

Low-frequency variability of atmospheric planetary waves Part I: observations and mechanisms

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Fluid temperature



If the rotation of the convective cells is reversed, the upward heat flux is the same

Planetary waves in Jan: geop. height 500 hPa





Planetary waves in Jan: eddy T and v-wind 850 hPa





ICTP School/Conference on the General Circulation: Variability of the planetary waves



The El Niño - Southern Oscillation

Walker and Bliss (1932); Bjerknes (1969)



ECMWF

TENSO teleconnections: rainfall and temperature

WARM EPISODE RELATIONSHIPS DECEMBER - FEBRUARY

COLD EPISODE RELATIONSHIPS DECEMBER - FEBRUARY



ICTP School/Conference on the General Circulation: Variability of the planetary waves



The North Atlantic Oscillation

Walker and Bliss (1932) Van Loon and Rogers (1978)



Positive NAO phase



Negative NAO phase



Teleconnection patterns: Wallace and Gutzler 1981

The Pacific / North American (PNA) pattern: correlations



FIG. 16. As in Fig. 12, but for base grid points (a) 20°N, 160°W; (b) 45°N, 165°W; (c) 55°N, 115°W; (d) 30°N, 85°W.

Teleconnection patterns: Wallace and Gutzler 1981

The Pacific / North American (PNA) pattern: composites





Teleconnection patterns: Wallace and Gutzler 1981

Eastern Atlantic (EA) pattern: composites





Understanding teleconnections: Horel and Wallace 1981



Correlation of 700hPa height with a) PC1 of Eq. Pacific SST c) SOI index

Schematic diagram of tropical-extratropical teleconnections during El Niño





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Atmospheric response to near-equatorial heating



a) Near-surface wind and vertical motion

b) Near-surface wind and sea-level pressure

c) Motion in x-z plane along the Equator



Response to a tropical heat source (15N):





Understanding telecon.: Sardeshmukh and Hoskins 1988

Definition of Rossby wave source generated by near-equatorial heating

specification of the Rossby wave source S. If we are interested in the generation of Rossby waves, we are interested in the equation

$$\left(\frac{\partial}{\partial t} + \mathbf{v}_{\phi} \cdot \nabla\right) \boldsymbol{\zeta} = S + F. \tag{2}$$

Here v_{ψ} is the rotational wind associated with ζ . Comparing with (1) we see that we should write the total wind v in terms of its rotational and divergent components

$$\mathbf{v} = \mathbf{v}_{\psi} + \mathbf{v}_{\mathbf{x}}$$

where $\nabla \cdot \mathbf{v}_{\chi} = D$. Then (1) may be written in the form of (2) but with

$$S = -\mathbf{v}_{\chi} \cdot \nabla \zeta - \zeta D. \tag{3}$$

Understanding telecon: Simmons Wallace Branstator 1983









A sequence of 5-day mean fields of 500 hPa geopotential height during boreal winter ...









different winter !







Lorenz (1963) model of R-B convection

- $dX/dt = \sigma (Y X)$
- dY/dt = -XZ + rX Y
- dZ/dt = X Y b Z

Unstable stationary states

- X = Y = Z = 0
- $X = Y = \pm [b(r-1)] \frac{1}{2}, Z = r-1$





- ${\bf q}$: barotropic or quasi-geostrophic potential vorticity
- $\partial_t q = -V_{\Psi} \cdot grad q D(q q^*)$
- steady state for instantaneous flow (multiple equilibria): $0 = -V_{\Psi} \cdot grad q - D(q - q^*)$
- steady state for time-averaged flow (regimes): $0 = - \langle V_{\psi} \rangle \cdot grad \langle q \rangle - D (\langle q \rangle - q^*)$ $- \langle V'_{\psi} \cdot grad q' \rangle$



Multiple steady states of low-order barotropic model with wave-shaped bottom topography



FIG. 4. Streamfunction fields of the stable first mode equilibria of a topographically forced flow for $k = 10^{-2}$, $L/a = \frac{1}{4}$, n = 2, $h_0/H = 0.2$ and $\psi_A^* = 0.2$; for the spectral model above resonance (a) and slightly below resonance (b); and for the grid-point model above resonance (c) and slightly below resonance (d). The nondimensional topographic heights are shown with light lines; the contour spacing is 0.05 units, with negative regions shaded.



Orographically forced models:

• **Charney and Straus 1980**: Form-grad instability, multiple equilibria and propagating planetary waves in baroclinic, orographically forced, planetary wave systems

• **Charney, Shukla and Mo 1981**: *Comparison of barotropic blocking theory with observation*

• **Legras and Ghil 1985**: *Persistent anomalies, blocking and variations in atmospheric predictability*

• **Benzi, Malguzzi, Speranza, Sutera 1986:** The statistical properties of the atmospheric general circulation: observational evidence and a minimal theory of bimodality

Thermally forced models:

• **Mitchell and Derome 1983:** Blocking-like solutions of the potential vorticity equation: their stability at equilibrium and growth at resonance

Weather regimes: Reinhold and Pierrehumbert 1983

Hemispheric weather regimes arising from equilibration of large-scale dynamical tendencies and "forcing" from transient baroclinic eddies







500 hPa geop. height

4-8 Feb. 1986



4-8 Feb. 1989







• **Green 1977**: The weather during July 1976: some dynamical consideration of the drought

• **Illari and Marshall 1983**: *On the interpretation of eddy fluxes during a blocking episode*

• **Shutts 1986**: A case study of eddy forcing during an Atlantic blocking episode

• Haines and Marshall 1987: Eddy-forced coherent structures as a propotype of atmospheric blocking



Regional regimes: Vautard and Legras 1988

Regional weather regimes arising from equilibration of large-scale dynamical tendencies and PV fluxes from transient baroclinic eddies





Low-frequency variability of atmospheric planetary waves Part II: impact on heat fluxes and transport

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Net heat flux at ocean surface from ERA-interim









The El Niño - Southern Oscillation



Variability in the phase of planetary waves with respect to the surface temperature distribution in winter

- Mitchell and Derome (1983)
- Shutts (1987)
- Marshall and So (1990)
- Equilibrated phase: planetary waves ridges with higher temperature in the lower-troposphere are located over the warm oceans; surface heat fluxes and friction are reduced
- Thermally forced phase: planetary wave throughs (cold air) located on the western side of the oceans, strong temperature advection compensated by increased surface fluxes

Empirical variability patterns: EOFs and COWL

COWL (cold ocean - warm land) index after Wallace et al. (1995)



• Thermal-balance Wave index (positive in COWL phase):

Zonal wanenumber-2 component of net surface heat flux (NSHF, positive upward) in the [45-70 N] latitudinal band: TW = NSHF [60W-30E] - NSHF [30E-120E] + NSHF [120E-150W] - NSHF [150W-60W] -

Molteni, King, Kucharski and Straus, *Climate Dynamics 2010*



Regression of Z500/NSHF onto TW index



Z 500 hPa JF 1950-2002

NCEP/NCAR Re-Analysis



Net surface heat flux (+ upwards)

Is the TW mode relevant to decadal variability ?



Is it relevant to changes in ocean circulation?





NASA GISS surface air/sea temperature



Atmospheric variability in the last two decades





Thter-decadal SST variability in DJF



C Decadal variability in surface fluxes and heat content

upward HF JF 2000/09 – 1990/99

HC_300m JF 2000/09 – 1990/99

-3.6 -2.4 -1.8 -1.2 -0.6 -0.3 0.3 0.6 1.2 1.8 2.4 3.6

u-stress (*100) JF 2000/09 - 1990/99



Surface fluxes from ERA-Interim

Ocean heat content in the top 300m from NEMOVAR-S4

Heat flux and heat content: winter and yearly means







- Variability in planetary waves may be induced either by internal atmospheric dynamics (quasi-stationary responses to orographic and thermal forcing) or by teleconnections with coupled modes of tropical variability.
- Poleward heat transport by atmospheric planetary waves is not strongly dependent on the specific phase of waves.
- On the other hand, the amount of heat transferred from the ocean to the atmosphere in northern mid-latitudes is dependent on the phase of the planetary waves with respect to the land-sea distribution.
- The reversal of the phase of the COWL-like (TW) pattern in the last decade is consistent with reduced heat transfer in the northern oceans and increased ocean heat content in these regions.