

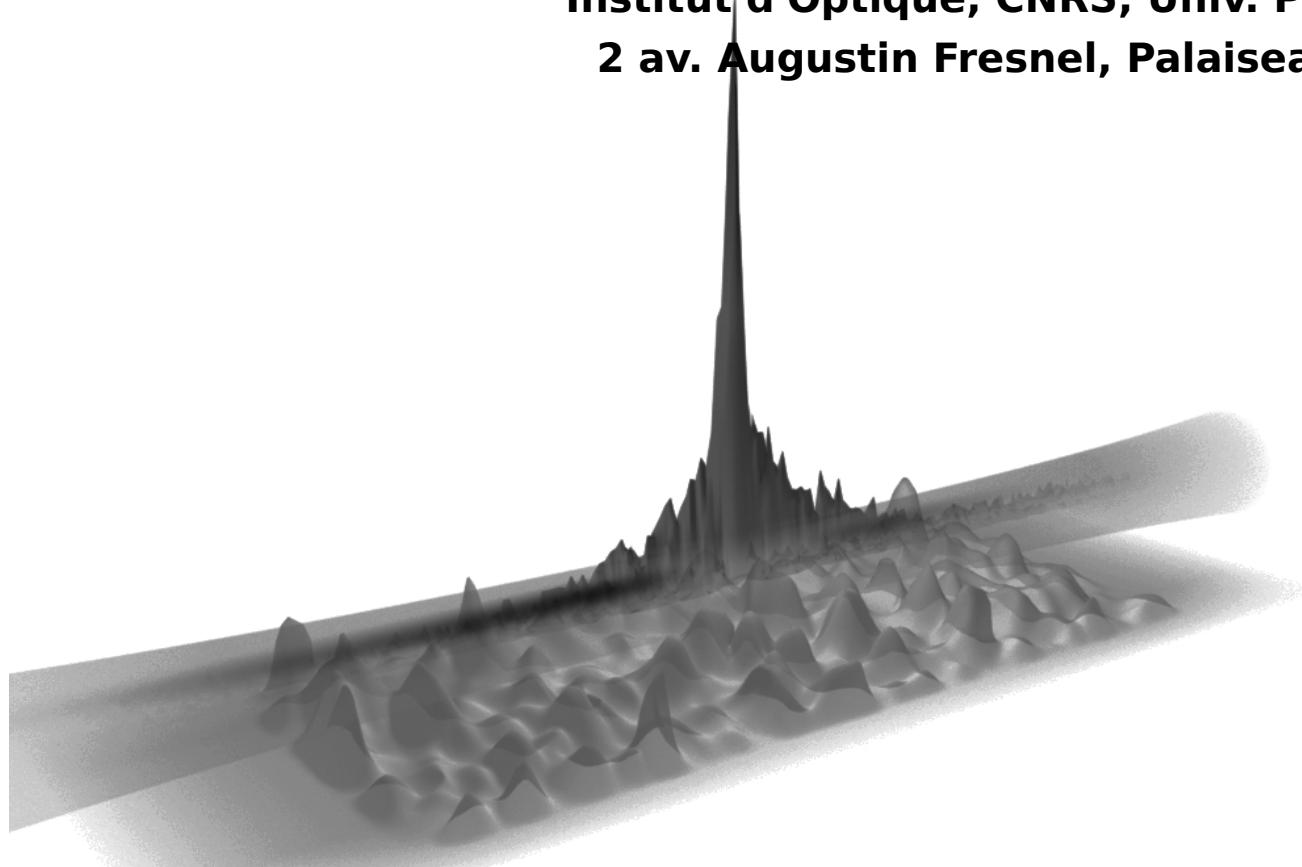
# **Ultracold Atoms in Controlled Disorder I:**

## **Anderson Localization in Dimension d=1**

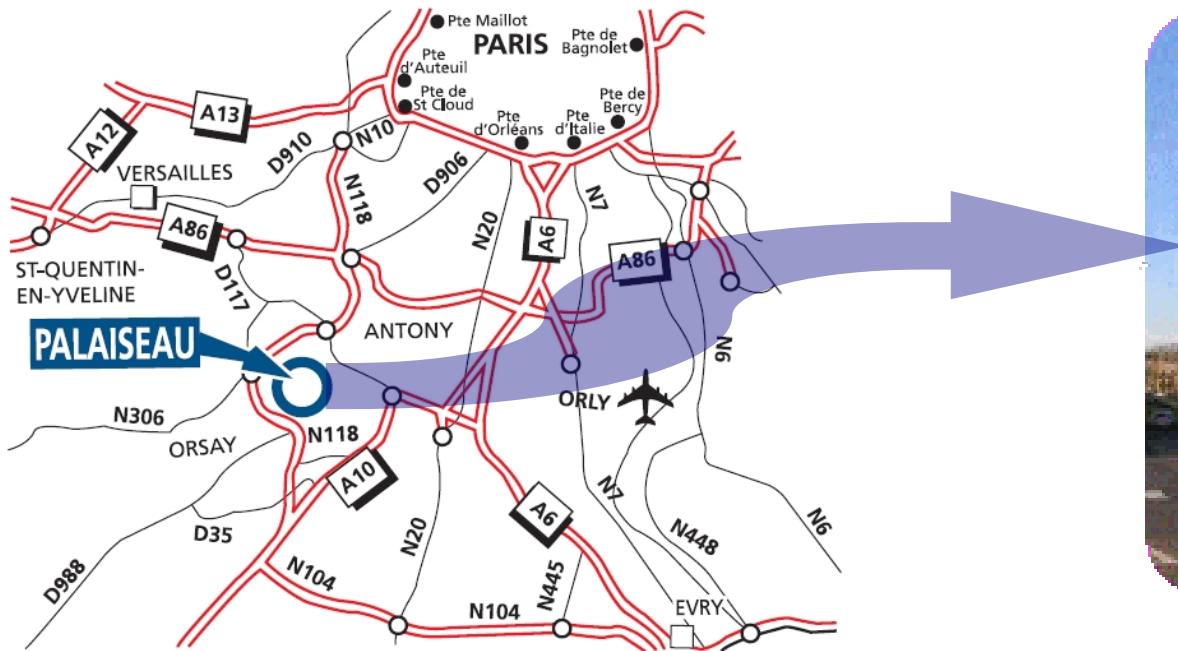
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**Laurent Sanchez-Palencia**

**Laboratoire Charles Fabry - UMR8501  
Institut d'Optique, CNRS, Univ. Paris-sud 11  
2 av. Augustin Fresnel, Palaiseau, France**



# Acknowledgements



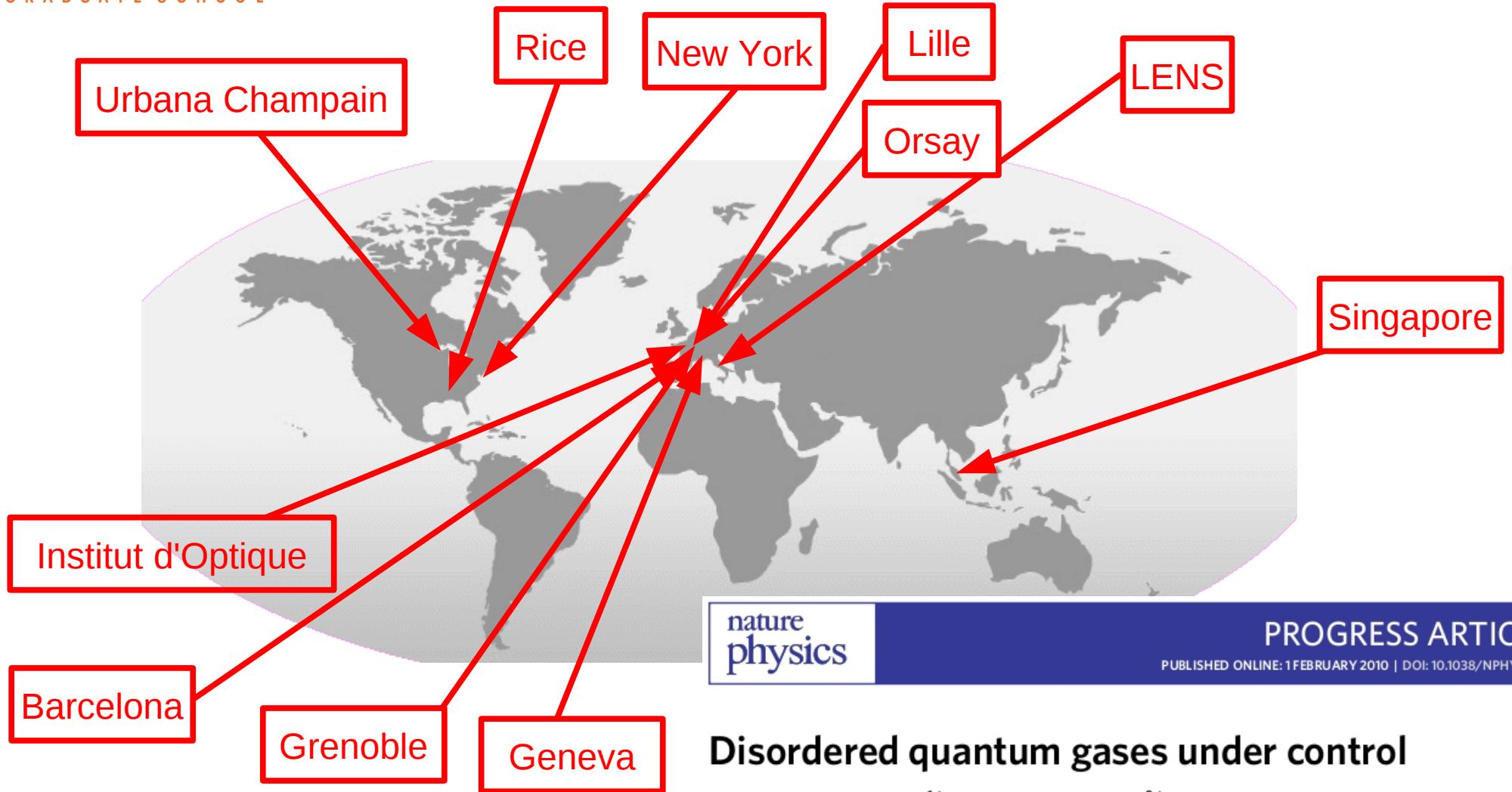
## « Theory of Ultracold Quantum Gases » team

- M. Piraud, S. Lellouch, G. Boeris, G. Carleo, LSP
- former members : L. Pezzé, T.-L. Dao, B. Hambrecht, P. Lugan

## Collaboration and stimulating discussions

- A. Aspect's experimental group, P. Bouyer, V. Josse, T. Bourdel *et al.*
- B. van Tiggelen, T. Giamarchi, M. Lewenstein, G.V. Shlyapnikov

# Disordered Quantum Gases



LSP & Lewenstein, Nat. Phys. **6**, 87 (2010)

special issue in Phys. Today 2009

Modugno, Rep. Prog. Phys. **73**, 102401 (2010)

Fallani *et al.*, Adv. At. Mol. Opt. Phys. **56**, 119 (2008)

Shapiro, J. Phys. A: Math. Theor. **45**, 143001 (2012)

## Disordered quantum gases under control

Laurent Sanchez-Palencia<sup>1\*</sup> and Maciej Lewenstein<sup>2\*</sup>

When attempting to understand the role of disorder in condensed-matter physics, we face considerable experimental and theoretical difficulties, and many questions are still open. Two of the most challenging ones—debated for decades—concern the effect of disorder on superconductivity and quantum magnetism. We review recent progress in the field of ultracold atomic gases, which should pave the way towards the realization of versatile quantum simulators, which help solve these questions. In addition, ultracold gases offer original practical and conceptual approaches, which open new perspectives to the field of disordered systems.

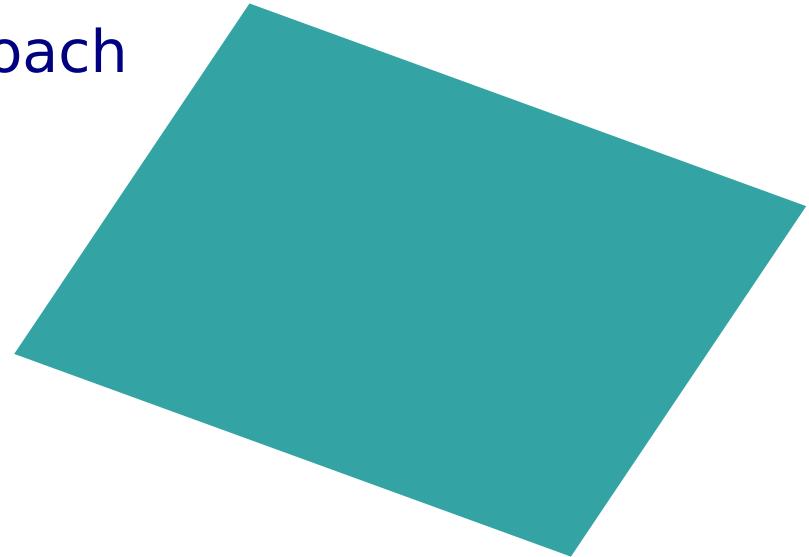
## Lecture #1

1. Basics of quantum transport theory in disordered media
2. Ultracold atoms and speckle potentials  
Specific features of a controlled disorder
3. Anderson localization in 1D speckle potentials
  - 3.1 Effective mobility edges
  - 3.2 Expansion of Bose-Einstein condensates (exp. vs theory)

- 💡 Anderson localization : microscopic approach

- Homogeneous (underlying) medium

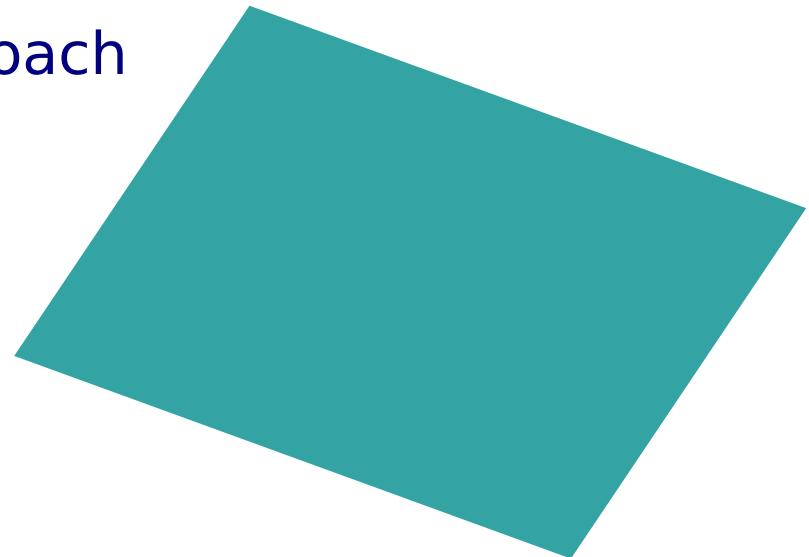
- $|\mathbf{k}\rangle$  states
- characterized by  $\epsilon(\mathbf{k})$  and  $\mathbf{v} = \partial\epsilon(\mathbf{k})/\partial\hbar\mathbf{k}$



- Anderson localization : microscopic approach

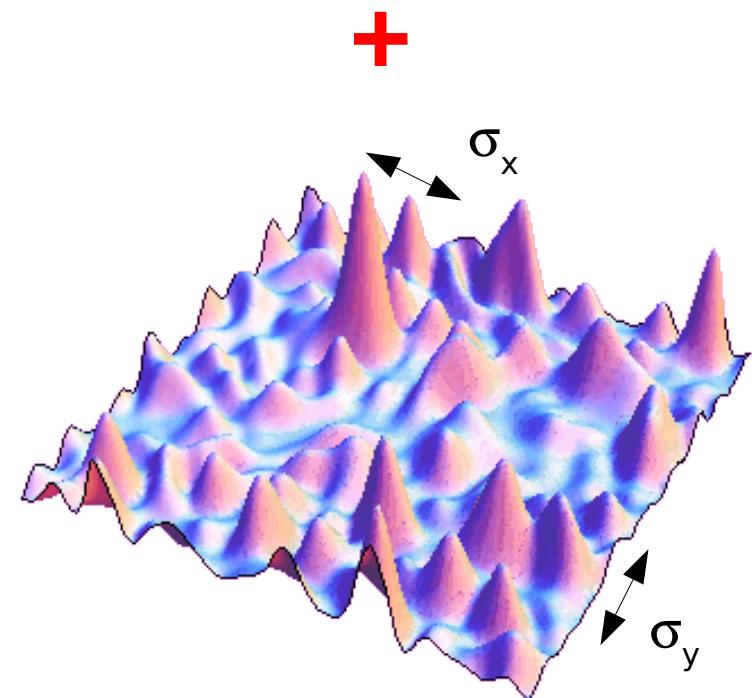
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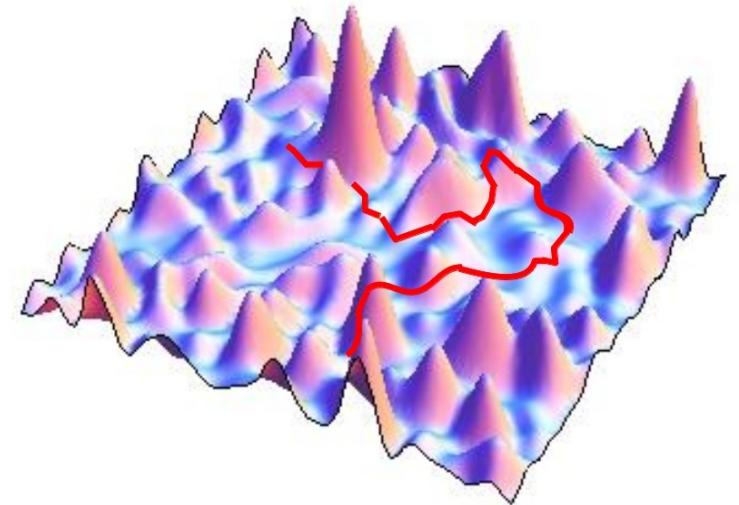
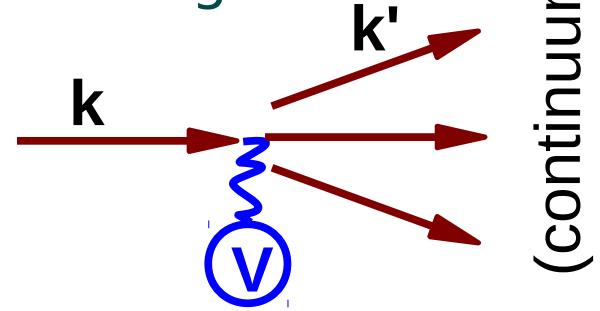


- Disorder

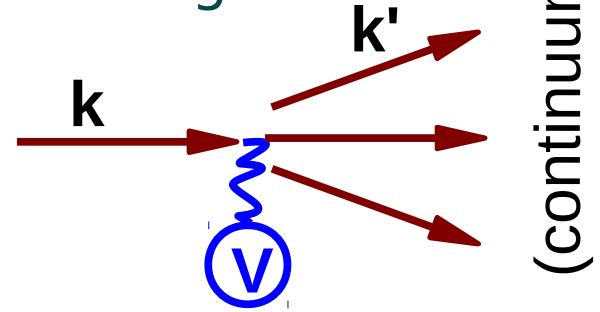
- $\bar{V}=0$  (sets the zero of energy)
- characterized by  $C_n(\mathbf{r}_1, \dots, \mathbf{r}_n) = \overline{V(\mathbf{r}_1) \dots V(\mathbf{r}_n)}$
- most importantly  $C_2(\mathbf{r}) = V_R^2 c_2(\{\mathbf{z}_j/\sigma_j\})$  and its Fourier transform  $C_2(\mathbf{k})$  (disorder power spectrum)
- If  $c_2$  is isotropic,  $\zeta_{j,l} = \sigma_j/\sigma_l$  (anisotropy factors)



- Single scattering



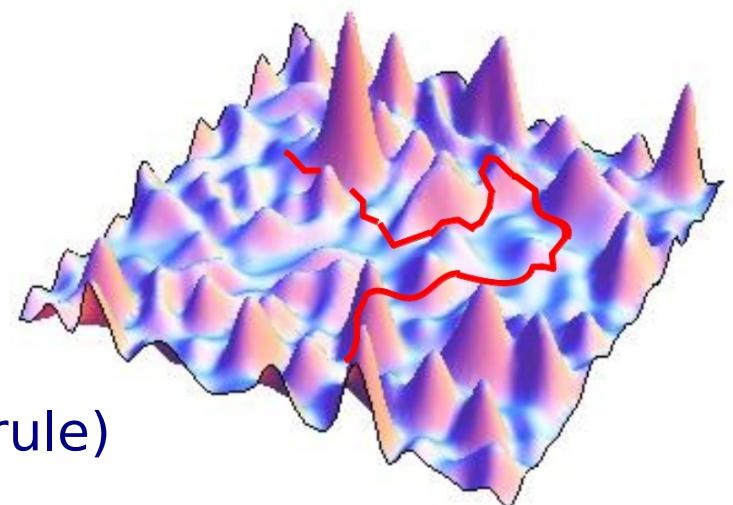
- Single scattering



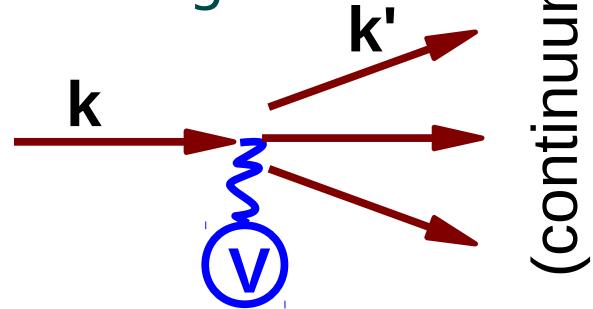
- finite life time  $\tau_s$  of  $|\mathbf{k}\rangle$  states (Fermi golden rule)

$$\tau_s(E, \mathbf{k}) = \frac{\hbar}{2\pi} \frac{1}{\langle \tilde{C}(\mathbf{k} - \mathbf{k}') \rangle_{\mathbf{k}'|E}}$$

$$\langle \dots \rangle_{\mathbf{k}'|E} = \int \frac{d\mathbf{k}'}{(2\pi)^d} \dots \delta [E - \epsilon(\mathbf{k}')] \quad$$



- Single scattering

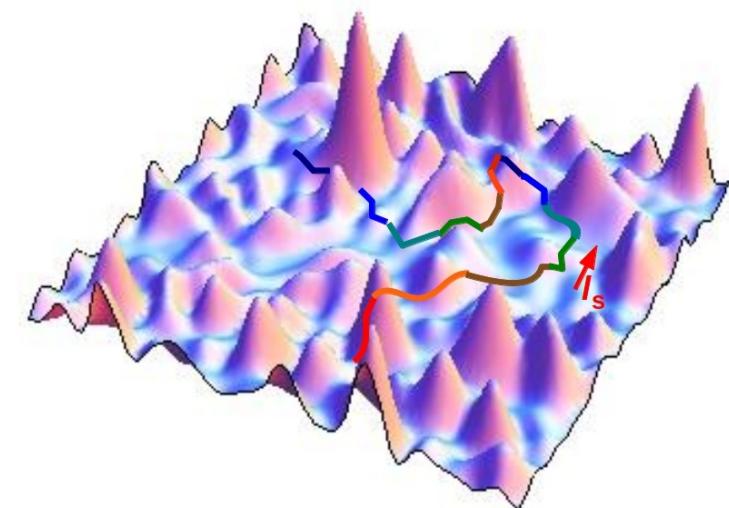


(continuum)

- finite life time  $\tau_s$  of  $|k\rangle$  states (Fermi golden rule)

$$\tau_s(E, \mathbf{k}) = \frac{\hbar}{2\pi} \frac{1}{\langle \tilde{C}(\mathbf{k} - \mathbf{k}') \rangle_{\mathbf{k}'|E}}$$

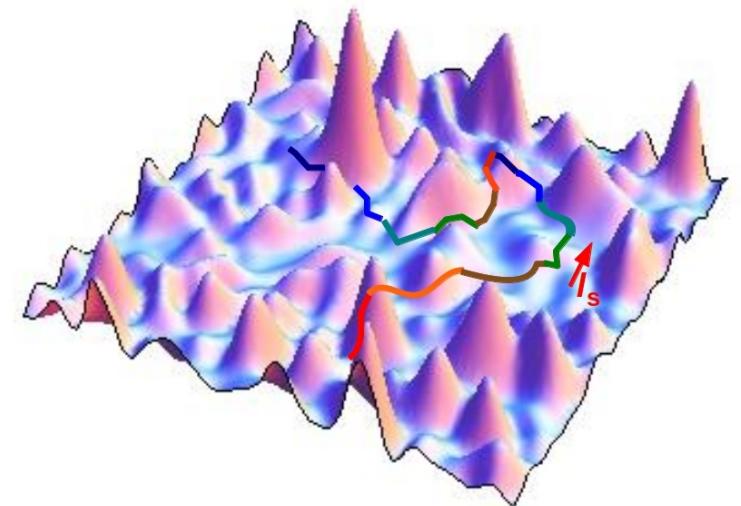
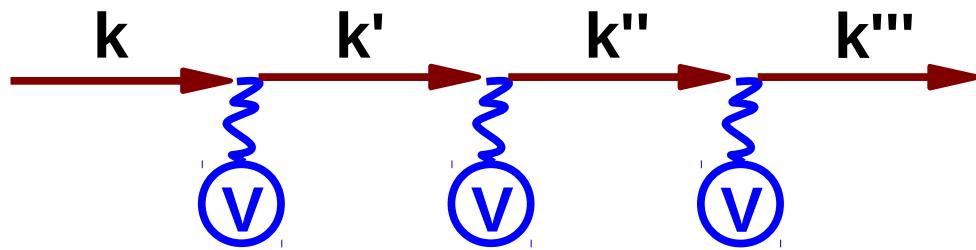
$$\langle \dots \rangle_{\mathbf{k}'|E} = \int \frac{d\mathbf{k}'}{(2\pi)^d} \dots \delta [E - \epsilon(\mathbf{k}')] \quad$$



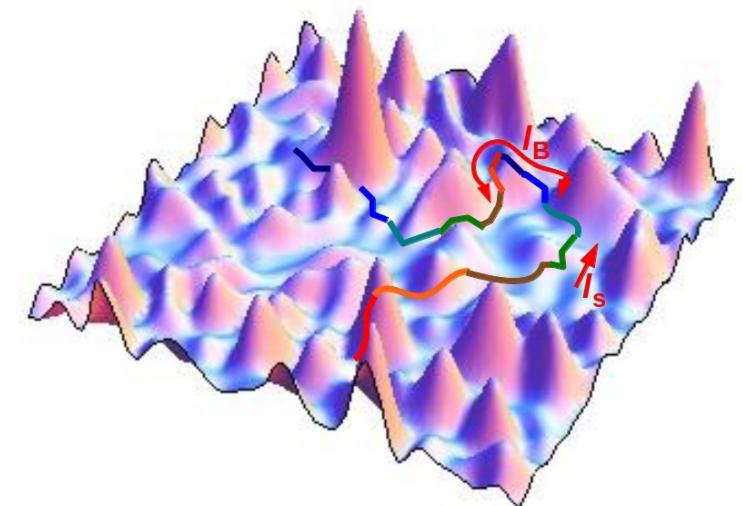
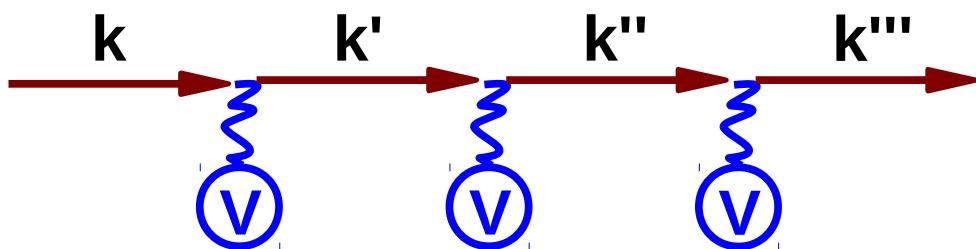
scattering mean free path  
 $l_s(\mathbf{k}) = |\mathbf{v}| \tau_s(\mathbf{k})$

- In general  $C(\mathbf{q})$  decreases with  $|\mathbf{q}|$  (maximum coupling  $\mathbf{k} \rightarrow \mathbf{k}$ , ie forward scattering)  $\rightarrow \tau_s$ : characteristic time of phase loss
- For white noise,  $C(\mathbf{q}) = cst$  and the scattering is isotropic  $\rightarrow \tau_s$  describes phase and direction loss

- Multiple scattering



- Multiple scattering

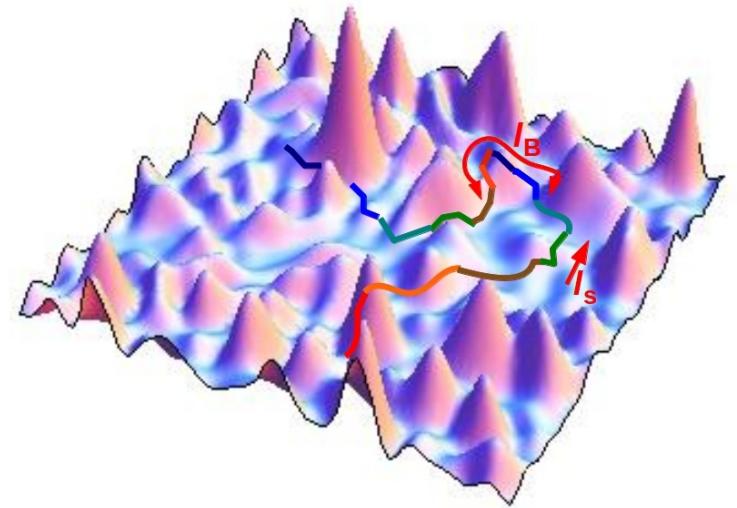
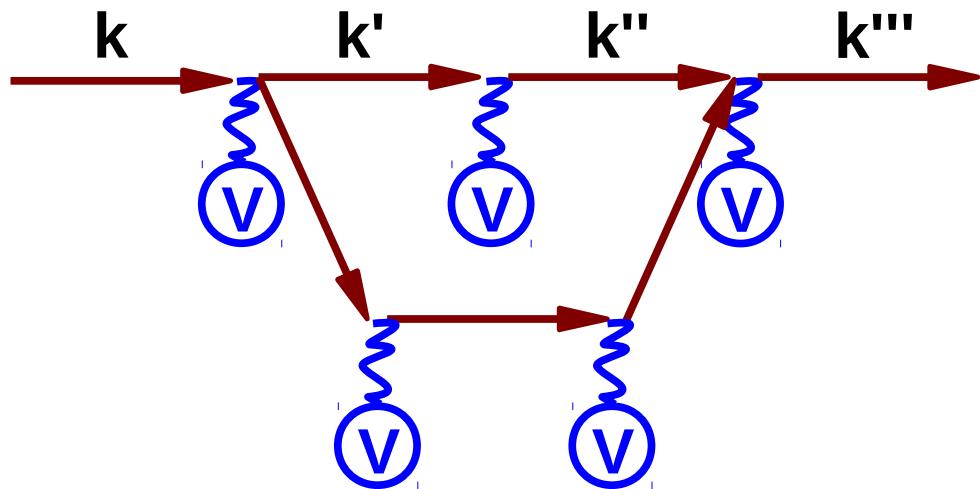


- deflection of  $|k\rangle$  states : random walk  
 $\rightarrow$  normal spatial diffusion,  $D_B(E)$

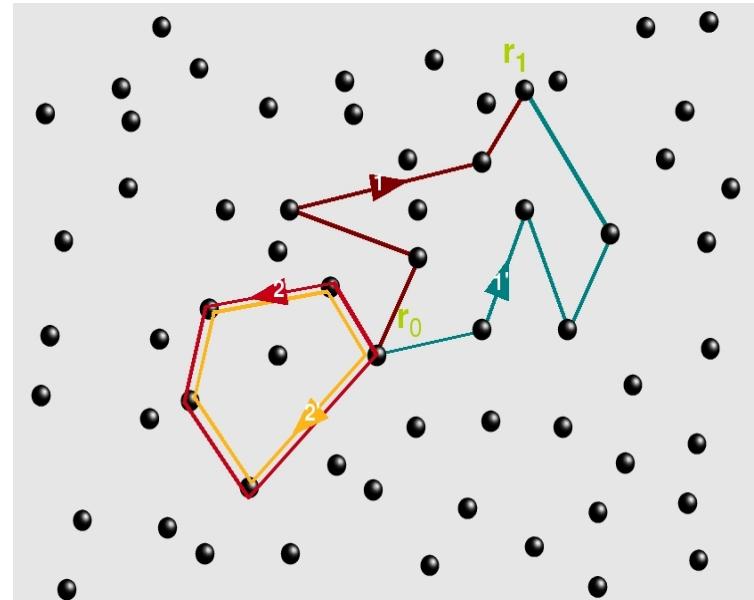
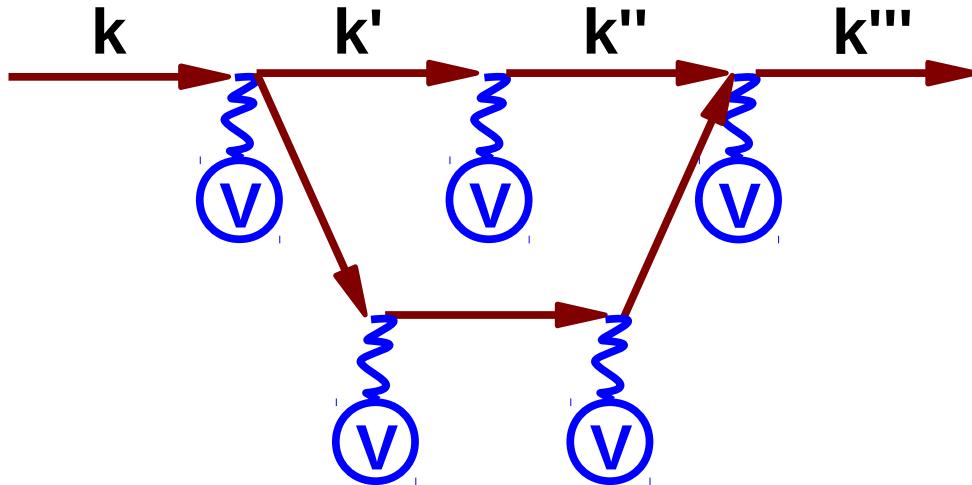
transport mean free path  
 $D_B(E) = (1/d) |\mathbf{v}| l_B(E)$

- In general  $C(\mathbf{q})$  decreases with  $|\mathbf{q}|$  (maximum coupling  $\mathbf{k} \rightarrow \mathbf{k}$ , ie forward scattering)  $\rightarrow l_B(E) > l_s(E)$
- For white noise,  $C(\mathbf{q}) = cst$  and the scattering is isotropic  $\rightarrow \tau_s$  describes phase and direction loss, and  $l_B(E) = l_s(E)$

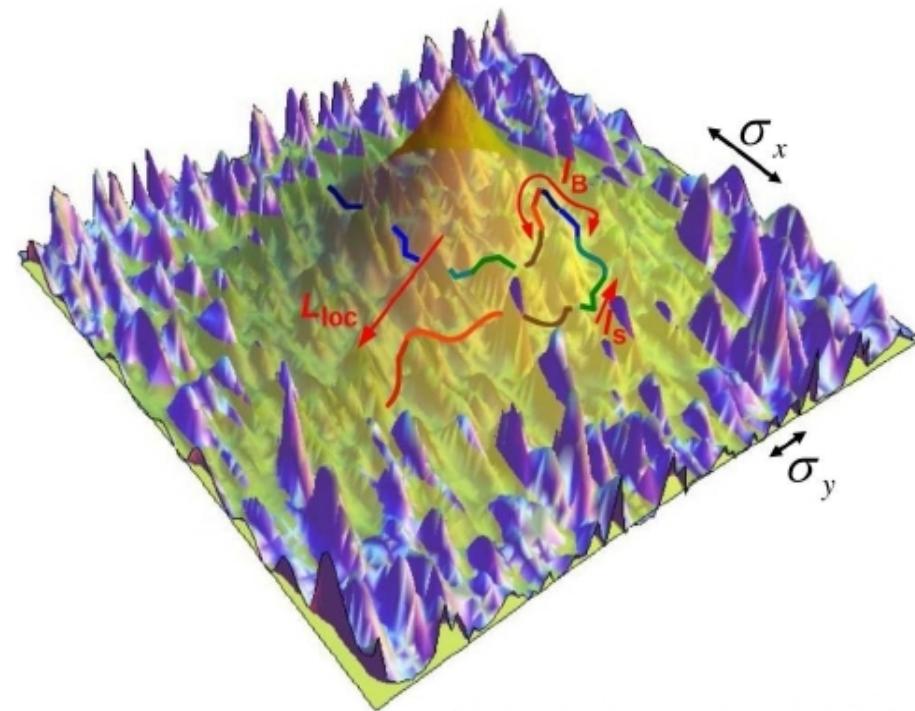
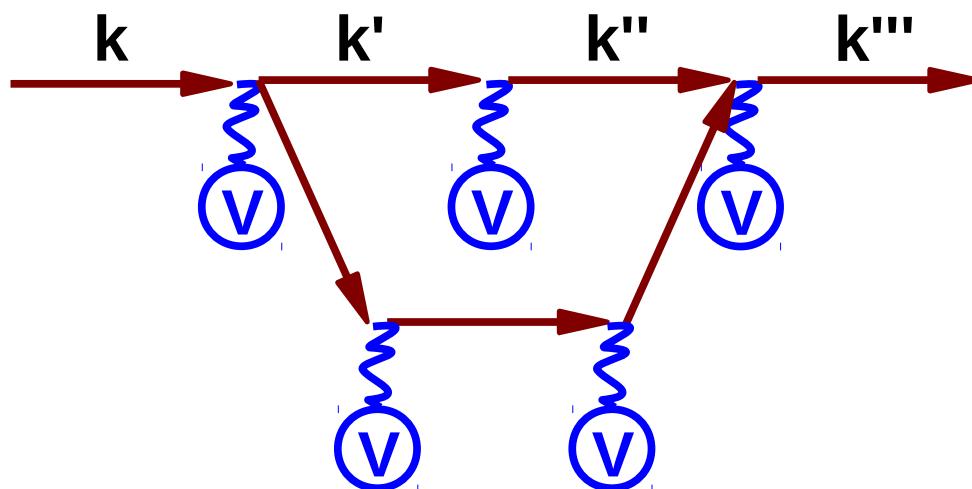
- Interference effects



- Interference effects



- Interference effects



decreases with  
 -  $d$  (return prob.)  
 -  $k l_B$  (interf. param.)

localization length  
 $\gamma(E) = L_{\text{loc}}(E)^{-1} = l_B^{-1} F_d(k l_B)$

- weak localization  $\rightarrow D_*(E) < D_B(E)$  : diffusion
- strong (Anderson) localization  $\rightarrow D_*(E) \sim -i\omega L_{\text{loc}}(E)^2$   
 exponential localization

## 💡 Anderson localization : scaling theory

- dimensionless conductance :  $g=0$  (insulator) ;  $g= \infty$  (metal)
- RG based on the fundamental « on-parameter scaling »  
hypothesis :  $d \ln(g) / d \ln(L) = \beta(g)$

P.W. Anderson, Phys. Rev. **109**, 1492 (1958)

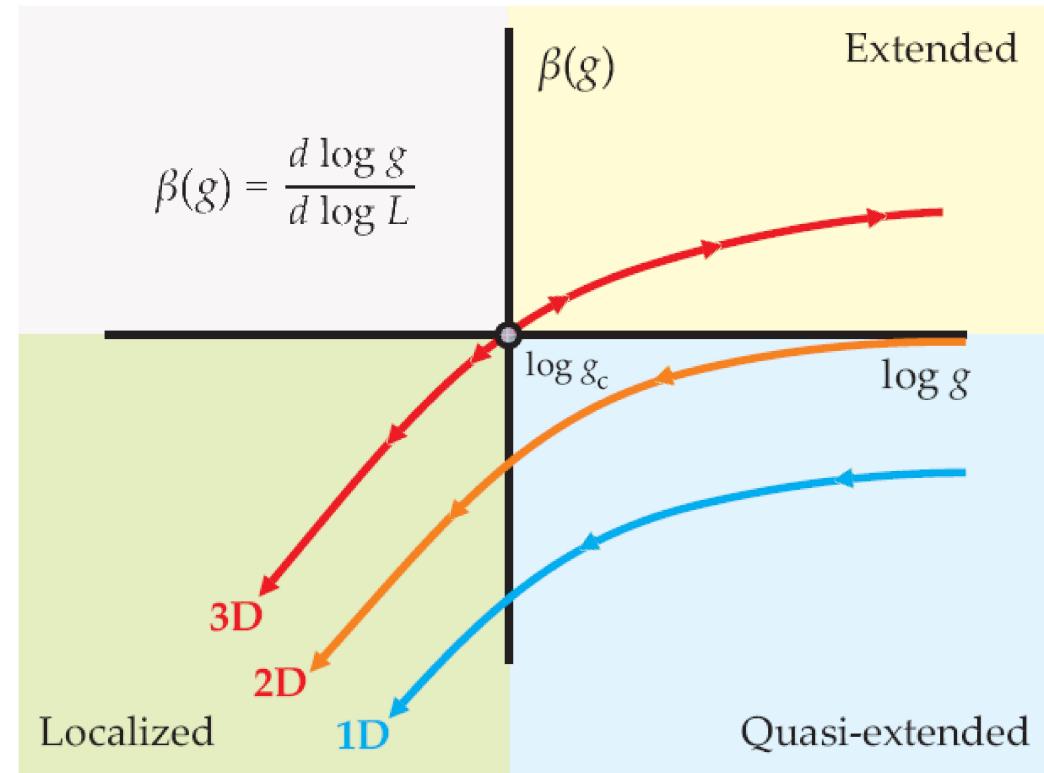
Edwards and Thouless, J. Phys. C **5**, 807 (1972)

Abrahams *et al.*, Phys. Rev. Lett. **42**, 673 (1979)

## Anderson localization : scaling theory

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1D : no diffusion & all states localized  
 2D : all states localized  
 3D : mobility edge,  $k l^* \sim 1$



P.W. Anderson, Phys. Rev. **109**, 1492 (1958)  
 Edwards and Thouless, J. Phys. C **5**, 807 (1972)  
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## ● Flexibility

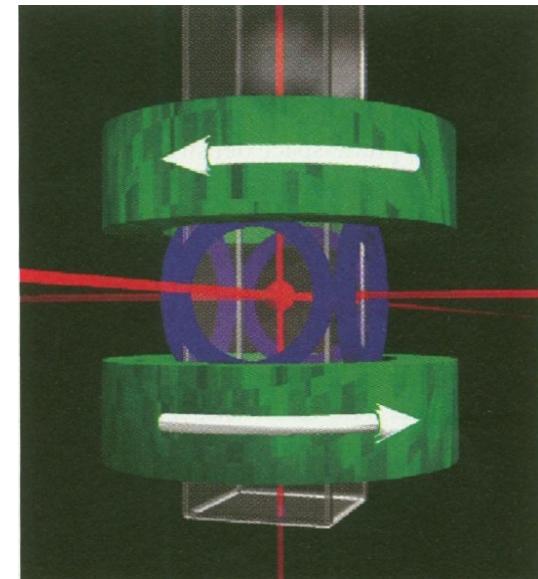
- control of parameters (dimensionality, shape, interactions, bosons/fermions/mixtures, ...)

## ● High-precision measurements

- complementary to condensed-matter tools
- variety of diagnostic tools (**direct imaging**, velocity distribution, oscillations, Bragg spectroscopy, ....)

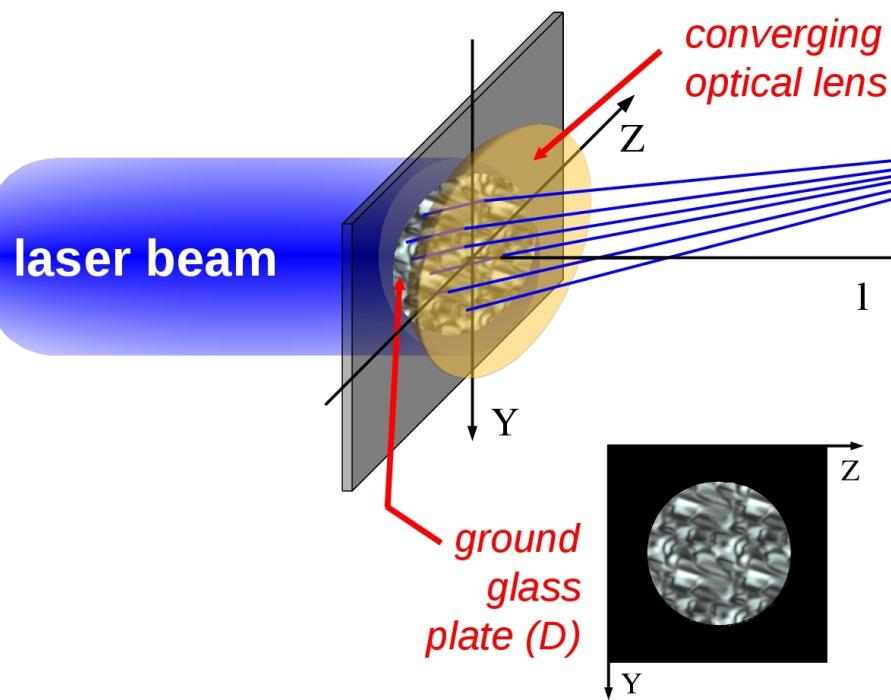
## ● Model systems

- parameters known ab-initio  $\Rightarrow$  direct comparison with theory
- towards quantum simulators  $\Rightarrow$  artificial, controlled systems to study fundamental questions

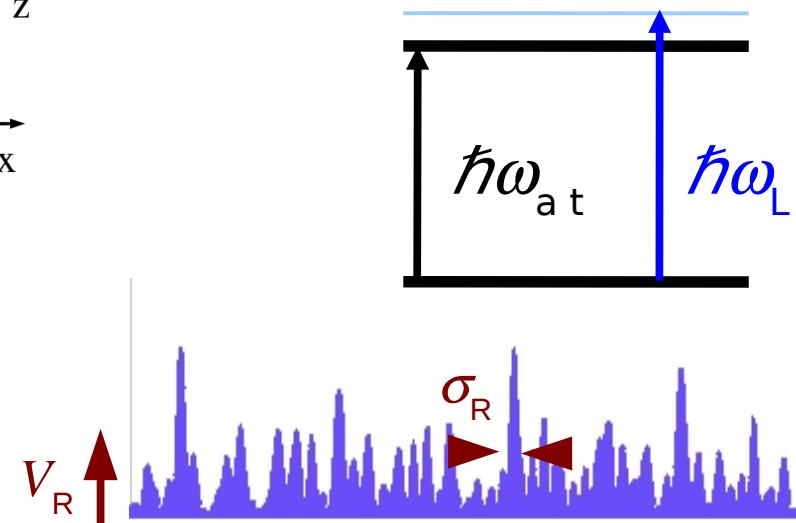


specific ingredients of ultracold atoms ...

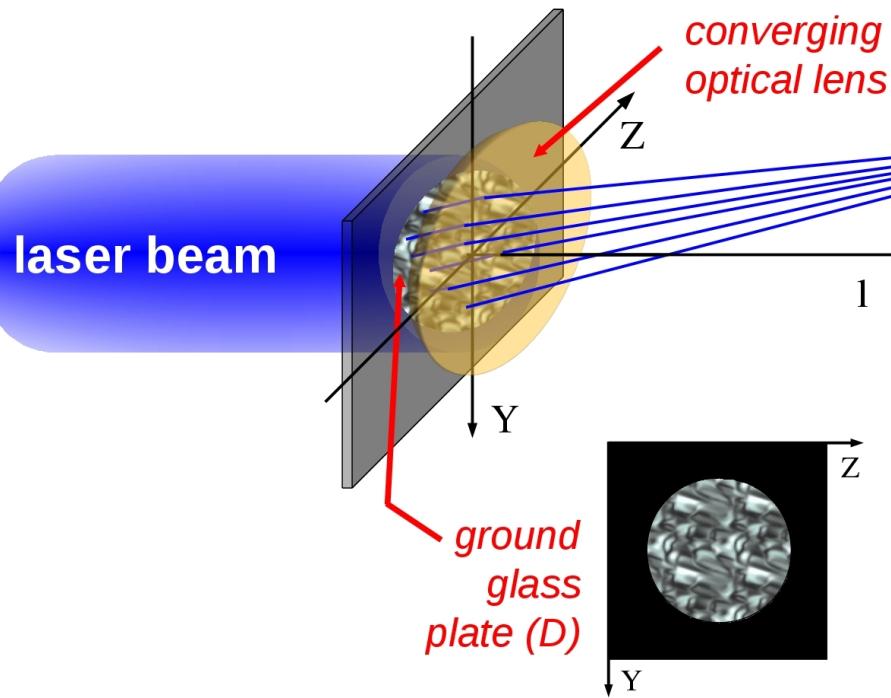
# Controlled Optical Disorder



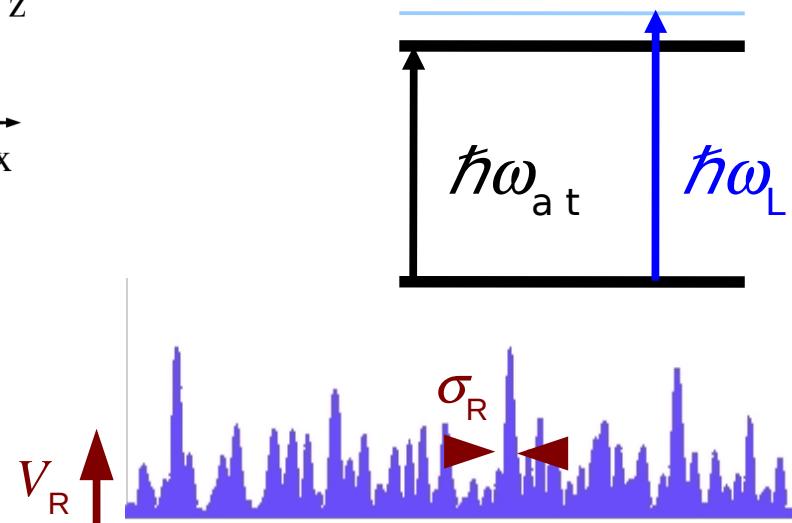
J. W. Goodman, *Speckle Phenomena in Optics* (2007)  
Clément et al., *New J. Phys.* **8**, 165 (2006)



# Controlled Optical Disorder



J. W. Goodman, *Speckle Phenomena in Optics* (2007)  
Clément et al., *New J. Phys.* **8**, 165 (2006)



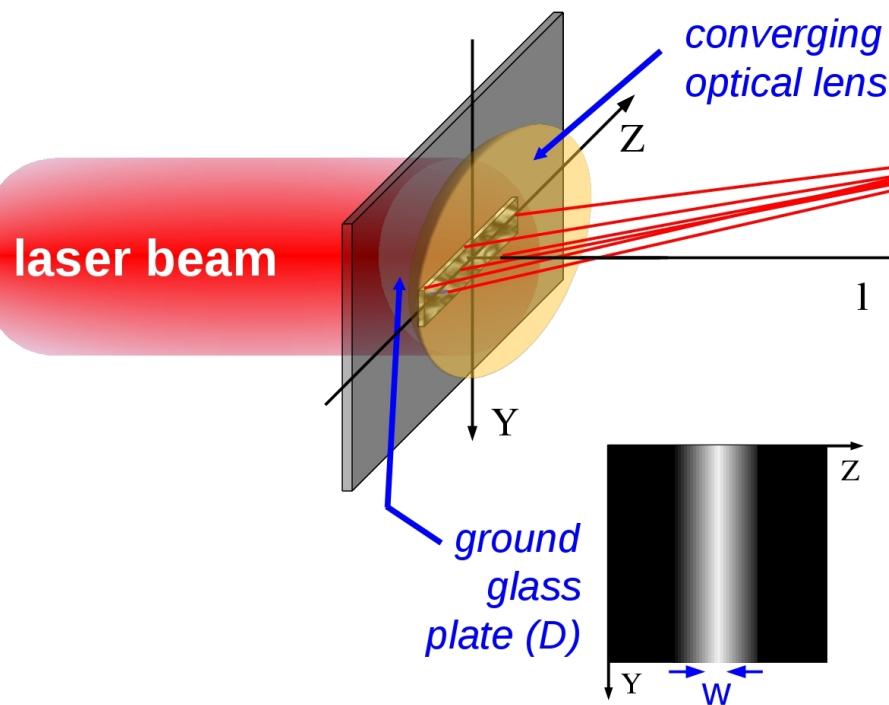
## Speckles: an original class of disorder

- $E(\mathbf{r})$  is a Gaussian random process
- $V(\mathbf{r}) \propto |E(\mathbf{r})|^2 - \overline{|E|^2}$  is **not Gaussian**  
is **not symmetric**
- $C_n(r_1, \dots, r_n) = \overline{V(r_1) \dots V(r_n)}$  all  
determined by  $C_E(\mathbf{r}) = \overline{E(\mathbf{r})^* E(\mathbf{0})}$

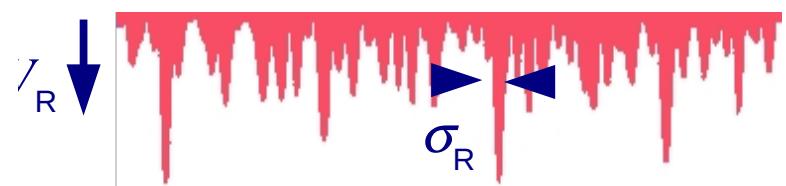
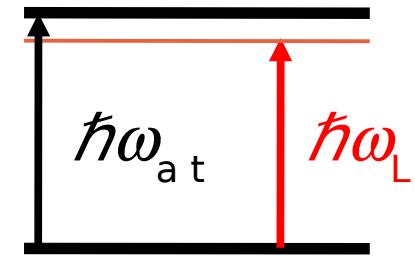
## Experimental control

- $$C_2(\mathbf{r}) = \overline{V(\mathbf{r})V(0)} - \overline{V}^2$$
- $V_R \propto I_L / \Delta$
  - $\sigma_R$ : correlation length  
(depends only on the intensity profile of the ground-glass plate)
  - $C_2(\mathbf{r})$  can be tailored almost at will

# Controlled Optical Disorder



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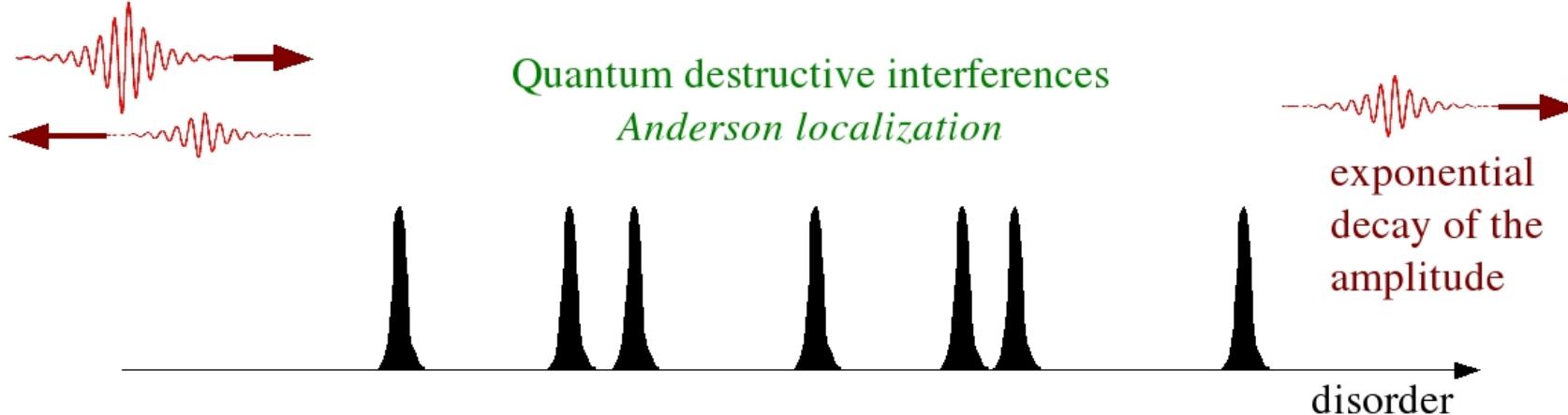
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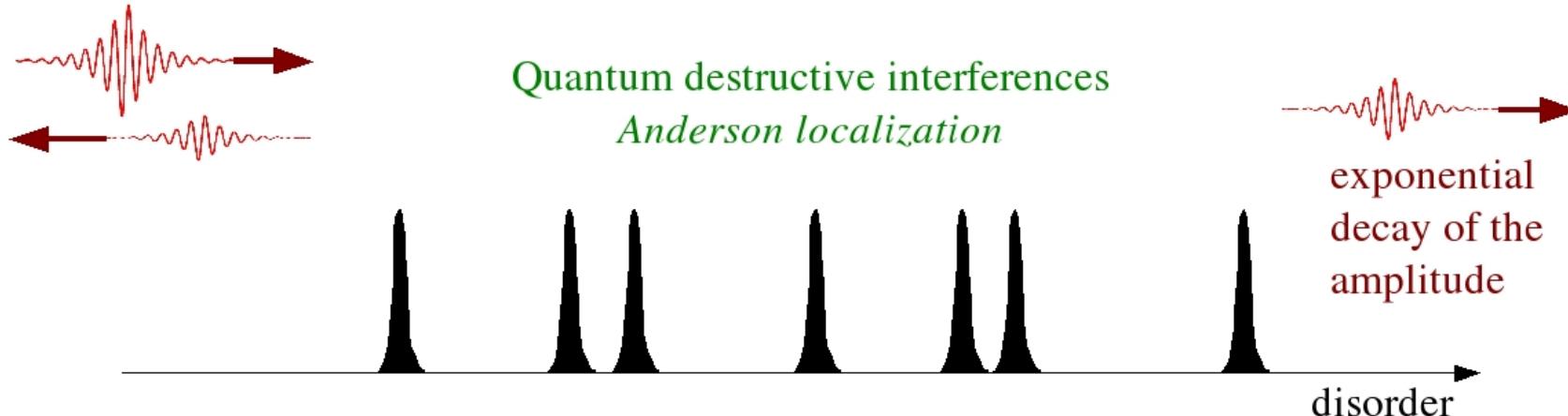
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## quantum physics

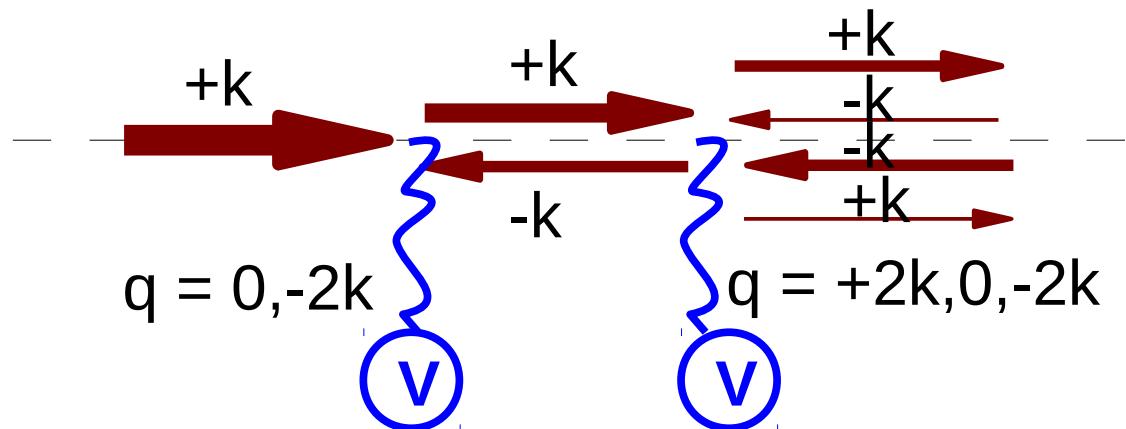


$$\ln(|\psi(z)|/|\psi(0)|) \sim -|z|/L_{\text{loc}}$$

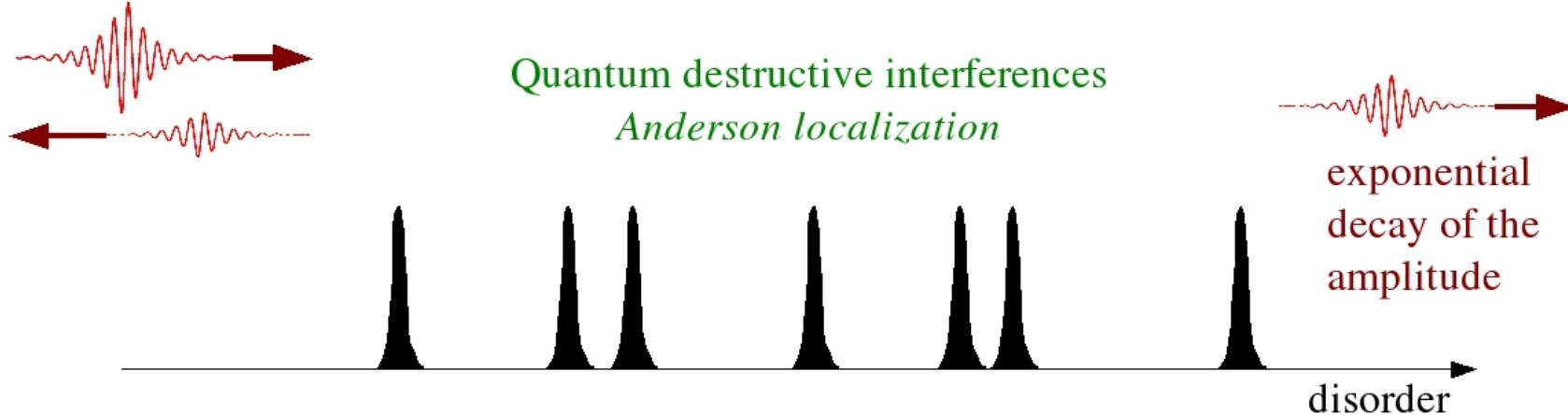
## quantum physics



$$\ln(|\psi(z)|/|\psi(0)|) \sim -|z|/L_{\text{loc}}$$



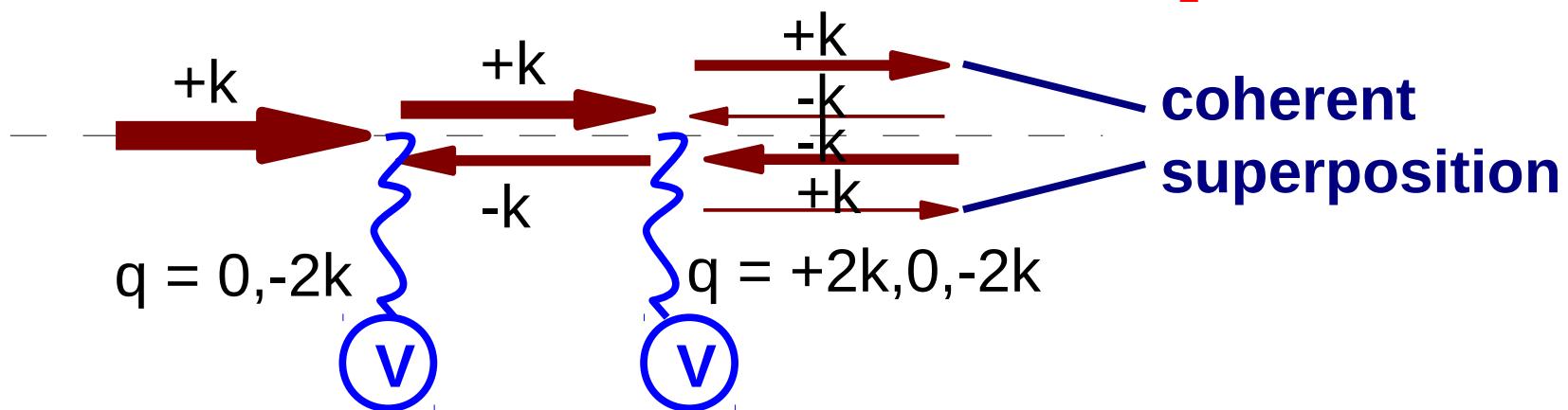
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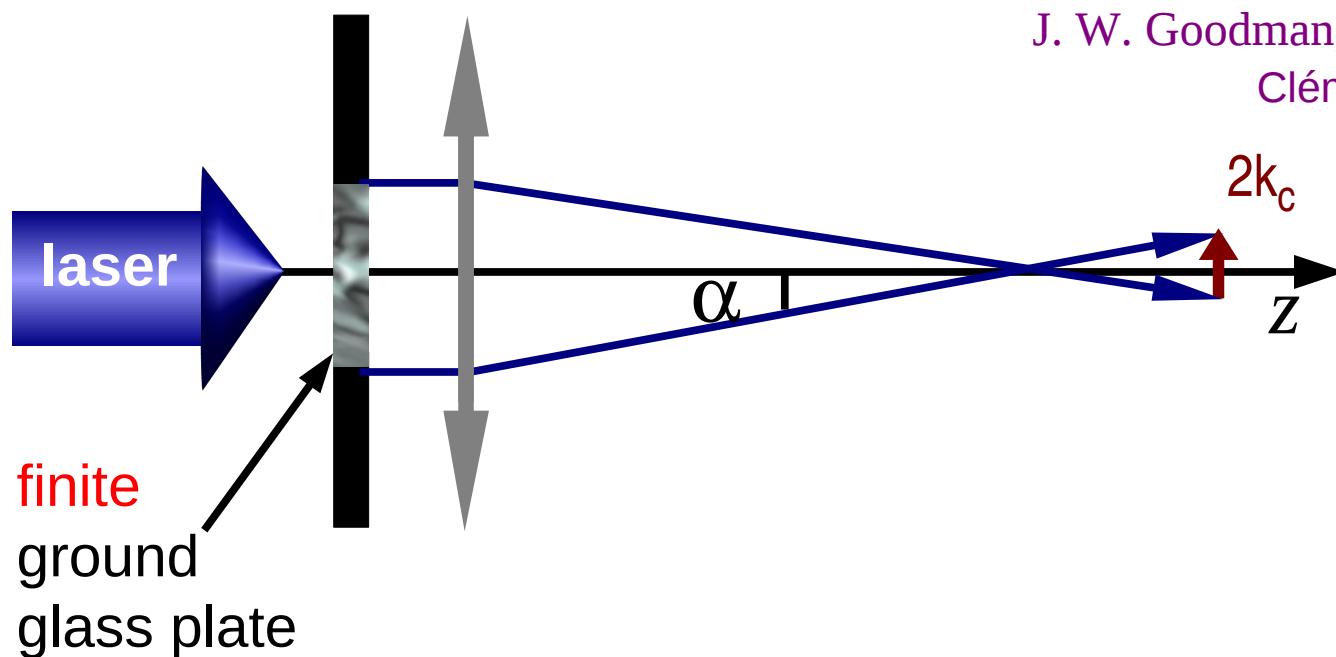


$$\ln(|\psi(z)|/|\psi(0)|) \sim -|z|/L_{\text{loc}}$$

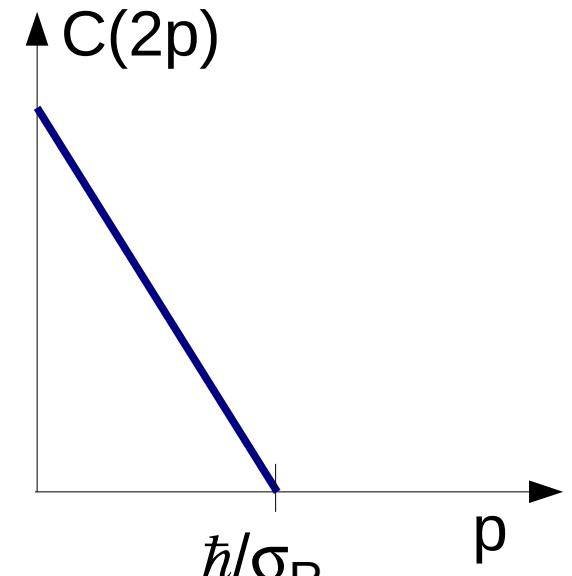
Quantity of interest:

$$C_2(2k) = \overline{V(+2k)V(-2k)} \neq 0$$





J. W. Goodman, *Speckle Phenomena in Optics* (2007)  
Clément et al., *New J. Phys.* **8**, 165 (2006)



Potential cut-off:

$$C_E(k)=0 \text{ for } k>k_c$$

$$C_2(2k)=0 \text{ for } k>k_c$$

$$k_c=(2\pi/\lambda_L)\sin\alpha \equiv \sigma_R^{-1}$$

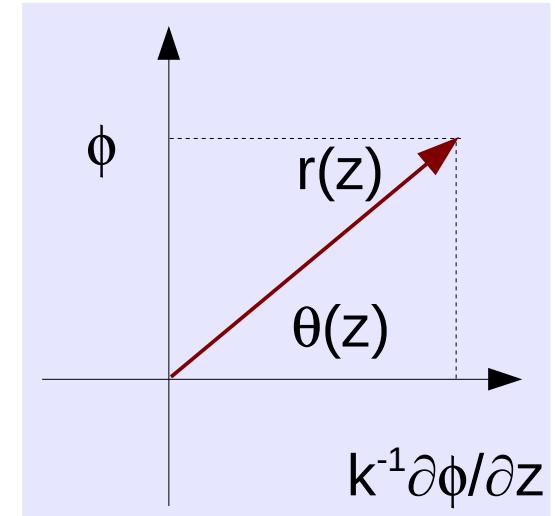
$$\tilde{C}(p) = \pi V_R^2 \sigma_R [1 - |p|\sigma_R/2\hbar] \oplus$$

## 💡 Perturbation theory in the phase formalism

$$\phi(z) = r(z)\sin[\theta(z)]; \quad \partial_z\phi = kr(z)\cos[\theta(z)]$$

$$\partial_z\theta(z) = k\{1 - [V(z)/E]\sin^2[\theta(z)]\}$$

$$\ln[r(z)/r(0)] = k \int_0^z dz' [V(z')/2E]\sin[2\theta(z')]$$



$k_E$  is defined by  
 $E = \hbar^2 k_E^{-2}/2m = p_E^{-2}/2m$

## 💡 Perturbation theory in the phase formalism

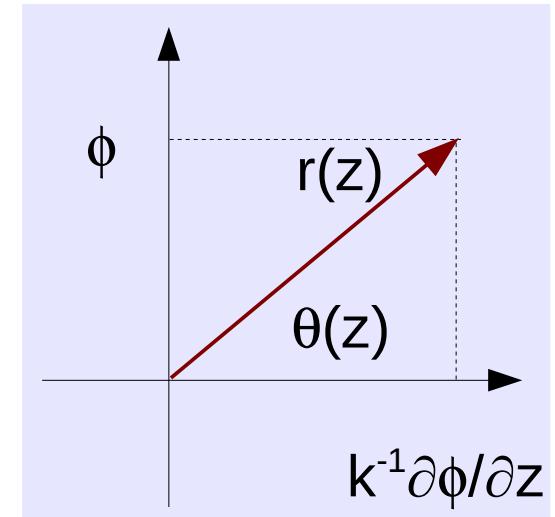
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$$\partial_z\theta(z) = k\{1 - [V(z)/E]\sin^2[\theta(z)]\}$$

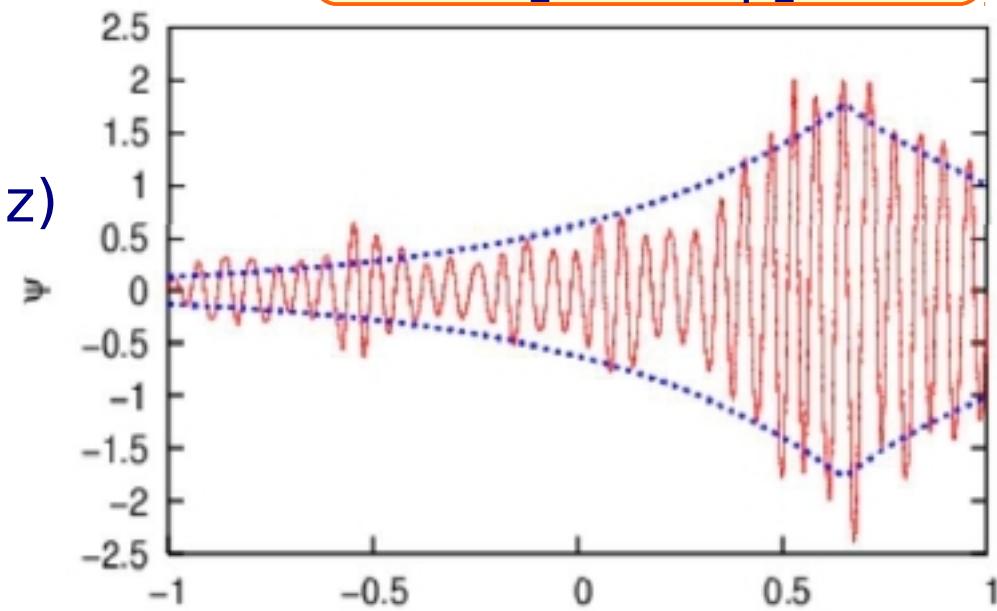
$$\ln[r(z)/r(0)] = k \int_0^z dz' [V(z')/2E] \sin[2\theta(z')]$$

- $\theta(z) = kz + \delta\theta^{(1)}(z) + \delta\theta^{(2)}(z) + \dots$
- Include  $\theta(z)$  into equation for  $r(z)$

$$\gamma(k) = \lim_{|z| \rightarrow \infty} \frac{\langle \ln[r(z)] \rangle}{|z|} = \sum_{n \geq 2} \gamma^{(n)}(k)$$



$k_E$  is defined by  
 $E = \hbar^2 k_E^2 / 2m = p_E^2 / 2m$



## Perturbation theory

$$\gamma^{(n)}(k) = \sigma_R^{-1} \left( \frac{\epsilon_R}{k\sigma_R} \right)^n f_n(k\sigma_R)$$

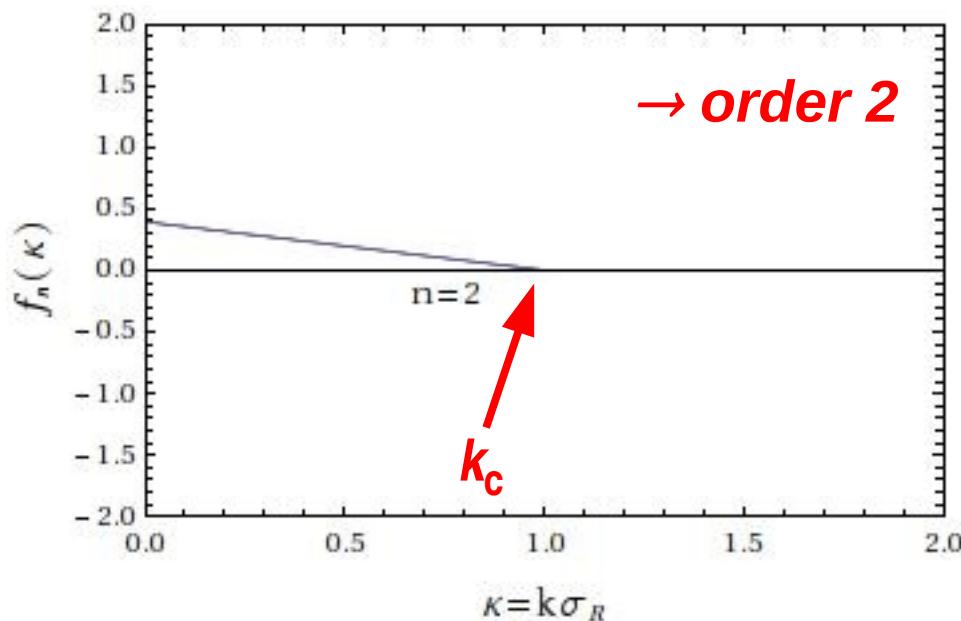
$$\text{where } \epsilon_R = 2m\sigma_R^2 V_R / \hbar^2$$

determined by the  
n-point correlation function

LSP et al., Phys. Rev. Lett **98**, 210401 (2007)

P. Lugan et al., Phys. Rev. A **80**, 023605 (2009)

E. Gurevich and O. Kenneth, Phys. Rev. A **79**, 063617 (2009)



## Order 2

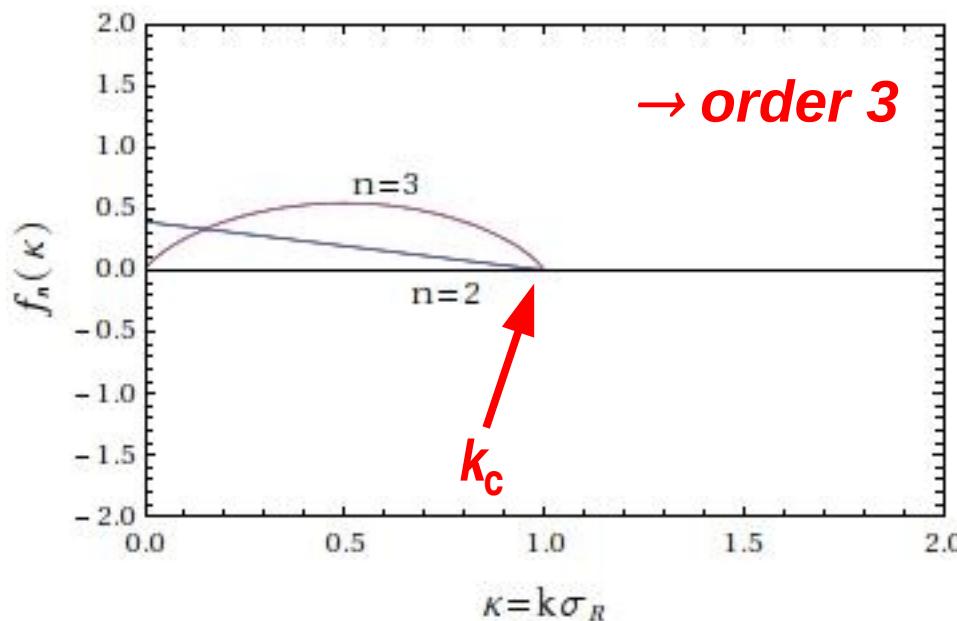
- cut-off at  $k_c \rightarrow$  defines an effective mobility edge
- no difference between  $V_R > 0$  and  $V_R < 0$

## Perturbation theory

$$\gamma^{(n)}(k) = \sigma_R^{-1} \left( \frac{\epsilon_R}{k\sigma_R} \right)^n f_n(k\sigma_R)$$

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## Order 3

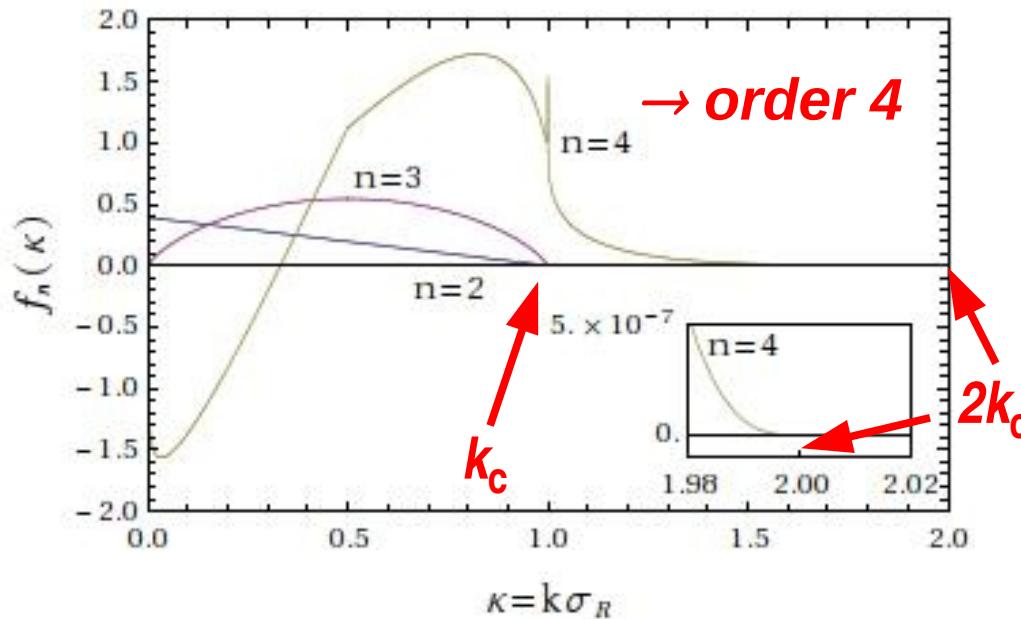
- same cut-off at  $k_c$
- opposite correction value of  $\gamma^{(3)}$  for  $V_R > 0$  and  $V_R < 0$
- this term can be relevant for experiments,  
see M. Piraud *et al.*, Phys. Rev. A 85, 063611 (2012)

## Perturbation theory

$$\gamma^{(n)}(k) = \sigma_R^{-1} \left( \frac{\epsilon_R}{k\sigma_R} \right)^n f_n(k\sigma_R)$$

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## Perturbation theory

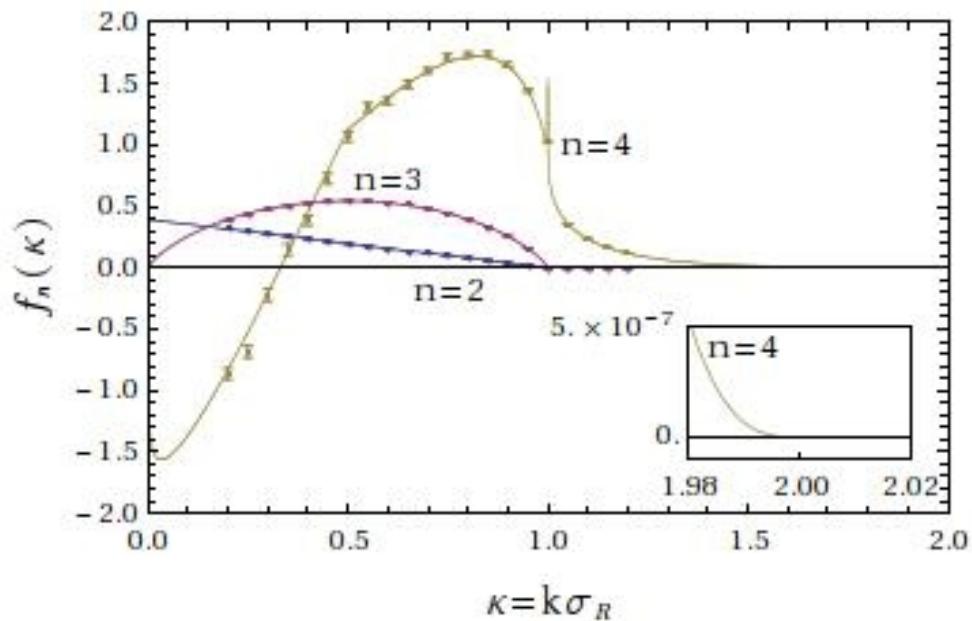
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$$\text{where } \epsilon_R = 2m\sigma_R^2 V_R / \hbar^2$$

determined by the  
n-point correlation function

### Order 4

- new cut-off at  $2k_c$  → defines a second effective mobility edge
- no difference between  $V_R > 0$  and  $V_R < 0$
- complicated behavior of  $\gamma^{(4)}$



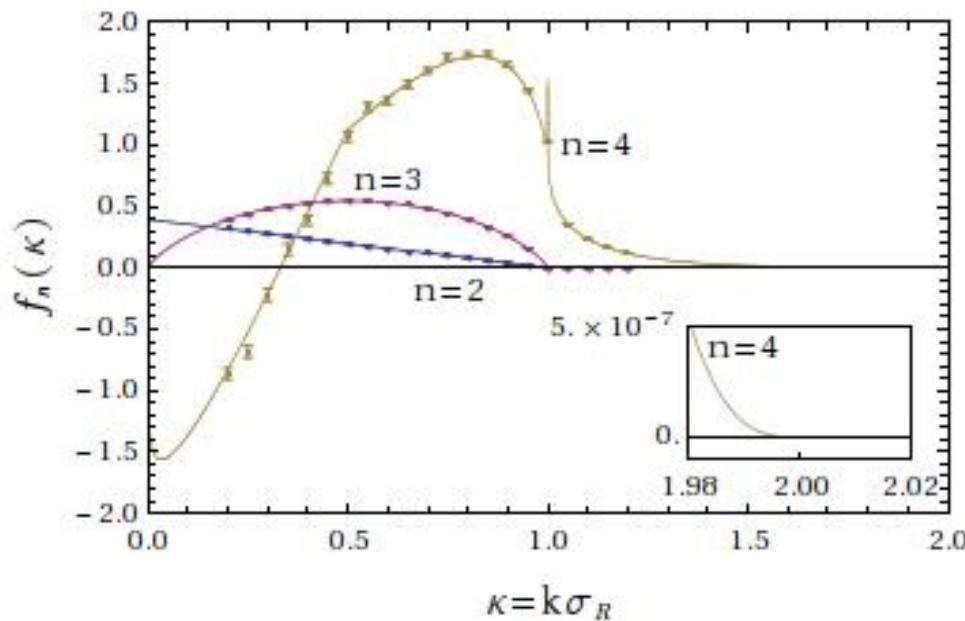
## Perturbation theory

$$\gamma^{(n)}(k) = \sigma_R^{-1} \left( \frac{\epsilon_R}{k\sigma_R} \right)^n f_n(k\sigma_R)$$

$$\text{where } \epsilon_R = 2m\sigma_R^2 V_R / \hbar^2$$

determined by the  
n-point correlation function

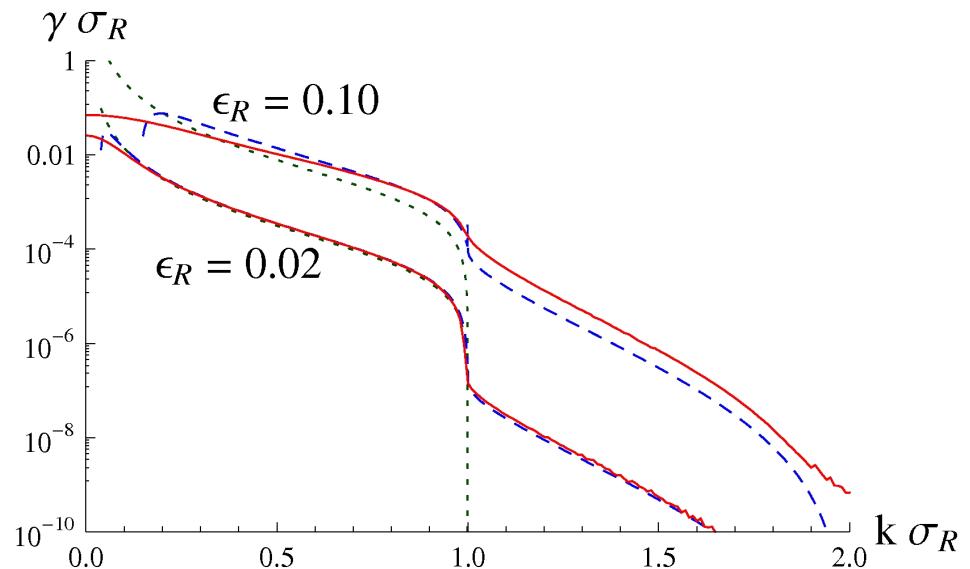
- Orders 2 to 4
- excellent agreement with exact numerical calculations (by D. Delande [P. Lugan *et al.*, Phys. Rev. A **80**, 023605 (2009)])



## Perturbation theory

$$\gamma^{(n)}(k) = \sigma_R^{-1} \left( \frac{\epsilon_R}{k\sigma_R} \right)^n f_n(k\sigma_R)$$

where  $\epsilon_R = 2m\sigma_R^2 V_R / \hbar^2$



## Lecture #1

1. Basics of quantum transport theory in disordered media
2. Ultracold atoms and speckle potentials  
Specific features of a controlled disorder
3. Anderson localization in 1D speckle potentials
  - 3.1 Effective mobility edges
  - 3.2 Expansion of Bose-Einstein condensates (exp. vs theory)

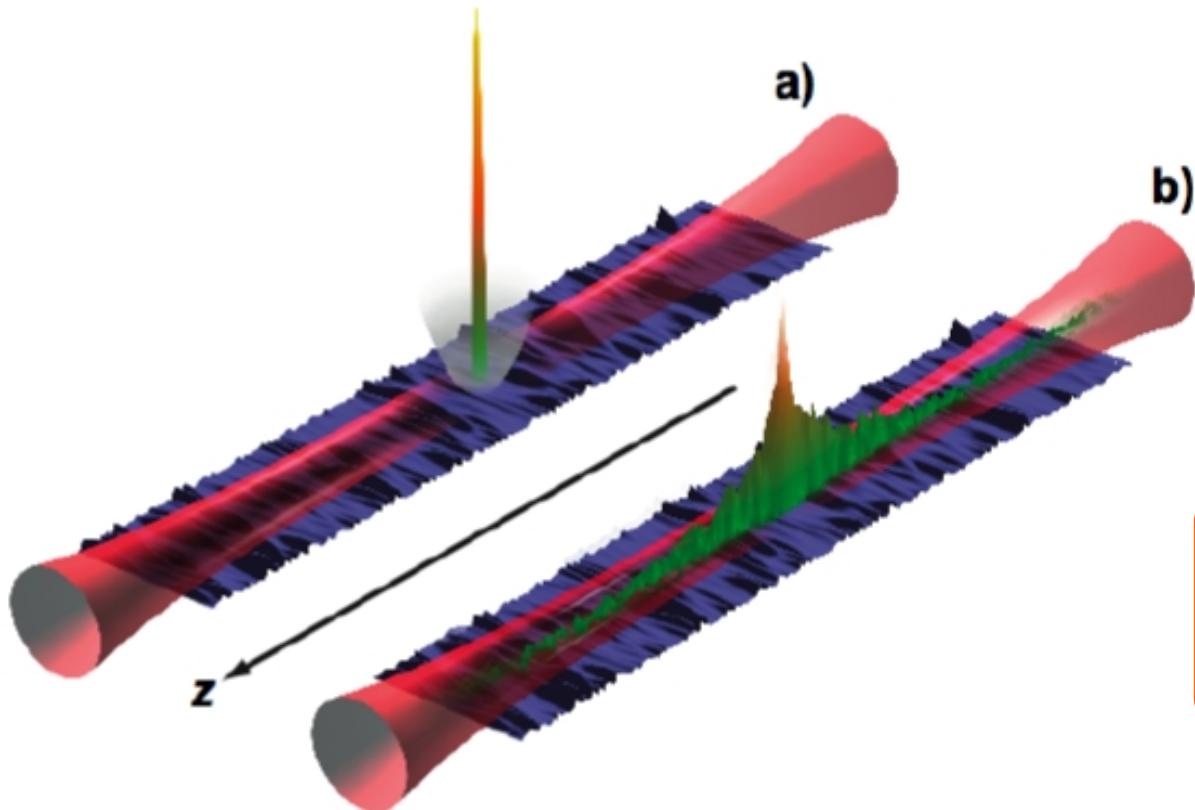
# Anderson Localization of an Expanding Matterwave in a 1D Speckle Potential

LSP *et al.*, Phys. Rev. Lett. **98**, 210401 (2007)

M. Piraud *et al.*, Phys. Rev. A **83**, 031603(R) (2011)

J. Billy *et al.*, Nature **453**, 891 (2008)

see also G. Roati *et al.*, Nature **453**, 895 (2008)

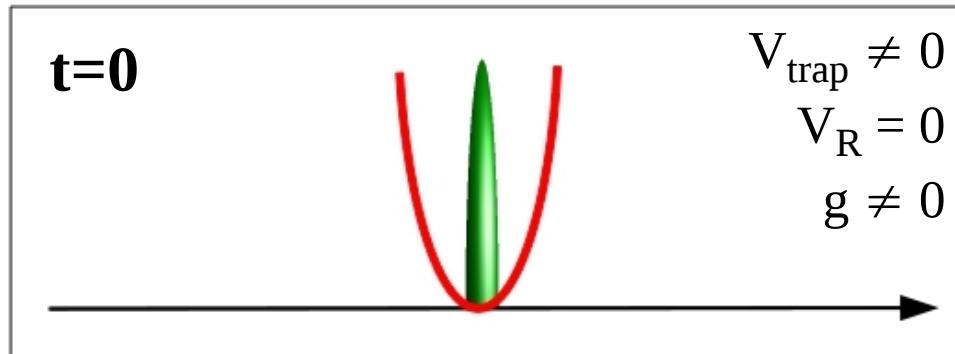


Superposition of matter waves with a broad energy distribution

$$n(z, t \rightarrow \infty) \simeq \int dE \mathcal{D}_E(E) P(z, t \rightarrow \infty | E)$$

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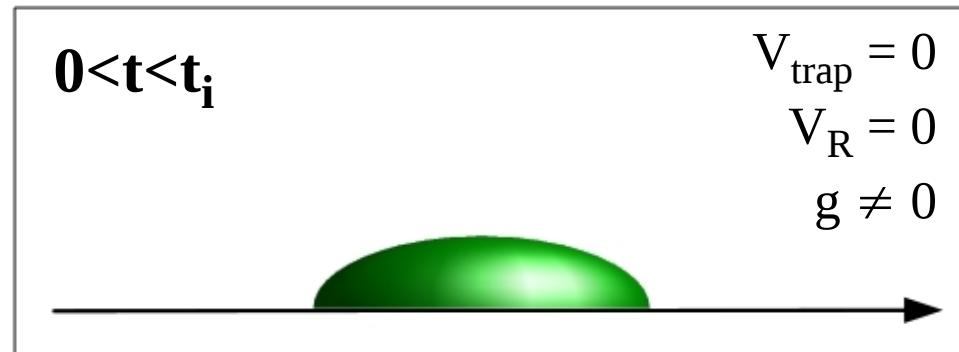
- ➊ Bose-Einstein condensate in a harmonic trap
- ➋ Thomas-Fermi regime :

$$n_0(z) \propto 1 - (z/L_{\text{TF}})^2$$

$$\theta(z) \equiv 0$$

LSP *et al.*, Phys. Rev. Lett. **98**, 210401 (2007)

M. Piraud *et al.*, Phys. Rev. A **83**, 031603(R) (2011)



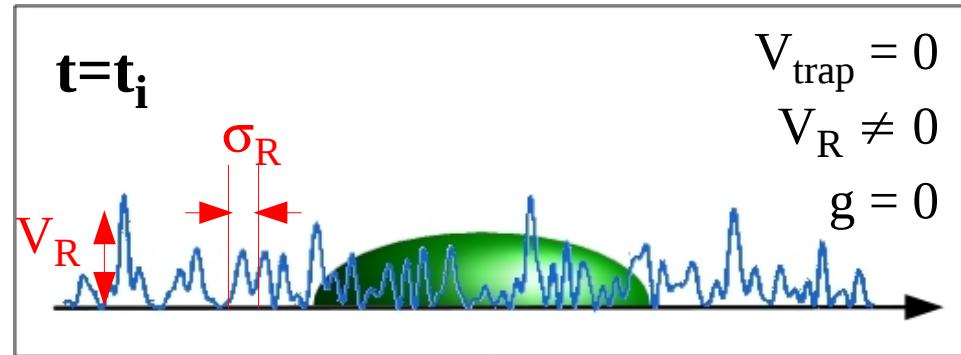
## Interaction-driven expansion of the BEC

- Scaling theory of expanding BEC (in TF regime)
- Momentum distribution :

$$\mathcal{D}_i(p) \propto [1 - (p\xi_{\text{in}}/\hbar)^2] \oplus \xi_{\text{in}} \equiv \hbar / \sqrt{4m\mu}$$

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LSP *et al.*, Phys. Rev. Lett. **98**, 210401 (2007)  
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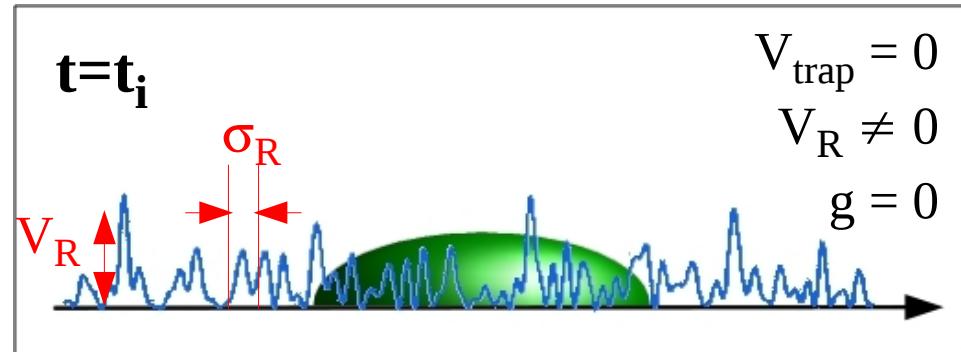


- ➊ Interactions off and disorder on
- ➋ « p » is no longer a good quantum number but « E » is !

$$\mathcal{D}_E(E) = \int dp \ A(p, E) \mathcal{D}_i(p)$$

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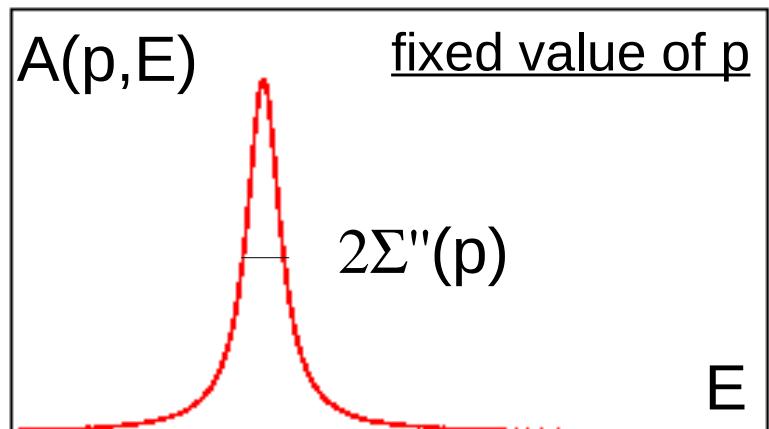


## ● Interactions off and disorder on

- «  $p$  » is no longer a good quantum number but «  $E$  » is !

$$\mathcal{D}_E(E) = \int dp \underbrace{A(p, E)}_{\text{spectral function}} \mathcal{D}_i(p)$$

spectral function



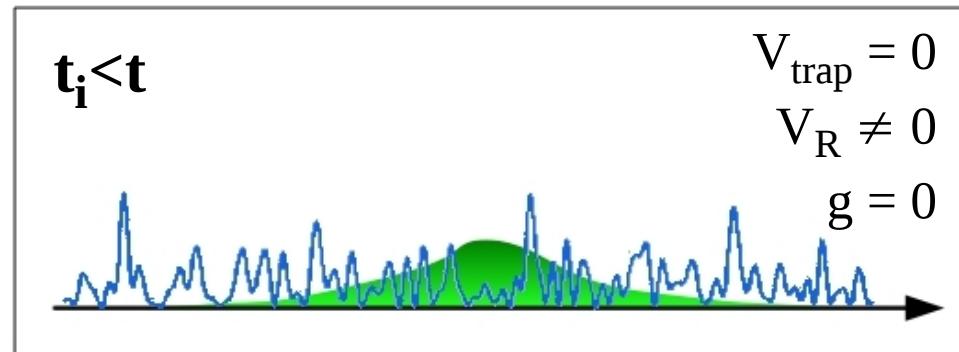
FT of the  
correlation  
function of the  
disorder

$$\Sigma''(p) \simeq -(m/2\hbar p) \left\{ \tilde{C}(0) + \tilde{C}(2p) \right\}$$

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LSP *et al.*, Phys. Rev. Lett. **98**, 210401 (2007)

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## Localization of the wave packet

$$P_\infty(z|E) = \frac{\pi^2 \gamma(E)}{8} \int_0^\infty du \ u \sinh(\pi u) \left[ \frac{1+u^2}{1+\cosh(\pi u)} \right] \\ \times \exp \left\{ -(1+u^2) \gamma(E) |z|/2 \right\},$$

<u>short distance :</u>	$\ln(P) \sim -2\gamma(E) z $
<u>long distance :</u>	$\ln(P) \sim -0.5\gamma(E) z $

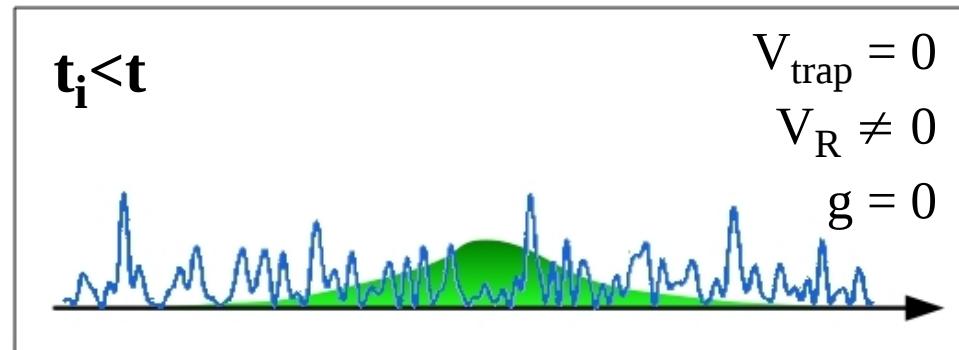
V.L. Berezinskii, Sov. Phys. JETP **38**, 620 (1974)

A.A. Gogolin *et al.*, Sov. Phys. JETP **42**, 168 (1976)

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$$\gamma(E) \simeq (m^2/2\hbar^2 p_E^2) \tilde{C}(2p_E)$$

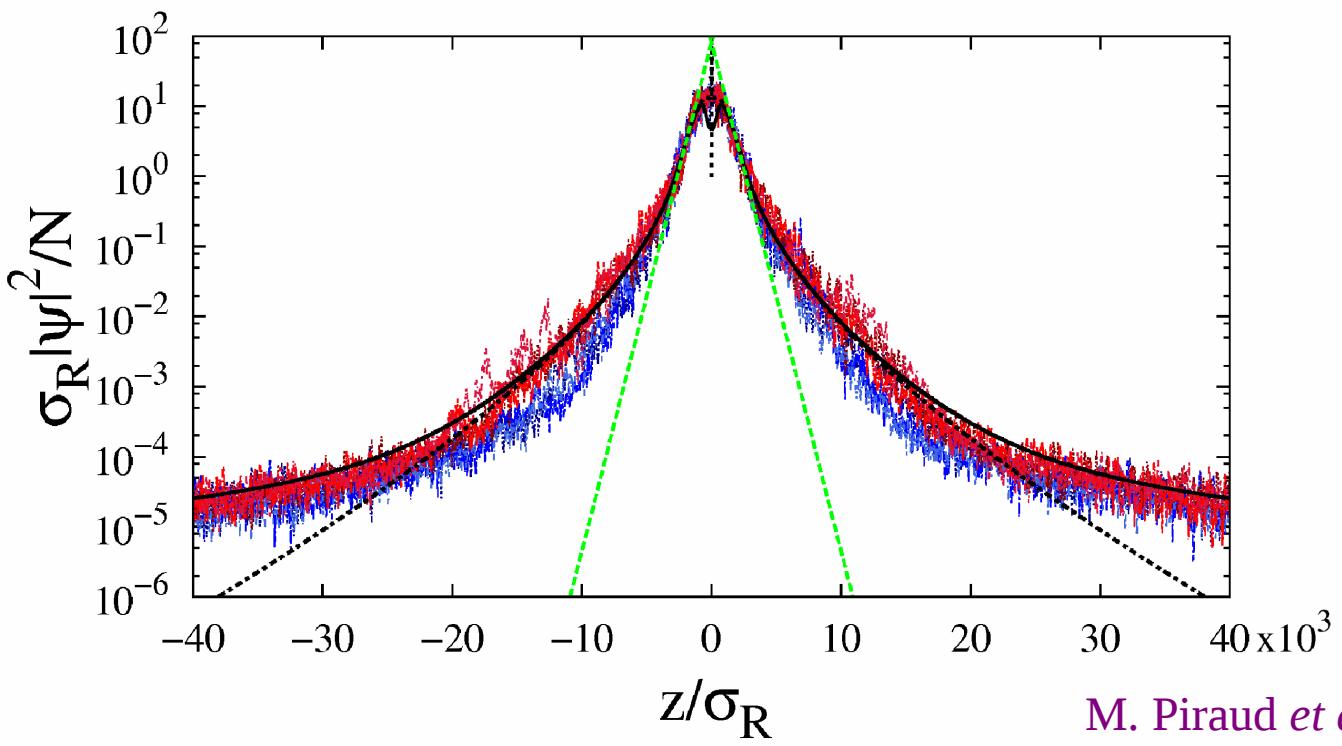


FT of the  
correlation  
function of the  
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V.L. Berezinskii, Sov. Phys. JETP **38**, 620 (1974)

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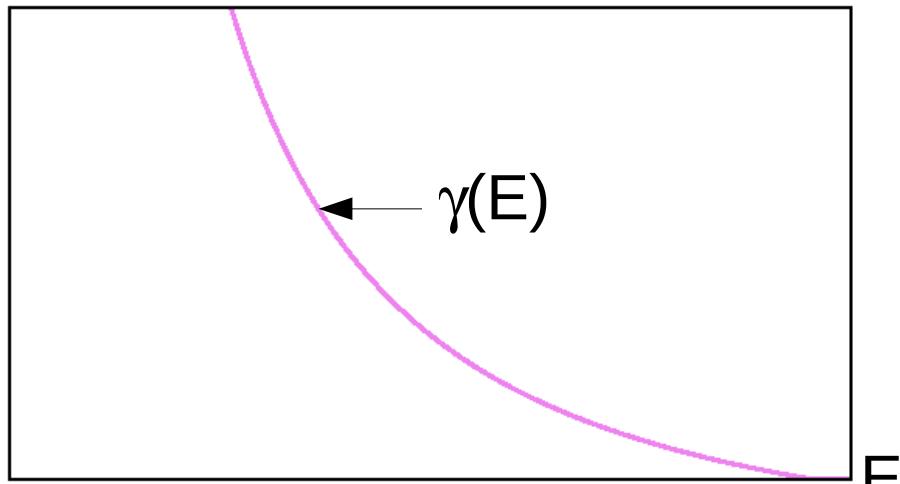
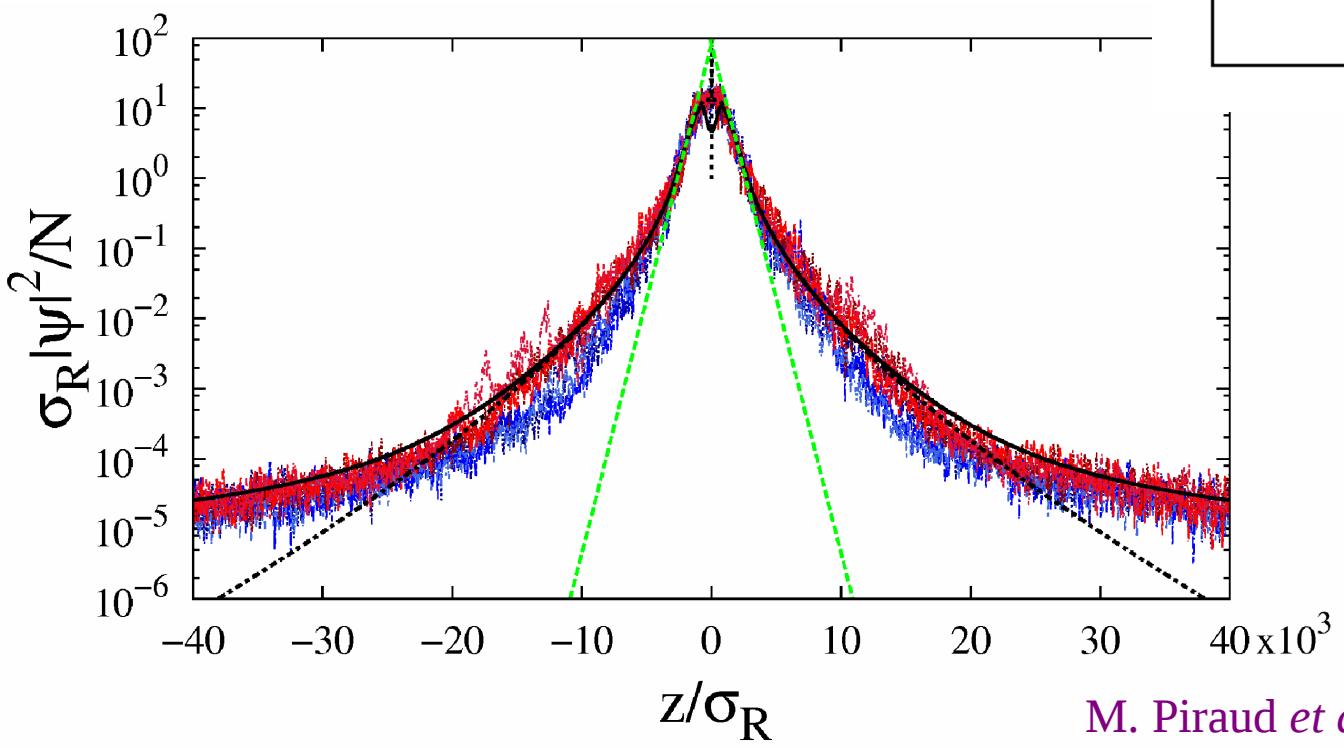
$n(z, t \rightarrow \infty) \simeq \int dE \mathcal{D}_E(E) P(z, t \rightarrow \infty | E)$  versus numerics



# Anderson Localization of an Expanding Matterwave in a 1D Speckle Potential

$$n(z, t \rightarrow \infty) \simeq \int dE \mathcal{D}_E(E) P(z, t \rightarrow \infty | E)$$

versus numerics

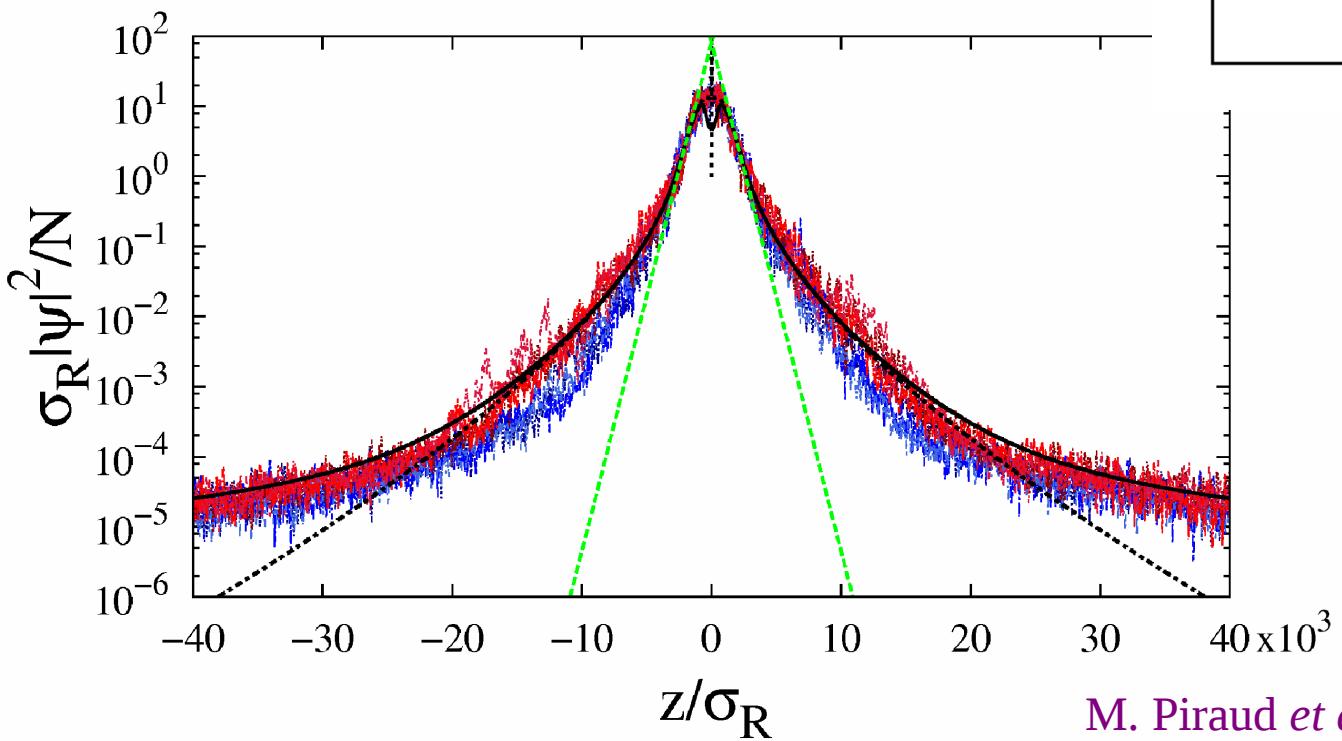


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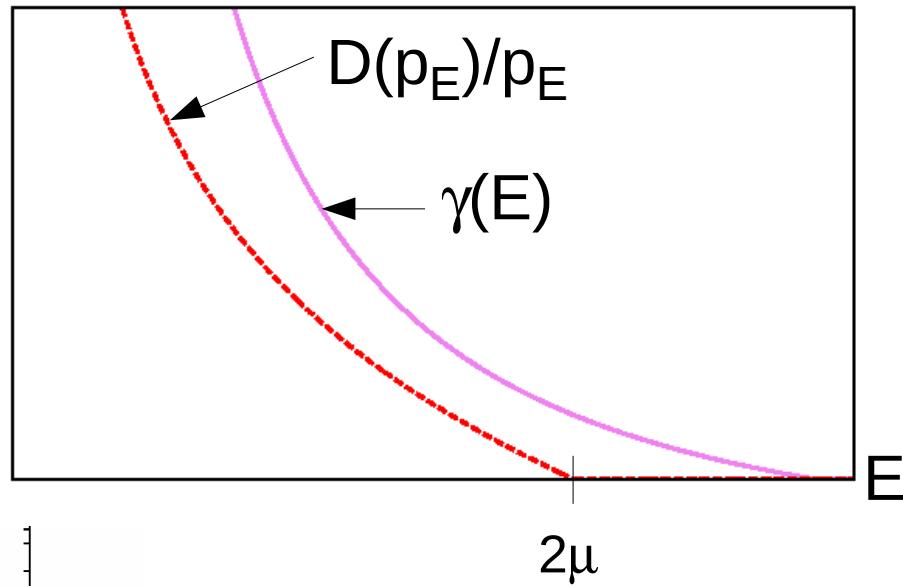
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$$\mathcal{D}_E(E) = \int dp A(p, E) \mathcal{D}_i(p) \simeq \frac{m}{p_E} \mathcal{D}_i(p_E)$$

(«on-shell» approximation)



versus numerics

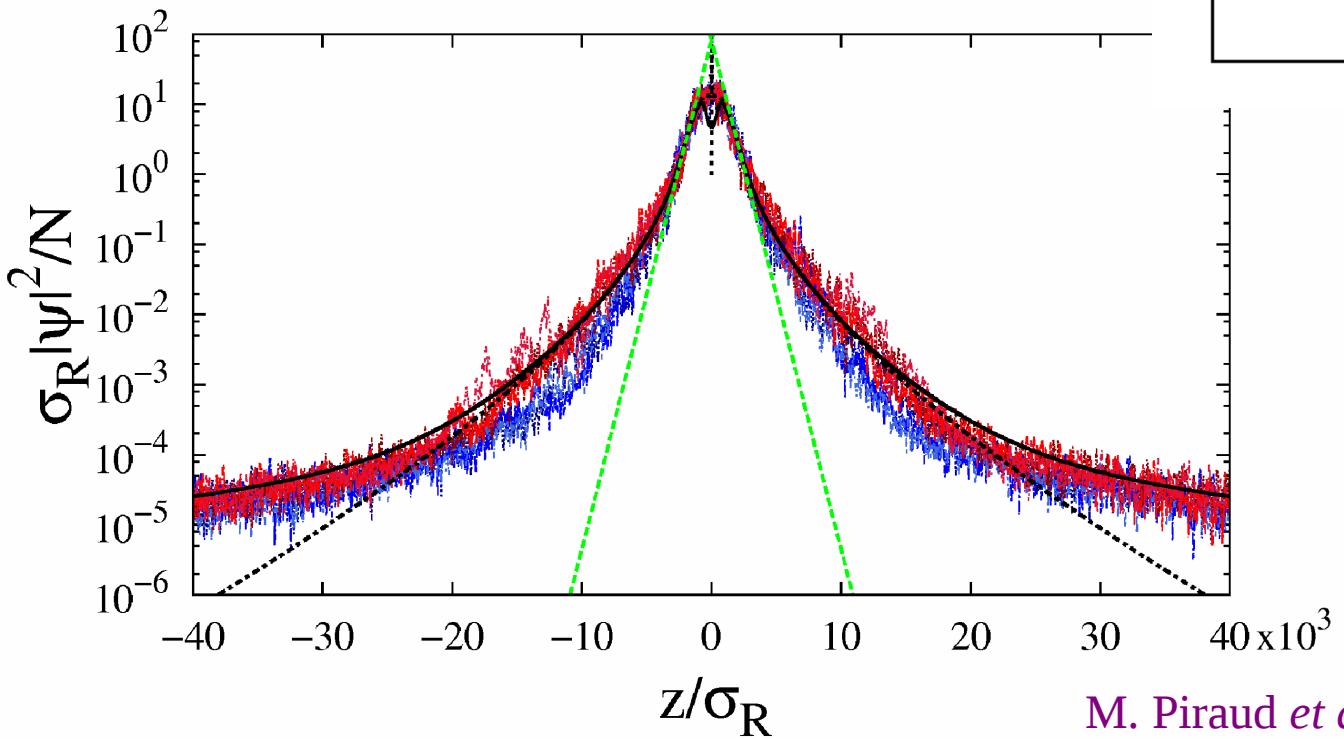


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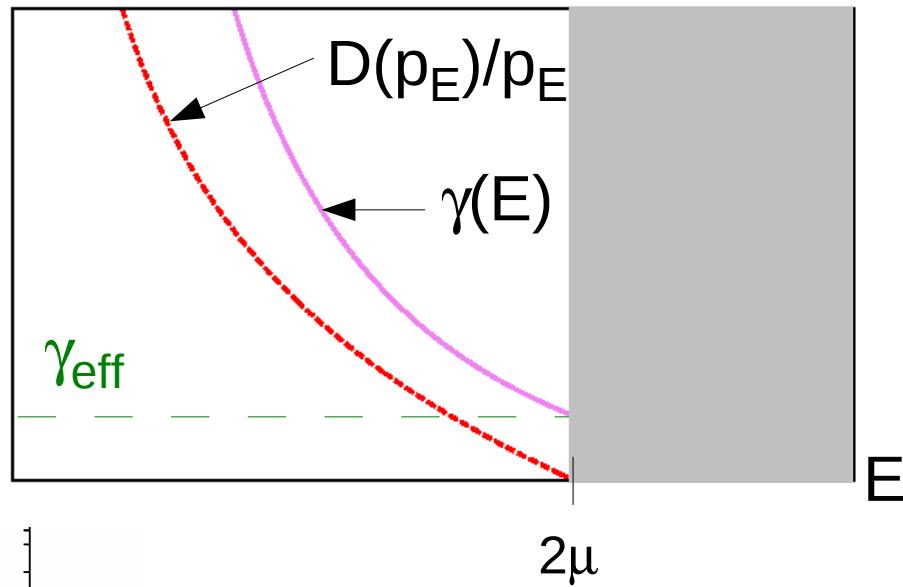
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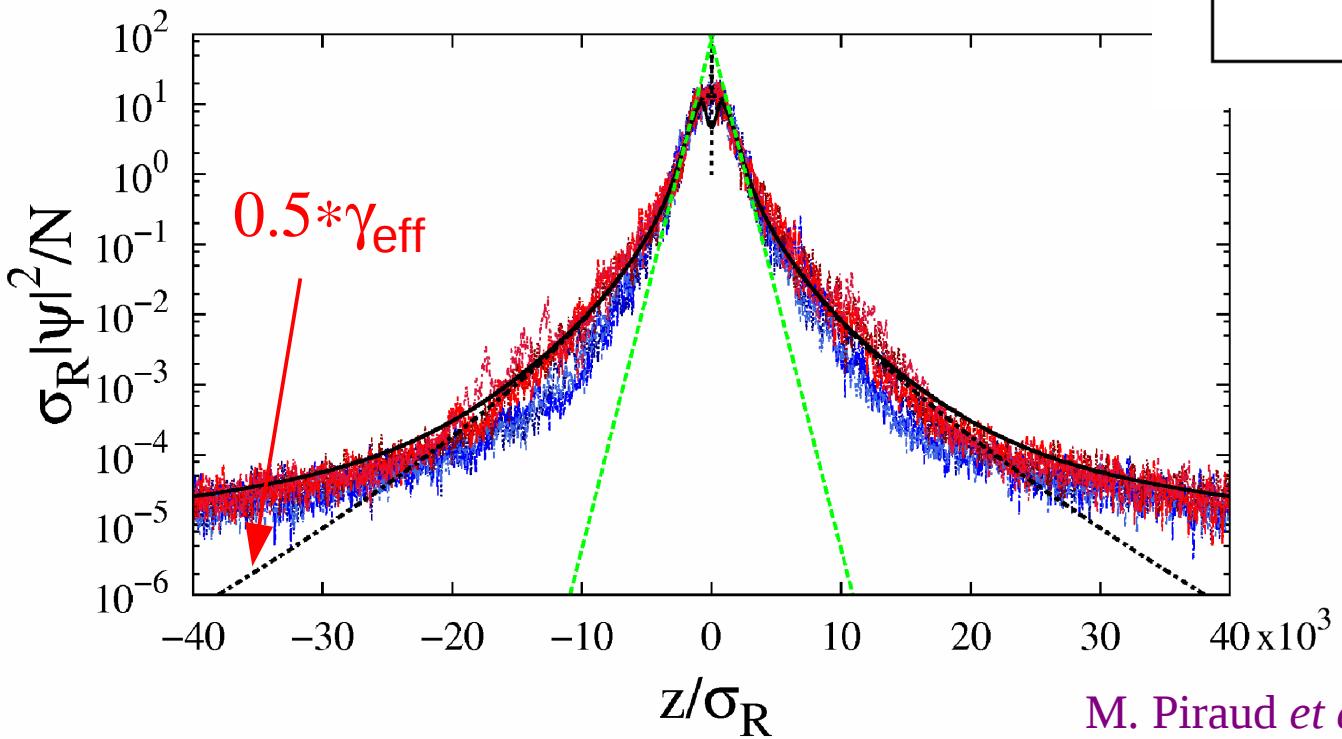


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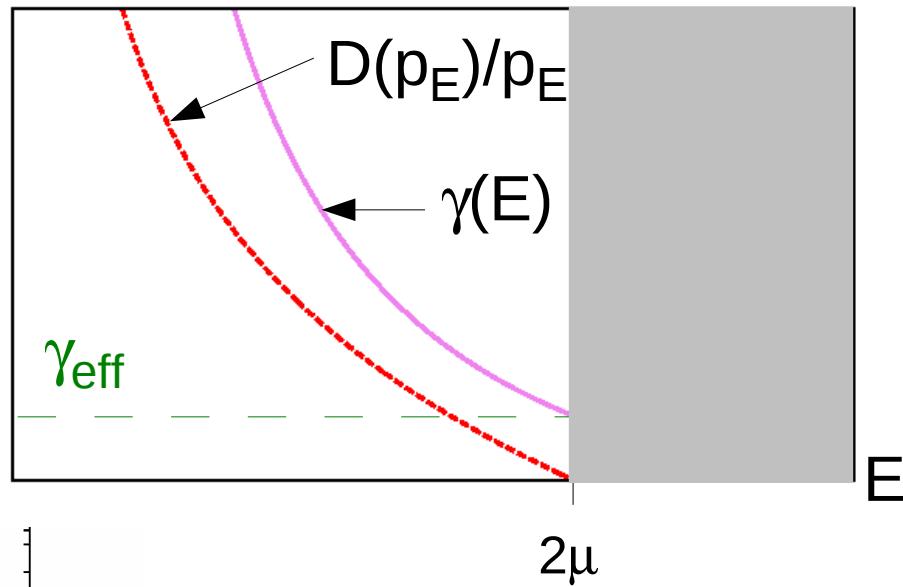
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versus numerics

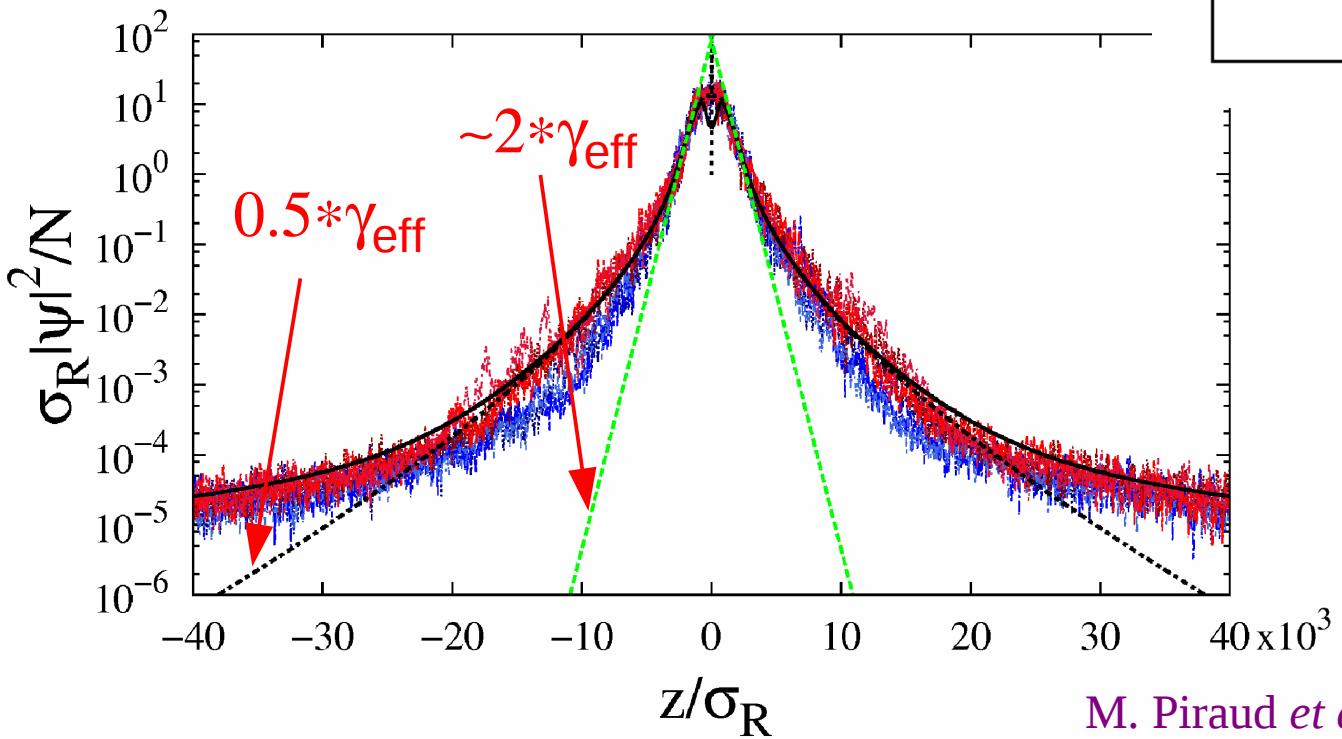


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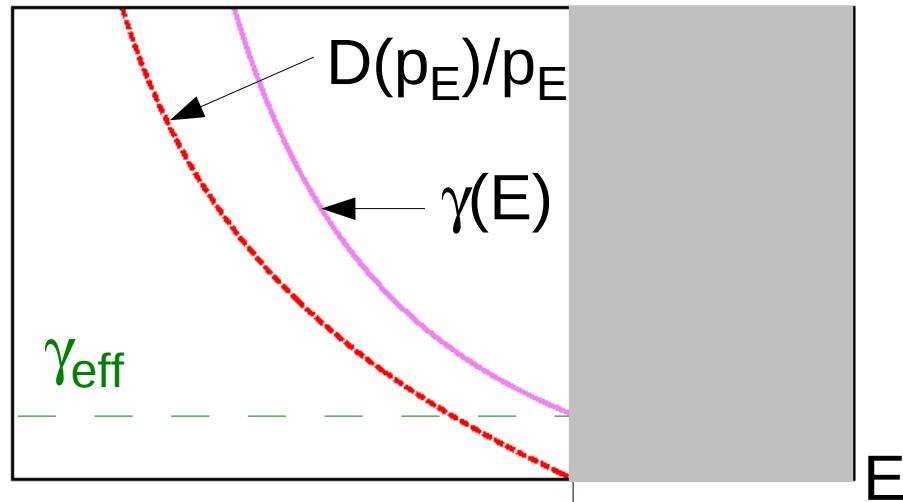
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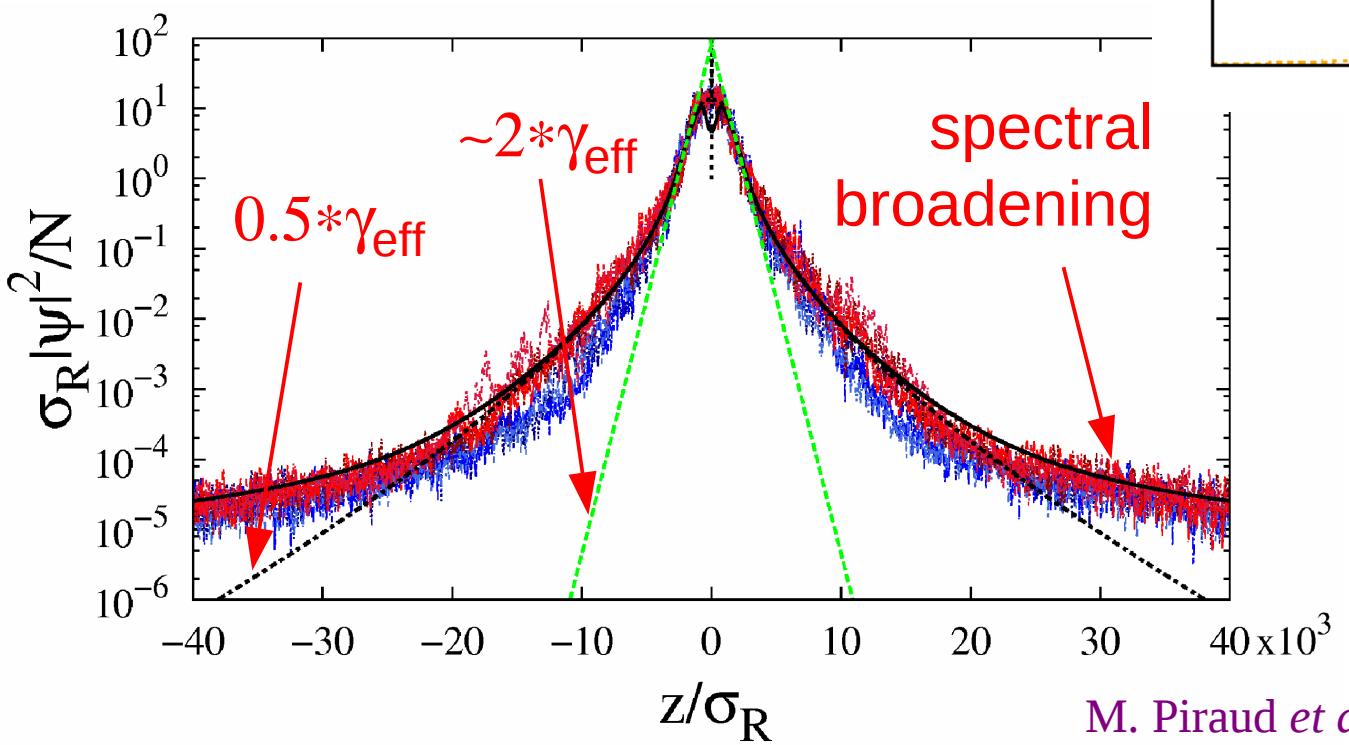
versus numerics



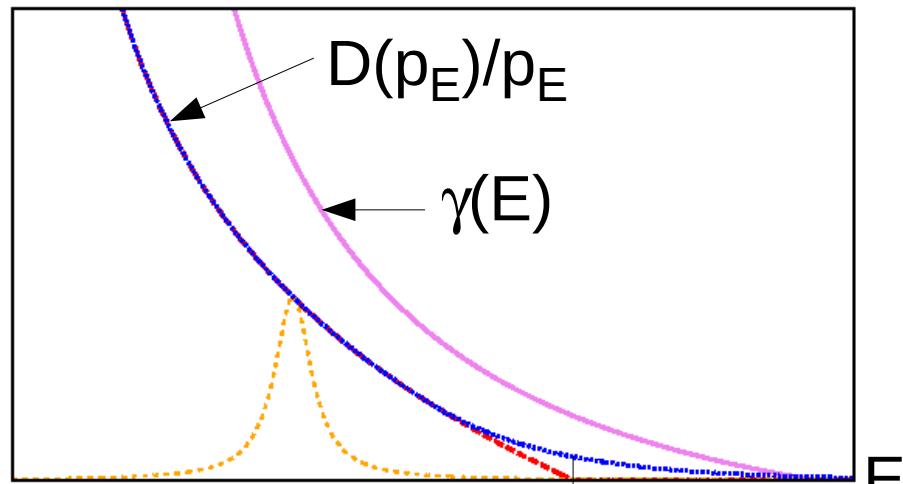
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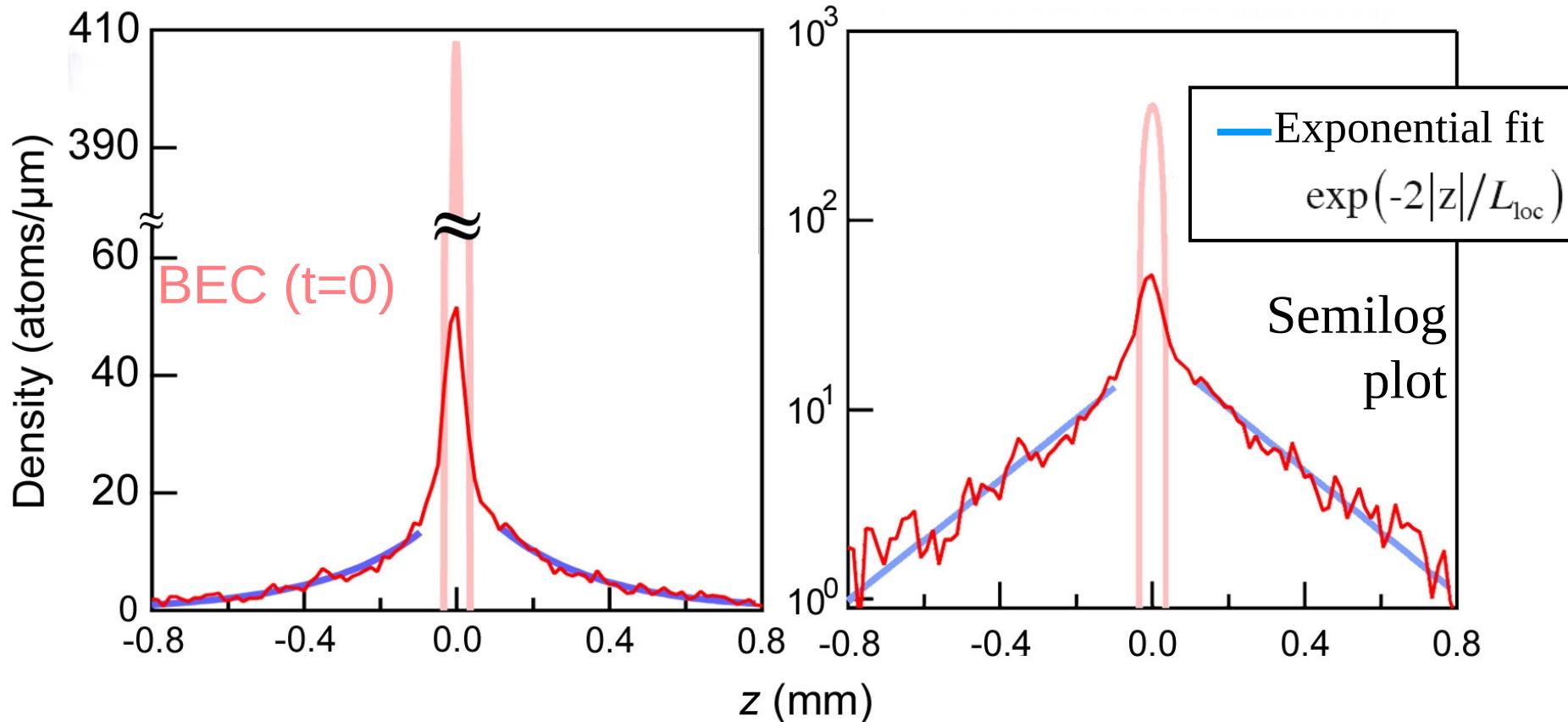


# Anderson Localization of an Expanding Matterwave in a 1D Speckle Potential

J. Billy *et al.*, Nature 453, 891 (2008)

see also G. Roati *et al.*, Nature 453, 895 (2008)

BEC parameters :  $N=1.7 \times 10^4$  atoms, ( $\xi_{\text{in}} \approx 1.5\sigma_R$ )  
 Weak disorder :  $V_R/\mu_{\text{in}} = 0.12 \ll 1$

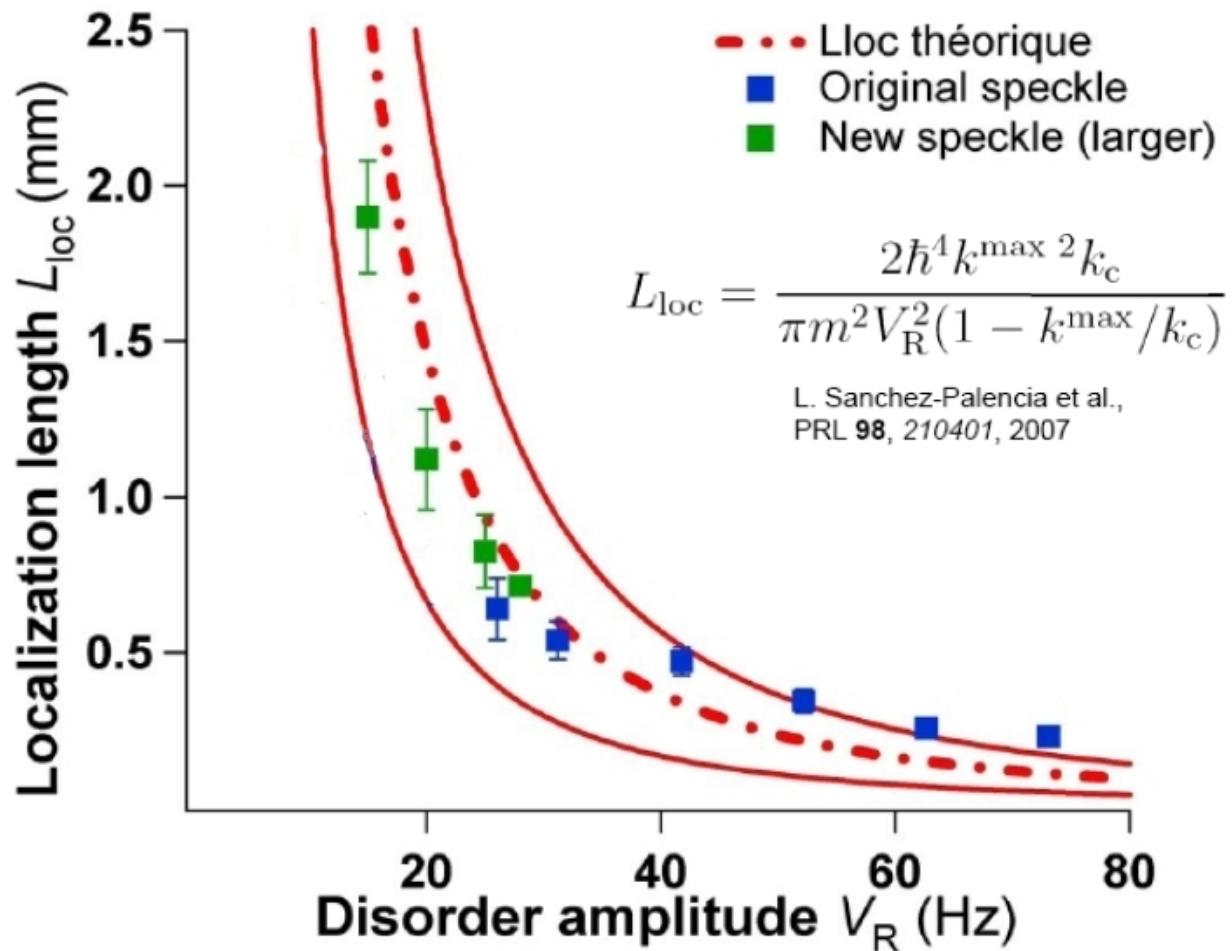


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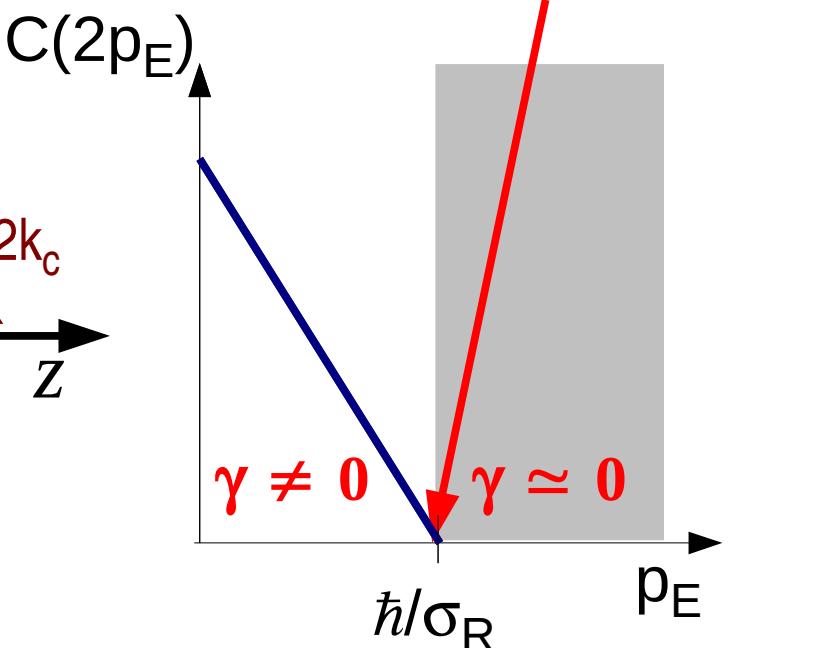
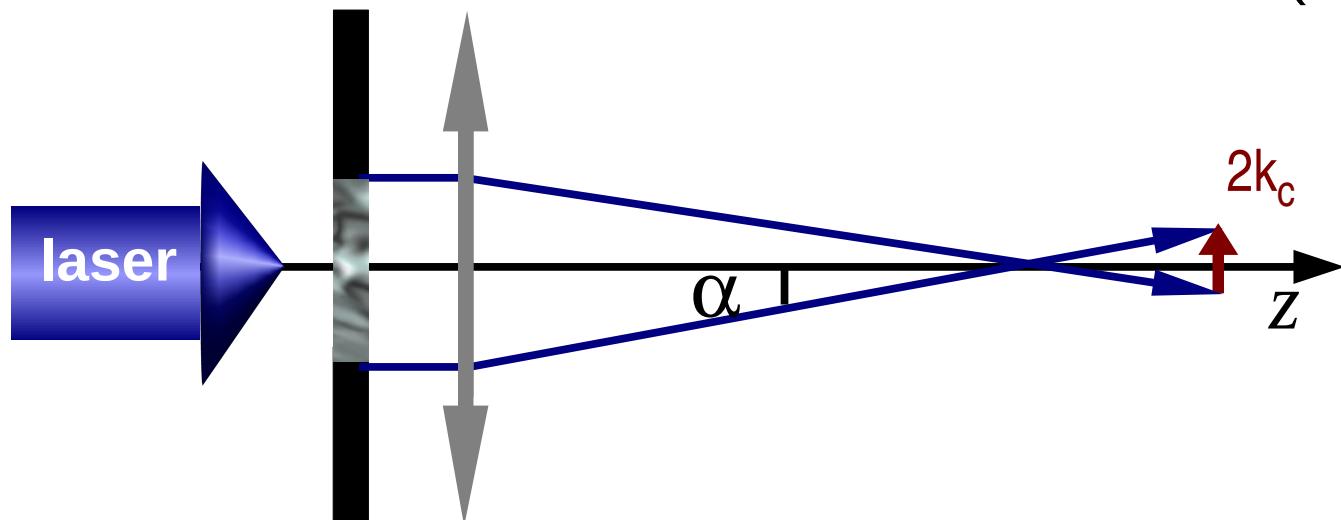
Fair agreement with  
theoretical prediction

$$\gamma(E) \simeq (m^2/2\hbar^2 p_E^2) \tilde{C}(2p_E)$$

FT of the correlation function of the disorder

## Correlation function of a 1D speckle potential

$$\tilde{C}(p) = \pi V_R^2 \sigma_R [1 - |p| \sigma_R / 2\hbar] \oplus$$



LSP *et al.*, Phys. Rev. Lett. **98**, 210401 (2007)

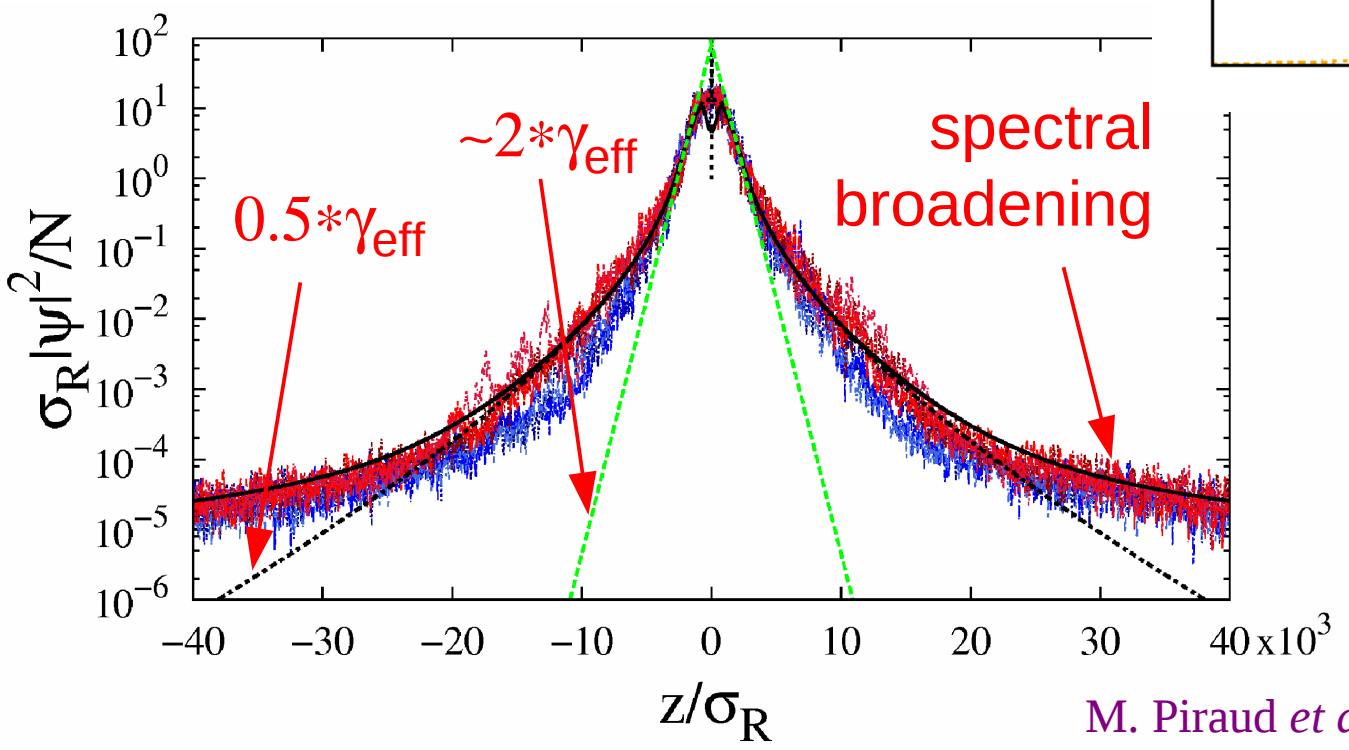
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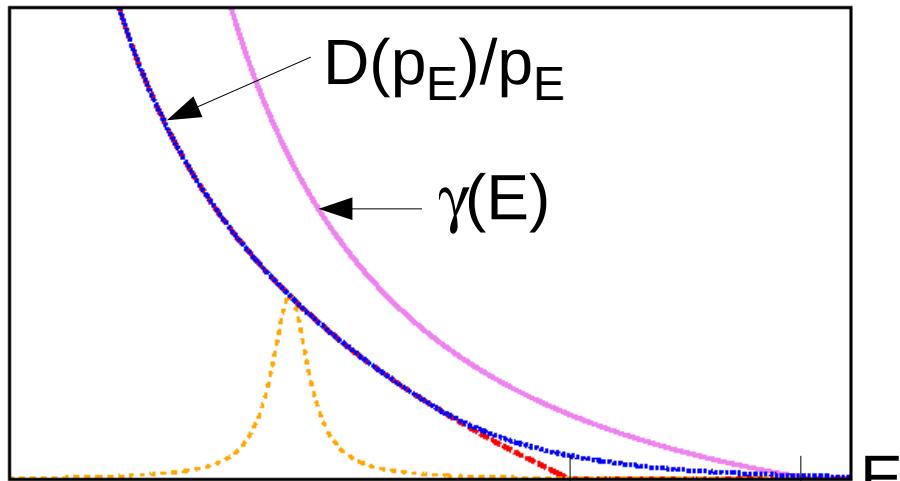
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versus numerics



$$2\mu = \hbar^2/2m\xi_{in}^2$$

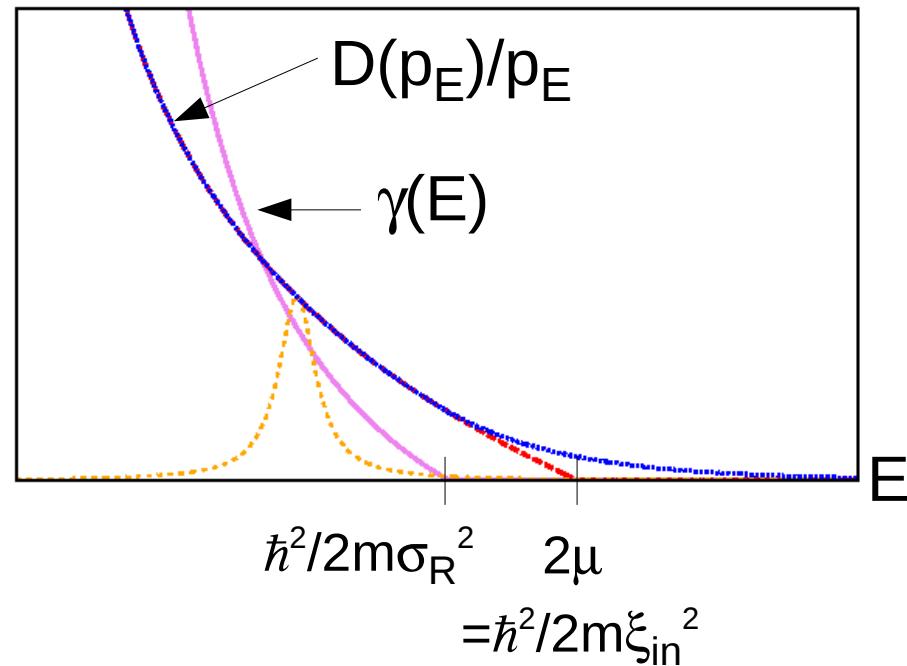
For  $\xi_{in} > \sigma_R$

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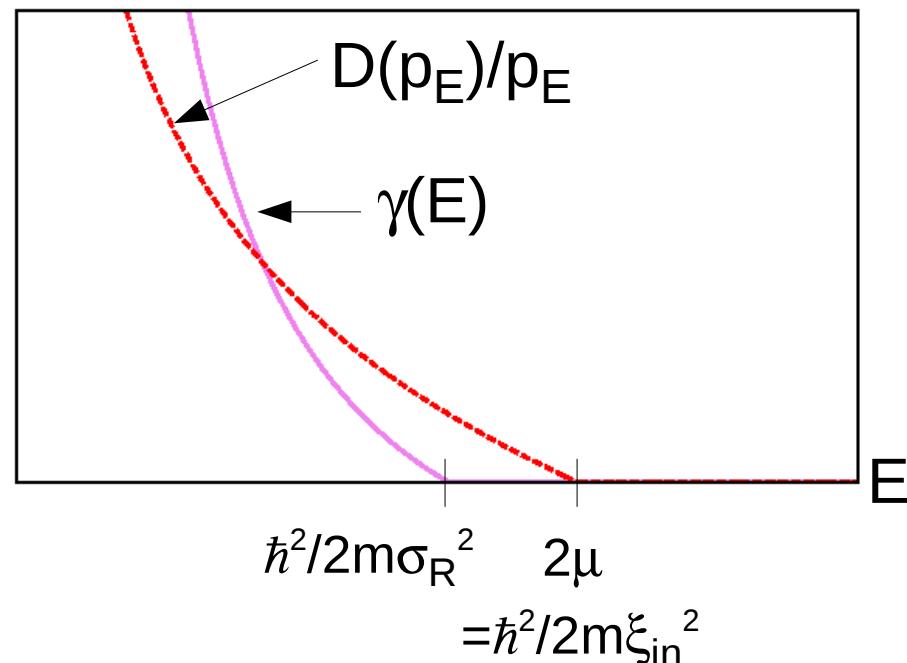
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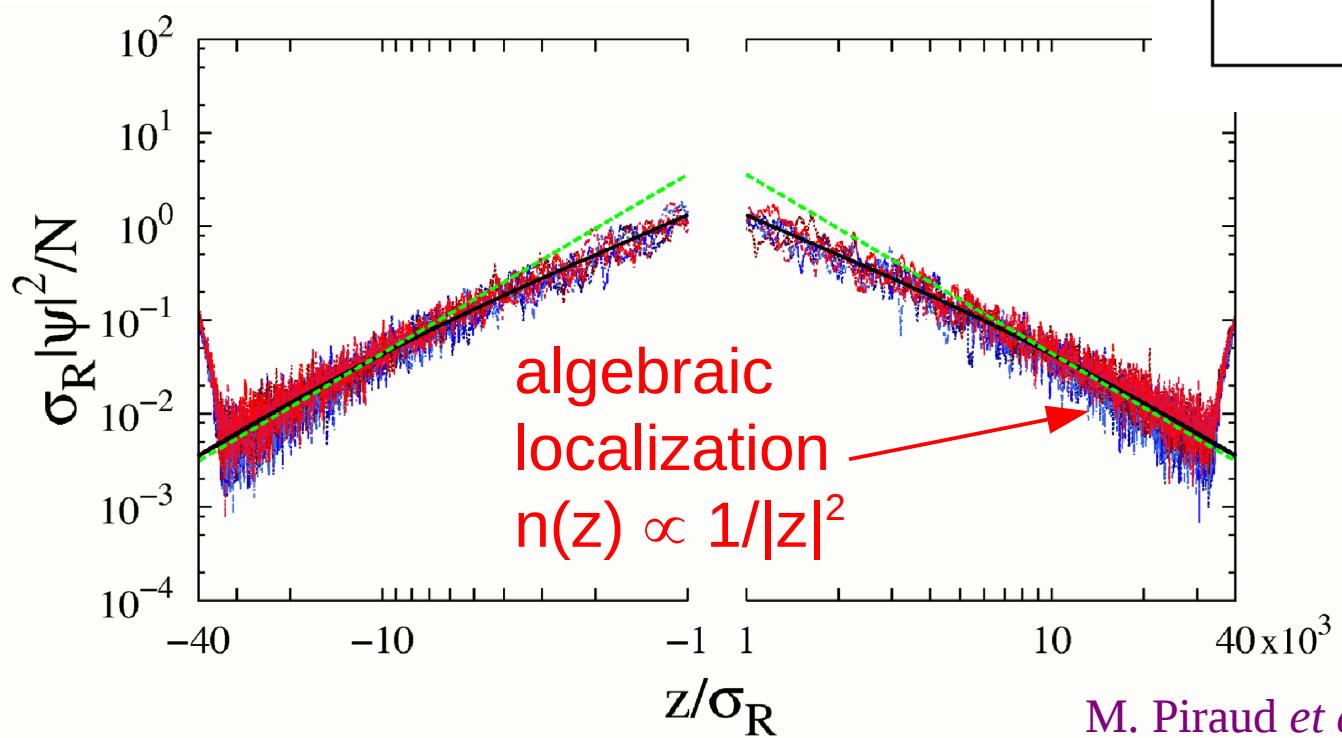
algebraic  
localization  
 $n(z) \propto 1/|z|^2$

For  $\xi_{in} < \sigma_R$

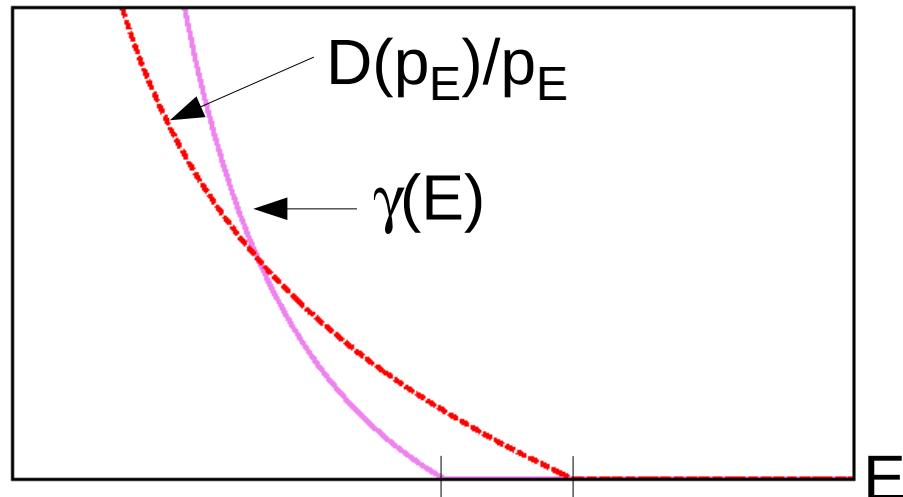
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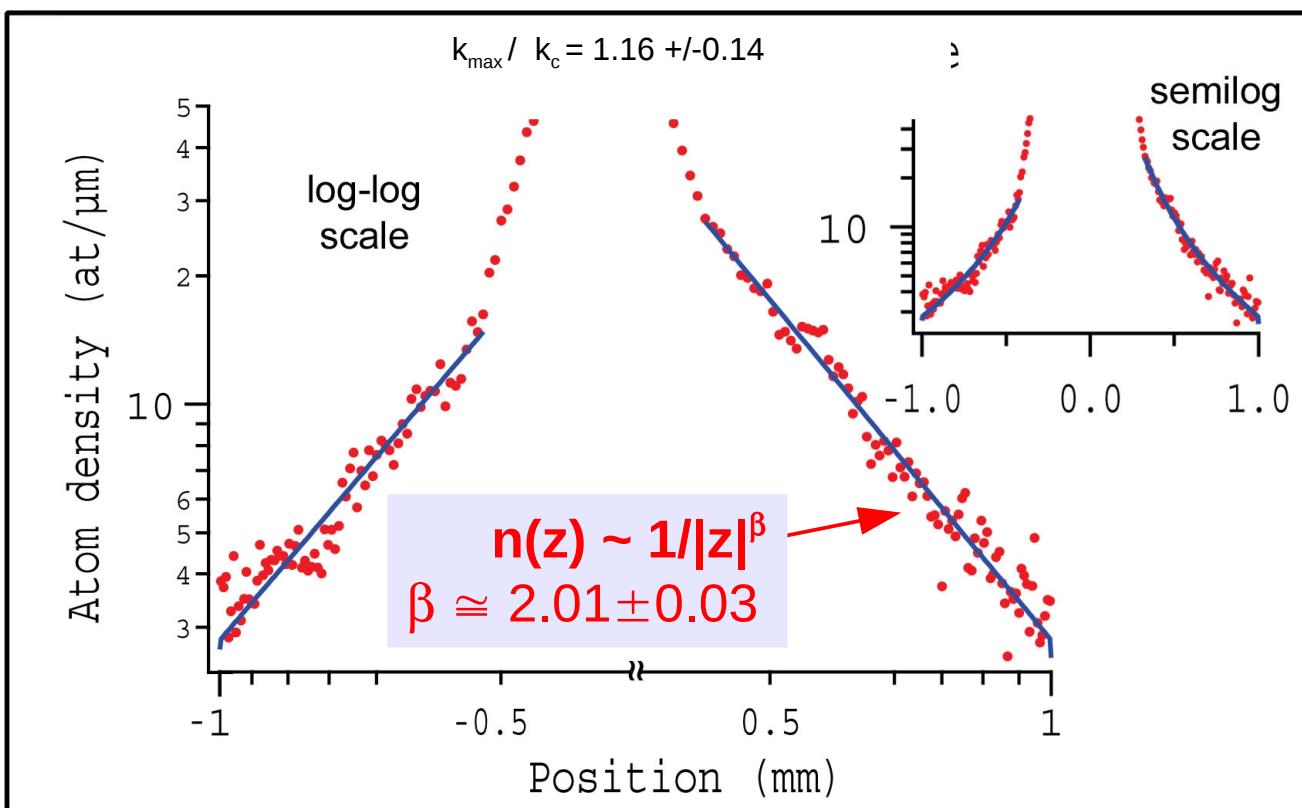
$$\begin{aligned} \hbar^2/2m\sigma_R^2 &= 2\mu \\ &= \hbar^2/2m\xi_{in}^2 \end{aligned}$$

For  $\xi_{in} < \sigma_R$

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J. Billy *et al.*, Nature 453, 891 (2008)

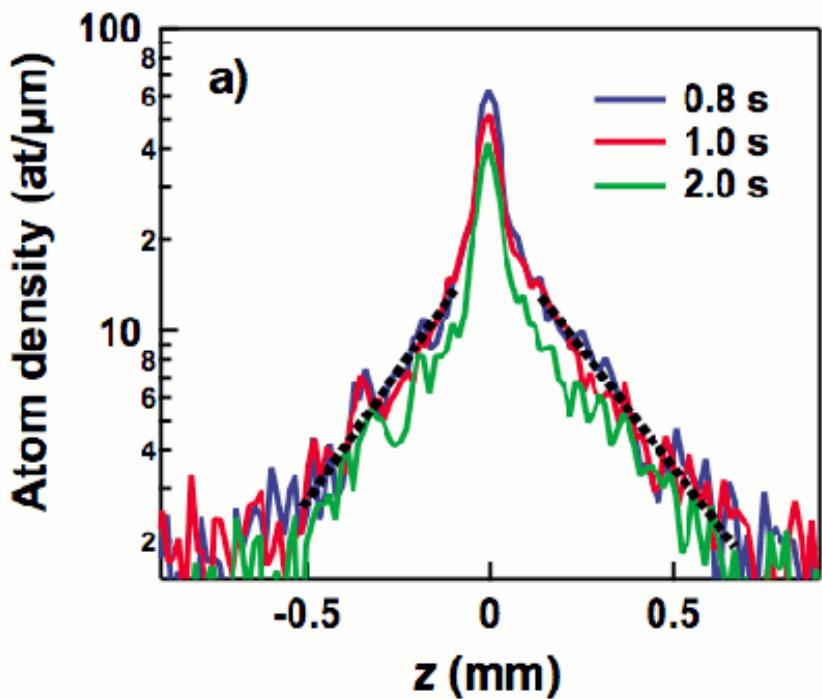
BEC parameters :  $N=1.7 \times 10^5$  atoms, ( $\xi_{\text{in}} \approx 0.8\sigma_R$ )  
 Weak disorder :  $V_R/\mu_{\text{in}} = 0.15 \ll 1$



# Anderson Localization in a 1D Speckle Potential : Conclusions

- ➊ Model of 1D Anderson localization

- ➊ Model of 1D Anderson localization
- ➋ For  $\xi_{in} > \sigma_R$ ,
  - short distance :  $\ln[n(z)] \sim -2\gamma_{\text{eff}}$  → observed



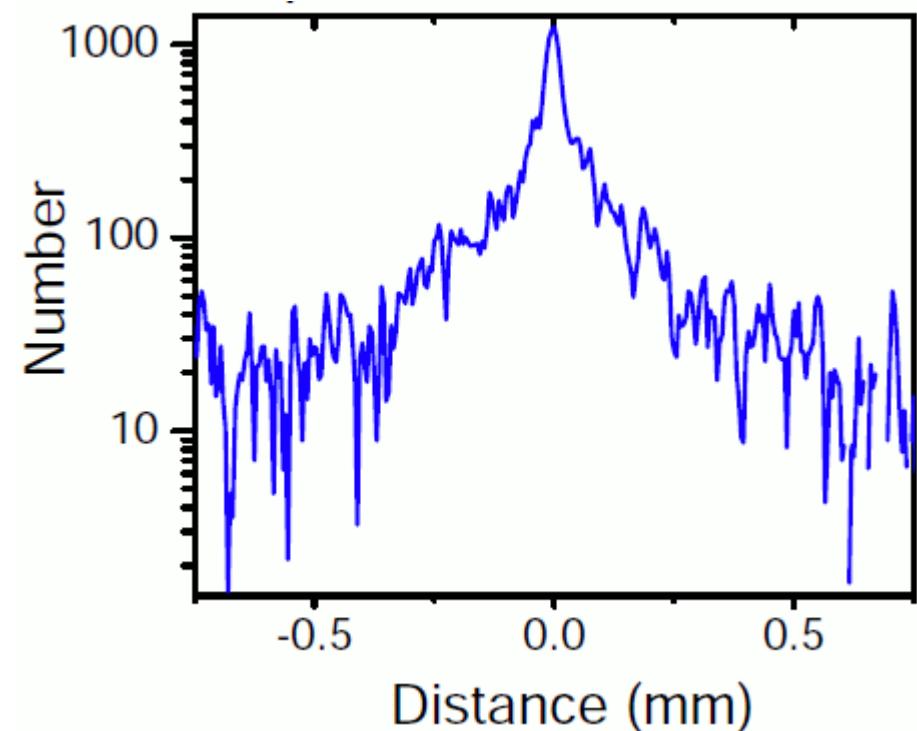
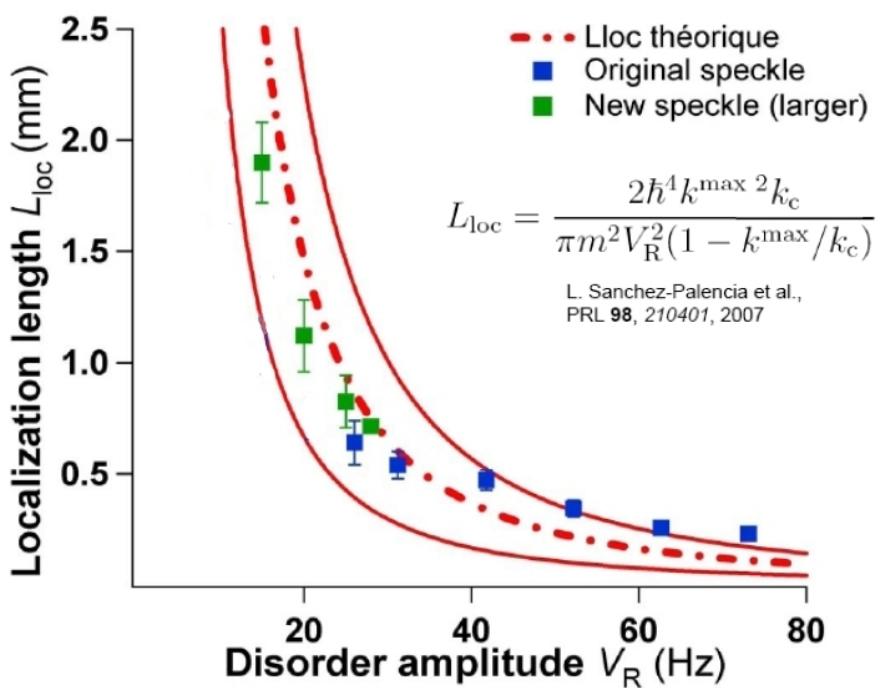
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- large distance :  $\ln[n(z)] \sim -0.5\gamma_{\text{eff}}$  → possible signatures (?)  
+ deviations from exponential decay



J. Billy et al., Nature 453, 891 (2008)

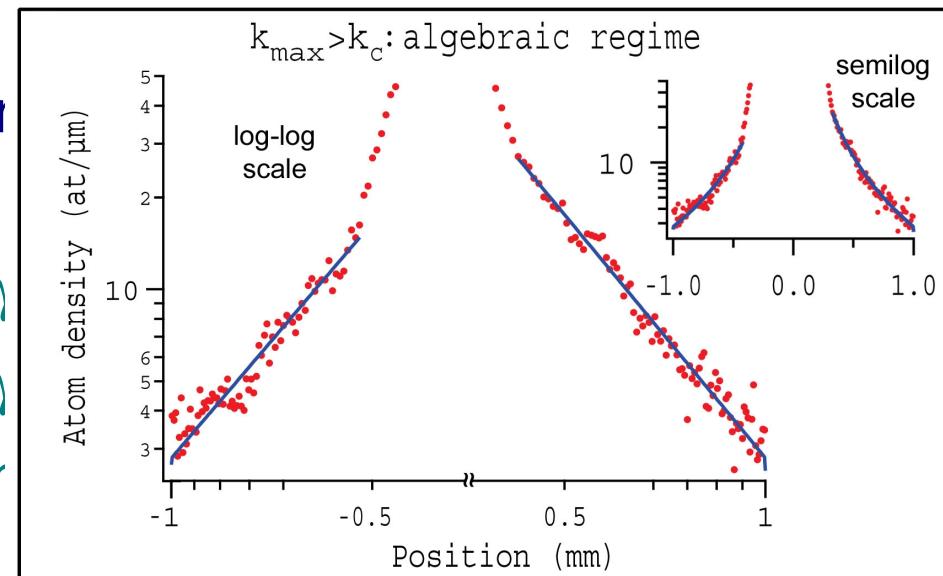
Courtesy of R. Hulet, unpublished (2009)

# Anderson Localization in a 1D Speckle Potential : Conclusions

- Model of 1D Anderson localization

- For  $\xi_{in} > \sigma_R$ ,

- short distance :  $\ln[n(z)] \sim -2|z|$
- large distance :  $\ln[n(z)] \sim -0.5|z|^2$   
+ deviations from theory



- For  $\xi_{in} < \sigma_R$ ,

J. Billy *et al.*, Nature 453, 891 (2008)

- central role of special finite-range correlations in 1D speckle potentials → effective mobility edge
- algebraic localization,  $n(z) \sim 1/|z|^2$  → observed
- not in contradiction with theorems [P. Lughan *et al.*, PRA (2009)] !

## Model of 1D Anderson localization

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  - not in contradiction with theorems [P. Lughan *et al.*, PRA (2009)] !
- Perspectives
  - tailored correlations [M. Piraud *et al.*, Phys. Rev. A **85**, 063611 (2012)]
  - extension to many-body localization

# Bibliography (I)

## ➊ Anderson localization in 1D disorder

- ➊ Lifshits, Gredeskul & Pastur, *Introduction to the Theory of Disordered Systems* (1988)
- ➊ V.L. Berezinskii, Sov. Phys. JETP **38**, 620 (1974)
- ➊ A.A. Gogolin *et al.*, Sov. Phys. JETP **42**, 168 (1976)

## ➋ Speckle potentials in ultracold atom systems

- ➋ J. W. Goodman, *Speckle Phenomena in Optics* (2007)
- ➋ D. Clément *et al.*, New J. Phys. **8**, 165 (2006)
- ➋ B. Shapiro, J. Phys. A: Math. Theor. **45**, 143001 (2012)

## ➌ Anderson localization in 1D speckle potentials

- ➌ LSP *et al.*, Phys. Rev. Lett. **98**, 210401 (2007)
- ➌ J. Billy *et al.*, Nature **453**, 891 (2008)
- ➌ P. Lugan *et al.*, Phys. Rev. A **80**, 023605 (2009)
- ➌ E. Gurevich and O. Kenneth, Phys. Rev. A **79**, 063617 (2009)
- ➌ M. Piraud *et al.*, Phys. Rev. A **83**, 031603(R) (2011)

## ➍ Localization in bichromatic lattices (not discussed here)

- ➍ S. Aubry and G. André, Ann. Israel Phys. Soc. **3**, 133 (1980)
- ➍ G. Roati *et al.*, Nature **453**, 895 (2008)
- ➍ G. Modugno, Rep. Prog. Phys. **73**, 102401 (2010)