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International Centre for Theoretical Physics**



**2252-S-3**

**Advanced Workshop on Non-Standard Superfluids and Insulators**

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**Disordered Commensurate Bosons**

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# Disordered Commensurate Bosons

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ICTP, Trieste, 19 July 2011

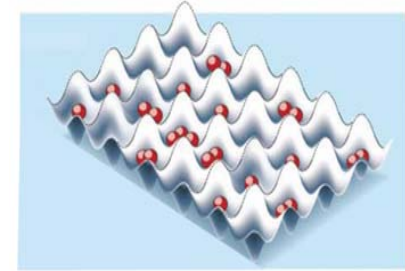


**ETH**



## Bose Hubbard model with bounded disorder at a commensurate filling

$$H = -t \sum_{\langle i,j \rangle} a_i^\dagger a_j + \frac{U}{2} \sum_i n_i (n_i - 1) + \sum_i \varepsilon_i n_i$$



$\varepsilon_i \in [-\Delta, \Delta]$  random on-site potential

$\nu \equiv \bar{n}_i = 1$  (or other integer)

Superfluid (SF)

Mott insulator (MI)    gapped insulator

Bose glass (BG)    compressible insulator

T. Giamarchi and H.J. Schulz,  
Europhys. Lett. **3**, 1287 (1987).

M.P.A. Fisher, P.B. Weichman,  
G. Grinstein, and D.S. Fisher,  
Phys. Rev. B **40**, 546 (1989).

**Q1:** Does disorder change the phase diagram at  $\Delta \ll U, t$  ?

**Q2:** Is disorder a relevant perturbation for SF-insulator transition?

## Experiment on disordered bosons in optical lattices

J. E. Lye, L. Fallani, M. Modugno, D. S. Wiersma, C. Fort, and M. Inguscio, Phys. Rev. Lett. 95, 070401 (2005).

L. Fallani, J. E. Lye, V. Guarrera, C. Fort, and M. Inguscio, Phys. Rev. Lett. 98, 130404 (2007).

L. Sanchez-Palencia, D. Clément, P. Lugan, P. Bouyer, G.V. Shlyapnikov, and A. Aspect, Phys. Rev. Lett. 98, 210401 (2007).

J. Billy, V. Josse, Z. Zuo, A. Bernard, B. Hambrecht, P. Lugan, D. Clement<sup>1</sup>, L. Sanchez-Palencia, P. Bouyer, and A. Aspect, Nature 453, 891 (2008).

G. Roati, C. D'Errico, L. Fallani, M. Fattori, C. Fort, M. Zaccanti, G. Modugno, M. Modugno, and M. Inguscio, Nature 453, 895 (2008).

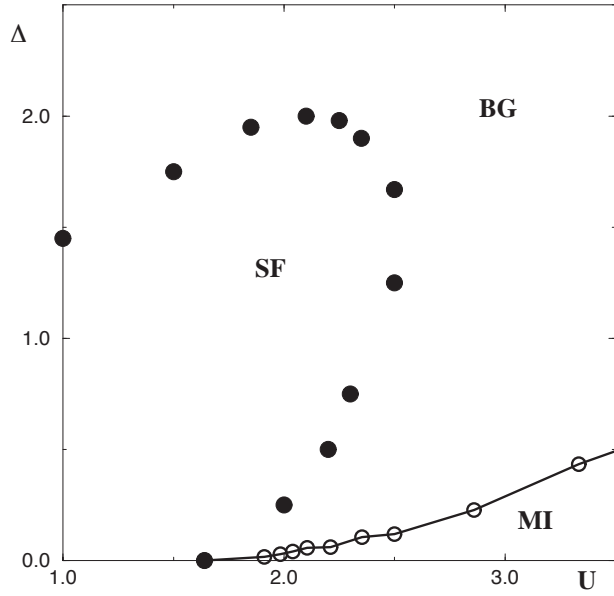
Y. P. Chen, J. Hitchcock, D. Dries, M. Junker, C. Welford, and R. G. Hulet, Phys. Rev. A 77, 033632 (2008).

M. White, M. Pasienski, D. McKay, S.Q. Zhou, D. Ceperley, and B. DeMarco, Phys. Rev. Lett. 102, 055301 (2009).

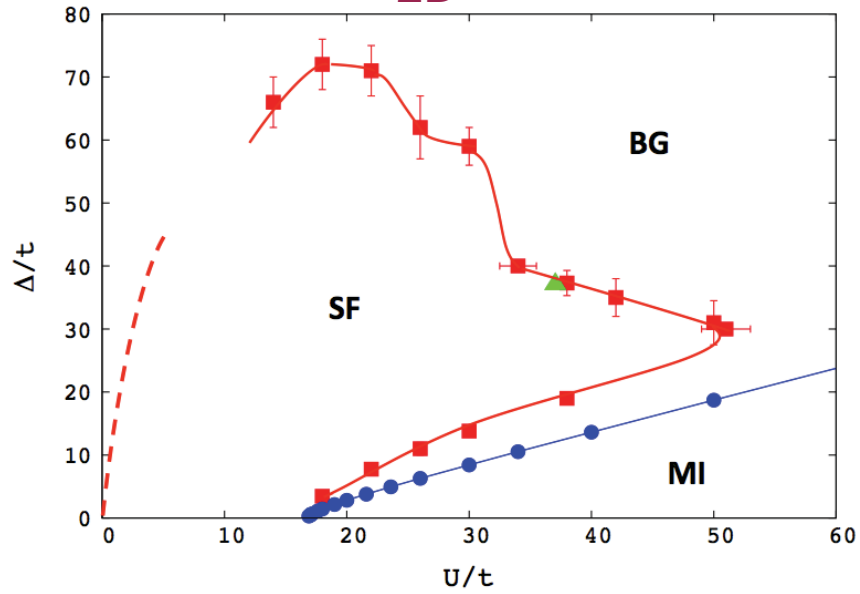
M. Pasienski, D. McKay, M. White, and B. DeMarco, Nature Physics 6, 677 (2010).

B. Deissler, M. Zaccanti, G. Roati, C. D'Errico, M. Fattori, M. Modugno, G. Modugno, and M. Inguscio, Nature Physics 6, 354 (2010)

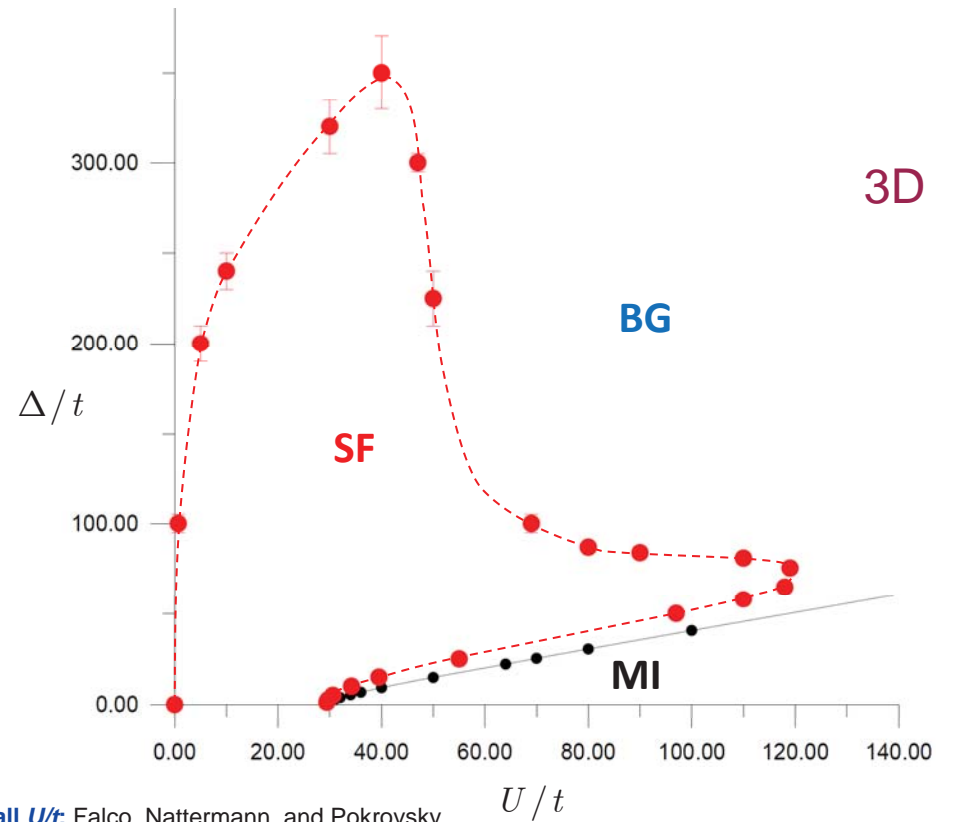
### 1D



### 2D

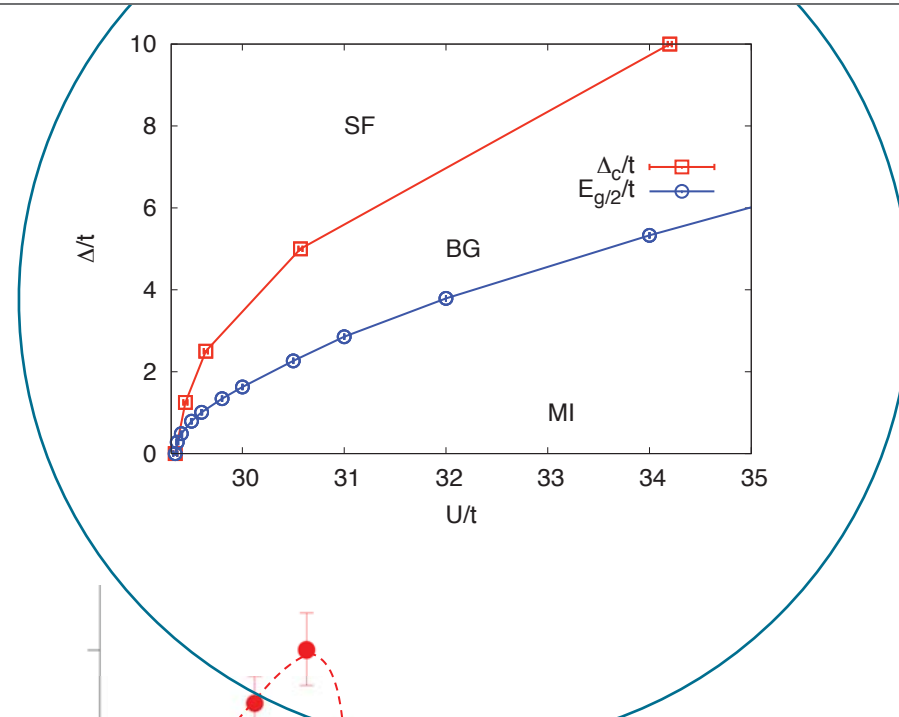
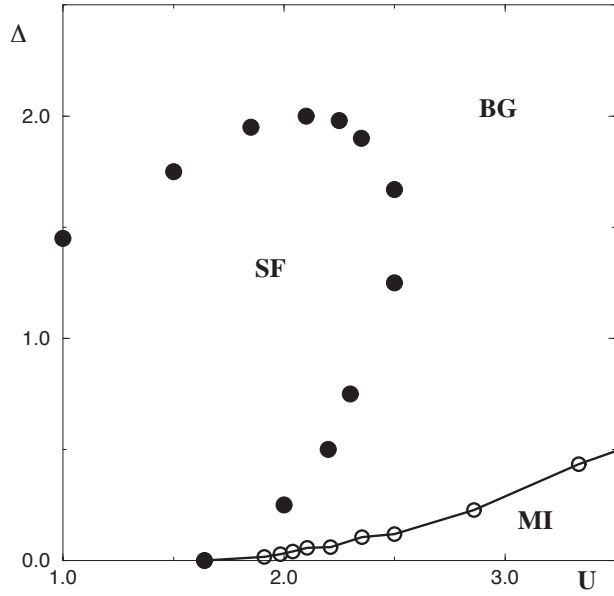


### 3D

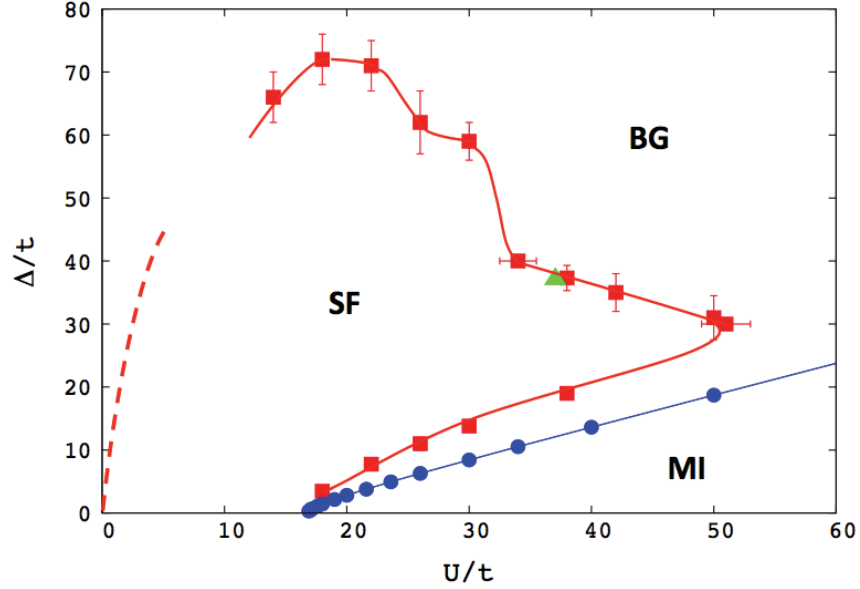


Small  $U/t$ : Falco, Nattermann, and Pokrovsky, Phys. Rev. B **80**, 104515 (2009).

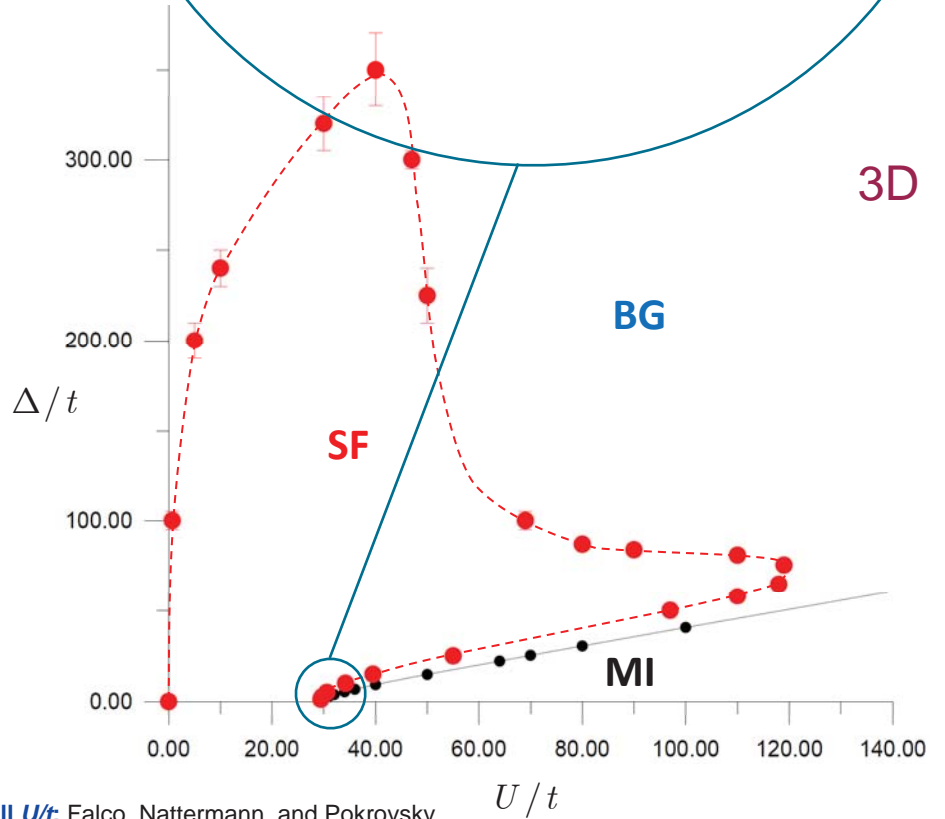
### 1D



### 2D

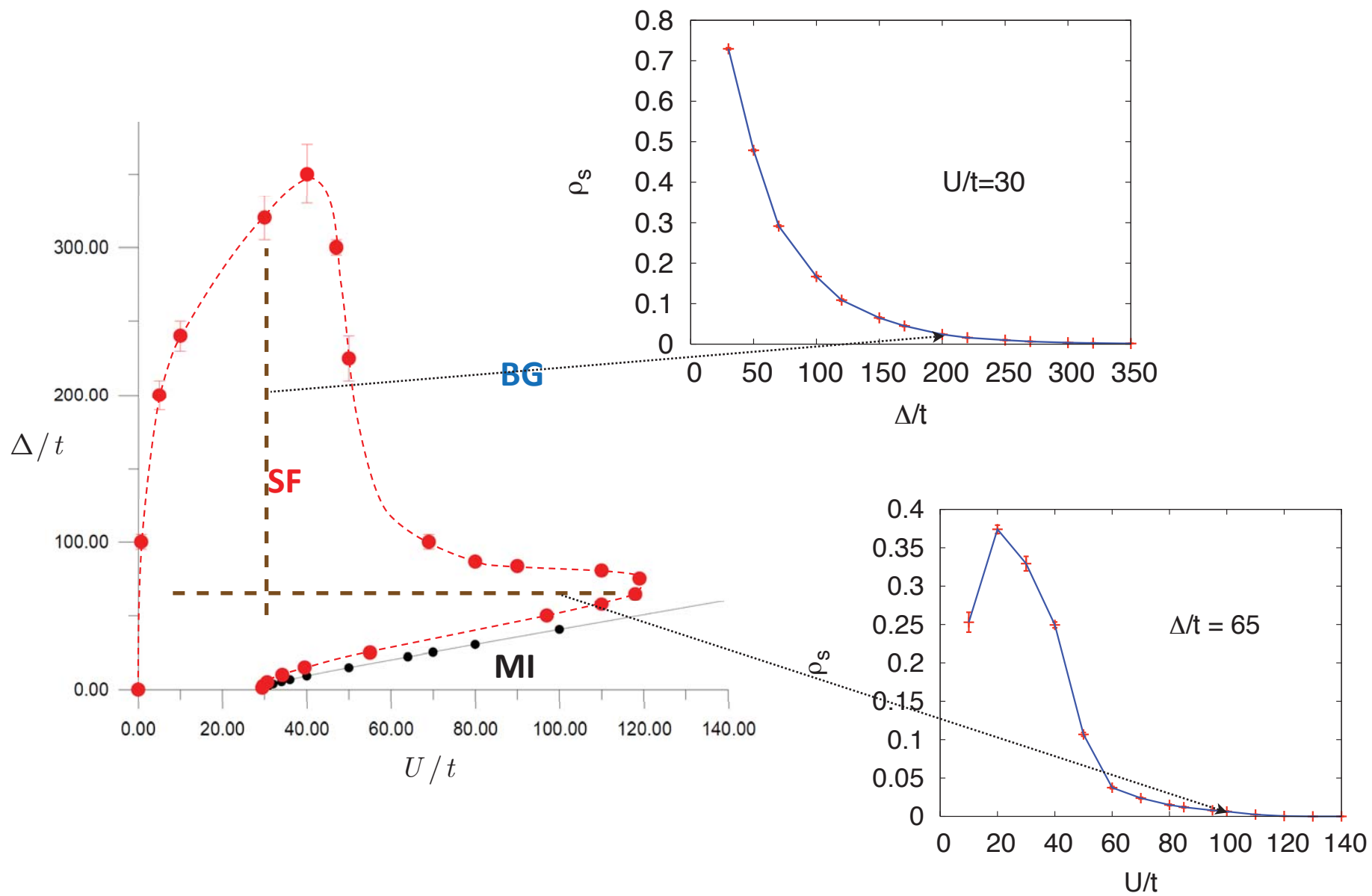


### 3D

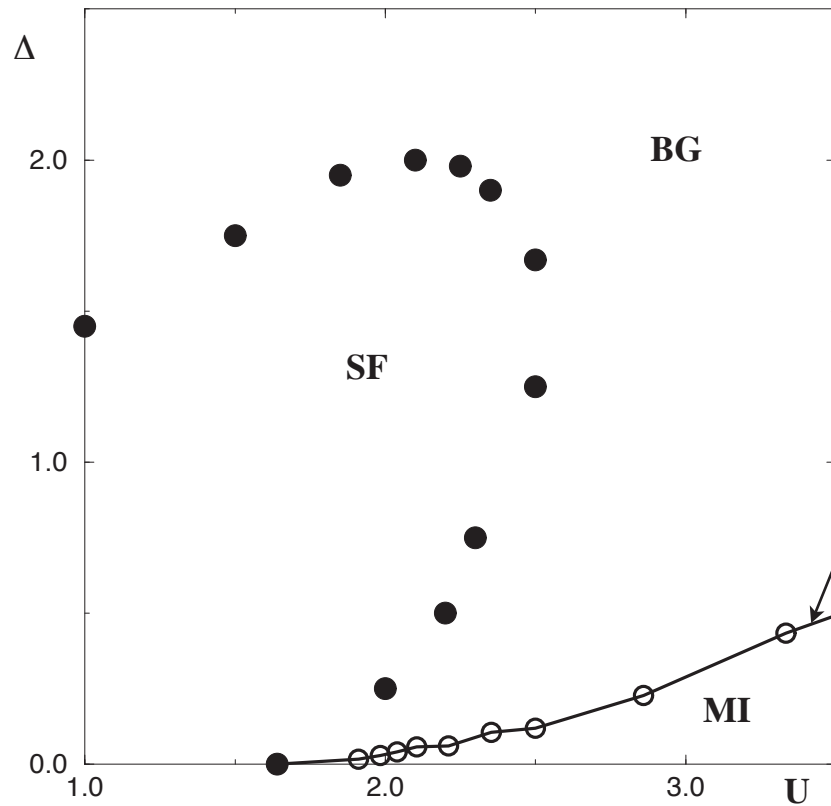


Small  $U/t$ : Falco, Nattermann, and Pokrovsky, Phys. Rev. B **80**, 104515 (2009).

Large numbers = fragile superfluidity !

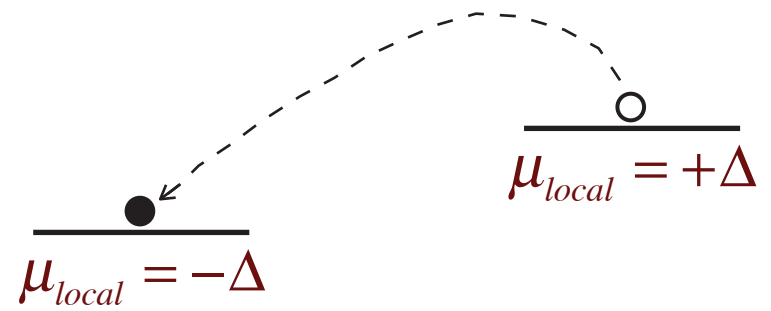


The line  $\Delta = E_{g/2}$



$E_{g/2}$  is the half of the gap in the *pure* system.

$$\Delta = E_{g/2}$$

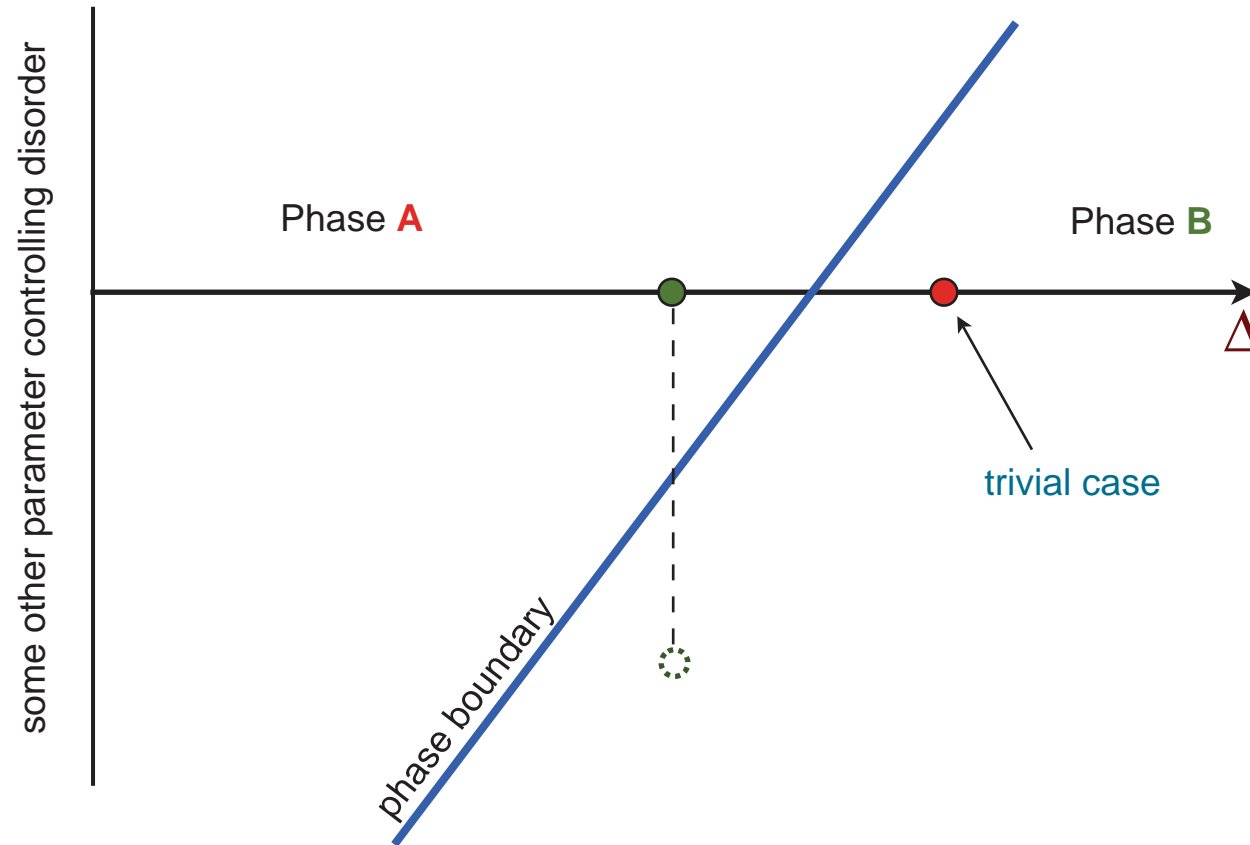




## Theorem of Inclusions

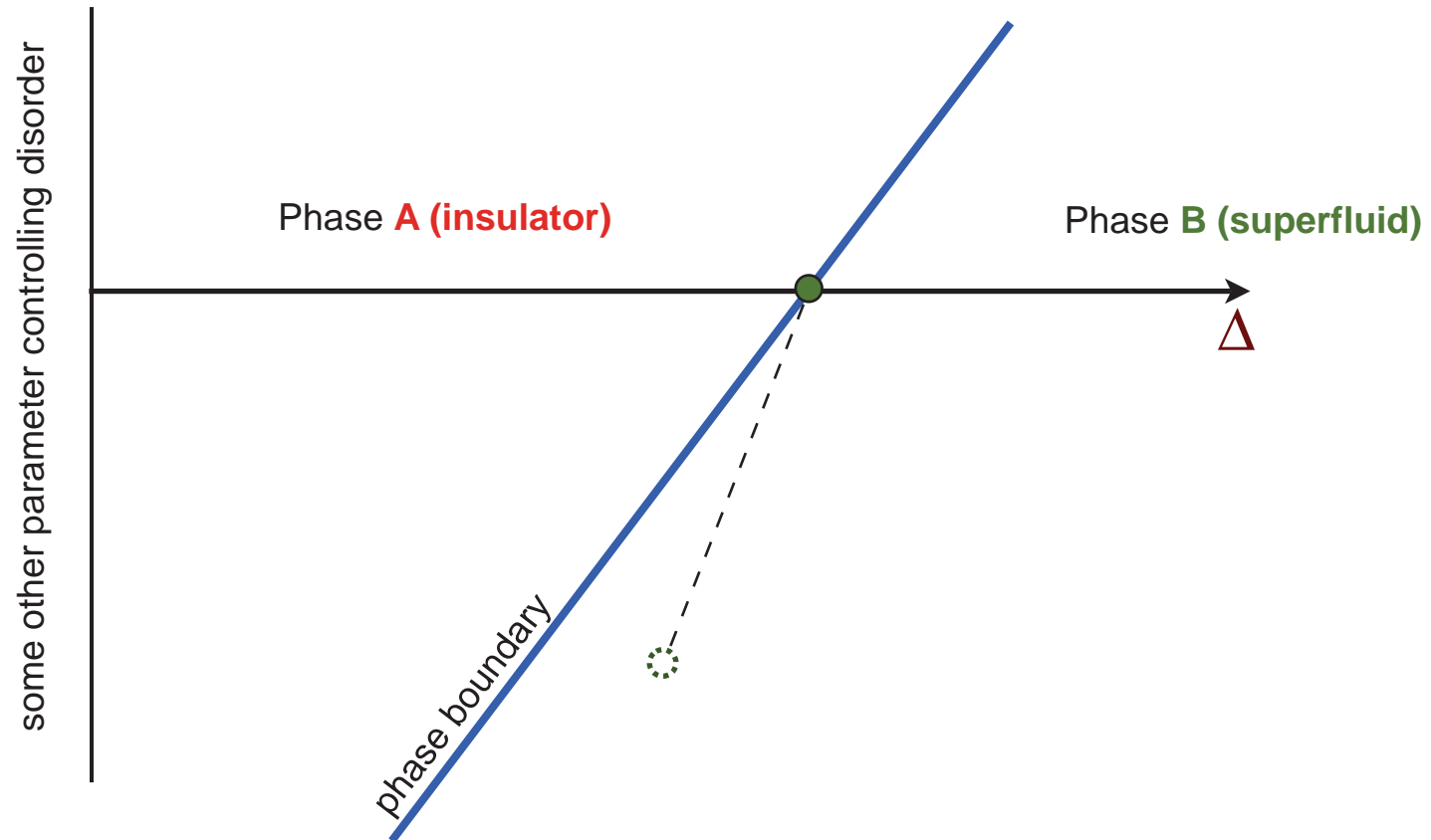
Def. ***Generic disorder***: Probability density for any (local) realization is finite.

## Theorem of Inclusions: Absence of SF-to-gapped-insulator transition



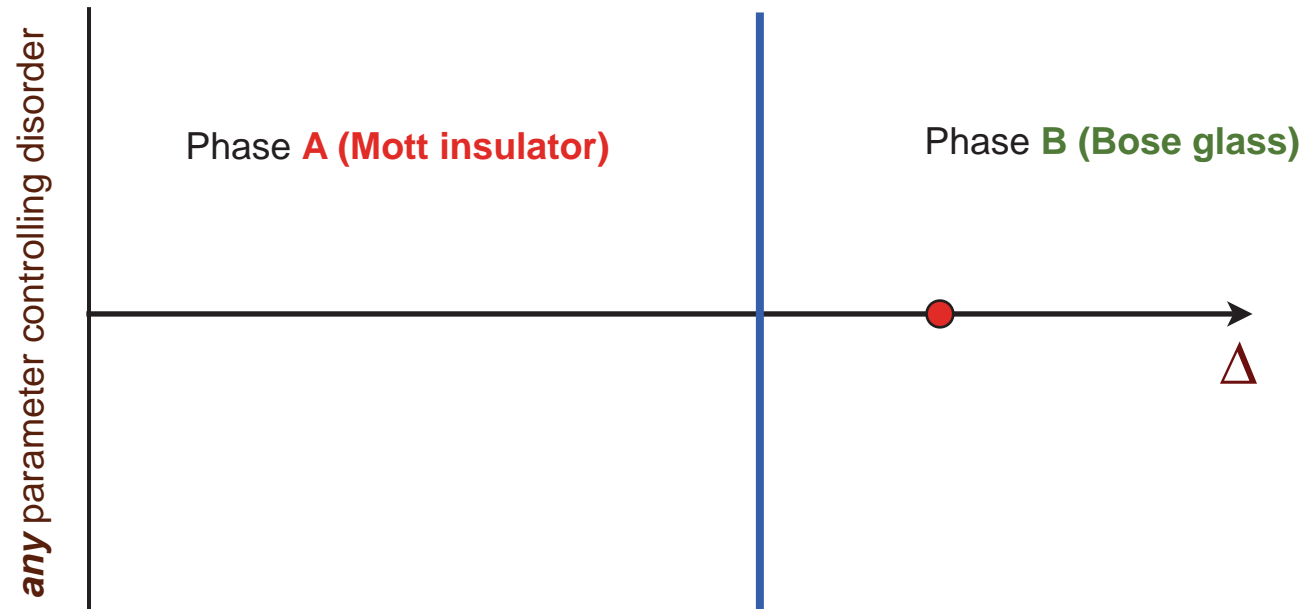
*In the presence of generic bounded disorder there exist rare, but arbitrarily large, regions of the competing phase across the transition line.*

## Enhanced version: Finite compressibility at the critical point



*There exist arbitrarily large regions with actual bound smaller than the global one by a finite value.*

An exception implied by the rule: MI to BG transition is of the Griffiths type



*The phase boundary line is generically vertical if, and only if, the transition is of the Griffiths type, i.e. driven by rare regions in which disorder emulates some regular external perturbation with the amplitude  $\Delta$ .*

Some important details on how disorder gets relevant in the vicinity of the SF-MI point

# 1D case: vortex instantons in (1+1) dimensions

$$S[\Phi] = \int dx d\tau \left\{ i\bar{n}(x) \partial_\tau \Phi + \frac{1}{2} \Lambda_s (\partial_x \Phi)^2 + \frac{1}{2} \kappa (\partial_\tau \Phi)^2 \right\}$$

$\kappa = dn / d\mu$   
compressibility

superfluid stiffness

$$S[\Phi] = \int dx dy \left\{ i\bar{n}(x) \partial_y \Phi + \frac{1}{2\pi K} (\nabla \Phi)^2 \right\}$$

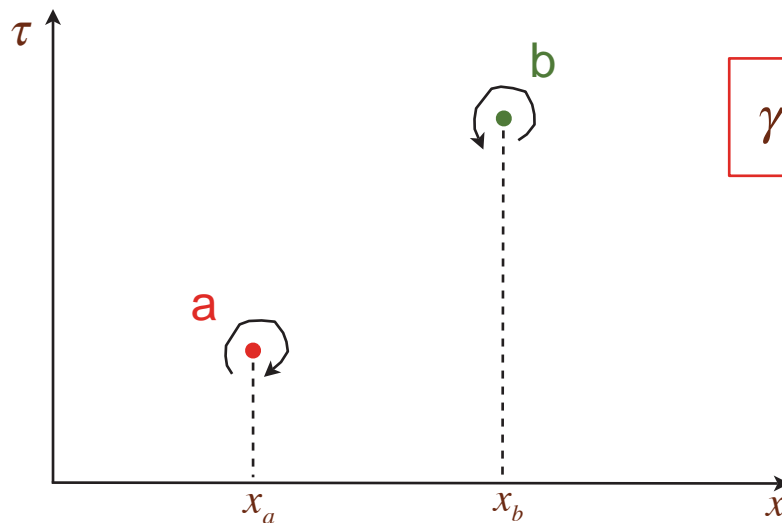
$$y = c\tau, \quad c = \sqrt{\Lambda_s / \kappa}, \quad K^{-1} = \pi \sqrt{\Lambda_s \kappa}$$

Luttinger-liquid parameter

$$S[\Phi] = \frac{1}{2\pi K} \int (\nabla \Phi)^2 dx dy - i \sum_j p_j \gamma(x_j)$$

$p_j = \pm 1, \pm 2, \pm 3, \dots$   
instanton charge

$\gamma(x) = 2\pi \int^x \bar{n}(x') dx'$   
instanton phase



$$\gamma_{ab} = 2\pi \int_{x_a}^{x_b} \bar{n}(x) dx$$

phase of the  
instanton pair

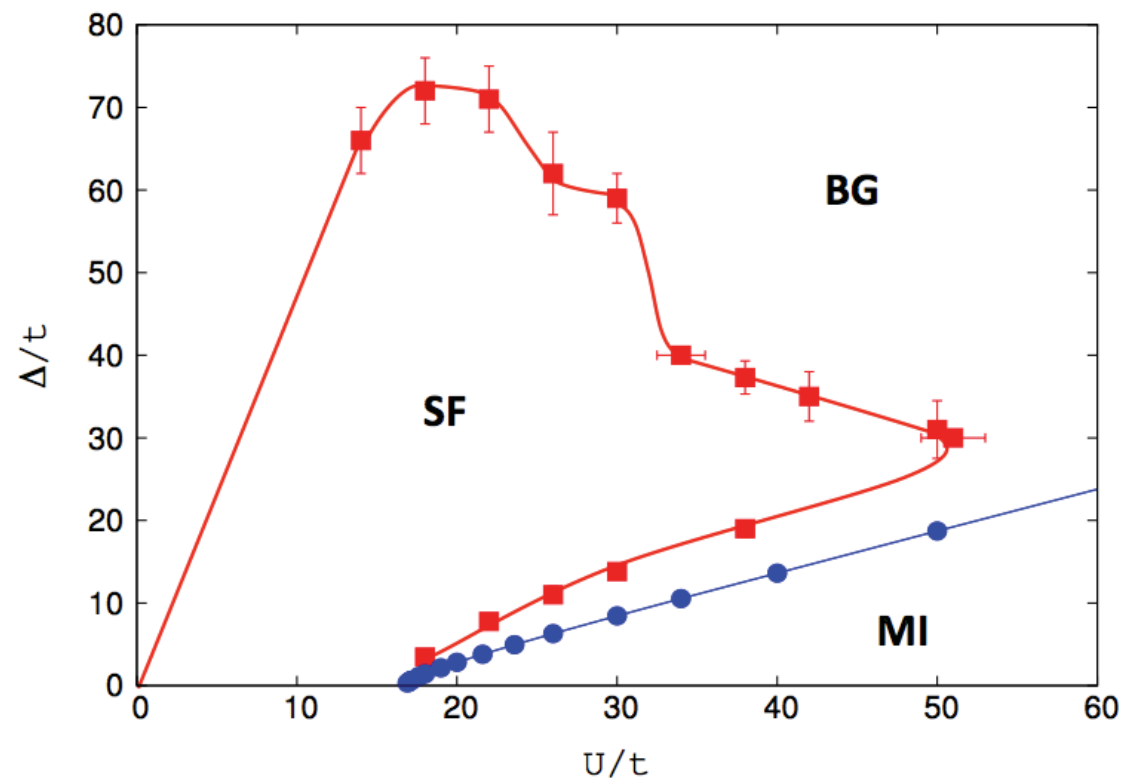
$$K_c = 1/2 \quad \text{SF-MI}$$

$$K_c = 2/3 \quad \text{SF-BG}$$

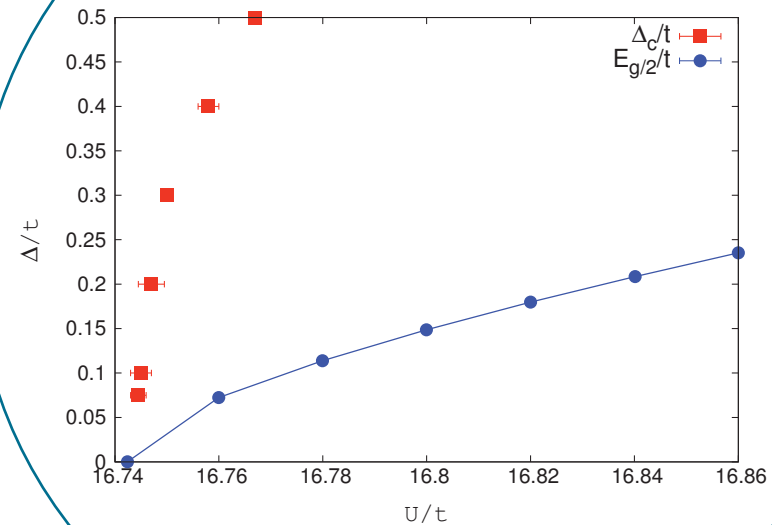
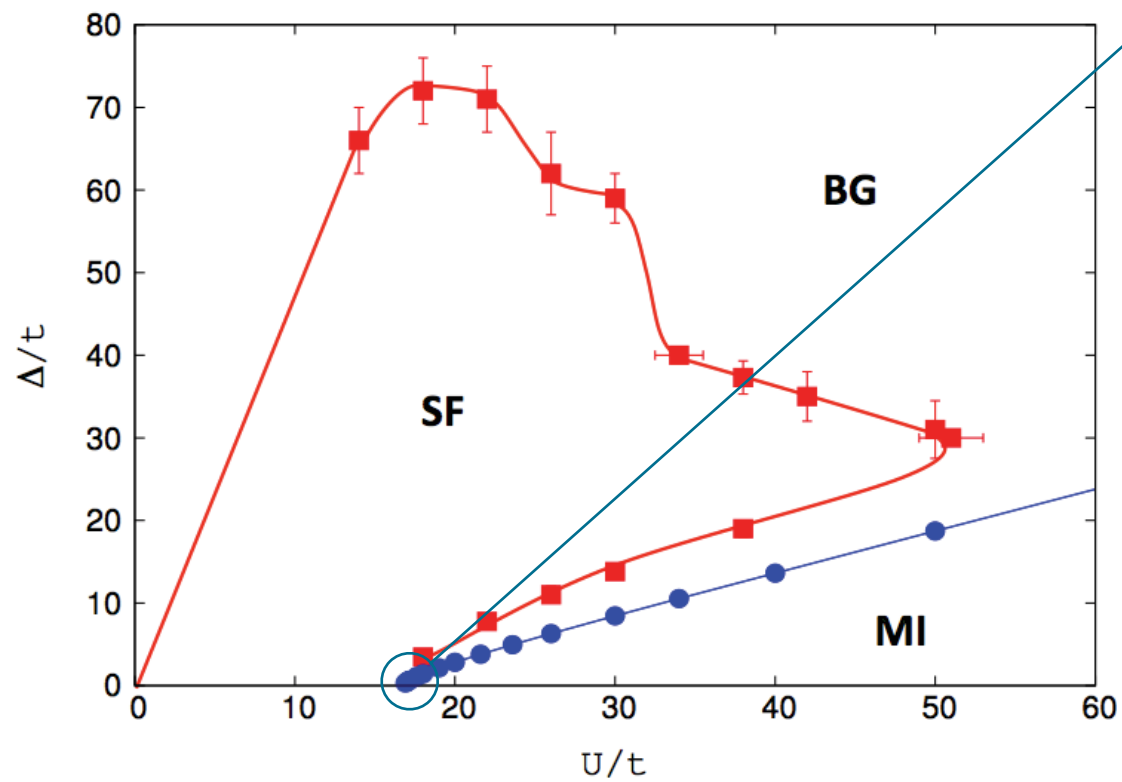
Shape of the SF-BG in the vicinity of the SF-MI point



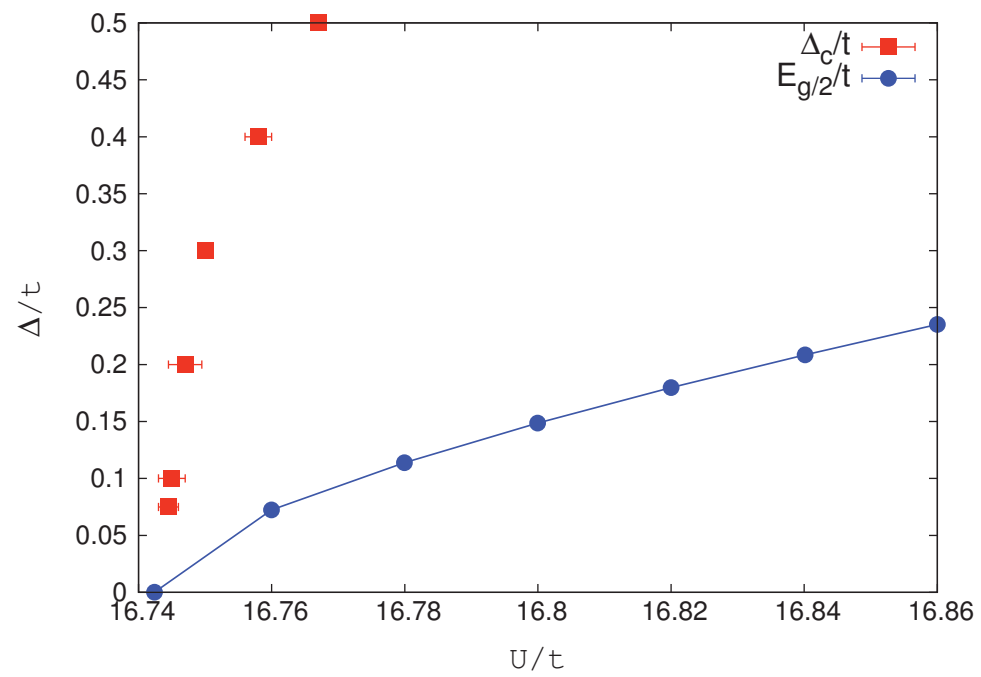
2D



2D



2D



## Main points

Exact numerics for Bose Hubbard model at unity filling in  $d=1,2$ , and 3.

Theorem of Inclusions: Absence of SF-to-gapped-insulator transitions

Enhanced version: Finite compressibility at the critical point

Extended version: MI to BG transition is of the Griffiths type

***Open question:*** Shape of the SF-BG curve in the vicinity of the SF-MI point in  $d=2,3$ .