



School on New Trends in Quantum Dynamics and Quantum Entanglement

14 - 18 February 2011

Quantum Channels with Memory

Nilanjana DATTA DPMMS, Mathematical Sciences Univ. of Cambridge U.K.

Memory Effects in Quantum Channels

Nilanjana Datta University of Cambridge, U.K.



- One of the most common and essential tasks of everyday life: transmission of information
- Examples of classical communications channels: telephones/mobile phones/computers
- A quantum communications channel : one which incorporates intrinsically quantum-mechanical effects
- Example of a quantum communications channel : -- an optical fibre
 - input to the channel : a photon in some quantum-mechanical state



UNIVERSITY OF CAMBRIDGE

- In Quantum Information Theory, information is carried by (or embodied in) *physical states of quantum-mechanical systems*:
- e.g. polarization states of a photon, spin states of electrons
- State space : Hilbert space *H* associated with the system,

(finite-dimensional Hilbert spaces)

e.g. $H \simeq C^2$ Qubit space

B(H): algebra of linear operators acting on H

• States: density matrices, ρ ; $\rho \ge 0$, $\operatorname{Tr} \rho = 1$

 $D(H) \subset B(H)$: set of all density matrices,

UNIVERSITY OF CAMBRIDGE

Quantum Channels

system

interactions

- Any allowed physical process that a quantum system can undergo is described by a :
 linear completely-positive, trace preserving (CPTP) map open open
- e.g. transmission through a quantum communications channel







- Trace preserving (TP): $\text{Tr } \rho' = \text{Tr} \rho = 1$
- *Positive:* $\rho' = \Phi(\rho) \ge 0$
- Completely positive (CP): $\Phi: D(H_A) \to D(H_B)$



 $\in D(H_{R}\otimes H_{F})$

 $(\Phi \otimes id_F)(\rho_{AE})$

= an allowed state of the composite system

 $(\Phi \otimes id)(\rho_{AE}) \geq 0$



"Church of the larger Hilbert Space"





Kraus Representation Theorem:

A quantum channel $\Phi: D(H_A) \rightarrow D(H_B)$

can be represented as follows:

$$\Phi(\rho) = \sum_{i=1}^{M} A_i \rho A_i^{\dagger}$$





- The biggest hurdle in the path of efficient transmission of information:
- -- Presence of noise in communications channels.
- Noise distorts the information sent through the channel.



• To combat the effects of noise: use error-correcting codes



To overcome the effects of noise:

instead of transmitting the original messages,

- -- the sender encodes her messages into suitable codewords
- -- these codewords are then sent through (multiple uses of)



• Error-correcting code: $C_n := (E_n, D_n)$:



• The idea behind the encoding:

- To introduce redundancy in the message so that upon decoding, Bob can retrieve the original message with a low probability of error:
- The amount of redundancy which needs to be added depends on the noise in the channel



Example

Memoryless binary symmetric channel (m.b.s.c.)



- it transmits single bits
- effect of the noise: to flip the bit with probability p

Repetition Code

• Encoding: $0 \longrightarrow 000$ $1 \longrightarrow 111$

codewords

• the 3 bits are sent through 3 successive uses of the m.b.s.c.

- Suppose
 O00 (Bob receives)
 m.b.s.c.
 Codeword
 - Decoding : (*majority voting*) 010 → 0 (Bob infers)



- Probability of error for the m.b.s.c. :
 - without encoding = p
 - with encoding = Prob (2 or more bits flipped) := q



Prove: q
-- in this case encoding helps!

(Encoding - Decoding) : Repetition Code.

UNIVERSITY OF CAMBRIDGE

Information transmission is said to be reliable if:
 -- the probability of error in decoding the output
 vanishes asymptotically in the number of uses of the channel







- To overcome the effects of noise:
- Alice encodes: messages ______ codewords ;
 codewords ______ (n uses of) the channel





□ codeword: $x^{(n)} = (x_1, x_2, ..., x_n); x_i \in \{0, 1\}^n$

• encoding:
$$\mathsf{E}_n : \mathsf{M} \mapsto \{0,1\}^n$$

• output:
$$y^{(n)} = (y_1, y_2, ..., y_n); y_i \in \{0, 1\}^n$$

• decoding: $\mathsf{D}_n: \{0,1\}^n \mapsto \mathsf{M}$

• Error-correcting code: $C_n := (E_n, D_n)$:

$$\frac{N^{(n)}:}{p(y^{(n)} | x^{(n)})}$$





• If $m' \neq m$ then an error occurs!

Information transmission is reliable:

Prob. of error $\rightarrow 0$ as $n \rightarrow \infty$

 Rate of info transmission = number of bits of message transmitted per use of the channel
 size of message (in bits) size of codeword (in bits)

- Aim: achieve reliable transmission whilst optimizing the rate
- Capacity: optimal rate of reliable information transmission



- Shannon in his Noisy Channel Coding Theorem:
- -- obtained an explicit expression for the capacity of a

memoryless classical channel



Memoryless (classical or quantum) channels

- action of each use of the channel is identical and it is independent for different uses
- -- i.e., the noise affecting states transmitted through the channel on successive uses is assumed to be uncorrelated.



Classical memoryless channel: a schematic representation

$$X \sim p(x) \qquad Y$$
input x
$$N \qquad output y$$

$$x \in J_X, \qquad p(y \mid x) \qquad y \in J_Y,$$

• channel: a set of conditional probs. $\{p(y | x)\}$

• Capacity
$$C(N) = \max_{\{p(x)\}} I(X : Y)$$

input distributions mutual information
 $I(X : Y) = H(X) + H(Y) - H(X, Y)$

Shannon Entropy $H(X) = -\sum p(x) \log p(x)$





• A classical channel has a unique capacity

BUT

a quantum channel has various different capacities

-- This is due to the greater flexibility in the use of a quantum channel

UNIVERSITY OF CAMBRIDGE

- The different capacities depend on:
 - the nature of the transmitted information

```
(classical or quantum)
```

the nature of the input states

(entangled or product states)

- the nature of the measurements done on the outputs (collective or individual)
- the presence or absence of any additional resource (e.g. prior shared entanglement between Alice & Bob)
- whether Alice & Bob are allowed to communicate classically with each other

UNIVERSITY OF CAMBRIDGE





C(Φ): Classical capacity of Φ
 = the maximum amount of classical info (in bits) that can be reliably transmitted per use of

 $C(\Phi) = \sup R$

--the supremum taken over all achievable rates



If Alice restricts her codewords to product states, i.e., if

$$i \to \rho_i^{(n)} = \rho_{i_1} \otimes \rho_{i_2} \otimes \dots \otimes \rho_{i_n}$$

And Bob does a collective measurement (POVM) on

 $\sigma_i^{(n)} \coloneqq \Phi^{(n)}(\rho_i^{(n)})$: the output of n uses of the channel

• Capacity : product state capacity $C^{(1)}(\Phi)$



Memoryless (classical and quantum) channels

- action of each use of the channel is identical and it is independent for different uses
- -- i.e., the noise affecting states transmitted through the channel on successive uses is assumed to be uncorrelated.

- Let $\Phi^{(n)}$: *n* successive uses of a quantum channel Φ
- For a memoryless channel:

$$\Phi^{(n)} = \Phi^{\otimes n}$$



Multiple uses of a memoryless channel

Consider a memoryless channel defined by

$$\Phi(\rho) = \sum_{i=1}^{M} A_{i} \rho A_{i}^{\dagger} \qquad \forall \rho \in D(H)$$

• Then the output of n uses of the channel is given by

$$\Phi^{(n)}\left(\rho^{(n)}\right) \equiv \Phi^{\otimes n}\left(\rho^{(n)}\right), \ \forall \ \rho^{(n)} \in \boldsymbol{D} \ (\boldsymbol{H}^{\otimes n})$$

where

$$\Phi^{(\otimes n)}\left(\rho^{(n)}\right) = \sum_{k_1,\dots,k_n}^{M} \left(A_{k_1} \otimes \dots \otimes A_{k_n}\right) \rho^{(n)}\left(A_{k_1}^{\dagger} \otimes \dots \otimes A_{k_n}^{\dagger}\right)$$



Transmission of classical info through a memoryless quantum

channel

$$\Phi^{(n)} = \Phi^{\otimes n}$$

• For product state inputs: $i \to \rho_i^{(n)} = \rho_{i_1} \otimes \rho_{i_2} \otimes \dots \otimes \rho_{i_n}$

Outputs = product states

$$\sigma_i^{(n)} \coloneqq \Phi^{\otimes n}(\rho_i^{(n)}) = \Phi(\rho_{i_1}) \otimes \Phi(\rho_{i_2}) \otimes \dots \otimes \Phi(\rho_{i_n})$$

- Product State Capacity $C^{(1)}(\Phi)$ given by the
 - Holevo-Schumacher-Westmoreland (HSW) Theorem



HSW Theorem

$$C^{(1)}(\Phi) = \max_{\{p_i, \rho_i\}} \chi(\{p_i, \Phi(\rho_i)\}) = \chi^*(\Phi)$$

Holevo Capacity

where

$$\chi\left(\{p_{i},\Phi(\rho_{i})\}\right) = S\left(\sum_{i} p_{i}\Phi(\rho_{i})\right) - \sum_{i} p_{i}S\left(\Phi(\rho_{i})\right)$$

 $S(\sigma) = -\text{tr}\sigma \log \sigma \quad : \text{ von Neumann entropy}$ Holevo χ - quantity: Let $\sigma_i := \Phi(\rho_i)$

$$\chi(\{p_{i},\sigma_{i}\}) = S(\sum_{i} p_{i}\sigma_{i}) - \sum_{i} p_{i}S(\sigma_{i})$$



• Holevo χ – quantity of an ensemble of states $\{p_i, \sigma_i\}$

$$\chi(\{p_{i},\sigma_{i}\}) \coloneqq S(\sum_{i} p_{i}\sigma_{i}) - \sum_{i} p_{i}S(\sigma_{i})$$

• Holevo Bound The maximum amount of classical info that Alice can send to Bob using $\{p_i, \sigma_i\}$ is $\leq \chi(\{p_i, \sigma_i\})$ $\chi(\{p_i, \sigma_i\}) \rightarrow S(\sigma)$ where $\sigma \coloneqq \sum_i p_i \sigma_i$ if the σ_i are pure $\because S(\sigma_i) = 0$



HSW Theorem

$$C^{(1)}(\Phi) = \max_{\{p_{i},\rho_{i}\}} \chi(\{p_{i},\Phi(\rho_{i})\}) = \chi^{*}(\Phi)$$

Holevo Capacity

tells us that the Holevo bound can be achieved --

IF Alice uses product state inputs

& Bob does a collective measurement

Optimal signal ensemble

$$\begin{array}{ccc} p_{av}^{(n)} \to 0 & p_{av}^{(n)} \not \to 0 \\ \hline \text{as } n \to \infty & \chi^*(\Phi) & \text{as } n \to \infty & R \text{ (rate)} \end{array}$$



• Classical capacity of a memoryless channel Φ :

(without the restriction of inputs being product states):

$$C(\Phi) = \lim_{n \to \infty} \frac{1}{n} \chi^* \left(\Phi^{\otimes n} \right)$$

regularised Holevo capacity

 $\chi^*(\Phi^{\otimes n})$ Holevo Capacity of the block $\Phi^{\otimes n}$ of n channels

 This generalization is obtained by considering inputs which are product states over blocks of n channels but which may be entangled within each block



 Classical capacity of a memoryless channel

$$C(\Phi) = \lim_{n \to \infty} \frac{1}{n} \chi^* \left(\Phi^{\otimes n} \right)$$

(Q) Can the classical capacity of a memoryless quantum channel be increased by using entangled states as inputs?

 This is related to the additivity conjecture of the Holevo capacity:

$$\chi^{*}(\Phi_{1} \otimes \Phi_{2}) = \chi^{*}(\Phi_{1}) + \chi^{*}(\Phi_{2}) \Longrightarrow \chi^{*}(\Phi^{\otimes n}) = n\chi^{*}(\Phi)$$
$$\Rightarrow C(\Phi) = \lim_{n \to \infty} \frac{1}{n} \chi^{*}(\Phi^{\otimes n}) = \lim_{n \to \infty} \frac{1}{n} \chi^{*}(\Phi) = \chi^{*}(\Phi)$$



 IF the Holevo capacity is additive for a memoryless quantum channel then using entangled inputs would not increase its classical capacity

• An interesting question:

Could entangled inputs increase the classical capacity of quantum channels with memory ?



For real-world communications channels, the assumption : noise is uncorrelated between successive uses of a channel cannot be justified!

Hence, memory effects need to be taken into account



quantum channels with memory

There are various examples of quantum channels with memory :

e.g. (1) one-atom maser or micromaser(2) spin-chain



• (1) One-atom maser or micromaser



- A stream of two-level atoms injected into an optical cavity.
- States of these input two-level atoms: signal states
- The atoms interact with the photons in the cavity
- If these photons have sufficiently long lifetimes, then the atoms entering the cavity feel the effect of the preceding atoms
- This introduces correlations between consecutive signal states



• (2) State transfer across a spin chain



- a spin chain : governed by a suitable Hamiltonian
- Spins at one end of the chain are prepared (by Alice) in the state which is required to be transmitted
- The spin chain is allowed to evolve for a specific amount of time under the action of the Hamiltonian; causing state to propagate
- The state is then retrieved from a set of spins at the other end of the spin chain - thus state transfer is achieved !


Alice

• (2) State transfer across a spin chain



When considered as a model for quantum communication:

assume : a reset of the spin chain occurs after the transmission of each signal e.g. by applying an external

magnetic field

memoryless quantum channel

 A continuous operation without reset might lead to higher transmission rates

quantum channel with memory



Exercises

- (1) Use the HSW theorem to prove that any quantum channel can be used to transmit classical information, as long as it is not a constant.
- (2) Use the HSW theorem to evaluate the product state capacity of a qubit depolarizing channel.

$$\Phi_{dep}(\rho) = (1-p)\rho + \frac{p}{3} \left[\sigma_x \rho \sigma_x + \sigma_y \rho \sigma_y + \sigma_z \rho \sigma_z \right]$$



Memory Effects in Quantum Channels

LECTURE II

Forgetful Channels

UNIVERSITY OF CAMBRIDGE

Lecture I - revisited

- IF Alice sends classical info through a quantum channel Φ
- using product state inputs
- Bob does a collective measurement
- Then capacity : Product-state capacity

HSW Theorem





$$C^{(1)}(\Phi) = \max_{\{p_i,\rho_i\}} \chi\left(\{p_i,\Phi(\rho_i)\}\right) = \chi^*(\Phi)$$

Holevo Capacity

Optimal signal ensemble



Regularised Holevo capacity



Quantum channels with memory = *quantum memory channels*

• Strategy : (i) start with a simple example of a memoryless quantum channel (ii) using it, construct a quantum memory channel Qubit depolarizing channel $\Phi_{dep} : B(H) \rightarrow B(H); \quad H \simeq C^2$

$$\Phi_{dep}(\rho) = (1-p)\rho + \frac{p}{3} \Big[\sigma_x \rho \sigma_x + \sigma_y \rho \sigma_y + \sigma_z \rho \sigma_z \Big] \quad \forall \rho \in D(H)$$
$$= \sum_{i=1}^{4} p_i \sigma_i \rho \sigma_i \qquad p_1 = (1-p); \quad p_2 = p_3 = p_4 = \frac{p}{3}$$

$$\sigma_{1} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}; \sigma_{2} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}; \sigma_{3} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}; \sigma_{4} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$
$$\sigma_{x} \qquad \sigma_{y} \qquad \sigma_{z}$$



$$\Phi_{dep}(\rho) = \sum_{i=1}^{4} p_i \sigma_i \rho \sigma_i$$

 $\Phi_{dep}^{(n)} = \Phi_{dep}^{\otimes n}$

- since the channel is memoryless
- input : $\rho^{(n)} \in D(H^{\otimes n})$; the output is given by

$$\Phi_{dep}^{\otimes n}(\rho^{(n)}) = \sum_{i=1}^{4} p_{i_1} p_{i_2} \dots p_{i_n} \left(\sigma_{i_1} \otimes \sigma_{i_2} \otimes \dots \otimes \sigma_{i_n} \right) \left(\rho^{(n)} \right) \left(\sigma_{i_1} \otimes \sigma_{i_2} \otimes \dots \otimes \sigma_{i_n} \right)$$

- Let $P_{i_1i_2...,i_n}$ = joint prob. of the n successive qubits being acted on by $\sigma_{i_1}, \sigma_{i_2}, ..., \sigma_{i_n}$ resply. $P_{i_1i_2...,i_n} = p_{i_1}p_{i_2}...p_{i_n}$
- this is in keeping with the notion that the noise acts independently on each successive use



Next:

- we consider an interesting generalization of this model
- which yields a model of a quantum memory channel

• This generalization involves a :

discrete-time Markov Chain



MARKOV CHAINS

- simplest mathematical models for random phenomena evolving in time
- It is a random process with the characteristic property that it retains no memory of where it has been in the past
- so only the current state of the process can influence where it goes next.

Discrete-time Markov Chain:

- time is discrete
- the instants of time are labelled by $n \in \mathbb{Z}^+ = \{0, 1, 2, ...\}$



An Example

Consider a fly hopping on the vertices of a triangle



$$I = \{1, 2, 3\}$$
 = state space of the MC

• Suppose the fly hops

- clockwise with prob.
 - anticlockwise with prob. $\frac{3}{-}$

So where it hops next depends only on where it is now

$$q_{12}$$
 = Prob(hops to 2 in next step | it is at 1) = $\frac{2}{3}$

 $q_{13} =$ Prob (hops to 3 in next step | it is at 1) $= \frac{1}{3}$ etc.

UNIVERSITY OF CAMBRIDGE





Invariant distribution

$$\lambda = \{\lambda_i\}_{i \in I}$$
; $\lambda = \lambda Q$; $\lambda_j = \sum_j \lambda_i q_{ij}$;

- many of the long-term properties of a MC depend on its invariant distribution
- A Markov Chain is defined by a sequence of random variables $X_0, X_1, ..., X_n; (X_n)_{n \ge 0}$; each X_i takes values in I $P(X_{n+1} = i_{n+1} | X_0 = i_0, X_1 = i, ..., X_n = i_n) = P(X_{n+1} = i_{n+1} | X_n = i_n)$ *Markov Property* $P(X_{n+1} = j | X_n = i) = q_{ij}$ = transition probability



Some properties of a Markov Chain

- Irreducibility:
- -- a Markov Chain is said to be irreducible if it is possible to go from any state to any other state in the chain



• Aperiodicity :

 -- a Markov Chain is said to be aperiodic if the return time to any state in the chain is not periodic (or if it has a period = 1)
 i.e., return can occur at irregular times





• n uses of a memoryless depolarising channel :

$$\Phi_{dep}^{\otimes n}(\rho^{(n)}) = \sum_{i=1}^{4} p_{i_{1}i_{2}....i_{n}} \left(\sigma_{i_{1}} \otimes \sigma_{i_{2}} \otimes ... \otimes \sigma_{i_{n}}\right) \left(\rho^{(n)}\right) \left(\sigma_{i_{1}} \otimes \sigma_{i_{2}} \otimes ... \otimes \sigma_{i_{n}}\right)$$

$$p_{i_{1}i_{2}....i_{n}} = p_{i_{1}}p_{i_{2}}....p_{i_{n}}$$
Now consider the case in which:
$$p_{i_{1}i_{2}....i_{n}} = \gamma_{i_{1}}q_{i_{1}i_{2}}....q_{i_{n-1}i_{n}}$$

$$q_{ij}, i, j = 1, 2, 3, 4:$$
 the elements of the transition matrix Q

of a discrete-time Markov Chain with state space $I = \{1, 2, 3, 4\}$

• Note : the states 1, 2, 3, 4, label the matrices

 $\sigma_1, \sigma_2, \sigma_3, \sigma_4$ of the depolarizing channel



In this case : the output after n uses

$$\Phi_{dep}^{(n)}(\rho^{(n)}) = \sum_{i_1,\dots,i_n} \gamma_{i_1} q_{i_1 i_2} \dots q_{i_{n-1} i_n} \left(\sigma_{i_1} \otimes \dots \otimes \sigma_{i_n}\right) \left(\rho^{(n)}\right) \left(\sigma_{i_1} \otimes \dots \otimes \sigma_{i_n}\right)$$

 $q_{i_{k-1}i_k} = P(k^{th}$ qubit acted on by $\sigma_{i_k} | (k-1)^{th}$ qubit acted on by $\sigma_{i_{k-1}}$)

 $\gamma_{i_1} = P(\text{the } 1^{st} \text{ qubit acted on by } \sigma_{i_k})$

- the noise acting on the k^{th} qubit depends on the noise acting on the $(k-1)^{th}$ qubit
- Note : the noise acting on successive qubits is correlated

model of a quantum memory channel - with Markovian

correlated noise

UNIVERSITY OF CAMBRIDGE

Macchiavello & Palma : -- introduced this model

studied the transmission of classical information through
 2 successive uses of this quantum memory channel with

$$q_{ij} = (1 - \mu)\gamma_j + \mu \delta_{ij} ; 0 \le \mu \le 1$$

i, j = 1, 2, 3, 4:

- with prob. μ the 2 qubits are acted on identically
- with prob. (1μ) the action of the channel on the 2 qubits is uncorrelated
 - μ : the degree of memory of the channel
 - $\mu = 0$: uncorrelated noise \bullet
 - $\mu = 1$: fully correlated noise (successive actions identical)



Macchiavello & Palma : -- showed that above a certain threshold value of the parameter μ

entangled inputs increase the Holevo χ – quantity

for 2 successive uses of the channel

• This suggests that above this value of μ :

One might be able to transmit a higher amount of classical information through this channel by using entangled input states

They did not, however, compute the capacity of the channel



- Next:
- we consider a more general quantum memory channel

with Markovian correlated noise

(of which the above model is a special case)

- and study its capacities
 - The model is constructed from ---- a finite set of memoryless quantum channels: $\{\phi_1, \phi_2, ..., \phi_M\}$ qubit channels $\forall i = 1, 2, ..., M, \phi_i : B(H) \rightarrow B(H); H \simeq C^2$



Quantum Channel with Markovian Correlated Noise

n uses of the channel

$$\Phi^{(n)}(\rho^{(n)}) = \sum_{i_1,\dots,i_n=1}^M \gamma_{i_1} q_{i_1 i_2} \dots q_{i_{n-1} i_n} \left(\phi_{i_1} \otimes \dots \otimes \phi_{i_n}\right) \left(\rho^{(n)}\right)$$

 q_{ij} elements of the transition matrix of a discrete-time Markov chain with finite state space I = {1, 2, ..., M}

$$\{\gamma_i\}$$
 = invariant distribution

- For each $i \in I$, ϕ_i CPT map on B(H):
- On each use of the channel, one of the given set of CPTP maps $\{\phi_1, \phi_2, ..., \phi_M\}$ acts on the qubit



Depending on the nature of the Markov Chain the channel Either : (1) forgetful or (2) not-forgetful

(1) Forgetful channel : a channel in which the correlation in the noise dies out with time

(2) not-forgetful channel : a channel with long-term memory

(Q) When is the quantum memory channel with Markovian correlated noise forgetful?



(A) If the underlying Markov chain is



In this case the Markov Chain has a unique invariant distribution and it satisfies the property called



UNIVERSITY OF CAMBRIDGE

lf

Quantum Channel with Markovian Correlated Noise

$$\Phi^{(n)}(\rho^{(n)}) = \sum_{i_1,...,i_n=1}^M \gamma_{i_1} q_{i_1 i_2} \dots q_{i_{n-1} i_n} \left(\phi_{i_1} \otimes ... \otimes \phi_{i_n}\right) \left(\rho^{(n)}\right)$$

satisfies "convergence to equilibrium"

$$i, j \in \mathbf{I}, \qquad q_{ij}^{(n)} \xrightarrow[n \to \infty]{} \gamma_j$$

• For *n* large enough, the prob. that the n^{th} qubit sent through the channel is acted upon by the memoryless channel ϕ_j does not depend on which memoryless channel ϕ_i acted on the first qubit.



In this case (of a forgetful channel) :

- The classical capacity of the channel is given by a formula which is very similar to that of a memoryless channel
- For a memoryless channel

$$C(\Phi) = \lim_{n \to \infty} \frac{1}{n} \chi^* (\Phi^{\otimes n})$$
$$= \lim_{n \to \infty} \frac{1}{n} \max_{\{p_i, \rho_i^{(n)}\}} \chi (\{p_i, \Phi^{\otimes n}(\rho_i^{(n)})\})$$

regularised Holevo capacity

For our forgetful channel

$$C(\Phi) = \lim_{n \to \infty} \frac{1}{n} \max_{\{p_i, \rho_i^{(n)}\}} \chi(\{p_i, \Phi^{(n)}(\rho_i^{(n)})\})$$



Why?

- The reason behind getting such a similar result: can be explained by a simple double-blocking argument
- We shall consider this argument in a more general setting.

 (I) Forgetful channels form an important subclass of ALL quantum channels with memory – (not only those with Markovian correlated noise)

- (II) For forgetful channels, expressions for each of the different capacities are similar to the corrs. capacity formulas for memoryless channels ---
- -- and can be understood by a double-blocking argument



General model for quantum channels with memory

- Thus far : we have studied only a small class of quantum memory channels those in which the memory is
- (i) classical and (ii) governed by an underlying Markov Chain
 - Bowen & Mancini : introduced a more general model for quantum memory channels in which the memory could even be quantum.
 - Kretschmann & Werner : studied this model exhaustively in the Heisenberg picture
 - -- they were the first to evaluate capacities of forgetful channels.



- In this model : a forgetful channel is one in which :
 The effect of the initializing memory dies away with time
- **Recall**: for the Markovian correlated noise model

condition for forgetfulness "convergence to equilibrium"

$$i, j \in \mathbf{I}, \quad q_{ij}^{(n)} \xrightarrow[n \to \infty]{} \gamma_j$$

- it ensures that the initializing memory dies out asymptotically
 - It is easy to evaluate the capacities of forgetful channel by reducing them to a memoryless setting via a double-blocking argument



The double-blocking argument

- Consider a strictly forgetful channel Φ
- one in which : the effect of the initializing memory dies away after a finite number of uses (say, m uses)
 - e.g. transmission of info over a quantum spin chain which is reset after every third use (m = 3).





Strictly forgetful channel





Strictly forgetful channel



- ignore the outputs of the first *m* channels of each such block
- actual encoding is done for the remaining l blocks
- Eventually let $l \rightarrow \infty$
- If we restrict inputs to products states of block length m+l

$$\operatorname{input} = \rho_1^{(m+l)} \otimes \rho_2^{(m+l)} \otimes \dots$$



- due to the strict forgetfulness of the channel:
 - -- the (relevant part of the) output state factorizes
- The whole set-up corrs. to a memoryless channel acting on a larger Hilbert space $\Phi^{(m+l)} \approx memoryless$ channel





The same double-blocking argument can be applied to channels which are forgetful (and not just strictly forgetful)

Classical Capacity

For a memoryless channel

$$C(\Phi) = \lim_{n \to \infty} \frac{1}{n} \chi^* (\Phi^{\otimes n})$$
$$= \lim_{n \to \infty} \frac{1}{n} \max_{\{p_i, \rho_i^{(n)}\}} \chi (\{p_i, \Phi^{\otimes n}(\rho_i^{(n)})\})$$

regularised Holevo capacity

For forgetful channels

$$C(\Phi) = \lim_{n \to \infty} \frac{1}{n} \max_{\{p_i, \rho_i^{(n)}\}} \chi(\{p_i, \Phi^{(n)}(\rho_i^{(n)})\})$$

UNIVERSITY OF CAMBRIDGE







Quantum Capacity

For a memoryless channel

$$Q(\Phi) = \lim_{n \to \infty} \frac{1}{n} \max_{\rho^{(n)}} I_c(\rho^{(n)}, \Phi^{\otimes n})$$

Regularised

Coherent information

• For forgetful channels

$$Q(\Phi) = \lim_{n \to \infty} \frac{1}{n} \max_{\rho^{(n)}} I_c(\rho^{(n)}, \Phi^{(n)})$$



LECTURE III

A channel with long-term memory (not-forgetful)

Coding Theorem for a Class of Quantum Channels with Long-Term Memory,

ND and Tony Dorlas,

J. Phys. A: Math. Theor. 40, 8147-8164 (2007).

UNIVERSITY OF CAMBRIDGE

A channel with long-term memory

- The correlation in the noise does not die out with time
- evaluating their capacities is a more challenging task
- Simplest example:

convex combinations of a finite number of memoryless

channels

 $\forall i=1,2,..,M, \quad \phi_i: D(H_A) \to D(H_B);$

 $\{\phi_1, \phi_2, \dots, \phi_M\}$

$$\Phi^{(n)}(\rho^{(n)}) = \sum_{i=1}^{M} \gamma_i \phi_i^{\otimes n} \left(\rho^{(n)}\right)$$

we uses of the channel: $\gamma_i > 0 \quad \forall i = 1, 2, ..., M, \quad \sum_{i=1}^{M} \gamma_i = 1$



A channel with long-term memory

$$\Phi^{(n)}(\rho^{(n)}) = \sum_{i=1}^{M} \gamma_i \phi_i^{\otimes n} \left(\rho^{(n)}\right)$$

$$H_A, H_B \simeq \mathbb{C}^2$$
 $\gamma_i > 0 \quad \forall i = 1, 2, ..., M, \quad \sum_{i=1}^M \gamma_i = 1$




with the Markovian correlated noise model:

- We note that (a) is a special case of (b):
 - The Markov Chain has M states



 $\left\{\phi_{1},\phi_{2},\ldots,\phi_{M}\right\}$

states of the MC



$$\Phi^{(n)}(\rho^{(n)}) = \sum_{i=1}^{M} \gamma_i \phi_i^{\otimes n} \left(\rho^{(n)}\right) \quad \dots \dots \dots (a)$$

with the Markovian correlated noise model:

- We note that (a) is a special case of (b):
 - The Markov Chain has M states

•
$$q_{ij} = \delta_{ij}$$

aperiodic but not irreducible

Convergence to equilibrium : so it is not forgetful





- Macchiavello and Palma considered: $q_{ij} = (1-\mu)\gamma_j + \mu\delta_{ij}; 0 \le \mu \le 1$
- Our choice $q_{ii} = \delta_{ii}$ corresponds to $\mu = 1$

(fully correlated noise -- successive actions identical)

Let us evaluate: the product state capacity of the channel

$$\Phi^{(n)}(\rho^{(n)}) = \sum_{i=1}^{M} \gamma_i \phi_i^{\otimes n} \left(\rho^{(n)} \right) \quad \dots \dots \dots (a)$$

• i.e., the classical capacity under the restriction of $C^{(1)}(\Phi)$ product-state inputs

UNIVERSITY OF CAMBRIDGE

Let us start by making a naïve guess:

• Recall : for a memoryless channel ϕ :

[HSW Theorem]

$$C^{(1)}(\phi) = \sup_{\{p_j, \rho_j\}} \chi\left(\{p_j, \phi(\rho_j)\}\right) = \chi^*(\phi) \dots (A)$$

$$p^{(n)}_{av} \to 0 \qquad p^{(n)}_{av} \not \to 0$$

$$as \ n \to \infty \qquad \chi^*(\Phi) \qquad as \ n \to \infty \qquad R \text{ (rate)}$$

• Any $R \leq \chi^*(\Phi)$ is achievable.



For a memoryless channel ϕ :

$$C^{(1)}(\phi) = \sup_{\{p_j, \rho_j\}} \chi(\{p_j, \phi(\rho_j)\}) = \chi^*(\phi) \dots (A)$$

[HSW Theorem]

So in this case, because the channel has <u>M</u> memoryless branches, one might naively expect:

$$C^{(1)}(\Phi) = \min_{1 \le i \le M} \chi^*(\phi_i)$$

= $\min_{1 \le i \le M} \max_{\{p_j, \rho_j\}} \chi(\{p_j, \phi_i(\rho_j)\})$(B)

BUT

• (B) is NOT TRUE ; min $\leftarrow \rightarrow$ max



• Theorem:

The product-state capacity of the long-term memory channel $\Phi^{(n)}(\rho^{(n)}) = \sum_{i=1}^{M} \gamma_i \phi_i^{\otimes n} \left(\rho^{(n)}\right) \quad \dots \dots \quad (a)$ is given by $C^{(1)}(\Phi) = \max_{\{p_j, \rho_j\}} \min_{1 \le i \le M} \chi\left(\{p_j, \phi_i(\rho_j)\}\right)$

whereas are guess was:

$$C^{(1)}(\Phi) = \min_{1 \le i \le M} \max_{\{p_{j}, \rho_{j}\}} \chi(\{p_{j}, \phi_{i}(\rho_{j})\})$$





Sketch of the proof



(Q) Is there any way in which Bob can find out which of the M memoryless branches the qubits have been sent through?

i.e., Can Bob distinguish between the outputs of the different memoryless branches ?





Sketch of the proof contd.



(A) Yes - provided

- Alice adds a preamble to her codewords &
- Bob does a collective measurement on the qubits he receives



Sketch of the proof contd.



Assume: ϕ_i , i = 1, 2, ..., M are all different

Else we do not need to distinguish between all of them

& we can introduce a compound prob. for each

set of identical branches.

• e.g. If
$$\phi_1 = \phi_2 = \tilde{\phi}$$
 prob. $\tilde{\gamma} := \gamma_1 + \gamma_2$



Sketch of the proof contd.



 ϕ_i , i = 1, 2, ..., M are all different

• For each pair $\phi_i, \phi_j: 1 \le i, j \le M$,

-- there exists states $\omega^{(ij)}$, such that

$$\phi_i(\omega^{(ij)}) \neq \phi_j(\omega^{(ij)}),$$





(Q) Can Bob distinguish the outputs of these 2 branches?

- Let ω be a state such that $\phi_1(\omega) \neq \phi_2(\omega)$
- To allow Bob to distinguish between the 2 branches,
 Alice adds a preamble to the input state $\rho_i^{(n)} \leftarrow i$ codeword









• Let us focus on the output of the first m qubits



state of the
first *m* qubits
that Bob
receives
$$\sigma_1^{\otimes m} = [\phi_1(\omega)]^{\otimes m} \text{ with probability } \gamma_1$$

with probability γ_2

(Q) Can Bob do a measurement to distinguish between $\sigma_1^{\otimes m} \& \sigma_2^{\otimes m}$?

(A) Yes. Consider the operator:

$$A^{(m)} = \gamma_1 \sigma_1^{\otimes m} - \gamma_2 \sigma_2^{\otimes m}$$

• Let $\Pi_1^{(m)}$: orthogonal projection onto the non-negative eigenspace of $A^{(m)}$

and

$$\Pi_2^{(m)} = \mathbf{1}^{(m)} - \Pi_1^{(m)}$$

UNIVERSITY OF CAMBRIDGE

• Let Bob does a projective measurement (a la Helstrom) described by the operators $\prod_{1}^{(m)} \& \prod_{2}^{(m)}$ on the state $\sigma_{j}^{\otimes m}$ that he receives: $\sigma_{j}^{\otimes m}$, j = 1, 2 with probs. $\gamma_{1} \& \gamma_{2}$ resply.

• For *m* large enough, by using Helstrom's strategy, Bob can indeed distinguish between $\sigma_1^{\otimes m} & \sigma_2^{\otimes m}$

-- with arbitrarily low probability of error.

Thus he can determine which memoryless branch the qubits have come through!



- Bob determines which branch the input has come through
 - from Bob's point of view : problem reduces to

decoding codewords sent through a memoryless channel

• So now he can do the appropriate decoding operation

on the remaining output state to infer Alice's message







Alice does not know what *i* is (no feedback)



Now one can understand why:

$$C^{(1)}(\Phi) = \max_{\{p_j, \rho_j\}} \min_{1 \le i \le M} \chi(\{p_j, \phi_i(\rho_j)\})$$
(A)

and

- For a memoryless channel ϕ_i

$$C^{(1)}(\phi_{i}) = \max_{\{p_{j}, \rho_{j}\}} \chi(\{p_{j}, \phi_{i}(\rho_{j})\})$$

- The input ensemble for which the max is achieved
 optimal signal ensemble
- IF Alice knew i apriori then she could encode her messages using the optimal signal ensemble for ϕ_i & obtain (B)

BUT Alice does NOT know i apriori.



- For ϕ_i , for any given input ensemble $\{p_{j}, \rho_{j}\}$ Max. amount of classical info that $= \chi(\{p_{j}, \phi_i(\rho_j)\})$
 - In our channel there are M memoryless branches:
 - . Max. amount of classical info that can be sent through it (for any given input ensemble $\{p_{j}, \rho_{j}\}$):

$$= \min_{1 \le i \le M} \chi \left(\{ p_{j,} \phi_i(\rho_j) \} \right)$$

• & this
$$\implies$$
 any rate $R \leq \max_{\{p_j, \rho_j\}} \min_{1 \leq i \leq M} \chi(\{p_j, \phi_i(\rho_j)\})$
is achievable
• & this $\implies C^{(1)}(\Phi) \geq \max_{\{p_j, \rho_j\}} \min_{1 \leq i \leq M} \chi(\{p_j, \phi_i(\rho_j)\})$



Theorem:
$$C^{(1)}(\Phi) = \max_{\{p_j, \rho_j\}} \min_{1 \le i \le M} \chi(\{p_j, \phi_i(\rho_j)\})$$

• We have proved Direct part (achievability) $C^{(1)}(\Phi) \ge \max_{\{p_j, \rho_j\}} \min_{1 \le i \le M} \chi(\{p_j, \phi_i(\rho_j)\})$

We also need to prove that : any rate

 $R \geq \max_{\{p_j, \rho_j\}} \min_{1 \leq i \leq M} \chi \left(\{p_j, \phi_i(\rho_j)\} \right) \quad \text{is not achievable}$

Weak Converse



- Ingredients needed to prove the Weak Converse:
 - Holevo bound
 - Subadditivity of the von Neumann entropy
 - Fano's inequality



- Recall: The quantum channel with Markovian correlated noise is forgetful IF the Markov Chain is
 (1) irreducible and (2) aperiodic
- The "not-forgetful" channel that we considered was aperiodic but not irreducuble
- Another example of a "not-forgetful" channel is one for which the Markov Chain is : irreducible but not aperiodic (i.e., memory governed by a periodic Markov Chain)
- E.g. 2-state Markov Chain :



• Transition Matrix $Q = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

irreducible, periodic (period=2).



• 2 states of the Markov Chain corrs. to 2 single qubit channels φ_1, φ_2 which act alternatively on successive inputs

$$\Phi^{(n)}(\rho^{(n)}) = \frac{1}{2} \begin{bmatrix} \phi_1 \otimes \phi_2 \otimes \phi_1 \otimes \dots + \phi_2 \otimes \phi_1 \otimes \phi_2 \otimes \dots \end{bmatrix} \begin{pmatrix} \rho^{(n)} \end{pmatrix}$$
n times
n times

In this case,

$$C^{(1)}(\phi_i) = \sup_{\{p_j, \rho_j\}} \frac{1}{2} \sum_{i=1}^2 \chi(\{p_j, \phi_i(\rho_j)\}) = \frac{1}{2} \sum_{i=1}^2 \chi^*(\phi_i)$$

= average of the Holevo capacities of the individual channels

• Similarly once can consider a channel where the underlying Markov Chain has a period L > 2



 ND and Tony Dorlas, Coding Theorem for a Class of Quantum Channels with Long-Term Memory, J. Phys. A: Math. Theor. 40, 8147-8164 (2007).

- ND, T.C.Dorlas and Y.M. Suhov, Entanglement Assisted Classical Capacity of a Class of Quantum Channels with Long-Term Memory,' Quantum Information Processing, 7, no. 6, 251--262 (2008).
- I._Bjelakovic and H. Boche,
 On quantum capacity of compound channels, arXiv:0808.1007.



• ND and Tony Dorlas,

Classical capacity of quantum channels with general Markovian correlated noise,

```
J.Stat.Phys. 134, pp 1173-1195, (2009).
```

 I._Bjelakovic and H. Boche,
 Entanglement transmission capacity of compound channels, arXiv:0904.3011.



Information Spectrum Method

- S.Verdu and T.S.Han, IEEE Trans. Inf. Theory 40, pp. 1147-1157 (1994);
- T.S.Han,

Information-Spectrum Methods in Information Theory (Springer-Verlag, 2002);

- H.Nagaoka and M.Hayashi, IEEE Trans. Inf. Theory {\bf 53}, 534--549 (2007)
- F.Buscemi and N.Datta, The quantum capacity of channels with arbitrarily correlated noise, IEEE Trans. inf. Th. 56, 1447-1460 (2010),

UNIVERSITY OF CAMBRIDGE



