

One dimensional quantum systems

T. Giamarchi

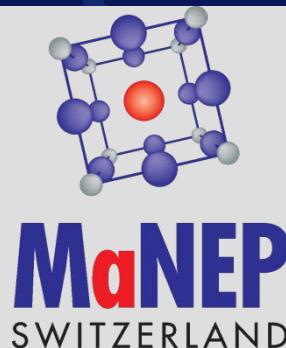
http://dpmc.unige.ch/gr_giamarchi/



**UNIVERSITÉ
DE GENÈVE**



FONDS NATIONAL SUISSE
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FONDO NAZIONALE SVIZZERO
SWISS NATIONAL SCIENCE FOUNDATION



C. Berthod

A. Lobos

A. Kantian

P. Bouillot

T. Ewart

A. Tokuno

A. Kozenkov

E. Agoritsas

T. Jarlborg

P. Chudzinski



Plan of the lectures (1)

■ Lecture 1:

- What are one dimensional systems
- Universal physics in one dimension (Luttinger liquid)
- Some realizations with cold atoms

Plan of the lectures (2)

- Lecture 2:
 - Effect of periodic potentials: Mott transition
 - Experimental tests of the Mot transition
 - Systems with internal degrees of freedom (spin)

Plan of the lectures (3)

■ Lecture 3:

- Disorder and 1D quantum systems (dirty bosons)
- Disorder and quasiperiodicity
- Loose ends and open questions

Lecture 1

References

TG, arXiv/0605472 (Salerno lectures)

TG, Quantum physics in one dimension, Oxford (2004)

M. Cazalilla et al.,
Rev. Mod. Phys. 83 1405 (2011)

How to describe solids ?

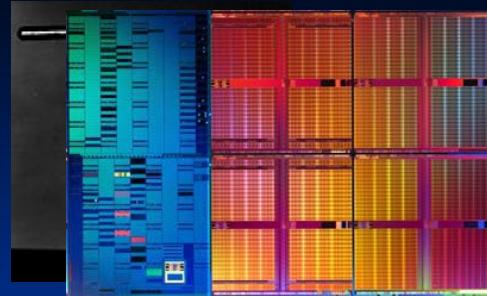
- Understood: free electrons
- Real systems : Coulomb interaction

$E \gg 10\,000\text{ K} !$

- Properties of realistic systems ?
- Free electron theory works quite well : Landau Fermi liquid



- Transistor



1956

- Supraconductivité



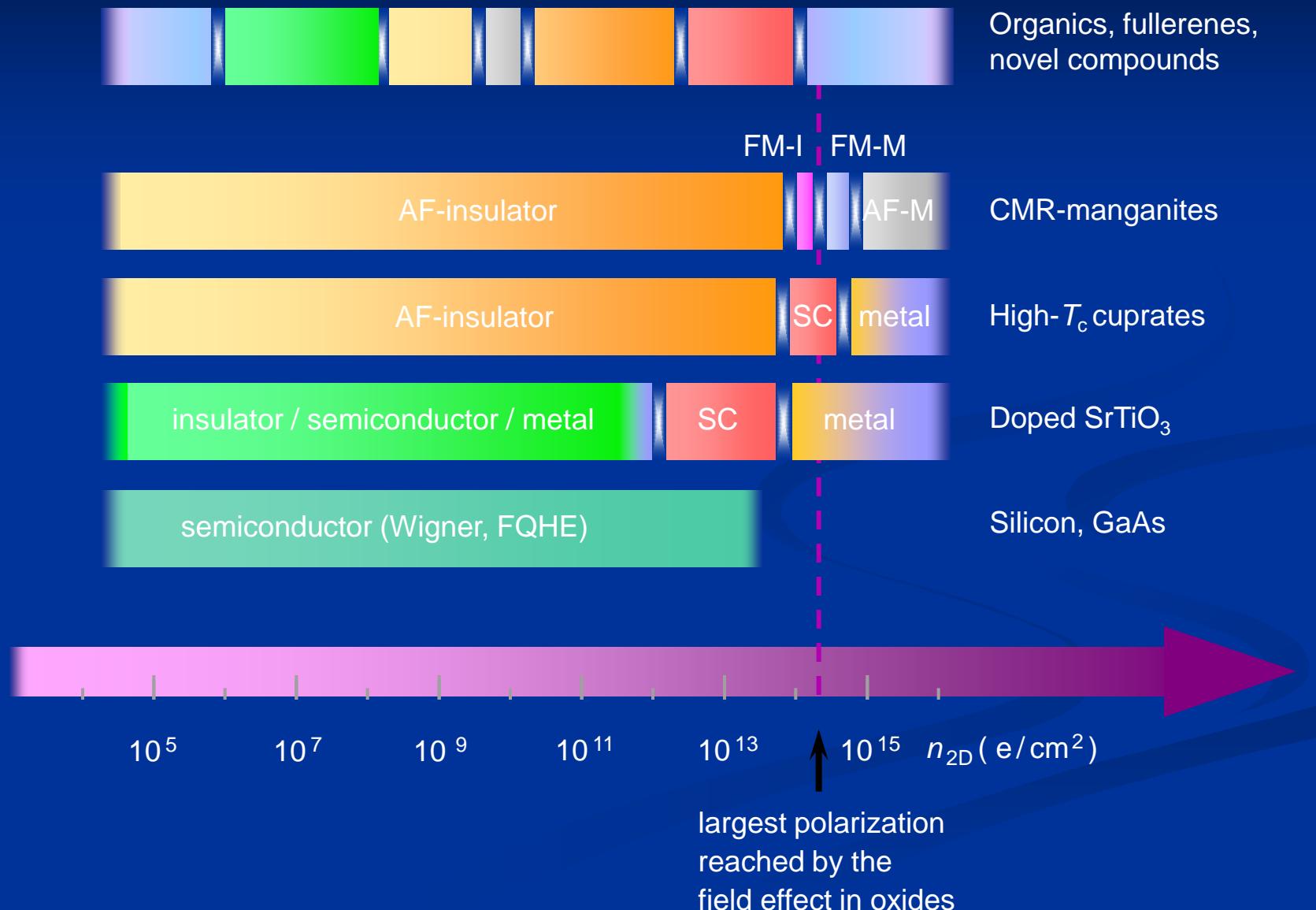
(1913), 1972,
1973, 1987, 2003

- Magnétorésistance géante



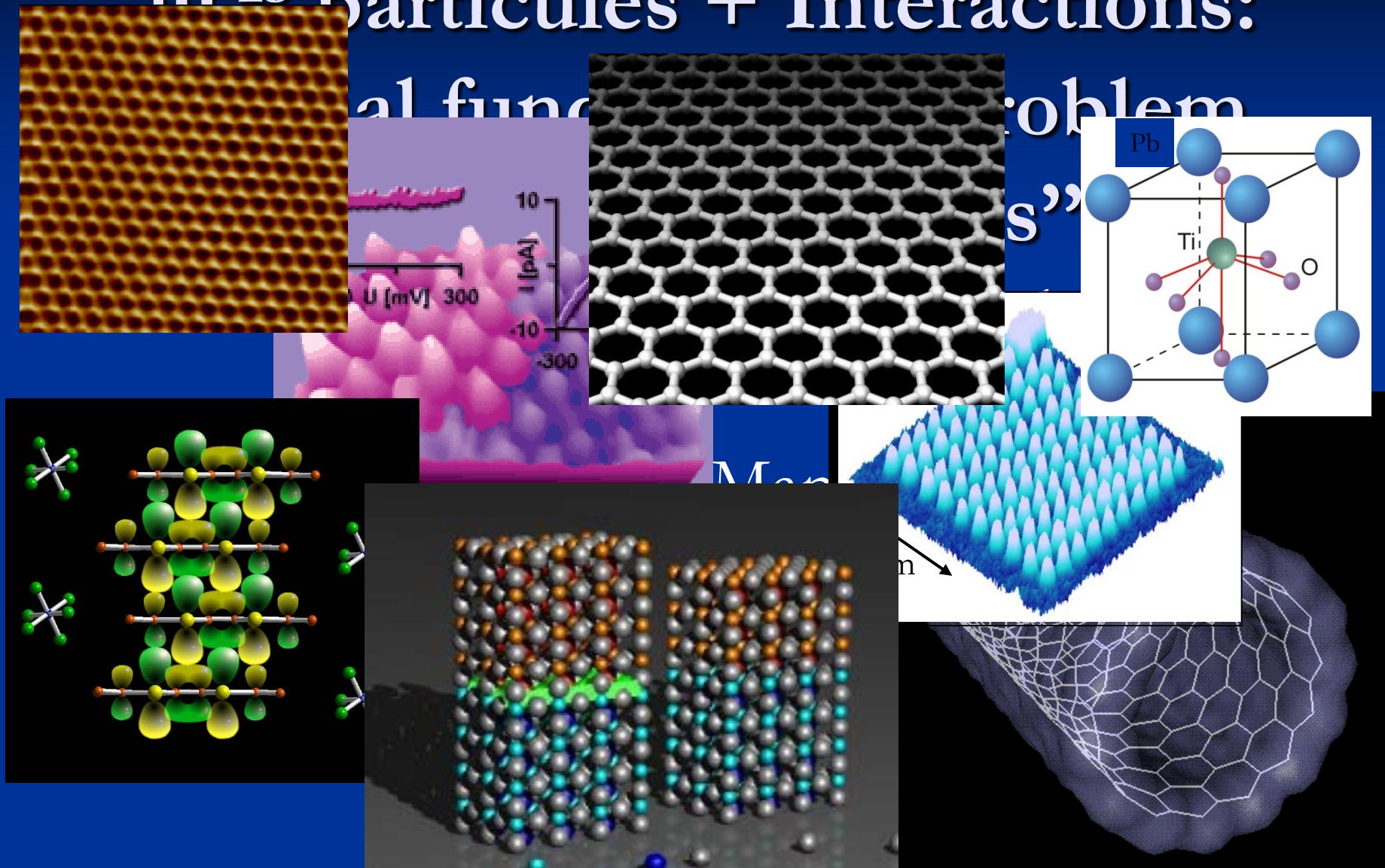
2007

Densities



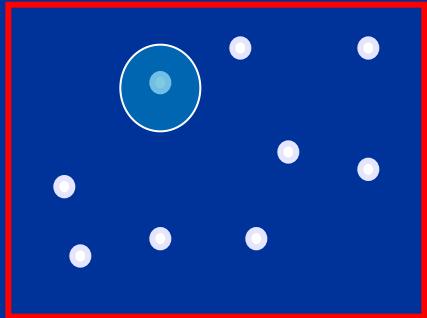
Need to understand interactions !

10^{23} particles + Interactions: real functionals problem



One dimension is specially interesting

- No individual excitation can exist (only collective ones)



- Strong quantum fluctuations

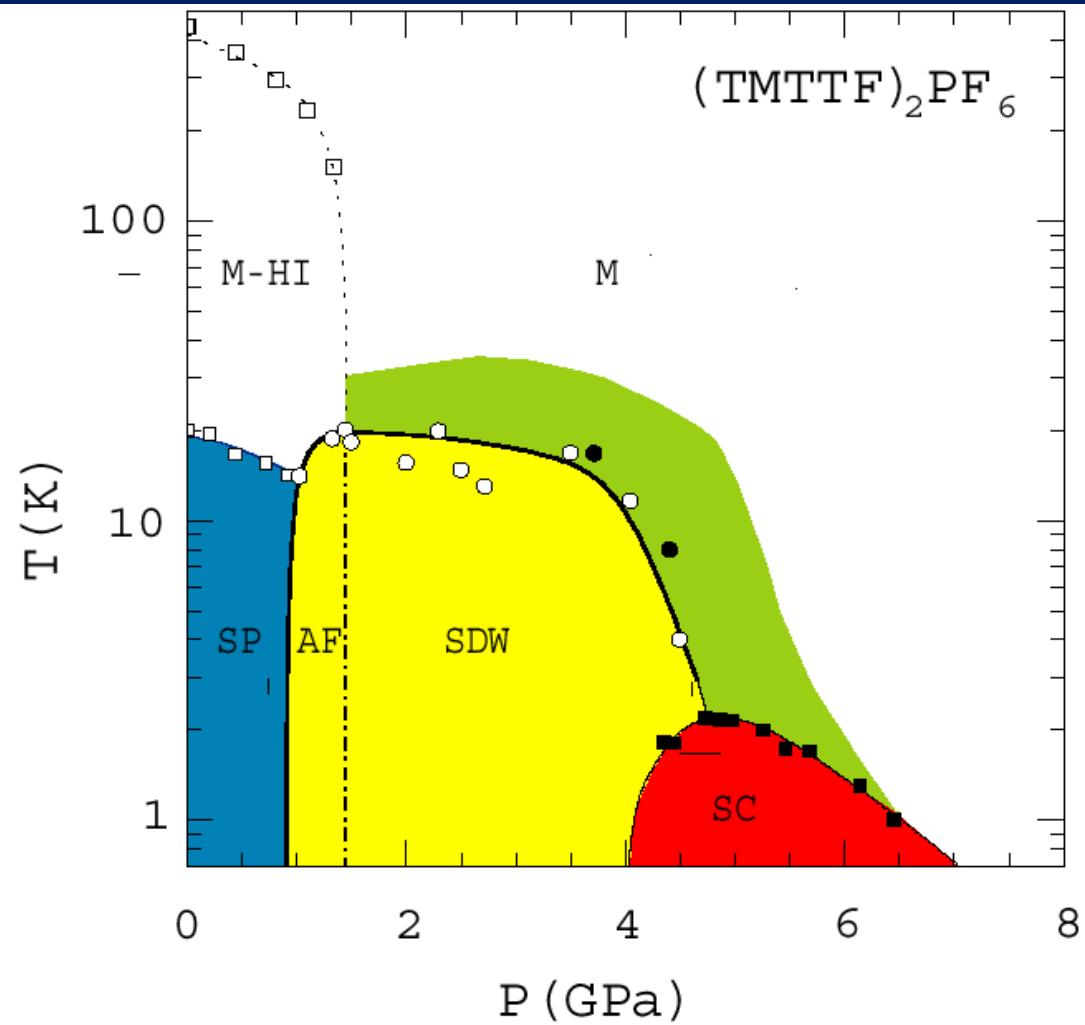
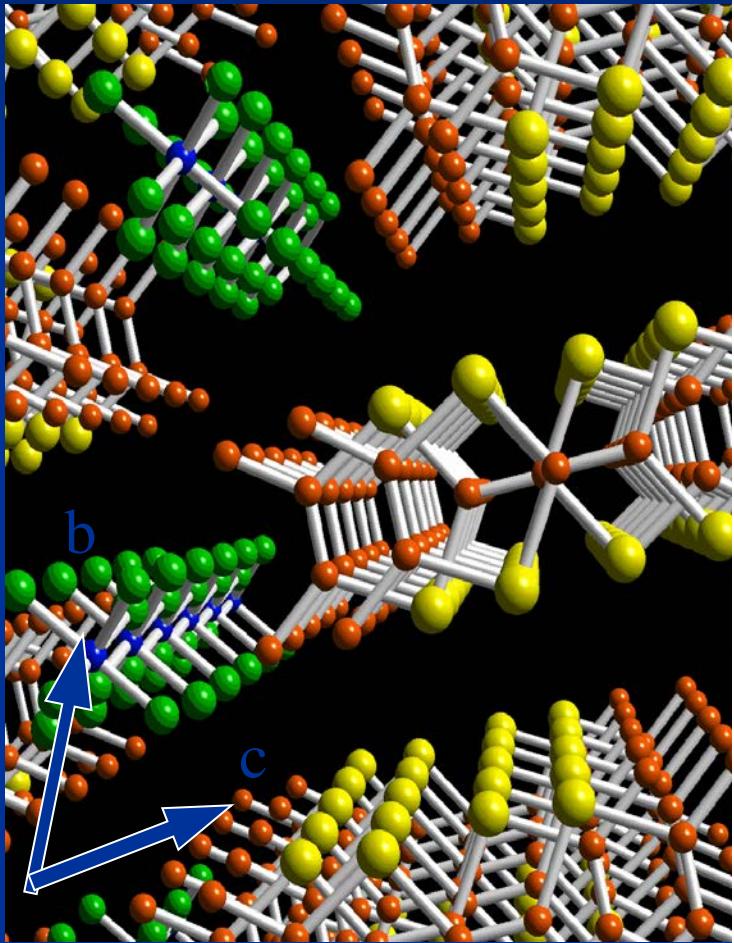
$$\psi = |\psi| e^{i\theta}$$

Difficult to order

How to realize

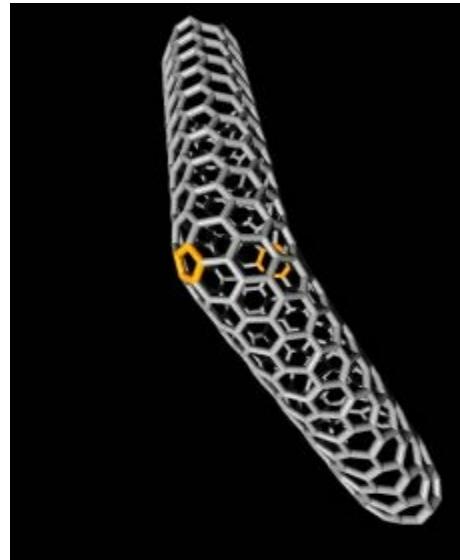
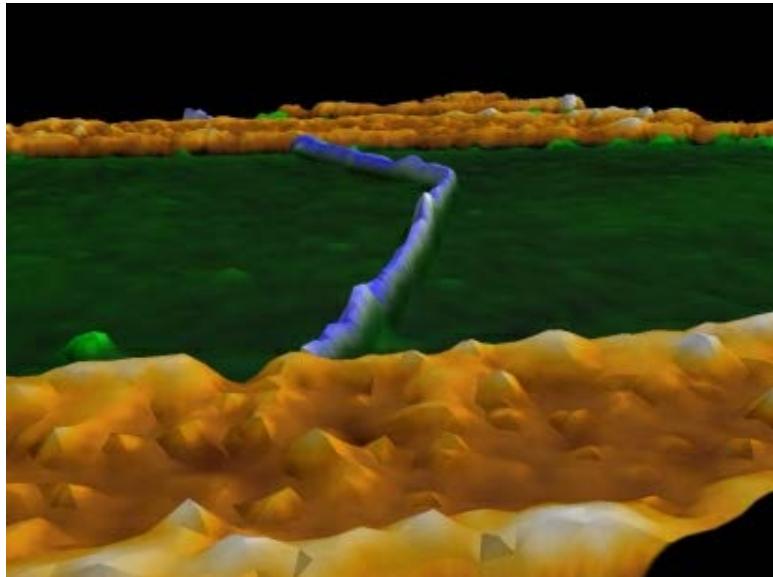
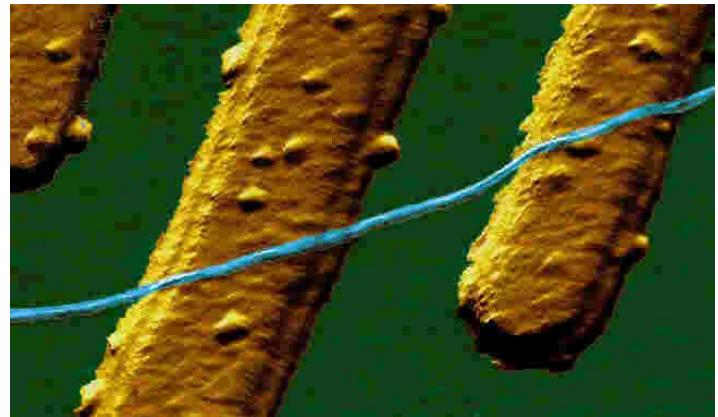


Organic conductors

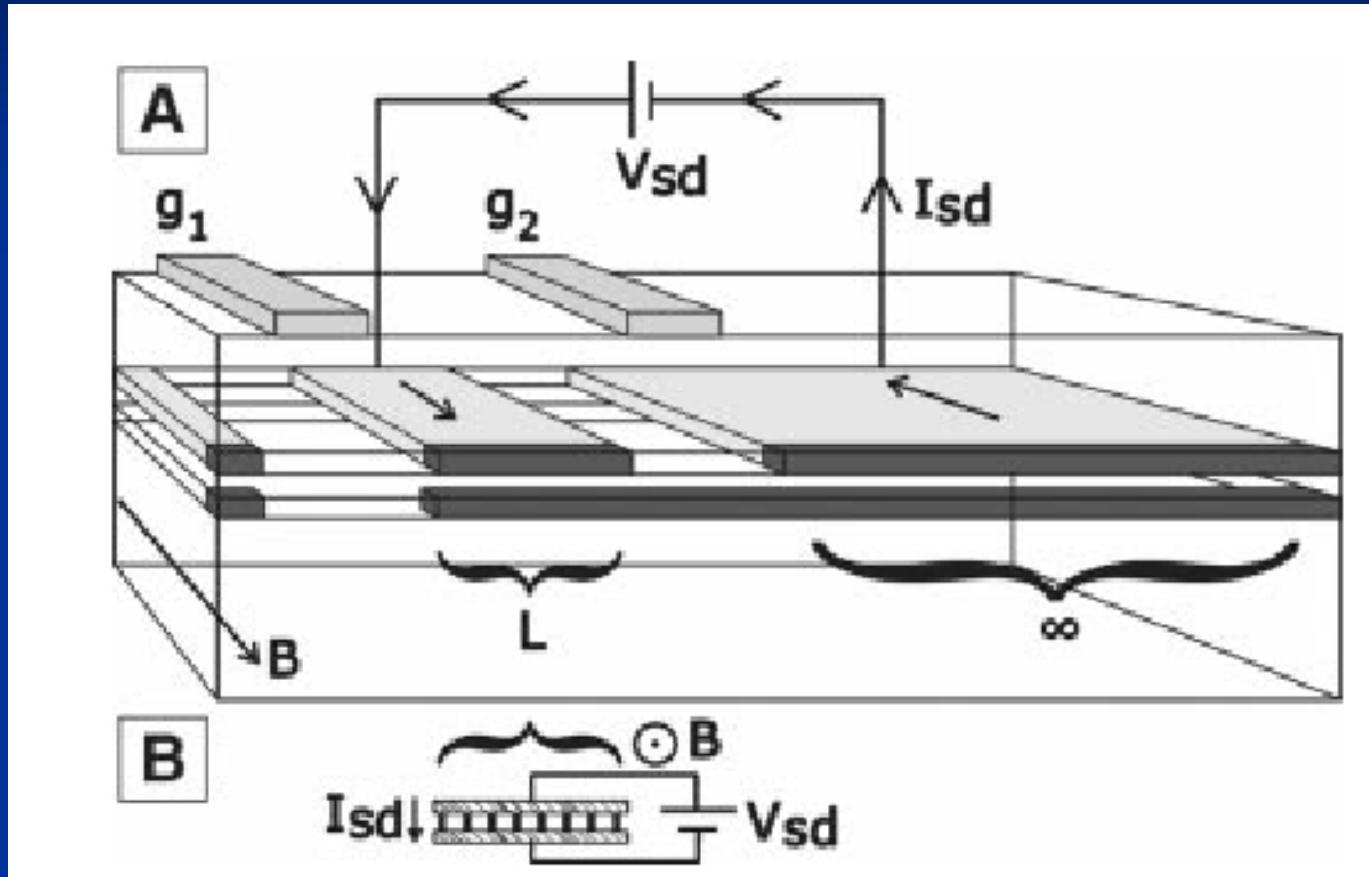


CARBON NANOTUBES

Cees Dekker



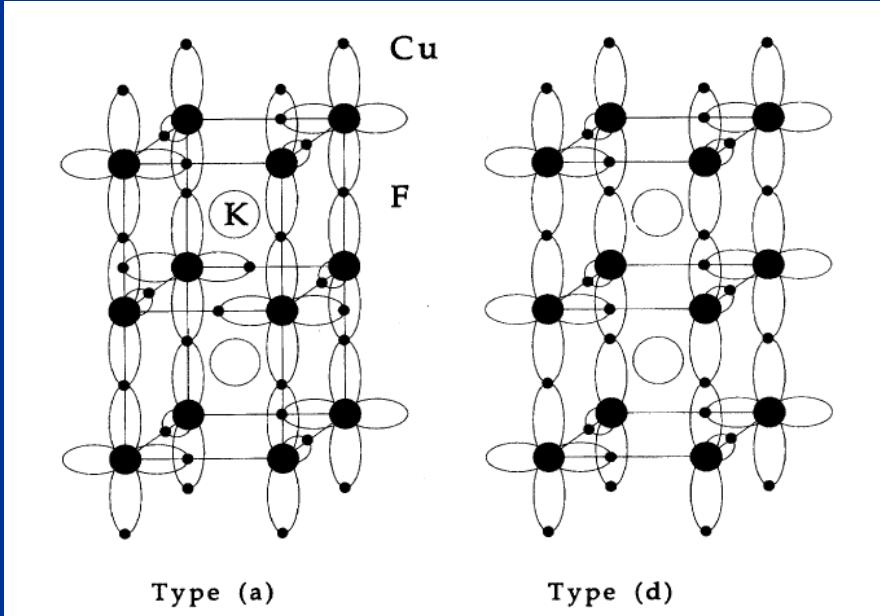
Quantum wires



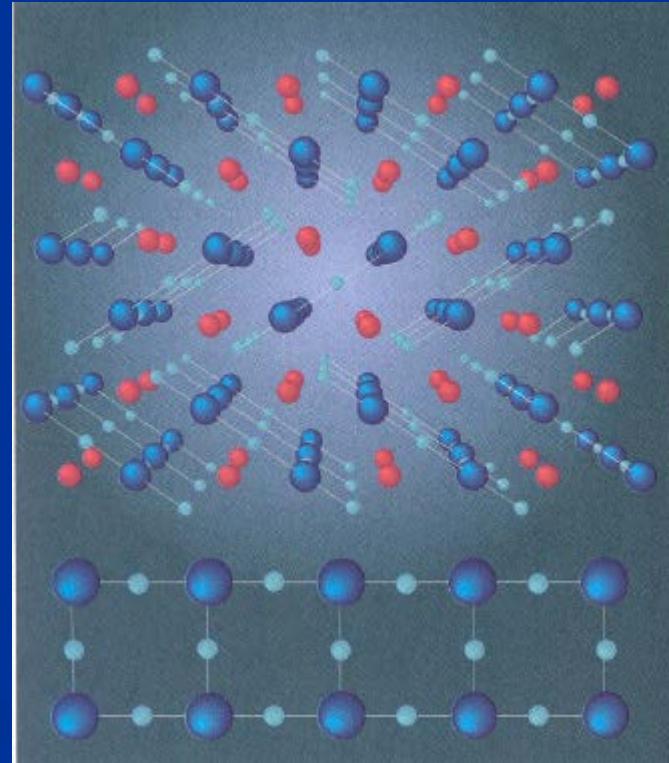
O.M Ausslander et al., Science 298 1354 (2001)

Spin chains and ladders

Spin chains

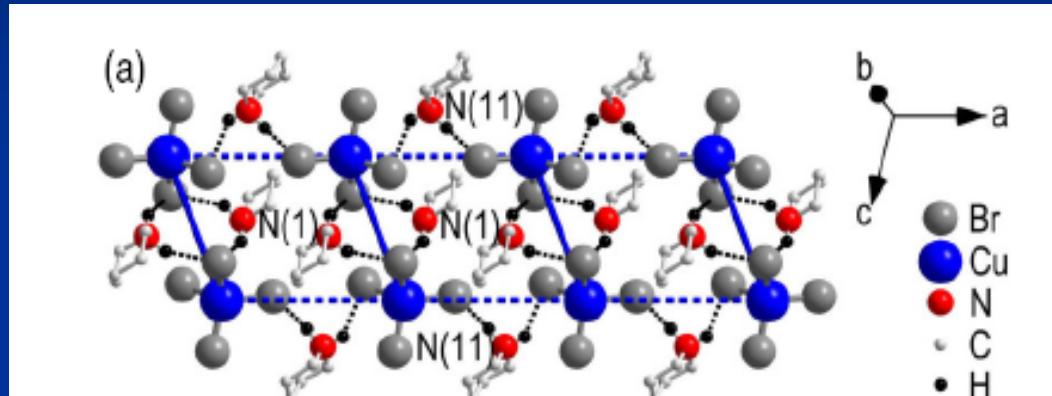


Ladder systems



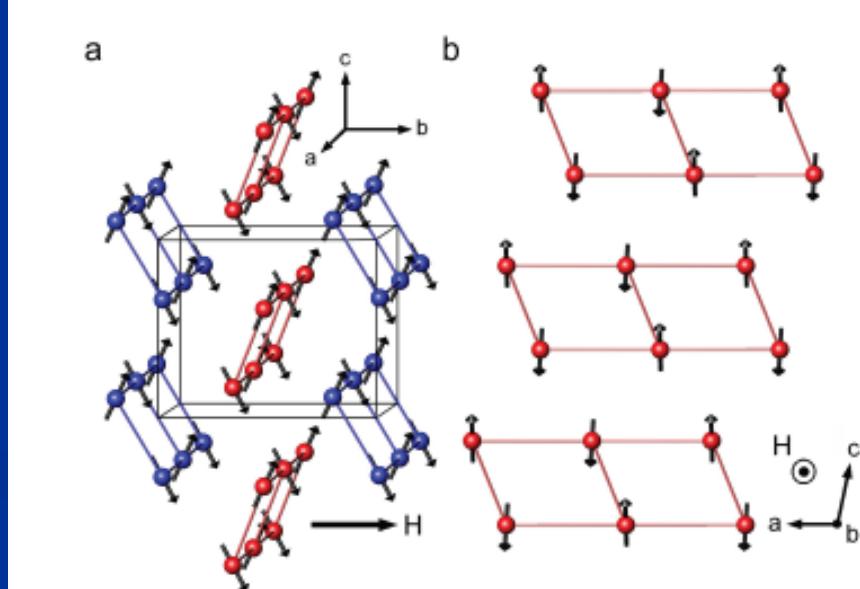
Spin dimer systems

B. C. Watson et al., PRL 86 5168 (2001)



M. Klanjsek et al.,
PRL 101 137207 (2008)

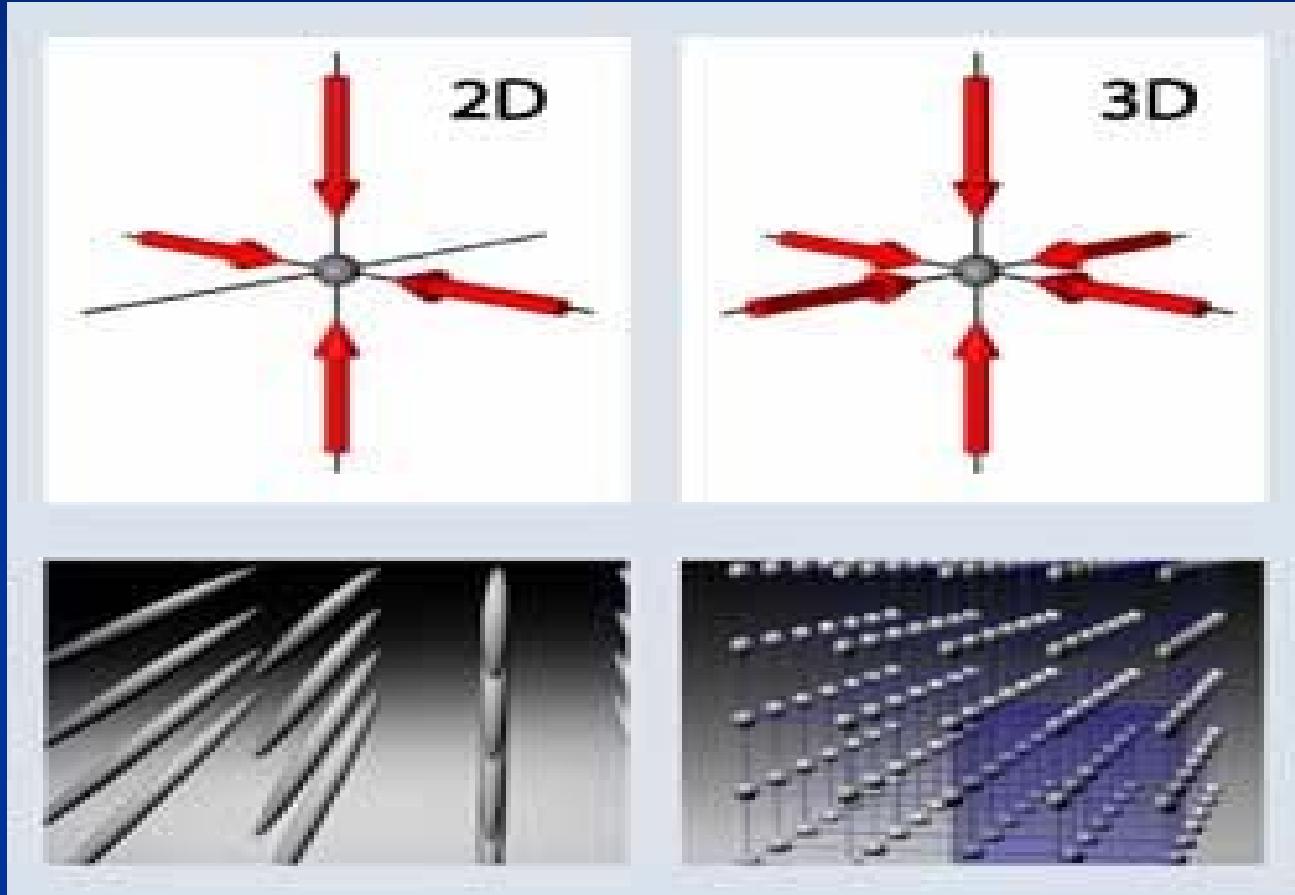
B. Thielemann et al.,
arXiv:0809.0440 (2008)



Cold atoms

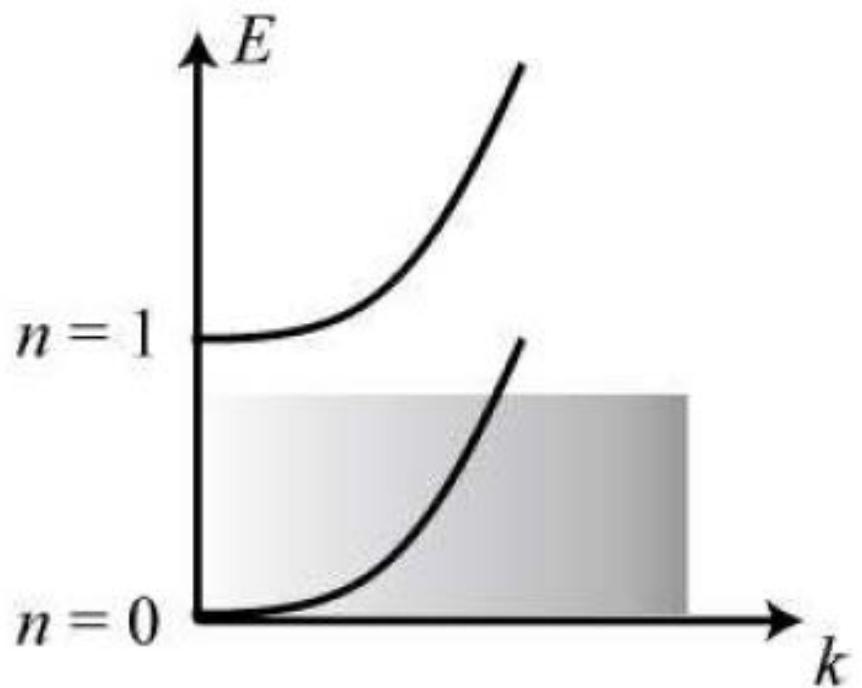
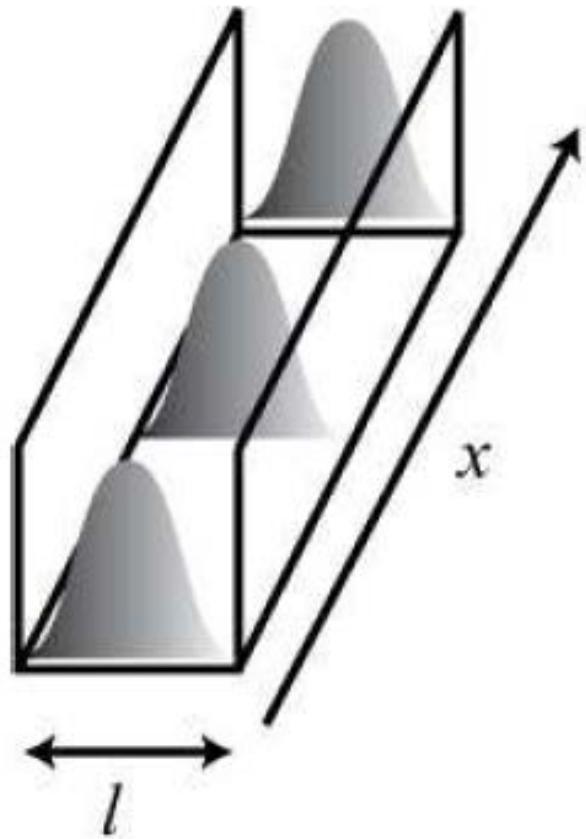


Control on the dimension



I. Bloch, Nat. Phys 1, 23 (2005)

What does 1D mean



How to treat



■ Exact Solutions

- Exact
- Limited to some special models
- Difficult to compute correlation functions

■ Numerics

- Very powerful
- Difficult to reach low energy physics (space/time)

■ Asymptotic methods

- Physically transparent
- Limited to low energy physics



Good excitations: collective ones

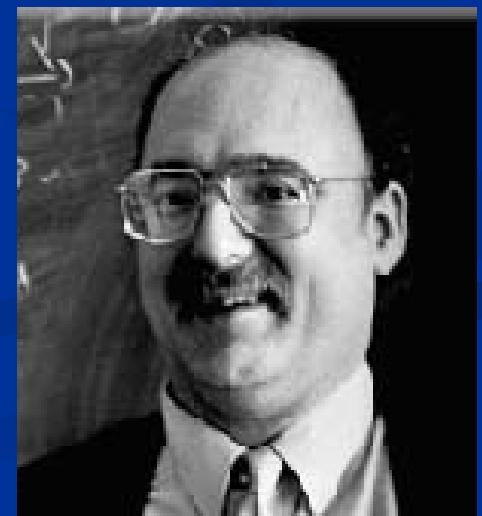
Tomonaga

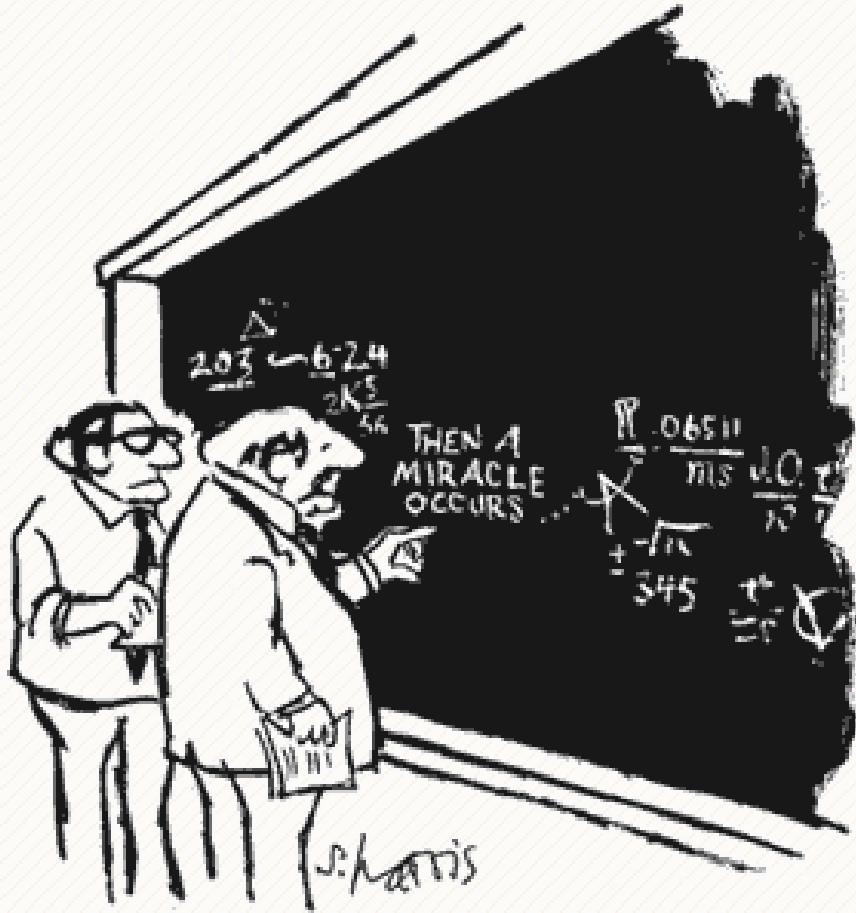
Luttinger

60': A. Larkin, I. Dzialoshinskii, L. Gorkov, Bichkov
E. Lieb, D. Mattis

70': A. Luther, V.J. Emery

80': F.D.M. Haldane





"I THINK YOU SHOULD BE MORE EXPLICIT HERE IN STEP TWO."

For more (among many refs) :

INTERNATIONAL SERIES OF MONOGRAPHS ON PHYSICS • 121

Quantum Physics in One Dimension

THIERRY GIAMARCHI



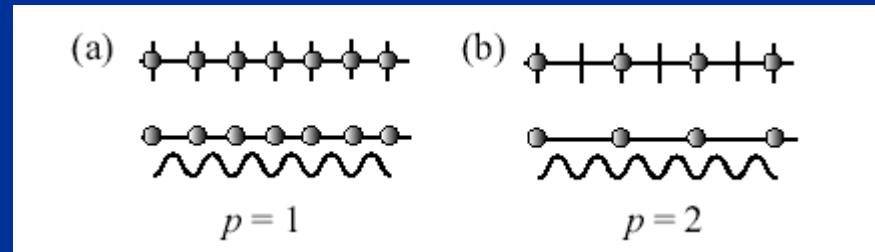
OXFORD SCIENCE PUBLICATIONS

Bosonization

- Continuum:

$$H = \int dx \frac{(\nabla\psi)^\dagger (\nabla\psi)}{2M} + \frac{1}{2} \int dx dx' V(x - x') \rho(x) \rho(x') - \mu \int dx \rho(x)$$

- Lattice:

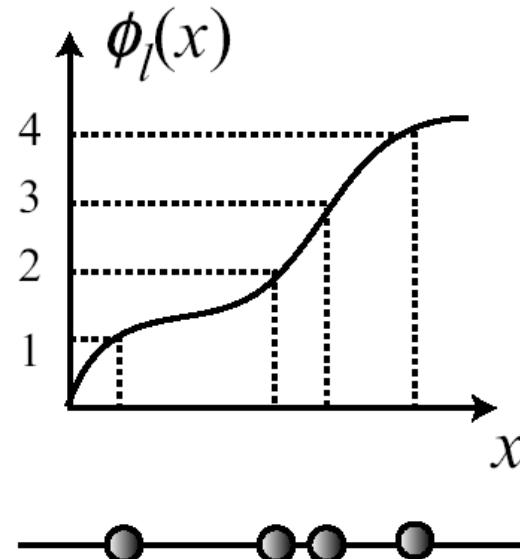
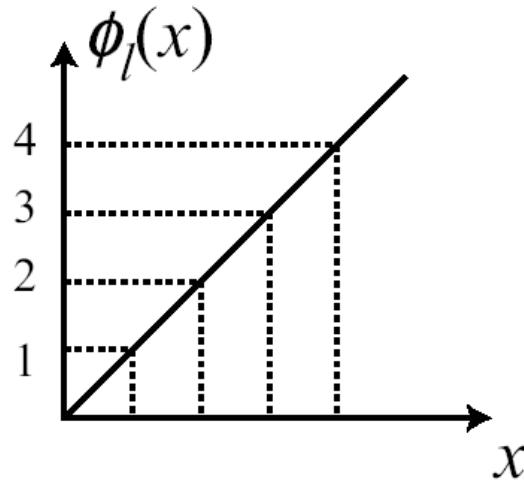


$$H = -J \sum_{\langle i,j \rangle} b_i^\dagger b_j + U \sum_i n_i(n_i - 1) - \mu \sum_i n_i$$

Labelling the particles

$$\begin{aligned}\rho(x) &= \sum_i \delta(x - x_i) \\ &= \sum_n |\nabla \phi_l(x)| \delta(\phi_l(x) - 2\pi n)\end{aligned}$$

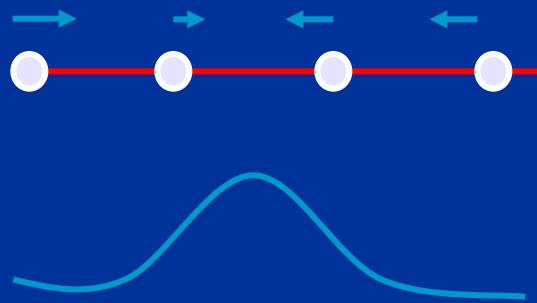
1D: unique way
of labelling



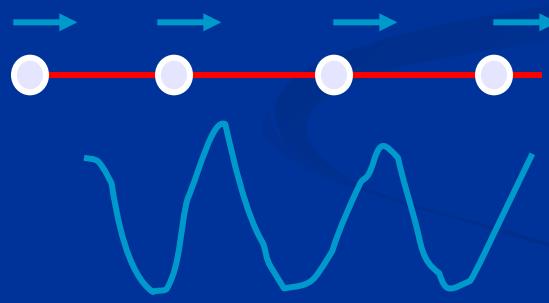
$$\phi_l(x) = 2\pi\rho_0x - 2\phi(x)$$

$$\rho(x) = \left[\rho_0 - \frac{1}{\pi} \nabla \phi(x) \right] \sum_p e^{i2p(\pi\rho_0x - \phi(x))}$$

$\phi(x)$ varies slowly



$$q \sim 0$$



$$q \sim 2\pi\rho_0$$

CDW

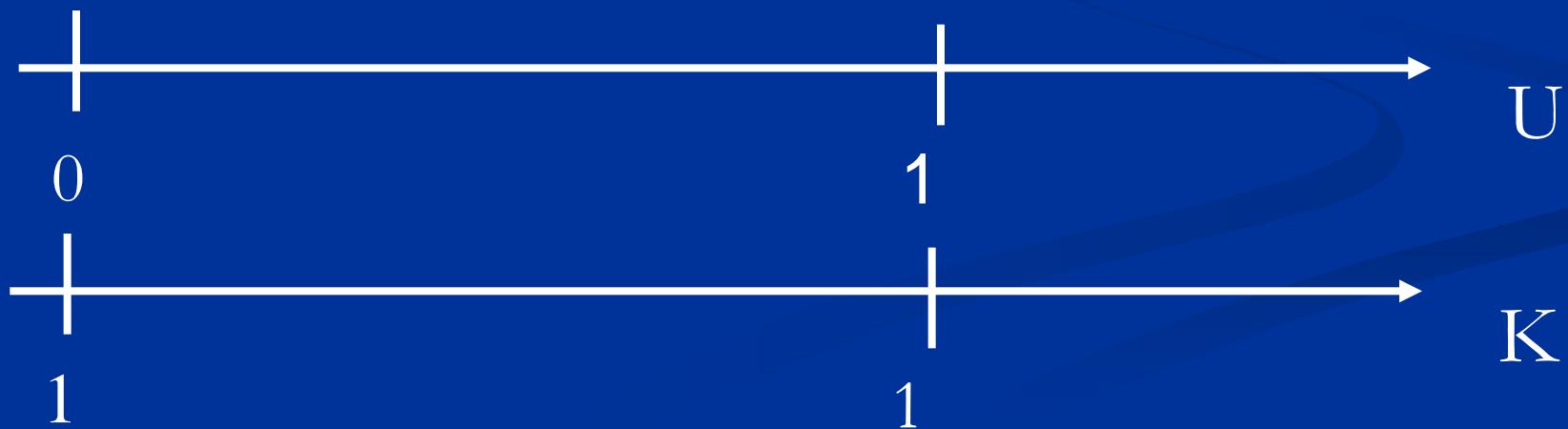
$$\psi^\dagger(x) = [\rho(x)]^{1/2} e^{-i\theta(x)}$$

θ : superfluid phase

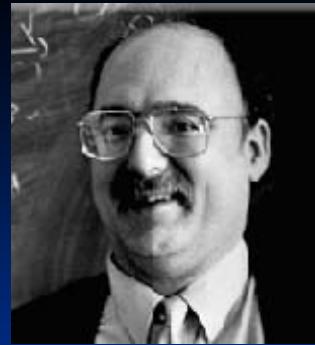
$$[\frac{1}{\pi} \nabla \phi(x), \theta(x')] = -i\delta(x - x')$$

Quantum
fluctuations

$$H = \frac{\hbar}{2\pi} \int dx \left[\frac{uK}{\hbar^2} (\pi\Pi(x))^2 + \frac{u}{K} (\nabla\phi(x))^2 \right]$$



Luttinger liquid concept

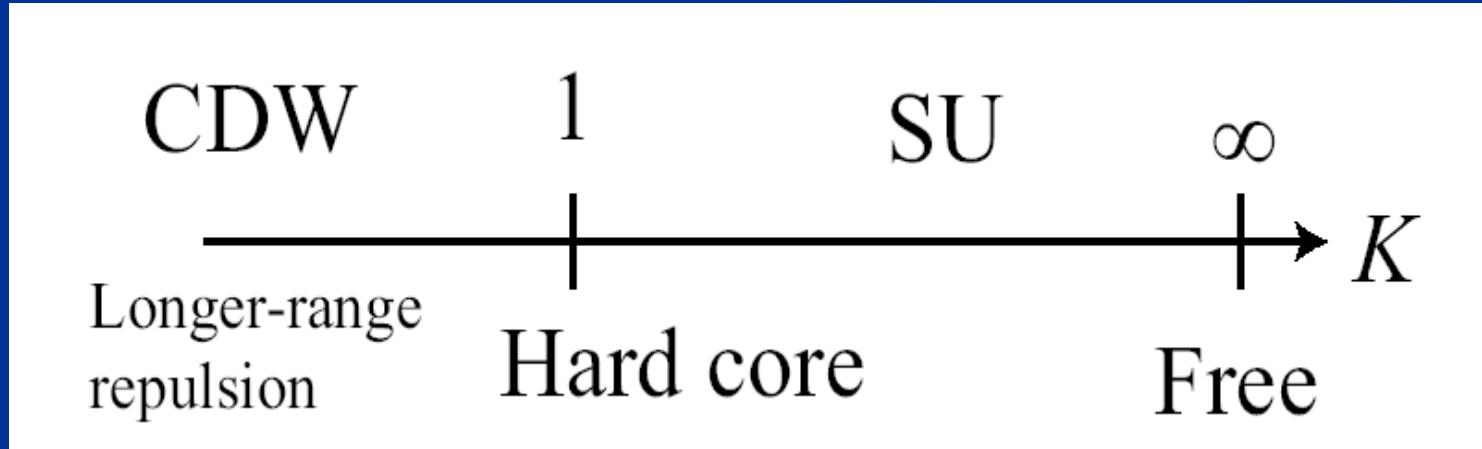


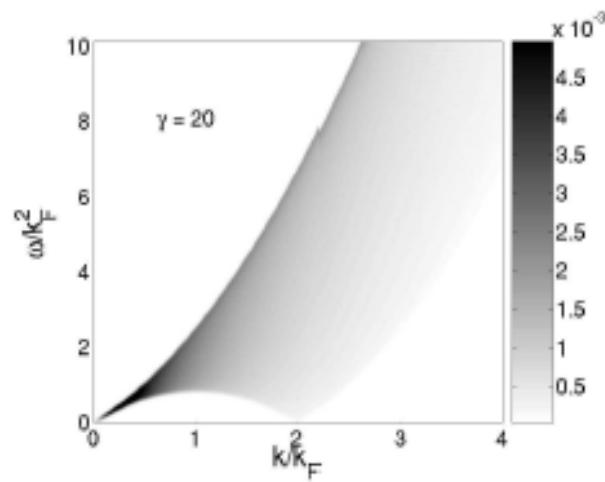
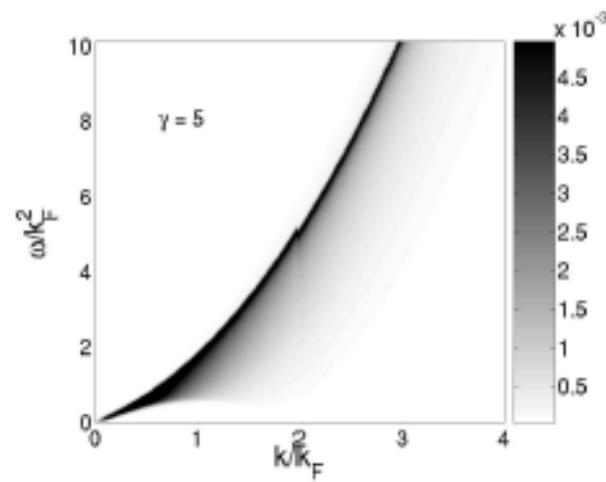
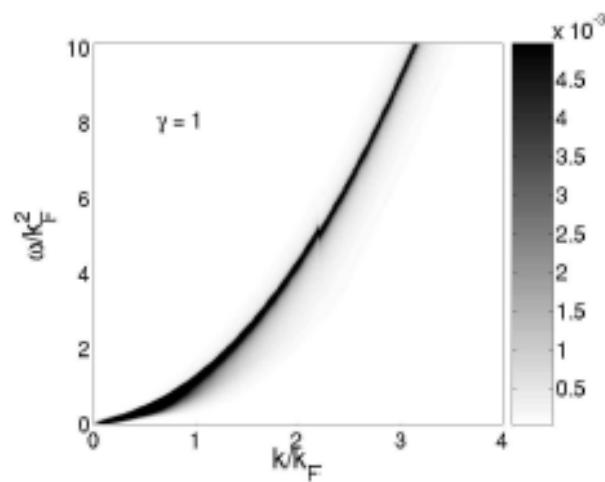
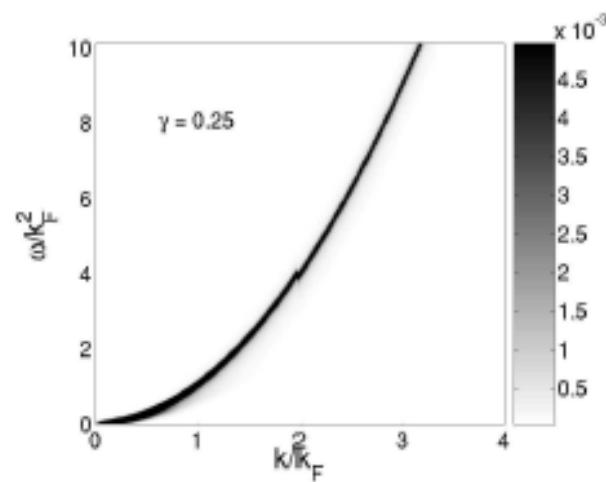
- How much is perturbative ?
- Nothing (Haldane):
provided the correct u, K are used
- Low energy properties: Luttinger liquid
(fermions, bosons, spins...)

Correlations

$$\langle \psi(r)\psi^\dagger(0) \rangle = A_1 \left(\frac{\alpha}{r}\right)^{\frac{1}{2K}} + \dots$$

$$\langle \rho(r)\rho(0) \rangle = \rho_0^2 + \frac{K}{2\pi^2} \frac{y_\alpha^2 - x^2}{(y_\alpha^2 + x^2)^2} + A_3 \cos(2\pi\rho_0 x) \left(\frac{1}{r}\right)^{2K} + \dots$$





S(q,!) J.S. Caux et al PRA 74 031605 (2006)

Finite temperature

Conformal theory



Calculation of Luttinger parameters

- Trick: use thermodynamics and BA or numerics
- Compressibility: u/K
- Response to a twist in boundary: $u K$
- Specific heat : T/u
- Etc.

Tonks limit



$U = 1$: spinless fermions

Not for $n(k)$: $\Psi_F \neq \Psi_B$

Free fermions: $\langle \rho(x)\rho(0) \rangle \propto \cos(2k_F x) \left(\frac{1}{x}\right)^2$

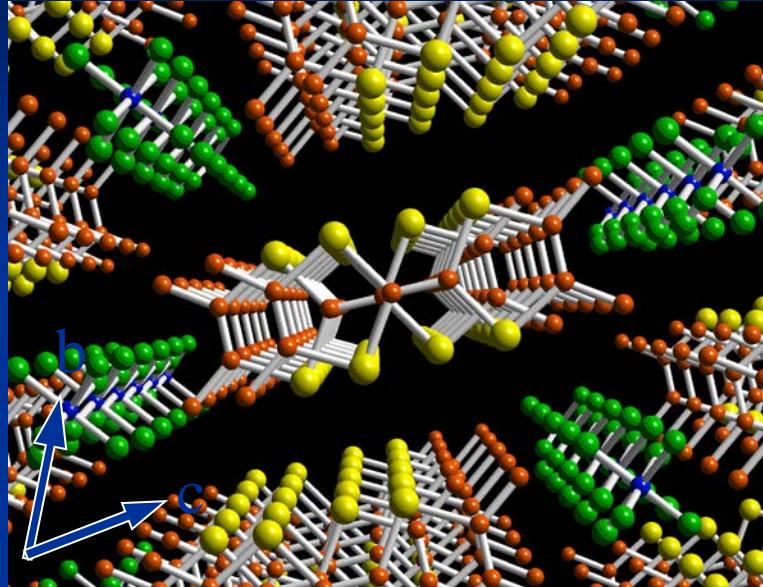
$K=1$

Note: $\langle \psi_B(x)\psi_B(0)^\dagger \rangle \propto \left(\frac{1}{x}\right)^{1/2}$

Tests of Luttinger liquids



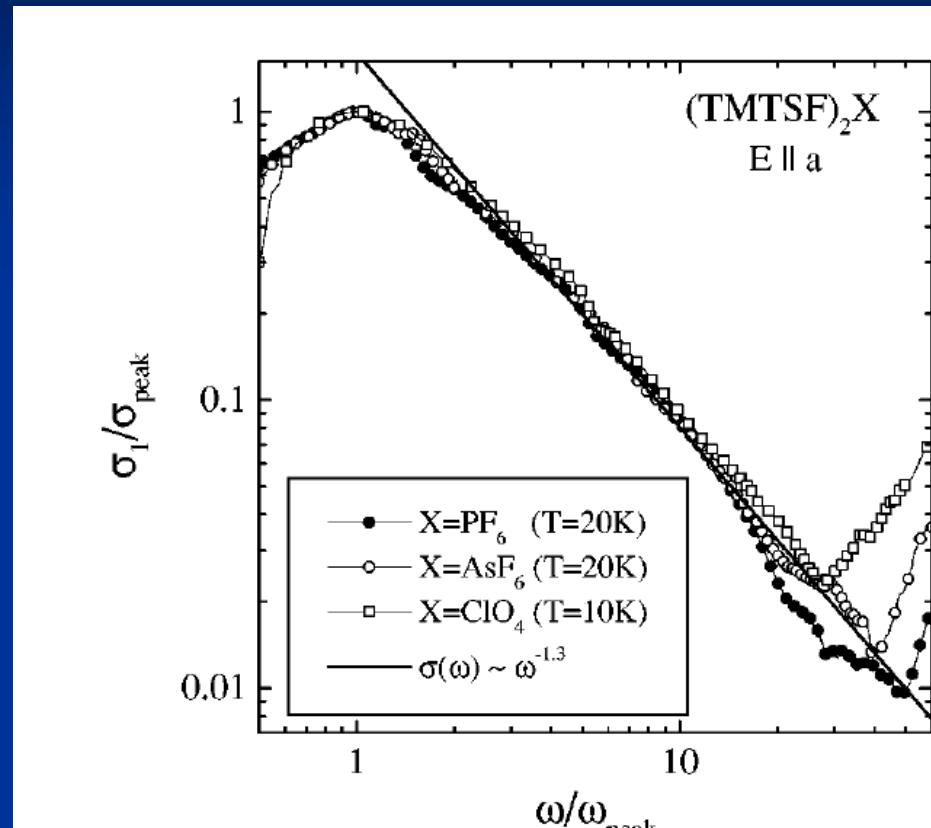
Organic conductors



$$\sigma(\omega) \sim \omega^{-\nu}$$

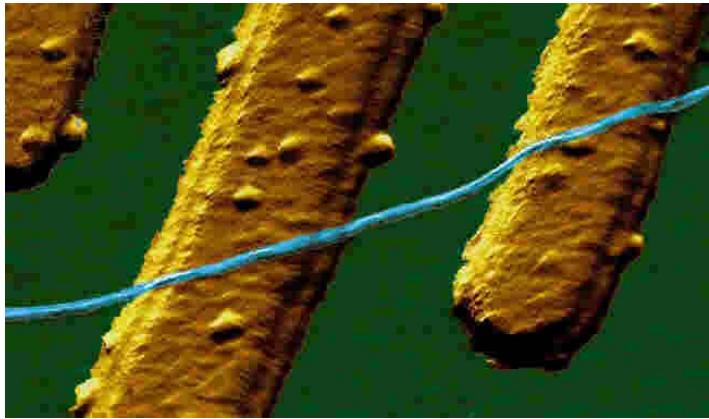
TG PRB (91) :

Physica B 230 (1996)

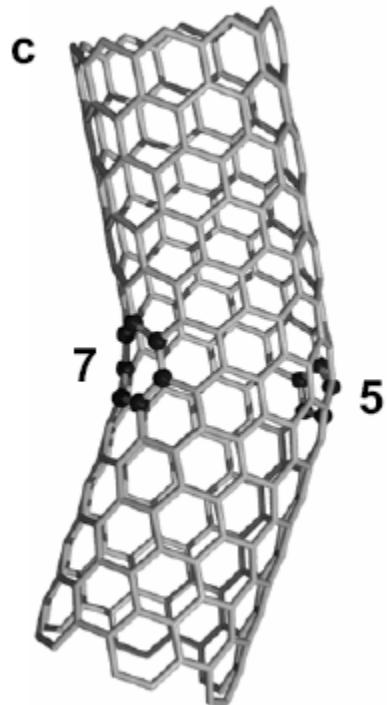
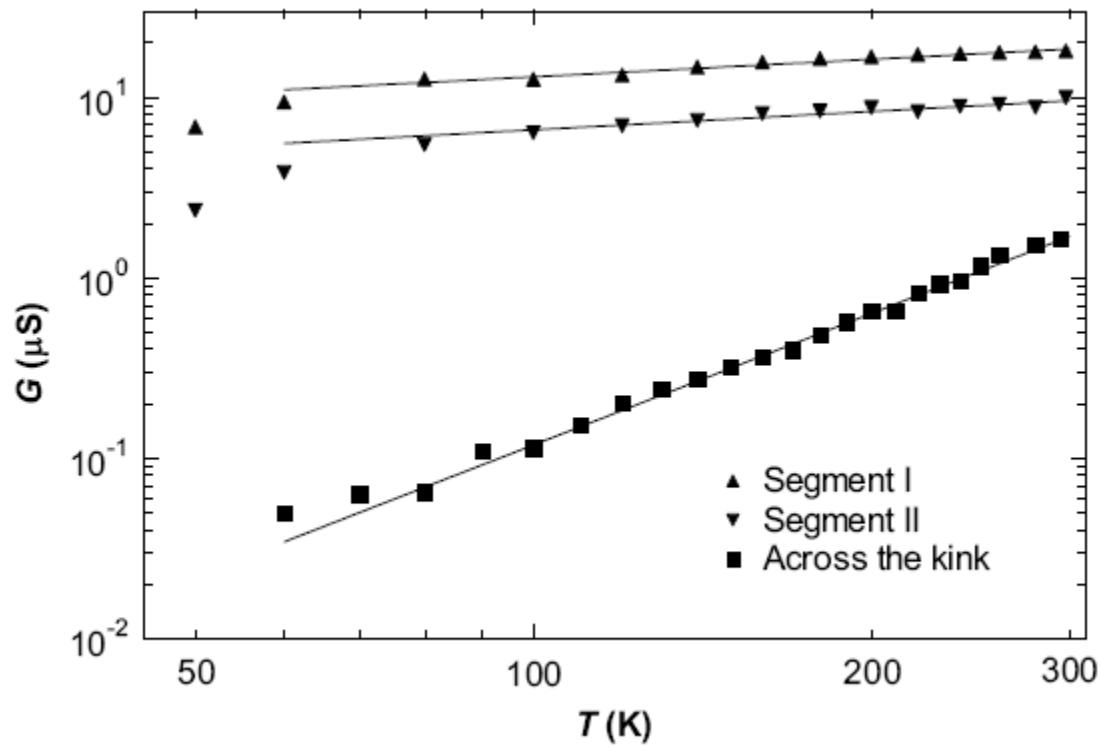


A. Schwartz et al. PRB 58 1261 (1998)

First observation of LL/powerlaw !!



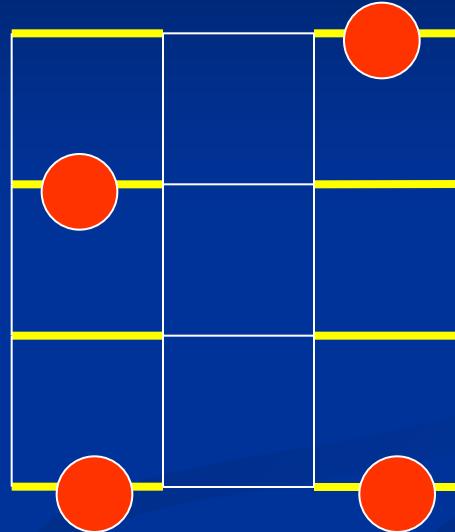
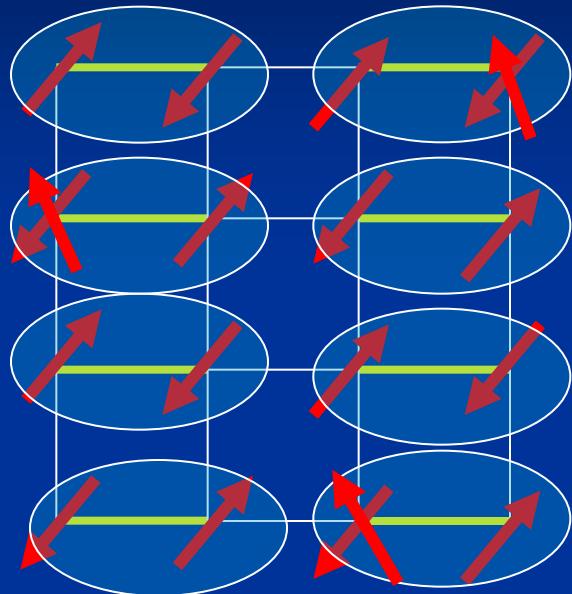
Z. Yao et al. Nature 402
273 (1999)



Magnetic insulators



triplon = hard core boson



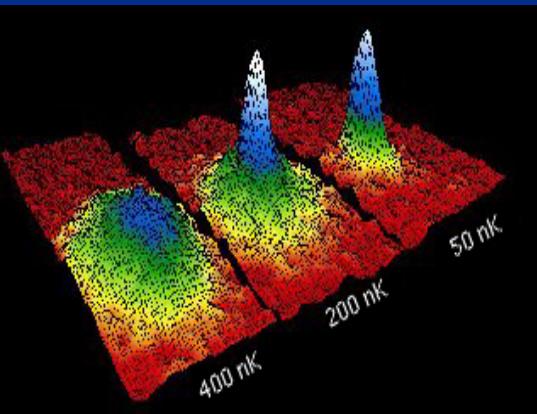
$h \sim h_c$ dilute limit: « free » bosons

Bose Einstein Condensation

(TG and A. M. Tsvelik PRB 59 11398 (1999))

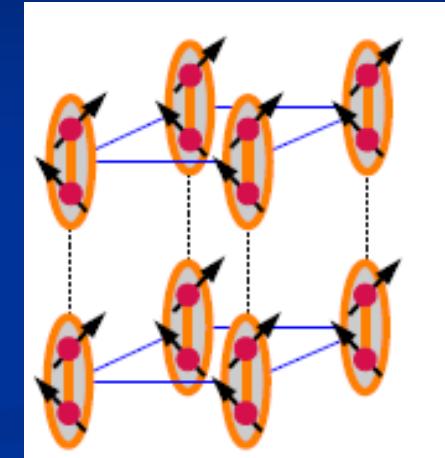
BEC vs BEC

Cold atoms



Dimers/Spins

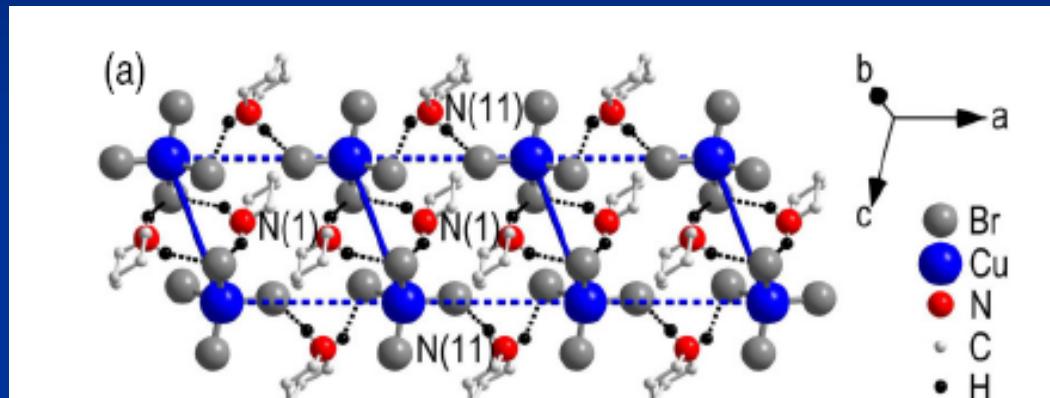
TG, Ch. Rüegg,
O. Tchernyshyov,
Nat. Phys. 4 198 (08)



Bose gas	Antiferromagnet
Particles	Spin excitations ($\Delta S^z = \pm 1$)
Boson number N	Spin component S^z
Charge conservation U(1)	Rotational invariance O(2)
Condensate wavefunction $\langle \psi(\mathbf{r}) \rangle$	Transverse magnetic order $\langle s_i^x + i s_i^y \rangle$
Chemical potential μ	Magnetic field B
Superfluid density ρ_s	Transverse spin stiffness
Mott insulating state	Integer magnetization plateau

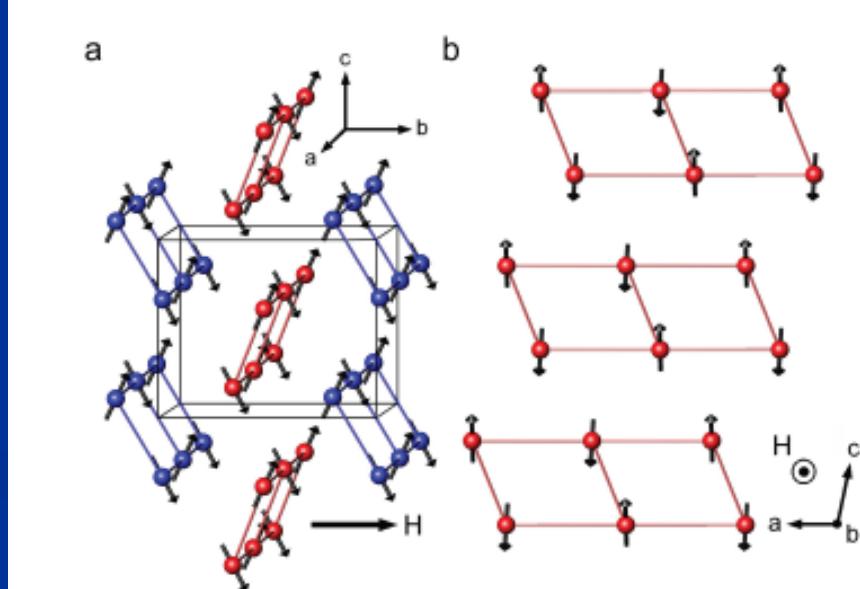
Spin dimer systems

B. C. Watson et al., PRL 86 5168 (2001)

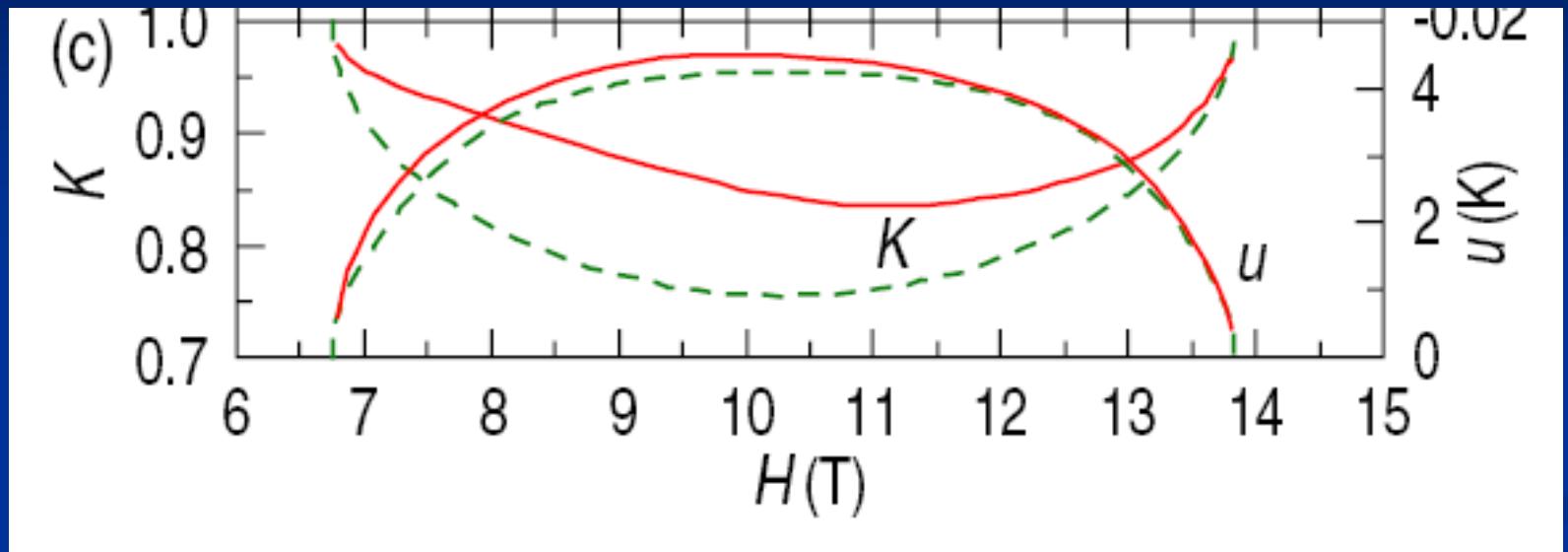


M. Klanjsek et al.,
PRL 101 137207 (2008)

B. Thielemann et al.,
arXiv:0809.0440 (2008)



Luttinger parameters



M. Klanjsek et al., PRL 101 137207 (2008)

Red : Ladder (DMRG)

Green: Strong coupling ($J_r \rightarrow 1$) (BA)

Correlation functions

M. Klanjsek et al., PRL 101 137207 (2008)

R. Chitra, TG PRB 55 5816 (97); TG, AM Tsvelik PRB 59 11398 (99)

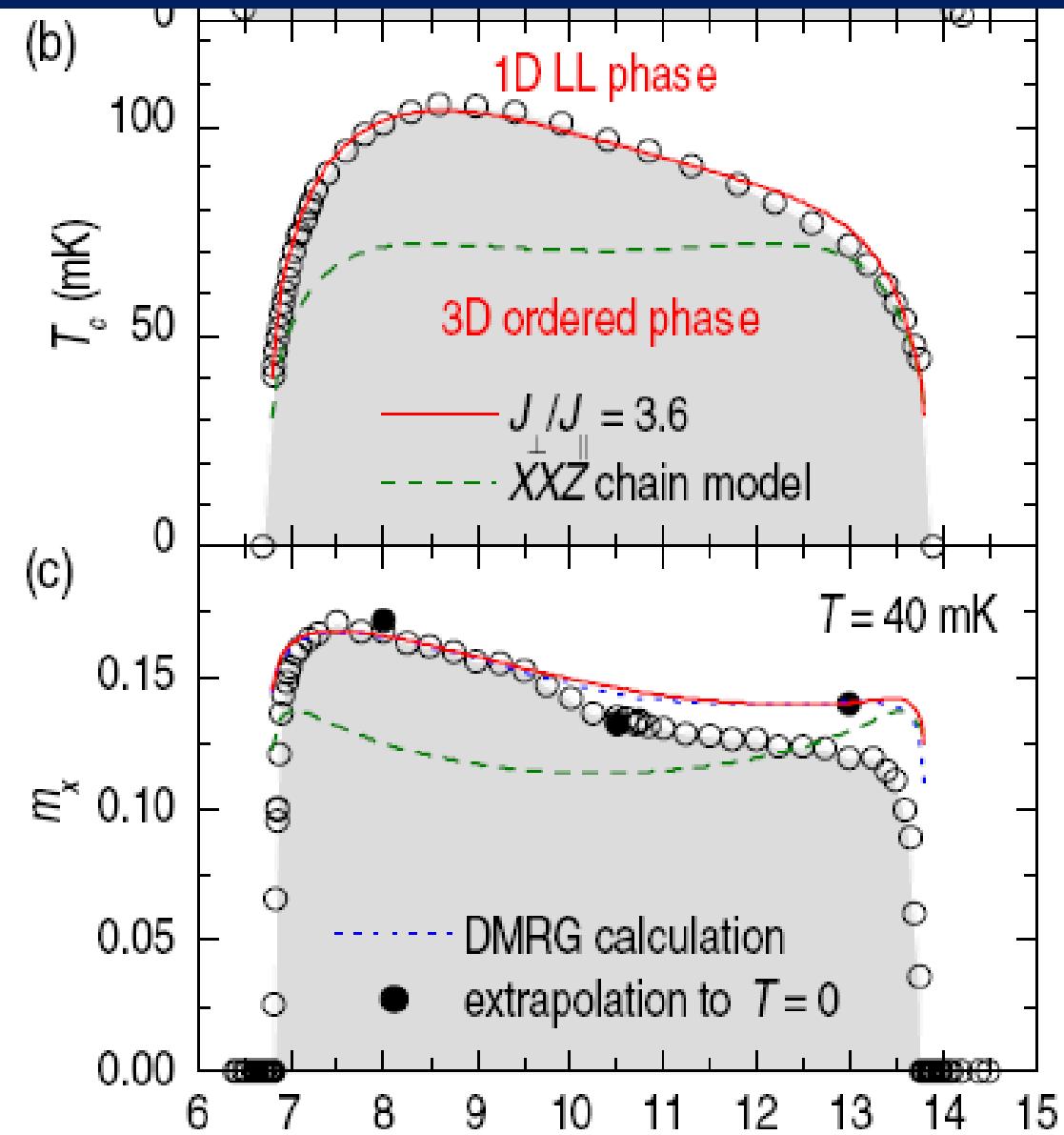
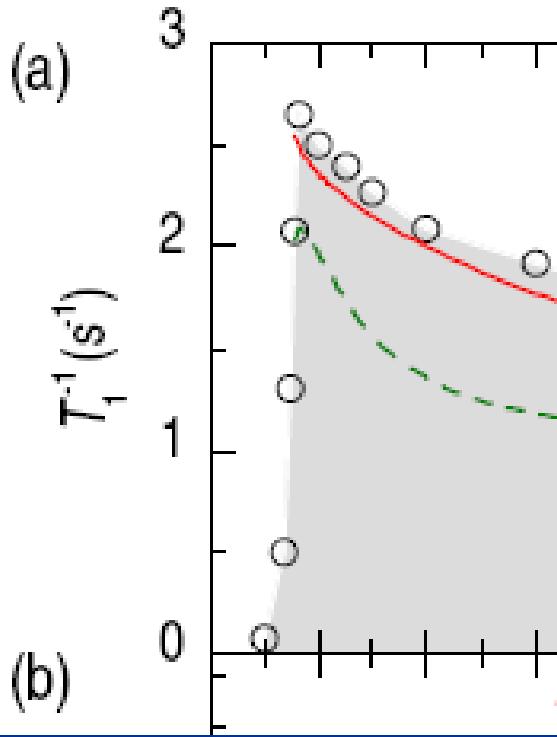
■ NMR relaxation rate:

$$T_1^{-1} = \frac{\hbar\gamma^2 A_\perp^2 A_0^x}{k_B u} \cos\left(\frac{\pi}{4K}\right) B\left(\frac{1}{4K}, 1 - \frac{1}{2K}\right) \left(\frac{2\pi T}{u}\right)^{(1/2K)-1},$$

■ Tc to ordered phase: $1/J' = \chi_{1D}(T_c)$

$$T_c = \frac{u}{2\pi} \left[\sin\left(\frac{\pi}{4K}\right) B^2\left(\frac{1}{8K}, 1 - \frac{1}{4K}\right) \frac{zJ' A_0^x}{2u} \right]^{2K/(4K-1)}.$$

NMR



M. Klanjsek et al.,

PRL 101 137207 (2008)

Cold atoms

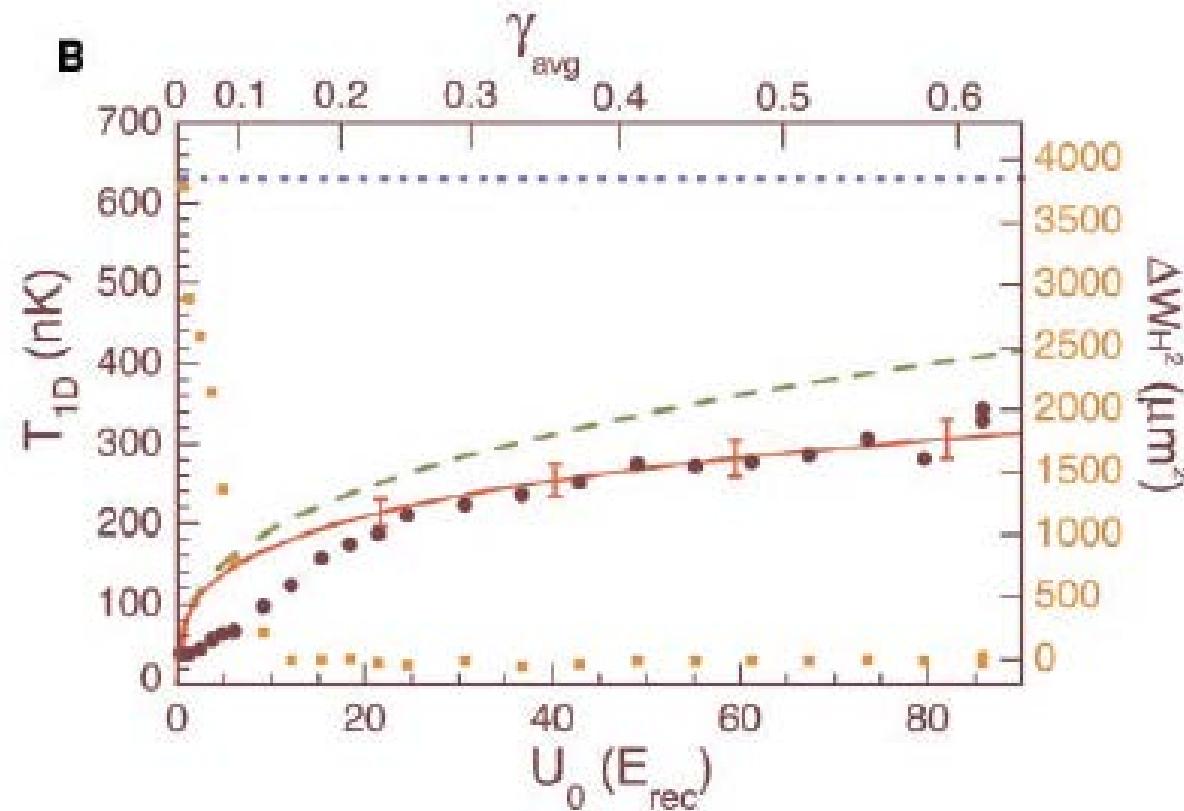
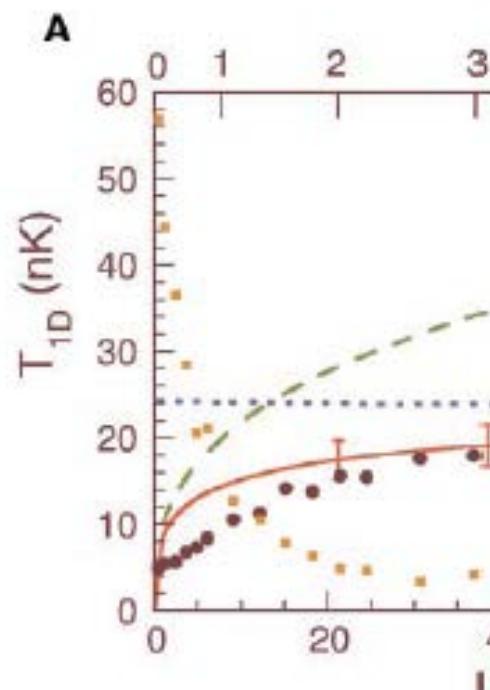


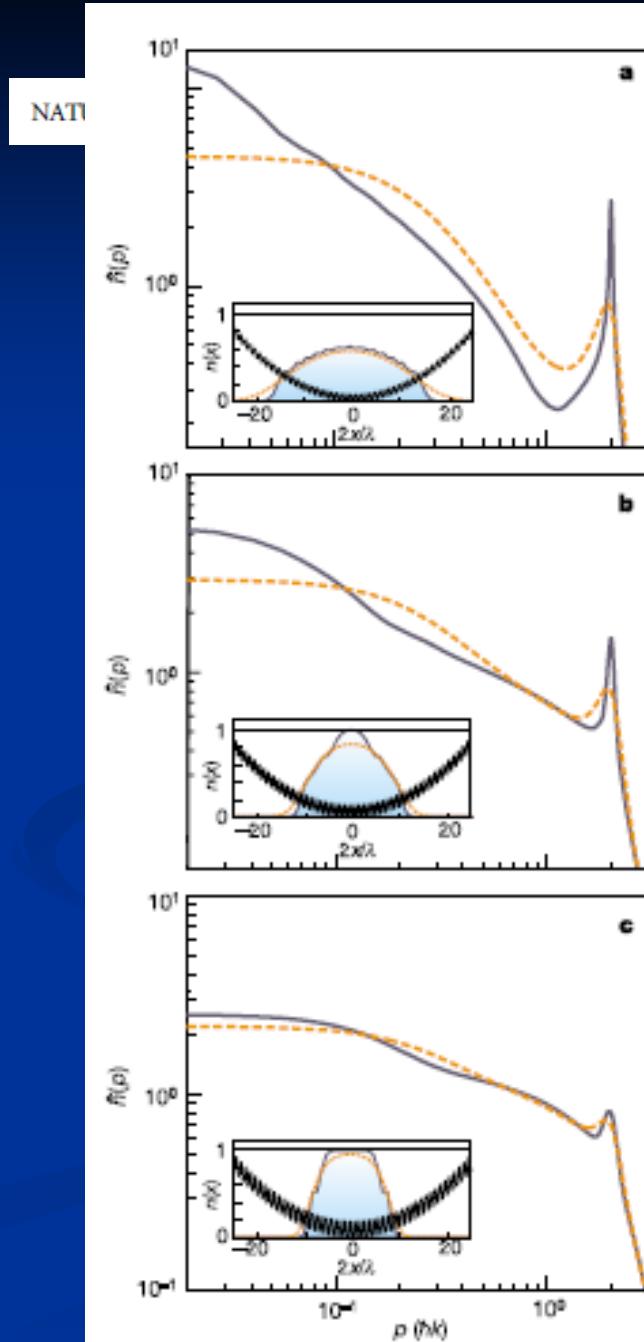
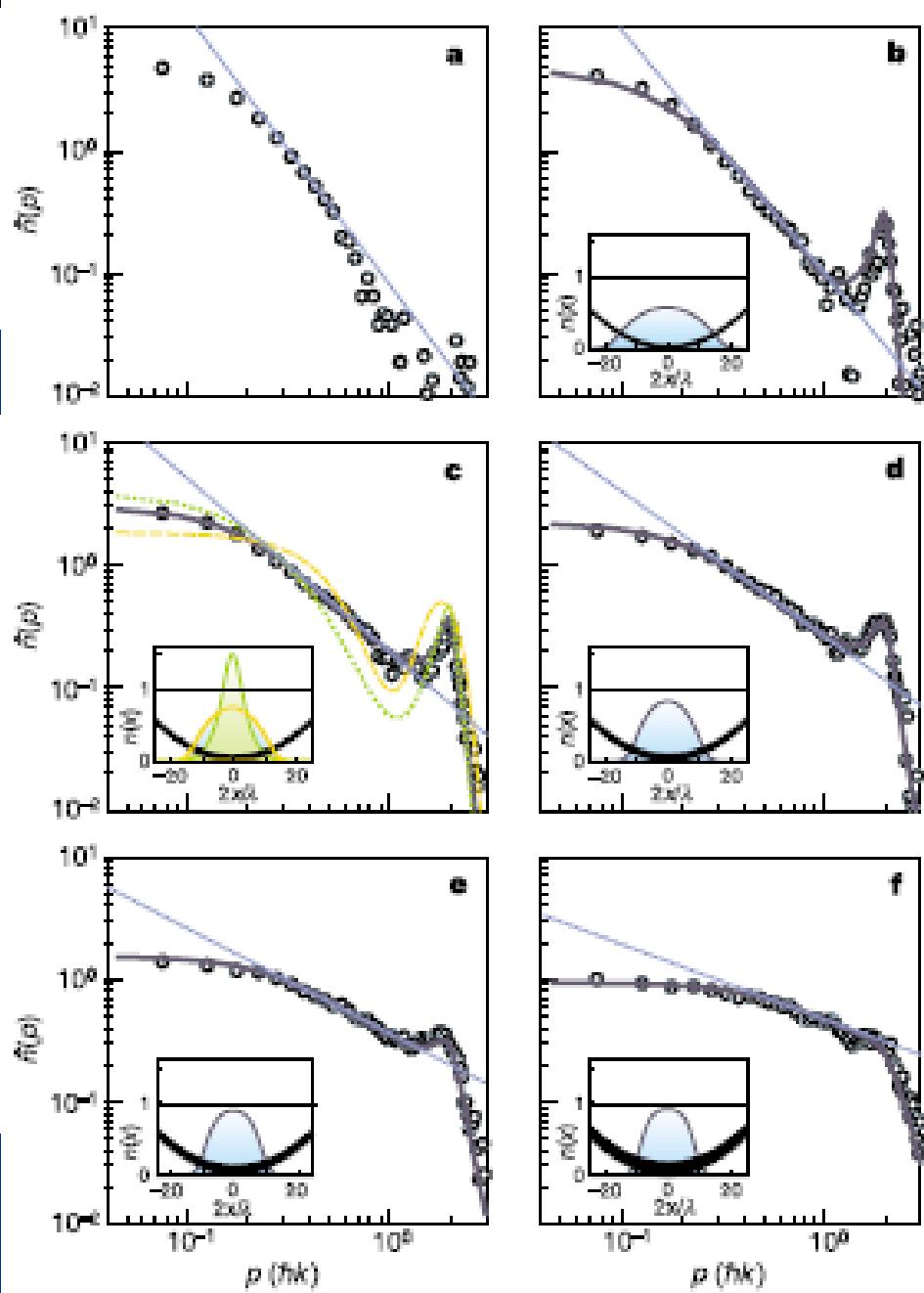
Observation of a One-Dimensional Tonks-Girardeau Gas

Toshiya Kinoshita, Trevor Wenger, David S. Weiss*

SCIENCE VOL 305 20 AUGUST 2004

1125

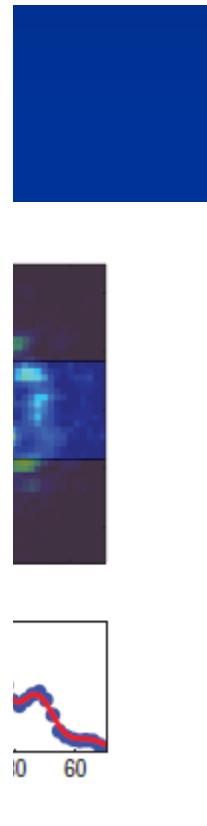
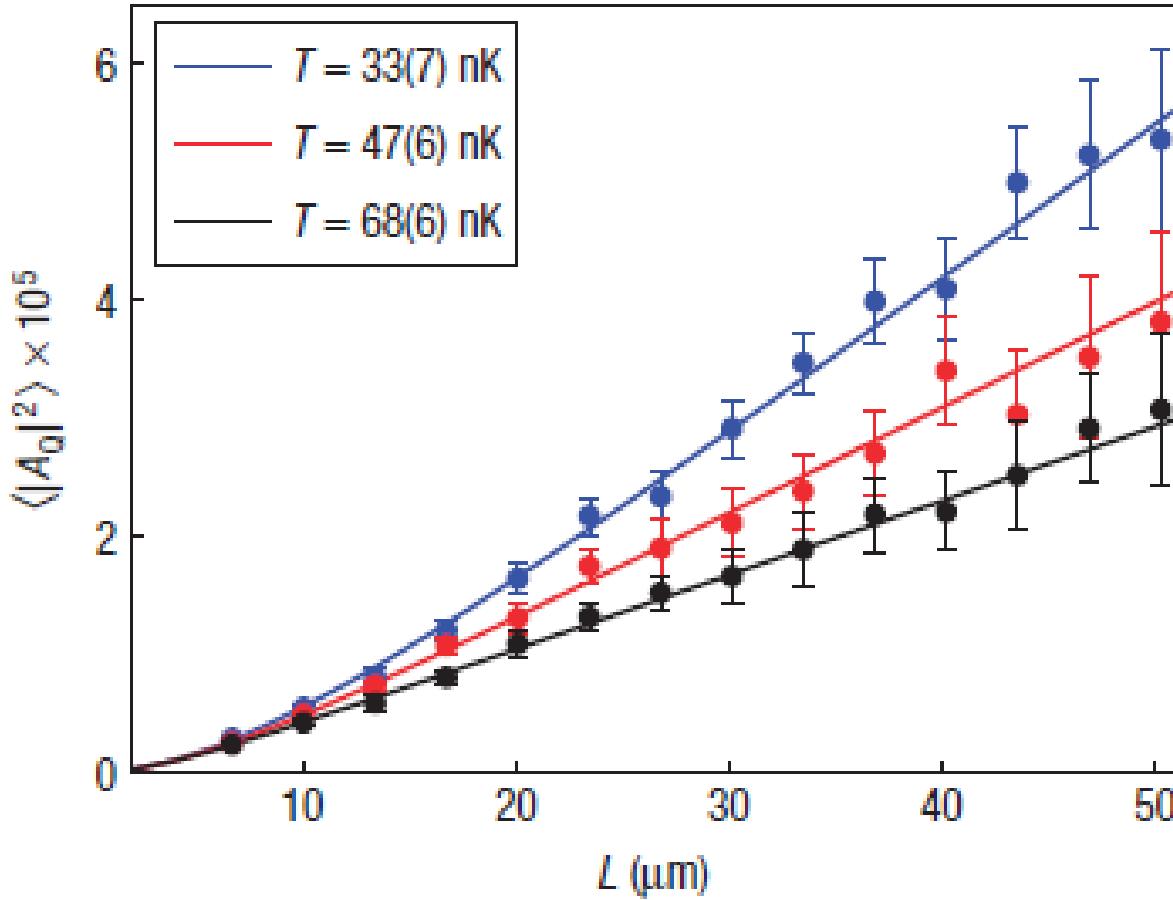




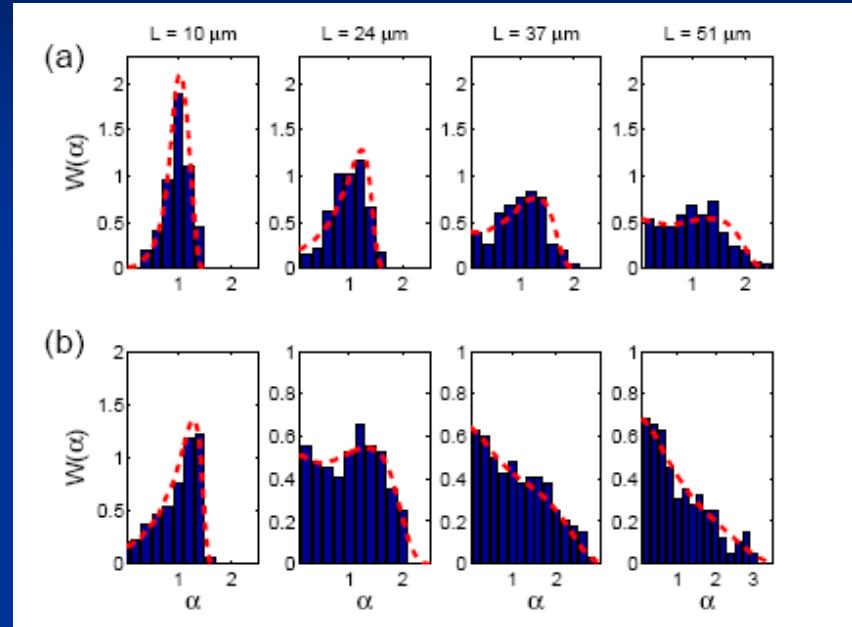
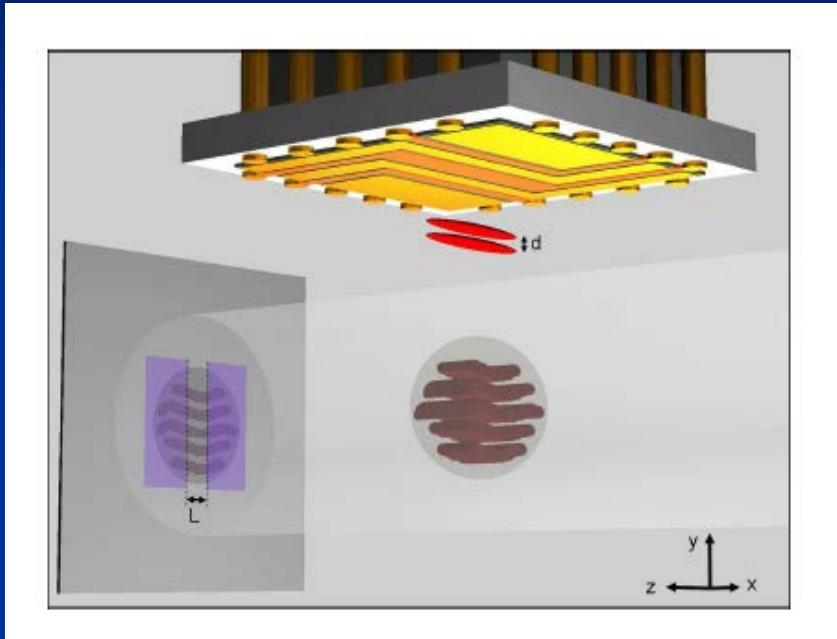
NATI

Probing quantum and thermal noise in an interacting many-body system

S. HOFFERBERTH^{1,2,3}, I. LESANOVSKY^{2,4}, T. SCHUMM¹, A. IMAMBEKOV^{3,5}, V. GRITSEV³, E. DEMLER³
AND J. SCHMIEDMAYER^{1,2*}



Cold atoms: Interferences



$$\int_0^L dr \langle \psi(r) \psi^\dagger(0) \rangle$$

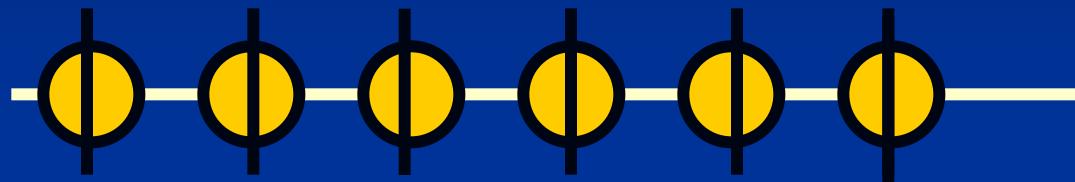
but K large (42)

S. Hofferberth et al. Nat. Phys 4
489 (2008)

hydrodynamic ?

Mott transition

Lattice: Mott transition

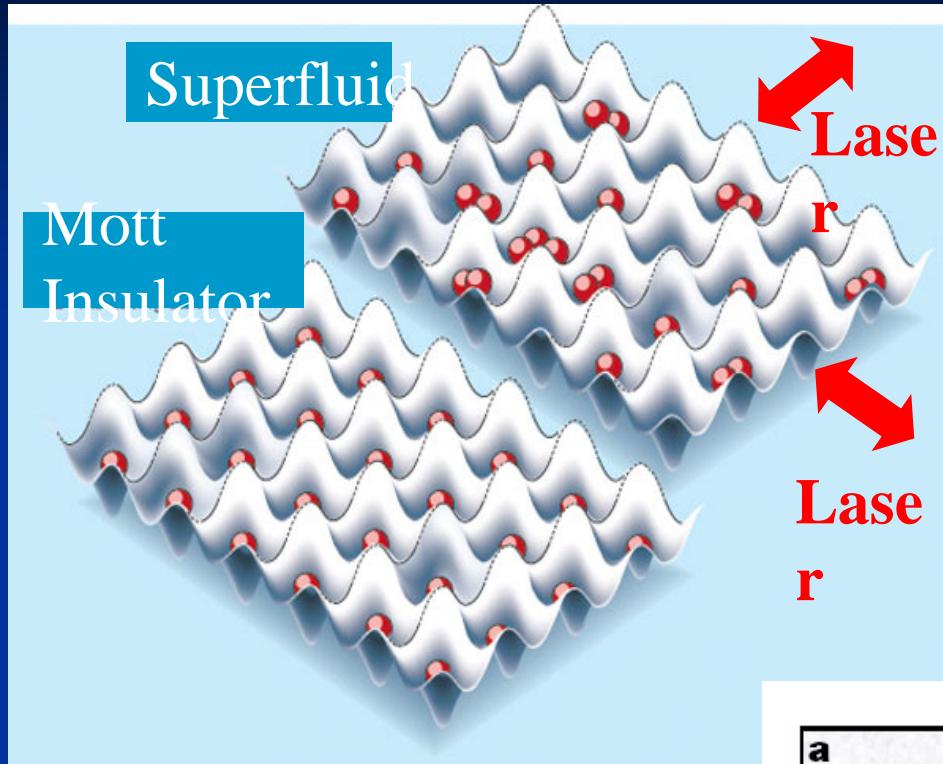


Costs U

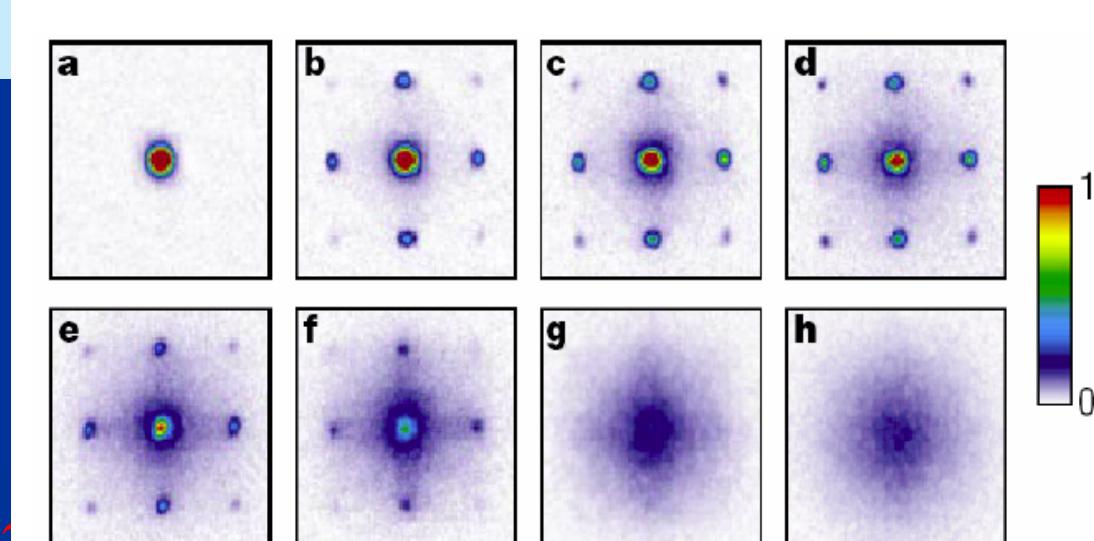
Quantum phase transition



Mott transition and cold atoms



Superfluid to Mott insulator
transition in a 3D optical
lattice
[D. Jaksch et al. PRL 81 (1998)]
[M Greiner et al. Nature, 415 (2002)]



Seminar **BINGO!**

To play, simply print out this bingo sheet and attend a departmental seminar.

Mark over each square that occurs throughout the course of the lecture.

The first one to form a straight line (or all four corners) must yell out

BINGO!!



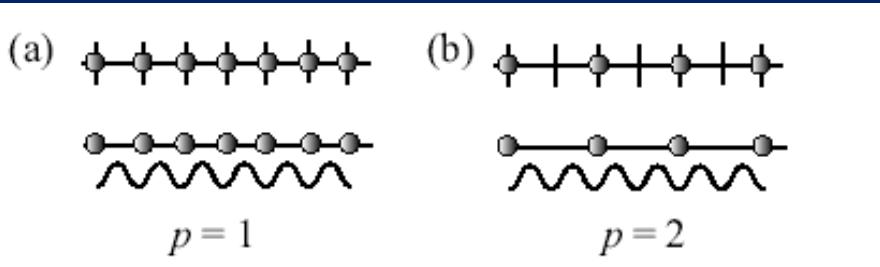
SEMINAR **B I N G O**

Speaker bashes previous work	Repeated use of "um..."	Speaker sucks up to host professor	Host Professor falls asleep	Speaker wastes 5 minutes explaining outline
Laptop malfunction	Work ties in to Cancer/HIV or War on Terror	"... et al."	You're the only one in your lab that bothered to show up	Blatant typo
Entire slide filled with equations	"The data clearly shows..."	FREE Speaker runs out of time	Use of Powerpoint template with blue background	References Advisor (past or present)
There's a Grad Student wearing same clothes as yesterday	Bitter Post-doc asks question	"That's an interesting question"	"Beyond the scope of this work"	Master's student bobs head fighting sleep
Speaker forgets to thank collaborators	Cell phone goes off	You've no idea what's going on	"Future work will..."	Results conveniently show improvement

JORGE CHAM © 2007

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How to treat?



$$H = \int dx V_0 \cos(Qx) \rho(x)$$

$$H = \int dx V_0 \cos(Qx) \rho_0 e^{i(2\pi\rho_0 x - 2\phi(x))}$$

- Incommensurate: $Q \neq 2 \pi \rho_0$

$$H = \int dx \cos(2\phi(x) + \delta x)$$

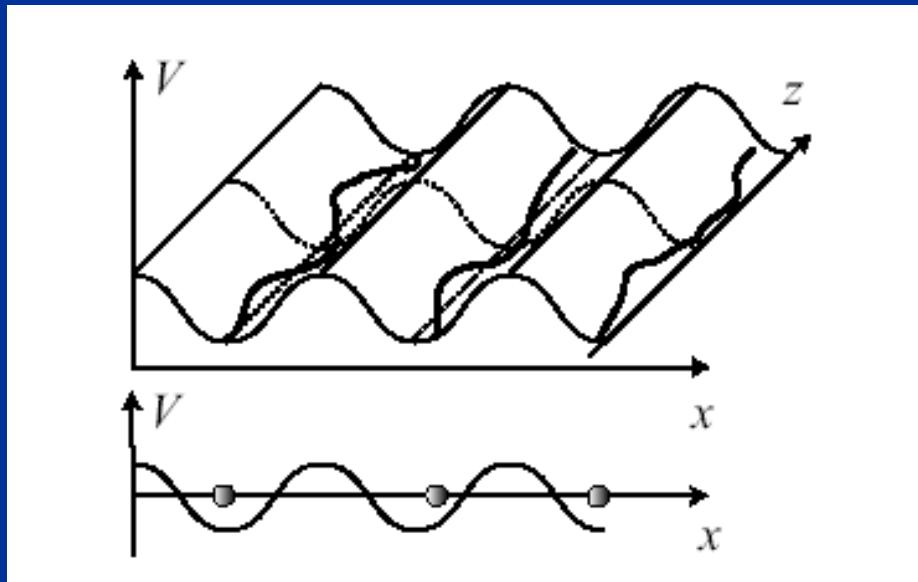
- Commensurate: $Q = 2 \pi \rho_0$

$$H = \int dx \cos(2\phi(x))$$

Competition

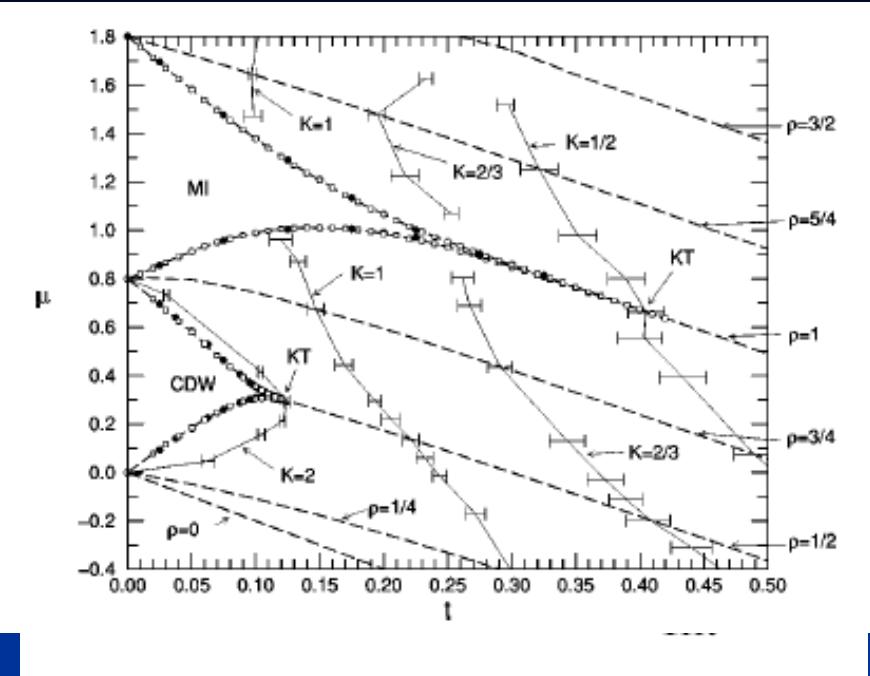
$$S_0 = \int \frac{dxd\tau}{2\pi K} [\frac{1}{u} (\partial_\tau \phi(x, \tau))^2 + u (\partial_x \phi(x, \tau))^2]$$

$$S_L = -V_0 \rho_0 \int dx d\tau \cos(2\phi(x))$$



Beresinski-
Kosterlitz-Thouless
transition

K=2



Mott insulator:
 ϕ is locked

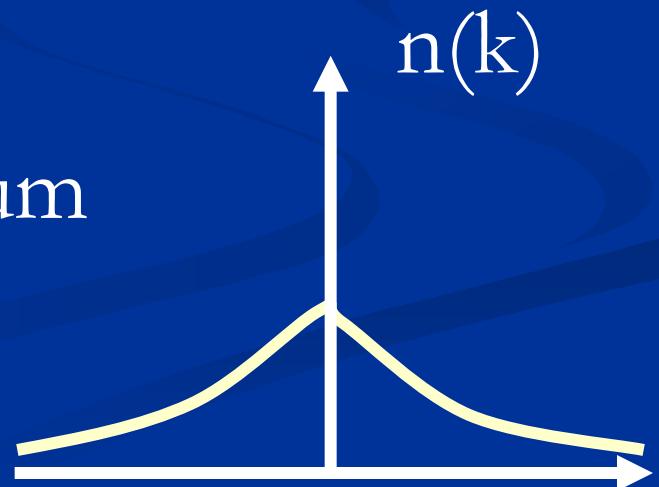
Density is fixed

TG, Physica B 230 975 (97)

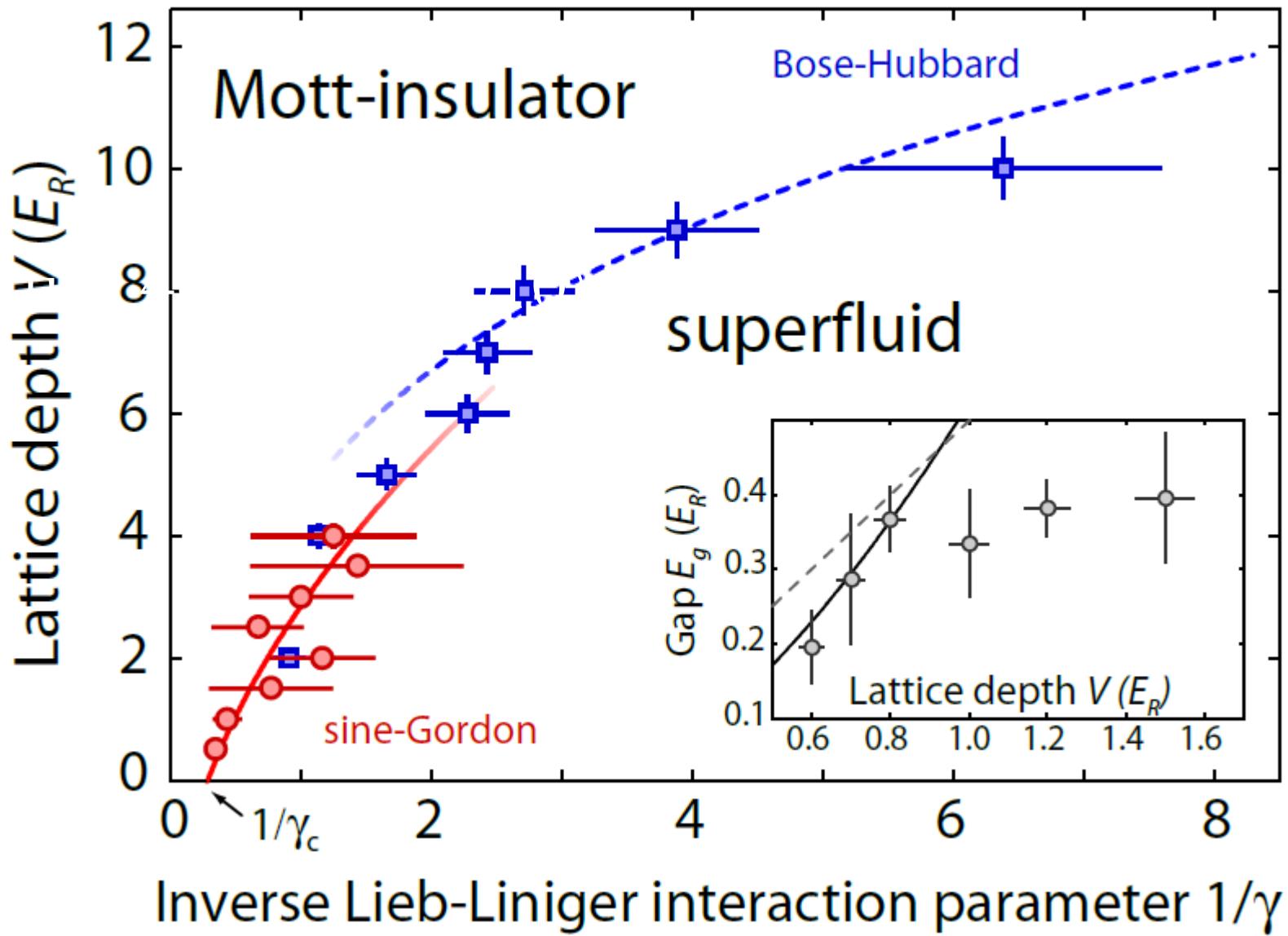
T. Kuhner et al. PRB 61 12474 (2000)

Gap in the excitation spectrum

$$G(r) \propto e^{-r/\xi}$$



Pinning quantum phase transition for a Luttinger liquid of strongly



Other 1D systems

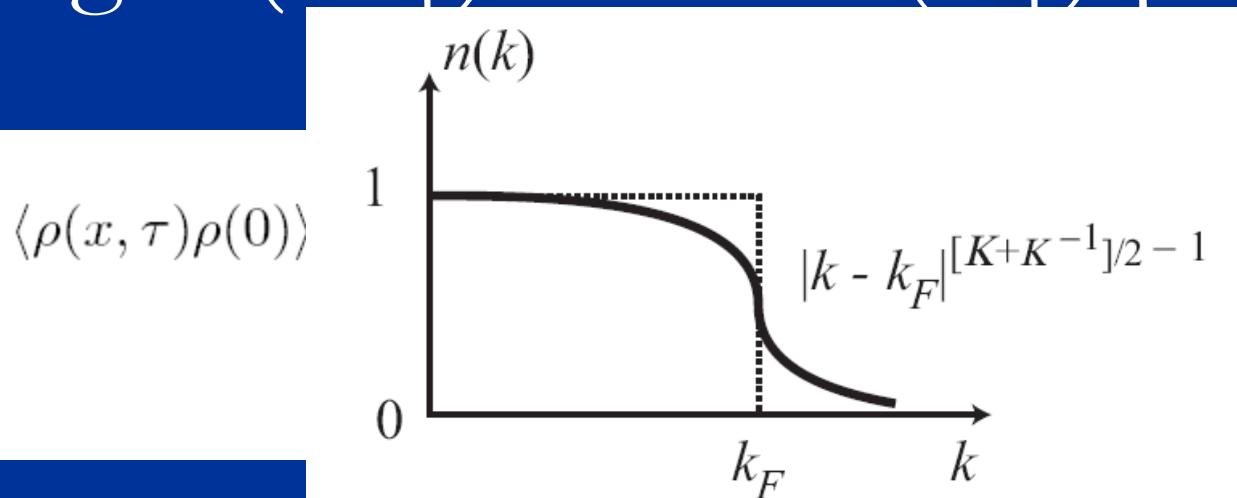


Fermions

$$\psi_F^\dagger(x) = \psi_B^\dagger(x) e^{i\frac{1}{2}\phi_l(x)}$$

$$\psi_F^\dagger(x) = [\rho_0 - \frac{1}{\pi} \nabla \phi(x)]^{1/2} \sum_p e^{i(2p+1)(\pi\rho_0 x - \phi(x))} e^{-i\theta(x)}$$

Right ($+k_F$) and left ($-k_F$) particles



$$\langle \rho(x, \tau) \rho(0) \rangle$$

$$\begin{aligned} & |k - k_F|^{[K+K^{-1}]/2 - 1} \rho_0 x \left(\frac{\alpha}{r}\right)^{2K} \\ & \pi \rho_0 x \left(\frac{\alpha}{r}\right)^{8K} + \dots \end{aligned}$$

Spins

Use boson or fermions mapping

$$S^+ = (-1)^i e^{i\theta} + e^{i\theta} \cos(2\phi)$$

$$S^z = \frac{-1}{\pi} \nabla \phi + (-1)^i \cos(2\phi)$$

Powerlaw correlation functions

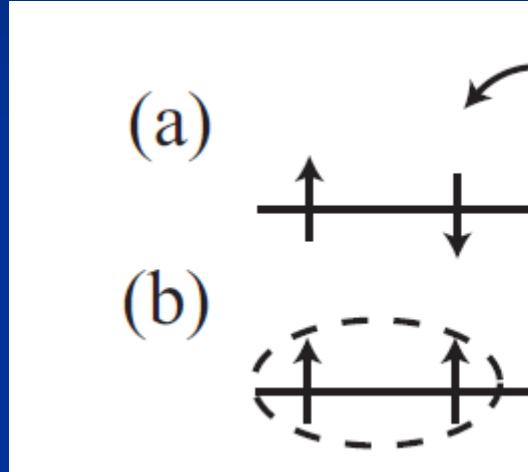
$$\langle S^z(x, 0) S^z(0, 0) \rangle = C_1 \frac{1}{x^2} + C_2 (-1)^x \left(\frac{1}{x}\right)^{2K}$$
$$\langle S^+(x, 0) S^-(0, 0) \rangle = C_3 \left(\frac{1}{x}\right)^{2K + \frac{1}{2K}} + C_4 (-1)^x \left(\frac{1}{x}\right)^{\frac{1}{2K}}$$

Non universal exponents K(h,J)

Fractionalization of excitations

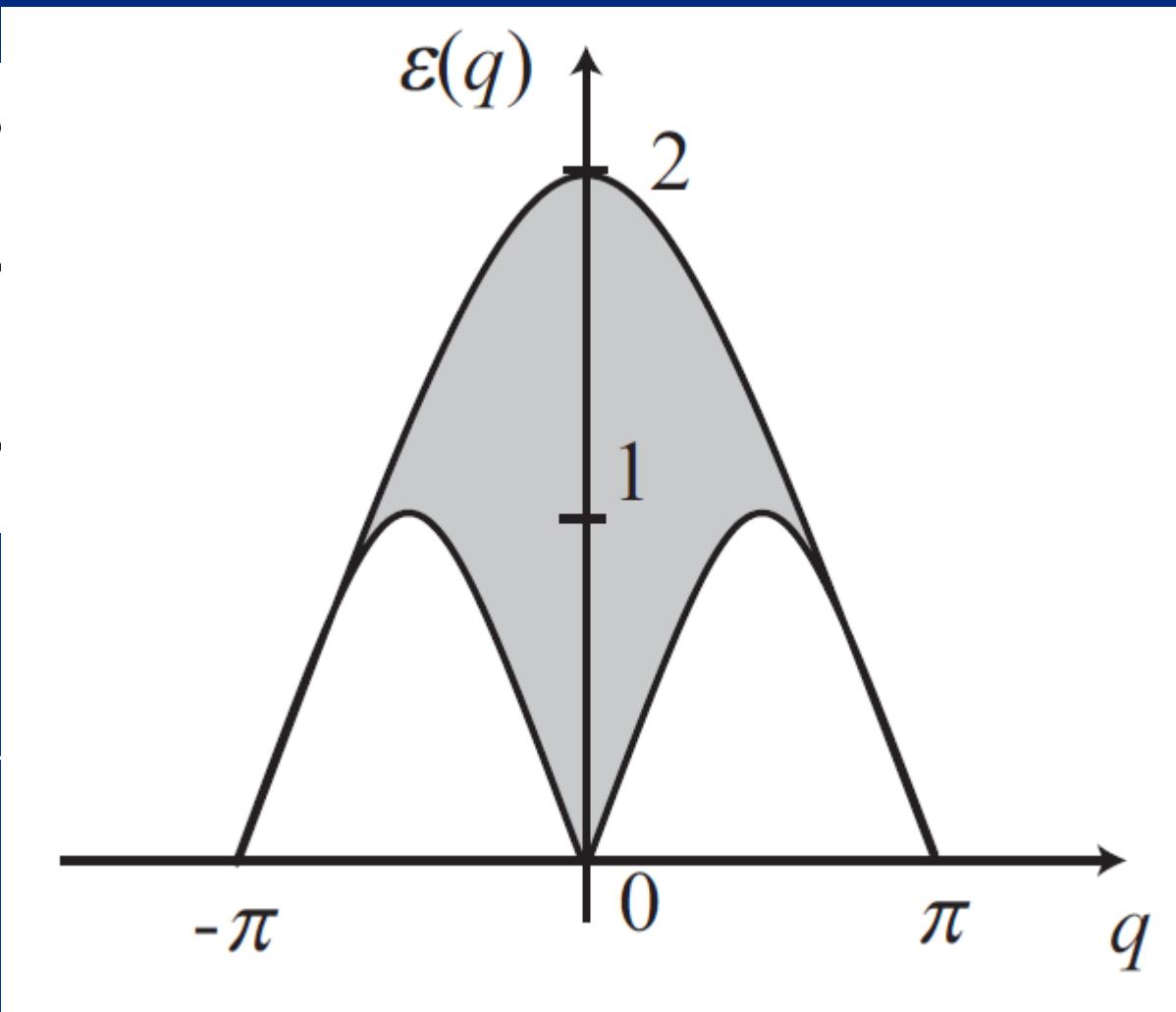


$$1 = \frac{1}{2} + \frac{1}{2}$$

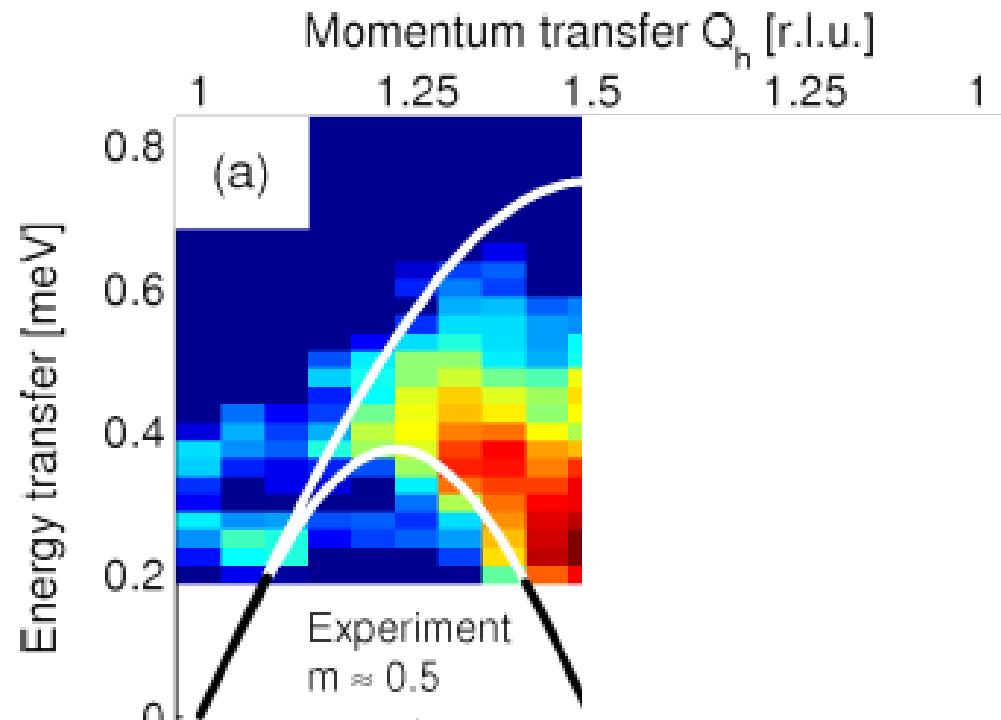


$$E(k) = \cos(k_1)$$

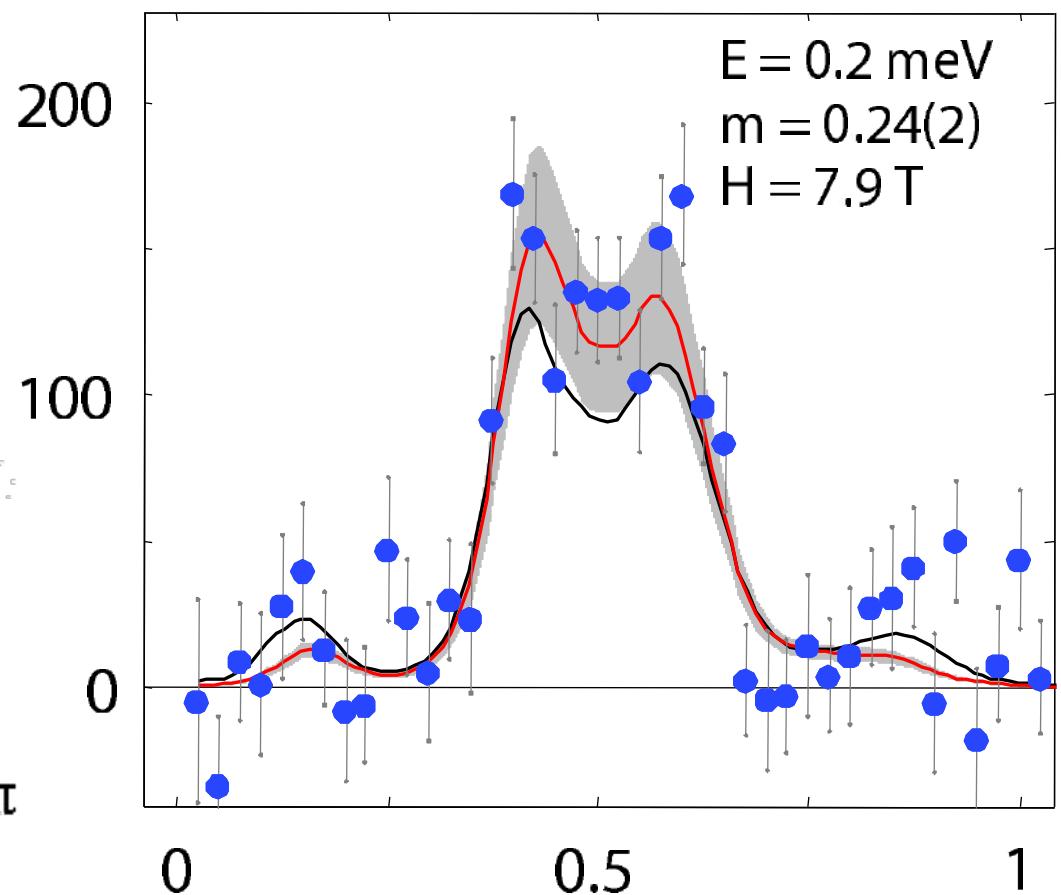
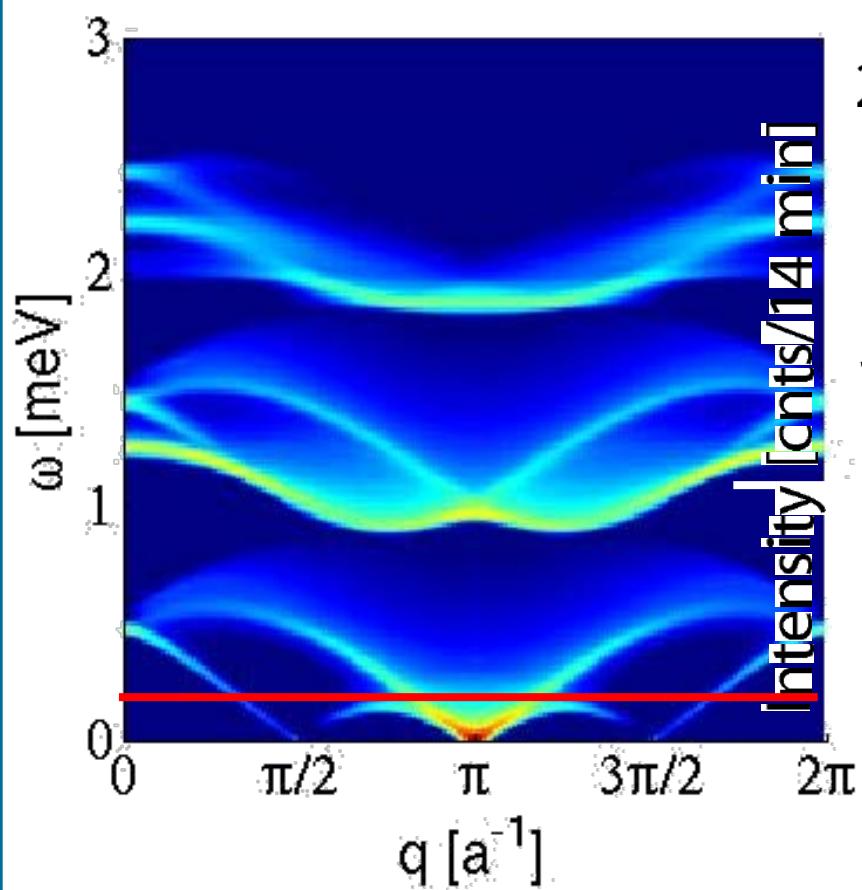
$$k = k_1 + k_2$$



Neutron scattering



B. Thielemann et al. PRL 102, 107204 (2009)

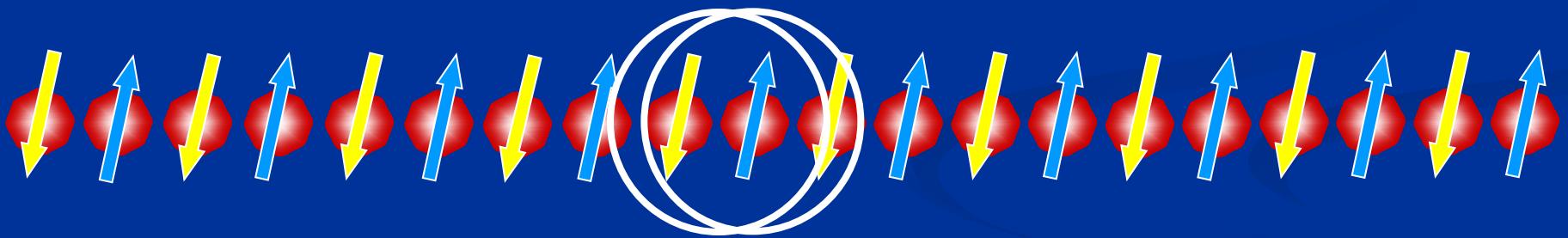


Systems with spins

Bosonize each spin separately

Spin-Charge Separation

Spin



Spinon

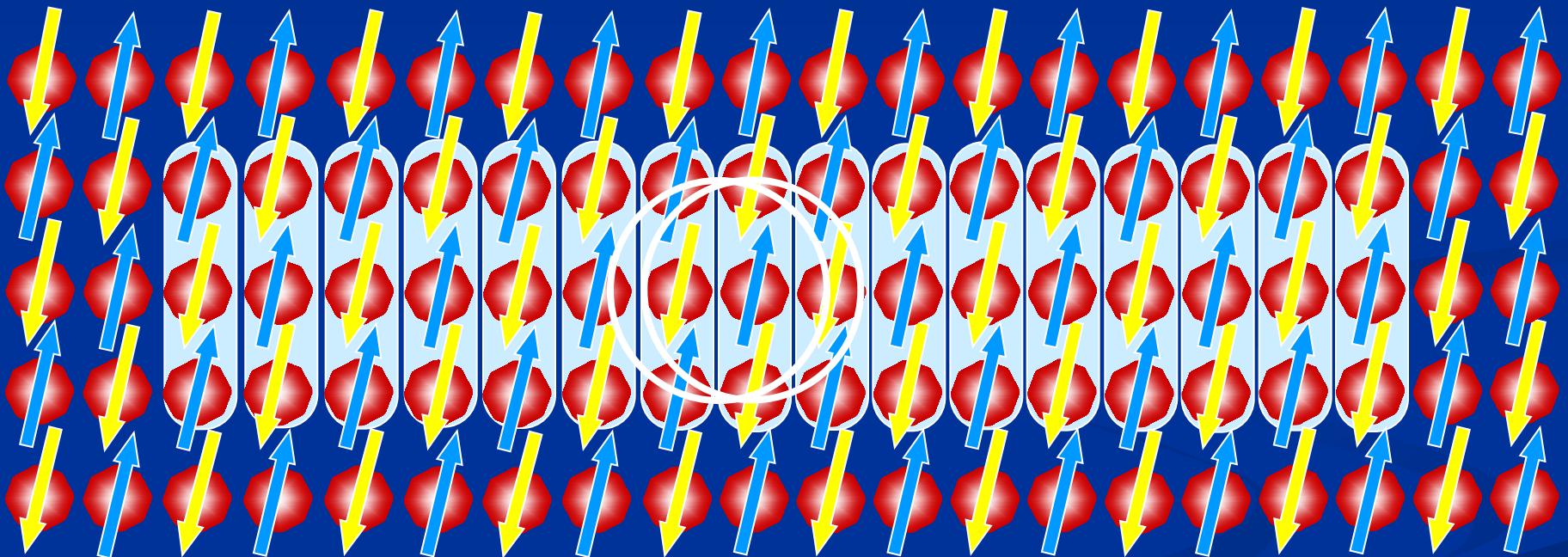
Charge

Holon

Spin-Charge Separation higher D ?

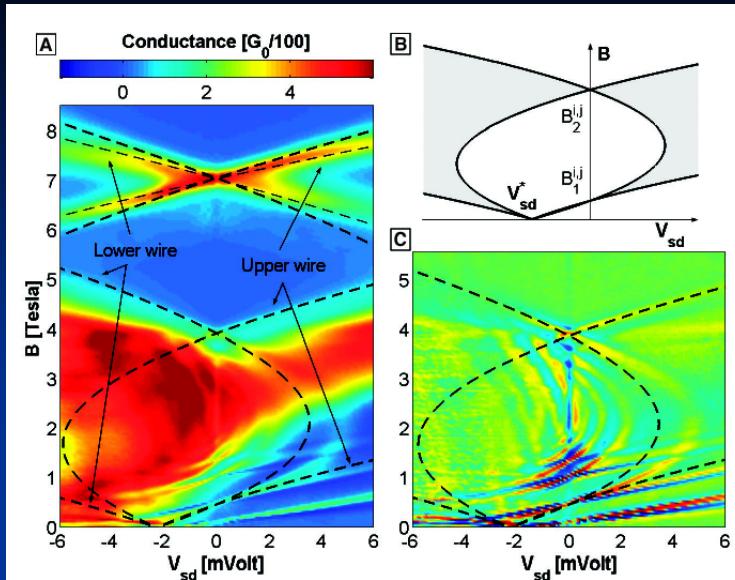
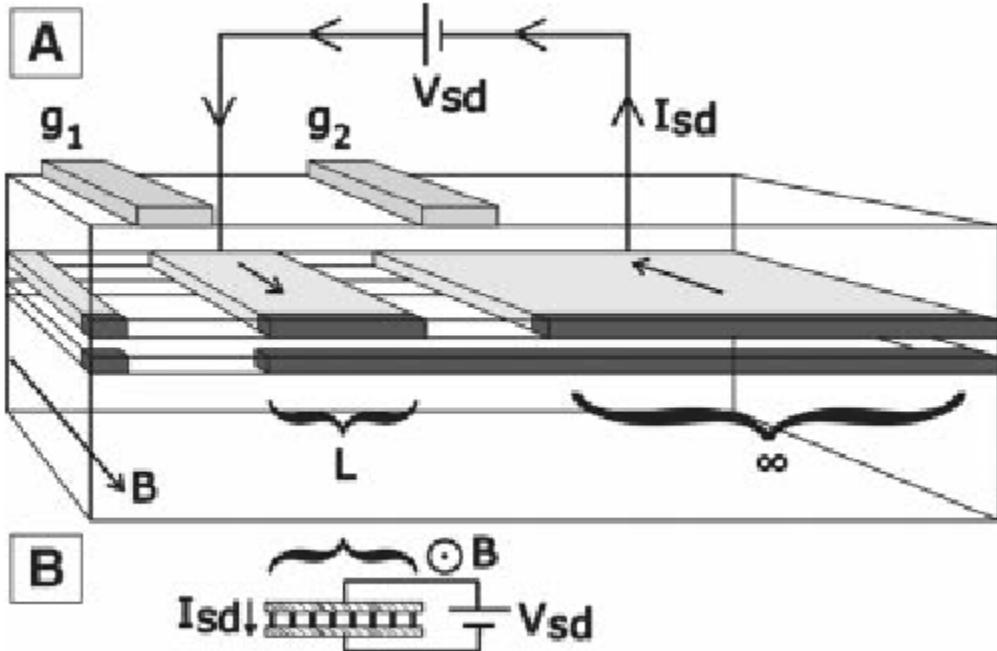
Spin

Charge



Energy increases with spin-charge separation

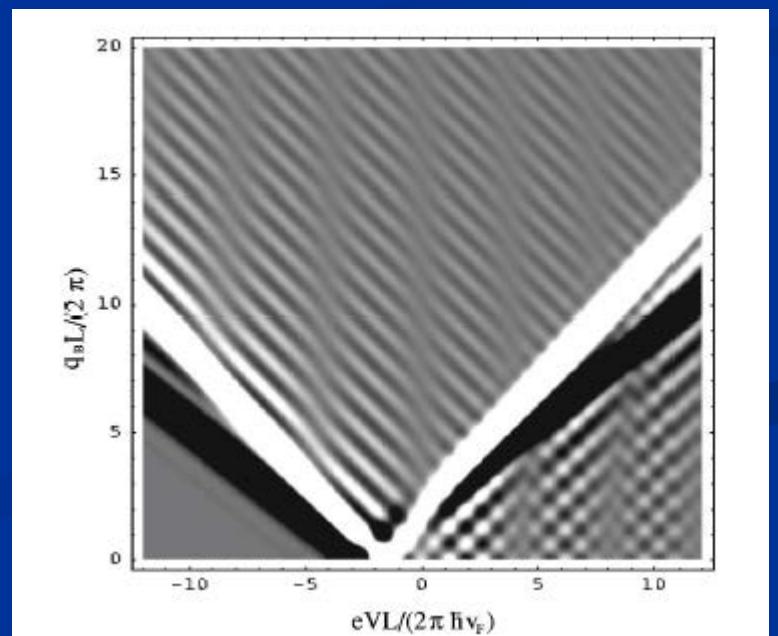
Confinement of spin-charge: « quasiparticle »



O.M Ausslander et al., Science
298 1354 (2001)

Y. Tserkovnyak et al., PRL 89
136805 (2002)

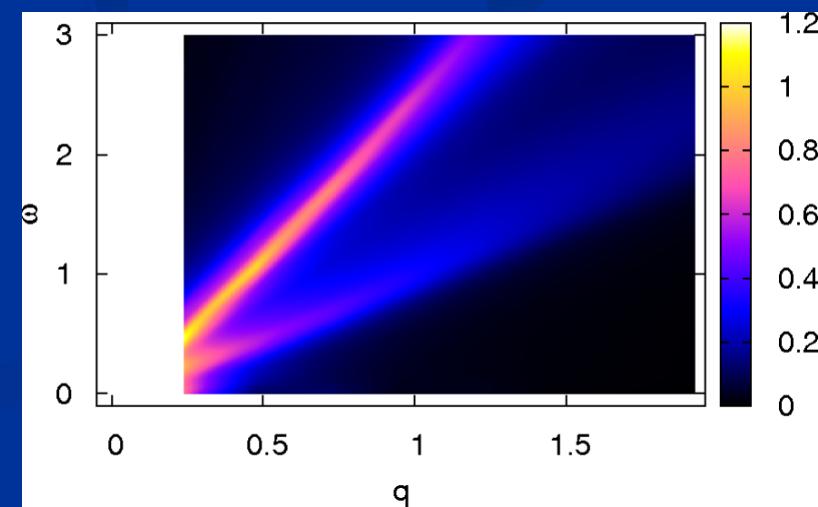
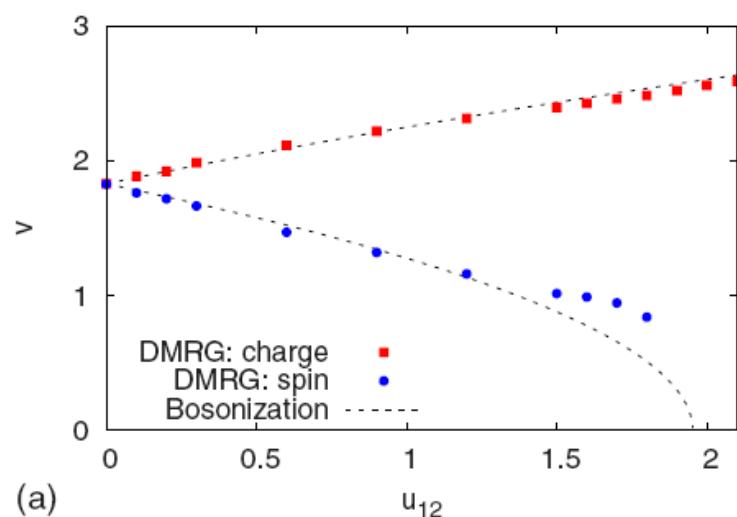
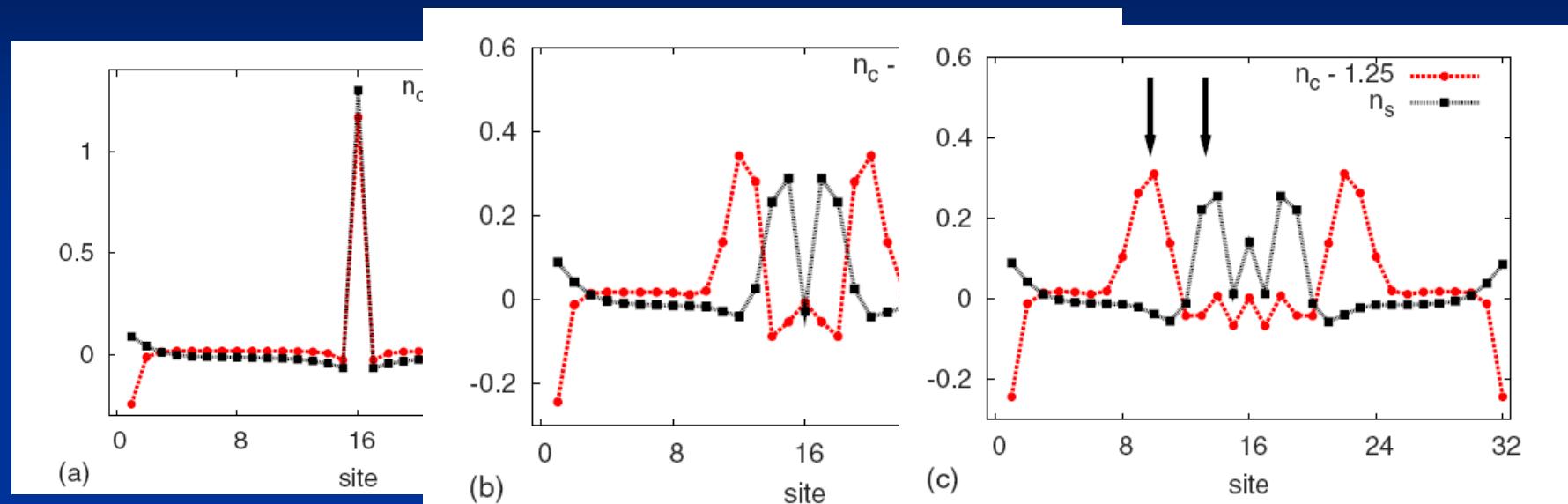
Y. Tserkovnyak et al., PRB 68
125312 (2003)



Proposal

- ^{87}Rb $|F=2, m_F=-1\rangle$ and $|F=2, m_F=1\rangle$ good potential candidates to observe spin-charge separation
- No problem of temperature (\neq fermions)
- Phase separation

A. Kleine, C. Kollath et al. PRA 77 013607 (2008);
 NJP 10 045025 (2008)



Beyond Luttinger liquids



To boldly go where no theorist
has gone before....



Out of equilibrium Luttinger liquids

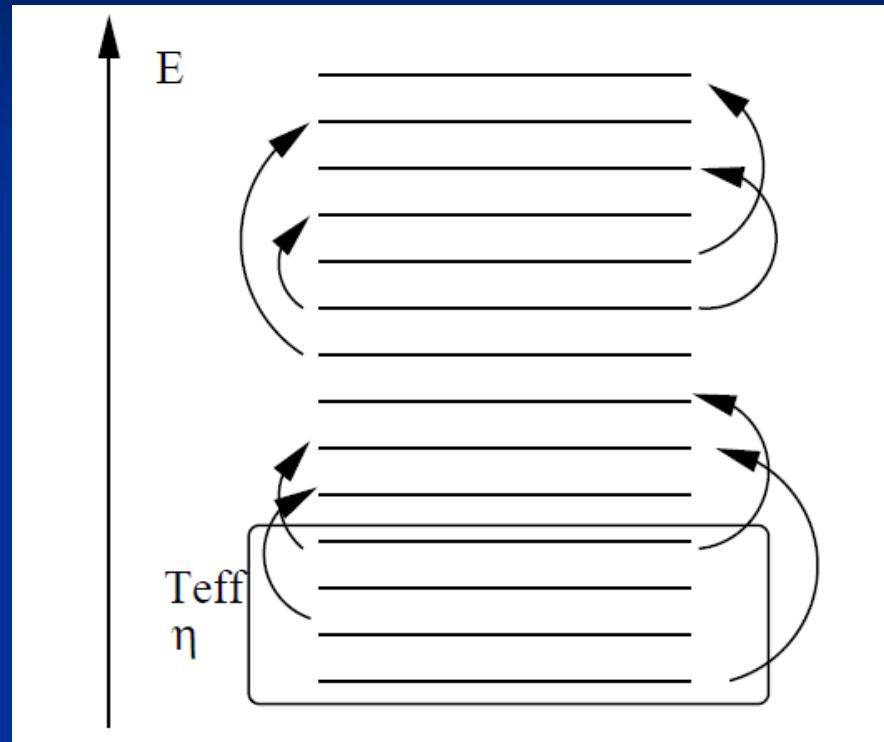
Quenches and thermalization in LL

A. Mitra, TG,

Phys. Rev. Lett. 107, 150602
(2011)

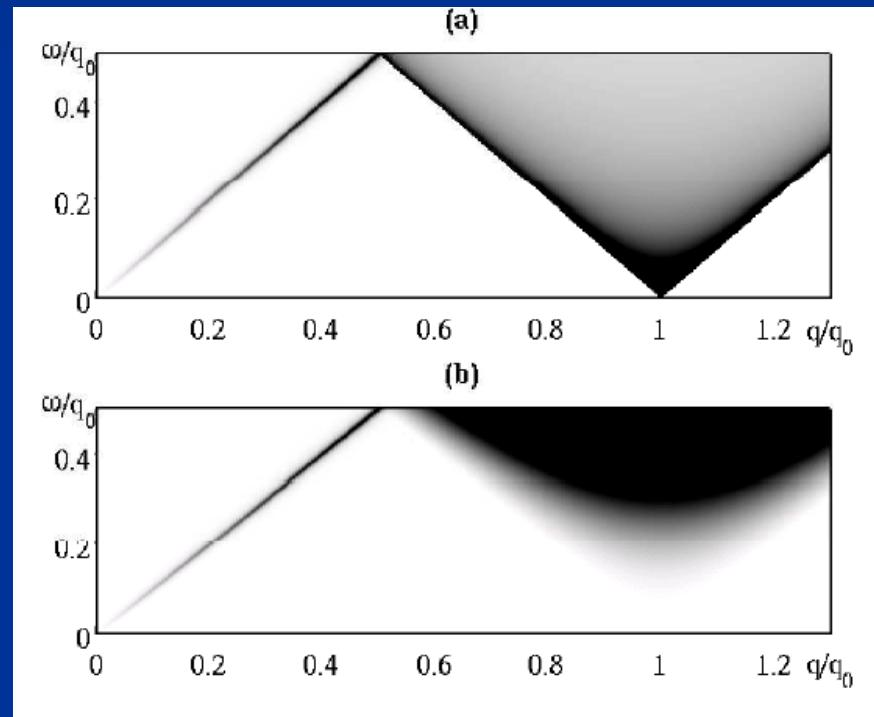
Phys. Rev. B 85, 075117 (2012)

J. Lancaster, TG, A Mitra,
Phys. Rev. B 84, 075143 (2011)



Out of equilibrium sine-Gordon
Thermalization for the low energy part

LL + external noise



E. Dalla Torre, E. Demler, TG, E. Altman, Nat. Phys. 6 806
(2010); PRB 85 184302 (2012)

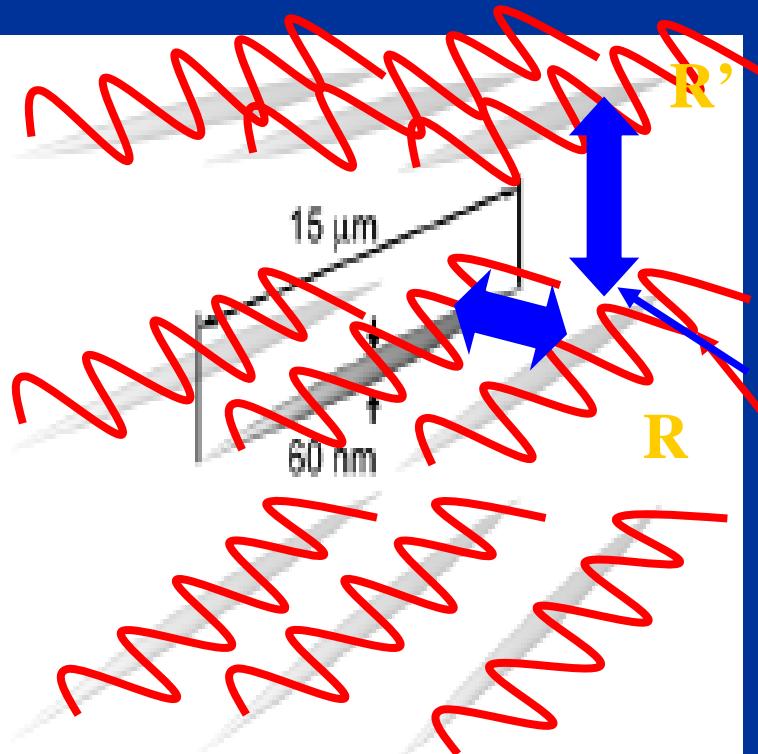
Coupled 1D chains: deconfinement



Mott vs. Josephson

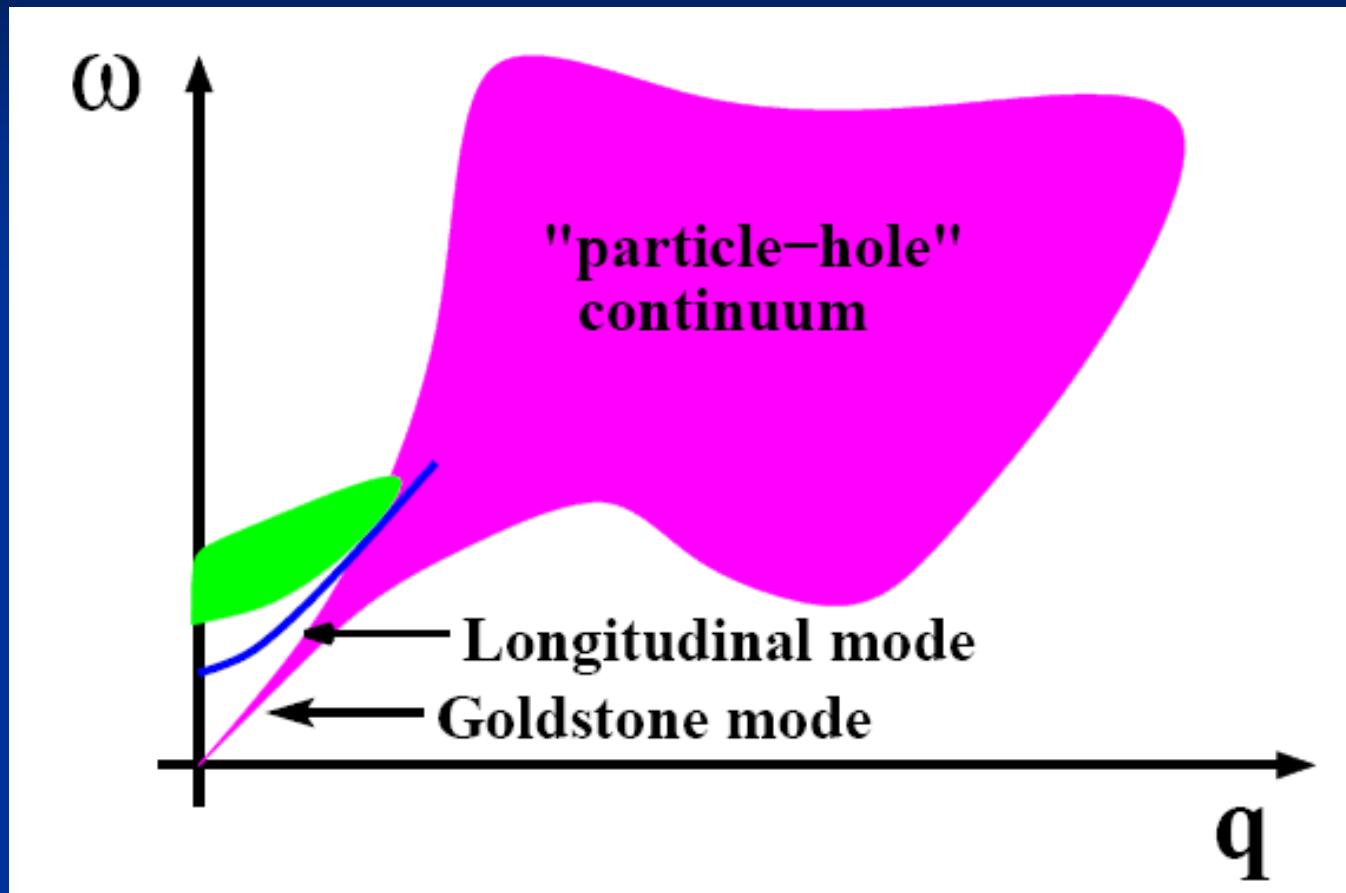
$$H_{\text{eff}} = \frac{\hbar v_s}{2\pi} \sum_{\mathbf{R}} \int_0^L dx \left[\frac{1}{K} (\partial_x \phi_{\mathbf{R}}(x))^2 + K (\partial_x \theta_{\mathbf{R}}(x))^2 \right]$$

$$\begin{aligned} & - \frac{\hbar v_s g_J}{2\pi a^2} \sum_{\langle \mathbf{R}, \mathbf{R}' \rangle} \int_0^L dx \cos(\theta_{\mathbf{R}}(x) - \theta_{\mathbf{R}'}(x)) \\ & + \frac{\hbar v_s g_u}{2\pi a^2} \sum_{\mathbf{R}} \int_0^L dx \cos(2\phi_{\mathbf{R}}(x) + \delta\pi x) \end{aligned}$$



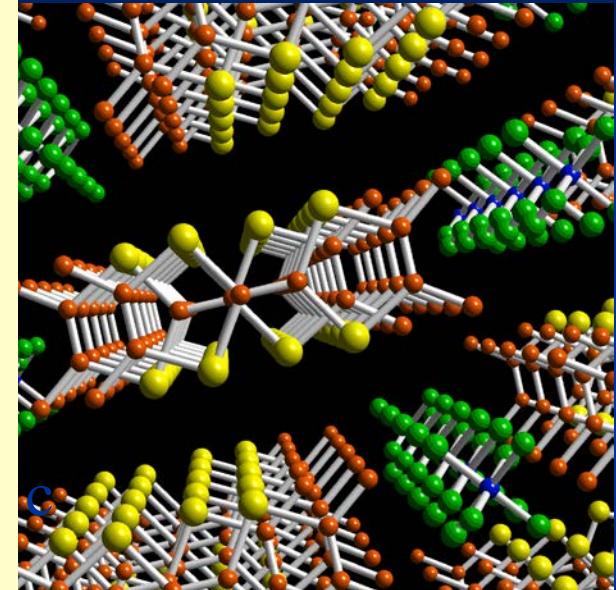
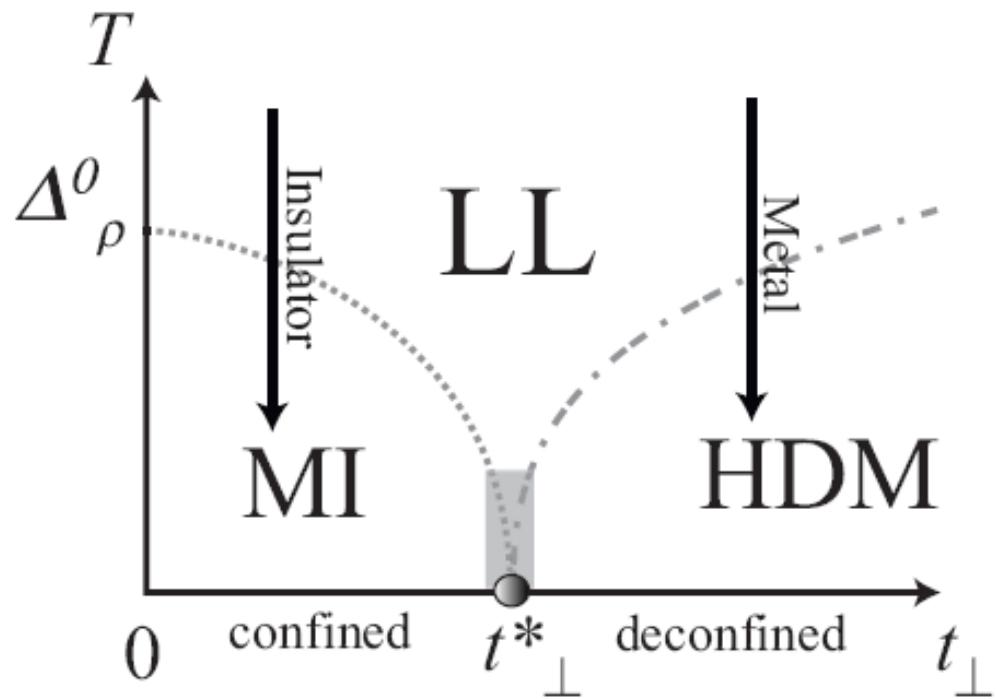
Josephson coupling: delocalizes atoms

“Mott” potential: localizes atoms



Higgs “amplitude” mode

Deconfinement

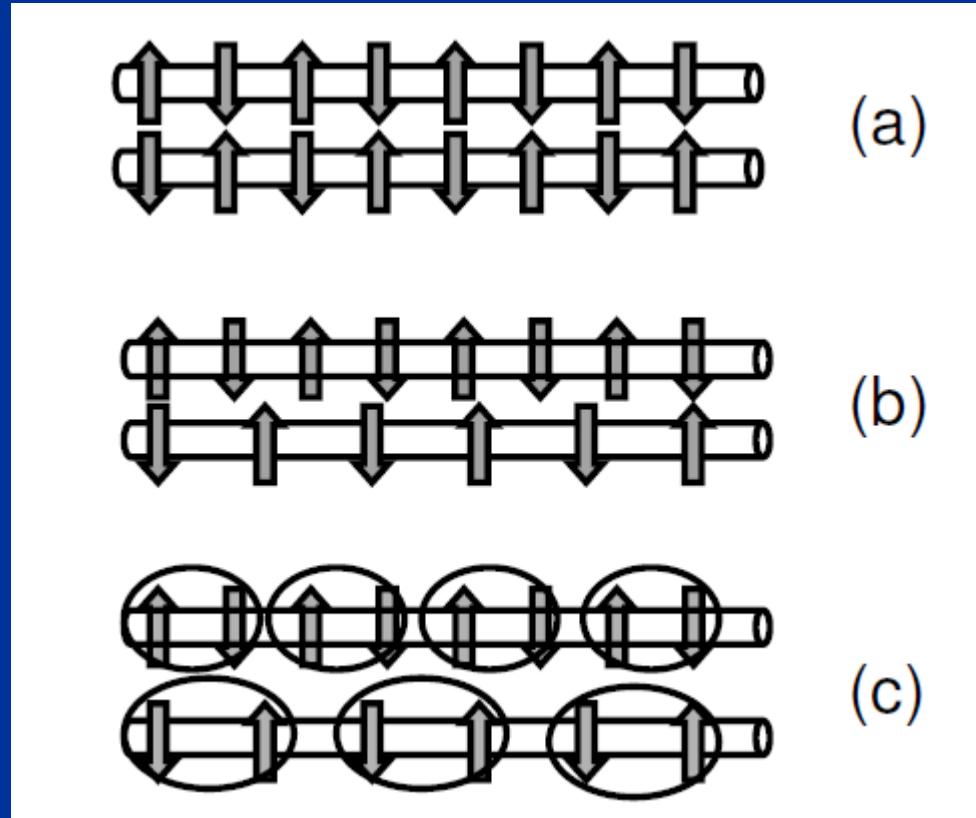


TG Chemical
Review 104 5037
(2004)

P. Auban-Senzier, D. Jérôme, C. Carcel and J.M. Fabre J de Physique IV, (2004)
A. Pashkin, D. Macéard et al., J. Phys. C, CL89 (2001)

Cold atoms

AF Ho, M.A. Cazalilla, TG, PRL **95**, 226402 (2005);
Arxiv/0604525



Triplet
superconductor
(repulsive
interactions)

Non-LL and impurity problems

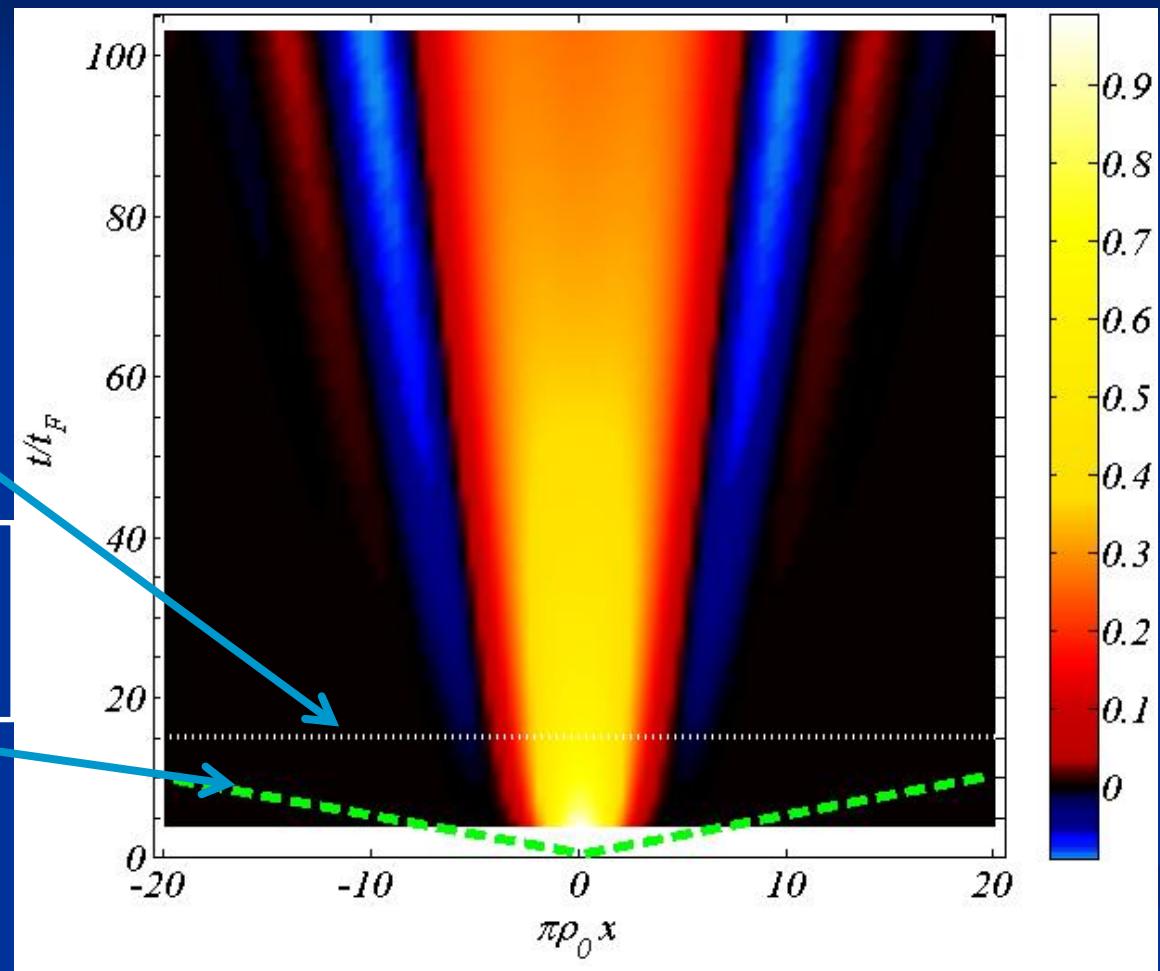


- Interacting 1D system (Luttinger liquid):
e.g. bosons -- ``spin up''
 - Add one particle of ``spin down'', or flip a spin
- $$G(x,t) = \langle S^+(x,t)S^-(0,0) \rangle$$
- Example: two component bose-Hubbard model

M. B. Zvonarev, V. V. Cheianov, TG, PRL 99 240404
 (2007)

Trapped/open
 regimes

Light cone of
 spinless bosons



$$G_{\perp}(x, t) \simeq \frac{1}{\sqrt{\ln(t/t_F)}} \exp \left\{ -\frac{1}{K} \frac{(\pi\rho_0 x)^2}{2 \ln(t/t_F)} \right\}.$$

$$G_{\perp} \simeq e^{-(x^2/2\ell^2)} t^{-\alpha} G_{\perp}^H, \quad \ell(t) = \frac{2K^{-(1/2)}}{\pi\rho_0} \frac{t/t_F}{\sqrt{\ln t/t_F}} \frac{m}{m_*}.$$

Theories:
Furusaka
Lamacra

Experi-
Innsbruc

Matveev,
Engardt,

e,



Driven impurity

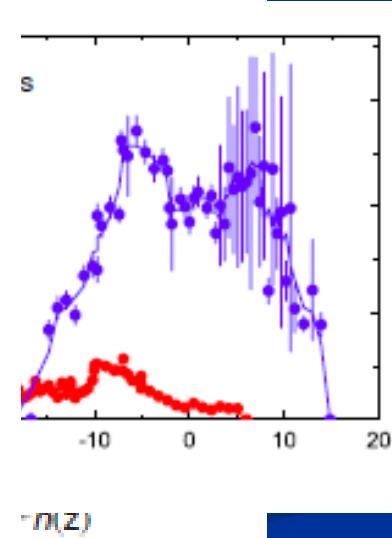
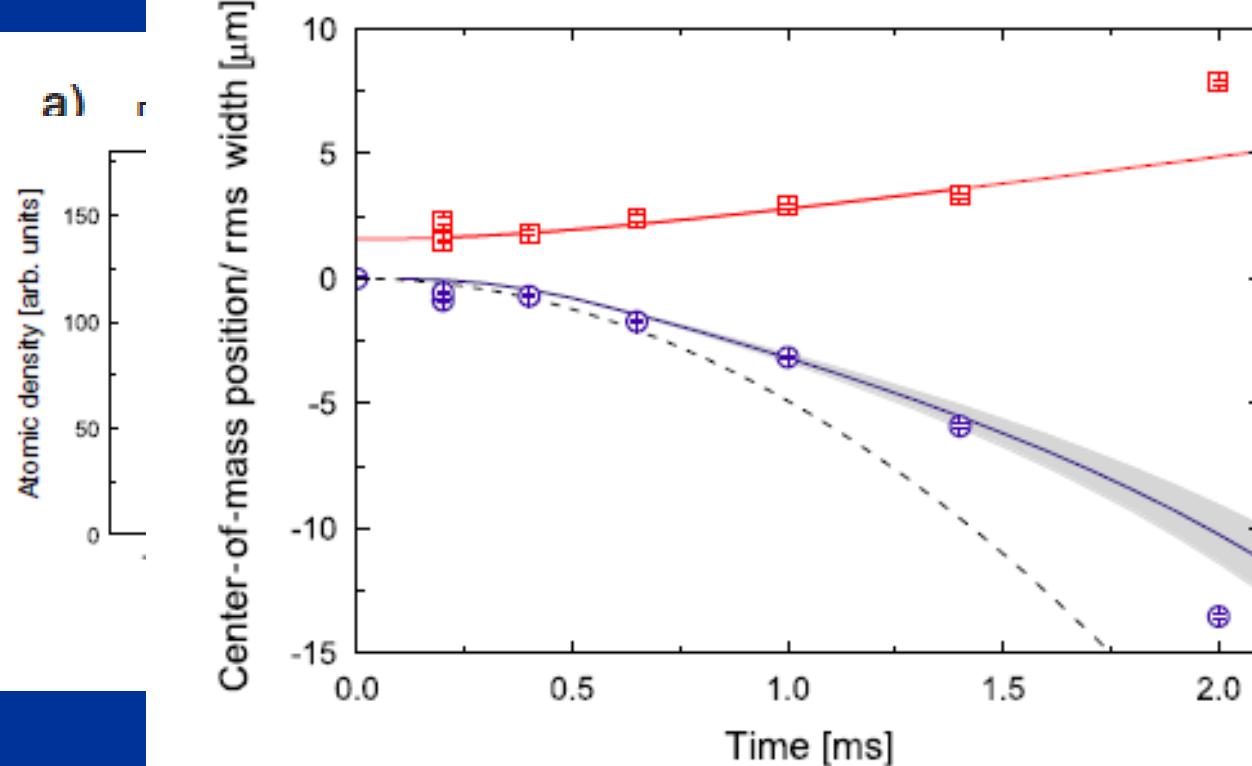
PRL 103, 150601 (2009)

PHYSICAL REVIEW LETTERS

week ending
9 OCTOBER 2009

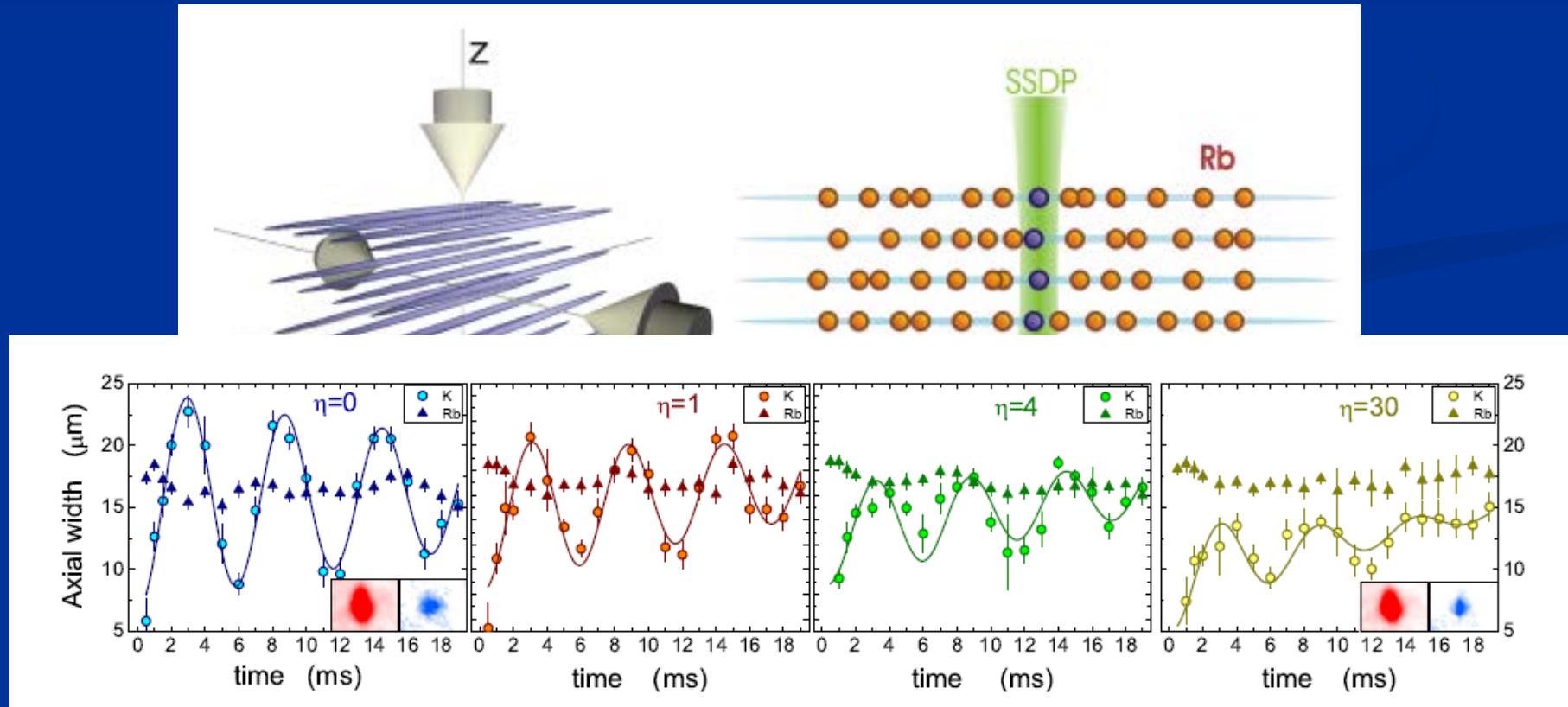
Quantum Transport through a Tonks-Girardeau Gas

Stefan Palzer, Christoph Zipkes, Carlo Sias,* and Michael Köhl



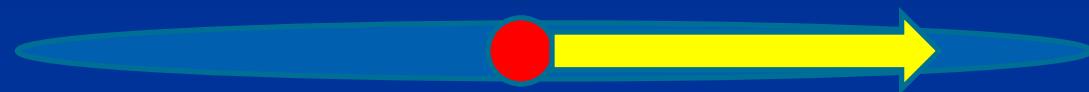
Diffusive impurity

J. Catani et al. PRA 85, 023623 (2012)



Driven impurity vs diffusion

- Normal transport



$$v = \mu F$$

$$v = f(F)$$



$$\langle x^2 \rangle \sim Dt$$

$$\langle x^2 \rangle \sim \log(t)$$

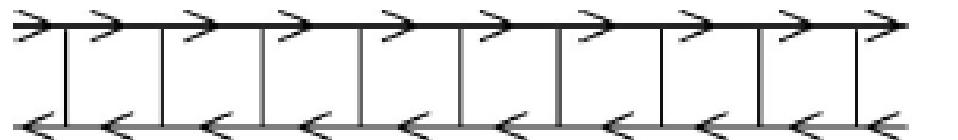
- Einstein relation: $\mu = D$

Artificial Gauge fields

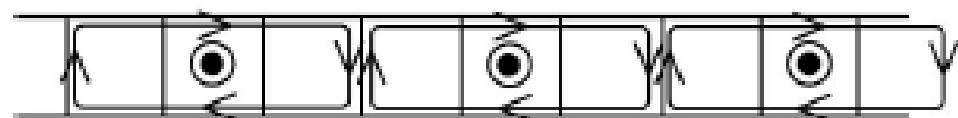


Meissner effect in a bosonic ladder

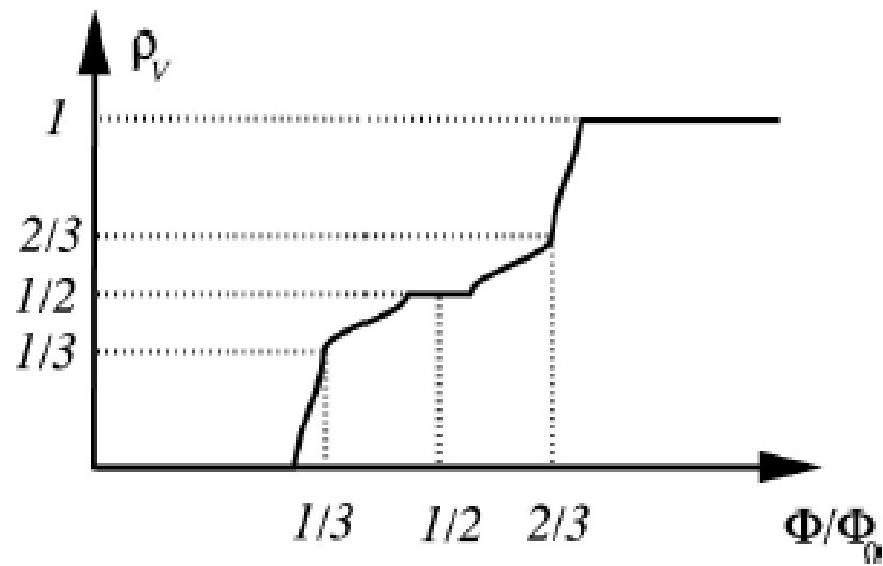
E. Orignac + TG



(a)



(b)



Disorder and interaction effects



Dirty bosons

TG + H. J. Schulz EPL 3 1287 (1987); PRB 37 325 (1988)

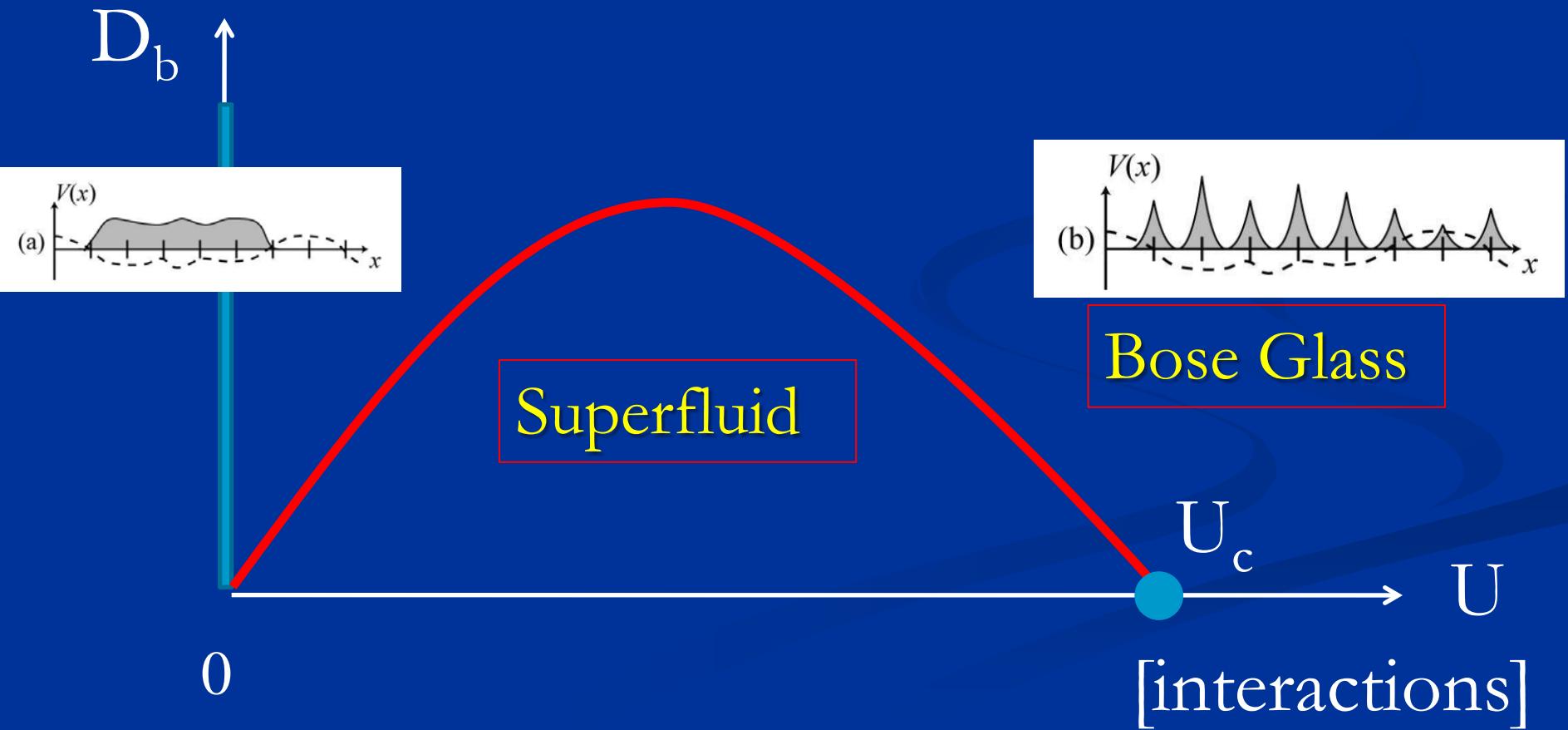
$$H_{\text{dis}} = \int dx V(x) \rho(x)$$

$$H_{\text{dis}} = \int dx V(x) \left[-\frac{1}{\pi} \nabla \phi(x) + \rho_0 (e^{i(2\pi\rho_0 x - 2\phi(x))} + \text{h.c.}) \right]$$

``Two'' fourier components of disorder

Bose glass phase

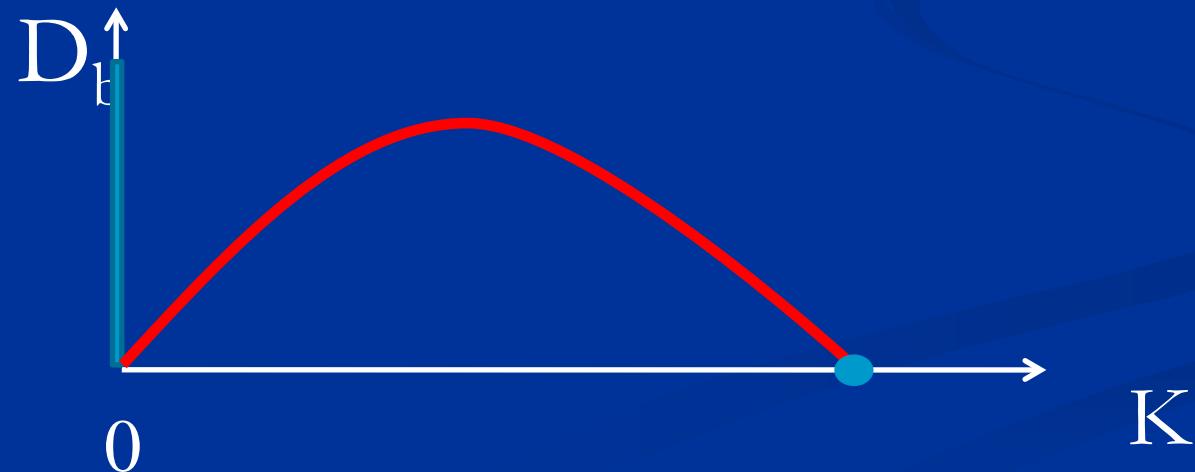
TG + H. J. Schulz EPL 3 1287 (87); PRB 37 325 (1988);
M.P.A. Fisher et al. PRB 40 546 (1989)



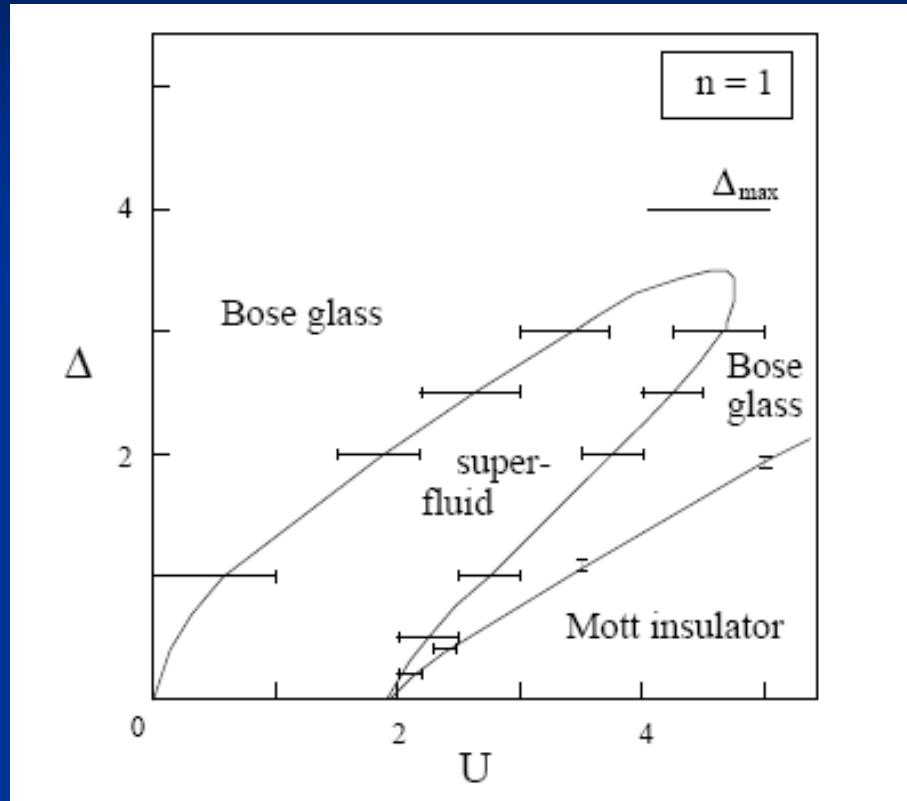
SU-BG transition in d=1

$$\langle \psi(r) \psi^\dagger(0) \rangle \sim \left(\frac{1}{r}\right)^{\frac{1}{2K}} \quad K \rightarrow 3/2$$

Universal exponent at the SU-BG transition !



Numerics



S. Rapsch, U. Schollwoeck,
W. Zwerger EPL 46 559
(1999);

G. Batrouni et al. PRL 65
1765 (90);
N. Prokofev et al. PRL 92
015703 (04);
O. Nohadani et al. PRL 95,
227201 (05)
K. G. Balabanyan et al. PRL
95, 055701 (05);
L. Pollet et al. PRL 103,
140402 (2009)

.....

Two loops RG for Bose glass

Z. Ristivojevic, A. Petkovic, P. Le Doussal, TG PRL (2012)

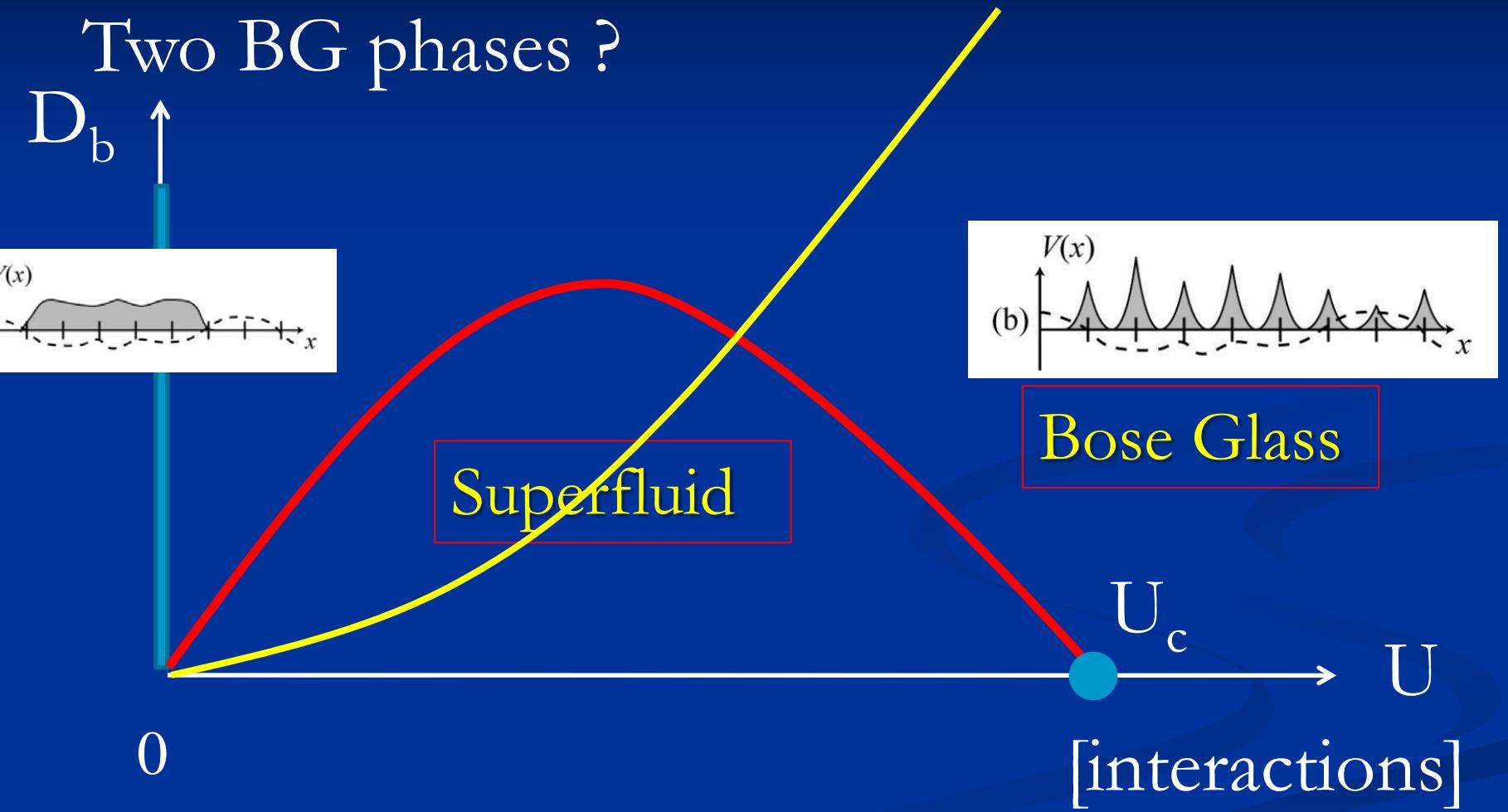
$$\Gamma_2^{(1)} = - \frac{d\mathcal{D}_R}{d\ell} = -2\mathcal{D}_R\delta_R + A\mathcal{D}_R^2 + \mathcal{O}(\mathcal{D}_R^2\delta_R), \quad , \quad (10)$$

$$\Gamma_2^{(2)} = - \frac{d\delta_R}{d\ell} = -9\mathcal{D}_R + B\mathcal{D}_R\delta_R + \mathcal{O}(\mathcal{D}_R^2), \quad \beta(x, \tau')] \quad , \quad (11)$$

$$\Gamma_2^{(3)} = - \frac{d}{d\ell}(aD_{fR}) = 0, \quad \frac{d}{d\ell}\left(\frac{v_R}{K_R}\right) = 0, \quad \beta_1]) \Big\}, \quad (12)$$

$$K = \frac{3}{2} + \delta_R$$

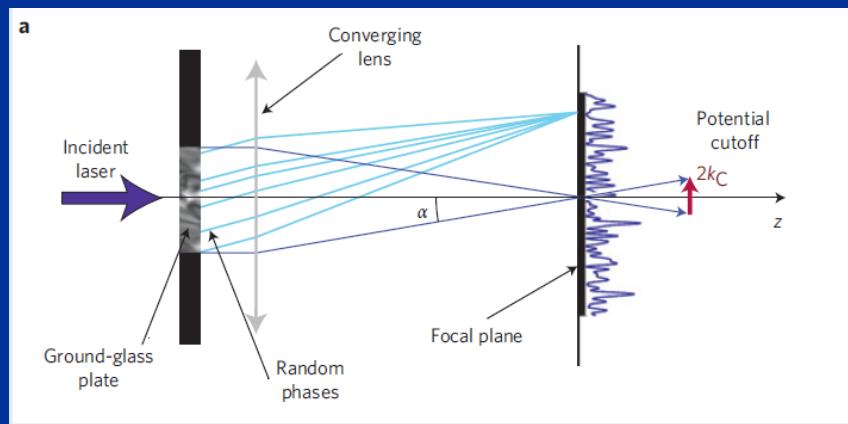
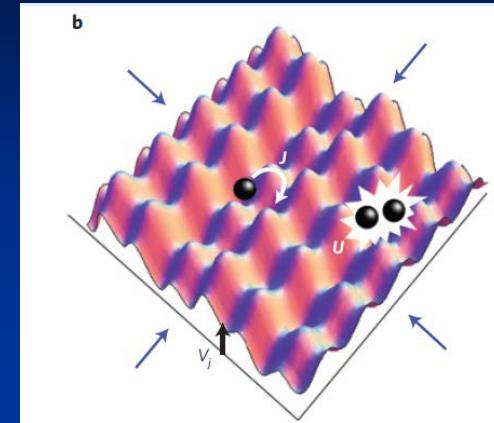
Phase diagram



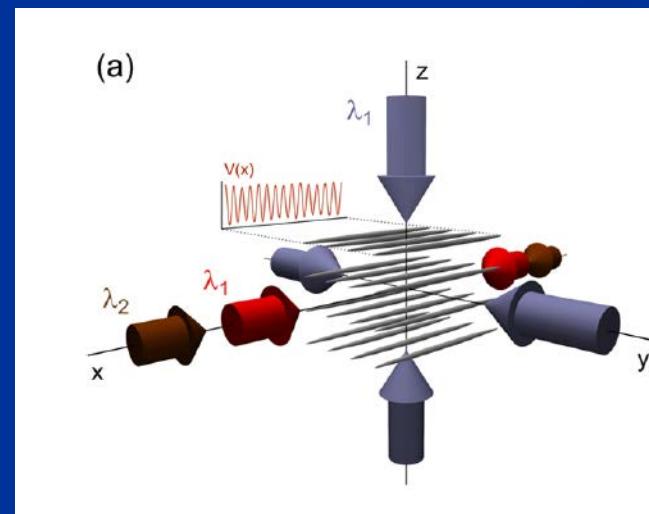
Order parameter ? Moments of distribution ?

Experiments

Cold atomic gases



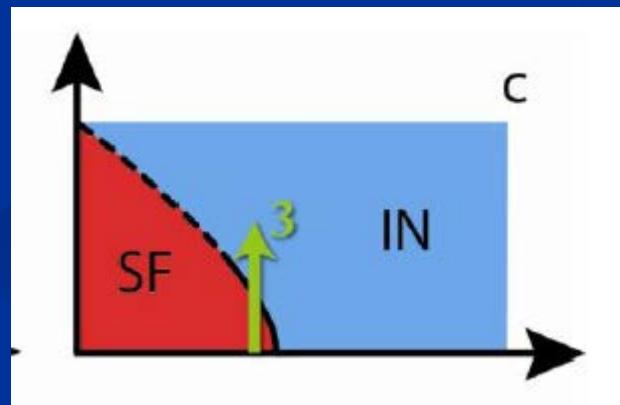
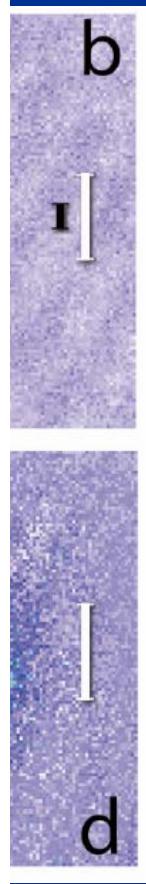
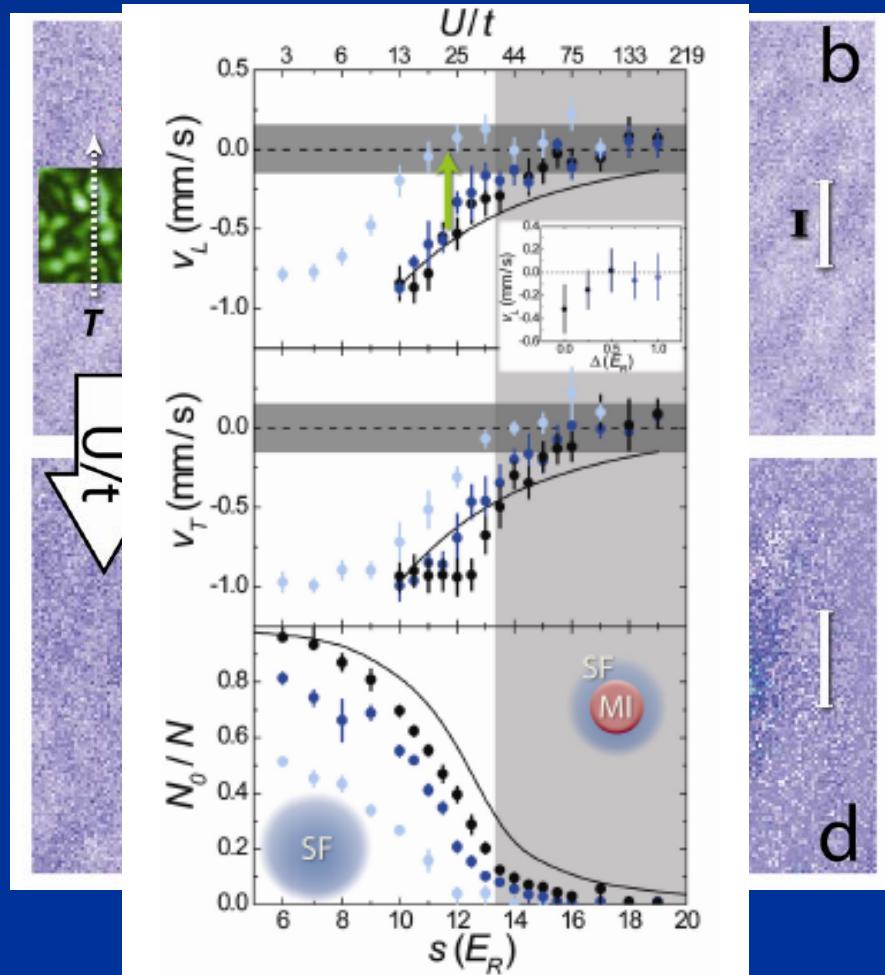
Speckle



Biperiodic lattices

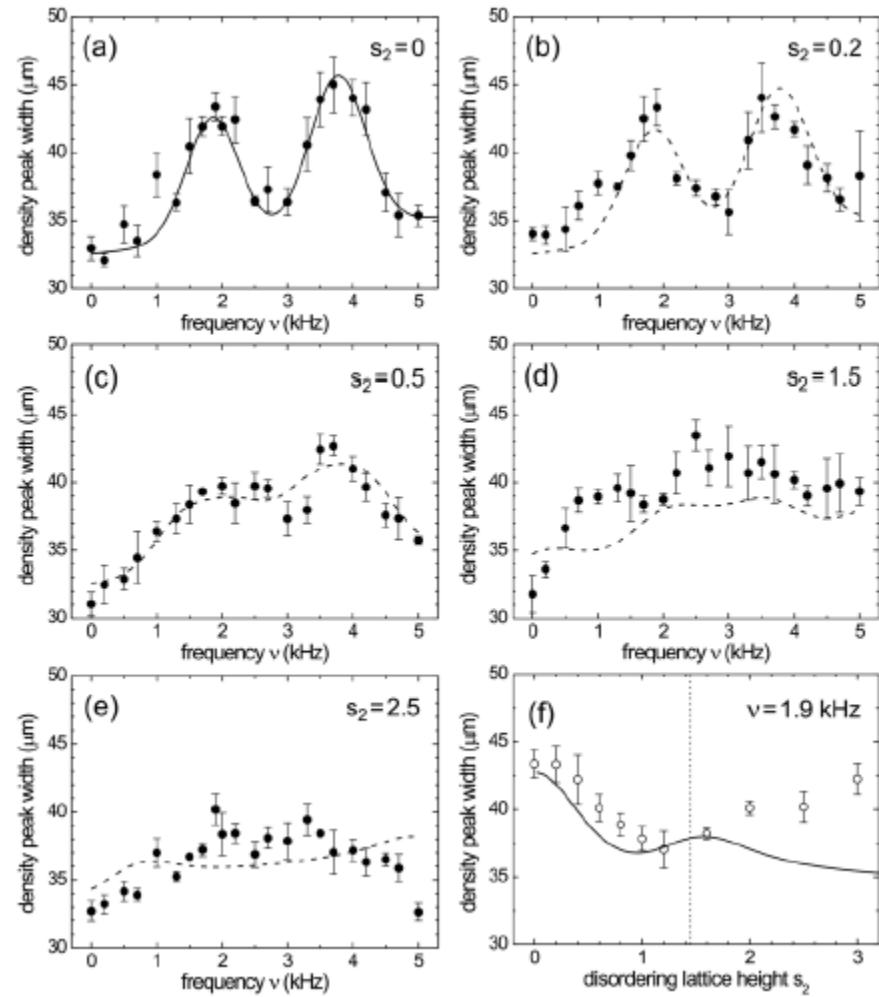
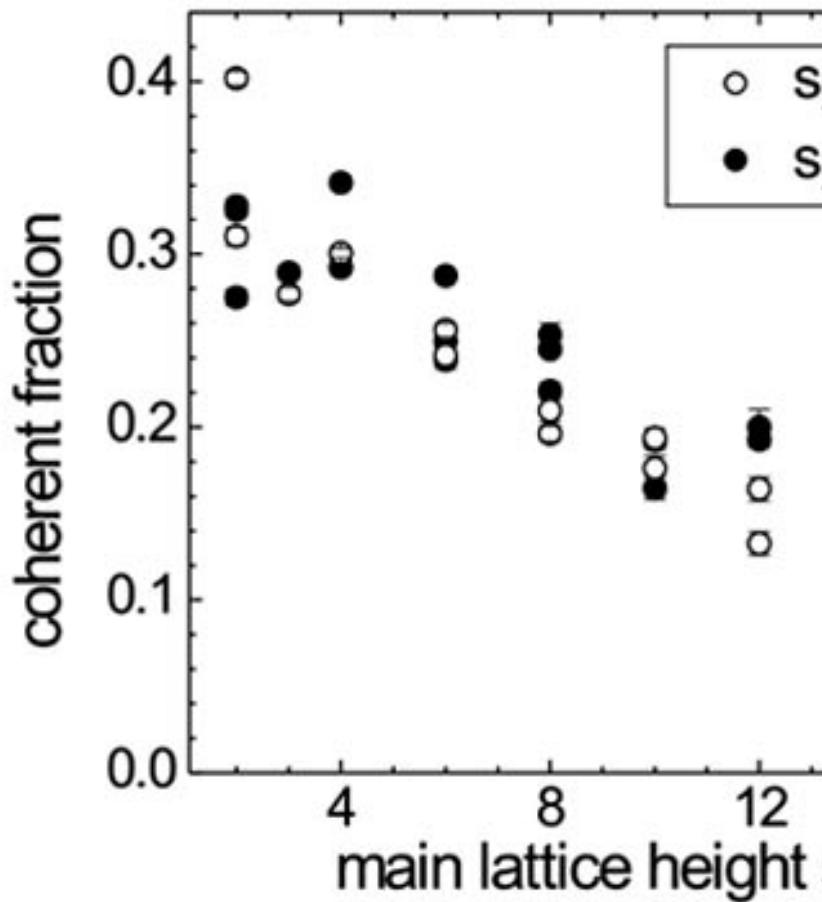
Speckle

Pasienski et al. Nat. Phys 6 677 (2010)



Quasi-periodic

Quasiperiodic (1D):



Same as true disorder ?

Renormalization treatment

J. Vidal, D. Mouhanna, TG PRL 83 3908 (1999); PRB 65 014201 (2001)

$$\frac{dK}{dl} = -K^2 \Xi(l),$$

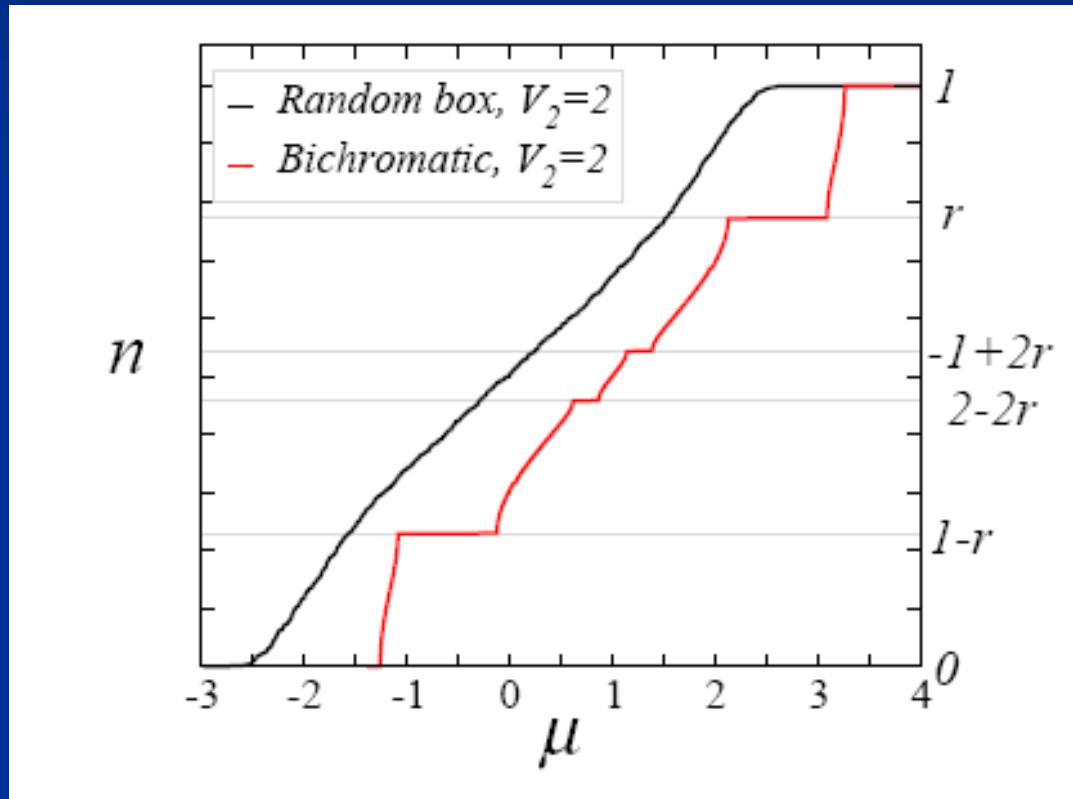
$$\frac{dy_Q}{dl} = (2 - K)y_Q,$$

y_Q : Fourier components
of potential

$$\Xi(l) = \frac{1}{2} \sum_Q y_Q^2 [J(Q^+ \alpha(l)) + J(Q^- \alpha(l))],$$

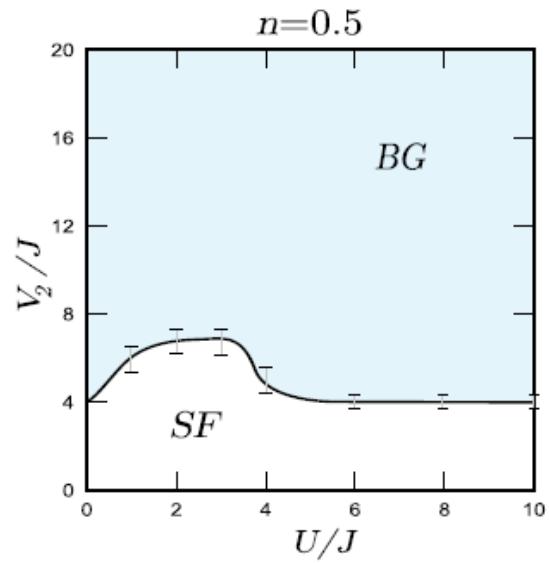
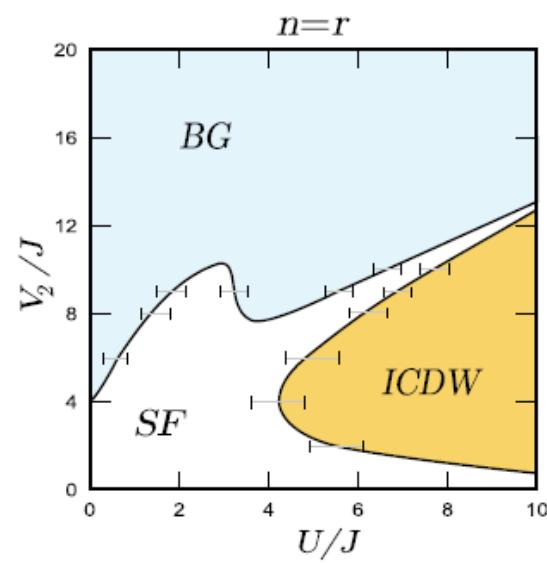
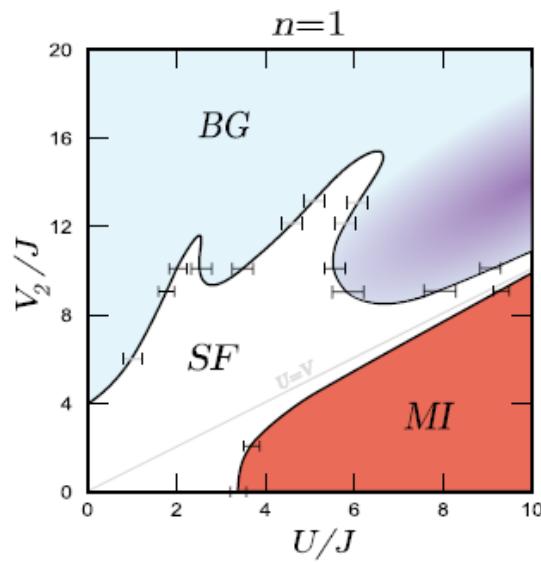
Numerics: DMRG

G. Roux et al. PRA 78 023628 (2008)



Related works: T. Roscilde, Phys. Rev. A 77, 063605 2008;
X. Deng et al PRA 78, 013625 (2008)

Phase diagram



Conclusions

- Tour of one dimensional physics
- Luttinger liquid theory provides a framework to study this physics, and to go beyond
- Many effects: lattice, disorder, long range forces, competition between orders, out of equilibrium physics, gauge fields,.....
- Cold atoms give access to novel one/quasi-one dimensional physics

Calvin and Hobbes

W. WILSON

WOW, IT REALLY
SNOWED LAST NIGHT!
ISN'T IT WONDERFUL?

EVERYTHING FAMILIAR HAS
DISAPPEARED! THE WORLD
LOOKS BRAND-NEW!

A NEW
YEAR...
A FRESH,
CLEAN
START!

IT'S LIKE HAVING A BIG
WHITE SHEET OF PAPER
TO DRAW ON!

A DAY
FULL OF
POSSIBILITIES!

IT'S A MAGICAL
WORLD, HOBES,
OL' BUDDY...

-LET'S GO
EXPLORING!