



The Abdus Salam
International Centre for Theoretical Physics



2246-34

**Workshop on Cosmic Rays and Cosmic Neutrinos: Looking at the
Neutrino Sky**

20 - 24 June 2011

Quantifying uncertainties in the high energy neutrino cross-section

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The high energy neutrino cross-section in the Standard Model and its uncertainty

Philipp Mertsch

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NUSKY 2011, ICTP, Trieste
24 June 2011



Science & Technology
Facilities Council



Why do we want predictions for the ν cross-section?

ν astronomy

- want to measure flux J_ν
- event rate $R \propto J_\nu \sigma_\nu$
- even considering the attenuation: $R/J_\nu \propto \sigma_\nu^{0.45}$

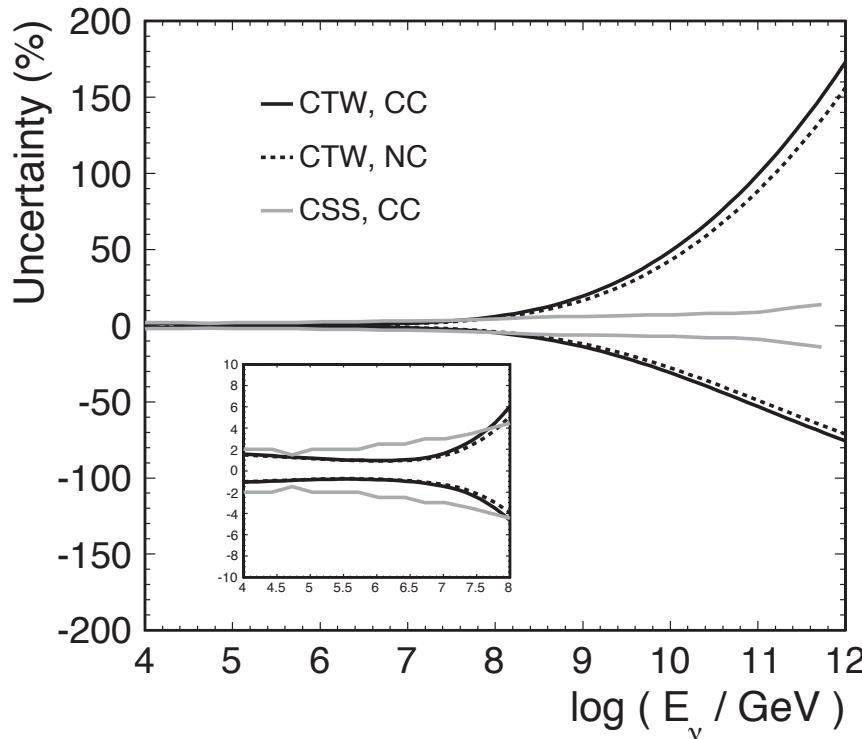
particle physics

- want to measure cross-section σ_ν
- test standard model at c.m. energies up to $\sim 10^3$ TeV:
gluon saturation, colour glass condensate, black holes?

- there are a few observables which are independent of σ
 \rightarrow measure flux **and** cross-section

Kusenko and Weiler, PRL 88, 161101 (2002); Anchordoqui *et al.*, PRD 74 (2006) 043008

How accurately can we predict the ν cross-section?

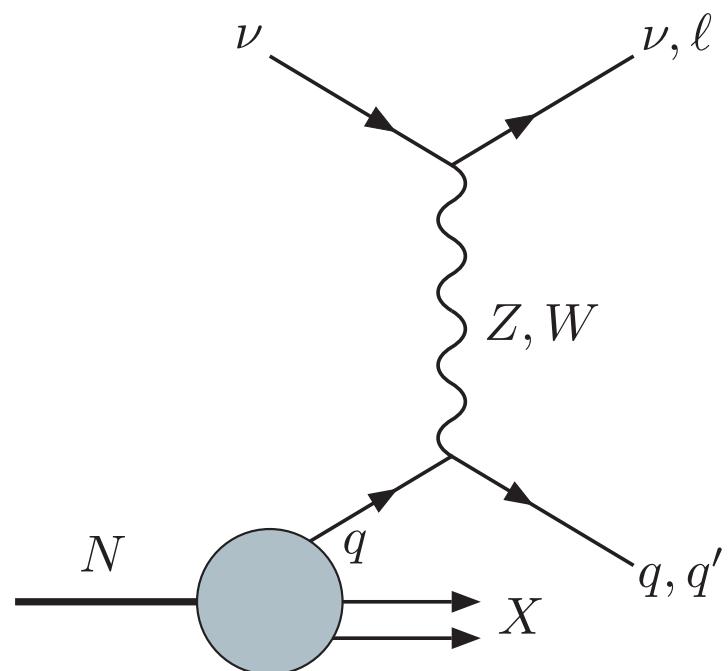


CSS: Cooper-Sarkar
and Sarkar,
JHEP 01 (2008) 075

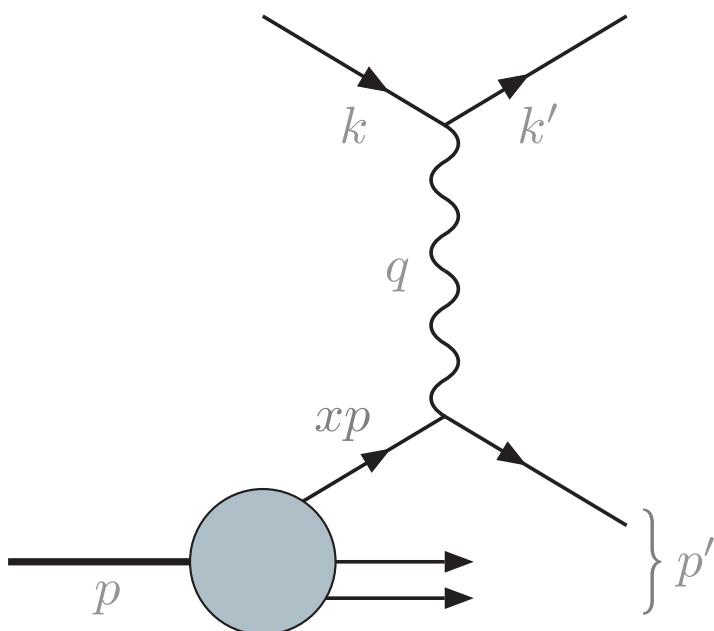
CTW: Connolly *et al.*,
arXiv:1102.0691

Does the uncertainty **really** blow up to $\mathcal{O}(1)$?

DIS



DIS



four Lorentz invariants:

- centre of mass energy \sqrt{s}
 $s = (p + k)^2$
- momentum transfer
 $Q^2 = -q^2 = -(k - k')^2$
- Bjorken scaling variable
 $x = Q^2 / (2p \cdot q)$
- inelasticity
 $y = p \cdot q / (p \cdot k)$

ν cross-section

Double differential cross-section

$$\frac{d^2\sigma(\nu(\bar{\nu})N)}{dx \ dQ^2} = \frac{G_F^2 M_W^4}{4\pi(Q^2 + M_W^2)^2 x} \sigma_r(\nu(\bar{\nu})N)$$

with reduced cross-section

$$\sigma_r(\nu(\bar{\nu})N) = [Y_+ F_2^\nu(x, Q^2) - y^2 F_L^\nu(x, Q^2) \pm Y_- x F_3^\nu(x, Q^2)]$$

where $Y_\pm = 1 \pm (1 - y)^2$.

Total cross-section

$$\sigma = \int dx \int dQ^2 \frac{d^2\sigma(\nu(\bar{\nu})N)}{dx \ dQ^2}$$

LO (quark parton model)

$q(x)$: probability density for quark q with momentum fraction x

structure functions

combination of quark PDFs,

$$F_2^\nu = x(u + d + 2s + 2b + \bar{u} + \bar{d} + 2\bar{c}),$$

$$F_L^\nu = 0,$$

$$xF_3^\nu = x(u + d + 2s + 2b - \bar{u} - \bar{d} - 2\bar{c}),$$

and similar for $\bar{\nu}$.

NLO

$q(x, Q^2)$ now scale-dependent, **no** probability density

structure functions

$$F_2 = \int_x^1 \frac{d\xi}{\xi} \left[\sum_i e_i^2 x q_i(\xi, Q^2) C_q \left(\frac{x}{\xi}, \alpha_s \right) + \left(\sum_i e_i^2 \right) x g(\xi, Q^2) C_g \left(\frac{x}{\xi}, \alpha_s \right) \right]$$

$$F_L = \dots$$

$$xF_3 = \dots$$

C_q, C_g : coefficient functions

PDF fitting: idea

problem

- ideally, would like to calculate PDFs from first principles
- however, interactions of partons are soft ($Q^2 \lesssim \Lambda_{\text{QCD}}^2$)
 - non-perturbative regime
 - lattice?

DGLAP evolution

- however, can calculate the evolution of PDFs in the perturbative regime ($Q^2 \gg \Lambda_{\text{QCD}}^2$)
 - assume parametric form at input scale and evolve to other scale

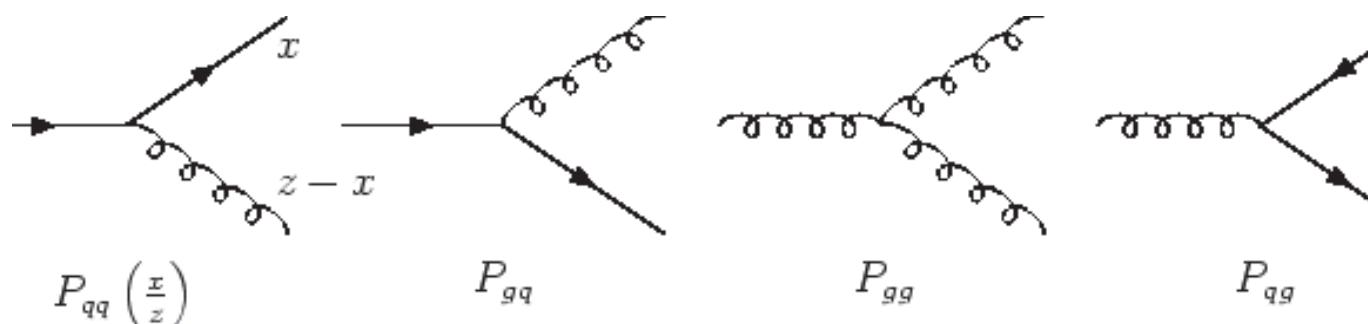
DGLAP evolution

$$\frac{\partial q^{\text{NS}}(x, Q^2)}{\partial \ln Q^2} = \frac{\alpha_s}{2\pi} (q^{\text{NS}} \otimes P_{qq})$$

$$\frac{\partial \Sigma(x, Q^2)}{\partial \ln Q^2} = \frac{\alpha_s}{2\pi} (\Sigma \otimes P_{qq} + g \otimes 2n_f P_{qg})$$

$$\frac{\partial \Sigma(x, Q^2)}{\partial \ln Q^2} = \frac{\alpha_s}{2\pi} (\Sigma \otimes P_{gq} + g \otimes 2n_f P_{gg})$$

Σ and q^{NS} are convenient linear combinations of quark PDFs.



PDF fitting: procedure

- chose parametrisation at input scale Q_0^2 , e.g.

$$xg = x^{\lambda_g} (1-x)^{\eta_g} P_g(x)$$
$$xS = x^{\lambda_S} (1-x)^{\eta_S} P_S(x)$$

...

- evolve to scale of measurement: $Q_0^2 \rightarrow Q^2$
- calculate F_2 , F_L and xF_3 functions and (differential) cross sections
- determine parameters λ_i , η_i , $P_i(x)$ by fitting to data

PDF errors

experimental uncertainties

- many experimental errors correlated
- correlation matrix diagonalised
 - linearly independent eigenvectors = variations of best-fit PDF
- can add errors from eigenvectors in quadrature

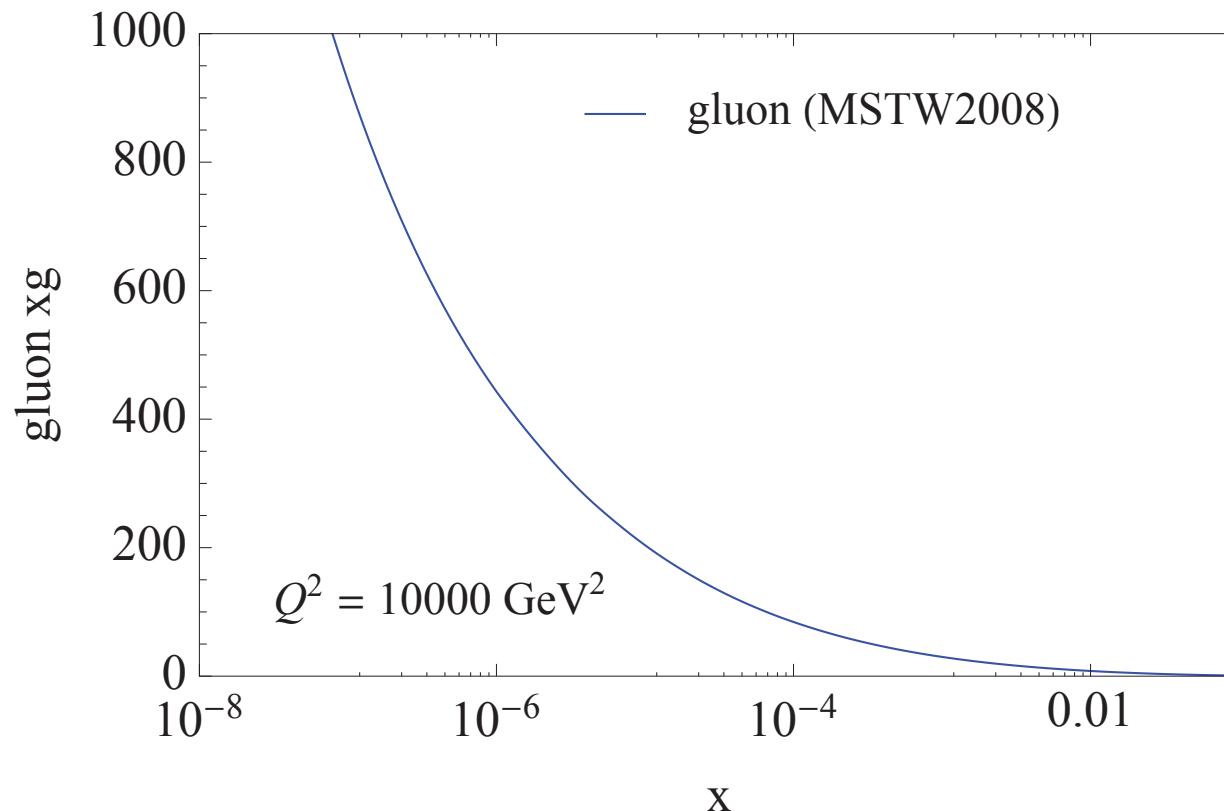
model/parameter uncertainties

- some parameters/model assumptions get fixed before fit
- vary these parameters within c.l. interval
 - variations of best-fit PDF

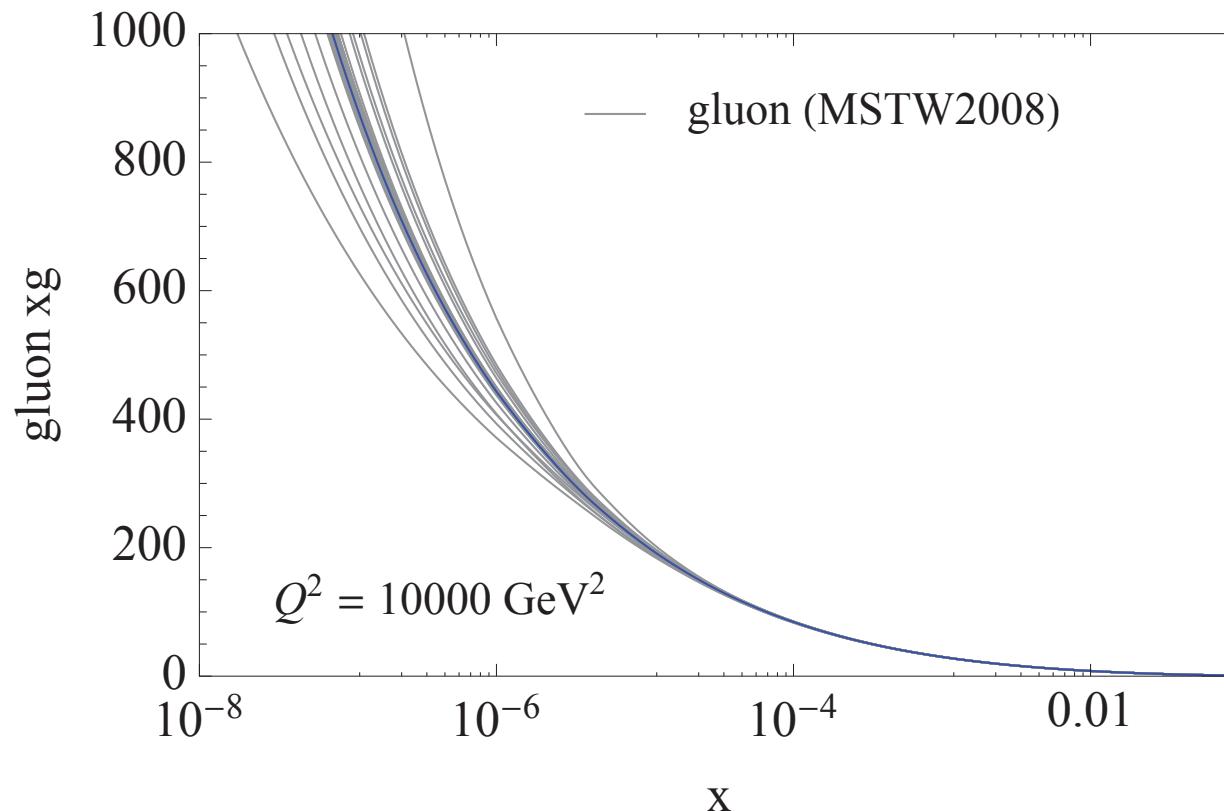
α_s uncertainties

- α_s determines how quickly PDFs rise at low x
- possibly large effect

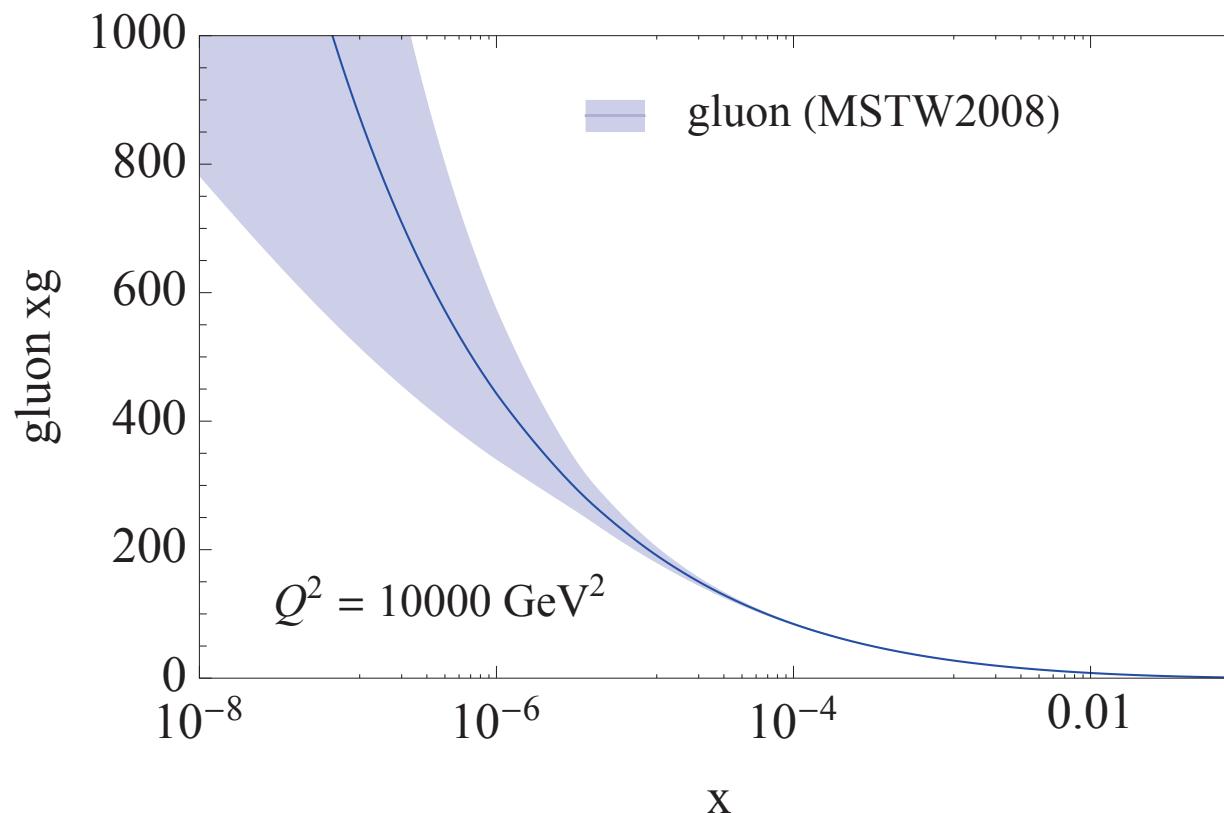
PDF errors



PDF errors



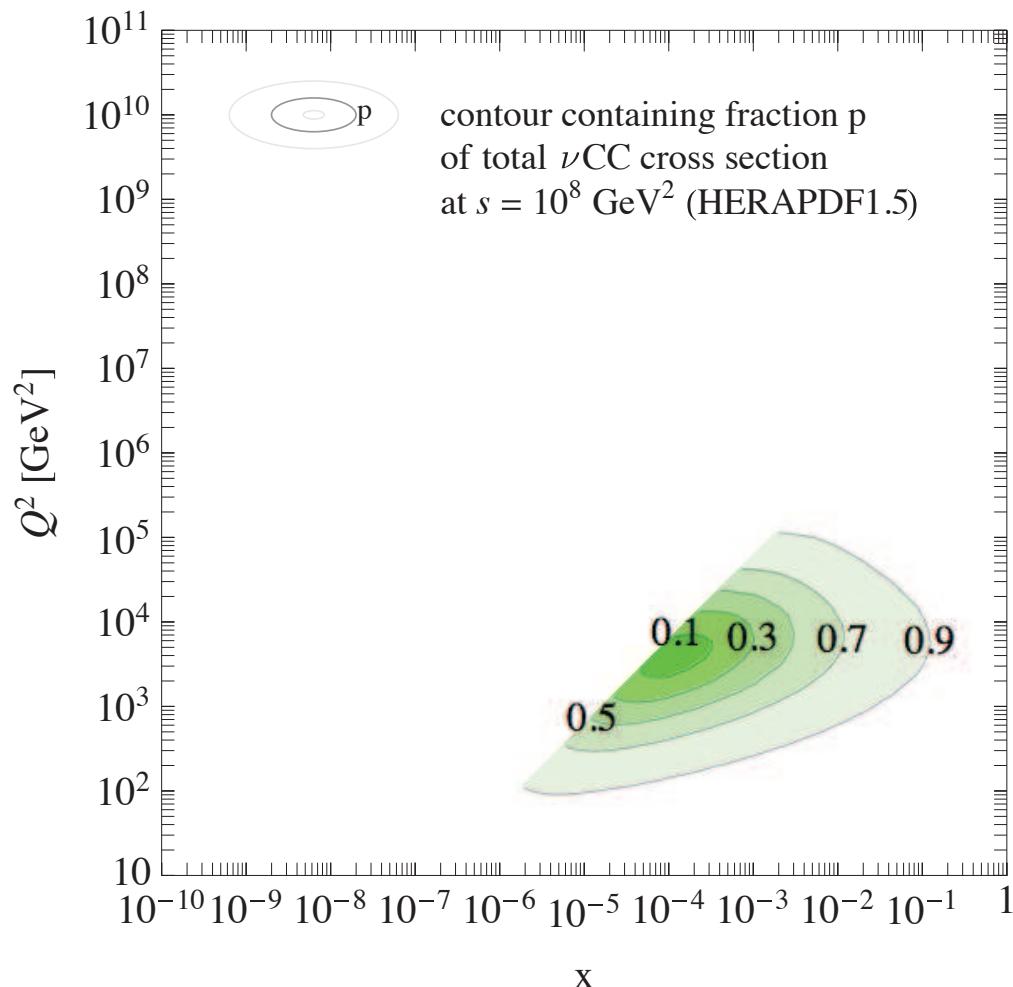
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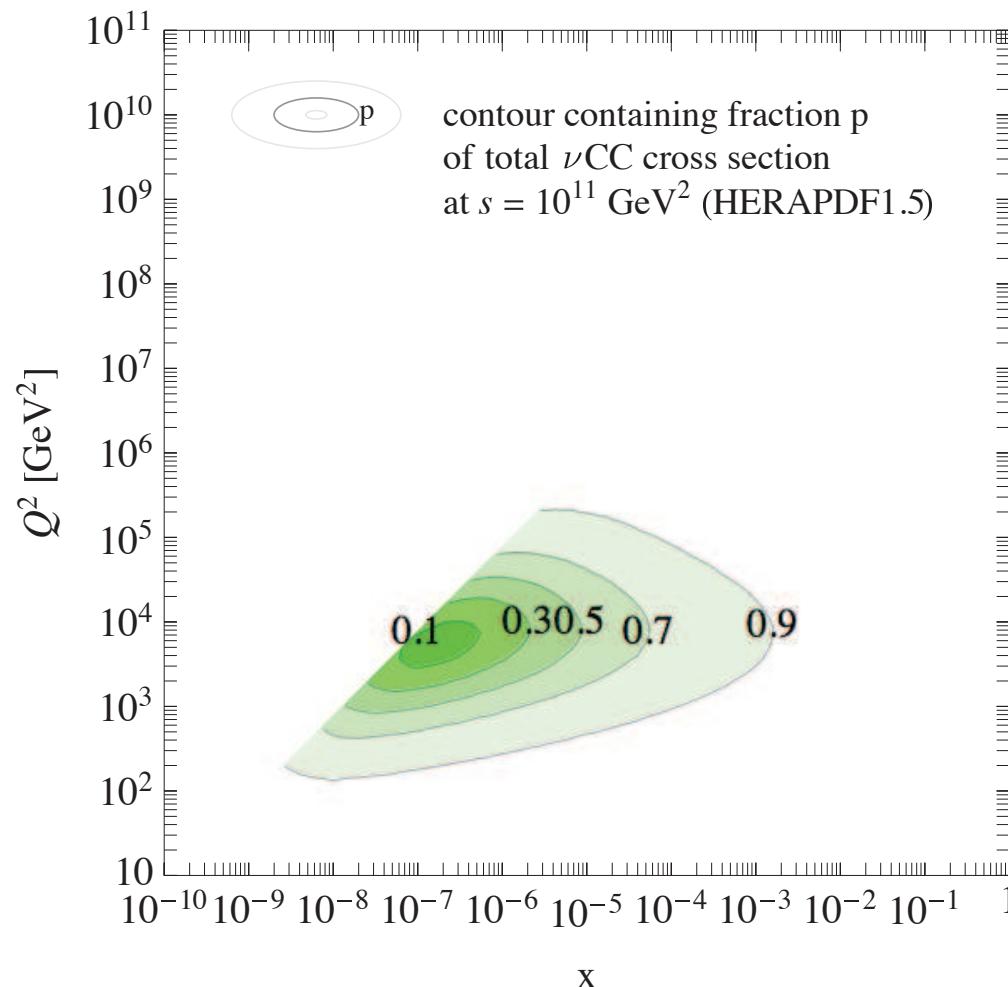
A detailed comparison

- use only up-to-date PDFs:
 - ▶ HERAPDF1.5 ✓
 - ▶ CT10 ✓
 - ▶ MSTW2008 ✗ (does not include combined HERA data)
- work **consistently** at NLO
- use only publicly available tools (e.g. LHAPDF)
- highlight different contributions to uncertainty **within DGLAP**:
 - ▶ experimental
 - ▶ parameters
 - ▶ model

The kinematic range



The kinematic range



Pitfalls

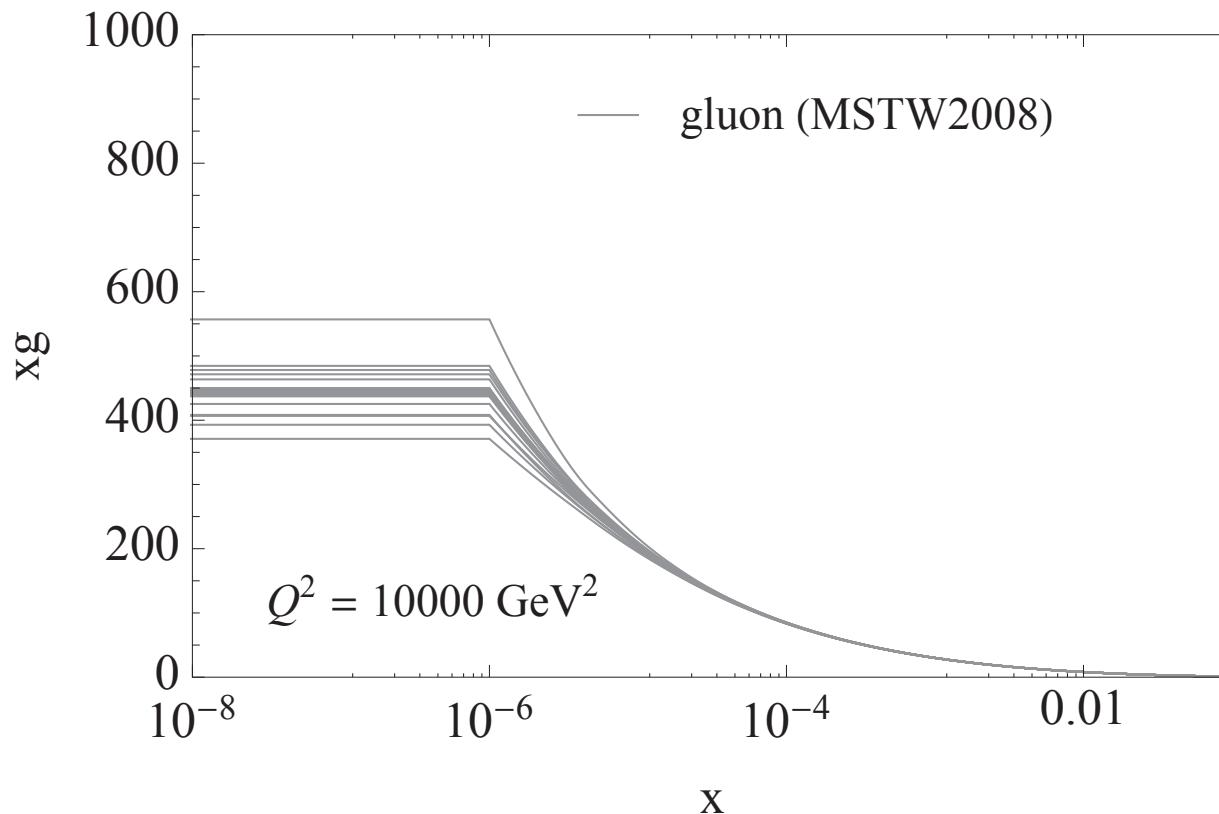
event generators, e.g. PYTHIA

- are for the most part LO
- using NLO PDFs: inconsistent \times

LHAPDF

- PDFs provided on a limited grid of points (x, Q^2)
- going beyond this grid: PDFs “freeze” \times

Pitfalls



Pitfalls

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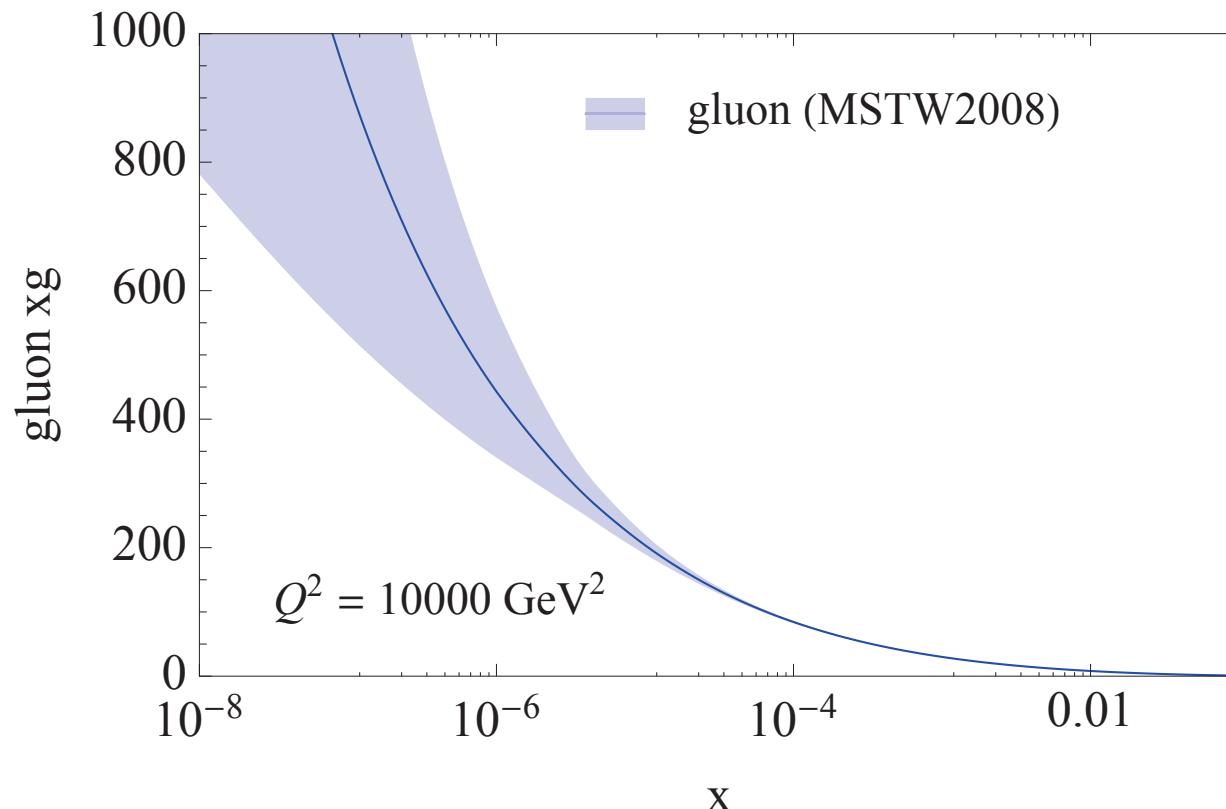
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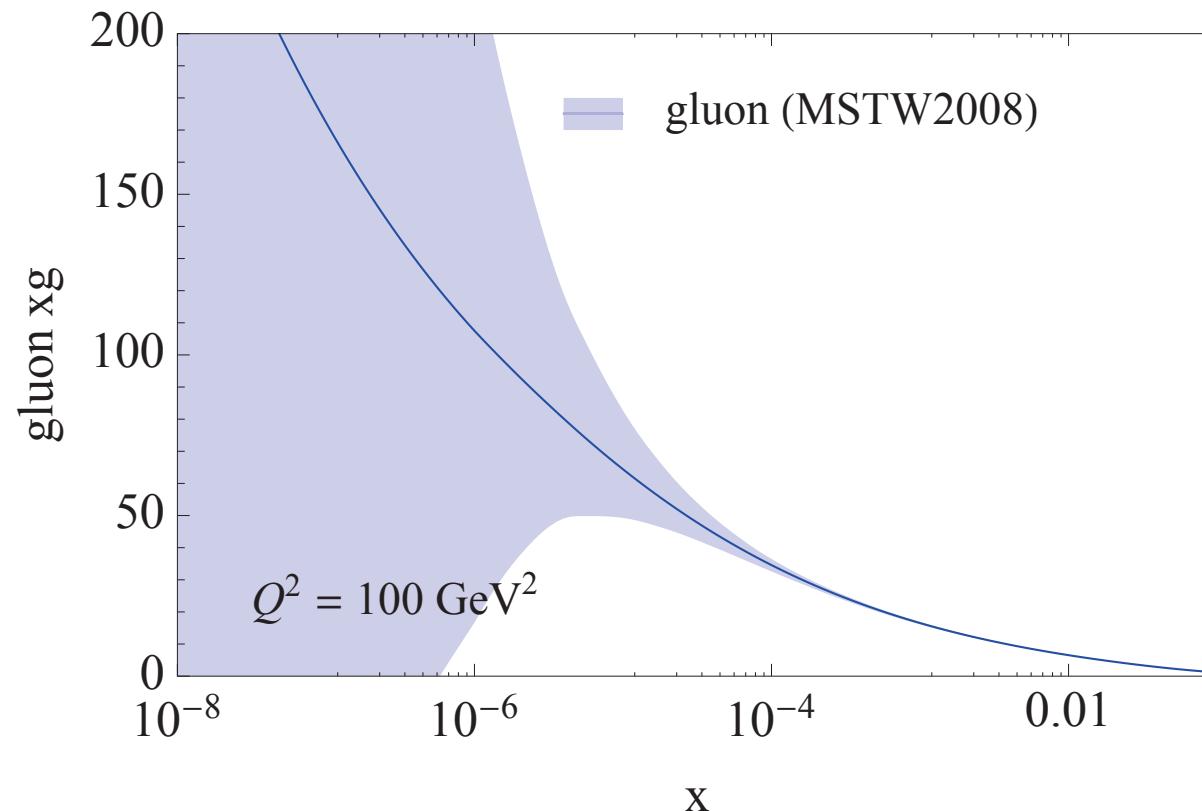
gluon parametrisation

- some groups choose a general parametrisation
- gluon PDF can go negative: meaning?

Example: MSTW2008 gluon momentum distribution



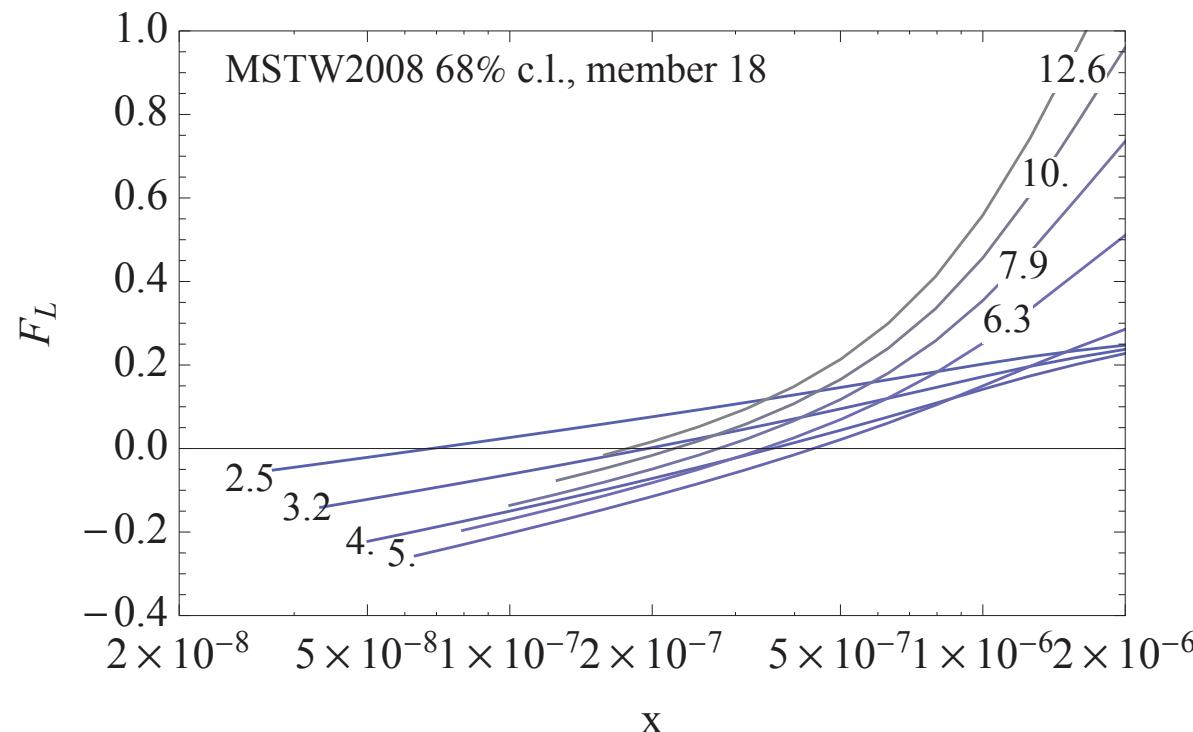
Example: MSTW2008 gluon momentum distribution



Could the gluon become negative?

at NLO, the gluon **could** become negative

however longitudinal structure function F_L **must** stay positive

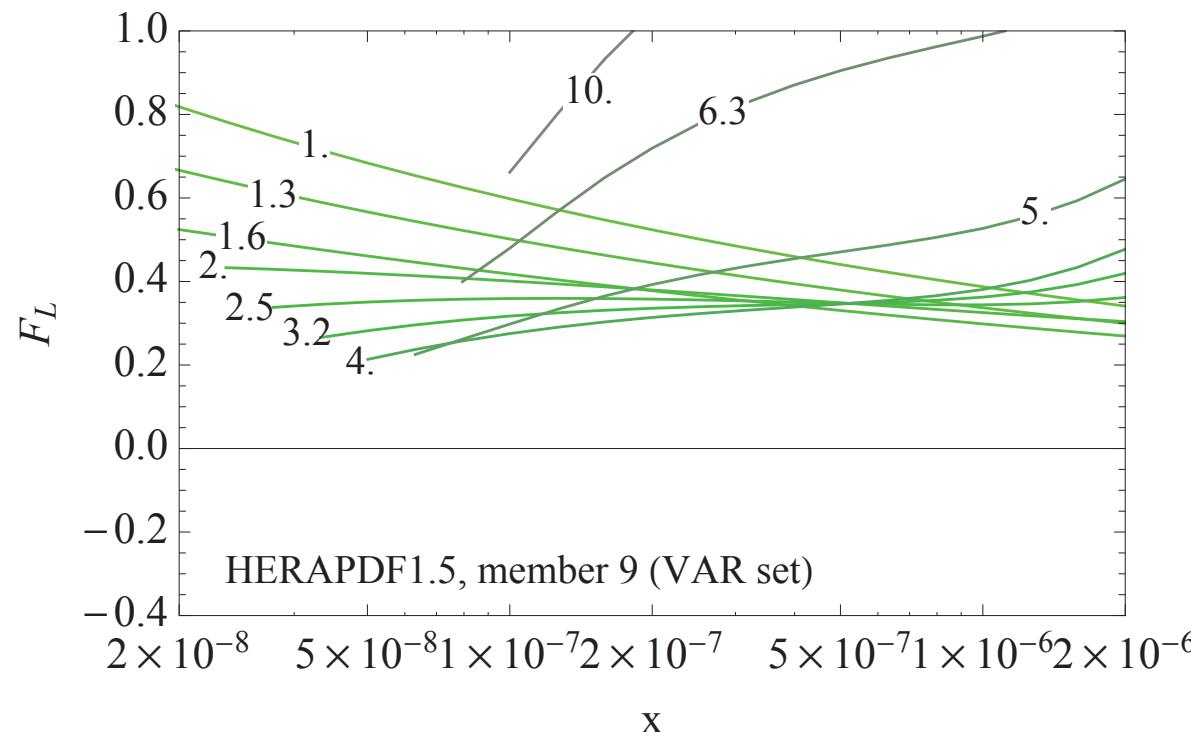


With MSTW2008, F_L **does** go negative!

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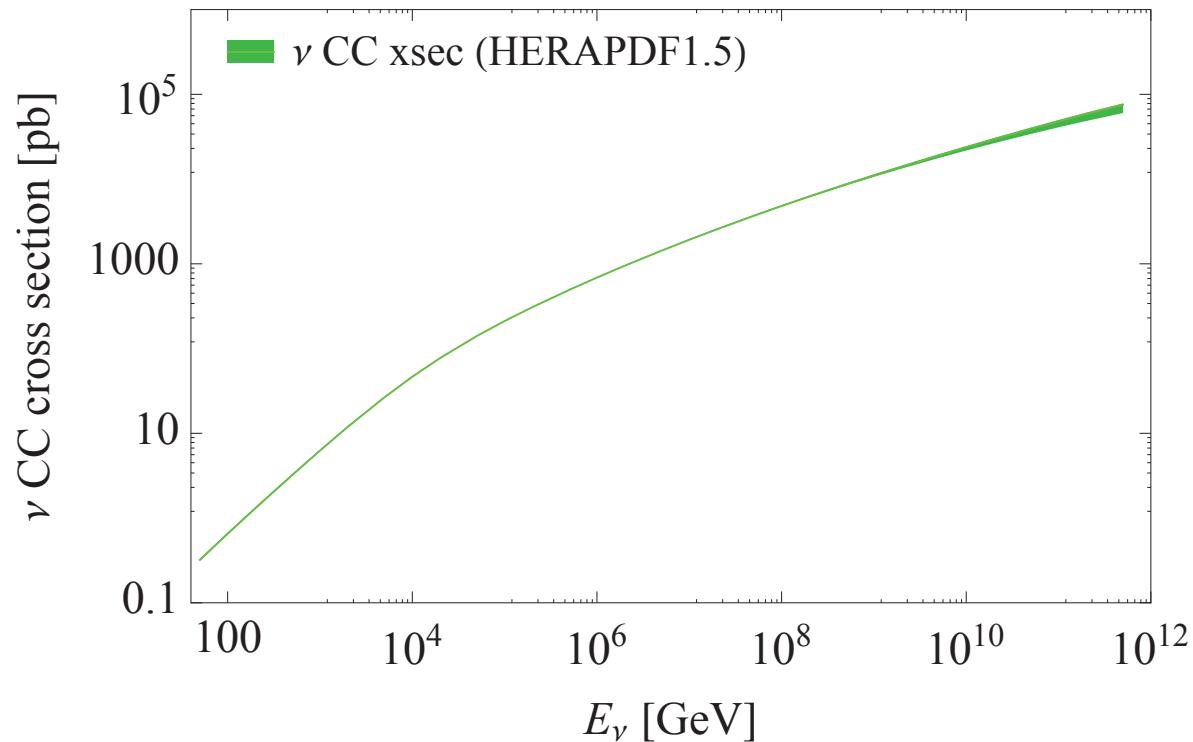
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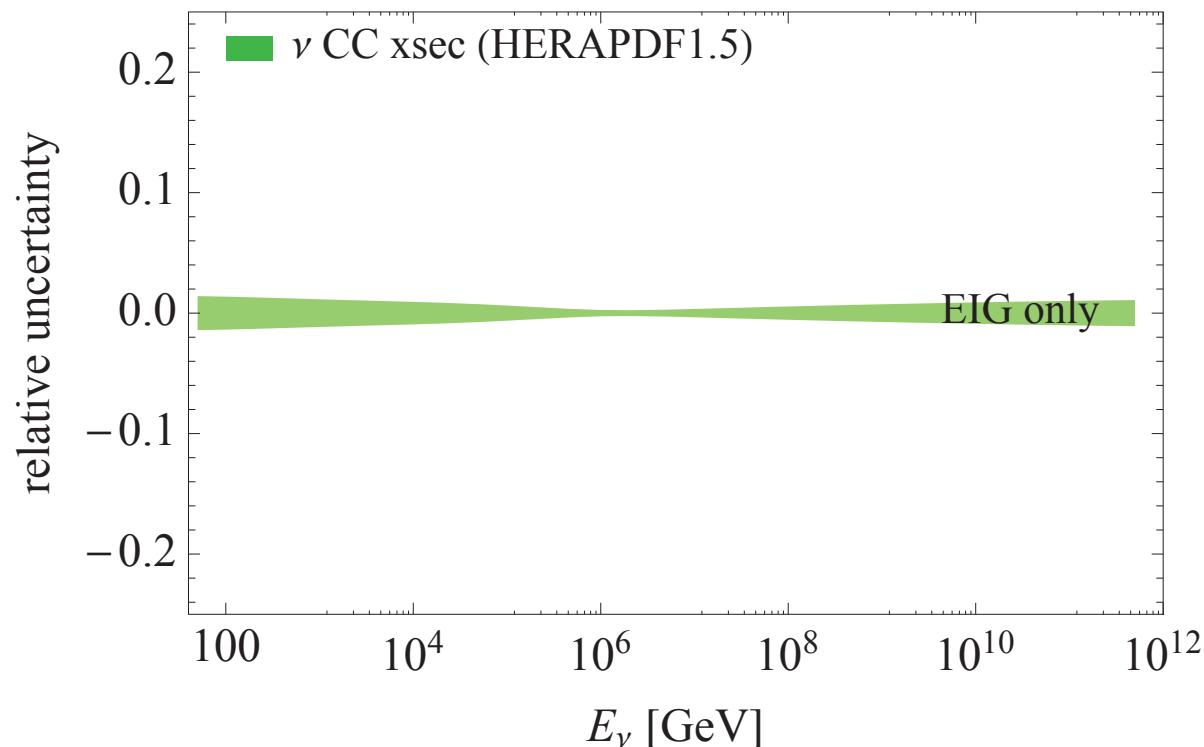


With HERAPDF1.5, F_L **does** stay positive!

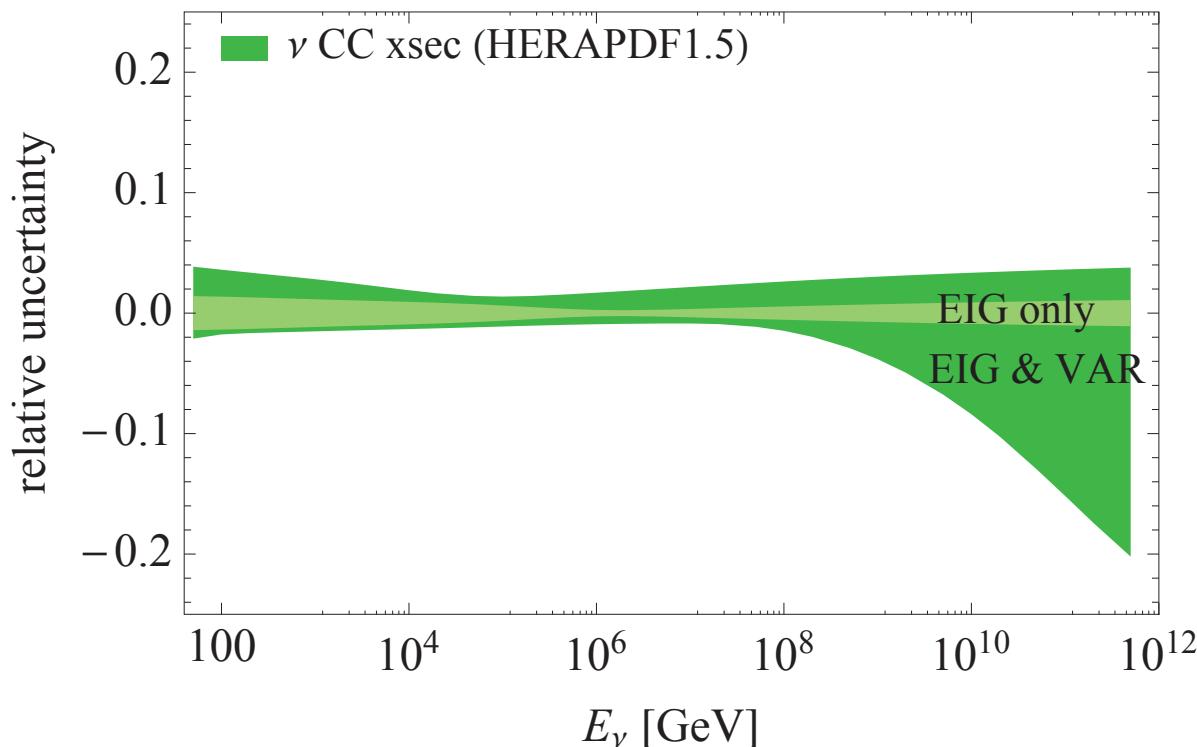
Total ν CC cross-section (HERAPDF1.5)



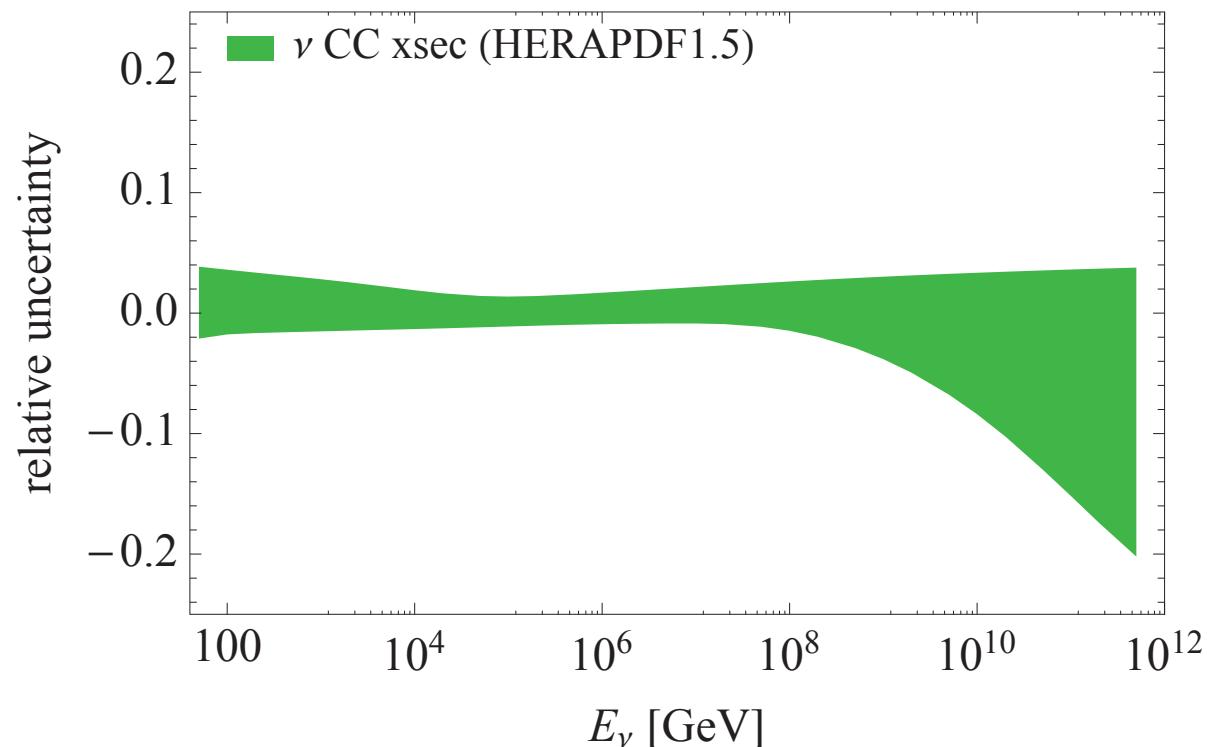
ν CC cross-section uncertainty (HERAPDF1.5)



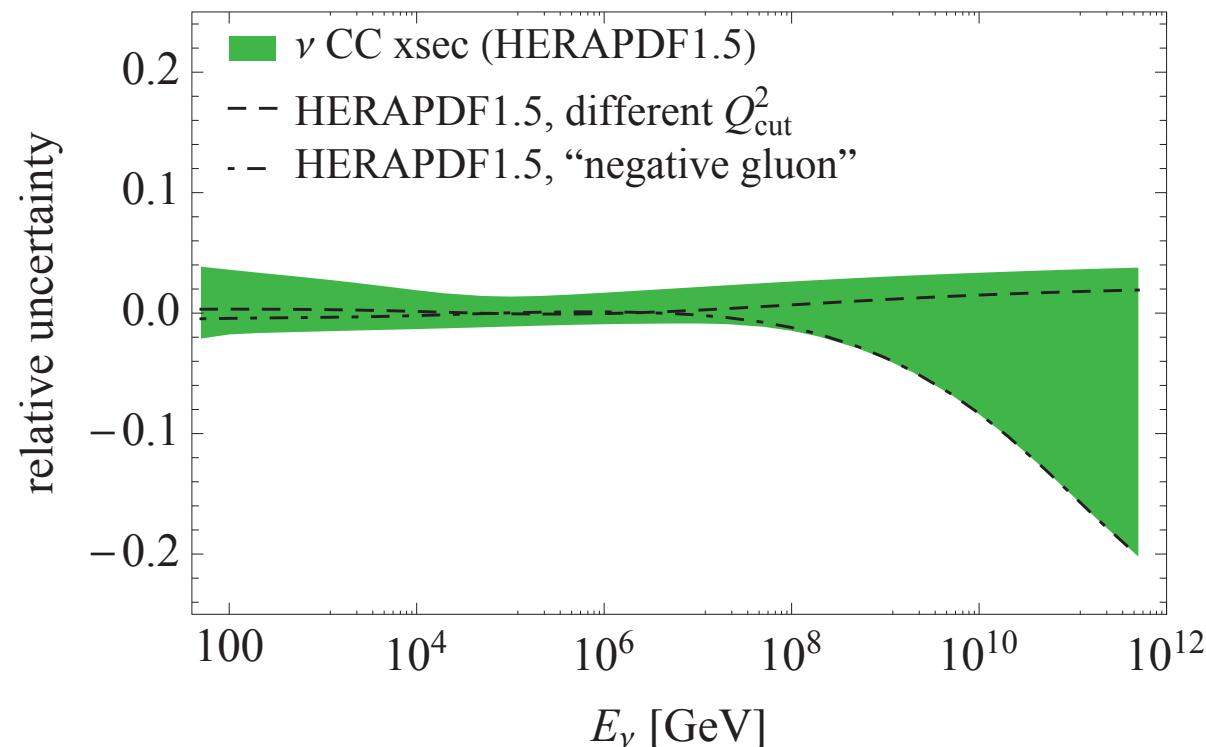
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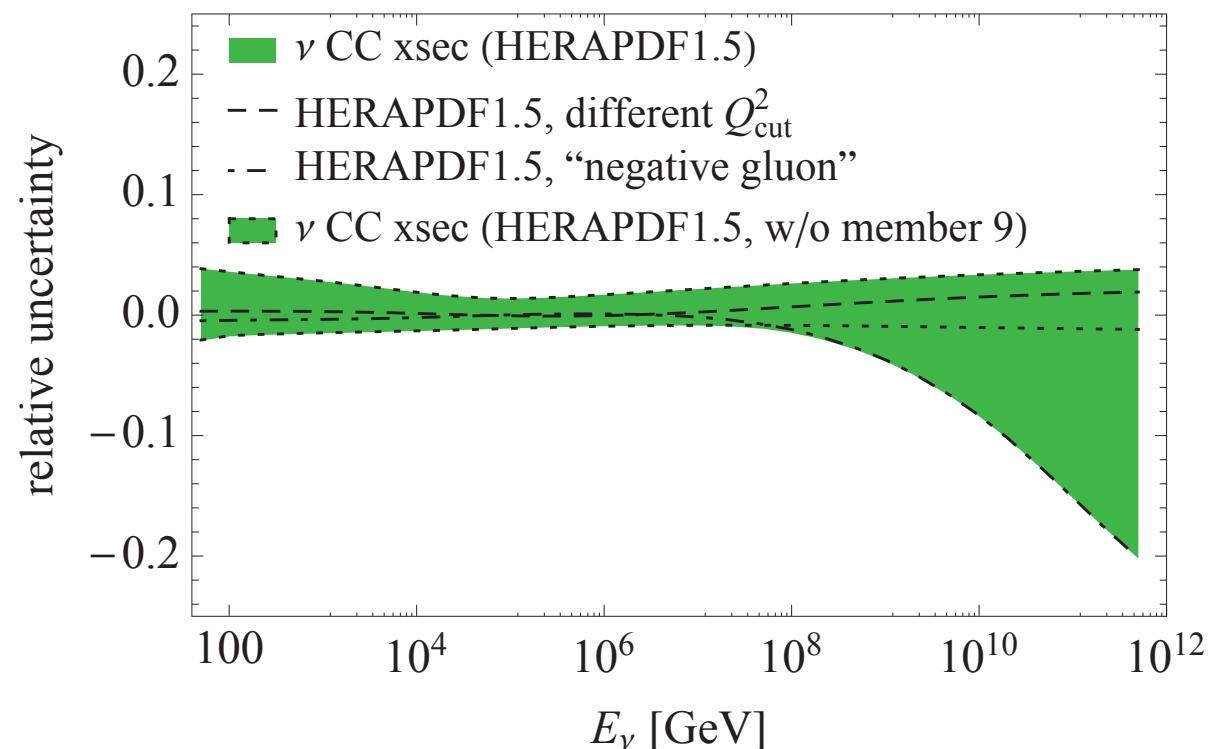
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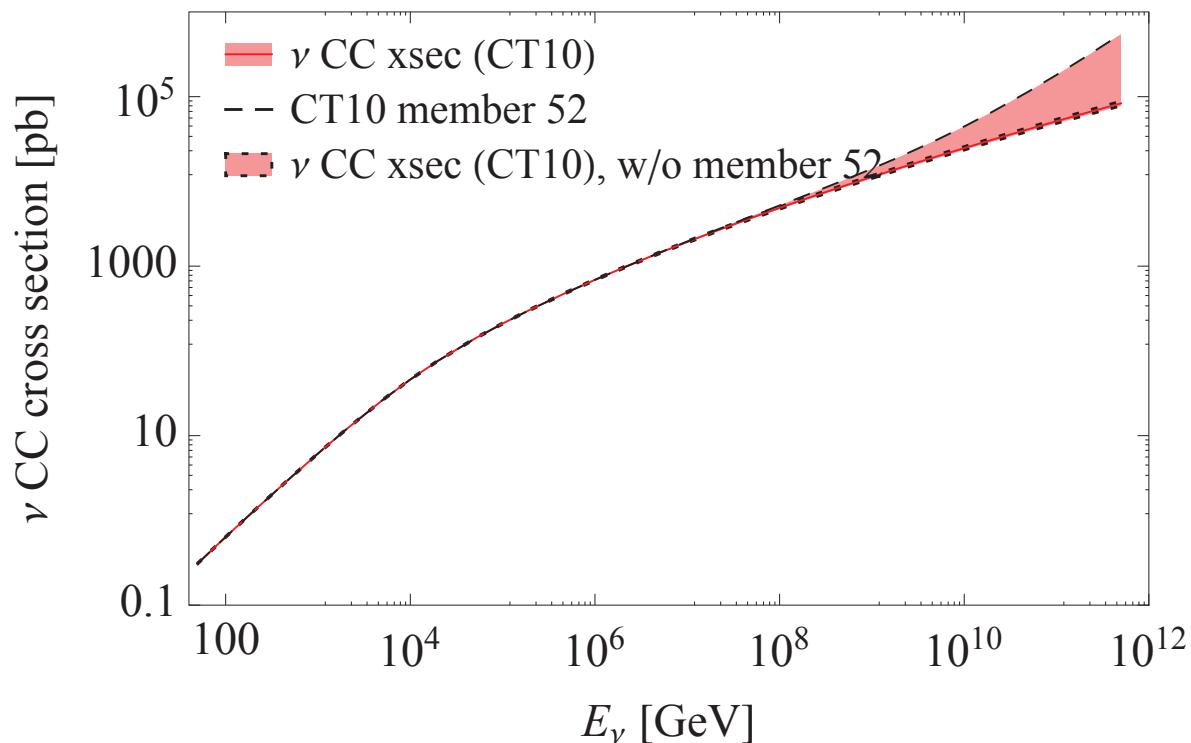
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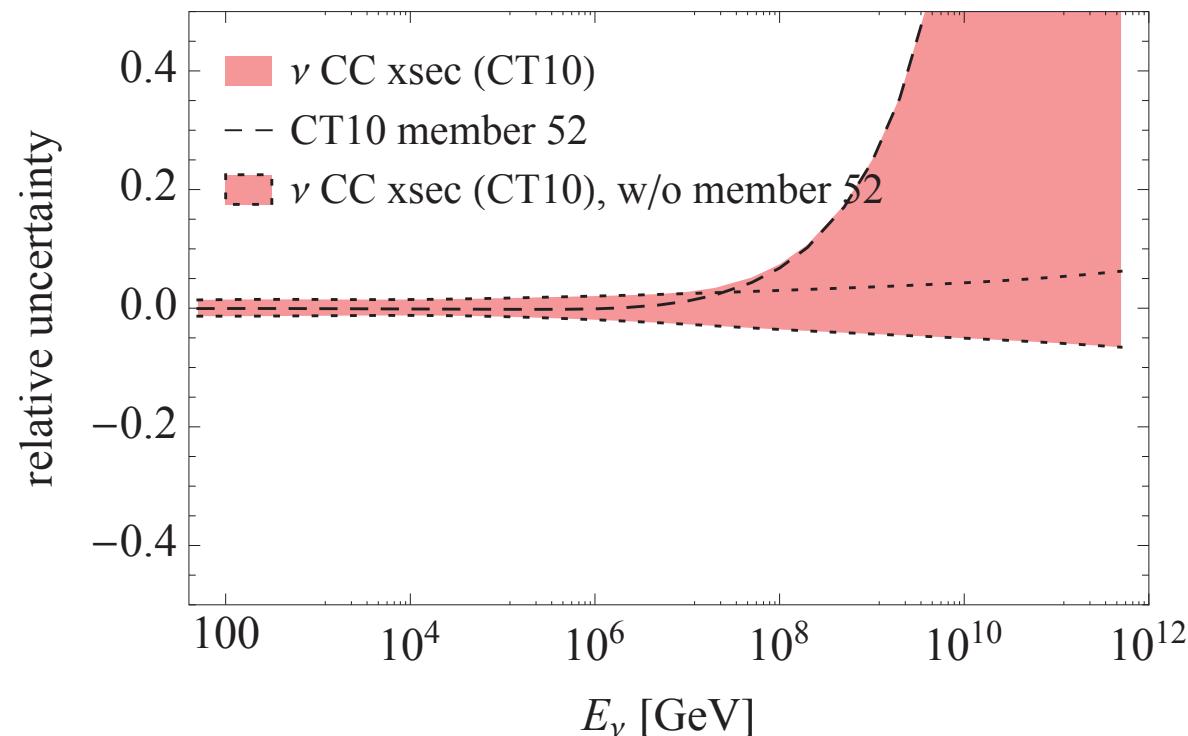


Total ν CC cross-section (CT10)

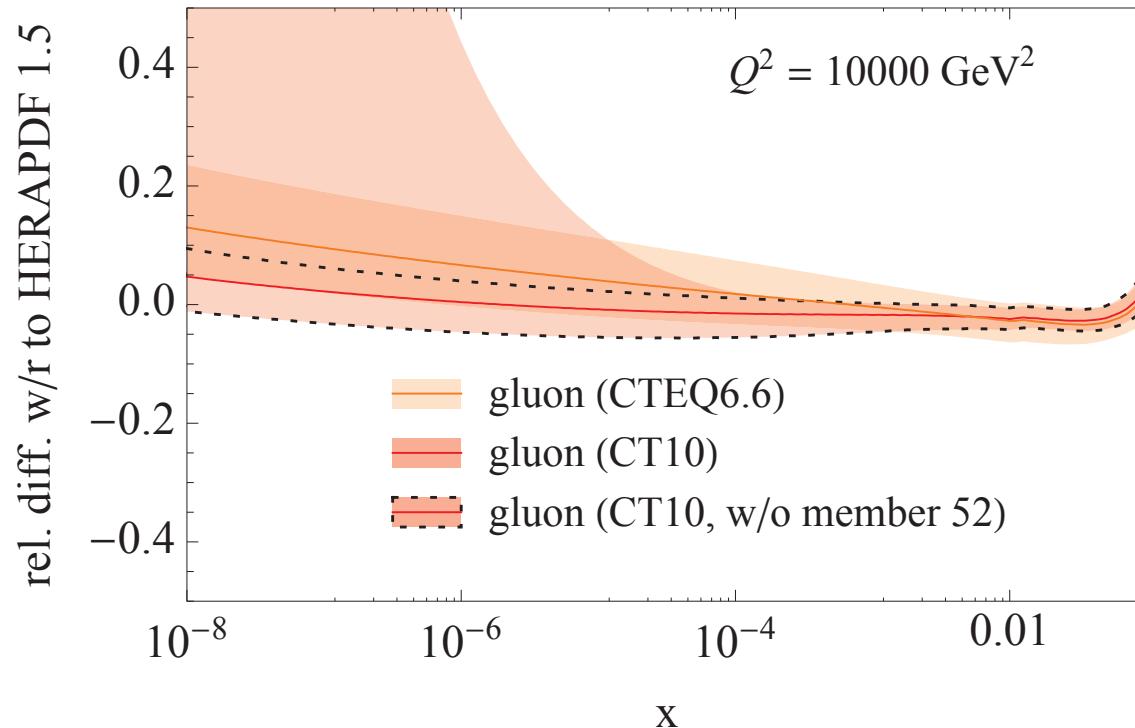


cross-section for member 52 rises $\propto E_\nu^{0.7}$; central member $\propto E_\nu^{0.3}$

ν CC cross-section uncertainty (CT10)

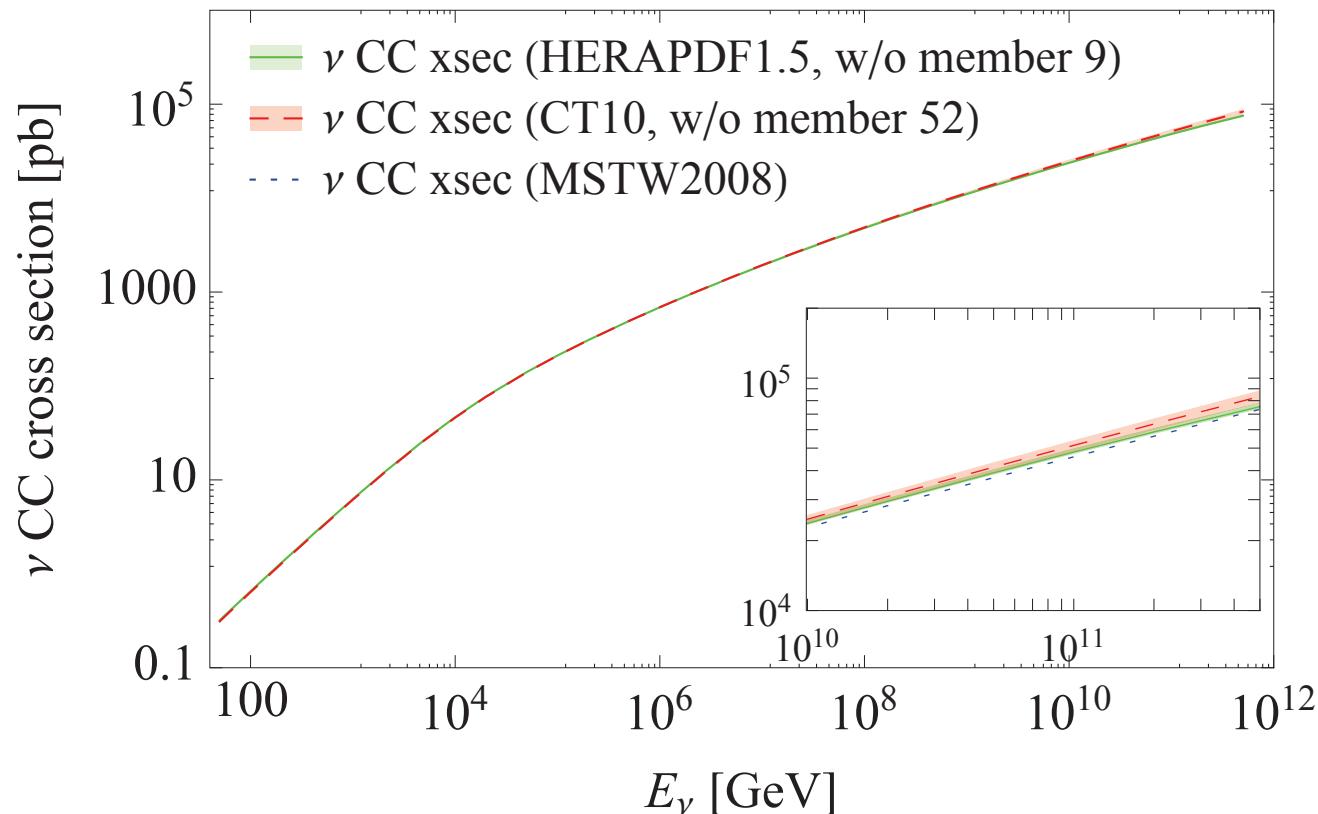


member 52 of CT10

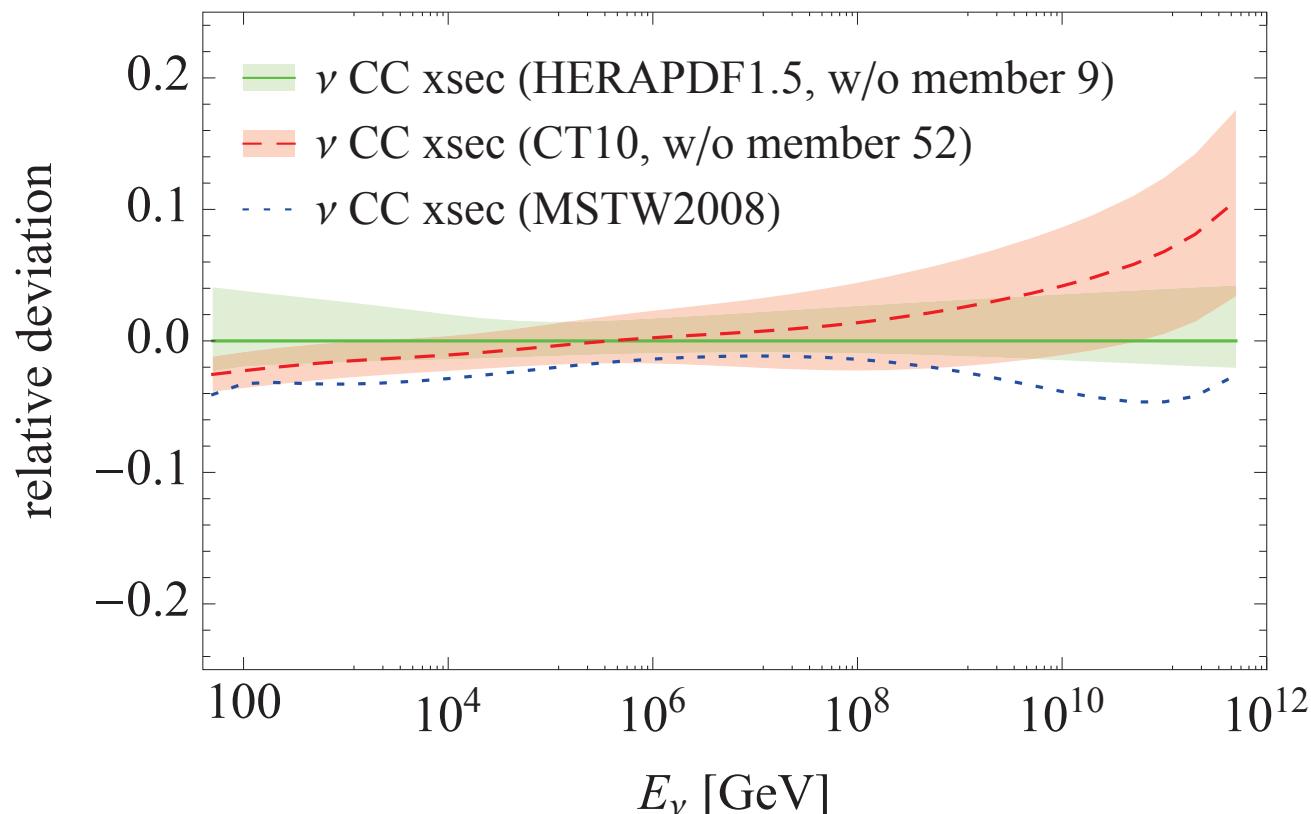


member 52 put in by hand

ν CC cross-section (excluding rogue members)



ν CC cross-section uncertainty (excluding rogue members)



Using these results . . .

more details in our paper

Cooper-Sarkar, Mertsch and Sarkar, arXiv:1106.3723

- more results: ν and $\bar{\nu}$, CC and NC
- comparison with event generators and other calculations
- tabulated total cross sections and uncertainties

future plans

updating the widely used event generator ANIS with the new cross-sections

A. Gazizov, M. P. Kowalski, Comput. Phys. Commun. **172** (2005) 203.

gluon at low x

$\ln 1/x$ resummation

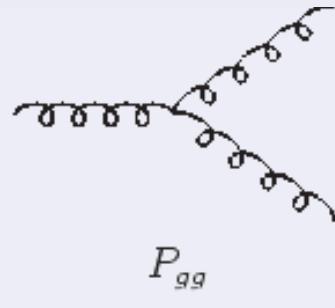
- DGLAP contains terms $\sim (\alpha_s \ln x_0/x)^n$
 - at low x this becomes larger than 1
- need to resum $\ln 1/x$ terms

Froissart bound

- DGLAP predicts $xg \propto x^{-\delta}$ at low x
- → $\sigma \propto s^\delta$ at large s
- however, unitarity demands $s, \sigma \propto (\ln s/s_0)^2$ at most

non-linear effects

- DGLAP eqns. are linear
- however, in DGLAP gluon and sea quark density large at small x
→ gluon saturation? gluon recombination?



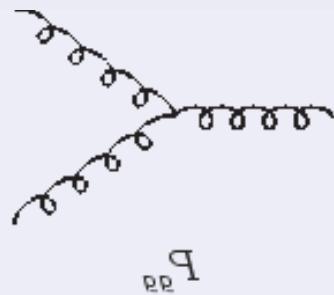
→ would tame the rise of σ

example

colour glass condensate

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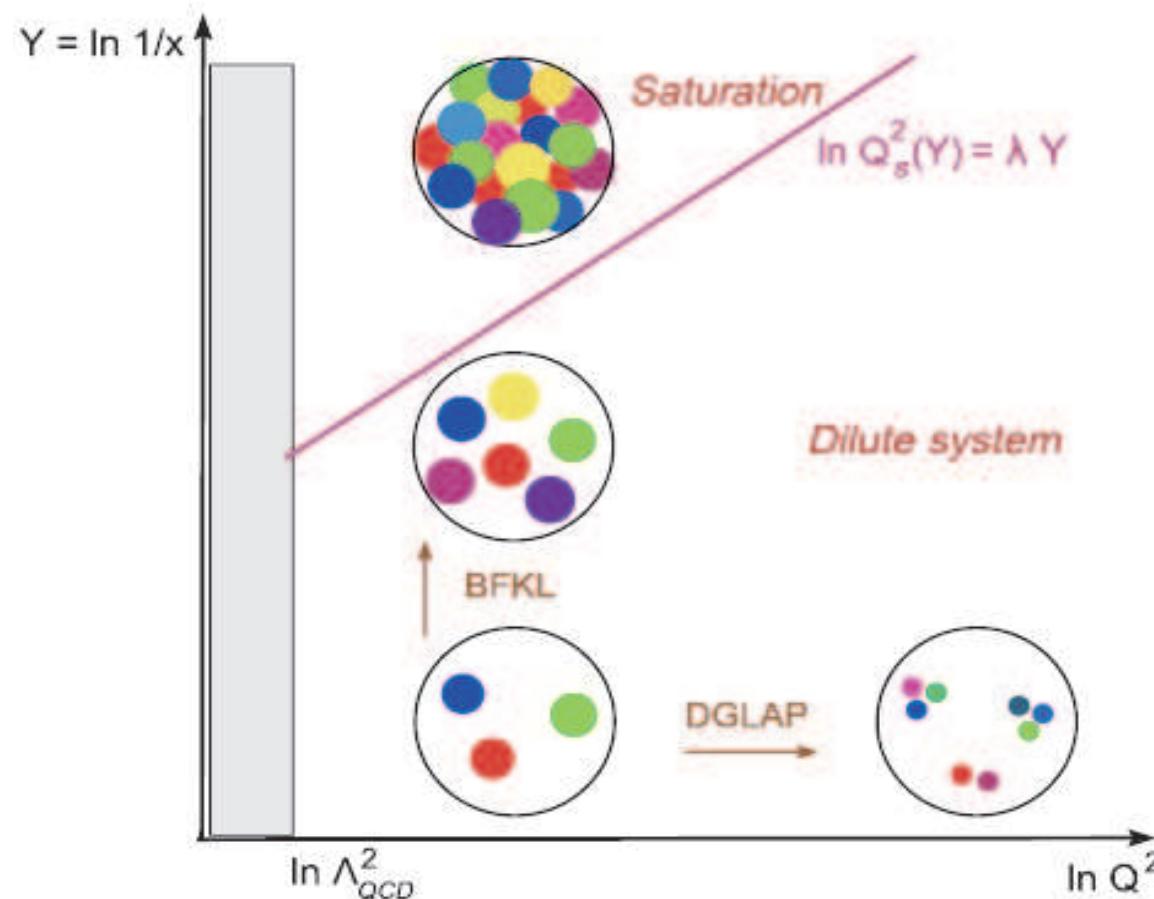


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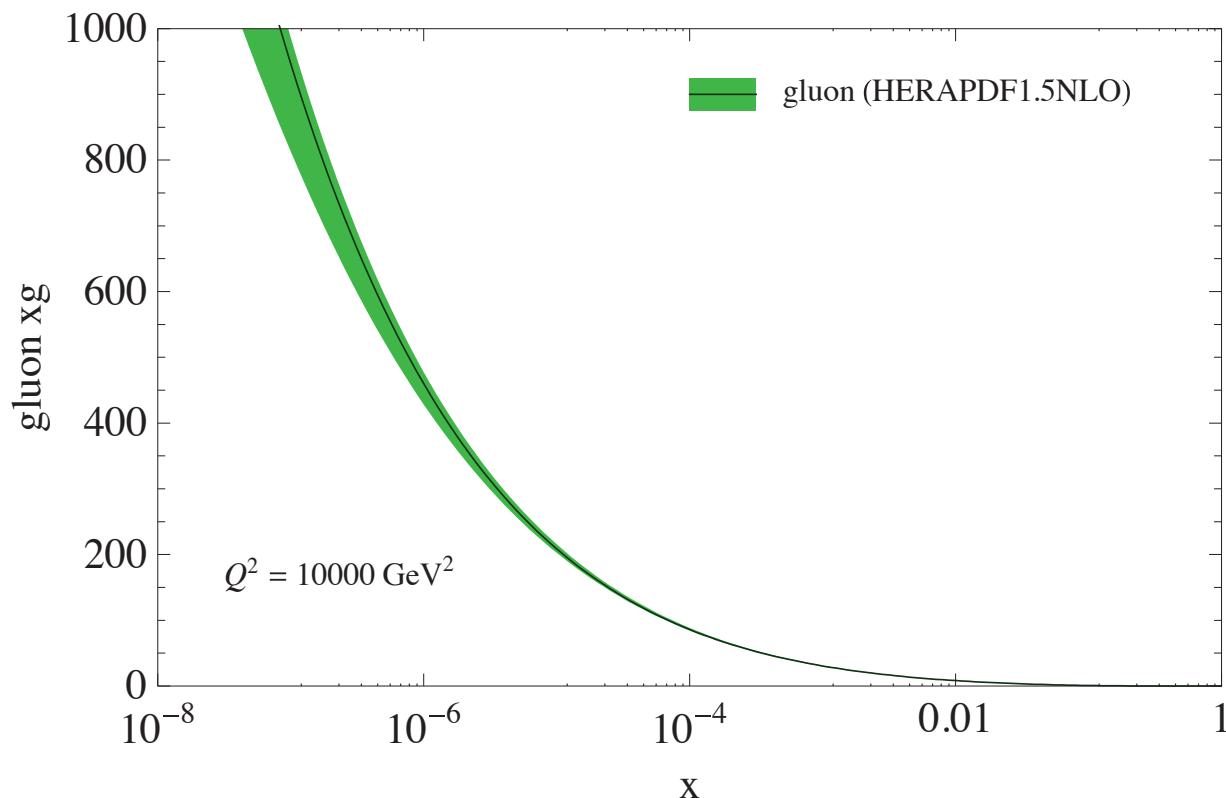


Conclusion

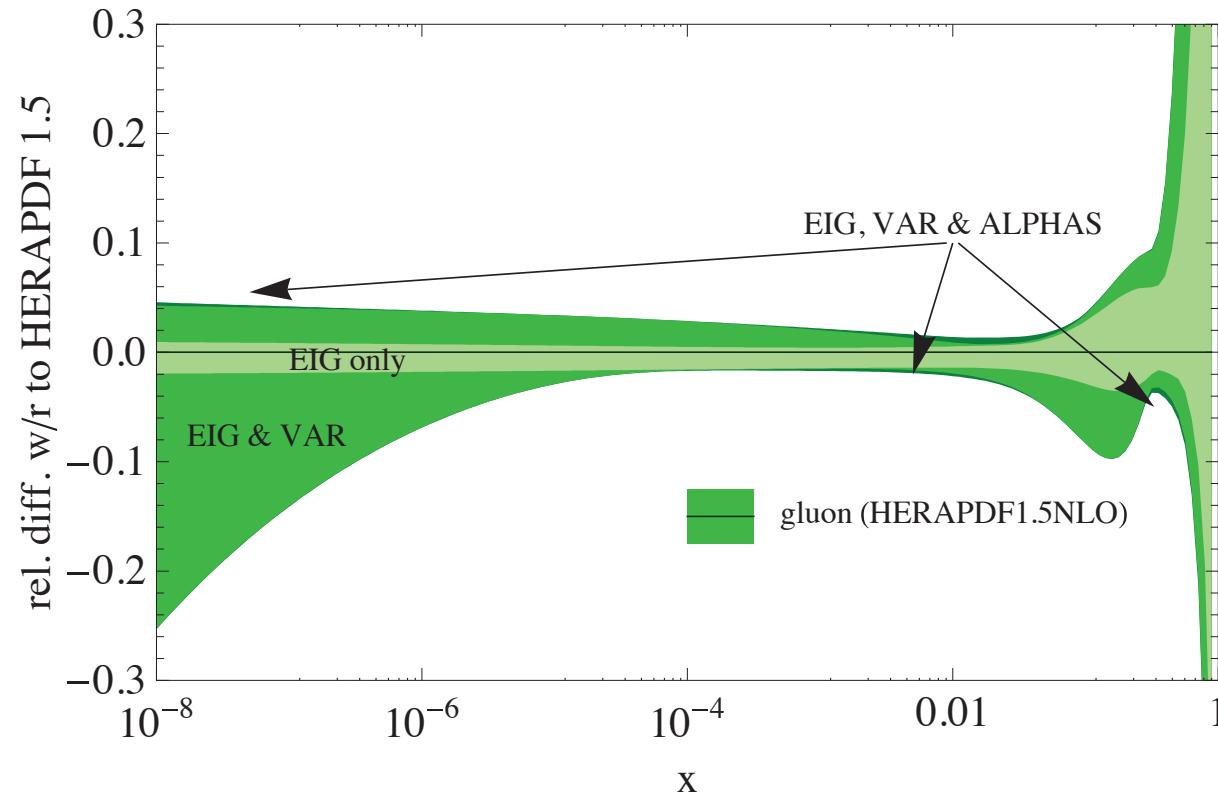
- cross-sections central values for
 - ▶ HERAPDF1.5
 - ▶ CT10
 - ▶ MSTW2008
- agree very well
- for HERAPDF1.5 and CT10 (under moderate assumptions) uncertainty is $\lesssim 10\%$, even at $E_\nu \sim 10^{20} \text{ GeV}$
- many pitfalls... e.g tabulated PDFs in LHAPDF “freeze” below some x value etc. → Don’t try to do this at home!
- Any measured deviation from these cross-sections would signal the need for new physics!

Backup slides

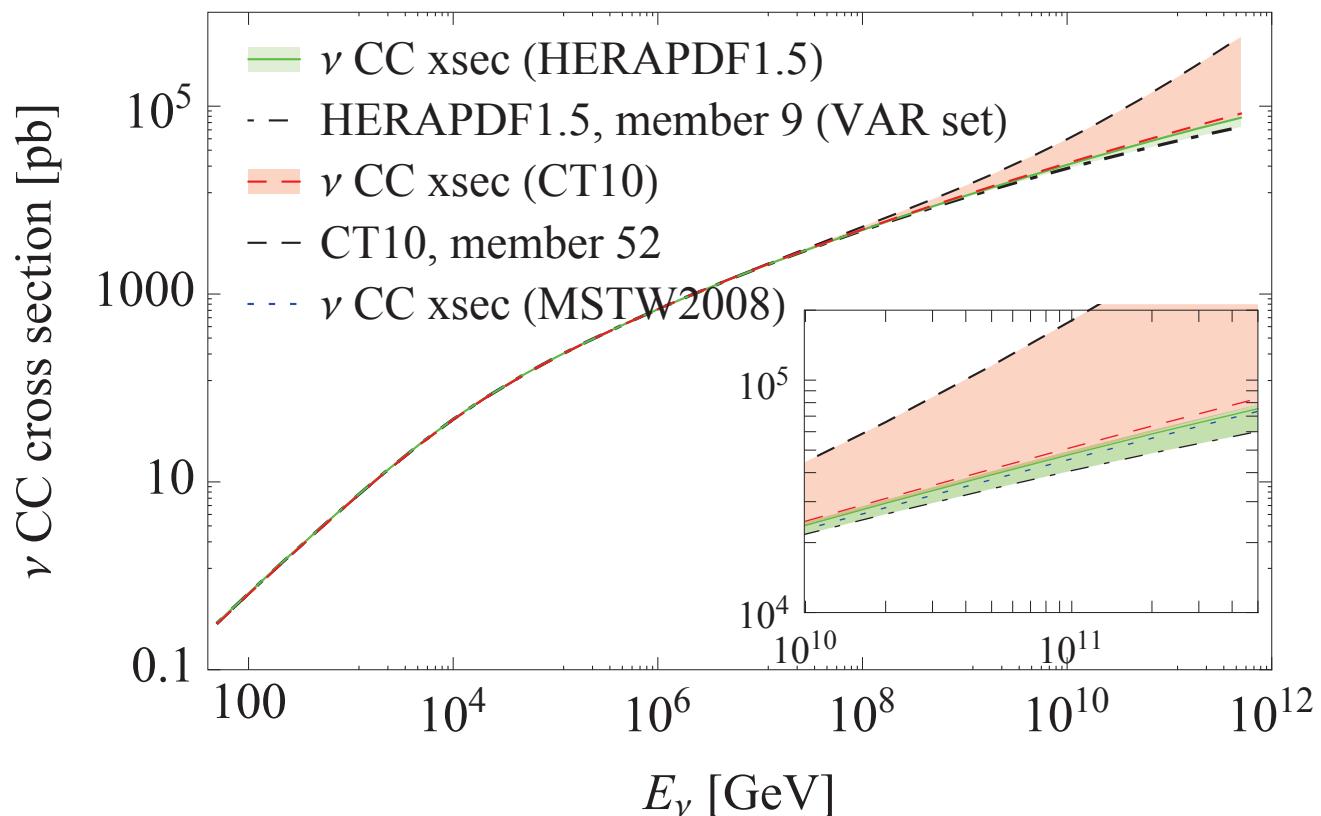
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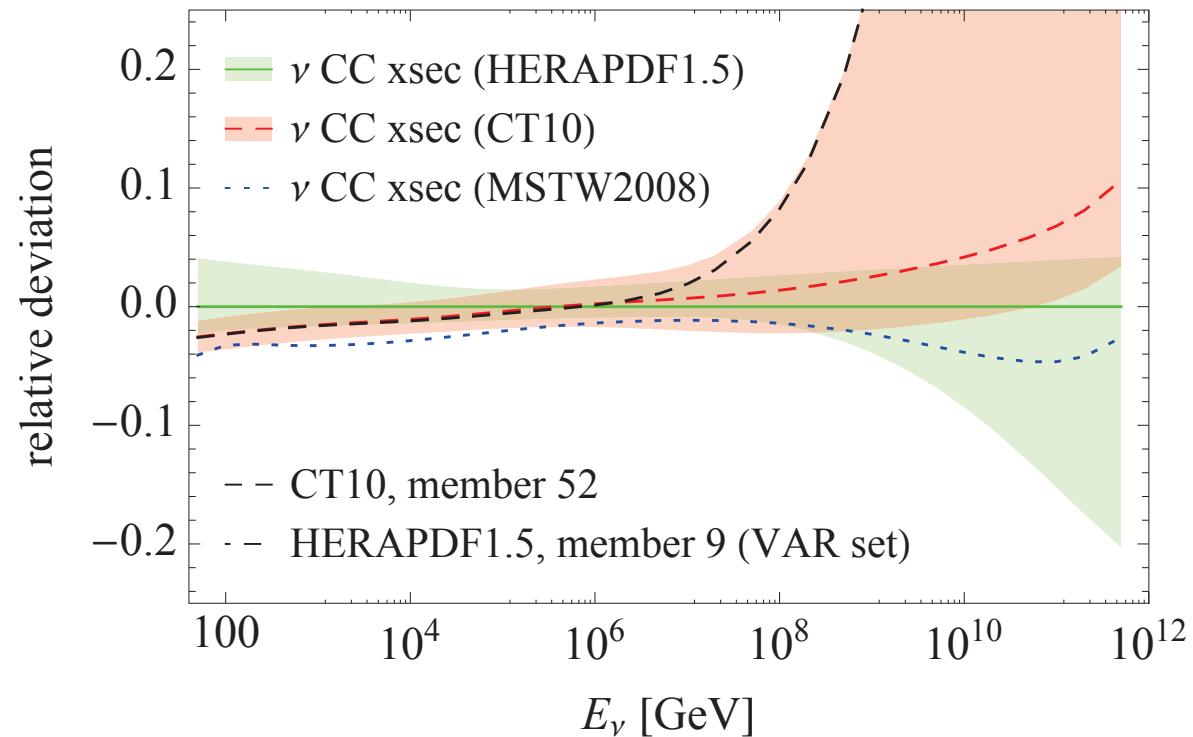
Example: HERAPDF1.5 gluon momentum distribution



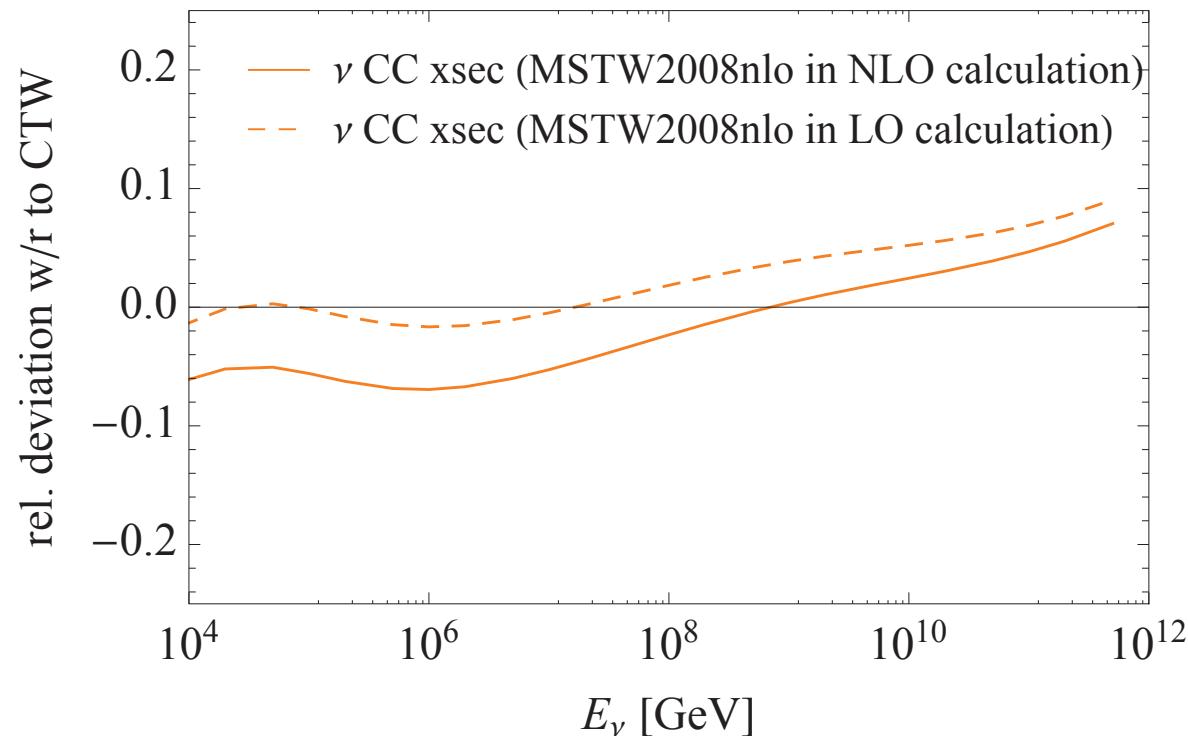
ν CC cross-section



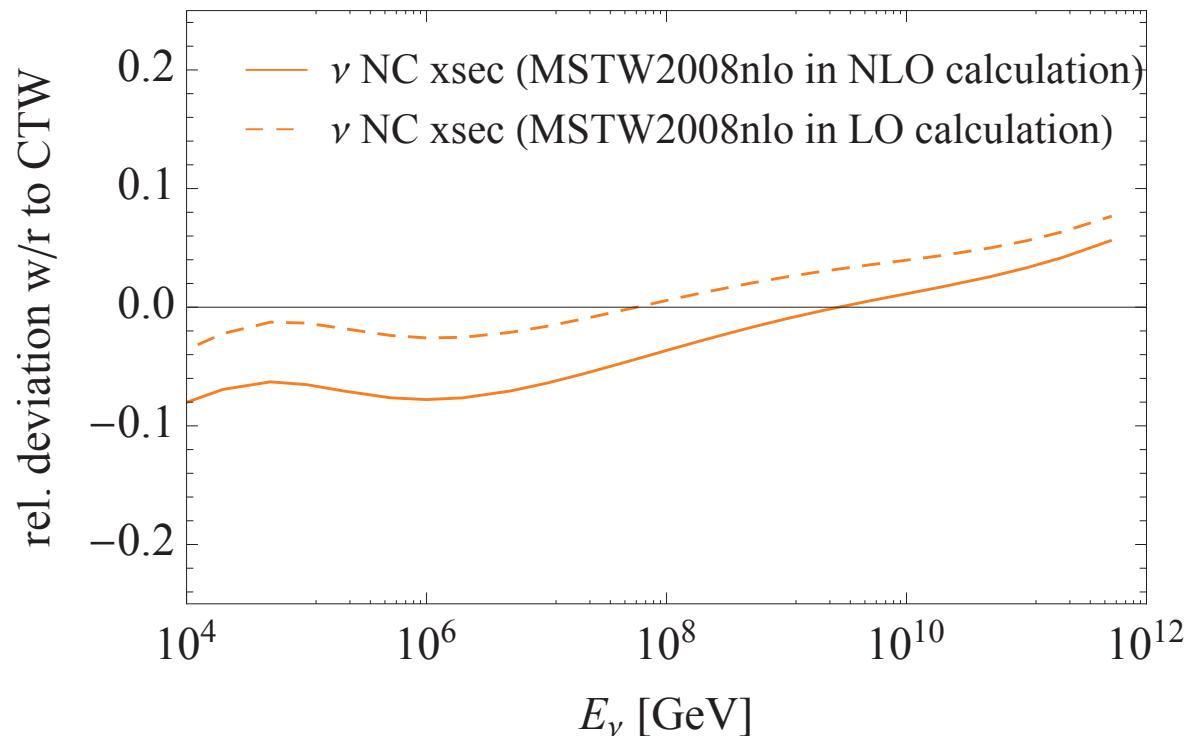
ν CC cross-section uncertainty



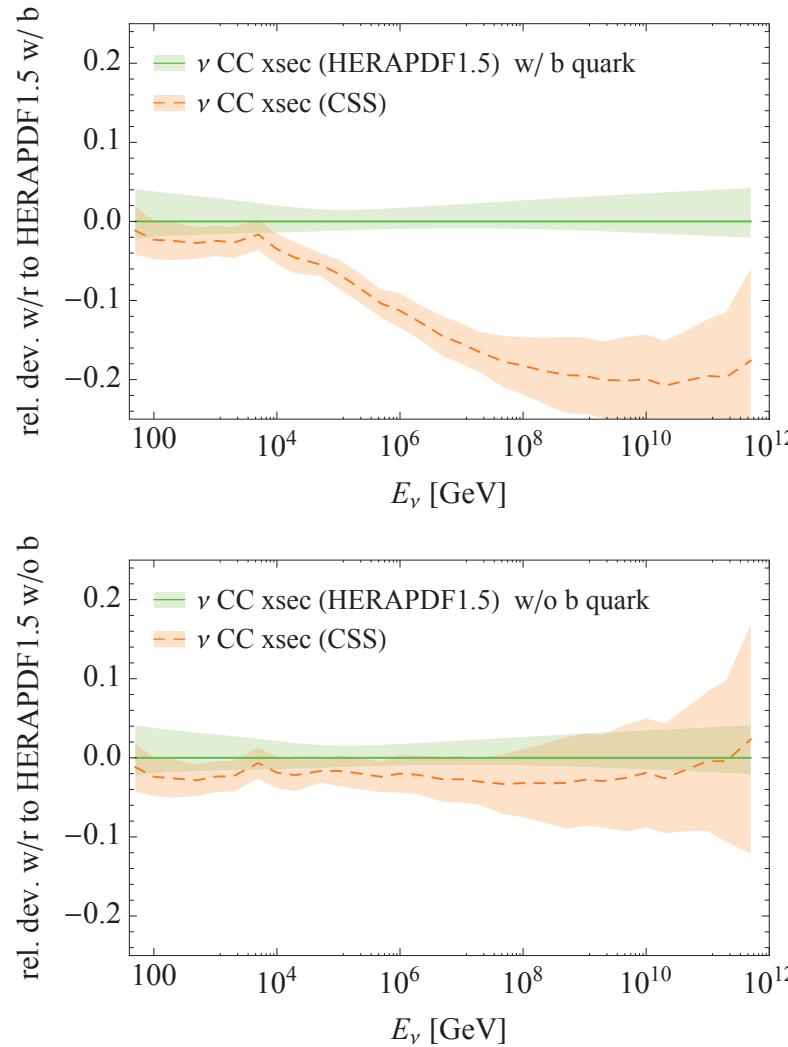
Comparison with CTW



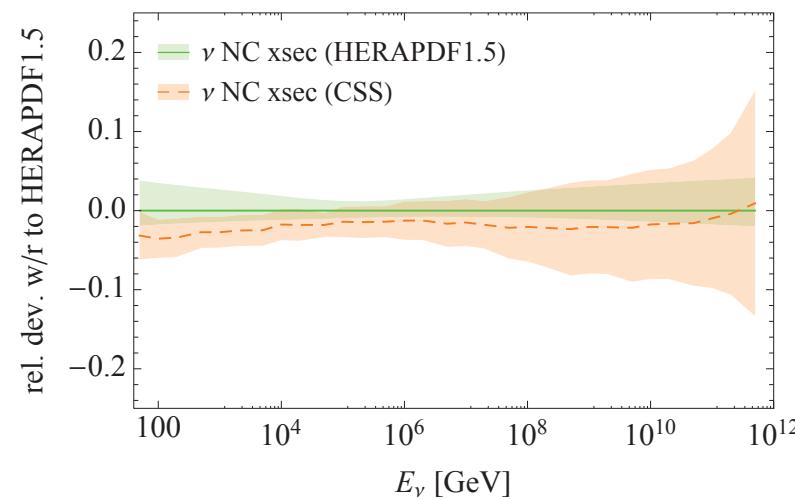
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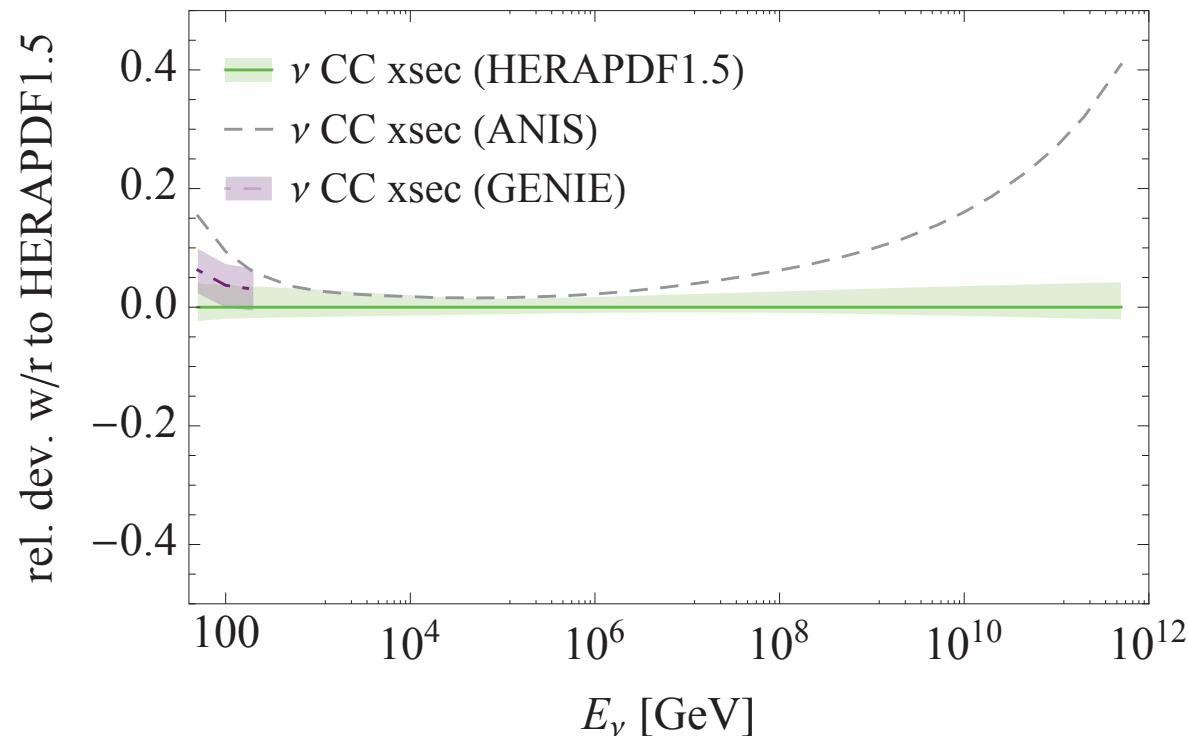
Comparison with CSS



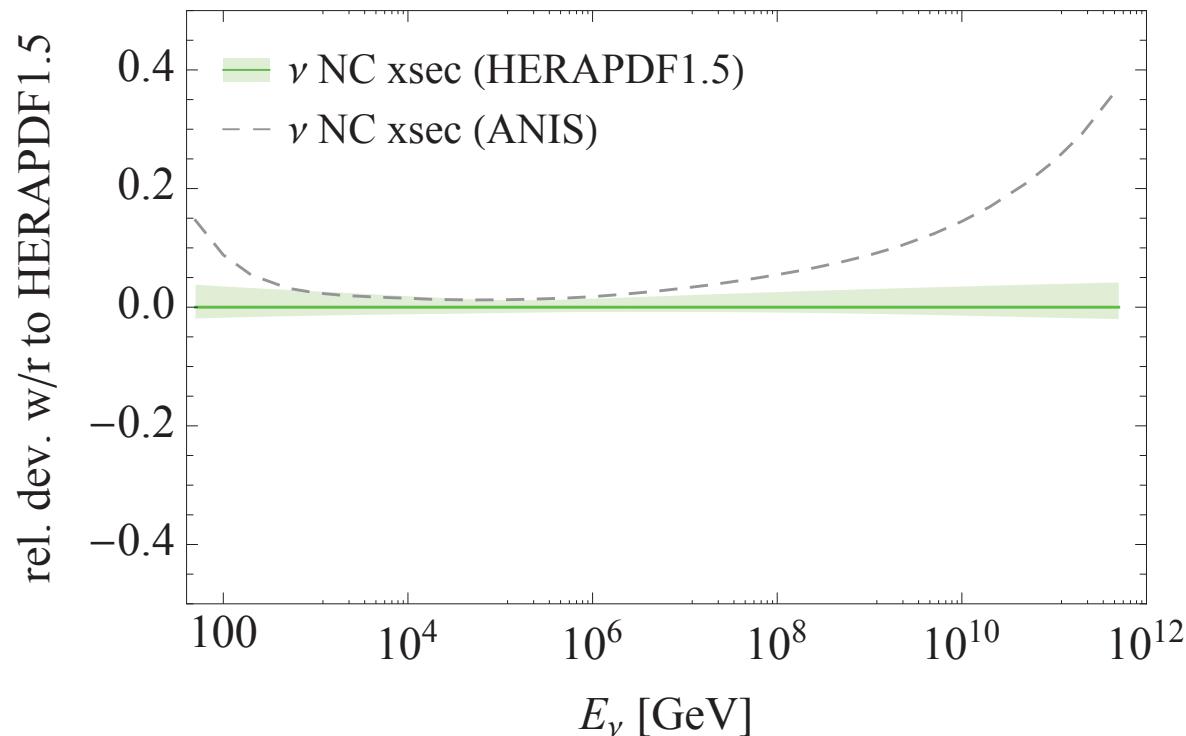
Comparison with CSS



Comparison with ANIS



Comparison with ANIS



The kinematic range

TBD: new figure w/o THERA!

