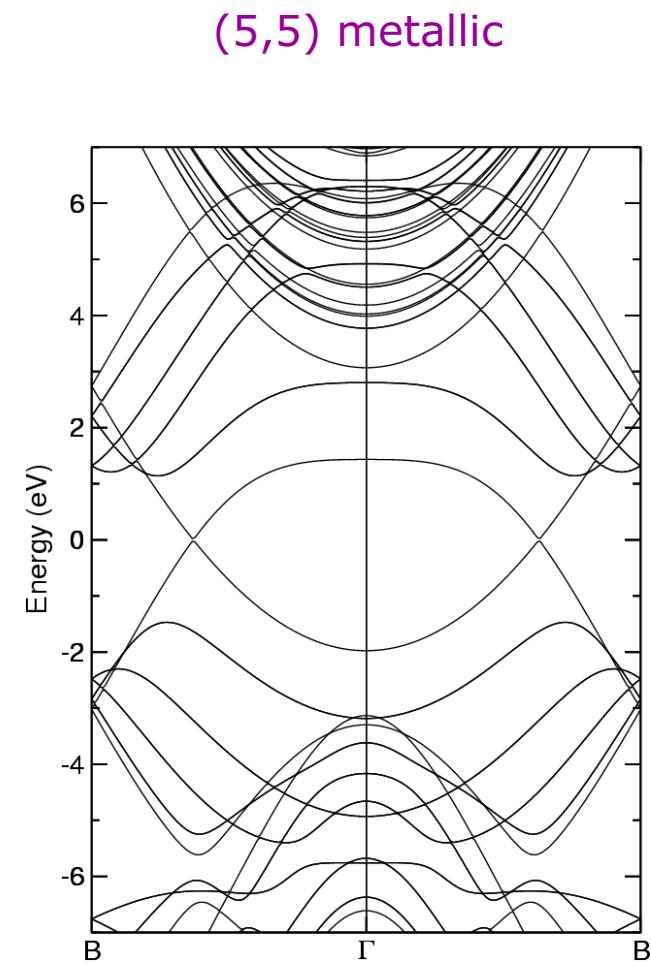
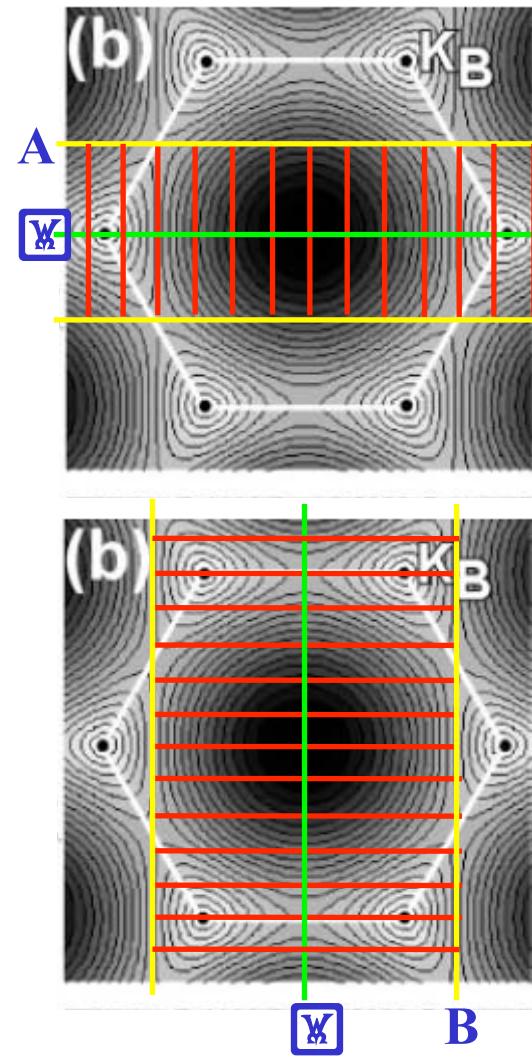
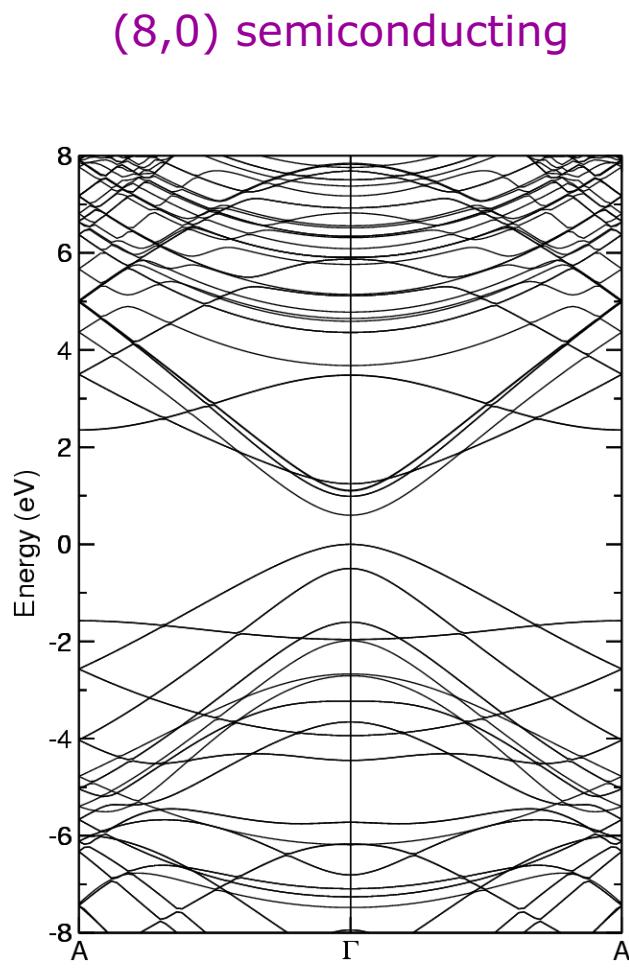
$$H = \begin{pmatrix} & & & & 0 \\ H_{01}^+ & H_{00} & H_{01} & & \\ & H_{01}^+ & H_{00} & H_{01} & \\ & & H_{01}^+ & H_{00} & H_{01} \\ 0 & & & & \end{pmatrix}$$



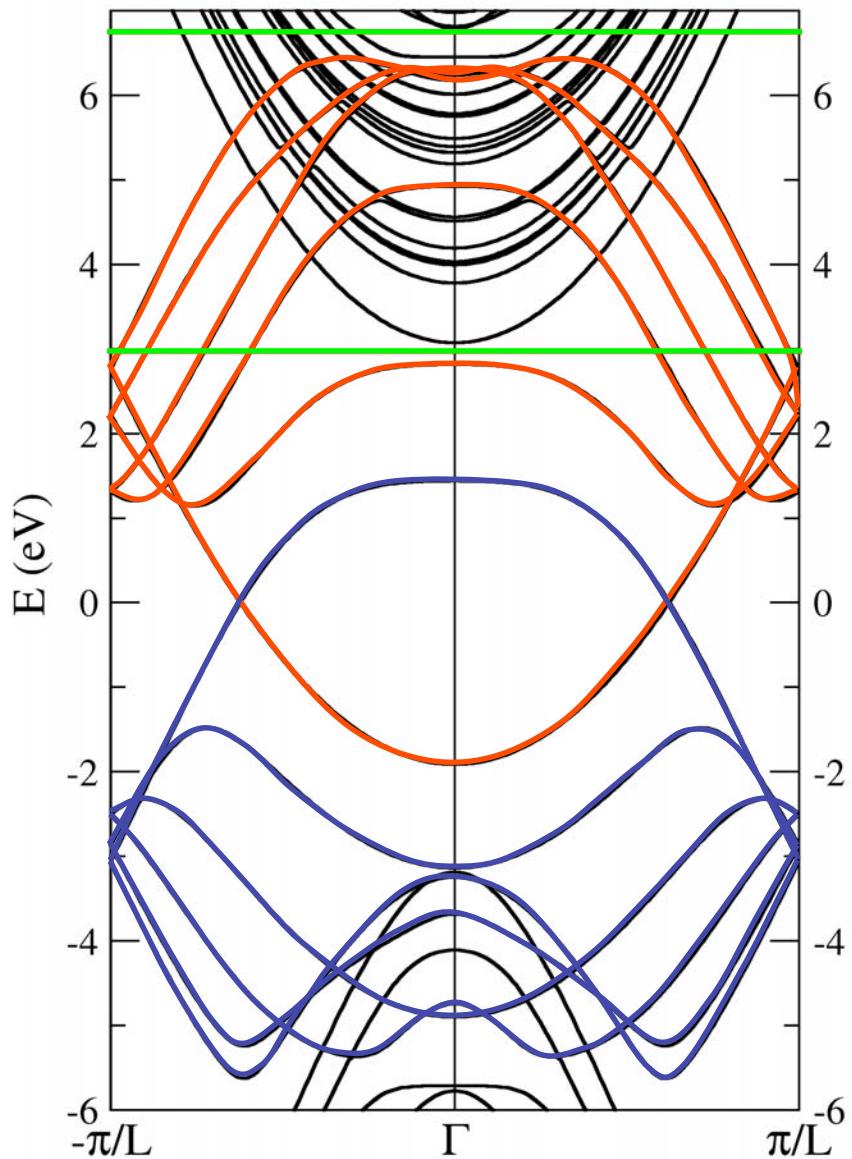
M. B. Nardelli, Phys. Rev. B 60, 7828 (1999)

A. Calzolari, N. Marzari, I. Souza, M. B. Nardelli, Phys. Rev. B 69, 035108 (2004)

Band Structure of (8,0) and (5,5) SWNT



Disentanglement: Conduction Bands in (5,5) SWNT



MLWFs from the Disentangled Subspace

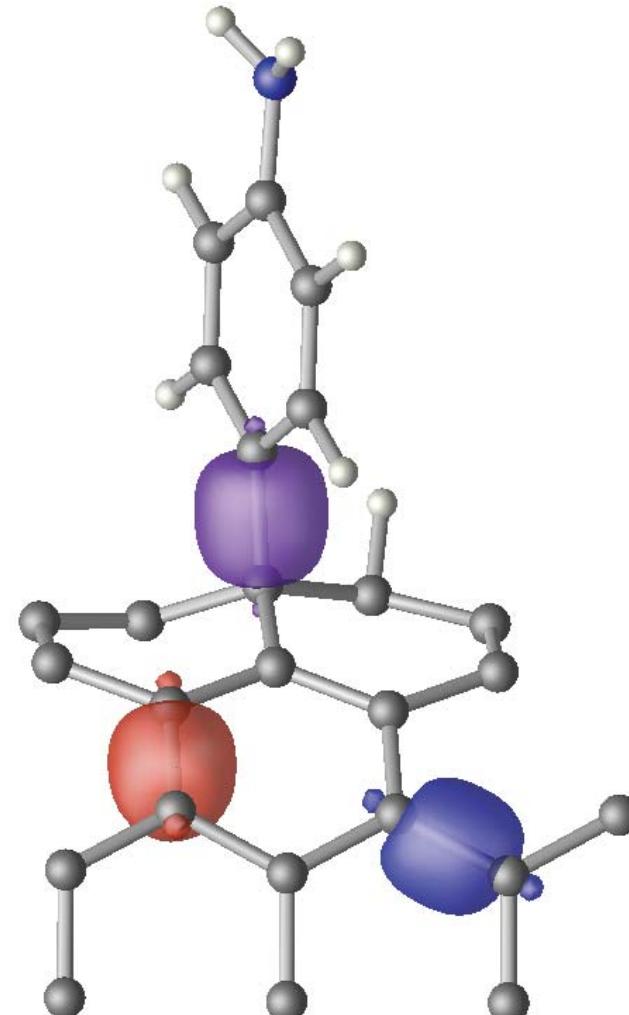
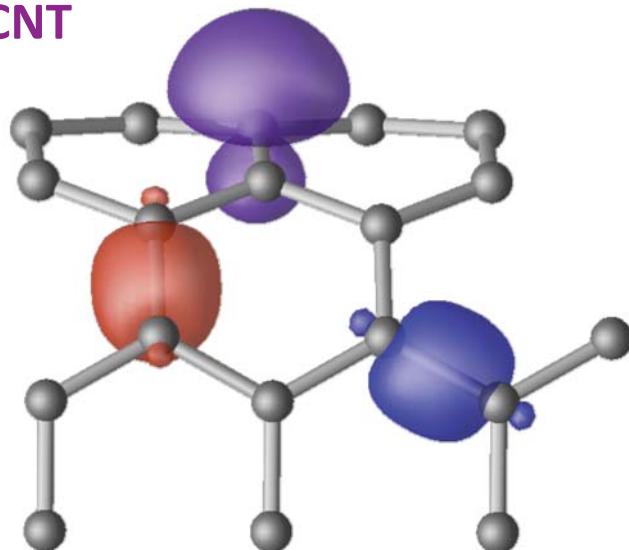
- Localization after disentanglement

(5,5) CNT + Nitrophenyl

$$w_n = \sum_{m=1}^N U_{mn} \psi_m$$

- MLWFs from the disentangled subspace
 - s bond orbitals + p orbitals

(5,5) CNT

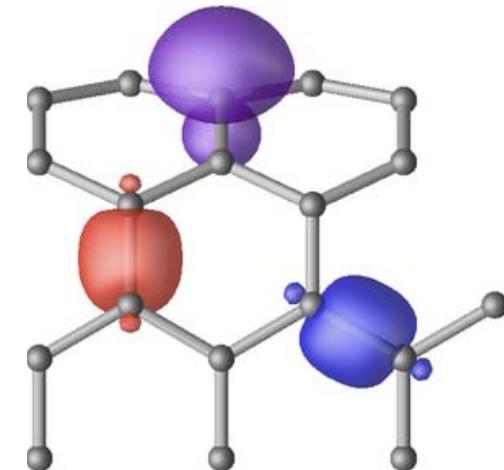


Max-loc WFs \leftrightarrow “Exact” Tight-Binding

Compact mapping of Bloch states into local orbitals

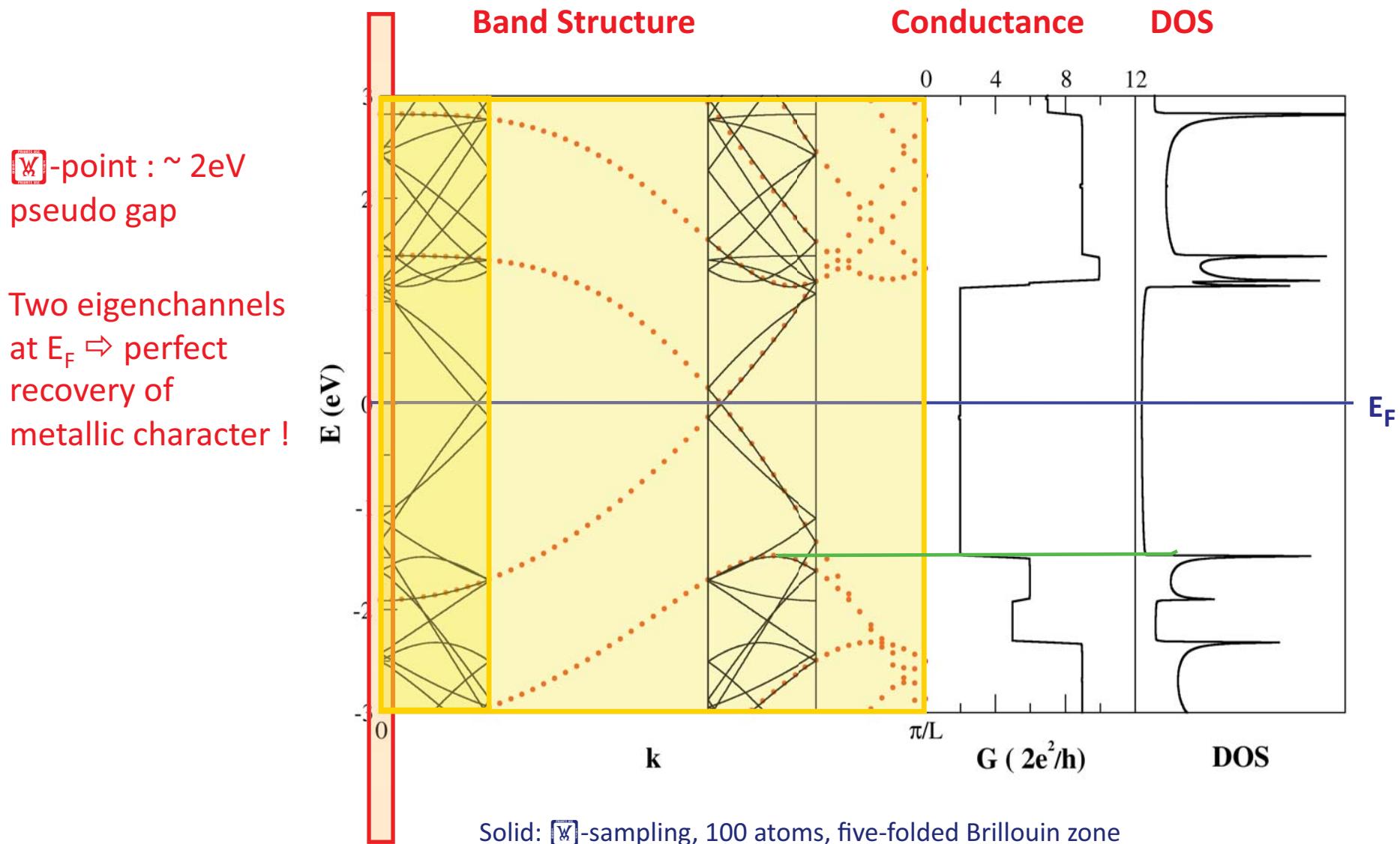
$$\omega_n(\mathbf{r} - \mathbf{R}) = \frac{V}{8\pi^3} \int_{BZ} e^{-i\mathbf{k}\cdot\mathbf{R}} \psi_{n\mathbf{k}}(\mathbf{r}) d\mathbf{k}$$

$$\psi_{n\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{N_R}} \sum_R e^{i\mathbf{k}\cdot\mathbf{R}} \omega_n(\mathbf{r} - \mathbf{R})$$



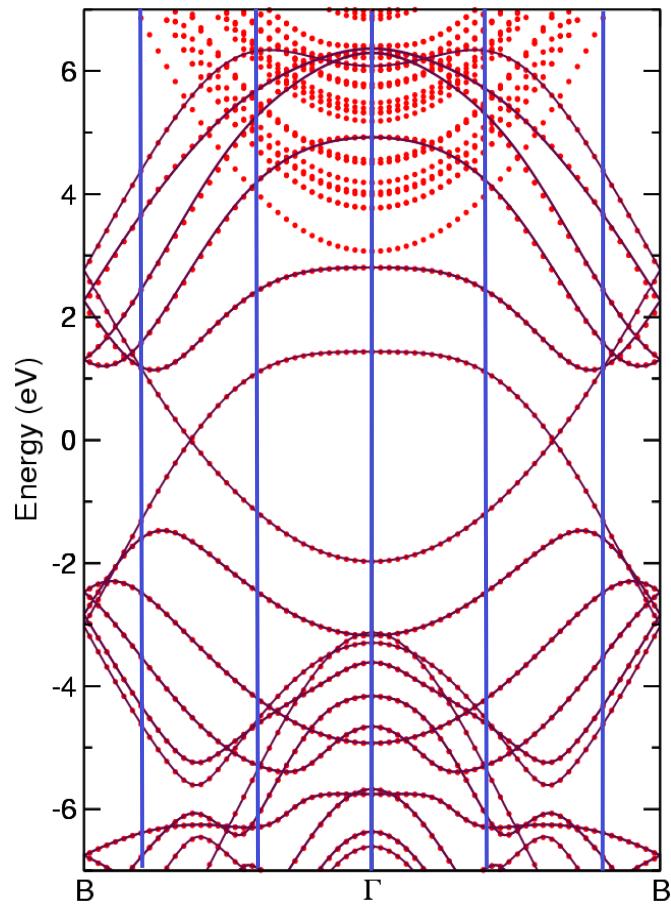
$$\langle \psi_{i\mathbf{k}} | \hat{H} | \psi_{j\mathbf{k}} \rangle = H_{ij}^{00} + e^{i\mathbf{k}\cdot\mathbf{R}} H_{ij}^{01} + e^{-i\mathbf{k}\cdot\mathbf{R}} H_{ij}^{0\bar{1}} \quad \Rightarrow \text{Diagonalize H Matrix}$$

Band Structure and Conductance of a (5,5) SWCNT

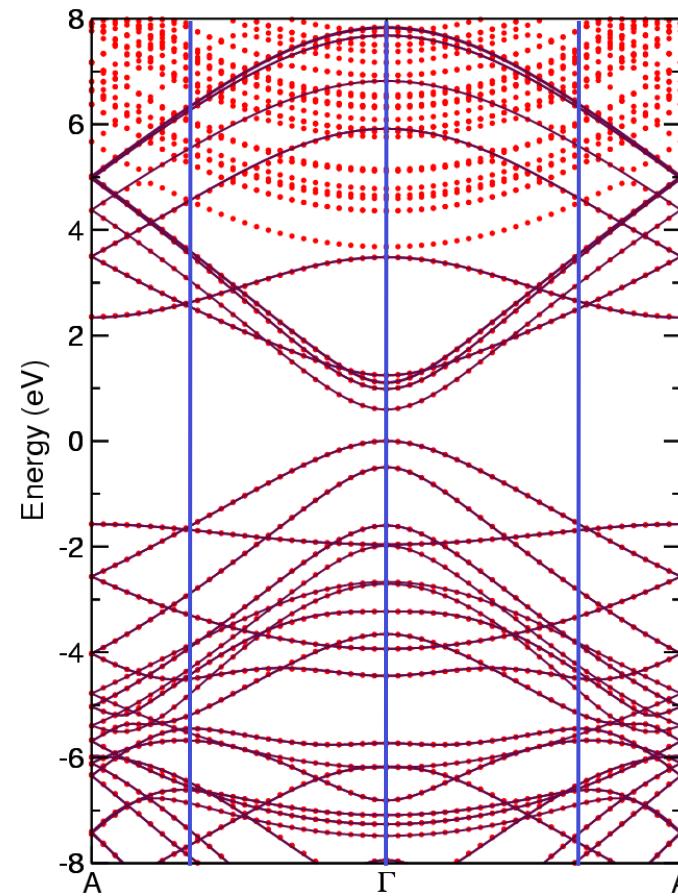


Exact Mapping onto a Tight-Binding Hamiltonian

(5,5) SWCNT

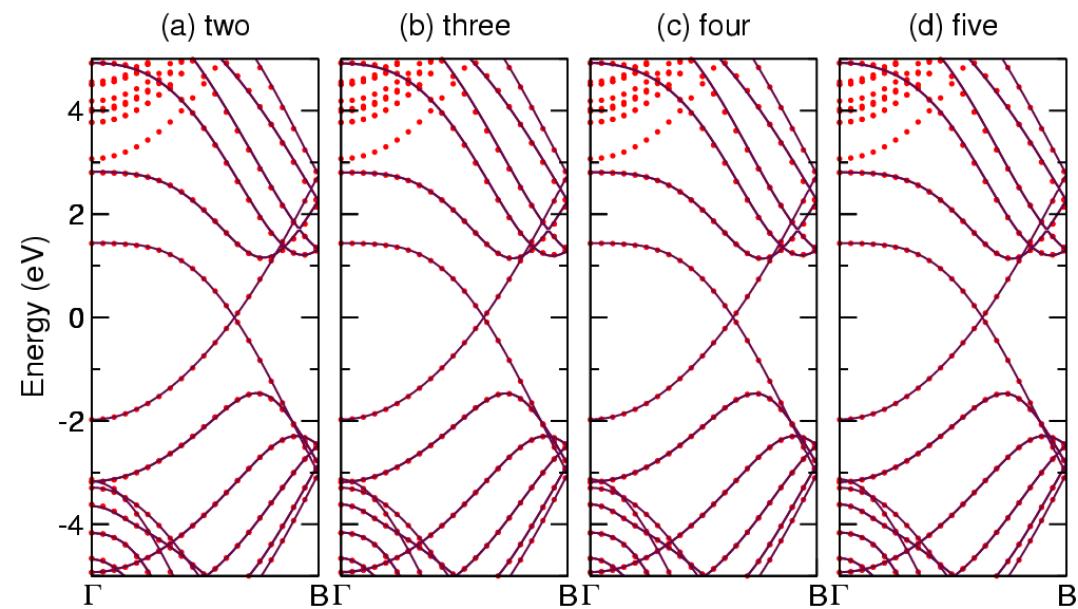
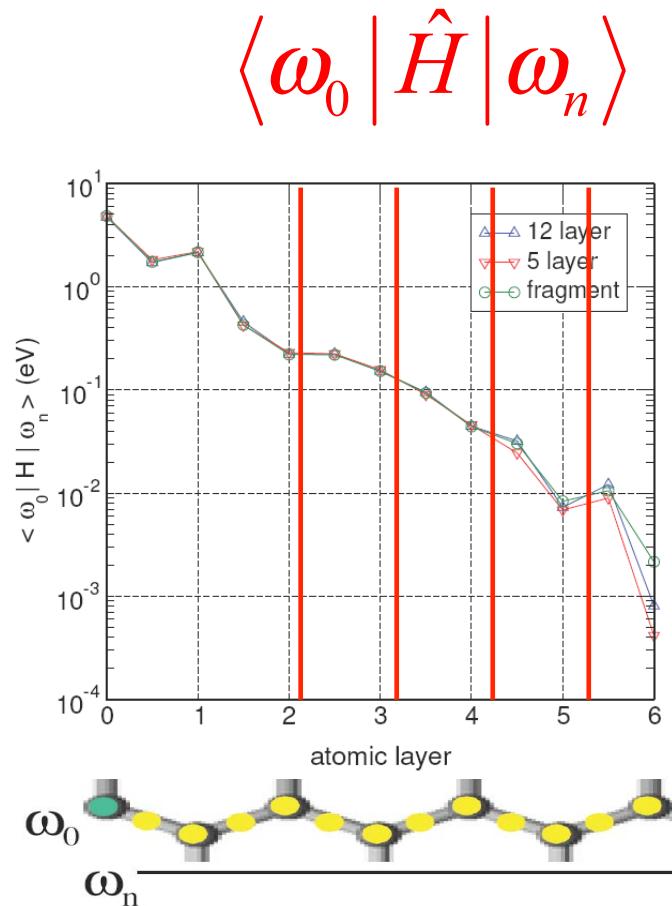


(8,0) SWCNT



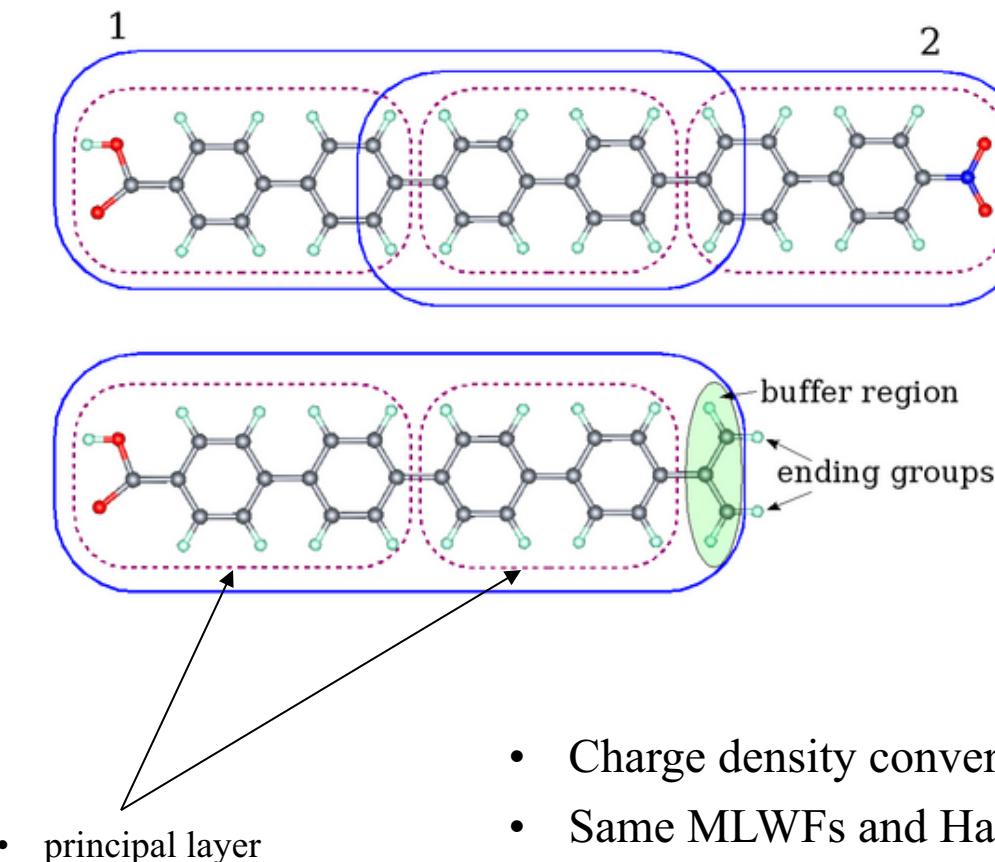
Exponential Decay

Even in a metal, we have smoothly connected manifolds – no relation with the physical decay of the density matrix



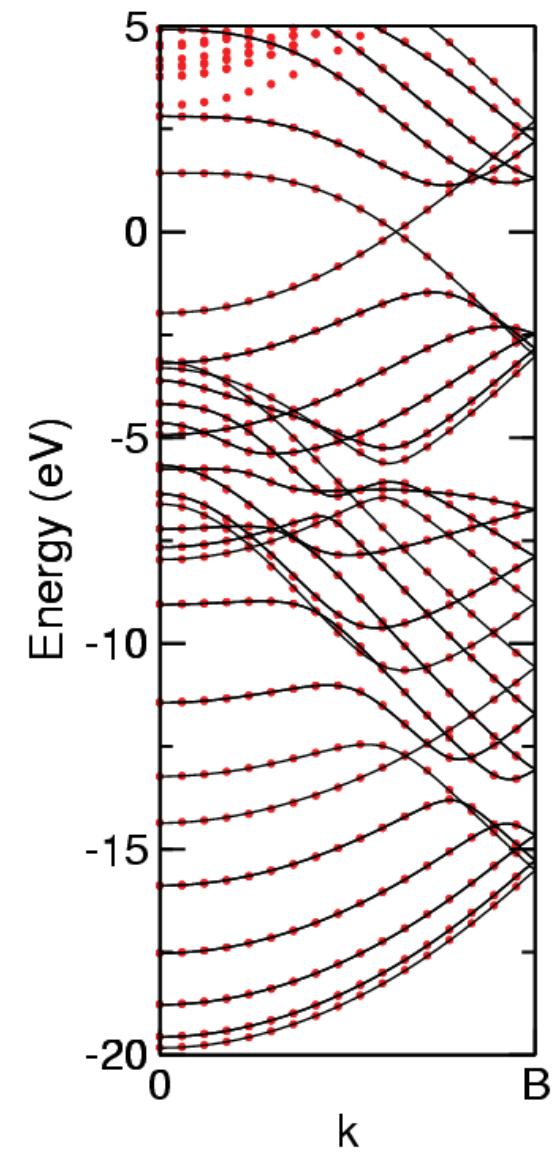
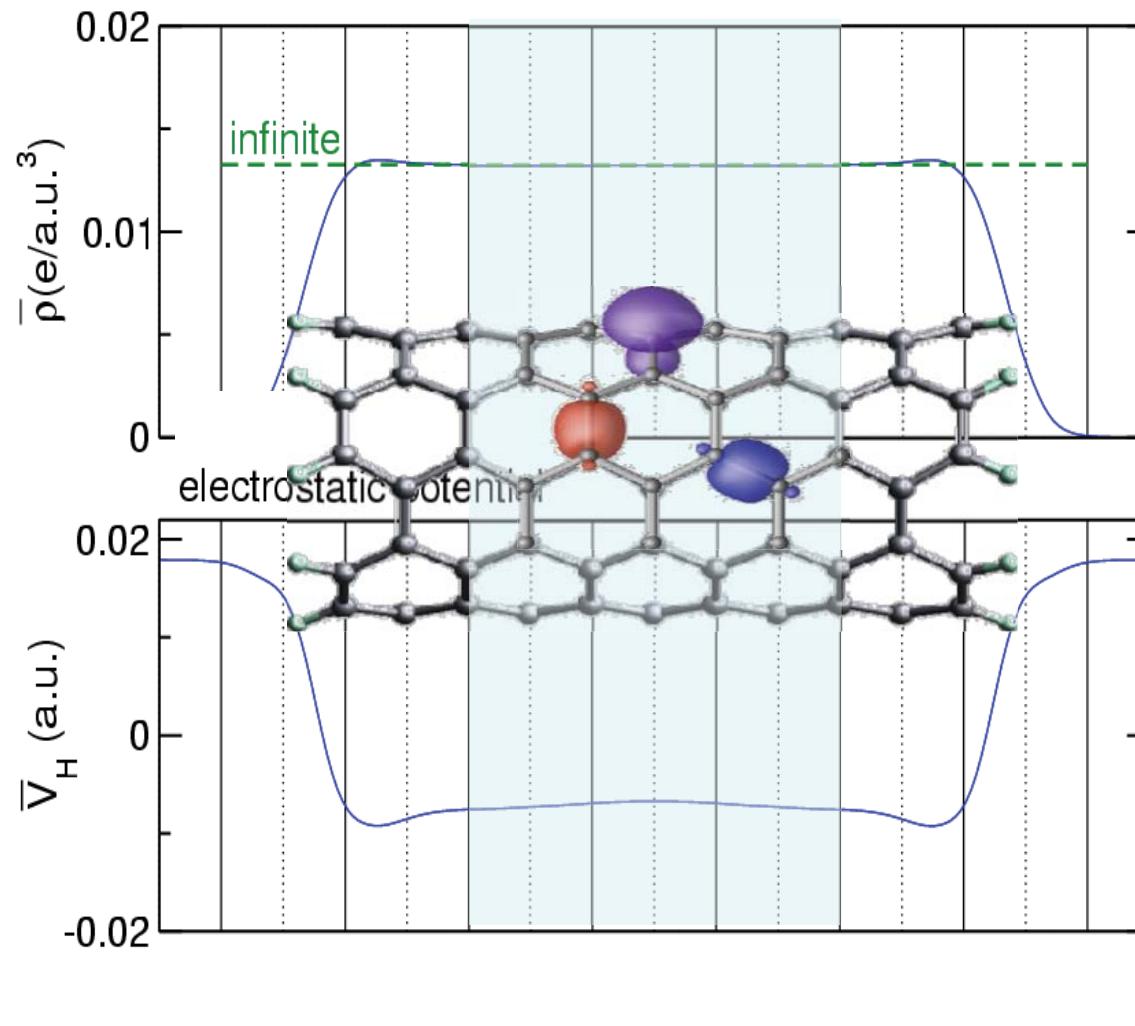
Large Scale Calculations

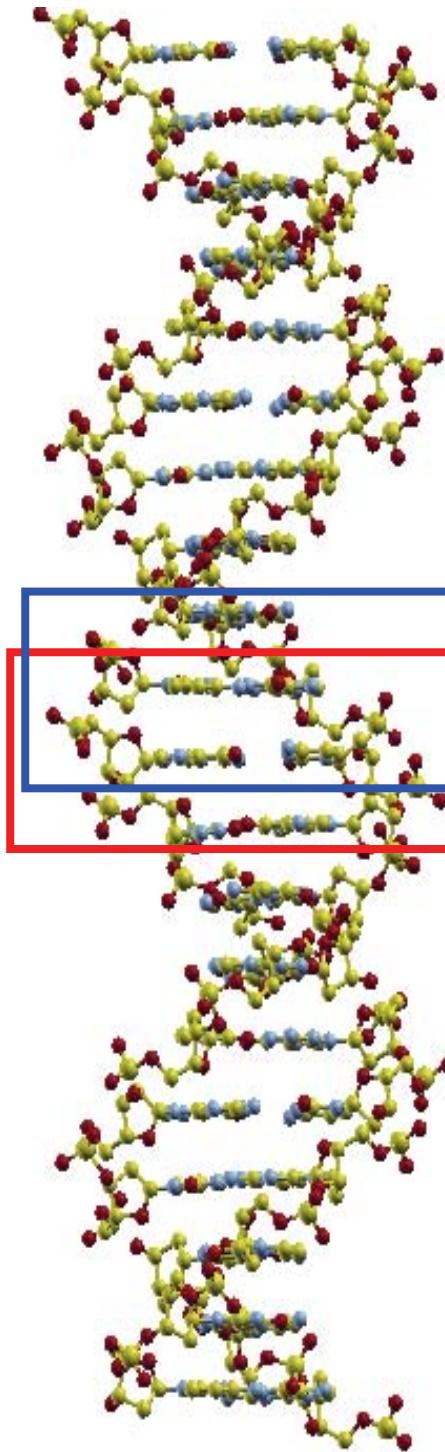
- Parameterization



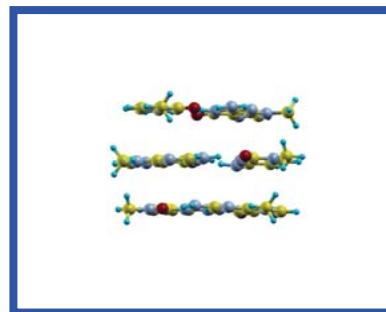
- principal layer
- Charge density convergence
- Same MLWFs and Hamiltonian matrix elements in overlapping region
- Electronic structure of long 1-D structure with a little cost

MLWFs Extraction From a Saturated Cluster

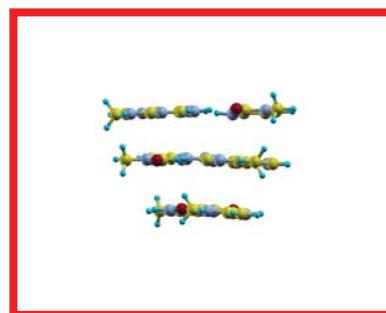




polyd(CpGp) B-DNA



$$\begin{aligned}A &= \mathbf{G} - \mathbf{C} \\B &= \mathbf{C} - \mathbf{G}\end{aligned}$$



First-principles DFT, \mathbb{W} -Point only Ultrasoft pseudopotentials GGA exchange-correlation functional

**B
A
B**

WANNIER FUNCTIONS

$$\mathbf{H} = \begin{pmatrix} \ddots & \vdots & 0 & 0 & 0 \\ \cdots & AA & AB & 0 & 0 \\ 0 & (AB)^\dagger & BB & BA & 0 \\ 0 & 0 & (BA)^\dagger & AA & \cdots \\ 0 & 0 & 0 & \vdots & \ddots \end{pmatrix}$$

**A
B
A**

WANNIER FUNCTIONS

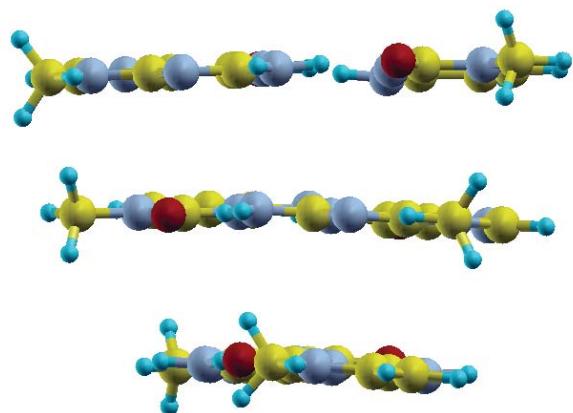
Optimal model Hamiltonian for infinite system

B-DNA Configuration

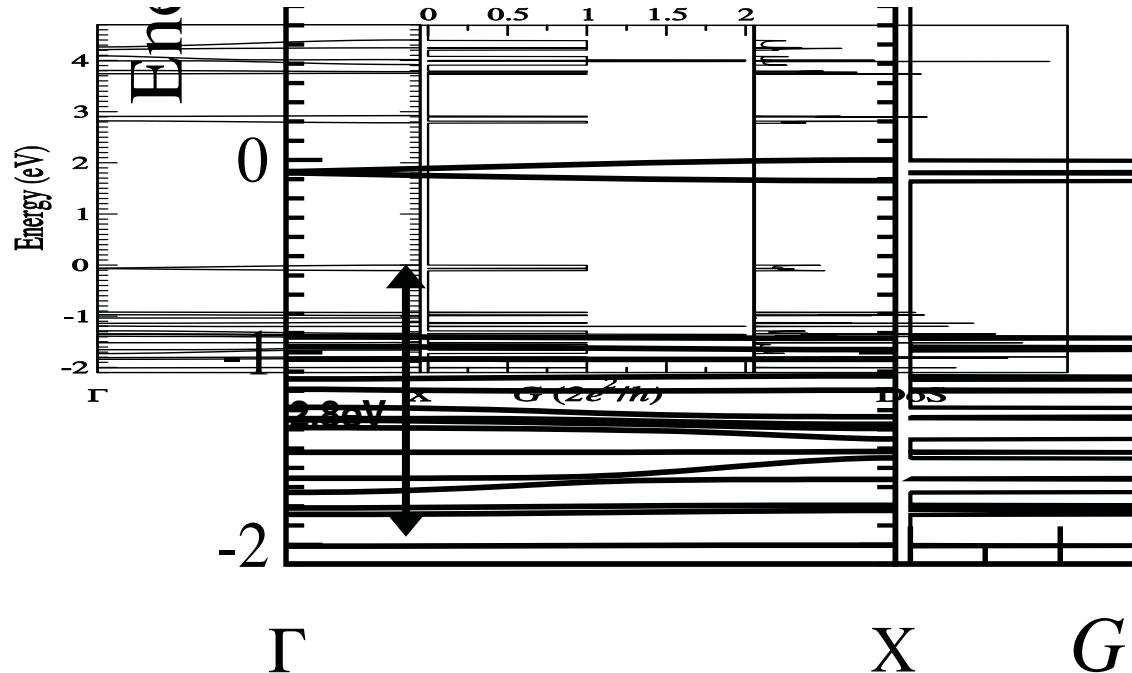
Dry, no backbone, 36° twist

Isolated fragments

- 105 atoms
- 330 electrons
- 18 Å cubic cell
- 186 WFs



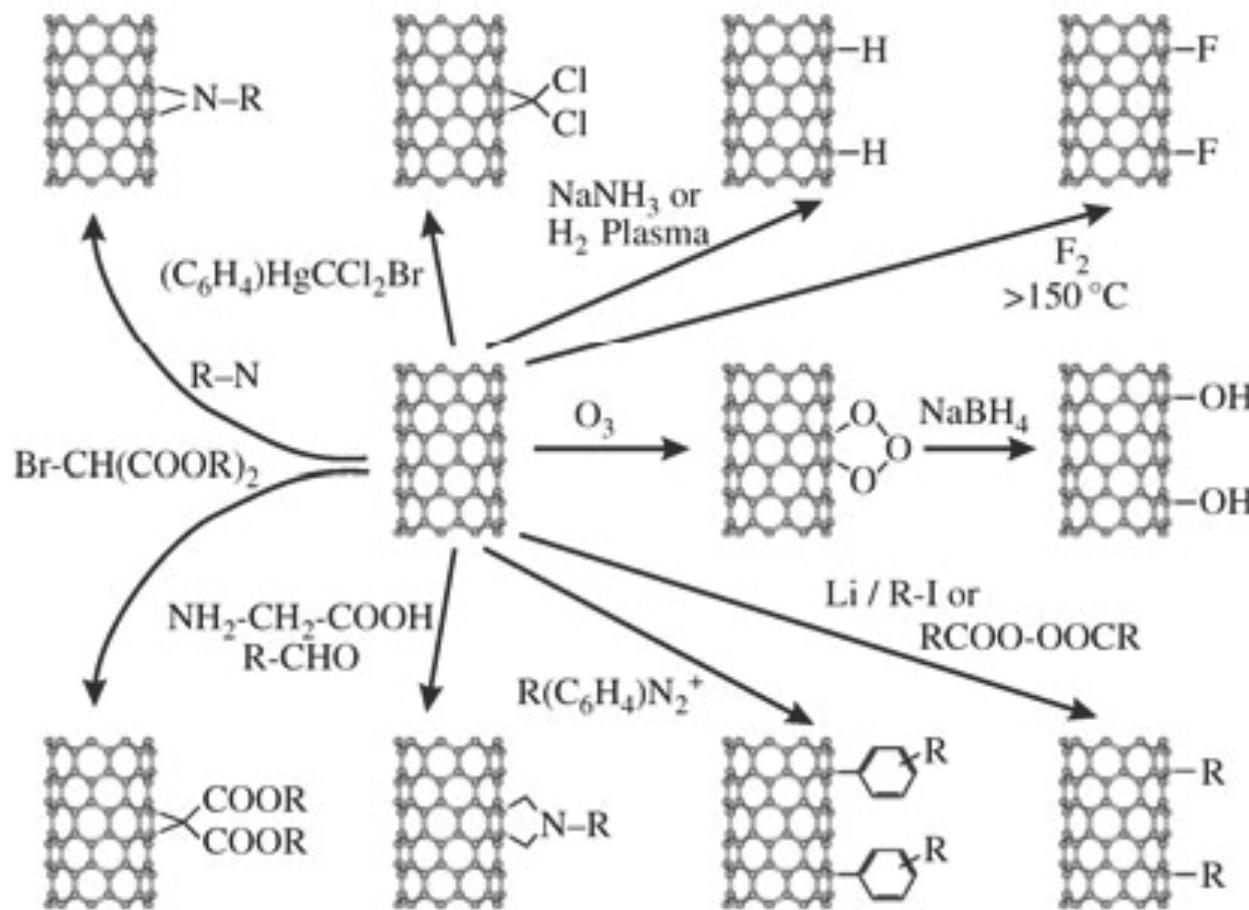
Equivalent to 350 atom calculation



LUMO bandwidth = 37 meV
HOMO bandwidth = 60 meV

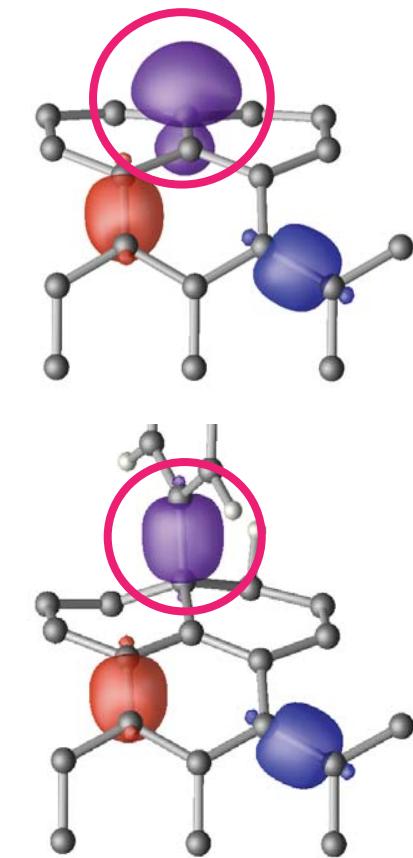
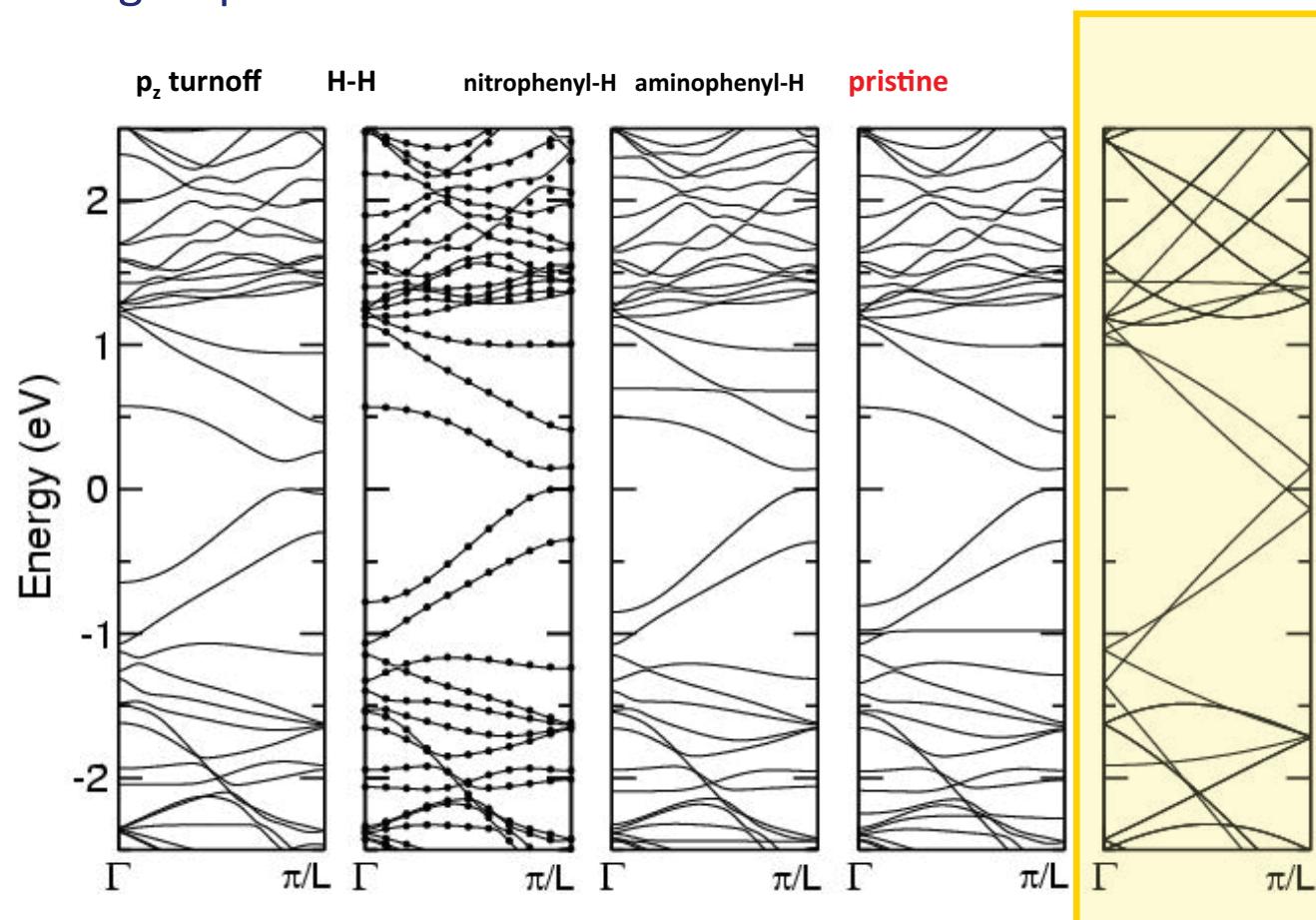
First-Principles Study of Functionalizations

- Currently available covalent functionalizations

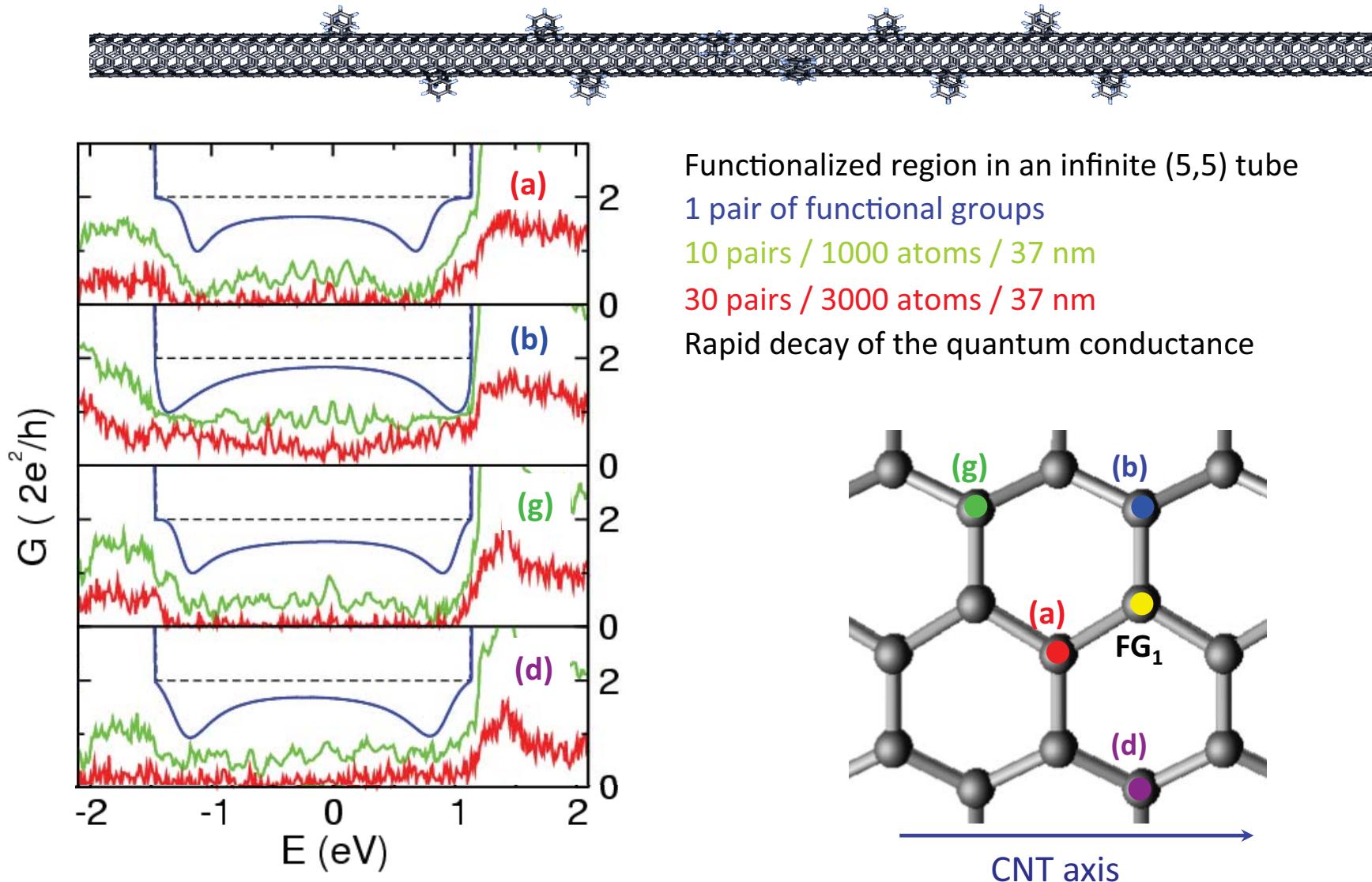


Aryl / Hydrogen

Band structure does not depend on the chemical nature of the functional groups



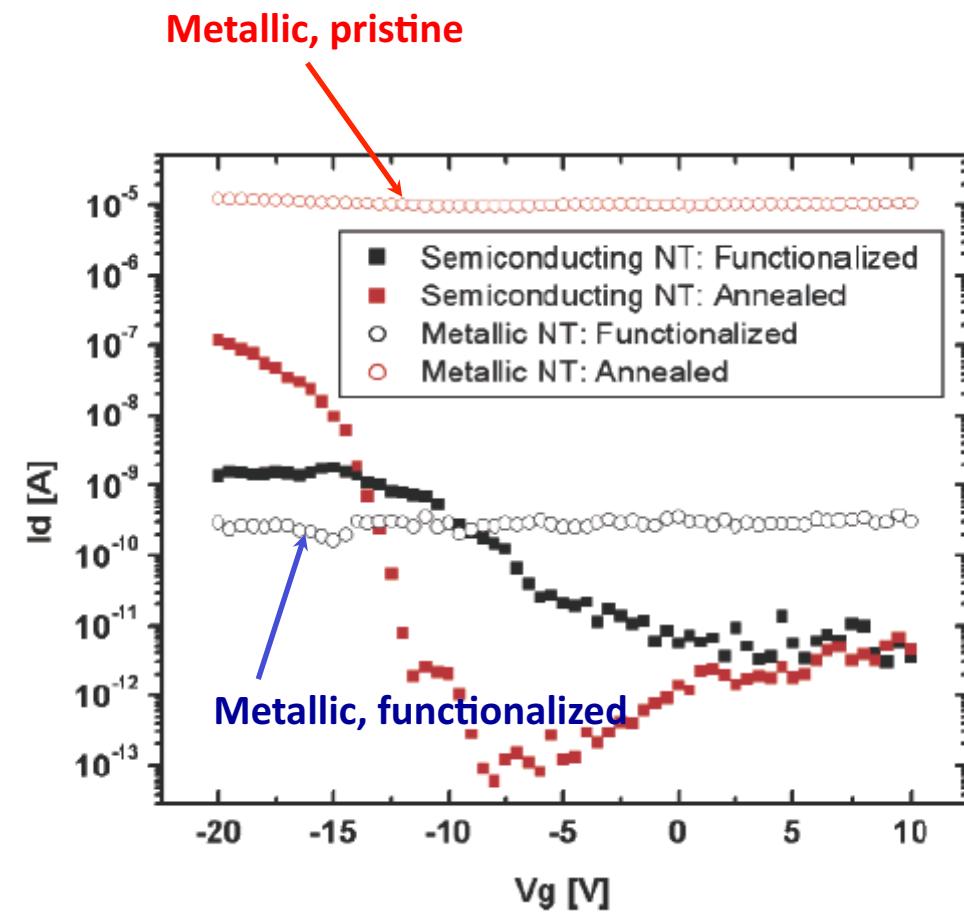
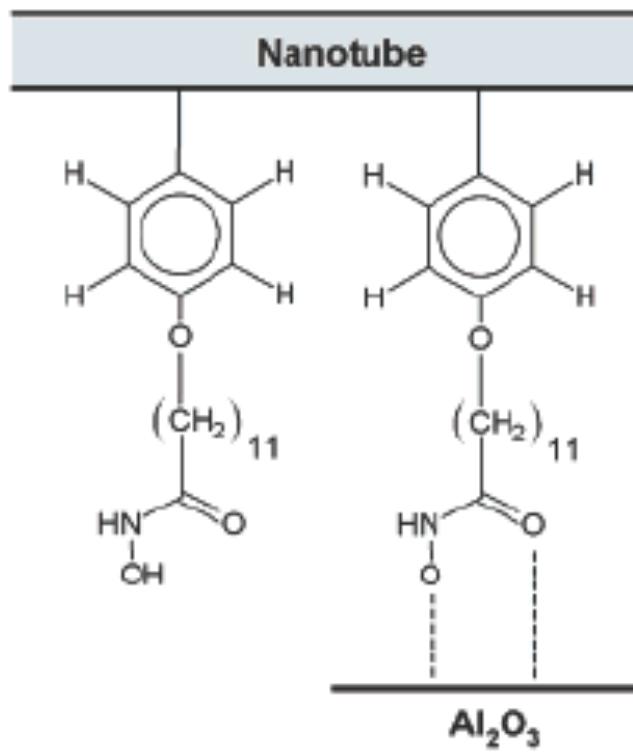
Aryl / Hydrogen - Quantum Conductance



Y.-S. Lee, M. Buongiorno Nardelli, and N. Marzari, Phys. Rev. Lett. 95, 076804 (2005)

Electrical Transport Measurements

IBM/Avouris: conductance decreases 5 orders of magnitude after aryl functionalization



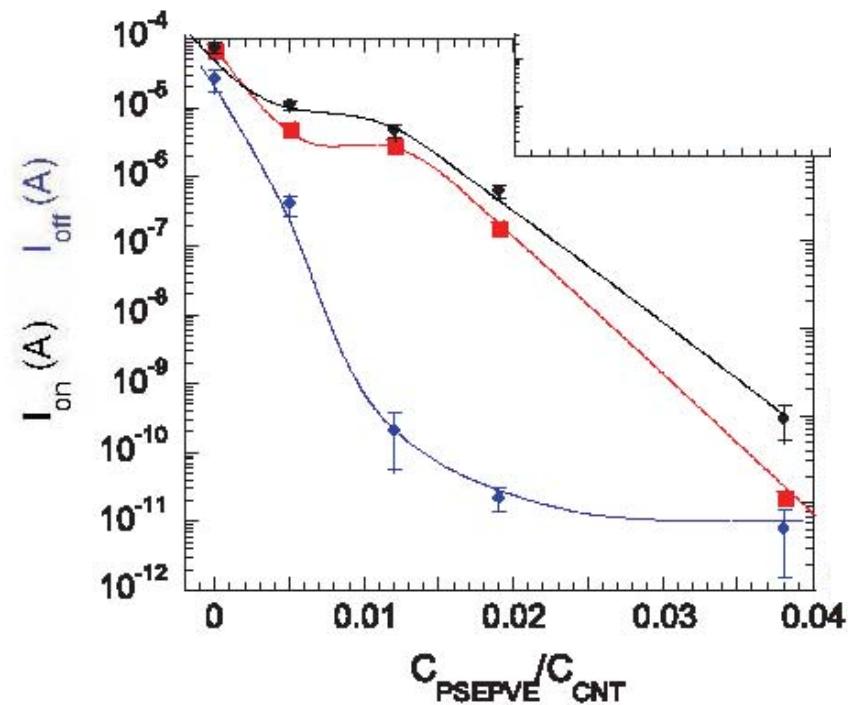
Electrical Transport Measurements

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Suppression of Metallic Conductivity of Single-Walled Carbon Nanotubes by Cycloaddition Reactions

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The high carrier mobility of films of semiconducting single-walled carbon nanotubes (SWNTs) is attractive for electronics applications, but the presence of metallic SWNTs leads to high off-current in transistor applications. The method presented here, cycloaddition of fluorinated olefins, represents an effective approach toward converting the “as grown” commercial SWNT mats into high-mobility semiconducting tubes with high yield and without further need for carbon nanotube separation. Thin-film transistors, fabricated from percolating arrays of functionalized carbon nanotubes, exhibit mobilities >100 square centimeters per volt-second and on-off ratios of 100,000. This method should allow for the use of semiconducting carbon nanotubes in commercial electronic devices and provide a low-cost route to the fabrication of electronic inks.



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- Transport: Young-Su Lee (MIT), Elise Li (MIT), Giovanni Cantele (University of Naples)



Conclusions

- Sparse Hamiltonians with chemical accuracy: Exact mapping of first-principles calculations onto tight binding Hamiltonians, using maximally-localized Wannier functions.
- Cycloadditions preserve the metallic conductance of carbon nanotubes even at high degrees of coverage
- Selected addends display fluxional tautomerism that
 - Can be directed with optical, chemical, or electrochemical means
 - Directly controls the conductance of the nanotubes