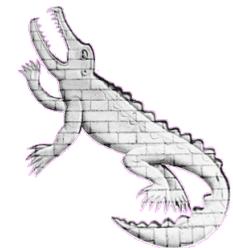




Time-evolving TD-DFT for time-dependent electronic problems

Emilio Artacho

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Miguel Pruneda
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Collaborators

Donostia International Physics Centre

Daniel Sanchez-Portal (implementation)

Andres Arnau

Inaki Juaristi

Pedro Echenique

& Discussions with Txema Pitarke

& Thanks to Peter Bauer (Linz, Austria)

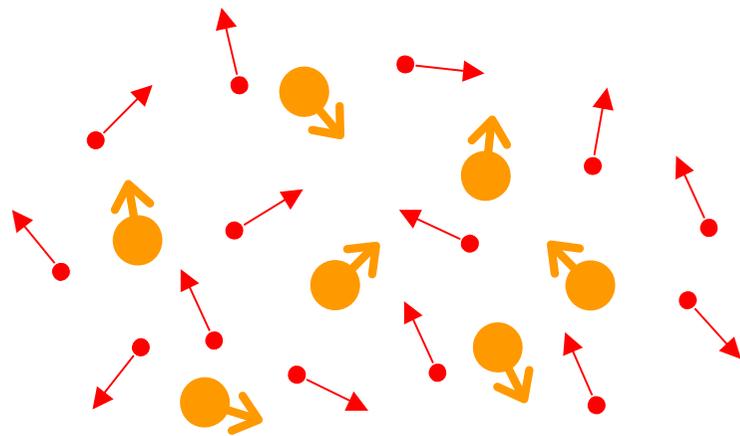
Jorge Kohanoff (Belfast)



Ahsan Zeb
U Cambridge

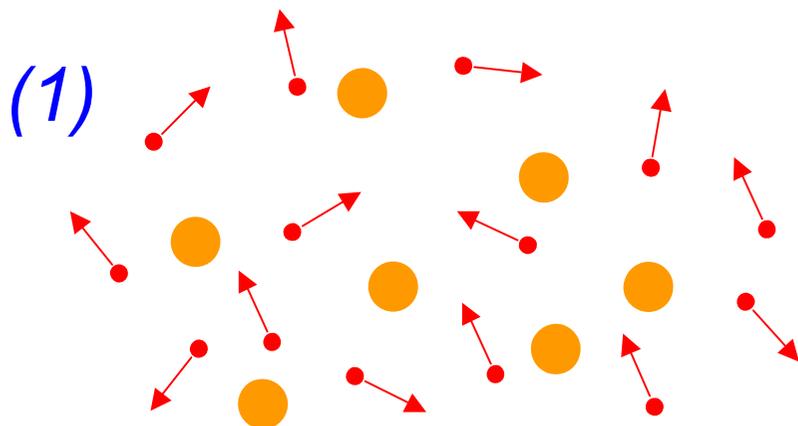


Adiabatic decoupling

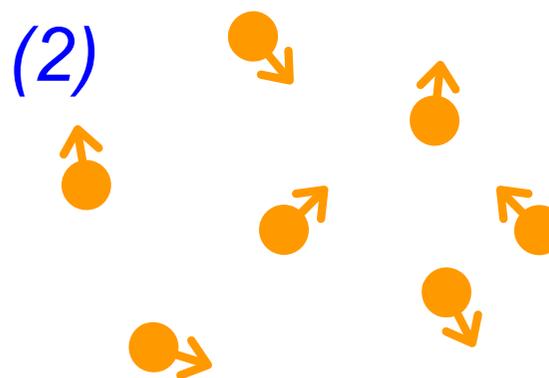


$$\frac{m_n}{m_e} \gg 1$$

\Rightarrow Nuclei are much slower than electrons



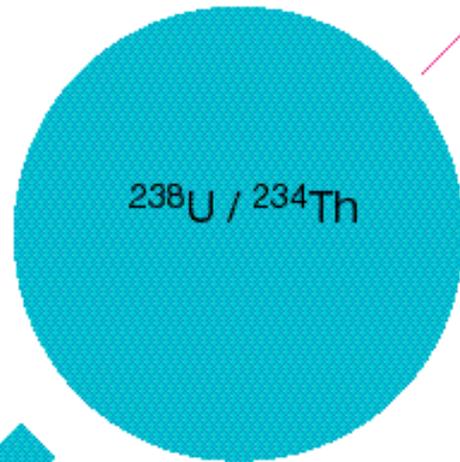
Quantum mechanics
Many electron problem:
Density Functional Theory



$F = m a$, evolution in
(discretised) time:
Molecular dynamics

Radiation damage

α -decay |



α -particle



$E \sim 5 \text{ MeV}$

$v \sim 7 \text{ a.u.}$

Recoil



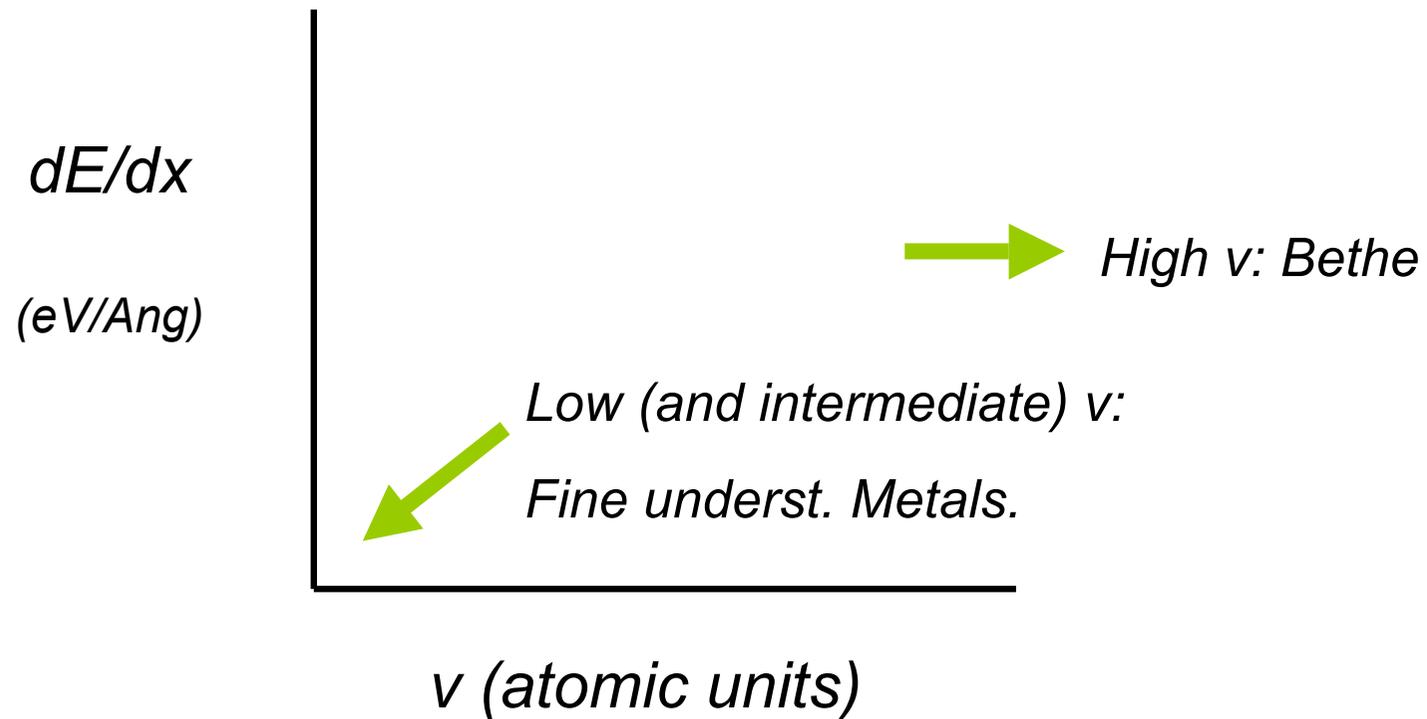
$E \sim 100 \text{ keV}$

$v \sim 0.1 \text{ a.u.}$

Other typical processes

- Neutrons, PKA
- Cosmic rays: Mostly protons $E > 1 \text{ GeV}$

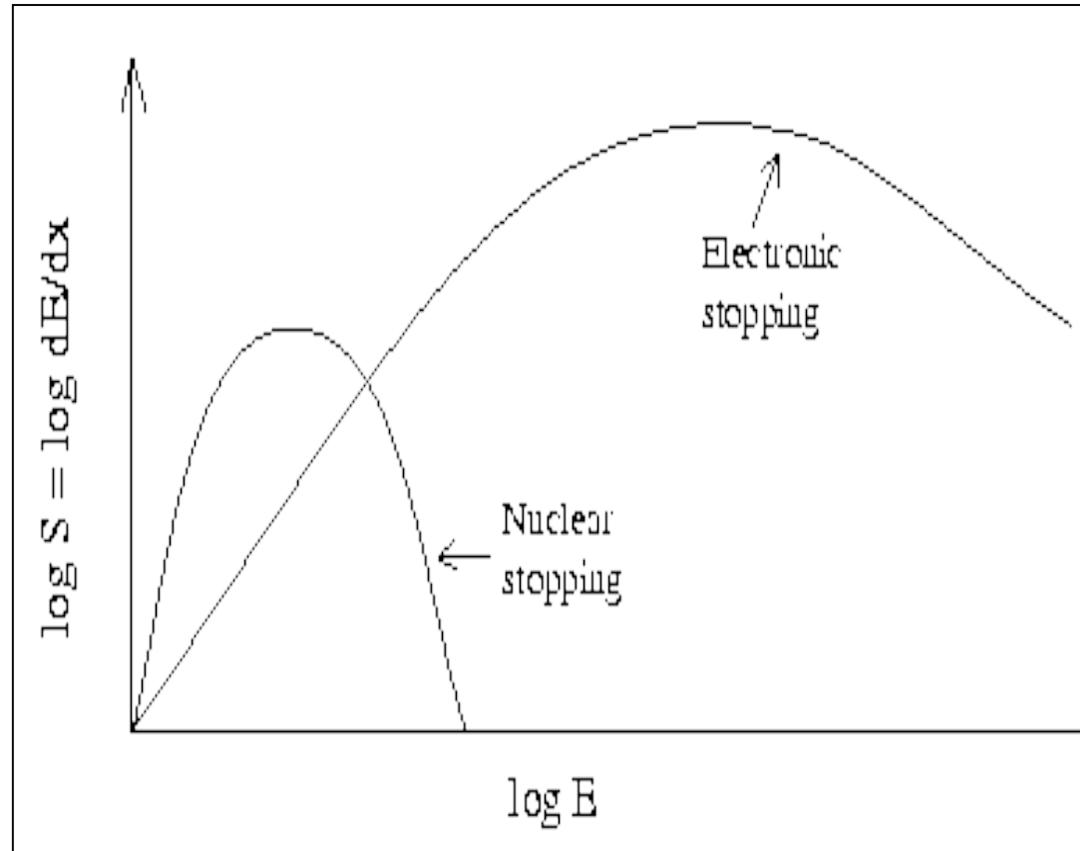
Energy transferred: Measured by stopping power



100 keV Recoiling Th nucleus: $v = 0.1$ a.u.

Electronic versus nuclear stopping

Metals:
 $S \sim v$
for $v \rightarrow 0$



Nuclear stopping dominates at low velocities

Clearly non-adiabatic systems

We cannot use the Born-Oppenheimer decoupling

Need to solve the coupled system of electrons and nuclei

Schroedinger equation for both

Dynamical problem (projectile moving) => TD-Schroedinger

Classical nuclei?

(see discussion of Ehrenfest approximation in Matthew Foulkes's talk)

TD-Schroedinger for electrons only

Time dependent DFT

Usual (stationary) DFT:

$$H\Psi(\{\vec{r}_i\}) = E\Psi(\{\vec{r}_i\}) \quad \rightarrow \quad h^{KS}\psi_n^{KS}(\vec{r}) = \varepsilon_n^{KS}\psi_n^{KS}(\vec{r})$$

Time-dependent DFT:

$$H\Psi(\{\vec{r}_i\}, t) = i\frac{\partial}{\partial t}\Psi(\{\vec{r}_i\}, t) \quad \rightarrow \quad h^{KS}\psi_n^{KS}(\vec{r}, t) = i\frac{\partial}{\partial t}\psi_n^{KS}(\vec{r}, t)$$

Neither forces on atoms (no MD), nor moving basis

Hohenberg - Kohn

$$\Psi(\{\vec{r}_i\}) \rightarrow \rho(\vec{r})$$

For our many-electron problem $\hat{H} = T + V_{ee} + \sum_{i=1}^N V_{ext}(\vec{r}_i)$

1. $E[\rho(\vec{r})] \equiv \int d^3\vec{r} V_{ext}(\vec{r}) \rho(\vec{r}) + F[\rho(\vec{r})] \geq E_{GS}$

(depends on nuclear positions) (universal functional)

2. $E[\rho_{GS}(\vec{r})] = E_{GS}$

Through 1-1 mapping
 $V_{ext}(r)$ and $n(r)$
Functional unknown!

Runge - Gross

PRL 84

$$\Psi(\{\vec{r}_i\}, t) \rightarrow \rho(\vec{r}, t)$$

For our many-electron problem $\hat{H} = T + V_{ee} + \sum_{i=1}^N V_{ext}(\vec{r}_i, t)$

1. $V_{ext}(\vec{r}, t) \Leftrightarrow \rho(\vec{r}, t)$ up to $c(t)$

1-1 correspondence => Physical properties are functionals of the density

2. $A[\Psi(\{\vec{r}_i, t\})] = \int_{t_0}^{t_1} dt \langle \Psi(t) | i \frac{\partial}{\partial t} - \hat{H}(t) | \Psi(t) \rangle$

Variational principle of the action instead of E

- Stationary point (not minimum)
- Initial value problem: $\Psi(t=0)$ defines

Kohn - Sham

Independent particles in an effective potential

They rewrote the functional as:

$$E[\rho] = T_0[\rho] + \int d^3\vec{r} \rho(\vec{r}) [V_{ext}(\vec{r}) + \frac{1}{2}\Phi(\vec{r})] + E_{xc}[\rho]$$

*Kinetic energy for system
with no e-e interactions*

Hartree potential

*The rest:
exchange
correlation*

*Equivalent to independent
particles under the potential*

$$V(\vec{r}) = V_{ext}(\vec{r}) + \Phi(\vec{r}) + \frac{\delta E_{xc}[\rho]}{\delta \rho(\vec{r})}$$

Runge - Gross

Independent particles in an effective potential

$$V_{xc} \equiv \frac{\delta E_{xc}[\rho]}{\delta \rho(\vec{r})} \quad \Rightarrow \quad V_{xc} \equiv \frac{\delta A_{xc}[\rho]}{\delta \rho(\vec{r})}$$

Where A_{xc} is obtained by subtracting the action of known bits (Hartree, Kinetic, etc) from the exact A

$$V_{eff}[\vec{r}, t, \rho(\vec{r}', t')] = V_{ext}(\vec{r}, t) + \Phi_{Hartree}(\vec{r}, t) + \frac{\delta A_{xc}[\rho(\vec{r}', t')]}{\delta \rho(\vec{r}', t')}$$

for $\vec{r}' \in \mathbb{R}^3$; $t' \leq t$

Just follow the t -dep of independent particles under V_{eff}

Approximations

Local Density Approximation (LDA)

$$V_{xc}[\rho] \approx V_{xc}(\rho(\vec{r})) \quad (\text{function parameterised for the homogeneous electron liquid as obtained from QMC})$$

Adiabatic Local Density Approximation (ALDA)

As above, plus t-dep is local in t

No history dependence

Same for Generalised Gradient Approximation (GGA)

$$V_{xc}[\rho] \approx V_{xc}(\rho(\vec{r}), \nabla \rho(\vec{r}))$$

(new terms parameterised for heterogeneous electron systems (atoms) as obtained from QC)

Time evolving

*Although TD-DFT was proposed in t -domain, most use is in frequency domain: **DFT for excitations***

*We describe t -dependent processes in
 t -TD-DFT*

As implemented by A Tsolakidis, D Sanchez-Portal and R M Martin, PRB 2002

The SIESTA method

Linear-scaling DFT based on NAOs (Numerical Atomic Orbitals)

P. Ordejon, E. Artacho & J. M. Soler , Phys. Rev. B 53, R10441 (1996)

- *Born-Oppenheimer* (relaxations, mol. dynamics)
- *DFT* (LDA, GGA)
- *Pseudopotentials* (norm conserving, factorised)
- *Numerical atomic orbitals as basis* (finite range)
- *Numerical evaluation of matrix elements* (3D grid)

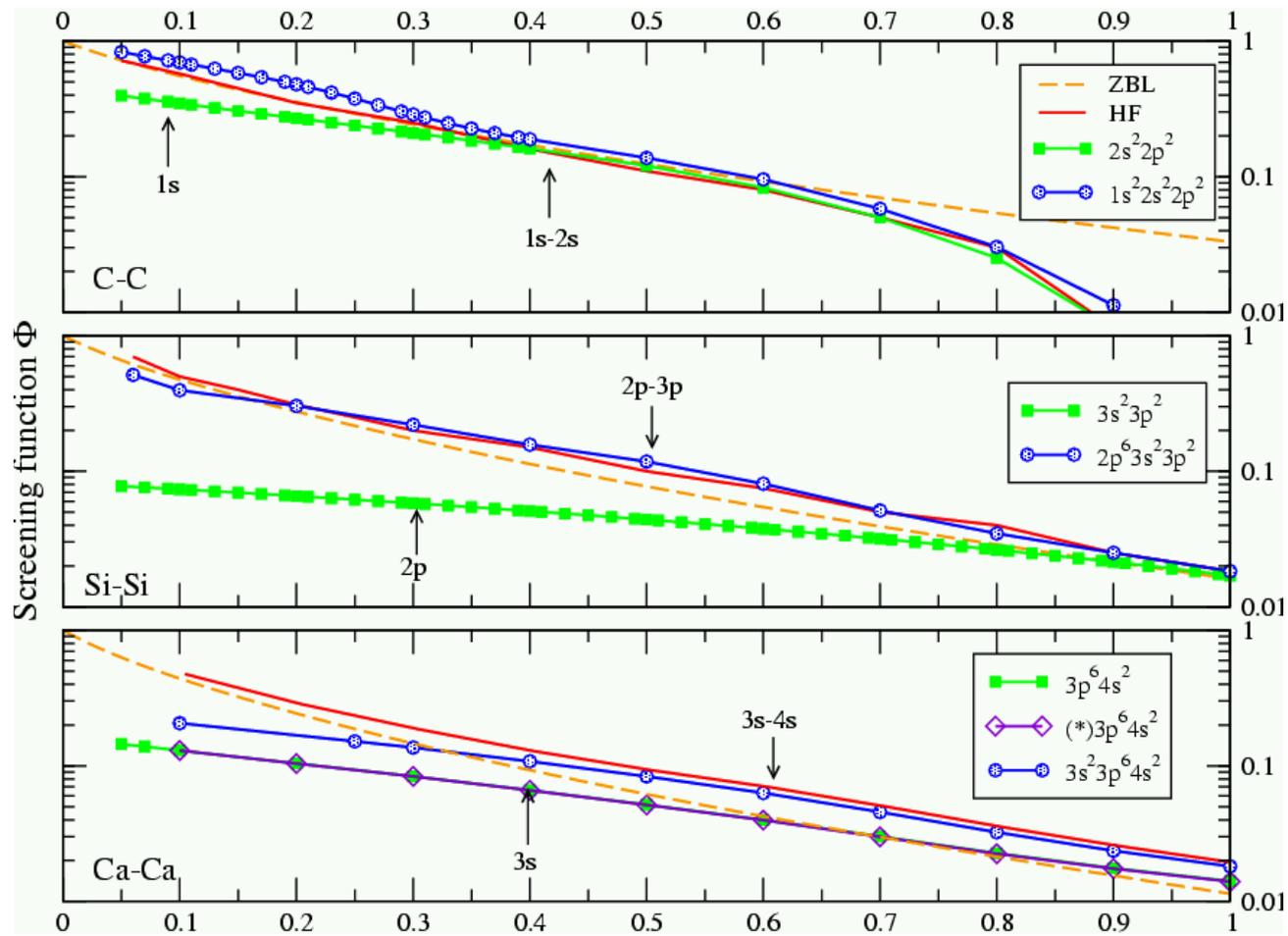
Implemented in the SIESTA program

J. M. Soler, E. Artacho, J. D. Gale, A. Garcia, J. Junquera, P. Ordejon & D. Sanchez-Portal, *J. Phys.: Condens. Matter* **14**, 2745 (2002)

Performance of pseudopotentials in high-energy collisions

In terms of the screening function

JM Pruneda & EA, PRB 2004



Real time evolution of the density

- **SIESTA (LCAO)**
$$\psi_i(r,t) = \sum_{\mu} c_i^{\mu}(t) \phi_{\mu}(r)$$
$$\rho(r,t) = \sum_{\mu,\nu} \rho_{\mu\nu}(t) \phi_{\mu}(r) \phi_{\nu}(r)$$

- *Evolution of the TD-KS equations:*

$$i \frac{\partial \psi}{\partial t} = H \psi$$

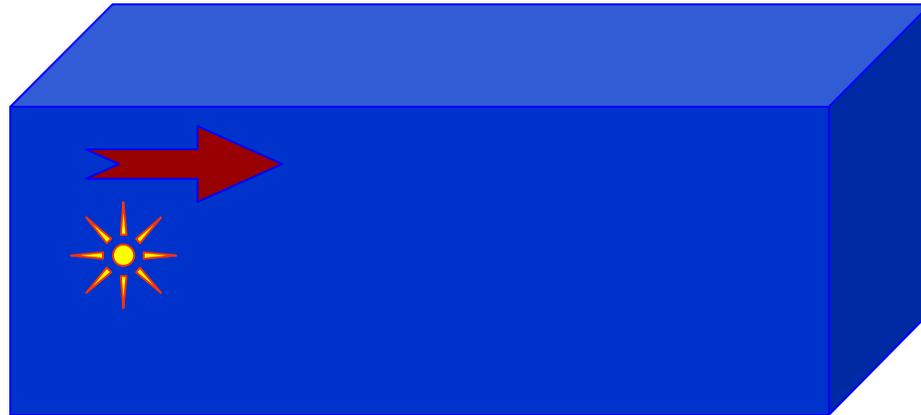
$$i \frac{\partial c}{\partial t} = S^{-1} H c$$

$$c(t_{n+1}) = \frac{1 - i S^{-1} H(t_n) \frac{\Delta t}{2}}{1 + i S^{-1} H(t_n) \frac{\Delta t}{2}} c(t_n)$$

Crank-Nicholson

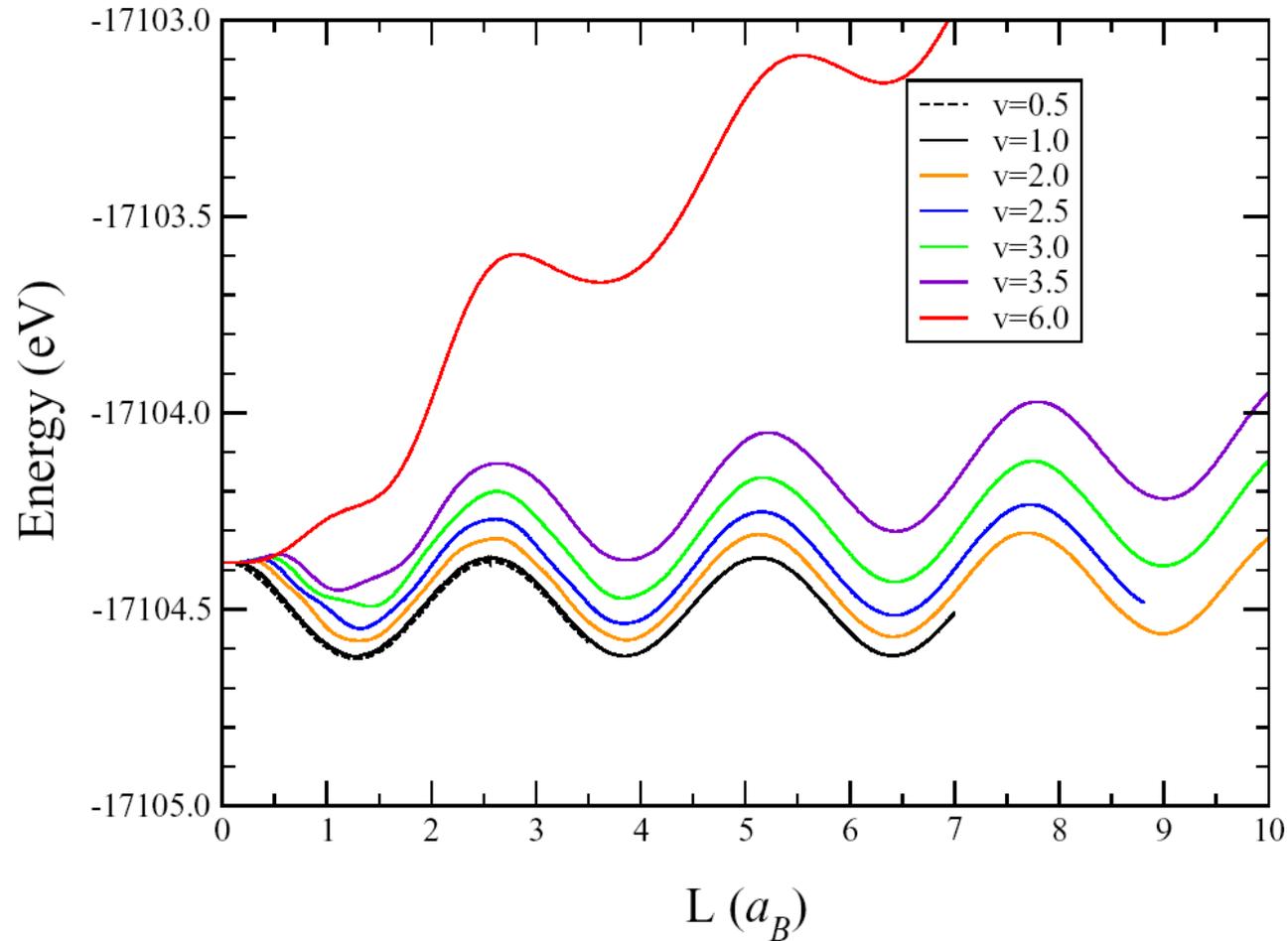
Our approach

- *Supercell of insulator's bulk*
- *Periodic boundary conditions*
- *Density functional theory*
- *Add external charge (potential)*



- *Move it and follow electron wave-functions with Time-Dependent DFT*

Energy as a function of distance: LiF



Quite stationary! Short transient, no obvious oscillation

Energy vs t

Implies that we take

$$\langle \Psi(t) | \hat{H}(t) | \Psi(t) \rangle = E[\rho(\vec{r}, t')]_{t' \leq t} \approx E_{KS}^{ALDA}(t)$$

*Goes to the right adiabatic limit, BUT
As far as I know, only justified a posteriori*

Summary

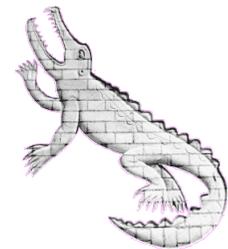
- *Need to go beyond adiabatic*
- *Starting with the electrons: A projectile as external potential gives a time-dependent Hamiltonian => TD-Schroedinger eq.*
- *Basics of TD-DFT*
- *Using time-evolving TD-DFT for stopping power of ions in matter*
- *Coupling to classical nuclei: see tomorrow's lecture by Matthew Foulkes*
- *We will see this applied to varied materials tomorrow*



Stopping power of projectiles shooting through matter
Predicting the rate of electron heating in radiation damage events

Emilio Artacho

Nanogune, Ikerbasque & DIPC, San Sebastian, Spain
Cavendish Laboratory, University of Cambridge





Miguel Pruneda
CIN2 - CSIC, Barcelona

Collaborators

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Daniel Sanchez-Portal (implementation)

Andres Arnau

Inaki Juaristi

Pedro Echenique

& Discussions with Txema Pitarke

& Thanks to Peter Bauer (Linz, Austria)

& Nuclear waste project at Cambridge

Jorge Kohanoff (Belfast)



Ahsan Zeb
U Cambridge



What to do with nuclear waste



**Shoot it at the sun. Send it to Earth's core.
What to do with nuclear waste?**

Government advisers consider 14 ways of getting rid of the troublesome legacy

Paul Brown, environment correspondent
Wednesday April 14, 2004
[The Guardian](#)

*Now: BNFL (UK) vitrifies it into
Borosilicate glass (~20%)*

Durability ~100 years

10ky – 1My needed!

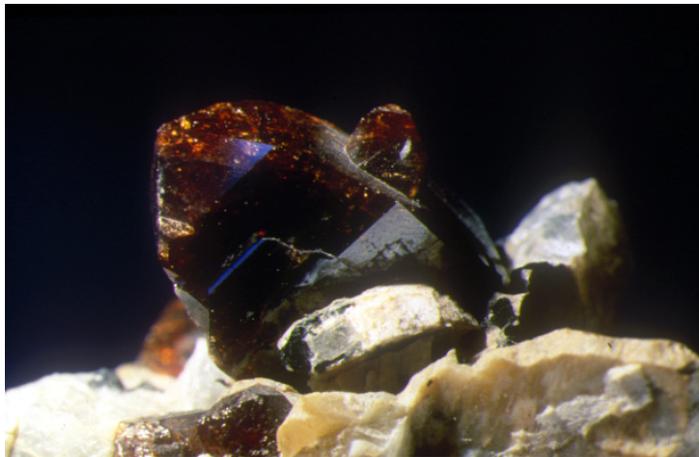
What to do with nuclear waste

Immobilisation by dilution in ceramics

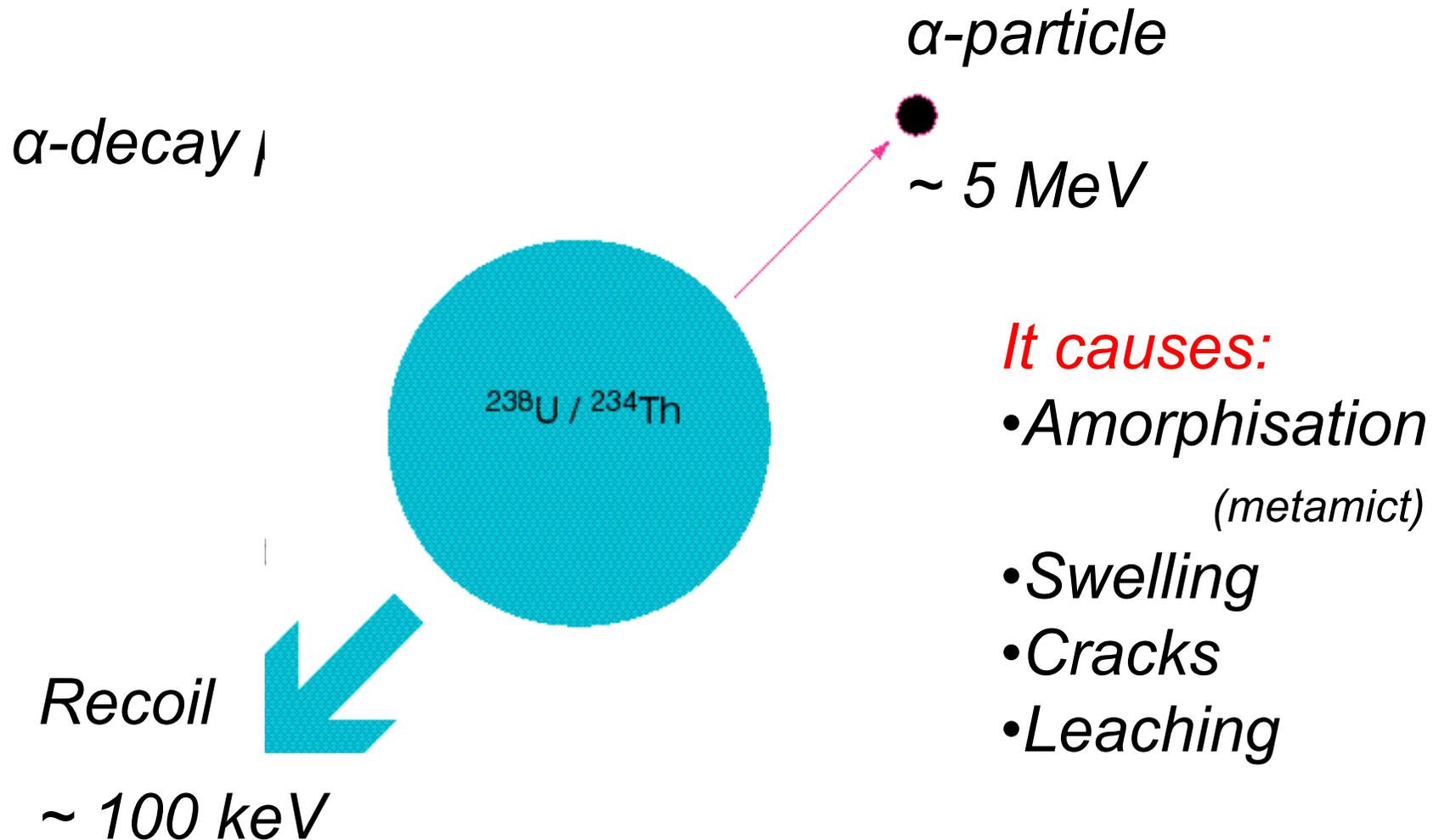
***SYN**thetic **ROCK**s with appropriate “minerals” to host high level nuclear waste*

Research in durability: resistance to radiation damage

Zircons have contained uranium for billions of years



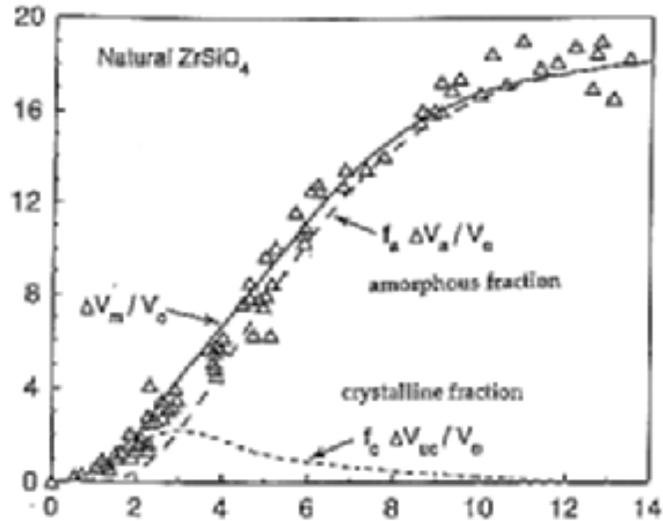
Durability: Radiation damage



Zircon: model study: old natural samples

Swelling in zircons

Volume change (%)

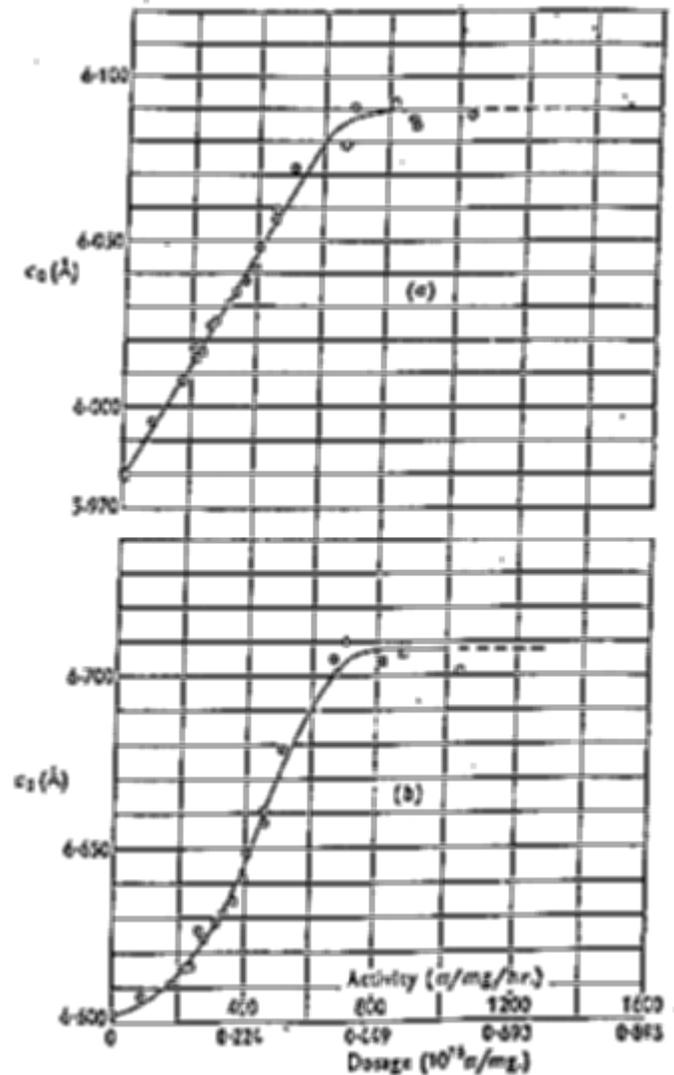


Dose (10^{16} decays/g)

Crystalline swelling:
lattice parameters vs dose

Total: ~20%

Crystalline: ~5% anisotropic



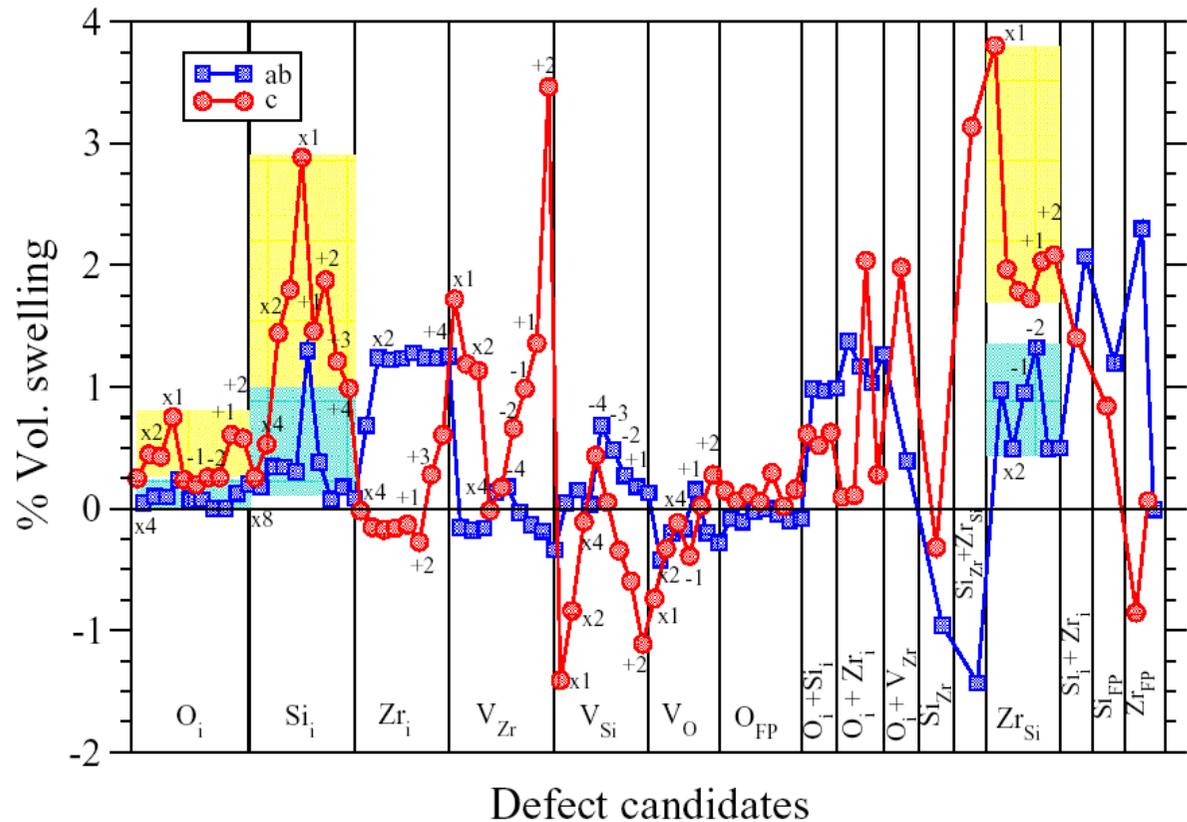
Intrinsic point defects and crystalline swelling in ZrSiO4

- Radiation cascades & defect accumulation $\sim 10^{21}$ defects/cm³
- Anisotropic swelling: $\sim 1.5\%$ in *ab*-plane *Si 0.2% swelling!!*
 $\sim 2\%$ in *c*-axis

• IR, Raman & NMR

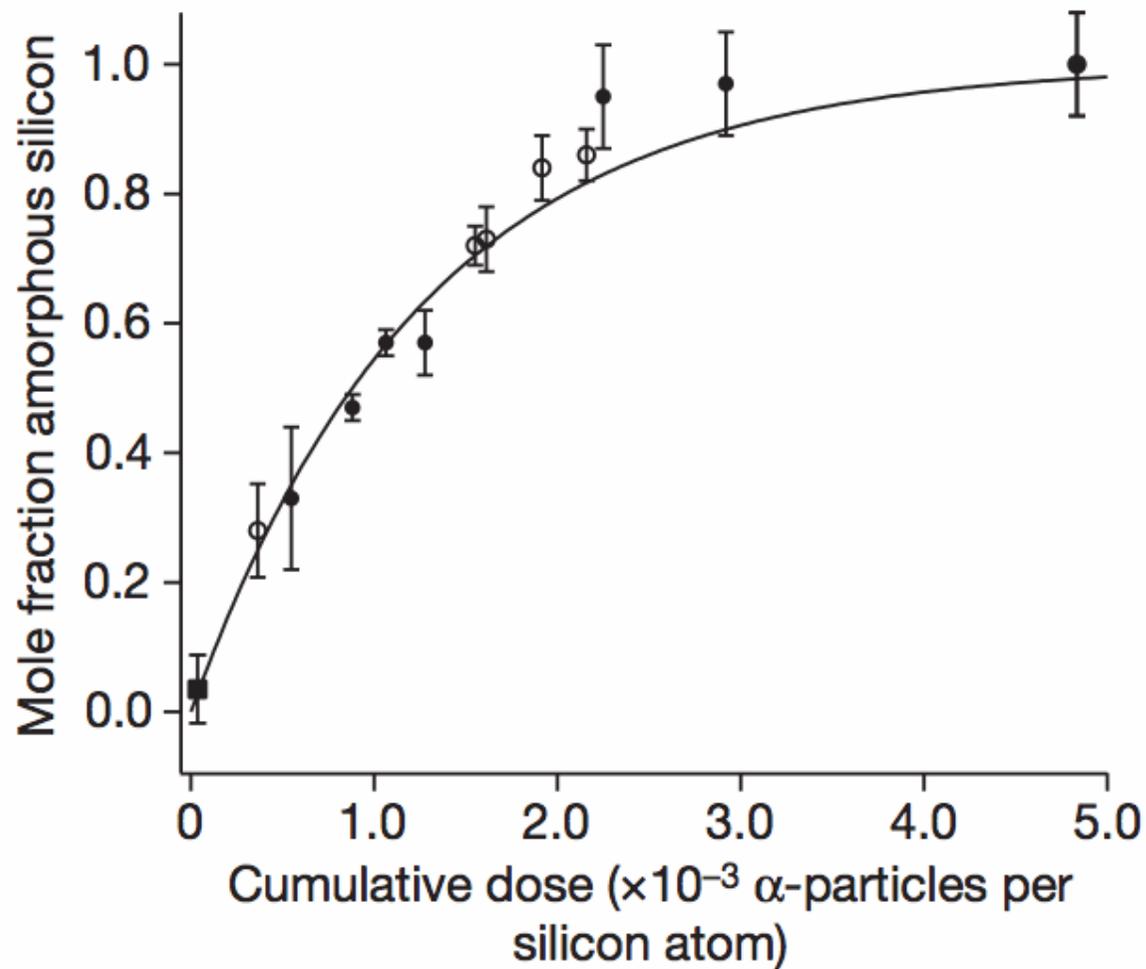
Possible defects:

O & Si interstitials
Zr_{Si} anti-site



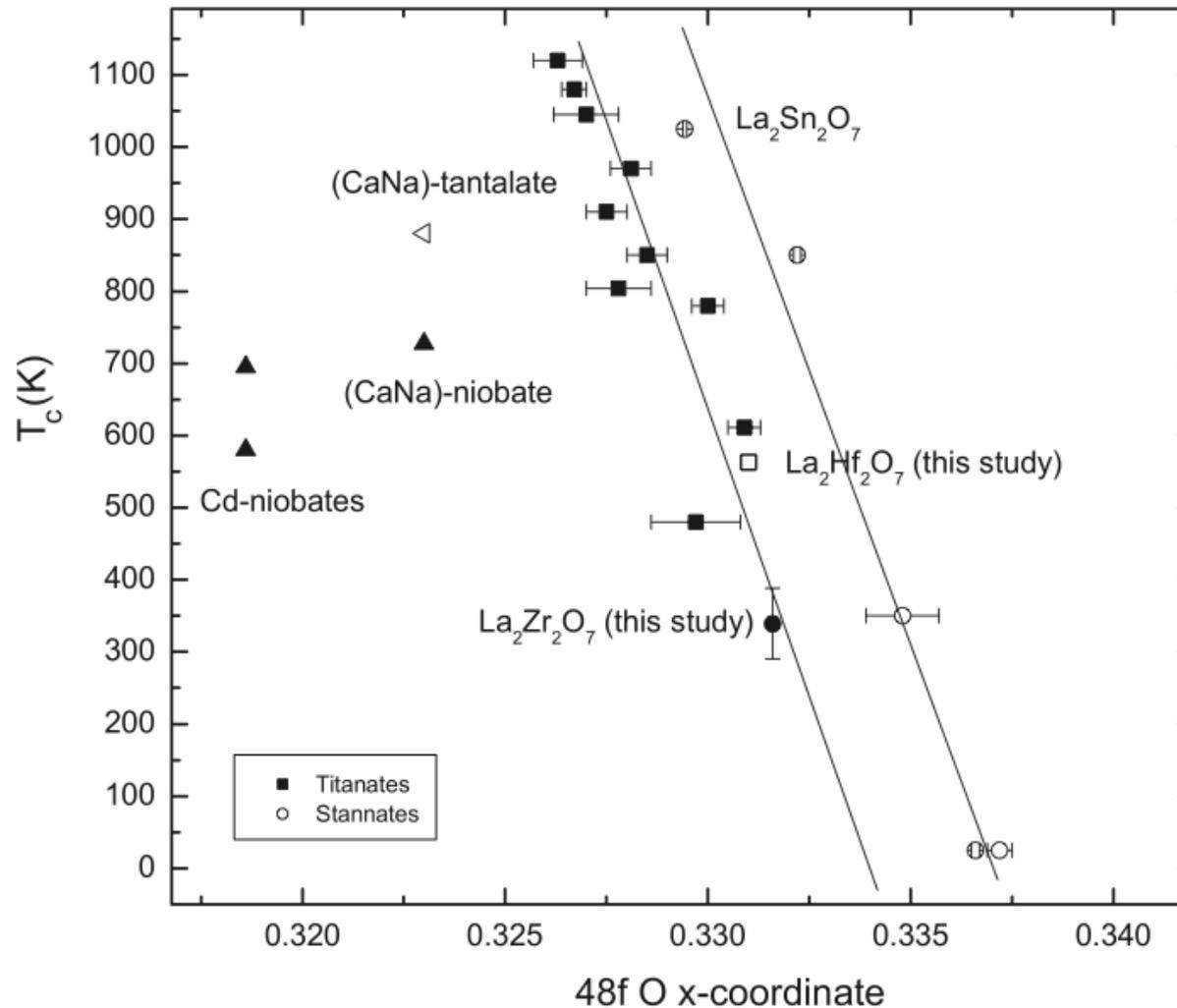
Experiments on live samples of Pu containing zircon ($ZrSiO_4$): NMR

I Farnan *et al*, Nature 2007



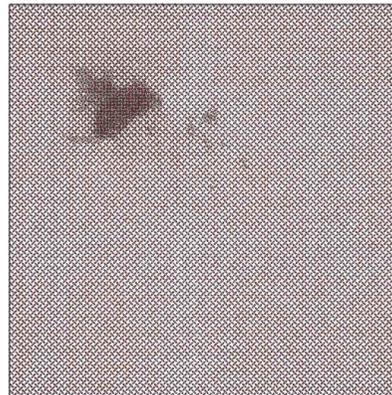
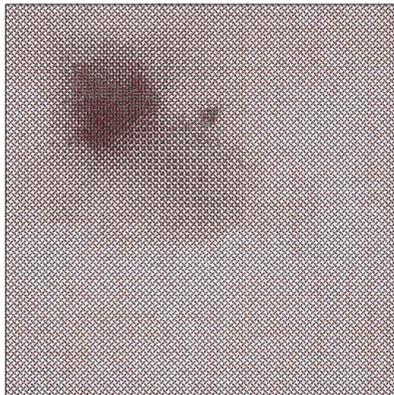
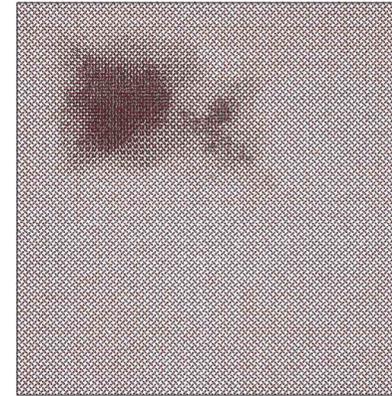
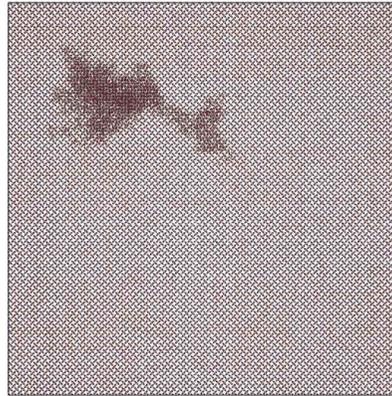
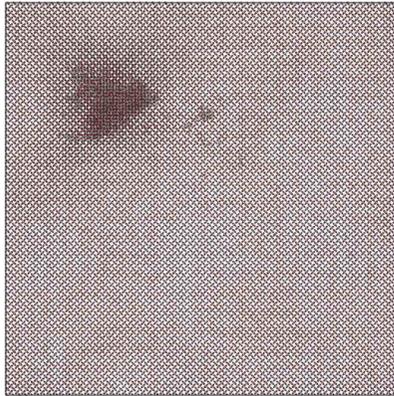
Experiments on irradiated samples of pyrochlores

GR Lumpkin *et al*, JPCM 2004



Large scale MD simulations based on empirical force fields

Rutile TiO₂



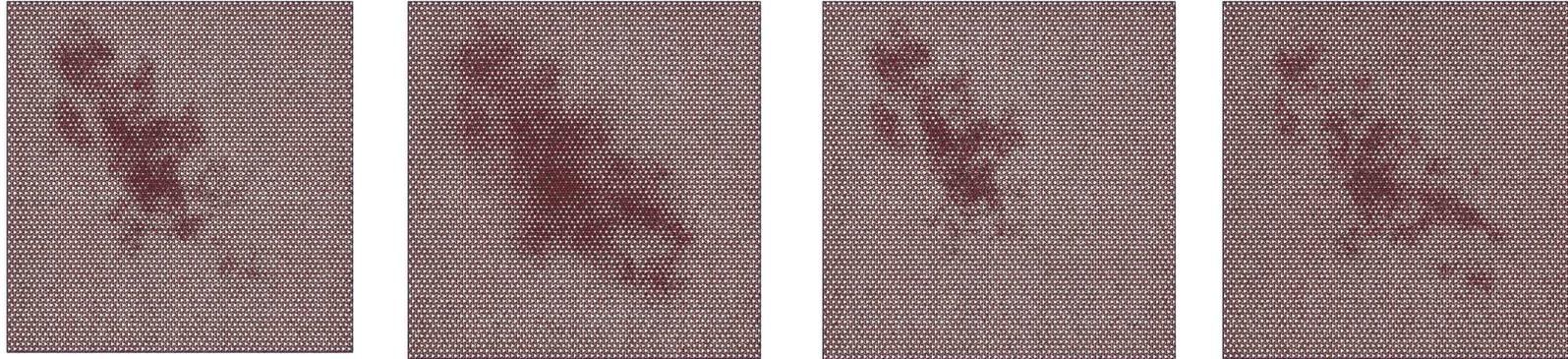
Kostya Trachenko



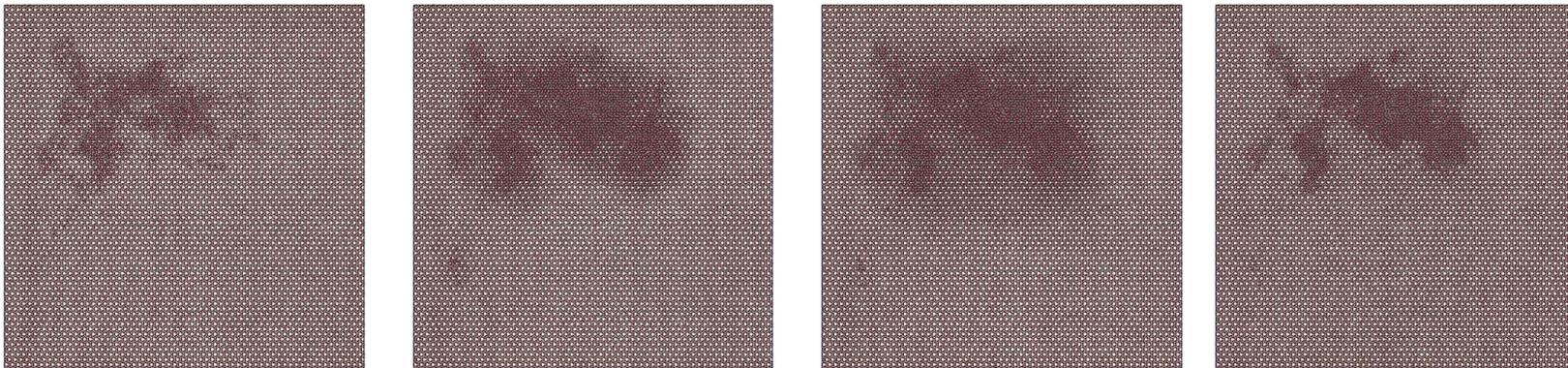
Martin Dove

Large scale MD simulations based on empirical force fields

Quartz SiO₂

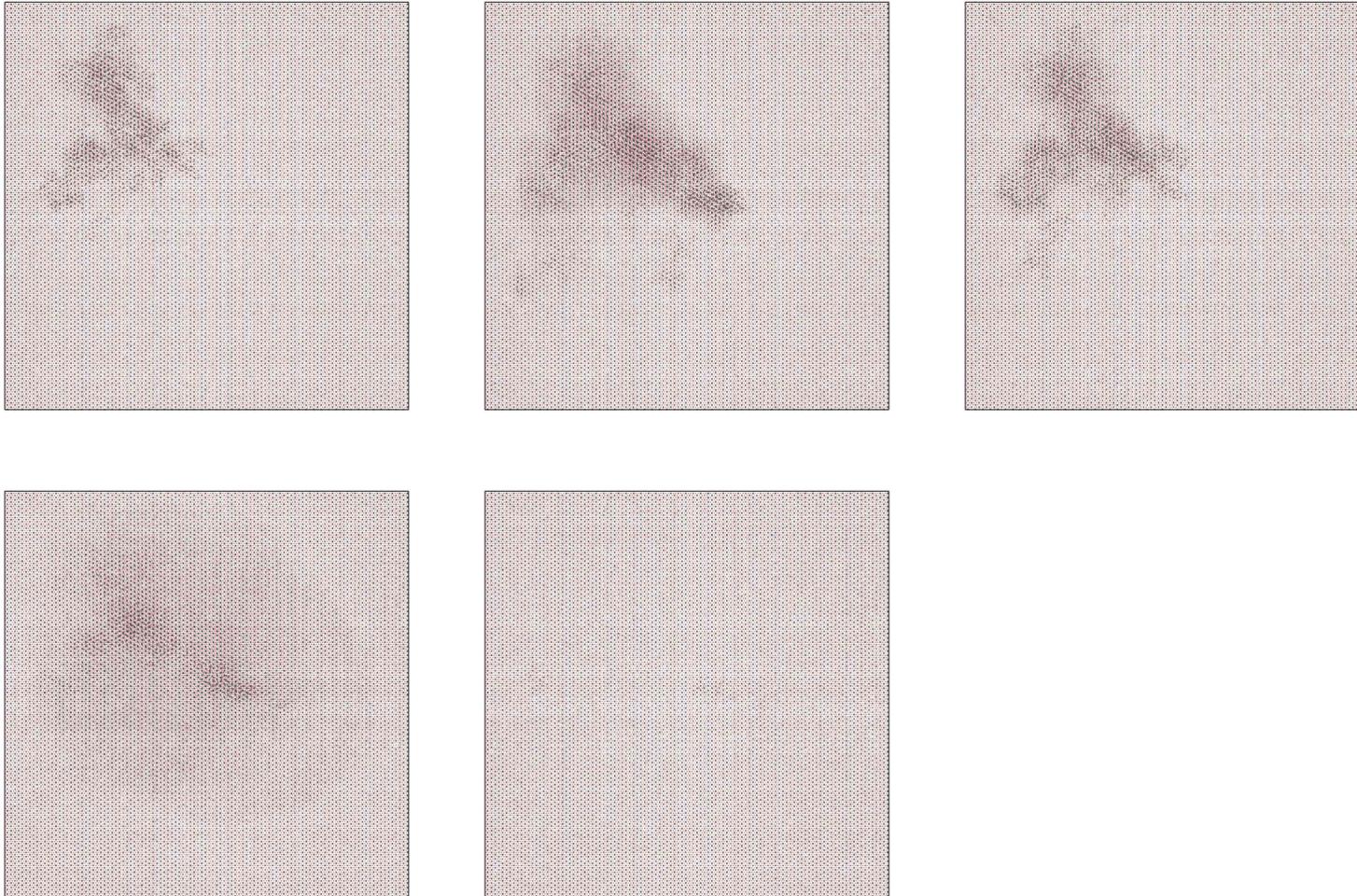


Quartz GeO₂



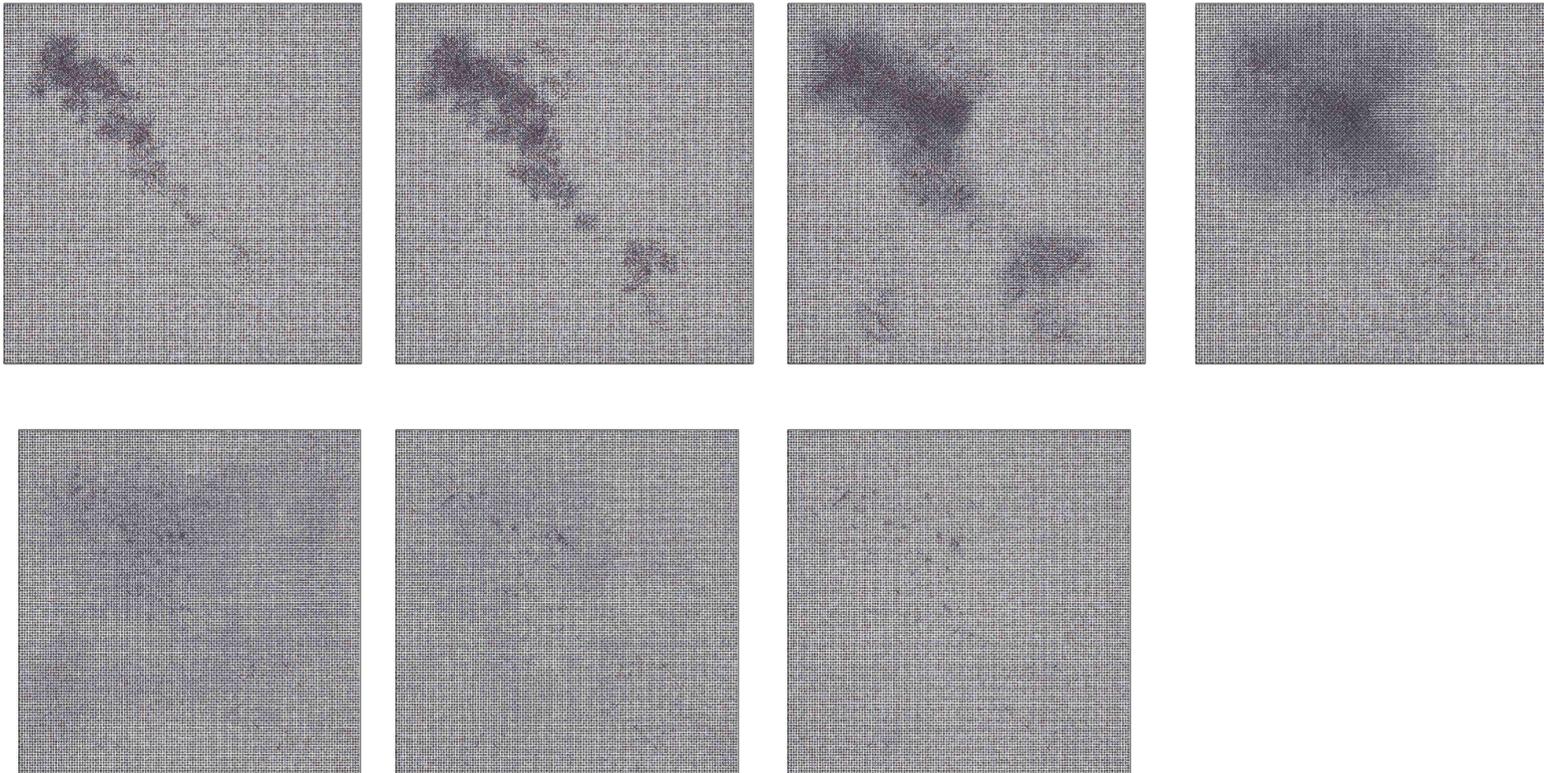
Large scale MD simulations based on empirical force fields

Corundum Al₂O₃



Large scale MD simulations based on empirical force fields

MgO



*Electrons heat up: effect on the material?
Effect on the simulations?*

*The ion moving in the solid transmits energy to electrons.
How much? How? Where? What consequences does it have?*

Materials soften if electronic subsystem substantially excited

Coupled electron-nuclei dynamics

FIRST: How much energy goes to electrons?

Not trivial either for exp or for th

Coupled dynamics of both electrons & nuclei

(realistic simulation demands ~ 2M atoms)

Multiscale: Different (decoupled?) timescales

Results from

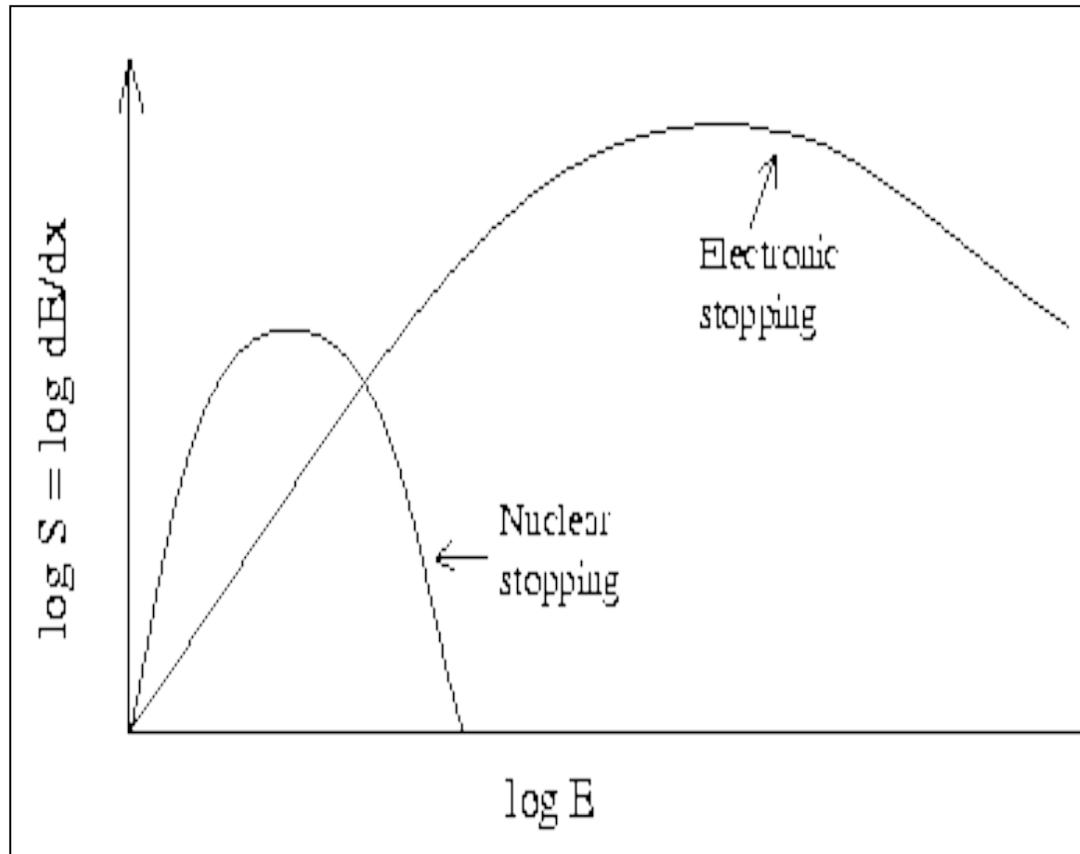
- this study (electronic excitation)*
- electron-phonon coupling +*
- heat conductivities for the electron
and phonon subsystems*

⇒ Continuum description of excess energy

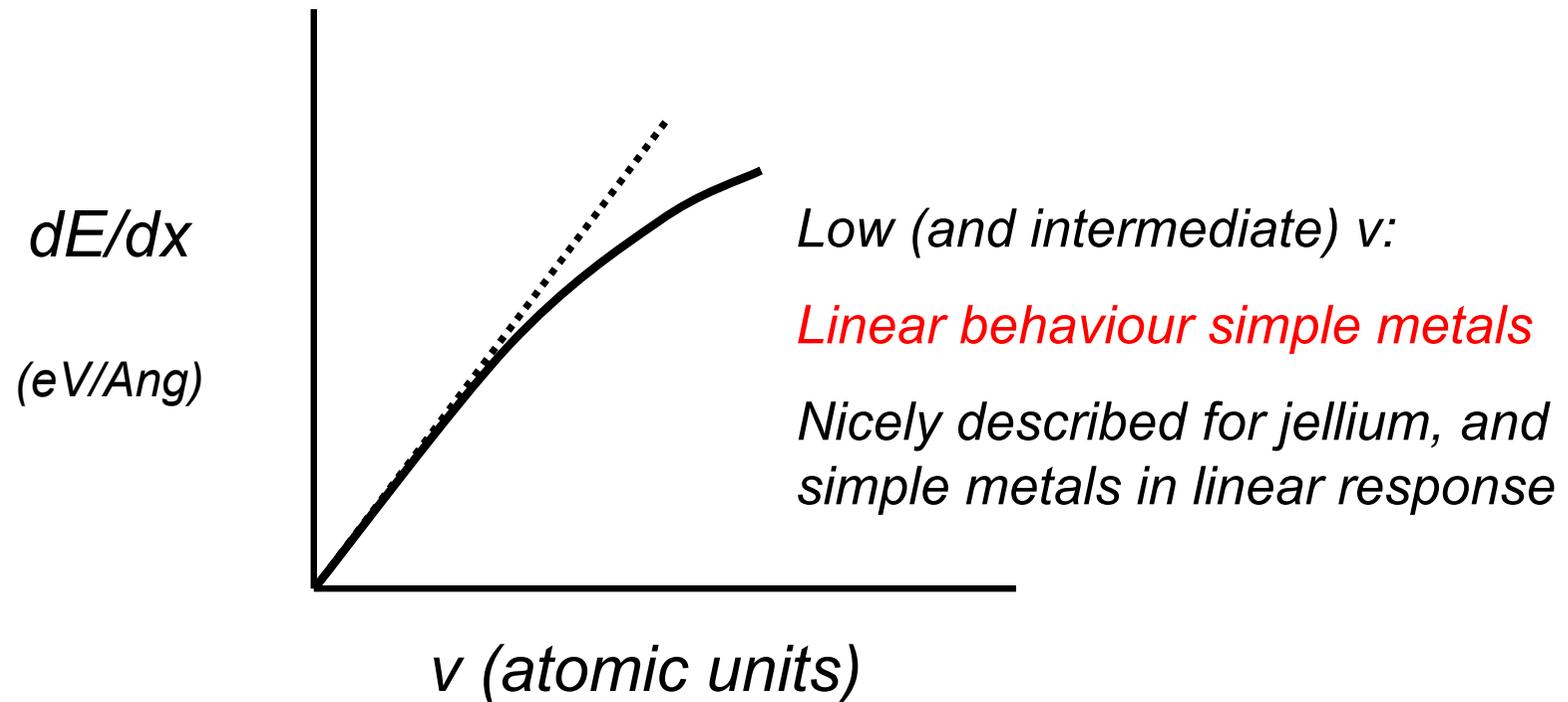
$$\Delta E_{el}(\mathbf{r}, t)$$

*First: how much energy is it pumped to
the electrons per unit time*

Electronic versus nuclear stopping



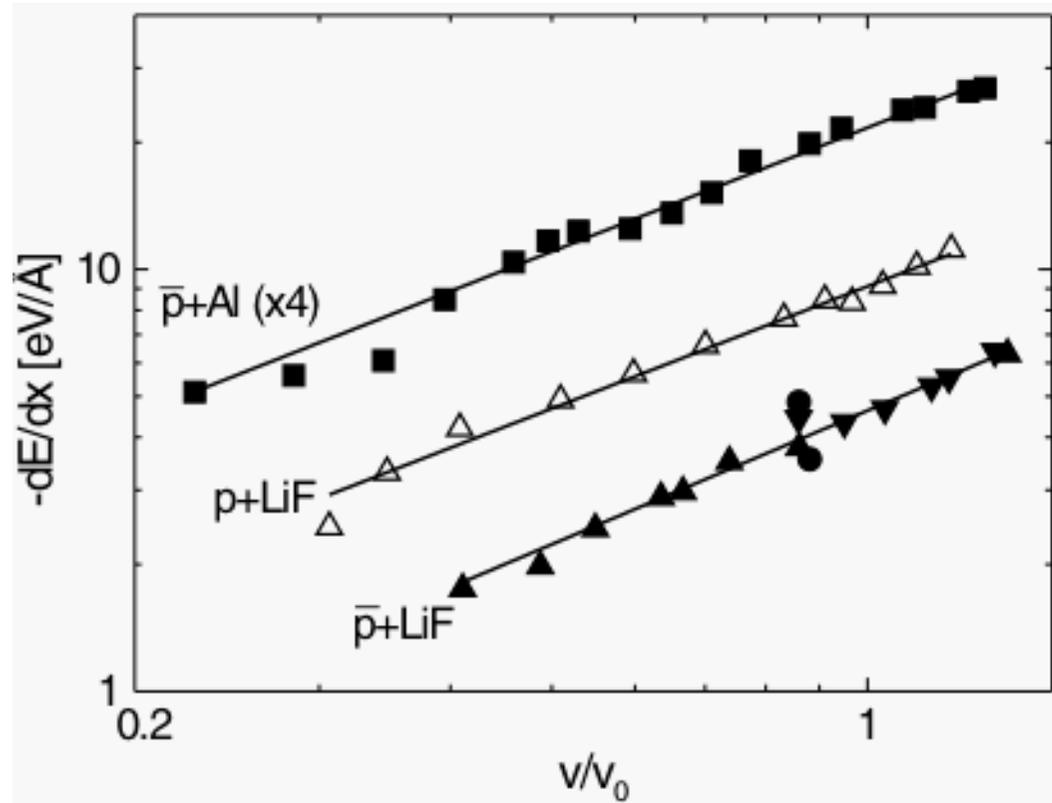
Electronic stopping power



Scale: few eV/Å at $v = 1$ a.u.

But what about insulators?

Protons and antiprotons into LiF thin films

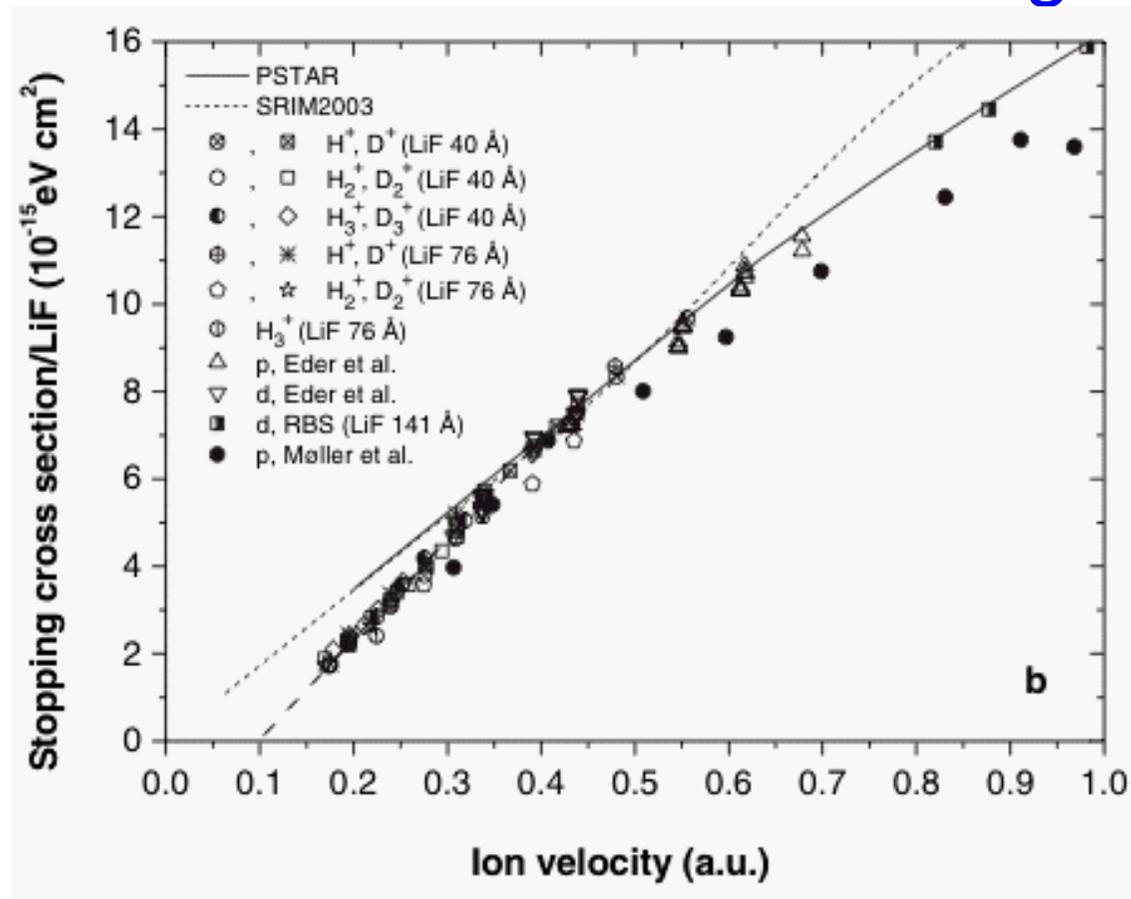


“Antiproton Stopping at Low Energies:
Confirmation of Velocity-Proportional Stopping Power”

SP Møller *et al.* PRL **88**, 193201 (2002) & PRL **93**, 042512 (2004)

Perfectly linear, no difference

Protons into LiF thin films again

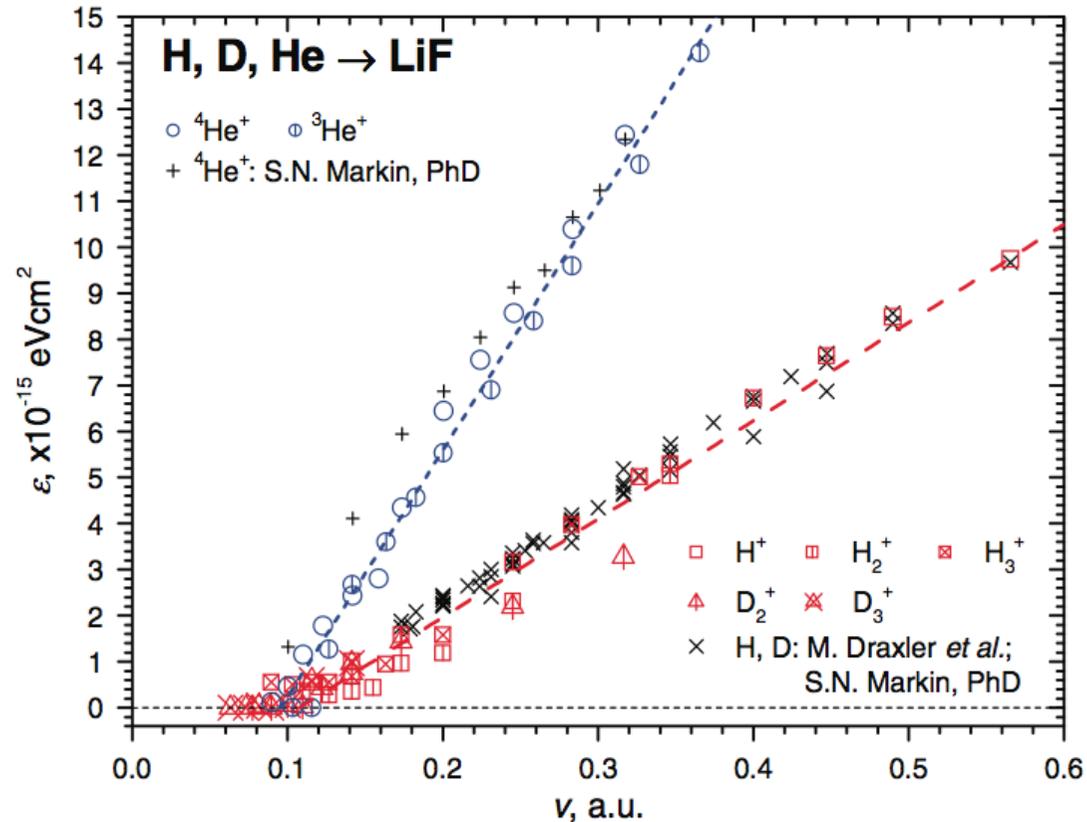


M. Draxler *et al*, PRL **95**, 113201 (2005)

Scale: $v = 0.1$ a.u. \Rightarrow Stopping ~ 1 eV/Ång

Threshold

Protons into LiF thin films again



S. Markin *et al*, PRL **103**, 113201 (2009)

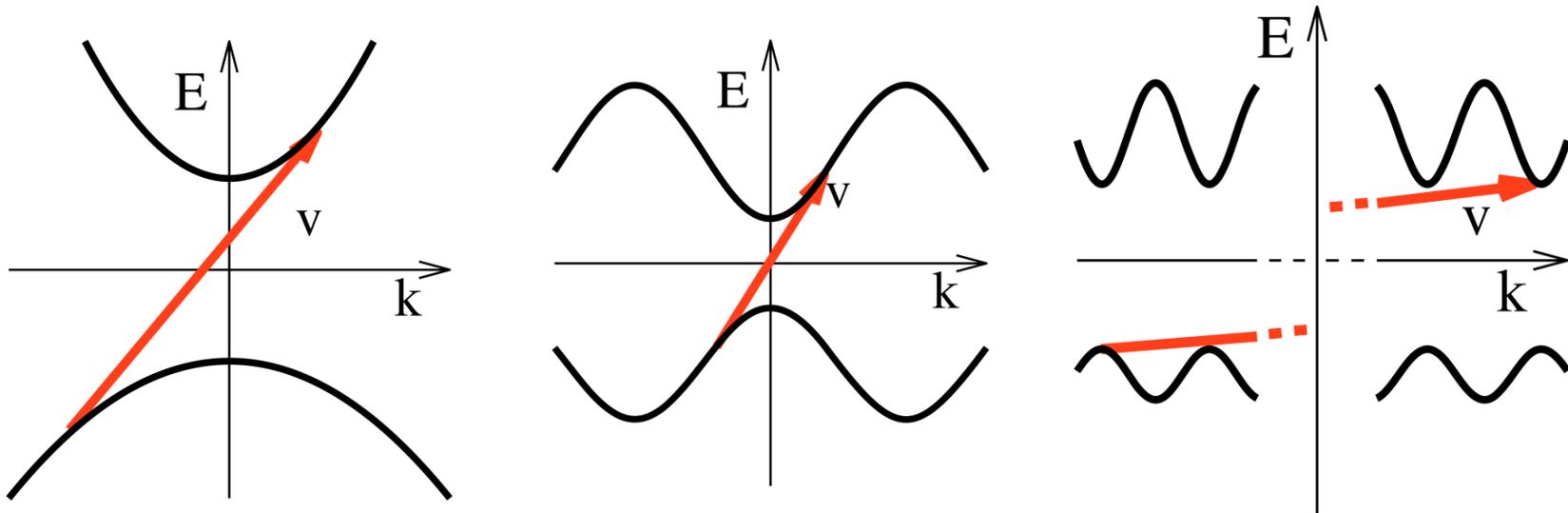
Scale: $v = 0.1 \text{ a.u.} \Rightarrow \text{Stopping} \sim 1 \text{ eV/\text{Ang}}$

Threshold

Threshold: what to expect?

TD Pertub. Th. (weak projectile potential)
e-h excitations such that

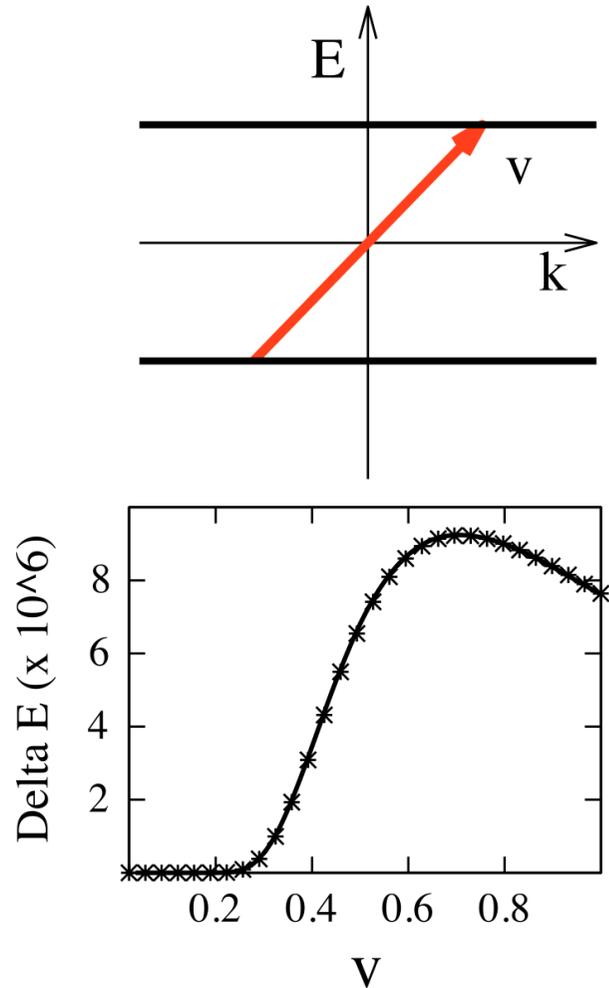
$$\Delta\varepsilon / \Delta k = v$$



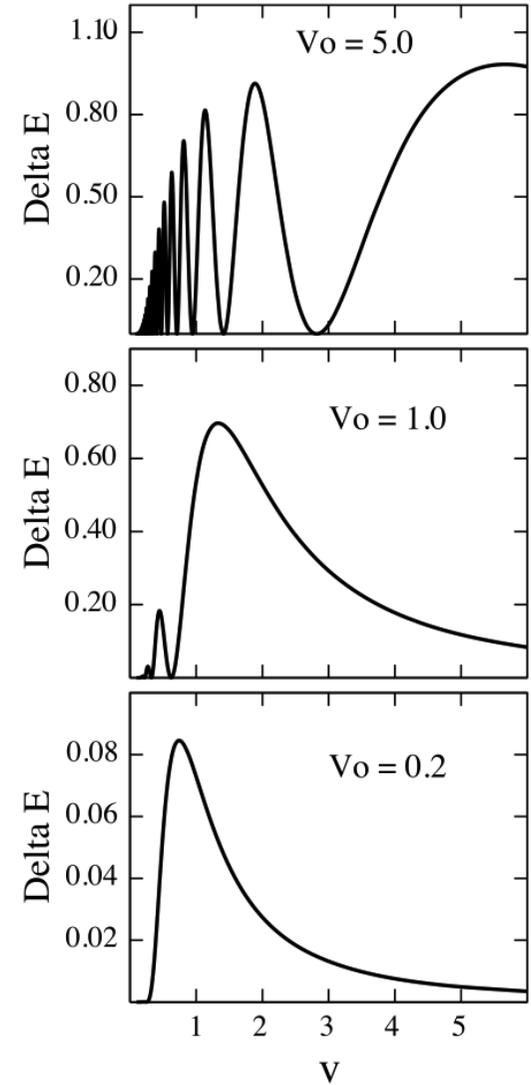
Strict threshold: $1/2 (m_e + m_h) v_c^2 = E_{gap}$

Excitation amplitude:
 $V(q=\Delta k)$

Flat-band limit: Simple model



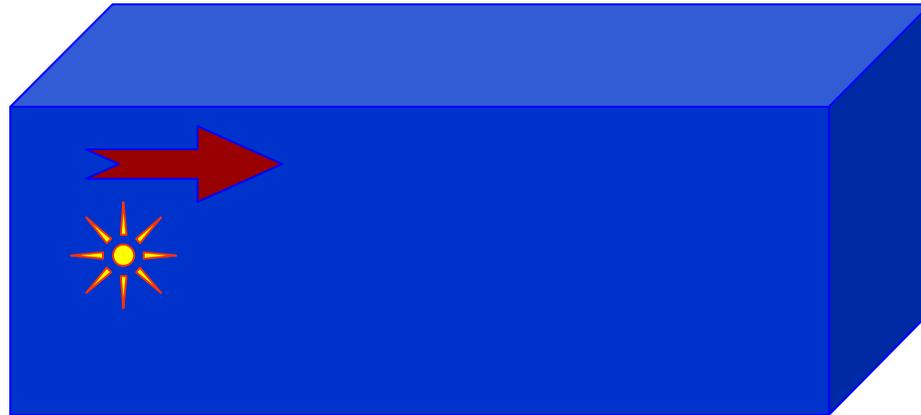
Gaussian perturbation scaling as V_0



$$S(v) \sim e^{-\left(\frac{1}{v}\right)^\eta}; \quad \eta \in \left[\frac{3}{2}, 2\right]$$

Realistic simulations: TD-DFT

- *Supercell of insulator's bulk*
- *Periodic boundary conditions*
- *Density functional theory*
- *Add external charge (potential)*



- *Move it and follow electron wave-functions with Time-Dependent DFT*

The SIESTA method

Linear-scaling DFT based on NAOs (Numerical Atomic Orbitals)

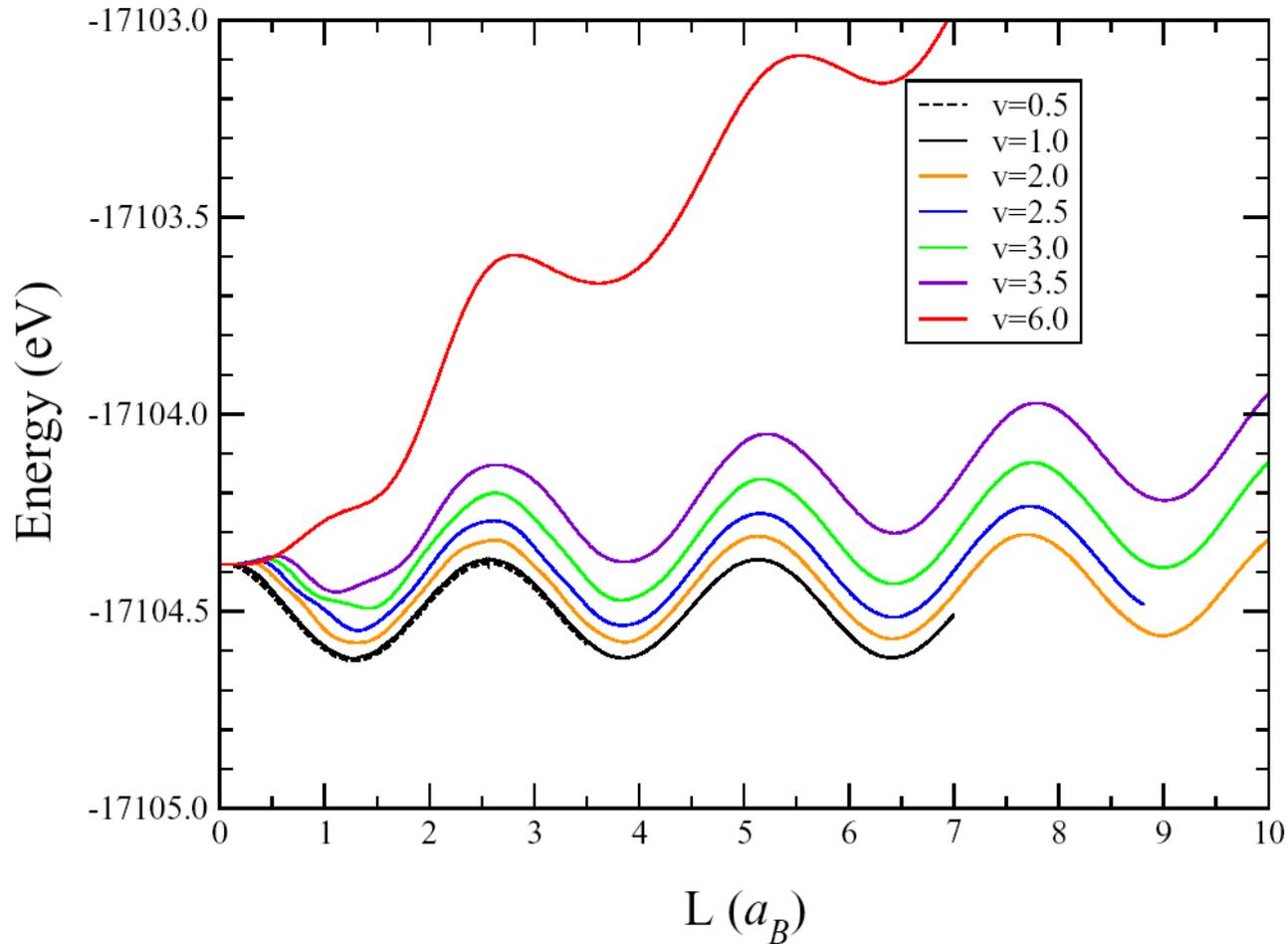
P. Ordejon, E. Artacho & J. M. Soler , Phys. Rev. B 53, R10441 (1996)

- *Born-Oppenheimer* (relaxations, mol. dynamics)
- *DFT* (LDA, GGA)
- *Pseudopotentials* (norm conserving, factorised)
- *Numerical atomic orbitals as basis* (finite range)
- *Numerical evaluation of matrix elements* (3D grid)

Implemented in the SIESTA program

J. M. Soler, E. Artacho, J. D. Gale, A. Garcia, J. Junquera, P. Ordejon & D. Sanchez-Portal, *J. Phys.: Condens. Matter* **14**, 2745 (2002)

Energy as a function of distance: LiF



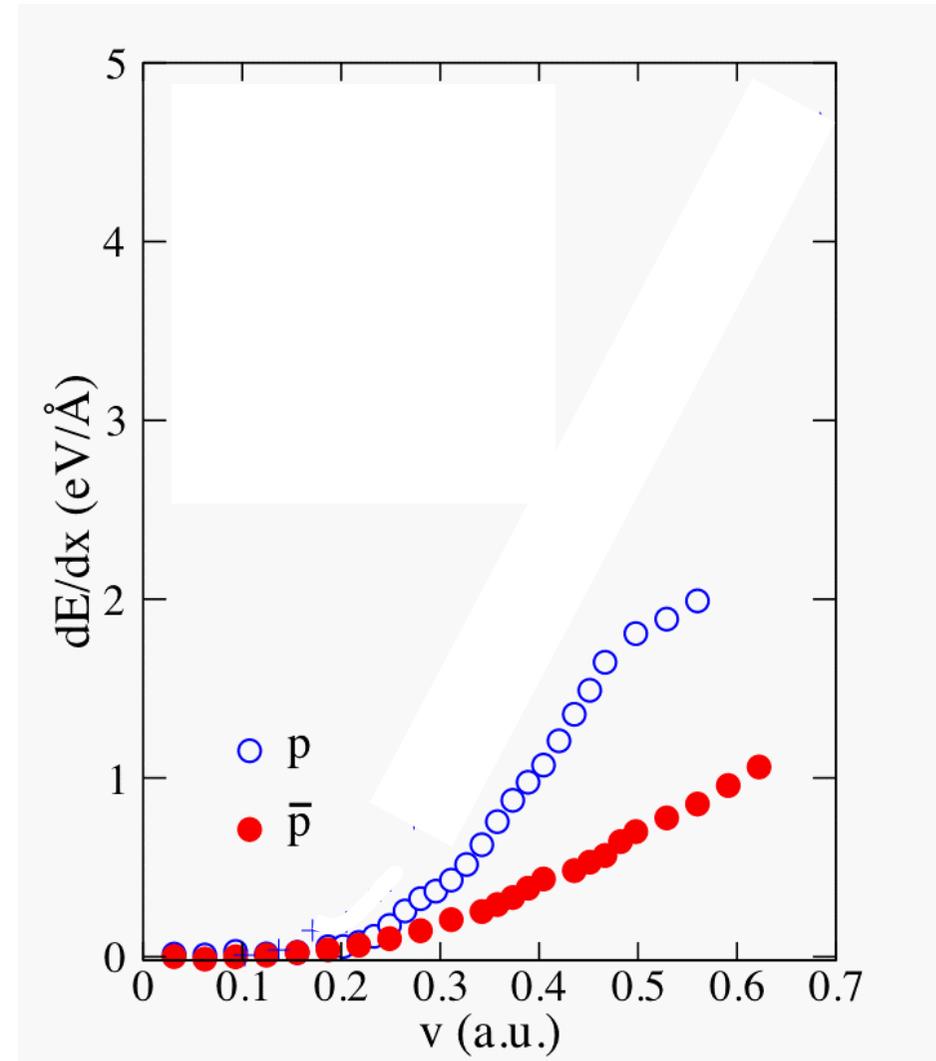
Quite stationary! Short transient, no obvious oscillation

Rate of energy transfer: electronic stopping power

*Protons and antiprotons
through LiF*

Threshold ~ 0.2 a.u. (exp ~ 0.1)

Ratio $SP_p/SP_a \sim 2.4$ (exp ~ 2.1)



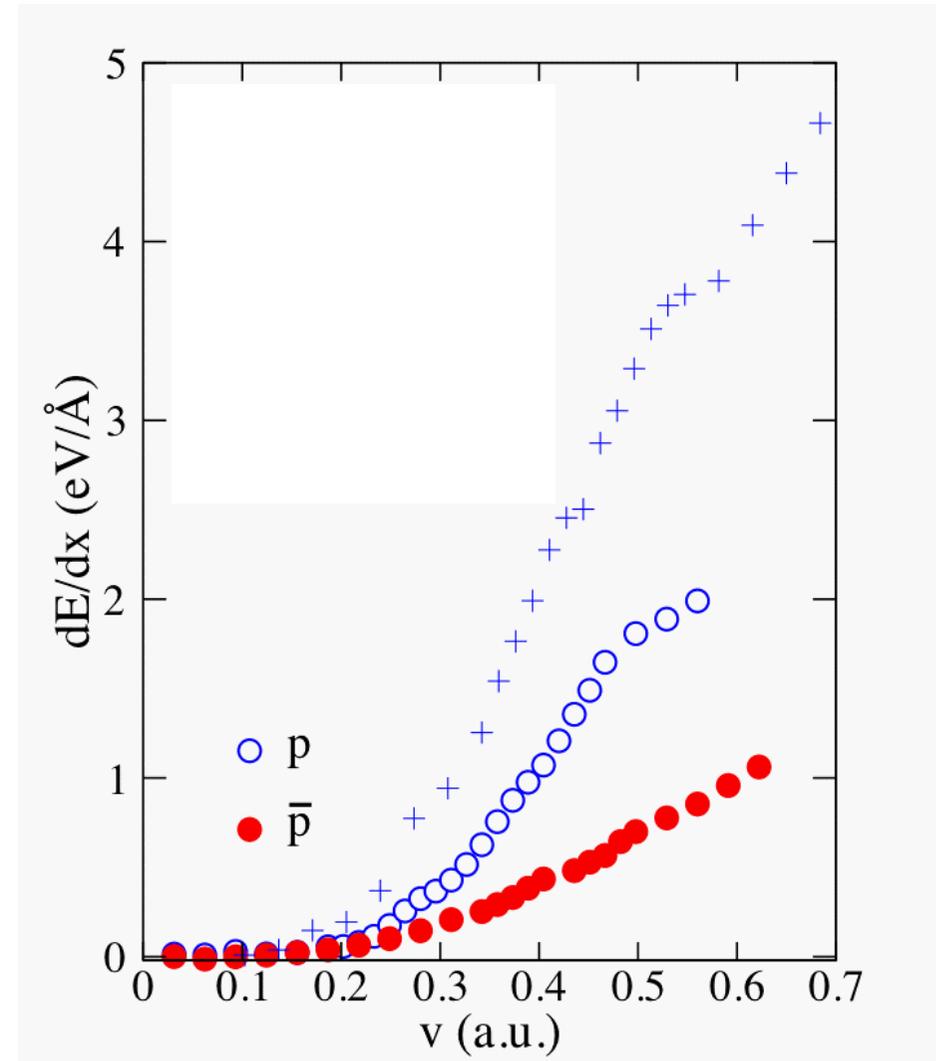
Rate of energy transfer: electronic stopping power

Protons and antiprotons through LiF

Threshold ~ 0.2 a.u. (exp ~ 0.1)

Ratio $SP_p/SP_a \sim 2.4$ (exp ~ 2.1)

*Absolute value: improve basis;
sp basis along trajectory (for p)*



Rate of energy transfer: electronic stopping power

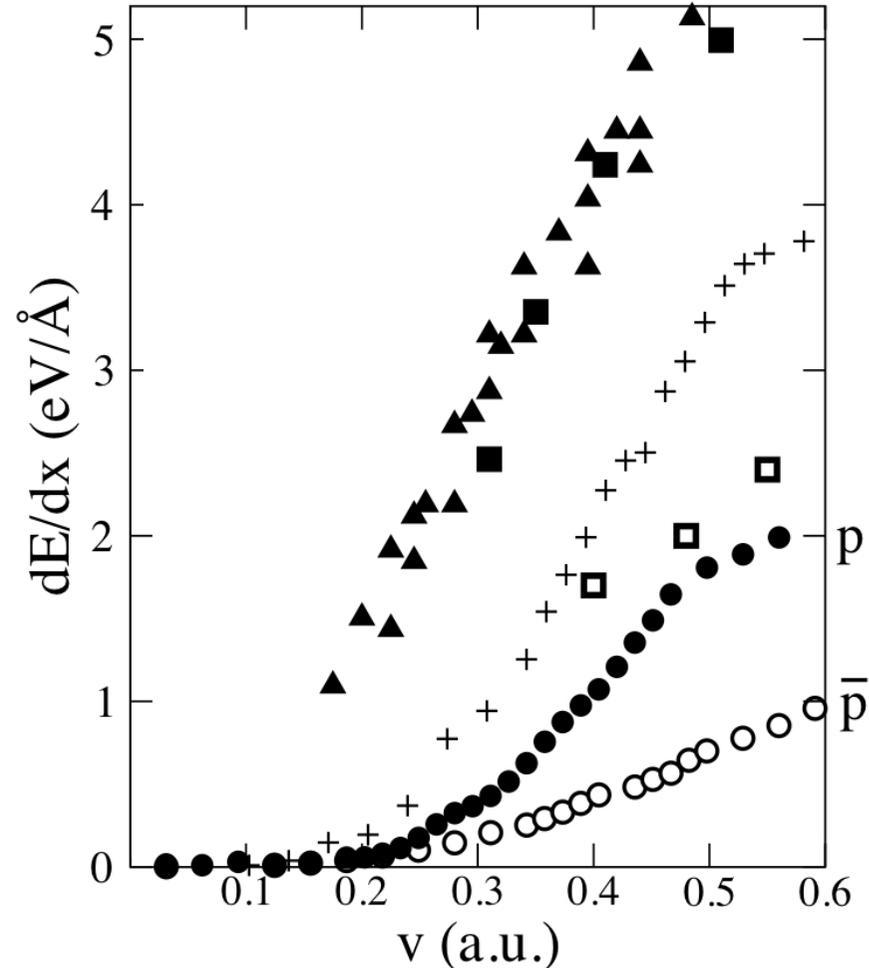
Protons and antiprotons through LiF

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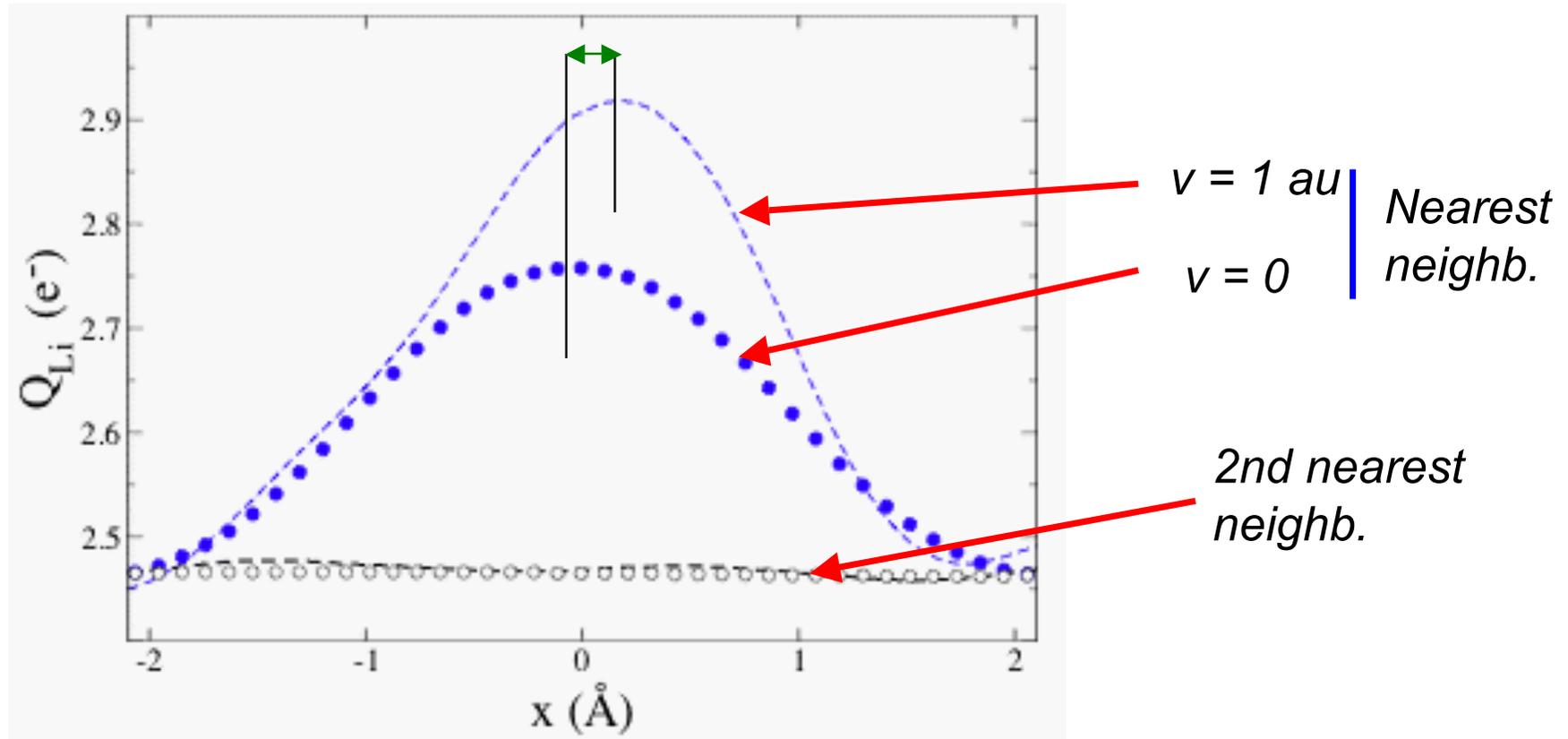
Ratio $SP_p/SP_a \sim 2.4$ (exp ~ 2.1)

Absolute value $\sim 20\text{-}50\%$ too small
as compared to experiment

But: This is channelling, exp is average



Evolution of the charge on nearby Li atoms



Position of projectile along trajectory ($x=0$ closest to nearest Li)

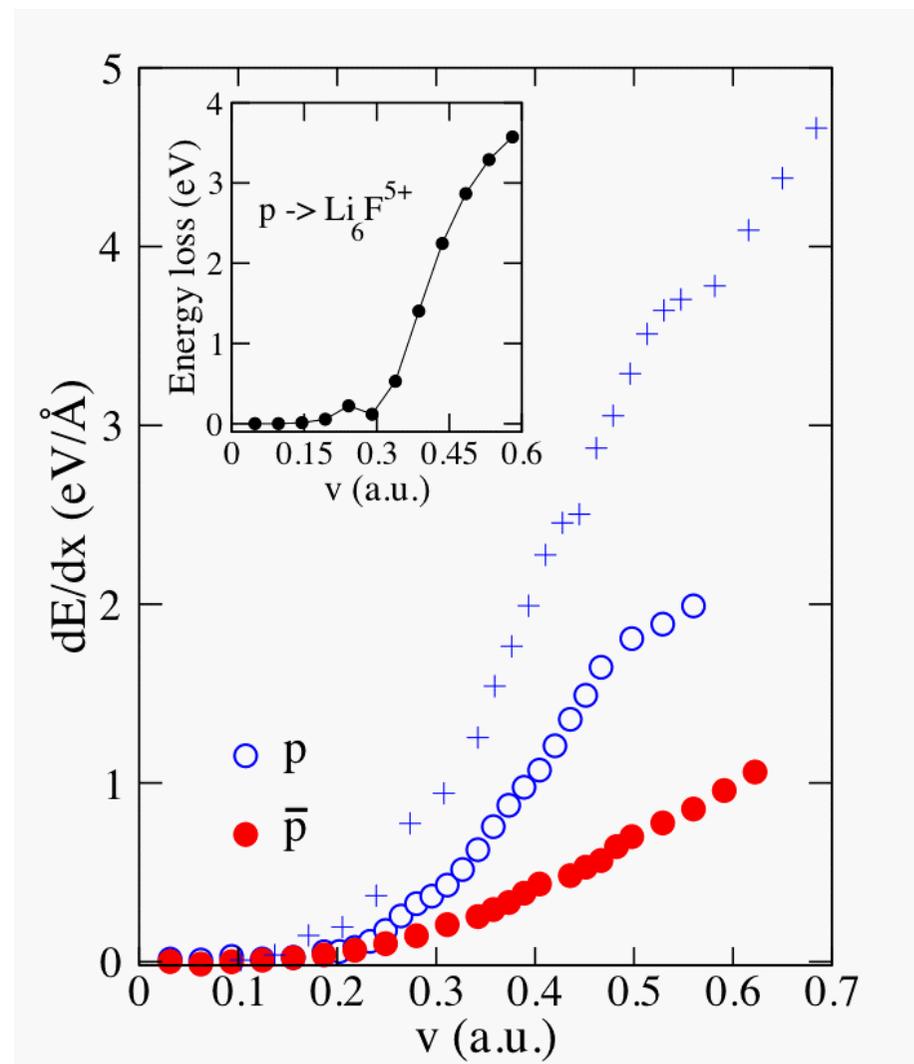
Screening of charge enhanced at finite v

Extremely short-ranged mechanism! Why?

Locality in the electronic stopping power

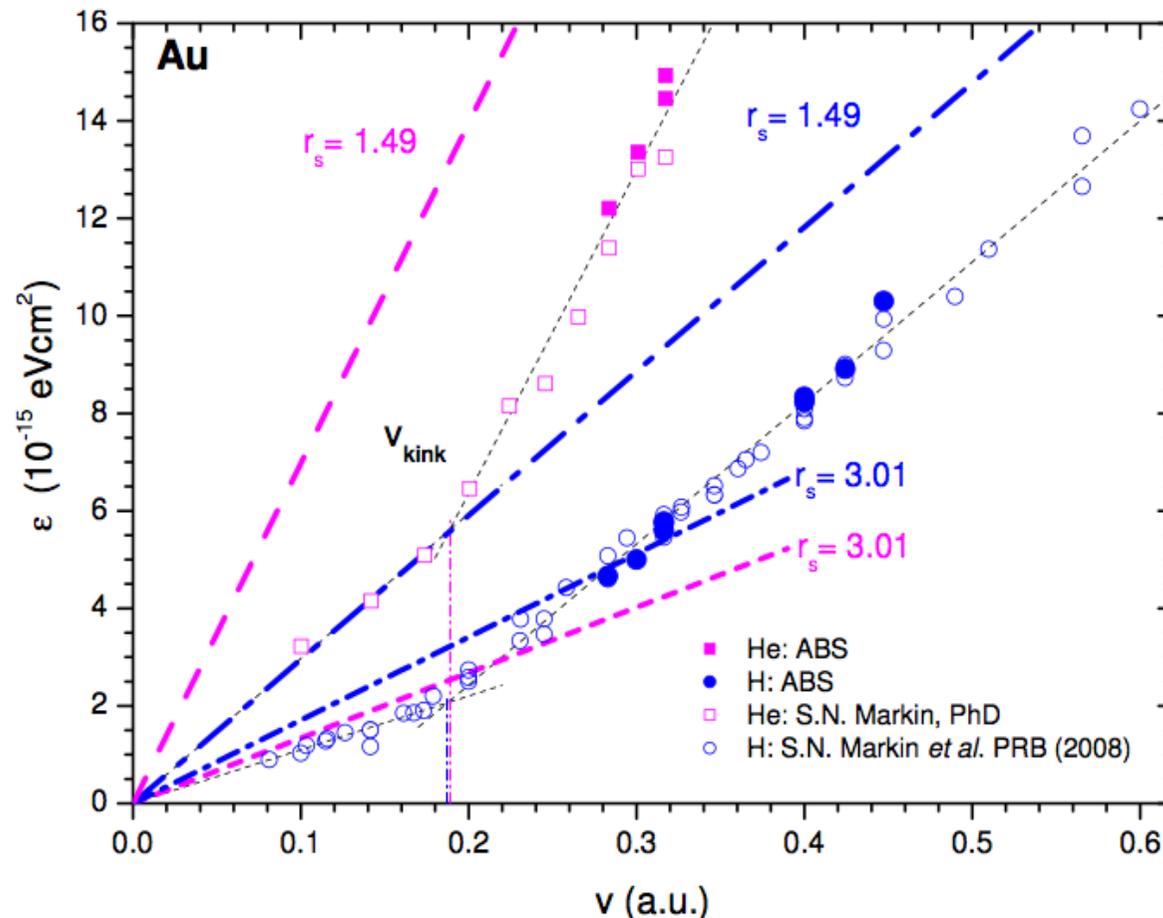
Protons in LiF

*Compare bulk with
small cluster Li_6F^{5+}*



What about metals beyond jellium?

Noble and transition metals



Electron heating by H and He projectiles in bulk Au

S. N. Markin, D. Primetzhofer, M. Spitz, and P. Bauer, PRB 2009

Summary

- *Using TD-DFT for obtaining the energy transfer from moving ions to electrons in insulators.*
- *New approach, lots of approximations, but complementary to previous studies*
- *Offers new kinds of information*
- *Need to couple with ion motion*

Important:

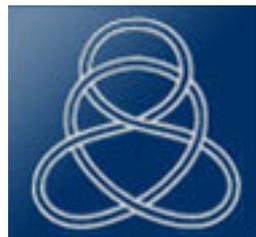
Recoiling Th-ion velocity: around threshold!

Stopping + Rad Dam + CEID

Funding



British Nuclear Fuels Ltd



Miller Institute for Basic
Research in Science

University of California, Berkeley