The Abdus Salam International Centre for Theoretical Physics

United Nations Educational, Scientific and Cultural Organization

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2146-35

Gribov-80 Memorial Workshop on Quantum Chromodynamics and Beyond'

26 - 28 May 2010

On Anomalous Quark Triangles

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Gribov-80 Memorial Workshop on Quantum Chromodynamics and beyond

May 26-28, 2010, ICTP, Trieste

On Anomalous Quark Triangles

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Anomalies, As A Manifestation Of The High Momentum Collective Motion In The Vacuum V.N. Gribov, 1981

Anomalies And A Possible Solution Of Problems Of Zero Charge And Infrared Instability V.N. Gribov, 1987

Luca Trentandue, Dmitri Kharzeev talks

Quark triangles show up in the muon anomalous magnetic moment

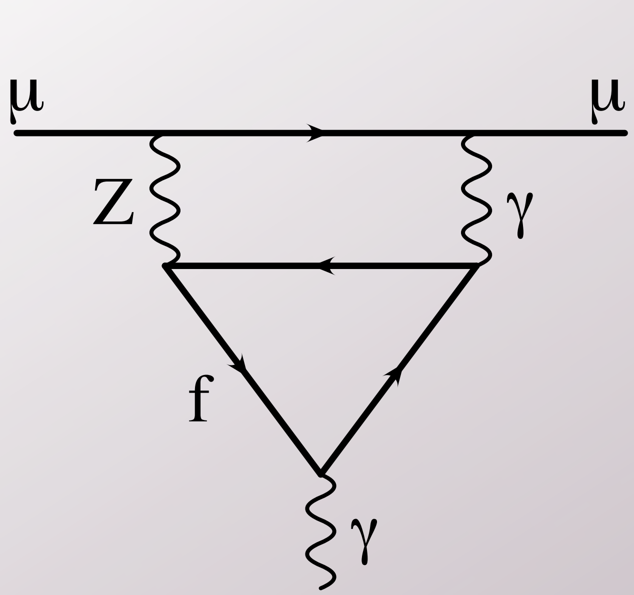
Refinements in electroweak contributions to the muon anomalous magnetic moment
Czarnecki, Marciano and AV, Phys. Rev. D 67, 073006 (2003)

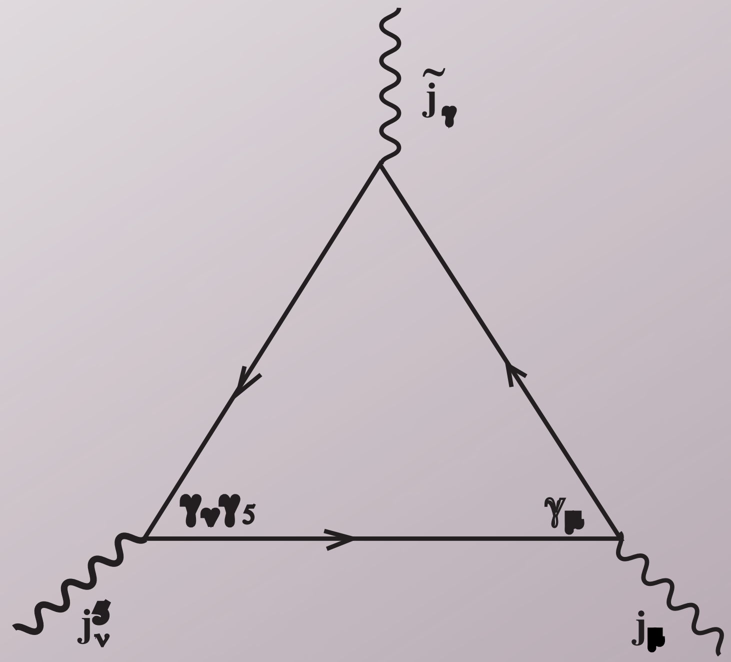
Perturbative and nonperturbative renormalization of anomalous quark triangles
AV, Phys. Lett. B 569, 187 (2003)

Triangle anomaly and the muon g-2 Czarnecki, Marciano and AV, Acta Phys. Polon. B 34, 5669 (2003)

Hadronic light-by-light scattering contribution to the muon anomalous magnetic moment revisited Melnikov and AV, Phys. Rev. D **70**, 113006 (2004)

Perturbative calculations





$$j_\mu = \bar{q}\, V \gamma_\mu q, \qquad j^5_\nu = \bar{q}\, A \gamma_\nu \gamma_5 q \qquad \tilde{j}_\gamma = \bar{q}\, \widetilde{V} \gamma_\mu q$$

$$T_{\mu\gamma\nu} = -\int \mathrm{d}^4x \mathrm{d}^4y \, \mathrm{e}^{iqx - iky} \, \langle 0 | \, T\{ j_\mu(x) \, \tilde{j}_\gamma(y) \, j^5_\nu(0) \} | 0 \rangle$$

Soft photon

$$T_{\mu\nu} = T_{\mu\gamma\nu} e^\gamma(k) = i \int \mathrm{d}^4x \, \mathrm{e}^{iqx} \, \langle 0 | \, T\{ j_\mu(x) \, j^5_\nu(0) \} | \, \gamma(k) \rangle$$

A general form of $T_{\mu\nu}$ contains two Lorentz structures

$$T_{\mu\nu} = -\frac{i}{4\pi^2}\left[w_T(q^2)\left(-q^2\tilde{f}_{\mu\nu} + q_\mu q^\sigma \tilde{f}_{\sigma\nu} - q_\nu q^\sigma \tilde{f}_{\sigma\mu}\right) + w_L(q^2)\, q_\nu q^\sigma \tilde{f}_{\sigma\mu}\right]$$

$$\tilde{f}_{\mu\nu} = \frac{1}{2}\,\epsilon_{\mu\nu\gamma\delta} f^{\gamma\delta}\,, \qquad f_{\mu\nu} = k_\mu e_\nu - k_\nu e_\mu$$

Rosenberg '63, Bell-Jackiw '69, Adler '69 calculated 1 loop

$$w_L^{1-\mathrm{loop}} = 2\,w_T^{1-\mathrm{loop}} = 2N_c\,\mathrm{Tr}\,A\,V\,\widetilde{V}\int_0^1 \frac{d\alpha\,\alpha(1-\alpha)}{\alpha(1-\alpha)Q^2 + m^2} \qquad Q^2 = -q^2$$

In the chiral limit $m = 0$

$$w_L^{1-\mathrm{loop}}[m=0] = 2\,w_T^{1-\mathrm{loop}}[m=0] = \frac{2N_c\,\mathrm{Tr}\,(A\,V\,\widetilde{V})}{Q^2}$$

Nonvanishing longitudinal part represents ABJ anomaly

$$q^\nu T_{\mu\nu} = \frac{i}{4\pi^2}\,Q^2 w_L\, q^\sigma \tilde{f}_{\sigma\mu} = \frac{i}{2\pi^2}\,N_c\,\mathrm{Tr}\,(A\,V\,\widetilde{V})\, q^\sigma \tilde{f}_{\sigma\mu}$$

Nonrenormalization, Adler-Bardeen theorem, implies that w_L stays intact when gluon interaction is switched on.

Nonrenormalization theorem for the transversal part

The relation

$$w_L[m=0] = 2\,w_T[m=0]$$

valid at the one-loop level gets no perturbative corrections. At $m=0$ the diagrams are symmetric under permutation $\mu \leftrightarrow \nu$, indices of the vector and axial currents.

For the symmetry to hold it is important that the part $q^2 \tilde{f}_{\mu\nu}$ produces just a constant in q term in $T_{\mu\nu}$. The singular part is symmetric and the constant term is fixed by the conservation of the vector current. (Independence on q is an alternative derivation of the AB theorem).

If Pauli-Vilars regularization is used to provide the vector current conservation then the antisymmetric part comes just from regulators.

Thus, the crossing symmetry relates the transversal and longitudinal parts and the AB theorem on the absence of perturbative corrections works for both.

For a general kinematics the relation was found in
'2004 Knecht, Peris, Perrottet and E. de Rafael and checked in
'2006 Jegerlehner and Tarasov .

What about nonperturbative corrections?
None in the longitudinal part ('t Hooft consistency condition) should present in the transversal part --
there is no massless spin one states.

Nonperturbative effects and OPE

$$\hat{T}_{\mu\nu} \equiv i \int \mathrm{d}^4x\, \mathrm{e}^{iqx}\, T\{j_\mu(x)\, j^5_\nu(0)\} = \sum_i c^i_{\mu\nu\gamma_1\ldots\gamma_i}(q)\, \mathcal{O}_i^{\gamma_1\ldots\gamma_i}$$

$$T_{\mu\nu} = \langle 0|\, \hat{T}_{\mu\nu}\, |\gamma(k)\rangle = \sum_i c^i_{\mu\nu\alpha_1\ldots\alpha_i}(q)\, \langle 0|\, \mathcal{O}_i^{\alpha_1\ldots\alpha_i}\, |\gamma(k)\rangle$$

$$\langle 0|\, \mathcal{O}_i^{\alpha\beta}\, |\gamma(k)\rangle = -\frac{i}{4\pi^2}\, \kappa_i\, \tilde{f}^{\alpha\beta}$$

$$\hat{T}_{\mu\nu} = \sum_i \left\{ c^i_T(q^2)\left(-q^2\mathcal{O}^i_{\mu\nu} + q_\mu q^\sigma \mathcal{O}^i_{\sigma\nu} - q_\nu q^\sigma \mathcal{O}^i_{\sigma\mu}\right) + c^i_L(q^2)\, q_\nu q^\sigma \mathcal{O}^i_{\sigma\mu} \right\}$$

$$w_{T,L}(q^2) = \sum c^i_{T,L}(q^2)\, \kappa_i$$

The leading d=2 operator

$$\mathcal{O}_F^{\alpha\beta} = \frac{1}{4\pi^2}\, \tilde{F}^{\alpha\beta} = \frac{1}{4\pi^2}\, \epsilon^{\alpha\beta\rho\delta}\partial_\rho A_\delta$$

$$c^F_L[\text{1-loop}] = 2c^F_T[\text{1-loop}] = \frac{2\, N_c}{Q^2}\, \mathrm{Tr}\, A\, V \widetilde{V}\left[1 + \mathcal{O}\left(\frac{m^2}{Q^2}\right)\right]$$

The next d=2 operator

$$\mathcal{O}_f^{\alpha\beta} = -i\, \bar{q}_f\, \sigma^{\alpha\beta}\gamma_5\, q^f \equiv \frac{1}{2}\, \epsilon^{\alpha\beta\gamma\delta}\bar{q}_f\, \sigma_{\gamma\delta}\, q^f$$

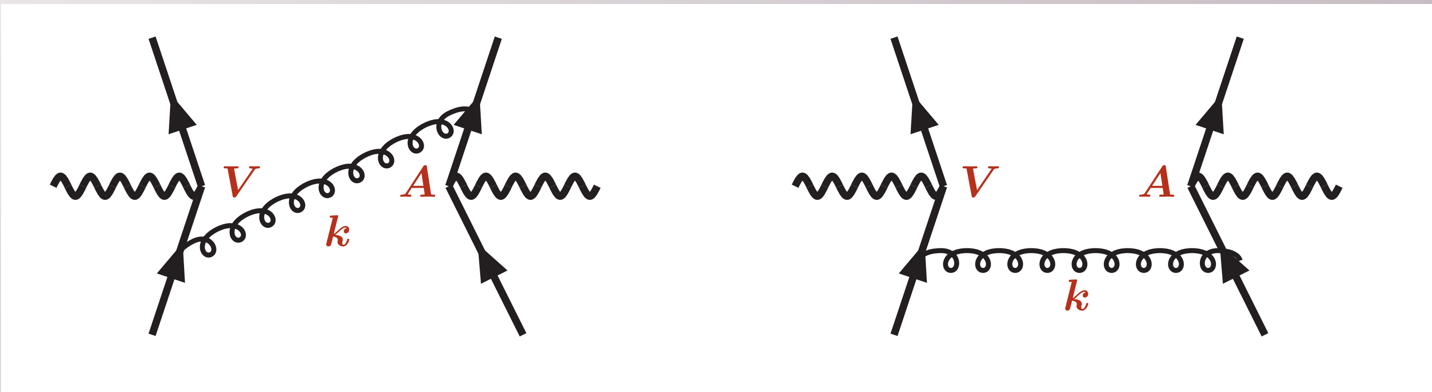
$$c_L^f = 2c_T^f = \frac{4\, A_f V_f\, m_f}{Q^4}$$

$$\Delta^{(d=3)} w_L = 2\, \Delta^{(d=3)} w_T = \frac{4}{Q^4} \sum_f A_f V_f m_f \kappa_f$$

$$\kappa_f = -4\pi^2 \widetilde{V_f}\, \langle \bar{q} q \rangle_0\, \chi$$

Quark condensate magnetic susceptibility χ Ioffe, Smilga

In the chiral limit the difference between longitudinal and transversal parts shows up at the level of d=6 four-fermion operators.



$$-\frac{8\pi\alpha_s Q_q}{k^6}\,\bar{q}\,t^a(\gamma_\alpha\hat{k}\gamma_\mu-\gamma_\mu\hat{k}\gamma_\alpha)q\otimes\bar{q}\,t^a(\gamma_\nu\hat{k}\gamma^\alpha-\gamma^\alpha\hat{k}\gamma_\nu)\gamma_5 q$$

Here the momentum q is substituted by k .

The model

$$w_{\mathrm{T}}[u,d]=\frac{1}{m_{a_1}^2-m_\rho^2}\left(\frac{m_{a_1}^2-m_\pi^2}{K^2+m_\rho^2}-\frac{m_\rho^2-m_\pi^2}{K^2+m_{a_1}^2}\right)$$

Applications

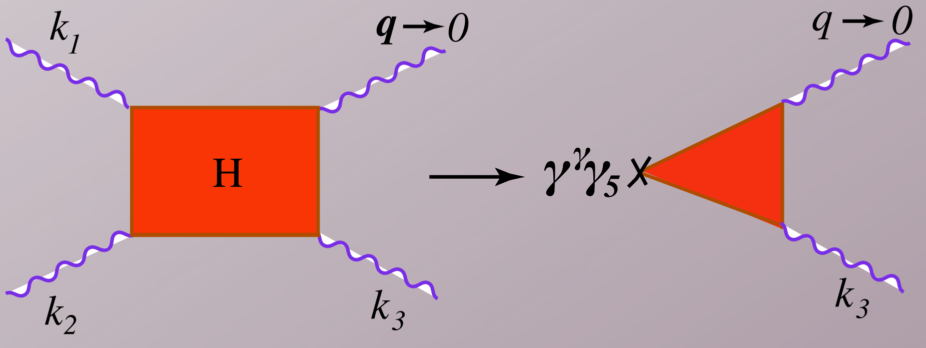
- Hadrons in the electroweak corrections to g-2

$$a_\mu^{\mathrm{EW}} = 154(1)(2) \times 10^{-11}$$

- Magnetic susceptibility of quark condensate

$$\chi = -\frac{N_c}{4\pi^2 F_\pi^2} = -\frac{1}{(335\ \mathrm{MeV})^2}$$

- Hadronic light-by-light in the muon g-2



$$a^{\mathrm{HLbL}} = (105 \pm 26) \times 10^{-11}$$

Thanks to

Julia, Luciano, Yuri, Peter, Daniele

and Nadia

for the perfect organization of the remarkable meeting