

## Summer School on Quantum Many-Body Physics of Ultra-Cold Atoms and Molecules

Lecture 1: Overview of tools in cold atom experiments

(11:30 am, July 12, 2012)

Lecture 2: In situ imaging: density, fluctuations, correlations and equation of state

(3:30 pm, July 12, 2012)

Lecture 3: Scale invariance, universality and quantum criticality of 2D quantum gases

(9 am, July 13, 2012)

## WORKSHOP ON QUANTUM SIMULATIONS WITH ULTRACOLD ATOMS

Talk: Exploring Universal Quantum Physics in few- and many-body atomic systems

(9:40am, July 17, 2012)

Cheng Chin, James Franck institute and Department of Physics, University of Chicago

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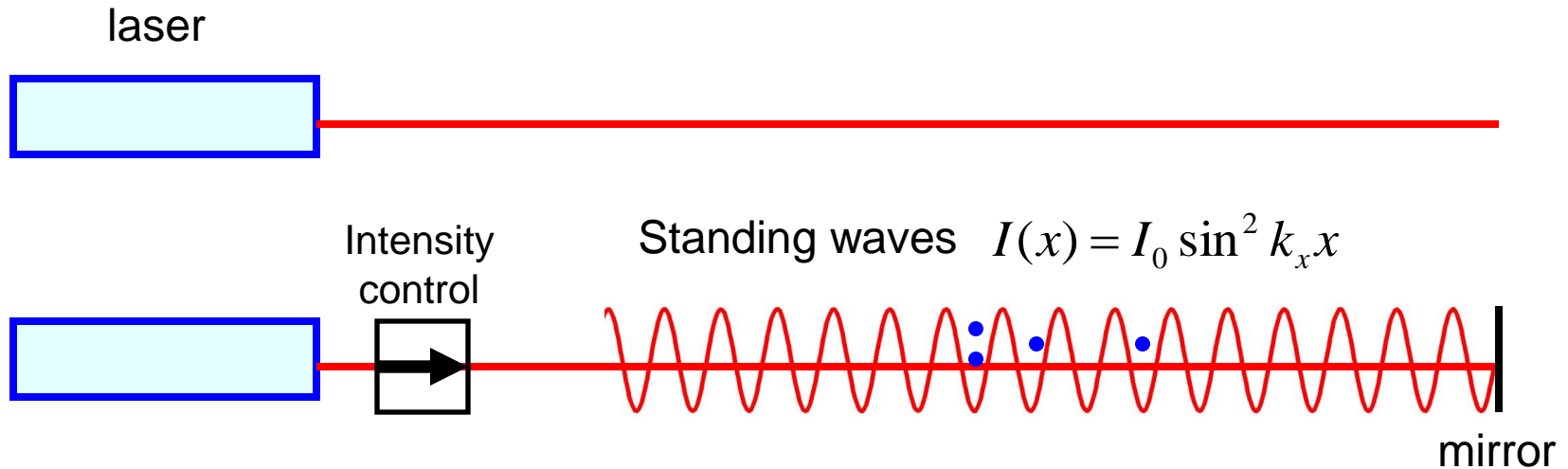
Lecture 3: Scale invariance, universality and quantum criticality of 2D quantum gases  
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## WORKSHOP ON QUANTUM SIMULATIONS WITH ULTRACOLD ATOMS

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# What is an optical lattice?



Atomic polarization

$$P = \epsilon E$$

Optical potential

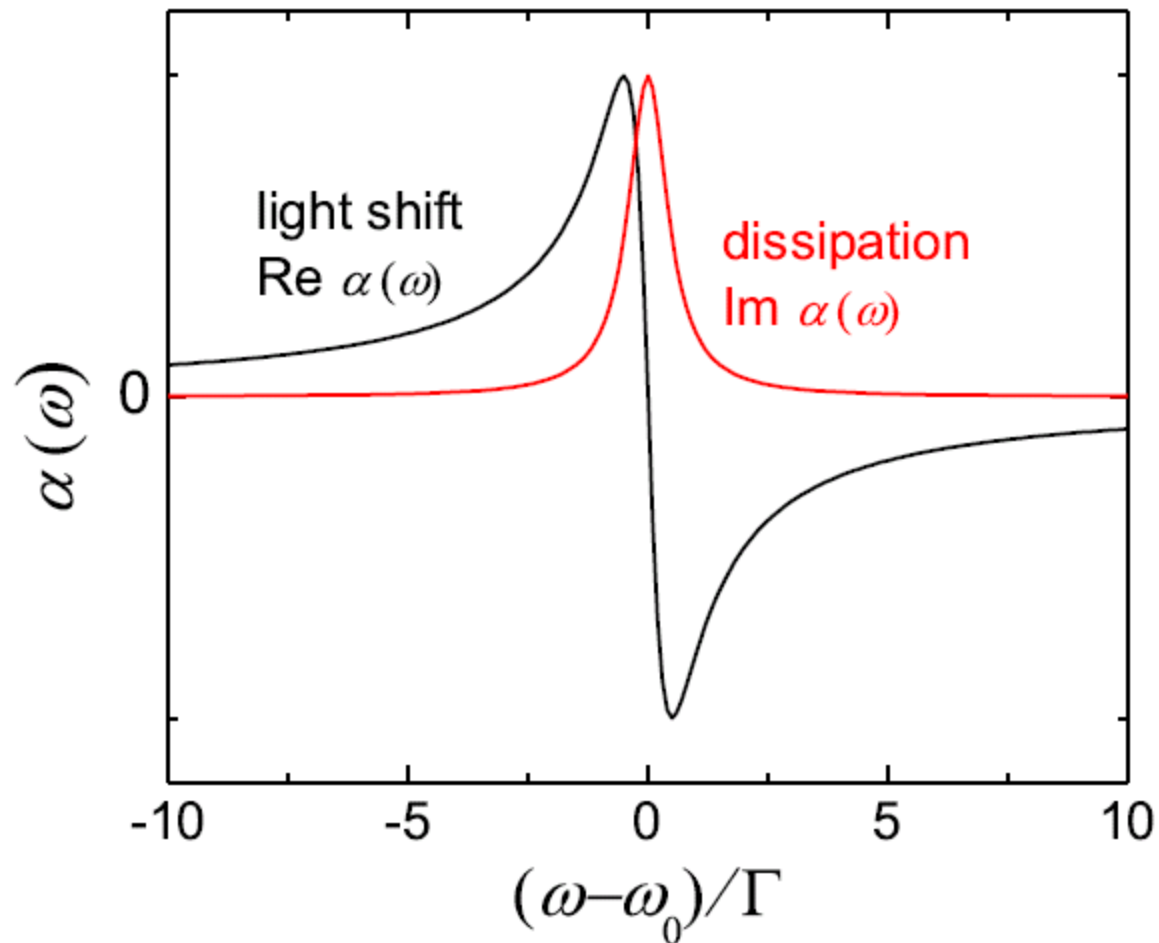
$$V = -\int P \cdot dE = -\frac{1}{2} \epsilon E^2 = -\frac{1}{2} \alpha_{AC} I$$

$\alpha_{AC}$  : AC polarizability

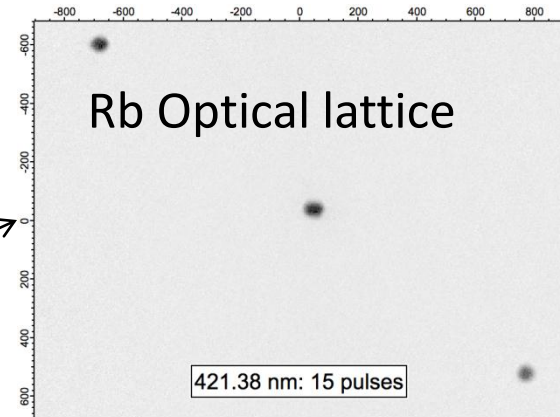
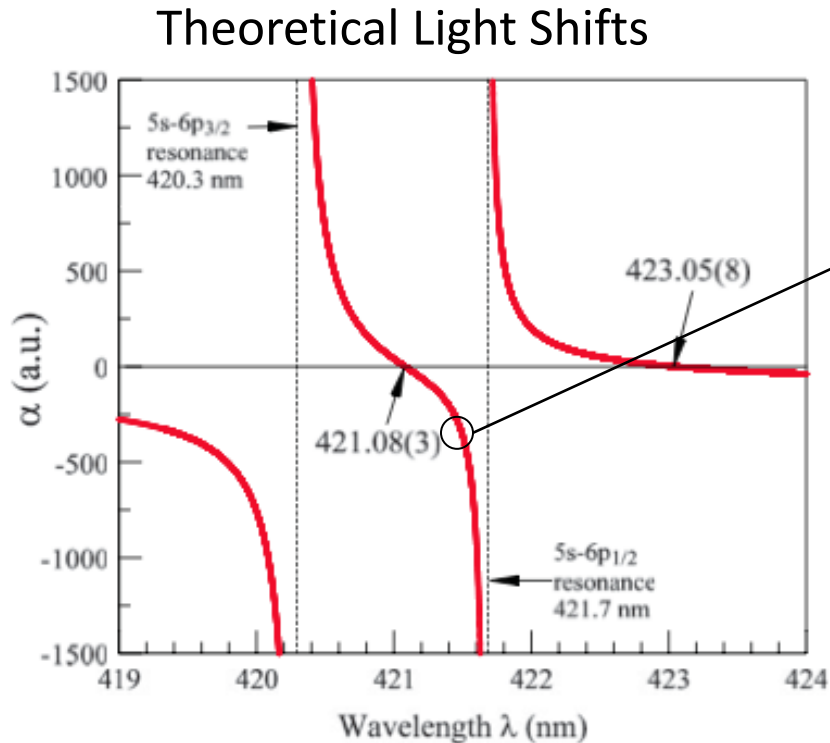
Optical lattice potential

$$V(x) = -\frac{1}{2} \alpha_{AC} I_0 \sin^2 k_x x$$

# AC polarizability



# Real atom: multiple transitions



Variations:

Magic wavelengths:

Precision measurement

Independent addressing

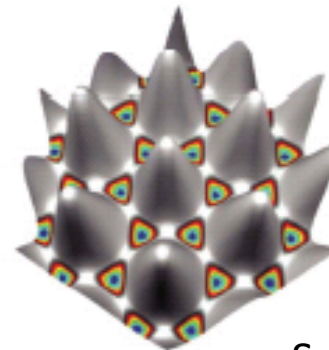
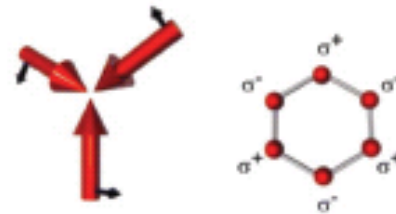
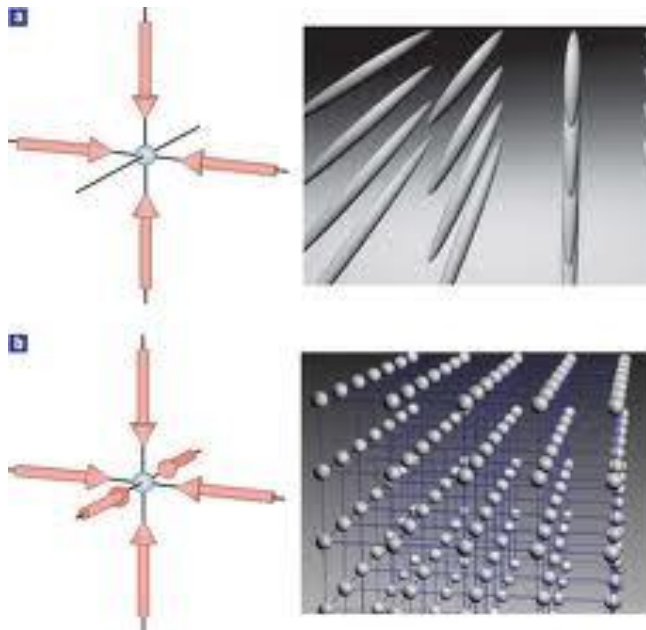
Spin-dependent lattice:

Lattice in higher dimensions

Arora et al. PRA 84 043401 (2011)

*Proto group, NIST & JQI*

# Optical lattice in higher dimensions

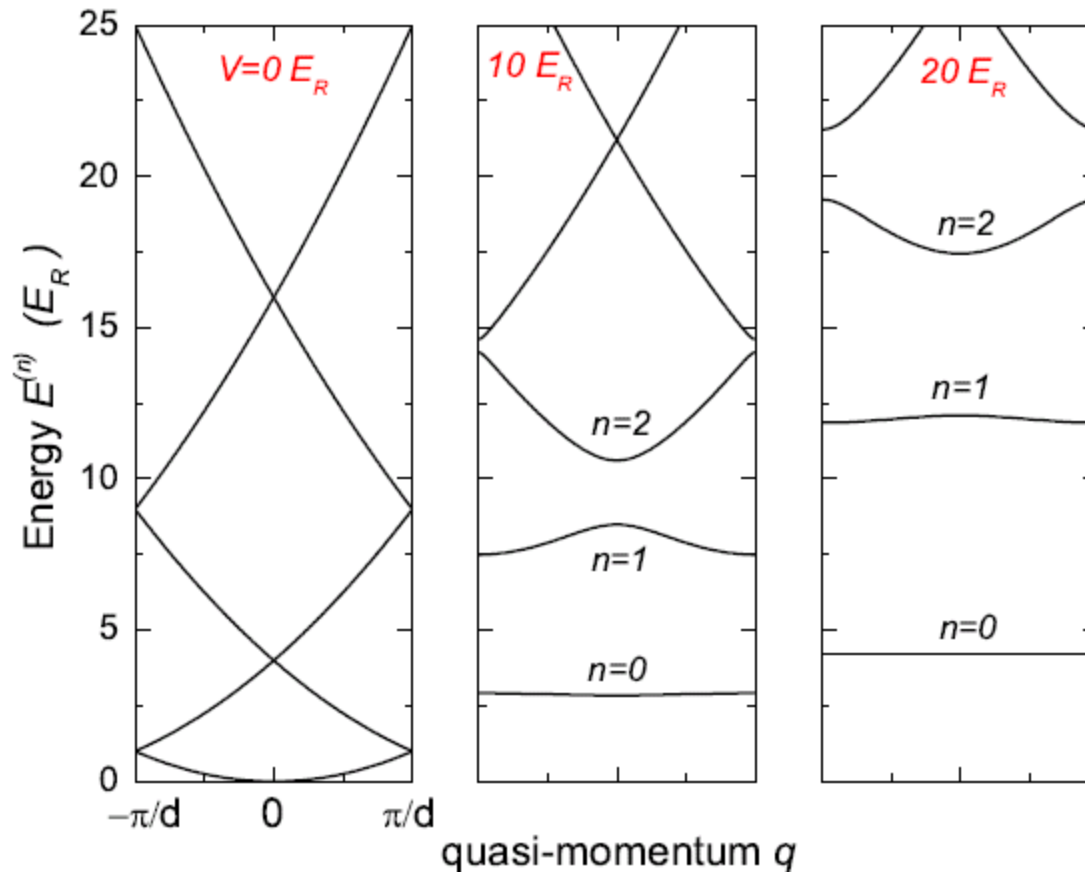


I. Bloch, Munich and MPQ

Sengstock, Hamburg

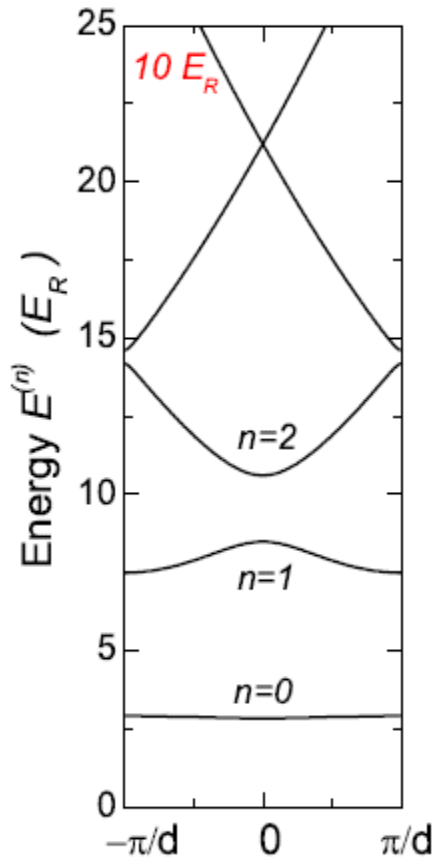
$$V(\vec{x}) = -\frac{1}{2}\alpha(\omega) \langle |Re[\sum_i \vec{E}_i e^{i(\vec{k}_i \cdot \vec{x} - \omega_i t + \phi_i)}]|^2 \rangle$$

# Band structure in optical lattice



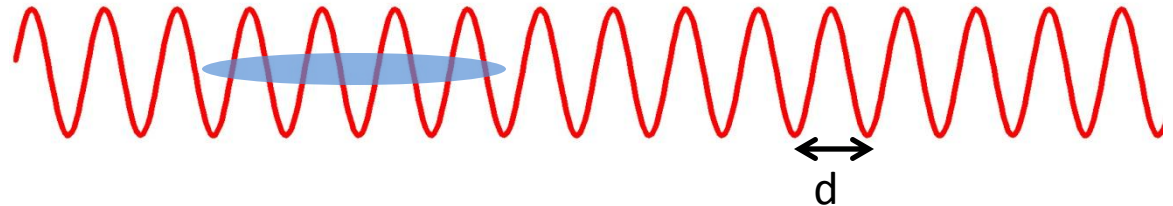
$$-\frac{\hbar^2}{2m} \partial_x^2 \psi_i(x) + V \sin^2(kx) \psi_i(x) = E_i \psi_i(x)$$

# Origin of the band structure: Bragg diffraction



$$e^{ik(x+d)} \quad e^{ik^*(x+d)} \quad \begin{matrix} \leftarrow \\ \rightarrow \end{matrix} \quad e^{ik(x+d)}$$

$$e^{ikx} \quad e^{ik^*x} \quad \begin{matrix} \leftarrow \\ \rightarrow \end{matrix} \quad e^{ikx}$$

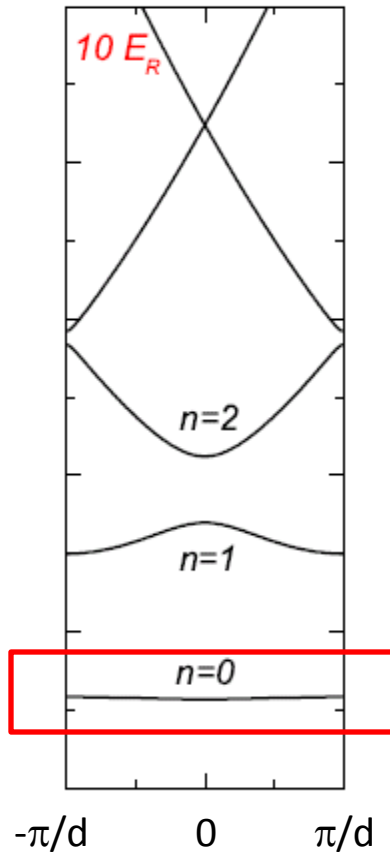


Condition for constructive interference:

$$(k+k^*)d = \pm 2\pi, \pm 4\pi, \pm 6\pi \dots$$



# Basic properties of particles in optical lattices



Eigen-states are delocalized Bloch waves  $\phi_q(x) = e^{iqx}f(x)$

$q$  is quasi-momentum

Eigen-energy  $E(q)$  is the dispersion

Localized wave packets, Wannier states, are defined as

$$w(x - x_i) = \sqrt{\frac{d}{2\pi}} \int_{-\pi/d}^{\pi/d} e^{iqx_i} \phi_q(x) dq$$

Hamiltonian in the Wannier state basis is

Quasi-momentum

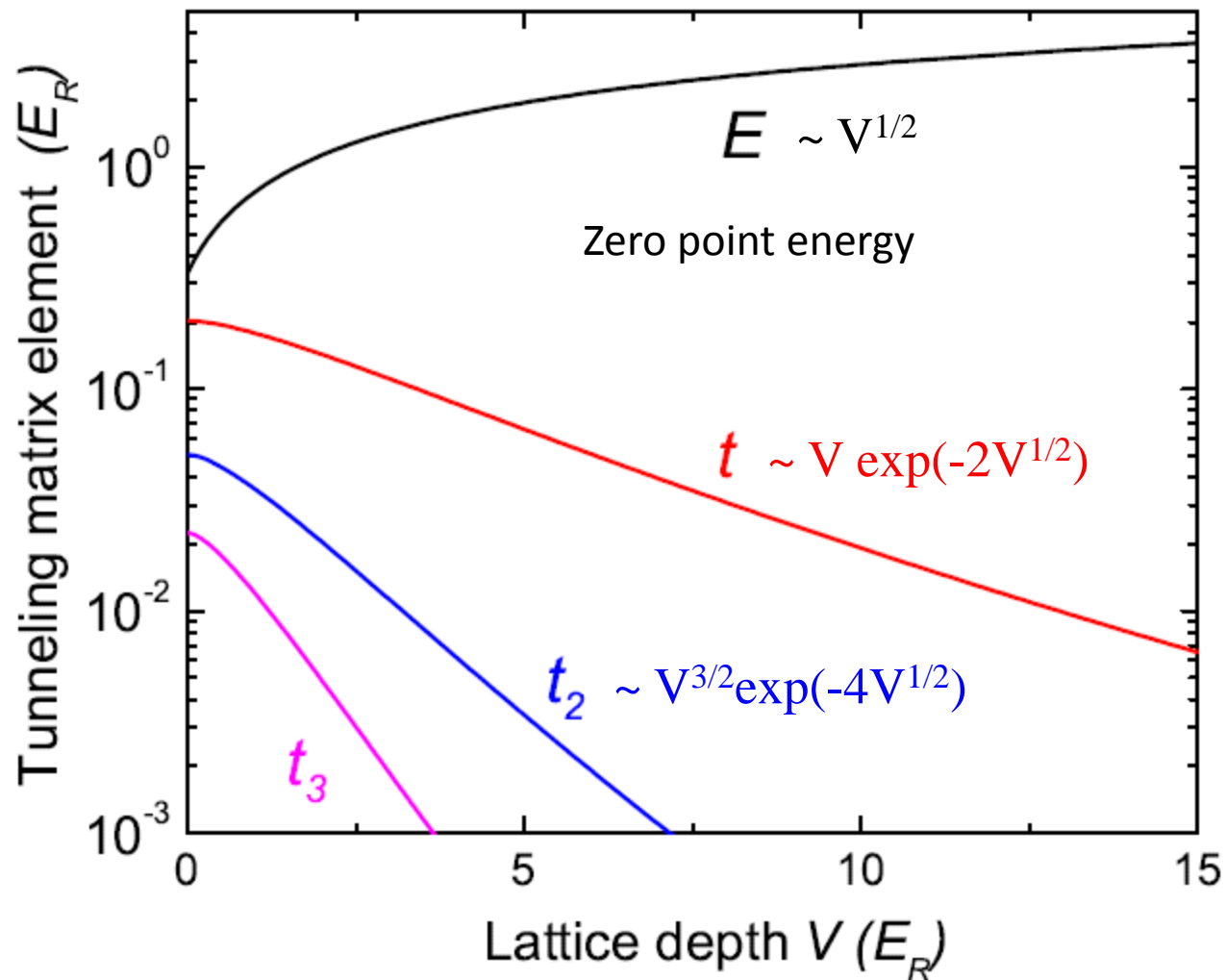
$$\langle w_i | H | w_j \rangle = \begin{pmatrix} E & -t & t_2 & -t_3 & \dots \\ -t & E & -t & t_2 & \dots \\ t_2 & -t & E & -t & \dots \\ -t_3 & t_2 & -t & E & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}$$

$t$ : nearest neighbor tunneling

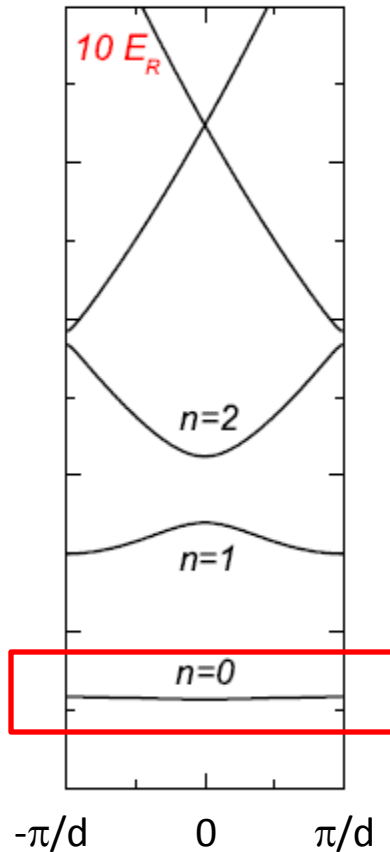
$t_2$ : next nearest neighbor tunneling

... are the Fourier expansion of  $E(q)$

# Energy scales in optical lattices



# Basic properties of particles in optical lattices



Eigen-states are delocalized Bloch waves  $\phi_q(x) = e^{iqx}f(x)$

$q$  is quasi-momentum

Eigen-energy  $E(q)$  is the dispersion

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Hamiltonian in the Wannier state basis is

$$\langle w_i | H | w_j \rangle = \begin{pmatrix} 0 & -t & t_2 & -t_3 & \dots \\ -t & 0 & -t & t_2 & \dots \\ t_2 & -t & 0 & -t & \dots \\ -t_3 & t_2 & -t & 0 & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}$$

Quasi-momentum

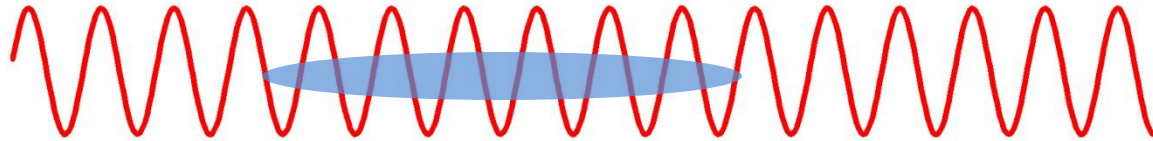
$t$ : nearest neighbor tunneling

$t_2$ : next nearest neighbor tunneling

... are the Fourier expansion of  $E(q)$

# Interaction of atoms in lattices

General wave function  $\hat{\psi} = \sum_i \hat{a}_i w(x - x_i)$



$$V = \frac{1}{2} \int \hat{\psi}^+(x_1) \hat{\psi}^+(x_2) V(x_1 - x_2) \hat{\psi}(x_1) \hat{\psi}(x_2) dx_1 dx_2$$

$$\approx \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1)$$

$$U = \frac{\hbar \omega}{\sqrt{2\pi}} \frac{a}{\ell}$$

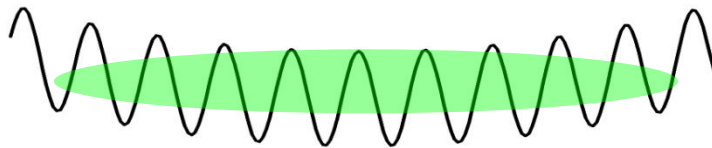
$V(x)$ : (contact) interaction  
 $n_i$ : particle number at site  $i$   
 $a$ : scattering length  
 $\omega$ : mean trap frequency  
 $\ell$ : mean oscillator length

# Bose-Hubbard model (*Fisher et al., PRB 1989, Greiner et al., Nature 2002*)

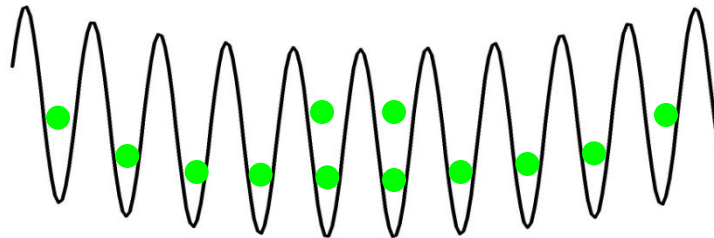
$$\hat{H} = - \underbrace{t \sum_{\langle i,j \rangle} (\hat{b}_i^\dagger \hat{b}_j + \hat{b}_j^\dagger \hat{b}_i)}_{\text{Tunneling}} + \underbrace{U \sum_i \frac{\hat{n}_i(\hat{n}_i - 1)}{2}}_{\text{Interaction}} - \underbrace{\sum_i \mu_i \hat{n}_i}_{\text{Trap potential}}$$

## Density distribution

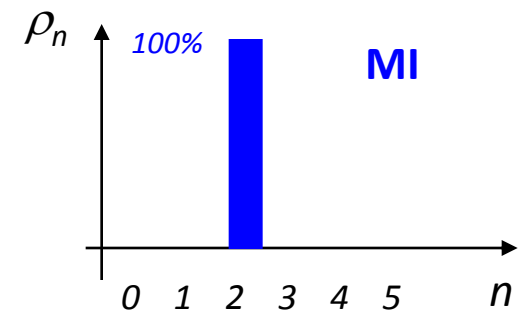
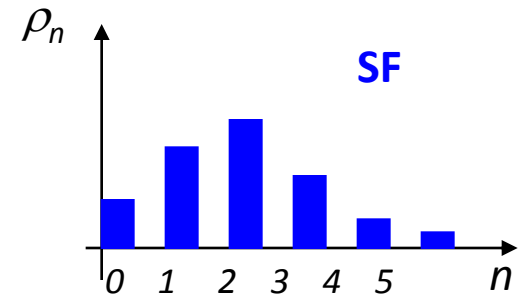
$t \gg U$   
**Superfluid**



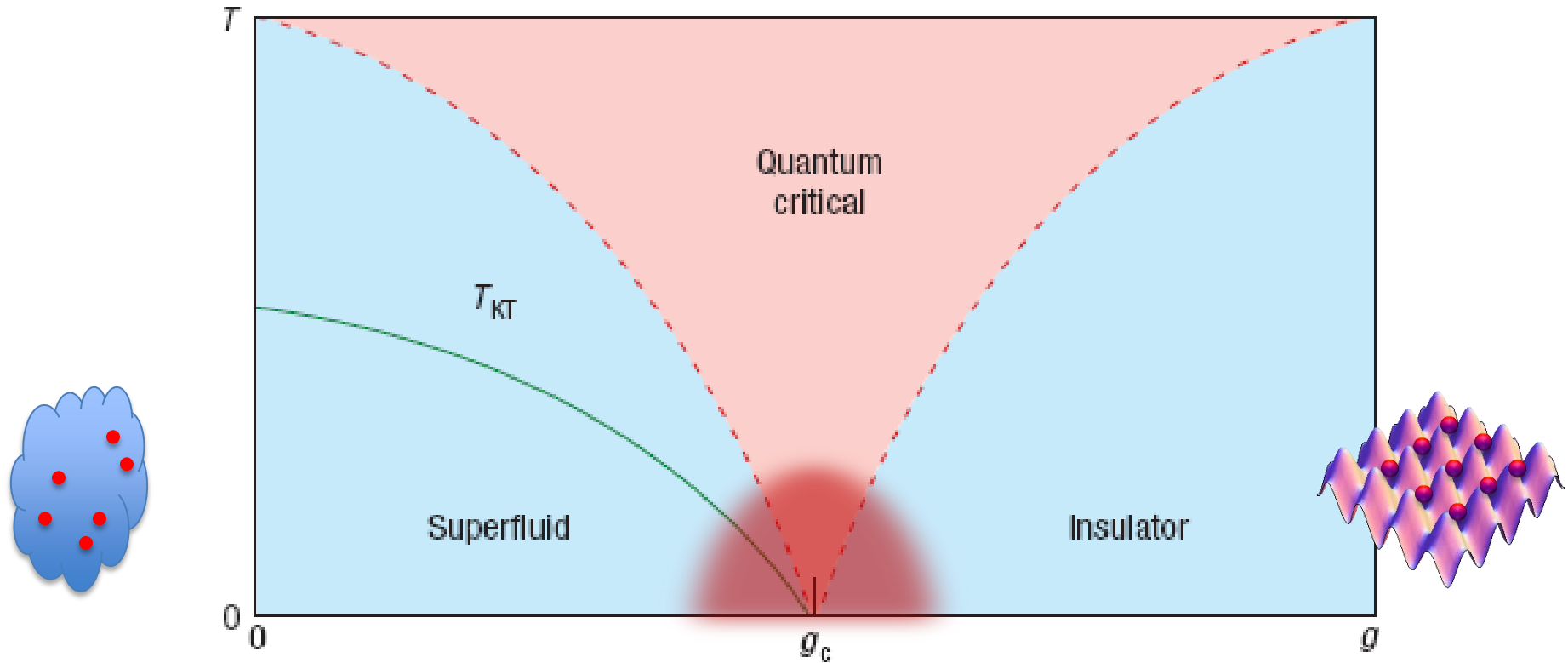
$t \ll U$   
**Mott insulator**



## Number statistics



# Quantum phase transition of bosons in 2D lattices



Ads-CFT duality: Sachdev, Nature physics (2007)

$$\frac{n - n_c}{T^{d/z+1-1/\nu z}} = h \left( \frac{\mu - \mu_c}{T^{1/\nu z}} \right)$$

Critical thermodynamics: Zhou and Ho, PRL (2010)

Hazzard and Mueller, PRA (2011)

# AdS-CFT duality

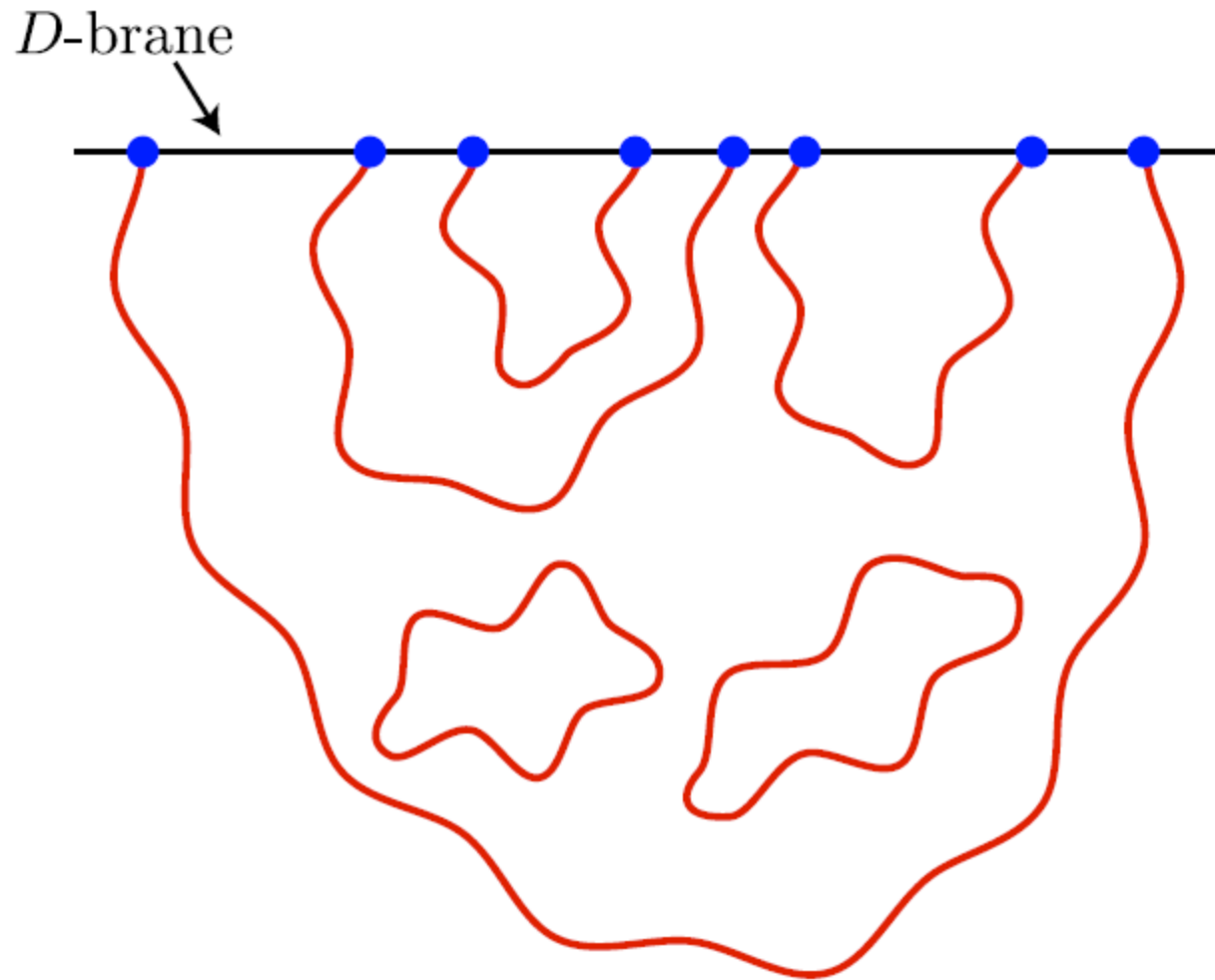
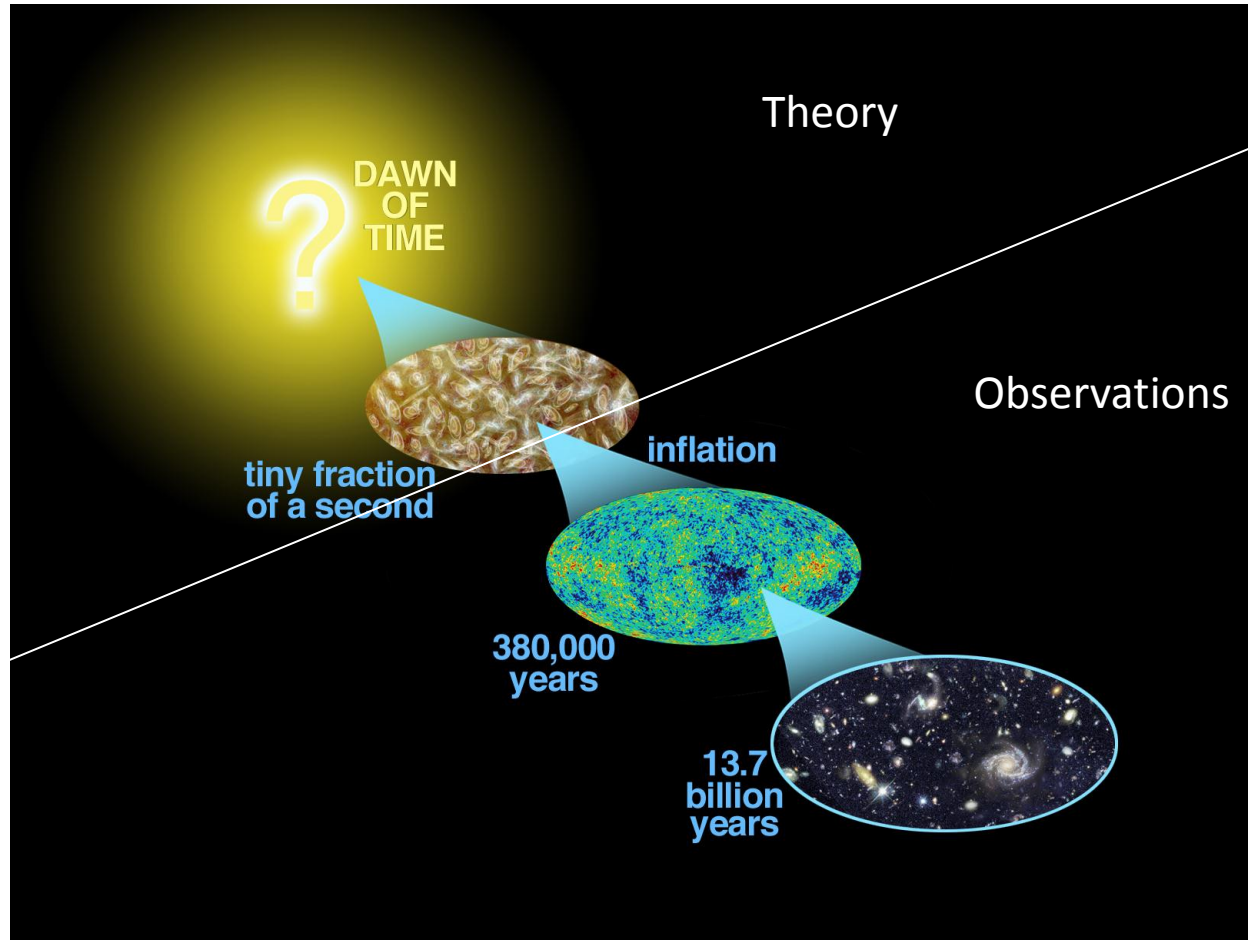


Figure: Subir Sachdev (2011)  
Juan Maldacena (1997)

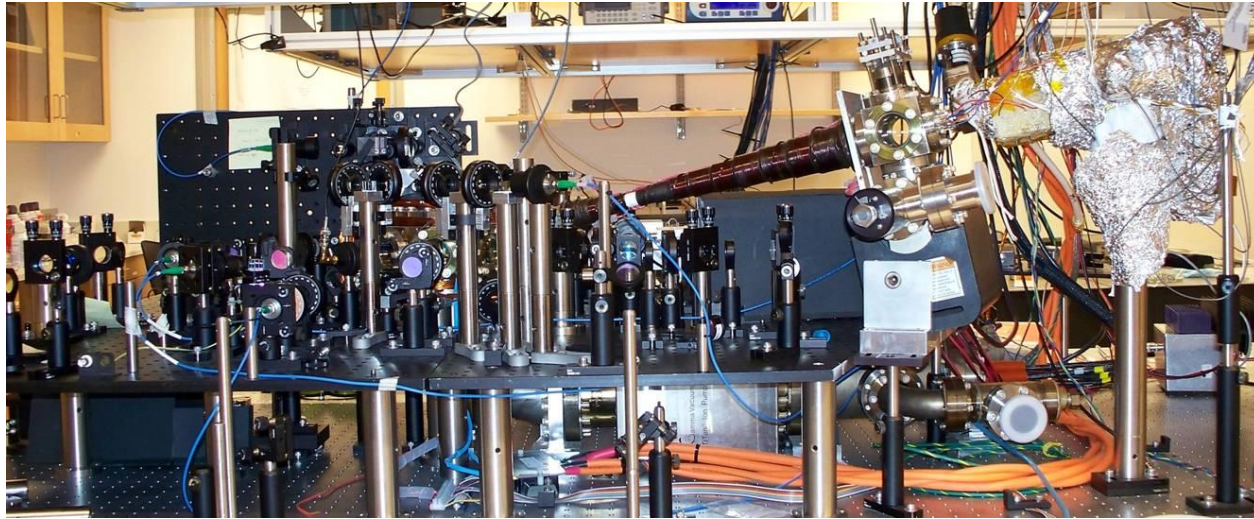
# Evolution of early universe



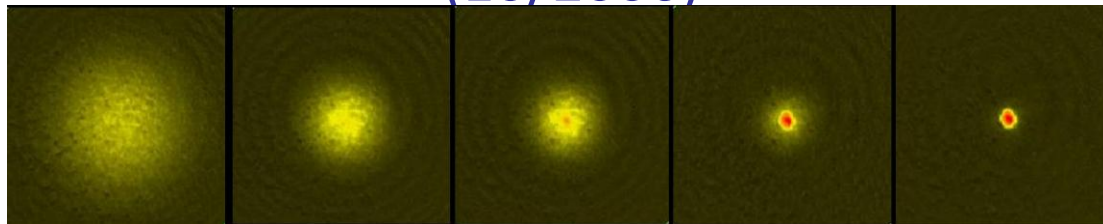
Source: NASA/WMAP Science Team



# Cesium-133 Bose-Einstein condensate apparatus



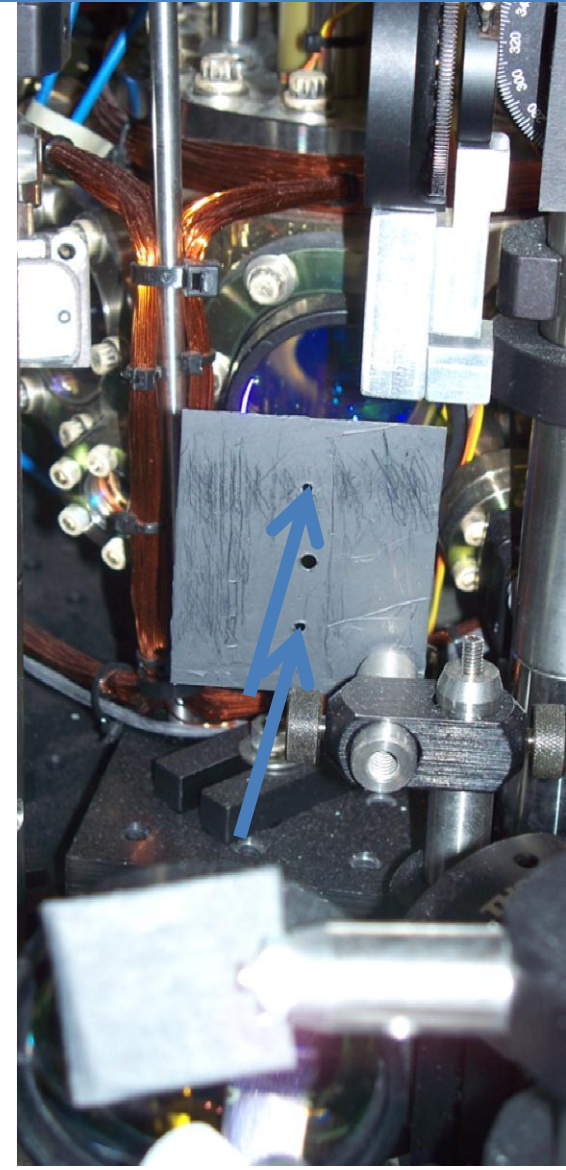
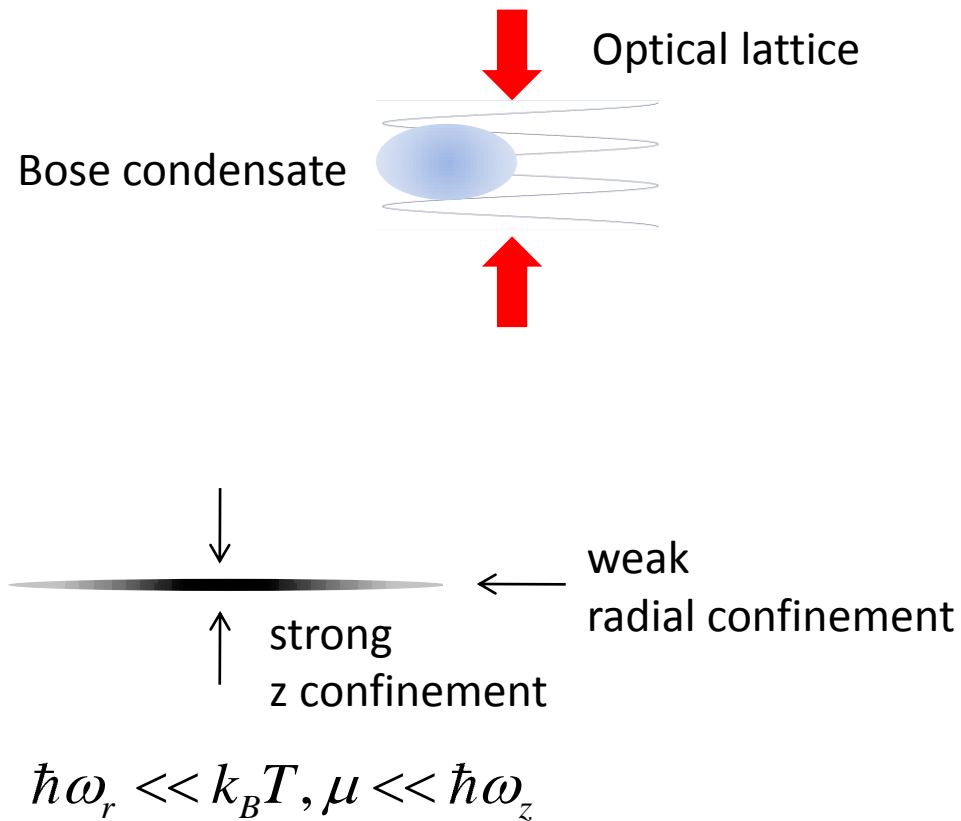
First cold atom apparatus at Univ. of Chicago  
(10/2006)



Fast evaporative cooling and Bose  
condensation (11/07/2007)

*Hung, Zhang, Gemelke, and Chin, PRA (2008)*

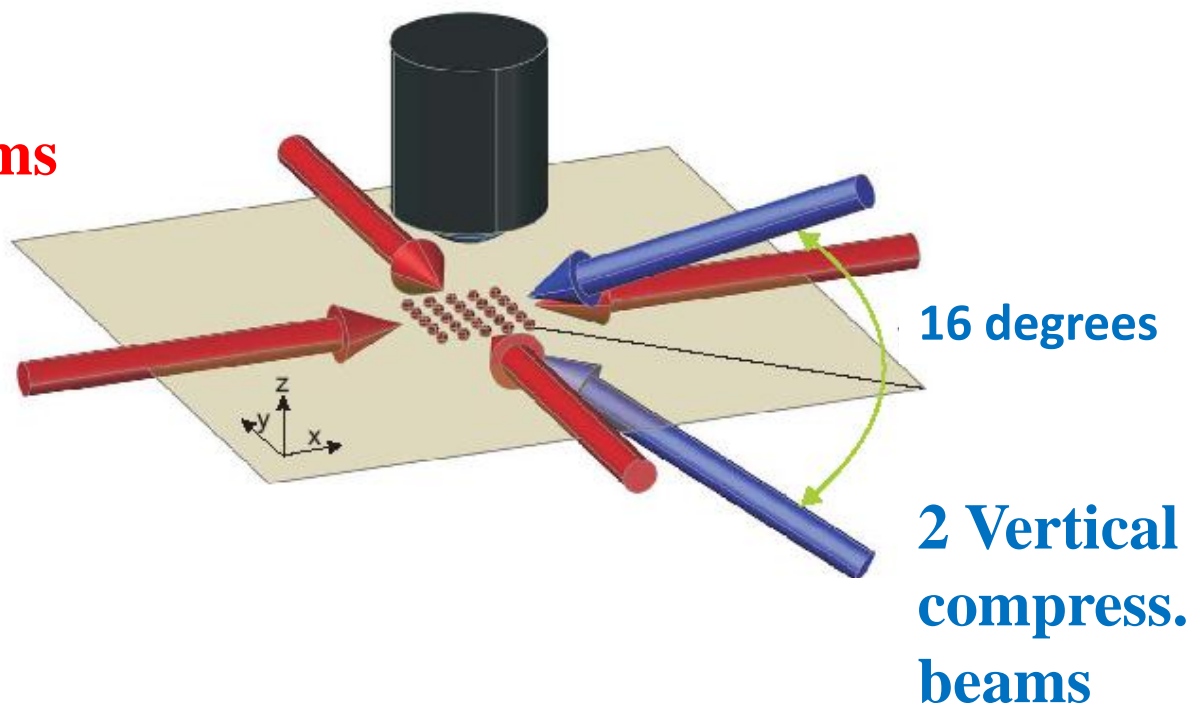
# Preparation of a 2D quantum gas



# Atoms in a 2D optical lattice

## Microscope objective

4 lattice beams

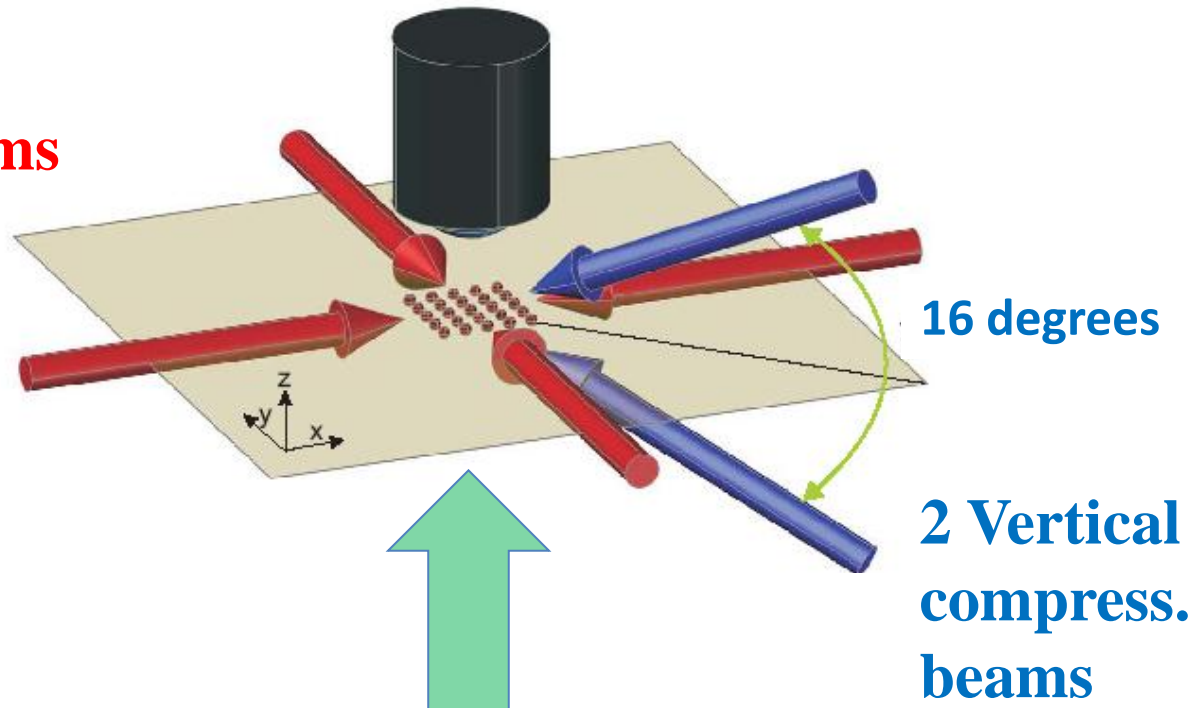


Gemelke, **Zhang**, Hung, Chin,  
*Nature* 460, 995 (2009)

# Perform in situ absorption imaging

## Microscope objective

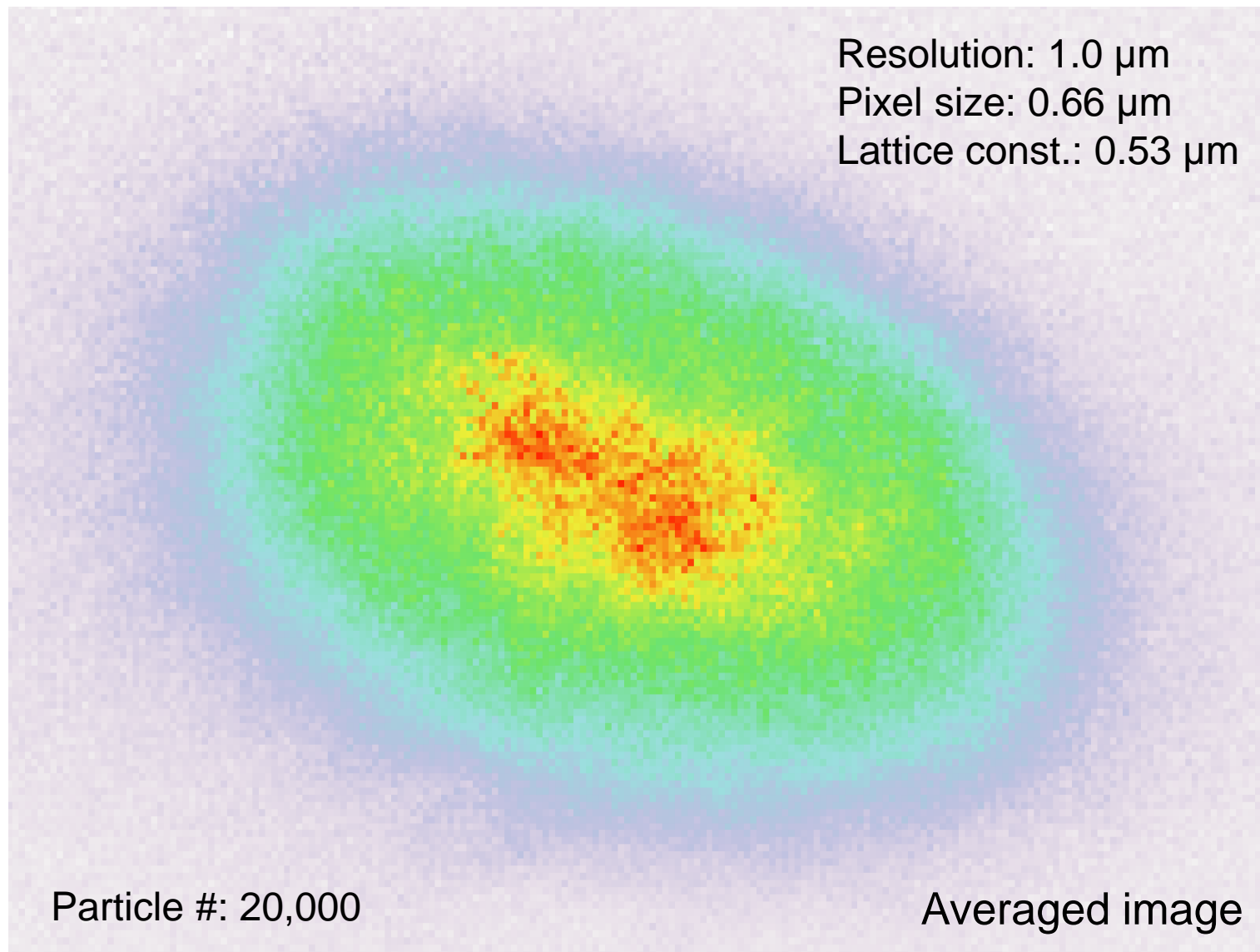
4 lattice beams



Imaging beam

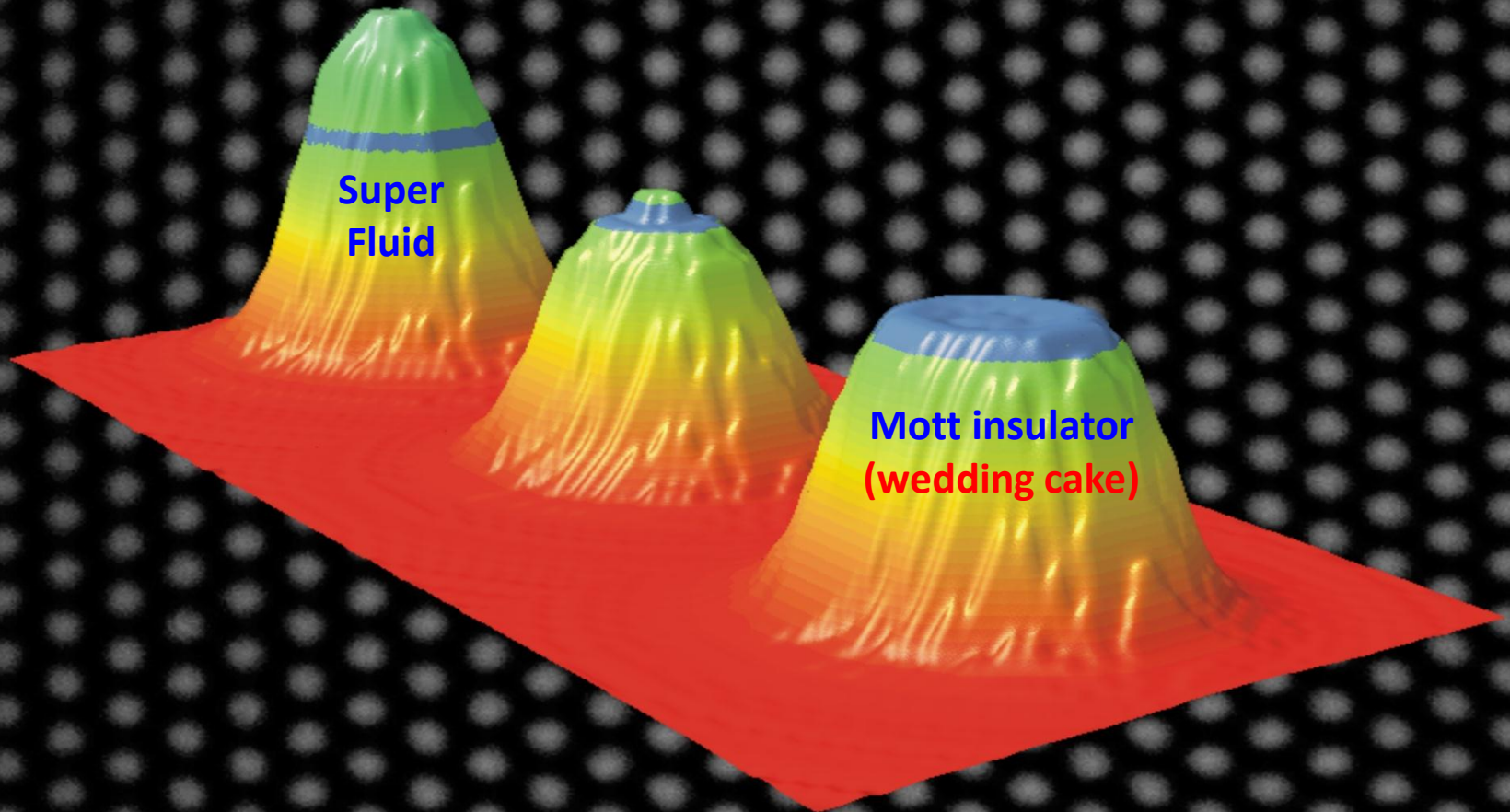
Gemelke, **Zhang**, Hung, Chin,  
*Nature* 460, 995 (2009)

## ***A closer look***

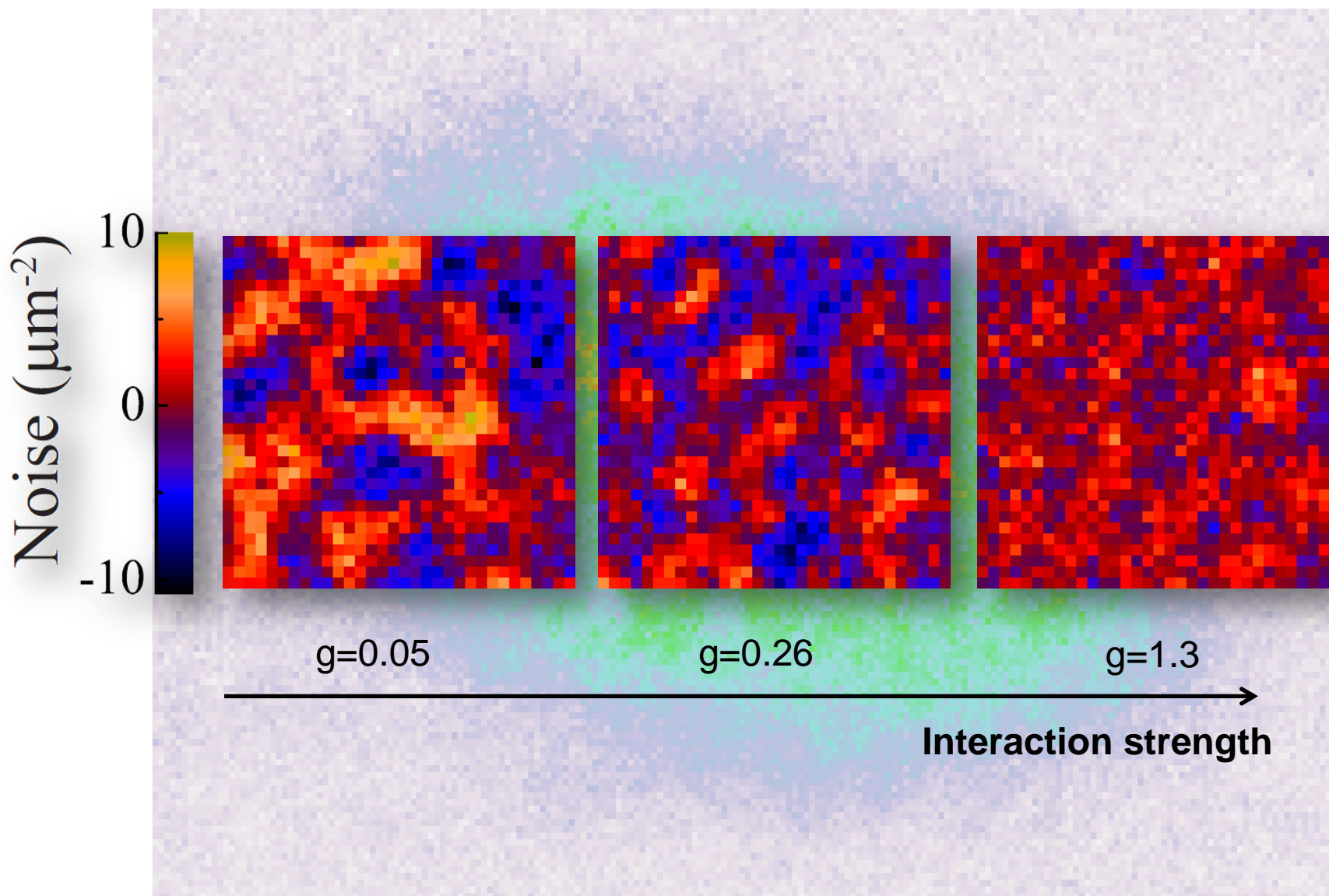




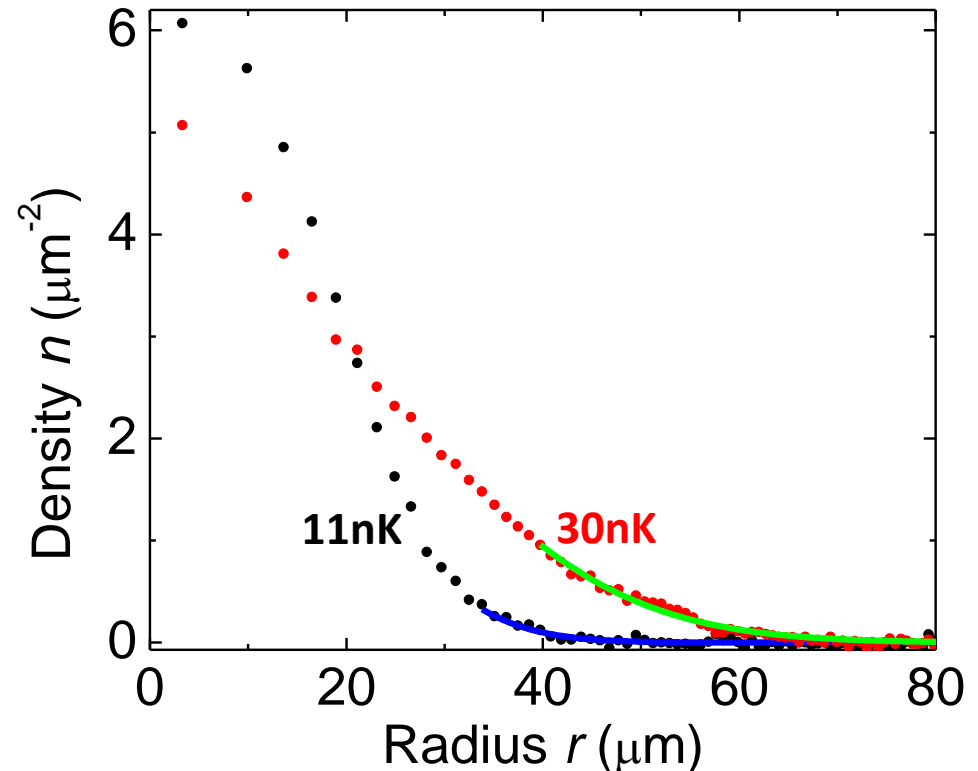
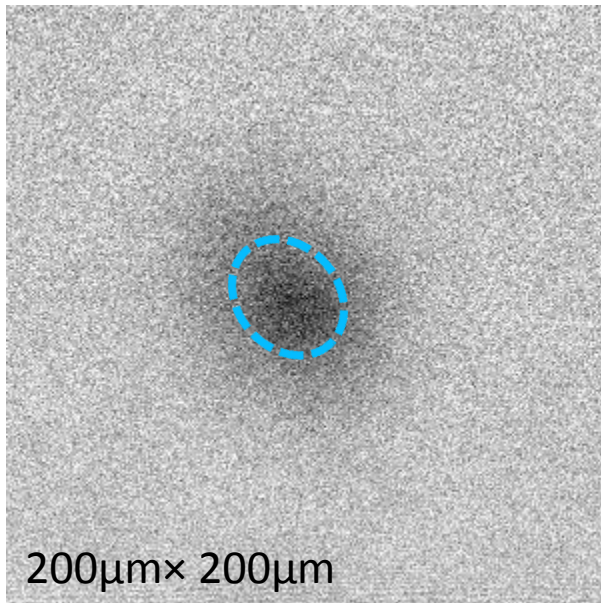
# Density profiles of atoms in 2D optical lattices



$$\delta n = n - \langle n \rangle$$



# Extract temperature $T$ and chemical potential $\mu_m$



Extracting  $T$  and  $\mu_m$  by fitting the low-density thermal tail based on a mean-field formula

$$n(r) = d^{-2} \sum_{l=1}^{\infty} \left[ I_0 \left( \frac{2lt}{k_B T} \right) \right]^2 \exp \left[ \frac{l(\mu_m - V(r) - 2U n d^2)}{k_B T} \right]$$

$d = 0.532\mu\text{m}$ ;  $t$ : tunneling;  $U$ : interaction;  $I_0$ : 0<sup>th</sup> Bessel

$$\mu = \mu_m - V(r)$$

Theory:

Ho & Zhou, Nature Physics 6, 131 ('10)

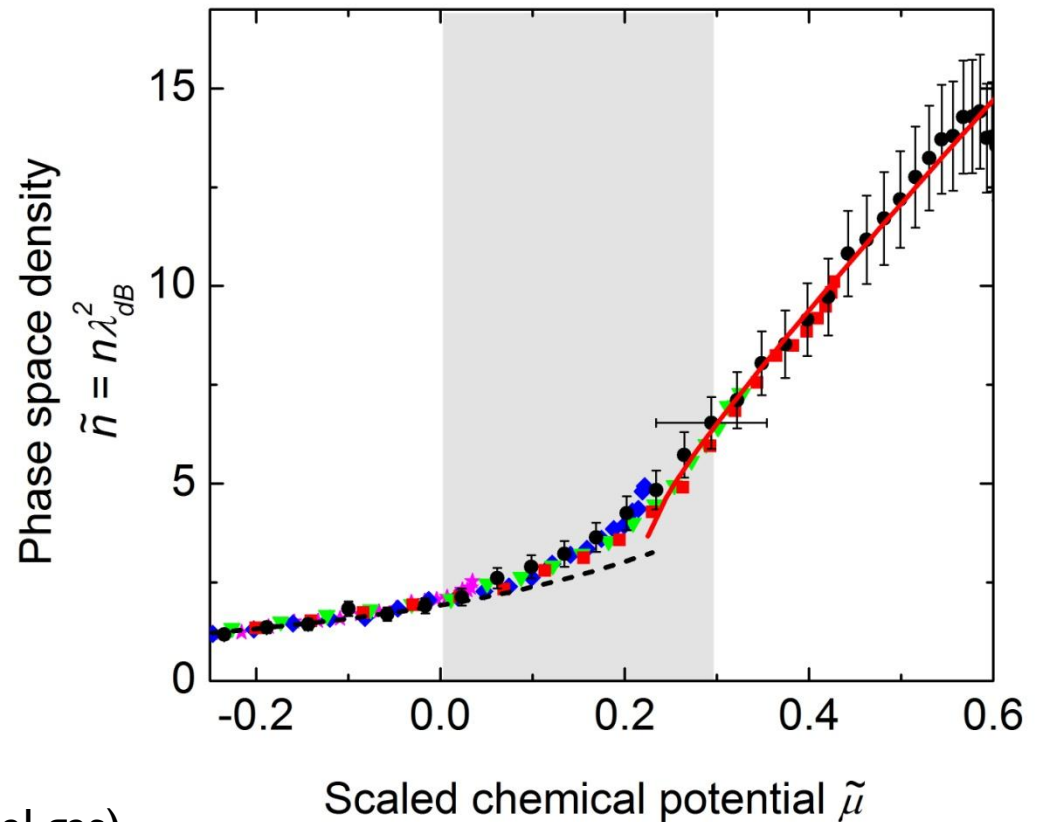
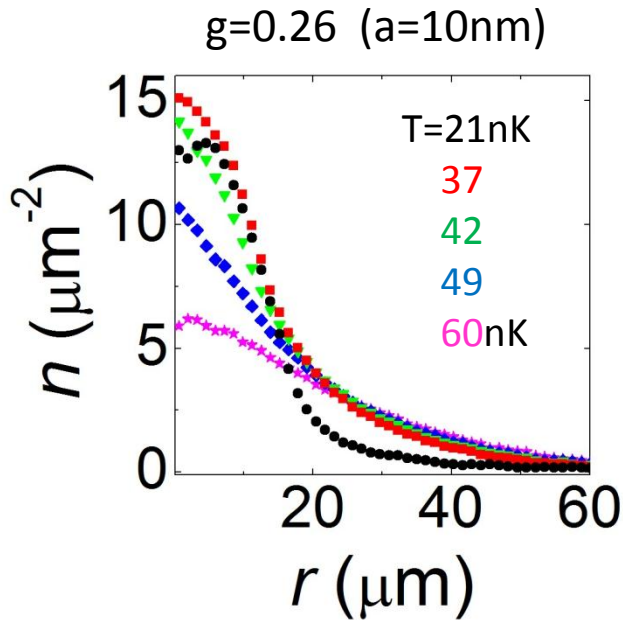
Exp:

Dalibard ('10), Zwierlein ('12), Chin ('11), ...



# Equation of state of a 2D gas And scale invariance

$$\tilde{n} = n\lambda_{dB}^2 = F\left(\frac{\mu}{kT}\right)$$



----- Mean-field (normal gas)

——— Mean-field (superfluid)

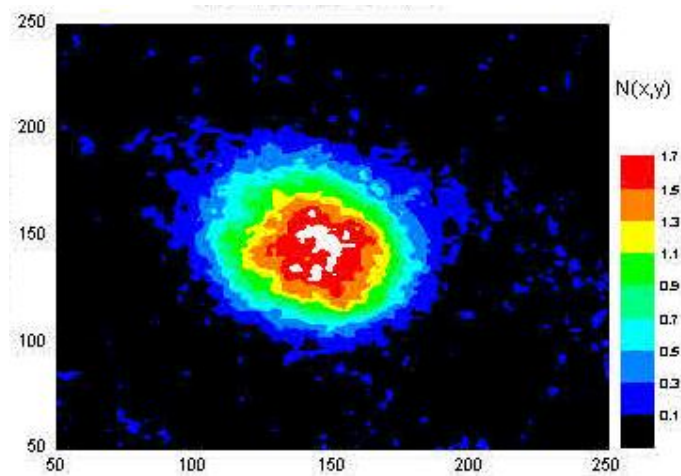
Prokof'ev and Svistunov, PRA (2002)

Chicago Experiment: Hung, Zhang, Gemelke and CC, Nature 470, 236-239 (2011)

ENS experiment: Yefsah et. Al., arXiv:1106.0188 (2011)

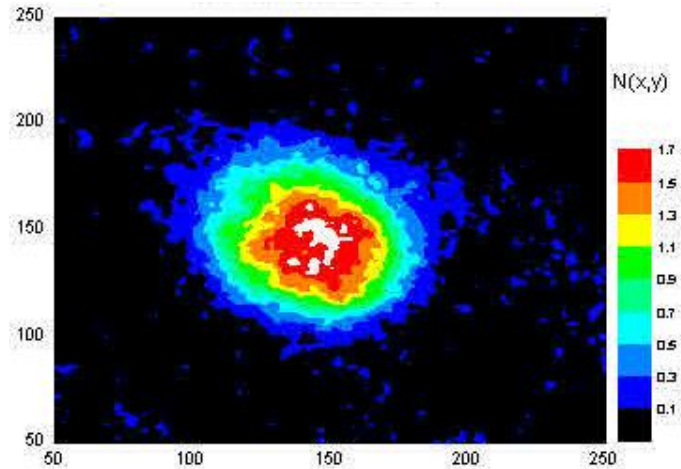
# Thermodynamics tomography (equilibrium)

Density  $n(x,y)$

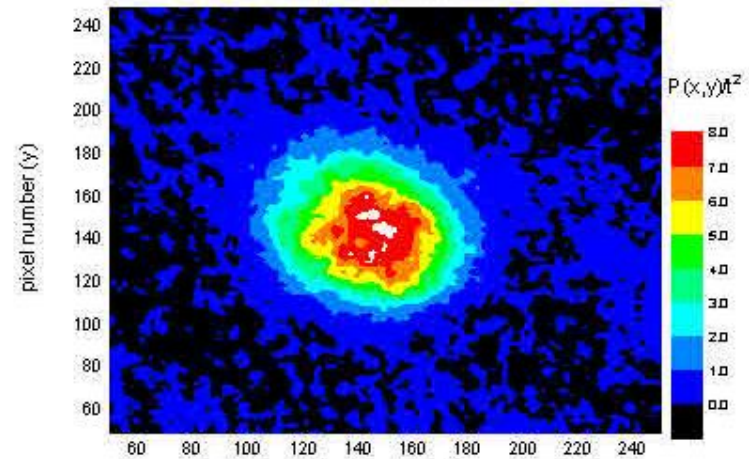


# Thermodynamics tomography

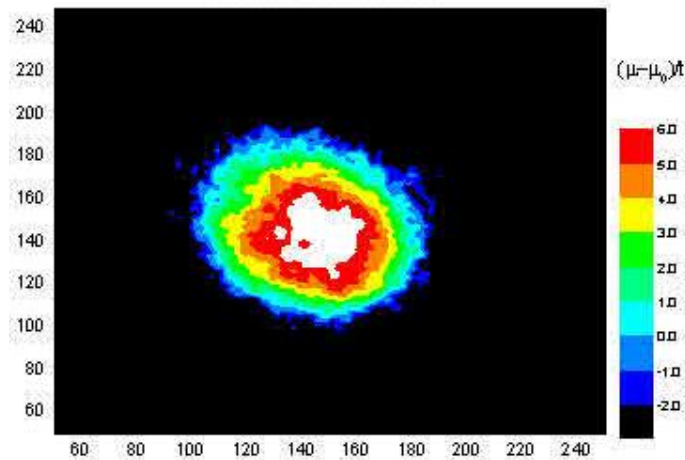
Density  $n(x,y)$



Pressure  $n(x,y)$



Chemical potential  $n(x,y)$



Entropy per particle  $S/N(x,y)$

