



The Abdus Salam
International Centre for Theoretical Physics



2333-15

Workshop on Science Applications of GNSS in Developing Countries (11-27 April), followed by the: Seminar on Development and Use of the Ionospheric NeQuick Model (30 April-1 May)

II April - I May, 2012

Scintillation Modeling

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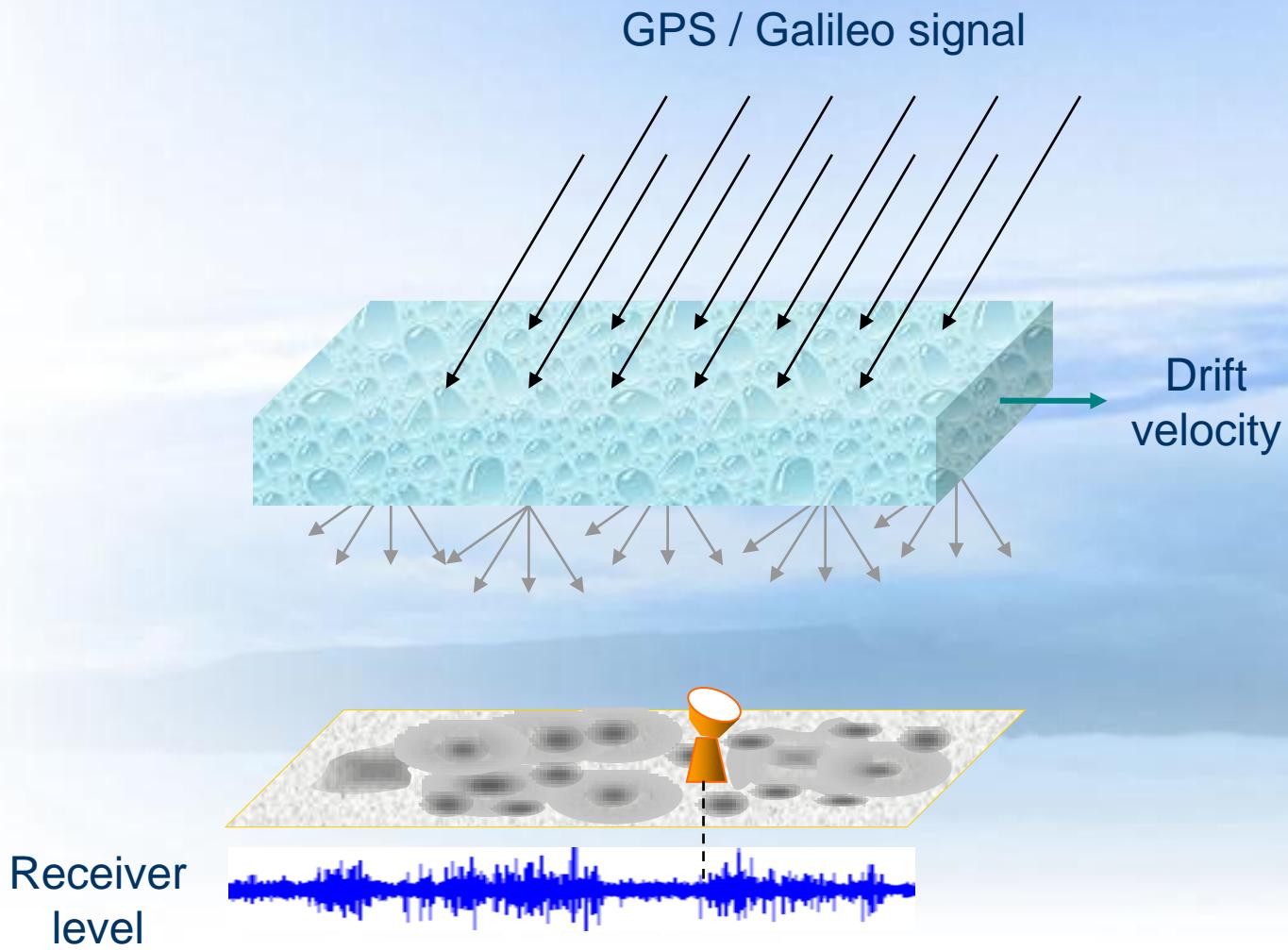
Low Latitudes Ionospheric Scintillation

Y. Béniguel, P. Hamel
IEEA, Paris, France

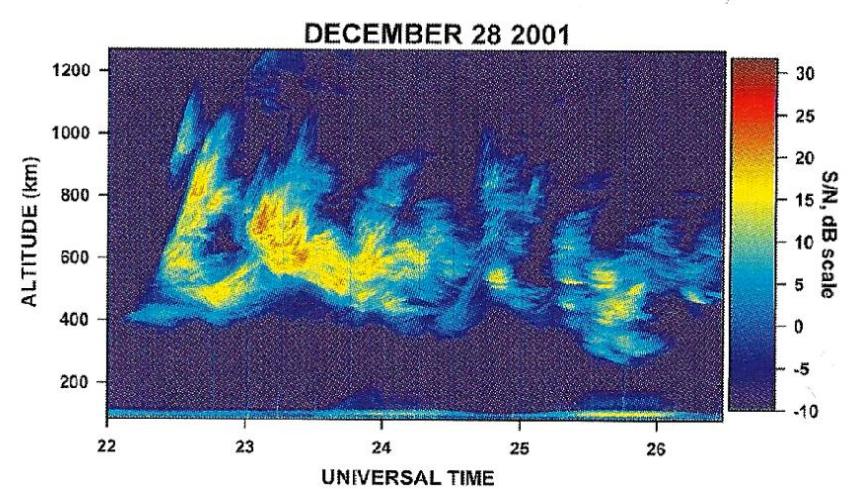
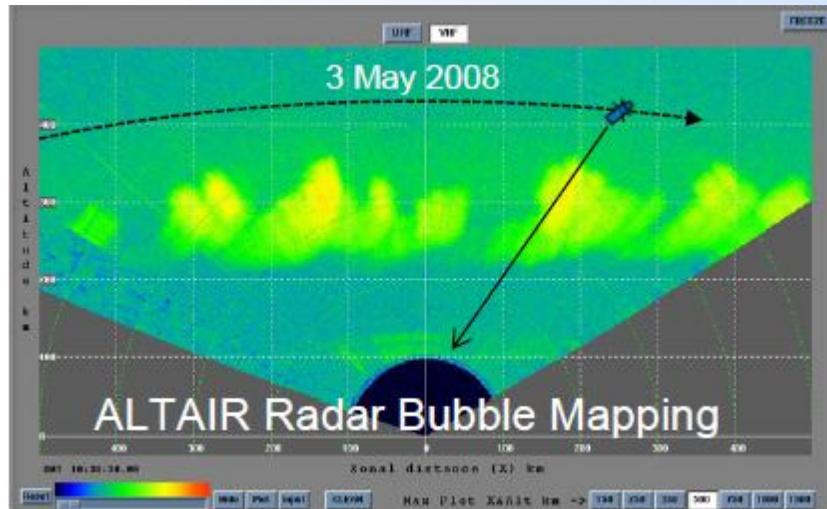
Outline

- Physical mechanism
- Observations
- Modelling & Comparisons
- Scattering function calculation (SAR observations)
- Conclusion

Physical Mechanism



Medium Radar Observations

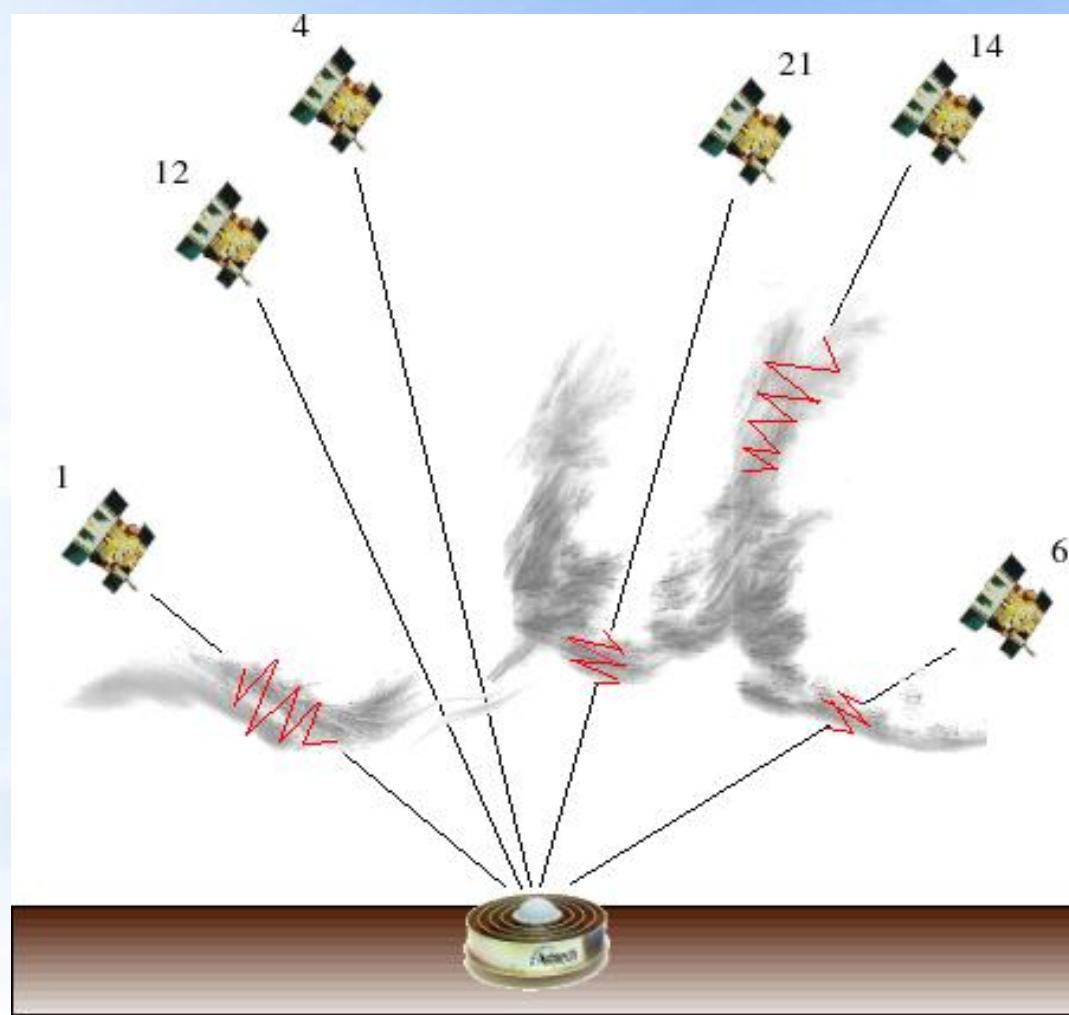


Observations at Kwajalen Islands
Courtesy K. Groves, AFRL

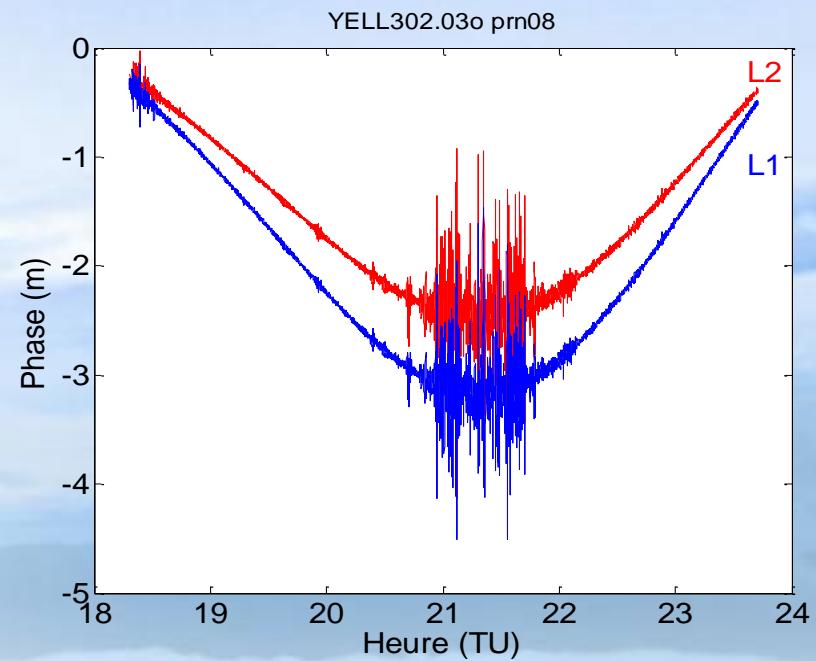
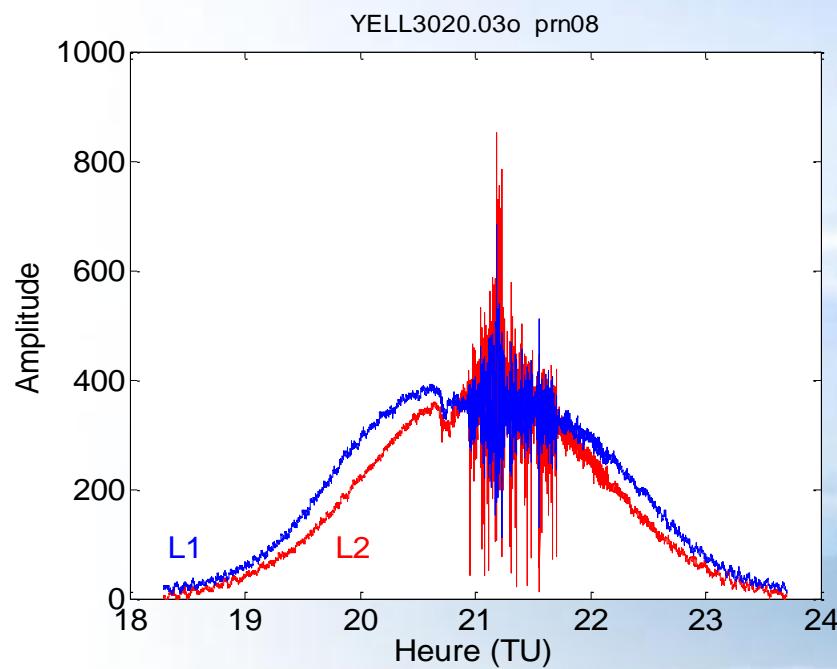
Observations in Brazil
Courtesy E. de Paula, INPE

The vertical extent may reach hundreds of kilometers

Medium Radar Observations



Signal at receiver level (Measurements)



Intensity

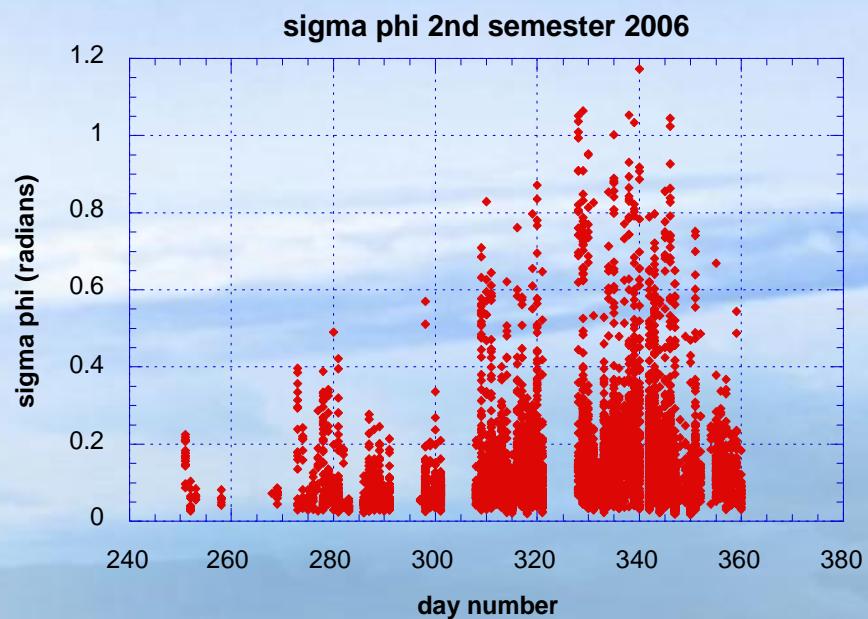
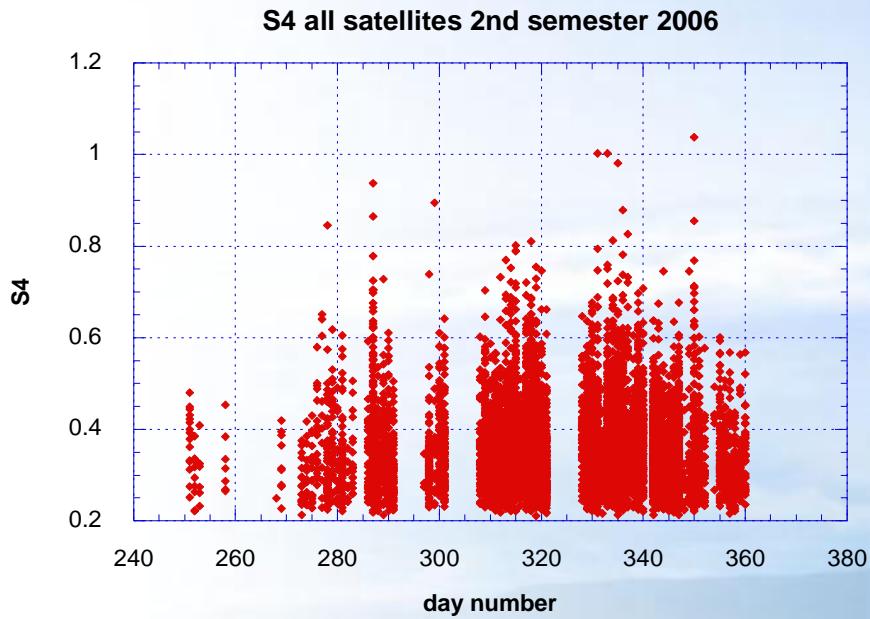
Phase

Observations at Low Latitudes

Béniguel Y., J-P Adam, N. Jakowski, T. Noack, V. Wilken, J-J Valette, M. Cueto, A. Bourdillon, P. Lassudrie-Duchesne, B. Arbesser-Rastburg, Analysis of scintillation recorded during the PRIS measurement campaign, Radio Sci., Vol 44, (2009), doi:10.1029/2008RS004090

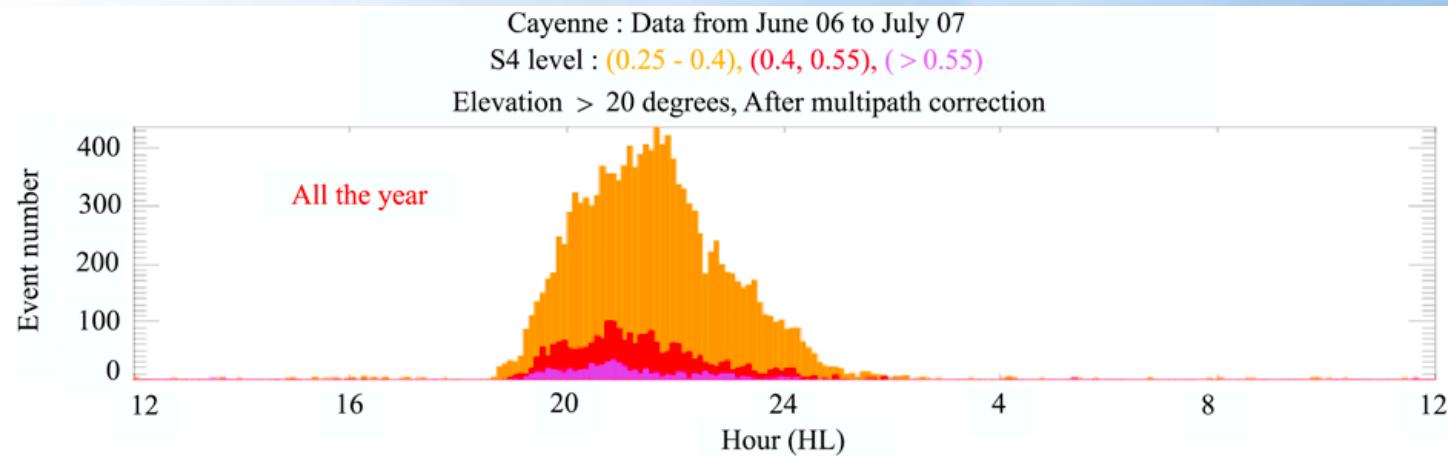
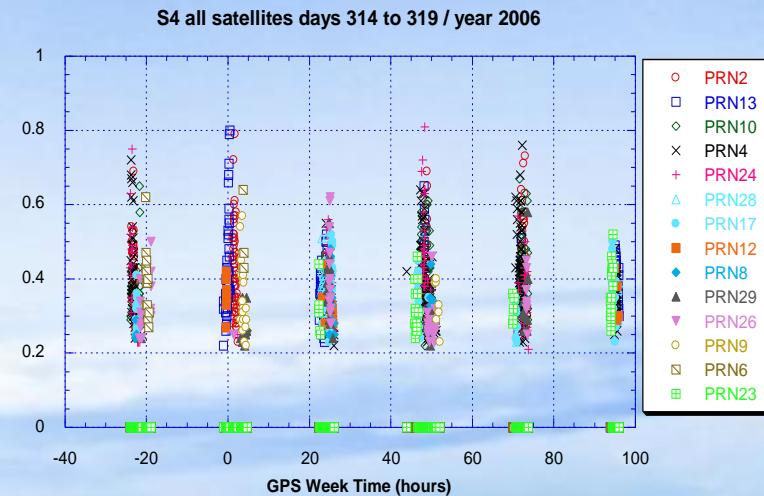
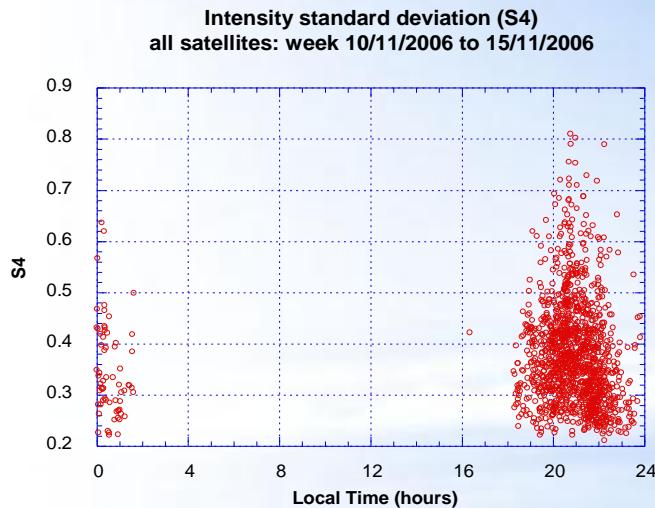
Prieto Cerdeira R., Y. Béniguel, "The MONITOR project: architecture, data and products", Ionospheric Effects Symposium, Alexandria VA, May 2011

Seasonal Dependency



Measurements in Guiana (4 N)

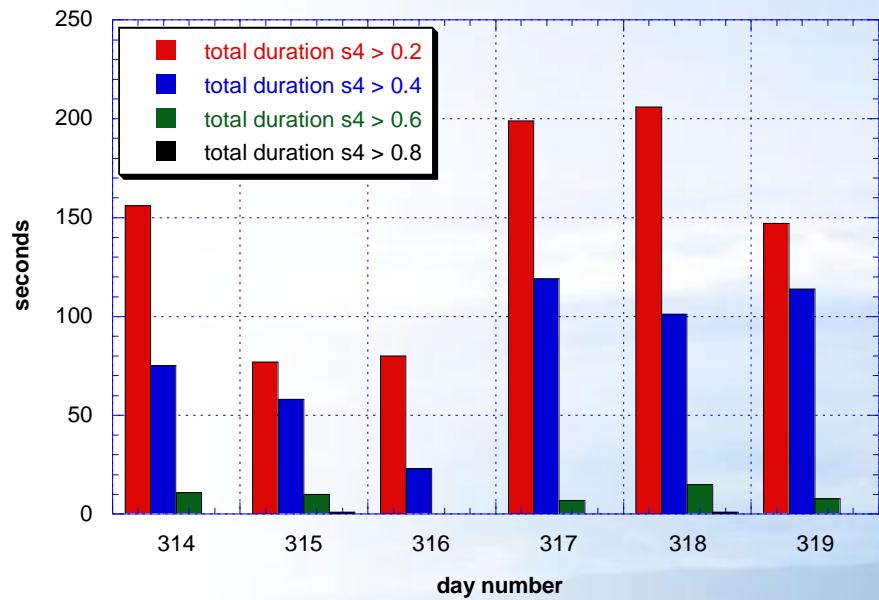
One week of measurements close to the equator



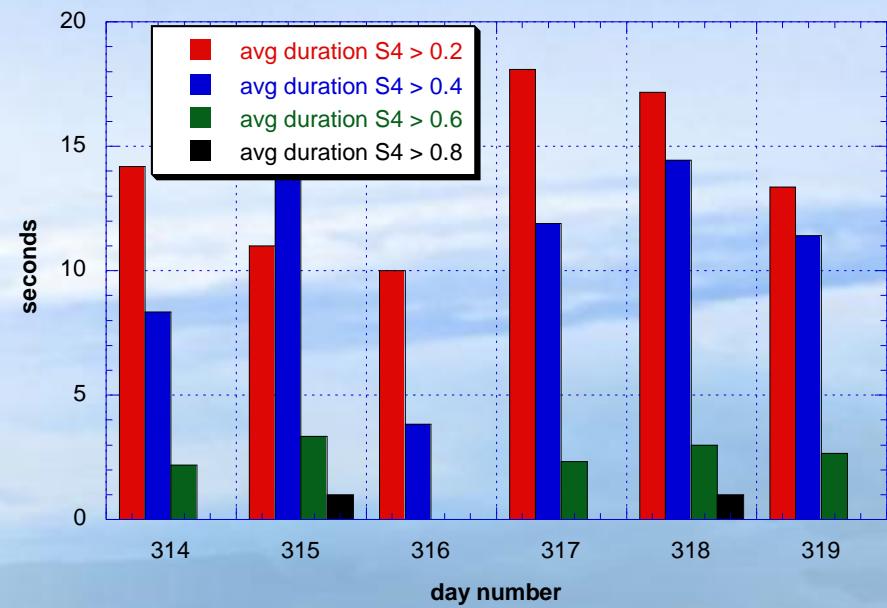
Local time : post sunset hours (CLS measurements, PRIS Campaign)

Fades Statistics

fades durations cumulated per day



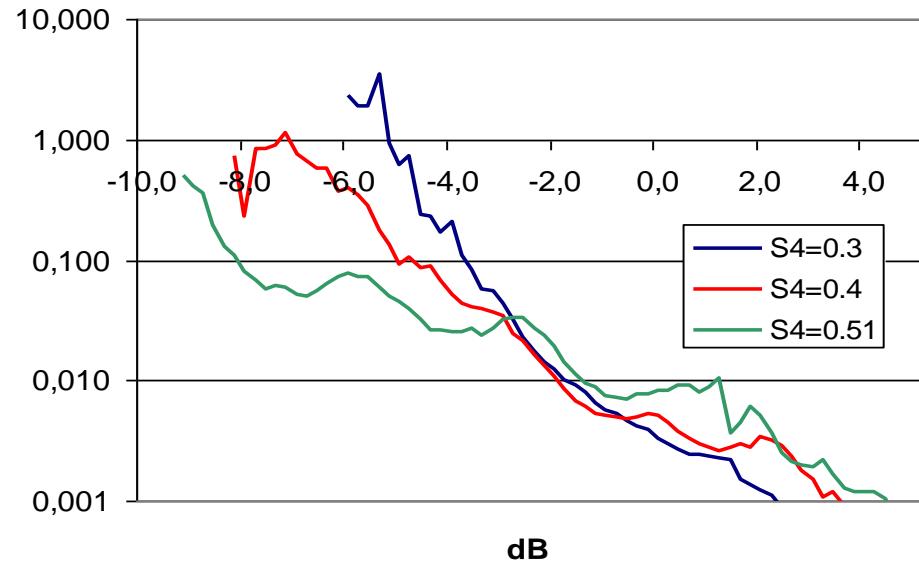
average durations per day



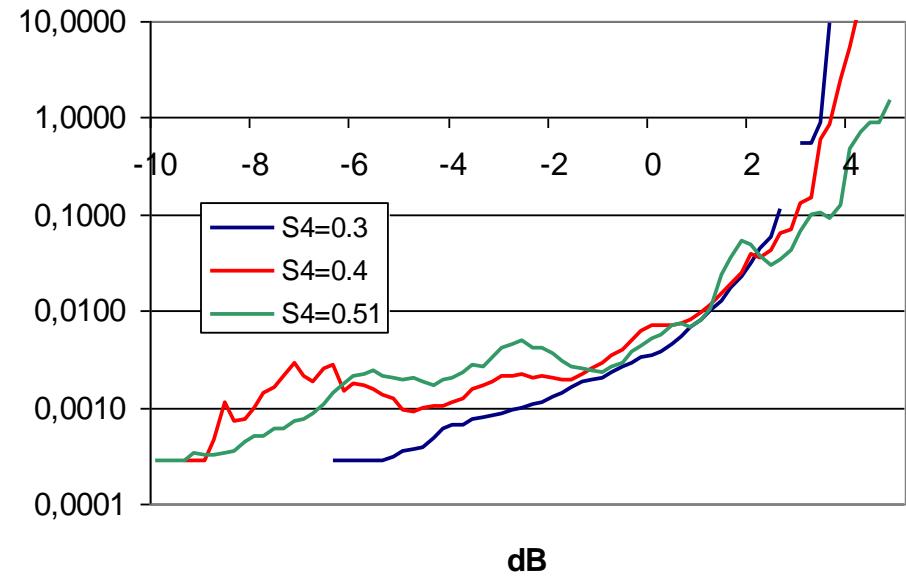
One week of measurements / november 2006

Fades Statistics

Time between fades

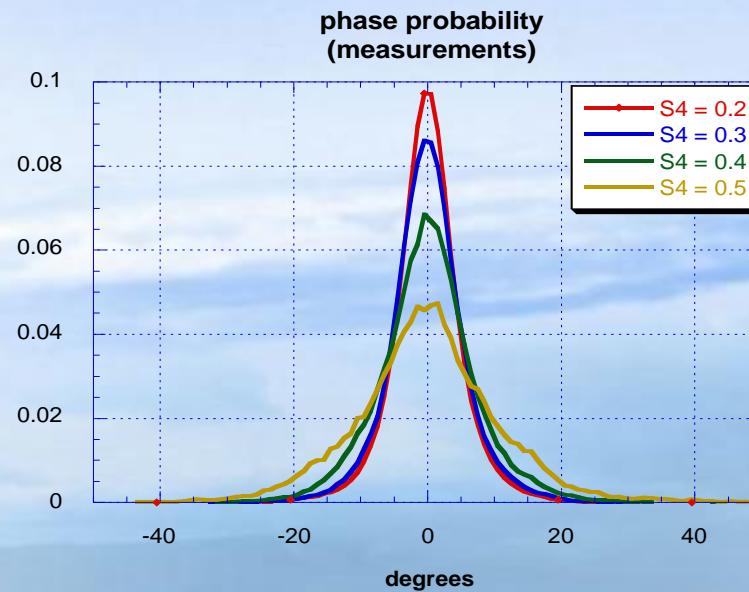
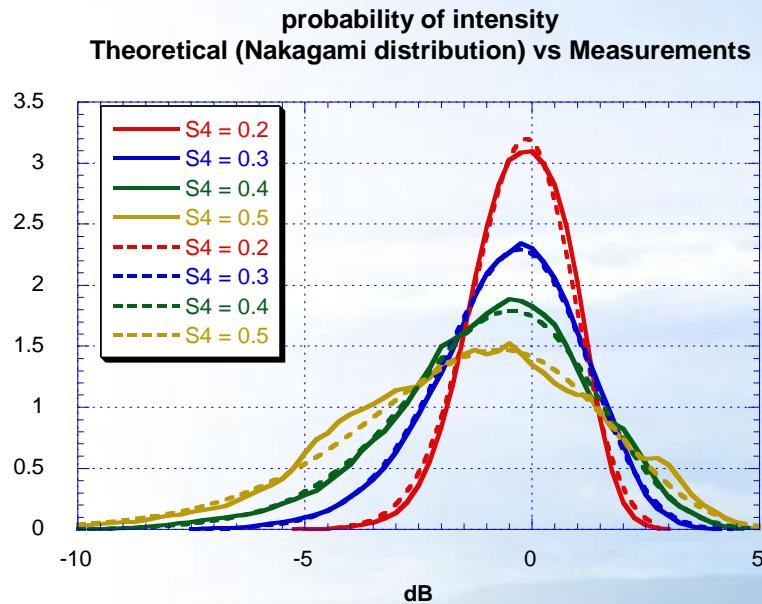


Fades duration



Intensity and Phase Distributions

one week of measurements



$$p(A) = \frac{2 m^m A^{2m-1}}{\Gamma(m)} \exp(-m A^2)$$

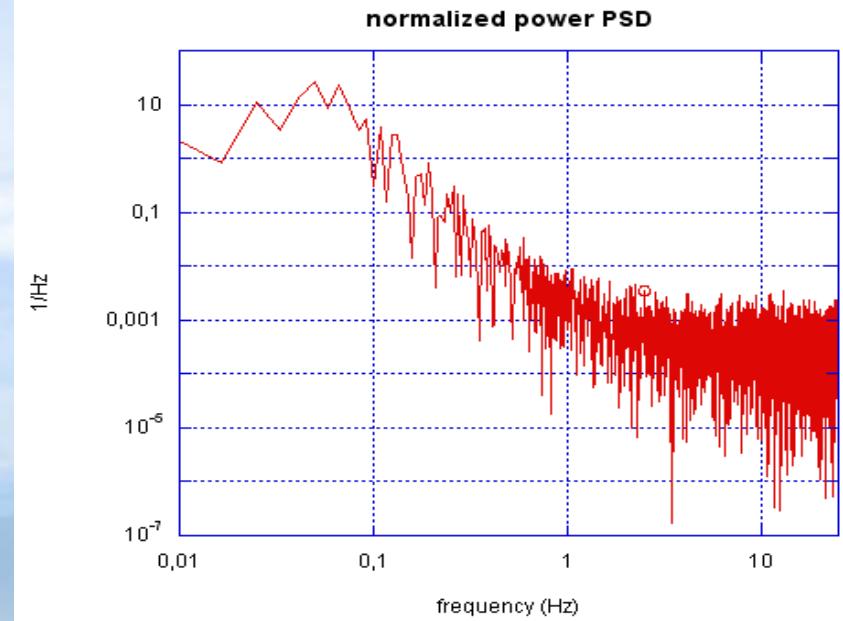
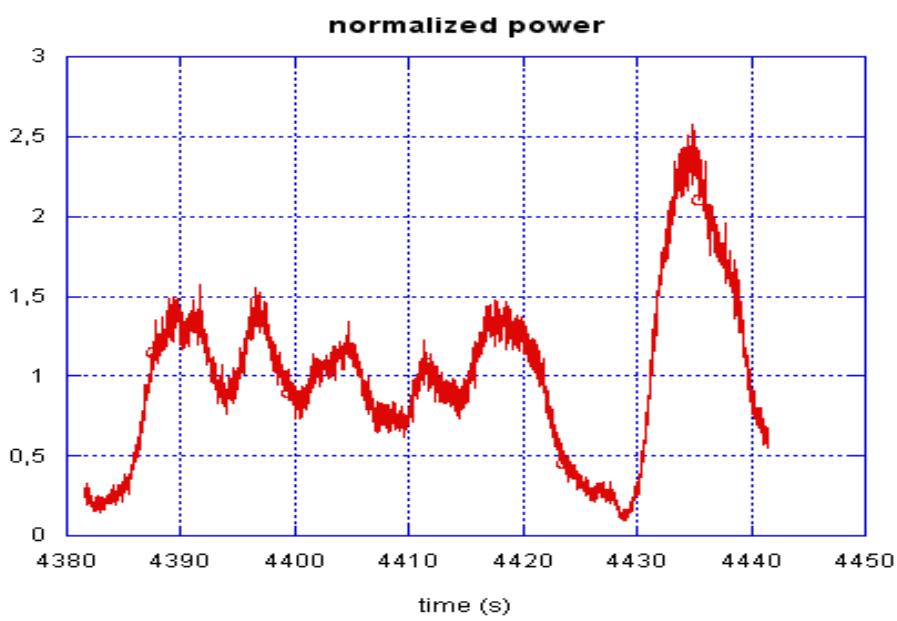
Gaussian distribution

with $m = 1 / S_4^2$

Spectrum Analysis

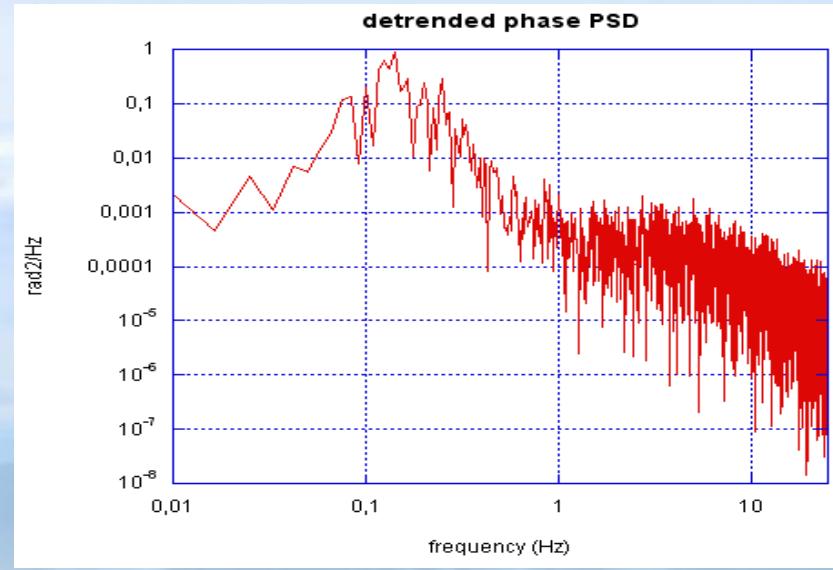
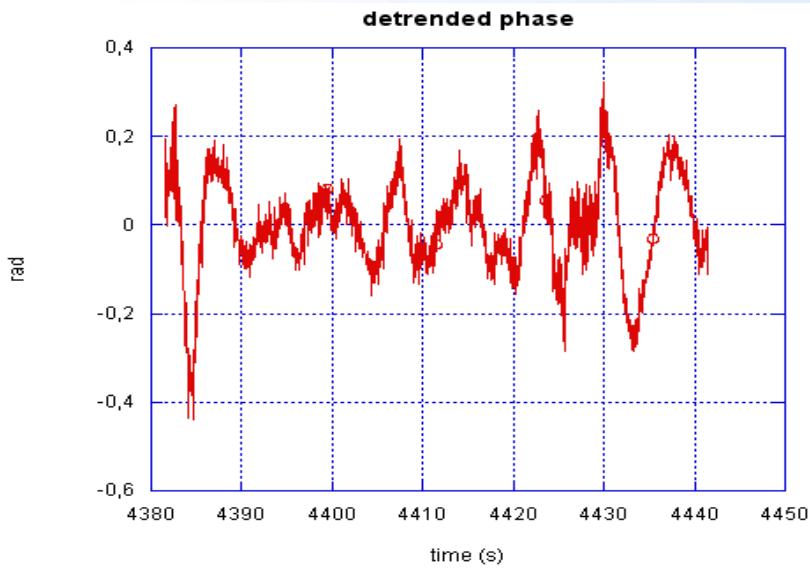
One Sample : Intensity

Sample characteristics : $S4 = 0.51$, $\sigma\phi = 0.11$



One Sample : Phase

Sample characteristics : $S_4 = 0.51$, $\sigma \phi = 0.11$

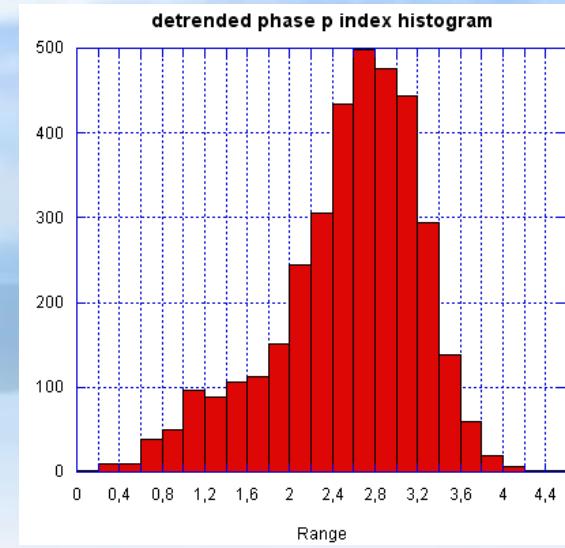
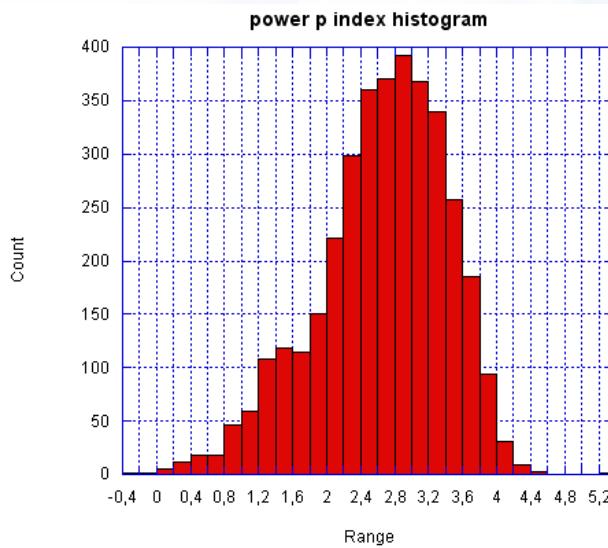


Spectrum Parameters

5 days RINEX files considered in the analysis

$S4 > 0.2$ & $\sigma\phi < 2$ (filter convergence)

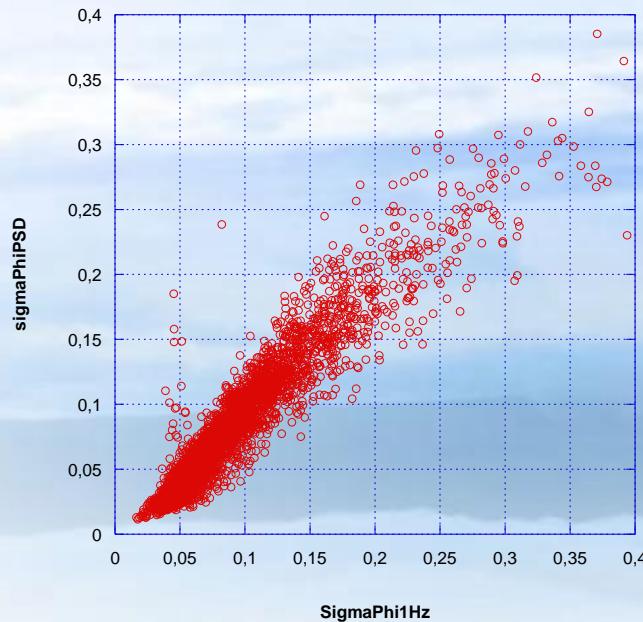
2 parameters to define the spectrum : T (1 Hz value) & p



Phase Variance : Time Domain vs Frequency Domain

$$\sigma\phi_i^2 = 2 \int_{fc}^{\infty} PSD(f) df = 2 \int_{fc}^{\infty} Tf^{-p} df = 2T \left[\frac{f^{-p+1}}{-p+1} \right]_{fc}^{\infty} = \frac{2T}{(p-1)fc^{p-1}} \quad (\text{if } p > 1)$$

Slope set to 2.8

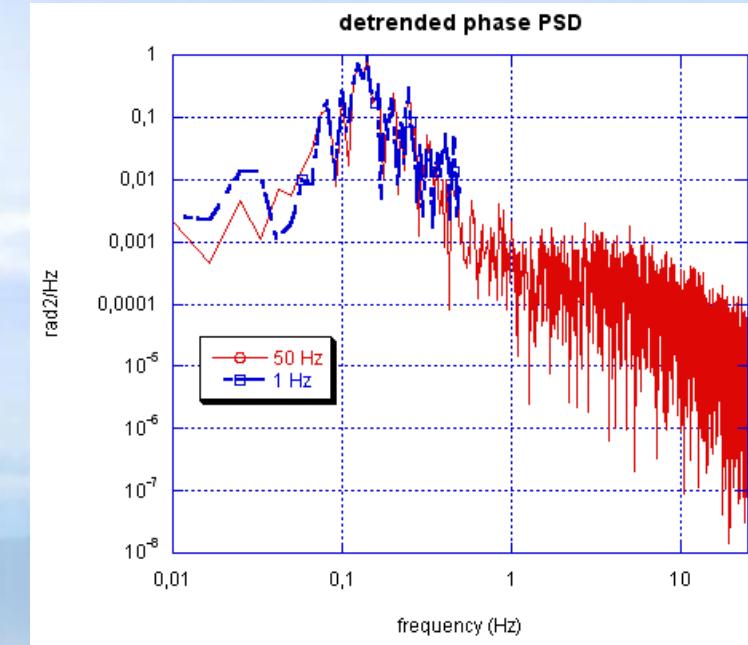
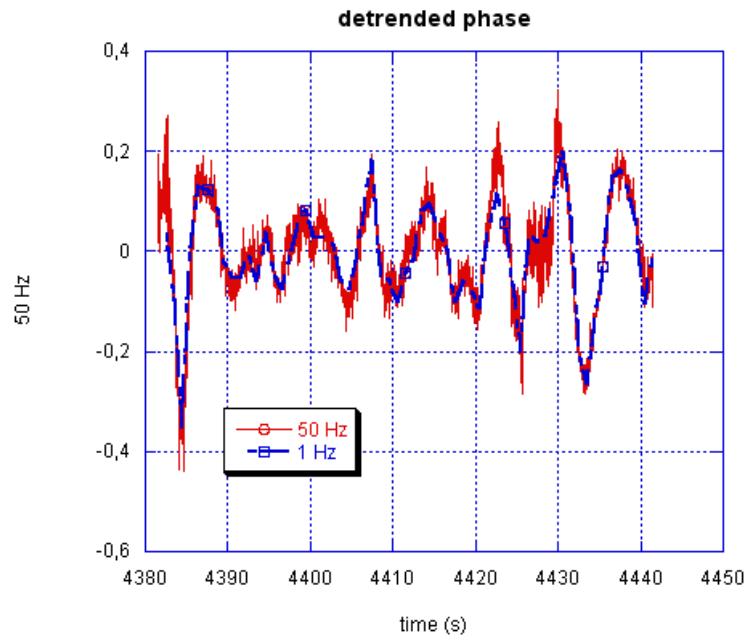


linear relationship

Under Sampling

1 Hz instead of 50 Hz

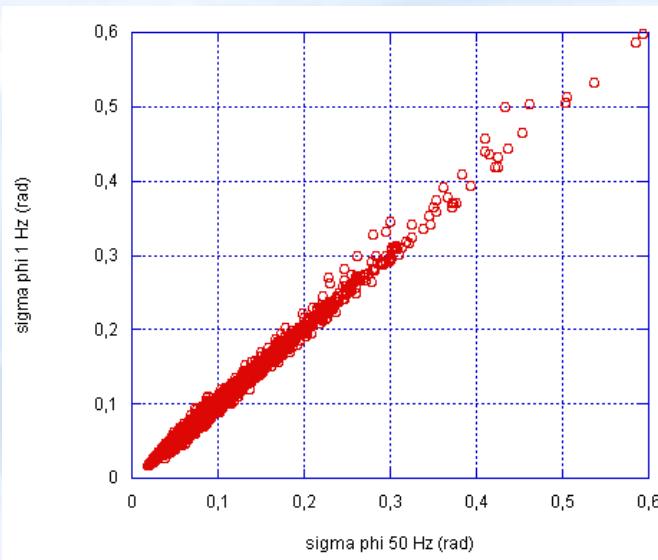
Cut off frequency → 0.5 Hz



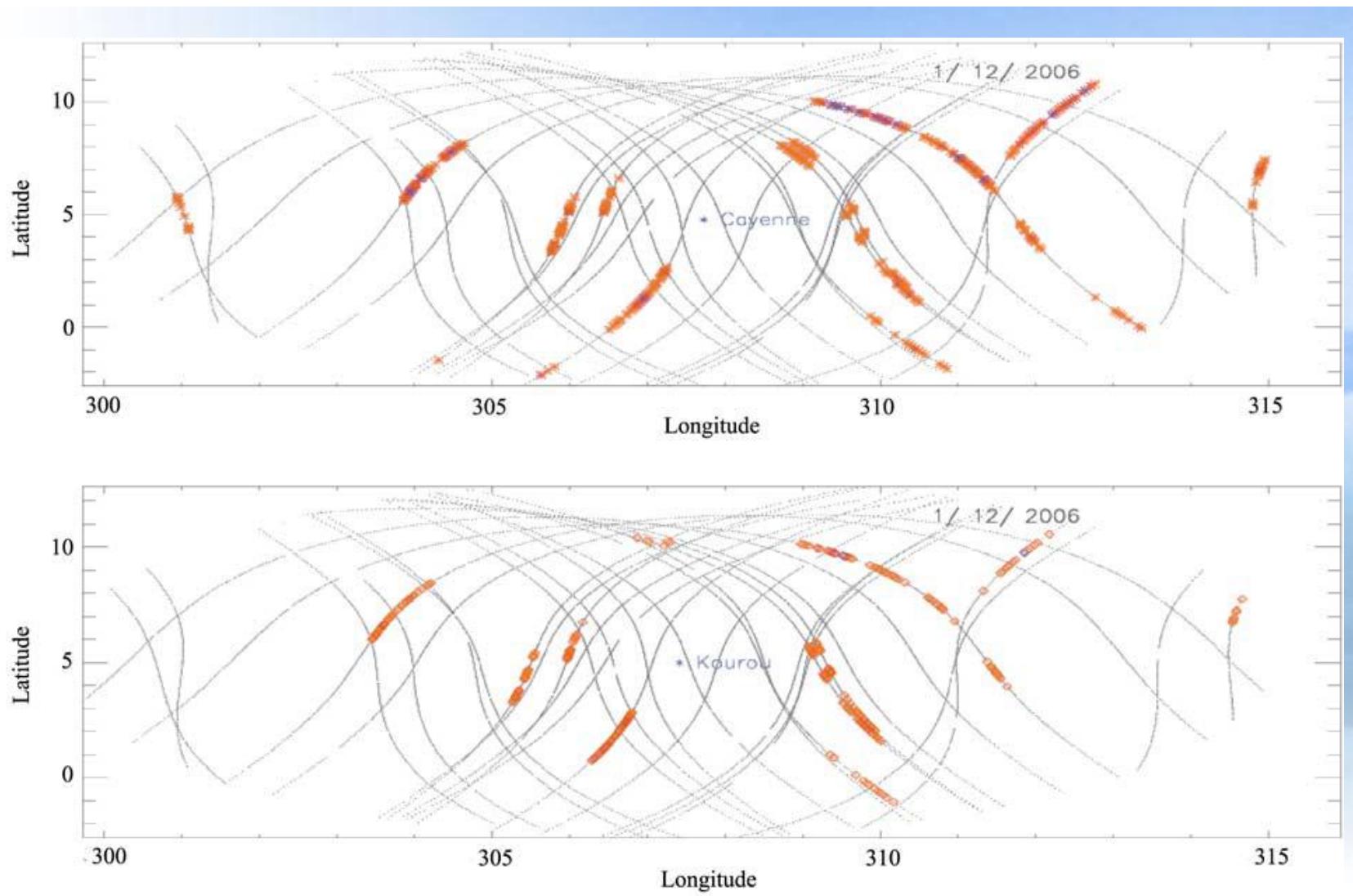
$$\sigma \phi i^2 = 2T \int_{0.1Hz}^{\infty} f^{-p} df \cong 2T \int_{0.1Hz}^{25Hz} f^{-p} df = 2T \left(\int_{0.1Hz}^{0.5Hz} f^{-p} df + \int_{0.5Hz}^{25Hz} f^{-p} df \right)$$

$$\int_{0.1Hz}^{0.5Hz} f^{-2.8} df = 33$$

$$\int_{0.5Hz}^{25Hz} f^{-2.8} df = 1.9$$



Scintillation indices at 50 Hz (top panel) vs 1 Hz (IGS data: bottom panel)



Modelling

Global Ionospheric Scintillation Model (GISM)

<http://www.ieea.fr/en/gism-web-interface.html>

Field received at ground level

Solution of the parabolic equation

Field Propagation Equation

Solution of the parabolic equation

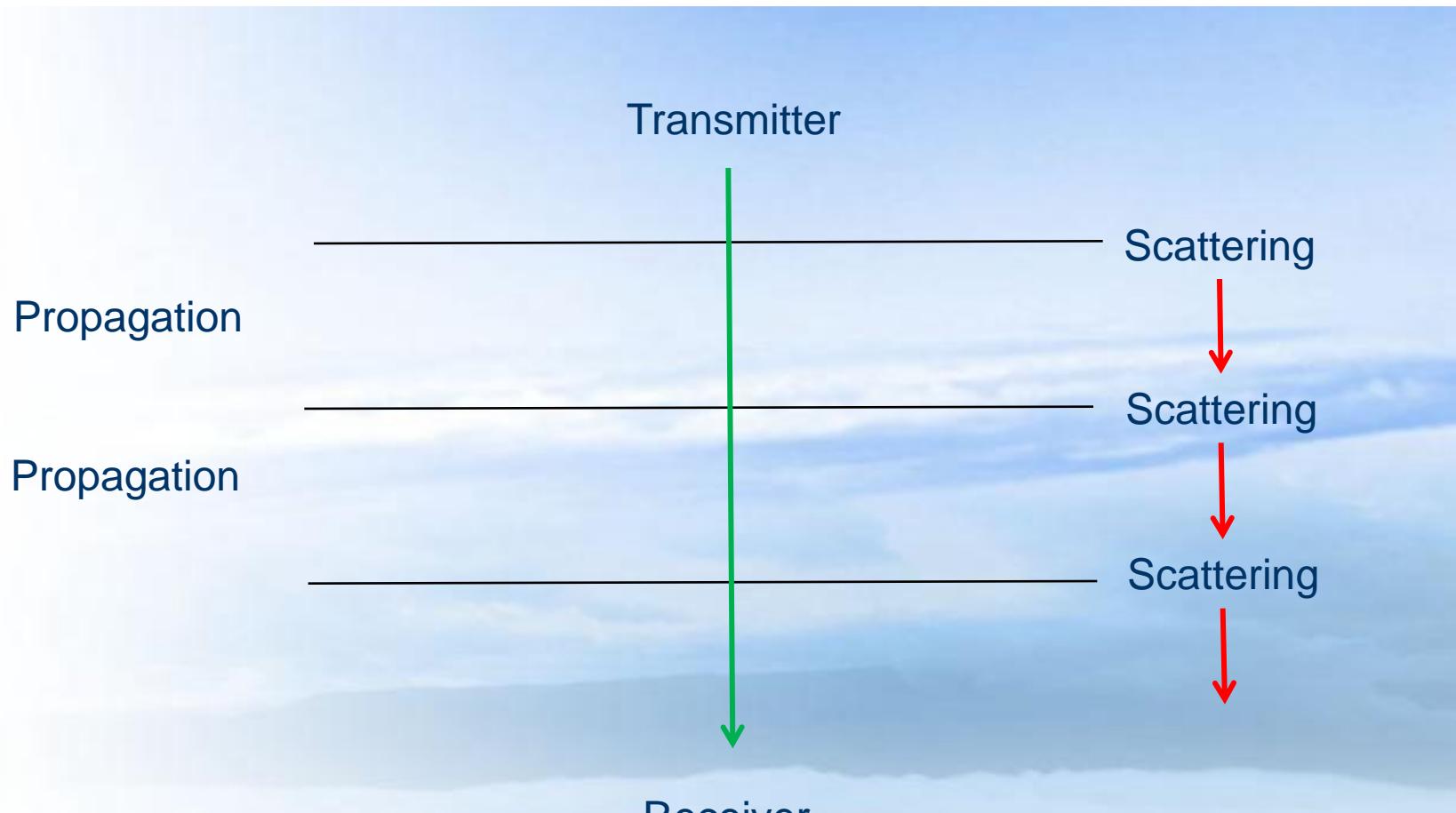
$$2jk \frac{\partial}{\partial z} \langle U(r) \rangle + \nabla_t^2 \langle U(r) \rangle + k^2 \langle \epsilon(r) U(r) \rangle = 0$$

$$2jk \frac{\partial}{\partial z} \langle U(r) \rangle + \nabla_t^2 \langle U(r) \rangle + j \frac{k^3}{4} A(0) \langle U(r) \rangle = 0$$

Using the phase index autocorrelation function

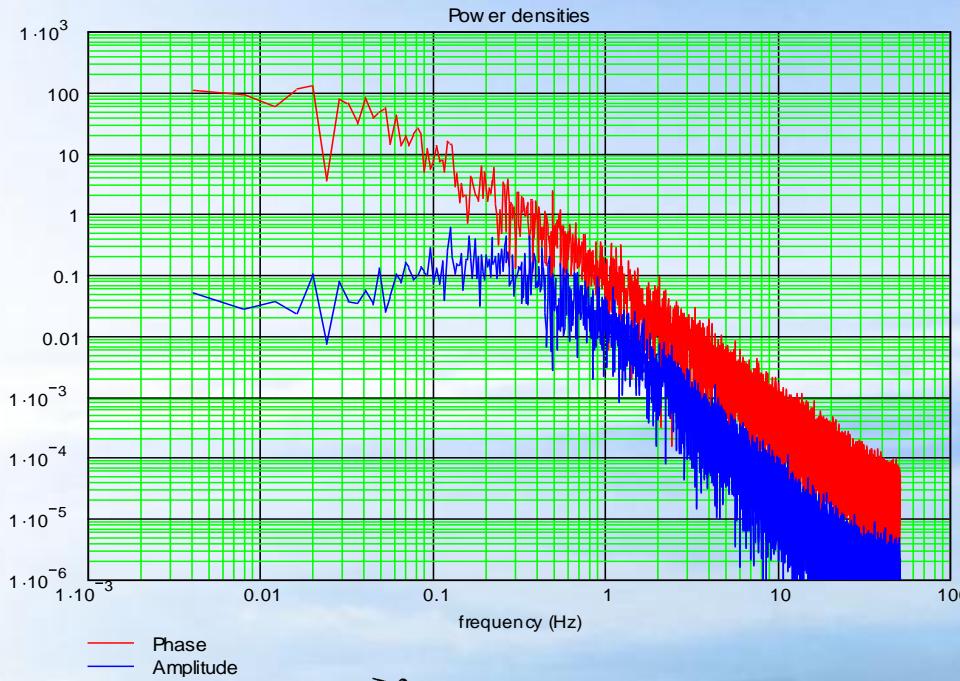
$$B(z, \rho) = \langle \epsilon(\rho_1) \epsilon(\rho_2) \rangle \quad A(\rho) = \int B(z, \rho) dz$$

Phase Screen Technique*



* Béniguel Y., P. Hamel, "A Global Ionosphere Scintillation Propagation Model for Equatorial Regions", Journal of Space Weather Space Climate, 1, (2011), doi: 10.1051/swsc/2011004

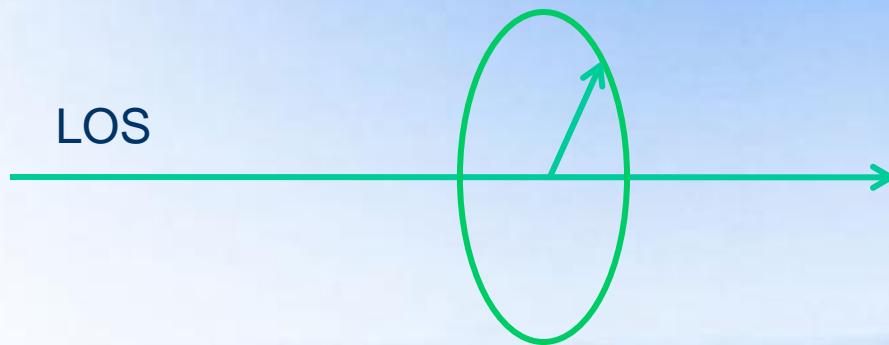
Medium's Characterisation*



$$\gamma_{\Phi}(K) = \frac{\mu r_e^{-2} L C_s \sigma_{Ne}^2}{(K^2 + q_0^2)^{1/2}} = \frac{C_p}{(K^2 + q_0^2)^{1/2}}$$

3 parameters : σ_{Ne} ; q_0 ; p

2D Analysis : Isotropic transverse medium



$$B_\Phi(\rho) = \frac{C_p}{4\pi} \iint \gamma_\Phi(K) \exp(-j\bar{K} \cdot \bar{\rho}) dK$$

→ $B_\Phi(\rho)_{\text{iso}} = \frac{\sigma_\Phi^2}{2^{(p-4)/2} \Gamma(p-2)/2} \rho q_0^{-((p-2)/2)} K_{((p-2)/2)} \rho q_0$

$$\sigma_\Phi^2 = B_\Phi(0) = \lambda r_e^{-2} L L_0 \sigma_{Ne}^2$$

1D Analysis : Isotropic transverse medium



$$B_{\Phi}(\rho) = \frac{C_p}{2\pi} \int \gamma_{\Phi}(k) \exp(-jk\rho) dk$$

→ $B_{\Phi}(\rho) = \frac{\sigma_{\Phi}^2}{2^{(p-3)/2} \Gamma(p-1)/2} \rho q_0^{-((p-1)/2)} K_{((p-1)/2)} \rho q_0^{-}$

$$\sigma_{\Phi}^2 = B_{\Phi}(0) = \lambda r_e^{-2} L L_0 \sigma_{Ne}^2$$

1D vs Isotropic (2D)

$$\mathbb{B}_\Phi(\rho)_{\text{is}} = \frac{\sigma_\Phi^2}{2^{(p-4)/2} \Gamma(p-2)/2} \langle \rho q_0 \rangle^{(p-2)/2} K_{((p-2)/2)} \langle \rho q_0 \rangle$$

$$\mathbb{B}_\Phi(\rho)_{\text{Ib}} = \frac{\sigma_\Phi^2}{2^{(p-3)/2} \Gamma(p-1)/2} \langle \rho q_0 \rangle^{(p-1)/2} K_{((p-1)/2)} \langle \rho q_0 \rangle$$

Slope $p \rightarrow p - 1$

Anisotropic vs Isotropic*



$$\gamma_{\Phi}(K) = \frac{4\pi r_e L C_s \sigma_{Ne}^2 a b}{(q^2 + q_0^2)^{1/2}}$$

$$\gamma_{\Phi}(K) = \frac{a b C_p}{((A K_{x\perp}^2 + B K_{x\perp} q_{y\perp} + C K_{y\perp}^2) + q_0^2)^{1/2}}$$

Additional geometric factor with respect to the 2D case

$$G = \frac{ab}{(AC - B^2/4)^{1/2}}$$

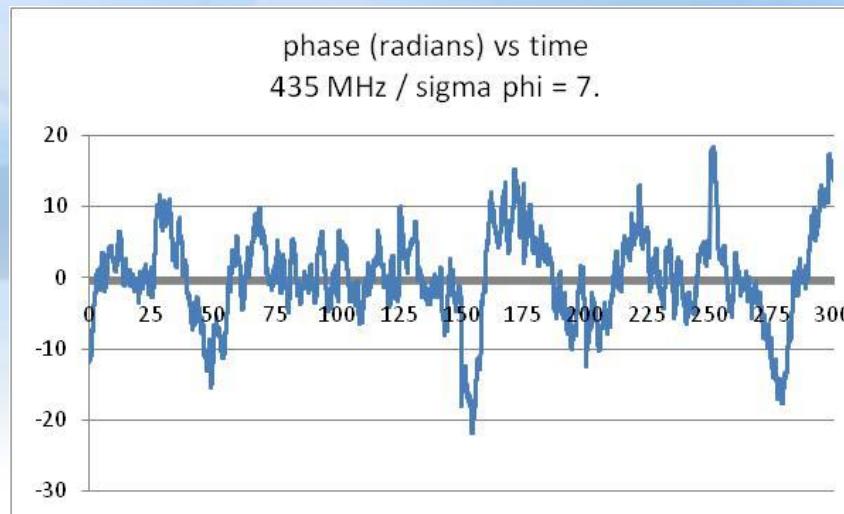
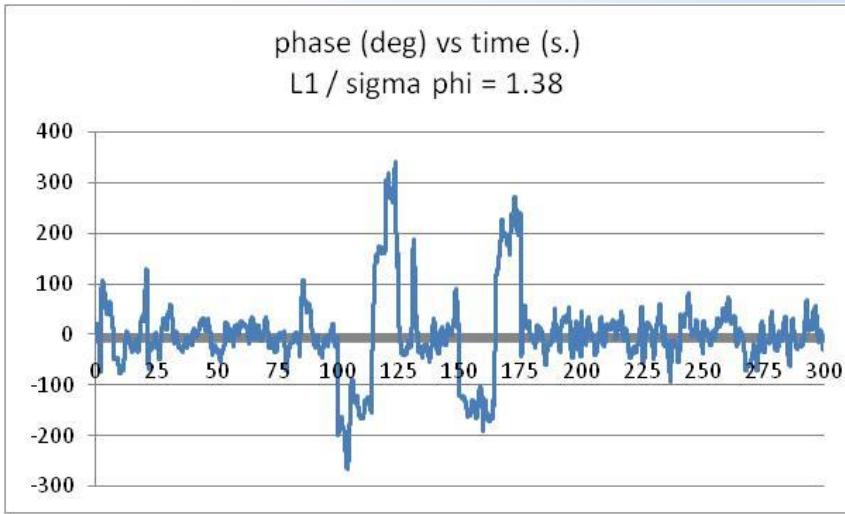
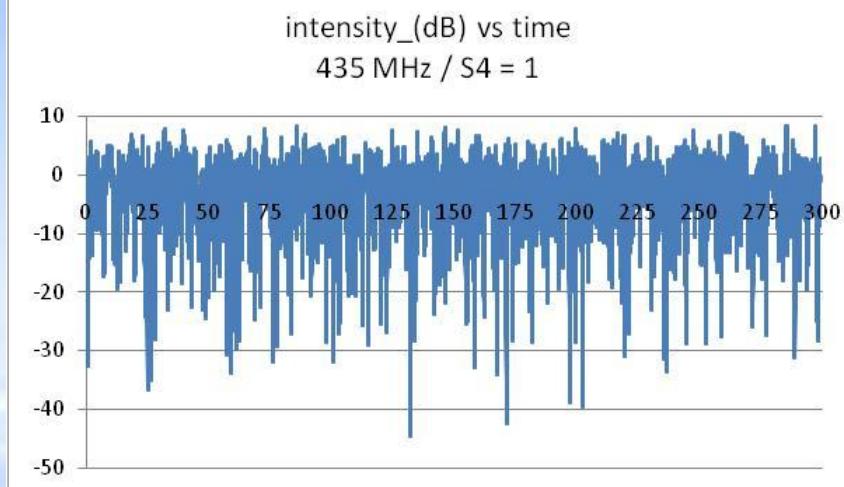
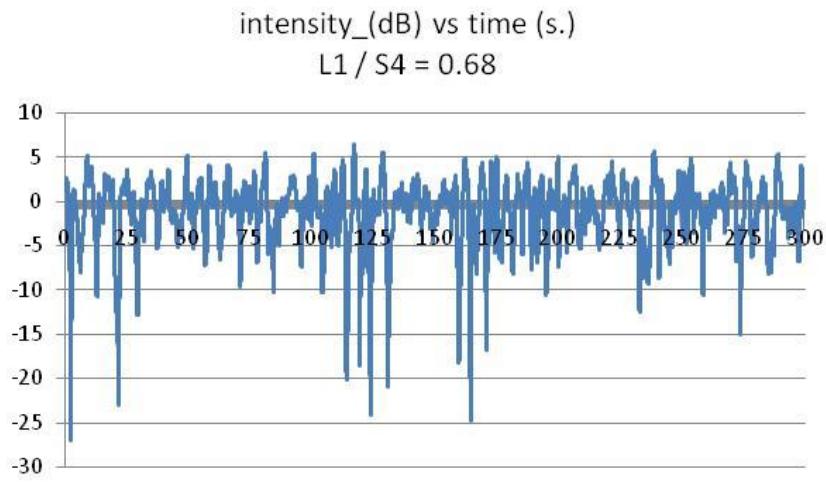
a, b ellipses axes

A, B, C trigonometric terms resulting from rotations related to variable changes

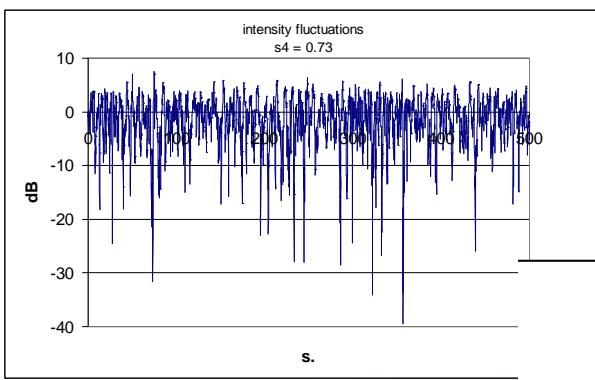
* Rino, Radio Sci. 1979

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Signal at receiver level

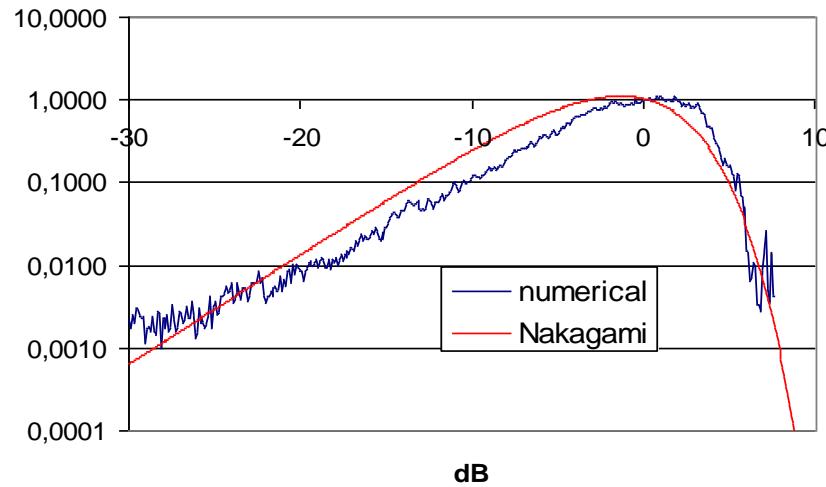


Probability of intensity Modelling



GISM output

probability of intensity fluctuations
 $s4 = 0.73$



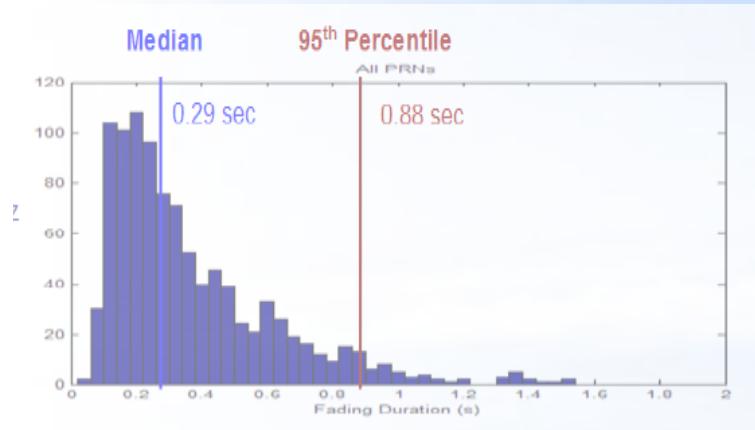
Nakagami law

$$p(A) = \frac{2 m^m A^{2m-1}}{\Gamma(m)} \exp(-m A^2) \quad \text{with}$$

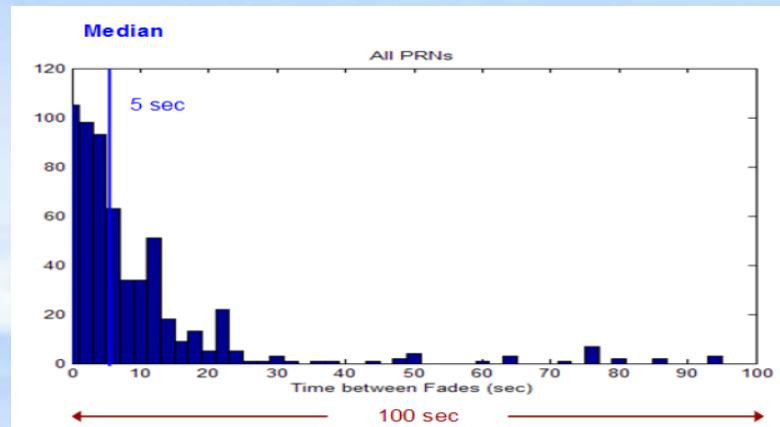
$$m = 1 / s_4^2$$

Fades Statistics

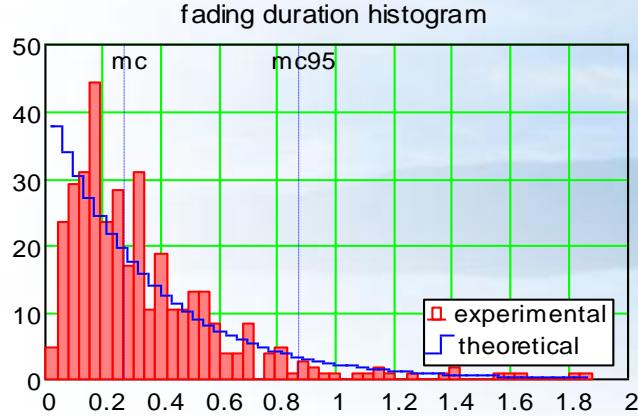
Example of equatorial scintillation in Ascension Island, in solarmax conditions (2001)



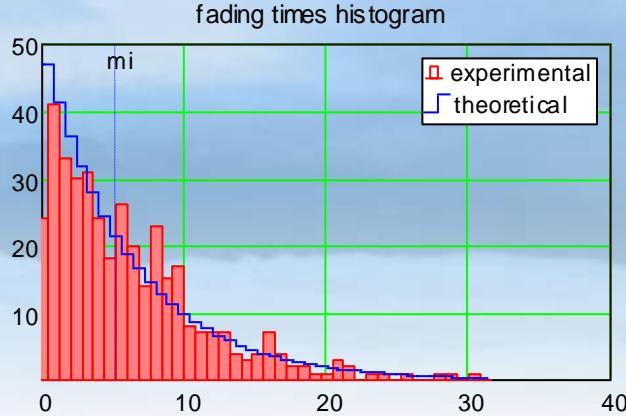
•Real data



fading times histogram



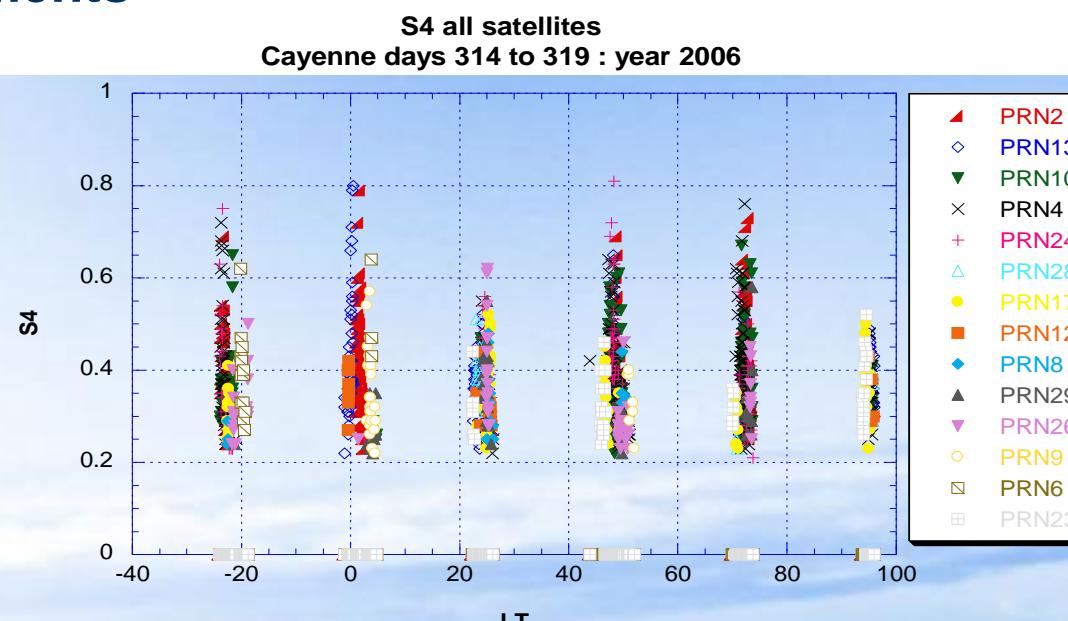
•GISM simulation



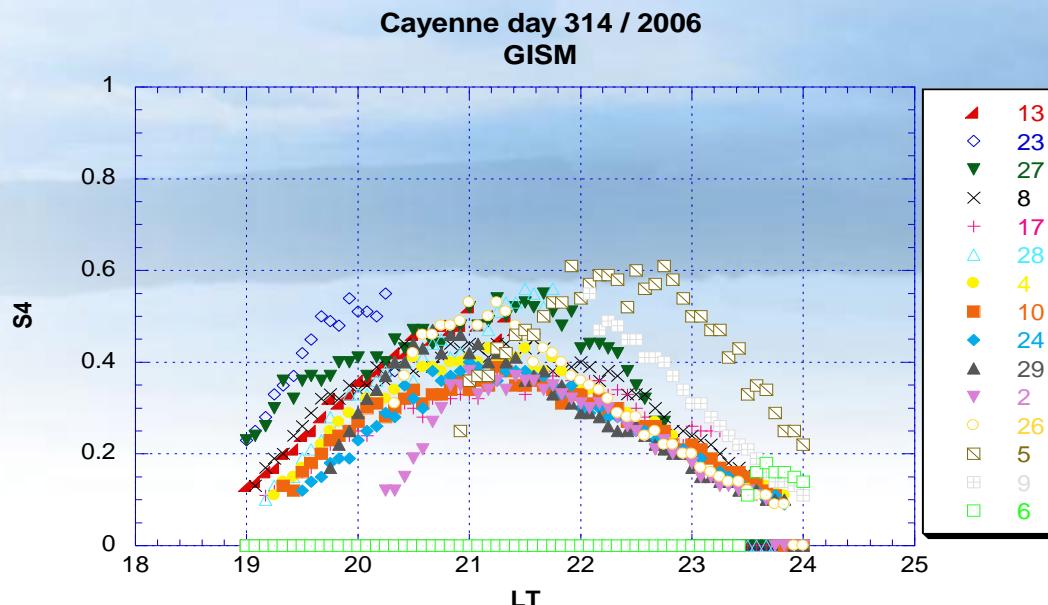
Modelling vs Measurements (Intensity)

Measurements →

Samples with $S4 < 0.2$
were ignored (noise level)



Modelling →



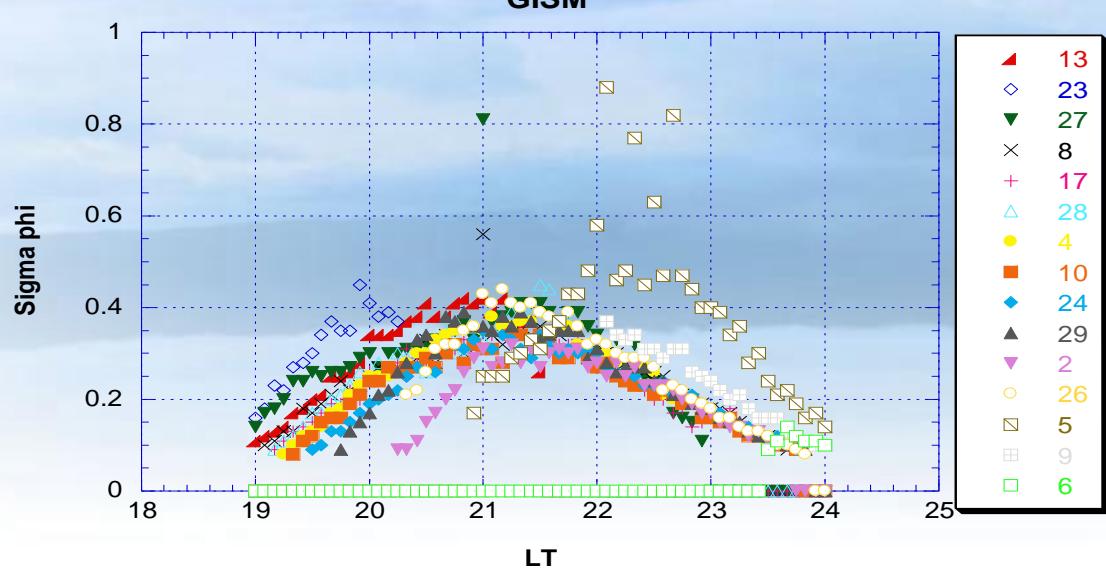
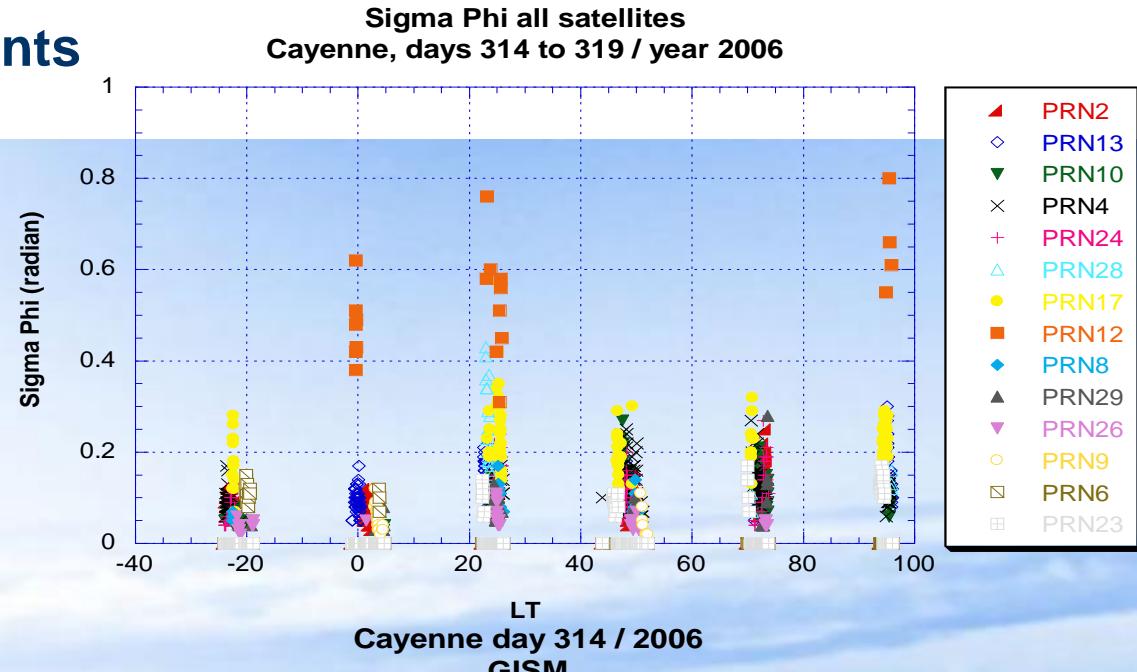
Modelling vs Measurements (Phase)

Measurements →

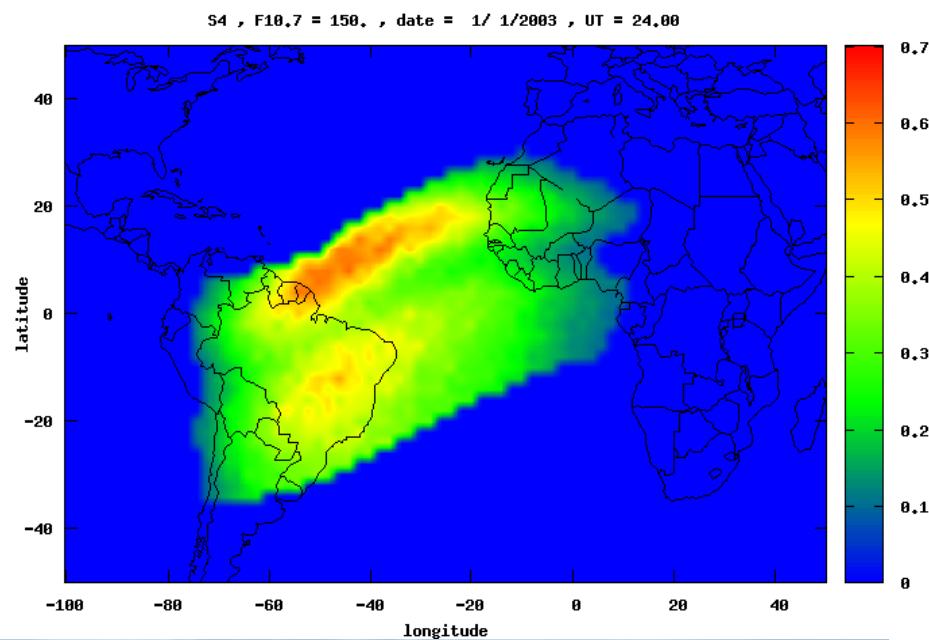
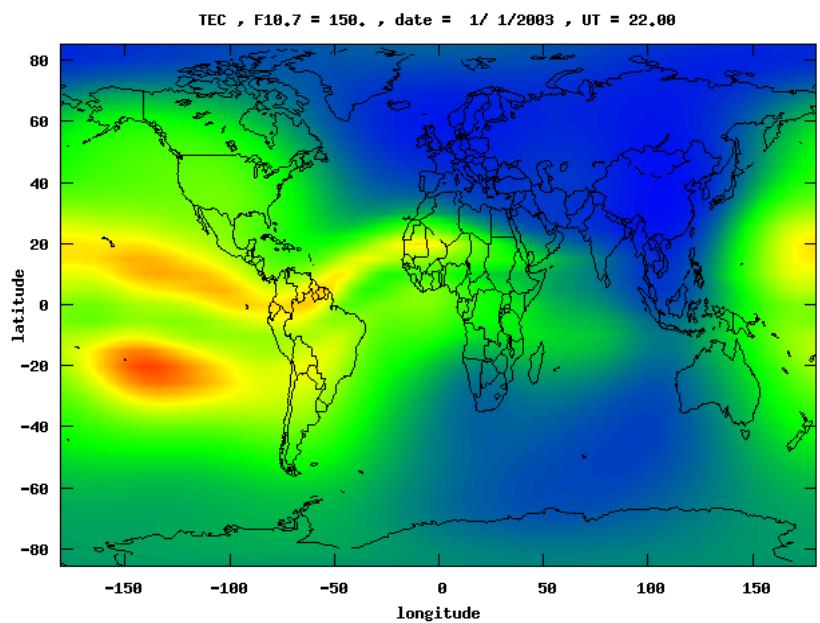
The phase RMS value is slightly lower than the S4 value

Some samples exhibit high values (both measurements and modelling) due to the phase jumps

Modelling →



Global Maps



TEC Map Modelling

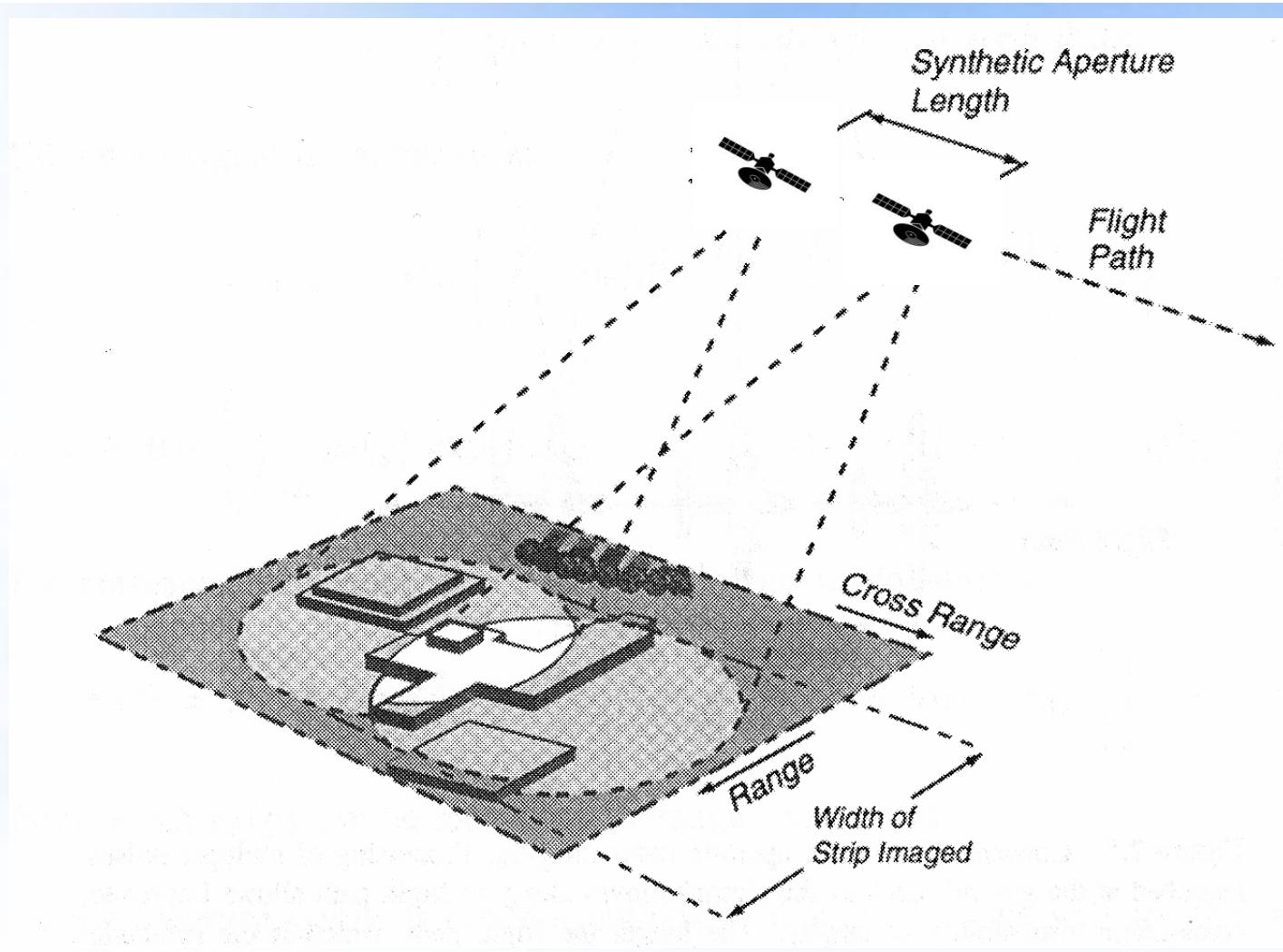
Scintillation Map Modelling



Radar Observations

Mutual Coherence Function

Correlation distance vs L_{SAR}



L_e Synthetic Aperture Length → 10 km

Ionosphere Effects

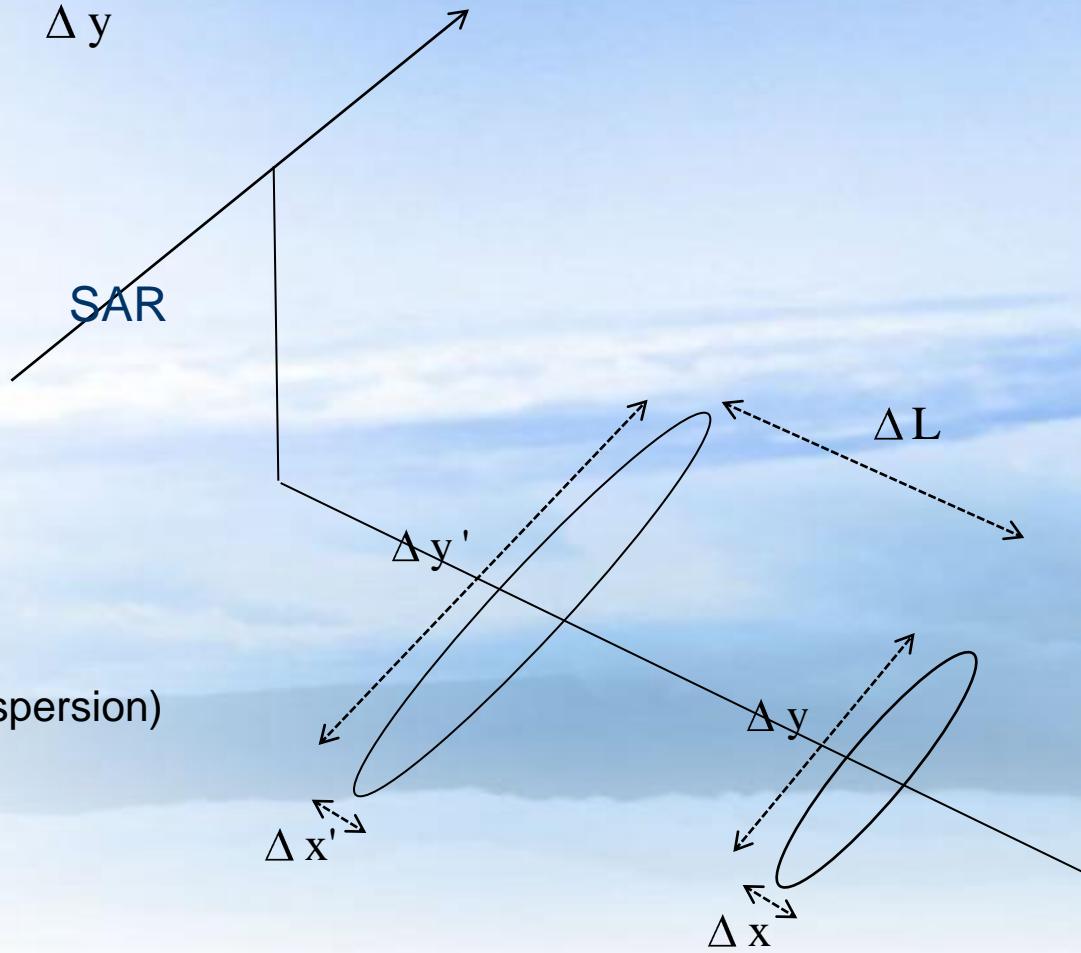
Free Space Resolution Δx Δy

Ionosphere Effects

$\Delta x'$ Pulse broadening (dispersion)

$\Delta y'$ Turbulence Effect

ΔL Group Delay



Two Points - Two Frequencies Coherence Function

$$\Gamma(z, k_1, k_2, \rho_1, \rho_2) = \langle U_1(z, k_1, \rho_1) U_2^*(z, k_2, \rho_2) \rangle$$

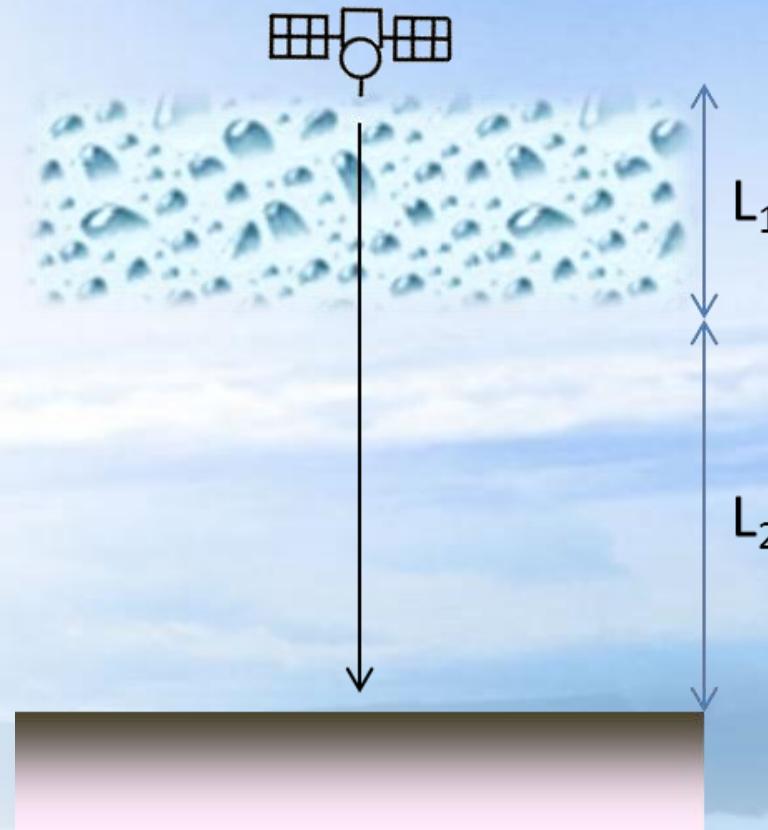
Using the parabolic equation

$$\left[\frac{\partial}{\partial z} - \frac{j}{2} \frac{k_d}{k_0^2} \nabla_d^2 + \frac{k_p^4}{8 k_0^2} \left[\frac{k_d^2}{k_0^2} A_\xi(0) + D_\xi(\rho) \right] \right] \Gamma(k_d, z, \rho) = 0$$

The structure function $D_\Phi(z, \rho) = 2 [B_\Phi(0) - B_\Phi(\rho)]$ is quadratic with respect to the distance

Same process than previously : propagation 1st & 3rd terms ; diffraction : 2nd & 3rd terms

Two Points - Two Frequencies Coherence Function



$$\rightarrow \Gamma(z, k_1, k_2, \rho_1, \rho_2) = \langle U_1(z, k_1, \rho_1) U_2^*(z, k_2, \rho_2) \rangle$$

Solution (one single screen)

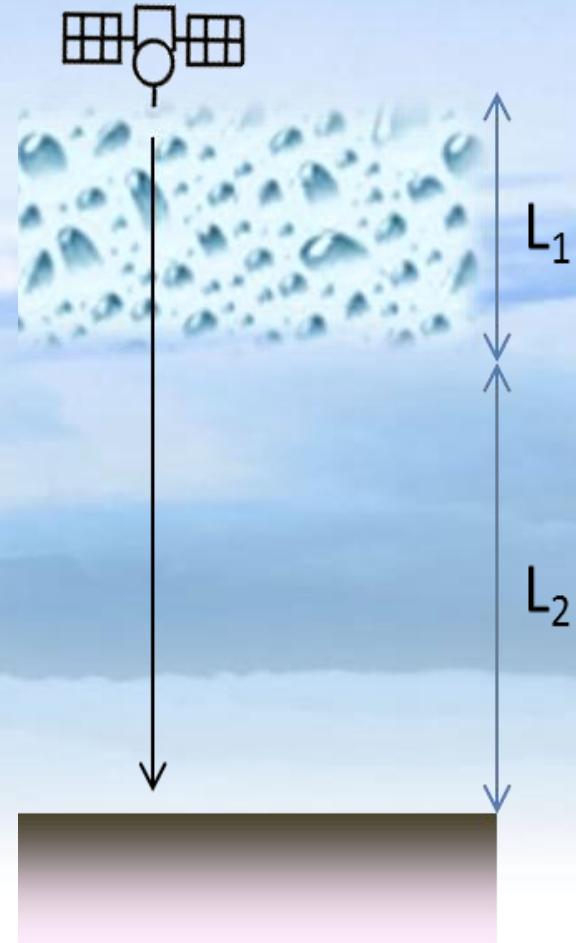
Scattering → 2 constants

$$B = \frac{\sigma_{\Phi}^2}{2 \omega^2}$$

$$S = \sigma_{\Phi}^2 \propto \frac{\log(L_0 / \ell_i)}{6 L_0^2}$$

Propagation → 1 constant

$$P = \frac{1}{2 c k^2} \left(\frac{1}{L_1 + L_2} - \frac{1}{L_1} \right)$$



* Nickisch, RS 92, Knepp & Nickisch, RS, 2010

One single Screen

Analytical Solution

$$\Gamma (\tau, K_x, z) = \frac{z}{2 \sqrt{BS}} \exp \left(- \frac{z^2 K_x^2}{4S} \right) \exp \left(- \frac{\epsilon + z^2 K_x^2 P}{4B} \right)$$

The equation is shown with two parts highlighted by arrows:

- A blue arrow points down from the term $\frac{z}{2 \sqrt{BS}}$ to the word "Maximum".
- A red arrow points up from the term $\exp \left(- \frac{z^2 K_x^2}{4S} \right)$.
- A yellow arrow points up from the term $\exp \left(- \frac{\epsilon + z^2 K_x^2 P}{4B} \right)$.

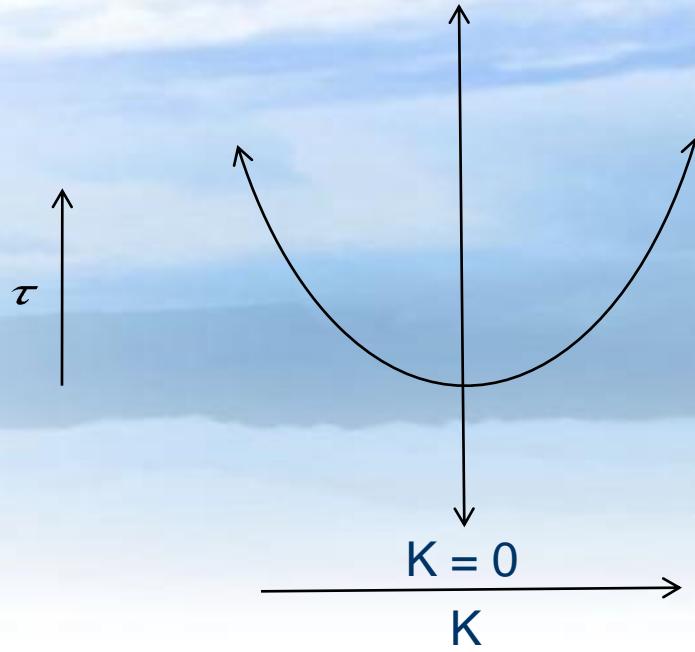
Spreading Extent

$$K = 0 \rightarrow \tau_1 = 2A \sqrt{B}$$

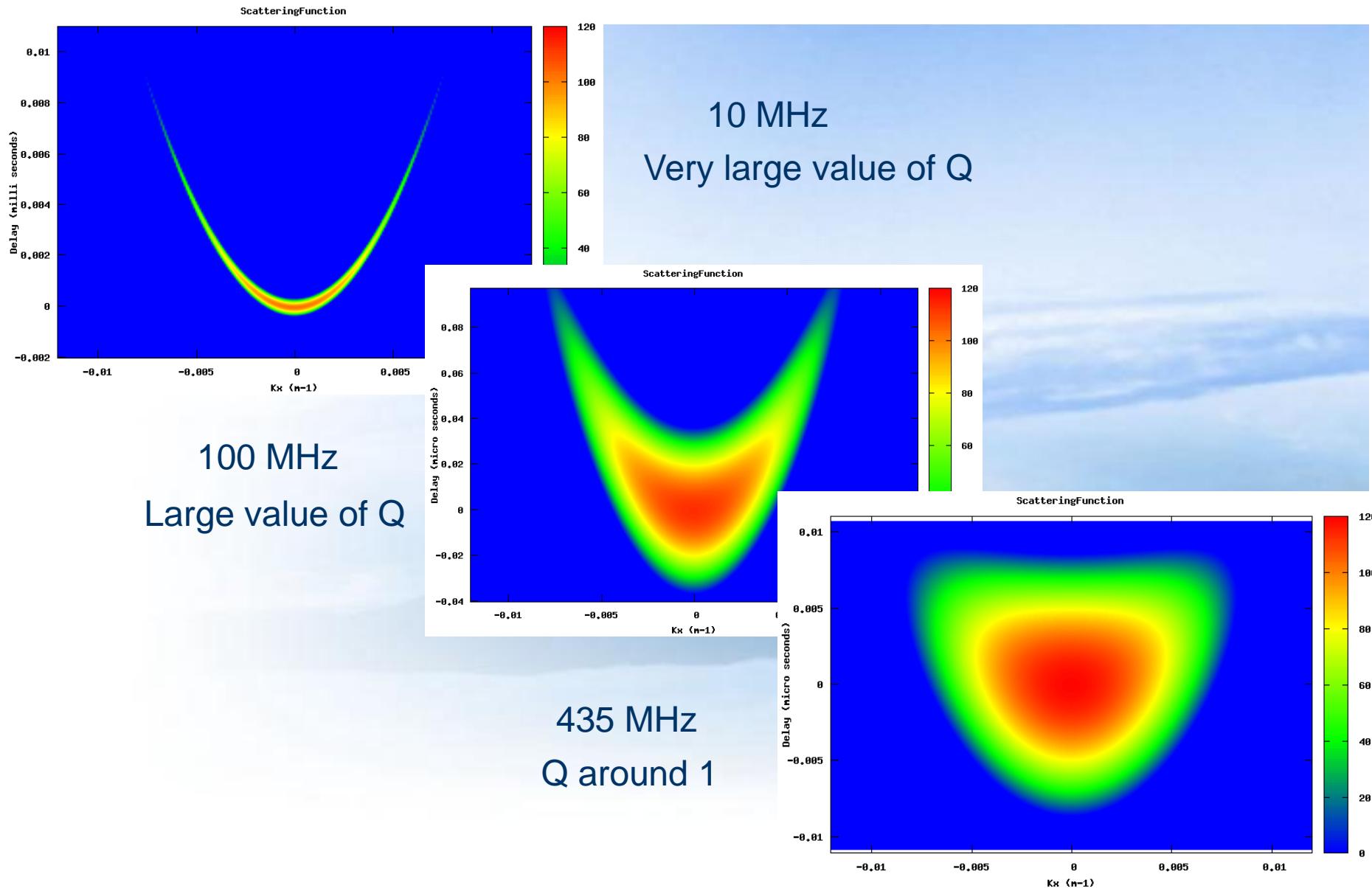
$$K = K_{\text{Max}} \rightarrow \tau_2 = 4 S P A$$

Spread factor

$$Q = \frac{\tau_2}{\tau_1} = \frac{2\sqrt{A} S P}{\sqrt{B}}$$



Spreading Extent



Several Screens / Refined Analysis

- The algorithm can easily be generalized
- Different statistical properties may be assigned to the different layers
- Numerical FFT (1D) shall be performed to get the coherence function

Ambiguity Function

$$\chi(r, r_0) = \sum_n \int g_n(t, r_n) f_n^*(t, r_{0n}) dt$$

g_n is the received signal and f_n is the matched filter

$$\chi_n(r_n, r_{0n}) = \frac{1}{(4\pi r_n)^2} \exp(j\Phi_0) \int \exp(j(\omega - \omega_0)\Phi_1 - (\omega - \omega_0)\Phi_2) d\omega$$

Value of Coherent field received

$$\langle \chi(r, r_0) \rangle = \frac{\exp(-\sigma_\Phi^2)}{(4\pi r_0)^2} \int_{-L_e/2}^{L_e/2} \exp\left(2jk_0\left(r_0 + \frac{\rho^2}{2r_0}\right)\right) d\rho = \frac{\exp(-\sigma_\Phi^2)}{(4\pi r_0)^2} \sin c(k_0 \rho L_e / r_0)$$

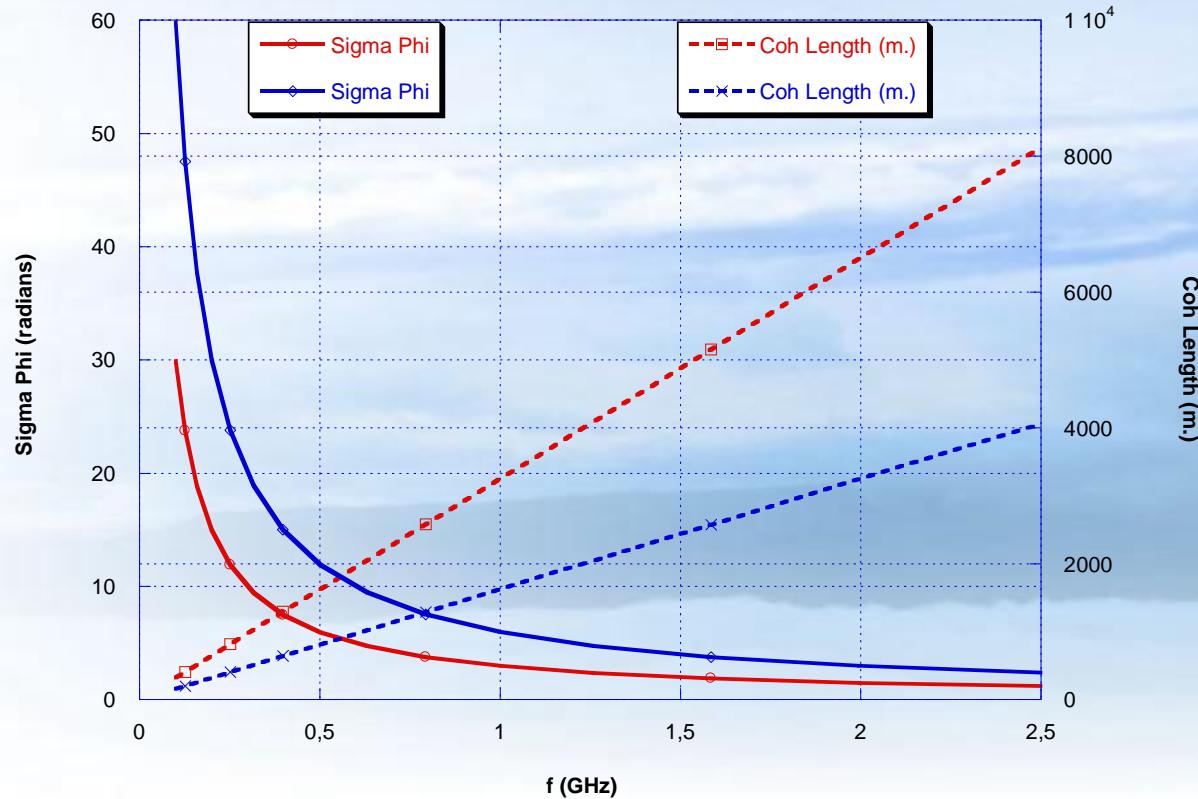
An attenuation factor on the coherent component has been included

Coherent Length

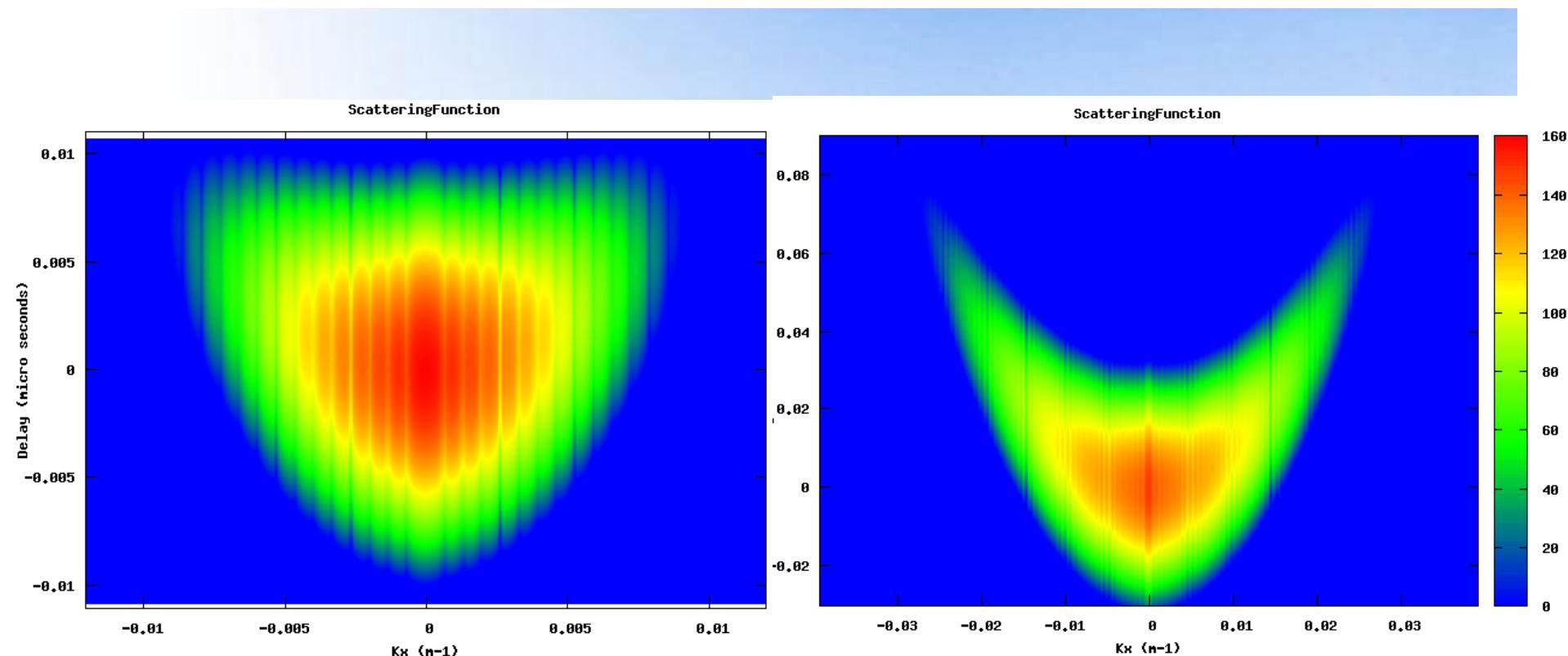
It is given by function

$$\Gamma(\omega_d, \rho, r) = \sqrt{\frac{D}{S}} \exp(-B\omega_d^2 - D(\rho/r_0))$$

Phase Standard Deviation and related Coherent Length
in the case of strong fluctuations (red lines) & very strong fluctuations (blue lines)



Scattering Function 400 MHz



Sigma Phi = 3.15

Sigma Phi = 10.

The Fourier transform (K_y space) of the antenna pattern is included in the calculation

Conclusion

- Reasonable agreement between the model (GISM) and the measurements
- The Mutual Coherence function has an analytical solution for one screen
- The algorithm can be extended for an arbitrary number of screens
- The azimuthal resolution of a SAR may be significantly decreased

Run GISM

<http://www.ieea.fr/en/gism-web-interface.html>