



**The Abdus Salam  
International Centre for Theoretical Physics**



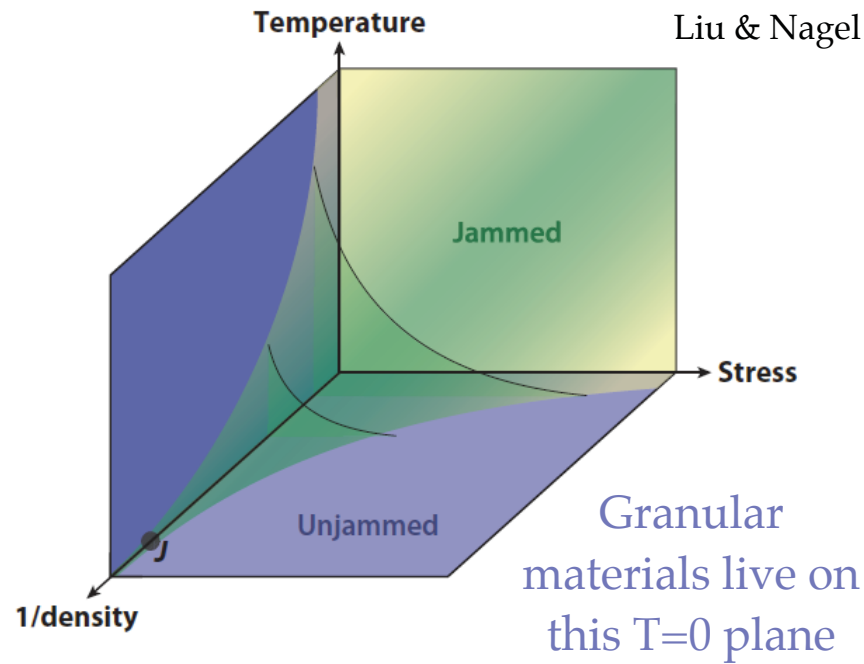
**2254-2**

**Workshop on Sphere Packing and Amorphous Materials**

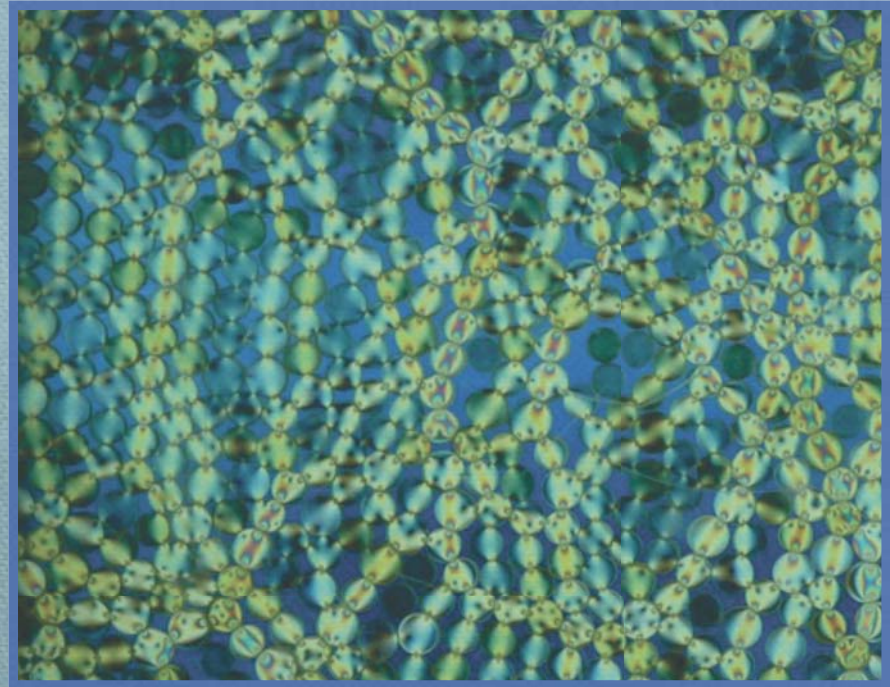
*25 - 29 July 2011*

**Growing Correlations at Unjamming of Sphere Packings**

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415 South Street  
Waltham, MA 02454  
U.S.A.*



J-Point is a non-equilibrium critical point

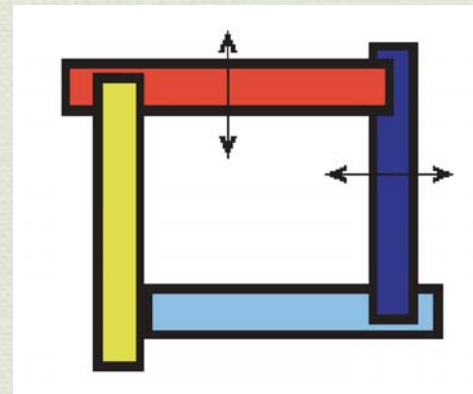


# Growing Correlations at Unjamming of sphere packings

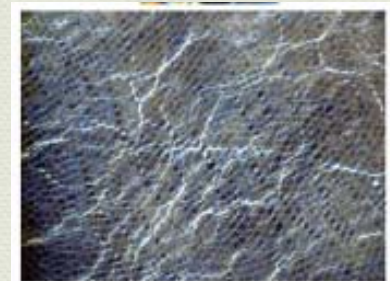
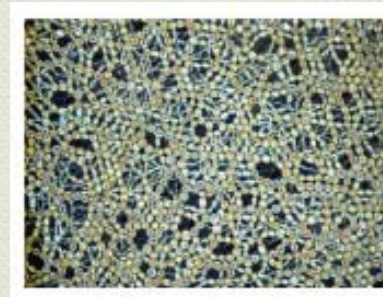


# Dry Granular Materials

- ❖ No cohesive interaction
- ❖ Grains are nearly rigid
- ❖ Unjammed side: Hard spheres
- ❖ Jamming : Increasing density leads to states that are jammed
- ❖ Deformable grains beyond Point J: states of mechanical equilibrium
- ❖ Unjamming: Falling apart because there is not enough imposed stress to impart rigidity



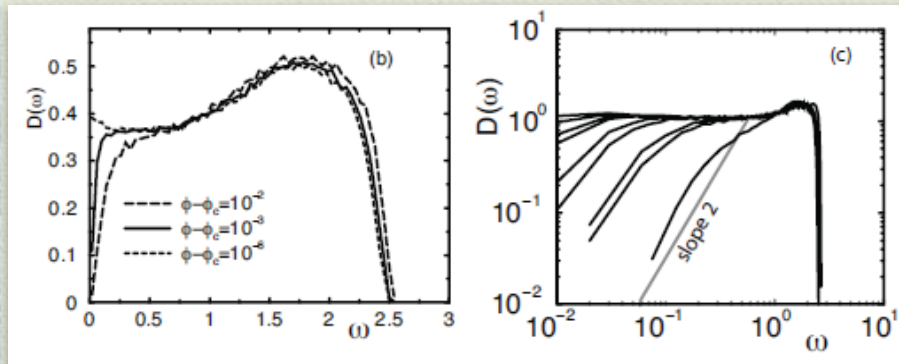
Put in  $N$  grains  
Change shape  
and/or size of box





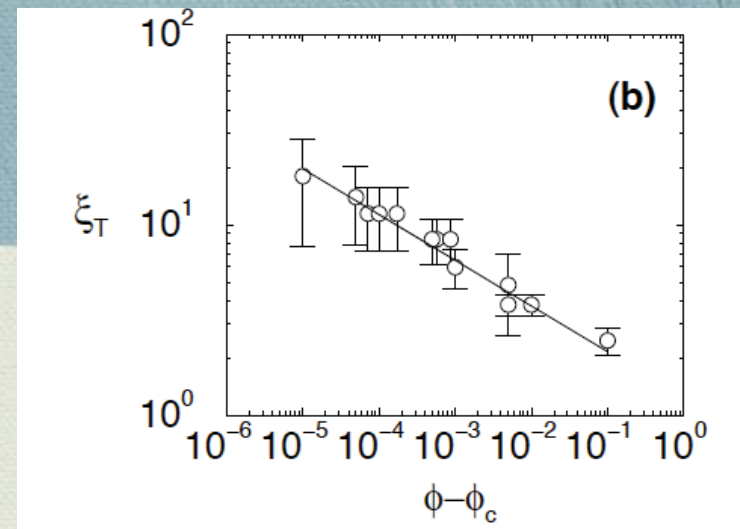
# Length Scales

## Soft Modes and $l^*$



O'hern et al, PRE 68 (2003) 01130

Silbert, Liu & Nagel, PRL 95, 098301 (2005)

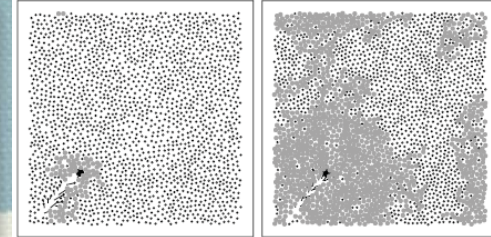
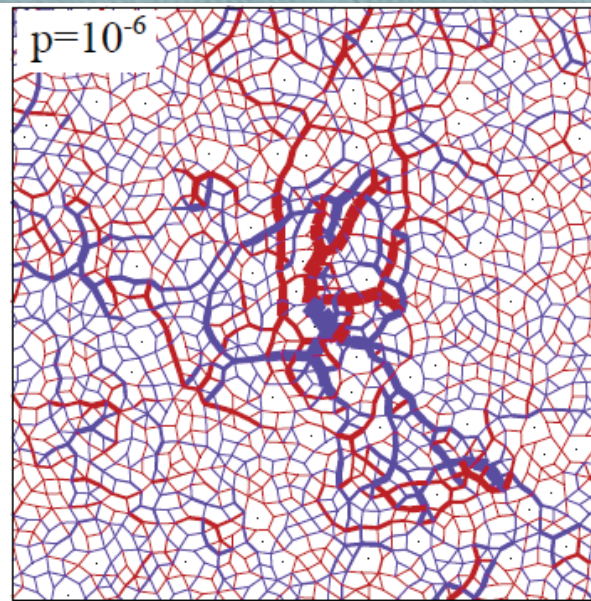
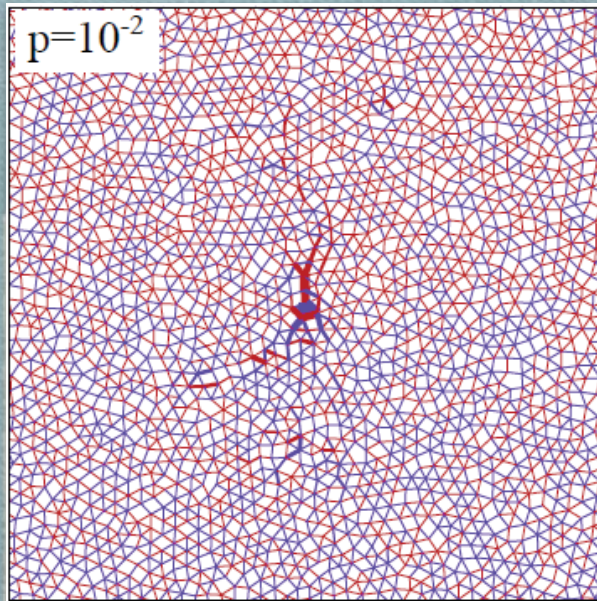


Exponent =  $1/4$

- ▶ Onset of unjamming associated with emergence of “anomalous modes.”
- ▶ Plateau defines crossover frequency  $\omega^*$ .
- ▶ Bulk-surface argument:  $\Rightarrow l^* \sim 1/\delta z \Rightarrow$  Exponent =  $1/2$  (M. Wyart, T. Witten)

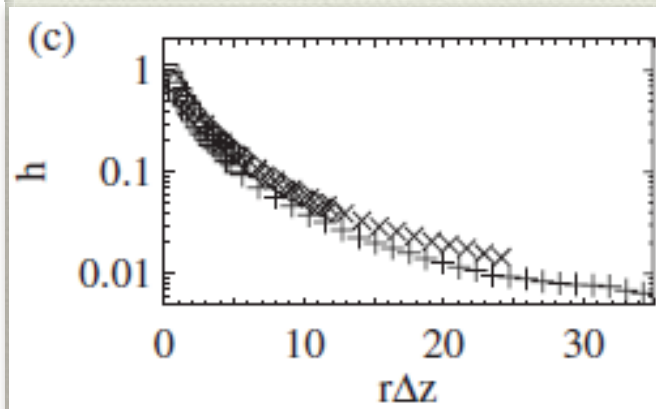
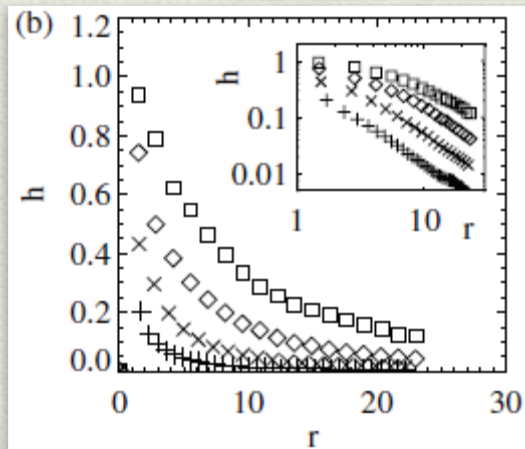


# MEASURING LENGTH SCALES FROM RESPONSE



J.A. Drocco et al., PRL **95** 088001 (2005)

A growing length scale is observed as a result of perturbations in simulation.



*$l^*$  scaling  
verified from  
force response*

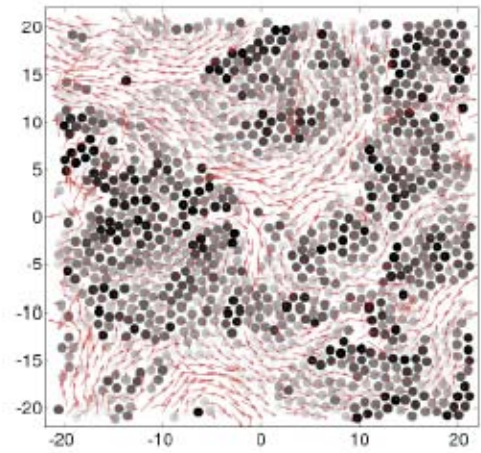
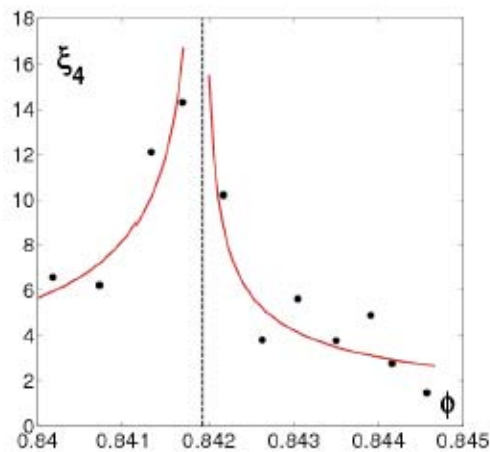
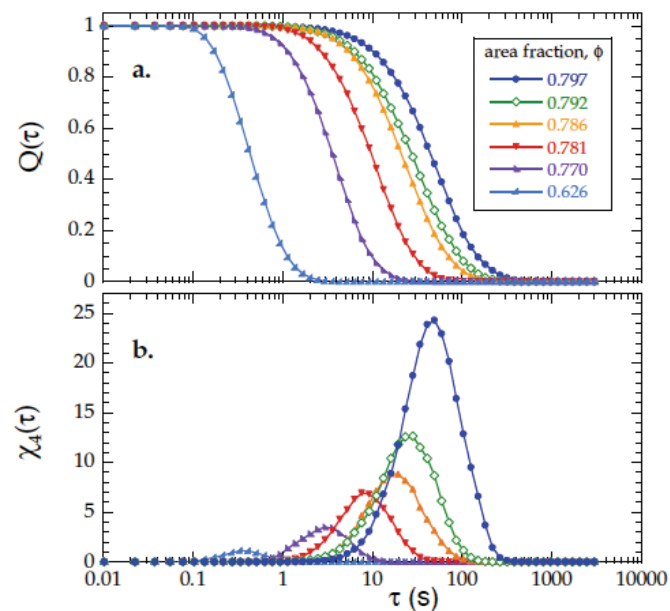
Ellenbroek W. G. et al, PRL **97** (2006) 258001



# Dynamical Heterogeneity

F. Lechenault et al., Europhys. Lett. 83, 46003 (2008)

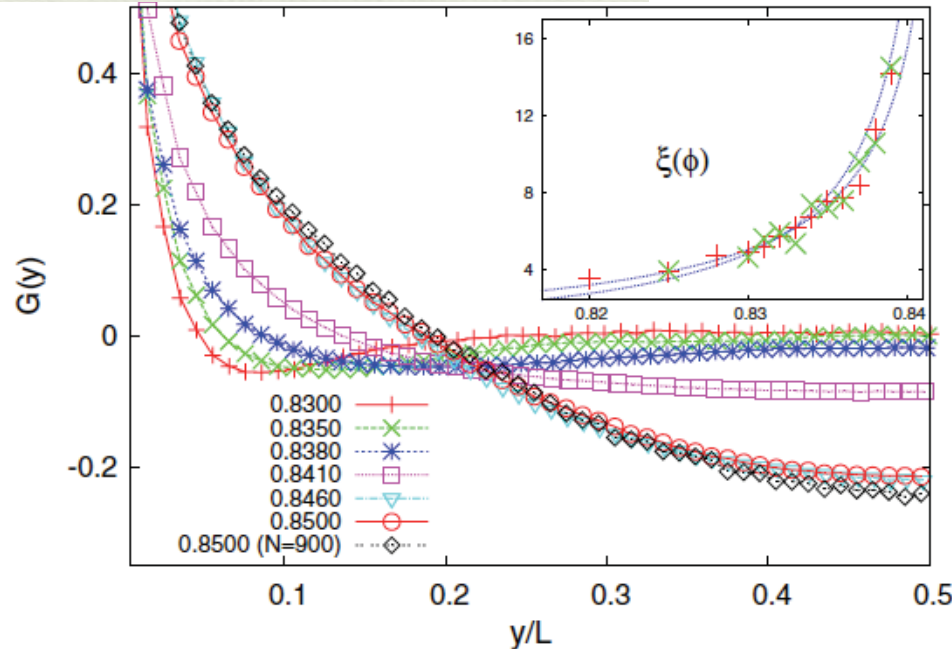
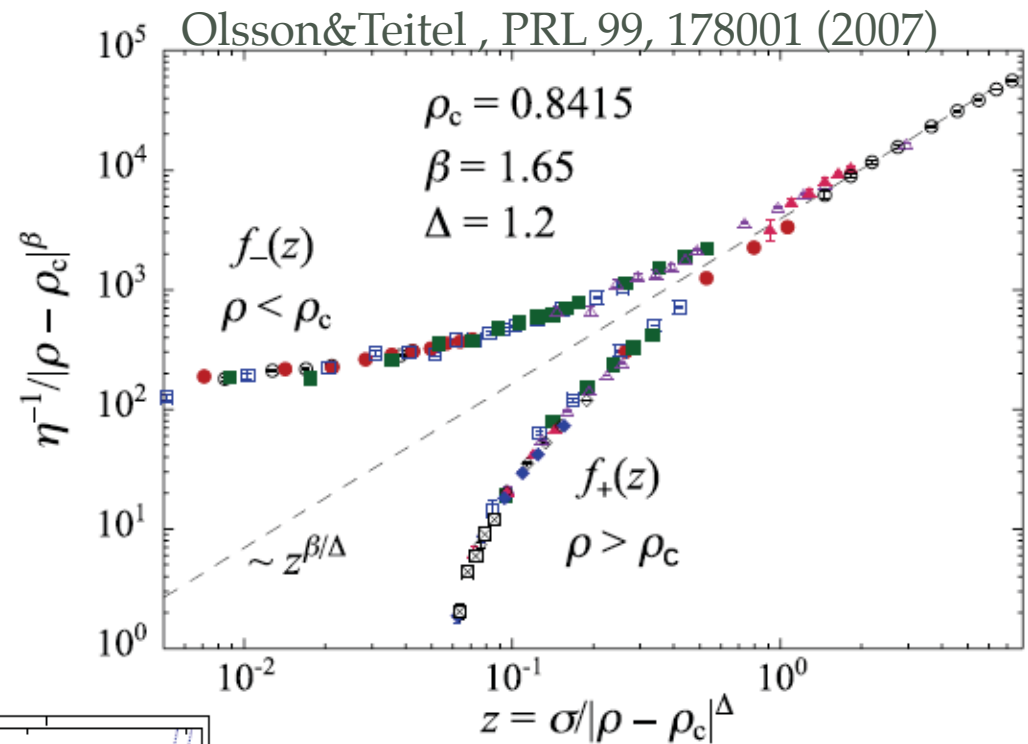
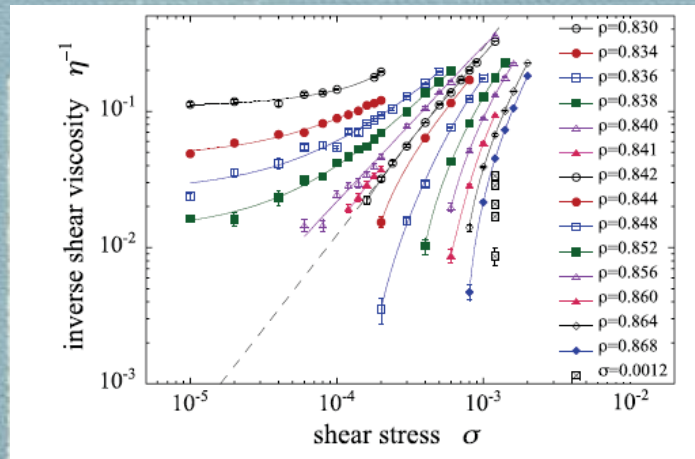
## Fluidized Monolayer



## Vibrated Monolayer

Abate and Durian, PRE 76, 021306 (2007)

# Scaling under shear



*Two-point correlations:  
no growing length scale  
on jammed side*

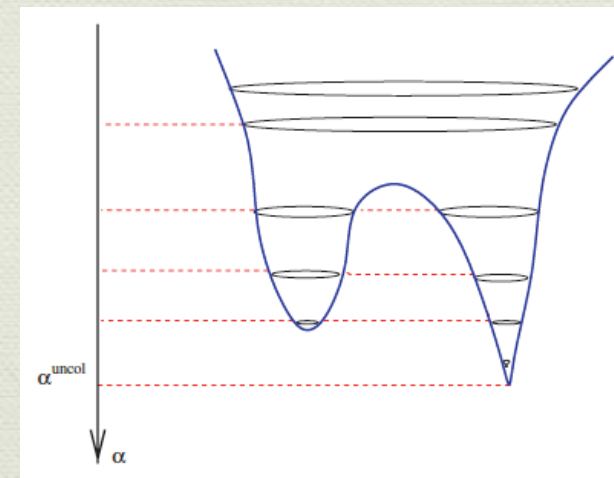
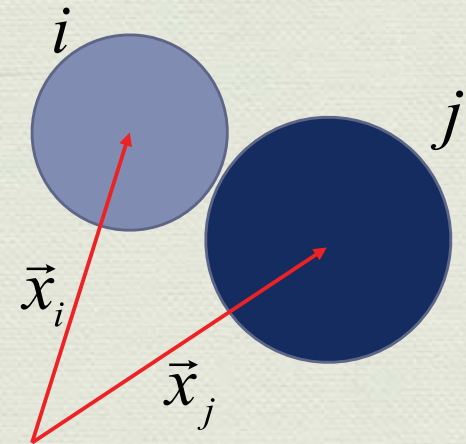
Heussinger & Barrat, PRL 102, 218303 (2009)



# Is there a Static Measure that detects a growing length scale?

- ▶ Two types of constraint satisfaction:
  - ▶ Jamming: Packing Problem
  - ▶ Unjamming: Mechanical stability constraints
- ▶ Many Solutions
- ▶ Entropy Crisis
- ▶ Mapping to Energy Landscape
- ▶ Connection to statistical mechanics of glassy systems

Krzakala & Kurchan, PRE 76, 021122 (2007)





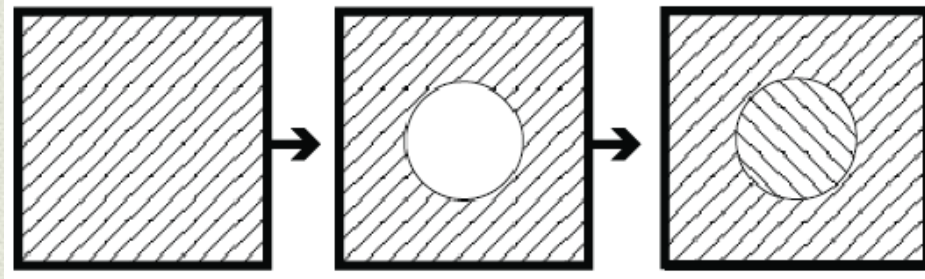
# POINT-TO-SET CORRELATIONS

Probes boundary effects:

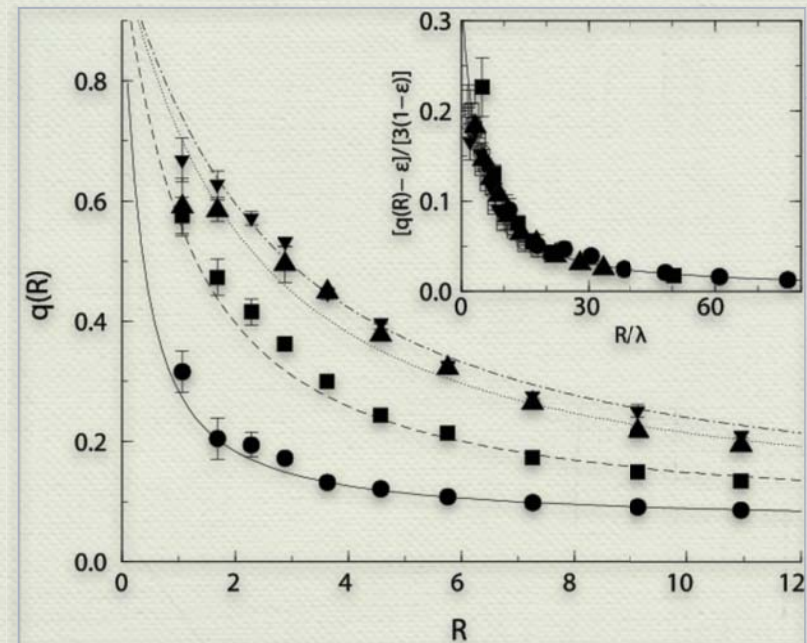
near a critical point, effects of boundaries should become long-ranged

Random-First-Order Theory: Entropic droplets

Lennard-Jones Liquids



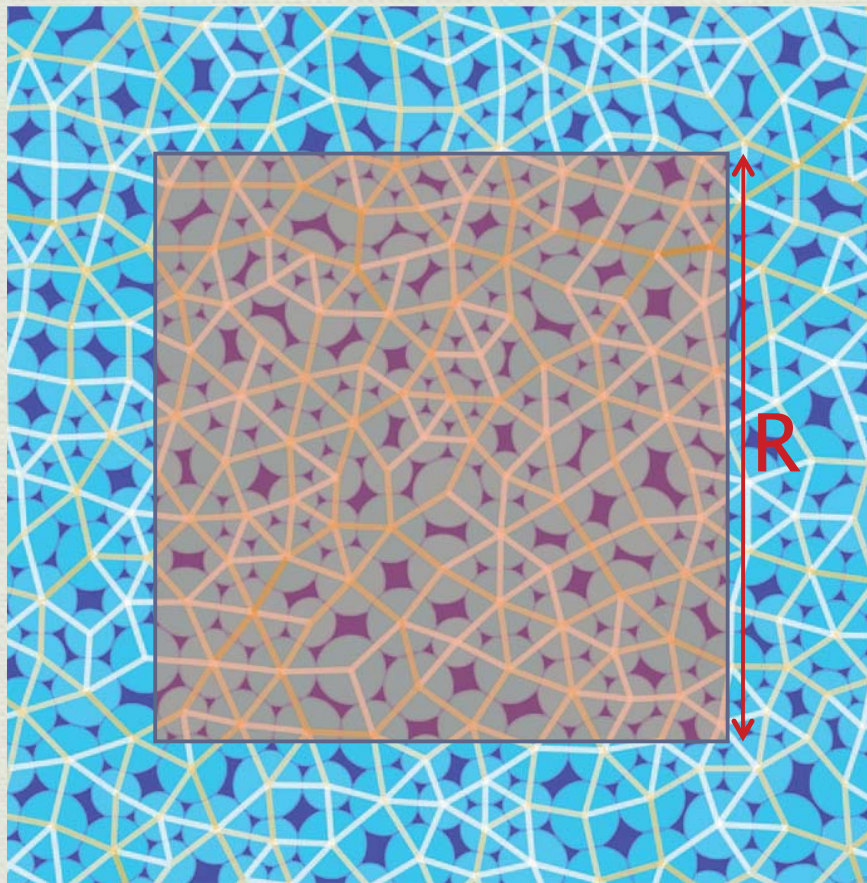
A. Cavagna et al., PRL 98 (2007) 187801,  
A. Montanari and G. Semerjian, J. Stat. Phys. 125 (2006)  
G. Biroli et al Nature Physics, 4, 771, (2008)





# Formulation of PTS In Force Space

M. Mailman, Thesis & M. Mailman and BC JSTAT L07002, (2011)



- ▶ Generate configurations at a given compression using quench protocol
- ▶ Keep grain configuration **fixed outside of  $B(R)$** .
- ▶ Find another configuration in ME by only reconfiguring **the interior**.

Force network image originated at: <http://jamming.research.yale.edu/multimedia/soft/gallery.html>R



# INDETERMINISM OF CONTACT FORCES: Force-network Ensemble (Leiden Group)

Brian P. Tighe, Jacco H. Snoeijer, Thijs J. H. Vlugt and Martin van Hecke, *Soft Matter*, 2010, **6**, 2908-2917

- ▶ Soft Spheres (M in d dimensions, z contacts / grain):
  - ▶ dM constraints from **mechanical equilibrium** (ME).

$$\sum_{j \in \text{contacts of } i} \left| \vec{f}_{ij} \right| \frac{\vec{r}_{ij}}{\left| \vec{r}_{ij} \right|} = 0$$

- ▶ Mz/2 equations associated with **contact forces**.

$$\vec{f}_{ij} = \vec{f}_{ij}(\vec{r}_i, \vec{r}_j; d_i, d_j)$$

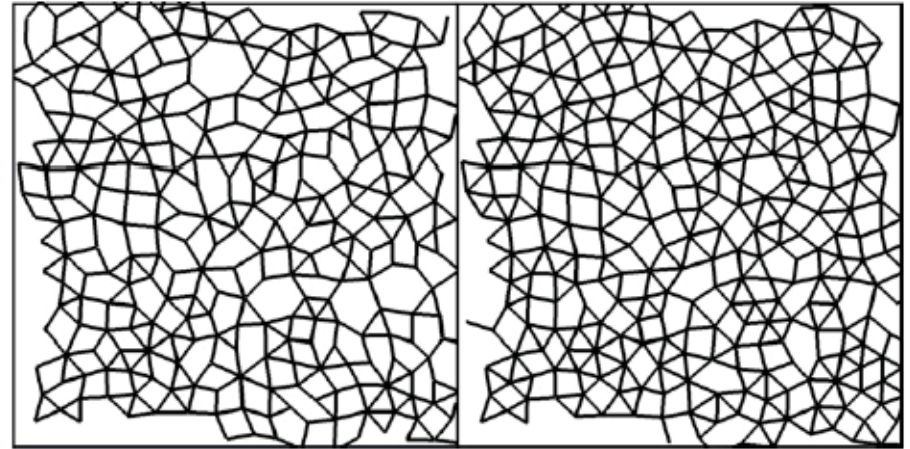
Force magnitude decouples from deformations if grains are hard enough.

- ▶ Force magnitudes are still constrained by the ME equations : Linear problem characterized by a geometry matrix, A



# Constructing force-networks

- ◆ Construct overcompressed packings of soft spheres using the O'Hern protocol (Packings are different from elastic network)
- ◆ Obtain the geometry matrix  $A$ , and the stress tensor
- ◆ Create multiple force networks on this geometry, keeping the same stress tensor



$$A\hat{g} = 0$$

*Random walk with reflecting boundary conditions in null space of matrix  $A$*

$$\vec{f} = \vec{f}(P_0) + \sum_{i=1}^{\delta z} c_i \hat{g}_i$$

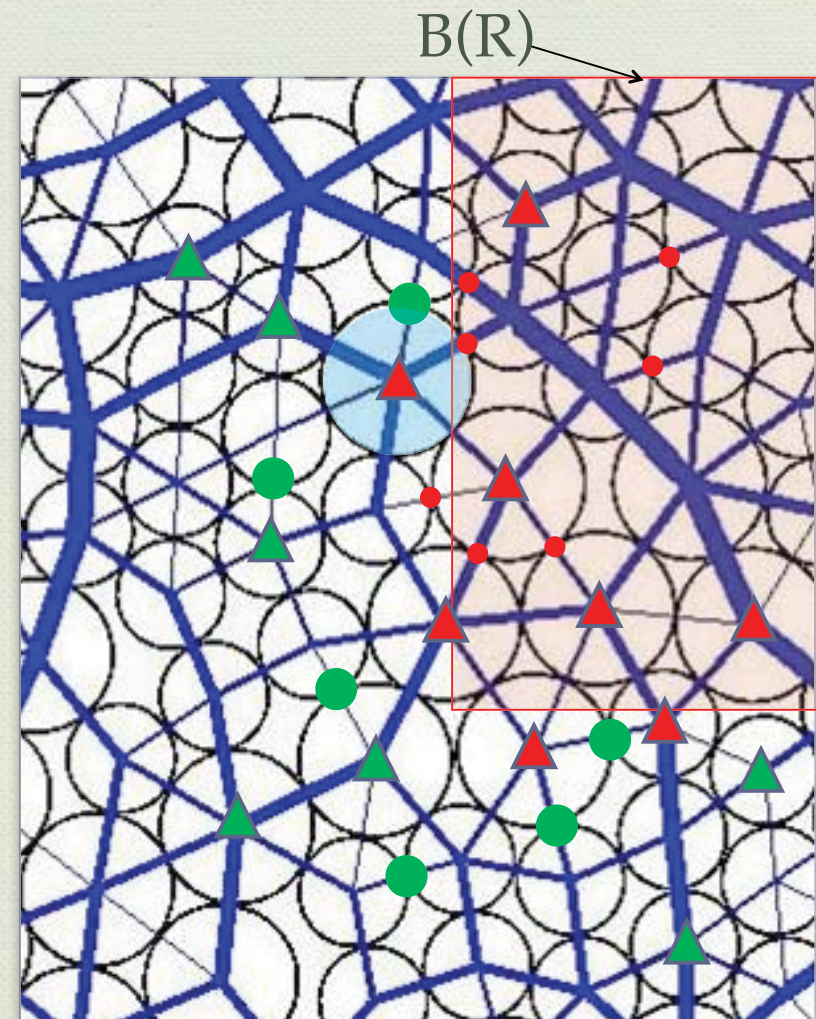


# BOUNDARY CONDITIONS AND THE CONSTRUCTION OF PTS

- ▶ Begin with static packing and  $\vec{f}_0$
- ▶ PTS is the overlap with an equivalent force network
- ▶ Freeze force variables outside of boundary

B(R):  $C(R) = \hat{f}_0 \cdot \hat{f}(R)$

- ▲ Grains contribute 2 ME constraints.
- ▲ No additional constraints.
- Fluctuating DOF.
- Frozen DOF.



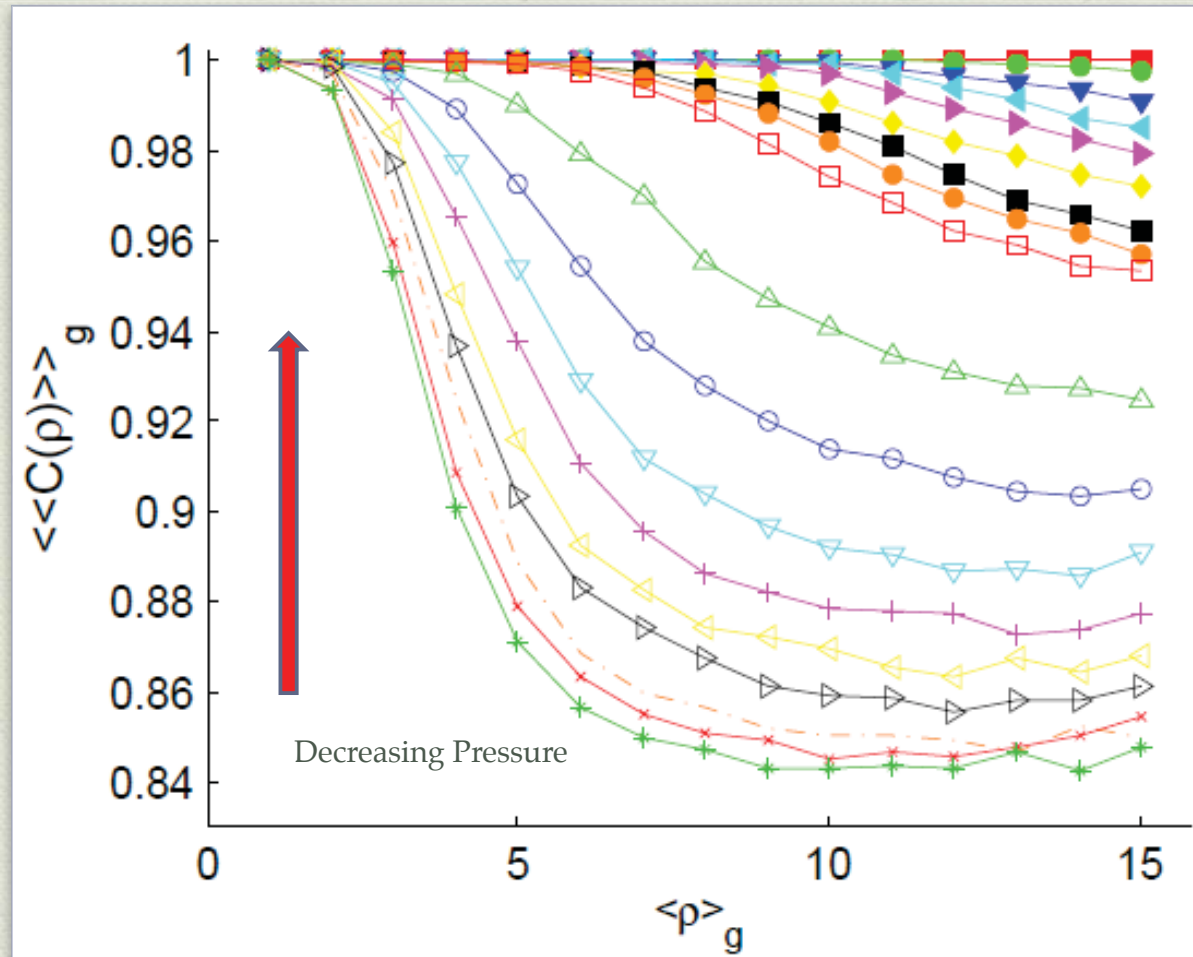


# RESULTS FOR PTS

- ▶ 900 grains, 2D, bidisperse.
- ▶ Packing fraction from  $10^{-3}$  to  $10^{-1}$ .
- ▶ 40 packing geometries.

$$\rho = R/D$$

Scaled Box size



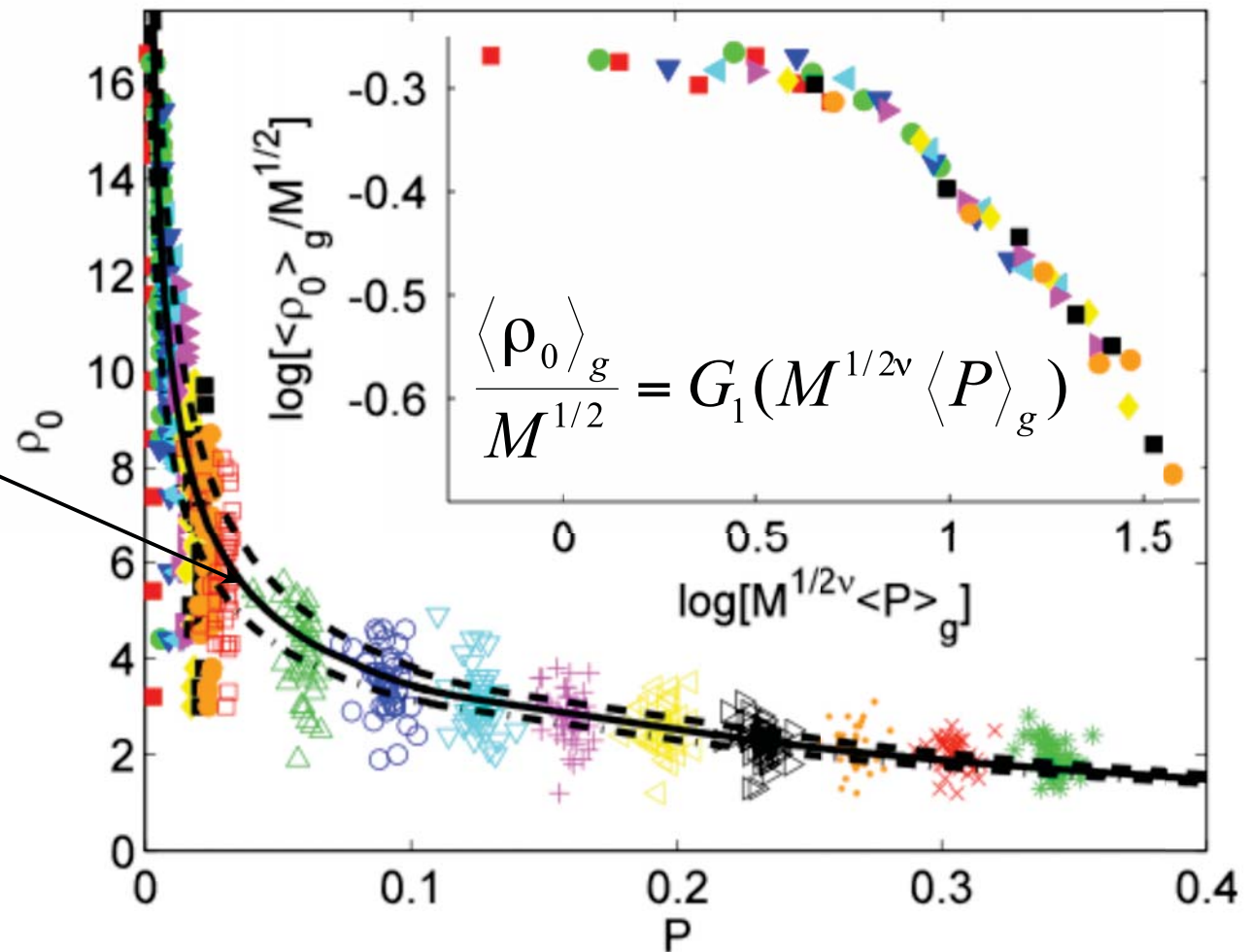


# FINITE SIZE SCALING

Fit:

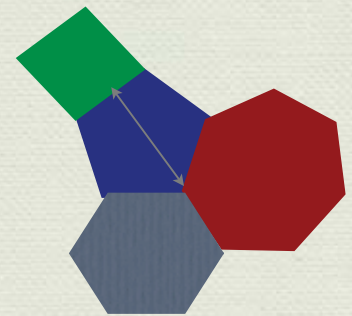
$$\nu = 0.46 \pm 0.01$$

Wyart  
argument



# Summary so far

- ◆ PTS correlation length in force space measures how far boundary forces propagate into the interior.
- ◆ PTS exhibits a growing length scale as packings are decompressed
- ◆ Finite size scaling hints at divergence
- ◆ PTS is the natural correlation function that should capture the length scale emerging from bulk-surface arguments of Wyart-Witten



*mosaic of  
force tiles*



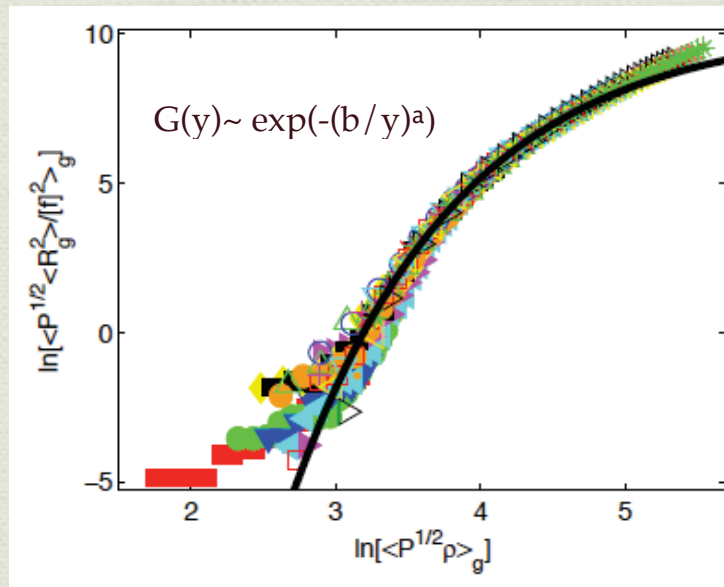
# CONFIGURATIONAL ENTROPY

- Relation to PTS

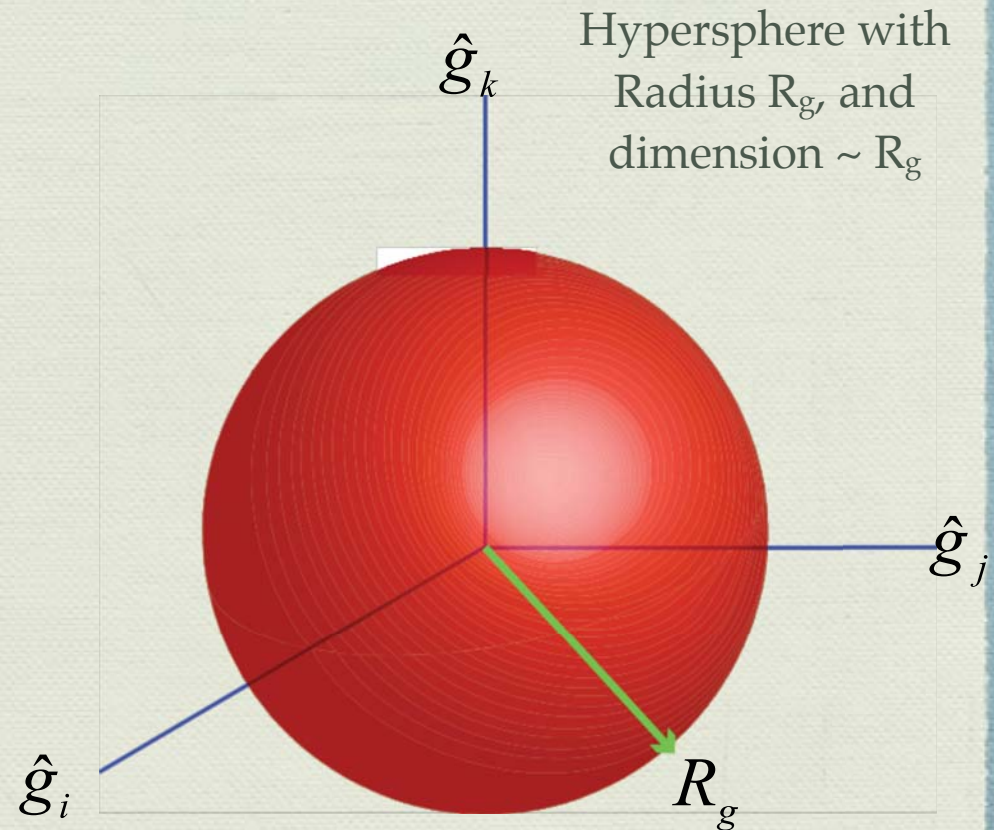
$$C(\rho) = p_{\alpha=\beta} q_1 + p_{\alpha \neq \beta} q_0$$

- Equal probability of each network is built into our sampling:

$$p_{\alpha} = \frac{1}{V_{\delta n}}$$



Logarithmic corrections to meanfield

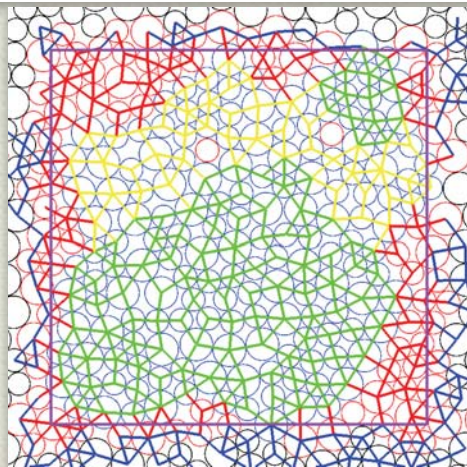
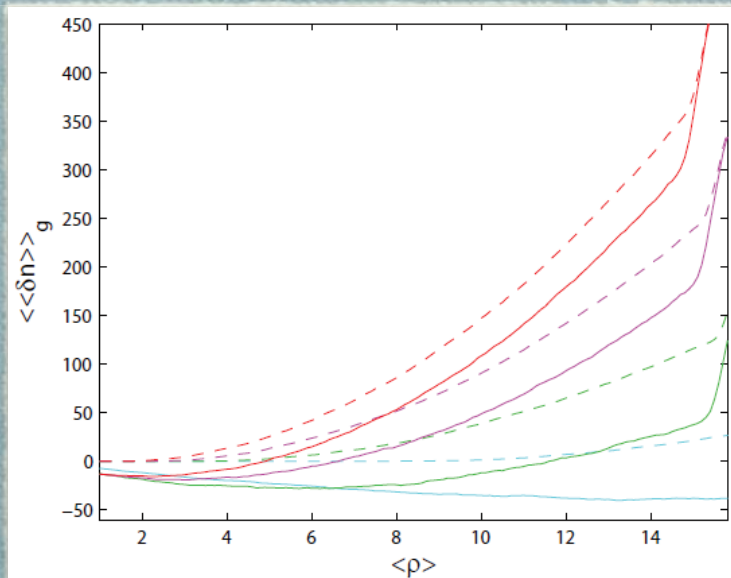


$$\langle \rho_0 \rangle_g^{-1} = \frac{\langle P \rangle_g^{1/2}}{b} \left( \ln \frac{G_0}{\eta_0} - \frac{\ln \langle P \rangle_g}{2} \right)^{1/\alpha}$$



# CORRECTIONS TO MEAN FIELD

Dashed line: nullity  
Solid line: counting

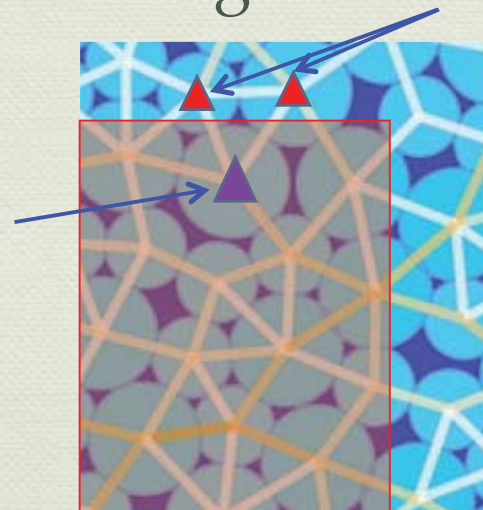


► Discrepancy between excess number of DOF and the nullity

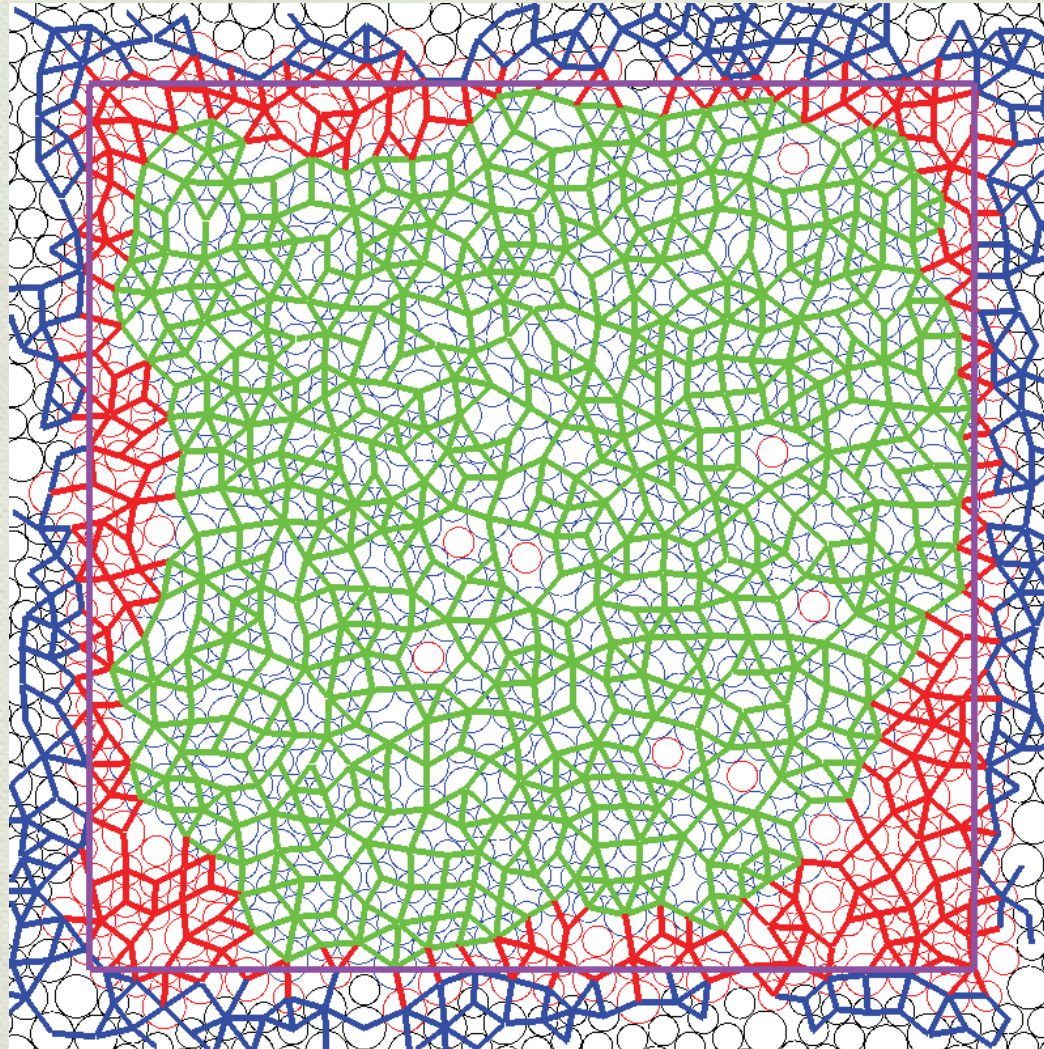
► Geometric origin:

2 constraints, 2 DOF

Effectively 2 constraints, 2 DOF

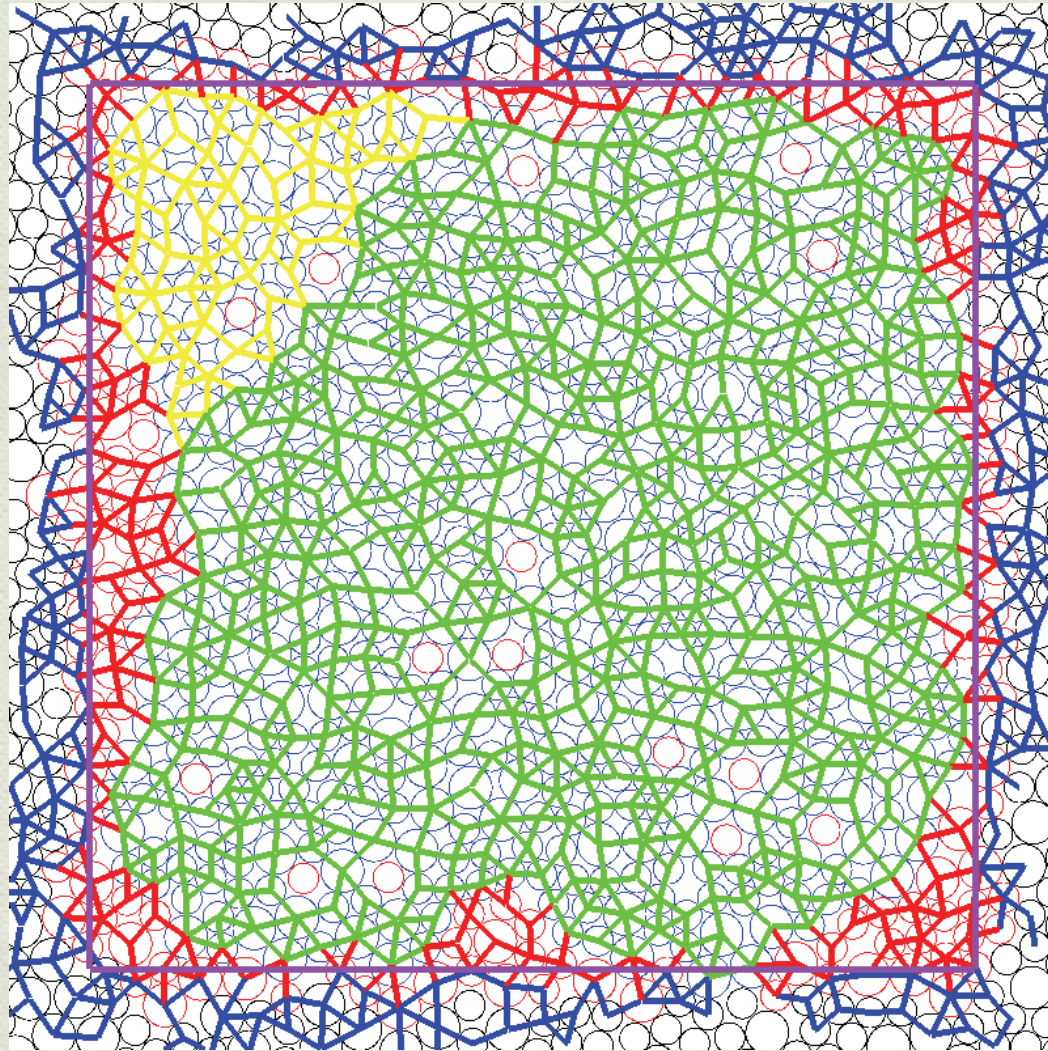






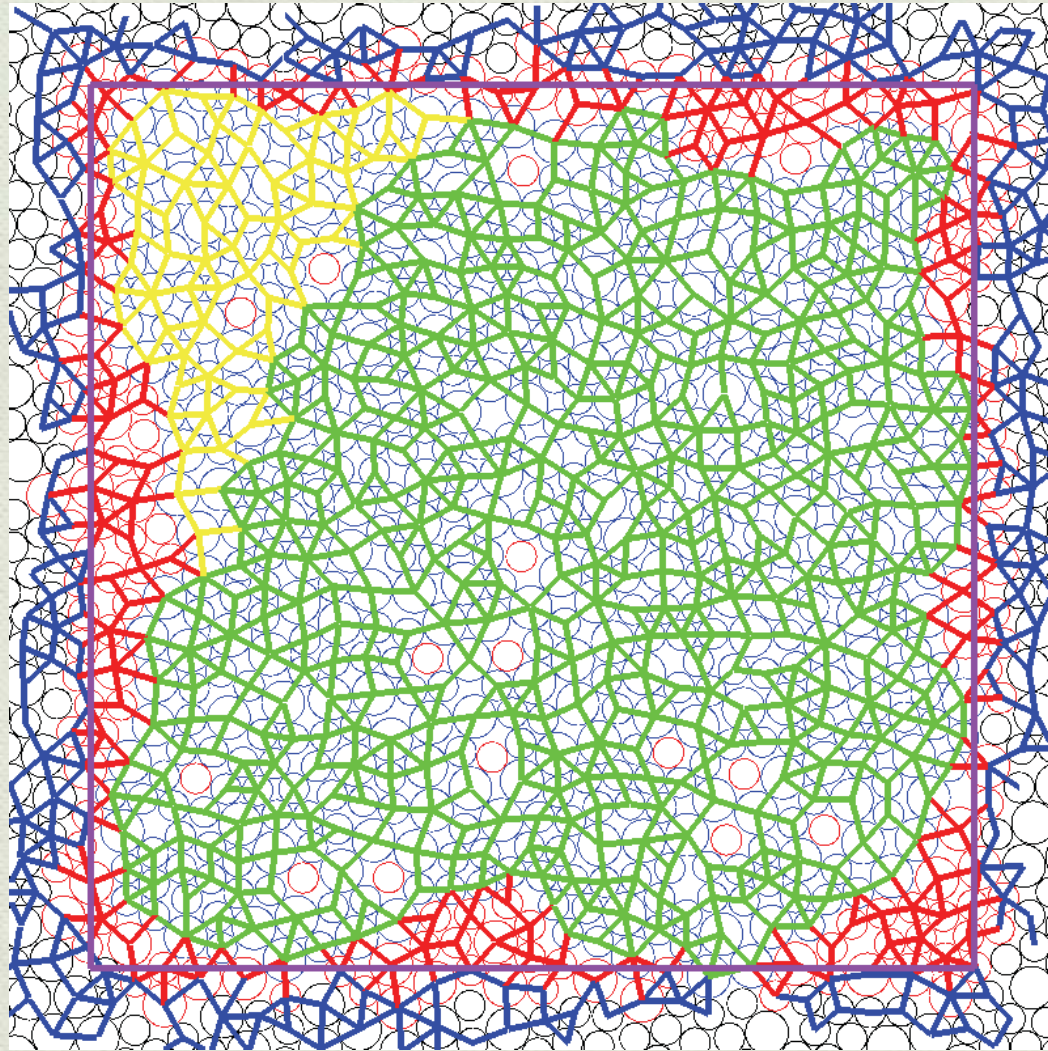
$$\delta\phi = 0.007$$





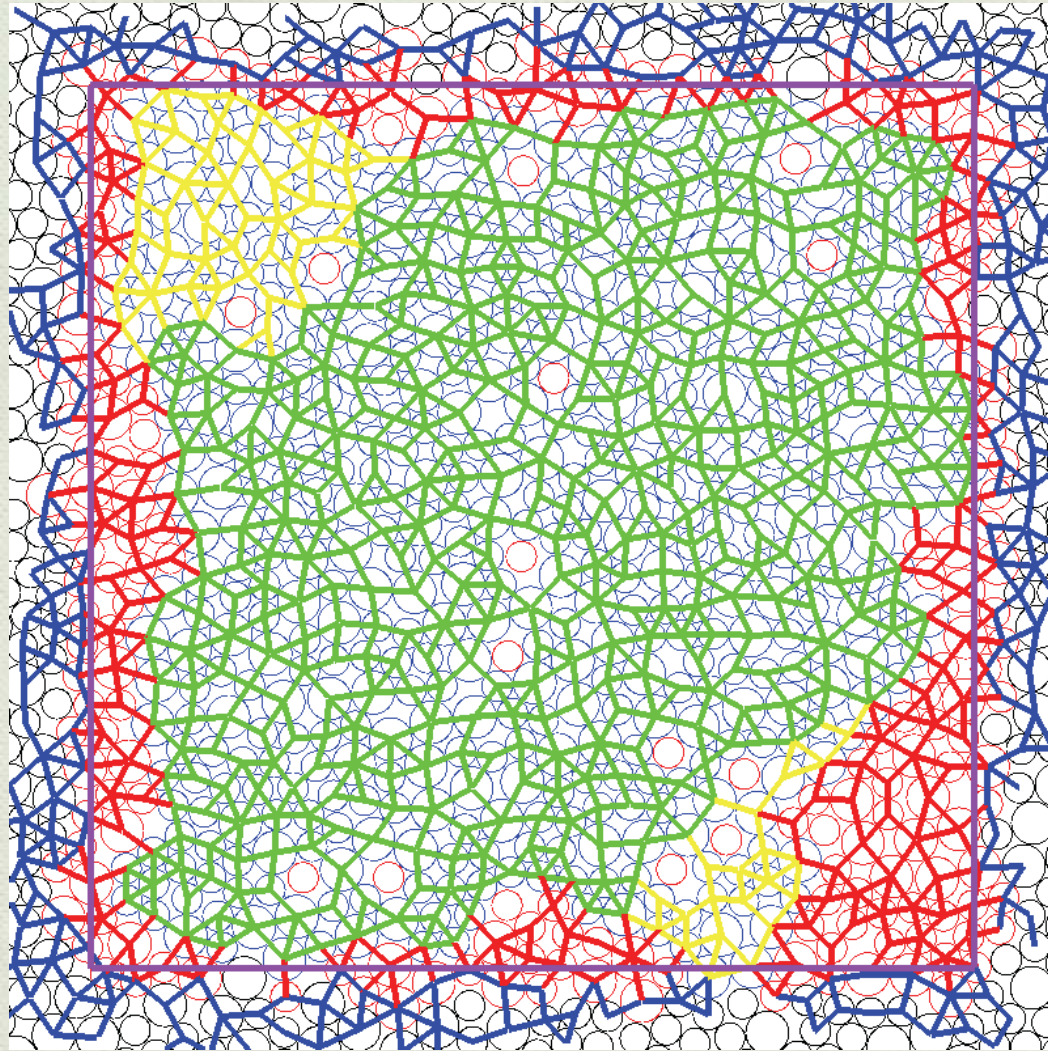
$$\delta\phi = 0.006$$





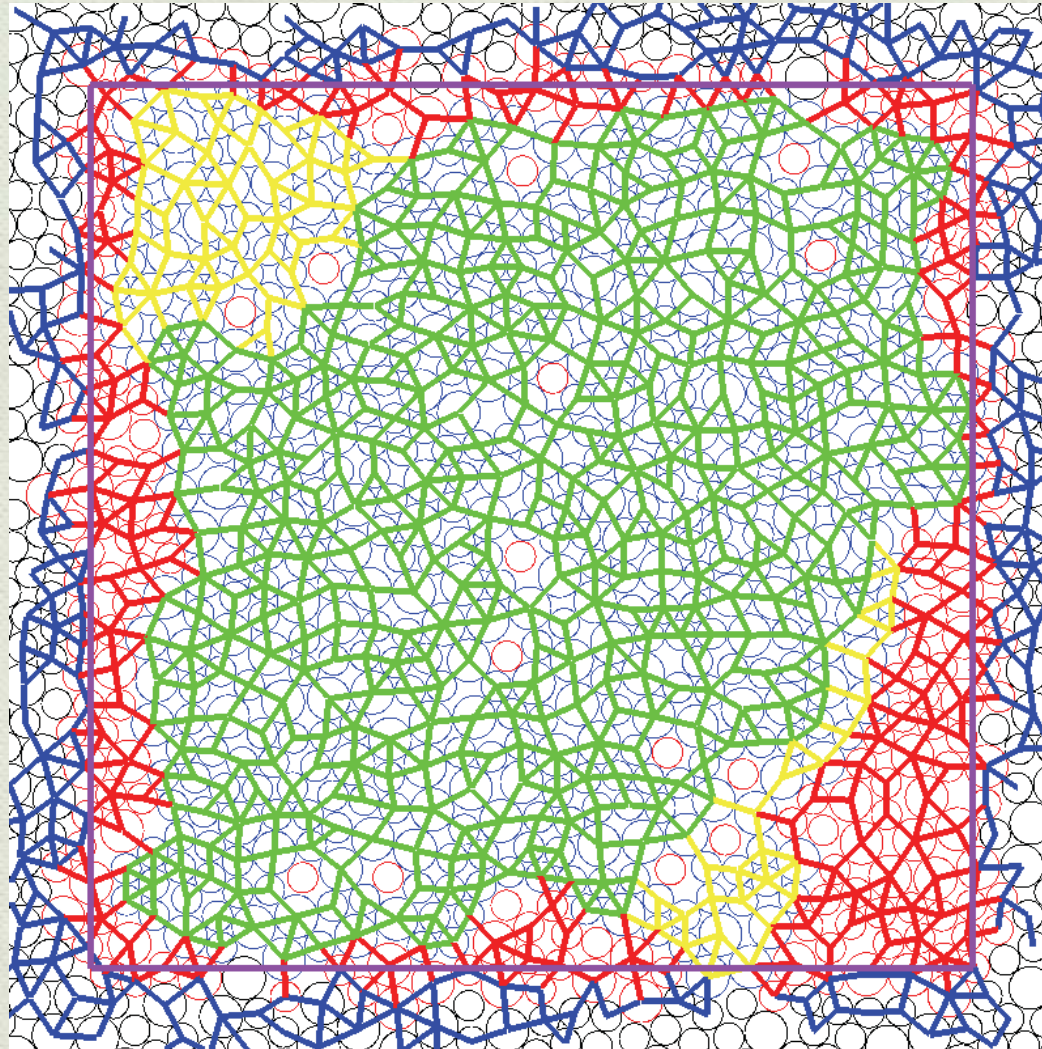
$$\delta\phi = 0.005$$





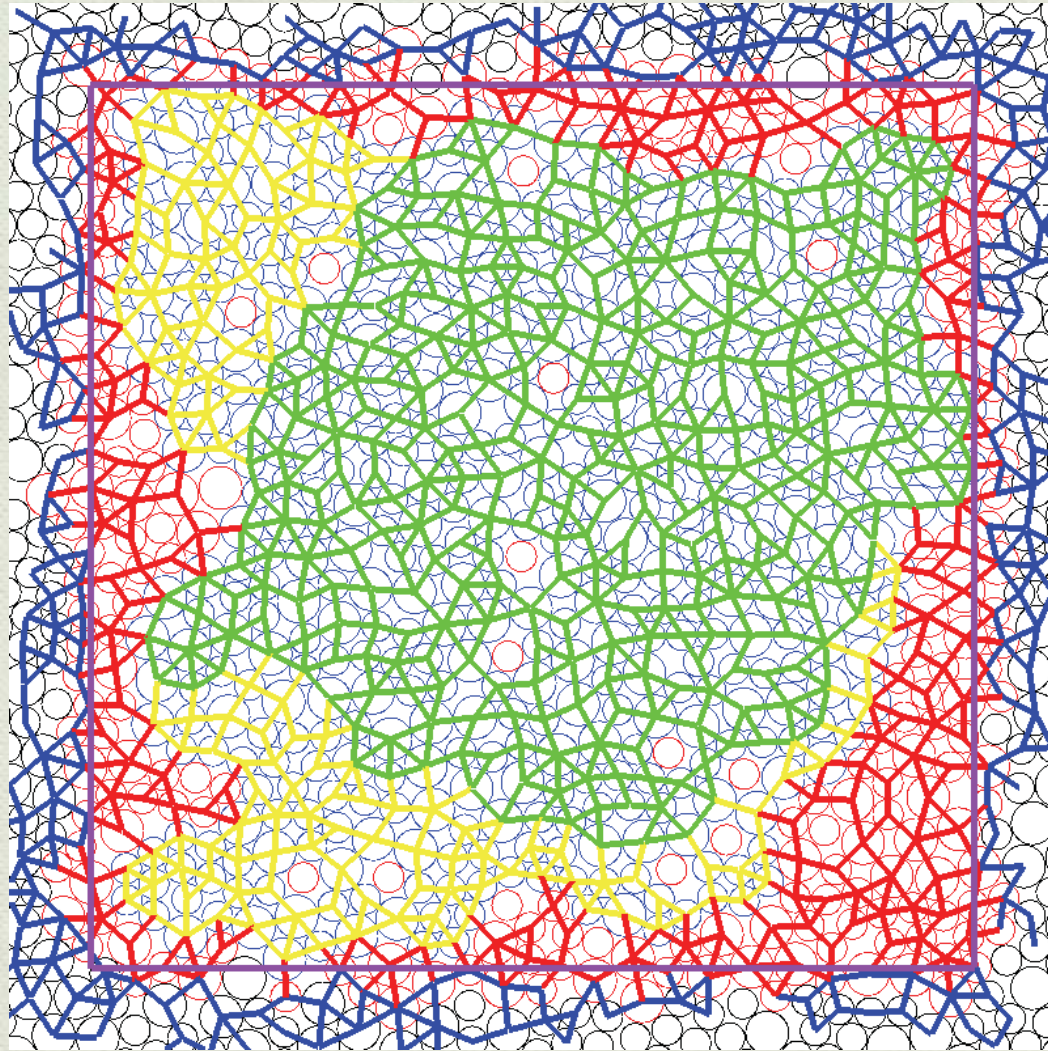
$$\delta\phi = 0.004$$





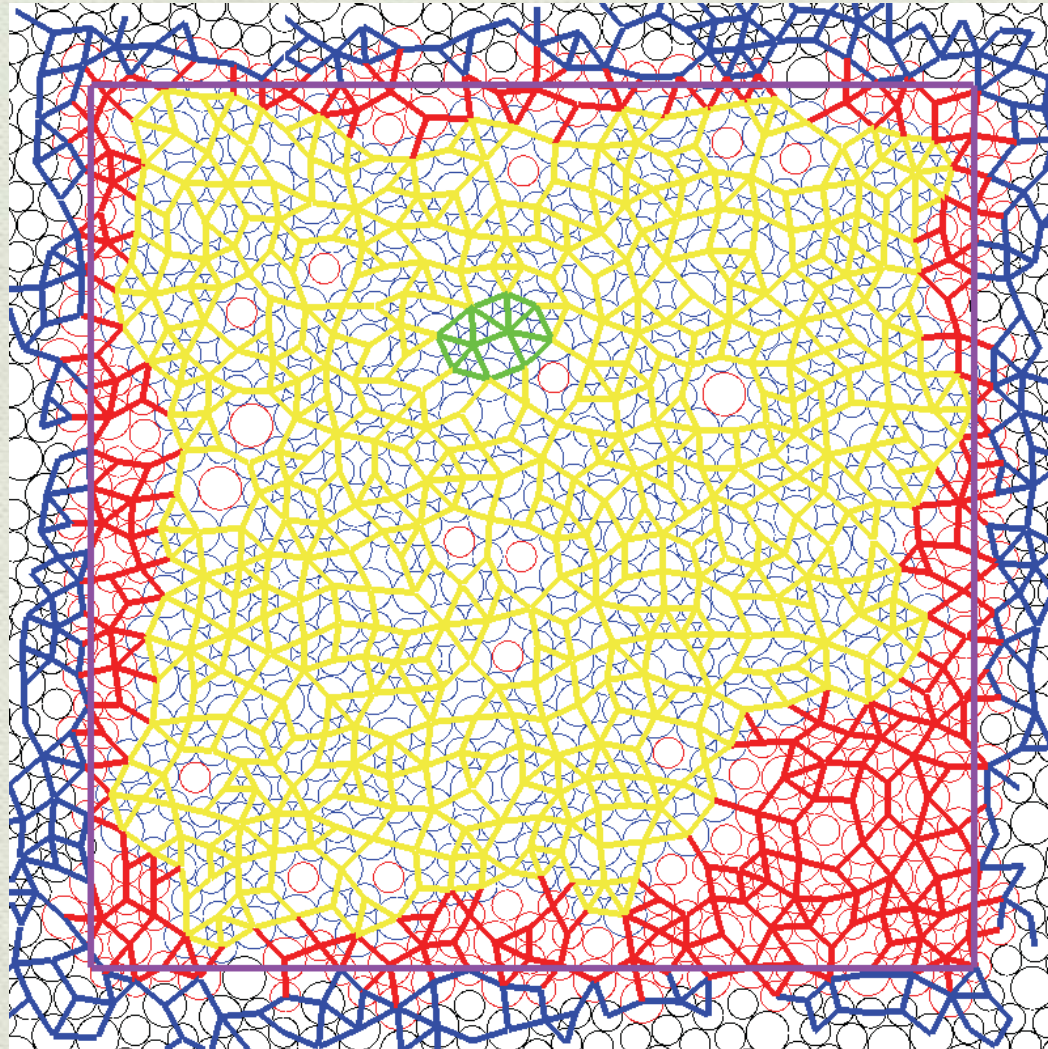
$$\delta\phi = 0.003$$





$$\delta\phi = 0.002$$



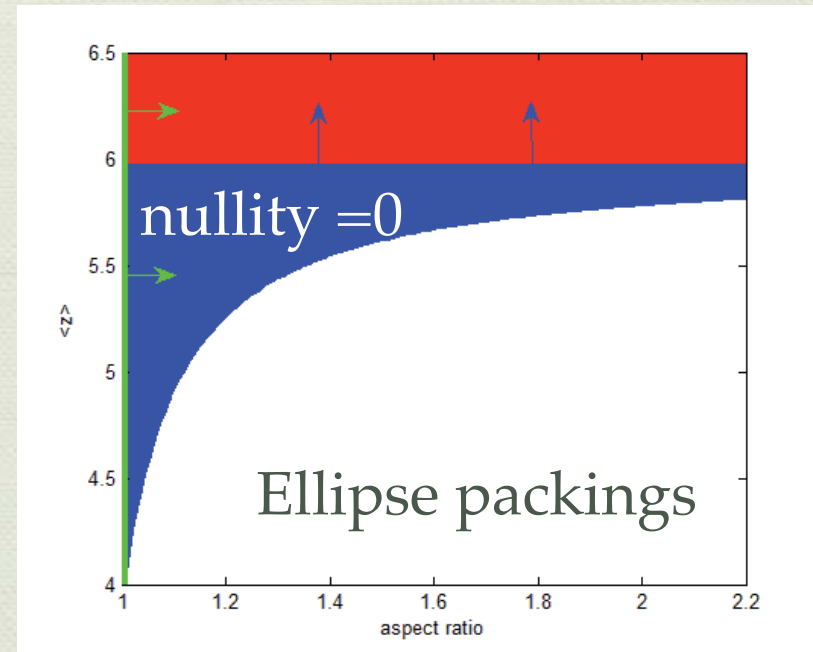


$$\delta\phi = 0.001$$



## Questions and Future Work

- ◆ Are the corrections to mean-field only relevant for one-sided spring networks?
- ◆ Stress fluctuations close to jamming, related to mesoscopic regions that fluctuate ?
- ◆ Relation to response, dynamical heterogeneities?
- ◆ Can the PTS length scale be measured experimentally?
- ◆ Anisotropic grains ? Hypostatic packings



## Acknowledgements:

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F. Krzakala

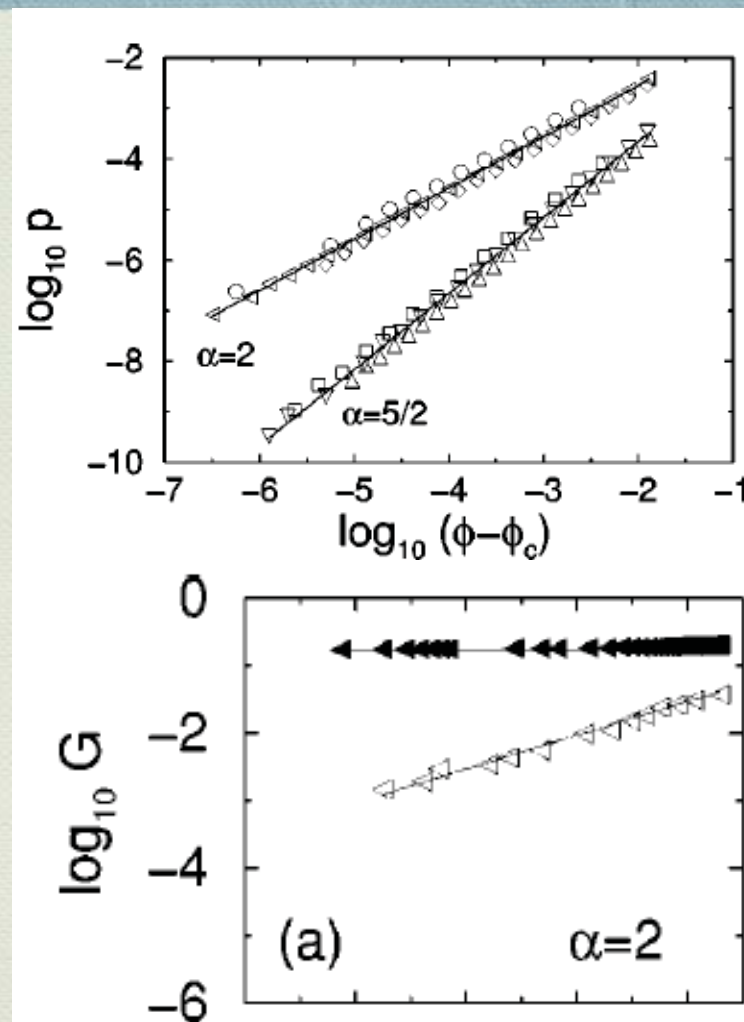
Funding: NSF



# Nature of Jamming/Unjamming Transition

## Signatures of a critical point

- ▶ Power law relationships
  - ▶ O'hern et al, PRE 68 (2003) 01130
- ▶ Long-ranged correlations on .
- ▶ Scaling
- ▶ Symmetry Breaking
- ▶ Possibility of First order transition

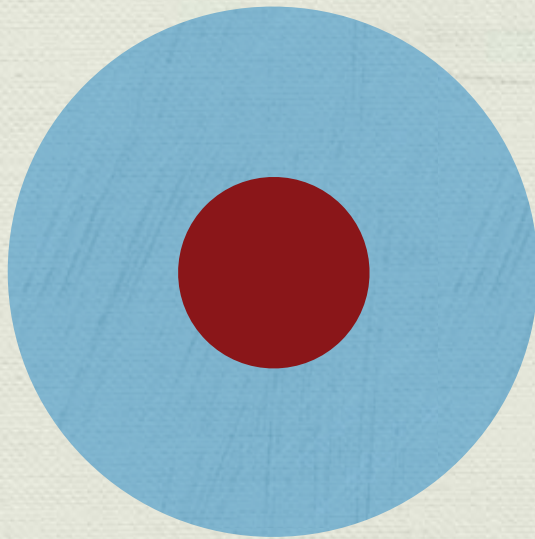


07/22/11



# RFOT: Entropic Nucleation

$$\tau \propto \exp \left[ B \left( \frac{\Delta F}{k_B T} \right)^\psi \right] \quad \xi^\theta \approx \frac{\Delta F}{k_B T} \quad \xi \approx \left( \frac{1}{TS_c} \right)^{\frac{1}{d-\theta}}$$



Entropy Gain:  $TS_c(T)\xi^d$

Surface Tension:  $\Gamma \xi^\theta : \theta \leq d-1$

Fragility ~ Related to Configurational Entropy

To get Vogel-Fulcher: Mosaic length scale has to diverge

Fragility is related to how strongly the mosaic length scale diverges