



*The Abdus Salam
International Centre for Theoretical Physics*



2146-3

**Gribov-80 Memorial Workshop on Quantum Chromodynamics and
Beyond'**

26 - 28 May 2010

CP Symmetry and Phase Transitions

Manuel Asorey
*Universidad de Zaragoza
Spain*

CP Symmetry and Phase Transitions

Manuel Asorey
Universidad de Zaragoza



Gribov–80
Trieste, May 2010

CP Symmetry

Matter-Antimatter Symmetry

- Electroweak interactions break CP symmetry
- Does QCD preserve CP symmetry?
- ϑ -Vacuum
- Strong CP problem
- Peccei-Quinn Theory: axions
- Non-perturbative breaking
- Vafa-Witten Theorem

QCD with ϑ term

- (Euclidean) partition function $\mathcal{Z}(\vartheta)$

$$\mathcal{Z}(\vartheta) = \int \mathcal{D}A e^{-S_{\text{ef}} - i\vartheta q} \quad S_{\text{ef}} \geq 0; \quad q \in \mathbb{Z}$$

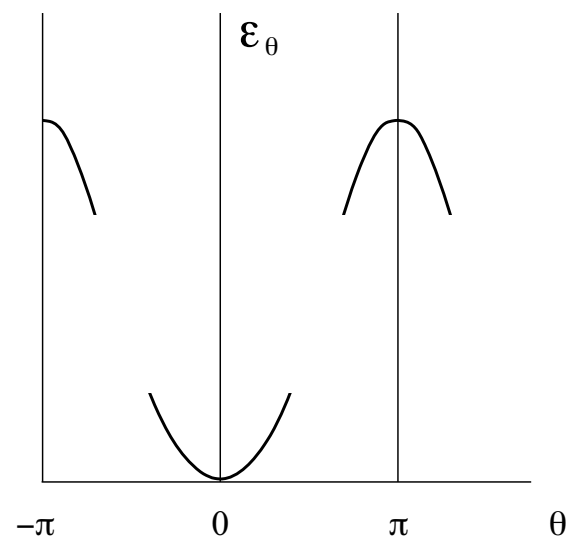
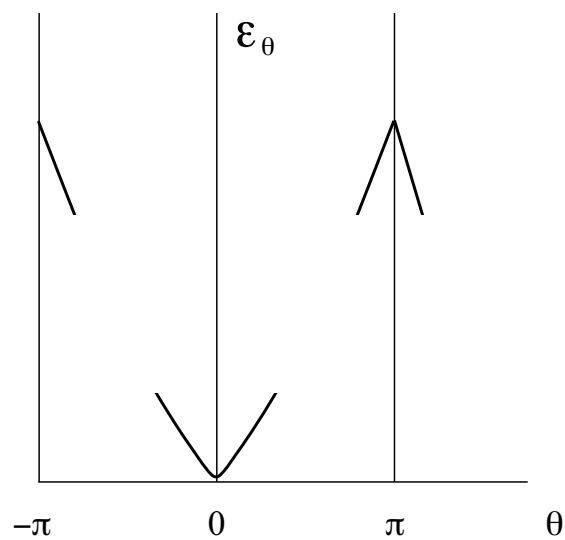
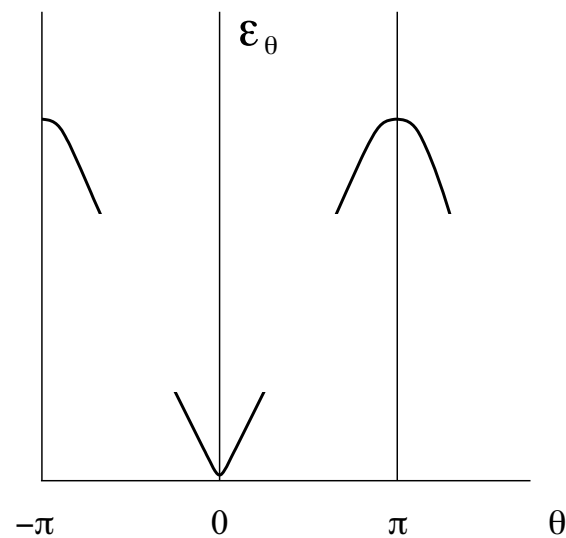
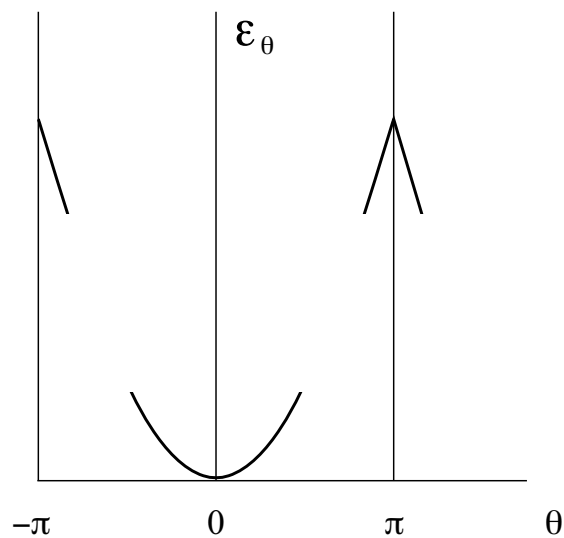
- $\mathcal{Z}(\vartheta) = \mathcal{Z}(\vartheta + 2\pi n)$ (periodicity)
- $\mathcal{Z}(\vartheta) = \mathcal{Z}(-\vartheta)$ (reflection symmetry)
- $\mathcal{Z}(\pi + \vartheta) = \mathcal{Z}(\pi - \vartheta)$ (Bragg symmetry)

- Energy density $\mathcal{E}_0(\vartheta)$

$$\mathcal{Z}(\vartheta) \sim e^{-VT\mathcal{E}_0(\vartheta)}$$

- $\mathcal{E}_0(\vartheta) = \mathcal{E}_0(\vartheta + 2\pi n)$ (periodicity)
- $\mathcal{E}_0(\vartheta) = \mathcal{E}_0(-\vartheta)$ (reflection symmetry)
- $\mathcal{E}_0(\pi + \vartheta) = \mathcal{E}_0(\pi - \vartheta)$ (Bragg symmetry)

CP symmetry and phase transitions



CP symmetry and phase transitions

1^{er} order transition:

$$\left. \frac{d\mathcal{E}(\vartheta)}{d\vartheta} \right|_{\vartheta=0,\pi} \neq 0$$

- Order parameter:

$$Q = \langle F\tilde{F}(x) \rangle \neq 0$$

Topological charge density

CP symmetry and phase transitions

2° order transition:

$$\left. \frac{d^2 \mathcal{E}(\vartheta)}{d\vartheta_-^2} \right|_{\vartheta=0,\pi} \neq \left. \frac{d^2 \mathcal{E}(\vartheta)}{d\vartheta_+^2} \right|_{\vartheta=0,\pi}$$

- Order parameter:

$$\chi_\tau = \int dx \left[\langle F\tilde{F}(0) F\tilde{F}(x) \rangle_0 - \langle F\tilde{F}(0) \rangle_0 \langle F\tilde{F}(x) \rangle_0 \right]$$

Topological susceptibility

χ_τ diverges if $Q \neq 0$

OS positivity and topological susceptibility

- Witten-Veneziano Formula

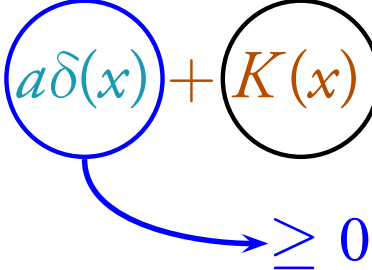
$$\chi_\tau \sim f_\pi^2 (m_{\eta'}^2 + m_\eta^2 - 2m_K^2) / 6$$

- Osterwalder-Schröder positivity

$$\langle F\tilde{F}(0) F\tilde{F}(x) \rangle \leq 0, \quad \text{for } x \neq 0$$

- Contact term

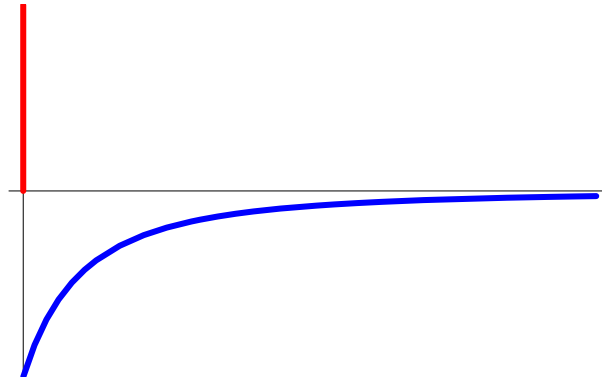
$$\langle F\tilde{F}(0) F\tilde{F}(x) \rangle \leq a\delta(x) + K(x) \rightarrow \leq 0$$



- Topological susceptibility positivity

$$a \geq \int dx K(x)$$

OS positivity and topological susceptibility



- Cluster property

$$\lim_{x \rightarrow \infty} \langle \tilde{F}\tilde{F}(0) \tilde{F}\tilde{F}(x) \rangle = \langle \tilde{F}\tilde{F} \rangle^2 \leq 0$$

- If CP symmetry is broken

$$\langle \tilde{F}\tilde{F} \rangle = iQ \neq 0$$

Vafa-Witten theorem

Parity symmetry is not spontaneously broken in QCD

- Basic VW argument: Order parameter: topological charge density

$$Q = \frac{1}{32\pi^2} F_{\mu\nu}^a \tilde{F}_{\mu\nu}^a$$

(Euclidean) partition function [with quarks]

$$\mathcal{Z}(\vartheta) = e^{-VT\mathcal{E}_0(\vartheta)} = \int \mathcal{D}A e^{-S_{\text{ef}} - i\vartheta q},$$

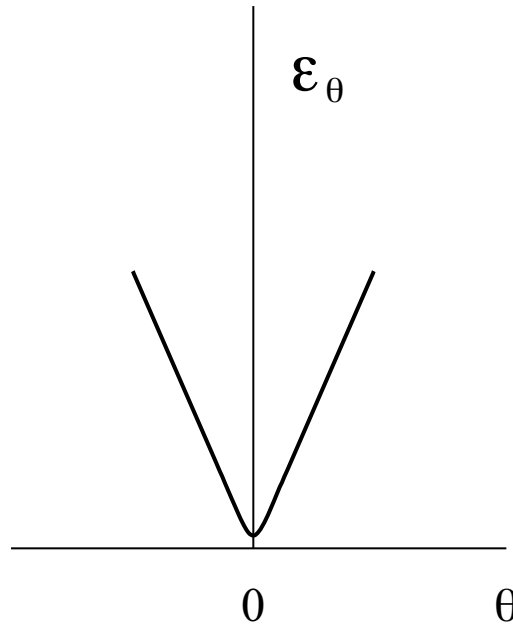
$$|\mathcal{Z}(\vartheta)| \leq \mathcal{Z}(0) \implies \mathcal{E}_0(\vartheta) \geq \mathcal{E}_0(0),$$

Assuming that $\mathcal{E}_0(\vartheta)$ is differentiable in $\vartheta = 0$

$$\mathcal{E}'_0(0) = \frac{i}{32\pi^2} \langle F_{\mu\nu}^a \tilde{F}_{\mu\nu}^a(0) \rangle_0 = \langle Q \rangle_{\vartheta=0} = 0.$$

Vafa-Witten theorem

- Possibility of cusp \Leftrightarrow 1st order phase transition



- The theorem can be proved with two additional assumptions
 - Renormalizability
 - Unitarity

Analiticity of the partition function

$$\mathcal{Z}_{\vartheta}(g) = \sum_{q=-\infty}^{\infty} e^{-q\vartheta_I + iq\vartheta_R} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(g,q)}$$

$$|\mathcal{Z}_{\vartheta}(g)| \leq \sum_{q=-\infty}^{\infty} e^{|q\vartheta_I|} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(g,q)} \leq \sum_{q=-\infty}^{\infty} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(g,q) + \frac{|\vartheta_I|}{8\pi^2} S_{\text{YM}}}$$

BPS bound

$$S_{\text{YM}}(q) = \frac{1}{2} \int F^{\mu\nu} F_{\mu\nu} \geq \frac{1}{2} \left| \int F^{\mu\nu} \tilde{F}_{\mu\nu} \right| = 8\pi^2 |q|$$

Analiticity of the partition function

$$\mathcal{Z}_{\vartheta}(g) = \sum_{q=-\infty}^{\infty} e^{-q\vartheta_I + iq\vartheta_R} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(g,q)}$$

$$|\mathcal{Z}_{\vartheta}(g)| \leq \sum_{q=-\infty}^{\infty} e^{|q\vartheta_I|} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(g,q)} \leq \sum_{q=-\infty}^{\infty} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(\tilde{g},q)} = \mathcal{Z}_0(\tilde{g})$$

$$\frac{1}{\tilde{g}^2} = \frac{1}{g^2} - \frac{|\vartheta_I|}{8\pi^2}$$

$\mathcal{Z}_{\vartheta}(g)$ is an analytic function in ϑ

Analiticity of the partition function

$$\mathcal{Z}_{\vartheta}(g) = \sum_{q=-\infty}^{\infty} e^{-q\vartheta_I + iq\vartheta_R} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(g,q)}$$

$$|\mathcal{Z}_{\vartheta}(g)| \leq \sum_{q=-\infty}^{\infty} e^{|q\vartheta_I|} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(g,q)} \leq \sum_{q=-\infty}^{\infty} \int_{c_2(A)=q} \mathcal{D}A e^{-S_{\text{ef}}(\tilde{g},q)} = \mathcal{Z}_0(\tilde{g})$$

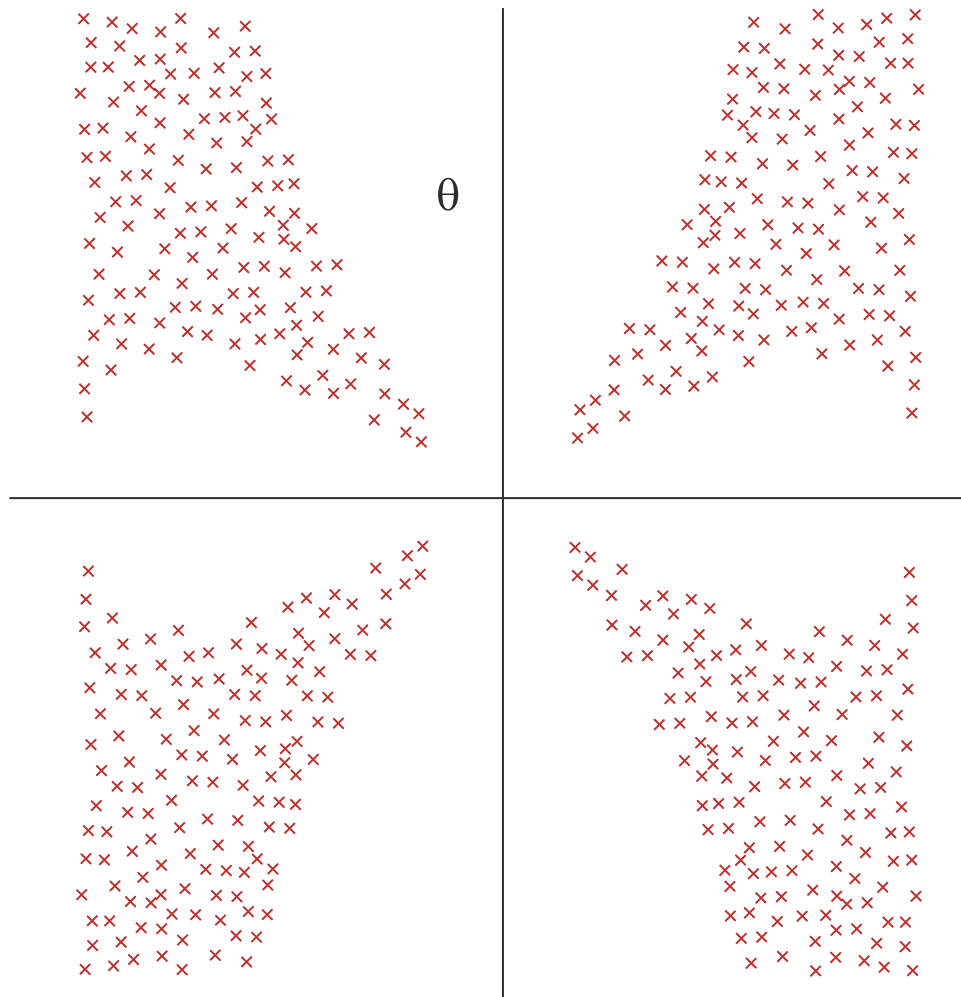
$$\frac{1}{\tilde{g}^2} = \frac{1}{g^2} - \frac{|\vartheta_I|}{8\pi^2}$$

[Lee-Yang Zeros in the partition function]

Lee-Yang Zeros

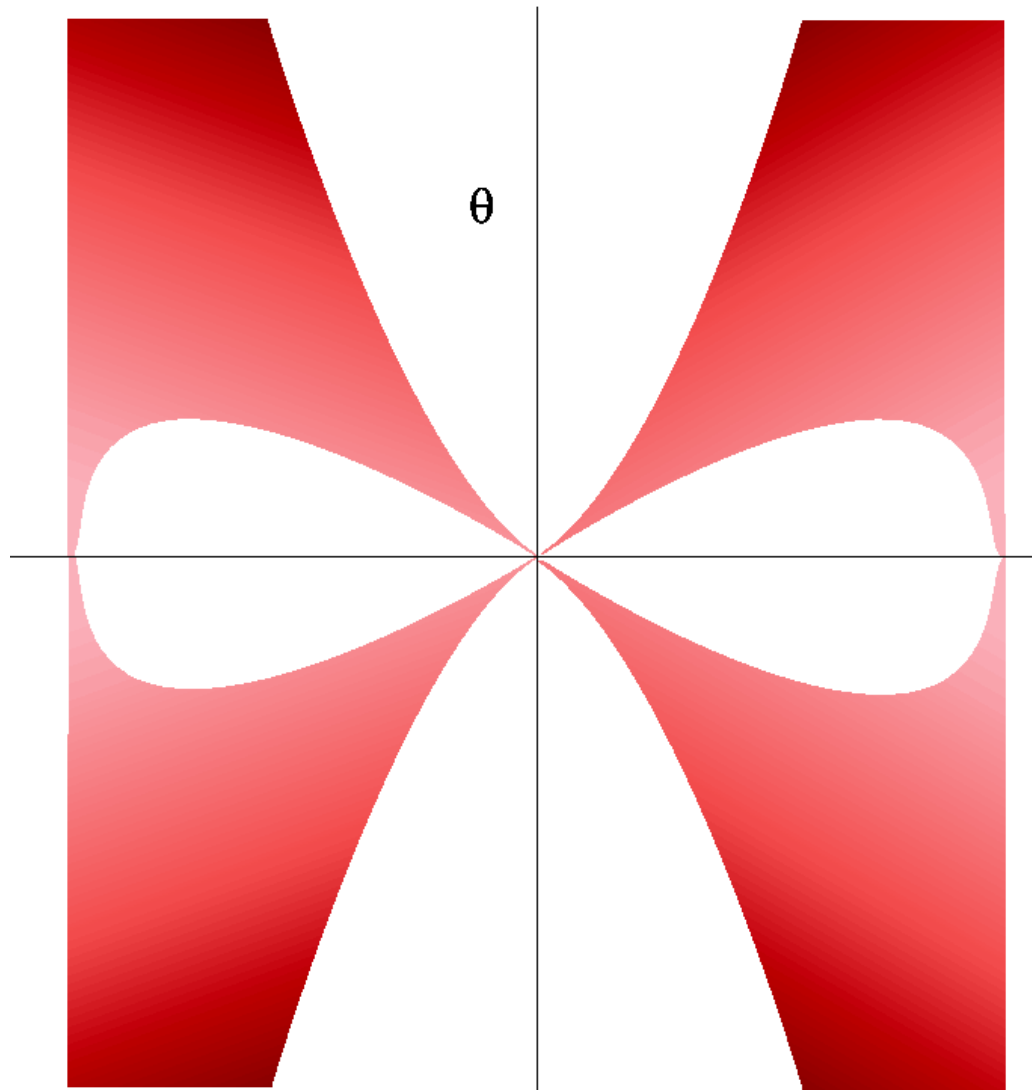
$$\mathcal{E}_V(\vartheta) = \frac{-2}{V} \operatorname{Re} \sum_{n=1}^{\infty} \rho_n \ln \left(1 - \frac{\sin^2 \frac{\vartheta}{2}}{\sin^2 \frac{\vartheta_n}{2}} \right) + \sum_{n=1}^{\infty} h_n(\vartheta, V)$$

ρ_n degree of degeneracy



Lee-Yang Zeros

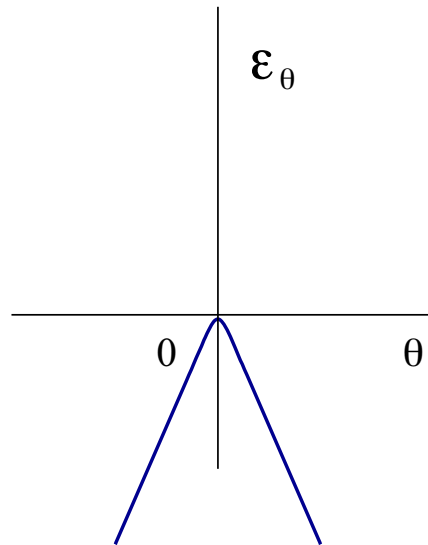
$$\mathcal{E}(\vartheta) = -2 \operatorname{Re} \int_0^\infty dt \rho(t) \ln \left(1 - \frac{\sin^2 \frac{\vartheta}{2}}{\sin^2 \frac{\vartheta(t)}{2}} \right) + \int_0^\infty dt h(t, \vartheta)$$



Vafa-Witten theorem

- Energy density $\mathcal{E}(\vartheta) \Leftrightarrow$ Electrostatic potential with negative charge density $-\rho > 0$

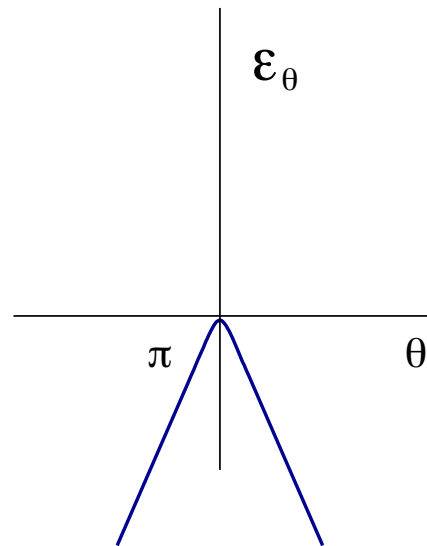
$$\mathcal{E}_V(\vartheta) \approx c|\vartheta|$$
$$c < 0$$



- Violation of Vafa-Witten inequality
 - CP symmetry is not preserved
 - There is no phase transition at $\vartheta = 0$

Vafa-Witten theorem

- The argument does not exclude the existence of a 1^{er} order phase transition at $\vartheta = \pi$



- CP symmetry can be spontaneously broken $\vartheta = \pi$
- 1^{er} order phase transition at $\vartheta = \pi$
- Even a 2nd order phase transition can appear at $\vartheta = 0$ if

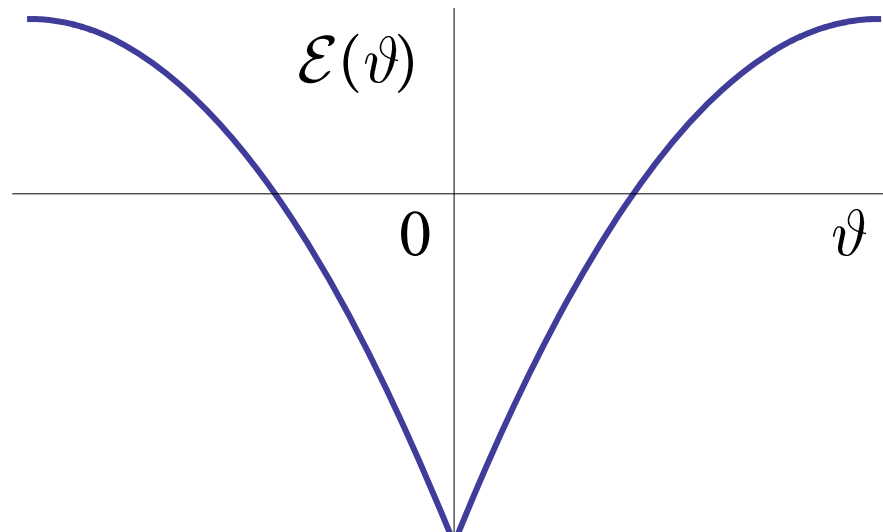
$$\rho(z) \sim \rho_1 z_2^\alpha \delta(z_1 - mz_2^\beta) \quad \text{with} \quad 0 < \alpha \leq 1$$

Vafa-Witten theorem

- Basic ingredients of the proof
 - Diamagnetic inequality
 - BPS bound

SCALAR FIELD THEORY IN **1+1** DIMENSIONS

$$S_{\vartheta}(\phi) = \frac{1}{2} \int d^2x \partial^\mu \phi^* \partial_\mu \phi - i\vartheta \int d^2x \varepsilon^{\mu\nu} \partial_\nu \phi^* \partial_\mu \phi$$



$\mathbb{C}P^{N-1}$ Sigma Models

2-D sigma model with $\mathbb{C}P^{N-1}$ target space

$$S[\Psi, \Psi^\dagger] = \frac{N}{2g^2} \int d^2x (D^\mu \Psi)^\dagger D_\mu \Psi$$

$$\Psi^\dagger \Psi = 1; \quad A_\mu = \frac{i}{2} \left(\Psi^\dagger \partial_\mu \Psi - \partial_\mu \Psi^\dagger \Psi \right)$$

- $U(1)$ gauge symmetry
- 1+1 D Poincaré symmetry
- $SU(N)$ global symmetry
- Conformal symmetry
- \mathcal{V} Vacuum

$\mathbb{C}P^{N-1}$ Sigma Models

2-D sigma model with $\mathbb{C}P^{N-1}$ target space

$$S[\Psi, \Psi^\dagger] = \frac{N}{2g^2} \int d^2x (D^\mu \Psi)^\dagger D_\mu \Psi - i \frac{\vartheta}{4\pi} \int \varepsilon_{\mu\nu} F_{\mu\nu}$$

- $U(1)$ gauge symmetry ✓
- 1+1 D Poincaré symmetry ✓
- $SU(N)$ global symmetry ✓
- Conformal symmetry ✗
- CP Invariance for $\vartheta = 0$ and $\vartheta = \pi$?

Exact Results

$\mathbb{C}P^1$ Sigma Model

1. $\vartheta = 0$ [Polyakov-Weigmann]

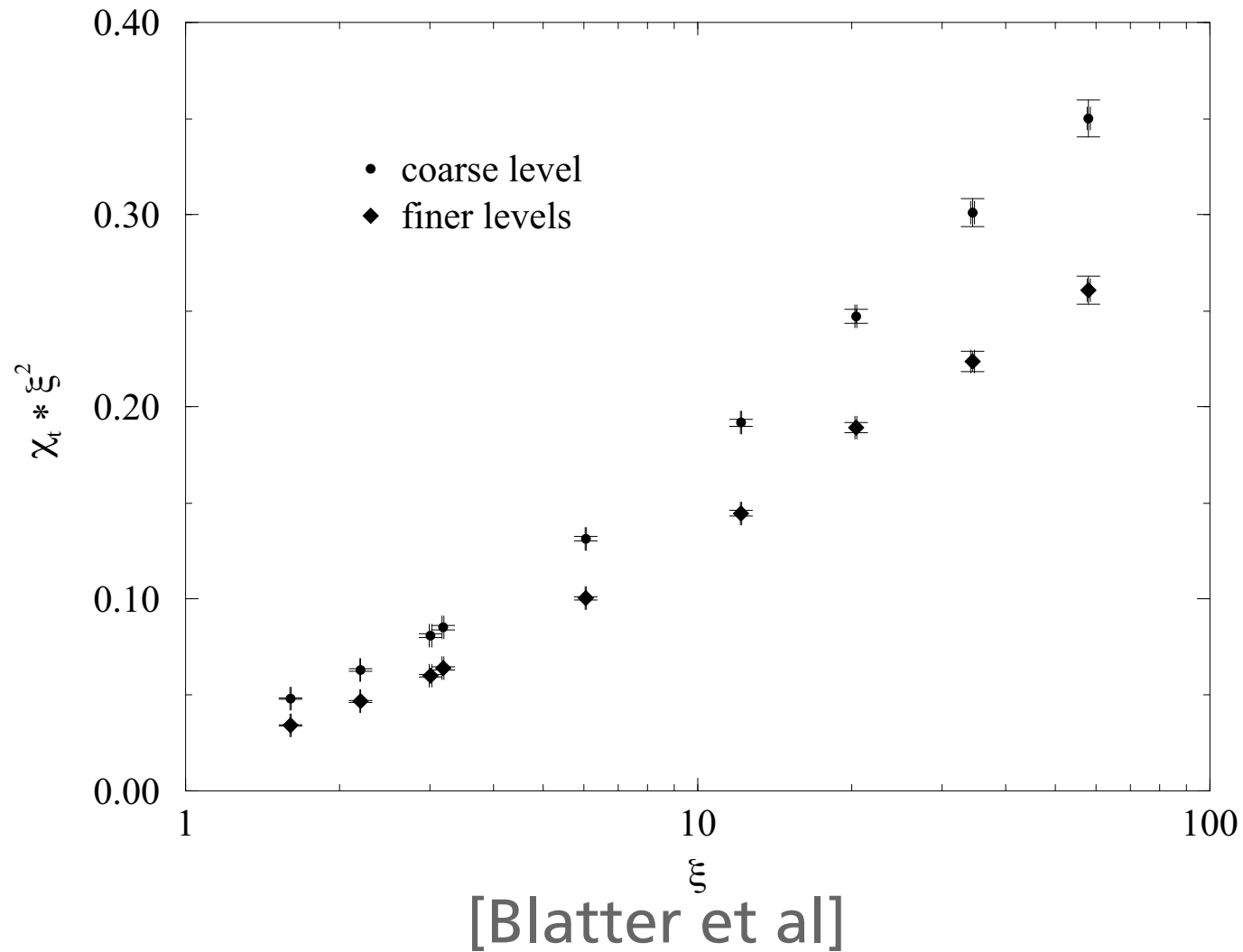
- Mass gap
- The lightless particle is in the **adjoint** $SU(2)$ representation
- No spontaneous **CP** symmetry breaking

2. $\vartheta = \pi$ [Zamolodchikov²]

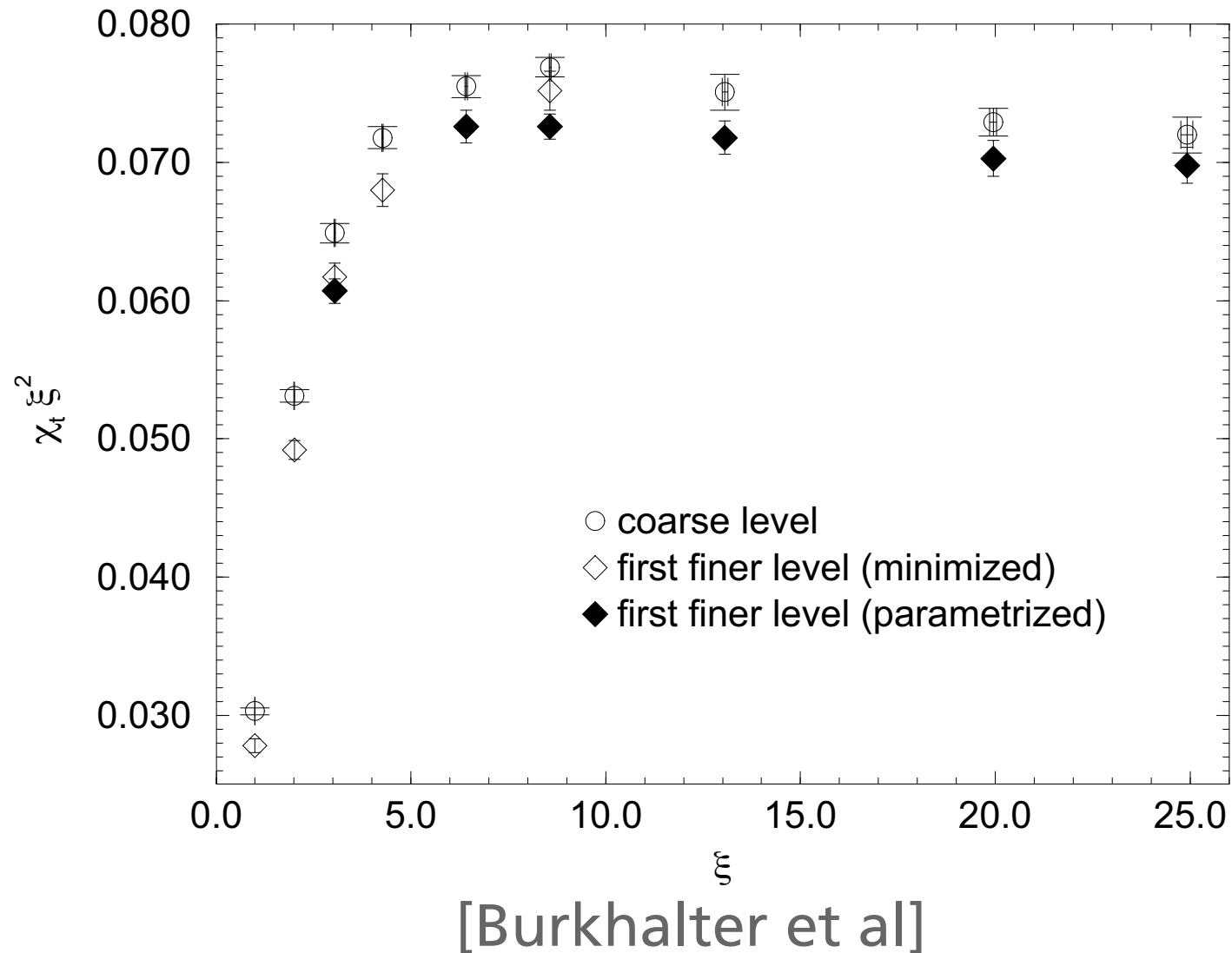
- Massless spectrum
- The basic massless particle is in the **fundamental** $SU(2)$ representation
- No spontaneous **CP** symmetry breaking

Topological Susceptibility in $\mathbb{C}P^1$

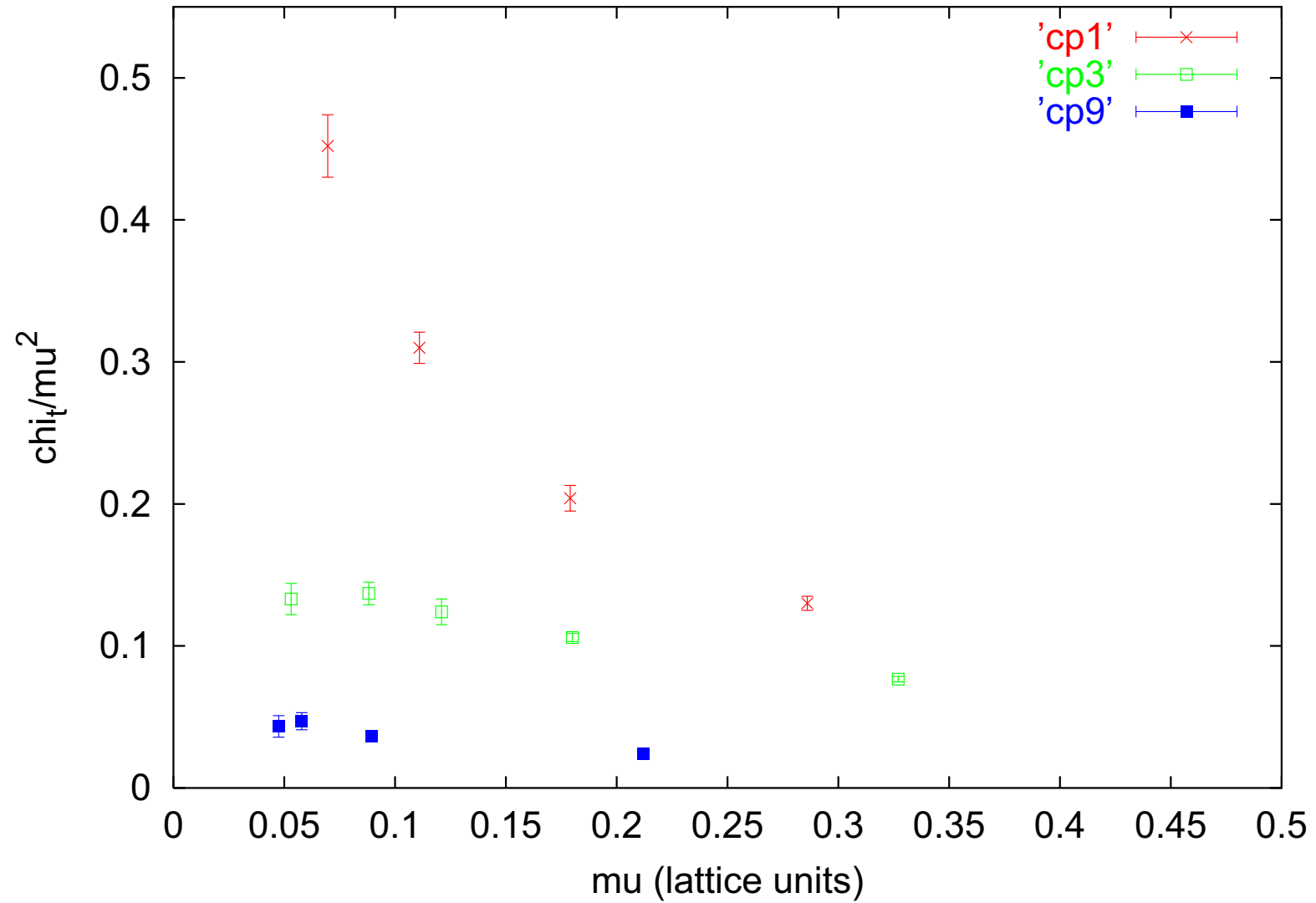
Lüscher's problem



Topological Susceptibility in $\mathbb{C}P^3$



Topological Susceptibility in $\mathbb{C}P^{N-1}$



[Ahmad et al]

CONCLUSIONS

- Vafa-Witten conjecture is correct [Theorem]
 - CP symmetry is preserved at $\vartheta = 0$
 - No 1^{er} order phase transitions $\vartheta = 0$
 - Topological susceptibility is positive thanks to a contact term
 - Analytic results are compatible with exact known results in the $\mathbb{C}P^1$ model. Phase transition at $\vartheta = 0$
 - They are also compatible with $N \gg 1$ limit results in $\mathbb{C}P^{N-1}$ models. CP symmetry is spontaneously broken at $\vartheta = \pi$
-

OPEN QUESTIONS

- Is topological susceptibility well defined in the $\mathbb{C}P^1$ sigma model at $\vartheta = 0$?
 - Does the theory confine in the $\vartheta = \pi$ vacuum?
 - There is a critical value N_c such that the $\mathbb{C}P^{N-1}$ model
 - undergoes a second order phase transition at $\vartheta = \pi$ for $N < N_c$ and
 - a first order phase transition for $N > N_c$
 - Is $N_c = 4$, or $N_c = 3$?
-