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International Centre for Theoretical Physics



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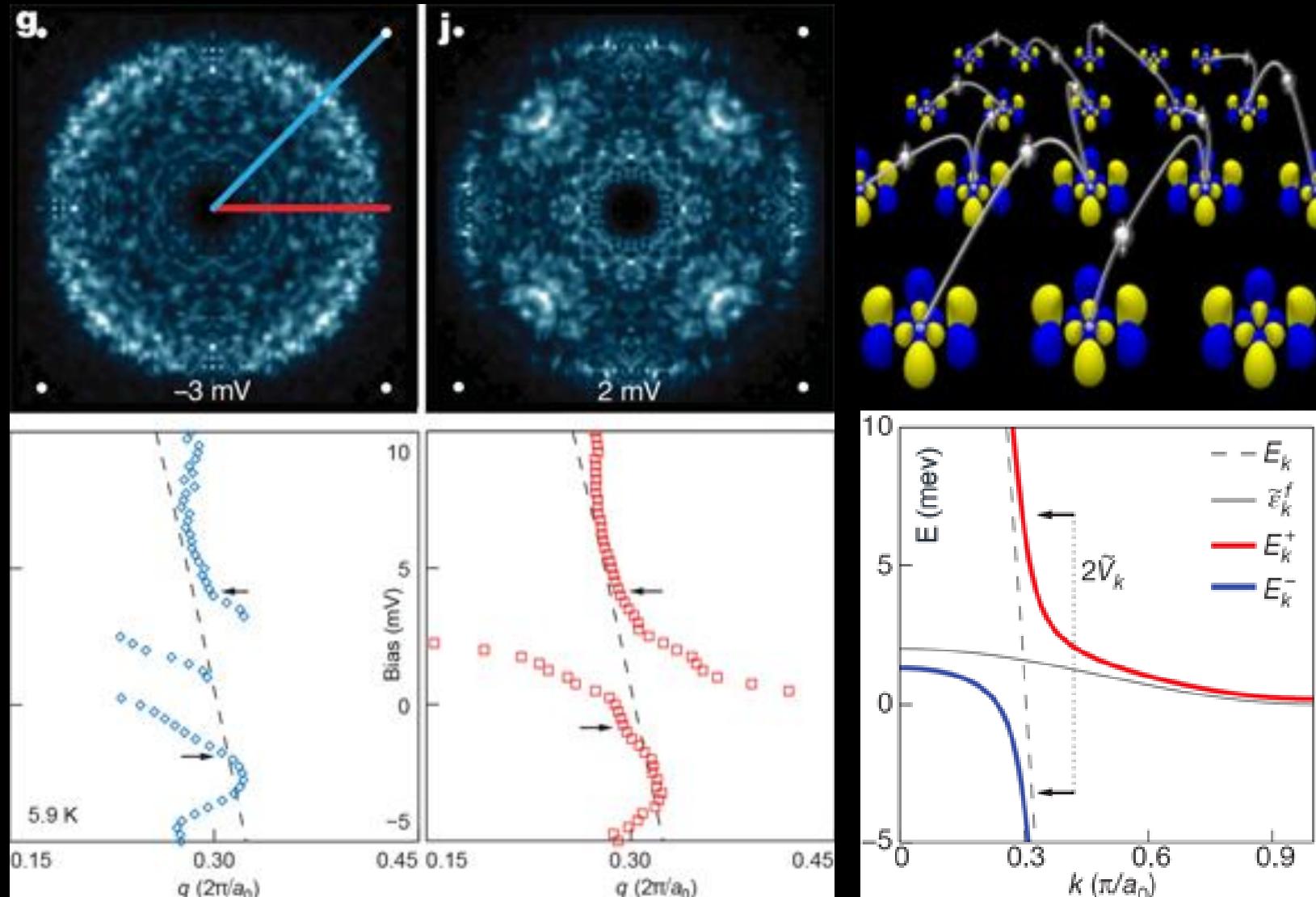
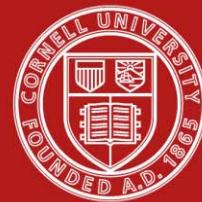
Workshop on Principles and Design of Strongly Correlated Electronic Systems

2 - 13 August 2010

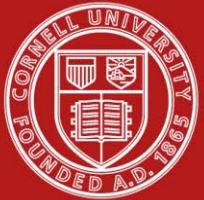
Imaging the Fano Lattice to "Hidden Order" Transition in URu₂Si₂

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Cornell University
Ithaca
NY 14853
U.S.A.*

Imaging the Fano Lattice to ‘Hidden Order’ Transition in URu_2Si_2



Collaboration



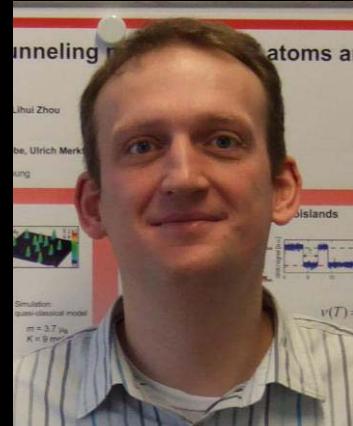
Dr. Andy Schmidt
Cornell /BNL



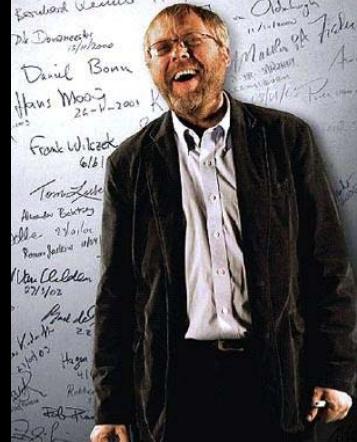
Mohammad Hamidian
Cornell/BNL



Dr. Peter Wahl
Cornell/
MPI Stuttgart



Dr. Focko Meier
Cornell/BNL



Dr. A.V. Balatsky
Los Alamos



Prof. Graeme Luke
McMaster



Travis Williams
McMaster



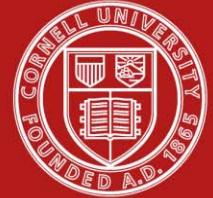
Dr. J .D. Garrett
(Brockhouse Institute)



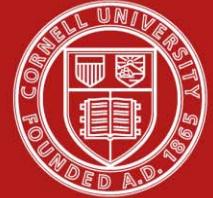
Prof. J.C. Davis
Cornell



Outline



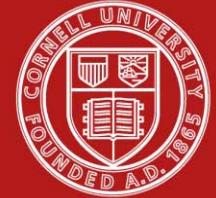
- Heavy Fermions, Kondo Effect / Kondo Lattice & Tunneling
- URu_2Si_2 : Introduction
- Spectroscopic Imaging STM & Heavy Fermion QPI
- URu_2Si_2 : Fano Lattice Imaging $T > T_0$
- URu_2Si_2 : Heavy Fermion QPI Imaging $T < T_0$
- New Perspectives from QPI on URu_2Si_2 ‘Hidden Order’
- Conclusions & Future Work



Heavy Fermions, Kondo Effect / Kondo Lattice & Tunneling

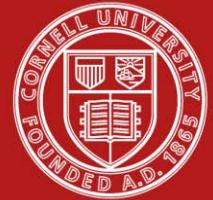


Heavy Fermions – Basics



- Electronic density of states up to 1000 times higher than copper at low temperatures
 - Seen in specific heat and magnetic susceptibility measurements
 - Heavy effective mass m^* - *INFERRRED*
- Partially filled *f-shell*
 - Matrix of localized magnetic moments immersed in a sea of conduction electrons

Single Impurity Kondo Effect



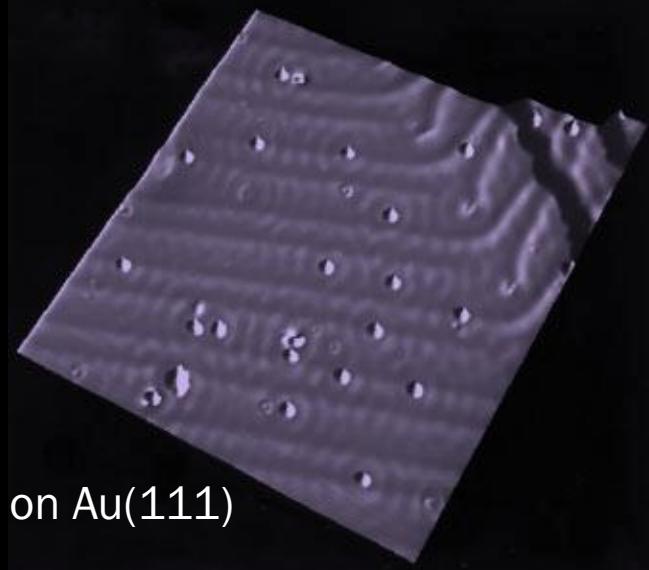
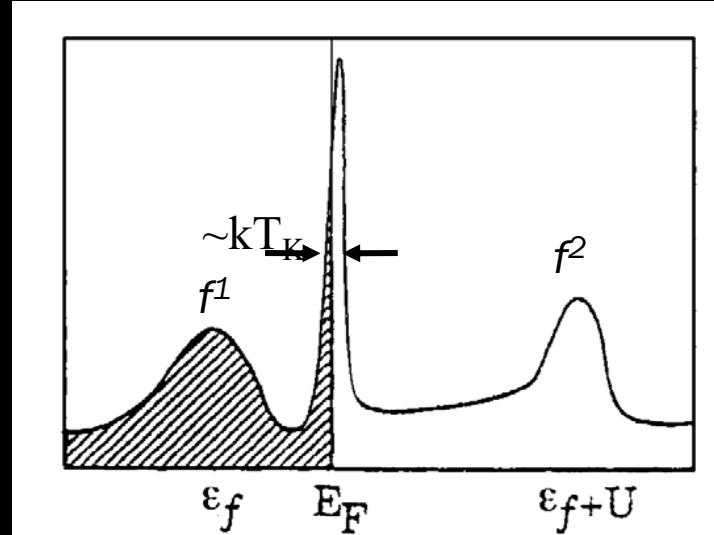
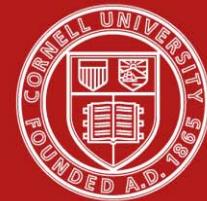
$T > T_K$

$T < T_K$

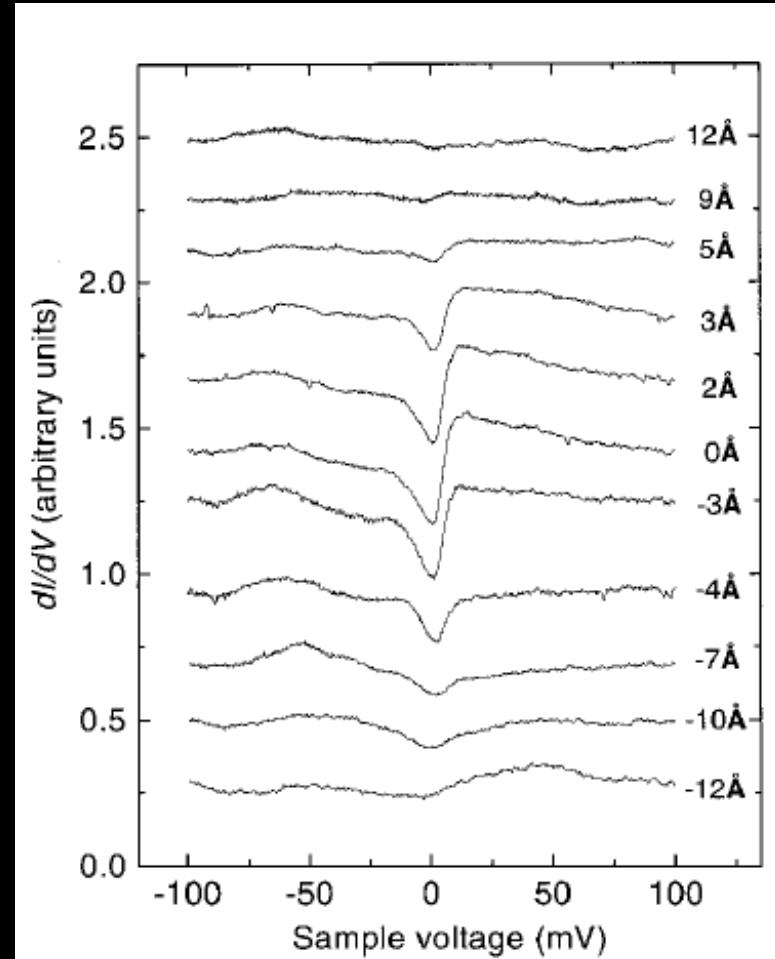
$$H = \sum_{k\sigma} \epsilon_{k\sigma} c_{k\sigma}^\dagger c_{k\sigma} + \sum_m \left(E_0 f_m^\dagger f_m + U f_m^\dagger f_m f_{-m}^\dagger f_{-m} \right) + \sum_{k,\sigma} (V_k e^{ik \cdot R_i} f_\sigma^\dagger c_{k\sigma} + H.C)$$



Fano Lineshape

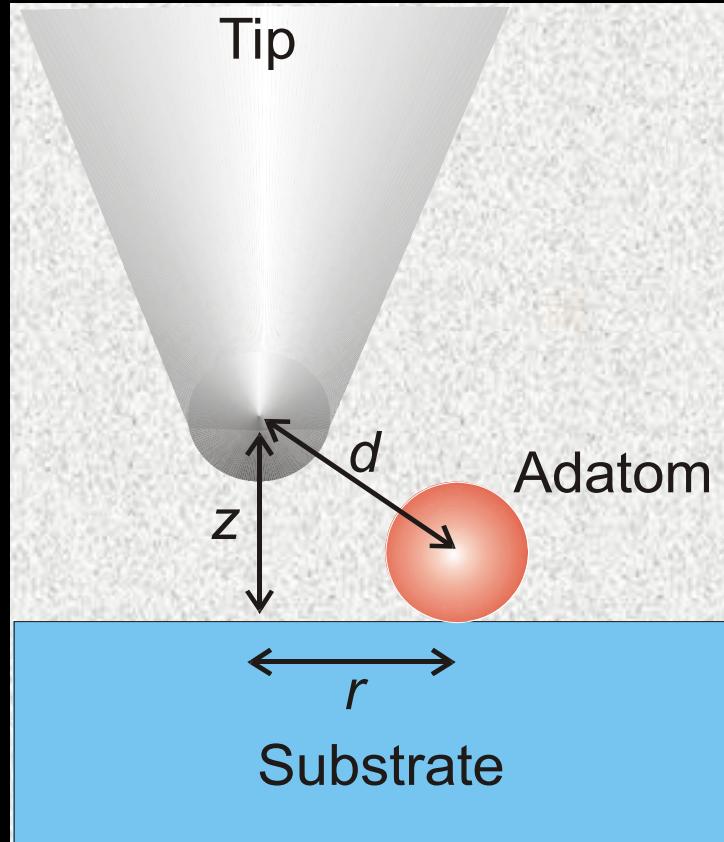
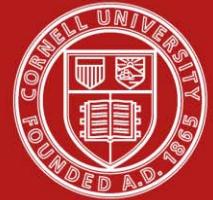


Co on Au(111)

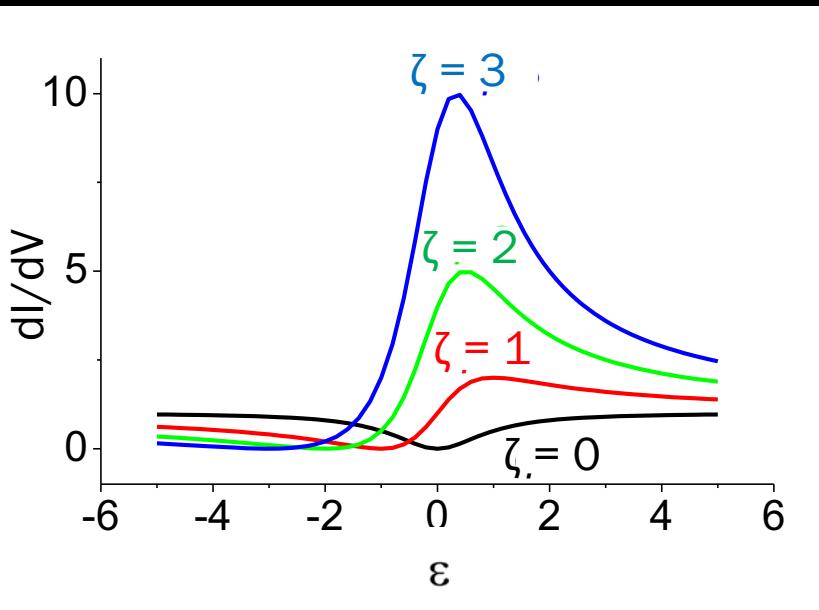


V. Madhavan *et al*, Science 280, 567 (1998)

Fano Lineshape



$$\frac{dI}{dV} \propto \frac{\zeta + \varepsilon'}{\varepsilon'^2 + 1}, \varepsilon' = \frac{(\varepsilon - \varepsilon_0)}{\Gamma/2}$$



ε_0 - energy of resonant state

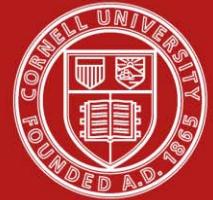
Γ - width of resonance

ζ - coupling ratio

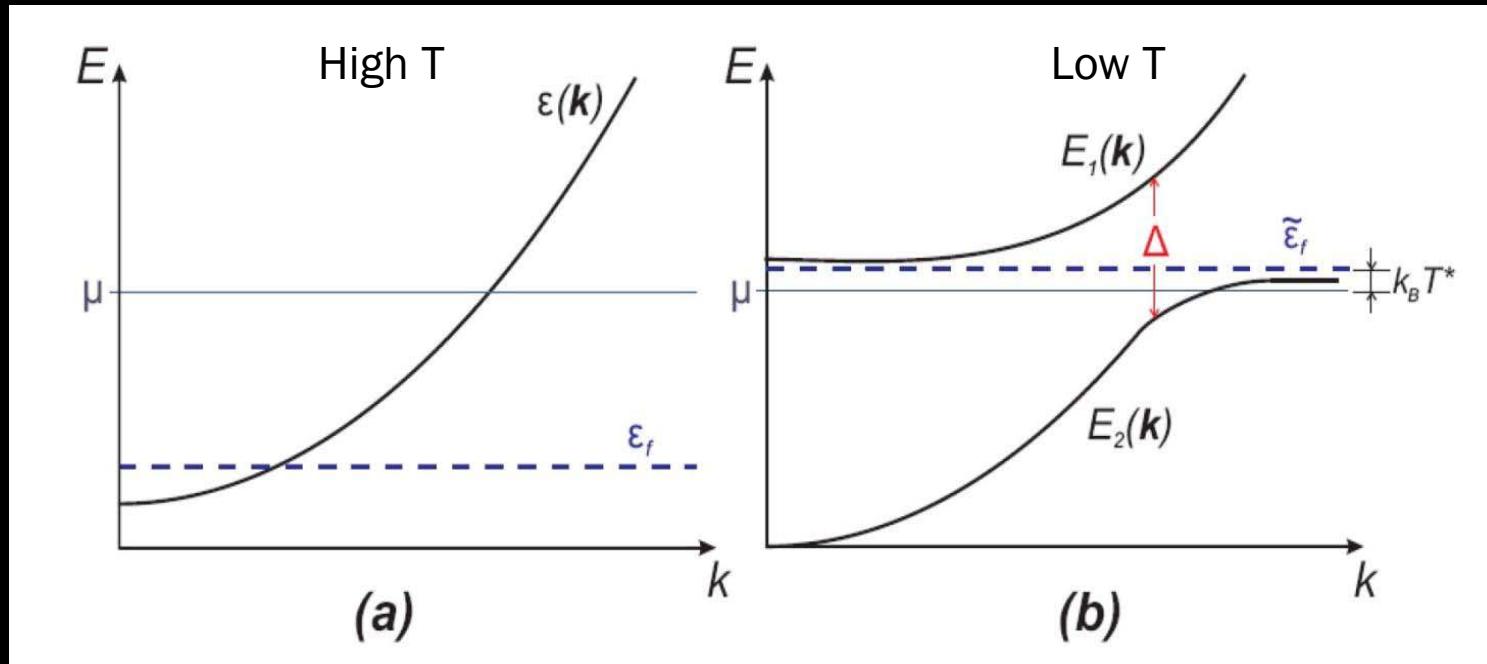
V. Madhavan *et al*, Science 280, 567 (1998)

M. Plihal und J.W. Gadzuk, Phys. Rev. B 63, 085404 (2001)

O. Újsághy, J. Kroha, L. Szunyogh und A. Zawadowski, Phys. Rev. Lett. 85, 2557 (2000)

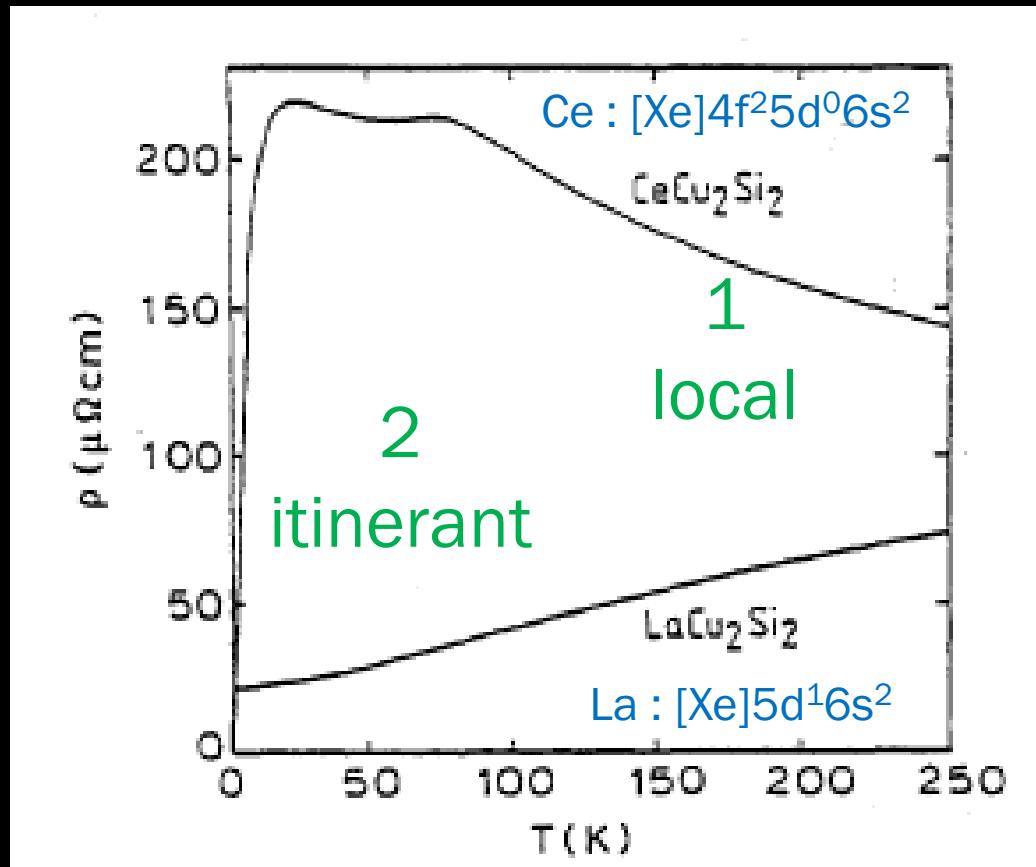
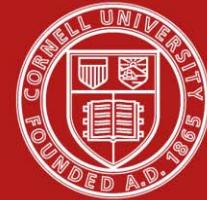


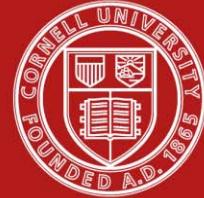
$$H = \sum_{k\sigma} \epsilon_{k\sigma} c_{k\sigma}^\dagger c_{k\sigma} + \sum_{i,m} \left(E_0 f_{im}^\dagger f_{im} + U f_{im}^\dagger f_{im} f_{i-m}^\dagger f_{i-m} \right) + \sum_{i,k,\sigma,m} \left(V_k e^{ik \cdot R_i} f_{im}^\dagger c_{k\sigma} + H.C \right)$$



$$E_{\pm} = \frac{E_k + \varepsilon_k^f}{2} \pm \left[\left(\frac{E_k - \varepsilon_k^f}{2} \right)^2 + |V_k|^2 \right]^{1/2}$$

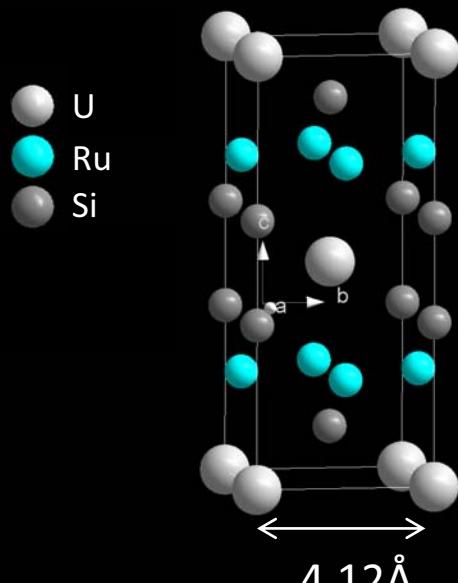
Path to Heaviness





URu₂Si₂

URu₂Si₂



$a=4.124\text{\AA}$ $c=9.582\text{\AA}$

Specific heat coefficient $\gamma = 65\text{mJ/mol/K}^2$

Effective mass $m^* = 25m_e - 50m_e$

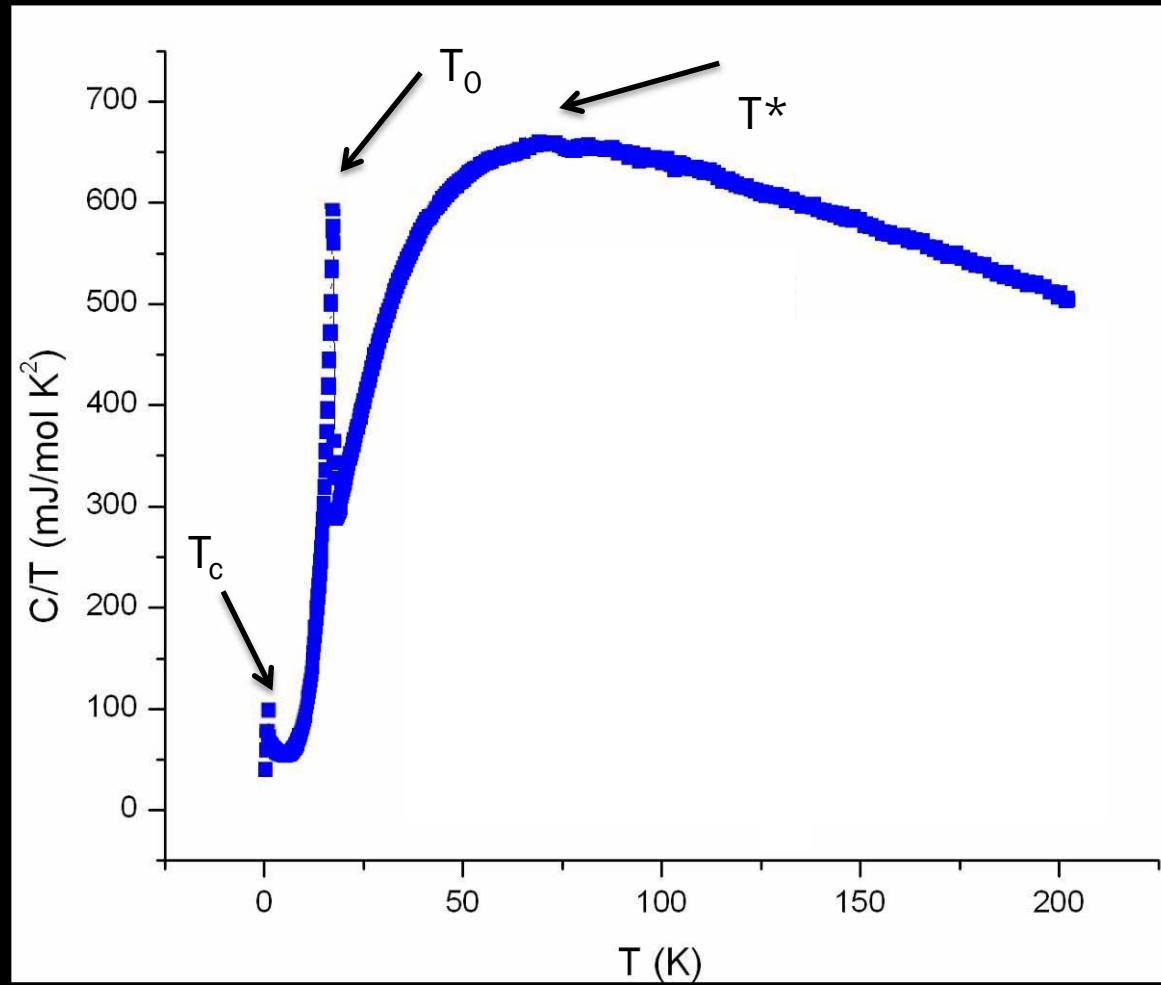
HF Coherence Temperature $T^* \sim 55\text{K}$

'Hidden Order' transition $T_N = 17.5\text{K}$

Superconducting transition $T_c = 1.5\text{K}$

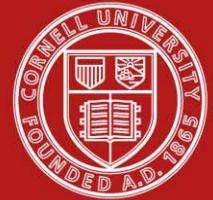
Palstra *et. al.* *PRL* **55**, 2727 (1985)

Maple *et. al.* *PRL* **56**, 185 (1986)





'Hidden Order' in URu_2Si_2



Energy-Gap $\sim 11\text{mV}$
(optical & specific heat)

Reorganization of band-structure and
magnetic excitation spectrum

- Palstra, T.T.M., Menovsky, A.A., & Mydosh, J.A. Superconducting and magnetic transitions in the heavy-fermion system URu_2Si_2 . *Phys. Rev. Lett.* **55**, 2727-2730 (1985).
- Broholm, C. *et al.* Magnetic excitations and order in the heavy-electron superconductor URu_2Si_2 . *Phys. Rev. Lett.* **58**, 1467-1470 (1987).
- Bonn, D.A. *et al.* Far-infrared properties of URu_2Si_2 . *Phys. Rev. Lett.* **61**, 1305-1308 (1988).
- Wiebe, C.R. *et al.* Gapped Itinerant spin excitations account for missing entropy in the hidden order state of URu_2Si_2 . *Nature Phys.* **3**, 96-99 (2007).
- Santander-Syro, A.F. *et al.* Fermi-surface instability at the 'hidden-order' transition of URu_2Si_2 . *Nature Phys.* **5**, 637-641 (2009).

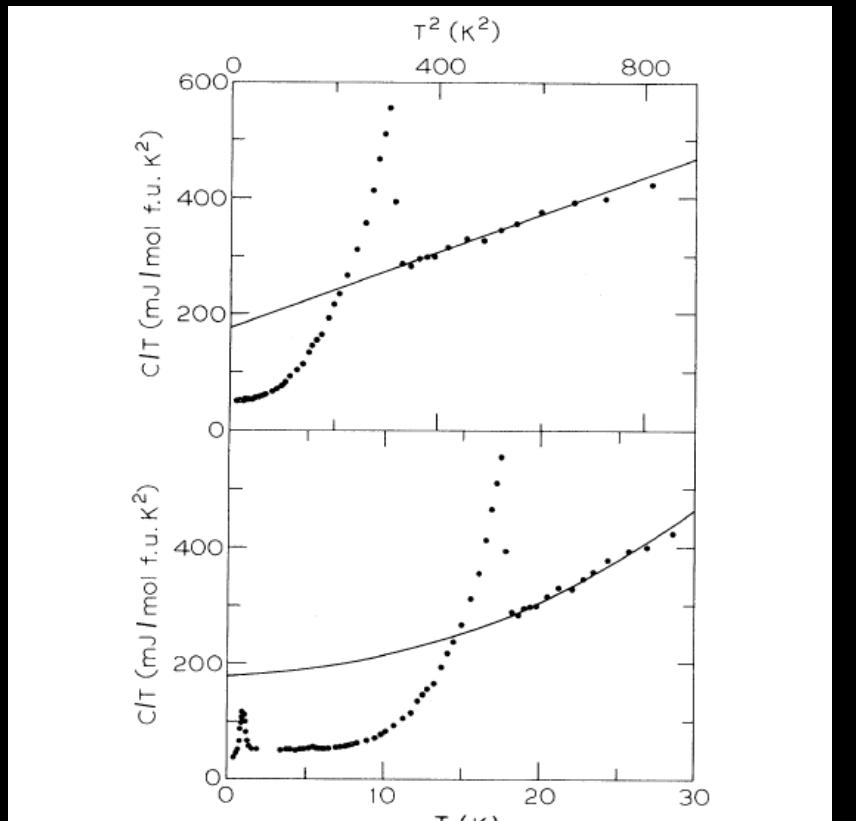
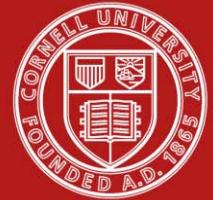


FIG. 1. Specific heat of URu_2Si_2 plotted as C/T vs T^2 (above) yielding γ and Θ_D , and as C/T vs T (below) showing the entropy balance.



Hypotheses for Identity of OP



Susceptibility of FL/FS Momentum Space

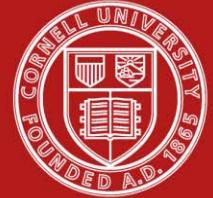
- Broholm, C. *et al.* Magnetic excitations in the heavy-fermion superconductor URu_2Si_2 . *Phys. Rev. B*. **43**, 809-822 (1991).
- Ikeda, H. & Ohashi, Y. Theory of unconventional spin density wave: a possible mechanism U-based heavy fermion compounds. *Phys. Rev. Lett.* **81**, 3723-3726 (1998).
- Chandra, P. *et al.* Hidden orbital order in the heavy fermion metal URu_2Si_2 . *Nature* **417**, 831-834 (2002).
- Varma, C.M. & Lijun, Z. Helicity order: Hidden order parameter in URu_2Si_2 . *Phys. Rev. Lett.* **96**, 036405-1-036405-4 (2006).
- Balatsky, A.V. *et al.* Incommensurate spin resonance in URu_2Si_2 . *Phys. Rev. B*. **79**, 214413 (2009).

'Altered' Kondo Effect / Real Space

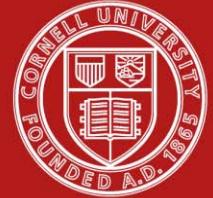
- Santini, P. Crystal field model of the mag properties of URu_2Si_2 . *Phys. Rev. Lett.* **73**, 1027-1030 (1994).
- Barzykin, V. & Gor'kov, L.P. Singlet magnetism in heavy fermions. *Phys. Rev. Lett.* **74**, 4301-4304 (1995).
- Haule K. & Kotliar G. Arrested Kondo effect and hidden order in URu_2Si_2 . *Nature Phys.* **5**, 796-799 (2009).
- Harima H., Miyake K., Flouquet J. Why the hidden order in URu_2Si_2 is still hidden - one simple answer. arXiv:1001.2369



Key Issues



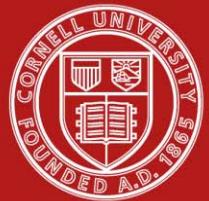
- What is the relationship between the initial Kondo Lattice and the ‘Hidden Order’ state?
- What are the alterations to real- and momentum-space electronic structure due to the onset of the ‘Hidden Order’?
- Can one distinguish between FS/FL susceptibility and local mechanisms?



Spectroscopic Imaging Scanning Tunneling Microscopy (SI-STM)



Our SI-STM Facilities



STM1 (9T/250mK)

Cornell



STM3 (4K→100K)

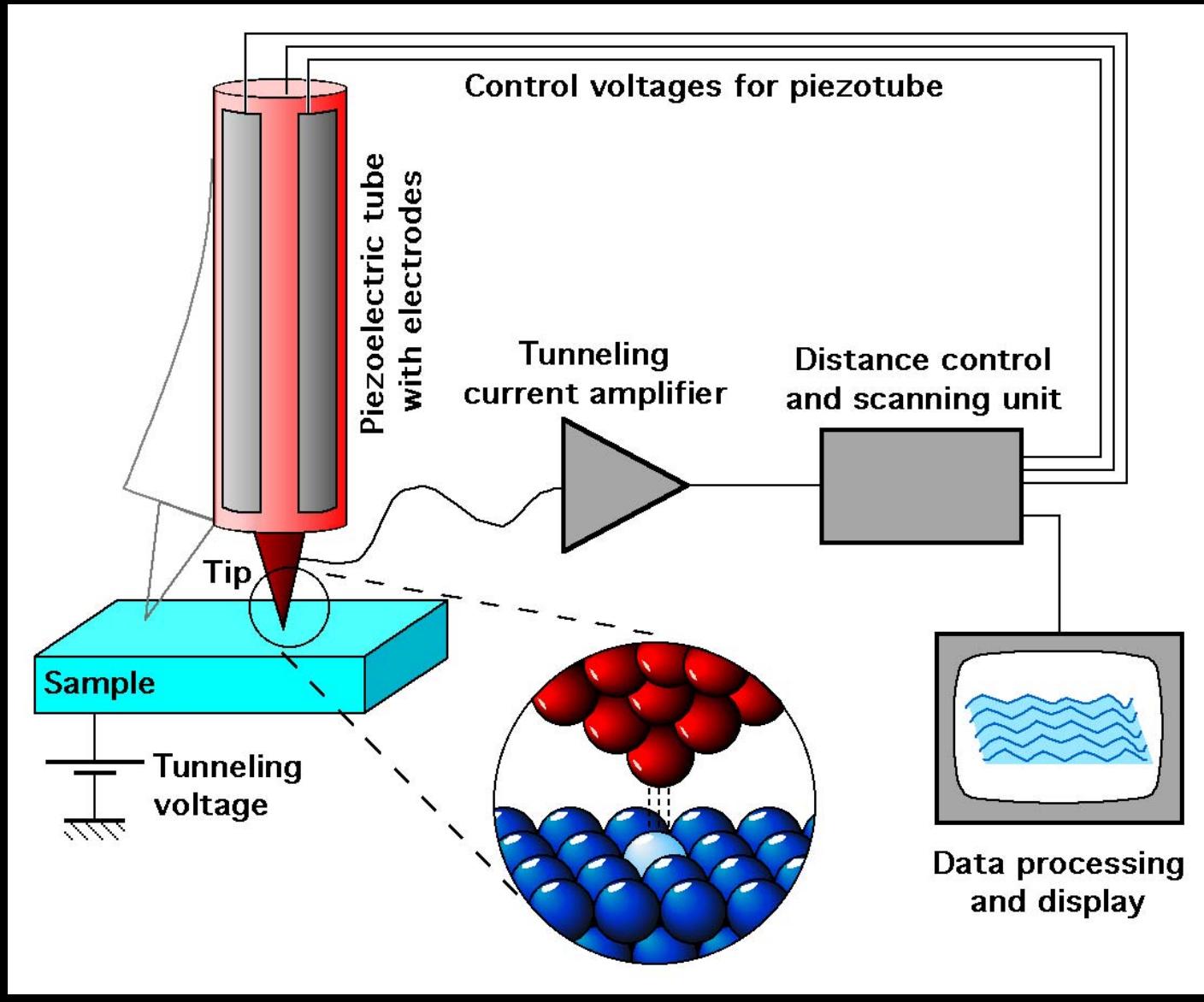
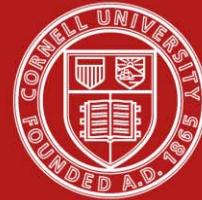
Brookhaven



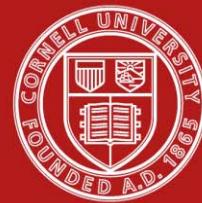
STM2(9T/10mK)

Cornell





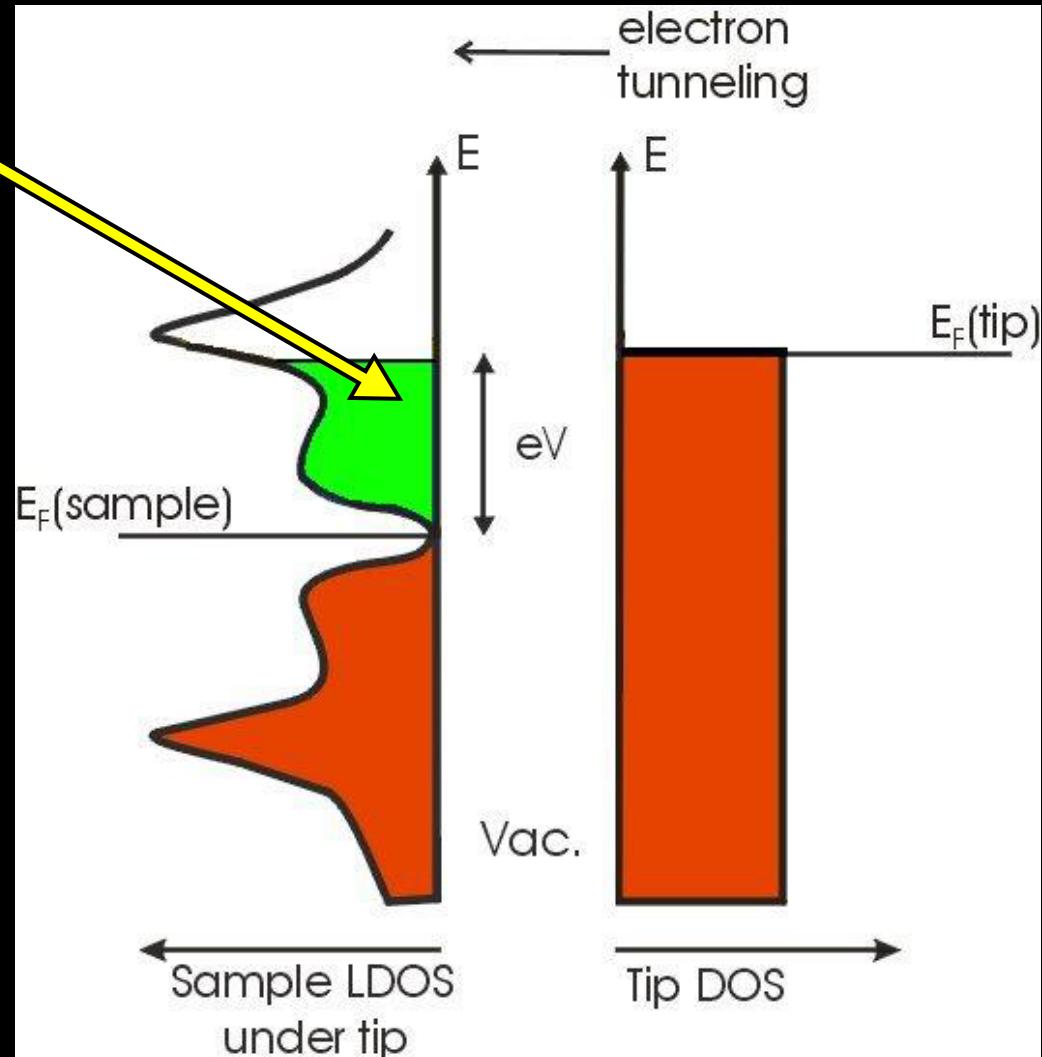
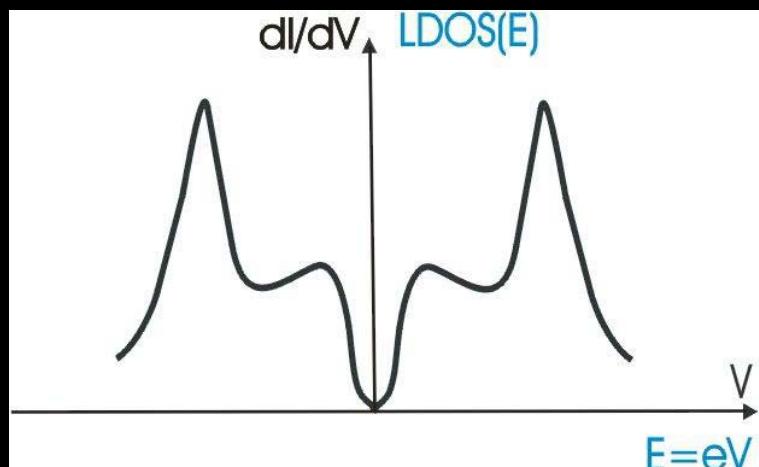
DOS(\vec{r}, E) from $dI/dV(\vec{r}, V)$ Spectra



$$I = C e^{-\frac{z(V)}{z_0}} \int_0^V LDOS(\vec{r}, E) dE$$

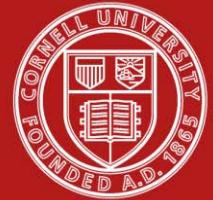
$$\Rightarrow \frac{dI}{dV} \propto LDOS(E = eV)$$

Point Spectrum

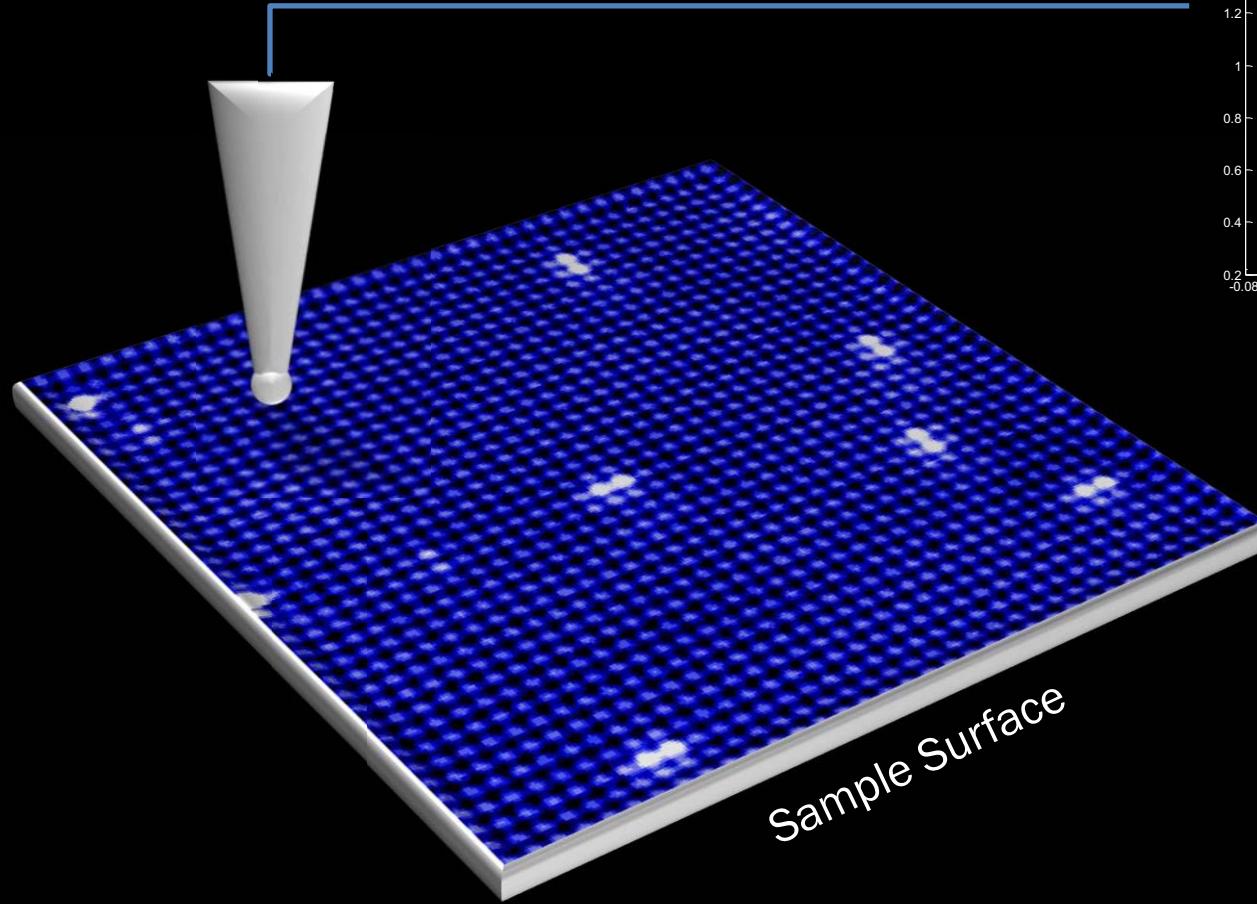




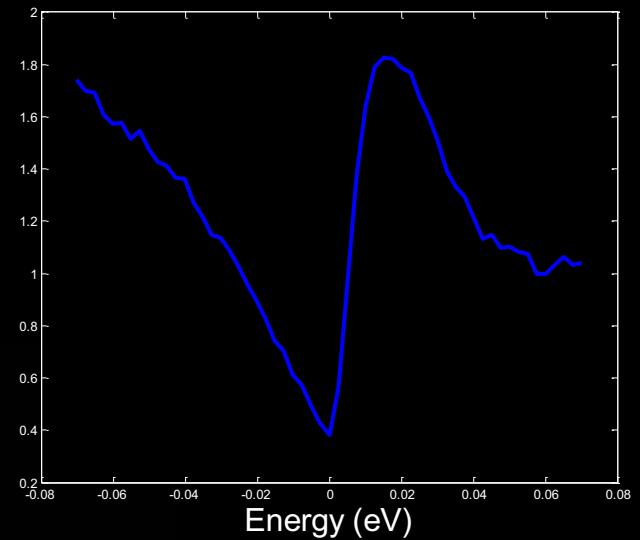
Spectroscopic Imaging



$$\frac{dI}{dV}(r, V) \equiv g(r, V) \propto \text{LDOS}(r, E = eV)$$

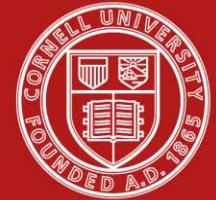


Sample Surface

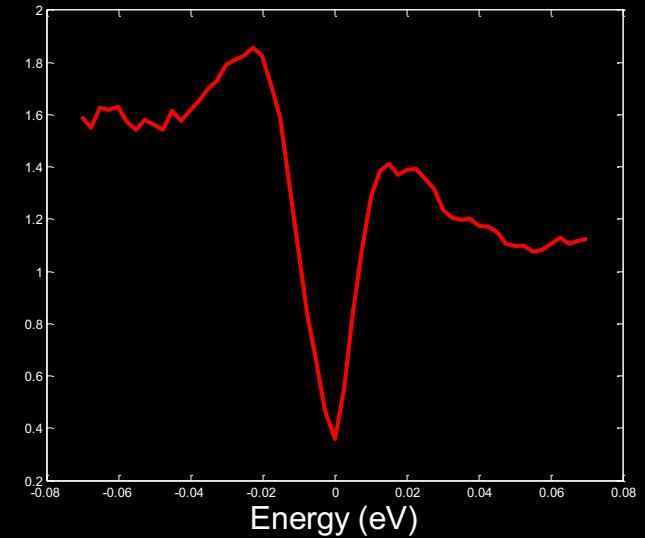
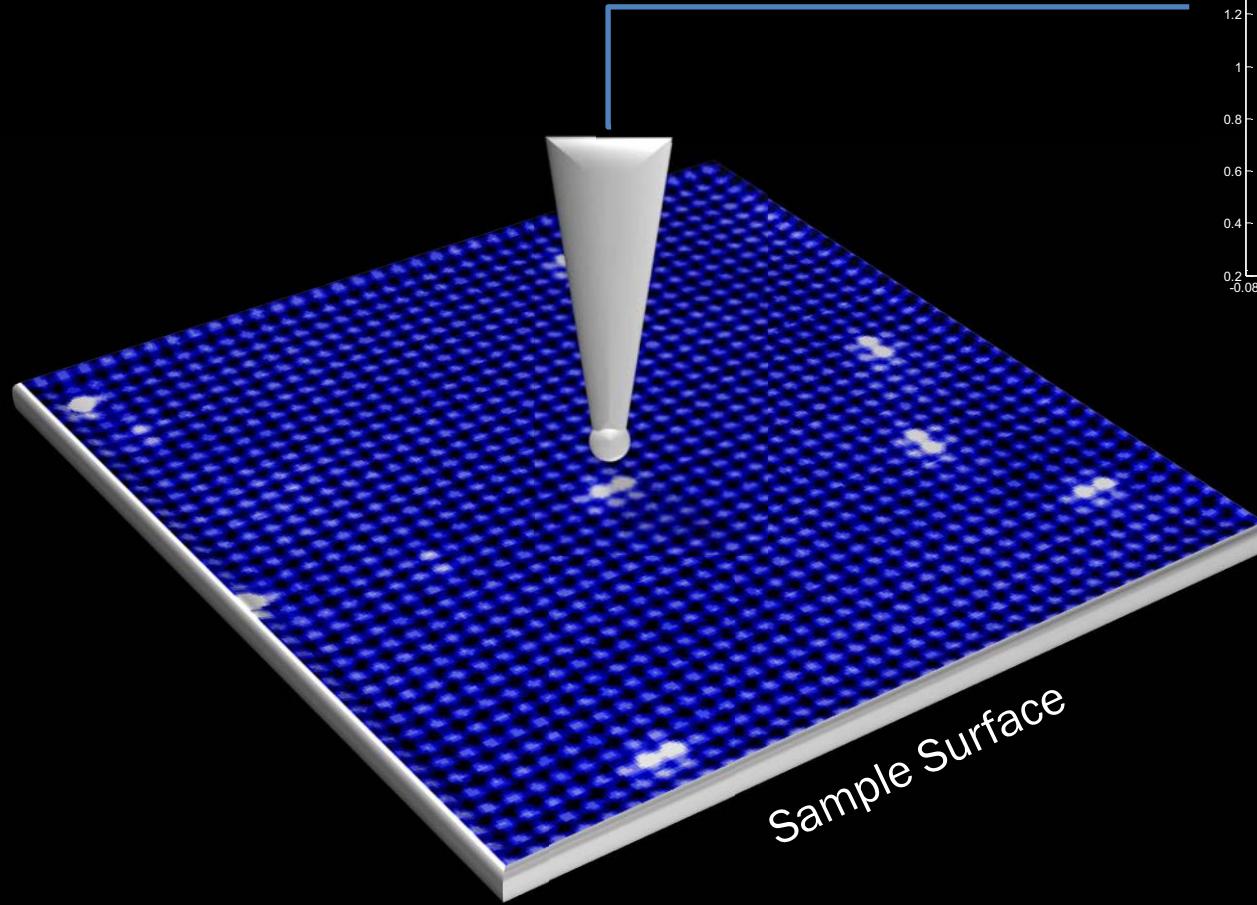


$\text{LDOS}(r = r_1, E)$

Spectroscopic Imaging



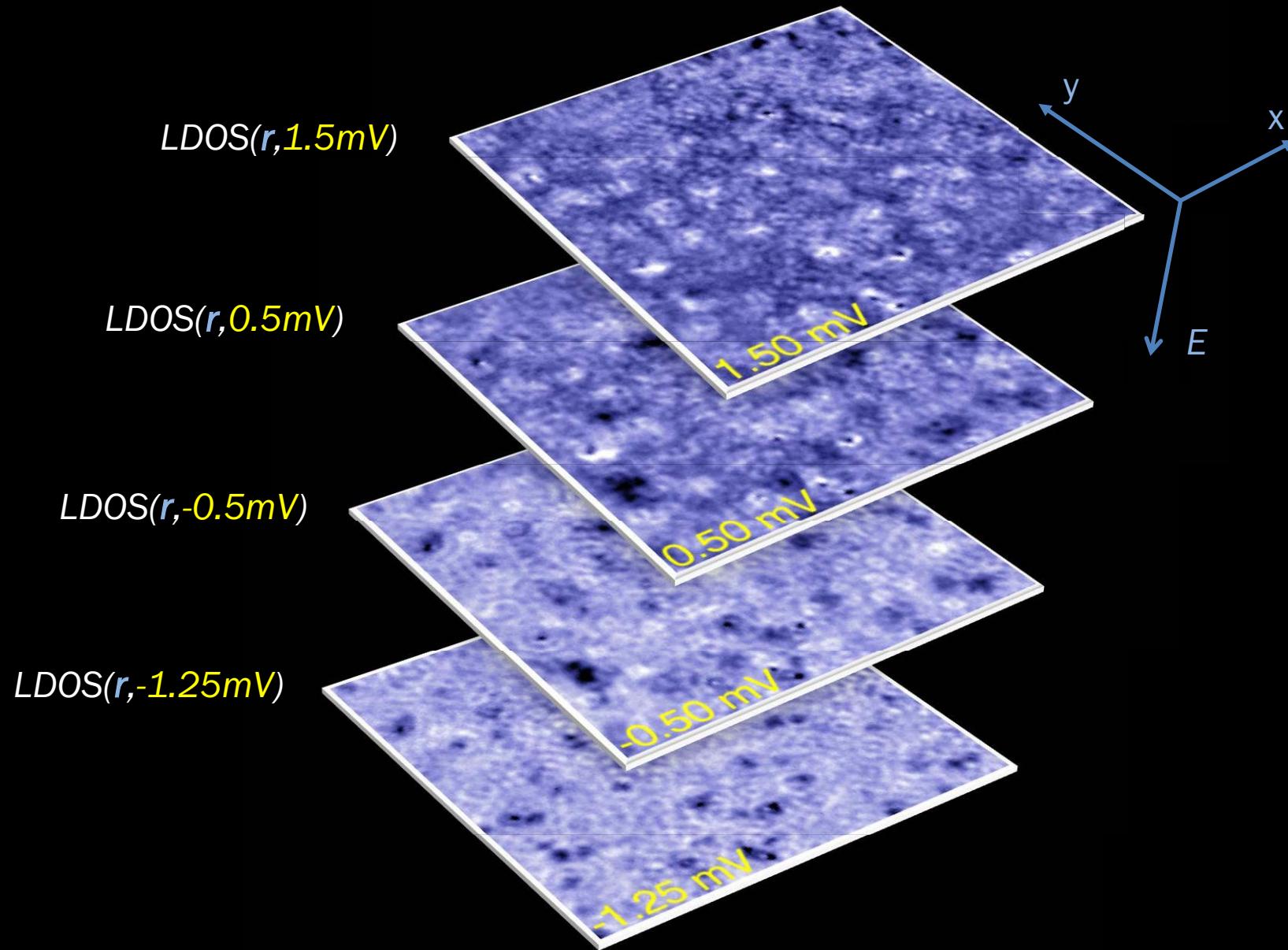
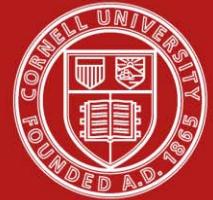
$$\frac{dI}{dV}(r, V) \equiv g(r, V) \propto \text{LDOS}(r, E = eV)$$



$\text{LDOS}(r = r_2, E)$

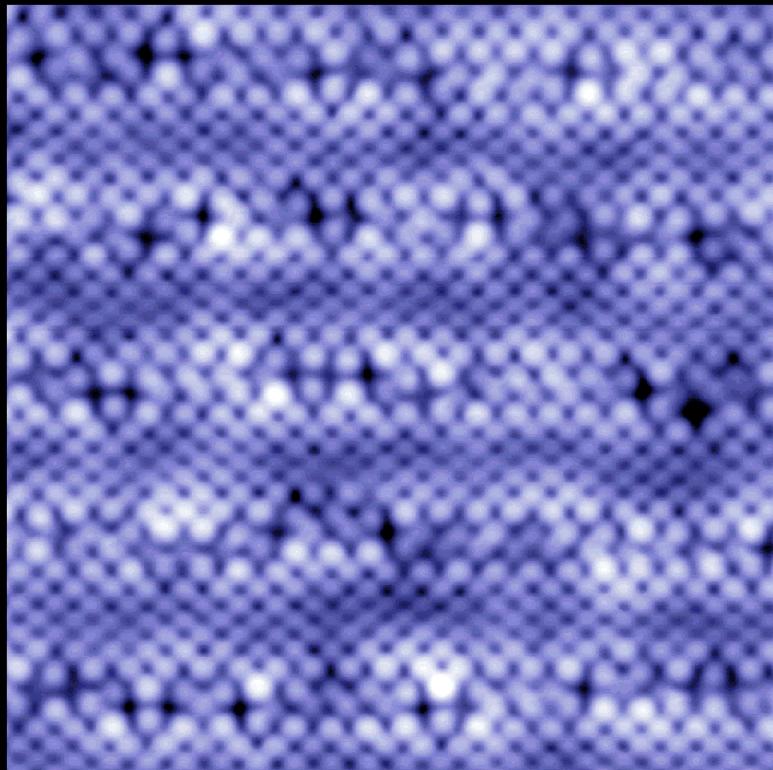
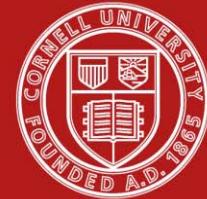


Spatially and Energy Resolved LDOS

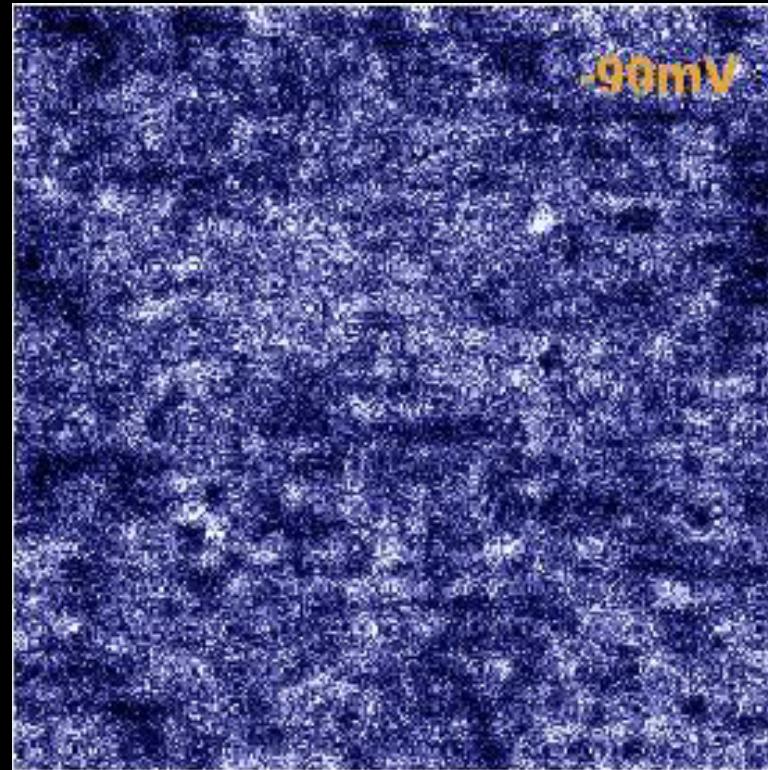




Example: Bi-2212



Topographic
Image

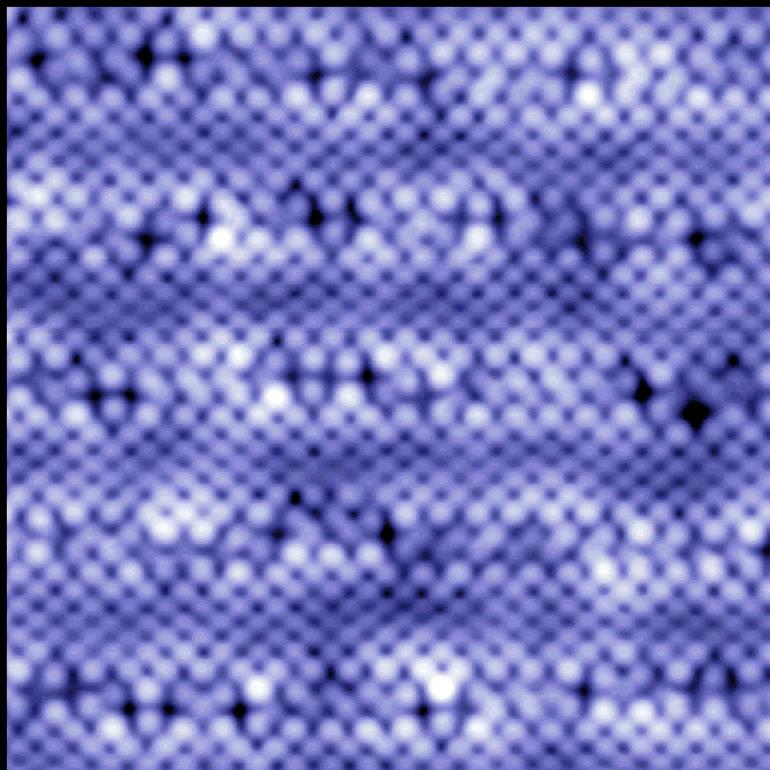
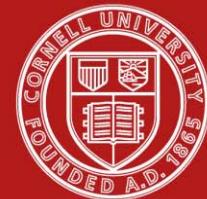


$g (r,V)$

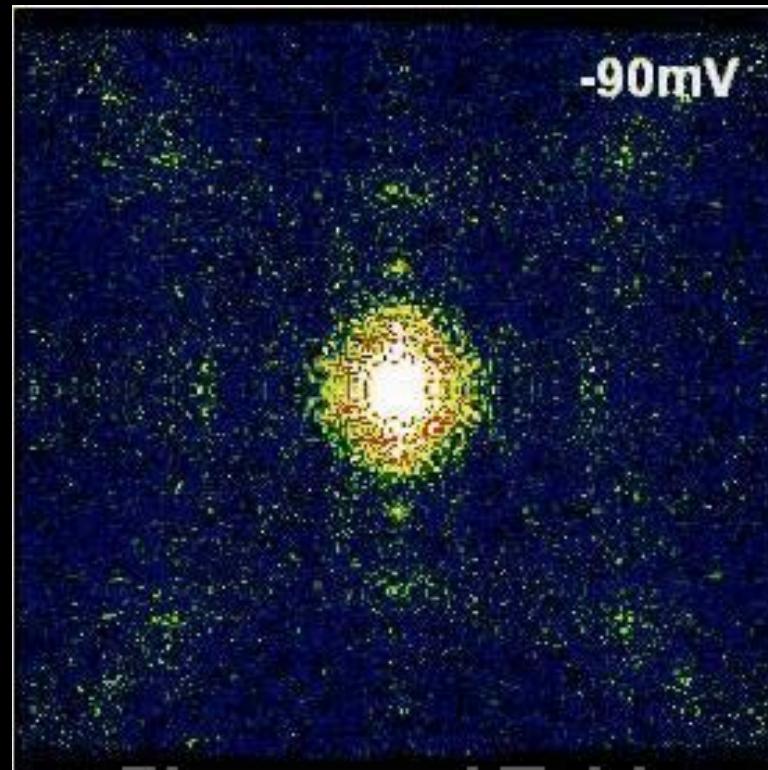
Nature doi:10.1038/nature09169 (2010)



QPI: Bi-2212



Topographic
Image

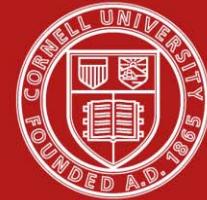


$g(q,V)$

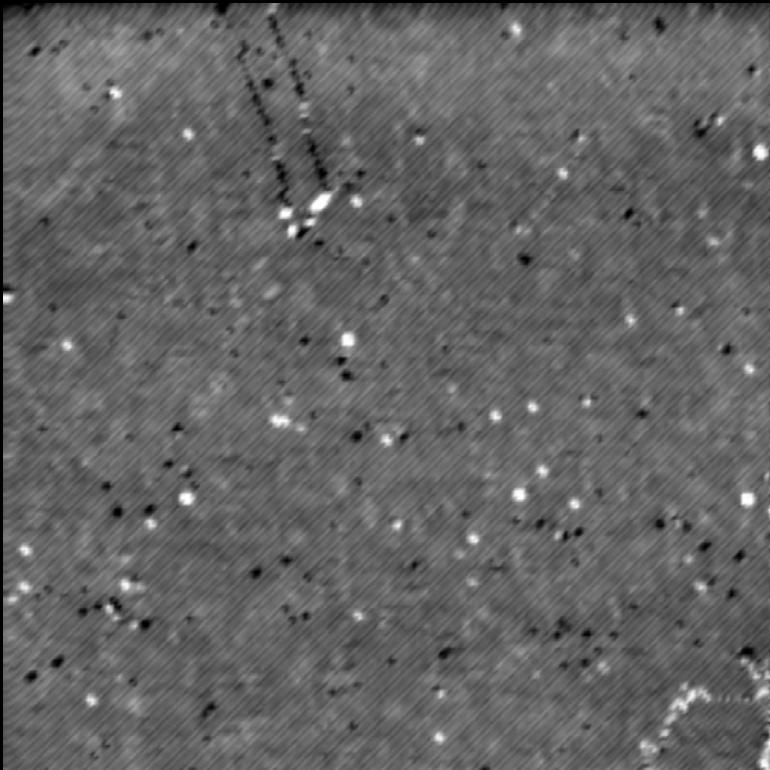
Nature doi:10.1038/nature09169 (2010)



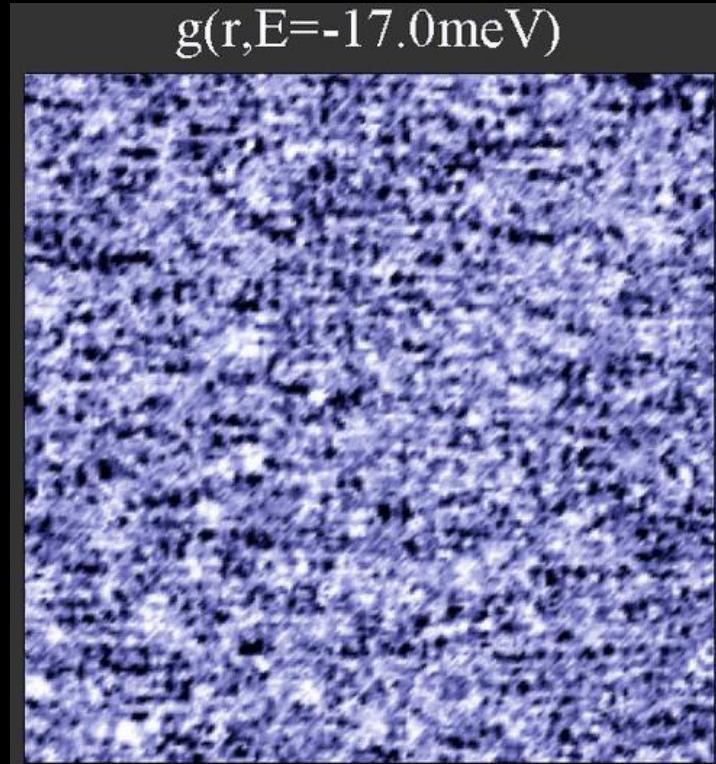
Example: Ca-122



Topography



Topographic
Image

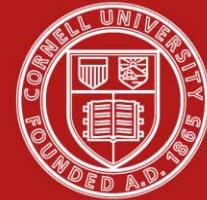


$g(r, V)$

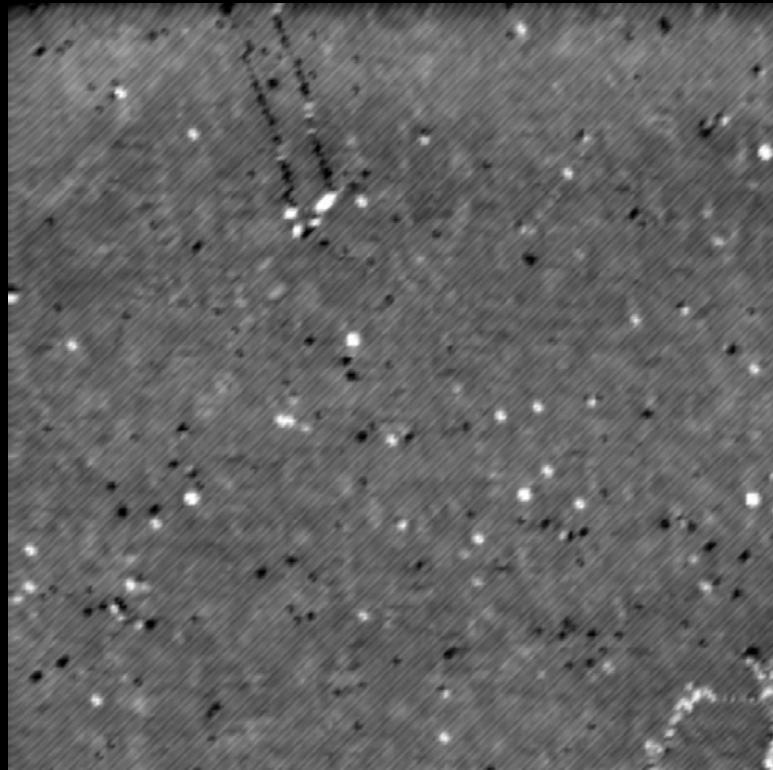
Science 327, 181 (2010)



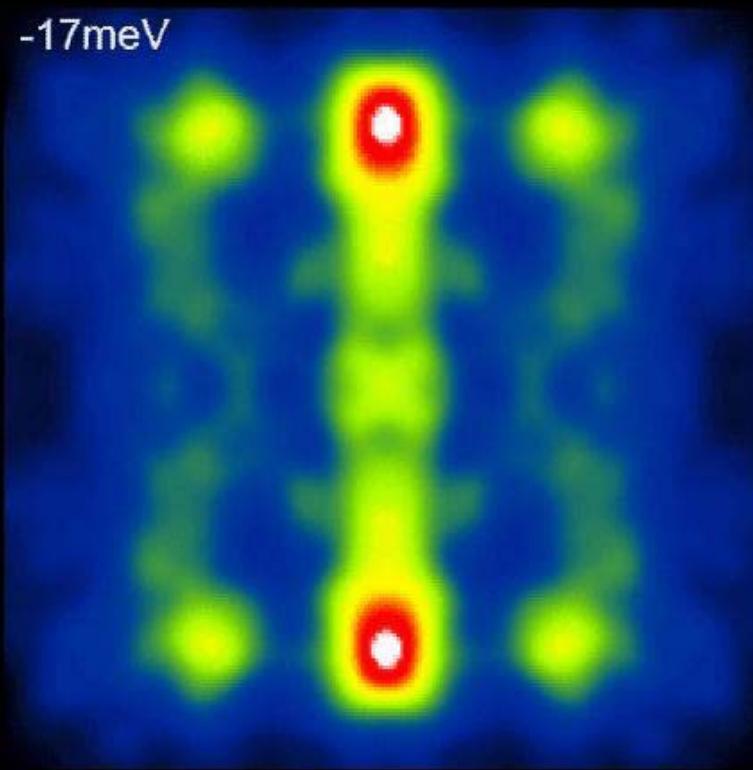
QPI: Ca-122



Topography



Topographic
Image

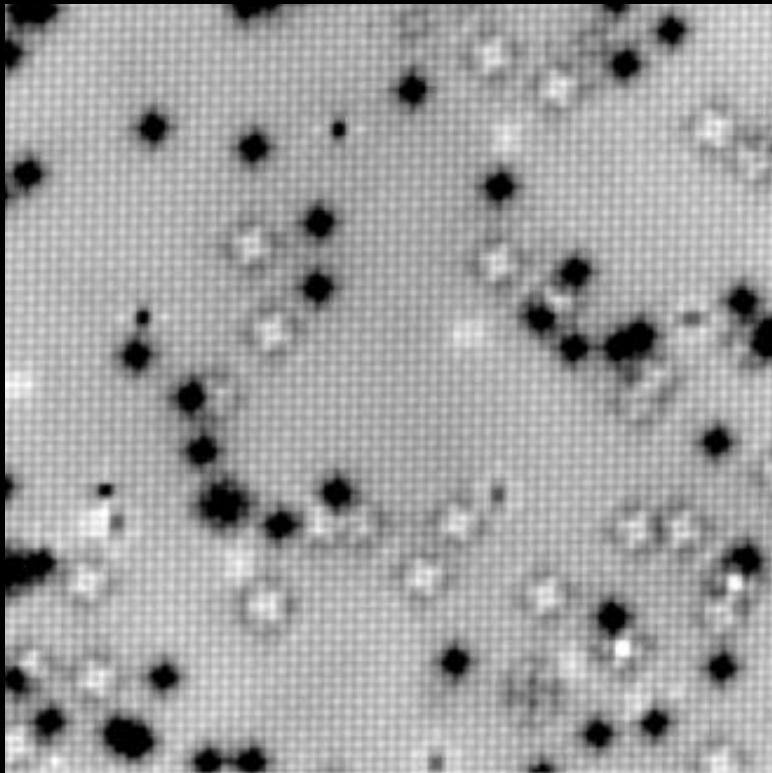
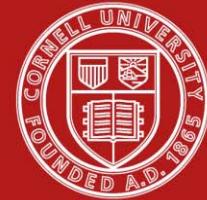


$g(q,V)$

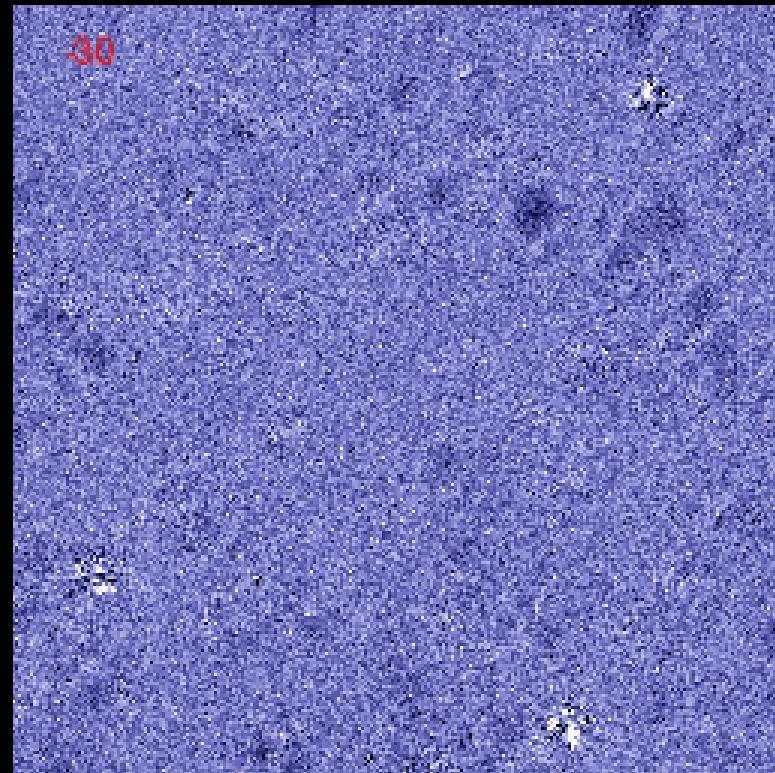
Science 327, 181 (2010)



Example: SrRuO₃-27



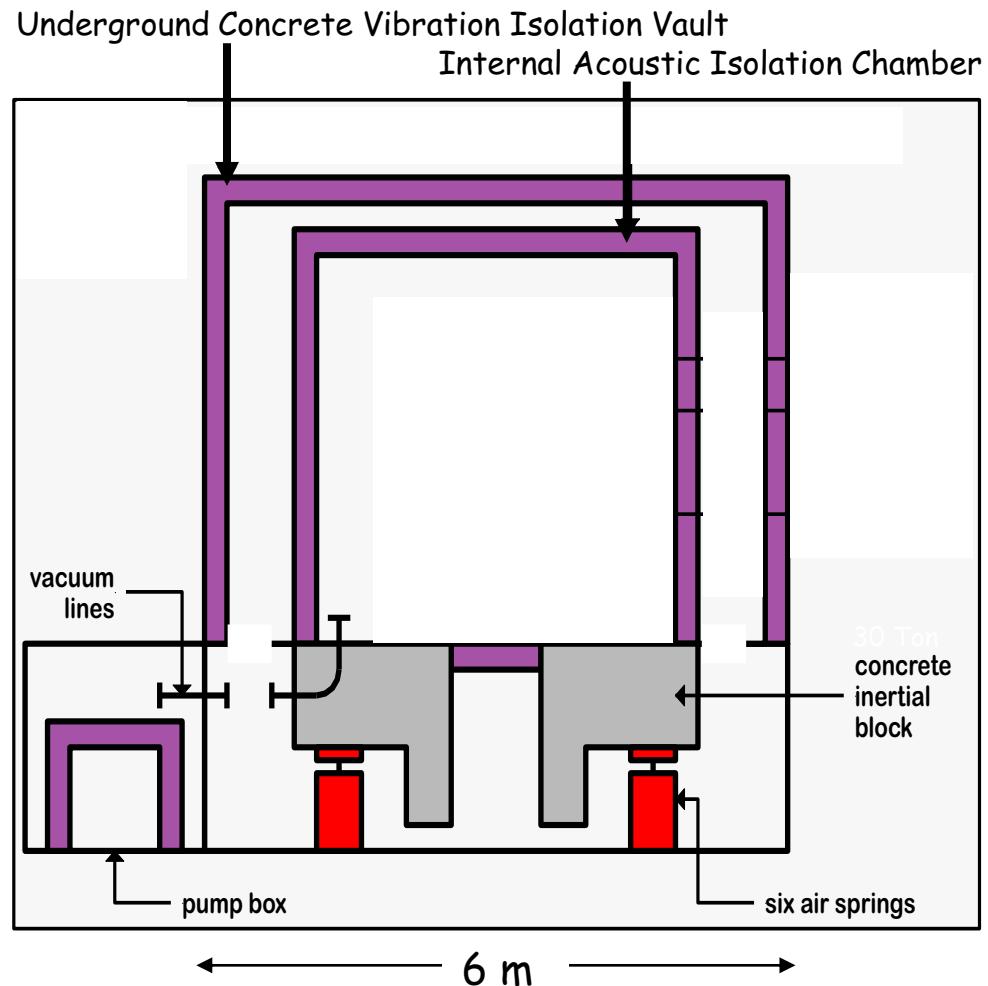
Topographic
Image



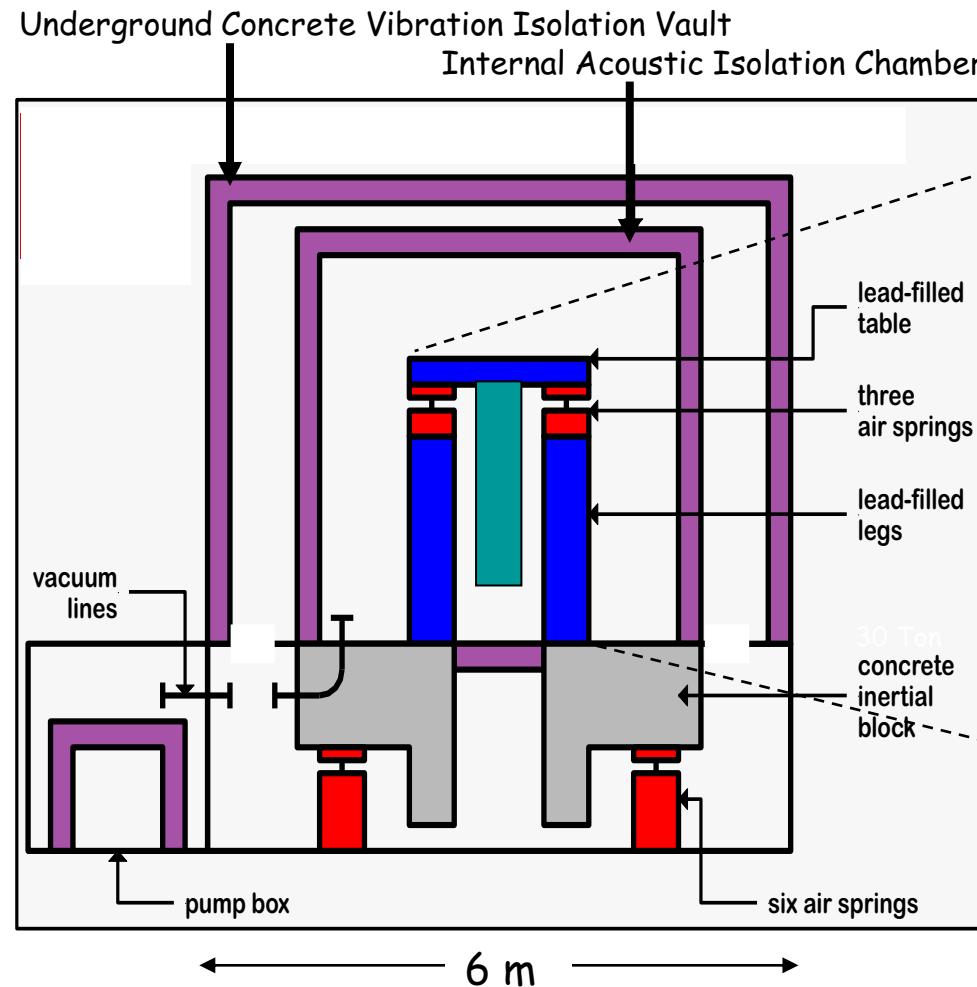
$g(r,V)$

Nature Physics 5, 800 (2009)

Ultra low vibration lab.

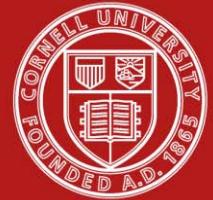


Rev. Sci. Inst. 70, 1459 (1999).

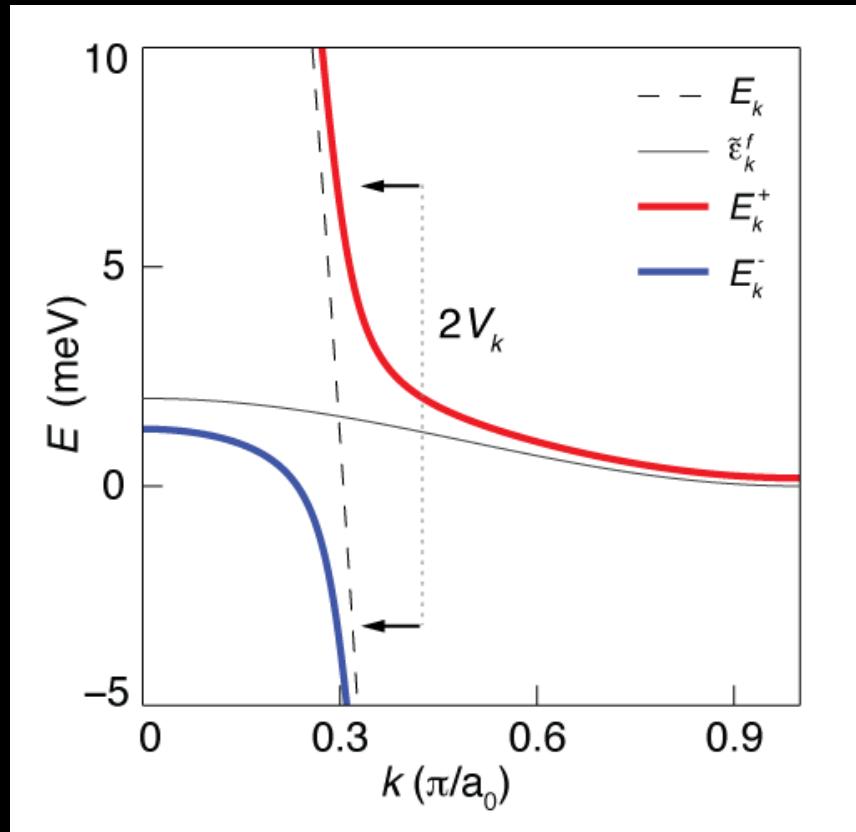


Rev. Sci. Inst. 70, 1459 (1999).

How/Why Heavy Fermion QPI?



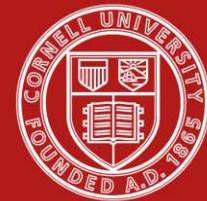
- Heavy Fermion many body state and bands are above E_F
- Heavy Fermion bands are extremely flat requiring $\sim 100 \mu\text{V}$ energy resolution or better
- → Sub-kelvin temperatures, low vib. & EM noise, plus high electronic sensitivity

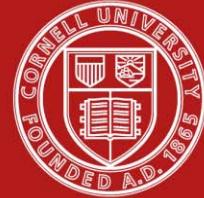


$$E_k^\pm = \frac{\tilde{\varepsilon}_k^f + E_k \pm \sqrt{(\tilde{\varepsilon}_k^f - E_k)^2 + 4|\tilde{V}_k|^2}}{2}$$



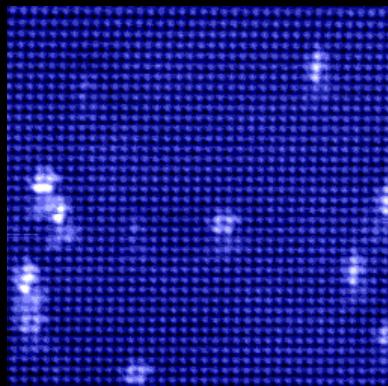
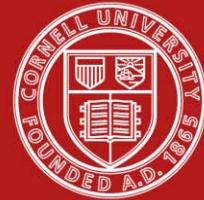
Dilution Fridge Spectroscopic Imaging-STM



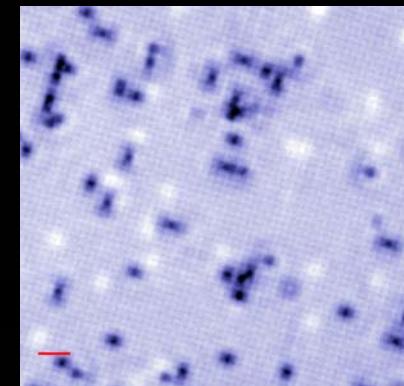
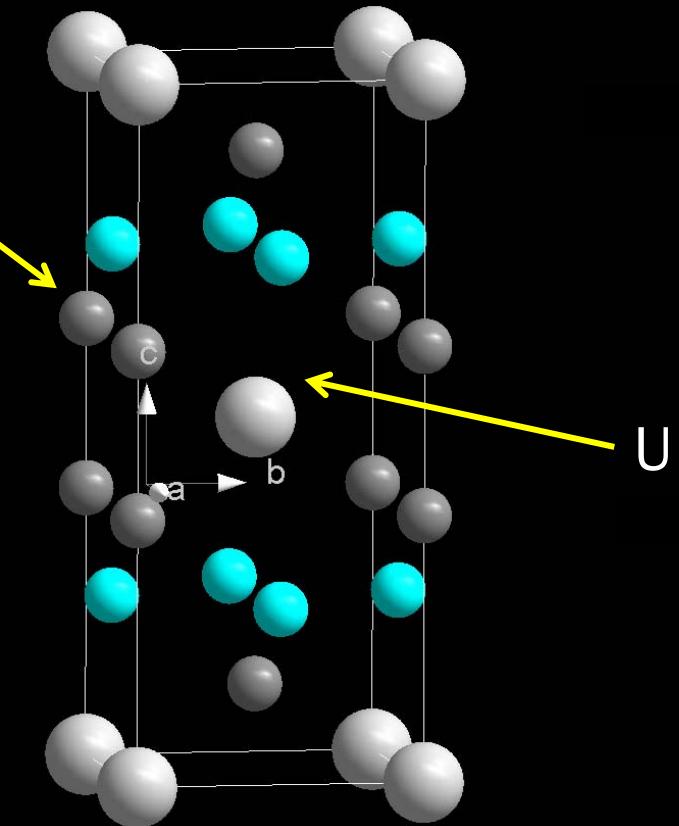


URu₂Si₂: Fano Lattice

Cleave Plane 1 - Si Surface



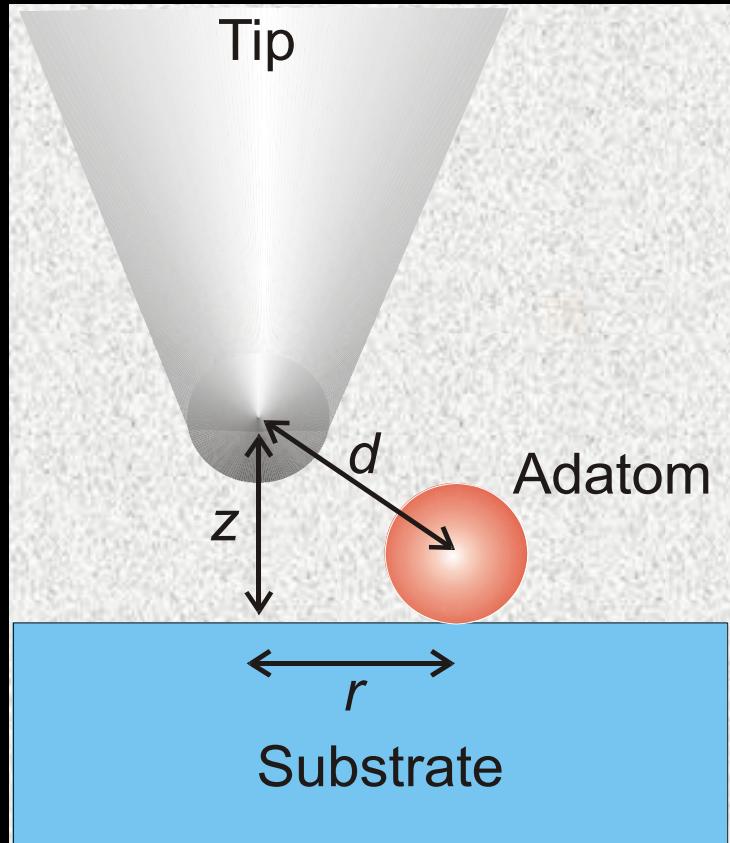
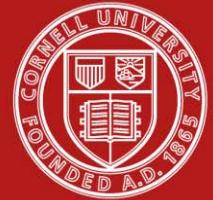
Si



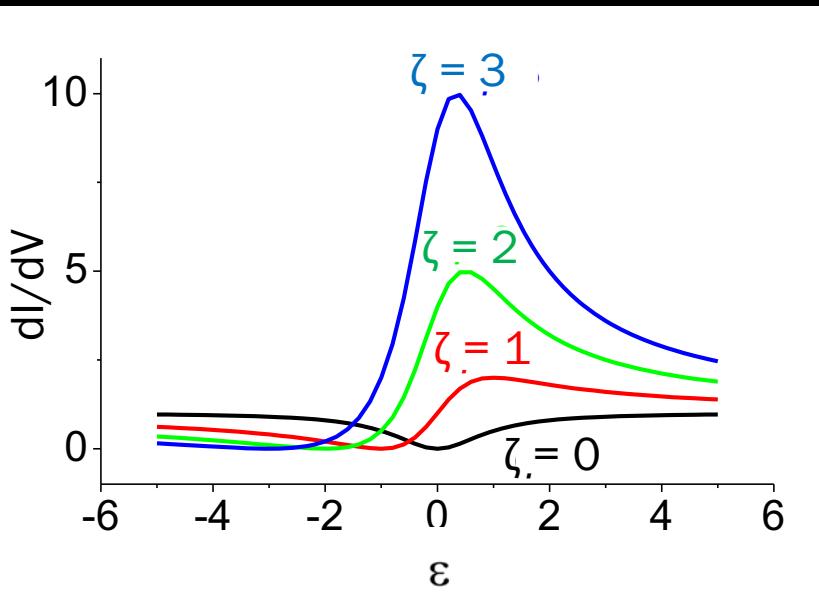
U
Ru
Si

$a=4.124\text{\AA}$; $c=9.582\text{\AA}$ (PRL65-3189)

Fano Lineshape



$$dI / dV(V) \propto \frac{\zeta + \varepsilon'^2}{\varepsilon'^2 + 1}, \quad \varepsilon' = \frac{(\varepsilon - \varepsilon_0)}{\Gamma/2}$$



ε_0 - energy of resonant state

Γ - width of resonance

ζ - coupling ratio

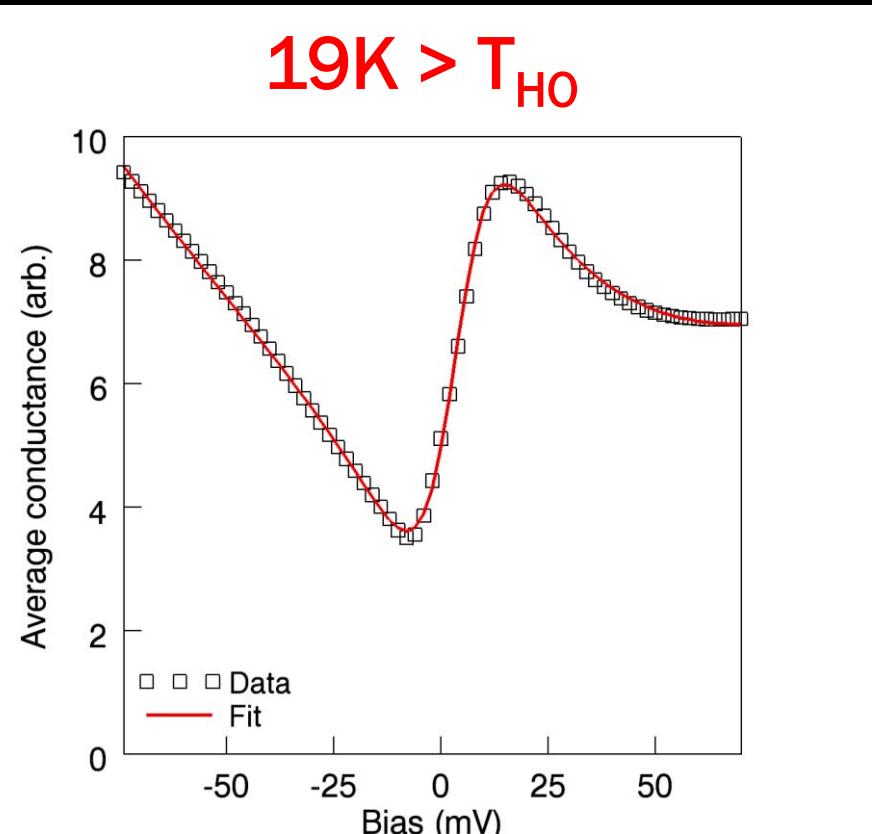
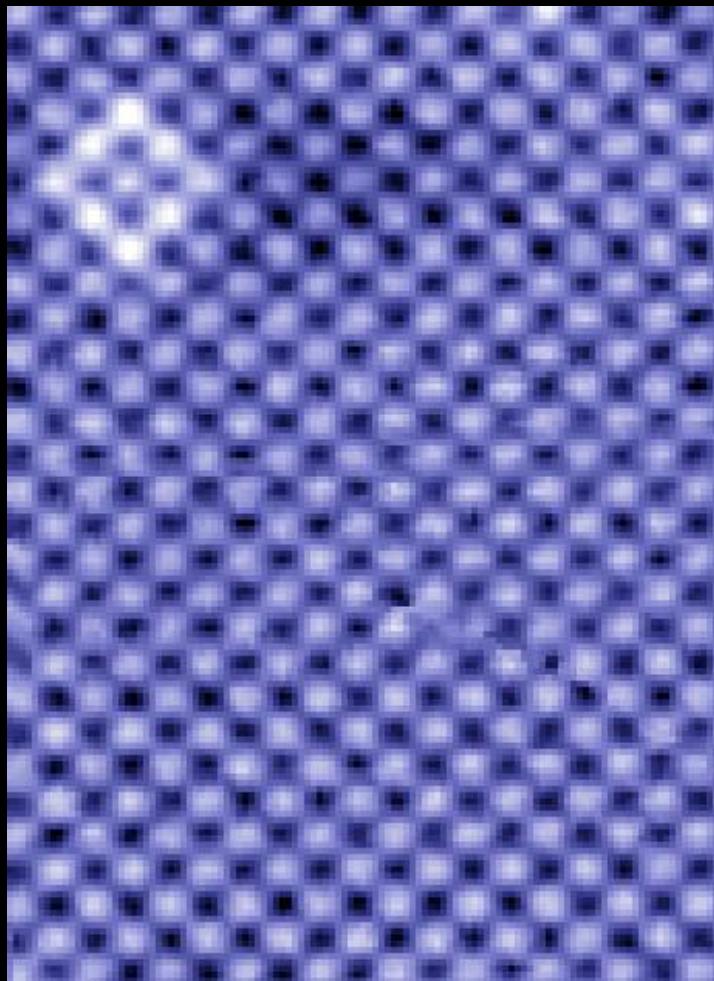
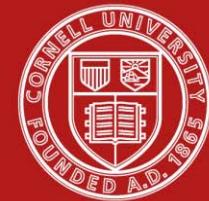
V. Madhavan *et al*, Science 280, 567 (1998)

M. Plihal und J.W. Gadzuk, Phys. Rev. B 63, 085404 (2001)

O. Újsághy, J. Kroha, L. Szunyogh und A. Zawadowski, Phys. Rev. Lett. 85, 2557 (2000)

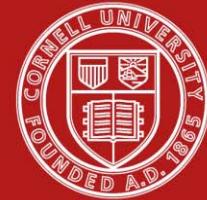


dI/dV Spectroscopy

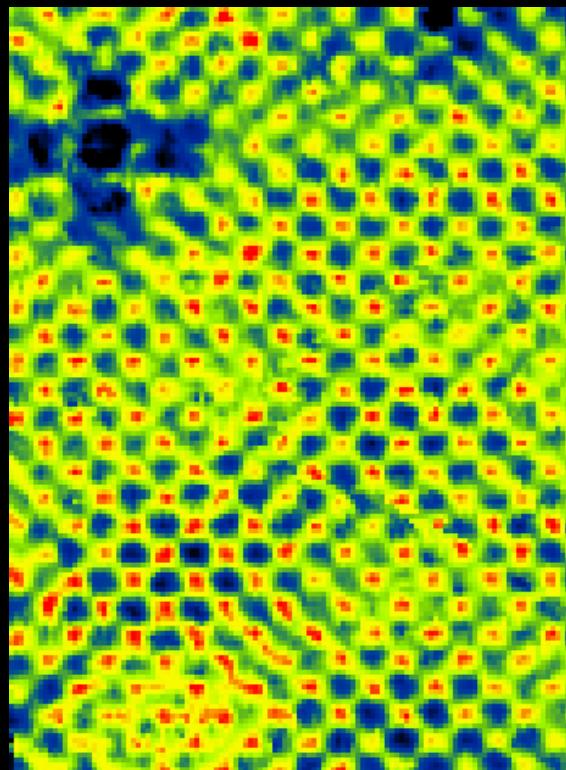




Imaging Fano Parameters

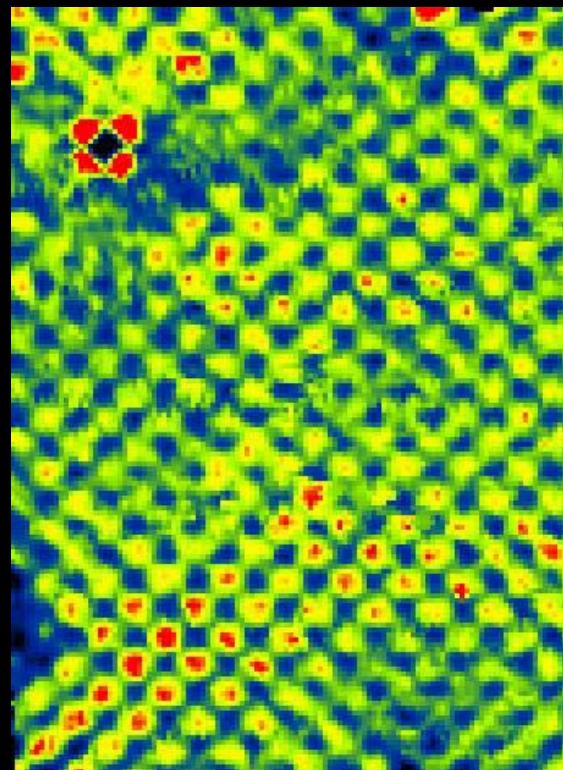


$\varepsilon_0(\mathbf{r})$



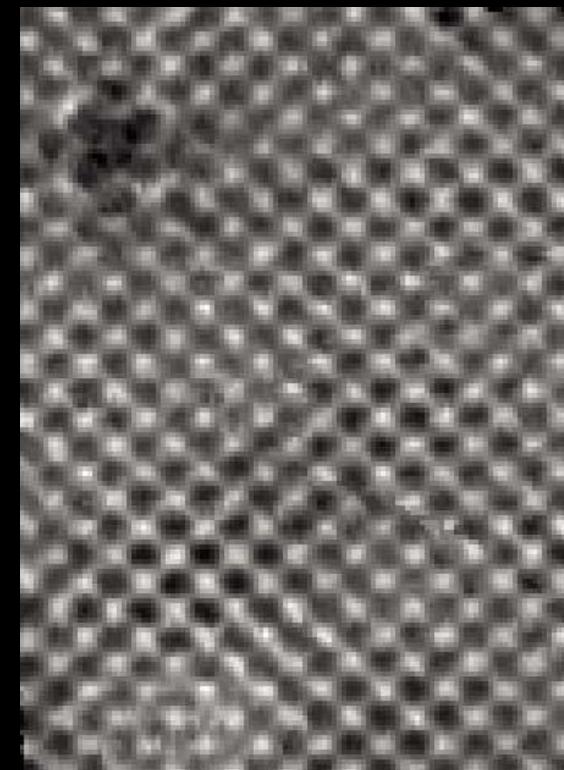
0.5 8 meV

$\Gamma(\mathbf{r})$



10 15.5 meV

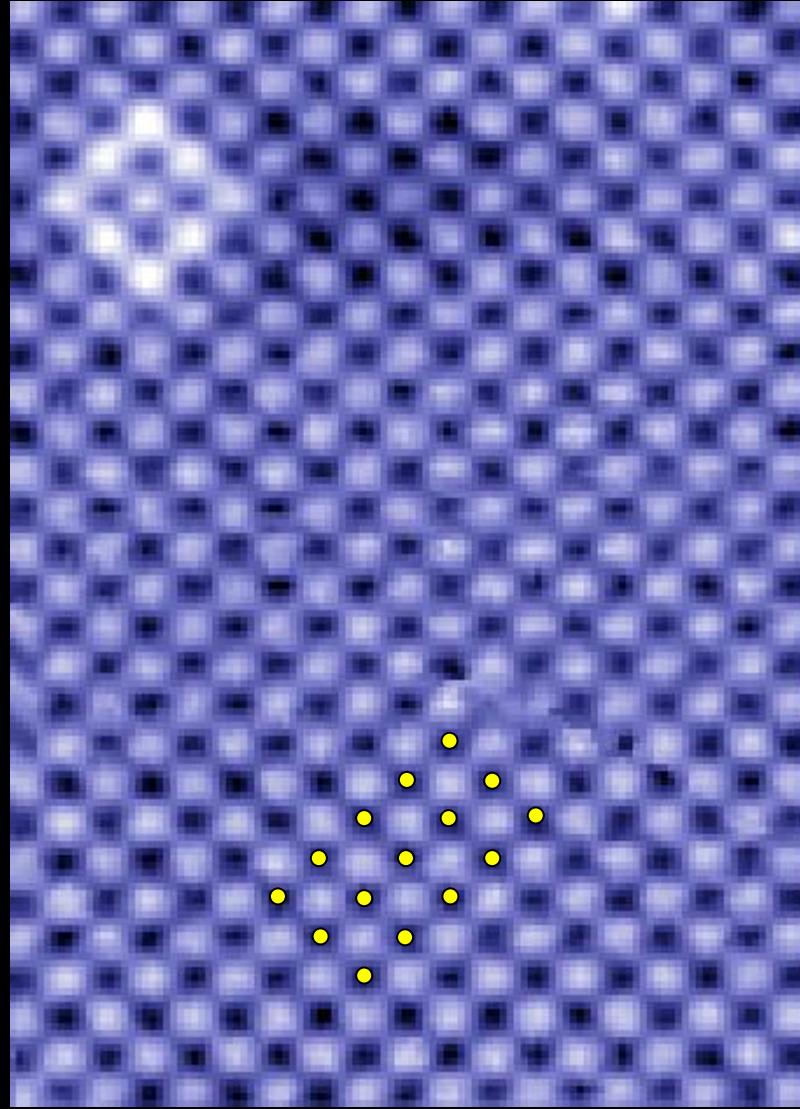
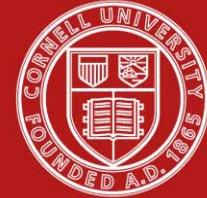
(\mathbf{r})



0.6 1.8



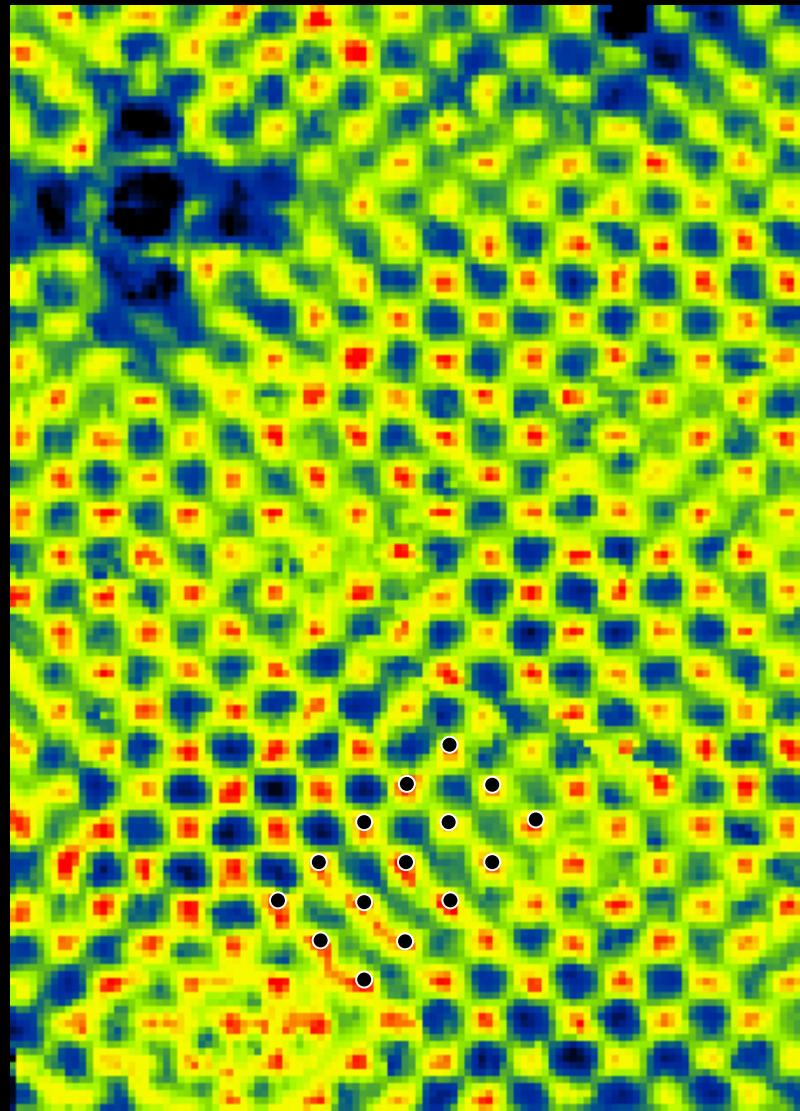
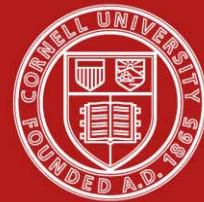
Topography



19K > T_{HO}



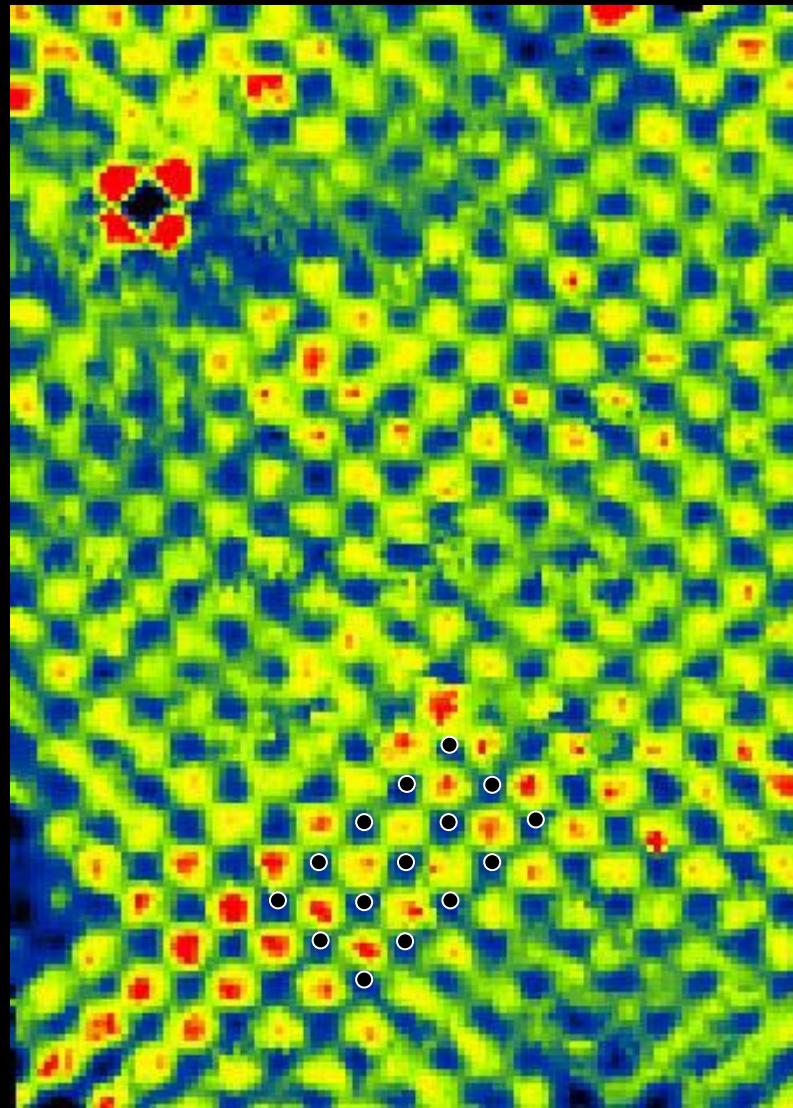
$\varepsilon_0(\mathbf{r})$



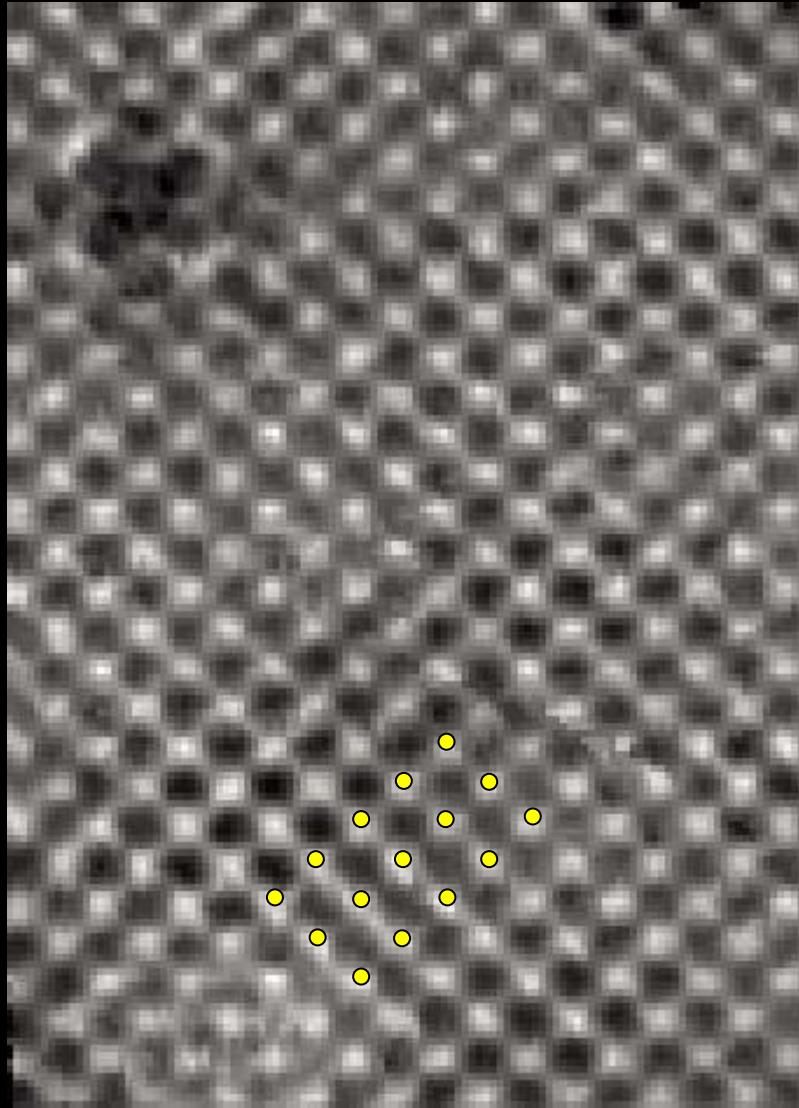
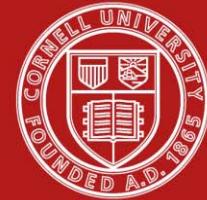
19K > T_{HO}

0.5 8 meV

$\Gamma(\mathbf{r})$



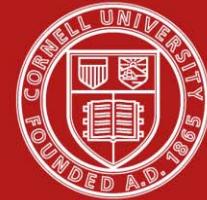
19K > T_{HO}

 $\zeta(r)$ 

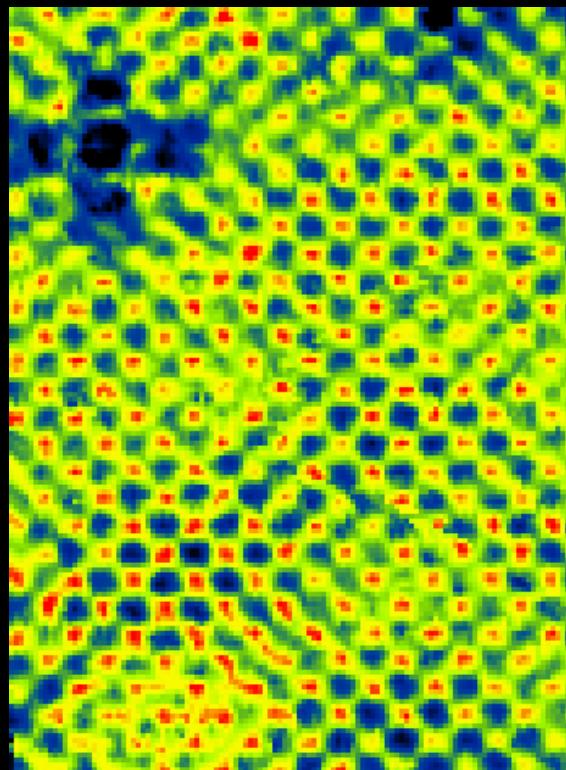
19K > T_{HO}



Visualization of a Fano Lattice

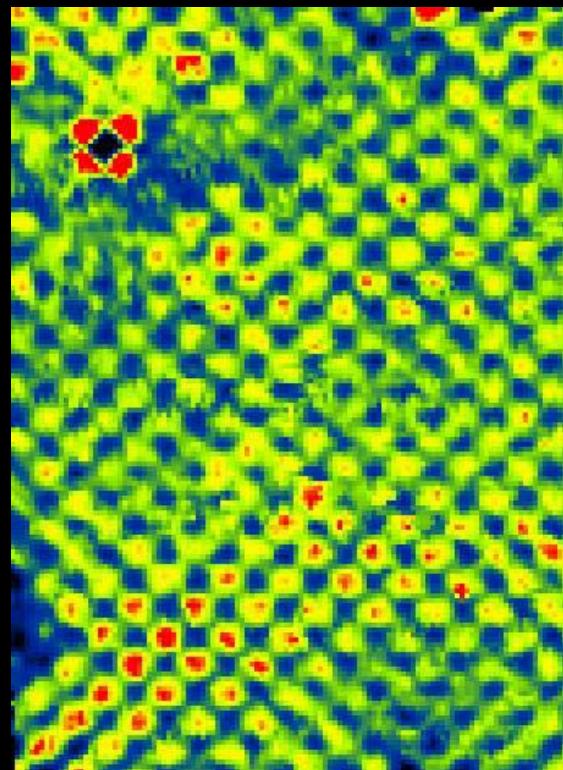


$\varepsilon_0(\mathbf{r})$



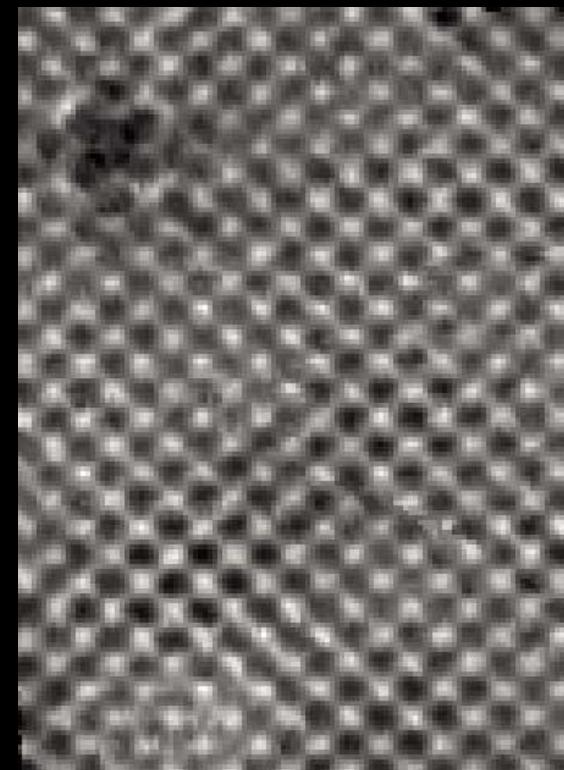
0.5 8 meV

$\Gamma(\mathbf{r})$



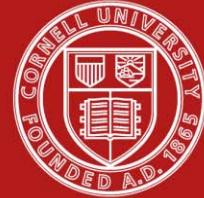
10 15.5 meV

$\zeta(\mathbf{r})$



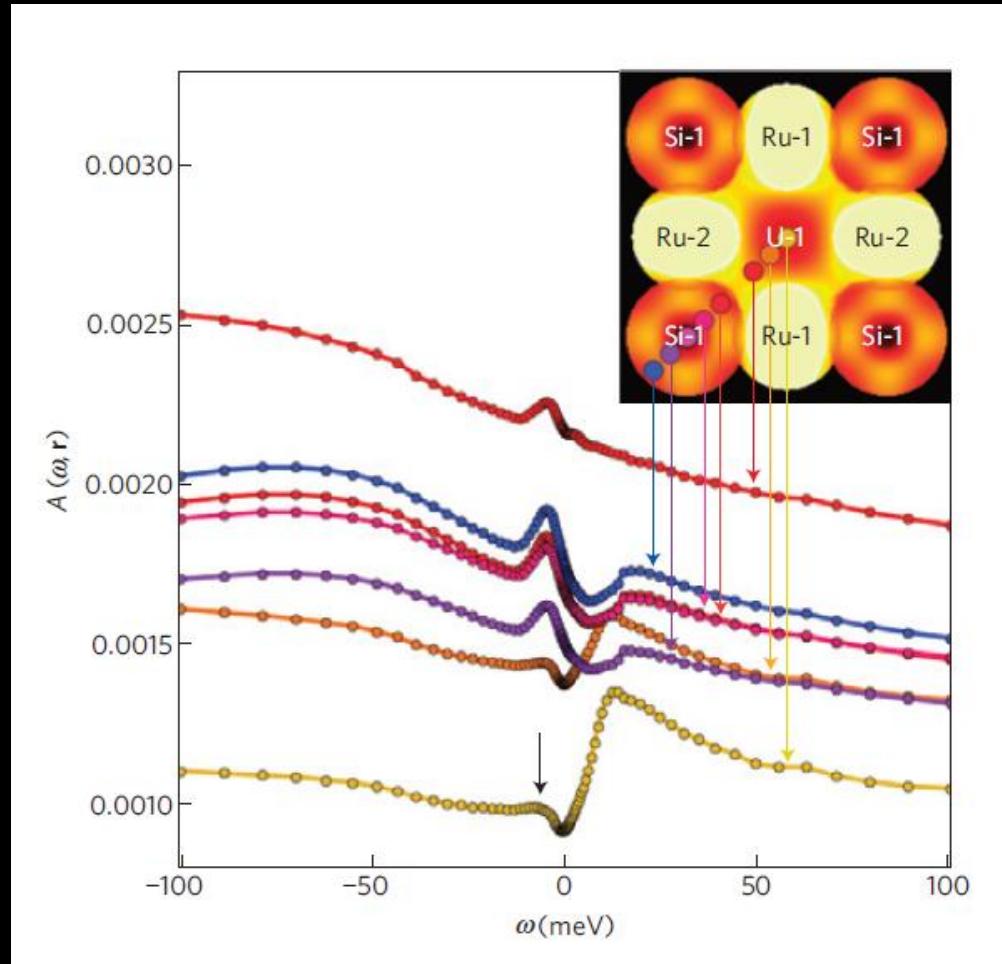
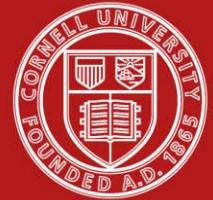
0.6 1.8

A. Schmidt & M. Hamidian *et al.* **Nature** **465**, 570 (2010)
P. Aynajian *et al.* **PNAS** **107**, 10383 (2010)

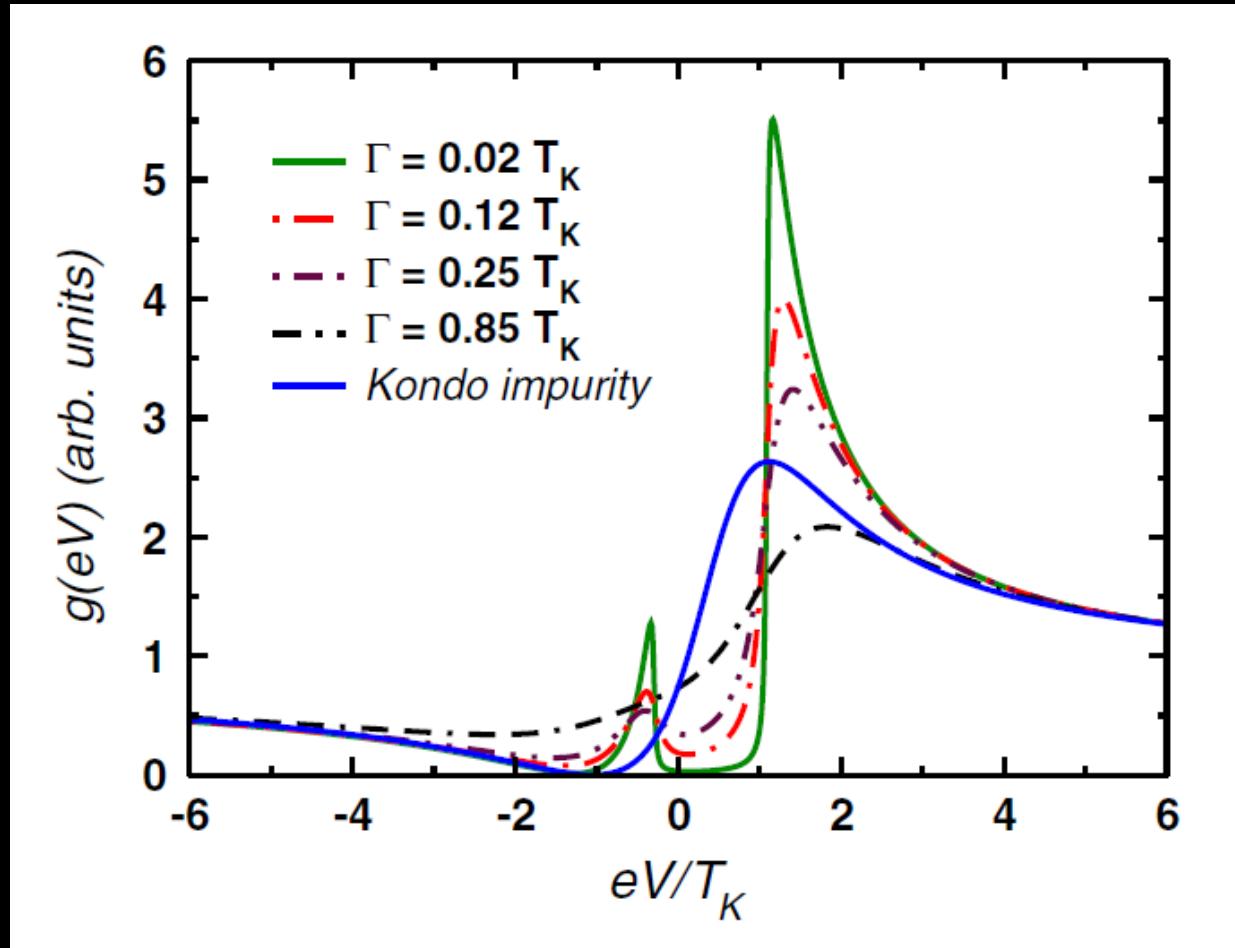


URu₂Si₂: DOS

Theory (DMFT)

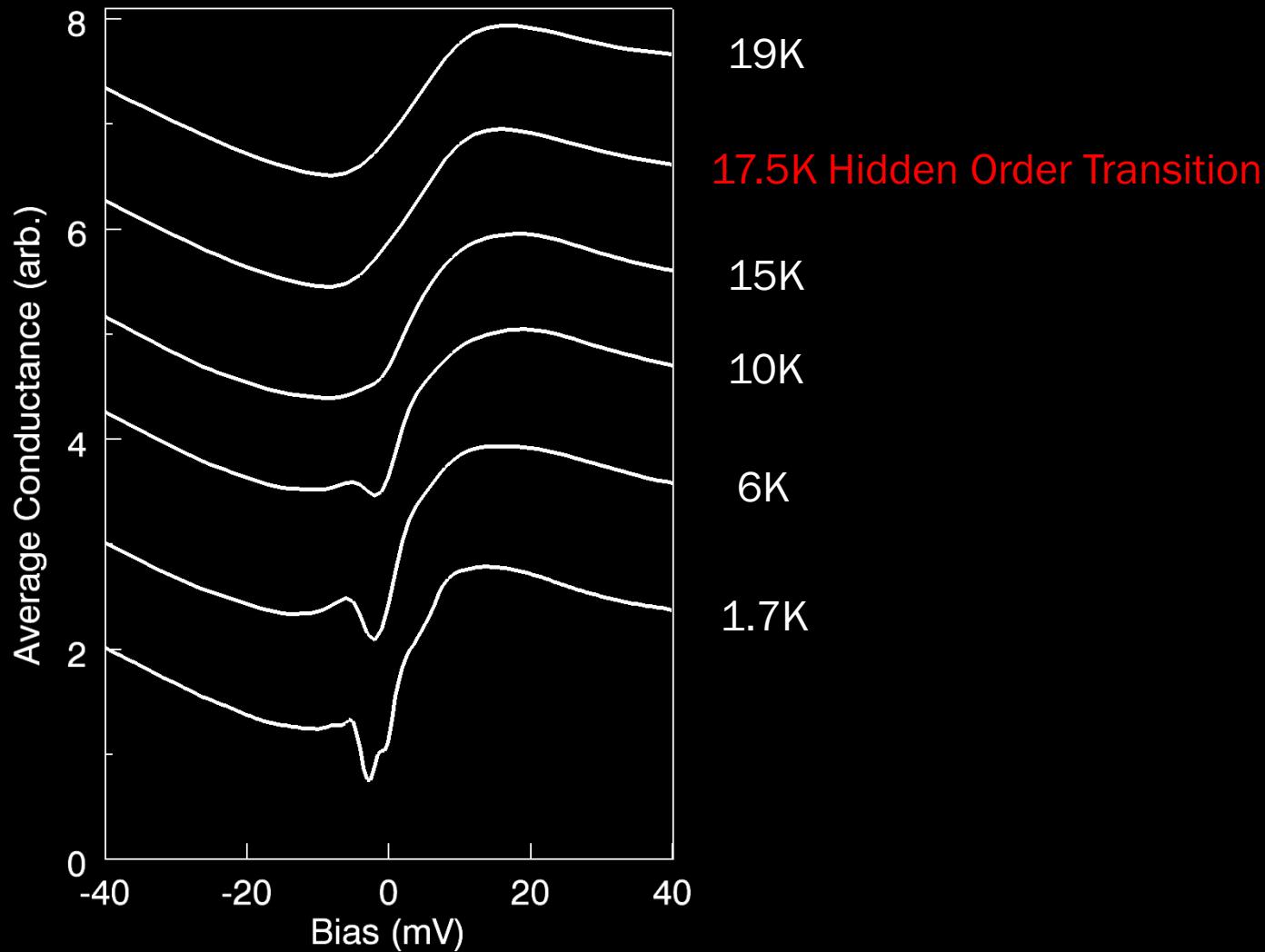
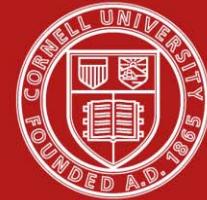


Haule, K. & Kotliar, G. *Nature Phys.* 6, 769 (2009)



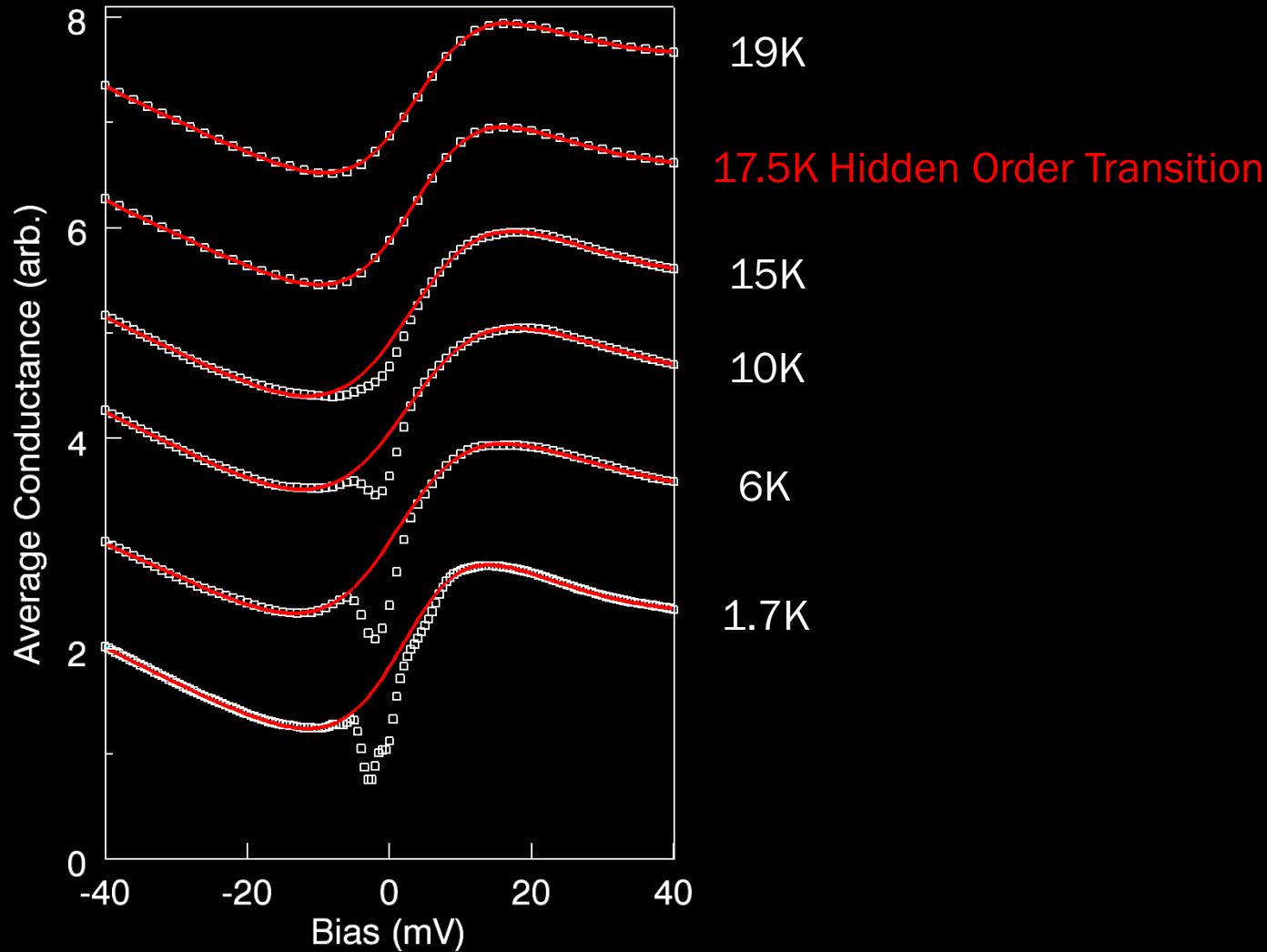
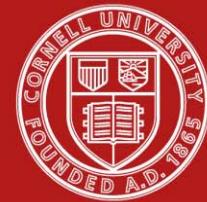


Temp. dependence of Fano Spectra



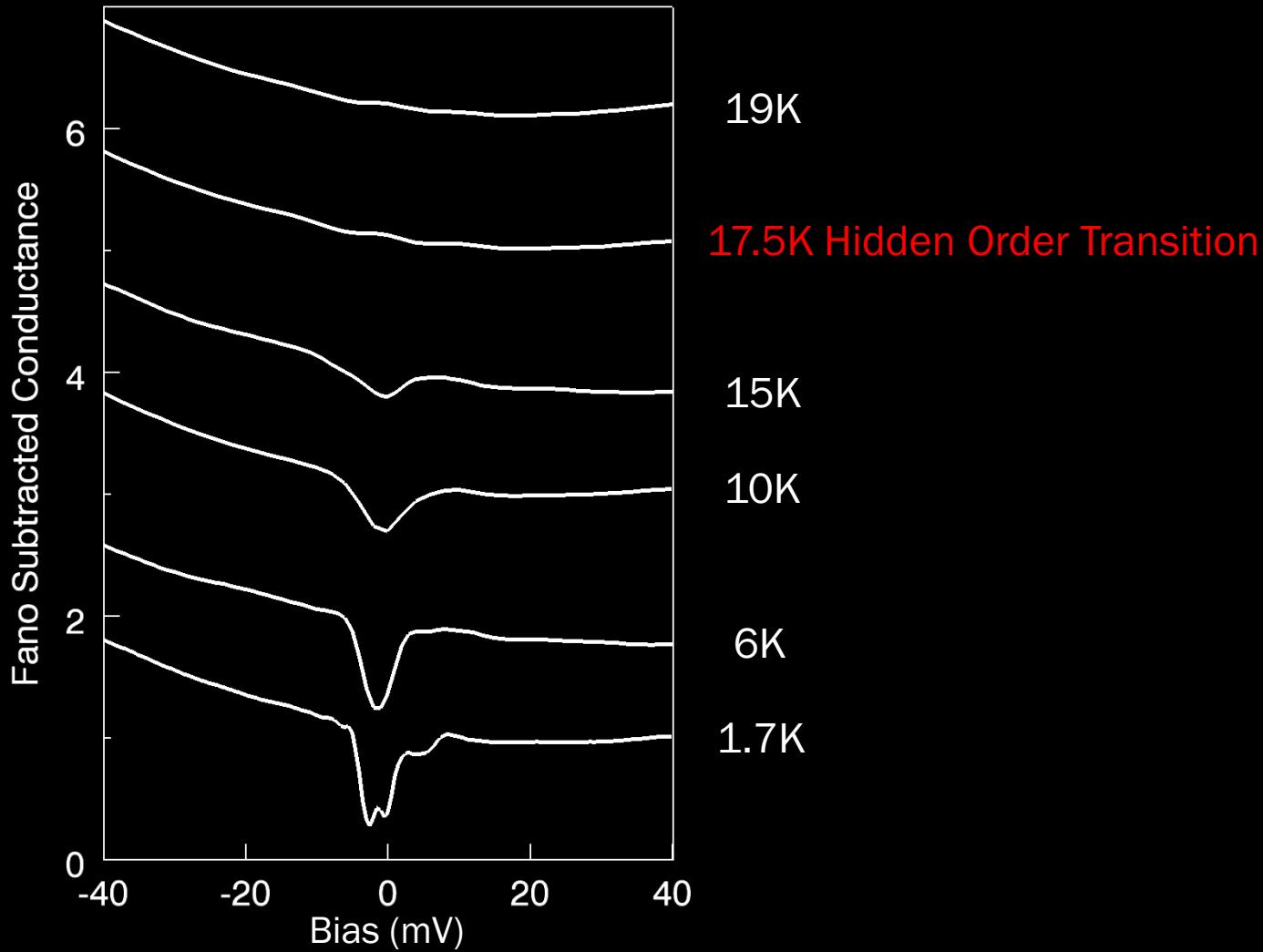
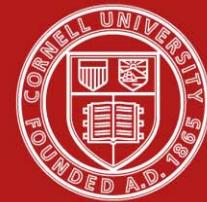


Temp. dependence of Fano Spectra

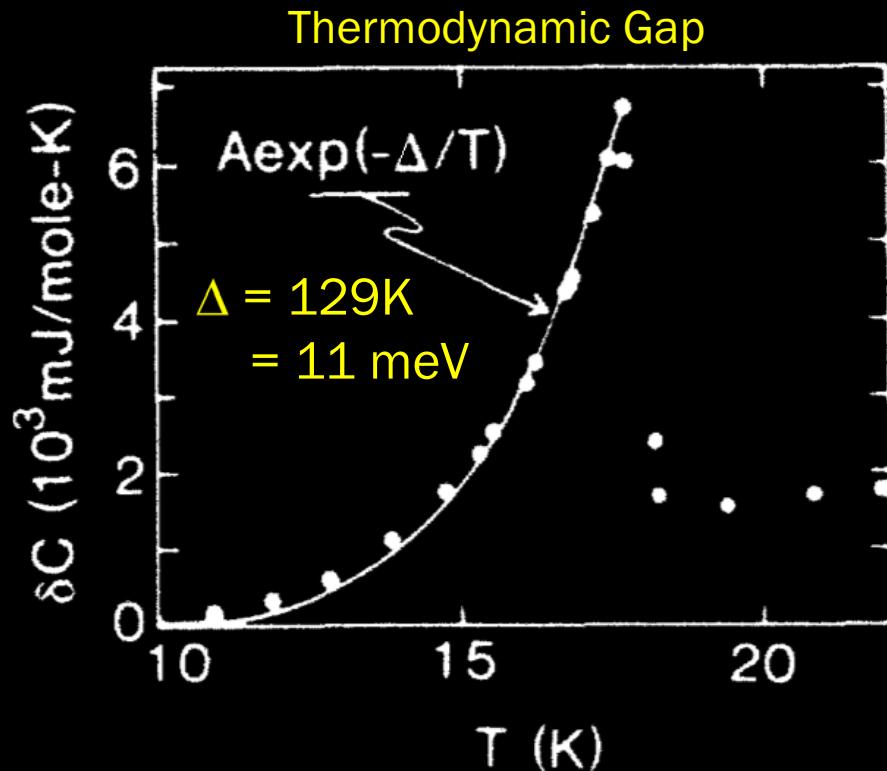
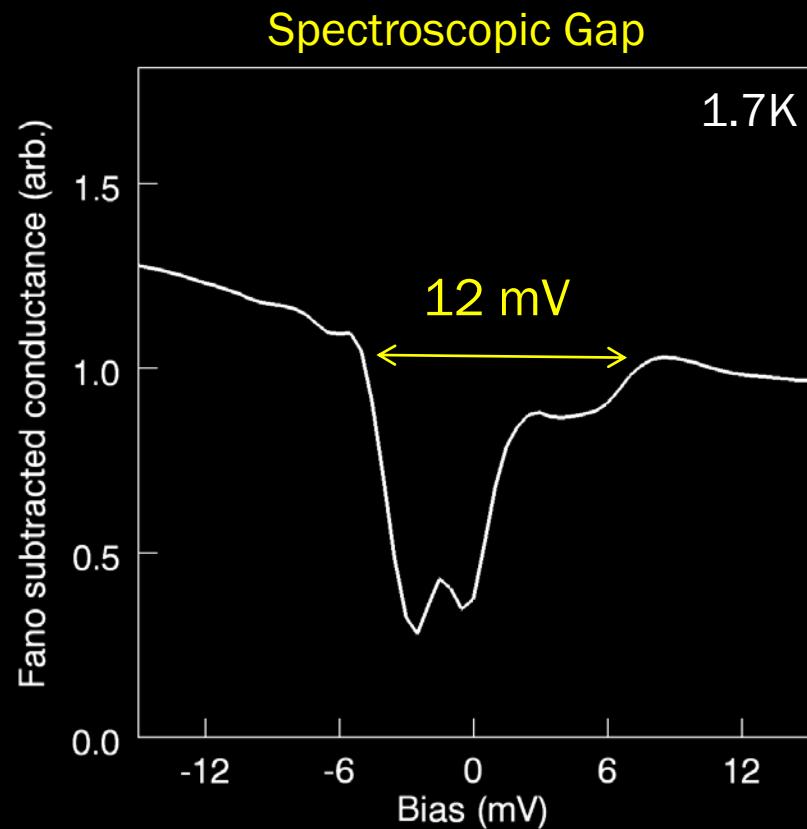
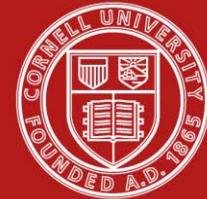




Temp. dependence of Fano Spectra



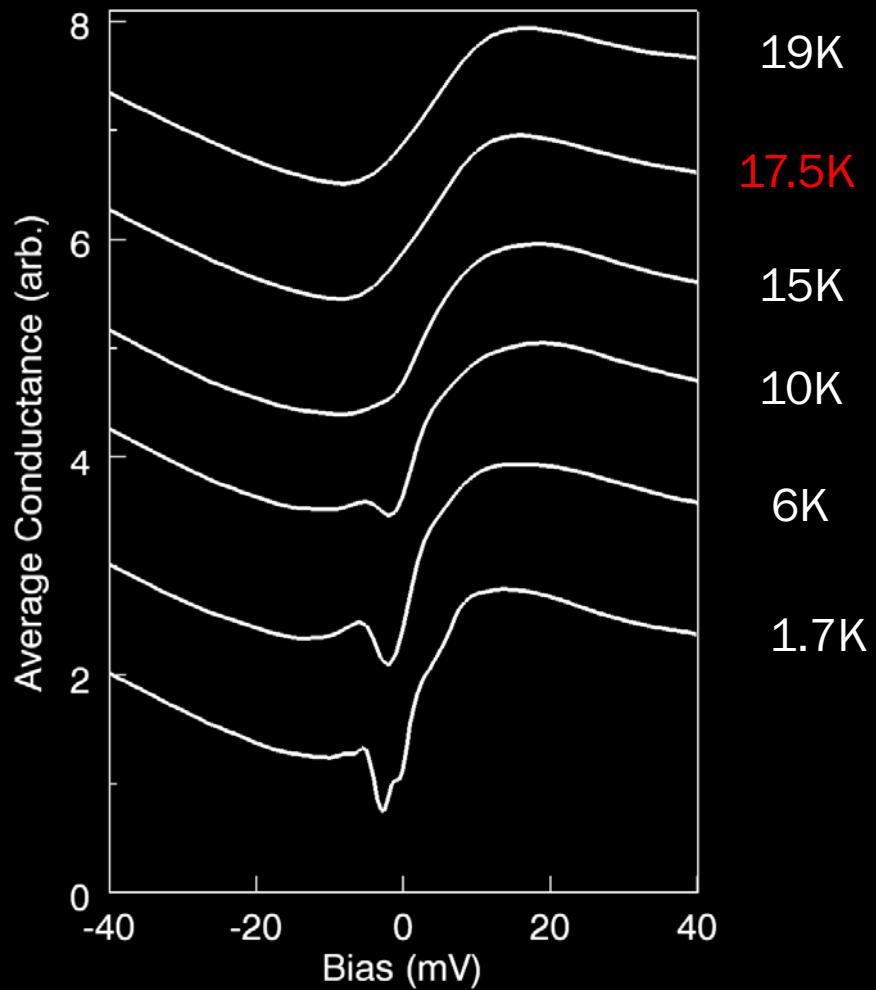
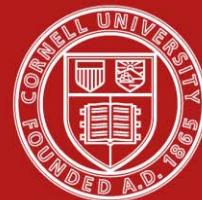
Consistent Spectroscopic Gap



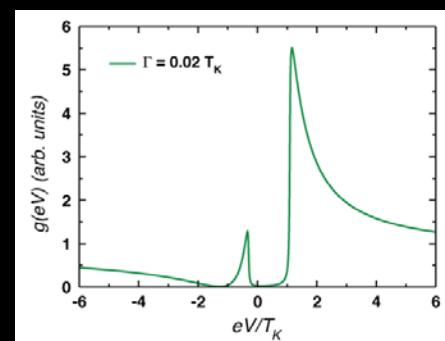
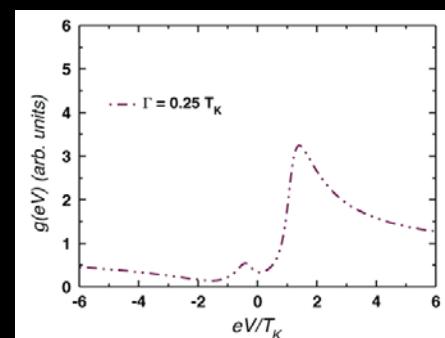
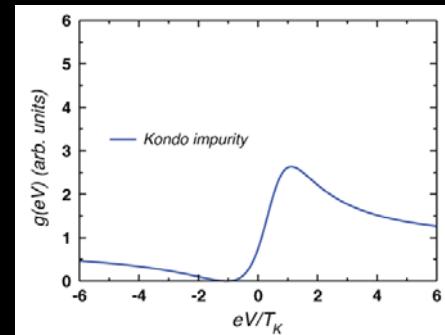
Maple, B. *et. al.* *PRL* **56**, 185 (1986)

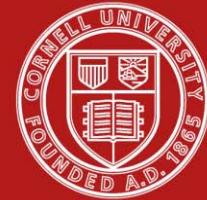
Bonn, D.A. *et al.* *PRL* **61**, 1305 (1988).

Comparison: Kondo Lattice Theory



Maltseva M. et al *PRL* **103**, 206402 (2009)

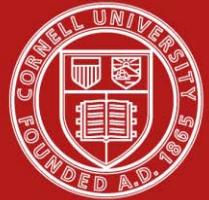




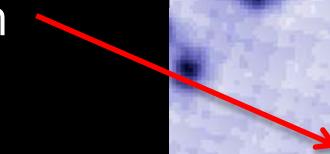
Thorium doped URu_2Si_2 : **for** **Quasiparticle Interference Imaging**



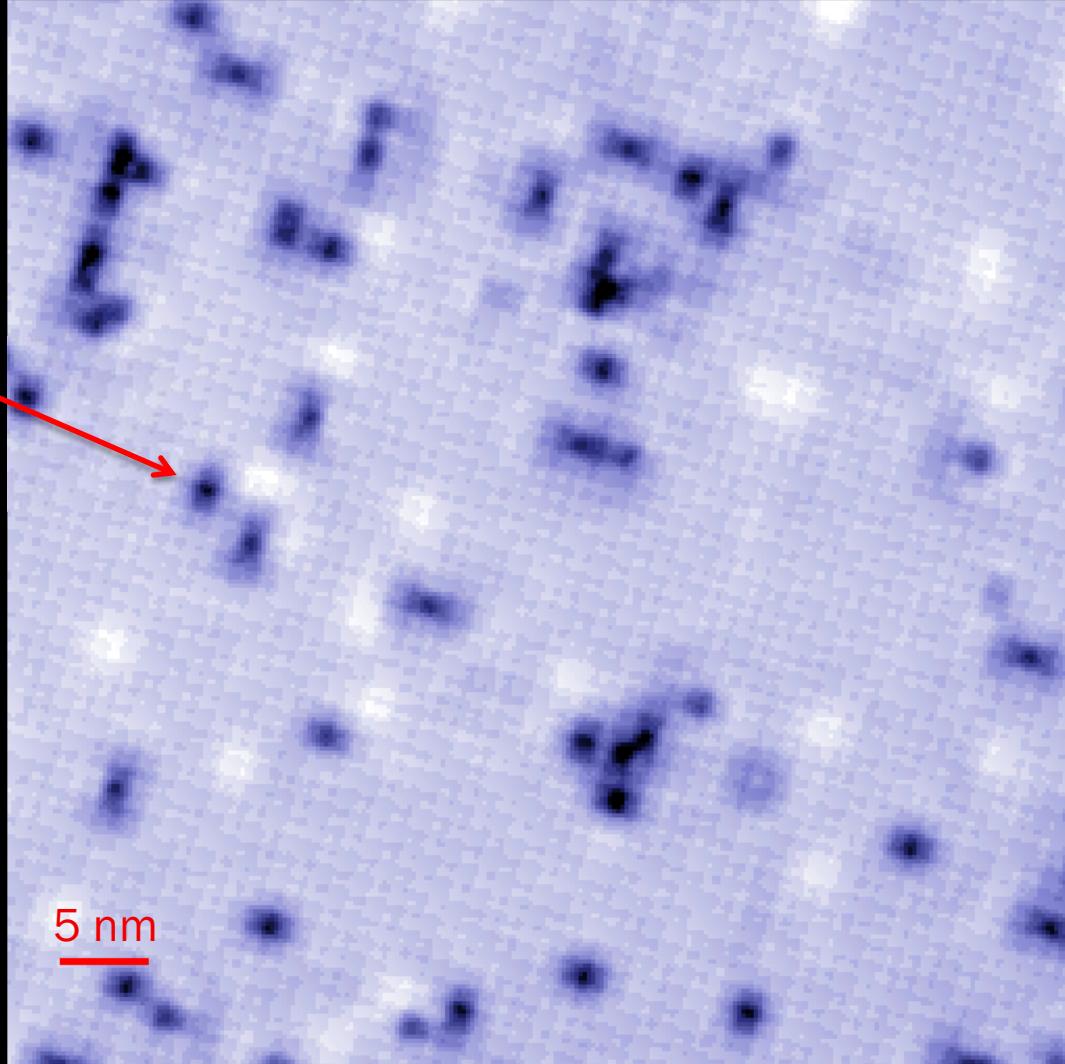
$U_{0.99}Th_{0.01}Ru_2Si_2$: U surface



Th



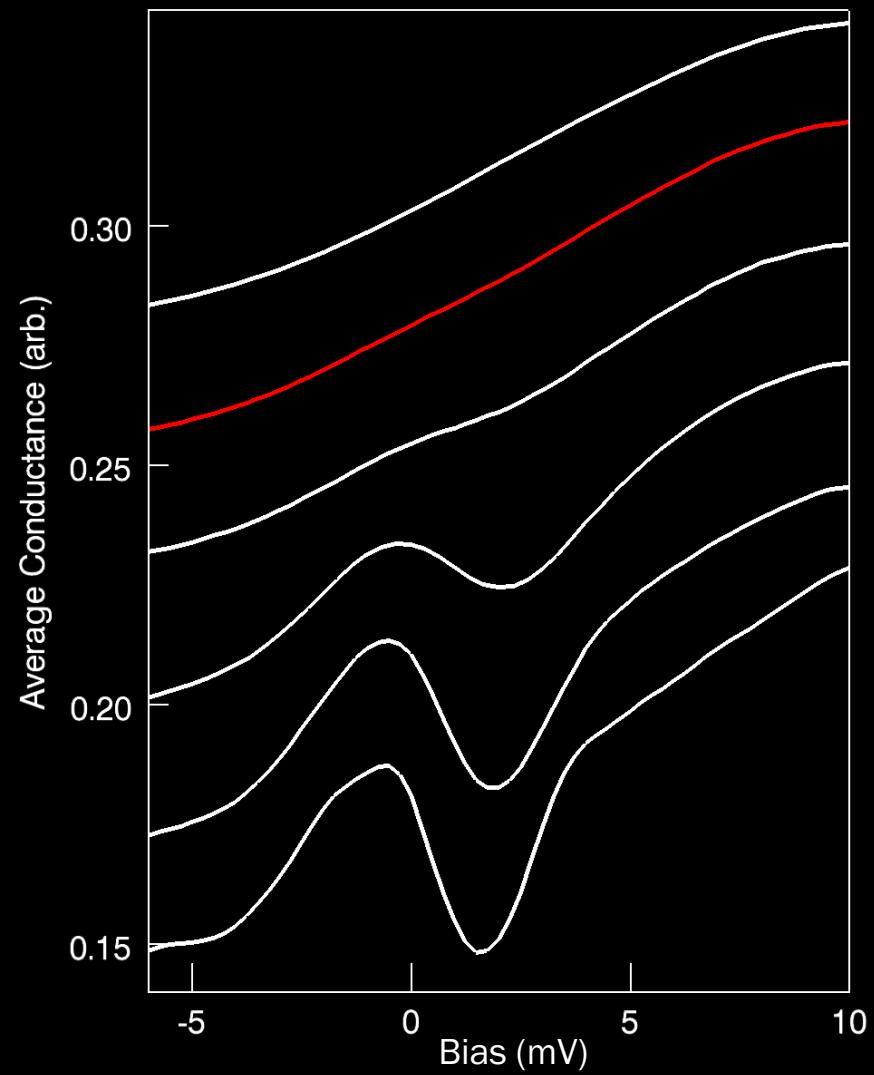
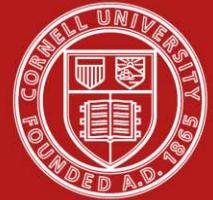
5 nm



Topography



$U_{0.99}Th_{0.01}Ru_2Si_2$: U surface



19K

16K Hidden Order Transition (drops with Th)

15K

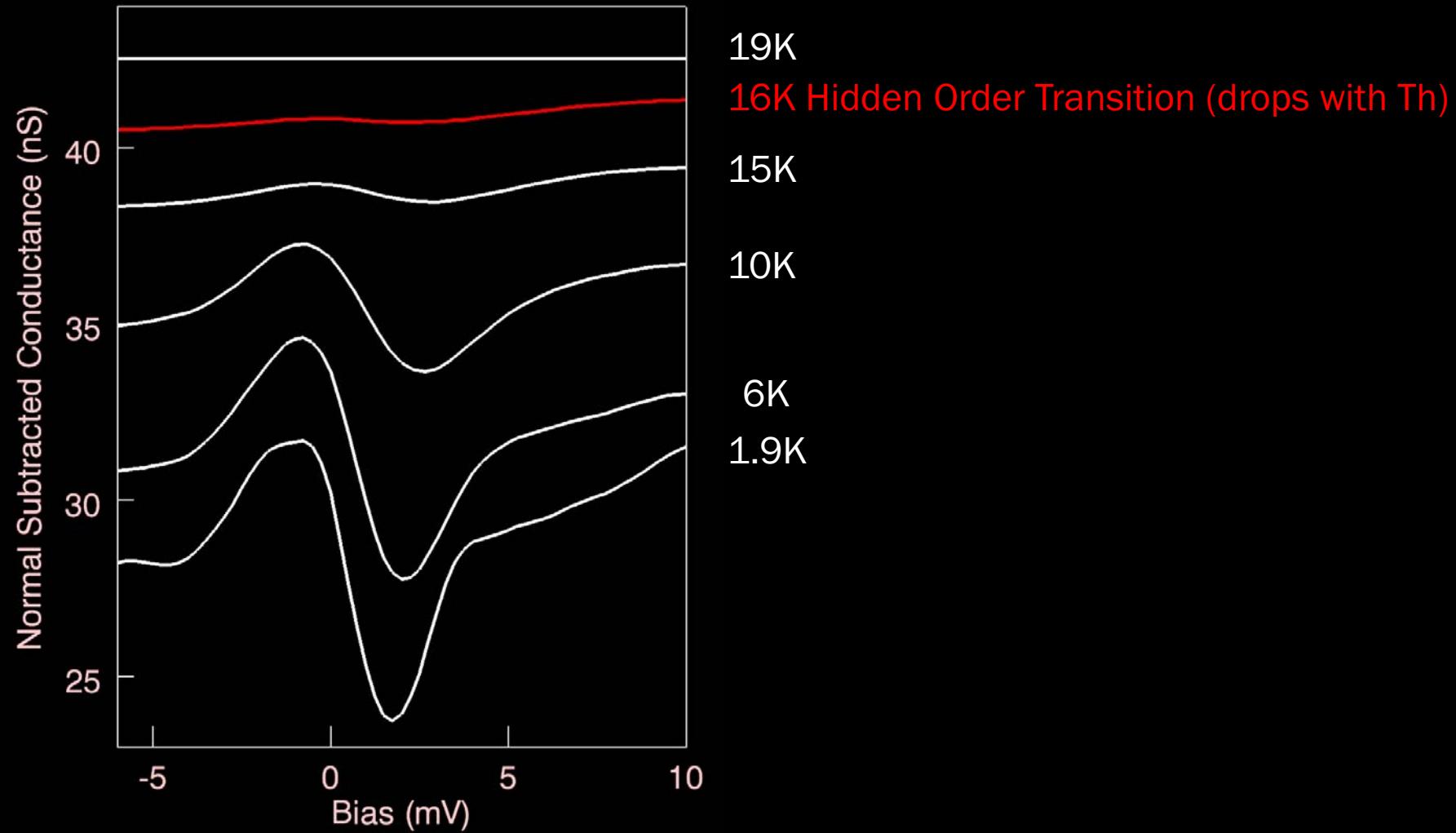
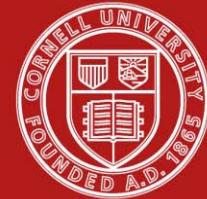
10K

6K

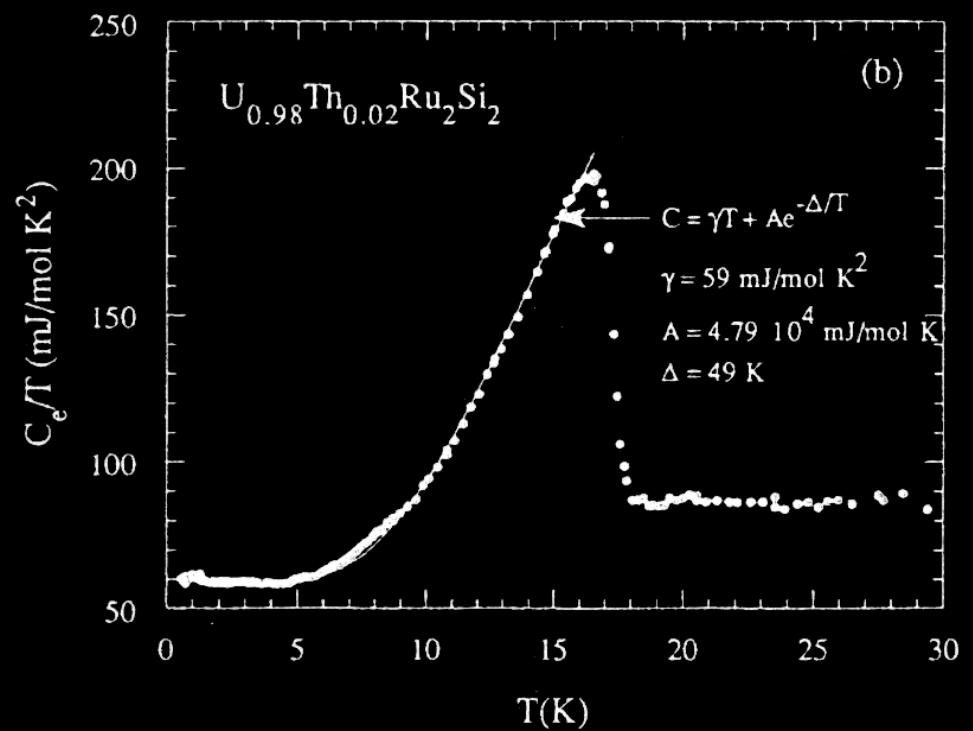
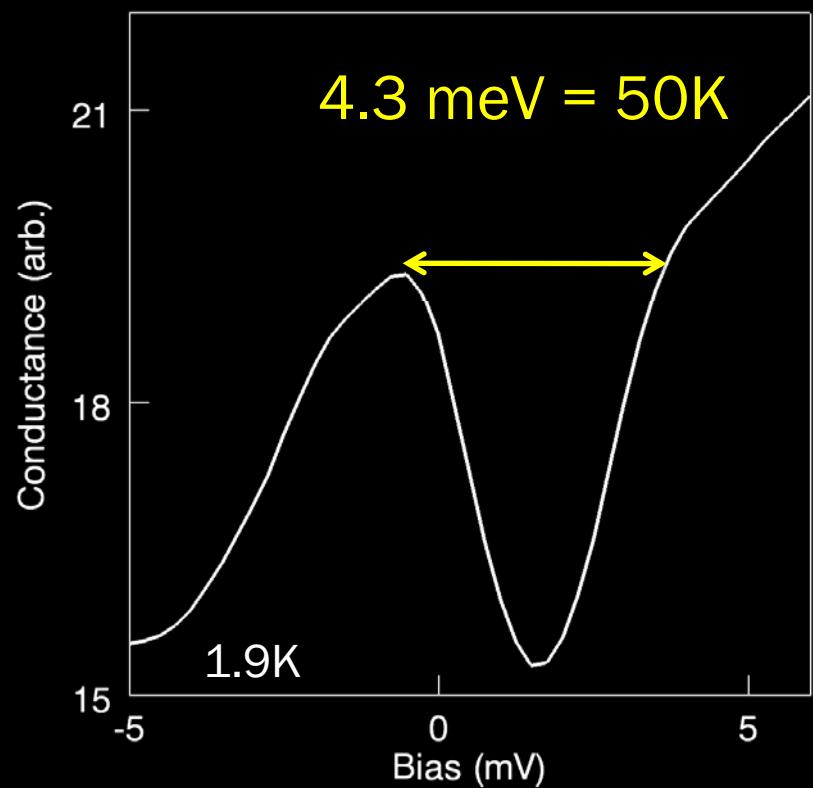
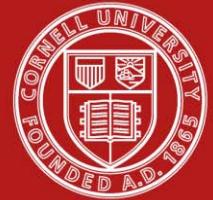
1.9K

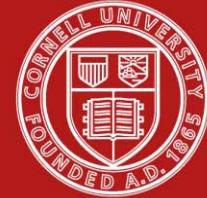


Transition within 1K of bulk value



Comparison with Specific Heat Gap





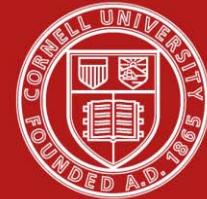
URu₂Si₂:

Temperature Dependent

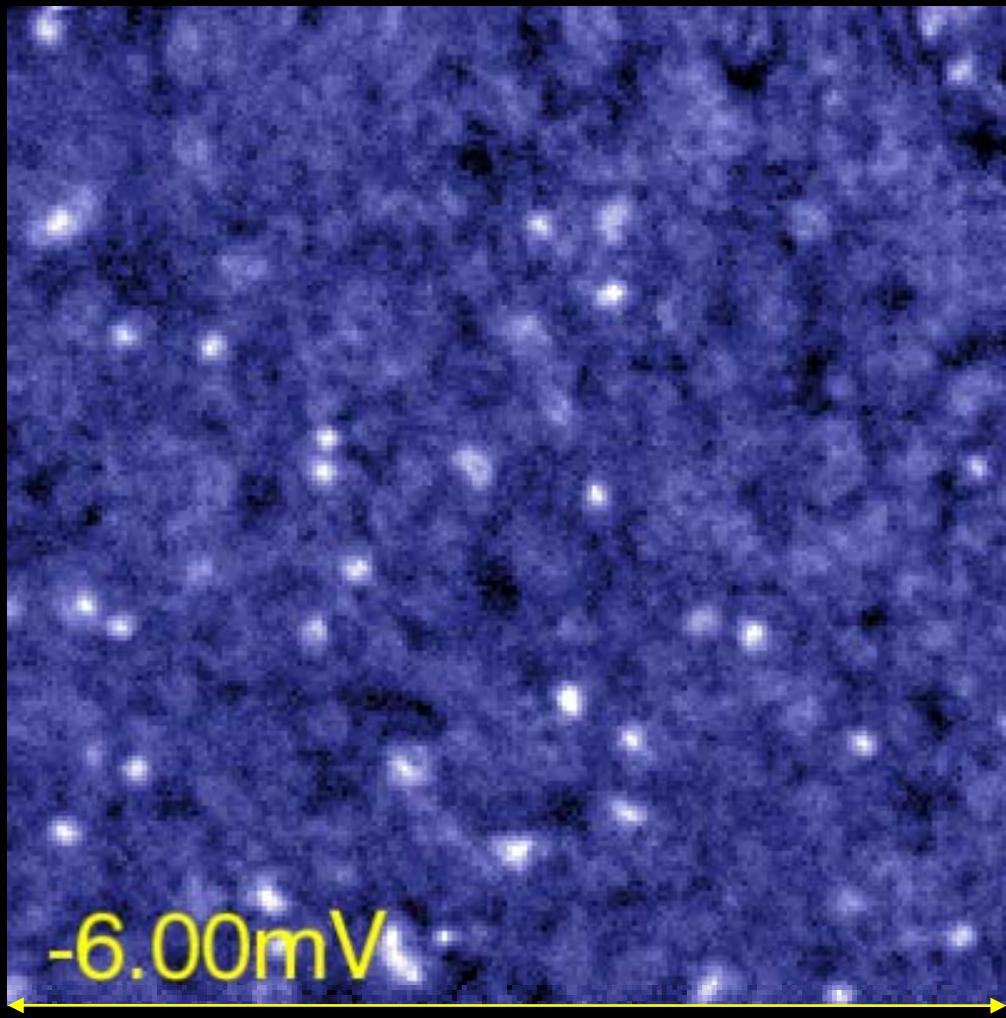
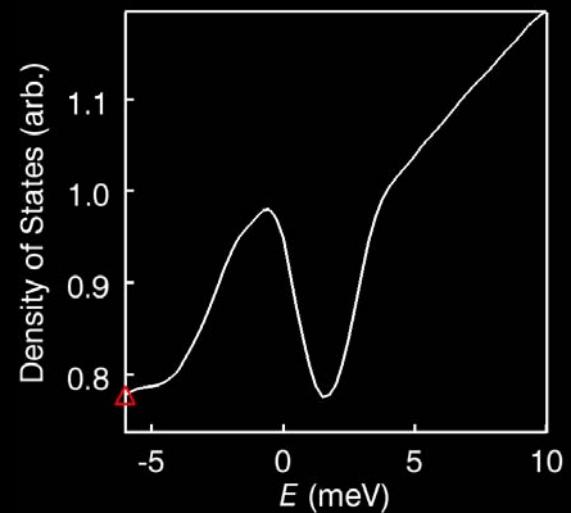
Quasiparticle Interference Imaging



1.9K Conductance Map $g(r,E)$

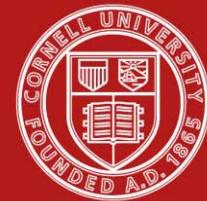


$U_{0.99}Th_{0.01}Ru_2Si_2$

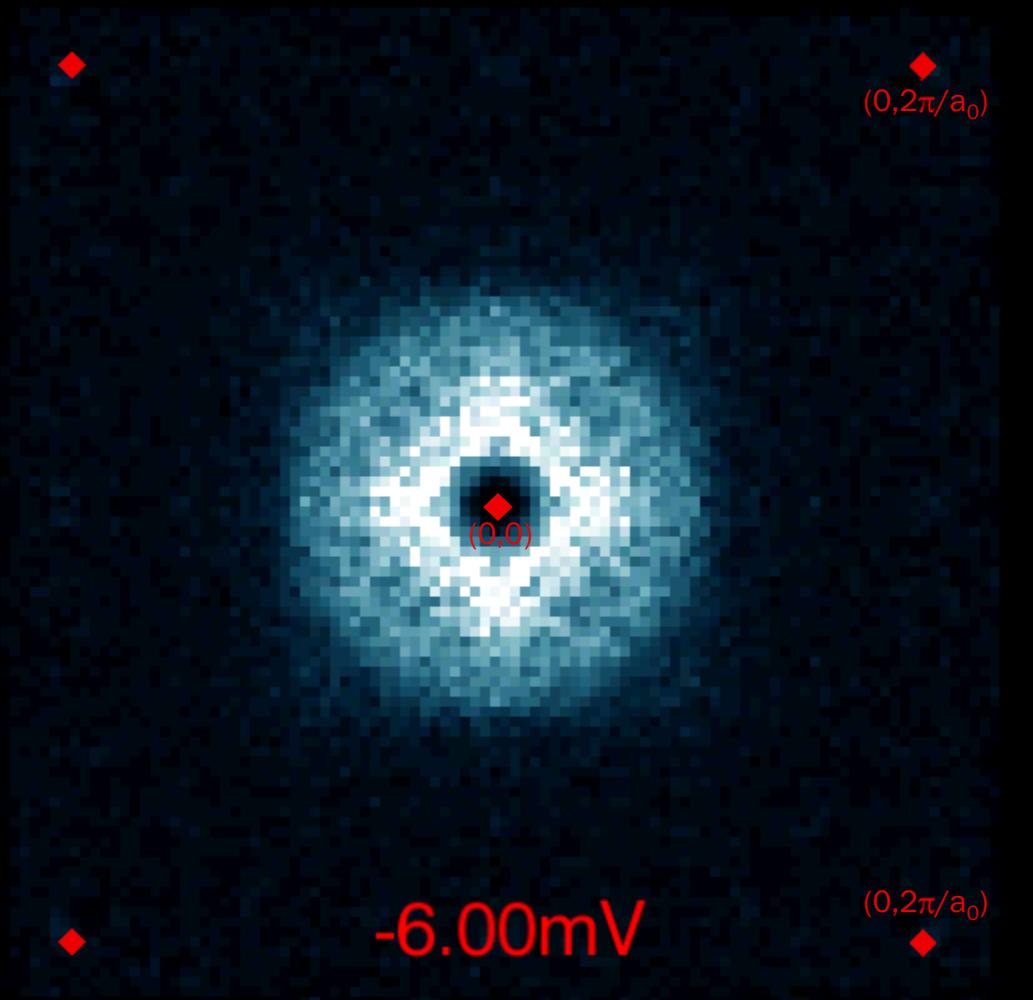
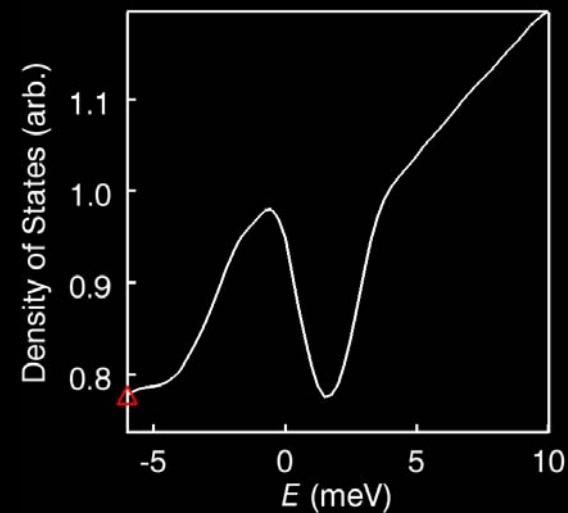




1.9K QPI Image $g(\mathbf{q}, E)$

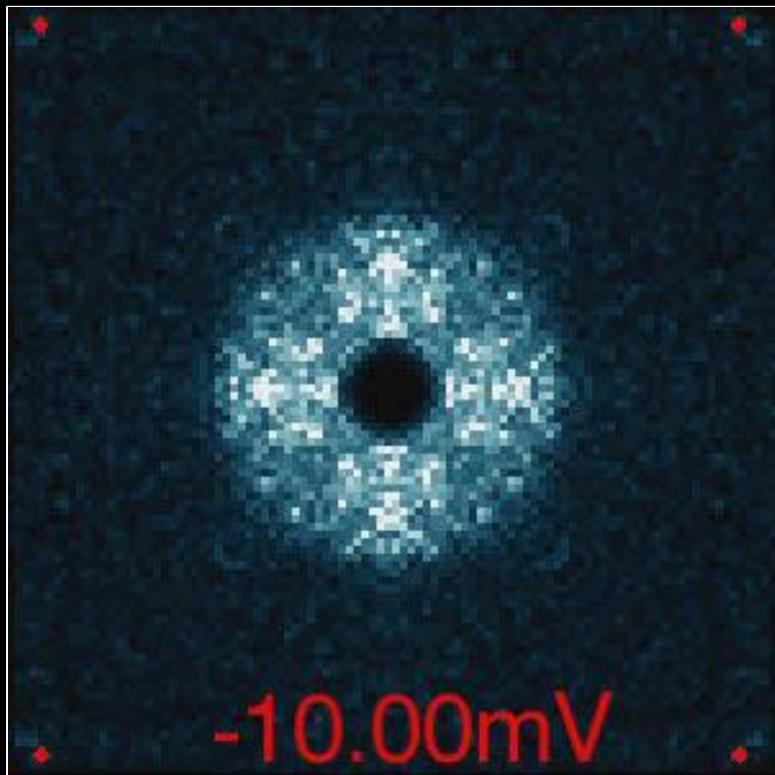
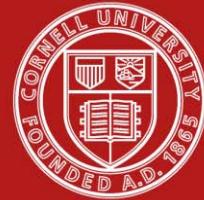


$\text{U}_{0.99}\text{Th}_{0.01}\text{Ru}_2\text{Si}_2$

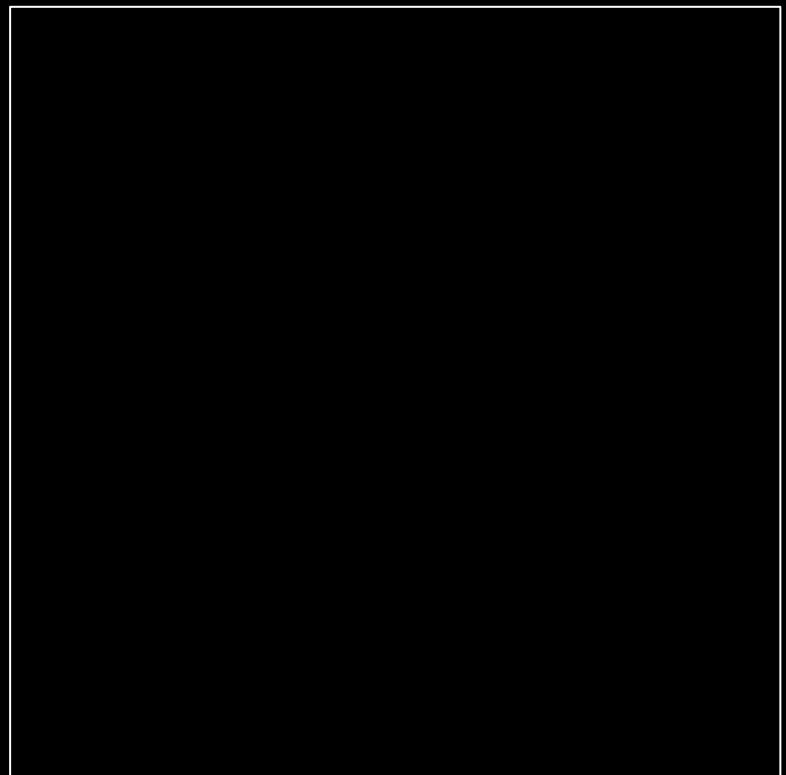




HO Transition in QPI $g(q,E)$



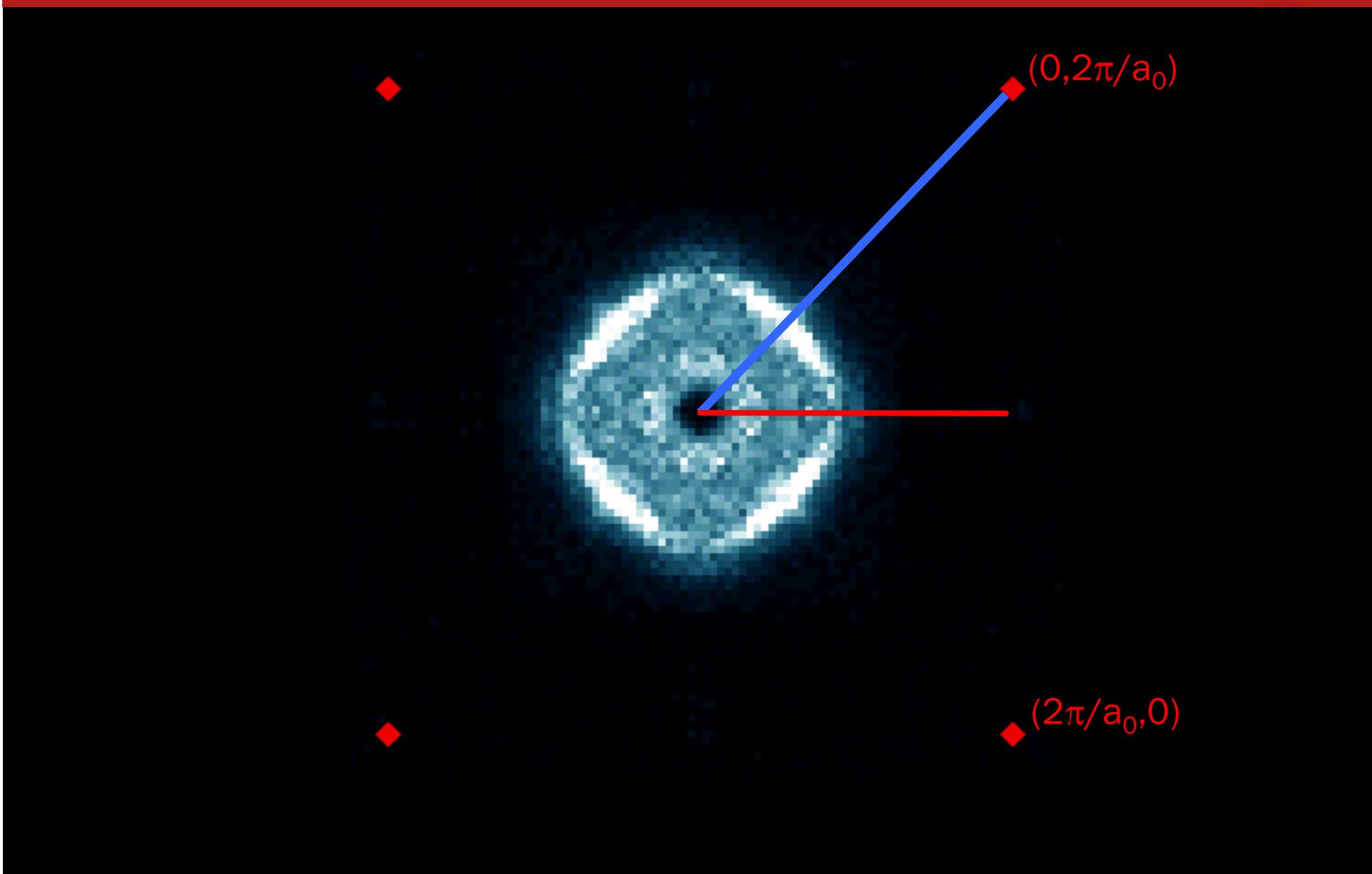
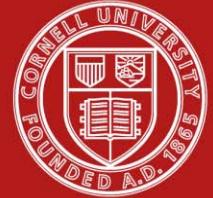
$T > T_c$ (19K)



$T \ll T_c$ (2K)



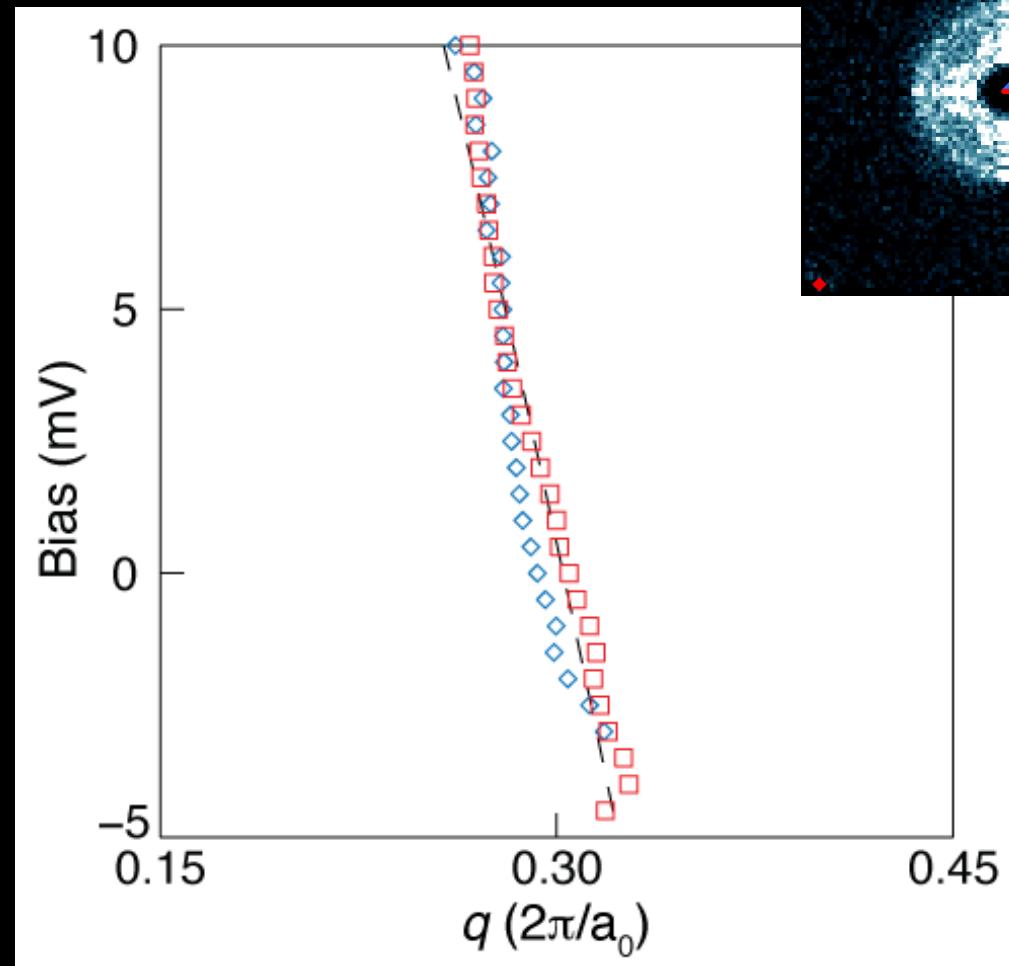
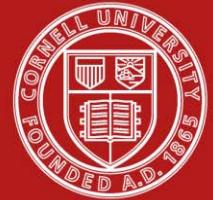
Tracking Heavy QPI

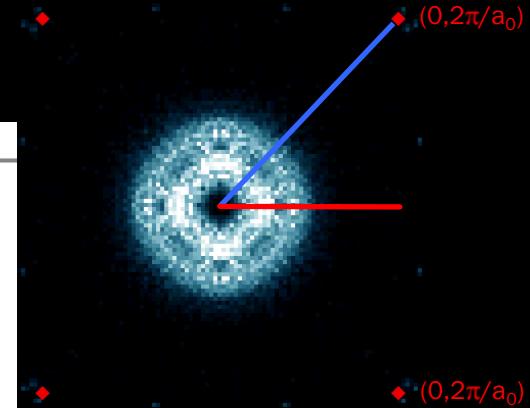
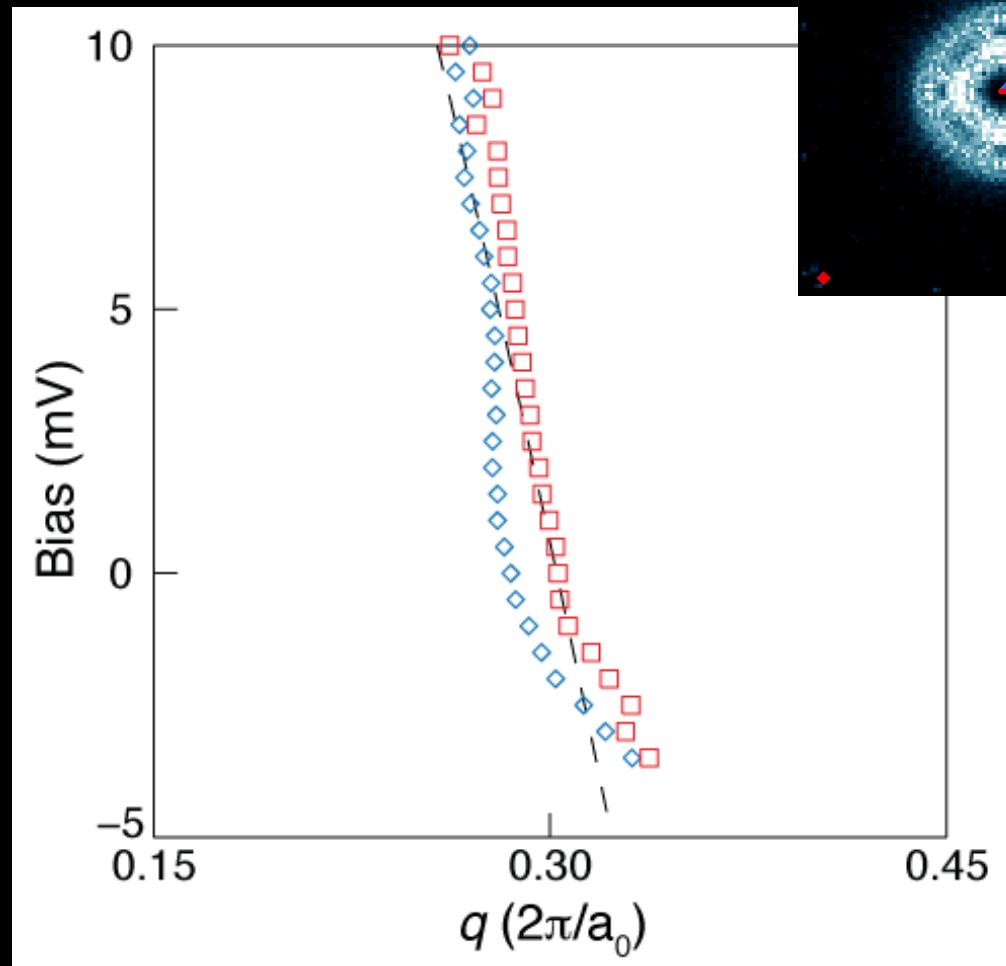


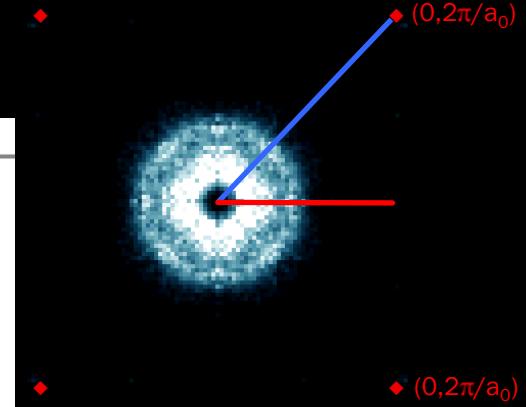
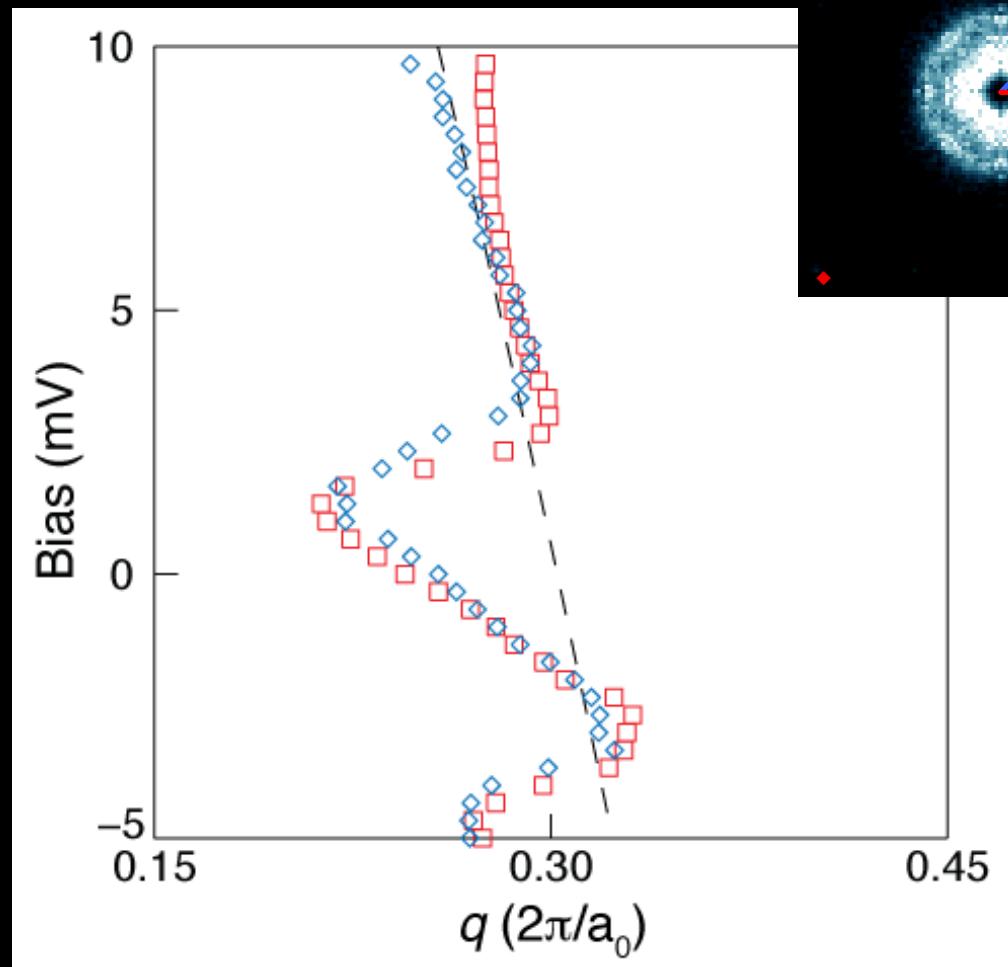


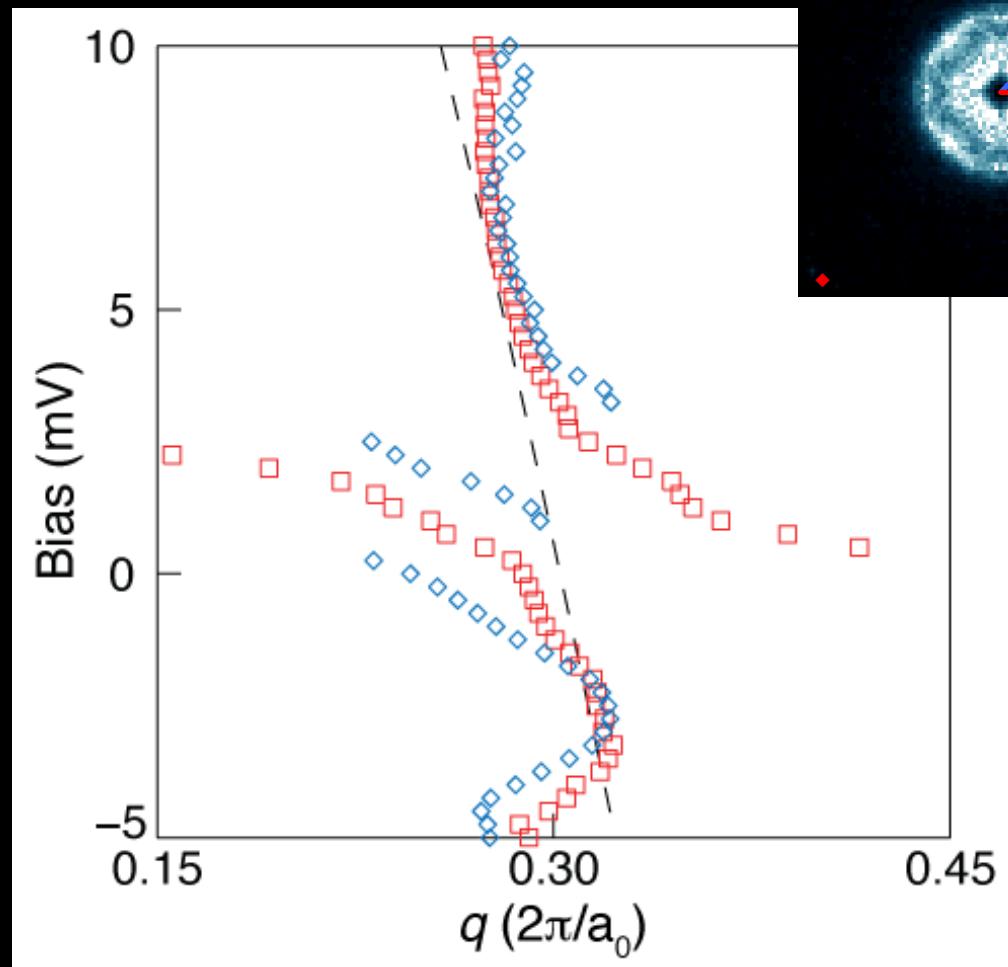
Heavy QPI

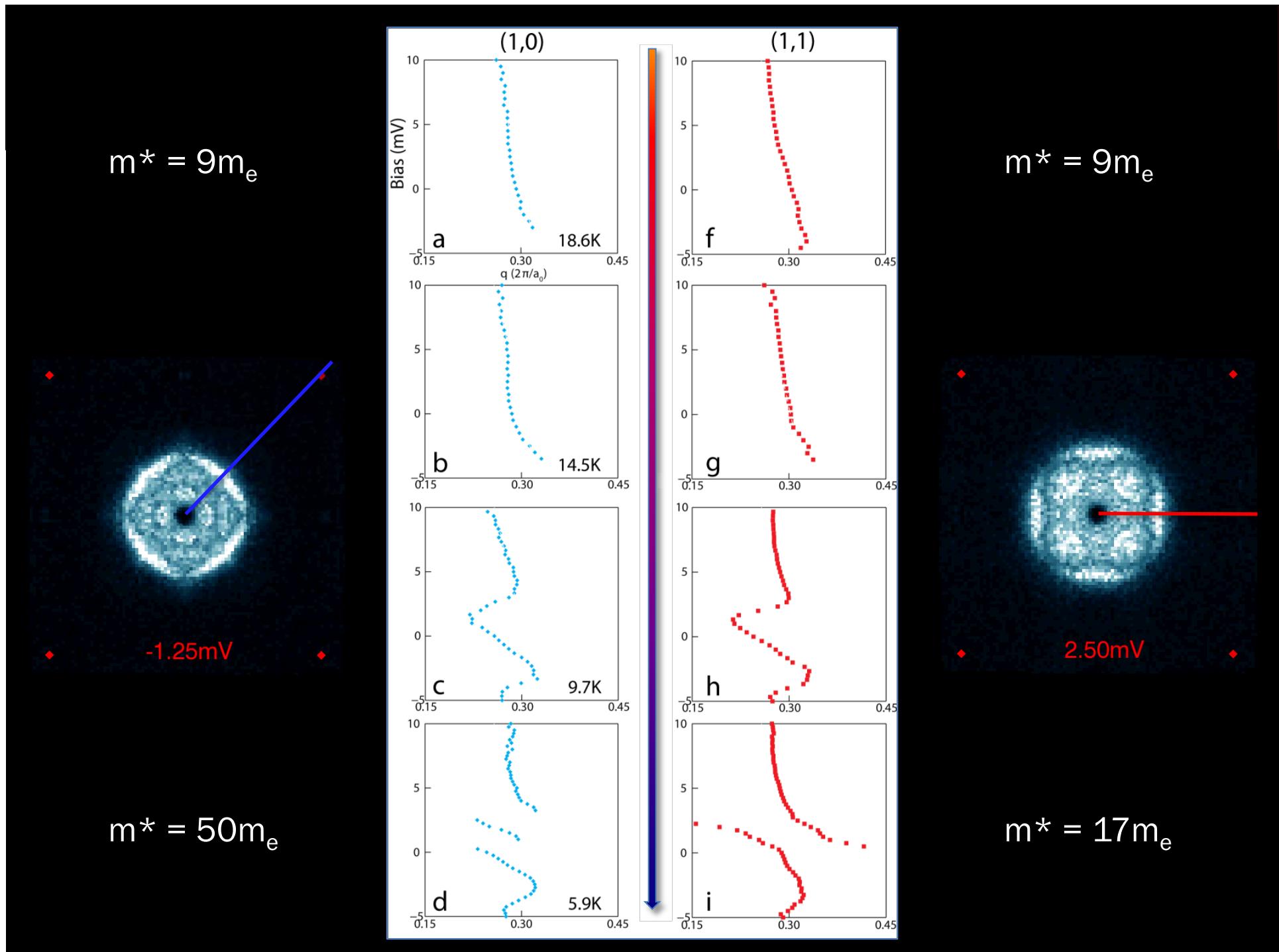
$18.6\text{K} > T_{\text{HO}}$



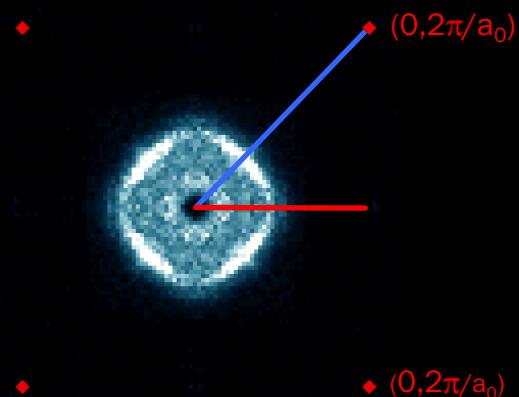
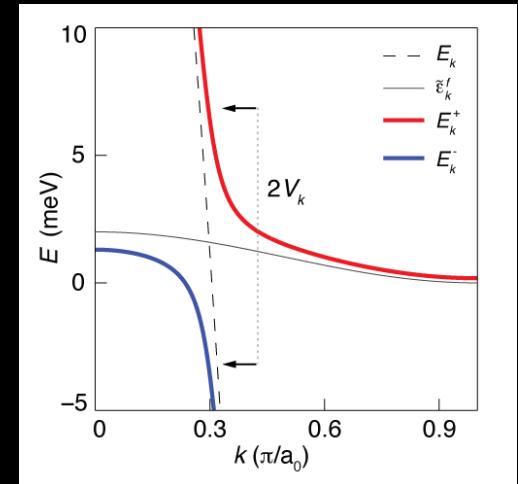
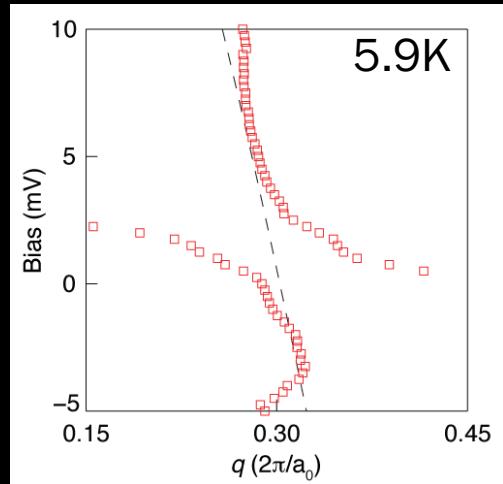
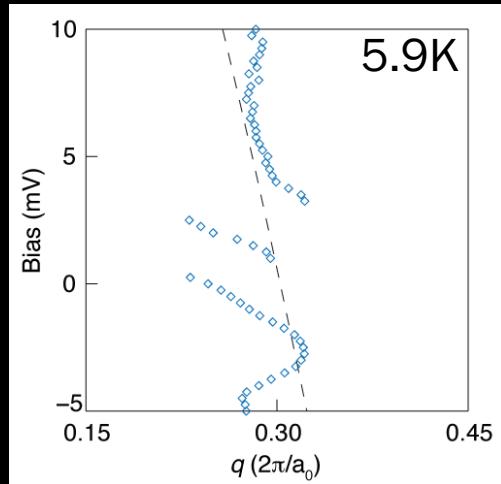
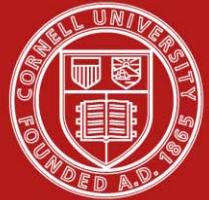






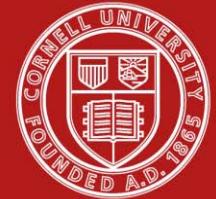


HO: Two New Heavy Bands

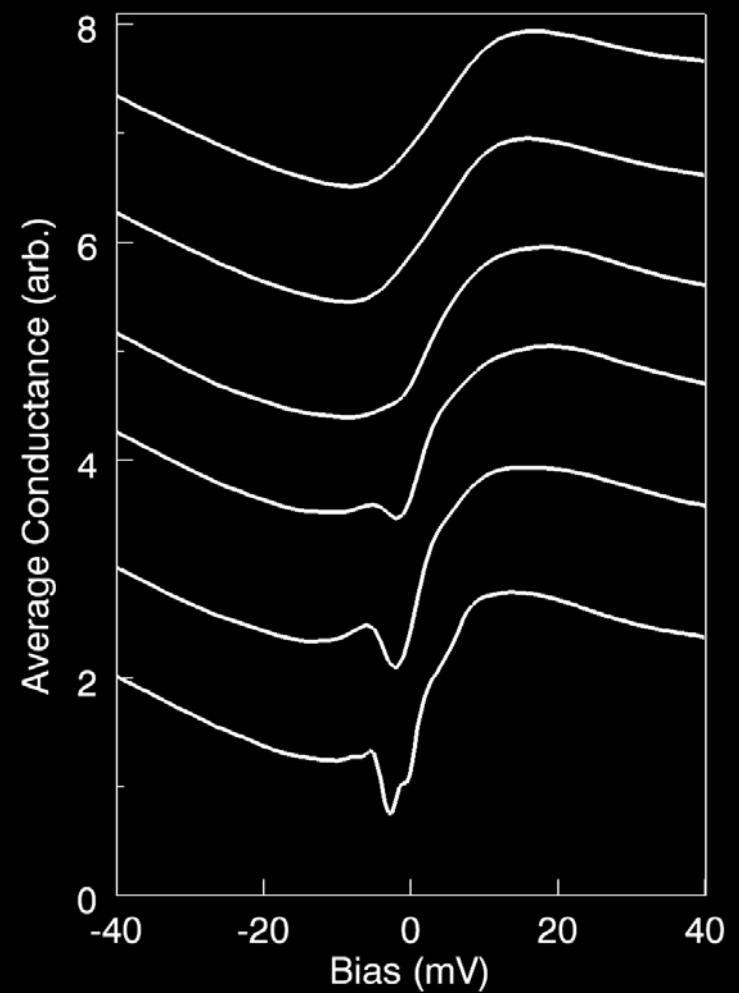


$$E_k^\pm = \frac{\tilde{\varepsilon}_k^f + E_k \pm \sqrt{(\tilde{\varepsilon}_k^f - E_k)^2 + 4|\tilde{V}_k|^2}}{2}$$

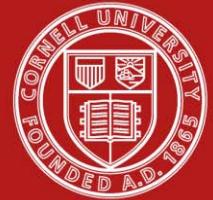
New Perspectives from QPI on HO



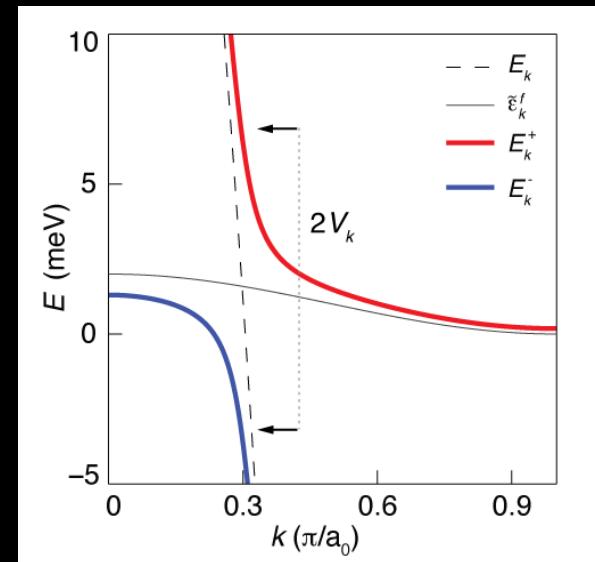
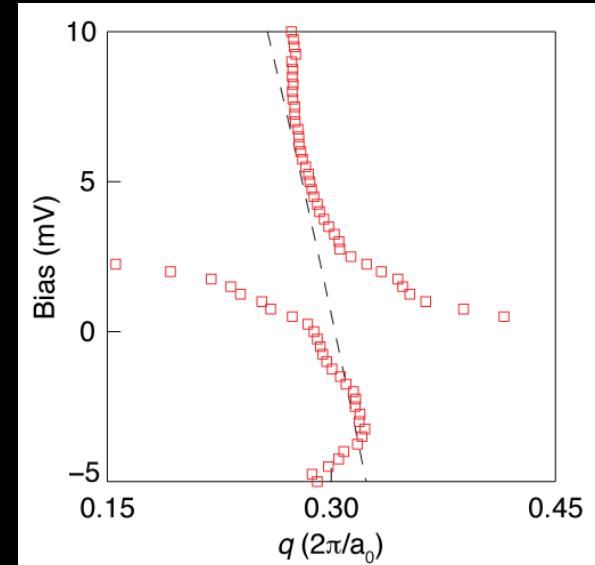
- Thermodynamics, ARPES, SI-STM are consistent with each other in HO phase
- No fixed \mathbf{Q} modulations, gap-edge state different \mathbf{k} -space locations below and above E_F , indirect gap does not cross $E_{F,\text{c}}$.
- DOS(E) emerging below T_0 looks quite like predicted gap for Kondo Lattice (!!)



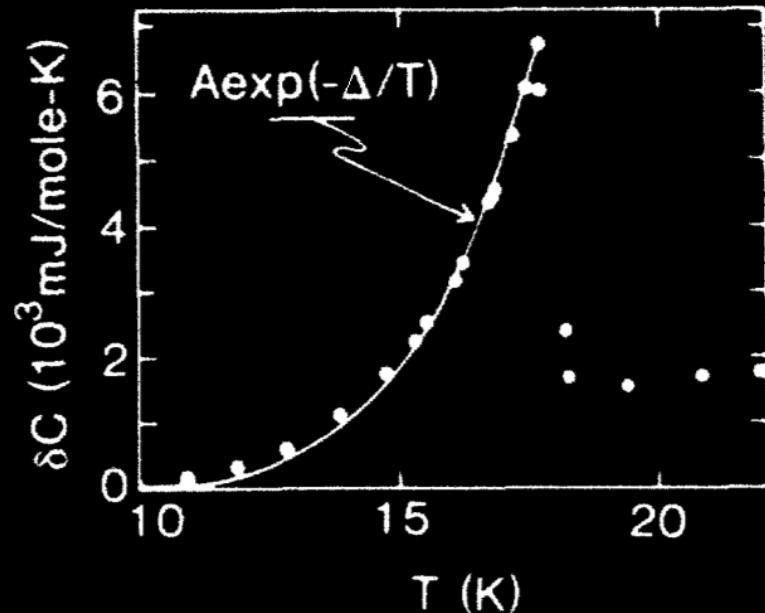
New Perspectives from QPI on HO

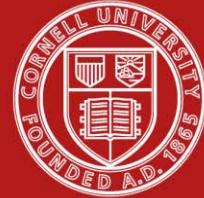


- Thermodynamics, ARPES, SI-STM are consistent with each other in HO phase
- No fixed \mathbf{Q} modulations, gap-edge state different \mathbf{k} -space locations below and above E_F , indirect gap does not cross $E_{F,\pm}$.
- DOS(E) emerging below T_0 looks quite like predicted gap for Kondo Lattice (!)
- A single light band is split into two new heavy bands below T_0
- These new bands appear remarkably like expectations for a Kondo Lattice (!)



- Thermodynamics, ARPES, SI-STM are consistent with each other in HO phase
- No fixed \mathbf{Q} modulations, gap-edge state different \mathbf{k} -space locations below and above E_F , indirect gap does not cross $E_{F,\downarrow}$
- DOS(E) emerging below T_0 looks quite like predicted gap for Kondo Lattice (!!)
- A single light band is split into two new heavy bands below T_0
- These new bands appear remarkably like expectations for a Kondo Lattice (!!)
- BUT... Mean-field-like, second order transition

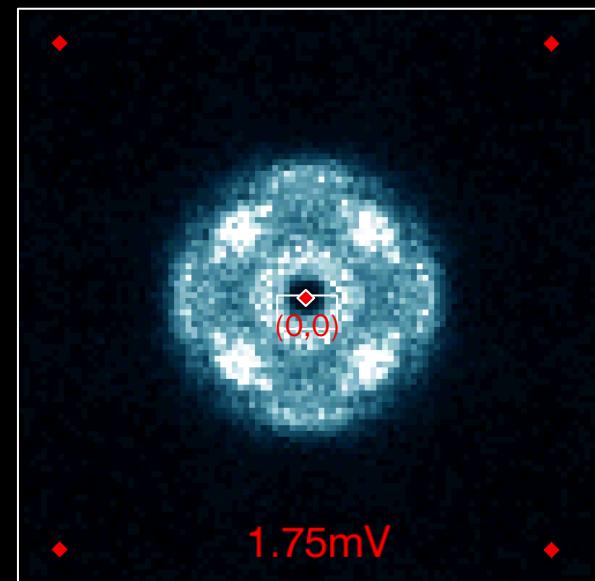
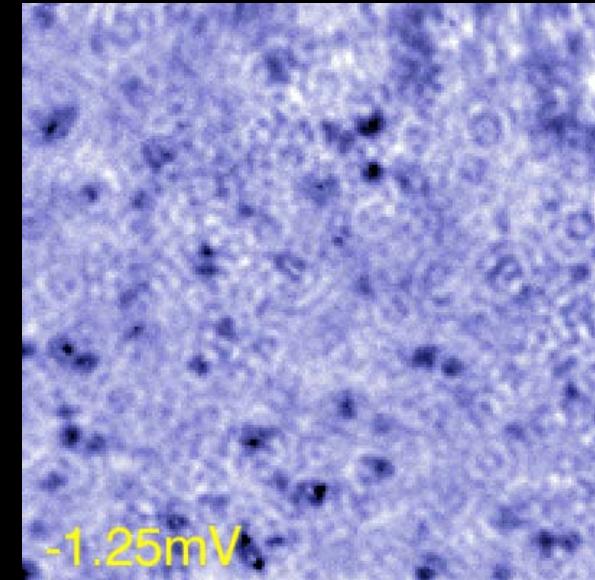
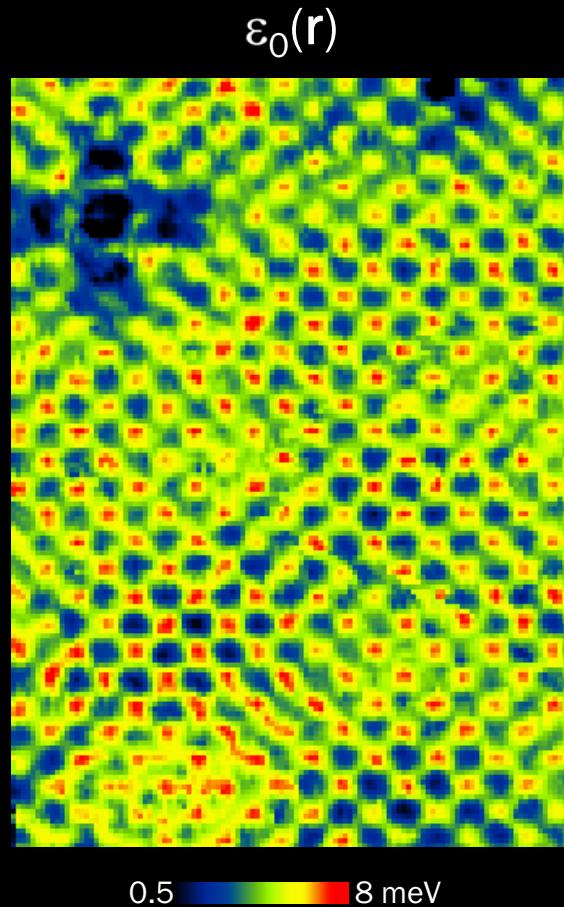
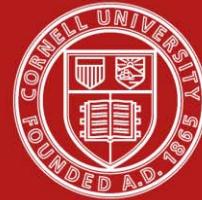




Future



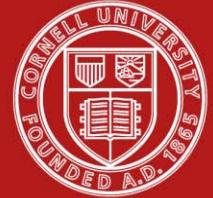
mK SI-STM/QPI for Heavy Fermions



Nature 465, 570 (2010)



Future Work



- SI-STM and QPI opens a new window onto the heavy fermion problem
 - Visualization Kondo Lattice formation/deformation
 - QPI carries symmetries of the Kondo interactions and allows the intricacies of the heavy bands to be measured
 - QPI of heavy f-electron superconductivity
 - The symmetry of URu₂Si₂ ‘hidden order’ within reach (?)

Collaboration



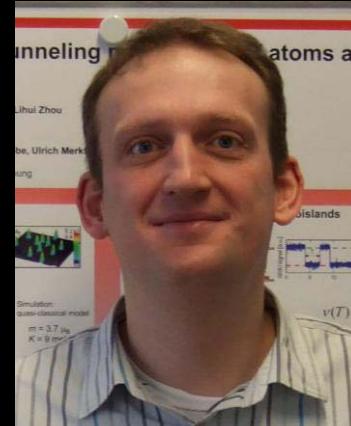
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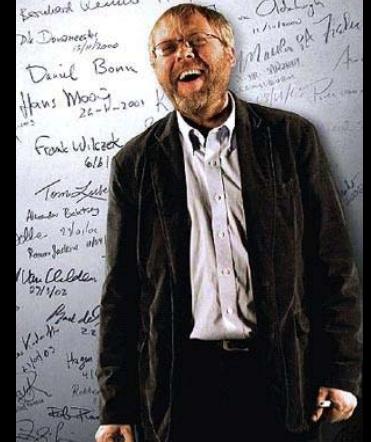
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Thank You

