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Could a Massless SU(5) Theory Underly the Standard Model? Big X-Section New LHC Physics is Predicted!

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Could a Massless SU(5) Theory Underly the Standard

Model ? {Big X-Section New LHC Physics is Predicted!} *

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Using supercritical RFT, I have argued that the uniquely unitary* Critical \mathbb{P} occurs \{via\ anomaly\ dynamics\}\ in\ "QCD_S" = QCD_{n_f=6}\ + \{a\ sextet\ quark\ doublet\} \rightarrow EW sym-breaking \{sextet\ pions\}\ \& Dark Matter \{stable\ sextet\ neutrons\}  \begin{pmatrix} Anomaly\ color\ factors \\ => big\ \mathbb{P}\ x-sections \\ for\ EW\ bosons\ \&\ sextet \\ nucleons\ -\ at\ the\ LHC\ ? \end{pmatrix}
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Remarkably, QCD_S embeds uniquely in "QUD" *- a massless SU(5) theory that, even more remarkably, may have (via massless fermion IR anomalies) a bound-state S-Matrix that reproduces the full Standard Model.

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* The RFT Critical P (alone) satisfies all high-energy unitarity constraints.
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* Presented at Gribov-80, May 2010, Trieste, Italy
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^{*} Quantum Uno/Unification/Unique/Unitary/Underlying Dynamics

QUD \longleftrightarrow SU(5) gauge theory with massless, left-handed, fermions in the ${f 5} \oplus {f 15} \oplus {f 40} \oplus {f 45}^*$ representation.

Uniquely discovered as 1) anomaly free 2) asymptotically free $\{\textit{just}\}\$ & 3) contains the EW symmetry-breaking sextet sector. Under $SU(3)\otimes SU(2)\otimes U(1)$

$$5 = (3,1,-\frac{1}{3}))^{\left\{3\right\}} + (1,2,\frac{1}{2})^{\left\{2\right\}} , \qquad 15 = (1,3,1) + (3,2,\frac{1}{6})^{\left\{1\right\}} + (6,1,-\frac{2}{3}) ,$$

$$40 = (1,2,-\frac{3}{2})^{\left\{3\right\}} + (3,2,\frac{1}{6})^{\left\{2\right\}} + (3^*,1,-\frac{2}{3}) + (3^*,3,-\frac{2}{3}) + (6^*,2,\frac{1}{6}) + (8,1,1) ,$$

$$45^* = (1,2,-\frac{1}{2})^{\left\{1\right\}} + (3^*,1,\frac{1}{3}) + (3^*,3,\frac{1}{3}) + (3,1,-\frac{4}{3}) + (3,2,\frac{7}{6}))^{\left\{3\right\}} + (6,1,\frac{1}{3}) + (8,2,-\frac{1}{2})$$

Not only does QUD contain QCD_S , both the triplet quark & lepton sectors $\{\text{not asked for}\}$ are amazingly close to the SM !!! There are three "generations" $-\{1\},\{2\},\{3\}$.

Very importantly, QUD is real $\{vector-like\}$ wrt $SU(3)xU(1)_{em}$. SU(2)xU(1) is not quite right but the lepton anomaly is correct => to be physically realistic,

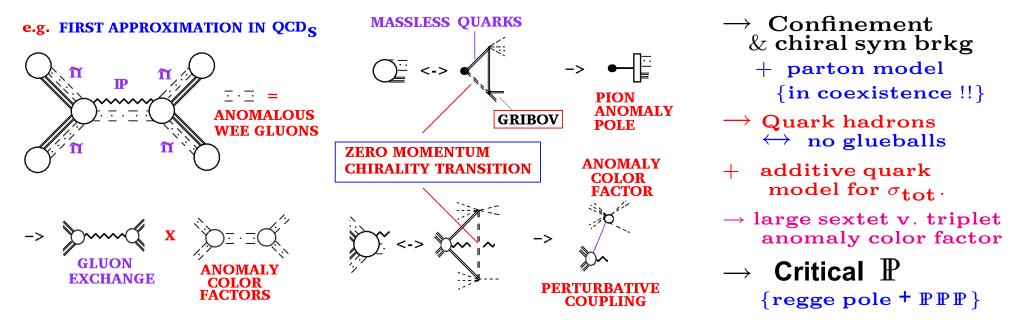
all elementary leptons & quarks must be confined & massless !!

As for QCD_S, access to the QUD bound-state S-Matrix is provided by multi-regge theory !!!

The multi-regge region involves multiple ∞ -momenta that $\{fundamentally\}$ allow wee partons to (simultaneously) play a vacuum role in both states & interactions.

I will build up the multi-regge QUD S-Matrix via IR divergences, by removing masses, & a k_{\perp} cut-off, in perturbative **reggeon diagrams**.

In massless* QCD $_S$ - Anomalies + IR fixed pt. \rightarrow IR divergence \rightarrow Wee gluons in all bound-states & interactions - coupled via zero-momentum quark chirality transitions,



^{*}Many massless mesons \rightarrow S-Matrix $\{???\}$. Masses require QUD.

S-Matrix IR chirality transitions play an even more fundamental role in QUD. (They are produced by the zero fermion mass limit of reggeon anomaly vertices.)

Although only an outline, for which {as will be obvious} much further development is needed and many details are missing, my construction will imply the following.

1. All elementary fermions are confined. Bound-states are formed by anomaly poles

the symmetry breaking $SU(5) \rightarrow SU(3) \otimes U(1)_{em}$ is due to zero-momentum chirality transitions.

- 2. Interactions are vector boson reggeons + anomalous wee gauge bosons.
- 3. Symmetry-breaking is an S-Matrix phenomenon \leftrightarrow no off-shell amplitudes (?) SU(5) is unbroken at large k_{\perp} but, although QUD lies in the "conformal window", the S-Matrix has only SM interactions & a spectrum of SM form.
- 4. All particles (including neutrinos) are bound-states with dynamical masses.

There is no Higgs !!

Because of it's uniqueness, QUD is either right or wrong - in it's entirety. {To be right} it must reproduce the full Standard Model S-Matrix!!

{Motivated by a unitary \mathbb{P} !!} QUD could provide a remarkably economic unification & even {perhaps} an origin* for the SM. Beyond the SM generations, there is only

- 1. A sextet quark sector that minimally, & naturally, solves two major mysteries Sextet pions $\to W^{\pm}, Z^0$ masses, & stable* sextet neutrons \to dark matter!
- 2. A "lepton-like" octet quark sector that is buried in all states as UV anomaly poles
 - → leptons & hadrons in SM generations
- 3. A pair of exotically charged quarks.

Nothing else!!

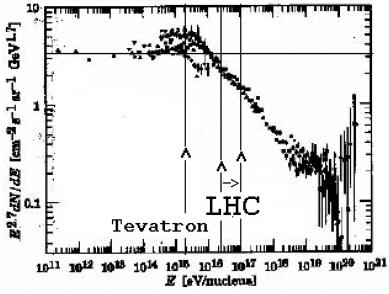
Although the physics is both novel & radical, the outcome is simple, consistent with established SM physics, & explains many puzzles. But, the multi-regge theory that I use to uncover it is so erudite that general interest will probably require (what would surely be)

A MAJOR EXPERIMENTAL DISCOVERY

^{*}Possibly, the SM (within QUD) is THE unique unitary particle S-Matrix!!

*Sextet protons are unstable - decaying to sextet neutrons.

COSMIC RAYS already suggest that new large x-section physics including



dark matter could appear at the LHC!!

The spectrum knee occurs between Tevatron and LHC energies. It is remarkably well-established, yet not understood. Although dark matter was unknown, a major threshold for neutral particles, unobserved in detectors, was initially suggested $\{\sim$ 40 years ago! $\}$. Underestimation of the energy would pile-up events as a "knee".

If the dark matter x-section is large at the LHC, a link to the knee is surely inevitable!!

For the sextet sector **three effects** should combine to produce a knee.

- 1. Prolific production of EW bosons increases $\langle p_{\perp} \rangle$ dramatically (& increases neutrino production) – leading to energy underestimation.
- 2. Threshold production of sextet neutron dark matter $\{\leftrightarrow \text{ inclusive } \mathbb{P}\}$.
- 3. Sextet neutrons as (UHE?) incoming cosmic rays with a (\mathbb{P}) threshold for atmospheric interaction not far below the knee.

X-sections must be **BIG** & at UHE must dominate σ_{tot} $\{ => \text{Dark Matter dominates } \}$

But, {most likely} \Longrightarrow very small mass neutrinos (no color/electric charge).

All particles are bound-states $\implies \alpha_{QUD}$ has no physical meaning.

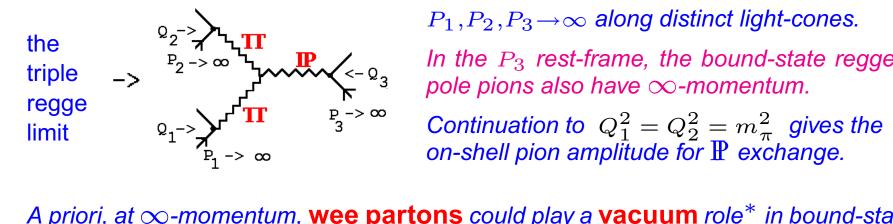
S-Matrix amplitudes are selected by an IR divergence ----

- Physical states & amplitudes all contain infinite sums of wee gauge bosons involving anomaly color factors {expressed, presumably, as integral formulae}.
 - \implies All interaction strengths are enhanced, with the SU(3) interaction strongly amplified by both color factors & the triple \mathbb{P} interaction $\{see\ later\}$.
- P anomaly color factors imply high-energy sextet x-sections are much larger than triplet x-sections => Dark Matter dominates early universe x-sections ??

My analysis of QUD anomaly dynamics 7 To give an outline, I begin with relies fundamentally on multi-regge theory. some basics - a regge pole gives

$$A^+(s,t) \underset{s \to \infty}{\sim} f(t) \xrightarrow{\frac{s^{\alpha(t)}}{sin \frac{\pi}{2}\alpha(t)}} \Longrightarrow t$$
- channel bound-state poles $\left(at \ \alpha(t) = 0, 2, \ldots\right)$ can be discovered

by calculating the $s \to \infty$ (∞ -momentum limit) in the cross-channel. In multi-regge limits multiple regge poles appear, e.g.



 $P_1, P_2, P_3 \rightarrow \infty$ along distinct light-cones.

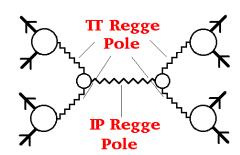
In the P_3 rest-frame, the bound-state regge pole pions also have as

A priori, at ∞ -momentum, wee partons could play a vacuum role* in bound-states if they are "universal". We will see that indeed, in QUD, multi-regge reggeon diagrams do produce "vacuum wee partons", but we must first introduce further ∞ -momenta !!

c.f. light-cone quantization using the perturbative vacuum.

In the "di-triple regge" (DTR) limit two triple-regge limits are separated by a further ∞ -momentum. Now, regge-pole $\pi's$ can scatter via the \mathbb{P} . All the $\pi's$ & the \mathbb{P} have ∞ -momentum in some frame

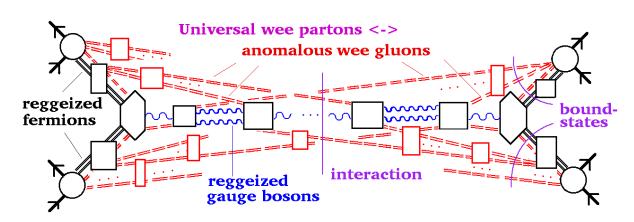
Hadron
DTR
Amplitude



=> both bound-states $(\pi's)$ & interactions (\mathbb{P}) could appear as parton (reggeon) states if universal wee reggeons carry vacuum properties $\{>> parton\ model\}$. We will see

that, in QUD, an initial divergence produces "anomalous wee gluons" universally !!

Multiple
wee parton
interactions
in a typical
initial QUD
DTR Amplitude



That SM states & interactions {& the Critical P} emerge from such diagrams as the complexity increases, is what has to be demonstrated!!

Reggeon Diagrams {a crash course} -

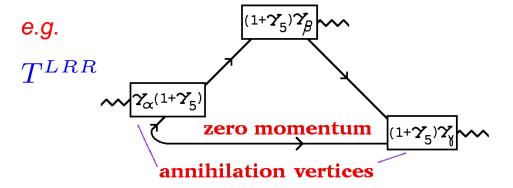
- 1. In (multi-)regge limits the large light-cone momenta are routed through feynman diagrams so that internal particles are maximally close to mass-shell & have large relative rapidities.
 - $ightarrow k_{\perp}$ diagram integral (ightarrow close to on-shell lines contracted) multiplied by rapidity logarithms.
- 2. Summing rapidity logarithms via reggeon propagators \rightarrow reggeon diagrams. { Infinitely many feynman diagrams \longleftrightarrow 1 reggeon diagram }
- 3. Internal particles with finite relative rapidity → couplings with more structure, including fermion loop interactions.
- 4. In a gauge theory **triangle anomalies** occur but {because a four-dimensional interaction is involved} only in **special vertices** coupling distinct reggeon channels **not*** in the single channel reggeon diagrams describing large k_{\perp} amplitudes.

^{*}SU(5) symmetry is not violated at "short distances" in QUD.

Reggeon anomaly vertices include axial-vector/vector/vector triangle diagrams T^{AVV} that, in QUD, must be defined at zero mass. At first sight,

- chirality is conserved at zero mass $\implies T^{AAA} = T^{AVV} = T^{RRR} + T^{LLL}$ \implies a conflict between the axial-vector anomaly & vector current conservation. But, regularization of γ_5 amplitudes is a major issue !!
- Fortunately, {it can be shown that} vector current conservation + axial anomaly
 unique massless anomaly pole amplitudes

IR pseudoscalar anomaly poles appear similiarly in both non chirality-violating & chirality-violating amplitudes via the triangle singularity.



Two annihilation vertices ↔ zero momentum chirality violation ⇒ the pseudoscalar pole can be a chiral Goldstone boson.

UV chirality transitions similarly produce anomaly poles as part of a Pauli–Villars subtraction when $\lambda_{\perp} \to \infty$.

I use QUD reggeon diagrams to construct small k_{\perp} DTR amplitudes*.

- I start with massive reggeons, via scalar VeV's, & with a k_{\perp} cut-off λ_{\perp} {VeV fermion masses, necessarily, identify particle/antiparticle pairs.}
- Because IR divergences produce wee partons in the massless theory, how the masses are removed is crucial. Anomalies play a key role {via the Gribov ambiguity}.

I take limits as follows -

- 1. Fermion masses $\{24 \& 5 \oplus 5^* \ VeV's\}$ are removed first, leaving chirality transitions in anomaly vertices that conserve SU(3) \otimes U(1) $_{em}$ but not SU(5).
- **2.** For gauge bosons, $\mathbf{5} \oplus \mathbf{5}^*$ VeV's \rightarrow smooth massless limit {complimentarity} Exponentiation of reggeization divergences confines* global SU(5), leaving IR finite interaction kernels. I take the limit in stages

$$ightarrow SU(2)_C$$
, $ightarrow SU(4)$, $\lambda_{\perp}
ightarrow \infty$, $ightarrow SU(5)$ (SU(2)_C $ightarrow SU(3)_C$)

asymptotic Supercritical Property of Critical Property (SU(2)_C $ightarrow SU(3)_C$)

 $^{^*}$ Very different from the large k_{\perp} use of QCD diagrams in BFKL physics.

^{*}Not true confinement! Multi-gluon singularities remain.

Because of λ_{\perp} , fermion loops do not have Ward identity zeroes when subsets of reggeon momenta vanish => many new divergences that exponentiate amplitudes to zero {including most left-handed boson couplings}.

The first color restoration \leftrightarrow (vector) SU(2) $_C \longrightarrow$ a wee gluon divergence that couples only via anomaly vertices & so does not exponentiate \longleftrightarrow

"anomalous wee gluons" $\left\{\begin{array}{l} \text{Sets of massless reggeized gluons, with all k_{\perp}'s scaled to zero wee gluons"} \\ \text{- with $I=0$ & "anomalous" color parity $\mathbf{C} \neq \tau = \text{signature.} \end{array}\right\}$

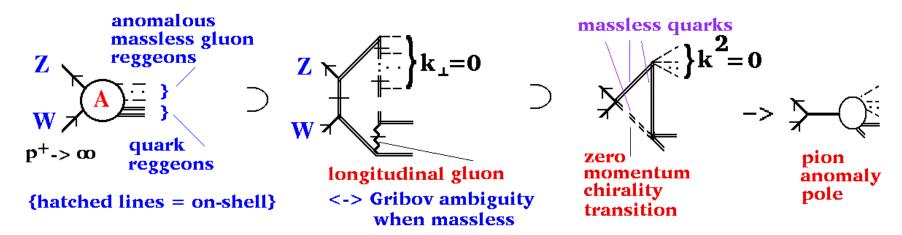
For SU(2), only $\tau = -C = -1$ is possible $\leftrightarrow 3, 5, ... \infty$ gluon reggeons.

IR fixed-point scaling \implies iteration of I=0 reggeon kernels reproduces the basic divergence with a

$$\left(\sum_{\substack{<\in \\ }} \mathbf{T}_N \stackrel{::}{\longrightarrow} \mathbf{K}_N^0 \stackrel{::}{\longrightarrow} \mathbf{T}_N \stackrel{::}{\longrightarrow} \left(\sum_{\stackrel{::}{\longrightarrow}} \mathbf{T}_N \stackrel{::}{\longrightarrow} \left(\sum_{\stackrel{::}{\longrightarrow}} \mathbf{K}_N^0 \stackrel{::}{\longrightarrow} \mathbf{T}_N \stackrel{::}{\longrightarrow} \mathbf{K}_N^0 \stackrel{::}{\longrightarrow} \mathbf{T}_N \stackrel{::}{\longrightarrow} \right) \right) \quad \text{factorized residue.} \\ \text{Factorizing off the divergence}$$

— universal wee gluon component of reggeon states & interactions.

The divergence also produces bound-state anomaly poles in vertices, e.g.



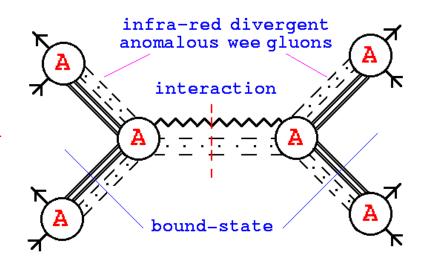
It is vital for S-Matrix symmetry breaking that an anomaly pole connects

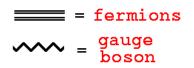
- 1. a physical fermion pair + wee gluon "vacuum component" ← Reggeon state with color zero projection → finite perturbative interactions, & {via the interaction}
- 2. an opposite chirality fermion pair one unphysical & with zero momentum. ↔ pseudoscalar Goldstone boson* associated with symmetry breaking.

Effectively, an anomaly pole bound—state is created by a zero momentum shift of the Dirac sea. By absorbing anomalous wee gluons, a physical fermion makes a symmetry—breaking chirality transition to an unphysical "hole state".

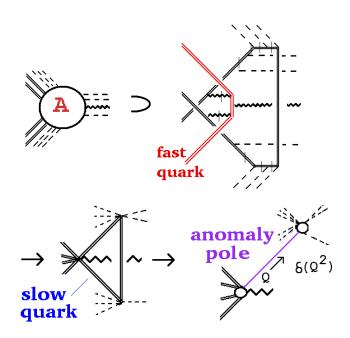
* An ∞ -momentum anomaly pole has physical Goldstone boson couplings.

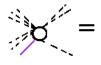
The simplest
DTR amplitudes
selected by the
anomalous wee
gluon divergence.





A = chirality transition anomaly pole vertex





WEE GLUON ANOMALY COLOR

FACTOR <-> sum over all wee gluons coupling to the slow quark loop

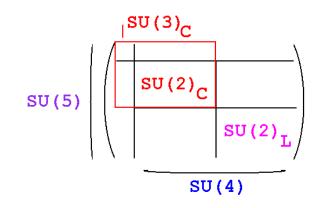
anomaly ∕ pole factorization ⇒ in QUD the P will couple more strongly to sextet states than to triplet states



Restoration of SU(2) $_C$ - gives anomaly pole chiral Goldstones (π_C 's) due to $5 \oplus 5^*$ chirality transitions, that are $q\bar{q}$ ("mesons") or qq & $\bar{q}\bar{q}$ ("nucleons").

The q's are **3's**, **6's**, & **8's** under $SU(3)_C$. **8's** are real wrt $SU(3)_C$, but contain complex chiral doublets wrt $SU(2)_C$. Via chirality transitions, the π_C 's are also reggeon states

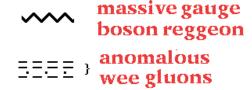
• Other reggeon states containing a π_C are also selected { will give leptons & SU(5) symmetry }



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\left\{\begin{array}{c} \begin{array}{c} \begin{array}{c} \\ \end{array} \end{array}\right\} \begin{array}{c} q_L \bar{q}_L \\ \end{array} \begin{array}{c} \text{pair} \\ \text{anomalous} \\ \text{wee gluons} \end{array}\right\} \iff \text{fr}_C
\begin{array}{c} \\ \\ \end{array} \begin{array}{c} \text{massive} \\ \text{gauge boson} \\ \text{reggeons} \\ \end{array} \begin{array}{c} \text{SU(2)}_C \\ \text{color} \\ \text{zero} \end{array}
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To avoid fermion loop exponentiation of the anomaly divergence the massive gauge boson reggeons must be vectors \leftrightarrow $SU(2)_C$ singlet gluons or photons.

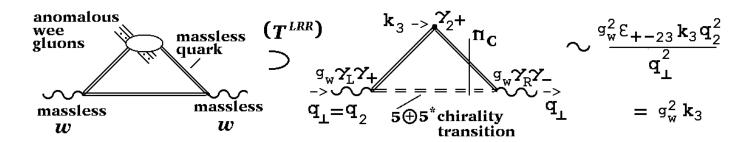
Interactions are even signature & are $SU(2)_C$ singlet massive vector exchange, together with anomalous wee gluons



- ullet Massive SU(3) $_C$ gluon exchange in the SU(2) $_C$ wee gluon background $\ \leftrightarrow \ \mathbb{P}$.
- ullet The massive gluon can be replaced by a massive γ , W^\pm or Z^0 .

Elementary left-handed exchanges ($W^{\pm} \& Z^{0}$) exponentiate to zero, **but**

5⊕5* chirality transitions
→ crucial wee-gluon interactions.

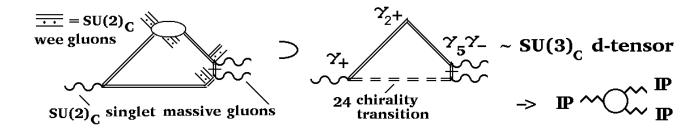


- ightarrow mass = $M_W^2 \sim g_W^2 \int dk k$ \leftrightarrow universal wee gluon integral.
- \leftrightarrow mixing with the $\pi_C{}'s$ $\{\pi_6 \text{ dominates after SU(3)}_{\mathbf{C}} \text{ is restored } \}$
- \rightarrow new quantum number $\{\rightarrow$ sextet flavor $\}$ & no exponentiation to zero.
- \rightarrow massive vector $W^{\pm} \& Z^{0}$

Odd-signature interactions are $\gamma \mathbb{P}$, $W^{\pm} \mathbb{P} \& Z^0 \mathbb{P}$ $\left\{ \begin{array}{ll} \rightarrow \text{ physical } \gamma, \mathbf{W}^{\pm}, \mathbf{Z}^0 \\ \text{ after } \mathbf{SU}(\mathbf{3})_{\mathbf{C}} \text{ restoration} \right\}. \end{array}$

24 chirality transitions

→ d-tensor triple P vertex.



- Wee gluons \rightarrow orthogonal γ -matrices $\leftrightarrow \gamma_5$.
 - \longrightarrow essential vertex for the SU(3) $_C$ interaction & the Critical $\mathbb P$.

SU(4) Restoration \longleftrightarrow SU(2) $_C$ singlet vector γ_4 becomes massless

- Other gauge bosons (that become massless) have left-handed couplings & appear only in interaction kernels {reggeon diagrams are exponentiated to zero.}
- γ_4 pairs have 1) even signature, 2) an SU(4) singlet projection, & 3) exponentiate to zero all amplitudes except when coupling to anomaly poles. $\{\leftrightarrow 1-e^{-\infty}\}$
- Anomalous wee bosons $\{\gamma_4 \text{ pairs} + SU(2)_C \text{ anomalous wee gluons}\}$ \rightarrow even signature γ - with SU(4) singlet projection.
- Anomalous wee bosons + massive gluon \rightarrow Supercritical \mathbb{P} .

Pseudoscalar anomaly poles coupling to γ_4 pairs are produced by

- 1. lepton pairs $(1, 2, \frac{1}{2})$ & $(1, 2, -\frac{1}{2})$.

 Chiral symmetry + **24** chirality transitions \rightarrow **pseudoscalars** $\pi_L^{\pm,0}$
- 2. $SU(2)_C$ singlets (8, 1, 1) & (8, 2, $-\frac{1}{2}$). Chiral symmetry + $\mathbf{5} \oplus \mathbf{5}^*$ chirality transitions \rightarrow pseudoscalars $\eta_8^{\pm,0}$.

Now, the anomalous wee boson divergence \implies physical bound-states

- 1. contain two pseudoscalar anomaly poles coupling to the divergence
- 2. have SU(4) singlet projections as reggeon states.
- 3. Fermion bound-states contain an additional elementary fermion.

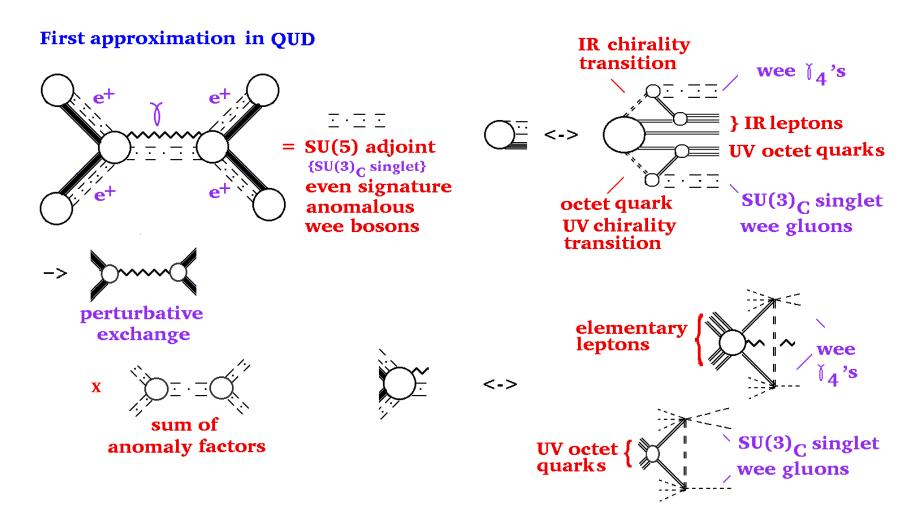
Leptons - π_L + π_8 + additional lepton \to 3 lepton generations. **Hadrons -** $\pi_{3,6}$ + η_8 \to mesons , + additional quark \to baryons. $\{$ More details after SU(5) restoration. $\}$

SU(5) Restoration ← SU(4) singlet vector becomes massless.

- $\begin{array}{ll} \textbf{1. Supercritical} \ \mathbb{P} & \{ \leftrightarrow SU(4) \ singlet + odd signature \ anomalous \ wee \\ & bosons \ with \ SU(3)_{\mathbf{C}} \ color \} \longrightarrow & \textbf{Critical} \ \mathbb{P} & \{ \ \leftrightarrow \ randomizing \\ & symmetry \ breaking \ ? \} \\ \end{array}$
- 2. Even signature anomalous wee bosons $\{with SU(3)_{\mathbb{C}} color zero\}$
 - \longrightarrow odd-signature massless γ & massive $W^{\pm}, \, Z^0.$
- 3. π_8 & η_8 combine in real octet SU(3) $_C$ representation (Π_8) \Longrightarrow
 - octet IR anomaly poles cancel in all amplitudes,
 - UV Π_8 poles appear as $\lambda_\perp \to \infty$ $\{\leftrightarrow \mathbf{Pauli-Villars\ subtractions}\}$
 - \implies Leptons & hadrons have IR components combined with UV $\,\Pi_{8}$'s.
- 4. Interactions are determined by IR components with states expressed as physical reggeons with an SU(5) singlet projection.
- 5. UV Π_8 anomaly poles \Longrightarrow IR fermions must be

$$(2, -\frac{1}{2})_L$$
 or $(2, \frac{1}{2})_R$ or $(1, 1)_L$ or $(1, -1)_R$ $\{SU(2) \otimes U(1)\}$

⇒ IR components of leptons & hadrons form SM generations.



Octet quarks coupled via $SU(3)_C$ color zero anomalous wee gluons {contained in SU(5) adjoint even—signature anomalous wee bosons} provide a UV light-cone contribution in all states & vertices.

The $SU(3) \times SU(2)_L \times U(1)$ lepton content, expressed as multi-fermion states, is

$$\begin{array}{l} \bullet \ \left(e^{-}, \nu \right) \ \leftrightarrow (1, 2, -\frac{1}{2}) \times \pi_{L}^{0} \times \Pi_{8} \\ \\ \leftrightarrow \ (1, 2, -\frac{1}{2}) \times (1, 2, -\frac{1}{2}) (1, 2, \frac{1}{2}) \times (8, 1, 1) (8, 2, -\frac{1}{2}) \\ \\ \leftrightarrow \ \mathbf{SU(5)} \ \mathbf{singlet/adjoint} - \mathbf{45}^{*} \times \mathbf{45}^{*} \times \mathbf{5} \times \mathbf{40} \times \mathbf{45}^{*} \\ \end{array}$$

•
$$(\mu^{-}, \nu) \leftrightarrow (1, 2, \frac{1}{2}) \times \pi_{L}^{-} \times \Pi_{8}$$

 $\leftrightarrow (1, 2, \frac{1}{2}) \times (1, 2, -\frac{1}{2}) (1, 2, -\frac{1}{2}) \times (8, 1, 1) (8, 2, -\frac{1}{2})$
 $\leftrightarrow SU(5) \text{ singlet/adjoint} - 5 \times 45^{*} \times 45^{*} \times 40 \times 45^{*}$

•
$$(\tau^-, \nu) \leftrightarrow (1, 2, -\frac{3}{2}) \times \pi_L^+ \times \Pi_8$$

 $\leftrightarrow (1, 2, -\frac{3}{2}) \times (1, 2, \frac{1}{2}) (1, 2, \frac{1}{2}) \times (8, 1, 1) (8, 2, -\frac{1}{2})$
 $\leftrightarrow \mathbf{SU}(\mathbf{5}) \mathbf{singlet/adjoint} - \mathbf{40} \times \mathbf{5} \times \mathbf{5} \times \mathbf{40} \times \mathbf{45}^*$

 e^+, μ^+, τ^+ can be obtained via charge conjugation.

QCD within QUD - the states are

- 1. triplet mesons & baryons
- 2. sextet "pions" & "nucleons" ($P_6 \& N_6$)
- **3. no** hybrid sextet/triplet quark states
- 4. no glueballs.

Consistent with, but fewer states than just confinement & chiral symmetry breaking.

- Sextet pions $\to W^{\pm} \& Z^{0} \implies$ sextet nucleons are the only new states.
- Sextet quarks have zero current mass $\implies N_6$ is stable \rightarrow DARK MATTER $\{electric\ charge \implies P_6\ is\ heavier\ -\ in\ contrast\ to\ the\ triplet\ sector\}$
- ullet Critical ${\mathbb P}$ & parton model, \longleftrightarrow no BFKL pomeron & no odderon.

Compared to conventional QCD, the states are fewer & the interaction is simpler - in agreement with experiment !!

 N_6 's have a very strong, very short range, QCD self-interaction $\{\rightarrow \text{clumps?}\}\$ & a QCD interaction with normal matter only at UHE. N_6 production will dominate UHE x-sections & early universe stable matter formation (& explain the CR knee!)

The bound-state mass spectrum will involve a combination of effects.

- 1. Perturbative reggeization is a very small effect, since α_{QUD} is so small.
- 2. Wee gluon anomaly interactions {which need to be catalogued} will mix the reggeon states & introduce anomaly color factors.
- 3. The SU(3) strong interaction will be emphasized by both **color factors & by the resulting high mass sector.** Electroweak charges will also play a role.
- 4. There is no symmetry that would conflict with the SM mass spectrum.

A wide range of scales should emerge & in bound-states all fermions will, effectively, have constituent masses. Connecting the η_6 to top production suggests

$$m_{q_6} \sim m_{top} \implies m_{N_6} \sim 500 \, GeV$$

In general, a better understanding of anomaly interactions & related wee gluon distributions is needed to determine if, & how many, parameters are involved.

Presumably, CP violation can be introduced via the anomalies, but is it essential?

QUD Virtues { beyond QCD}

It would be hard to over-emphasize the scientific $\{\&\ aesthetic\}$ importance of an underlying massless field theory for the Standard Model. In addition,

- 1. QUD is self-contained with only SM Interactions !!
- 2. The only new physics is the strong interaction sextet sector giving EW symmetry breaking, dark matter, & unification without supersymmetry !!
- 3. The parity properties of the strong & weak interaction are naturally explained.
- 4. Particles & fields are truly distinct. Hadrons & leptons have equal status. No off-shell amplitudes & no Higgs field ←→ all symmetries & masses are S-Matrix properties.
- 5. The QUD S-Matrix could be the only "non-perturbative" part of field theory needed with infinite momentum physics retaining a "perturbative" diagrammatic description.
- 6. Anomaly interactions mix the reggeon states &, presumably, introduce parameters.
- Anomaly color factors should produce a wide range of scales & masses that could produce the SM spectrum - there is no conflicting symmetry.
- 8. Small neutrino masses should be due to the underlying small coupling.
- 9. There is no proton decay, but the SU(5) Weinberg angle should hold!
- 10. Perturbatively, QUD is an asymptotically free, massless, fixed-point theory that should have no perturbative renormalons \implies it has no vacuum energy & would induce Einstein gravity with zero cosmological constant.