

Defects on Nematic Shells

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Summary

Nematic Shells

Mathematical Model

Levi-Civita Parallel Transport

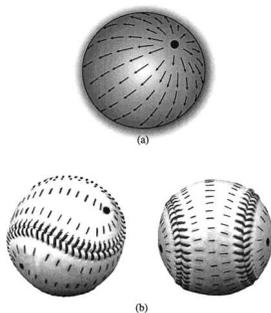
Surface Nematic Elasticity

Valence Control

Nematic Shells

A *thin* film of nematic liquid crystal deposited on a *colloidal* particle is a *nematic shell*, which is treated here as a *two-dimensional* order texture.

tetravalent tennis balls



Slide 2 *Defects* in the order texture are potential *hot spots* where ligands may adhere bridging one particle to another in a *metamaterial* structure. NELSON (2002)

Mathematical Model

Molecules are envisaged as *ribbons* lying flat on a closed surface \mathcal{S} and freely gliding on it.

planar degenerate anchoring

This picture is perhaps too naive: a degenerate planar anchoring *disluted* in a thin layer is a more realistic microscopic picture.

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- ℓ molecular director
- \mathcal{S} closed orientable surface
- ν outer unit normal
- $\ell \cdot \nu \equiv 0$

order tensor

$$\mathbf{Q} := \langle \ell \otimes \ell - \frac{1}{2} \mathbf{P} \rangle$$

$$\mathbf{P} := \mathbf{I} - \nu \otimes \nu$$

$\langle \cdot \rangle$ ensemble average

\mathbf{Q} is a fully biaxial tensor

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$$\mathbf{Q} = \lambda (\mathbf{n} \otimes \mathbf{n} - \mathbf{n}_\perp \otimes \mathbf{n}_\perp) \quad \mathbf{n}_\perp := \nu \times \mathbf{n}$$

Defects occur whenever $\lambda = 0$ and so $\mathbf{Q} = \mathbf{0}$

order bounds

$$0 \leq \lambda \leq \frac{1}{2}$$

q-representation

$$\mathbf{Q} = q_1(\mathbf{e} \otimes \mathbf{e} - \mathbf{e}_\perp \otimes \mathbf{e}_\perp) + q_2(\mathbf{e} \otimes \mathbf{e}_\perp + \mathbf{e}_\perp \otimes \mathbf{e})$$

$\mathbf{e}, \mathbf{e}_\perp$ tangent unit vector fields

$$\mathbf{e}_\perp := \boldsymbol{\nu} \times \mathbf{e} \quad \mathbf{e} \cdot \boldsymbol{\nu} \equiv 0$$

$$\mathbf{n} = \cos \varphi \mathbf{e} + \sin \varphi \mathbf{e}_\perp$$

$$\mathbf{n}_\perp = -\sin \varphi \mathbf{e} + \cos \varphi \mathbf{e}_\perp$$

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$$\cos 2\varphi = \frac{q_1}{\sqrt{q_1^2 + q_2^2}} \quad \sin 2\varphi = \frac{q_2}{\sqrt{q_1^2 + q_2^2}}$$

Levi-Civita Parallel Transport

For *planar* director fields, the *topological charge* of a defect is directly related to the *winding number*.

For *non-flat* fields, an intrinsic distortion is due to the curvature of the underlying surface.

Along a curve \mathcal{C} on \mathcal{S}

$$\boldsymbol{\nu}' = \boldsymbol{\Omega}_\parallel \times \boldsymbol{\nu}$$

$$\boldsymbol{\Omega}_\parallel = \boldsymbol{\nu} \times (\nabla_s \boldsymbol{\nu}) \mathbf{t}$$

' arc-length derivative

$\nabla_s \boldsymbol{\nu}$ surface gradient of $\boldsymbol{\nu}$

\mathbf{t} unit tangent to \mathcal{C}

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principal curvatures

$$\nabla_s \boldsymbol{\nu} = \sigma_1 \mathbf{e}_1 \otimes \mathbf{e}_1 + \sigma_2 \mathbf{e}_2 \otimes \mathbf{e}_2$$

$H := \frac{1}{2}(\sigma_1 + \sigma_2)$ mean curvature

$K := \sigma_1 \sigma_2$ Gaussian curvature

intrinsic distortion

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Transporting a tangent unit vector \mathbf{u}

$$\mathbf{u}' = \boldsymbol{\Omega}_\parallel \times \mathbf{u}$$

$$\Delta \vartheta_{\mathcal{C}} = \int_{\mathcal{S}_{\mathcal{C}}} K da$$

$\mathcal{S}_{\mathcal{C}}$ surface on \mathcal{S} enclosed by \mathcal{C}

topological charge

$$\mathbf{n}' = \boldsymbol{\Omega} \times \mathbf{n}$$

$$\boldsymbol{\Omega} = \boldsymbol{\Omega}_\nu \boldsymbol{\nu} + \boldsymbol{\Omega}_\parallel$$

$$\boldsymbol{\Omega}_\nu := \mathbf{n}' \cdot \mathbf{n}_\perp \quad \mathbf{n}_\perp := \boldsymbol{\nu} \times \mathbf{n}$$

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For the field \mathbf{n}

$$2\pi m_{\mathbf{n}} = \oint_{\mathcal{C}} \mathbf{n}' \cdot \mathbf{n}_\perp ds + \int_{\mathcal{S}_{\mathcal{C}}} K da \quad \forall \mathcal{C}$$

$2m_{\mathbf{n}} \in \mathbb{Z}$ \mathbf{n} and $-\mathbf{n}$ give the same \mathbf{Q}

$$2\pi m_{\mathbf{n}} = \oint_{\mathcal{C}} \mathbf{n}' \cdot \mathbf{n}_{\perp} ds + \int_{\mathcal{S}_{\mathcal{C}}} K da$$

For another tangent unit vector field \mathbf{e}

$$2\pi m_{\mathbf{e}} = \oint_{\mathcal{C}} \mathbf{e}' \cdot \mathbf{e}_{\perp} ds + \int_{\mathcal{S}_{\mathcal{C}}} K da$$

In the q -representation

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$$\mathbf{Q} = q_1(\mathbf{e} \otimes \mathbf{e} - \mathbf{e}_{\perp} \otimes \mathbf{e}_{\perp}) + q_2(\mathbf{e} \otimes \mathbf{e}_{\perp} + \mathbf{e}_{\perp} \otimes \mathbf{e})$$

$$m_{\mathbf{n}} = m_{\mathbf{e}} + \frac{1}{4\pi} \oint_{\mathcal{C}} \frac{q_1 q_2' - q_1' q_2}{q_1^2 + q_2^2} ds$$

$$m_{\mathbf{n}} = m_{\mathbf{e}} + \frac{1}{2\pi} \oint_{\mathcal{C}} \varphi' ds$$

It is easily shown that the topological charge is **additive** in the contour enclosure.

total charge

If \mathcal{S} splits into the union of N patches, each containing a defect of \mathbf{n}

$$2\pi m_i = \oint_{\mathcal{C}_i} \mathbf{n}' \cdot \mathbf{n}_{\perp} ds + \int_{\mathcal{S}_i} K da \quad \mathcal{S} = \bigcup_{i=1}^N \mathcal{S}_i$$

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$$2\pi \sum_{i=1}^N m_i = \int_{\mathcal{S}} K da$$

Euler characteristics

$$\int_{\mathcal{S}} K da = 2\pi \chi(\mathcal{S}) = 2(1 - g(\mathcal{S}))$$

$$\chi(\mathcal{S}) = F - E + V$$

F number of faces
 E number of edges of **any** tessellation of \mathcal{S}
 V number of vertices
 $g(\mathcal{S})$ number of **handles**

\mathcal{S}	g	χ
sphere	0	2
torus	1	0

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Poincaré Theorem $\sum_{i=1}^N m_i = \chi(\mathcal{S})$

Surface Nematic Elasticity

The measure of surface **order distortion** is $\nabla_s \mathbf{Q}$

$$\iota_1 := Q_{ij;k} Q_{ij;k} = |\nabla_s \mathbf{Q}|^2$$

$$\iota_2 := Q_{ij;k} Q_{ik;j}$$

$$\iota_3 := Q_{ij;j} Q_{ik;k} = (\text{div}_s \mathbf{Q})^2$$

; surface derivative

null-Lagrangian ?

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$$\int_{\mathcal{S}} (\iota_2 - \iota_3) da = - \int_{\mathcal{S}} K (\text{tr} \mathbf{Q}^2) da$$

elastic free-energy density

$$f_e := c_1 \iota_1 + c_2 \iota_2 + c_{24}(\iota_2 - \iota_3)$$

$$= \frac{1}{2}(k_1 + k_3)Q_{ij;k}Q_{ij;k} + \frac{1}{2}(k_1 - k_3)Q_{ij;k}Q_{ik;j} - \frac{1}{2}k_{24}K \operatorname{tr} \mathbf{Q}^2$$

$$k_1 = c_1 + c_2 \quad k_3 = c_1 - c_2 \quad k_{24} = c_{24}$$

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one-constant approximation

$$k_1 = k_3 = \frac{1}{2}k$$

$$f_e = \frac{1}{2}(k|\nabla_s \mathbf{Q}|^2 - k_{24}K(\operatorname{tr} \mathbf{Q}^2))$$

$$-1 \leq \frac{k_{24}}{k} \leq 1$$

internal potential

$$f_p(\mathbf{Q}) = \frac{A}{2} \operatorname{tr} \mathbf{Q}^2 + \frac{C}{4} (\operatorname{tr} \mathbf{Q}^2)^2$$

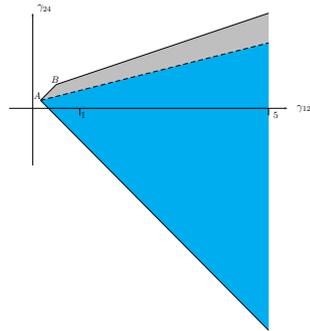
$$A = A_0 \frac{T - T_c}{T_c} \quad A_0 > 0$$

T temperature
 T_c critical temperature

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Ericksen's inequalities

f_e is **positive definite** in the admissible region



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$$\gamma_{12} = \frac{2c_1 + c_2}{2c_1 - c_2 - c_{24}}$$

$$\gamma_{24} = \frac{c_{24}}{2c_1 - c_2 - c_{24}}$$

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condensation order

$$\lambda_c = \sqrt{\frac{A_0}{2C} \left(\frac{T_c - T}{T_c} \right)} \quad T < T_c$$

nematic coherence length

$$\xi_0 := \sqrt{\frac{k}{A_0}}$$

reduced temperature

$$t := \frac{T - T_c}{T_c}$$

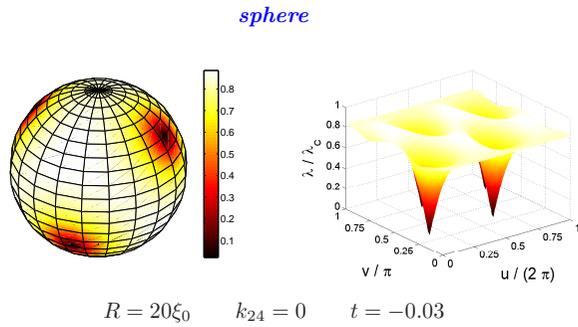
free-energy functional

$$\mathcal{F} := \frac{1}{2} \int_{\mathcal{D}} \left\{ k|\nabla_s \mathbf{Q}|^2 + (A - k_{24}K)(\operatorname{tr} \mathbf{Q}^2) + \frac{C}{2} (\operatorname{tr} \mathbf{Q}^2)^2 \right\} da$$

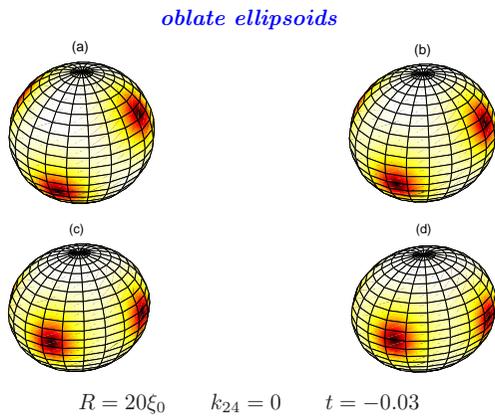
Valence Control

Both k_{24} and K appear capable of controlling both number and location of defects

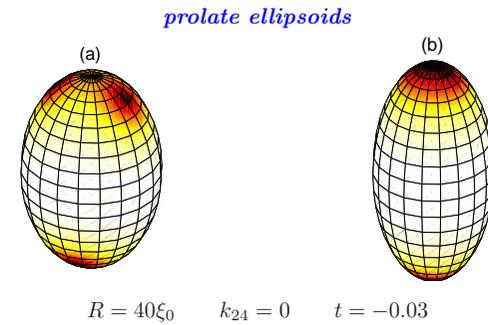
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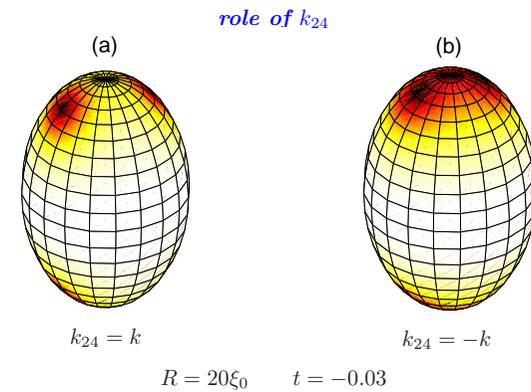


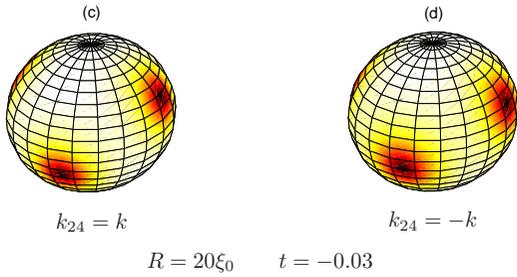
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Defects migrate towards the region the largest K , to the point of possibly *merging*. This confirms simulations of [BATES, SKAČEJ & ZANNONI \(2010\)](#), which however did not show defect merging.

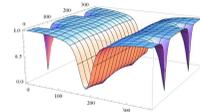
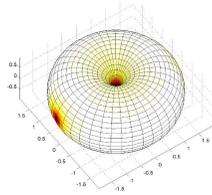
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For $K > 0$ a *negative* k_{24} enhances the migration of defects towards the regions with the largest K .



torus

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Credits

Co-workers

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Discussion

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More information



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