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**Preparatory School to the Winter College on Optics and the Winter College on  
Optics: Advances in Nano-Optics and Plasmonics**

*30 January - 17 February, 2012*

**Non linear optics of chiral materials**

C. Sibilia

*University Roma La Sapienza  
Roma  
Italy*



**SAPIENZA**  
UNIVERSITÀ DI ROMA

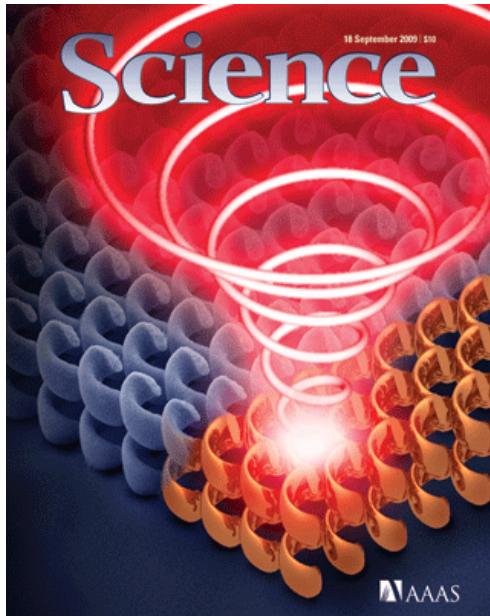
# Nonlinear Optics of chiral materials

**Concita Sibilia**

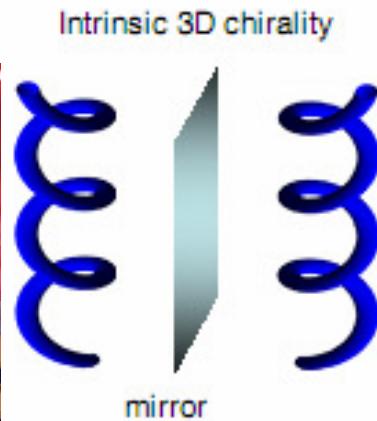
Dipartimento SBAI-Università' di Roma La Sapienza



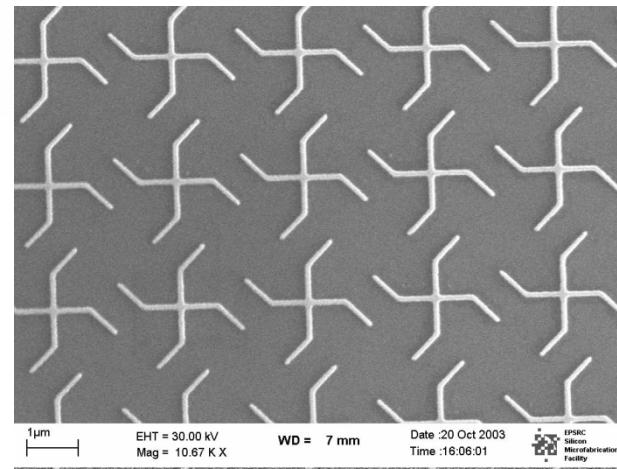
# 3D and 2D Chirality



3D chiral metamaterial



Absence of  
mirror rotation  
axis



2D chiral metamaterial

IOP PUBLISHING  
J. Opt. A: Pure Appl. Opt. 11 (2009) 074009 (7pp)

JOURNAL OF OPTICS A: PURE AND APPLIED OPTICS  
doi:10.1088/1464-4258/11/7/074009

## Extrinsic electromagnetic chirality in metamaterials

E Plum, V A Fedotov and N I Zheludev<sup>1</sup>

Optoelectronics Research Centre, University of Southampton, Southampton SO17 1BJ, UK

E-mail: [niz@orc.soton.ac.uk](mailto:niz@orc.soton.ac.uk)

In the last years a lot of effort was devoted to the study on chiral metamaterials. Both 3D- and 2D-chiral metamaterials were studied.

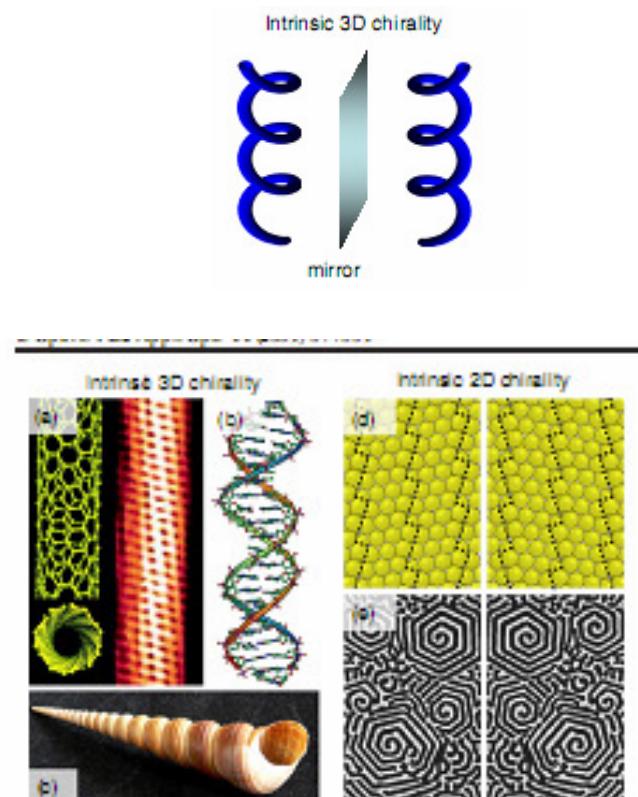


Figure 2. Examples of *intrinsically* chiral artificial and natural objects. 3D chirality: (a) single-wall carbon nanotube (structure and tunnelling electron microscope image) [1], (b) DNA double helix [2] and (c) seashell from the Turritellidae family [3]. 2D chirality: (d) surface (643) of a gold monocrystal, (e) surface of a metal–metal anhydride [4].

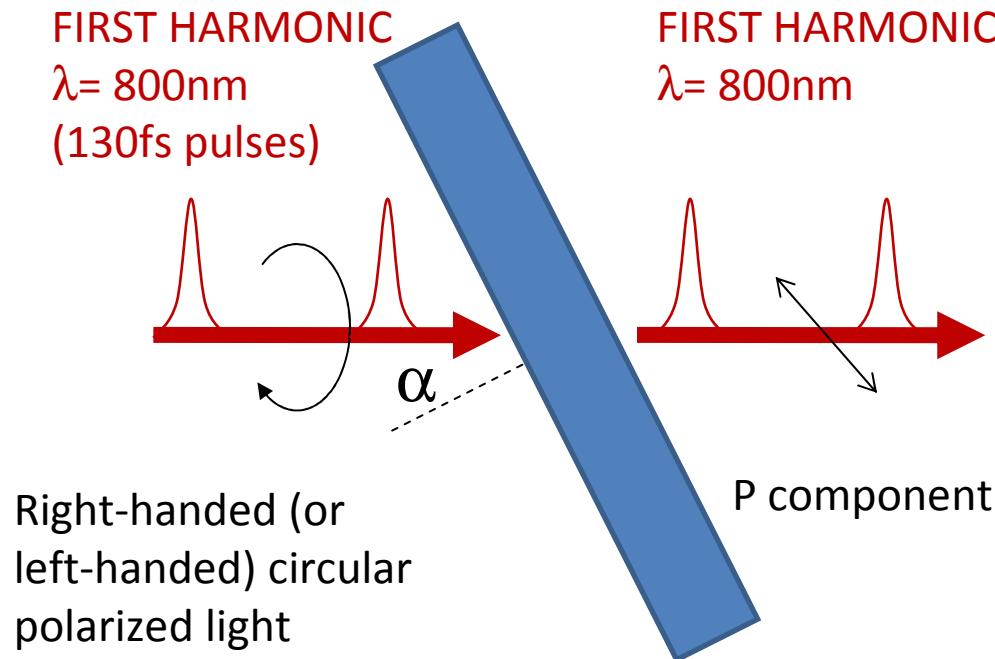
## From N.Zheludev

The term ‘chirality’ (or handedness from the Greek  $\chi\epsilon\rho$  ‘hand’) usually refers to an intrinsic sense of twist (or helicity) of a three-dimensional object, such as a helix, which has two different mirror forms called enantiomers (see figure 1(a)). 3D chirality appears to be a signature of life forms and can be found in many important organic molecules and proteins, as well as in inorganic structures (see figures 2(a)–(c)). 3D chirality results in such fundamental polarization effects as optical activity (a medium’s ability to rotate the polarization state of light) and the associated phenomenon of circular dichroism, which manifests itself as a difference in transmission levels for left and right circularly polarized light (see figure 3(a)). Since its discovery in 1811 by Dominique Arago, optical activity has played an important role in analytical chemistry, crystallography, molecular biology and the food industry, and is now used as a test for detecting life forms in space missions. The recognition of chirality as a source of negative refraction of light [5–11] that is needed for the creation of a perfect lens [12] has recently inspired intense work in developing optically active microwave and photonic chiral metamaterials [13–18] and has led to the observation of



# Linear measurements

The measurements performed as a function of the incidence angle  $\alpha$



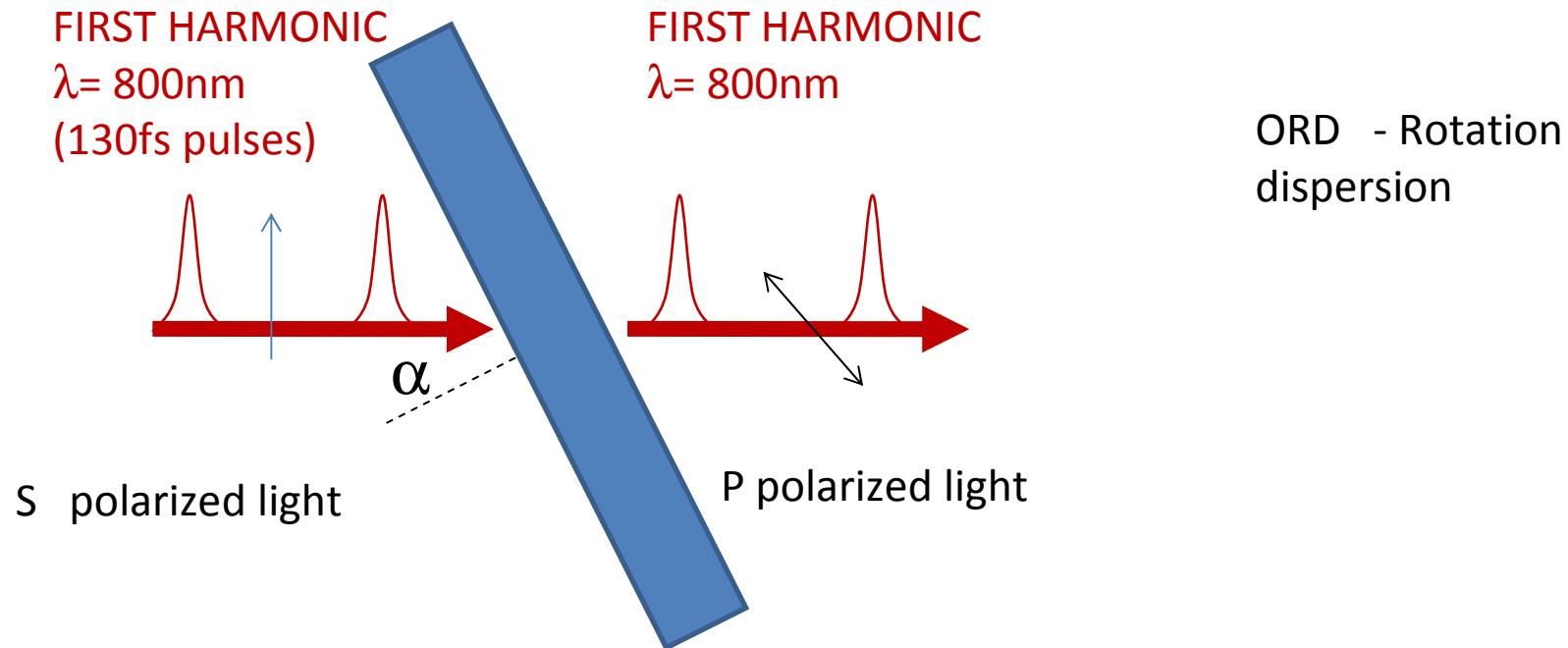
circular dichroism (CD)

$$CD = \frac{I_L^\omega - I_R^\omega}{(I_L^\omega + I_R^\omega)/2}$$



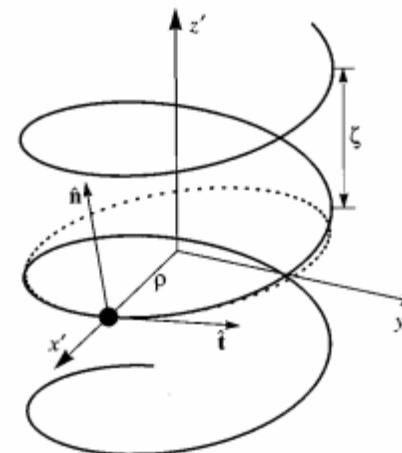
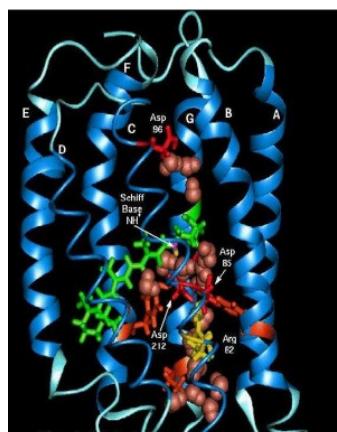
# Linear measurements

The measurements performed as a function of the incidence angle  $\alpha$





As the electrons of chiral molecules are displaced from their equilibrium by the application of the electromagnetic field, they are forced helical like paths



$$H = -\mu \cdot E - m \cdot B - Q : \nabla E + \dots \text{weak}$$



# Nonlinear optics, chirality, magneto-optics: a serendipitous road [Invited]

Andre Persoons<sup>1,2,3,\*</sup>

<sup>1</sup>Department of Chemistry, University of Leuven, B-3001 Leuven, Belgium

<sup>2</sup>College of Optical Sciences, University of Arizona, Tucson, Arizona 85721, USA

<sup>3</sup>[apersoons@optics.arizona.edu](mailto:apersoons@optics.arizona.edu)

[\\*andre.persoons@fys.kuleuven.be](mailto:andre.persoons@fys.kuleuven.be)

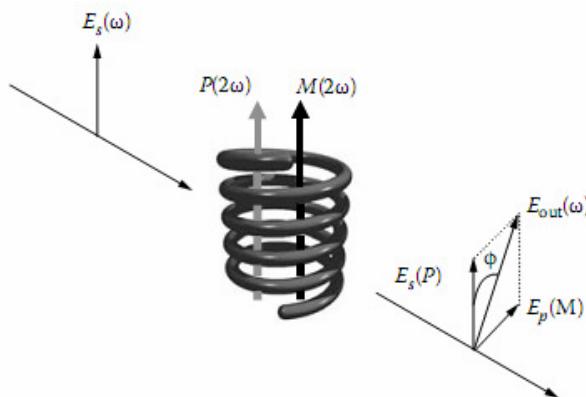
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1 May 2011 / Vol. 1, No. 1 / OPTICAL MATERIALS EXPRESS 5

## Nonlinear Optical Activity



$$\chi_{ijk}^{(2)} = N \sum_{ijk} \langle (I \cdot i)(J \cdot j)(K \cdot k) \rangle \beta_{ijk}$$

PHYSICAL REVIEW

VOLUME 130, NUMBER 3

1 MAY 1963

## Nonlinear Optical Properties of Solids: Energy Considerations\*

P. S. PERSHAN

Division of Engineering and Applied Physics, Harvard University, Cambridge, Massachusetts

(Received 17 December 1962)



## Nonlinear Optical Frequency Polarization in a Dielectric

ERIC ADLER

*IBM Watson Laboratory and Physics Department, Columbia University, New York, New York*

(Received 25 July 1963)

A perturbation calculation of optical frequency polarization quadratic in the Maxwell field is made for a dielectric in which the electrons are localized on units of the crystal. The result is expressed in a power series in  $K\langle r \rangle \sim 10^{-8}$  where  $\langle r \rangle$  is the size of the unit and  $K$  is the wave number of the field. The term of zeroth order is the electric-dipole term which vanishes in a crystal with a center of symmetry. The term first order in  $K\langle r \rangle$  is separated into electric-quadrupole and magnetic-dipole contributions by introducing a special gauge for the electromagnetic potentials. Higher power terms are neglected.

PRL 106, 103902 (2011)

PHYSICAL REVIEW LETTERS

week ending  
11 MARCH 2011

### Circular Dichroism in Biological Photonic Crystals and Cubic Chiral Nets

M. Saba,<sup>1</sup> M. Thiel,<sup>2</sup> M. D. Turner,<sup>3</sup> S. T. Hyde,<sup>4</sup> M. Gu,<sup>3</sup> K. Grosse-Brauckmann,<sup>5</sup> D. N. Neshev,<sup>6</sup>  
K. Mecke,<sup>1</sup> and G. E. Schröder-Turk<sup>1,\*</sup>

## Second-Order Nonlinear Optical Imaging of Chiral Crystals

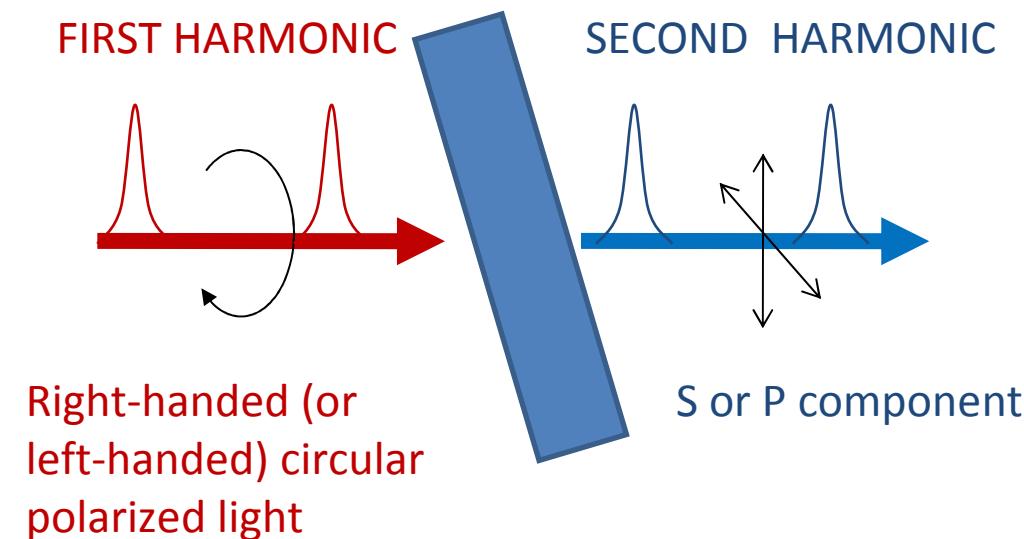
David J. Kissick, Debbie Wanapun, and Garth J. Simpson

Department of Chemistry, Purdue University, West Lafayette, Indiana 47907;



# SHG as chirality detector

It is possible to investigate chirality via second harmonic generation experiments



One useful parameter is:  
Second harmonic generation-circular dichroism (SHG-CD)

$$SHG - CD = \frac{I_L^{2\omega} - I_R^{2\omega}}{(I_L^{2\omega} + I_R^{2\omega})/2}$$

ORD

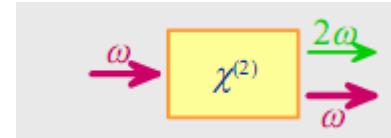
$$R(2\omega) = \frac{E_s(2\omega)}{E_p(2\omega)} = \frac{f_s E_p^2(\omega) + g_s E_s^2(\omega) + h_s E_p(\omega) E_s(\omega)}{f_p E_p^2(\omega) + g_p E_s^2(\omega) + h_p E_p(\omega) E_s(\omega)}$$



- Second order

$$P(t) = \chi^{(2)} E^2(t) = \chi^{(2)} E_\omega^2 e^{-i2\omega t}$$

→ second-harmonic generation



## Multipolar Symmetries in SecondOrder Nonlinear Optics

- electric-dipole and higher-multipole nonlinearities –
- isotropic and centrosymmetric materials



# Nonlinear susceptibility expansion

$$\chi(E) = \chi^{(1)} + \chi^{(2)}E + \chi^{(3)}EE + \dots$$

$$P(E) = \epsilon_0 (\chi^{(1)} + \chi^{(2)}E + \chi^{(3)}E^2 + \dots) E$$

$$\vec{P}_{eff}(2\omega) = \vec{P}_D(2\omega) - \frac{1}{\omega} \vec{k} x \vec{M}(2\omega)$$

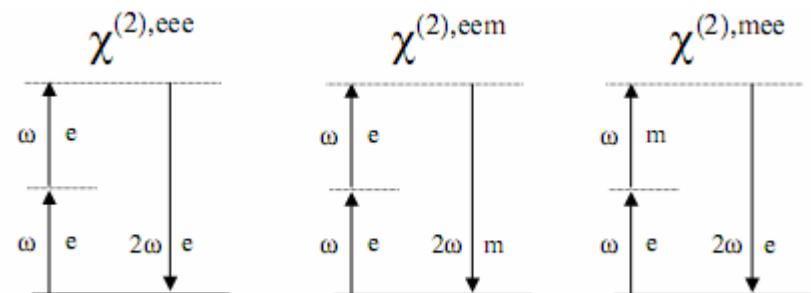
$$\vec{P}_D(2\omega) = \sum_{j,k} \chi_{ijk}^{(2)eee} E_j(\omega) E_k(\omega) + \chi_{ijk}^{(2)eem} E_j(\omega) B_k(\omega)$$

Dipole dipole source

$$\vec{M}(2\omega) = \sum_{j,k} \chi_{ijk}^{(2)mee} E_j(\omega) E_k(\omega)$$

Nonlinear magnetization

Magnetic dipole source





- Second-harmonic generation

$$P_i = \chi_{ijk}^{eee} E_j E_k + \chi_{ijk}^{eem} E_j B_k + \chi_{ijkl}^{eeQ} E_j \nabla_k E_l$$

axial      4th rank

- Magnetic and quadrupole tensors

- symmetry properties are different from those of the electric-dipole tensor

→ electric-dipole-forbidden  
effects can occur



# $\chi(2)$ Nonlinear Interaction

bulk

$$\eta = \frac{P^{(2\omega)}(L)}{P^{(\omega)}(0)} \propto \sin c^2 \left( \frac{1}{2} \Delta k L \right) \iff \Delta k = 0 \iff n_o^{2\omega} < n_e^\omega$$

Phase matching ==  
equal PHASE velocity

For nonhomogeneous media- Phc

P.M. (equal PHASE velocity),

For SH  $2K_1 = K_2$  ....  $n(w) = n(2 w)$  ---if there is *geometrical dispersion*

Thin films

No P.M. Interference between bound and free components ( surface and bulk)

At the nanoscale

No P.M.

Field localization and overlap



*J. Phys. Chem. B* 2009, 113, 13437–13445

13437

## Measurement of the Second-Order Hyperpolarizability of the Collagen Triple Helix and Determination of Its Physical Origin

Ariane Deniset-Besseau,<sup>†</sup> Julien Duboisset,<sup>‡</sup> Emmanuel Benichou,<sup>‡</sup> François Hache,<sup>†</sup> Pierre-François Brevet,<sup>‡</sup> and Marie-Claire Schanne-Klein<sup>\*‡</sup>

*Laboratoire d'Optique et Biosciences, Ecole Polytechnique, CNRS, INSERM U696, 91128 Palaiseau, France, and Laboratoire de Spectroscopie Ionique et Moléculaire, CNRS, Université Claude Bernard Lyon 1, 69622 Villeurbanne, France*

*Received: May 19, 2009; Revised Manuscript Received: July 16, 2009*

$$\chi_{ijk}^{(2)} = N \sum_{ijk} \langle (I \cdot i)(J \cdot j)(K \cdot k) \rangle \beta_{ijk}$$

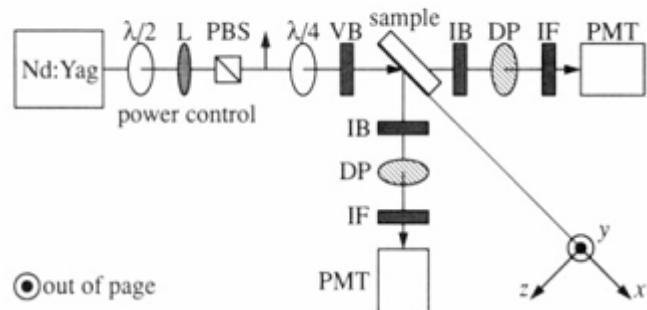
## Nonlinear Optical Activity and Biomolecular Chirality

Thierry Verbiest, Martti Kauranen, Andre Persoons, Marjo Ikonen, Jouni Kurkela, and Helge Lemmetyinen

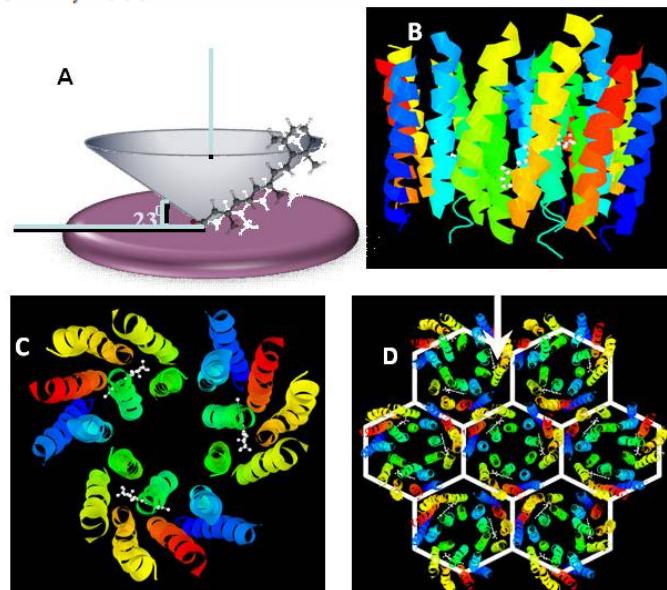
*J. Am. Chem. Soc.*, 1994, 116 (20), 9203-9205 • DOI: 10.1021/ja00099a040 • Publication Date (Web): 01 May 2002

Downloaded from <http://pubs.acs.org> on March 17, 2009

9204 *J. Am. Chem. Soc.*, Vol. 116, No. 20, 1994



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Biophysical Journal Volume 73 December 1997 3164–3170

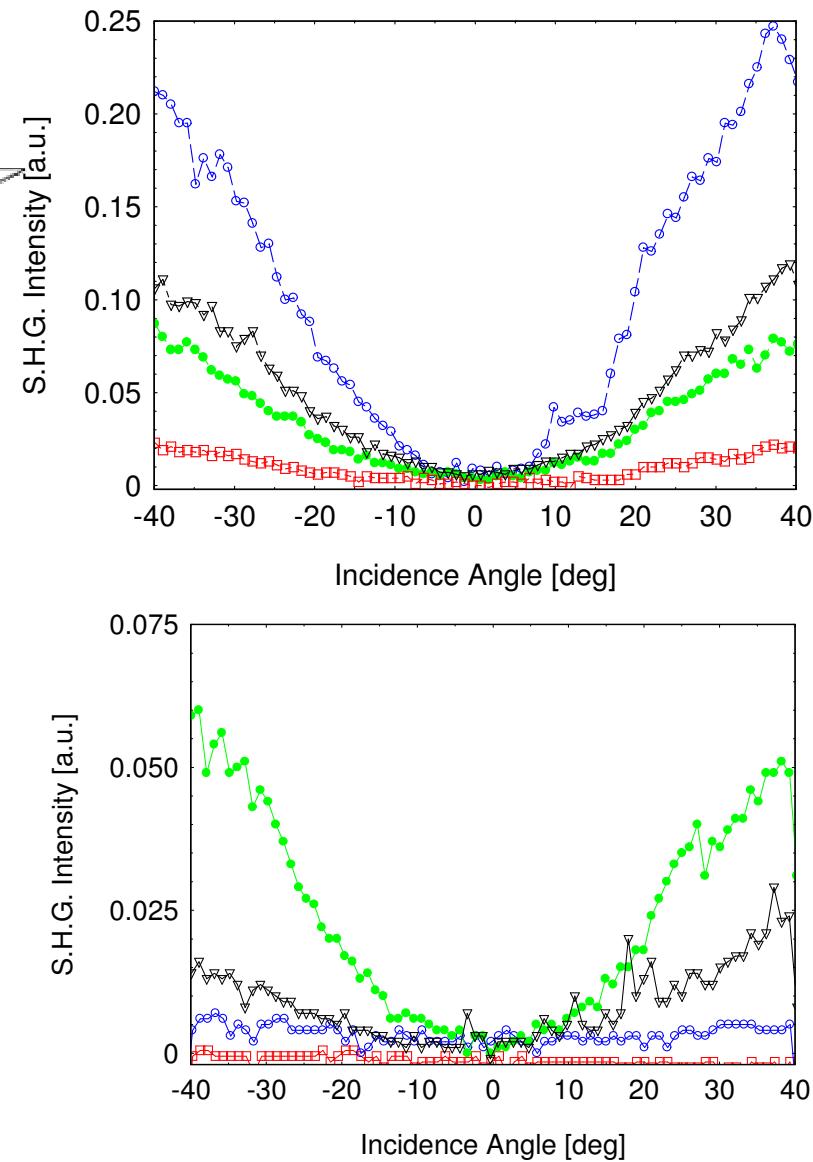
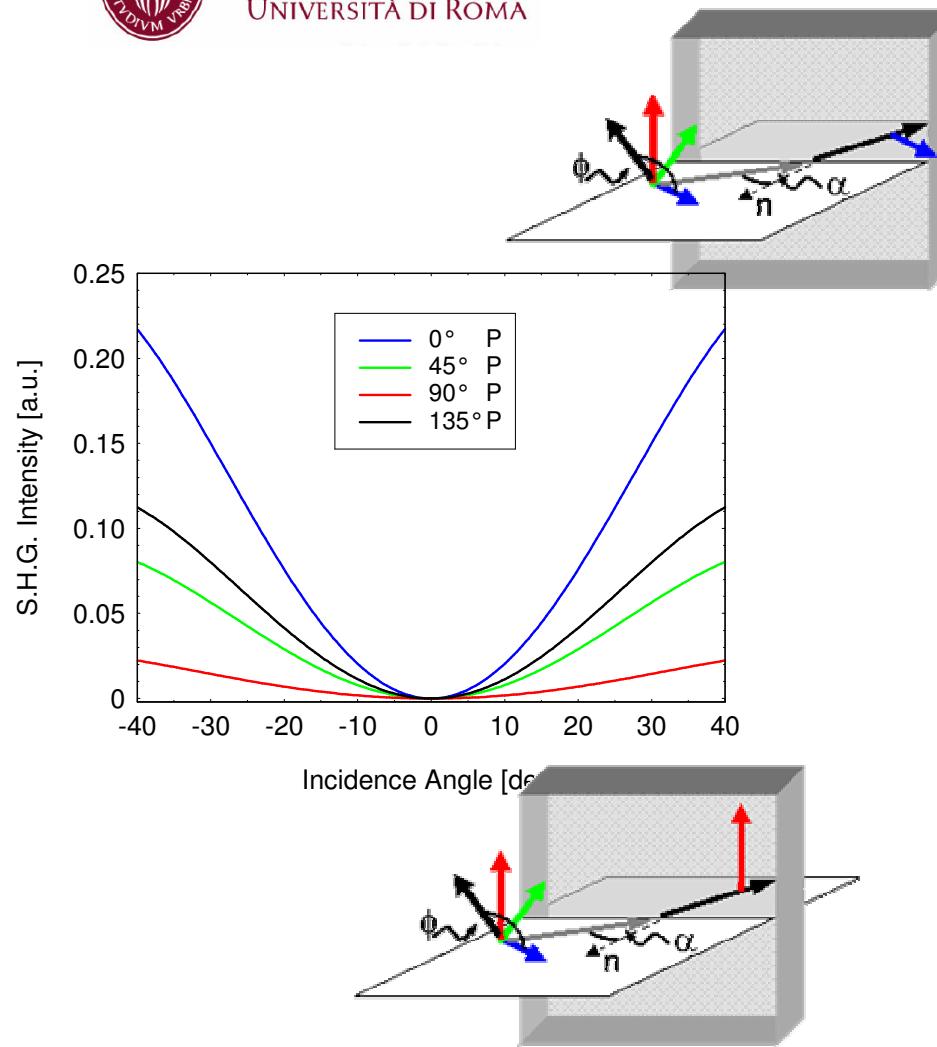
### Optical Rotation of the Second Harmonic Radiation from Retinal in Bacteriorhodopsin Monomers in Langmuir-Blodgett Film: Evidence for Nonplanar Retinal Structure

V. Volkov,\* Yu. P. Svirko,<sup>#§</sup> V. F. Kamalov,\* L. Song,\* and M. A. El-Sayed\*

\*Department of Chemistry and Biochemistry, Georgia Institute of Technology, Atlanta, Georgia 30332 USA; <sup>#</sup>Department of Physics, University of Southampton, Southampton SO17 1BJ, England; and <sup>§</sup>General Physics Institute, Moscow 117942, Russia



## Maker Fringes on BR

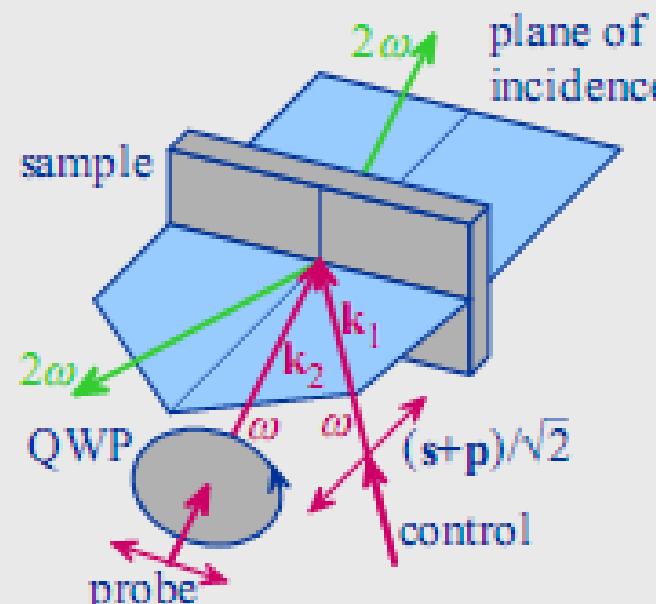
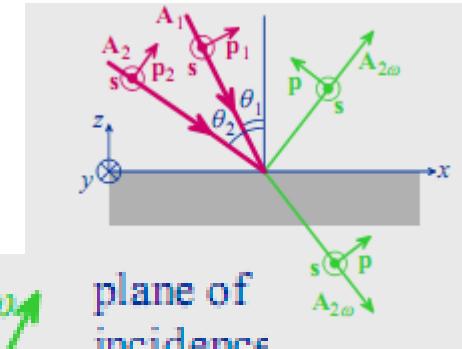


$d_{15} = d_{24} = 2.02 \text{ pm/V}$ ,  $d_{31} = d_{32} = 2.44 \text{ pm/V}$ ,  $d_{33} = 4.70 \text{ pm/V}$ , as well as the chiral coefficients  $d_{25} = -d_{14} = 0.34 \text{ pm/V}$ .



## Two beams

- **Control beam**
  - polarization fixed
- **Probe beam**
  - polarization varied
- **SHG signals**
  - reflected and transmitted





Transmission  
Coefficient at  $\omega$  for pump 1 & 2

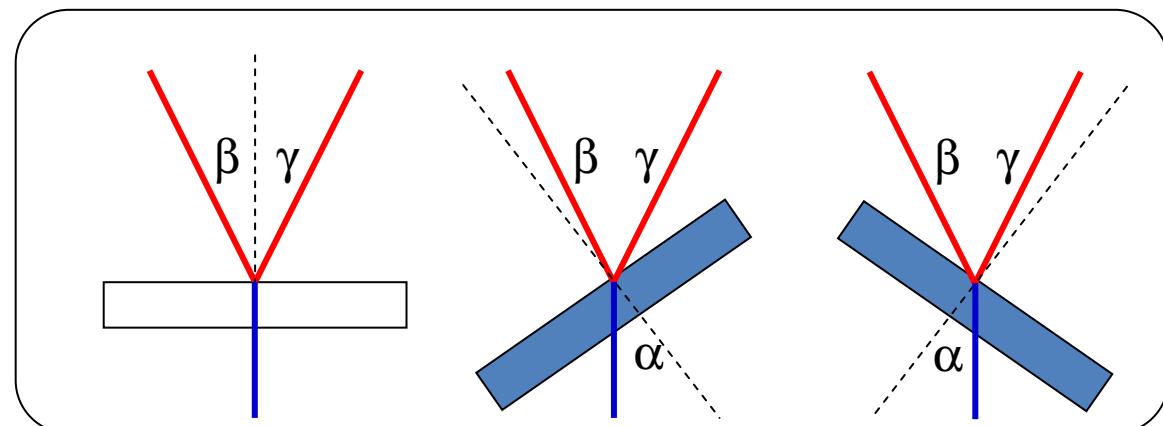
$$W_{\omega_1+\omega_2}(\alpha) = \left( \frac{512\pi^3}{A} \right) (t_{\omega_1})^2 \cdot (t_{\omega_2})^2 \cdot T_{\omega_1+\omega_2} \cdot W_{\omega_1} \cdot W_{\omega_2} \frac{\sin^2(\Psi_{SHG}(\alpha))}{[n_{\omega_1} \cdot n_{\omega_2} - n_{\omega_1+\omega_2}^2]^2} (d_{eff}(\alpha))^2$$

Power of the field at  $\omega$  for pump  
1 & 2

Transmittance  
at  $2\omega$

$$\Psi_{SHG} = \left( \frac{\pi L}{2} \right) \left( \frac{2}{\lambda} \right) [n_{\omega_1} \cdot \cos(\alpha' \omega_1) + n_{\omega_2} \cdot \cos(\alpha' \omega_2) - 2n_{\omega_1+\omega_2} \cdot \cos(\alpha' \omega_1 + \omega_2)]$$

- ✓ The difference between the two incidence angles is set to a fixed value ( $18^\circ$ )

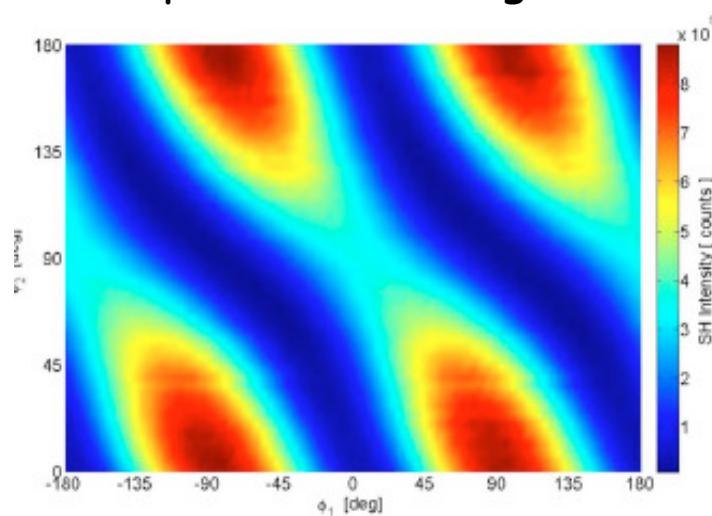




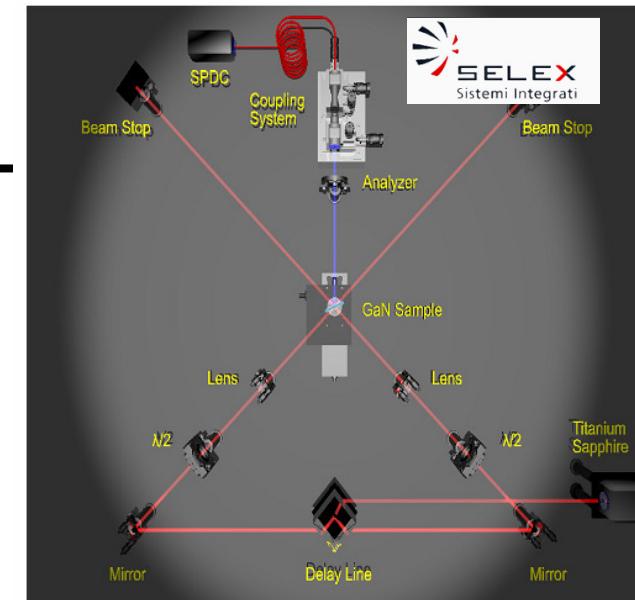
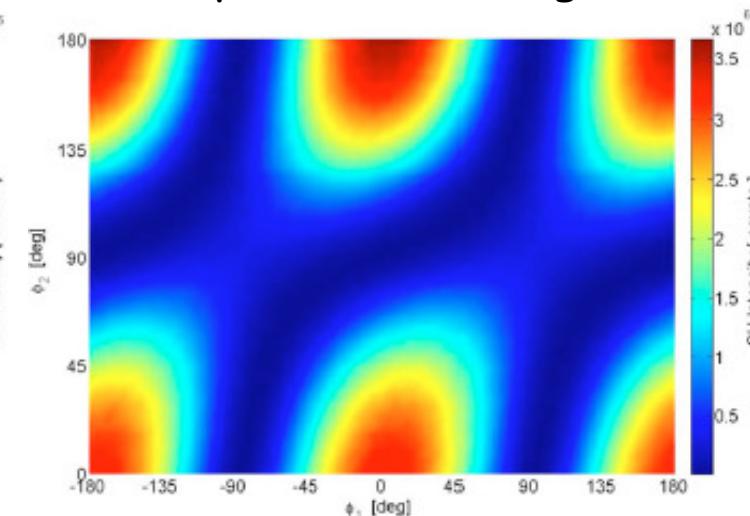
## Polarization chart

- At a fixed incidence angle, the generated signal is investigated by varying the polarization state of both fundamental beams.
- Example: GaN slab, ~300nm thick

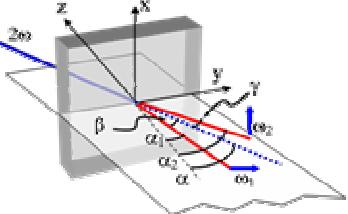
S-polarized SH signal



P-polarized SH signal

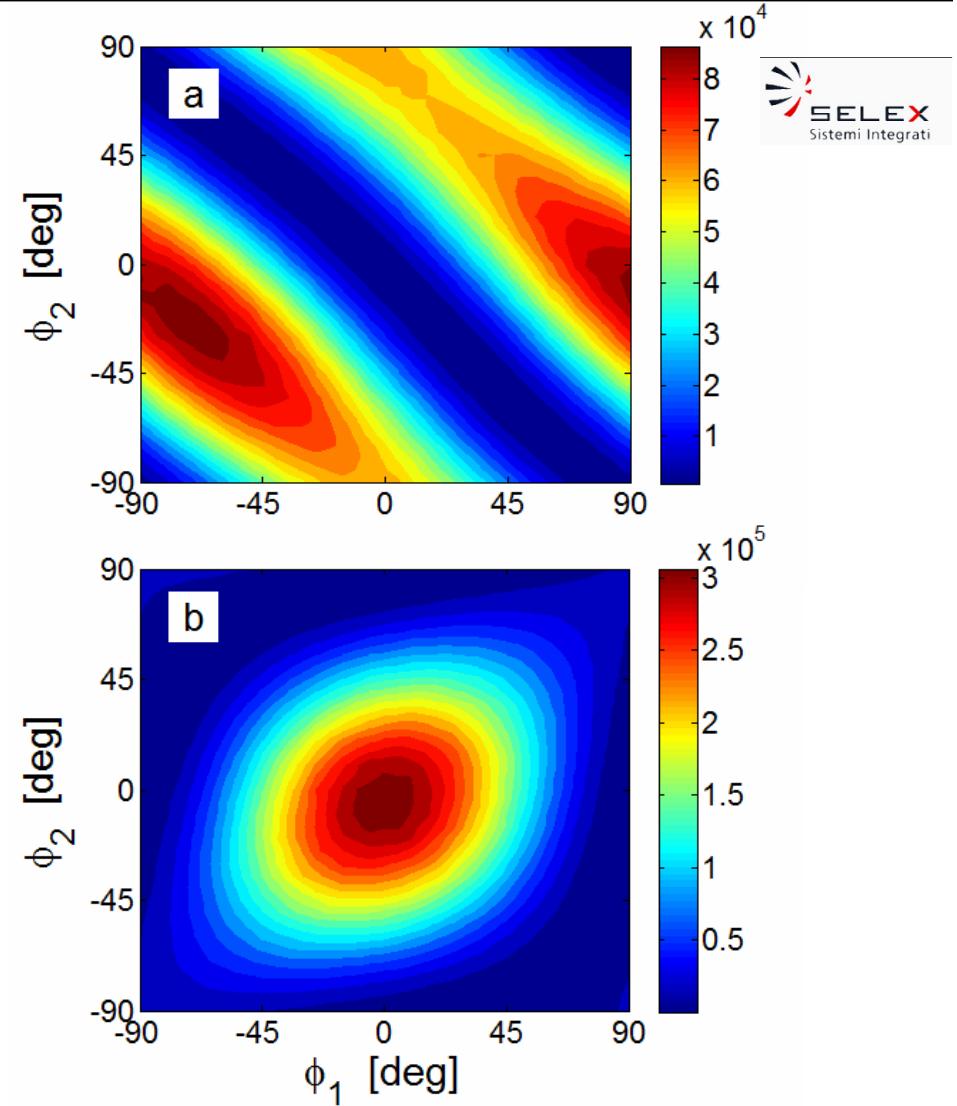
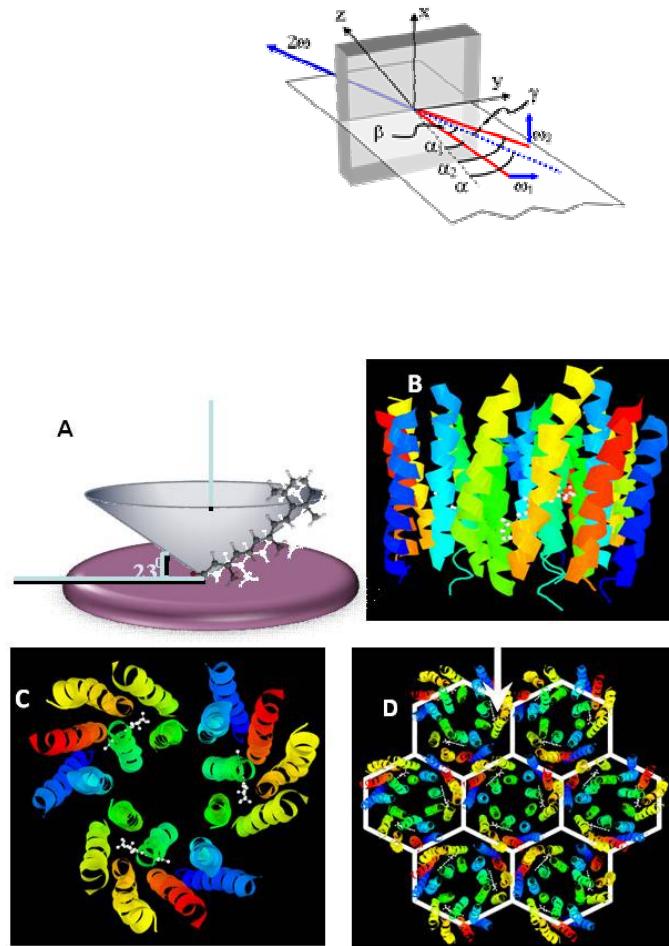


F.A.Bovino  
Selex-SI-  
Genova



M.C.Larciprete, F.A.Bovino, M.Giardina, A.Belardini, M.Centini, C.Sibilia, M.Bertolotti, A. Passaseo and V.Tasco, "Mapping the nonlinear optical susceptibility by noncollinear second-harmonic generation", Optics Letters, Vol. 34, pp.2189-2191 (2009)

# Polarization Chart in noncollinear SHG-BR



F.A Bovino et al- Selex-SI-Italy



# Power of SH Signal

$$W_{\omega_1+\omega_2}(\alpha) = \left( \frac{512\pi^3}{A} \right) (t_{\omega_1})^2 \cdot (t_{\omega_2})^2 \cdot T_{\omega_1+\omega_2} \cdot W_{\omega_1} \cdot W_{\omega_2} \frac{\sin^2(\Psi_{SHG}(\alpha))}{[n_{\omega_1} \cdot n_{\omega_2} - n_{\omega_1+\omega_2}^2]^2} (|d_{eff}^{(2)}(\alpha)|)^2$$

$$\begin{aligned} d_{eff}^{(2)} = & (\hat{p}_{out})(d_{il}^{(eee)})(\hat{p}_{in}^{1e})(\hat{p}_{in}^{2e}) + \\ & (\hat{p}_{out})(d_{il}^{(eem1)})(\hat{p}_{in}^{1e})(\hat{p}_{in}^{2m}) + (\hat{p}_{out}^e)(d_{il}^{(eem2)})(\hat{p}_{in}^{1m})(\hat{p}_{in}^{2e}) + \\ & (\hat{p}_{out})xK(d_{il}^{(mee)})(\hat{p}_{in}^{1e})(\hat{p}_{in}^{2e}) \end{aligned}$$

$\hat{p}_{out}$  : polarization versor of output sh signal

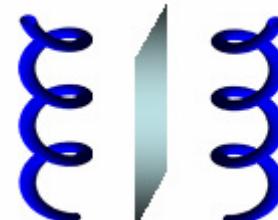
$\hat{p}_{in}^1$  : polarization versor of input pump 1 signal

$\hat{p}_{in}^2$  : polarization versor of input pump 2 signal

3p

$$d_{il} = \frac{\hat{\chi}_{ijk}^{(2)}}{2}$$

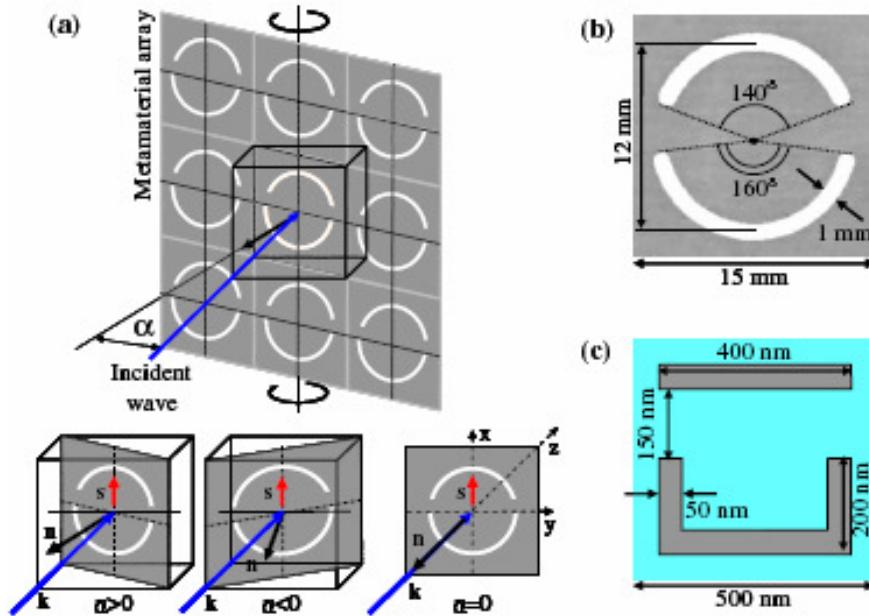
Intrinsic 3D chirality



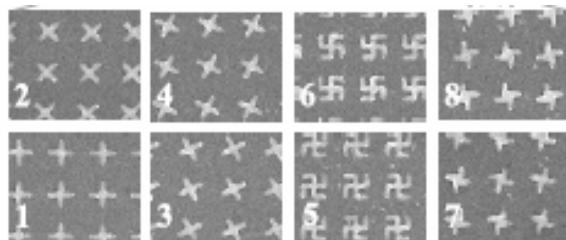
$$\begin{pmatrix} 0 & 0 & 0 & d_{14}^{eee} & d_{15}^{eee} & 0 \\ 0 & 0 & 0 & d_{24}^{eee} & d_{25}^{eee} & 0 \\ d_{31}^{eee} & d_{32}^{eee} & d_{33}^{eee} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & d_{14}^{eem} & d_{15}^{eem} & 0 \\ 0 & 0 & 0 & d_{24}^{eem} & d_{25}^{eem} & 0 \\ d_{31}^{eem} & d_{32}^{eem} & d_{33}^{eem} & 0 & 0 & d_{36}^{eem} \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & d_{14}^{ree} & d_{15}^{ree} & 0 \\ 0 & 0 & 0 & d_{24}^{ree} & d_{25}^{ree} & 0 \\ d_{31}^{ree} & d_{32}^{ree} & d_{33}^{ree} & 0 & 0 & 0 \end{pmatrix}$$

# Extrinsic Chirality

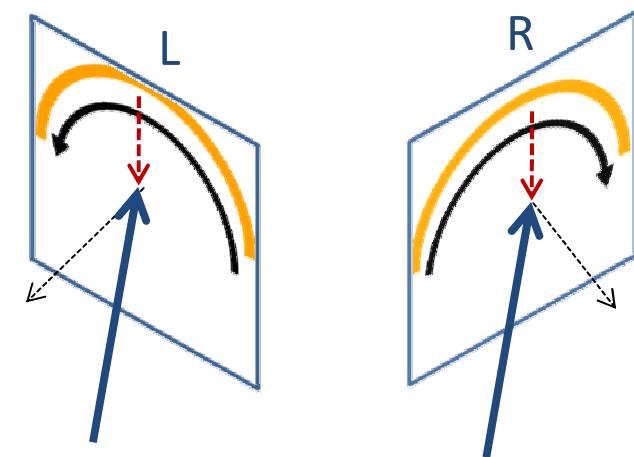
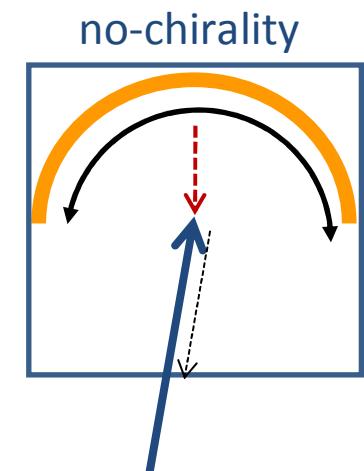
E. Plum, PRL 102,113902 (2009)



S.N. Volkov, PRA 79, 043819 (2009)



**Extrinsic chirality:**  
A non-chiral  
element breaks the  
symmetry when is  
tilted with respect  
the incidence light.  
The element thus  
assumes a chiral  
behavior



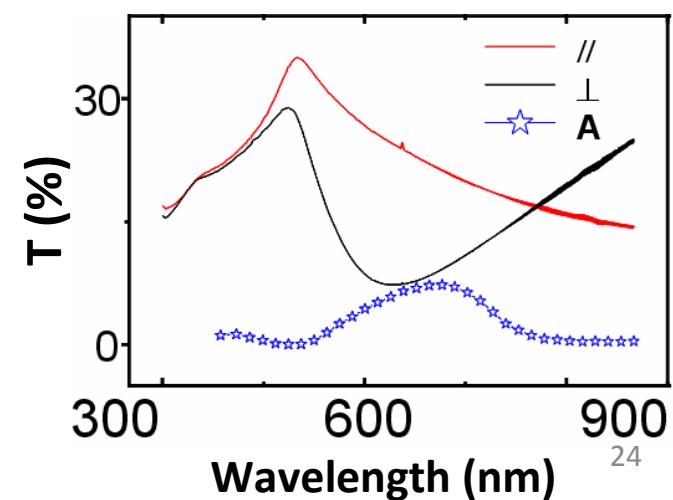
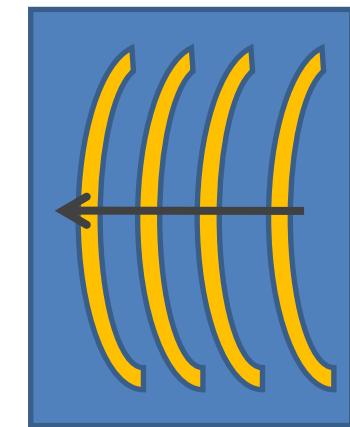
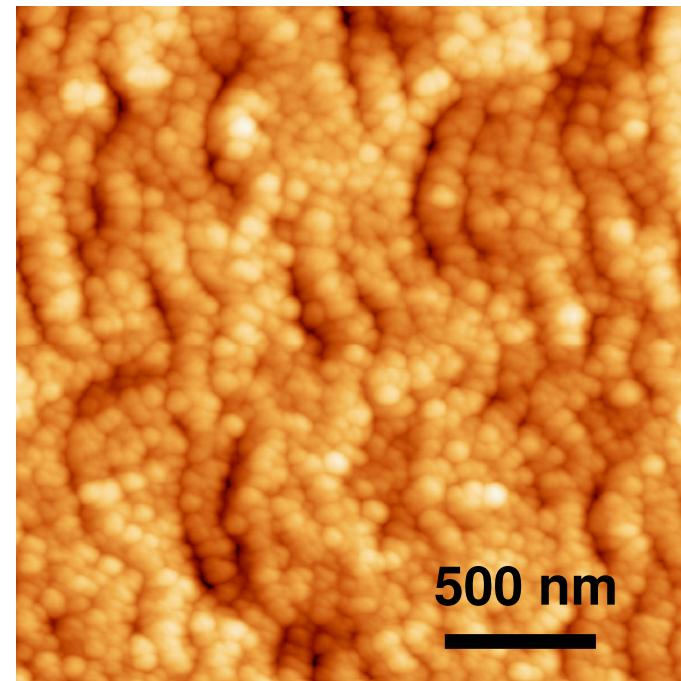
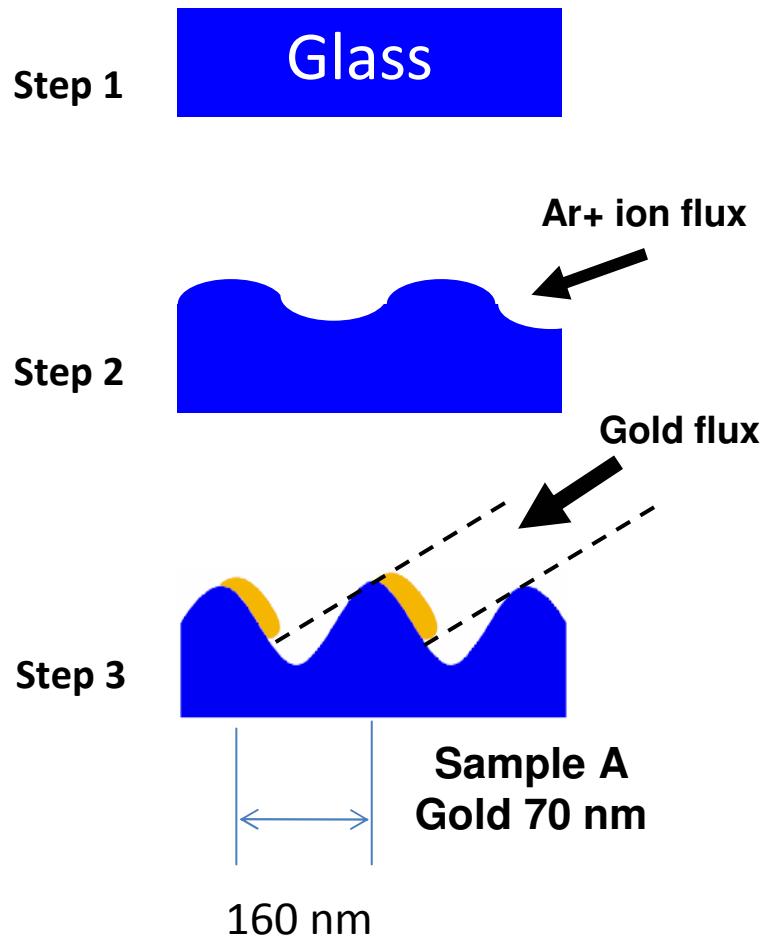


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## Extrinsic Chirality

# Test sample fabrication – Sample A

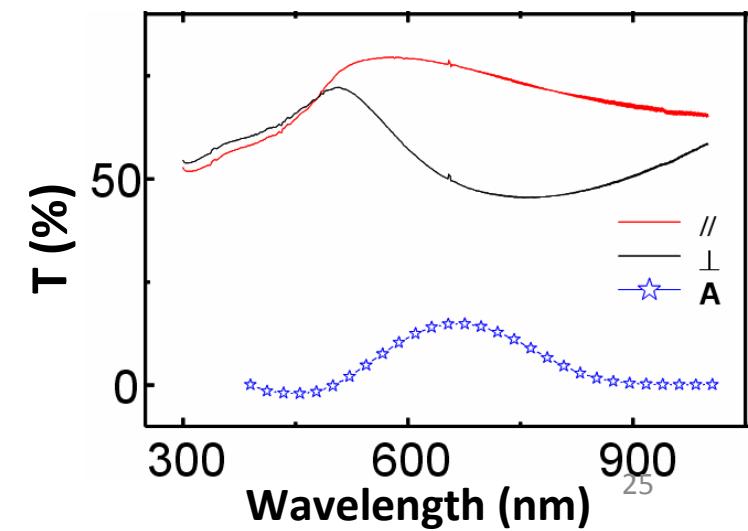
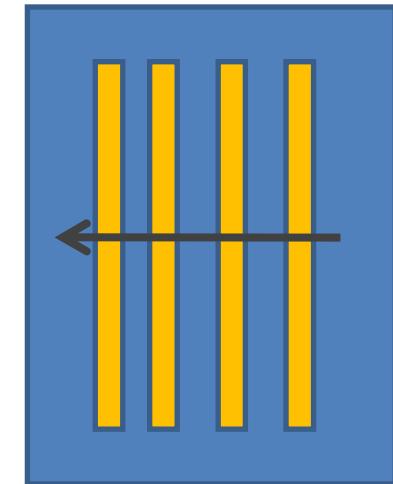
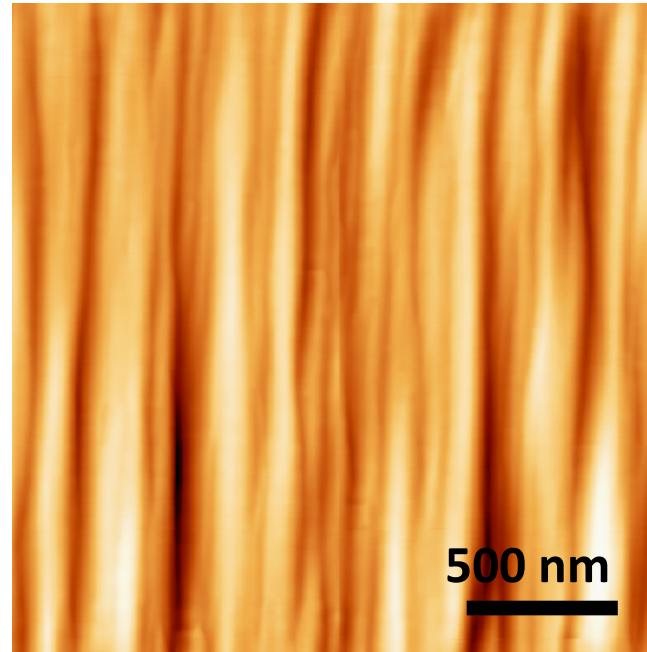
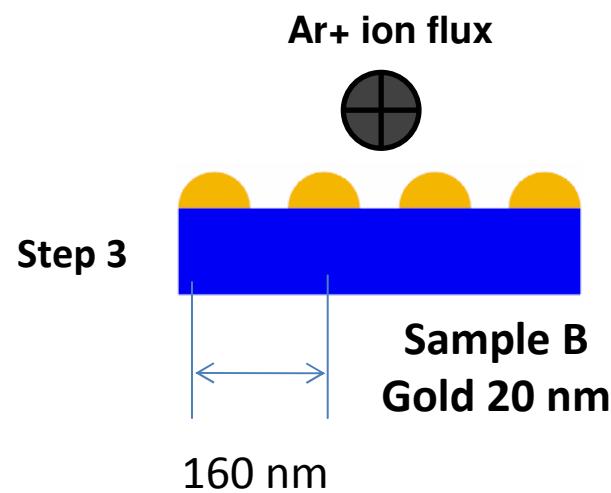
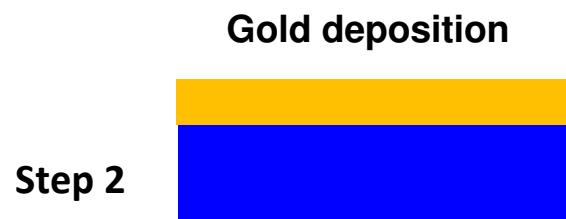
Low cost self assembling procedure



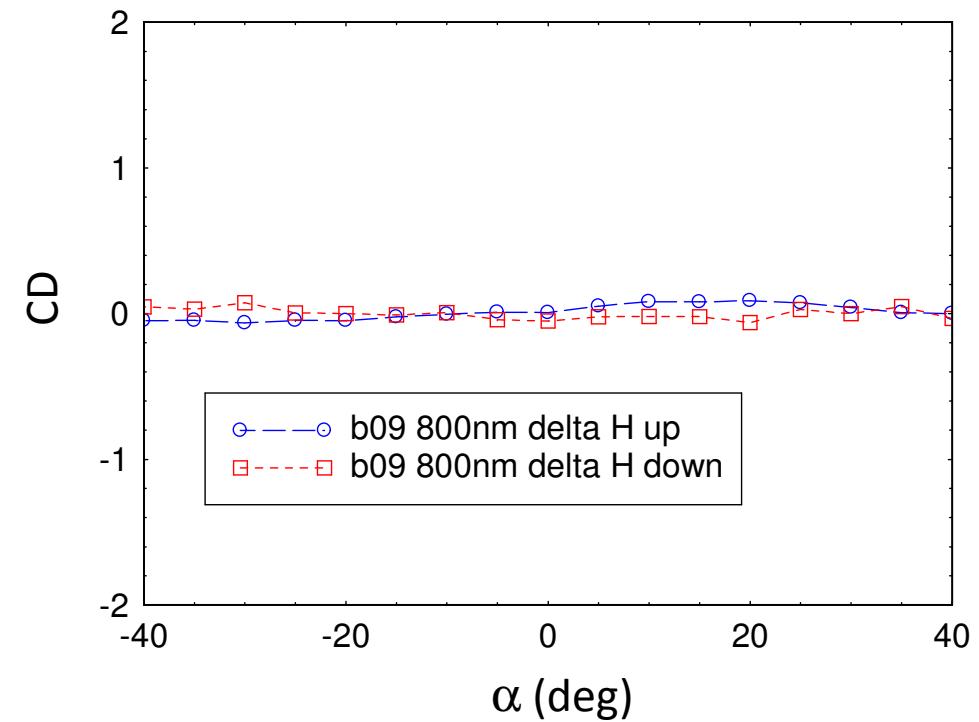
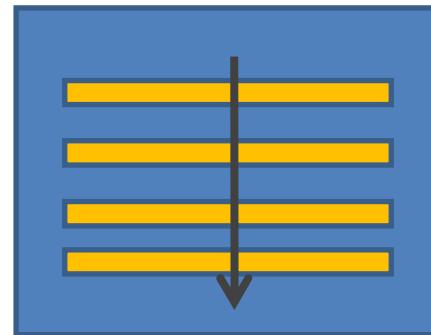
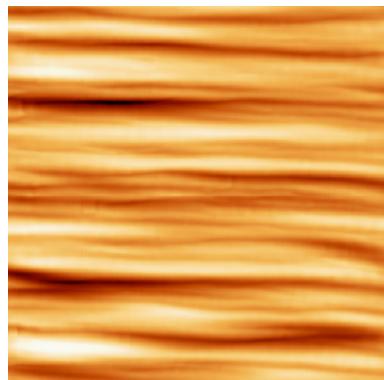
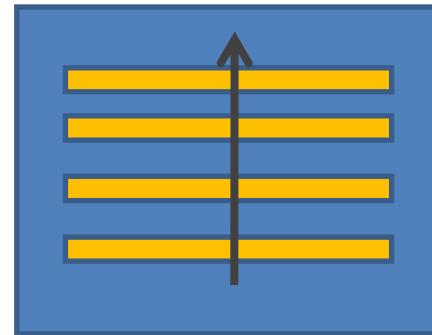
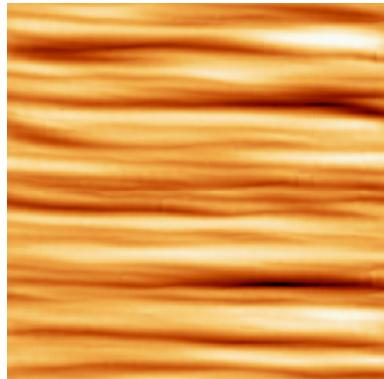
# Reference sample fabrication – Sample B



Low cost self assembling procedure

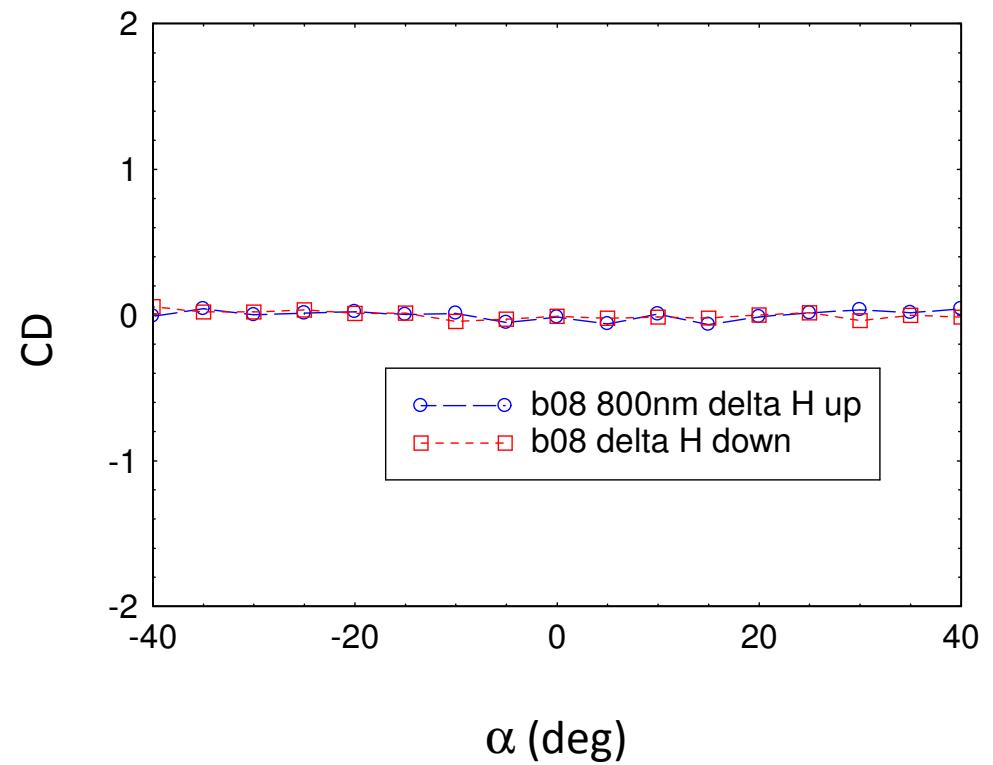
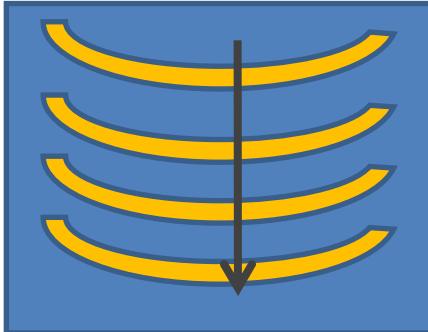
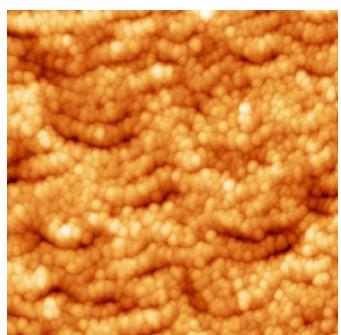
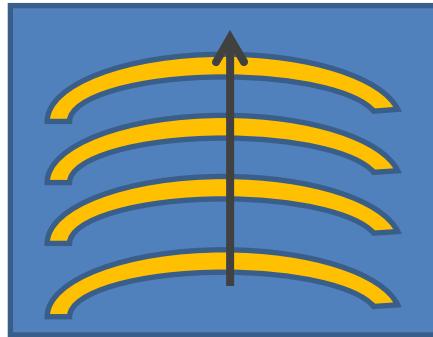
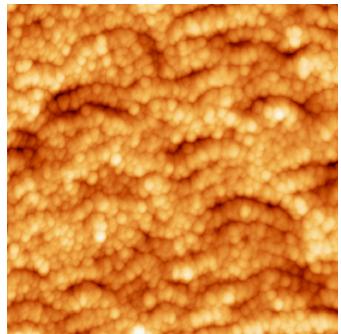


# Linear measurements on sample B - P pol





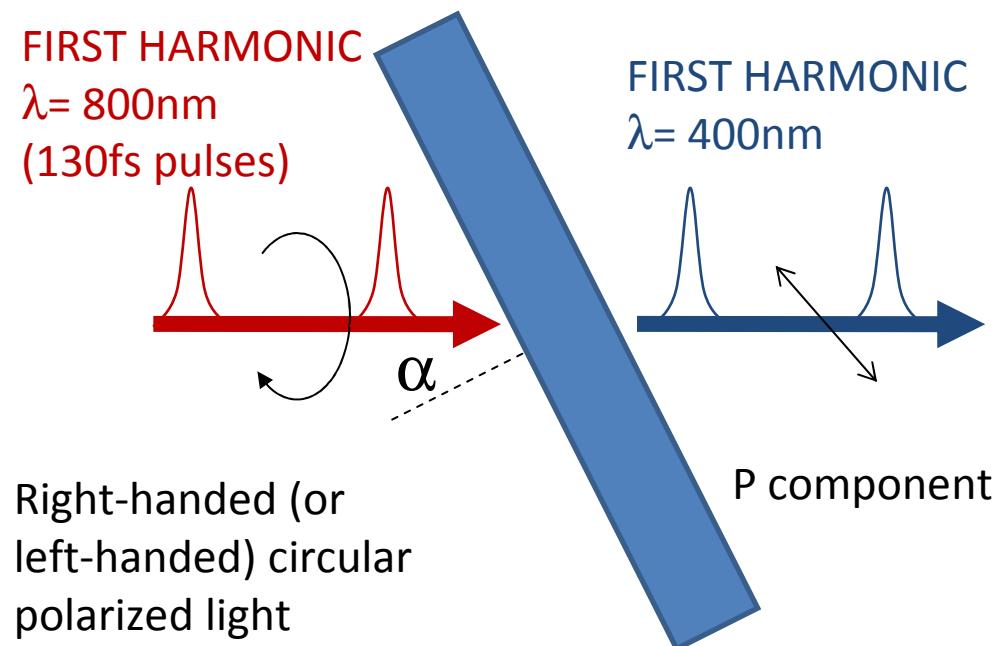
# Linear measurements on sample A - P pol





# SHG measurements – P pol

The measurements were performed as a function of the incidence angle  $\alpha$

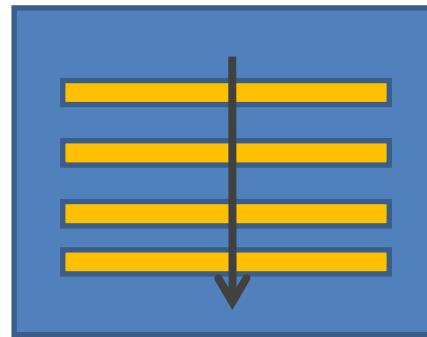
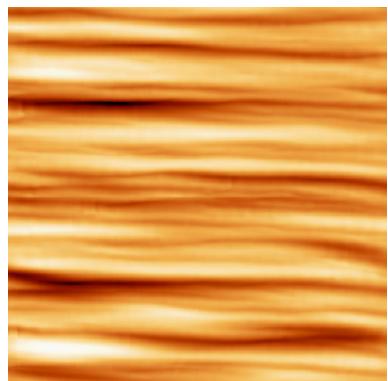
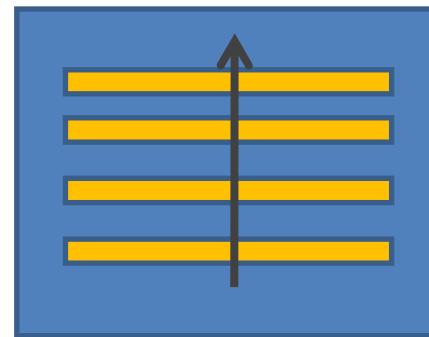
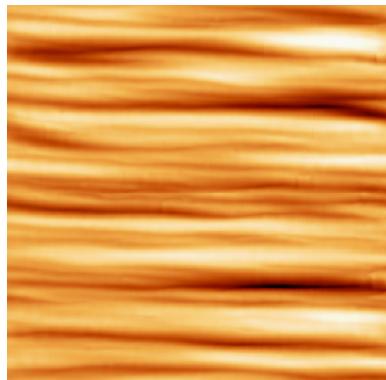


Second harmonic  
generation - circular  
dichroism (SHG-CD)

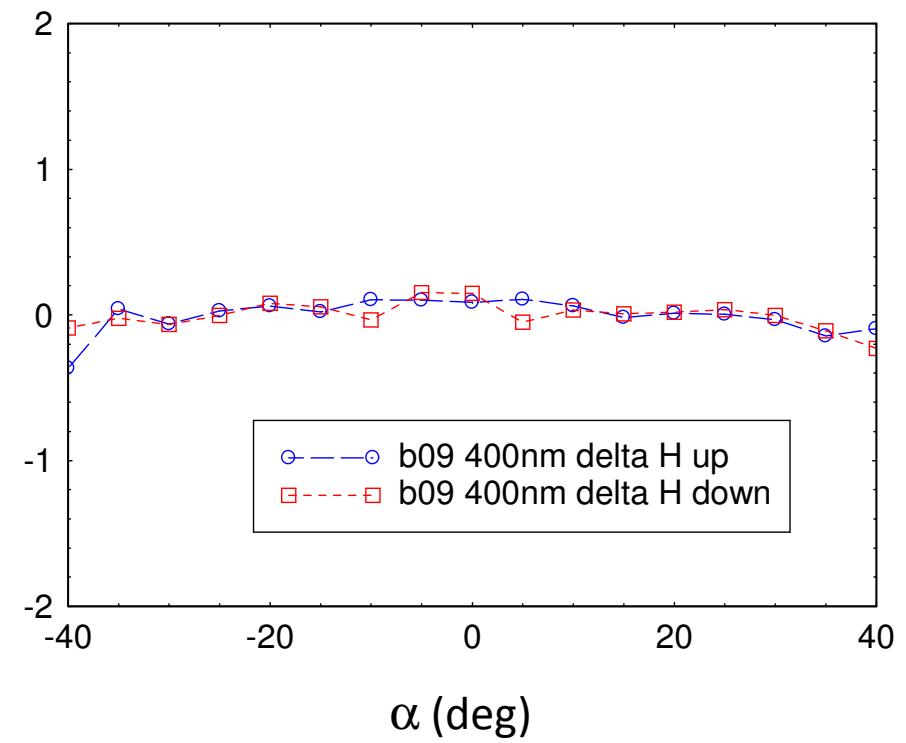
$$SHG - CD = \frac{I_L^{2\omega} - I_R^{2\omega}}{(I_L^{2\omega} + I_R^{2\omega})/2}$$



# SHG measurements on sample B – P pol

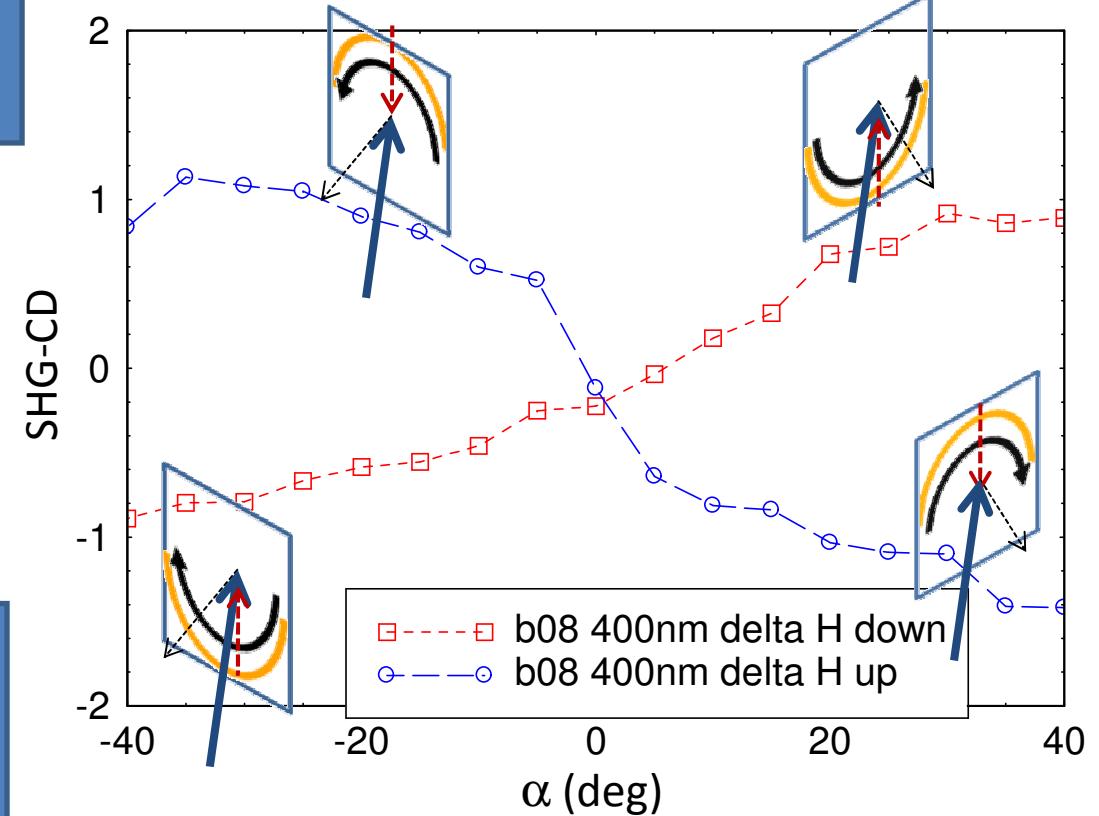
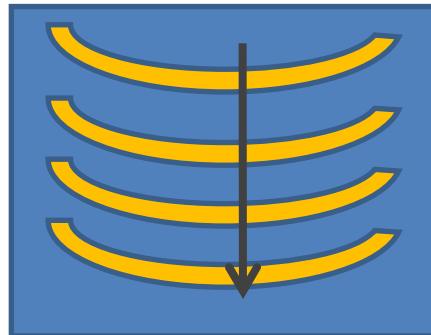
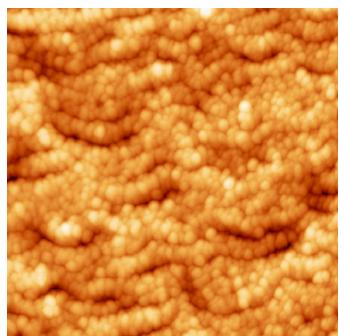
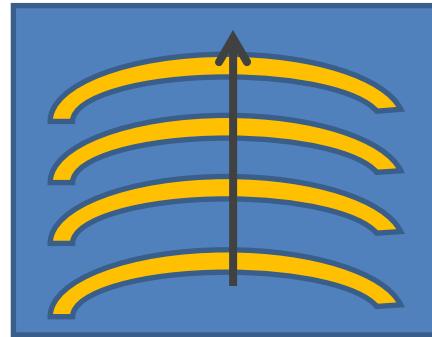
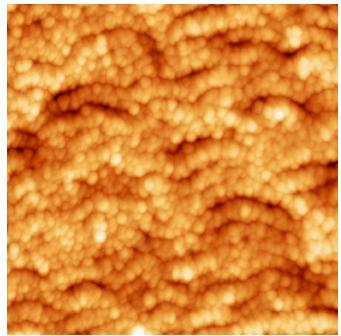


SHG-CD





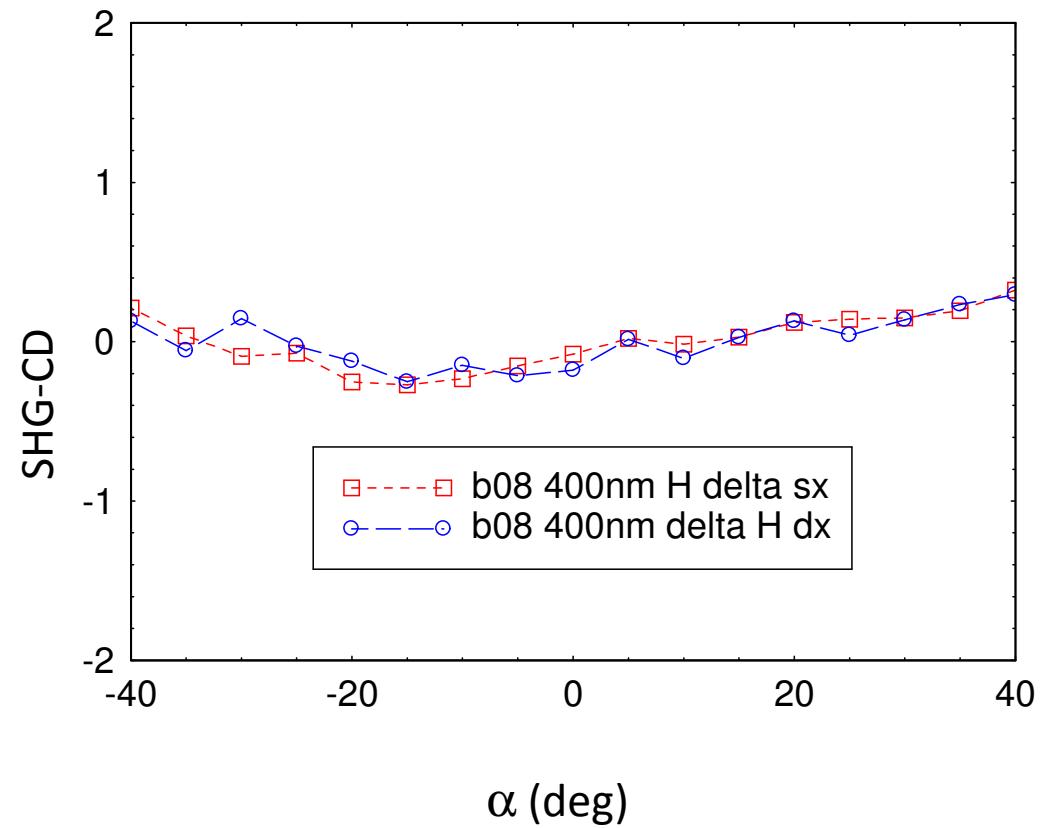
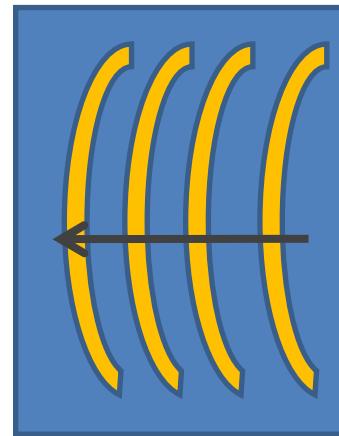
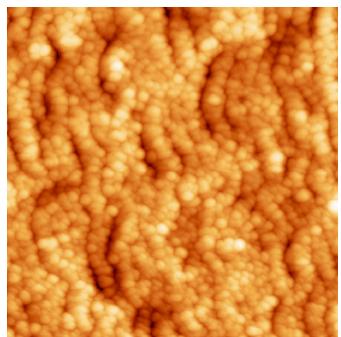
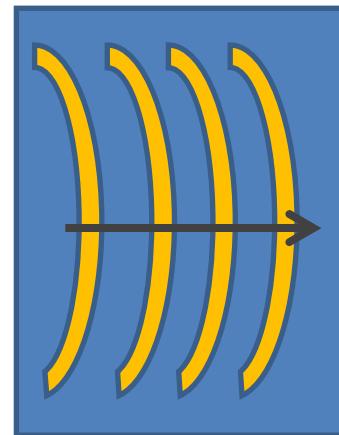
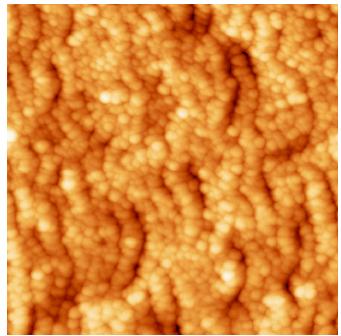
# SHG measurements on sample A – P pol



A.Belardini et al PRL  
2011

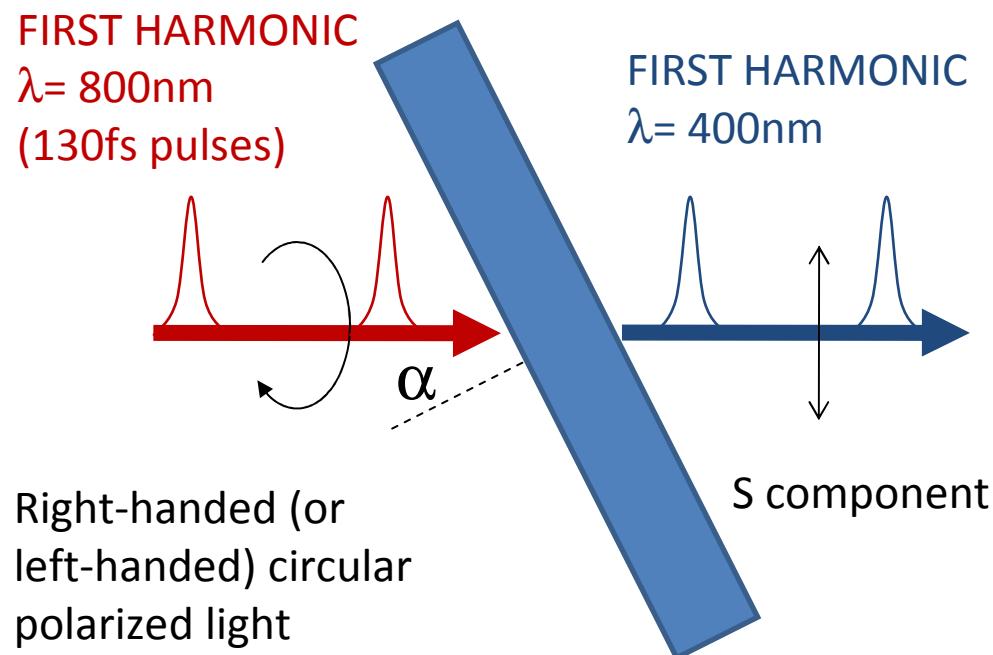


# SHG measurements on sample A – P pol





# SHG measurements – S pol

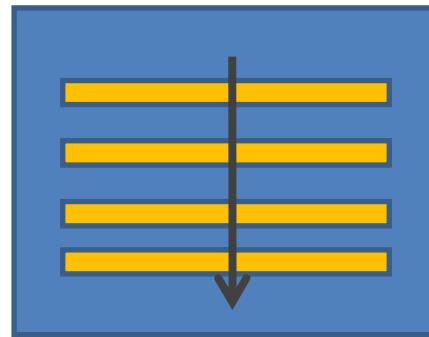
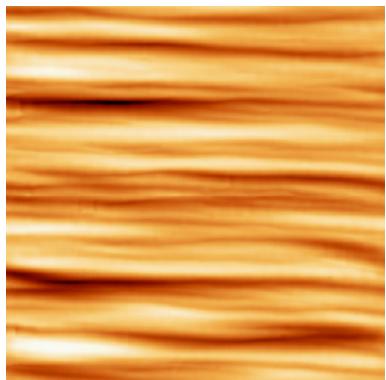
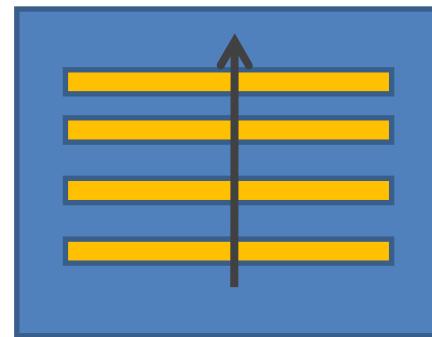
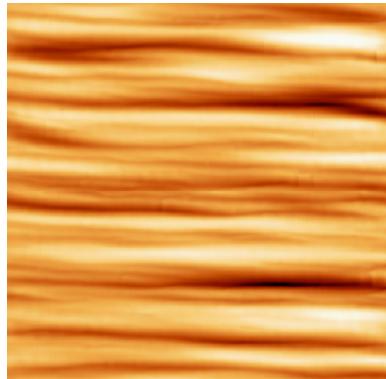


Second harmonic generation - circular dichroism (SHG-CD)

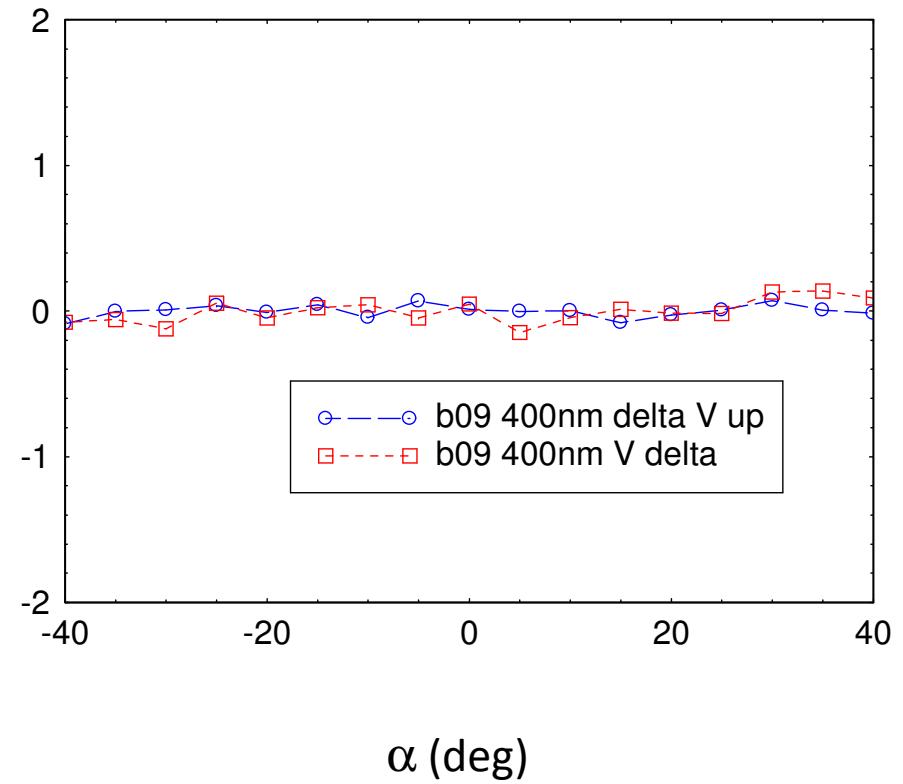
$$SHG - CD = \frac{I_L^{2\omega} - I_R^{2\omega}}{(I_L^{2\omega} + I_R^{2\omega})/2}$$



# SHG measurements on sample B – S pol

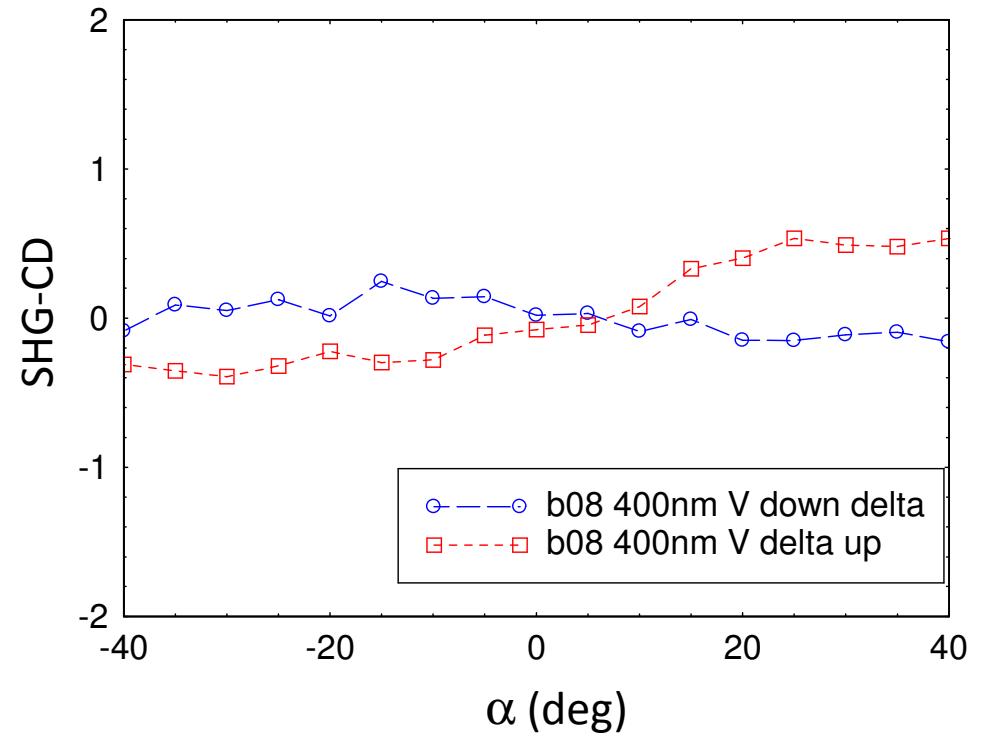
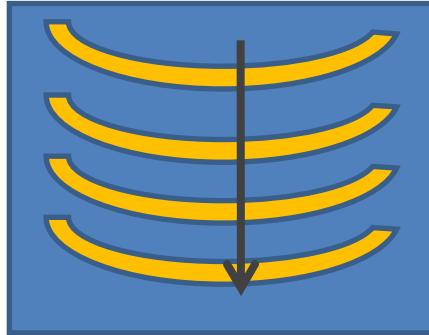
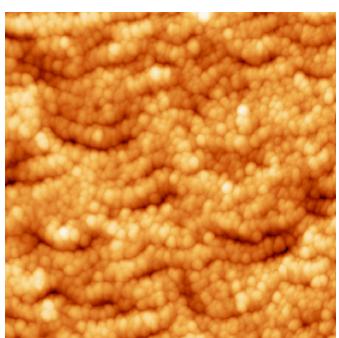
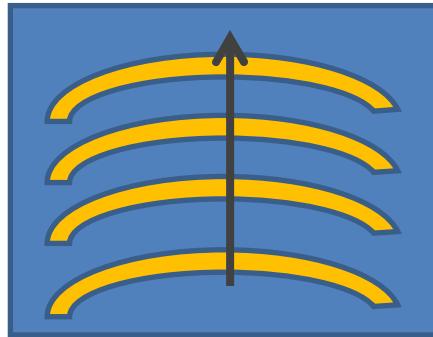
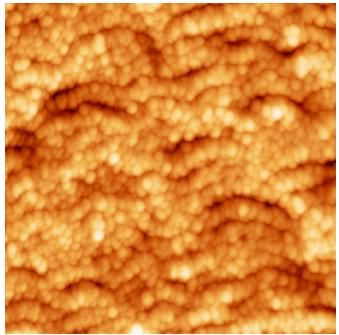


SHG-CD





# SHG measurements on sample A – S pol

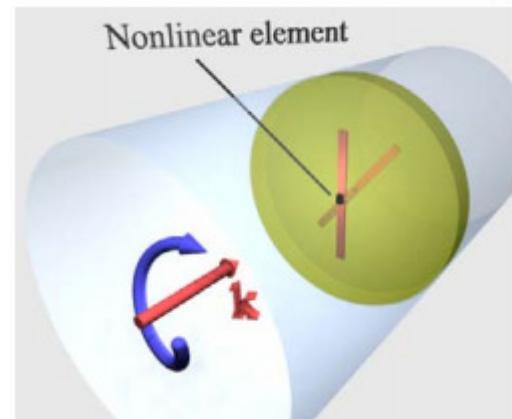
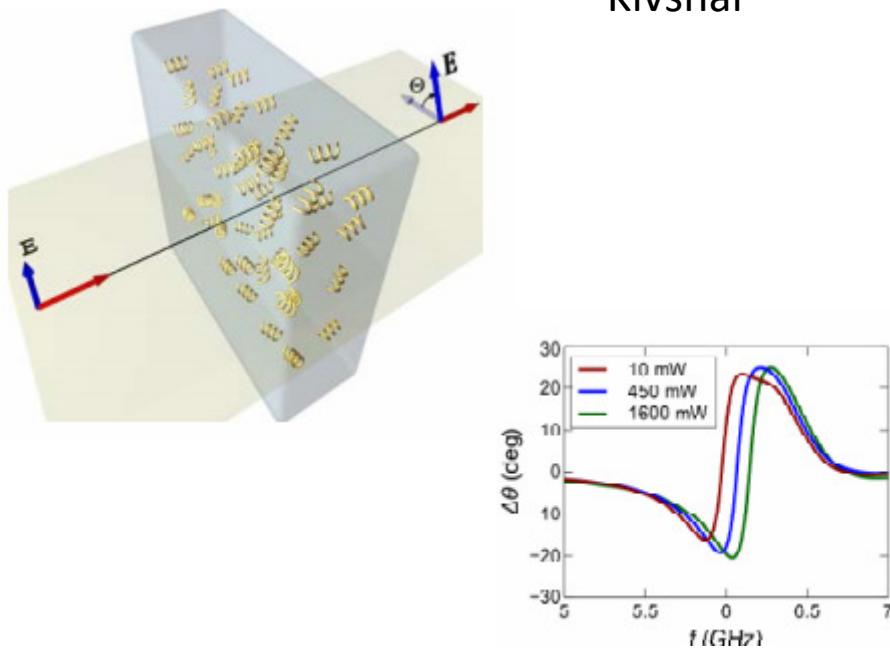




Chiral Metamaterials  
Unlocking Nonlinear Optical Activity  
Intensity dependent effects

AOS News Volume 25 Number 1 2011

by David A. Powell, Ilya V. Shadrivov, Vassili A. Fedotov, Nikolay I. Zheludev and Yuri S. Kivshar



Optical diode for  
circular polarized  
light



## Optical chirality without optical activity: How surface plasmons give a twist to light

Aurélien Drezet, Cyriaque Genet, Jean-Yves Laluet and Thomas  
W. Ebbesen

ISIS, Louis Pasteur University, 8 allée Gaspard Monge, 67000, Strasbourg, France

#97343 - \$15.00 USD      Received 12 Jun 2008; revised 16 Jul 2008; accepted 16 Jul 2008; published 5 Aug 2008  
(C) 2008 OSA      18 August 2008 / Vol. 16, No. 17 / OPTICS EXPRESS 12559

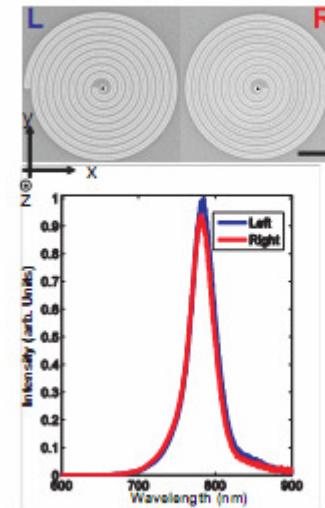


Fig. 1. Chiral plasmonic metasubstrates. On the top panel: scanning electron micrographs of the left (L) and right (R) handed enantiomer (mirror image) planar chiral structures investigated. The scale bar is 3  $\mu\text{m}$  long. The parameters characterizing the structures are the following: hole diameter  $d = 500$  nm, film thickness  $\lambda = 100$  nm, relative index  $n = 1.93$ .



- 3D, 2D-chirality, Extrinsic chirality
- SHG as chirality detection method
- Nonlinear measurements
- Intensity dependence



## Acknwoledgments

SBAI

A.Belardini  
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A.Passaseo CNR- Lecce –Italy  
C.Gergeley – Univ. Montpellier2 – Montpellier-Fr  
F.Buatier- University of Genova -Italy



Thank you !!!



**SAPIENZA**  
UNIVERSITÀ DI ROMA

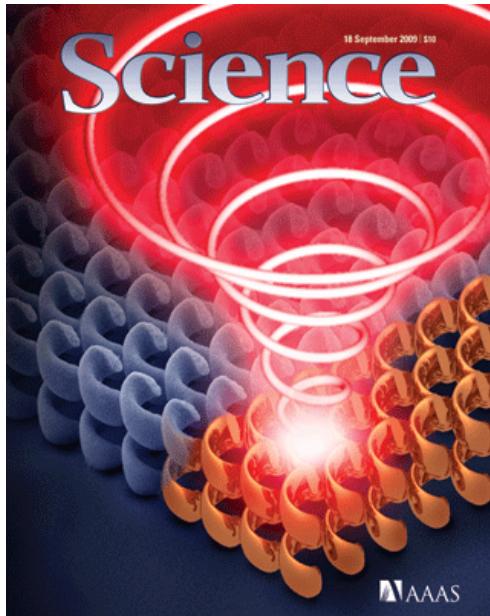
# Nonlinear Optics of chiral materials

**Concita Sibilia**

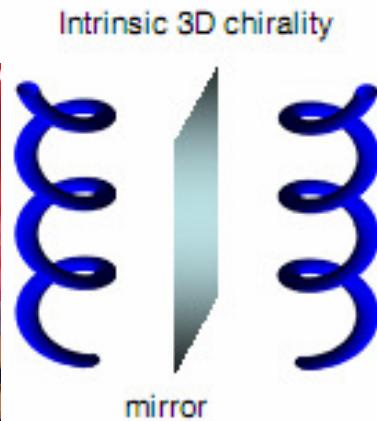
Dipartimento SBAI-Università' di Roma La Sapienza



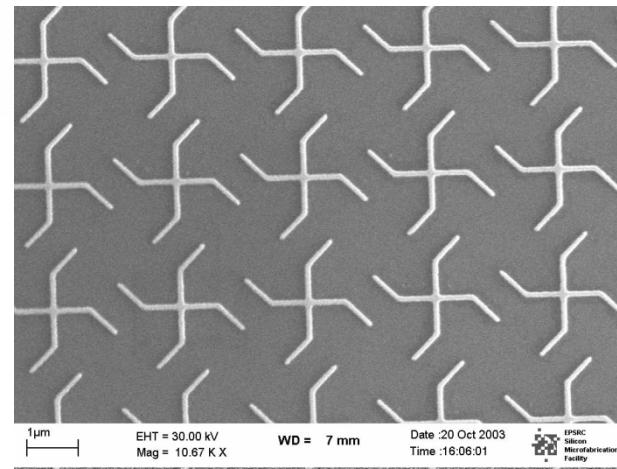
# 3D and 2D Chirality



3D chiral metamaterial



Absence of  
mirror rotation  
axis



2D chiral metamaterial

## Extrinsic electromagnetic chirality in metamaterials

E Plum, V A Fedotov and N I Zheludev<sup>1</sup>

Optoelectronics Research Centre, University of Southampton, Southampton SO17 1BJ, UK

E-mail: [niz@orc.soton.ac.uk](mailto:niz@orc.soton.ac.uk)

In the last years a lot of effort was devoted to the study on chiral metamaterials. Both 3D- and 2D-chiral metamaterials were studied.

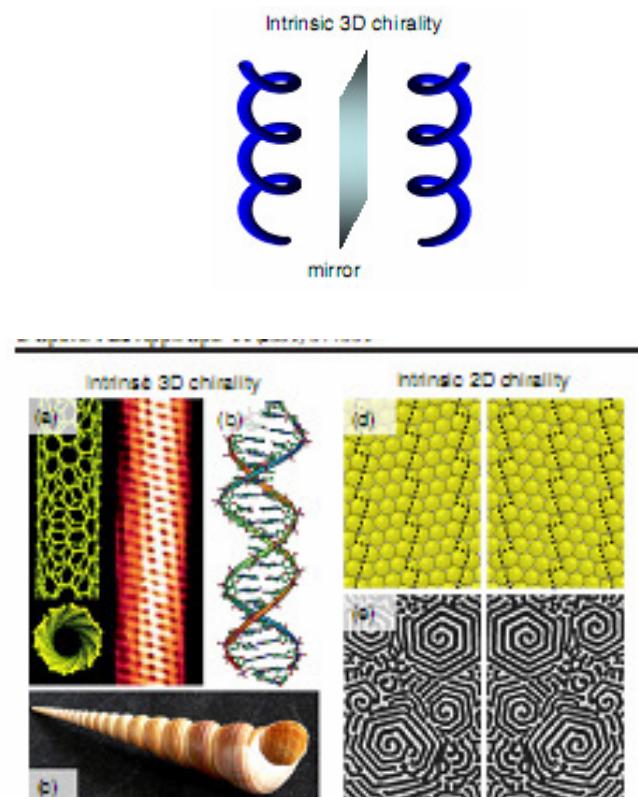


Figure 2. Examples of *intrinsically* chiral artificial and natural objects. 3D chirality: (a) single-wall carbon nanotube (structure and tunnelling electron microscope image) [1], (b) DNA double helix [2] and (c) seashell from the Turritellidae family [3]. 2D chirality: (d) surface (643) of a gold monocrystal, (e) surface of a metal–metal anhydride [4].

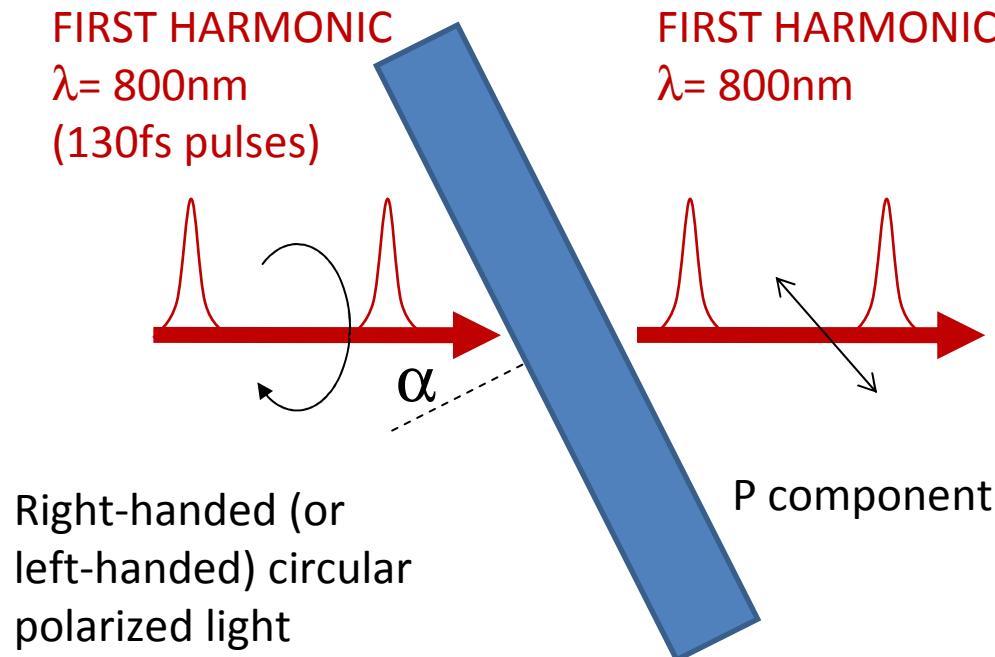
## From N.Zheludev

The term ‘chirality’ (or handedness from the Greek  $\chi\epsilon\rho$  ‘hand’) usually refers to an intrinsic sense of twist (or helicity) of a three-dimensional object, such as a helix, which has two different mirror forms called enantiomers (see figure 1(a)). 3D chirality appears to be a signature of life forms and can be found in many important organic molecules and proteins, as well as in inorganic structures (see figures 2(a)–(c)). 3D chirality results in such fundamental polarization effects as optical activity (a medium’s ability to rotate the polarization state of light) and the associated phenomenon of circular dichroism, which manifests itself as a difference in transmission levels for left and right circularly polarized light (see figure 3(a)). Since its discovery in 1811 by Dominique Arago, optical activity has played an important role in analytical chemistry, crystallography, molecular biology and the food industry, and is now used as a test for detecting life forms in space missions. The recognition of chirality as a source of negative refraction of light [5–11] that is needed for the creation of a perfect lens [12] has recently inspired intense work in developing optically active microwave and photonic chiral metamaterials [13–18] and has led to the observation of



# Linear measurements

The measurements performed as a function of the incidence angle  $\alpha$



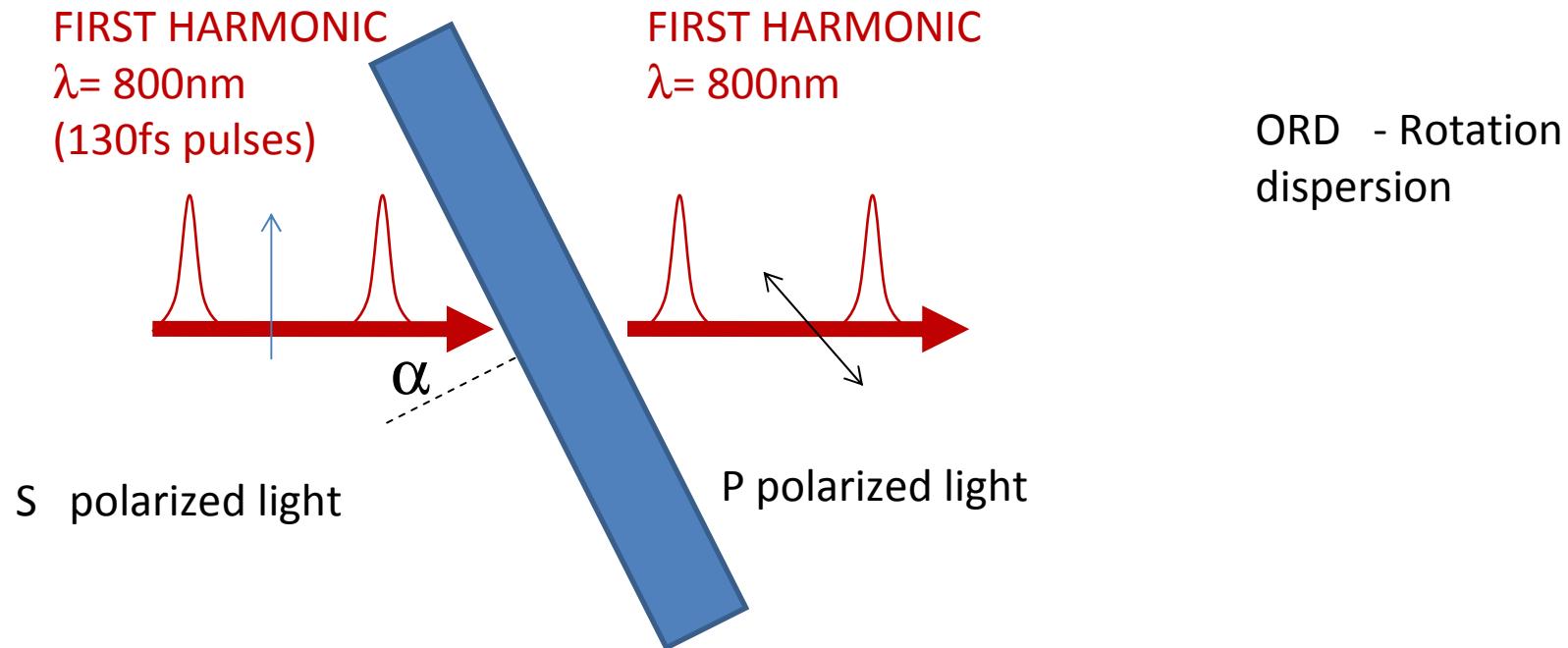
circular dichroism (CD)

$$CD = \frac{I_L^\omega - I_R^\omega}{(I_L^\omega + I_R^\omega)/2}$$



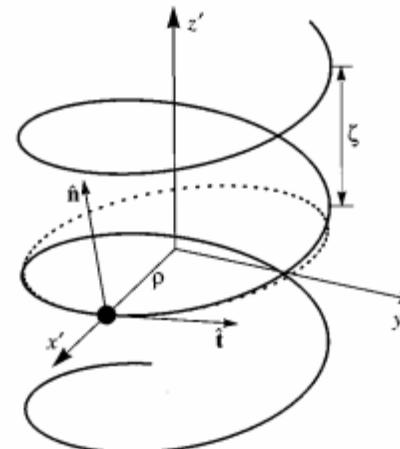
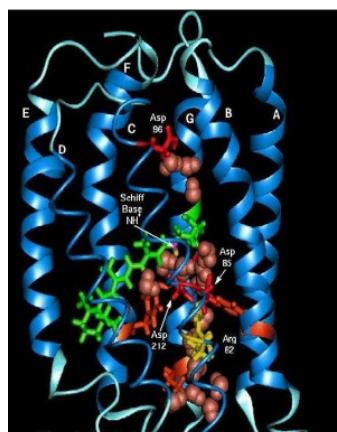
# Linear measurements

The measurements performed as a function of the incidence angle  $\alpha$





As the electrons of chiral molecules are displaced from their equilibrium by the application of the electromagnetic field, they are forced helical like paths



$$H = -\mu \cdot E - m \cdot B - Q : \nabla E + \dots \text{weak}$$



# Nonlinear optics, chirality, magneto-optics: a serendipitous road [Invited]

Andre Persoons<sup>1,2,3,\*</sup>

<sup>1</sup>Department of Chemistry, University of Leuven, B-3001 Leuven, Belgium

<sup>2</sup>College of Optical Sciences, University of Arizona, Tucson, Arizona 85721, USA

<sup>3</sup>[apersoons@optics.arizona.edu](mailto:apersoons@optics.arizona.edu)

[\\*andre.persoons@fys.kuleuven.be](mailto:andre.persoons@fys.kuleuven.be)

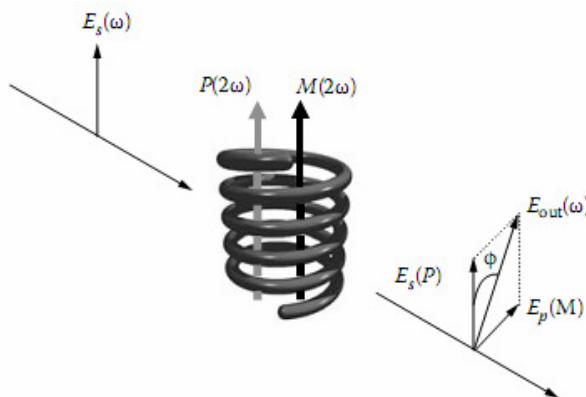
#143177 - \$15.00 USD

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1 May 2011 / Vol. 1, No. 1 / OPTICAL MATERIALS EXPRESS 5

## Nonlinear Optical Activity



$$\chi_{ijk}^{(2)} = N \sum_{ijk} \langle (I \cdot i)(J \cdot j)(K \cdot k) \rangle \beta_{ijk}$$

PHYSICAL REVIEW

VOLUME 130, NUMBER 3

1 MAY 1963

## Nonlinear Optical Properties of Solids: Energy Considerations\*

P. S. PERSHAN

Division of Engineering and Applied Physics, Harvard University, Cambridge, Massachusetts

(Received 17 December 1962)



## Nonlinear Optical Frequency Polarization in a Dielectric

ERIC ADLER

*IBM Watson Laboratory and Physics Department, Columbia University, New York, New York*

(Received 25 July 1963)

A perturbation calculation of optical frequency polarization quadratic in the Maxwell field is made for a dielectric in which the electrons are localized on units of the crystal. The result is expressed in a power series in  $K\langle r \rangle \sim 10^{-8}$  where  $\langle r \rangle$  is the size of the unit and  $K$  is the wave number of the field. The term of zeroth order is the electric-dipole term which vanishes in a crystal with a center of symmetry. The term first order in  $K\langle r \rangle$  is separated into electric-quadrupole and magnetic-dipole contributions by introducing a special gauge for the electromagnetic potentials. Higher power terms are neglected.

PRL 106, 103902 (2011)

PHYSICAL REVIEW LETTERS

week ending  
11 MARCH 2011

### Circular Dichroism in Biological Photonic Crystals and Cubic Chiral Nets

M. Saba,<sup>1</sup> M. Thiel,<sup>2</sup> M. D. Turner,<sup>3</sup> S. T. Hyde,<sup>4</sup> M. Gu,<sup>3</sup> K. Grosse-Brauckmann,<sup>5</sup> D. N. Neshev,<sup>6</sup>  
K. Mecke,<sup>1</sup> and G. E. Schröder-Turk<sup>1,\*</sup>

## Second-Order Nonlinear Optical Imaging of Chiral Crystals

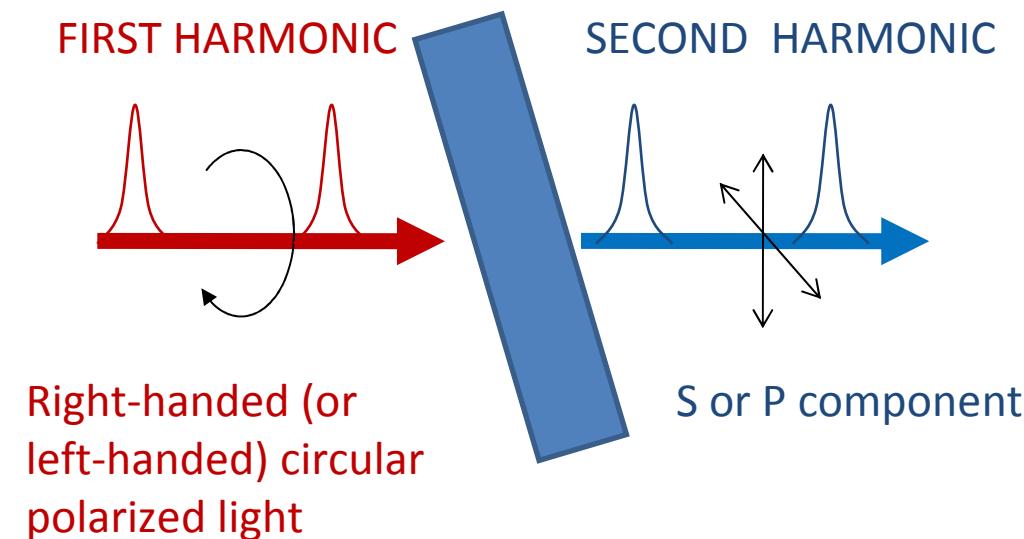
David J. Kissick, Debbie Wanapun, and Garth J. Simpson

Department of Chemistry, Purdue University, West Lafayette, Indiana 47907;



# SHG as chirality detector

It is possible to investigate chirality via second harmonic generation experiments



One useful parameter is:  
Second harmonic generation-circular dichroism (SHG-CD)

$$SHG - CD = \frac{I_L^{2\omega} - I_R^{2\omega}}{(I_L^{2\omega} + I_R^{2\omega})/2}$$

ORD

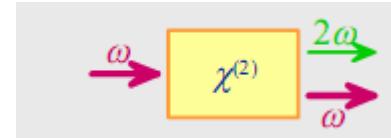
$$R(2\omega) = \frac{E_s(2\omega)}{E_p(2\omega)} = \frac{f_s E_p^2(\omega) + g_s E_s^2(\omega) + h_s E_p(\omega) E_s(\omega)}{f_p E_p^2(\omega) + g_p E_s^2(\omega) + h_p E_p(\omega) E_s(\omega)}$$



- Second order

$$P(t) = \chi^{(2)} E^2(t) = \chi^{(2)} E_\omega^2 e^{-i2\omega t}$$

→ second-harmonic generation



## Multipolar Symmetries in SecondOrder Nonlinear Optics

- electric-dipole and higher-multipole nonlinearities –
- isotropic and centrosymmetric materials



# Nonlinear susceptibility expansion

$$\chi(E) = \chi^{(1)} + \chi^{(2)}E + \chi^{(3)}EE + \dots$$

$$P(E) = \epsilon_0 (\chi^{(1)} + \chi^{(2)}E + \chi^{(3)}E^2 + \dots) E$$

$$\vec{P}_{eff}(2\omega) = \vec{P}_D(2\omega) - \frac{1}{\omega} \vec{k} x \vec{M}(2\omega)$$

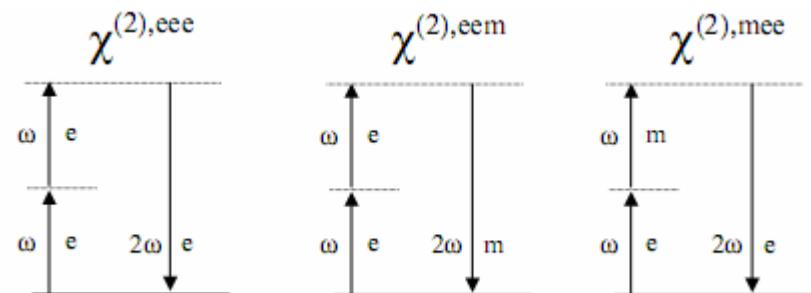
$$\vec{P}_D(2\omega) = \sum_{j,k} \chi_{ijk}^{(2)eee} E_j(\omega) E_k(\omega) + \chi_{ijk}^{(2)eem} E_j(\omega) B_k(\omega)$$

Dipole dipole source

$$\vec{M}(2\omega) = \sum_{j,k} \chi_{ijk}^{(2)mee} E_j(\omega) E_k(\omega)$$

Nonlinear magnetization

Magnetic dipole source





- Second-harmonic generation

$$P_i = \chi_{ijk}^{eee} E_j E_k + \chi_{ijk}^{eem} E_j B_k + \chi_{ijkl}^{eeQ} E_j \nabla_k E_l$$

axial      4th rank

- Magnetic and quadrupole tensors

- symmetry properties are different from those of the electric-dipole tensor

→ electric-dipole-forbidden  
effects can occur



# $\chi(2)$ Nonlinear Interaction

bulk

$$\eta = \frac{P^{(2\omega)}(L)}{P^{(\omega)}(0)} \propto \sin c^2 \left( \frac{1}{2} \Delta k L \right) \iff \Delta k = 0 \iff n_o^{2\omega} < n_e^\omega$$

Phase matching ==  
equal PHASE velocity

For nonhogenous media- Phc

P.M. (equal PHASE velocity),

For SH  $2K_1 = K_2$  ....  $n(w) = n(2 w)$  ---if there is *geometrical dispersion*

Thin films

No P.M. Interference between bound and free components ( surface and bulk)

At the nanoscale

No P.M.

Field localization and overlap



*J. Phys. Chem. B* 2009, 113, 13437–13445

13437

## Measurement of the Second-Order Hyperpolarizability of the Collagen Triple Helix and Determination of Its Physical Origin

Ariane Deniset-Besseau,<sup>†</sup> Julien Duboisset,<sup>‡</sup> Emmanuel Benichou,<sup>‡</sup> François Hache,<sup>†</sup> Pierre-François Brevet,<sup>‡</sup> and Marie-Claire Schanne-Klein<sup>\*‡</sup>

*Laboratoire d'Optique et Biosciences, Ecole Polytechnique, CNRS, INSERM U696, 91128 Palaiseau, France, and Laboratoire de Spectroscopie Ionique et Moléculaire, CNRS, Université Claude Bernard Lyon 1, 69622 Villeurbanne, France*

*Received: May 19, 2009; Revised Manuscript Received: July 16, 2009*

$$\chi_{ijk}^{(2)} = N \sum_{ijk} \langle (I \cdot i)(J \cdot j)(K \cdot k) \rangle \beta_{ijk}$$

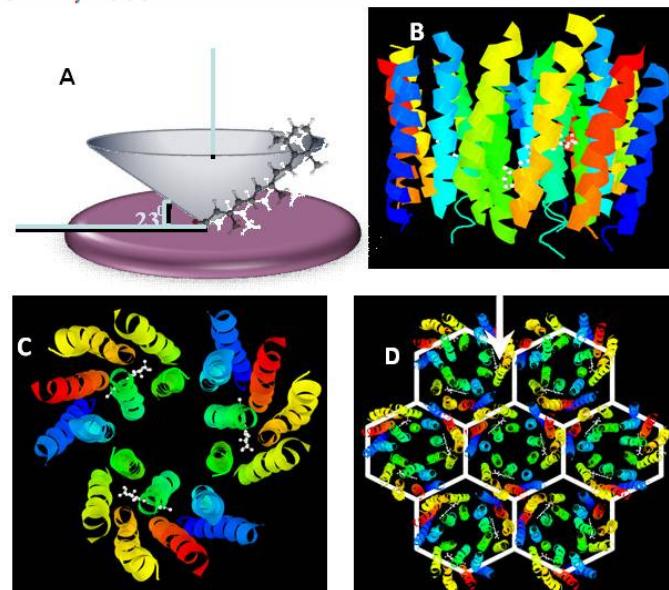
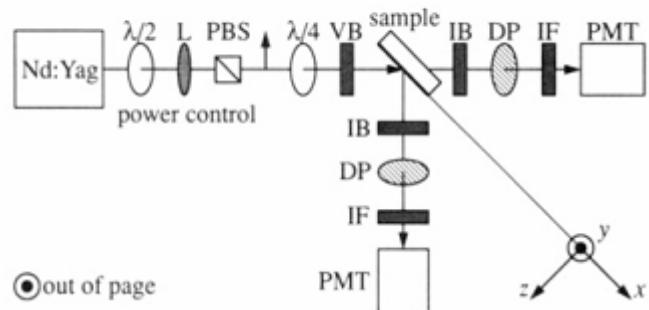
## Nonlinear Optical Activity and Biomolecular Chirality

Thierry Verbiest, Martti Kauranen, Andre Persoons, Marjo Ikonen, Jouni Kurkela, and Helge Lemmetyinen

*J. Am. Chem. Soc.*, 1994, 116 (20), 9203-9205 • DOI: 10.1021/ja00099a040 • Publication Date (Web): 01 May 2002

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9204 *J. Am. Chem. Soc.*, Vol. 116, No. 20, 1994



3164

Biophysical Journal Volume 73 December 1997 3164-3170

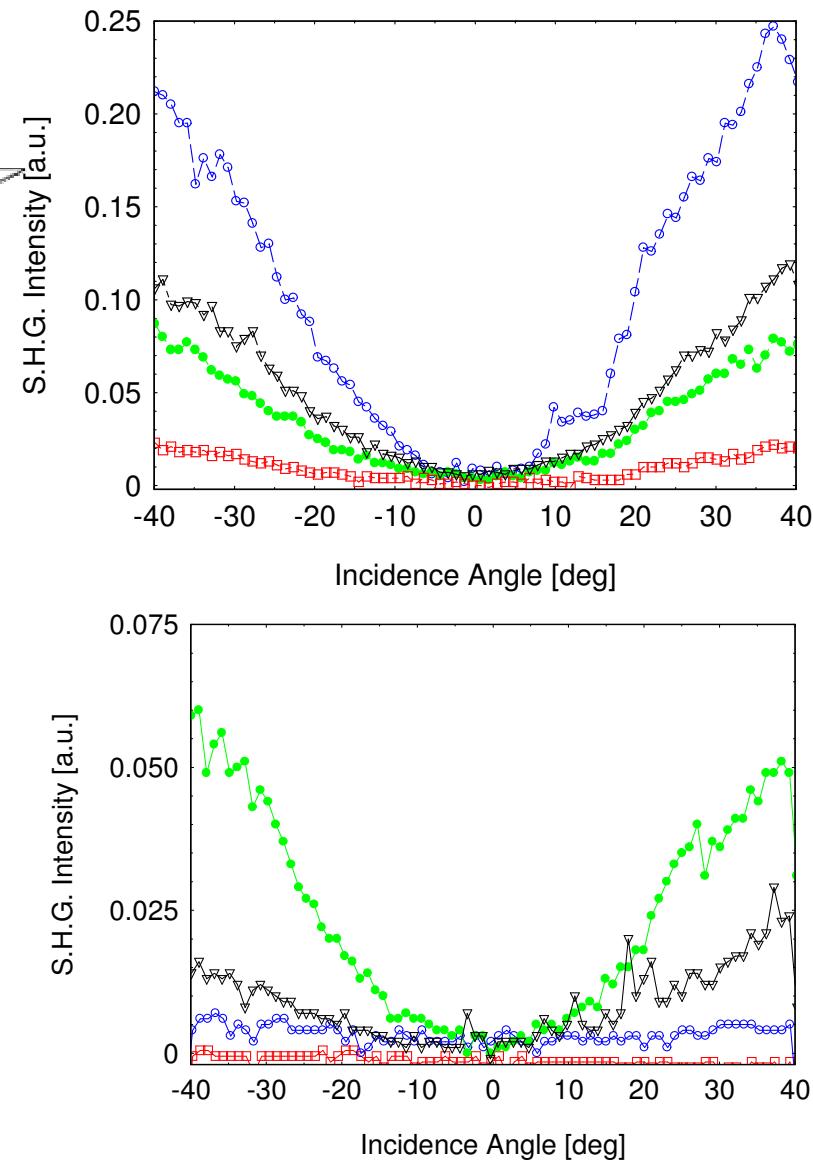
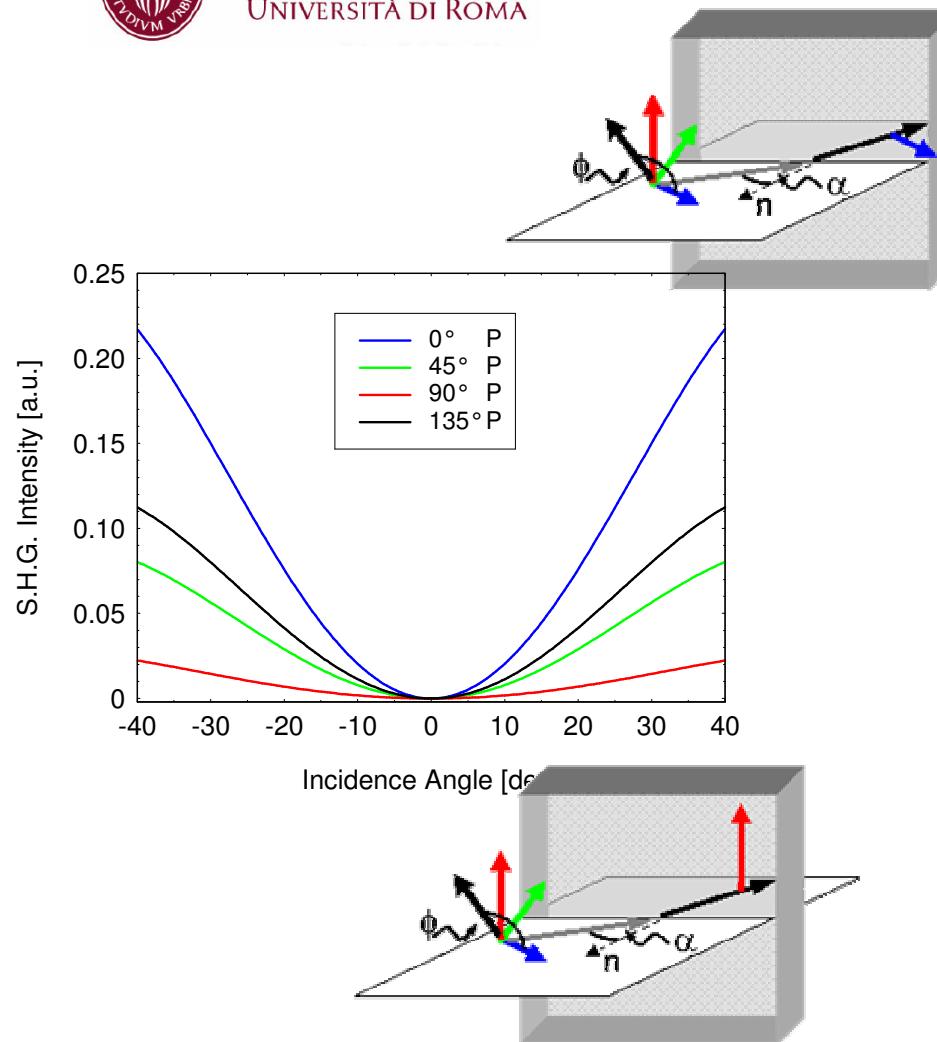
### Optical Rotation of the Second Harmonic Radiation from Retinal in Bacteriorhodopsin Monomers in Langmuir-Blodgett Film: Evidence for Nonplanar Retinal Structure

V. Volkov,\* Yu. P. Svirko,<sup>#§</sup> V. F. Kamalov,\* L. Song,\* and M. A. El-Sayed\*

\*Department of Chemistry and Biochemistry, Georgia Institute of Technology, Atlanta, Georgia 30332 USA; <sup>#</sup>Department of Physics, University of Southampton, Southampton SO17 1BJ, England; and <sup>§</sup>General Physics Institute, Moscow 117942, Russia



## Maker Fringes on BR

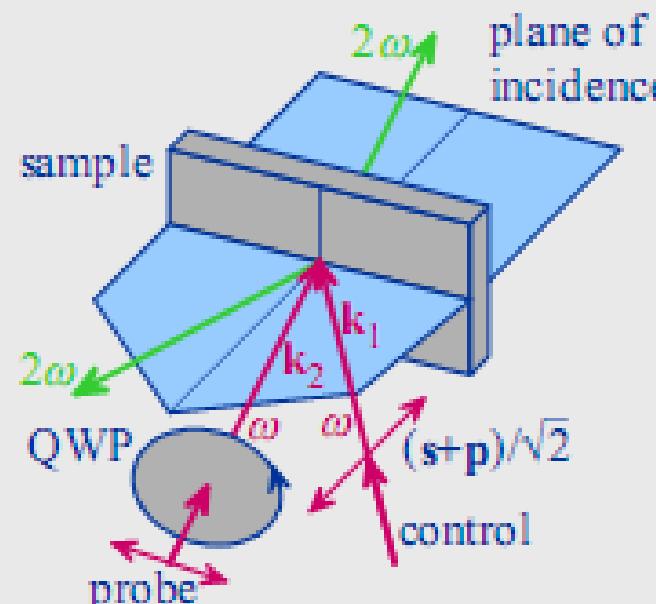
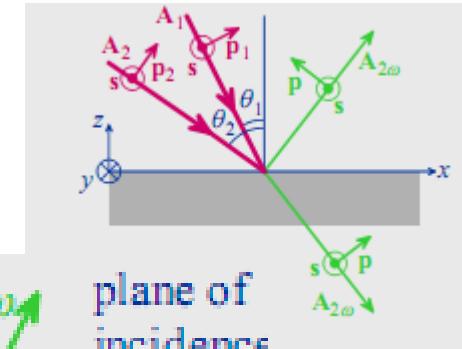


$d_{15} = d_{24} = 2.02 \text{ pm/V}$ ,  $d_{31} = d_{32} = 2.44 \text{ pm/V}$ ,  $d_{33} = 4.70 \text{ pm/V}$ , as well as the chiral coefficients  $d_{25} = -d_{14} = 0.34 \text{ pm/V}$ .



## Two beams

- **Control beam**
  - polarization fixed
- **Probe beam**
  - polarization varied
- **SHG signals**
  - reflected and transmitted





Transmission  
Coefficient at  $\omega$  for pump 1 & 2

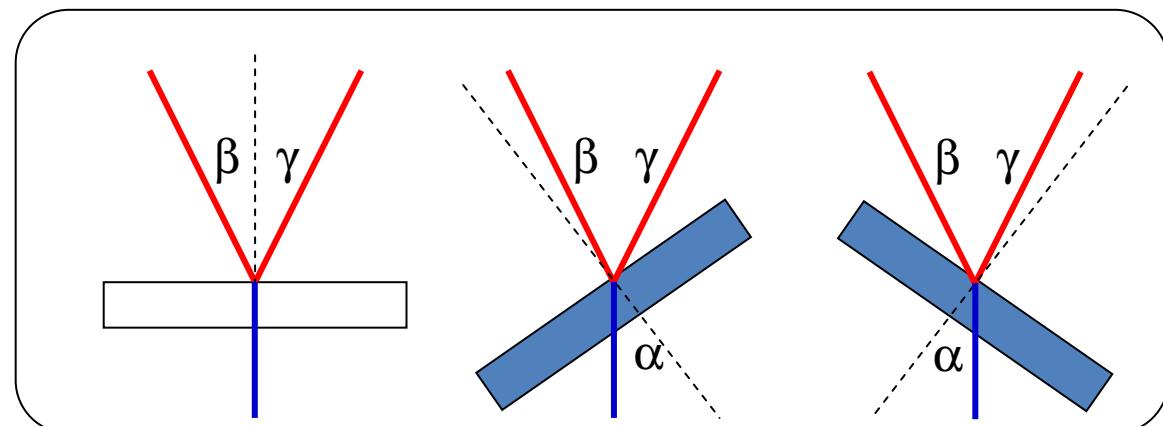
$$W_{\omega_1+\omega_2}(\alpha) = \left( \frac{512\pi^3}{A} \right) (t_{\omega_1})^2 \cdot (t_{\omega_2})^2 \cdot T_{\omega_1+\omega_2} \cdot W_{\omega_1} \cdot W_{\omega_2} \frac{\sin^2(\Psi_{SHG}(\alpha))}{[n_{\omega_1} \cdot n_{\omega_2} - n_{\omega_1+\omega_2}^2]^2} (d_{eff}(\alpha))^2$$

Power of the field at  $\omega$  for pump  
1 & 2

Transmittance  
at  $2\omega$

$$\Psi_{SHG} = \left( \frac{\pi L}{2} \right) \left( \frac{2}{\lambda} \right) [n_{\omega_1} \cdot \cos(\alpha' \omega_1) + n_{\omega_2} \cdot \cos(\alpha' \omega_2) - 2n_{\omega_1+\omega_2} \cdot \cos(\alpha' \omega_1 + \omega_2)]$$

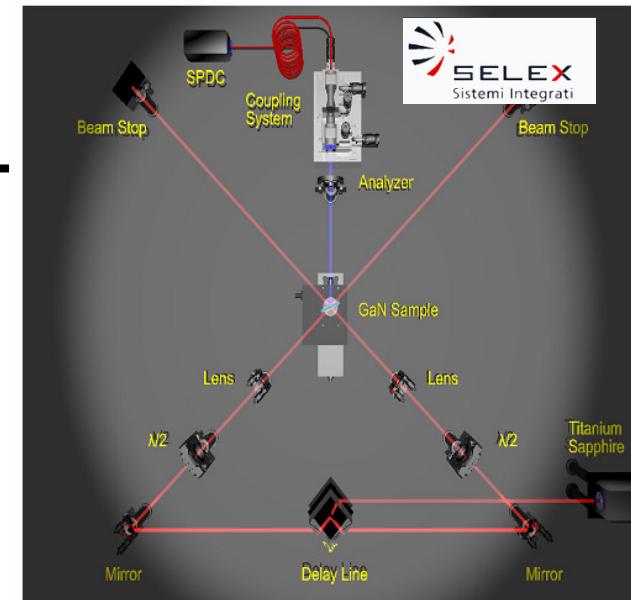
- ✓ The difference between the two incidence angles is set to a fixed value ( $18^\circ$ )



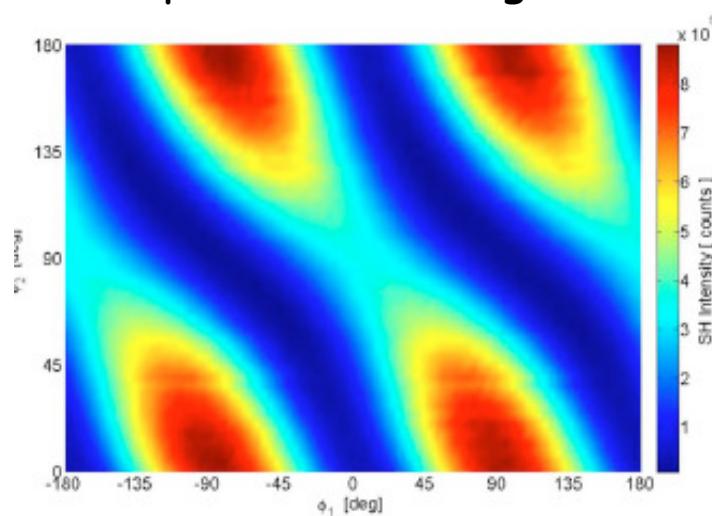


## Polarization chart

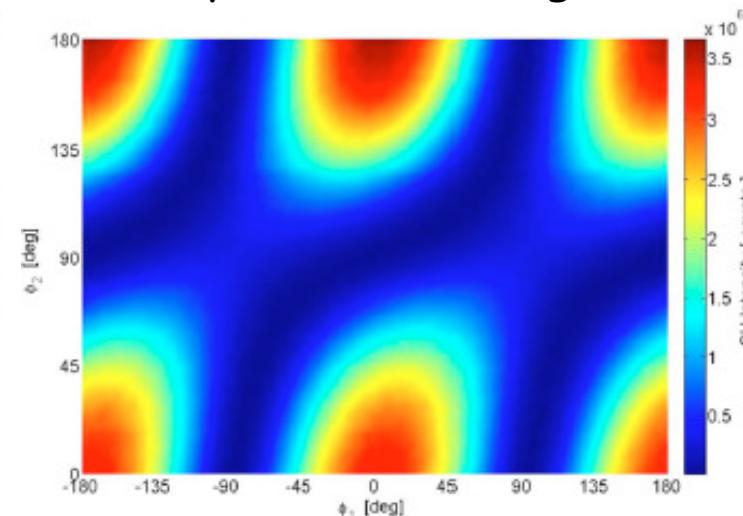
- At a fixed incidence angle, the generated signal is investigated by varying the polarization state of both fundamental beams.
- Example: GaN slab, ~300nm thick



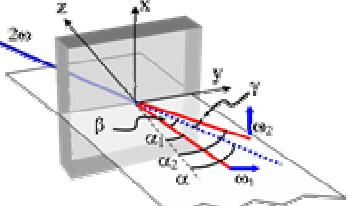
S-polarized SH signal



P-polarized SH signal

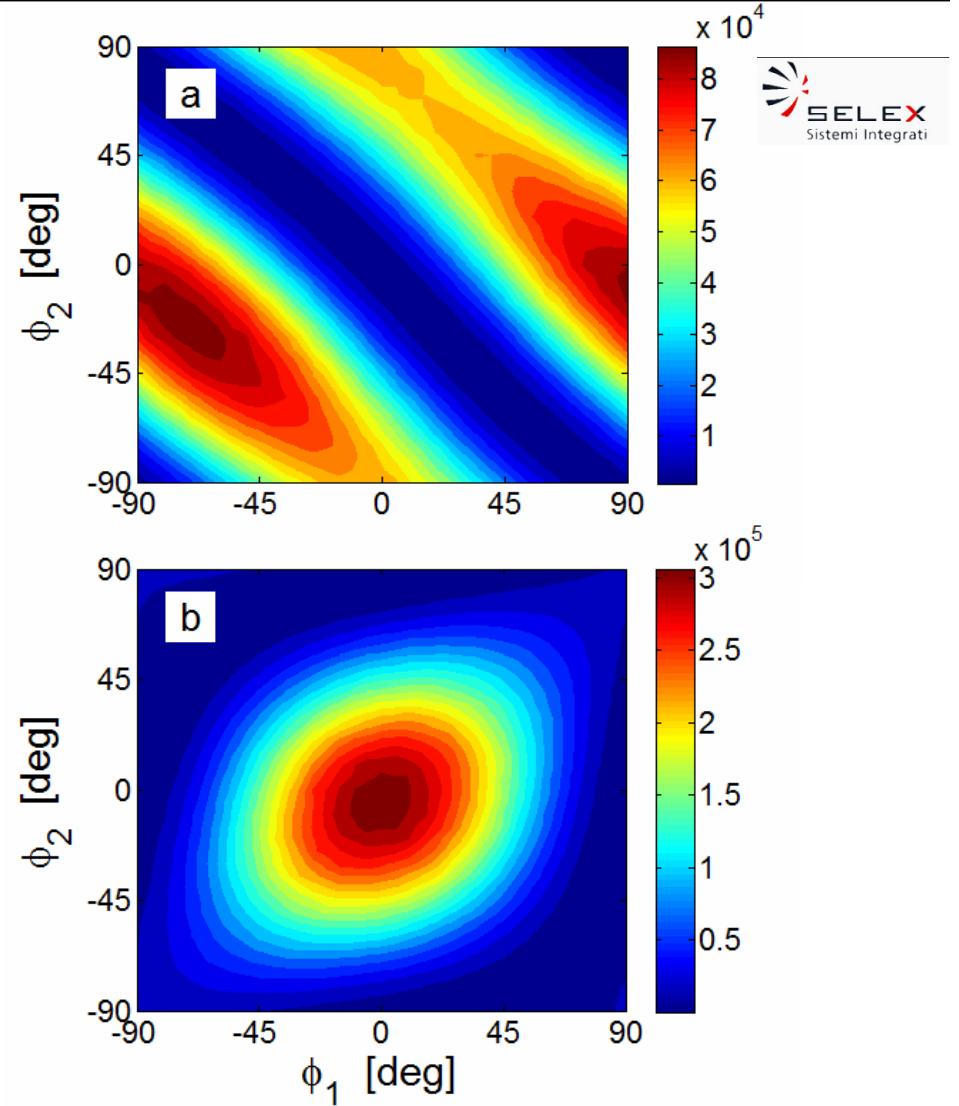
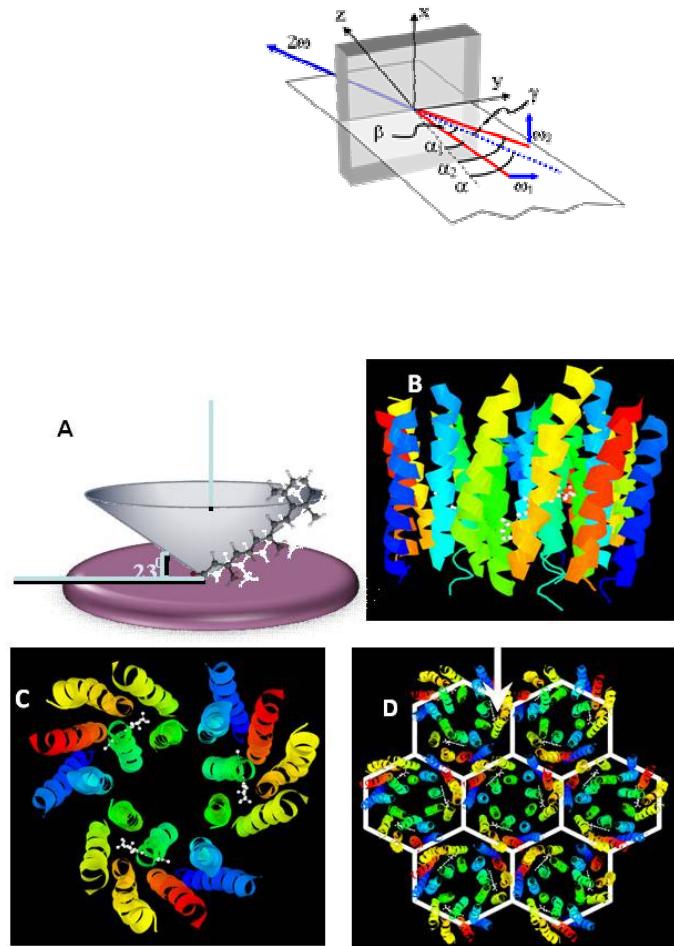


F.A.Bovino  
Selex-SI-  
Genova



M.C.Larciprete, F.A.Bovino, M.Giardina, A.Belardini, M.Centini, C.Sibilia, M.Bertolotti, A. Passaseo and V.Tasco, "Mapping the nonlinear optical susceptibility by noncollinear second-harmonic generation", Optics Letters, Vol. 34, pp.2189-2191 (2009)

# Polarization Chart in noncollinear SHG-BR



F.A Bovino et al- Selex-SI-Italy



# Power of SH Signal

$$W_{\omega_1+\omega_2}(\alpha) = \left( \frac{512\pi^3}{A} \right) (t_{\omega_1})^2 \cdot (t_{\omega_2})^2 \cdot T_{\omega_1+\omega_2} \cdot W_{\omega_1} \cdot W_{\omega_2} \frac{\sin^2(\Psi_{SHG}(\alpha))}{[n_{\omega_1} \cdot n_{\omega_2} - n_{\omega_1+\omega_2}^2]^2} (|d_{eff}^{(2)}(\alpha)|)^2$$

$$\begin{aligned} d_{eff}^{(2)} = & (\hat{p}_{out})(d_{il}^{(eee)})(\hat{p}_{in}^{1e})(\hat{p}_{in}^{2e}) + \\ & (\hat{p}_{out})(d_{il}^{(eem1)})(\hat{p}_{in}^{1e})(\hat{p}_{in}^{2m}) + (\hat{p}_{out}^e)(d_{il}^{(eem2)})(\hat{p}_{in}^{1m})(\hat{p}_{in}^{2e}) + \\ & (\hat{p}_{out})xK(d_{il}^{(mee)})(\hat{p}_{in}^{1e})(\hat{p}_{in}^{2e}) \end{aligned}$$

$\hat{p}_{out}$  : polarization versor of output sh signal

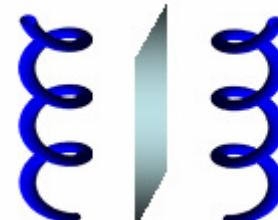
$\hat{p}_{in}^1$  : polarization versor of input pump 1 signal

$\hat{p}_{in}^2$  : polarization versor of input pump 2 signal

3p

$$d_{il} = \frac{\hat{\chi}_{ijk}^{(2)}}{2}$$

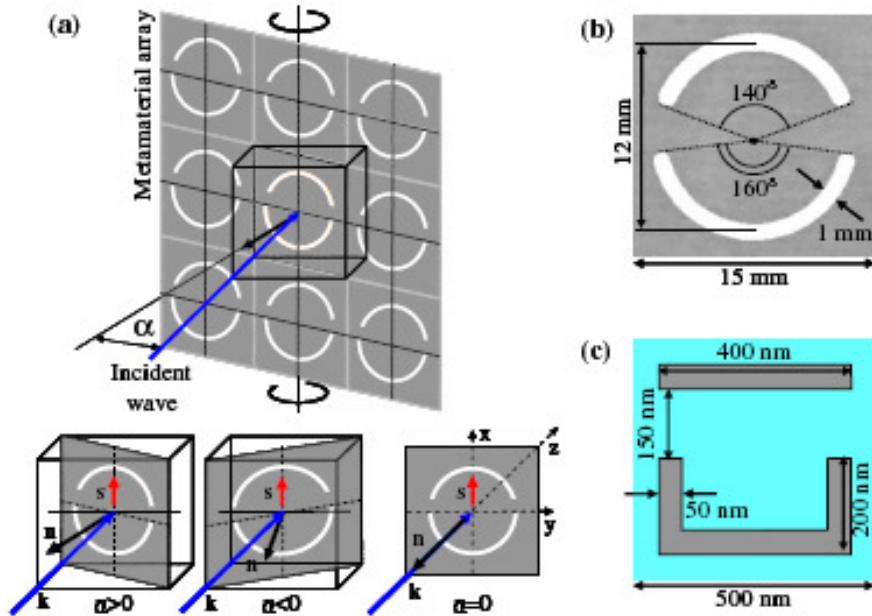
Intrinsic 3D chirality



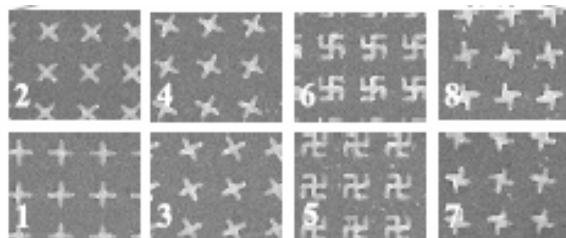
$$\begin{pmatrix} 0 & 0 & 0 & d_{14}^{eee} & d_{15}^{eee} & 0 \\ 0 & 0 & 0 & d_{24}^{eee} & d_{25}^{eee} & 0 \\ d_{31}^{eee} & d_{32}^{eee} & d_{33}^{eee} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & d_{14}^{eem} & d_{15}^{eem} & 0 \\ 0 & 0 & 0 & d_{24}^{eem} & d_{25}^{eem} & 0 \\ d_{31}^{eem} & d_{32}^{eem} & d_{33}^{eem} & 0 & 0 & d_{36}^{eem} \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & d_{14}^{ree} & d_{15}^{ree} & 0 \\ 0 & 0 & 0 & d_{24}^{ree} & d_{25}^{ree} & 0 \\ d_{31}^{ree} & d_{32}^{ree} & d_{33}^{ree} & 0 & 0 & 0 \end{pmatrix}$$

# Extrinsic Chirality

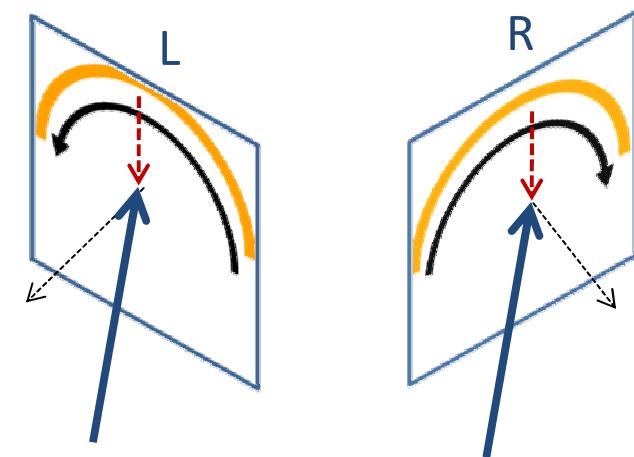
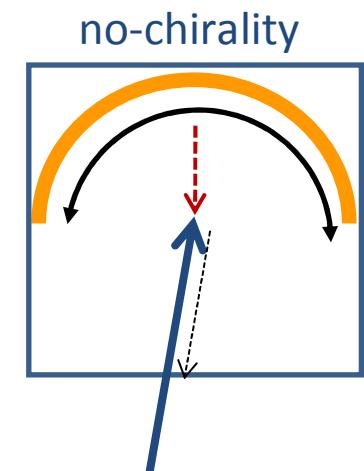
E. Plum, PRL 102,113902 (2009)



S.N. Volkov, PRA 79, 043819 (2009)



**Extrinsic chirality:**  
A non-chiral  
element breaks the  
symmetry when is  
tilted with respect  
the incidence light.  
The element thus  
assumes a chiral  
behavior





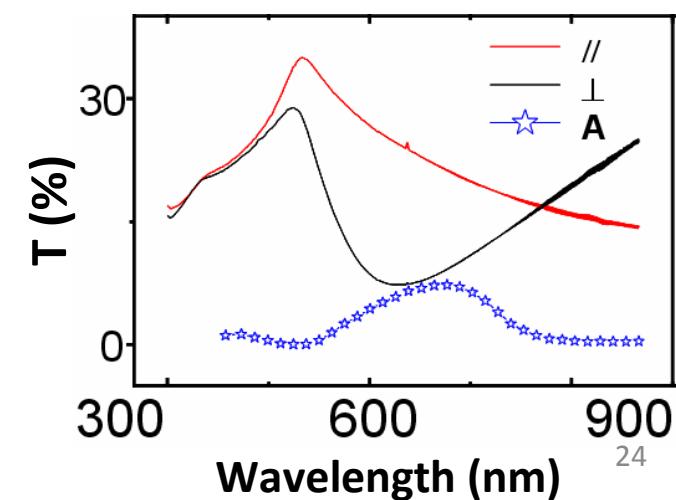
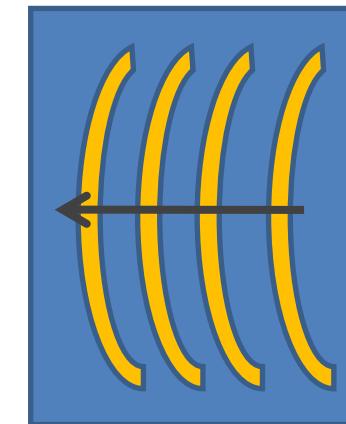
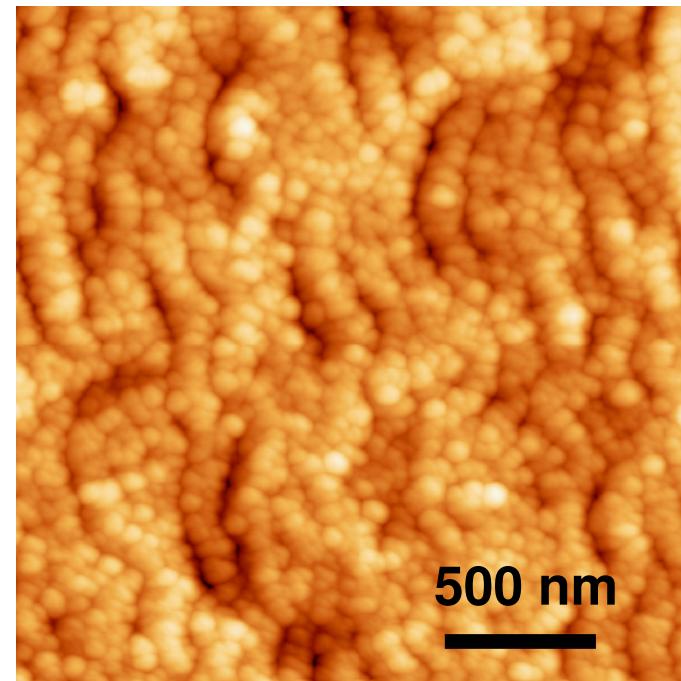
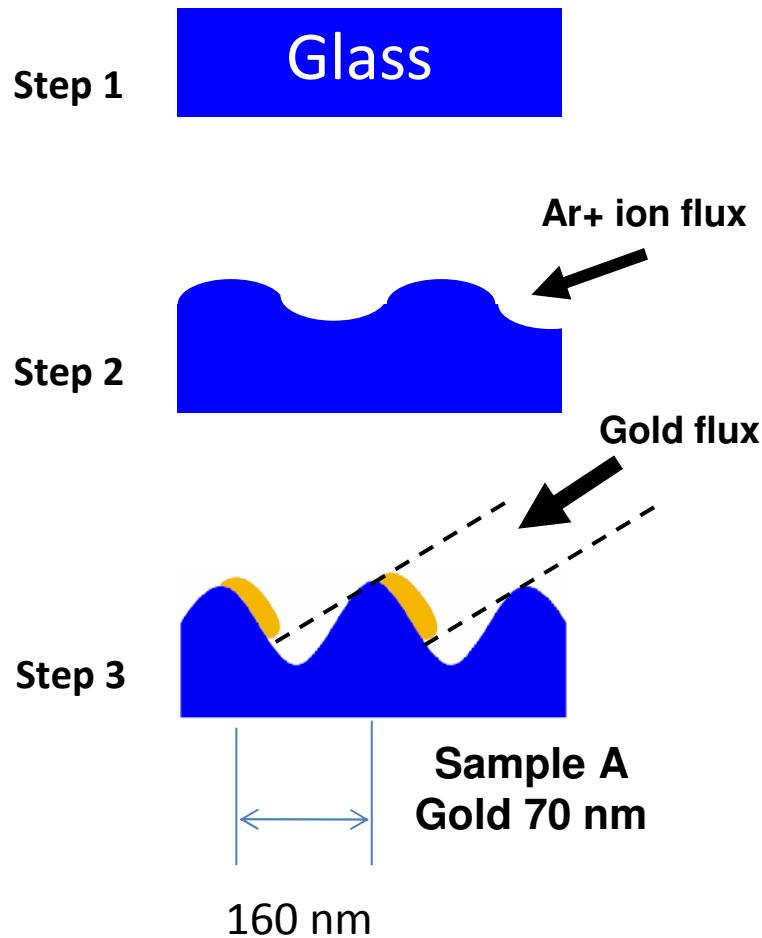
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## Extrinsic Chirality



# Test sample fabrication – Sample A

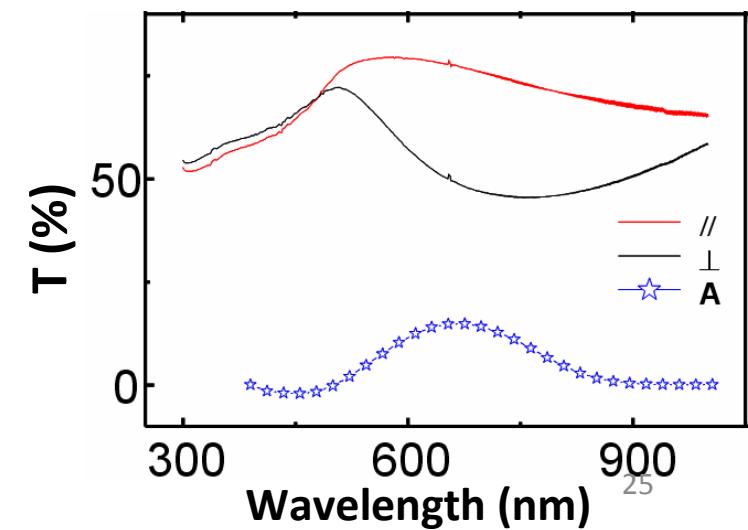
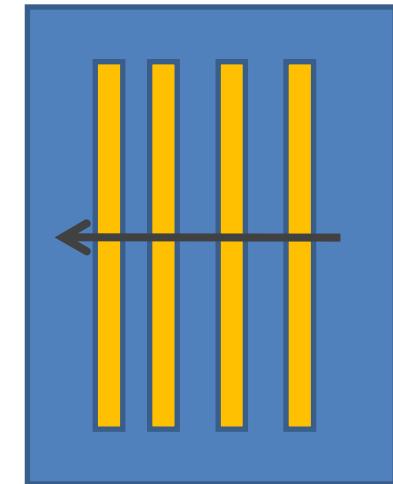
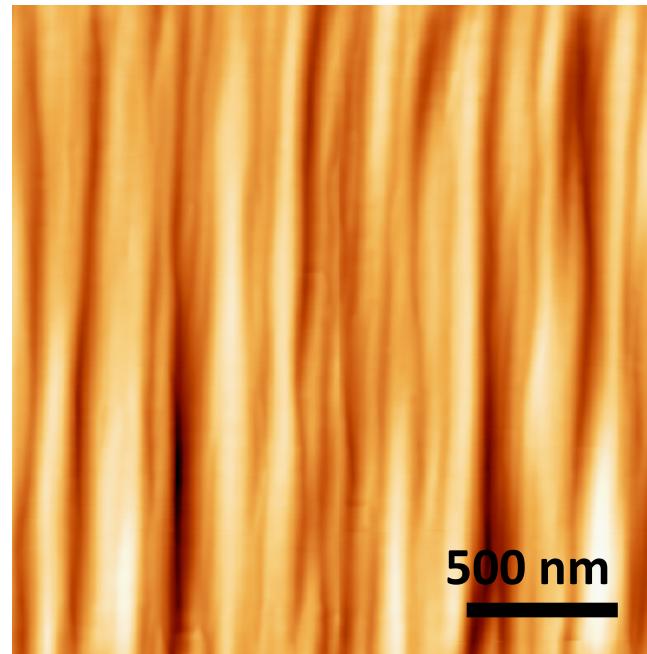
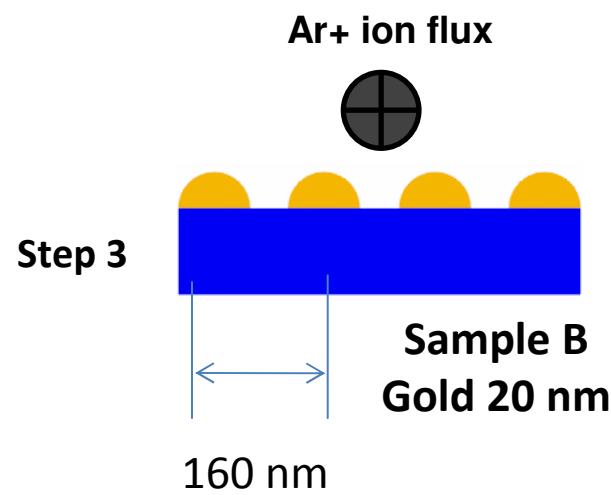
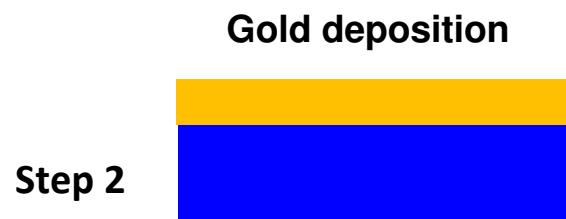
Low cost self assembling procedure



# Reference sample fabrication – Sample B

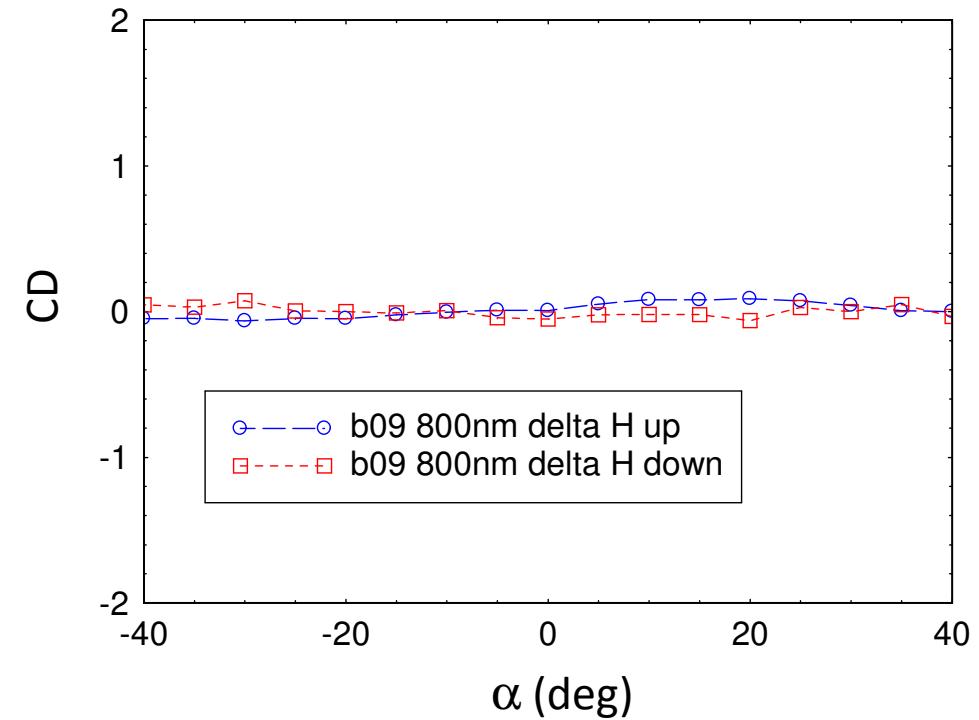
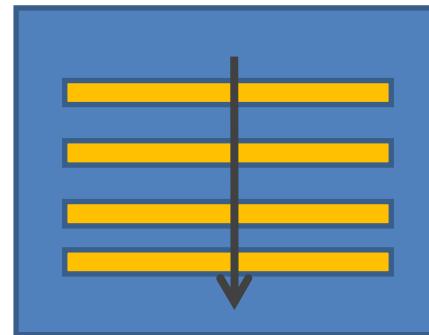
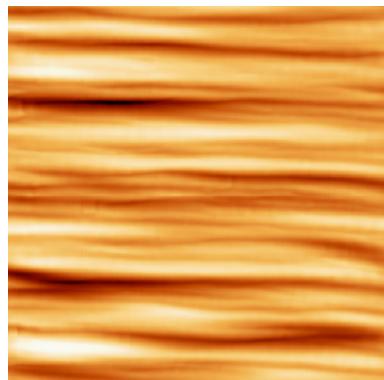
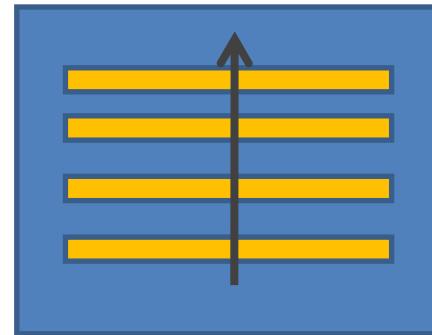
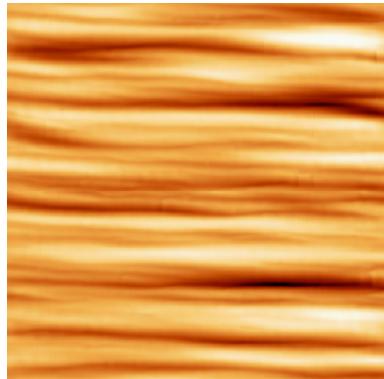


Low cost self assembling procedure



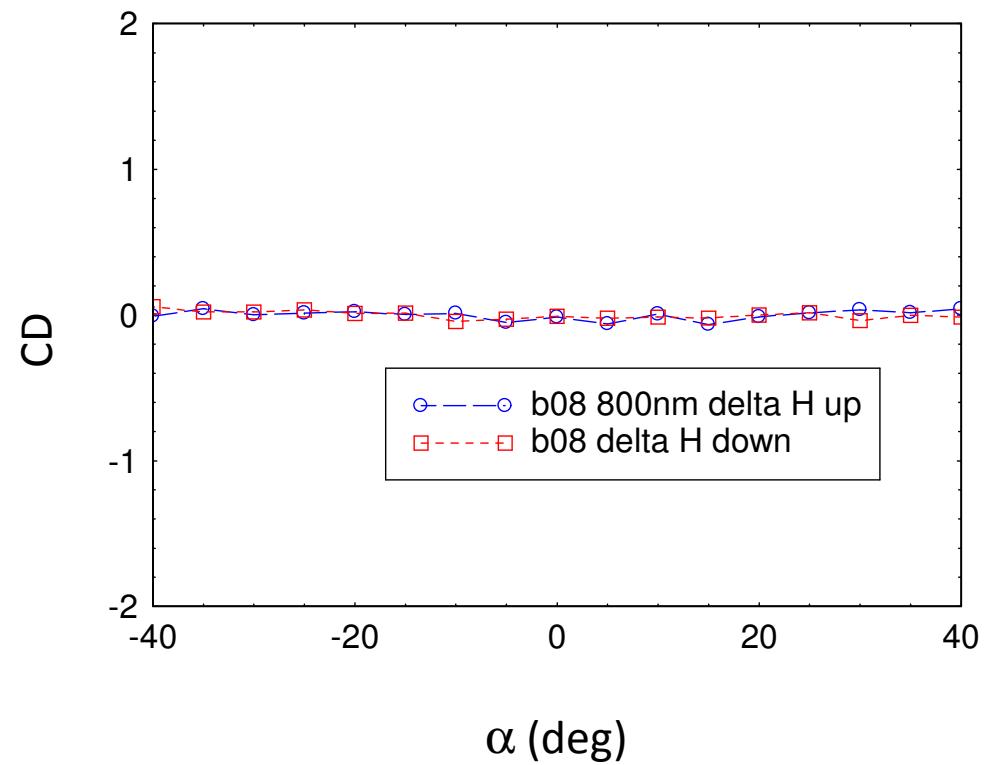
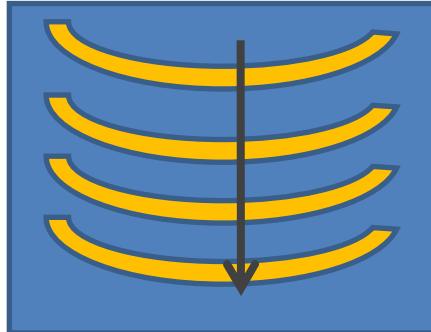
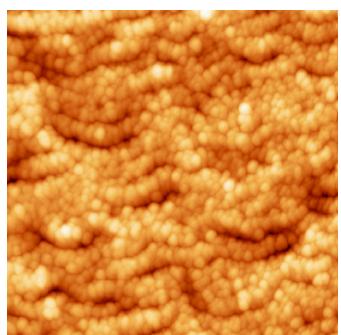
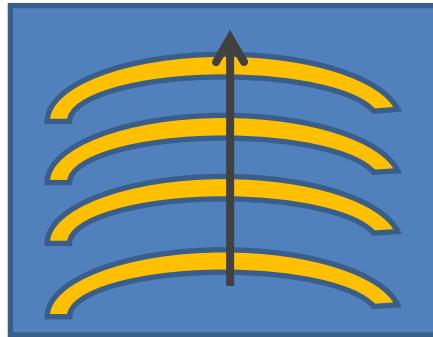
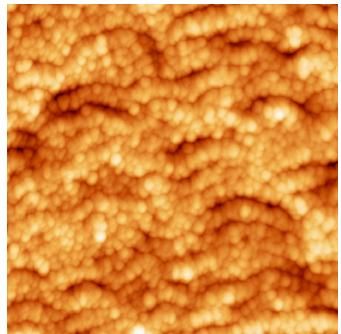


# Linear measurements on sample B - P pol





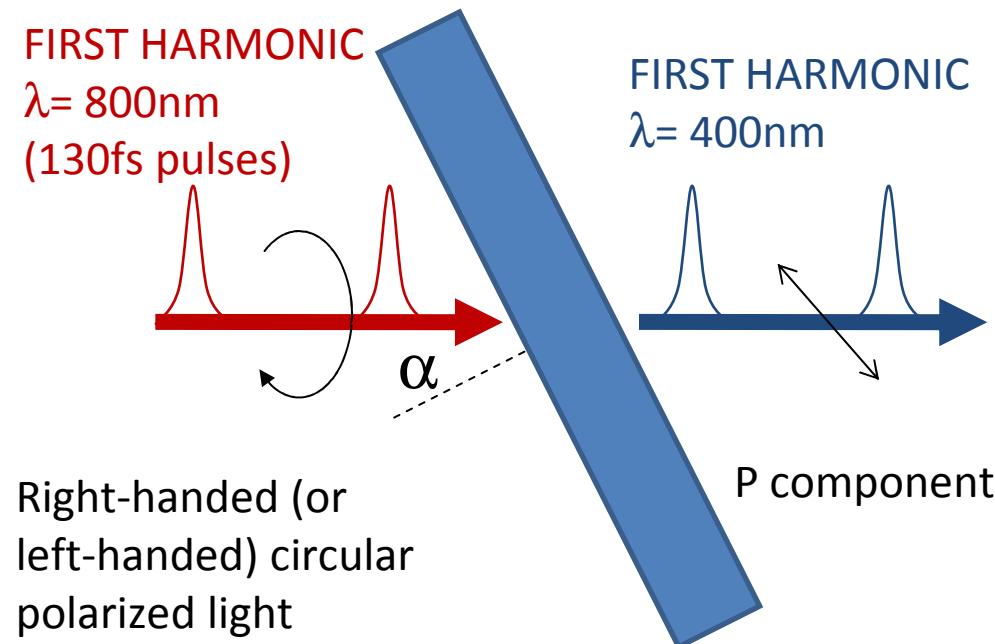
# Linear measurements on sample A - P pol





# SHG measurements – P pol

The measurements were performed as a function of the incidence angle  $\alpha$

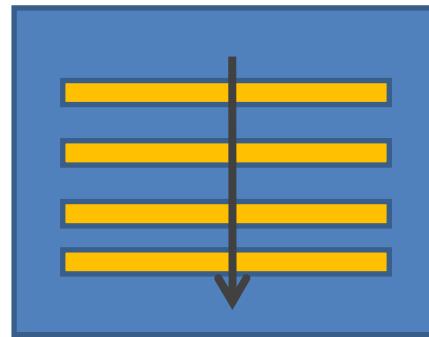
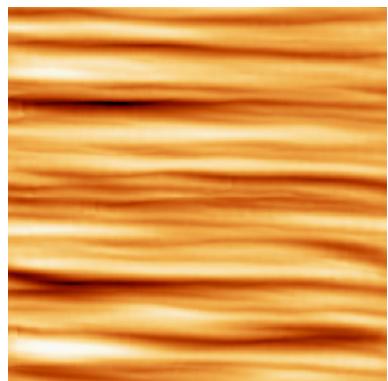
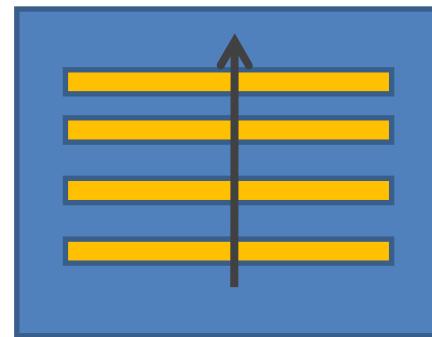
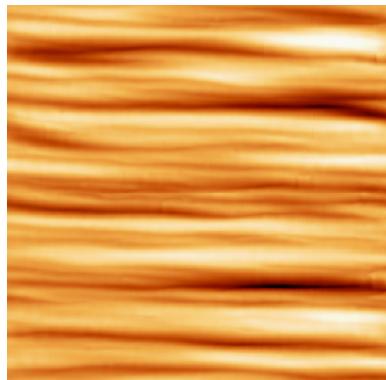


Second harmonic  
generation - circular  
dichroism (SHG-CD)

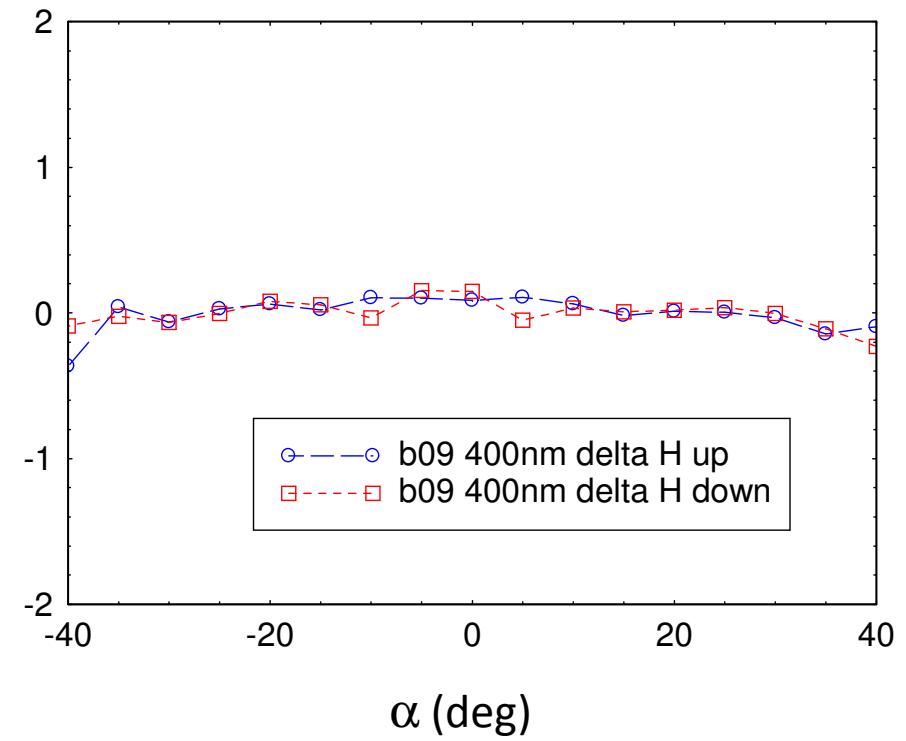
$$SHG - CD = \frac{I_L^{2\omega} - I_R^{2\omega}}{(I_L^{2\omega} + I_R^{2\omega})/2}$$



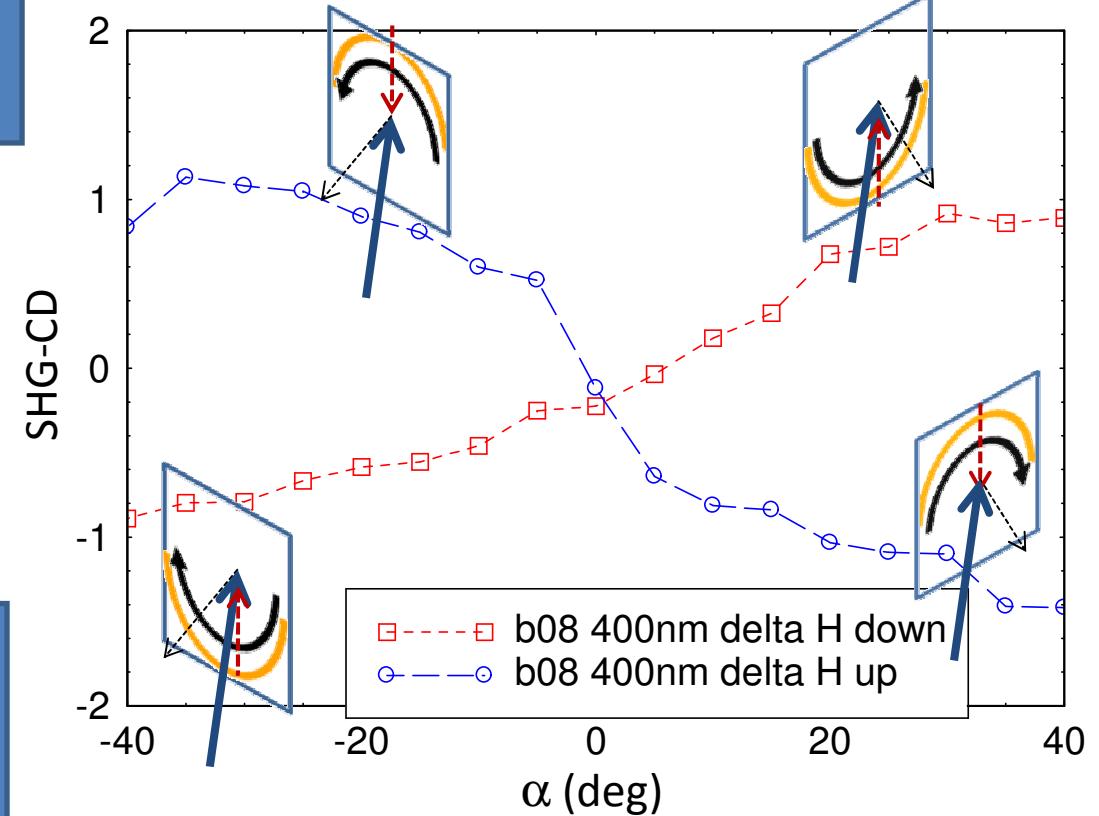
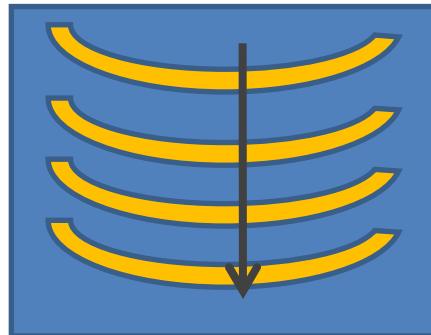
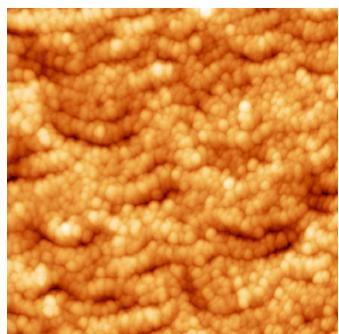
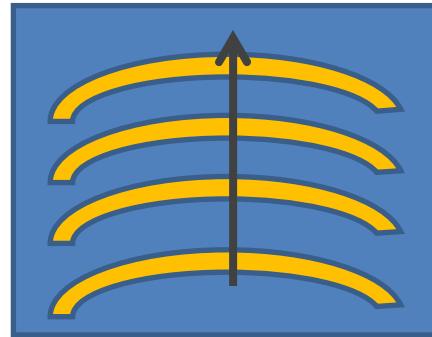
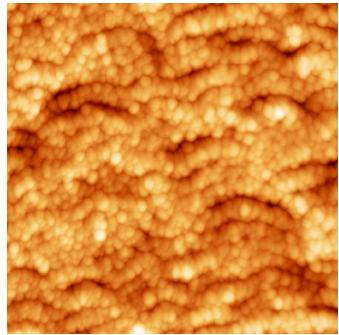
# SHG measurements on sample B – P pol



SHG-CD



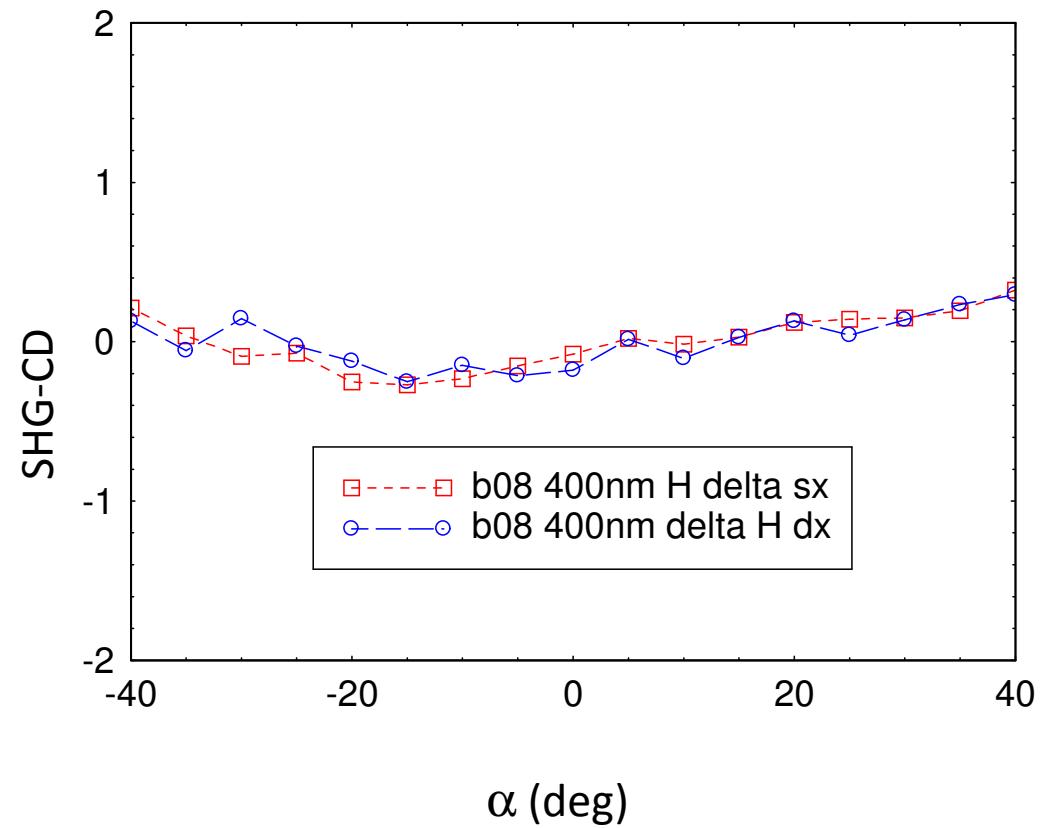
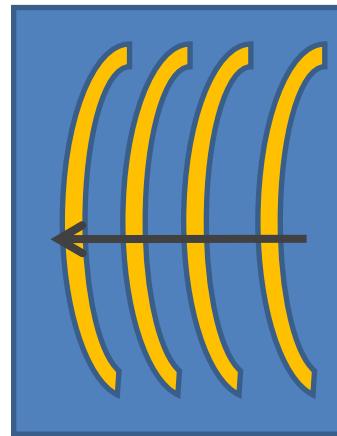
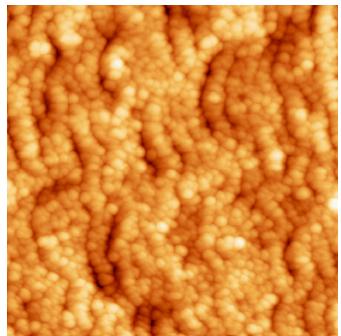
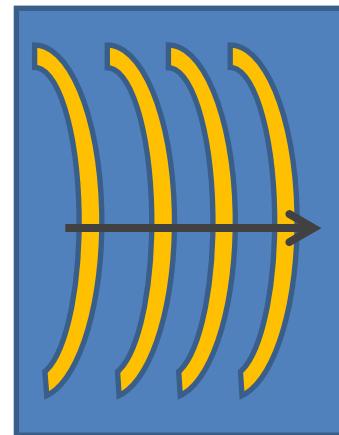
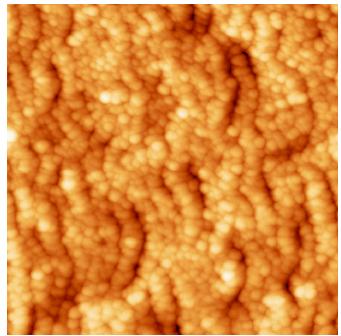
# SHG measurements on sample A – P pol



A.Belardini et al PRL  
2011

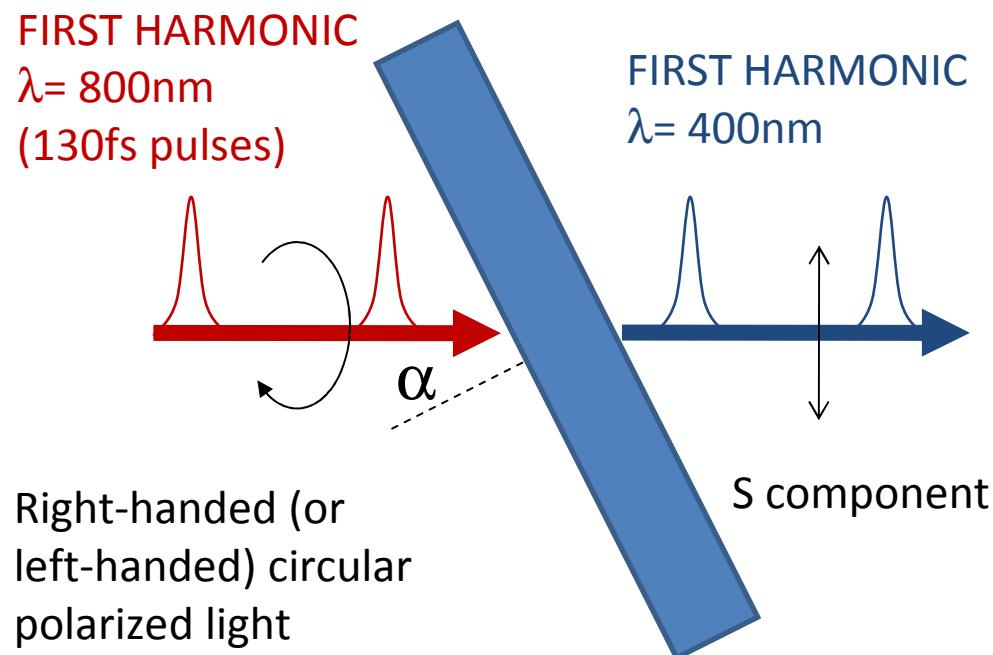


# SHG measurements on sample A – P pol





# SHG measurements – S pol

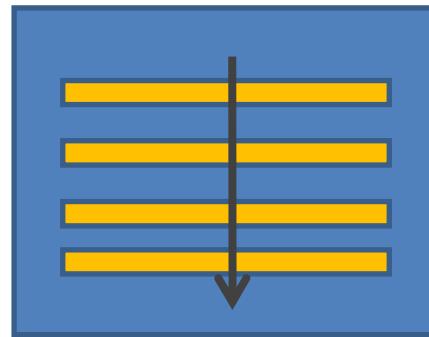
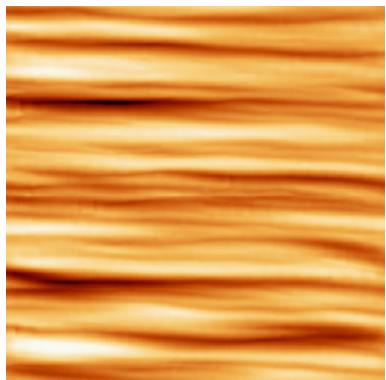
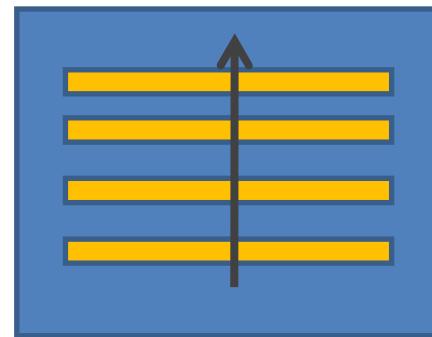
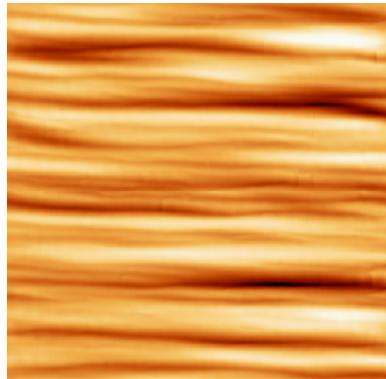


Second harmonic generation - circular dichroism (SHG-CD)

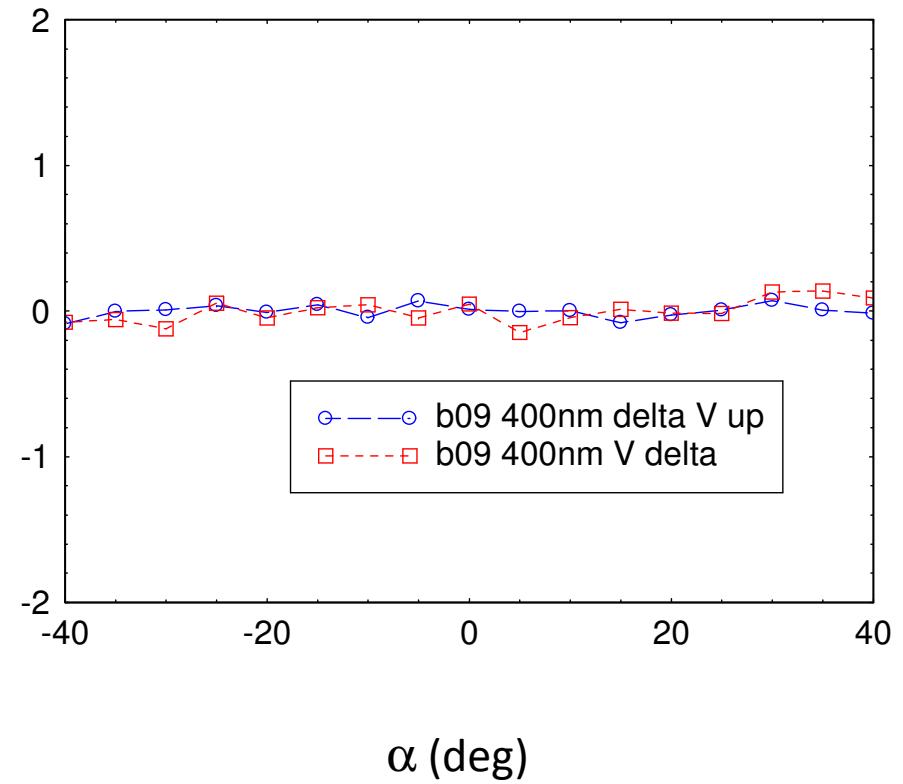
$$SHG - CD = \frac{I_L^{2\omega} - I_R^{2\omega}}{(I_L^{2\omega} + I_R^{2\omega})/2}$$



# SHG measurements on sample B – S pol

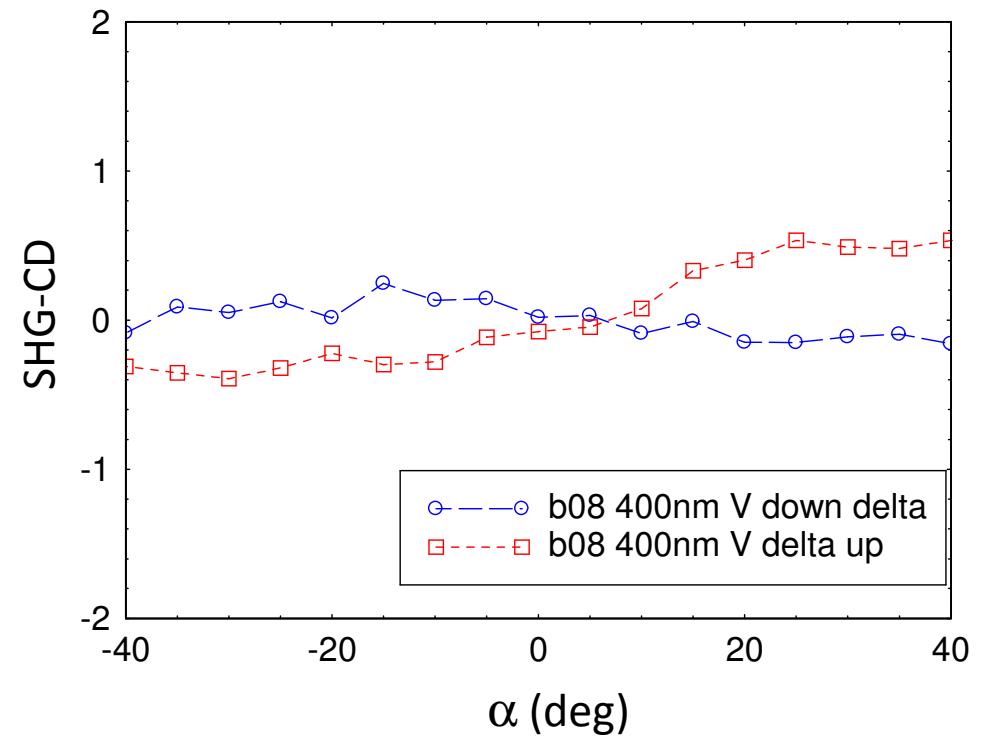
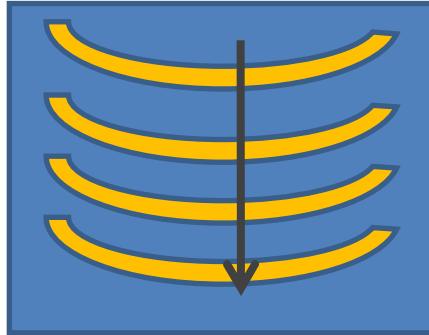
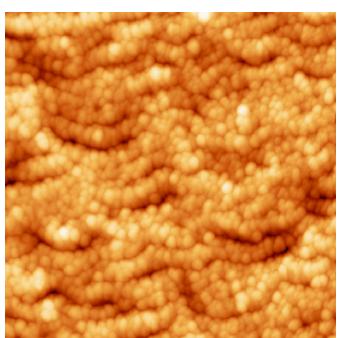
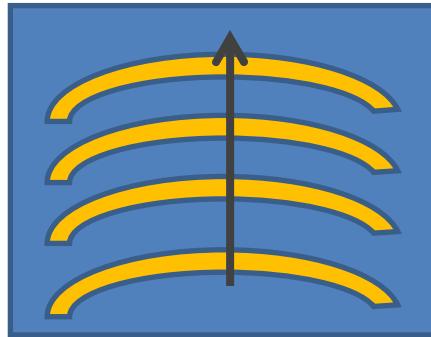
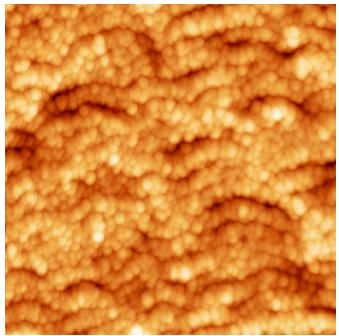


SHG-CD





# SHG measurements on sample A – S pol

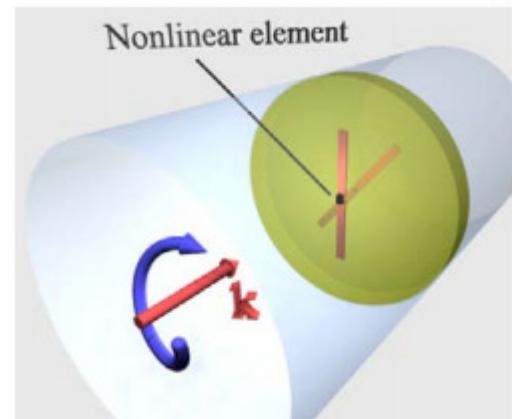
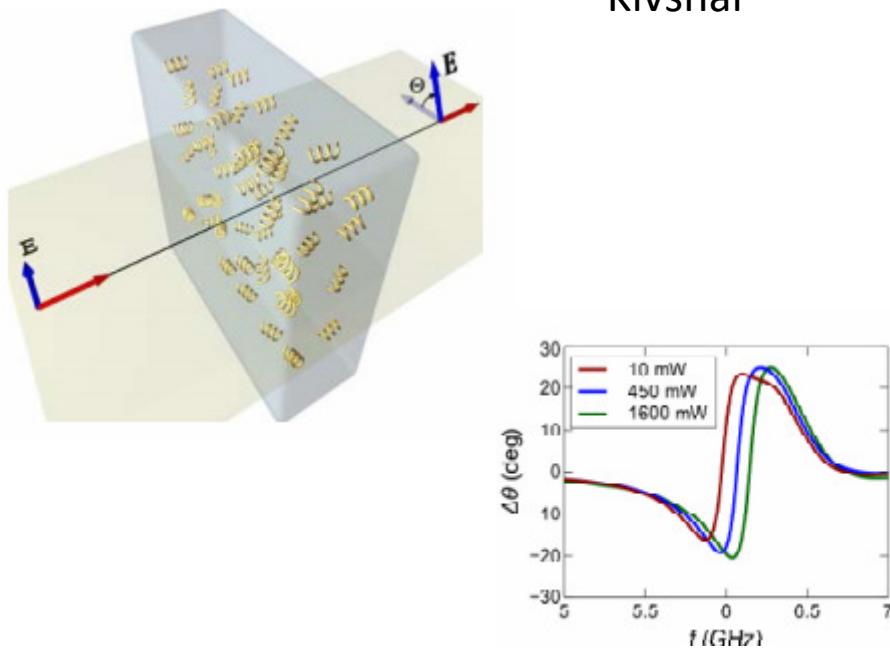




Chiral Metamaterials  
Unlocking Nonlinear Optical Activity  
Intensity dependent effects

AOS News Volume 25 Number 1 2011

by David A. Powell, Ilya V. Shadrivov, Vassili A. Fedotov, Nikolay I. Zheludev and Yuri S. Kivshar



Optical diode for  
circular polarized  
light



## Optical chirality without optical activity: How surface plasmons give a twist to light

Aurélien Drezet, Cyriaque Genet, Jean-Yves Laluet and Thomas  
W. Ebbesen

ISIS, Louis Pasteur University, 8 allée Gaspard Monge, 67000, Strasbourg, France

#97343 - \$15.00 USD      Received 12 Jun 2008; revised 16 Jul 2008; accepted 16 Jul 2008; published 5 Aug 2008  
(C) 2008 OSA      18 August 2008 / Vol. 16, No. 17 / OPTICS EXPRESS 12559

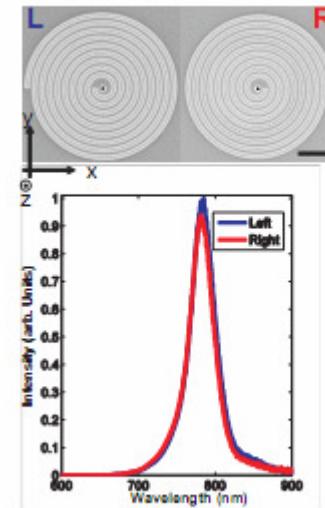


Fig. 1. Chiral plasmonic metasubstrates. On the top panel: scanning electron micrographs of the left (L) and right (R) handed enantiomer (mirror image) plasmonic chiral structures investigated. The size bar is 3  $\mu\text{m}$  long. The parameters characterizing the structures are the following: hole diameter  $d = 500$  nm, film thickness  $\Lambda = 100$  nm, dielectric constant  $\kappa = 7.95$ .



- 3D, 2D-chirality, Extrinsic chirality
- SHG as chirality detection method
- Nonlinear measurements
- Intensity dependence



## Acknwoledgments

SBAI

A.Belardini  
A.Benedetti  
M.Bertolotti  
M.Centini  
MC Larciprete

F.A.Bovino –Selex-SI-Genova-Italy  
A.Passaseo CNR- Lecce –Italy  
C.Gergeley – Univ. Montpellier2 – Montpellier-Fr  
F.Buatier- University of Genova -Italy



Thank you !!!