



*The Abdus Salam  
International Centre for Theoretical Physics*



**2246-15**

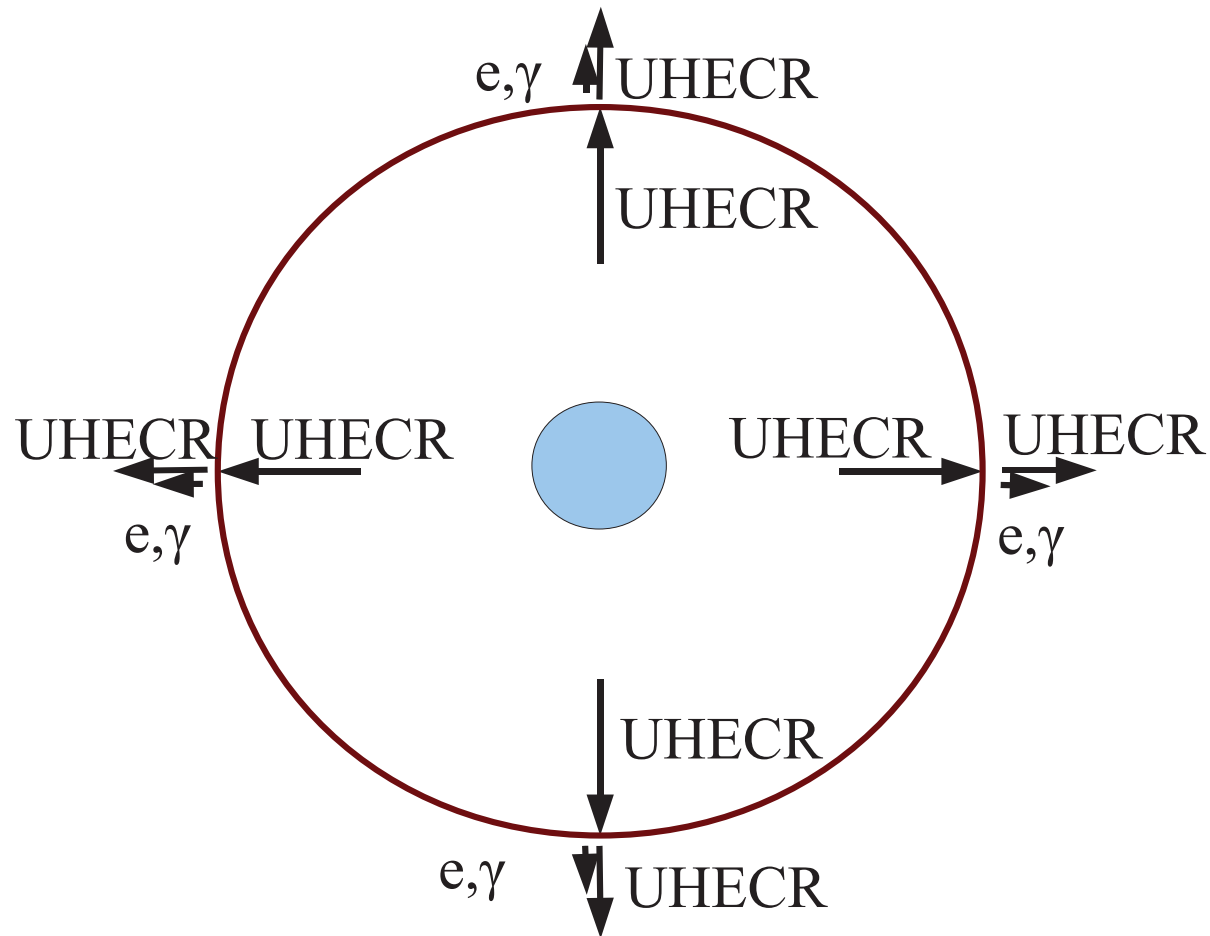
**Workshop on Cosmic Rays and Cosmic Neutrinos: Looking at the  
Neutrino Sky**

*20 - 24 June 2011*

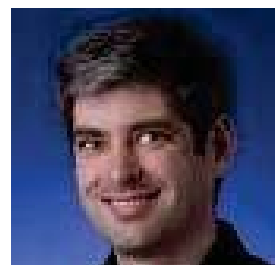
**Probing local UHE sources using CR nuclei and gamma-rays**

Andrew TAYLOR  
*University of Geneva  
Switzerland*

# Probing Local UHECR Sources with Nuclei + Gamma-Rays



# Probing Local UHECR Sources with Nuclei + Gamma-Rays



# Talk Structure

## Using Nuclei

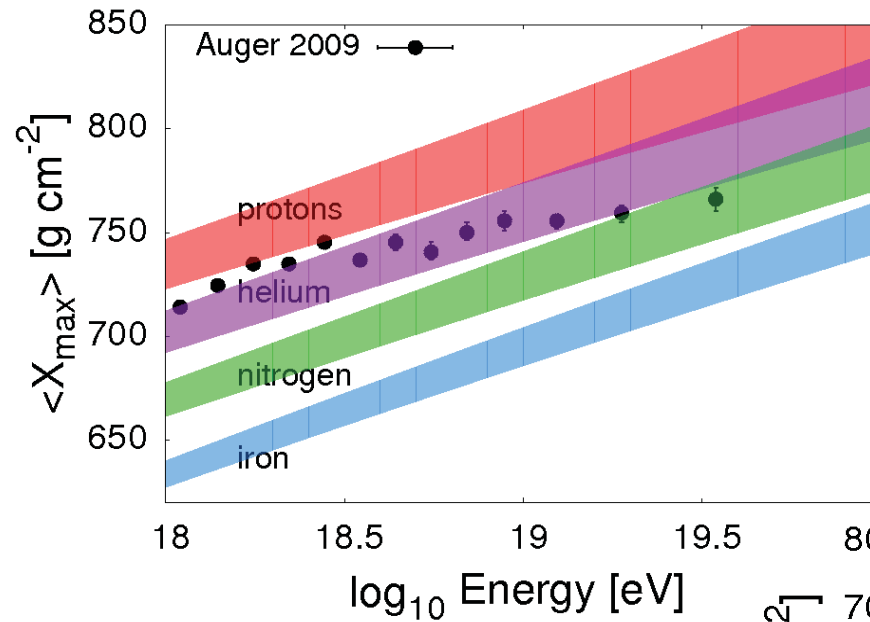
- 1) Measurements of UHECR by the PAO
  - 2) Implications for UHECR source requirements
- .....

## Using Gamma-Rays

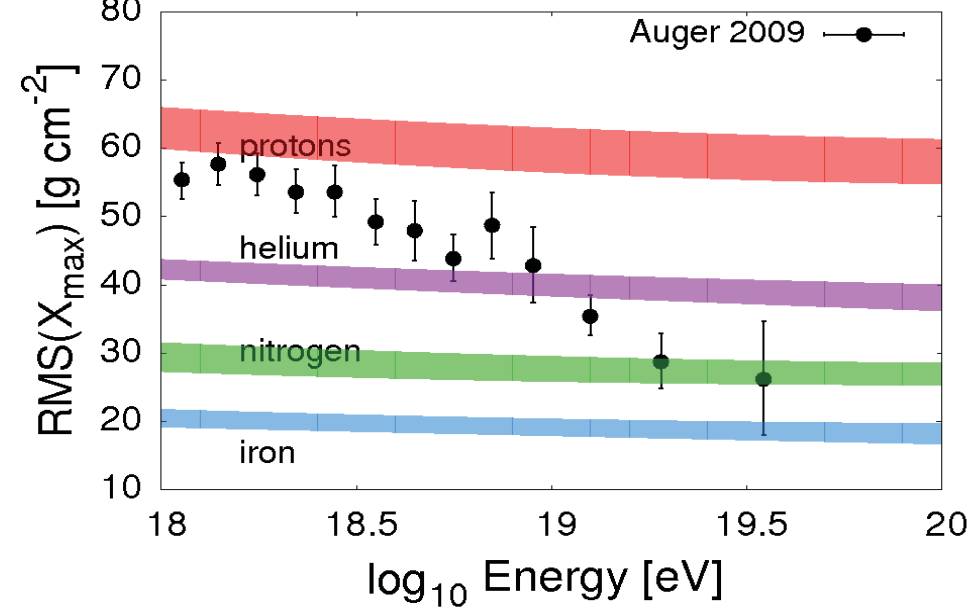
- 3) Their birth + death in regions surrounding the source
- 4) What current PAO Gamma-Ray limits already tell us about the sources

# **Part 1: Nuclei**

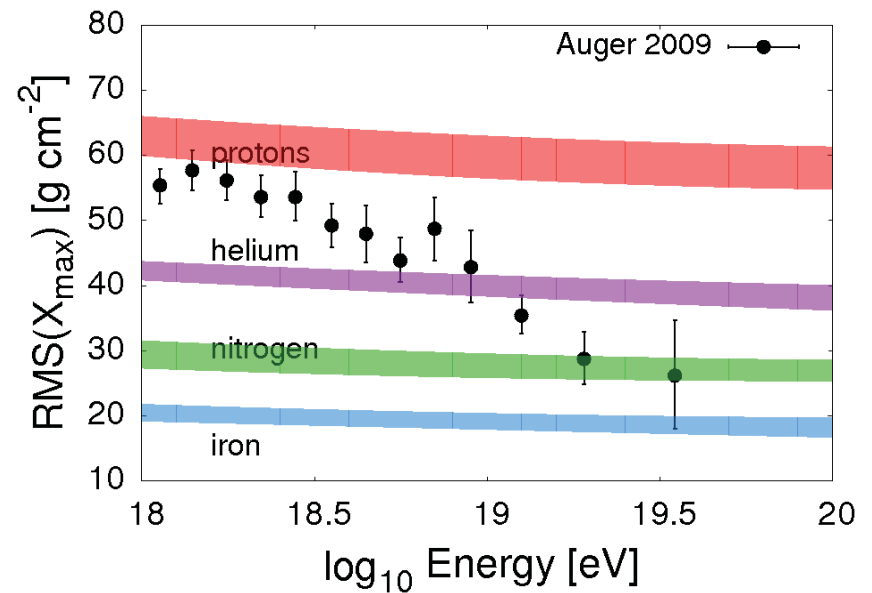
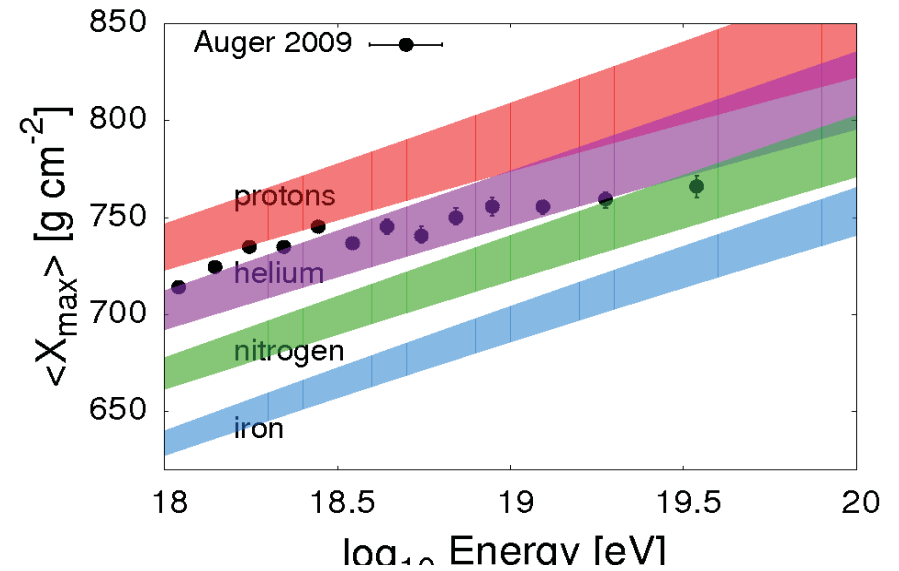
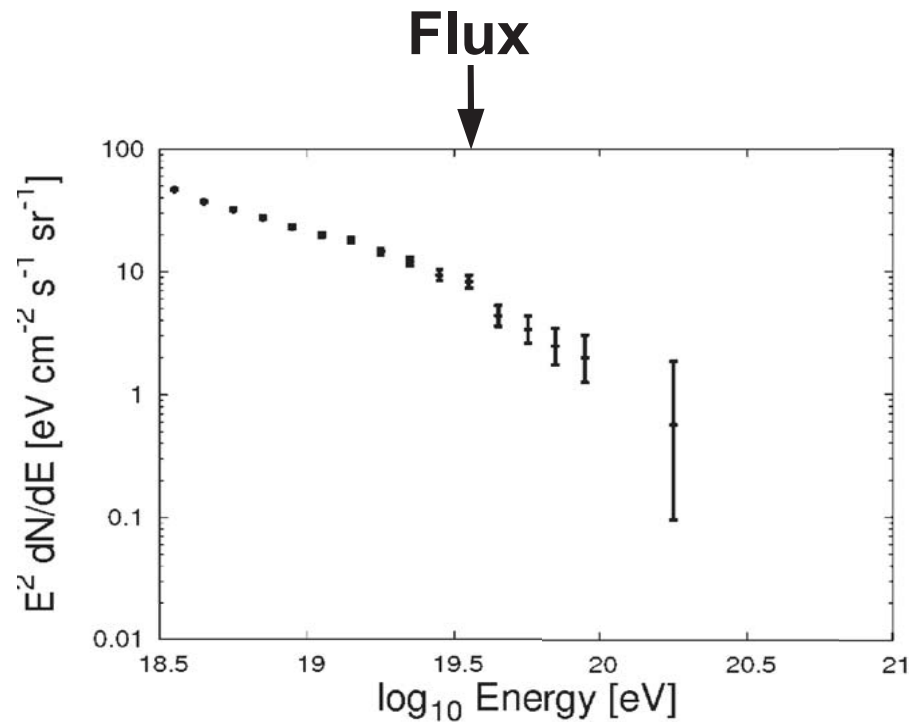
# MEASUREMENTS of UHECR by the PAO



$\leftarrow \langle X_{\max} \rangle$   
 $\text{RMS}(X_{\max})$

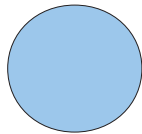
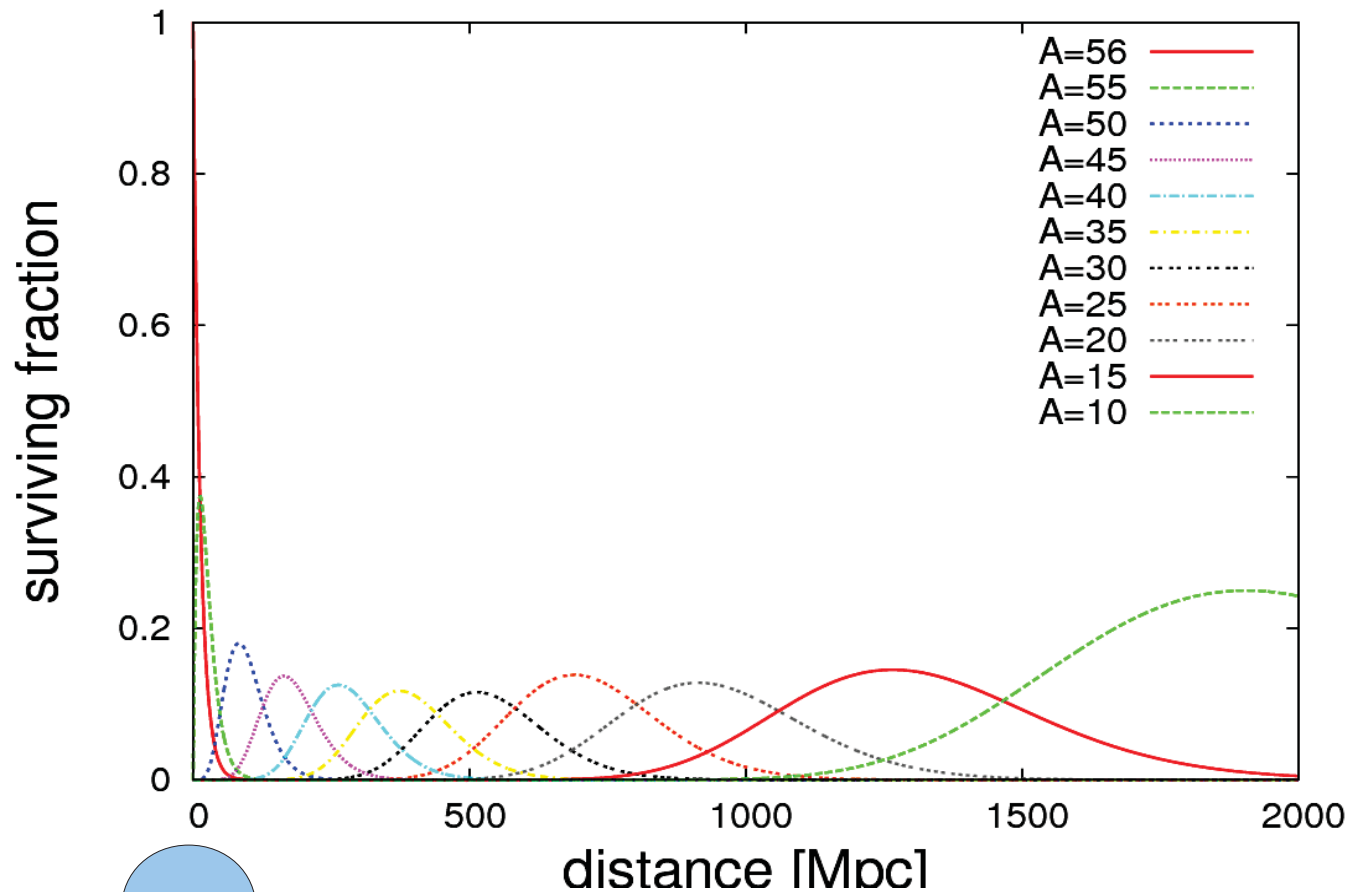


# MEASUREMENTS of UHECR by the PAO



# Nuclei Propagation Away from their Source + their Transmutation

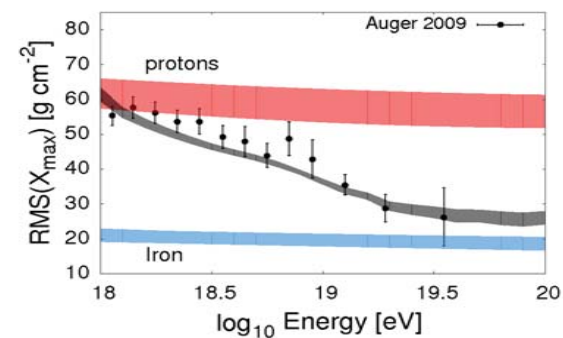
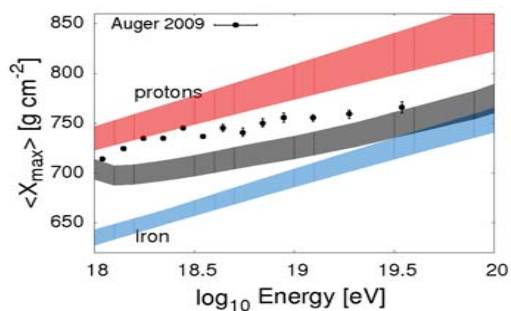
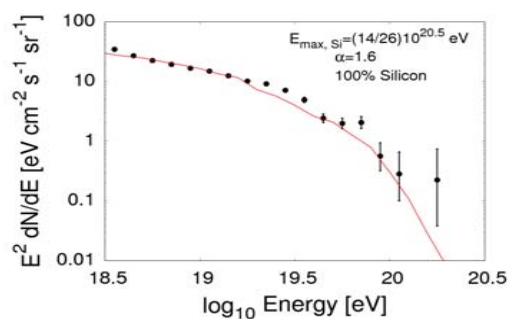
## Lorentz factor of nuclei ~conserved



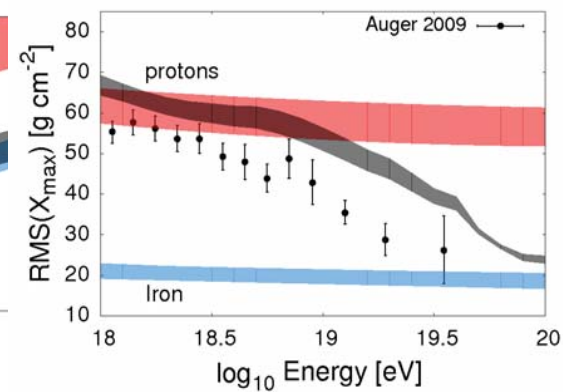
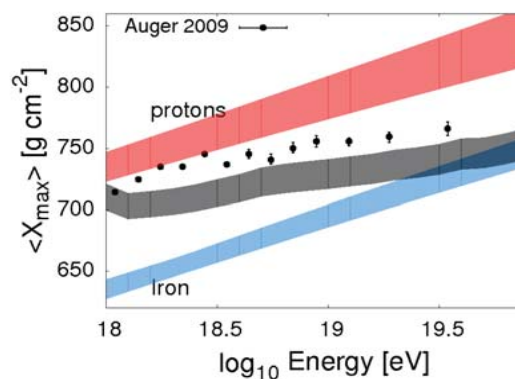
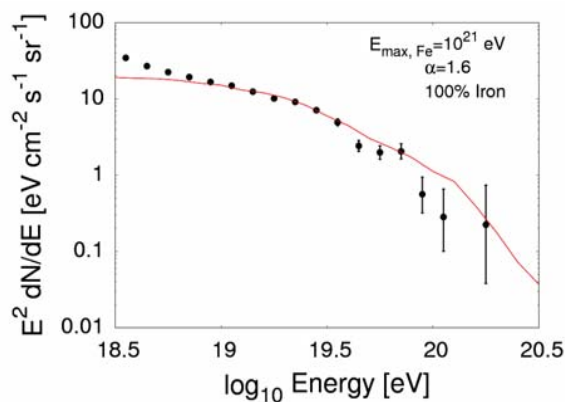


# Nuclei Propagation Away from their Source + their Transmutation

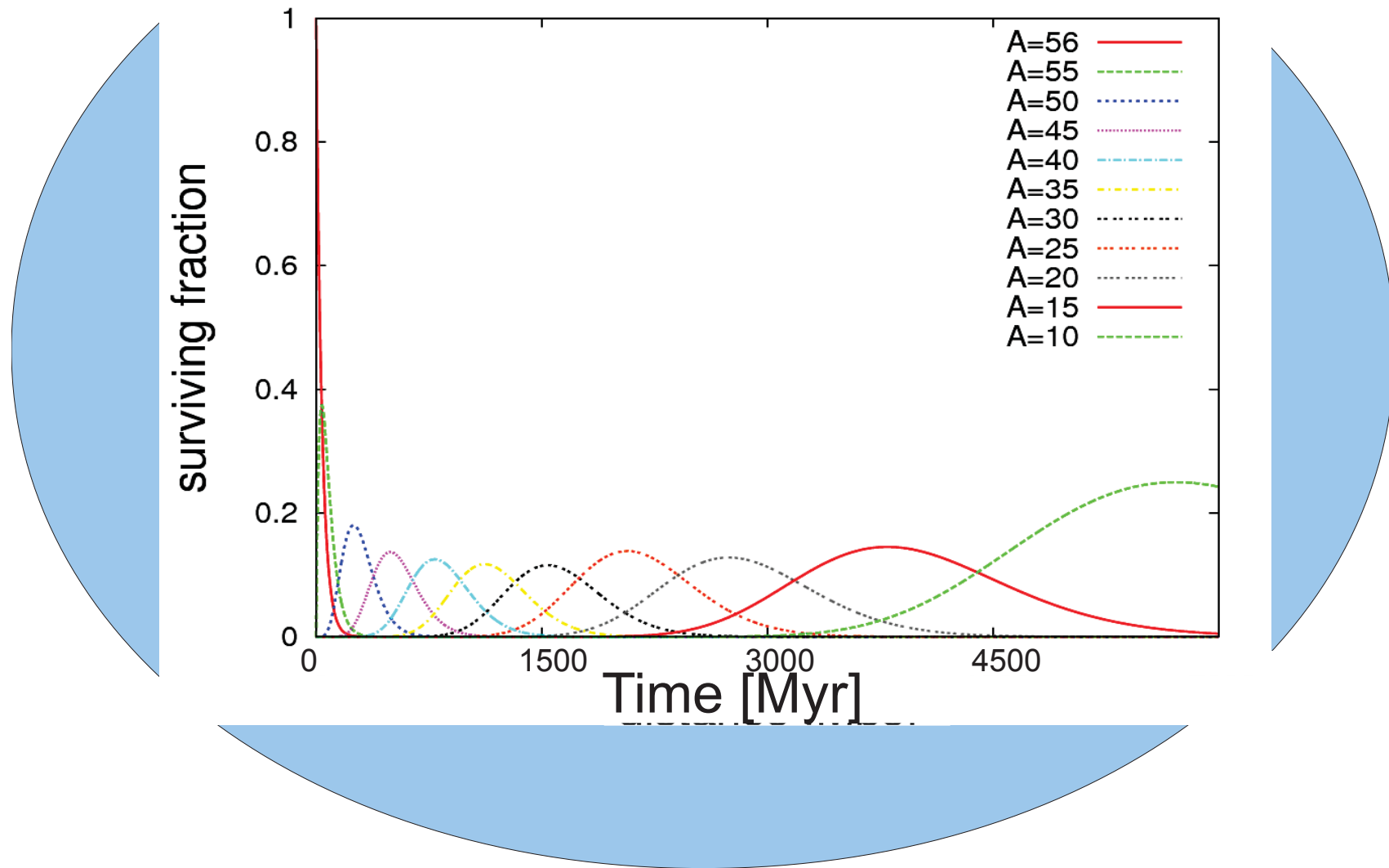
Silicon only?



Iron only?

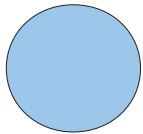


# Nuclei Propagation Within their Source + their Transmutation



# IMPLICATIONS for UHECR Sources

$$f = \frac{t_{\text{trap}}}{t_{\text{int.}}^{\text{CR}\gamma}}$$



$$t_{\text{int.}}^{\text{CR}\gamma} \approx \frac{1}{n_{\gamma} \sigma_{\text{CR}\gamma} c}$$

$$n_{\gamma} = \frac{L_{\gamma}}{c 4\pi R^2 \epsilon_{\gamma}}$$

$$t_{\text{trap}} \approx \frac{R^2}{2D} = \frac{3R^2}{2R_{\text{Larmor}}}$$

$$f^{\text{CR}\gamma} = \frac{3L_{\gamma} \sigma_{\text{CR}\gamma} Z B}{8\pi \epsilon_{\gamma} E_{\text{CR}}}$$

# IMPLICATIONS for UHECR Sources

$$f^{\text{CR}\gamma} = \frac{3L_{\gamma}\sigma_{\text{CR}\gamma}ZB}{8\pi\epsilon_{\gamma}E_{\text{CR}}} = \frac{s_1}{s_2}$$

Photo-disintegration threshold:

$$2E_{\text{CR}}\epsilon_{\gamma} > Am_p c^2 E_{\text{bind.}}, \text{ where } m_p c^2 E_{\text{bind.}} = 10^{16} \text{ eV}^2$$

Since,	$L_{\gamma}[10^{40} \text{ erg s}^{-1}] = 2 \times 10^{41} \text{ eV cm}^{-1}$
	$\sigma_{\text{CR}\gamma}[\text{A mb}] = A \times 10^{-27} \text{ cm}^2$
	$B[1 \text{ G}] = 300 \text{ eV cm}^{-1}$

$$\frac{L_{\gamma}\sigma_{\text{CR}\gamma}B}{A} = 6 \times 10^{16} \text{ eV}^2, \text{ ergo.... } f^{\text{CR}\gamma} = 50 \frac{Z}{26}$$

# IMPLICATIONS for UHECR Sources

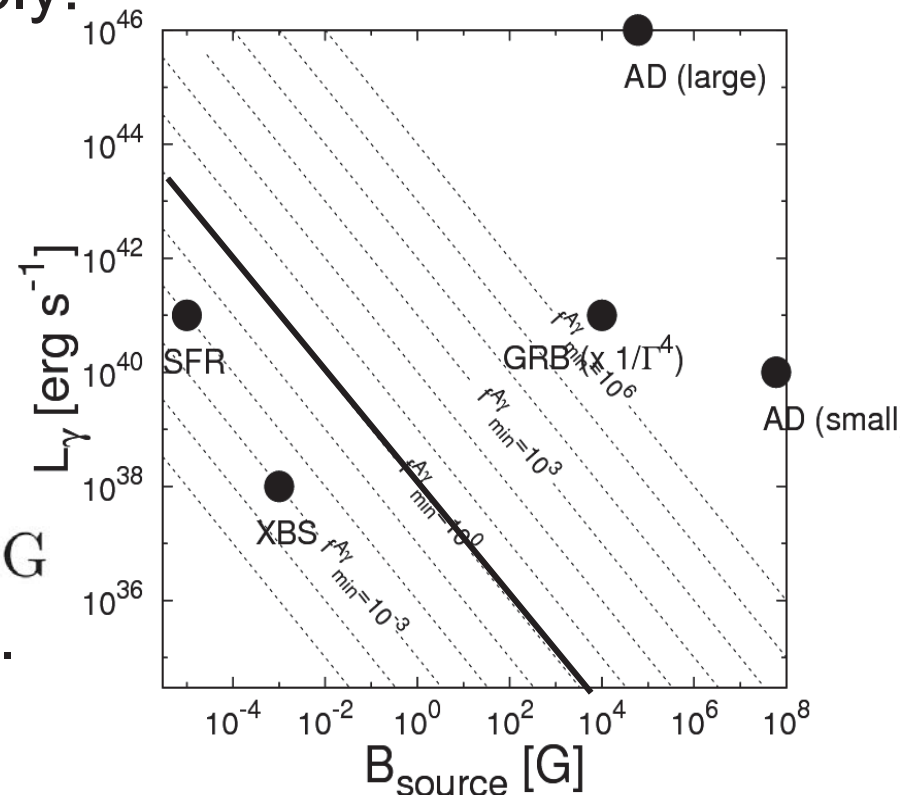
Since, 
$$\frac{L_{\gamma}^{\text{Edd.}} \sigma_{\text{CR}\gamma} B^{\text{Edd.}}}{A} = 4 \times 10^{23} \left( \frac{M}{M_{\odot}} \right)^{1/2} \text{ eV}^2$$

Only heavily sub-Eddington  
power objects need apply!

If magnetic + photon  
luminosity are in  
equipartition:

$$L_{\gamma} \approx \beta R^2 B^2$$

Requiring,  $B < 4 \times 10^{-5} \text{ G}$   
to ensure safe passage.



# IMPLICATIONS for UHECR Sources

Since, 
$$\frac{L_{\gamma}^{\text{Edd.}} \sigma_{\text{CR}\gamma} B^{\text{Edd.}}}{A} = 4 \times 10^{23} \left( \frac{M}{M_{\odot}} \right)^{1/2} \text{ eV}^2$$

Only heavily sub-Eddington  
power objects need apply!

If magnetic + photon luminosity  
are in equipartition:

$$L_{\gamma} \approx \beta R^2 B^2$$

**WARNING:**

$$R_{\text{Larmor}}(10^{20} \text{ eV Fe}) \approx 0.1 \text{ kpc}$$

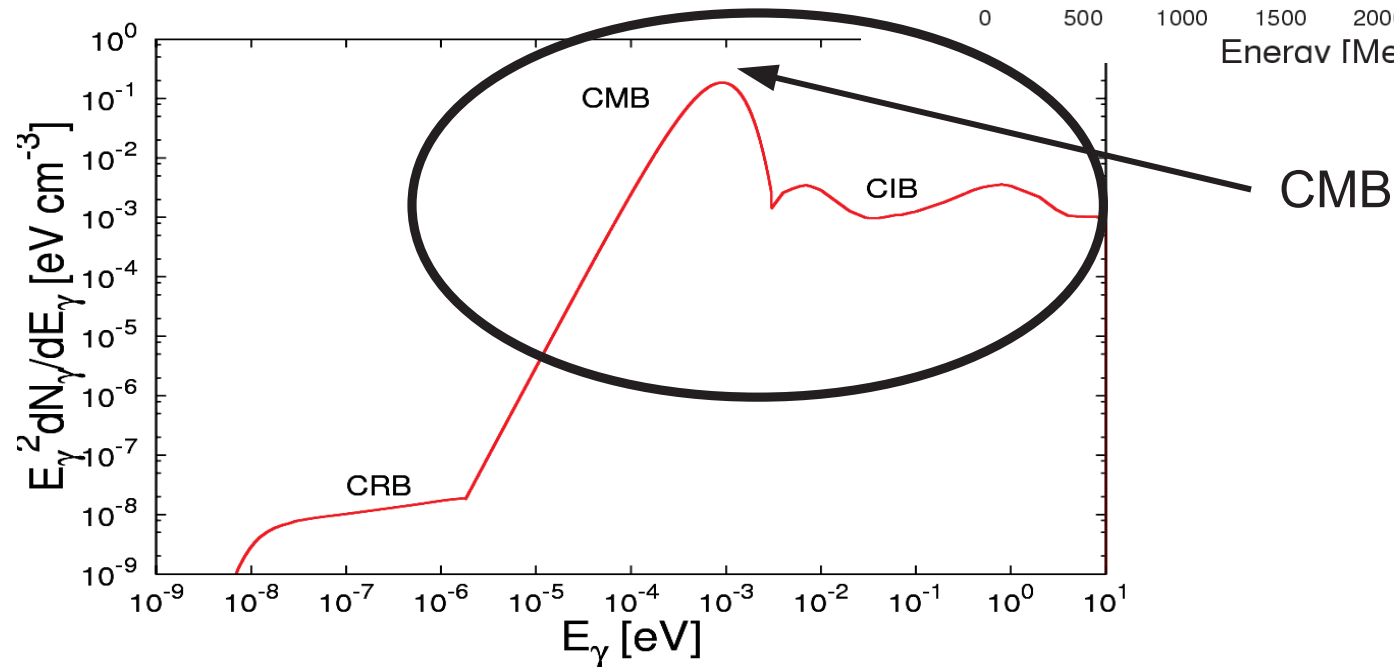
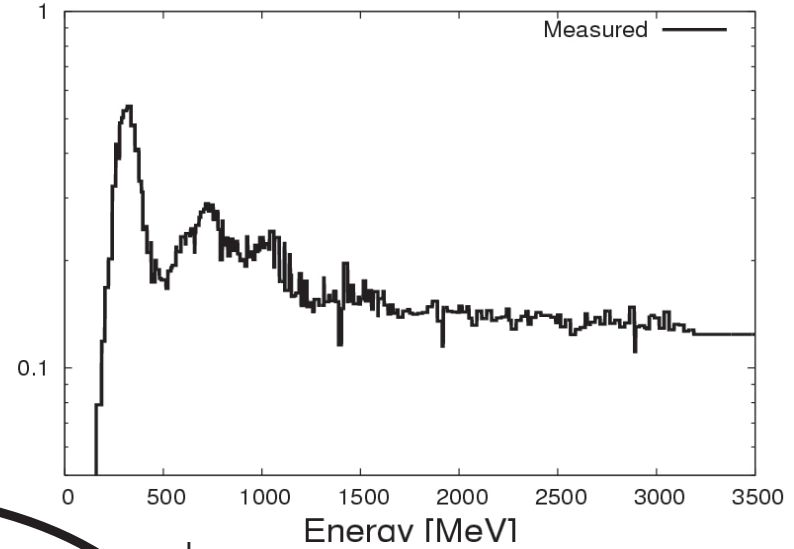
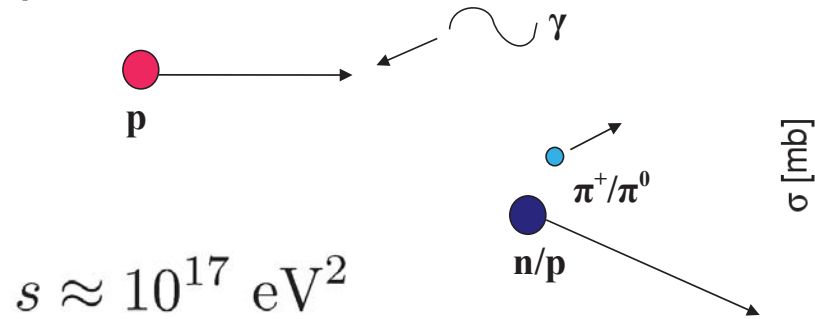
$$t_{\text{acc}} \approx R_{\text{Larmor}} / c\beta^2$$

Requiring,  $B < 4 \times 10^{-5} \text{ G}$

## **Part 2: Gamma-Rays**

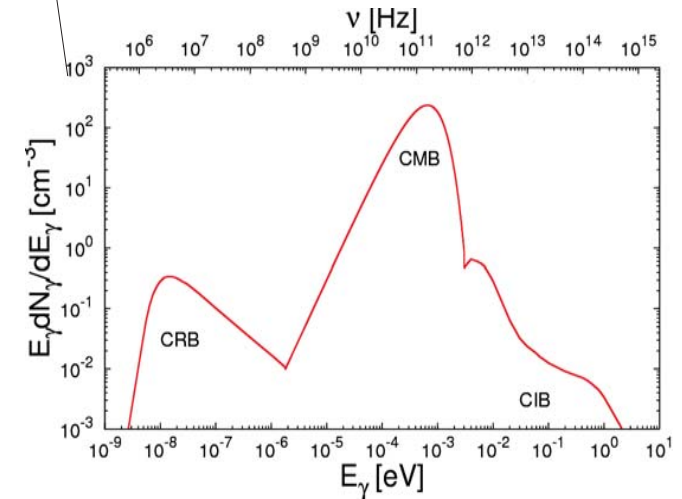
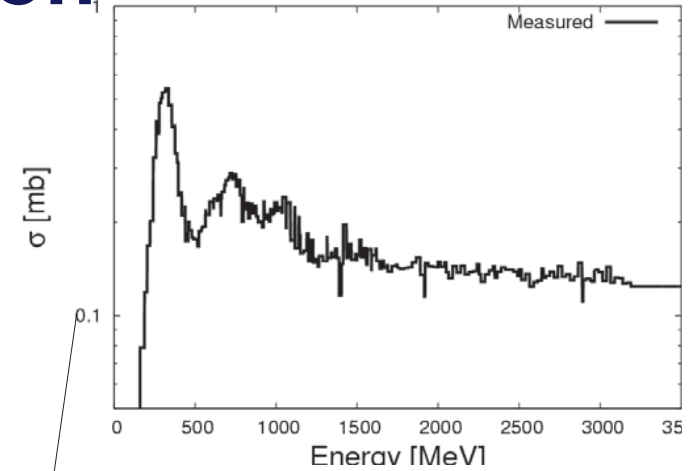
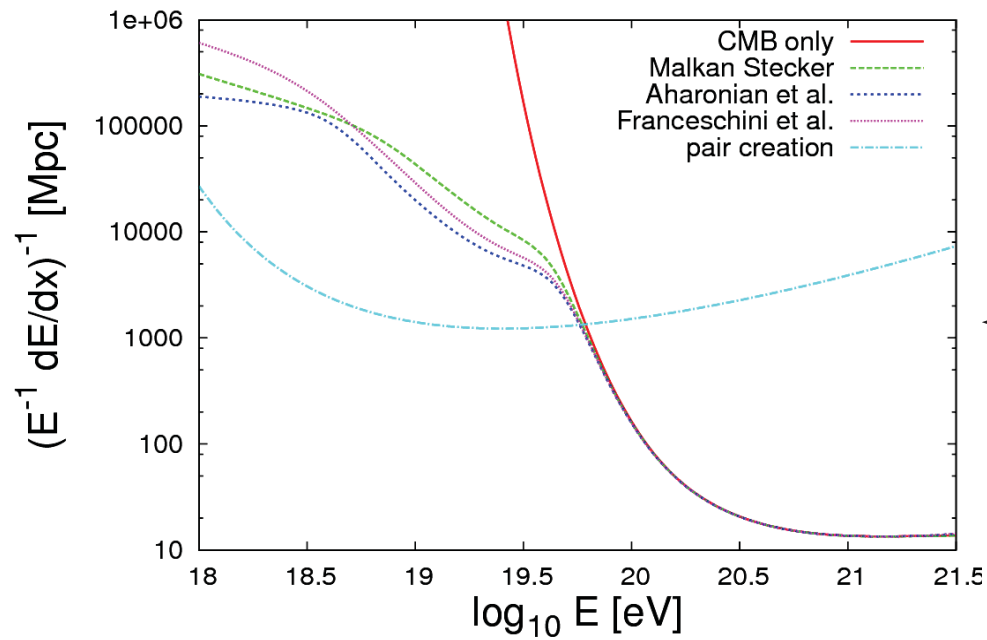
# Gamma-Ray Production

protons

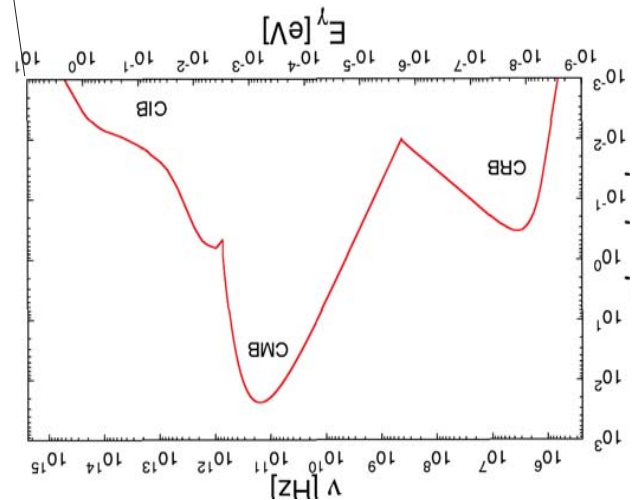
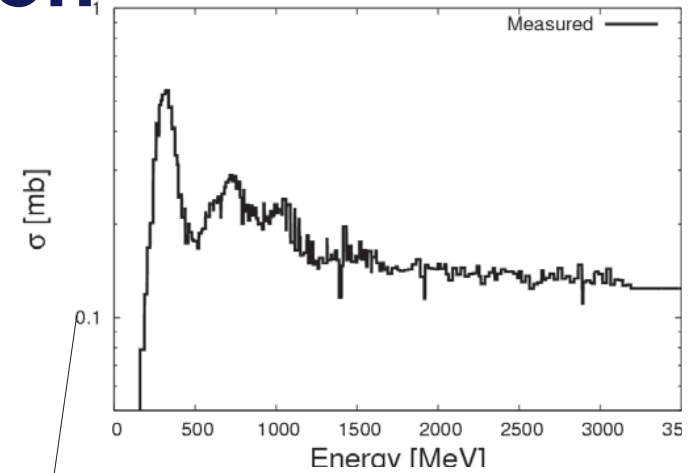
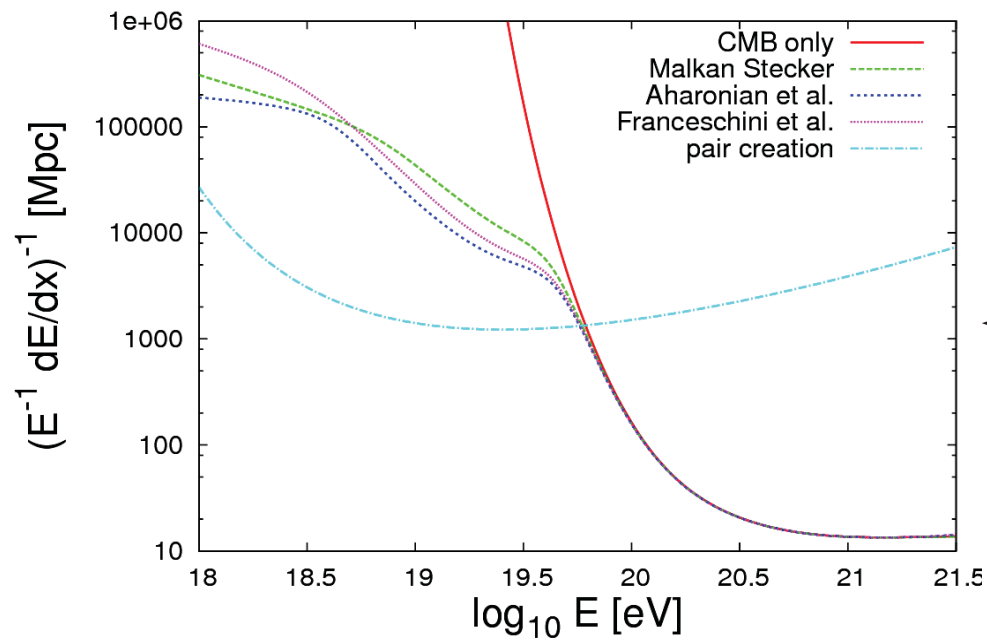




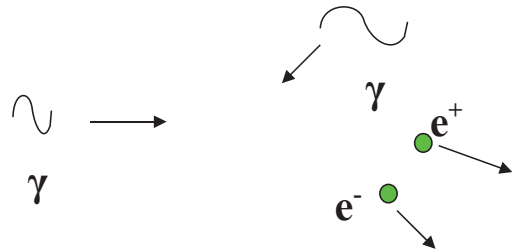
# Gamma-Ray Production



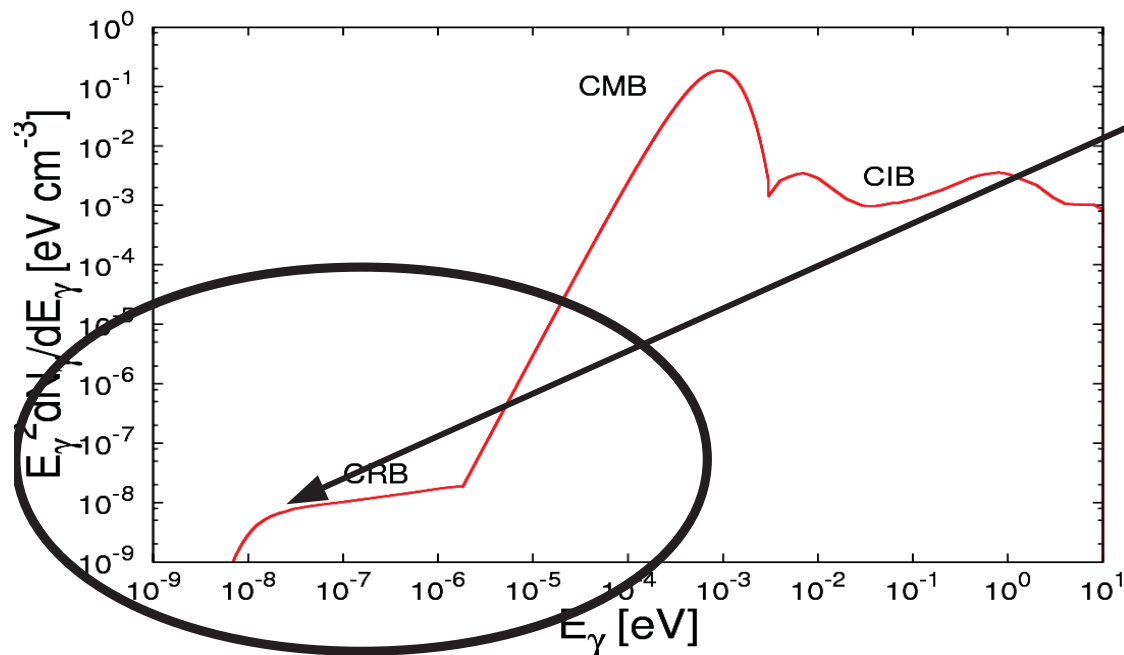
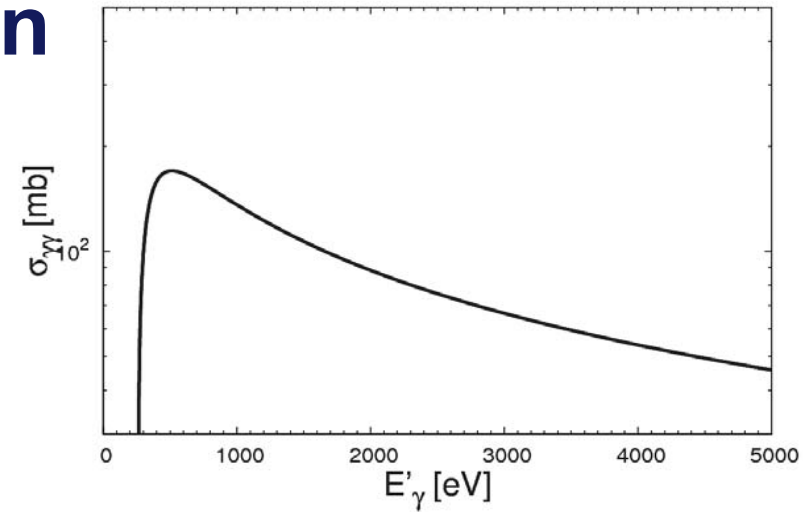
# Gamma-Ray Production



# Gamma-Ray Interaction

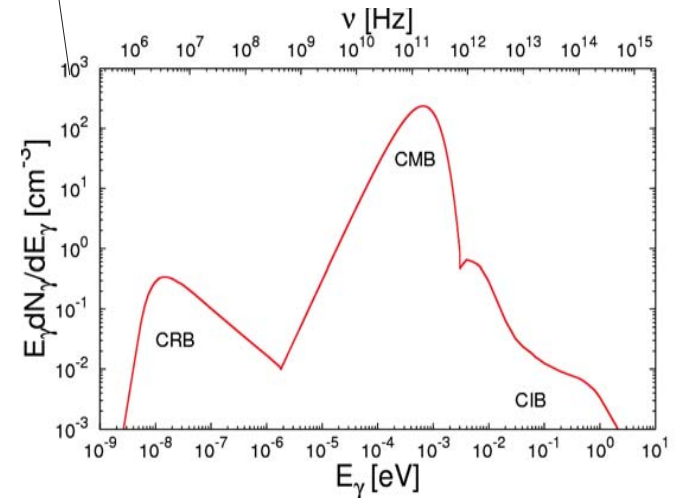
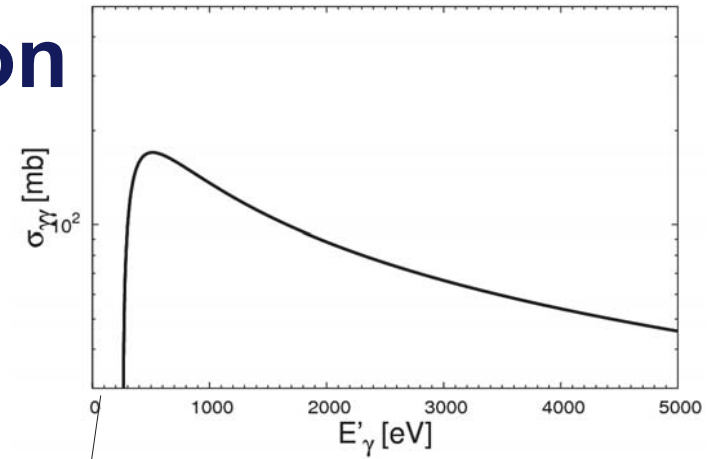
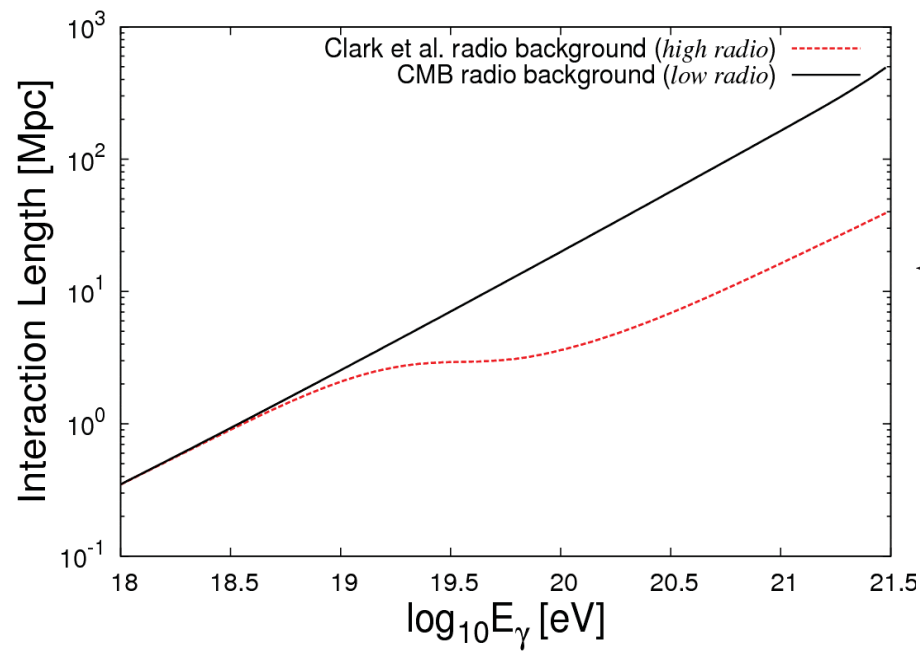


$$s \approx 10^{12} \text{ eV}^2$$

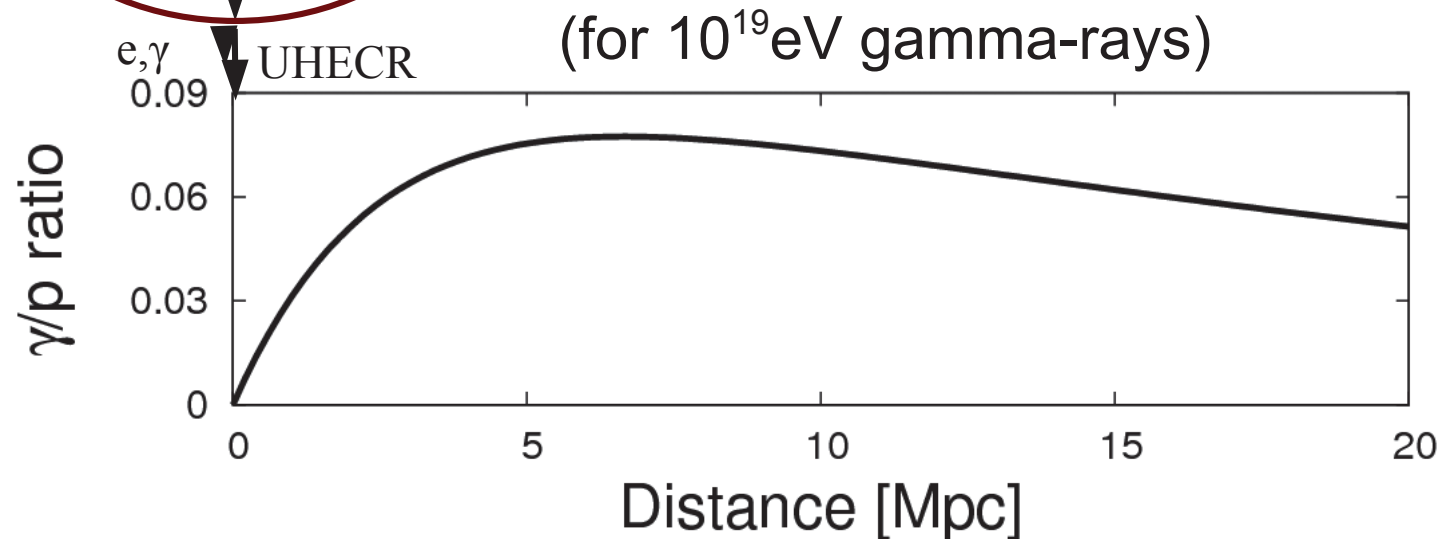
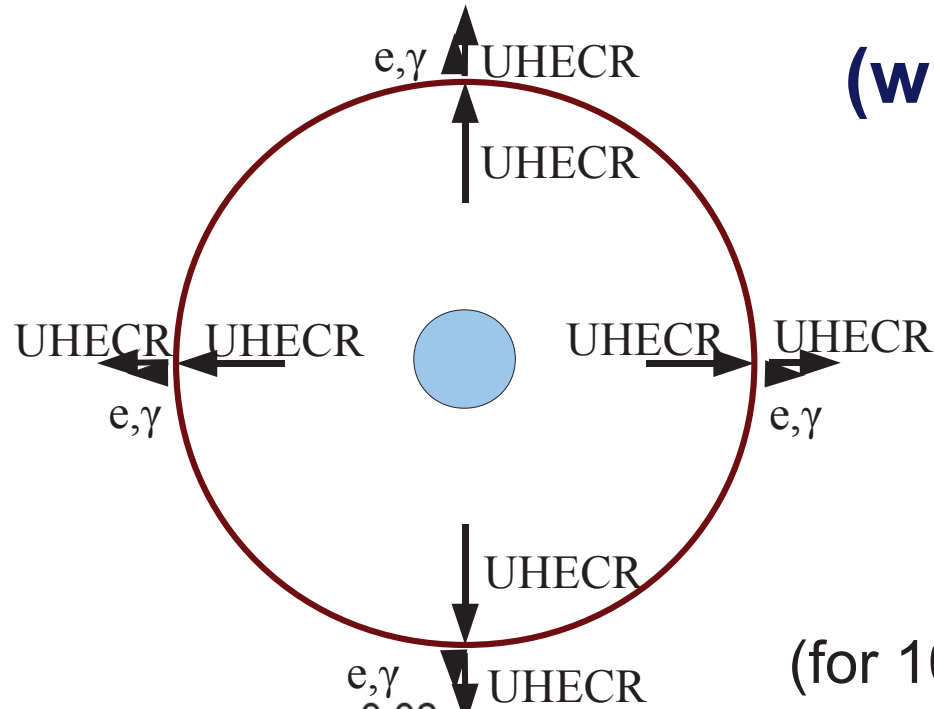


Radio  
Background

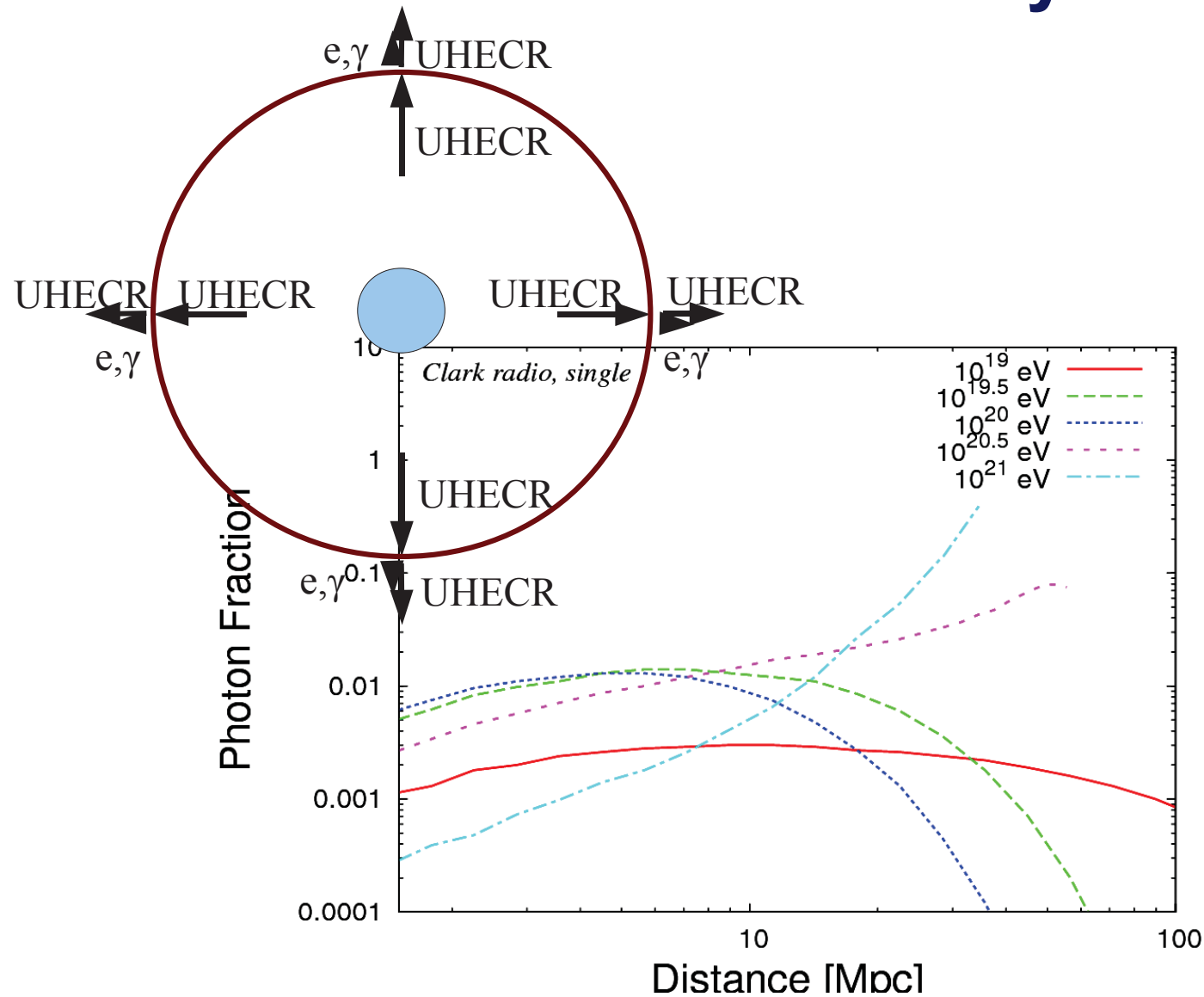
# Gamma-Ray Interaction



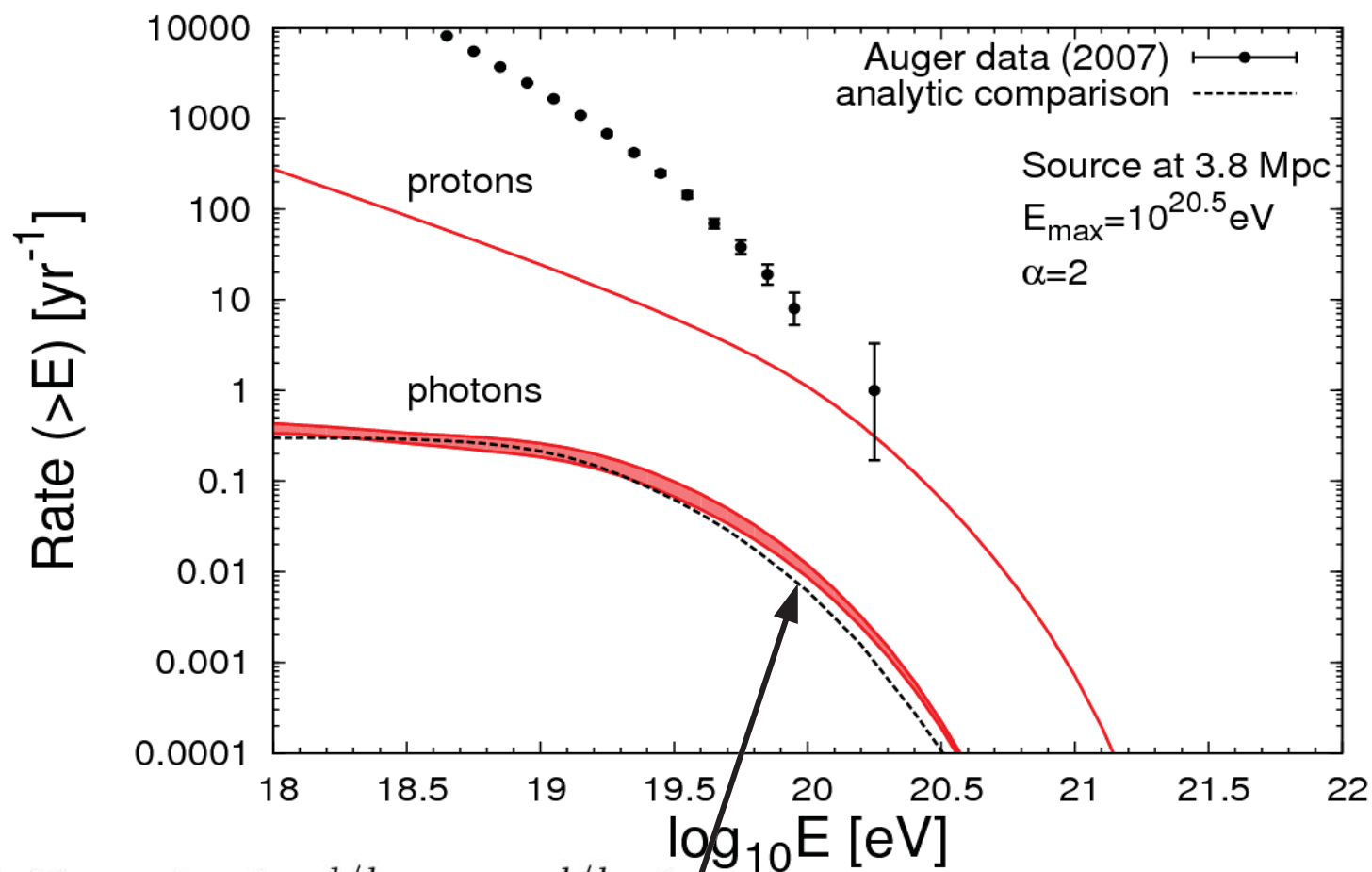
# The Halo Around Heavenly Bodies (which acc. UHECR)



# The Halo Around Heavenly Bodies



# The Halo Around Cen A



$$\frac{N_{\gamma}(l)}{N_p(0)} = \frac{l_{\gamma\gamma}(e^{-l/l_{p\gamma}} - e^{-l/l_{\gamma\gamma}})}{(l_{p\gamma} - l_{\gamma\gamma})}$$

# Conclusion

The dominance of nuclei at the highest energies provides useful new information about the nature of UHECR sources

Regions close to luminous objects are excluded as UHECR sources, favouring slow acceleration scenarios

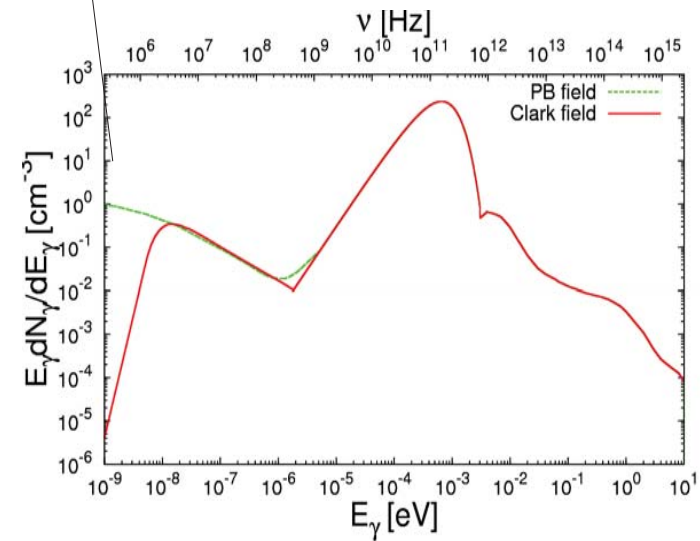
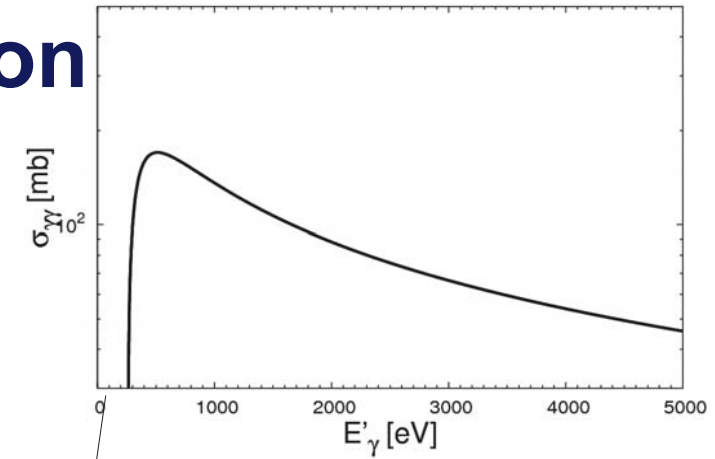
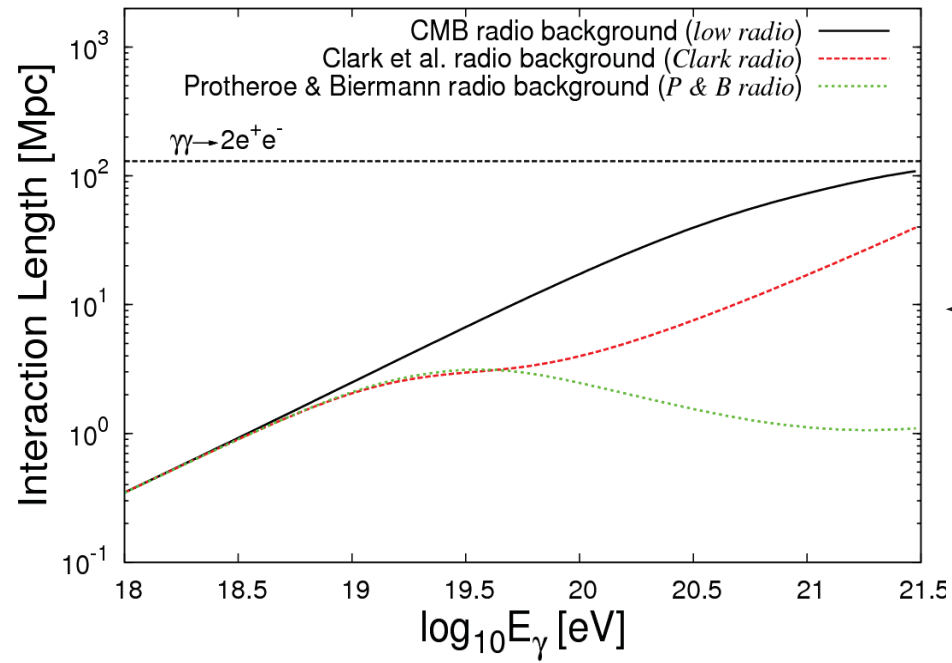
UHE photons can provide a useful probe of local sources

Applied to Cen A we expect an UHE photon in 5 years, if 2 UHECR in the PAO 57 UHECR set originated from Cen A.

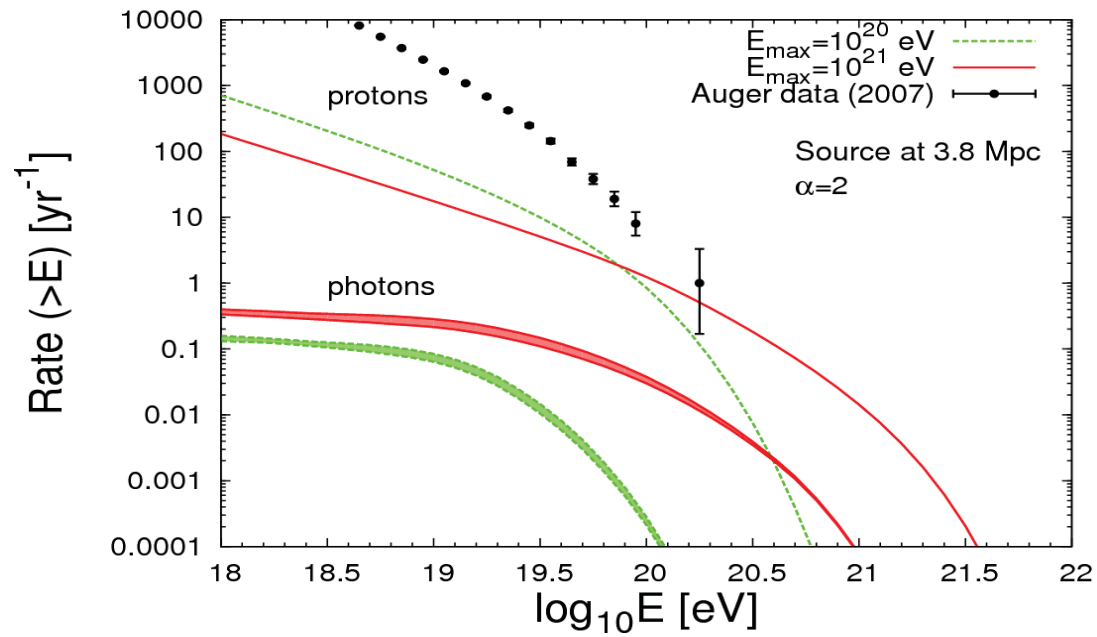
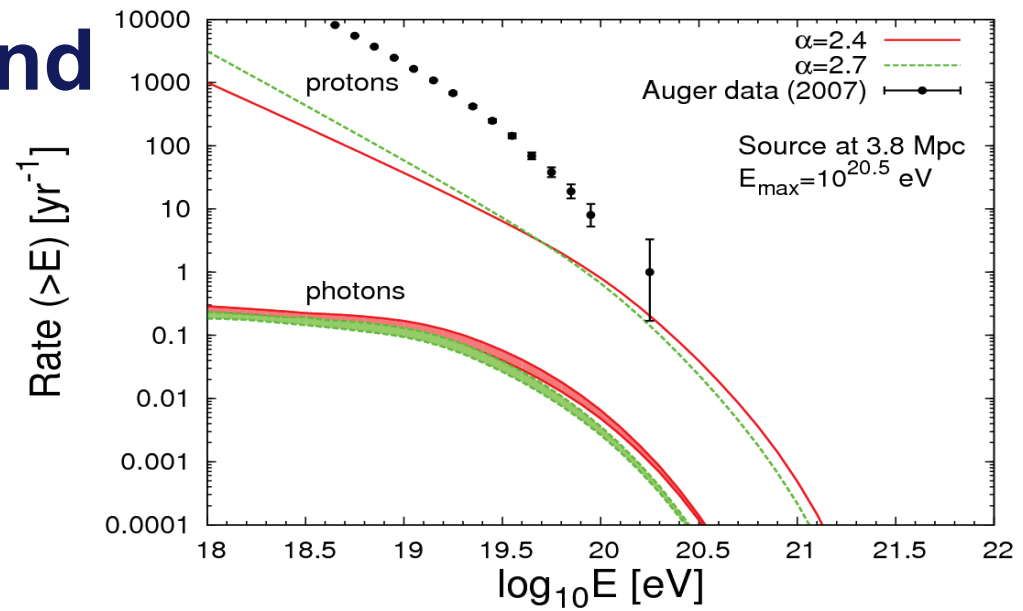


# **Extra Slides**

# Gamma-Ray Interaction



# The Halo Around Cen A

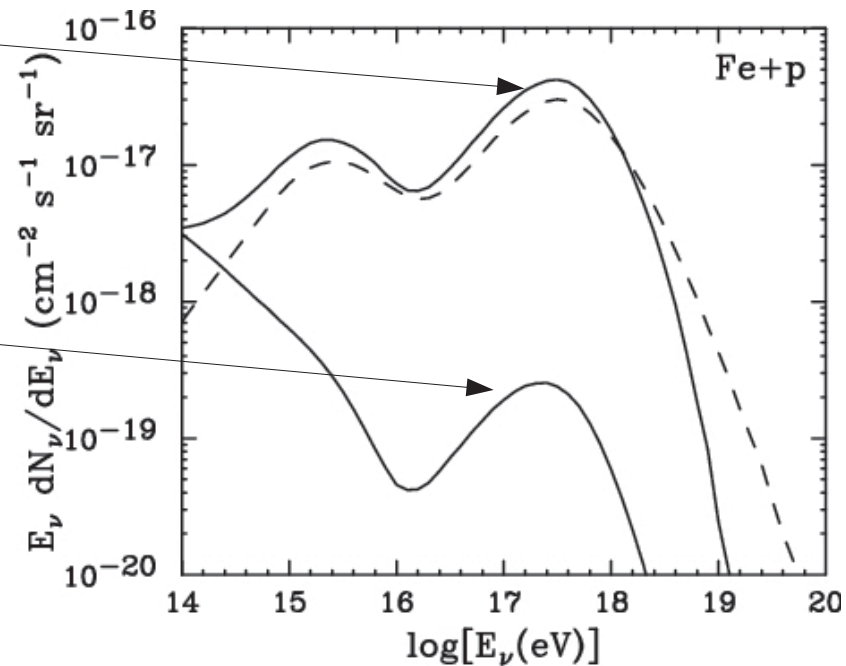


# The Cosmogenic Neutrino Flux

The high energy ( $>10^{17}$  eV) flux quoted as the “Guaranteed flux” value

lowest value compatible with all the data

Smaller value obtained since best agreement found for a dominant Fe fraction with  $E_{\text{max}} = 10^{21}$  eV



# Ratio Between Photo-Pion and Photo-Disintegration Rates (2)

with,

$$\sigma_{p,\gamma} = 1.0 \text{ mb}, E_{p,\gamma} = 31 \text{ MeV}, \Delta_{p,\gamma} = 100 \text{ MeV}$$

and

$$\sigma_{A_{56},\gamma} = 81 \text{ mb}, E_{A_{56},\gamma} = 18 \text{ MeV}, \Delta_{A_{56},\gamma} = 8 \text{ MeV}$$

therefore

$$\begin{aligned} R_{A_{56},\gamma}(\Gamma) &\approx \frac{\sigma_{A_{56},\gamma}}{\sigma_{p,\gamma}} R_{p,\gamma}(15\Gamma) \\ &= 160 R_{p,\gamma}(15\Gamma) \end{aligned}$$