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Innovations in Strongly Correlated Electronic Systems: School and Workshop

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From Renormalization Group to Emergent Gravity: holographic description of quantum many-body systems

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From Renormalization Group to Emergent Gravity :

holographic description of quantum many-body systems

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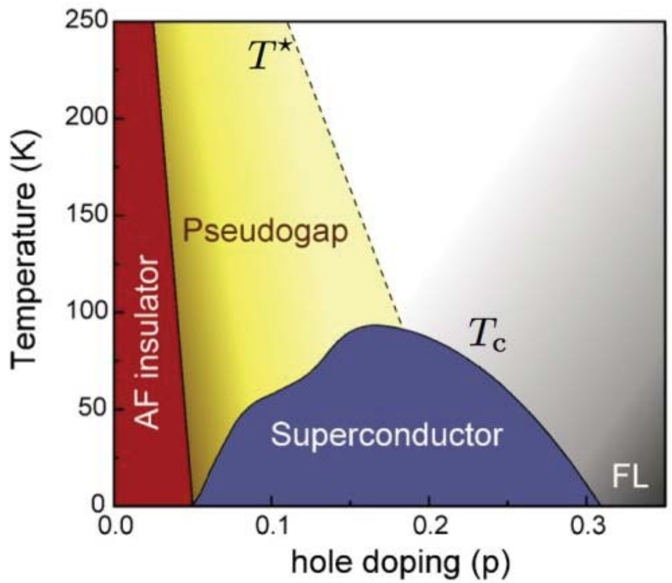
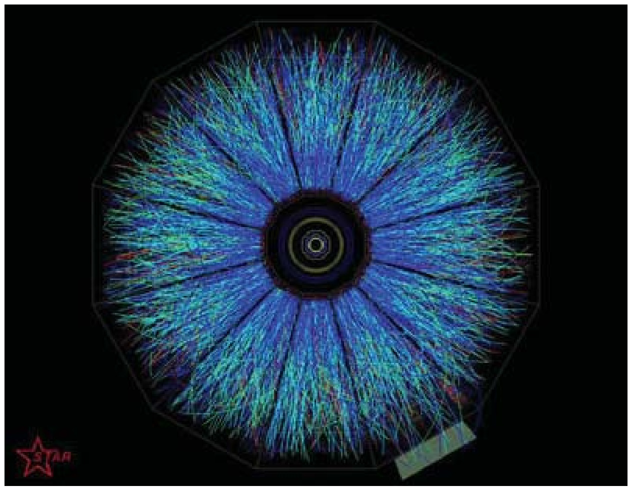
Strongly Correlated Systems

- Interactions change physical properties qualitatively
- Ubiquitous in physics
  - Condensed matter systems : high Tc, heavy fermion systems near QCP, …
  - QCD; quark-gluon plasma
- Long distance physics is described by strongly interacting QFT
- One may identify a QFT that captures the low energy physics, but it is hard to solve it
- No systematic way to incorporate strong correlations



[image from http://www.toulouse.lncmi.cnrs.fr/]



Duality : change of variable

- One theory can be mapped into another theory
  - 3D XY model = U(1) gauge theory w/ charged boson
  - Electromagnetic duality in 4D gauge theory
- Sometimes, strongly coupled theories are mapped into weakly coupled theories under duality
  - Original `particles' remain strongly coupled and have short life time, yet they are organized into long-lived (weakly coupled) collective excitations
  - Non-trivial dynamical information contained in the change of variable
  - Dual variable may carry new (sometimes fractional) quantum numbers : fractionalization
  - Dual variable may live in different space : holography

Holography : AdS/CFT correspondence

[Maldacena]

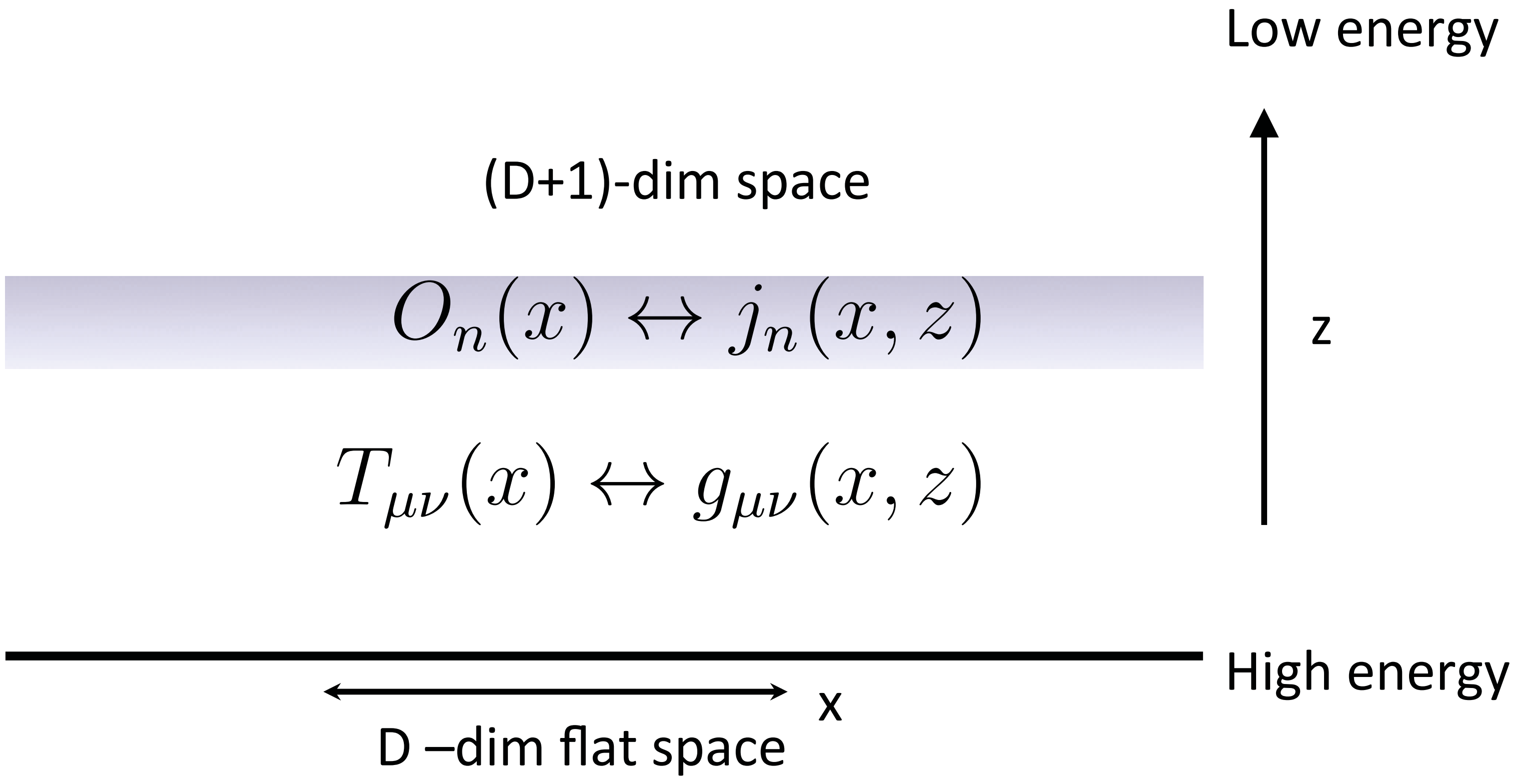
- D-dim QFT is dual to (D+1)-dim gravitational theory
  - N=4 SU(N) gauge theory in 4D = IIB superstring theory in $AdS_5 x S^5$
  - Weak coupling description for strongly coupled QFT for a large N
- Believed to be a general framework for a large class of QFT's

[Das, Jevicki; Gopakumar; Heemskerk, Penedones, Polchinski; Lee; Faulkner, Liu, Rangamani; Douglas, Mazzucato, Razamat,…]

- The correspondence has been used to reproduce many phenomena in condensed matter systems : hydrodynamics, superconductivity, non-Fermi liquid (mostly phenomenological with a few exceptions)
- No first-principle derivation for the conjecture : no systematic way to derive the dual theory for a given QFT / quantum many-body system

Dictionary in the duality

[Gubser, Klebanov, Polyakov; Witten]



$$\int D\phi(x) e^{iS_D[\phi(x)] + i\int J_n(x) O_n} = \int Dj(x,z) e^{iS_{D+1}[j(x,z)]}\Big|_{j_n(x,z=0)=J_n(x)}$$

$$\to e^{iS_{D+1}[\bar{j}(x,z)]}\Big|_{\bar{j}_n(x,z=0)=J_n(x)}$$

This talk :

- Q : Can one explicitly construct the gravitational theory from a general QFT/quantum many-body system ?

- A : Yes. Quantize beta function !

Step 0 : Partition function is functional of spacetime dependent sources

$$Z[J(x)] = \int D\phi \; e^{i\int dx \mathcal{L}}$$

$$\mathcal{L} = -J_n(x)O_n + J_{mn}(x)O_mO_n + ...$$

- O_n : set of `fundamental' operators
- Any local operator allowed by symmetry can be written as a polynomial of fundamental operators and their derivatives

e.g.

$$(\phi\partial_{\mu_1}\partial_{\mu_2}...\partial_{\mu_i}\phi) \quad \text{for Ising field,}$$

$$\sum_a(\phi_a\partial_{\mu_1}\partial_{\mu_2}...\partial_{\mu_i}\phi_a) \quad \text{for vector field}$$

$$\mathrm{tr}(\phi\partial_{\mu_1}\partial_{\mu_2}...\partial_{\mu_i}\phi\partial_{\nu_1}\partial_{\nu_2}...\partial_{\nu_j}\phi...) \quad \text{for matrix field}$$

Step 1 : remove non-fundamental operators by introducing auxiliary fields

$$Z[J(x)] = \int Dj_n^{(1)} Dp_n^{(1)} D\phi \; e^{i \int dx \mathcal{L}'}$$

$$\mathcal{L}' = j_n^{(1)}(p_n^{(1)} - O_n) - J_n p_n^{(1)} + J_{nm} p_n^{(1)} p_m^{(1)} + \ldots$$

- $J_n^{(1)}$: Lagrangian multiplier that plays the role of dynamical source that enforces the constraint $p_n^{(1)} = O_n$
- $P_n^{(1)}$: dynamical operator

Step 2 : Integrate out high energy mode

$$\phi_< : |k| < \Lambda e^{-dz}, \quad \phi_> : \Lambda e^{-dz} < |k| < \Lambda$$

$$Z[J(x)] = \int Dj_n^{(1)} Dp_n^{(1)} D\phi_< \, e^{i\int dx \mathcal{L}''}$$

$$\mathcal{L}'' = J_{nm} p_n^{(1)} p_m^{(1)} + p_n^{(1)} (j_n^{(1)} - J_n) + dz \mathcal{L}_c[j^{(1)}]$$

$$-(j_n^{(1)} + dz A_n[j^{(1)}]) O_n + dz B_{nm}[j^{(1)}] O_n O_m$$

- Casimir energy
- Quantum correction for fundamental operators
- Quadratic term of fundamental operators

Step 3 : remove non-fundamental operators by introducing a second set of auxiliary fields

$$Z[J(x)] = \int Dj_n^{(1)} Dp_n^{(1)} Dj_n^{(2)} Dp_n^{(2)} D\phi_< \, e^{i\int dx \mathcal{L}'''}$$

$$\mathcal{L}''' = J_{nm} p_n^{(1)} p_m^{(1)} + p_n^{(1)} (j_n^{(1)} - J_n) + dz \mathcal{L}_c[j_n^{(1)}]$$

$$+ j_n^{(2)} (p_n^{(2)} - O_n) - (j_n^{(1)} + dz A_n[j^{(1)}]) p_n^{(2)} + dz B_{nm}[j^{(1)}] p_n^{(2)} p_m^{(2)}$$

- Low energy fields has only fundamental operators
- Quadratic term in p_n

Step 4 : repeat 2-3 again and again

$$Z[J(x)] = \int \prod_{l=1}^{\infty} Dj_n^{(l)}(x) Dp_n^{(l)}(x)\ e^{i\int d^D x \mathcal{L}}$$

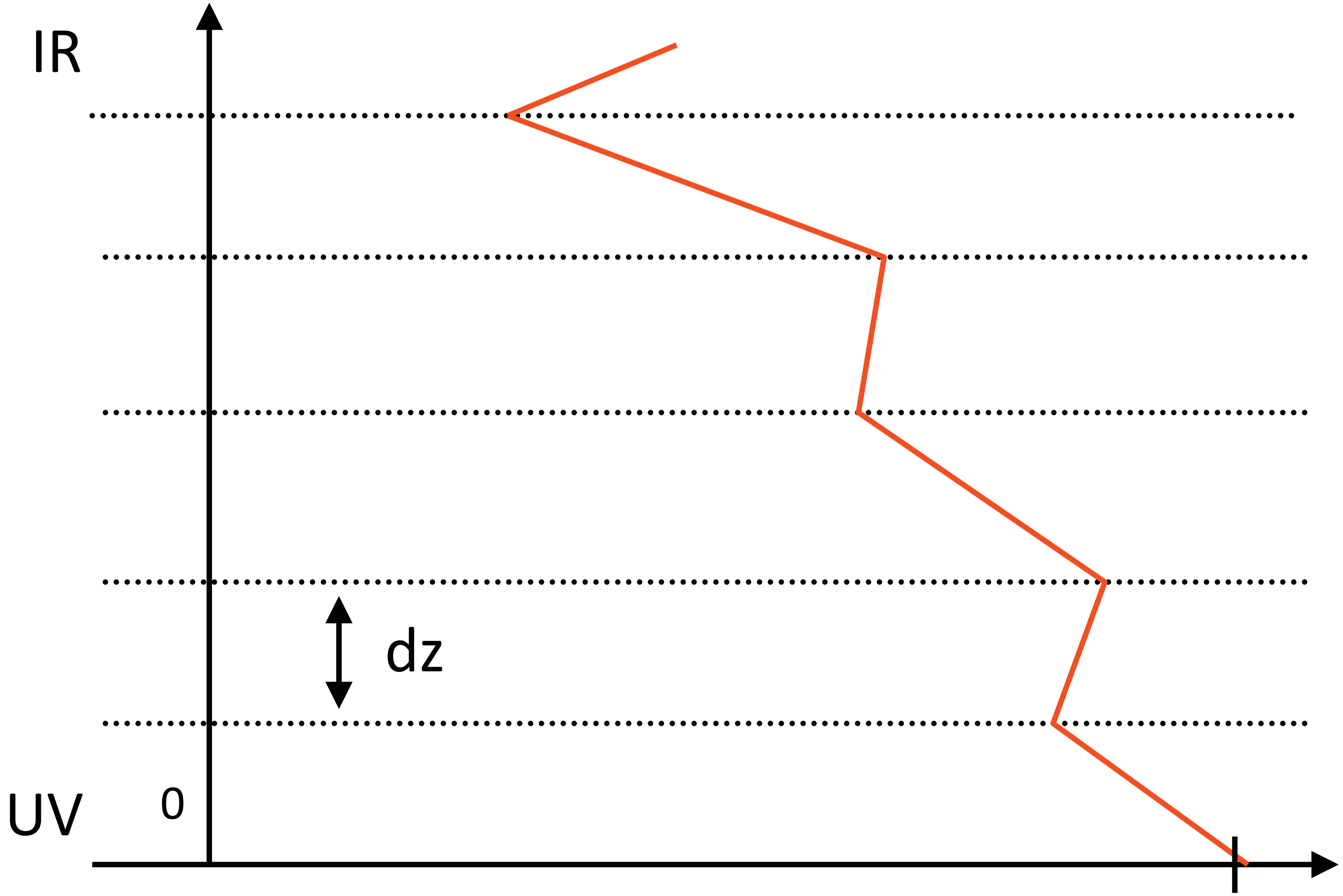
$$\mathcal{L} = J_{nm} p_n^{(1)} p_m^{(1)} + \sum_{i=1}^{\infty} \Big[p_n^{(i)} (j_n^{(i)} - j_n^{(i-1)}) + dz \mathcal{L}_c [j_n^{(i)}]$$

$$- dz A_n [j^{(i)}] p_n^{(i+1)} + dz B_{nm} [j^{(i)}] p_n^{(i+1)} p_m^{(i+1)} \Big]$$

$$j_n^{(0)}(x) = J_n(x)$$

- A set of dynamical sources and dynamical operators are introduced at each step of RG at the expense of decimating high energy mode bit by bit

Extra dimension as a length scale



$$j_n^{(l)},\ \ p_n^{(l)}$$

Boundary condition $\quad j_n^{(0)}(x) = J_n(x)$

Continuous extra dimension

$$j_n^{(l)}(x) \to j_n(x,z)$$
$$p_n^{(l)}(x) \to p_n(x,z)$$

$$z = l dz$$

$$Z[J(x)] = \int Dj_n(x,z) Dp_n(x,z)\ e^{iS_{D+1}}\Big|_{j(x,0)=J_n(x)}$$

$$S_{D+1} = \int dx J_{nm} p_n(x,0) p_m(x,0) + \int_0^\infty dz \int d^D x \Big[p_n(x,z)\partial_z j_n(x,z)$$
$$+\mathcal{L}_c[j_n(x,z)] - A_n[j(x,z)]p_n(x,z) + B_{nm}[j(x,z)]p_n(x,z)p_m(x,z)\Big]$$

- The length scale becomes an extra coordinate [Verlinde]

Key features

- An exact change of variables
- D-dimensional partition function can be written as (D+1)-dimensional functional integration for dynamical sources and operator fields
- General scheme : can be applied to any QFT
  - For general theory, the holographic description is not useful
  - When all symmetry allowed operators are large(large N), the holographic theory become classical

$$(\phi_a \partial_{\mu_1} \partial_{\mu_2} ... \partial_{\mu_i} \phi_a), \quad \mathrm{tr}(\phi \partial_{\mu_1} \partial_{\mu_2} ... \partial_{\mu_i} \phi)$$

- Holographic theory always includes gravity
  - Energy momentum tensor ➔spin-2 sources

Quantum beta function

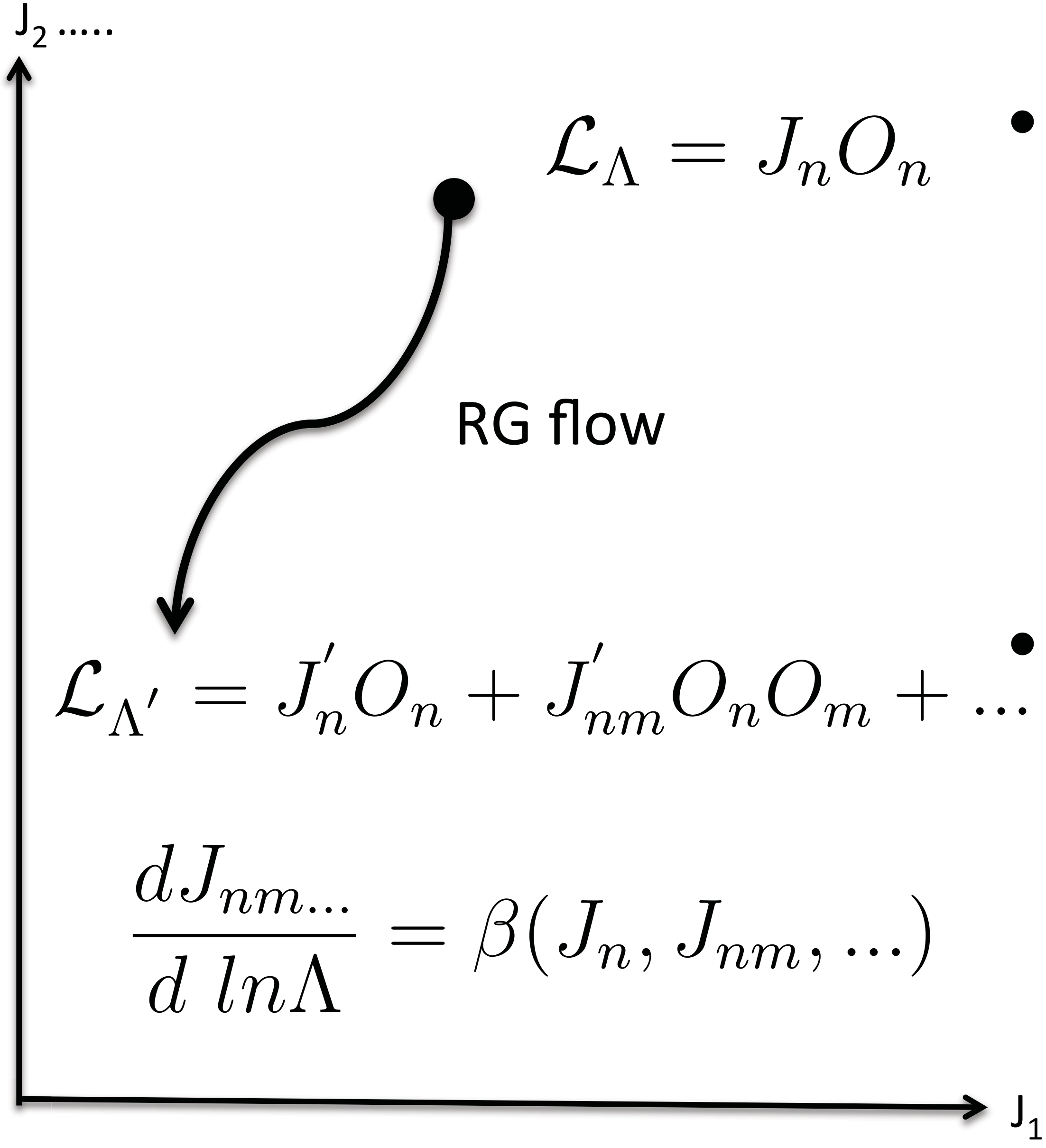
$$Z = \lim_{T\to\infty} < \Psi_f | e^{-iT\hat{H}} | \Psi_i >$$

Wavefunction for D - dimensional spacetime dependent sources

- It is useful to view the scale parameter z as `time’
- Dynamical sources and dynamical operators are conjugate to each other $[\hat{j}_l(x), \hat{p}_m(x^{'})] = \delta_{l,m}\delta(x - x^{'})$
- Partition function is written as a transition amplitude of D-dimensional quantum wavefunction of coupling constants
- The Hamiltonian generates scale transformation for dynamical couplings
- The Heisenberg equation : quantum beta function

$$\frac{d\hat{j}_n}{d\ z} = [\hat{H}, \hat{j}_n]$$

Conventional RG



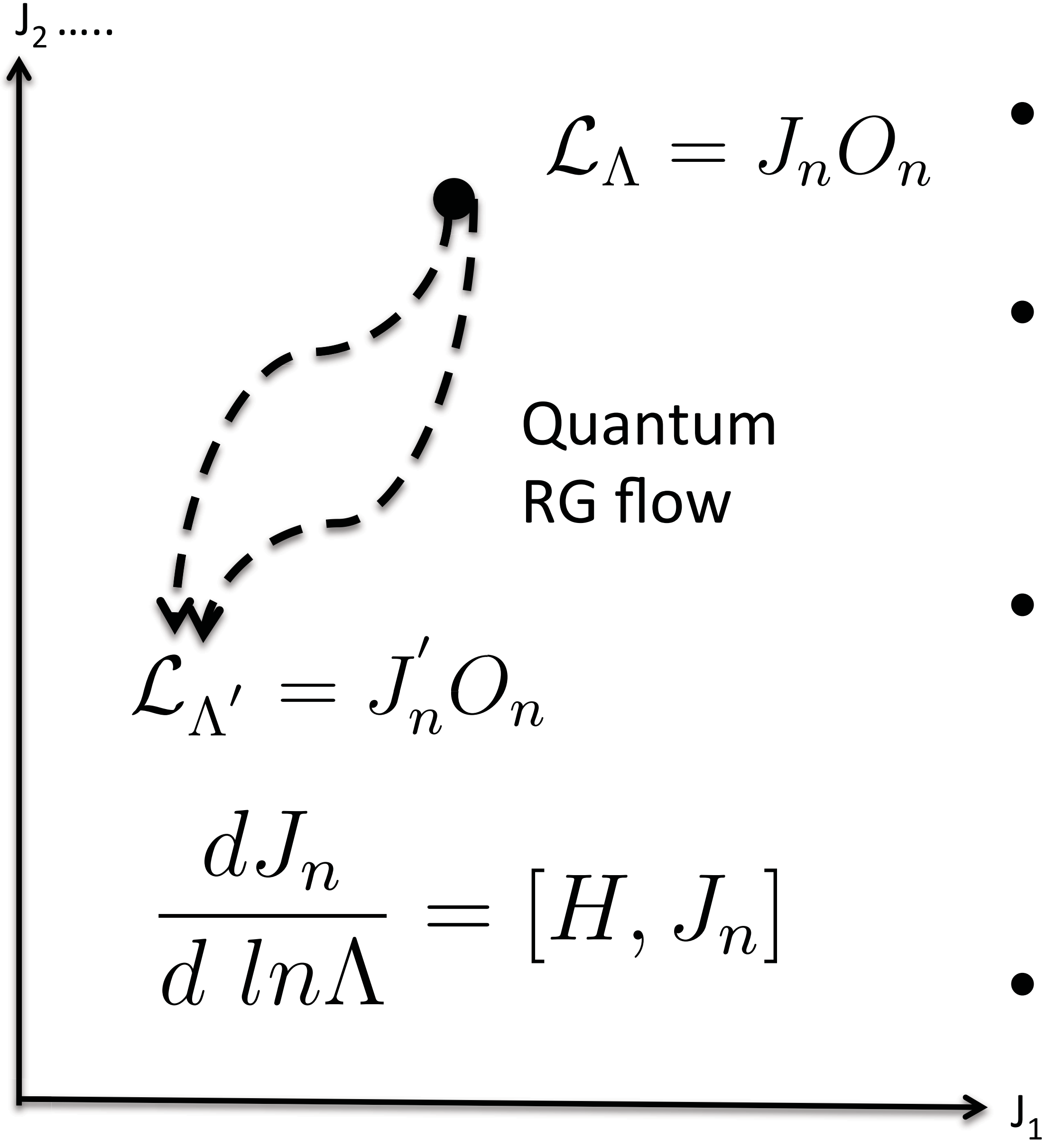
$$\mathcal{L}_{\Lambda} = J_n O_n$$

$$\mathcal{L}_{\Lambda'} = J_n^{'} O_n + J_{nm}^{'} O_n O_m + \ldots$$

$$\frac{dJ_{nm\ldots}}{d\ ln\Lambda} = \beta(J_n, J_{nm}, \ldots)$$

- Even though one starts with the fundamental operators at a given scale, non-primary operators are generated
- For a given initial condition, there is a unique trajectory = RG trajectory is classical (no fluctuations)

Holography



- Only fundamental operators appear
- The sources for fundamental operators become dynamical
- Quantum fluctuations in the RG trajectory : sources become operators!
- RG flow is governed by a quantum `Hamiltonian’

Local RG prescription

- Spacetime dependent coarse graining

$$-\phi\left(\nabla^2+\frac{\nabla^4}{\Lambda^2}+\dots\right)\phi \to -\phi\left(\nabla^2+\frac{\nabla^4}{\Lambda(x)'^2}+\dots\right)\phi$$

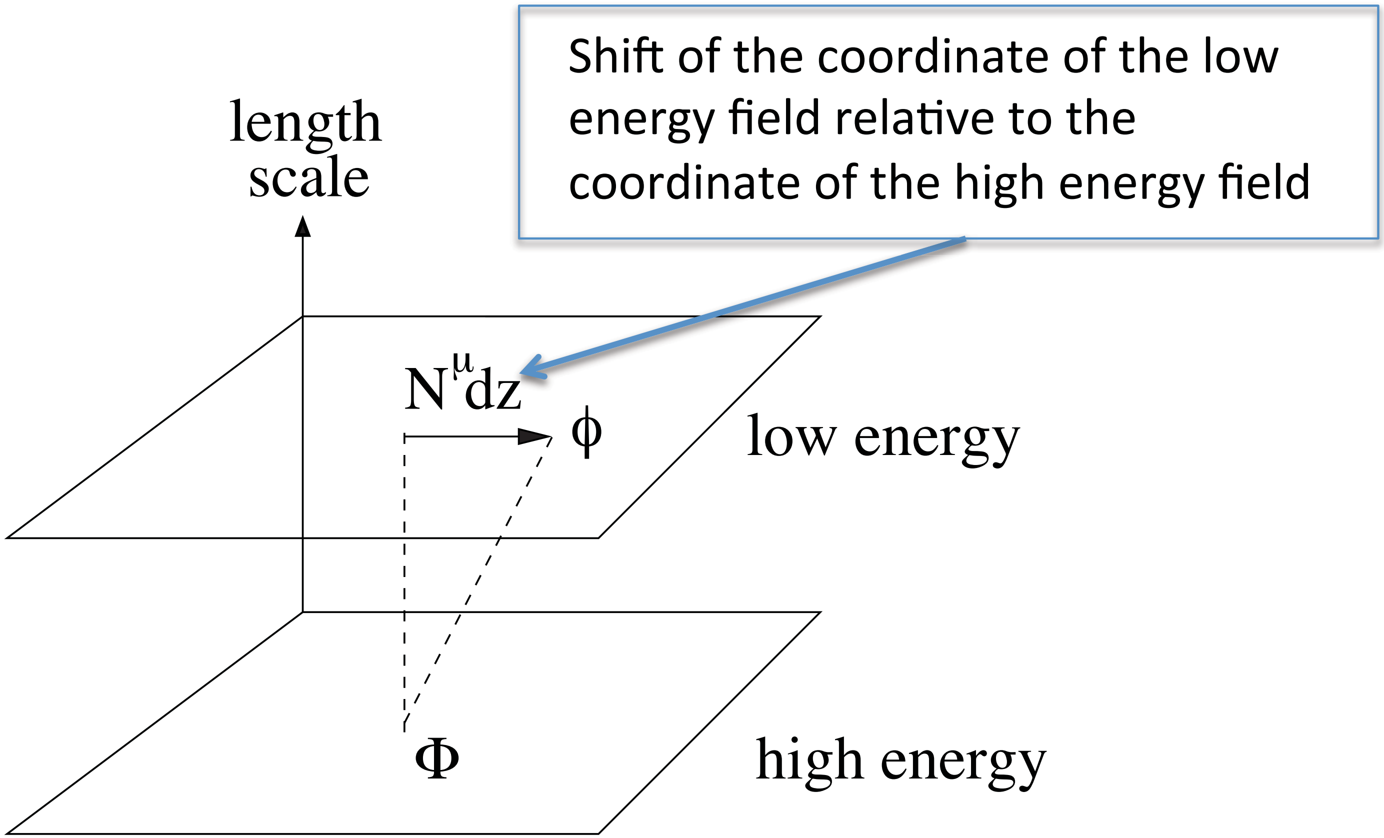
$$\Lambda(x)' = \Lambda e^{-\alpha(x)dz}$$

Speed of coarse graining

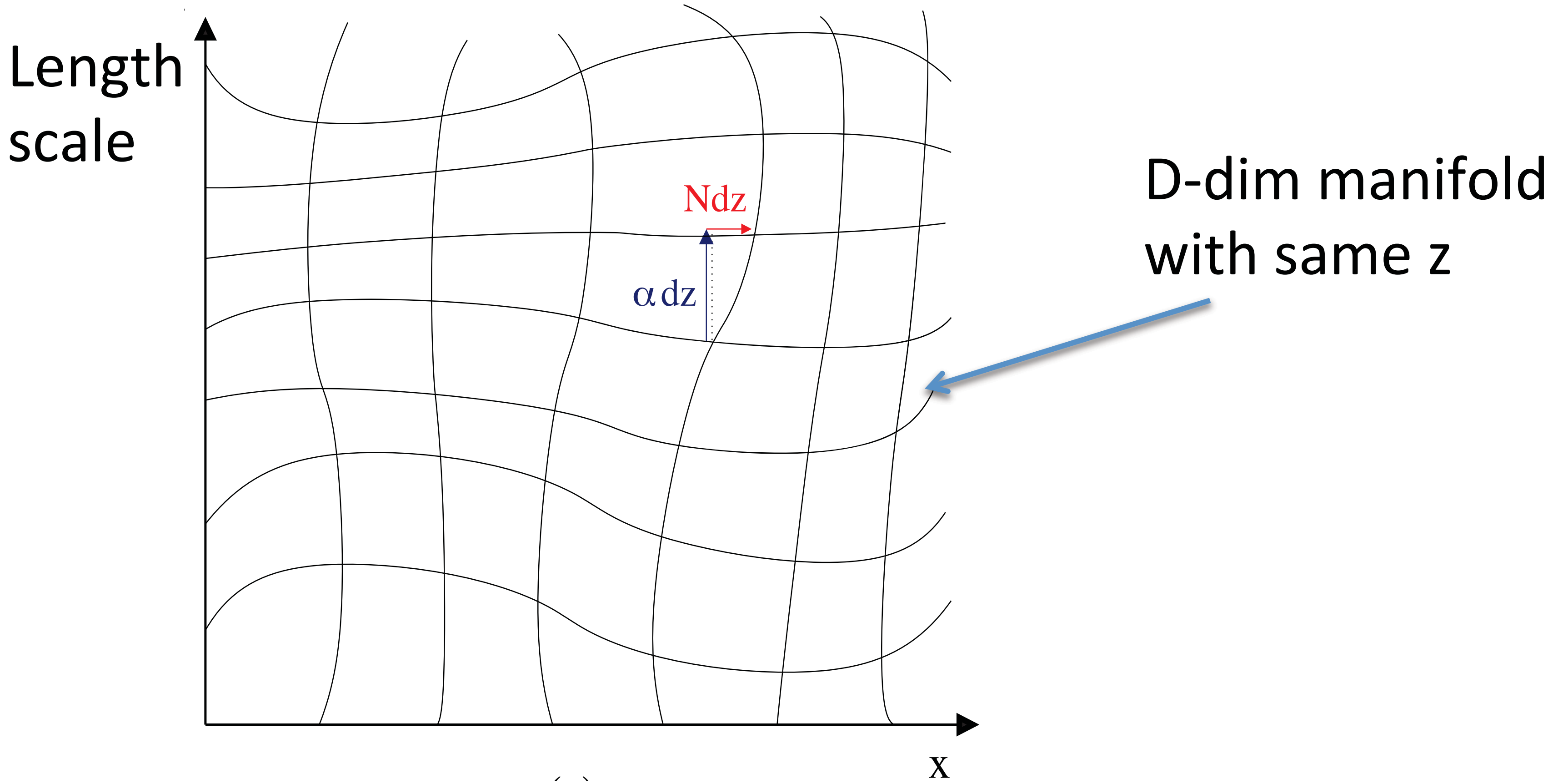
- By construction, Z is independent of $\alpha(x,z)$
- Choosing different RG scheme $\alpha(x,z)$: choosing different gauge

Shift

- One does not have to choose the coordinate of the low energy field as the coordinate of the high energy mode



Diffeomorphism = Freedom to choose different local RG schemes



D-dimensional matrix field theory

single-trace operators

$$O_{[q+1;\{\mu_j^i\}]} = \frac{1}{N}\mathrm{tr}\left[\Phi\left(\partial_{\mu_1^1}\partial_{\mu_2^1}..\partial_{\mu_{p_1}^1}\Phi\right)\left(\partial_{\mu_1^2}\partial_{\mu_2^2}..\partial_{\mu_{p_2}^2}\Phi\right)\dots\left(\partial_{\mu_1^q}\partial_{\mu_2^q}..\partial_{\mu_{p_q}^q}\Phi\right)\right]$$

ϕ : N x N tracelss symmetric real matrix field

Spacetime dependent sources

$$Z[\mathcal{J}] = \int D\Phi \ \exp\left[iN^2\int d^Dx\left(-\mathcal{J}^m O_m + V[O_m;\mathcal{J}^{\{m_i\},\{\nu_j^i\}}]\right)\right]$$

multi-trace deformation

$$V[O_m;\mathcal{J}^{\{m_i\},\{\nu_j^i\}}] = \sum_{q=1}^{\infty}\mathcal{J}^{\{m_i\},\{\nu_j^i\}}O_{m_1}\left(\partial_{\nu_1^1}..\partial_{\nu_{p_1}^1}O_{m_2}\right)\left(\partial_{\nu_1^2}..\partial_{\nu_{p_2}^2}O_{m_3}\right)\dots\left(\partial_{\nu_1^q}..\partial_{\nu_{p_q}^q}O_{m_{q+1}}\right)$$

(D+1)-dimensional gravity

$$Z[\mathcal{J}] = \int \mathcal{D}J(x,z)\mathcal{D}\pi(x,z) \; e^{i\left(S_{UV}[\pi(x,0)]+S[J(x,z),\pi(x,z)]+S_{IR}[J(x,\infty)]\right)}\Big|_{J(x,0)=\mathcal{J}(x)}$$

Bulk action : $S = N^2 \int d^D x dz \left[(\partial_z J^n)\pi_n - \alpha(x,z)\mathcal{H} - N^\mu(x,z)\mathcal{H}_\mu \right]$

Hamiltonian constraint :

$$\mathcal{H} = \tilde{A}^{\mu\nu}[J(x)]\pi_{[2,\mu\nu]} - \frac{\tilde{B}^{\mu\nu\lambda\sigma}[J(x)]}{\sqrt{|G|}}\pi_{[2,\mu\nu]}\pi_{[2,\lambda\sigma]} - \sqrt{|G|}\left\{ C_0[J(x)] + C_1[J(x)]\mathcal{R} \right\} + \ldots,$$

Momentum constraint :

$$\mathcal{H}_\mu = -2\nabla^\nu \pi_{[2,\mu\nu]} - \sum_{[q,\{\mu_j^i\}]\neq[2,\mu\nu]} \left[\sum_{a,b} \nabla_\nu \left(J^{[q,\{\mu_1^1\mu_2^1\ldots\mu_{b-1}^a\nu\mu_{b+1}^a\ldots\}]}\pi_{[q,\{\mu_1^1\mu_2^1\ldots\mu_{b-1}^a\mu\mu_{b+1}^a\ldots\}]} \right) + (\nabla_\mu J^{[q,\{\mu_j^i\}]})\pi_{[q,\{\mu_j^i\}]} \right].$$

[SL, arXiv:1204.1780]

Summary

- D-dimensional QFT can be explicitly mapped into a (D+1)-dimensional quantum theory of gravity based on a local RG
- Quantum beta function
- Example of emergent gravity
- Condensed matter example ?