



**The Abdus Salam  
International Centre for Theoretical Physics**



**2146-1**

**Gribov-80 Memorial Workshop on Quantum Chromodynamics and  
Beyond'**

*26 - 28 May 2010*

**Gribov Reggeon calculus in QCD, SUSY and gravity**

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# Gribov Reggeon calculus in QCD, SUSY and gravity

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# 1 Regge ansatz for amplitudes

Regge kinematics

$$s = 4E^2 \gg -t = \vec{q}^2$$

$t$ -channel partial wave expansion

$$A^p(s, t) = s \int_{a-i\infty}^{a+i\infty} \frac{d\omega}{2\pi i} ((-s)^\omega - ps^\omega) f_\omega^p(t)$$

Regge pole hypothesis

$$f_\omega^p(t) = \frac{\gamma^2(t)}{\omega - \omega_p(t)}$$

Linear Regge trajectories

$$\omega_p(t) = \Delta - \alpha' \vec{q}^2$$

## 2 Mandelstam cuts

### Regge asymptotics

$$A_{Regge}^p(s, t) = \xi_p(t) s^{1+\omega_p(t)} \gamma^2(t), \quad \xi_p(t) = e^{-i\pi\omega_p(t)} - p$$

### Mandelstam cut contribution

$$A_{Mand}^p(s, t) = \xi_p s \int \frac{d^2k}{(2\pi)^2} \Phi^2(k, q-k) s^{\omega_{p_1}(-k^2)} s^{\omega_{p_2}(-(q-k)^2)}$$

### Gribov signature conservation rule

$$p = p_1 p_2$$

### Total cross sections and Pomeron

$$\sigma = \frac{1}{s} \Im A(s, 0), \quad A(s, t) \approx i s s^{\Delta_P - \alpha'_P q^2}, \quad \Delta_P \ll 1$$

### 3 Gribov Pomeron calculus

Multi-particle unitarity conditions in  $t$ -channel

$$\Im_t f_\omega(t) \sim \sum_n \int d\Omega_n |f_\omega^{(n)}|^2$$

Separation of particle clusters in their rapidities

$$0 < y_1 < y_2 < \dots < y_k < \ln s, \quad 1 \ll y_k - y_{k-1} \ll \ln s$$

Non-relativistic reggeon propagators

$$G_0 = \frac{1}{E + \Delta - \frac{k^2}{2m}}, \quad E = -\omega, \quad \alpha' = \frac{1}{2m}$$

Gribov effective Pomeron action

$$S = \int dy d^2\rho \left( \phi^* (\partial_y - \Delta) \phi + \frac{1}{2m} |\partial_\mu \phi|^2 + \lambda \phi^* \phi^2 + \dots \right)$$

# 4 Gluon reggeization in QCD

## QCD Born amplitude

$$M_{AB}^{A'B'}(s, t)|_{Born} = g T_{A'A}^c \delta_{\lambda_{A'} \lambda_A} \frac{2s}{t} g T_{B'B}^c \delta_{\lambda_{B'} \lambda_B}$$

## Leading logarithmic approximation

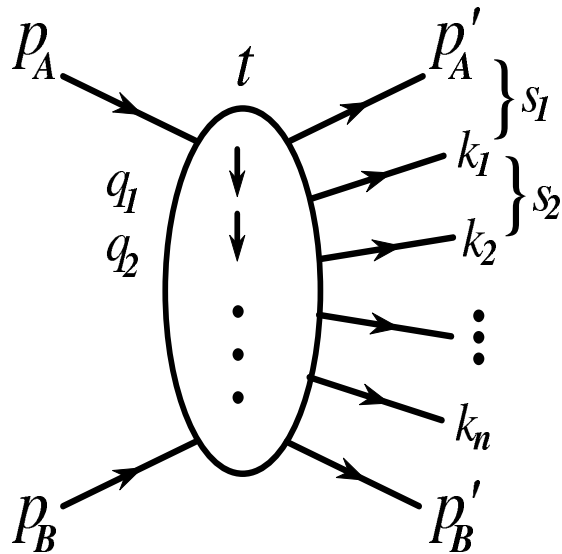
$$M_{AB}^{A'B'}(s, t) = M_{AB}^{A'B'}(s, t)|_{Born} s^{\omega(t)},$$

$$\alpha_s \ln s \sim 1, \quad \alpha_s = \frac{g^2}{4\pi} \ll 1$$

## Gluon Regge trajectory in LLA

$$\omega(-|q|^2) = -\frac{\alpha_s N_c}{4\pi^2} \int d^2k \frac{|q|^2}{|k|^2 |q-k|^2} \approx -\frac{\alpha_s N_c}{2\pi} \ln \frac{|q|^2}{\lambda^2}$$

## 5 Amplitudes in multi-Regge kinematics



$$M_{2 \rightarrow 2+n}^{BFKL} \sim \frac{s_1^{\omega_1}}{|q_1|^2} g T_{c_2 c_1}^{d_1} C(q_2, q_1) \frac{s_2^{\omega_2}}{|q_2|^2} \dots g T_{c_{n+1} c_n}^{d_n} C(q_{n+1}, q_n) \frac{s_{n+1}^{\omega_{n+1}}}{|q_{n+1}|^2},$$

$$\omega_r = -\frac{\alpha_s N_c}{2\pi} \left( \ln \frac{|q_r^2|}{\mu^2} - \frac{1}{\epsilon} \right), \quad C(q_2, q_1) = \frac{q_2 q_1^*}{q_2^* - q_1^*}, \quad \sigma_t = \sum_n \int d\Gamma_n |M_{2 \rightarrow 2+n}|^2$$

# 6 Analyticity, unitarity and bootstrap

Steinmann relations for overlapping channels

$$\Delta_{s_r} \Delta_{s_{r+1}} M_{2 \rightarrow 2+n} = 0$$

Dispersion representation for  $M_{2 \rightarrow 3}$  in the Regge ansatz

$$M_{2 \rightarrow 3} = c_1 (-s)^{j(t_2)} (-s_1)^{j(t_1) - j(t_2)} + c_2 (-s)^{j(t_1)} (-s_2)^{j(t_2) - j(t_1)}$$

Dispersion representation for  $M_{2 \rightarrow 4}$  in the Regge ansatz

$$\begin{aligned} M_{2 \rightarrow 4} = & d_1 (-s)^{j_3} (-s_{012})^{j_2 - j_3} (-s_1)^{j_1 - j_2} + d_2 (-s)^{j_1} (-s_{123})^{j_2 - j_1} (-s_3)^{j_3 - j_2} \\ & + d_3 (-s)^{j_3} (-s_{012})^{j_1 - j_3} (-s_2)^{j_2 - j_1} + d_4 (-s)^{j_1} (-s_{123})^{j_3 - j_1} (-s_2)^{j_2 - j_3} \\ & + d_5 (-s)^{j_2} (-s_1)^{j_1 - j_2} (-s_3)^{j_3 - j_2}, \quad j_r = j(t_r) \end{aligned}$$

Bootstrap relation in LLA (BFKL (1975-1978))

$$\pi \omega(t_1) M_{2 \rightarrow 2+n} = \sum_r \mathfrak{S}_{s_{0r}} M_{2 \rightarrow 2+n} = \sum_t M_{2 \rightarrow 2+t} M_{2+t \rightarrow 2+n}$$



# 7 BFKL equation (1975)

Balitsky-Fadin-Kuraev-Lipatov equation

$$E \Psi(\vec{\rho}_1, \vec{\rho}_2) = H_{12} \Psi(\vec{\rho}_1, \vec{\rho}_2), \quad \sigma_t \sim s^\Delta, \quad \Delta = -\frac{\alpha_s N_c}{2\pi} E_0$$

BFKL Hamiltonian

$$H_{12} = \frac{1}{p_1 p_2^*} (\ln |\rho_{12}|^2) p_1 p_2^* + \frac{1}{p_1^* p_2} (\ln |\rho_{12}|^2) p_1^* p_2 + \ln |p_1 p_2|^2 - 4\psi(1),$$

$$\rho_{12} = \rho_1 - \rho_2, \quad \rho_r = x_r + iy_r, \quad \Delta = 4\alpha N_c \ln 2 / \pi$$

Möbius invariance and eigenvalues (L. (1986))

$$\rho_k \rightarrow \frac{a\rho_k + b}{c\rho_k + d}, \quad m = \gamma + n/2, \quad \tilde{m} = \gamma - n/2, \quad \gamma = 1/2 + i\nu,$$

$$E = \psi(m) + \psi(1 - m) + \psi(\tilde{m}) + \psi(1 - \tilde{m}) - 4\psi(1)$$

# 8 Effective action approach

Locality in the rapidity space

$$y = \frac{1}{2} \ln \frac{\epsilon_k + |k|}{\epsilon_k - |k|}, \quad |y - y_0| < \eta, \quad \eta \ll \ln s$$

Gluon and Reggeized gluon fields

$$v_\mu(x) = -iT^a v_\mu^a(x), \quad A_\pm(x) = -iT^a A_\pm^a(x), \quad \delta A_\pm(x) = 0$$

Effective action for reggeized gluons (L., 1995)

$$S = \int d^4x (L_{QCD} + Tr(V_+ \partial_\mu^2 A_- + V_- \partial_\mu^2 A_+)) ,$$

$$V_+ = -\frac{1}{g} \partial_+ P \exp \left( -\frac{g}{2} \int_{-\infty}^{x^+} v_+(x') d(x')^+ \right) = v_+ - gv_+ \frac{1}{\partial_+} v_+ + \dots$$

## 9 Pomeron in $N = 4$ SUSY

BFKL kernel in two loops (F., L. and C.,C. (1998))

$$\omega = 4\hat{a} \chi(n, \gamma) + 4\hat{a}^2 \Delta(n, \gamma), \quad \hat{a} = g^2 N_c / (16\pi^2),$$

Hermitian separability in  $N = 4$  SUSY (K.,L. (2000))

$$\Delta(n, \gamma) = \phi(M) + \phi(M^*) - \frac{\rho(M) + \rho(M^*)}{2\hat{a}/\omega}, \quad M = \gamma + \frac{|n|}{2},$$

$$\rho(M) = \beta'(M) + \frac{1}{2}\zeta(2), \quad \beta'(z) = \frac{1}{4} \left[ \Psi'\left(\frac{z+1}{2}\right) - \Psi'\left(\frac{z}{2}\right) \right]$$

Maximal transcendentality (K.,L. (2002))

$$\phi(M) = 3\zeta(3) + \Psi''(M) - 2\Phi(M) + 2\beta'(M) \left( \Psi(1) - \Psi(M) \right),$$

$$\Phi(M) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k+M} \left( \Psi'(k+1) - \frac{\Psi(k+1) - \Psi(1)}{k+M} \right)$$

# 10 Pomeron and reggeized graviton

BFKL Pomeron in a diffusion approximation

$$j = 2 - \Delta - D\nu^2, \quad \gamma = 1 + \frac{j-2}{2} + i\nu$$

Constraint from the energy-momentum conservation

$$\gamma = (j-2) \left( \frac{1}{2} - \frac{1/\Delta}{1 + \sqrt{1 + (j-2)/\Delta}} \right)$$

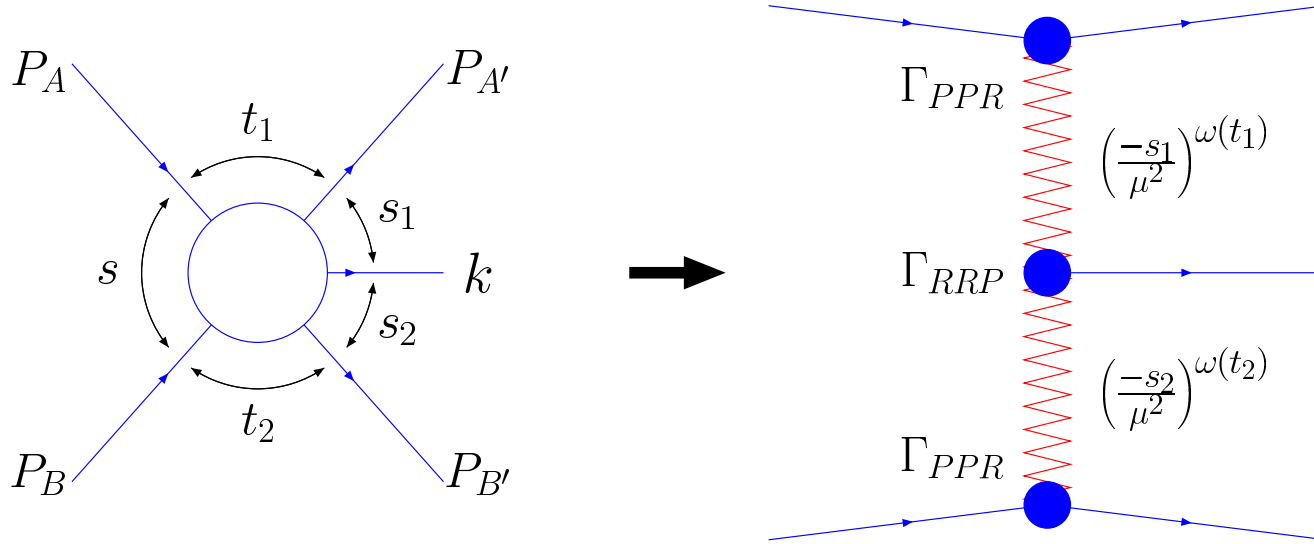
AdS/CFT relation for the graviton Regge trajectory

$$j = 2 + \frac{\alpha'}{2} t, \quad t = E^2/R^2, \quad \alpha' = \frac{R^2}{2} \Delta$$

Large coupling asymptotics for  $\Delta$  (KLOV, BPST)

$$\gamma = -\sqrt{2\pi(j-2)} \hat{a}^{1/4}, \quad j = 2 - \Delta, \quad \Delta = \frac{1}{2\pi} \hat{a}^{-1/2}$$

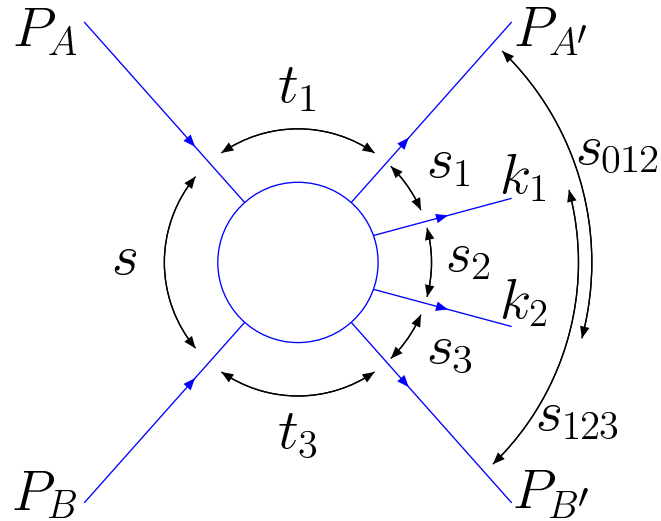
# 11 BDS production amplitude (BLS)



$$\ln \Gamma_{21} = -\frac{1}{2} \left( \omega(t_1) + \omega(t_2) - \int_0^a \frac{da'}{a'} \left( \frac{\gamma_K(a')}{4\epsilon} + \beta(a') \right) \right) \ln \frac{-k_{\perp}^2}{\mu^2} -$$

$$\frac{\gamma_K(a)}{16} \left( \ln^2 \frac{-k_{\perp}^2}{\mu^2} - \ln^2 \frac{-t_1}{-t_2} - \zeta_2 \right) - \frac{1}{2} \int_0^a \frac{da'}{a'} \ln \frac{a}{a'} \left( \frac{\gamma_K(a')}{4\epsilon^2} + \frac{\beta(a')}{\epsilon} + \delta(a') \right)$$

## 12 Regge factorization violation (BLS)



$$M_{2 \rightarrow 4}^{BDS} |_{s, s_2 > 0; s_1, s_3 < 0} = C \Gamma_1 \left( \frac{-s_1}{\mu^2} \right)^{\omega(t_1)} \Gamma_{21} \left( \frac{-s_2}{\mu^2} \right)^{\omega(t_2)} \Gamma_{32} \left( \frac{-s_3}{\mu^2} \right)^{\omega(t_3)} \Gamma_3,$$

$$C = \exp \left[ \frac{\gamma_K(a)}{4} i\pi \left( \ln \frac{t_1 t_2}{(\vec{k}_1 + \vec{k}_2)^2 \mu^2} - \frac{1}{\epsilon} \right) \right]$$

# 13 Mandelstam cuts in $j_2$ -plane

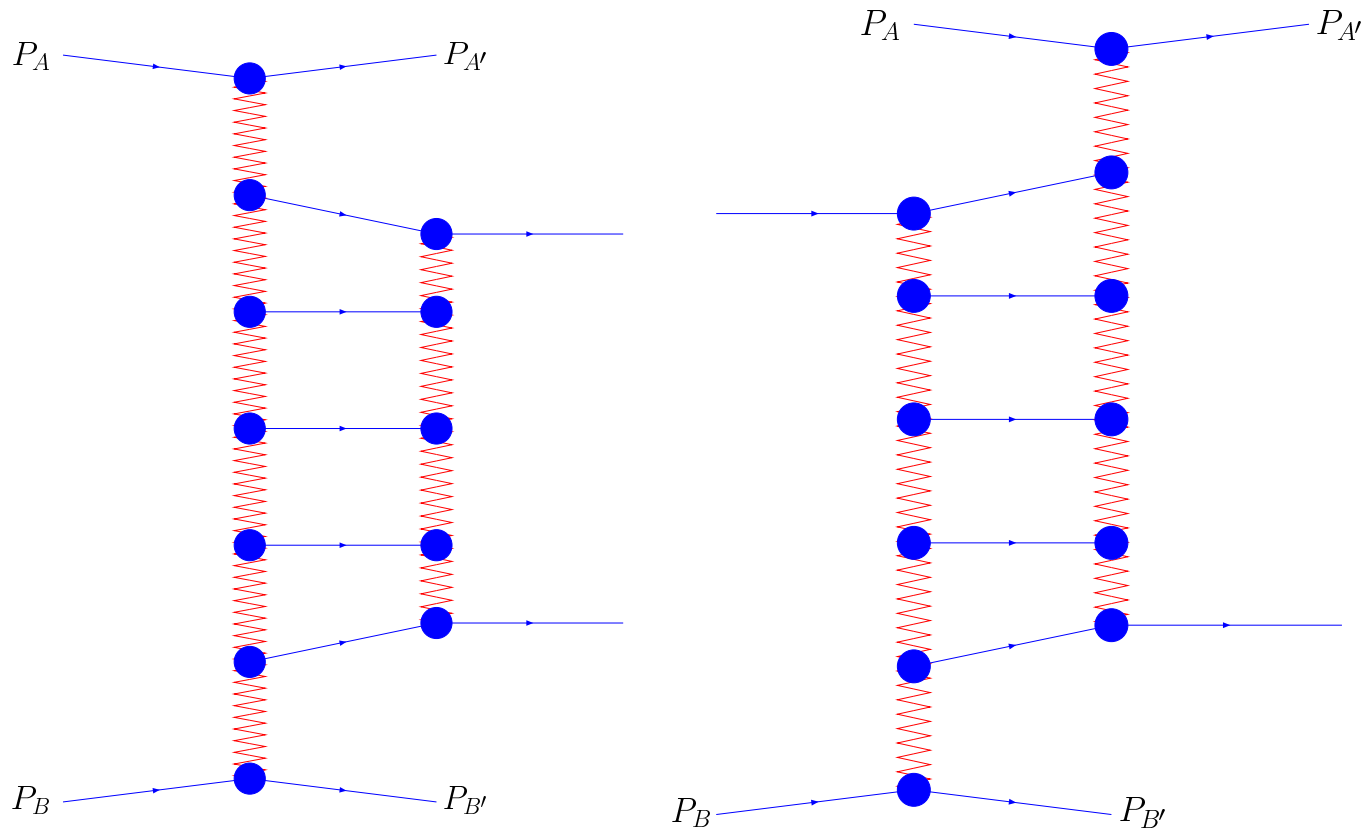


Figure 1: BFKL ladders in  $M_{2 \rightarrow 4}$  and  $M_{3 \rightarrow 3}$

# 14 Correction to the BDS result (BLS)

Factorization of infrared divergencies in LLA

$$M_{|s, s_2 > 0; s_1, s_3 < 0}^{2 \rightarrow 4} = (1 + i\Delta_{2 \rightarrow 4}) M_{2 \rightarrow 4}^{BDS},$$

Analytic result in LLA in the region  $a \ln s_2 \sim 1$

$$\Delta_{2 \rightarrow 4} = \frac{a}{2} \sum_{n=-\infty}^{\infty} (-1)^n \int_{-\infty}^{\infty} \frac{d\nu}{\nu^2 + \frac{n^2}{4}} (V^*)^{i\nu - \frac{n}{2}} V^{i\nu + \frac{n}{2}} \left( s_2^{\delta(\nu, n)} - 1 \right)$$

Duality transformation (L.L. (1999))

$$V = \frac{q_3 k_1}{k_2 q_1} \rightarrow \frac{z_{03} z_{0'1}}{z_{0'3} z_{01}}$$

Functions of 4-dimensional anharmonic ratios

$$i\Delta_{2 \rightarrow 4} = \frac{a^2}{4} Li_2(\chi) \ln \frac{\chi t_2 s_{13}}{s_3 t_1} \ln \frac{\chi t_2 s_{02}}{t_3 s_1} + \dots, \quad \chi = 1 - \frac{s s_2}{s_{012} s_{123}}$$



# 15 Multi-gluon states in octet channels

Holomorphic hamiltonian for n-gluon composite states

$$h = \ln(Z_1^2 \partial_1) - 2\psi(1) + \ln \partial_{n-1} + \sum_{k=1}^{n-2} h_{k,k+1}, \quad p_k = Z_{k-1,k}, \quad Z_0 = 0, \quad Z_n = \infty$$

Pair hamiltonian of the spin chain

$$h_{1,2} = \ln(Z_{12}^2 \partial_1) + \ln(Z_{12}^2 \partial_2) - 2 \ln Z_{12} - 2\psi(1)$$

Monodromy matrix (L. (1993))

$$t(u) = L_1(u)L_2(u)\dots L_{n-1}(u) = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix}$$

Integrals of motion and Baxter equation for integrable spin chain

$$[D(u), h] = 0, \quad D(u)Q(u) = (u - i)^{n-1}Q(u - i)$$

# 16 Effective action for gravity

Locality in the rapidity space

$$y = \frac{1}{2} \ln \frac{\epsilon_k + |k|}{\epsilon_k - |k|}, \quad |y - y_0| < \eta, \quad \eta \ll \ln s$$

Metric tensor and reggeized graviton fields

$$d^2 S = \sum_{\mu\nu} g_{\mu\nu} dx^\mu dx^\nu, \quad \delta A^{++}(x) = \delta A^{--}(x) = 0$$

Effective action for reggeized gravitons

$$S = \int d^4 x \left( -\frac{\sqrt{-g}}{2\kappa} R + G_{++} A^{++} + G_{--} A^{--} \right)$$

Effective metric tensor

$$G_{++} = g_{++} + (\partial_\sigma g_{++}) \frac{1}{4\partial_+^2} \partial_\sigma g_{++} - (\partial_\sigma g_{++}) \frac{1}{\partial_+} g_{+\sigma} - g_{+\sigma}^2 + \dots$$

# 17 Discussion

1. Locality in the rapidity space
2. Steinmann relations and bootstrap.
3. Remarkable properties of NLLA in  $N = 4$  SUSY.
4. Breakdown of the BDS ansatz.
5. Mandelstam cuts in planar amplitudes at  $n \geq 6$ .
6. Integrability of planar scattering amplitudes
7. Effective action for the reggeized gravitons