

**2369-23**

**CIMPA/ICTP Geometric Structures and Theory of Control**

*1 - 12 October 2012*

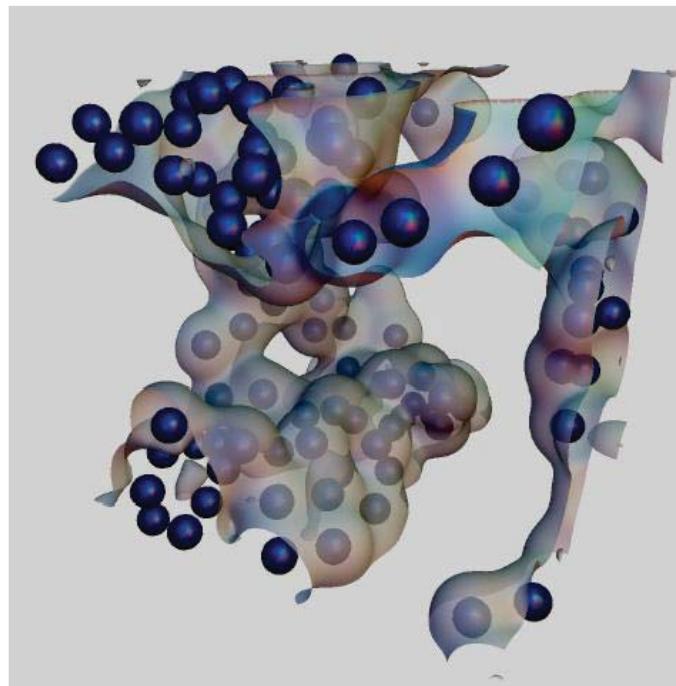
**New trends for warm dense matter research works and their applications**

Hitoki Yoneda

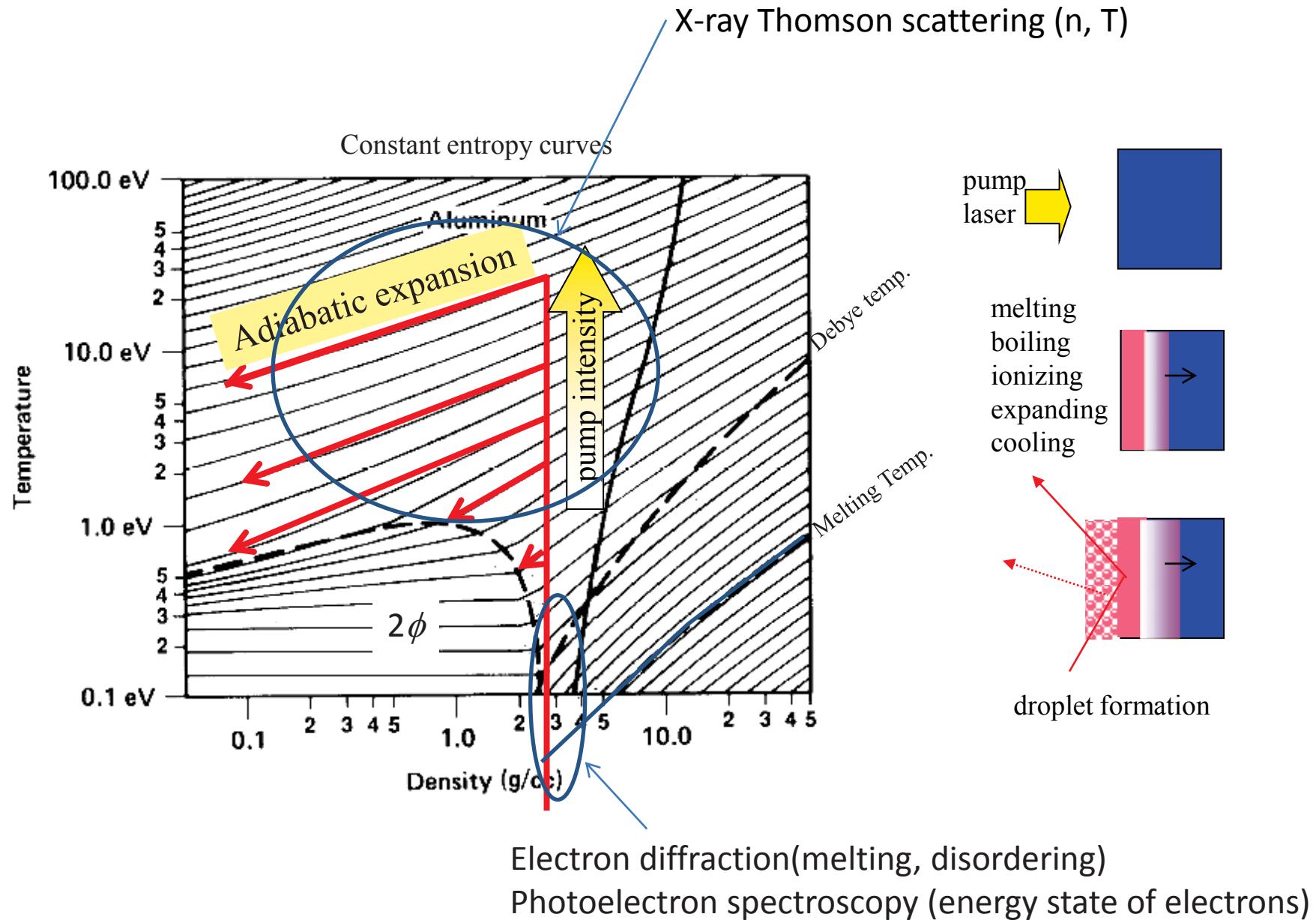
*Institute for Laser Science  
Tokyo  
Japan*

New trends for warm dense matter  
research works  
and their applications

We should measure directly the energy density state of electrons, disordering of ions, transport coefficient (electron, diffusion, thermal....), sound velocity, viscosity, and specific heat.



QMD simulation view of wdm Al, Michael Desjarlais, 2008 WDM school

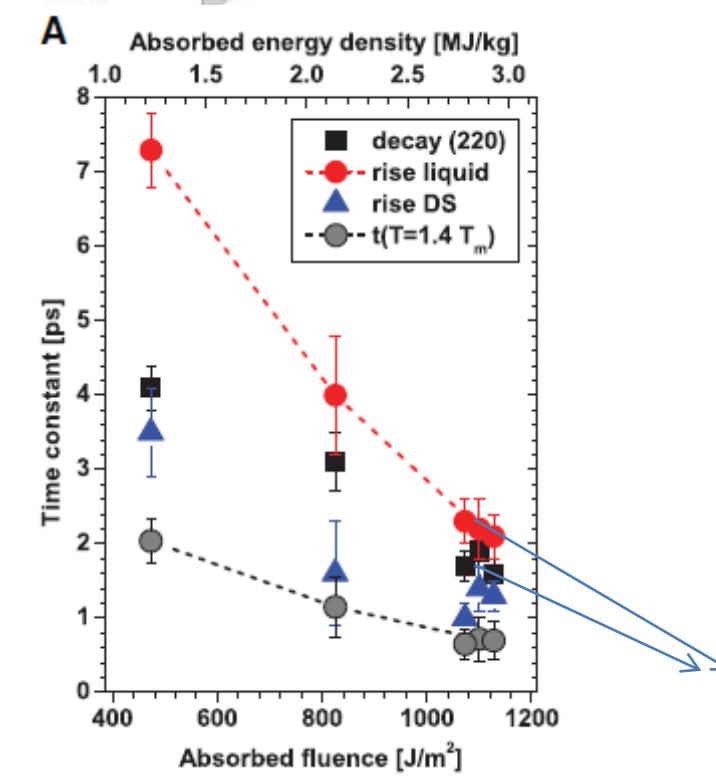
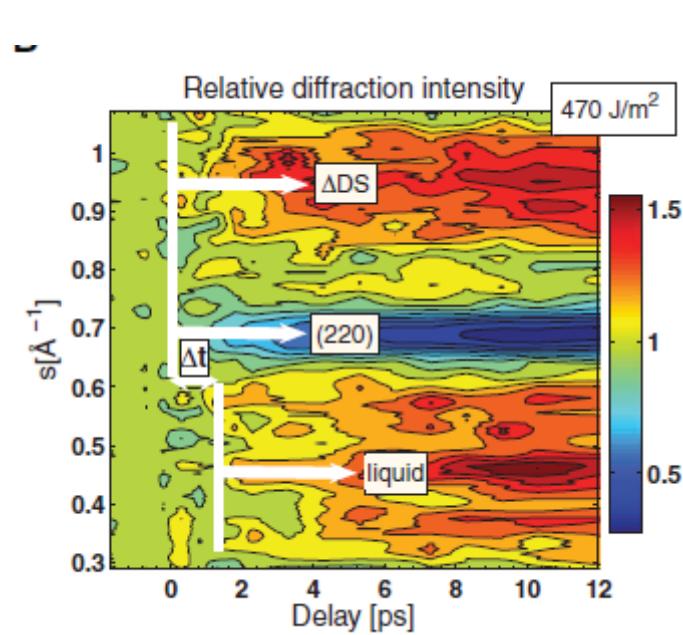
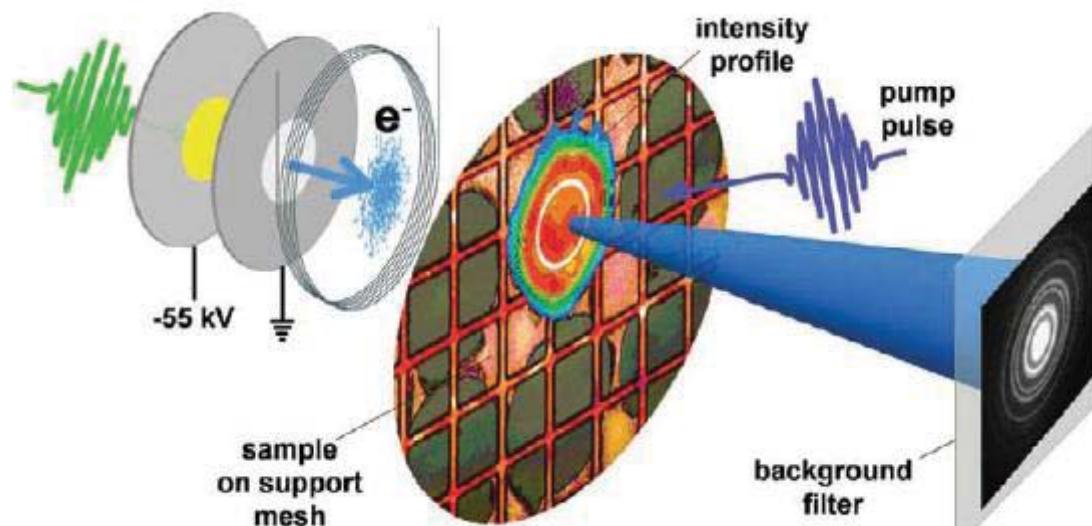


# New measurement method

- Electron diffraction experiment
- Angle-resolved photoemission spectroscopy
- X ray scattering and transmission experiment

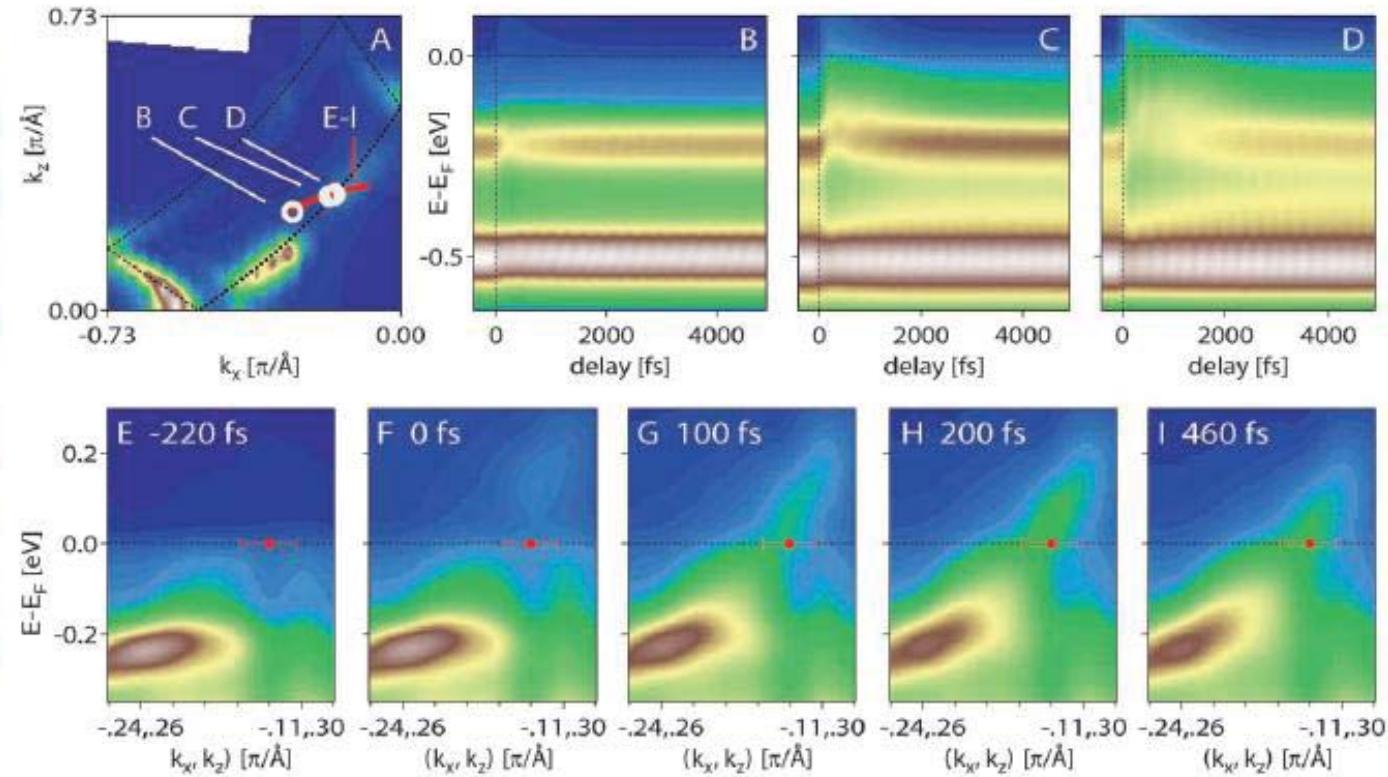
## Electron diffraction experiment

Ralph Ernstorfer, SCIENCE, 2009 VOL 323

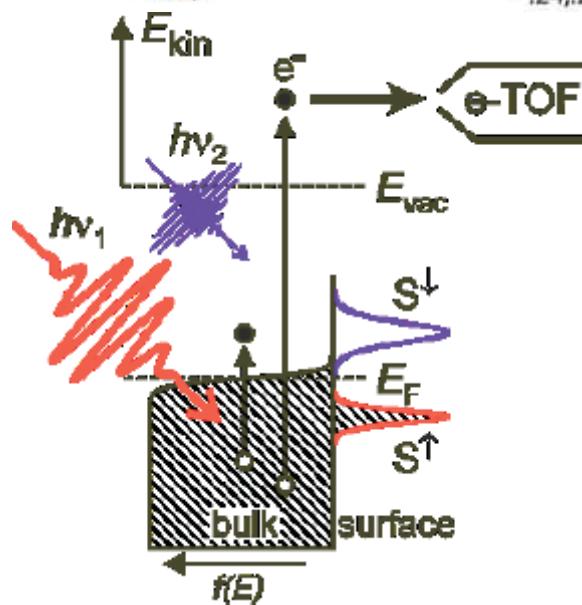


## Angle-resolved photoemission spectroscopy

**Fig. 3.** (A) Detail of the FS plot in Fig. 1A' with indicated positions (white circles) of time-resolved data shown in (B) to (D) for fixed  $k$  as a function of time delay. Indicated cut position (red line) of photoelectron intensity is shown as a function of energy, and position [(E) to (I)] for a momentum scan is shown as a function of time delays. All data were collected at 100 K and  $F = 2 \text{ mJ/cm}^2$ .  $k_F$  is marked in (E) to (I) (red dot). Error bars indicate the distance to the neighboring sample points, which is a good estimate for the error of  $k_F$ .



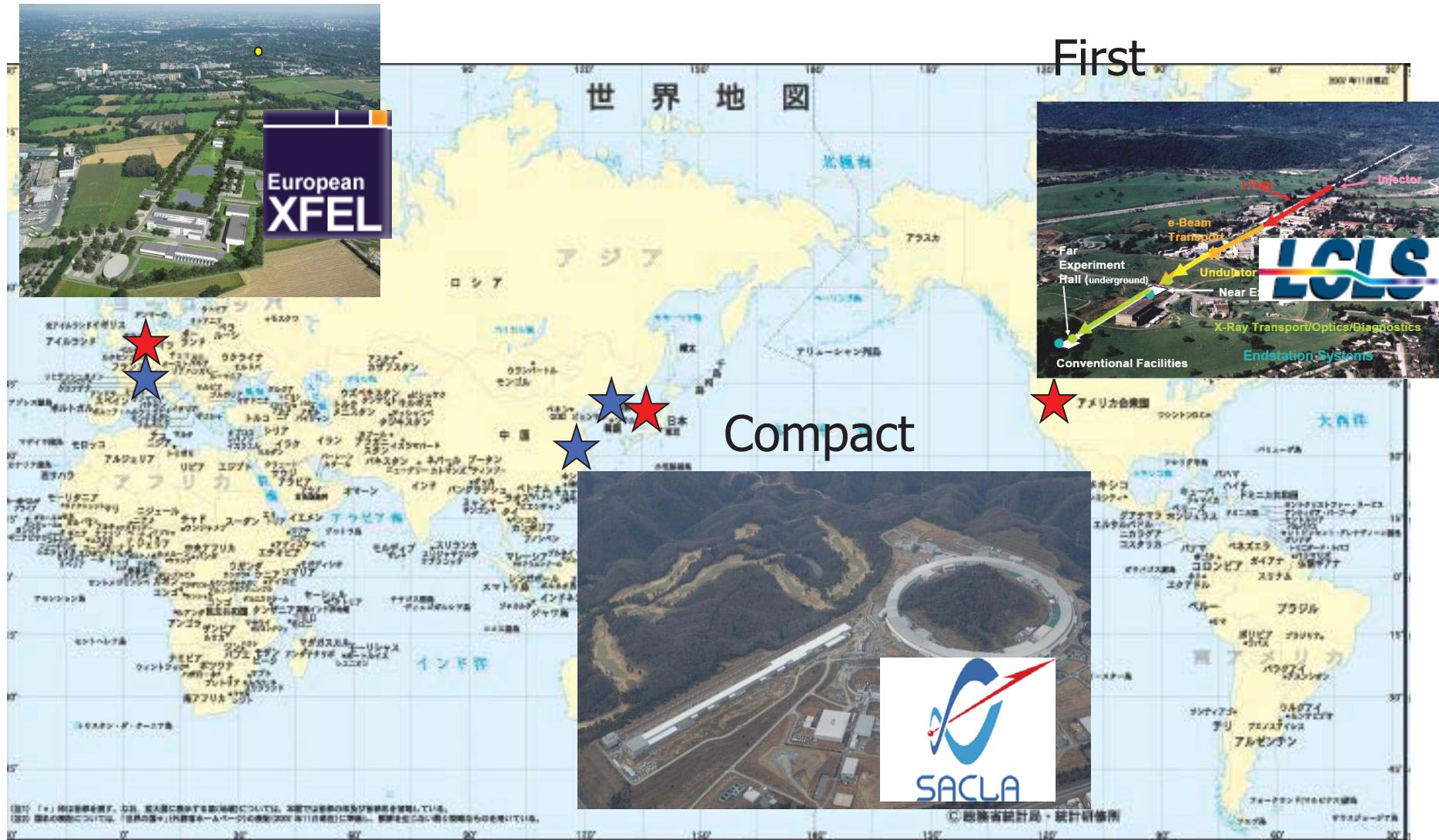
F. Schmitt, et al., *Science* **321**, 1649 (2008);



Martin Wolf, European XFEL SCS Workshop,  
PSI, 2 - 4 June 2009

# XFEL sources in the world

# High rep rate

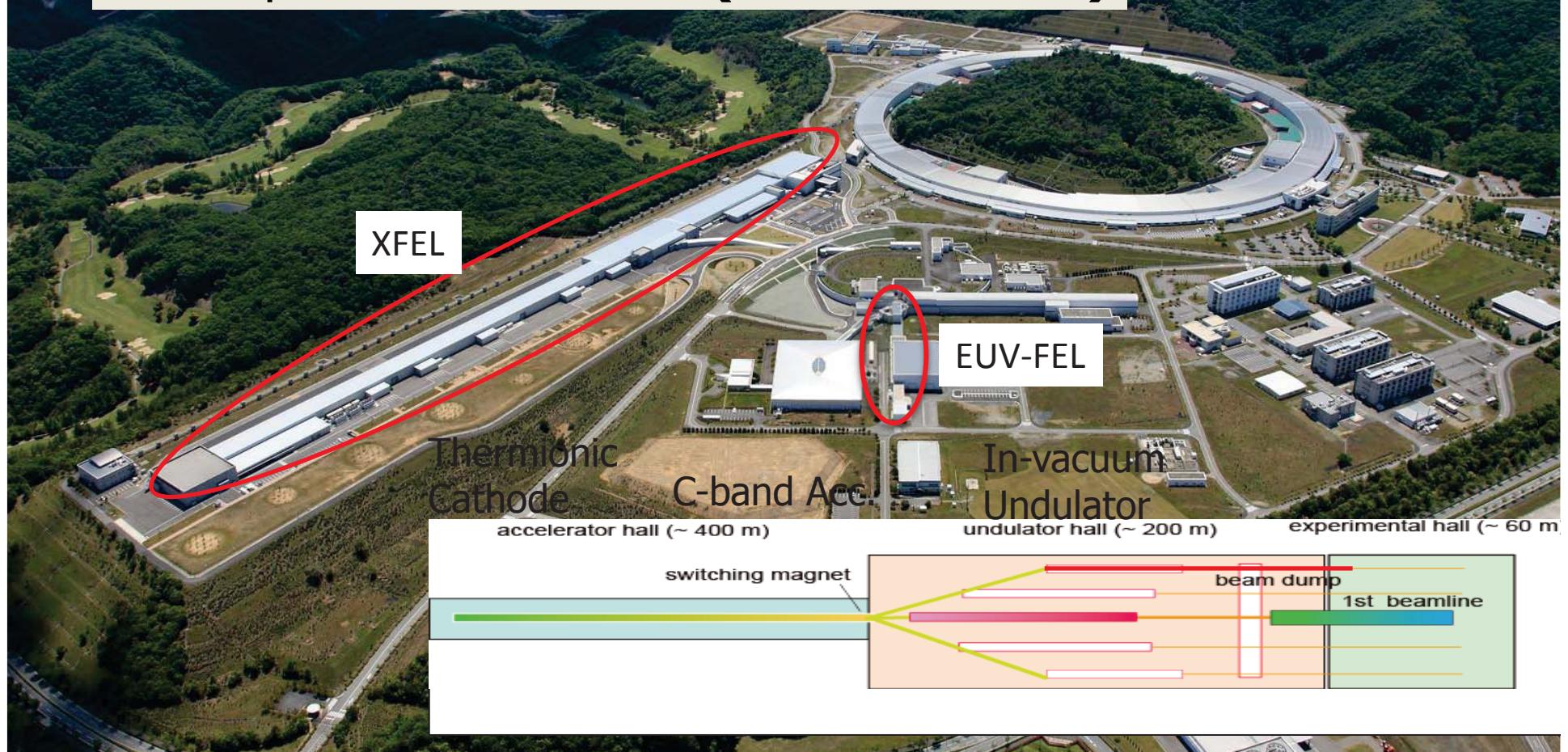


# SACLA & EUV FEL @SPring-8

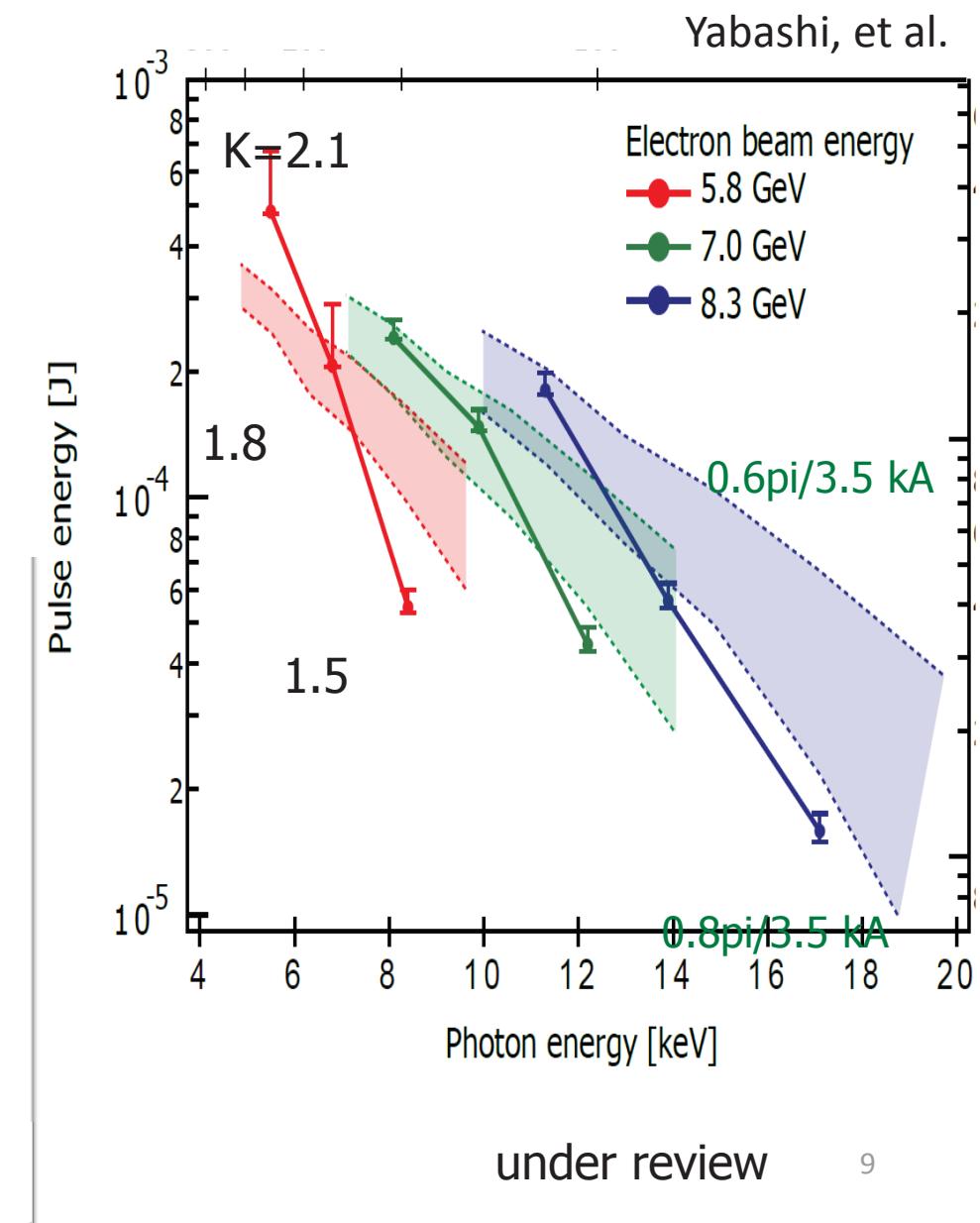
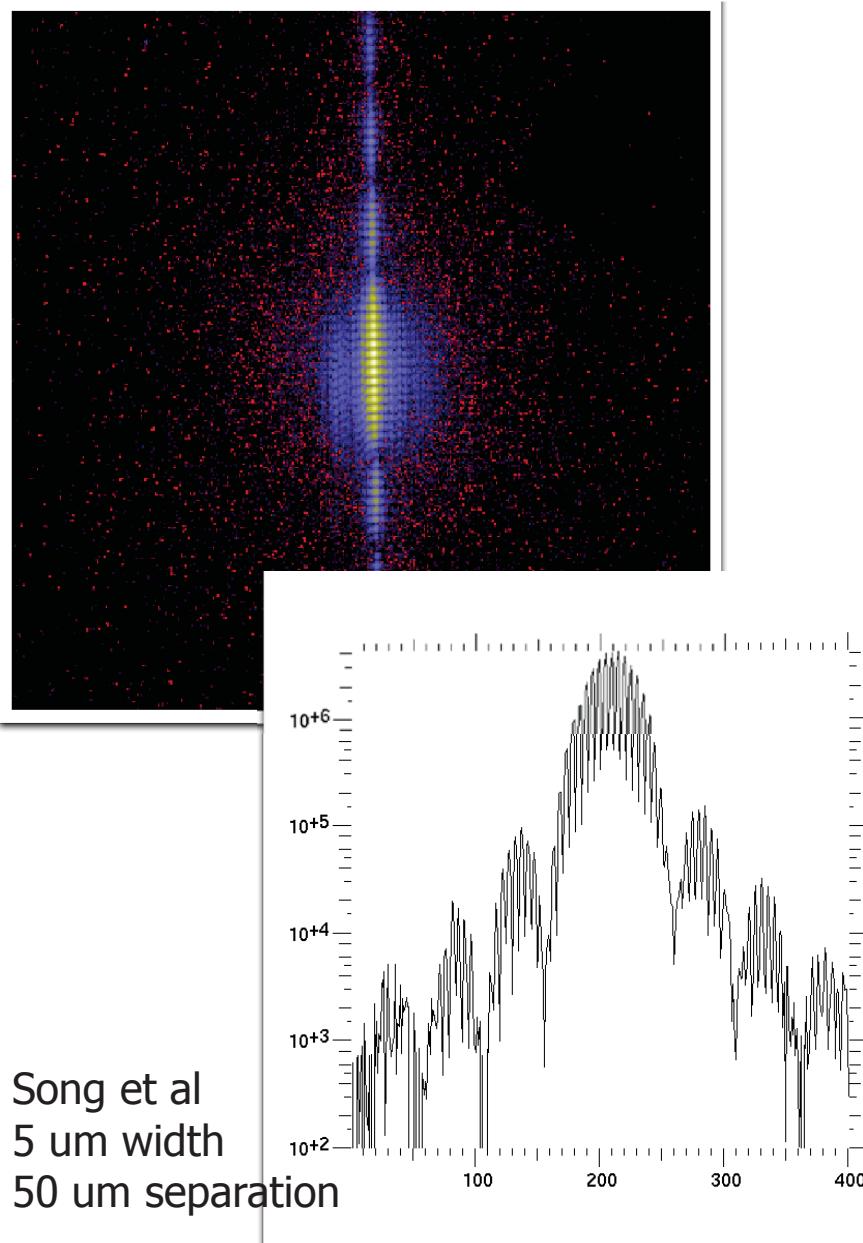
Compact XFEL

Construction: FY2006~2010

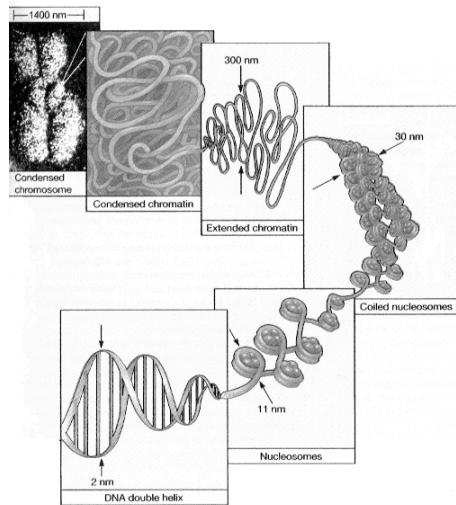
User Operation: FY2011~ **(7 March 2012)**



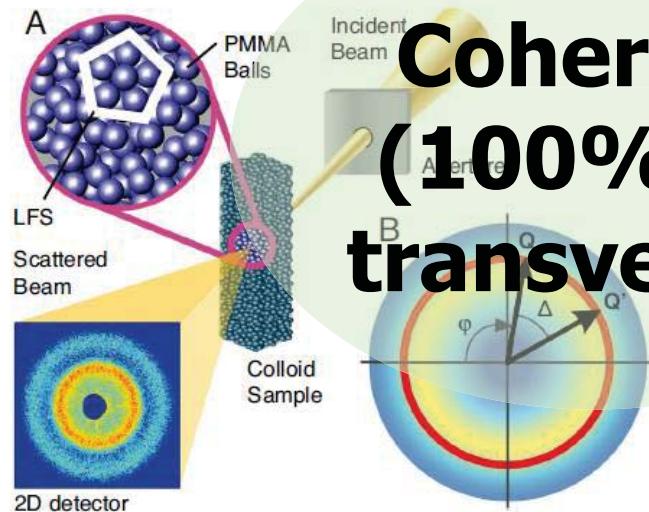
# Coherence & Wavelength & Pulse energy



# XFEL explores new worlds of science



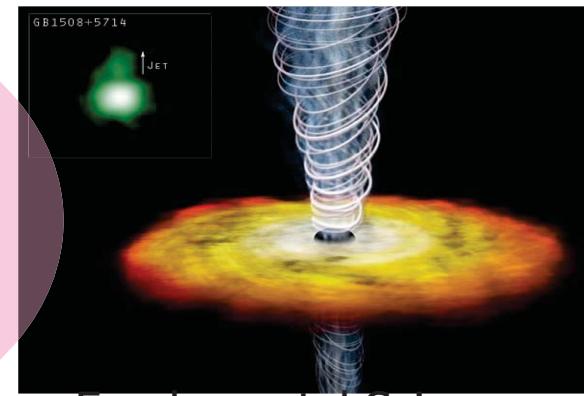
Biology & Medicine: Investigate non-crystalline materials



**Coherent  
(100% in  
transverse)**

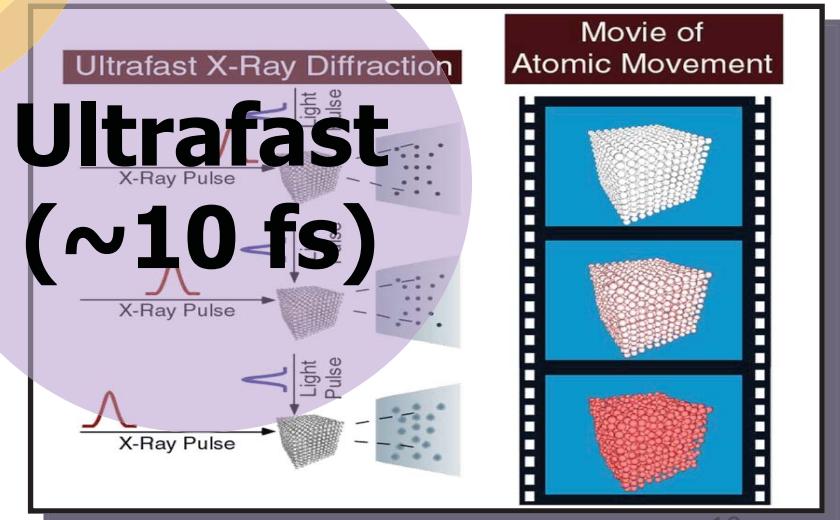
**Brilliance  
( $\times 10^9$ )**

**XFEL**



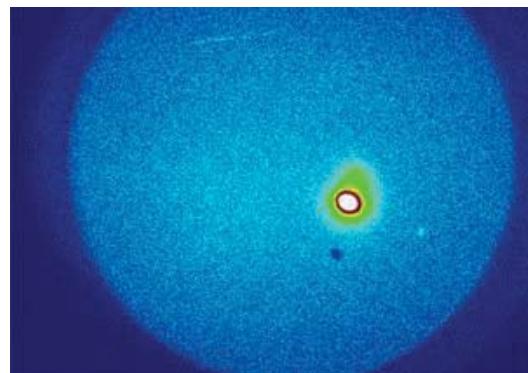
Fundamental Science:  
Create Extreme State

Environmental & Energy Science:  
Probe ultrafast reactions



# How intense of X-ray?

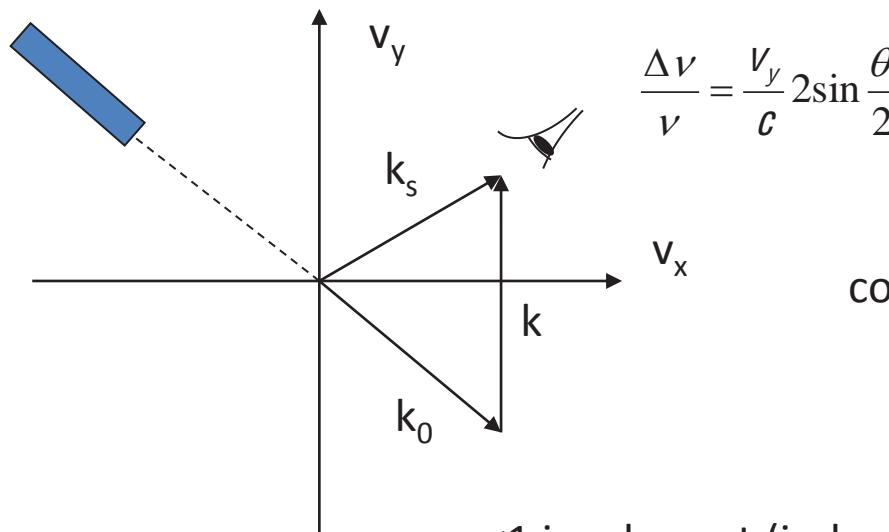
- X-ray Free electron laser (from 2011)  $I \sim 10^{21} \text{ W/cm}^2$   
5~20keV, ~mJ, ~10fs, focus diameter ~50nm(min. 7nm), 60Hz
- EUV Free electron laser (since 2006)  $I \sim 10^{15} \text{ W/cm}^2$   
45~170nm, ~20μJ, ~60fs, focus diameter ~6μm, 30Hz



# Diagnostics for Warm dense matter with X-ray

## Thomson Scattering

$$\alpha = \frac{\lambda_{Laser}}{4\pi\lambda_D \sin(\theta/2)} = 1/k_{Laser}\lambda_D \sin(\theta/2) = \frac{\lambda_{Laser}}{4\pi \sin(\theta/2)} \left( \frac{4\pi n e^2}{k T_e} \right)^{1/2}$$



collective or incoherent?

$\alpha < 1$  incoherent (independent electron  
 $\Rightarrow$  electron velocity distribution function)

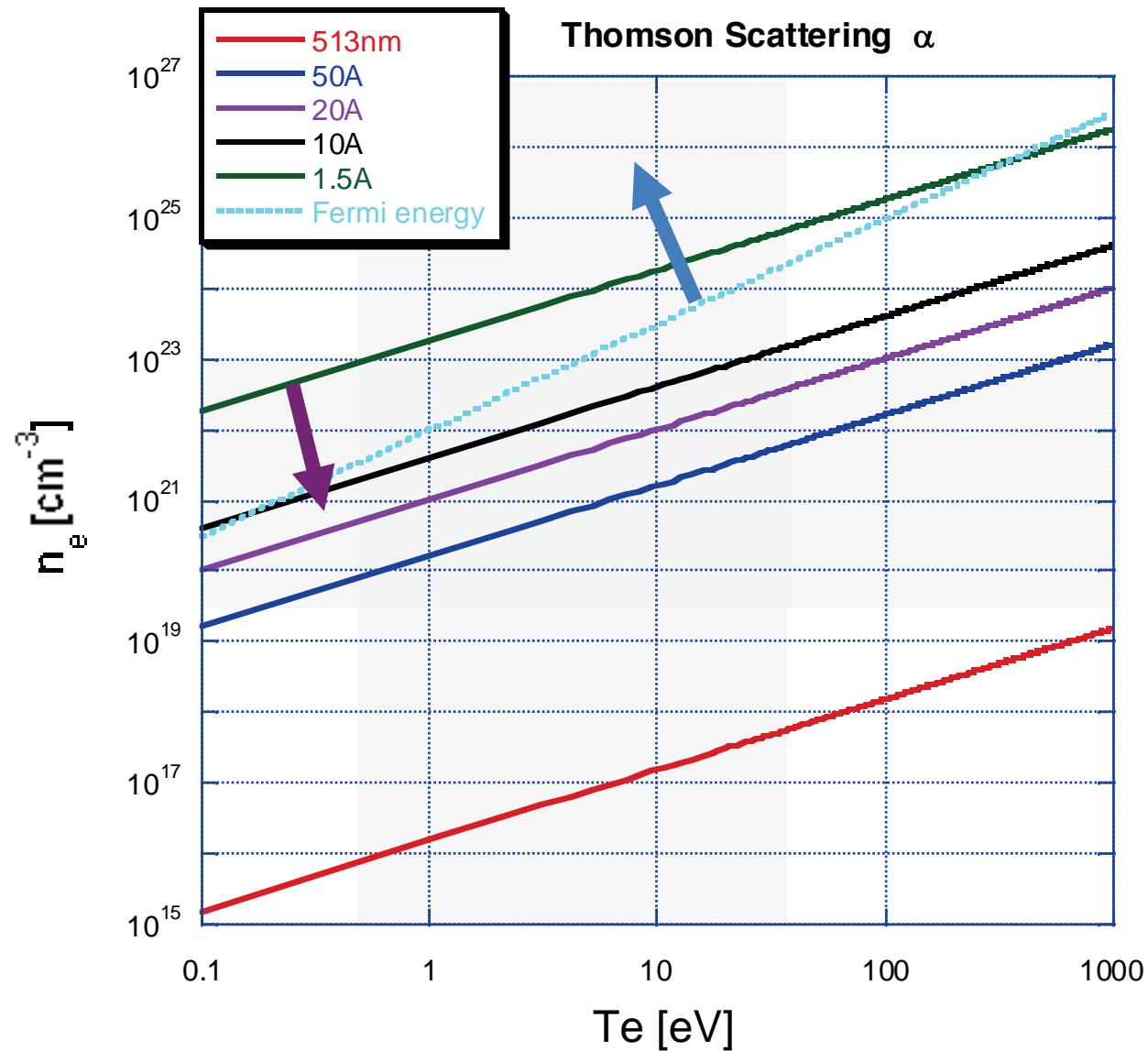
$\alpha > 1$  collective (ion acoustic wave, electron plasma wave)

dispersion relation:  $\omega_0^2 = \omega_{pe}^2 + 3(kT_e/m)k^2$

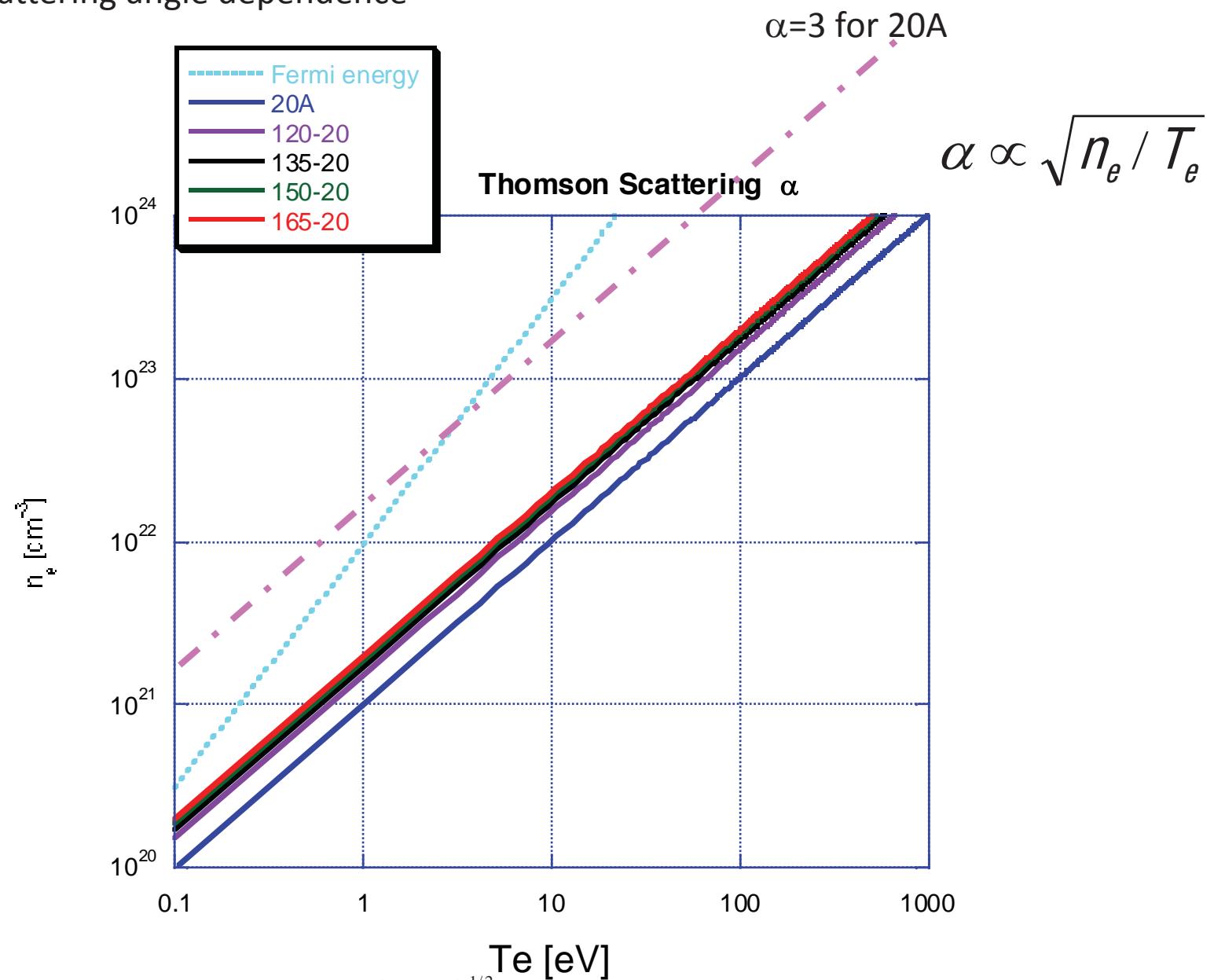
spectral width  $\sim$  ion Doppler width

At  $\alpha \sim 1$ , we can decide both of temperature and density.

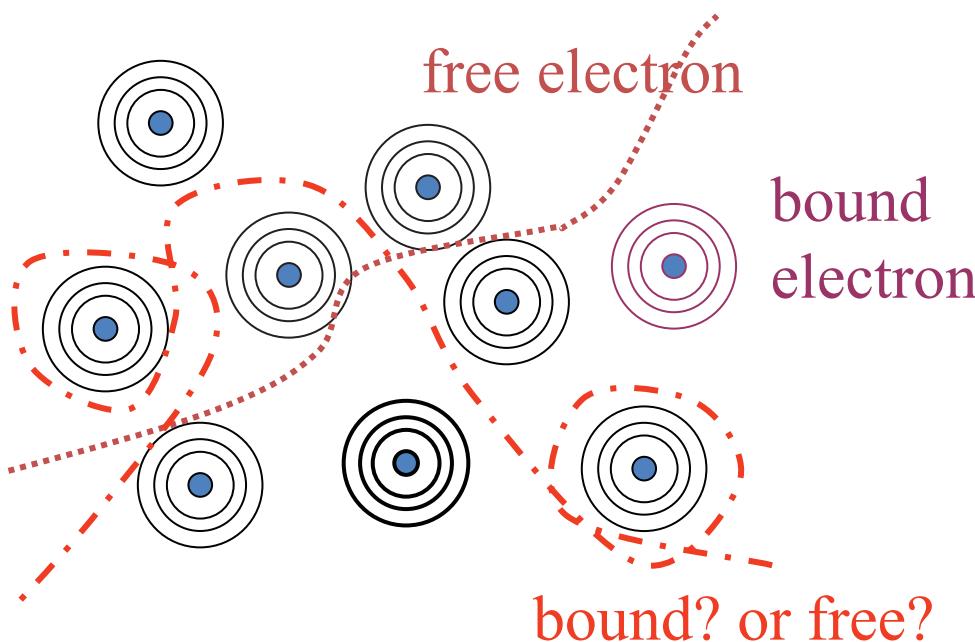
$\alpha=1$  line



## Scattering angle dependence



$$\alpha = \frac{\lambda_{\text{Laser}}}{4\pi\lambda_D \sin(\theta/2)} = 1/k_{\text{Laser}}\lambda_D \sin(\theta/2) = \frac{\lambda_{\text{Laser}}}{4\pi \sin(\theta/2)} \left( \frac{4\pi n e^2}{k T_e} \right)^{1/2}$$



## Scattering cross section

$$S(k, \omega) = |f_i(k) + q(k)|^2 S_{ii}(k, \omega) + Z_f S_{ee}^0(k, \omega) + Z_c \int \tilde{S}_{ce}(k, \omega - \omega') S_s(k, \omega') d\omega'$$

density correlations of electrons <= ion motion

free electron scattering

inelastic scattering with bound electrons

ion form factor

screening cloud of free and valence electrons

ion-ion density correlation function

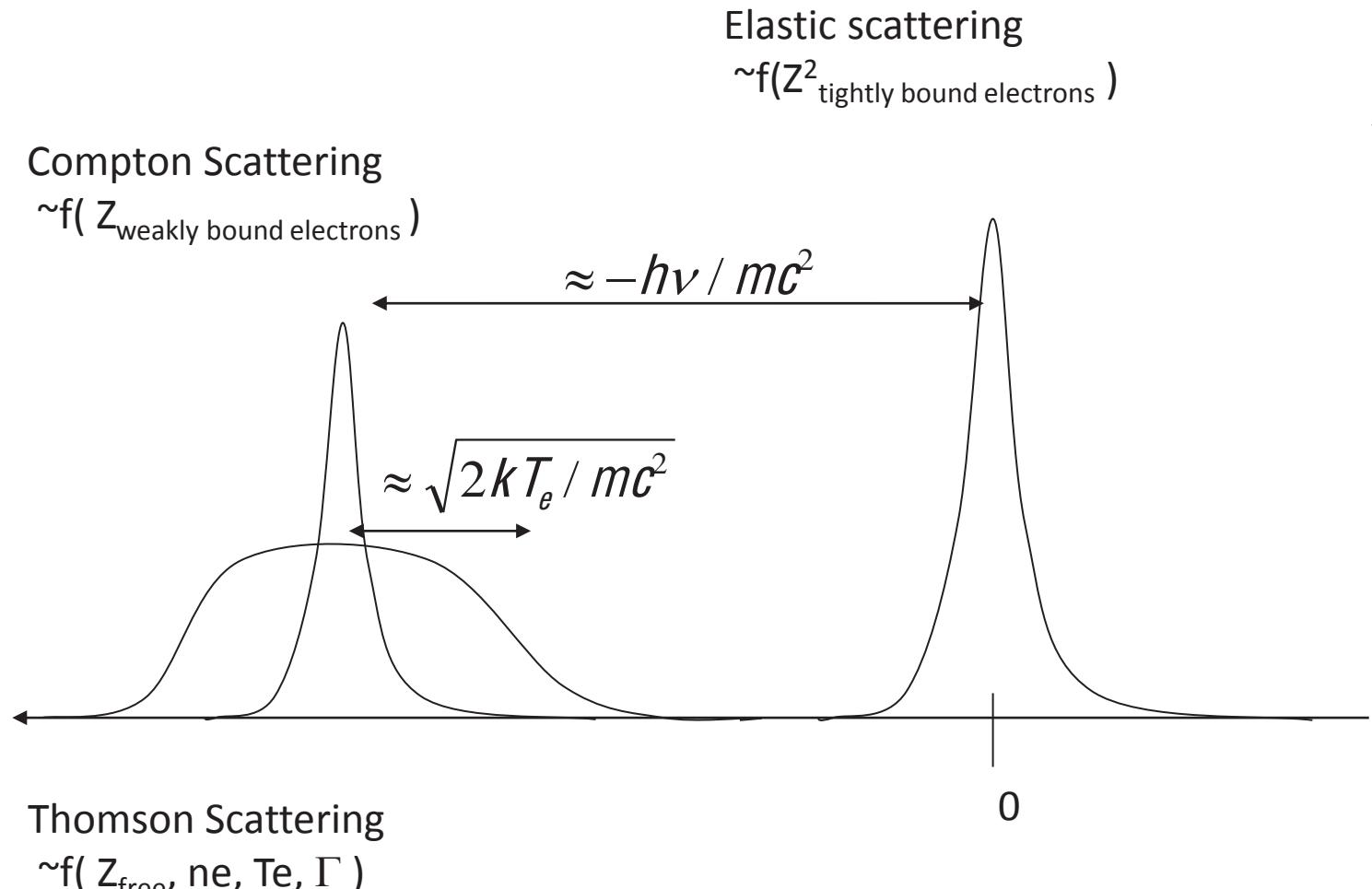
charge for free electron

high frequency part of e-e correlation function

charge for bound electron

bound-free transition

self-motion of ions



$$h\nu=1\text{keV}, T_e=10\text{eV}, \text{shift}=2\times 10^{-5}, \text{width}=6\times 10^{-3}$$

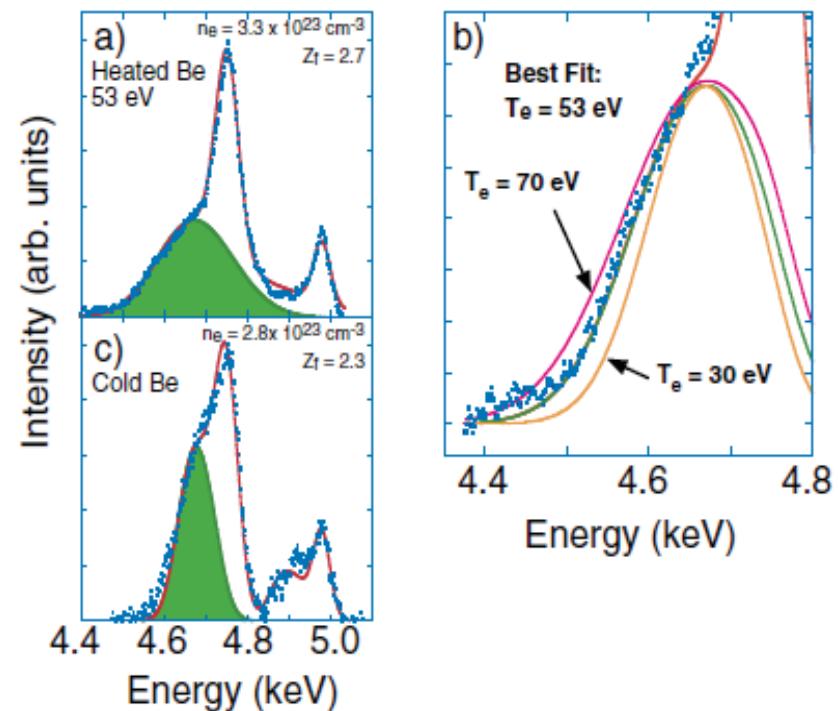
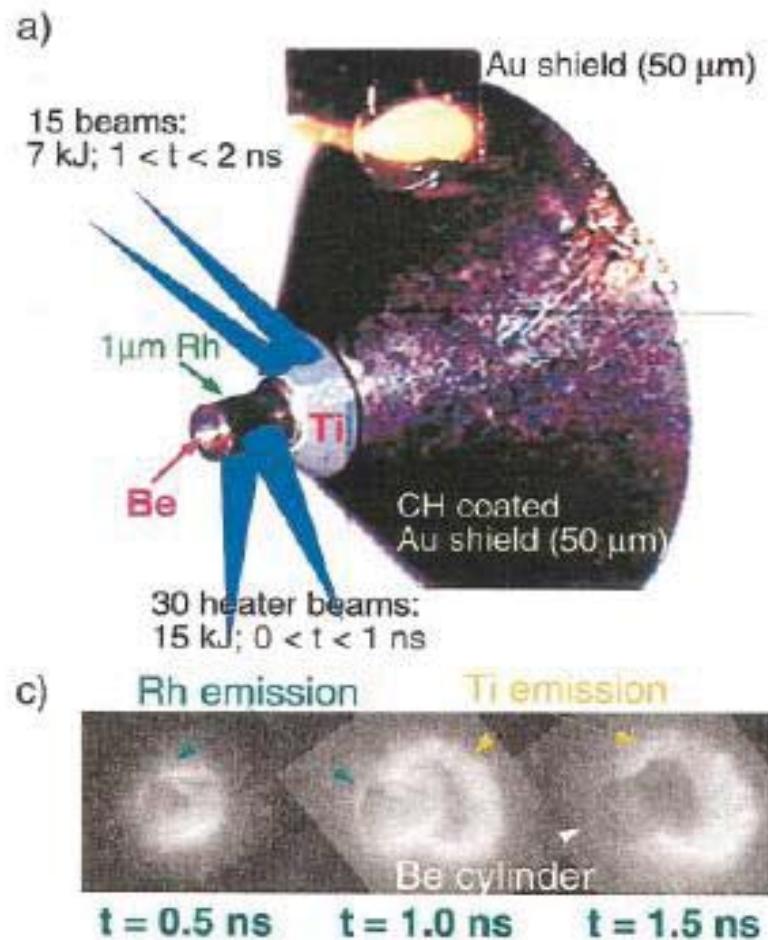


FIG. 3 (color). (a) Experimental x-ray scattering data (blue dots) from the heated Be plasma with a theoretical fit yielding  $T_e = 53$  eV and  $n_e = 3.3 \times 10^{23} \text{ cm}^{-3}$ . (b) Red wing of the scattering spectrum together with the best fit and calculated spectra for  $T_e = 30$  eV and  $T_e = 70$  eV indicating that the electron temperature is determined with 10%–20% accuracy. (c) Experimental x-ray scattering data (blue dots) from the cold Be indicating  $T_e \approx 2$  eV and  $n_e = 2.8 \times 10^{23} \text{ cm}^{-3}$ . The green shaded area corresponds to the Compton down-shifted feature as obtained from the fit.

## Electronic structure measurements of dense plasmara.

G. Gregori, et al.,

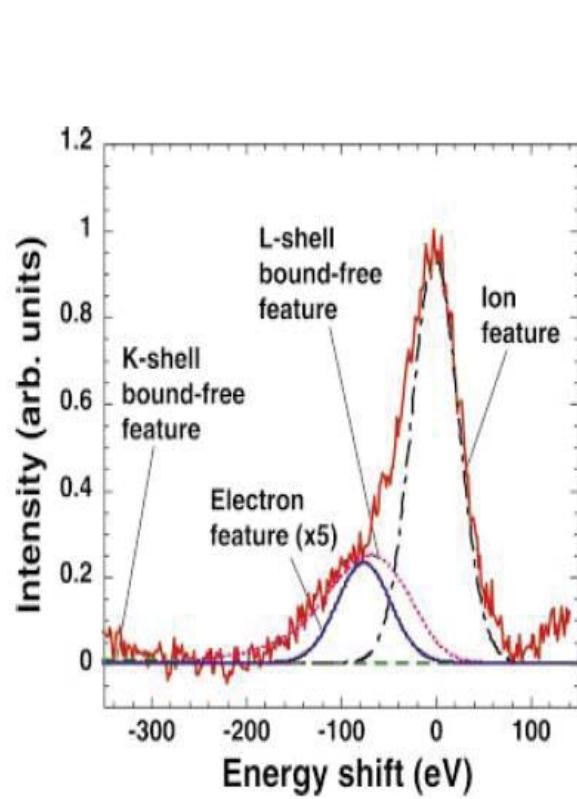


FIG. 6. (Color) Spectrum of the unheated foam with separated contributions from each different scattering mechanism. Convolution with instrument response is added each terms. The probe radiation is the Ti He- $\alpha$  line at 4.75 keV, and the scattered x rays are collected at  $\sim 130^\circ \pm 5^\circ$  scattering angle with best fit parameters  $T_e < 5$  eV and  $Z_f = 0.26$ . The ionization energy for isolated neutral carbon is  $E_B^0 = 11$  eV for L-shell electrons and  $E_B^0 = 286$  eV for K-shell electrons.

4.75keV Ti He- $\alpha$  ( $d\lambda/\lambda < 0.005$ )

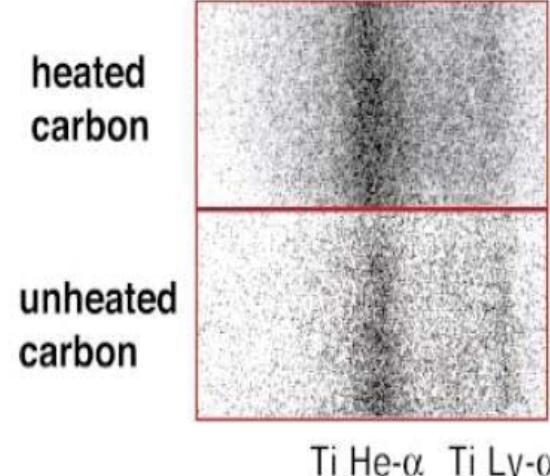
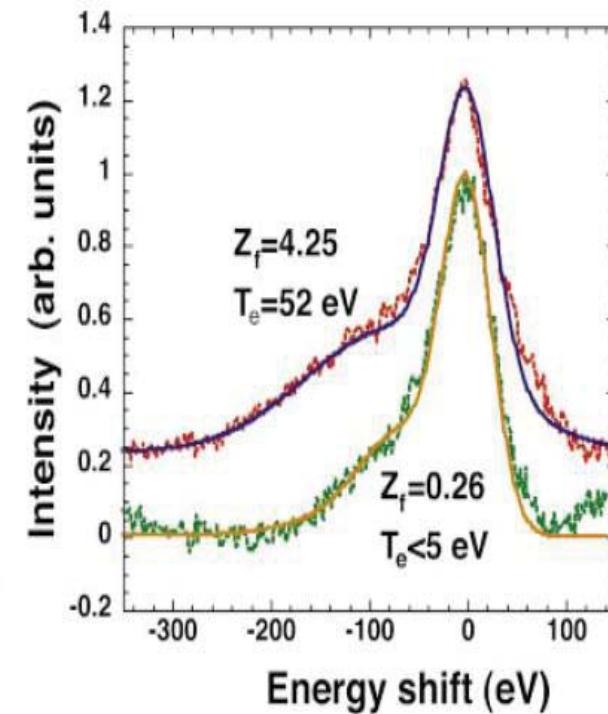


FIG. 4. (Color) Experimental x-ray scattering data from a heated carbon foam ( $0.72 \text{ g/cm}^3$ ) and a cold (unheated) carbon foam. The raw data, as measured by the detector are shown on the left side. Lineouts and best fits are reported on the right panel. The probe radiation is the Ti He- $\alpha$  line at 4.75 keV, and the scattered x rays are collected at  $\sim 130^\circ \pm 5^\circ$  scattering angle. Best fit parameters and corresponding spectra are also plotted in the figure. For the high temperature foam,  $\alpha = 0.17$ ,  $T_F = 10.4$  eV, and  $\Gamma = 0.2$ ; while for the cold foam  $\alpha = 0.13$ ,  $T_F = 1.6$  eV, and  $\Gamma = 0.9$ .



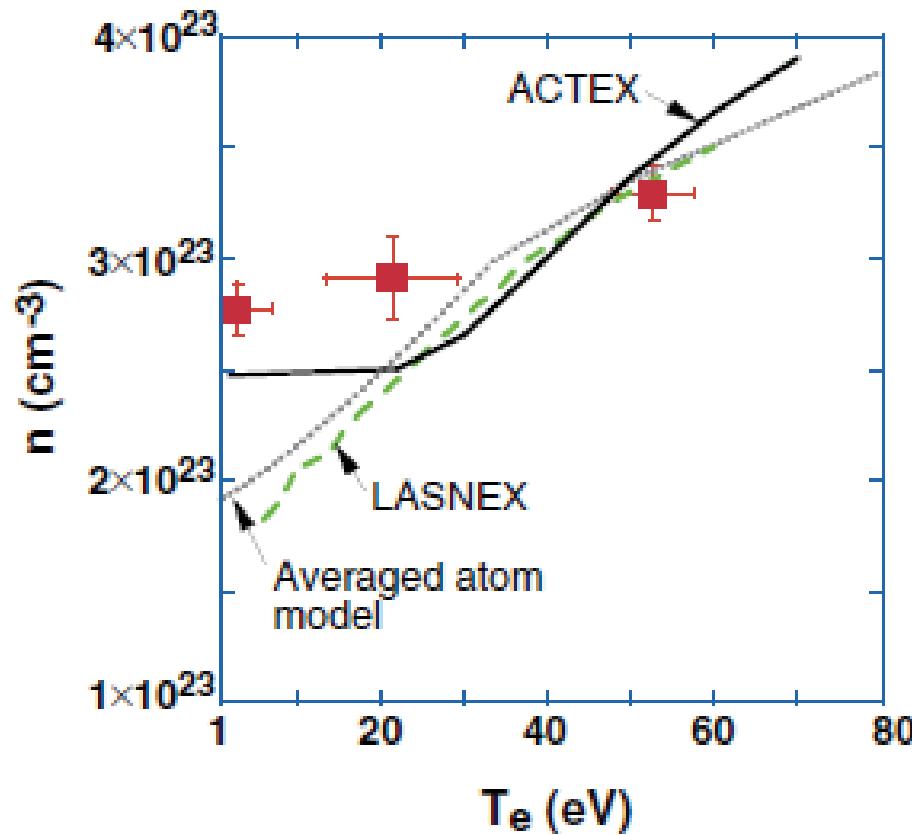
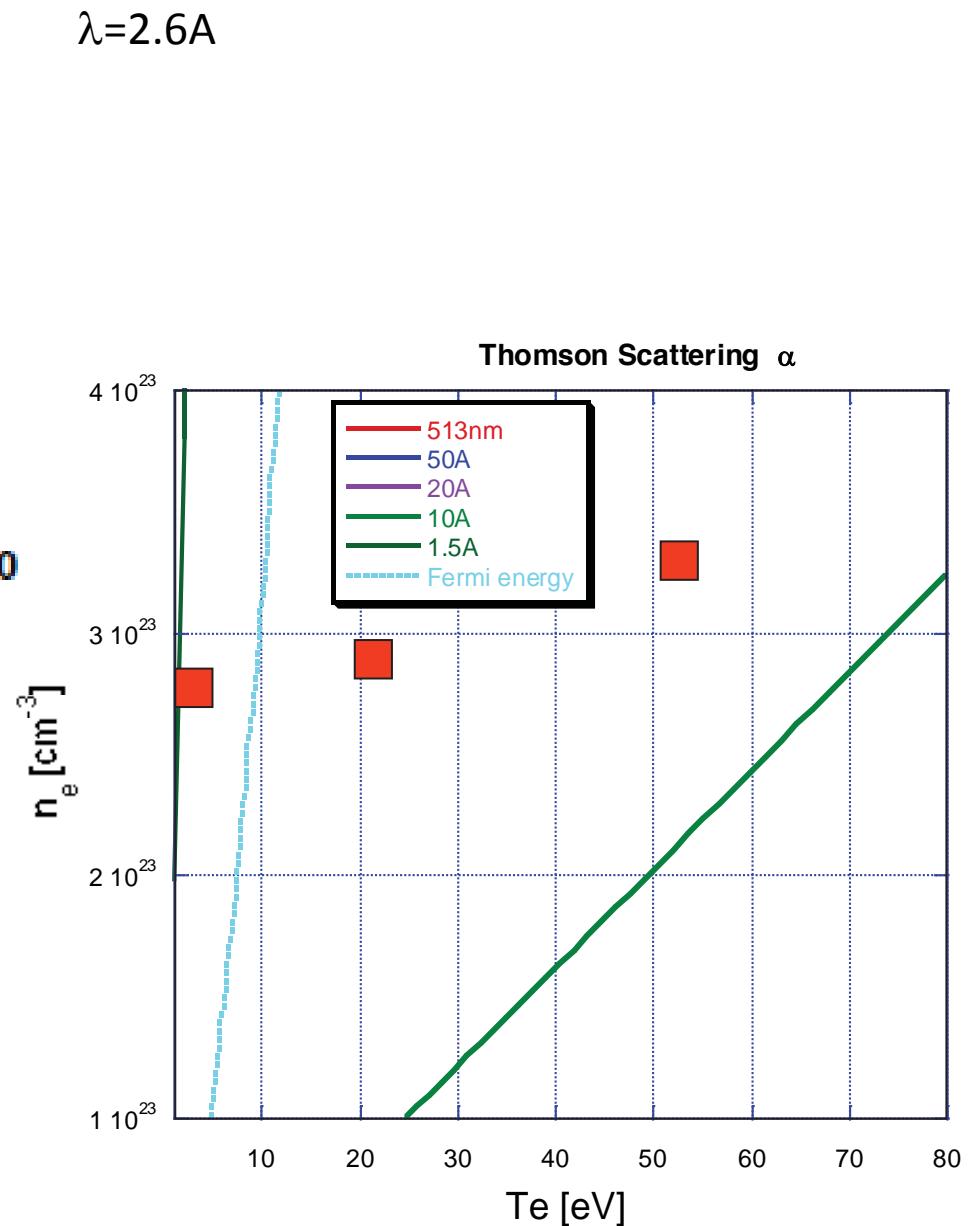


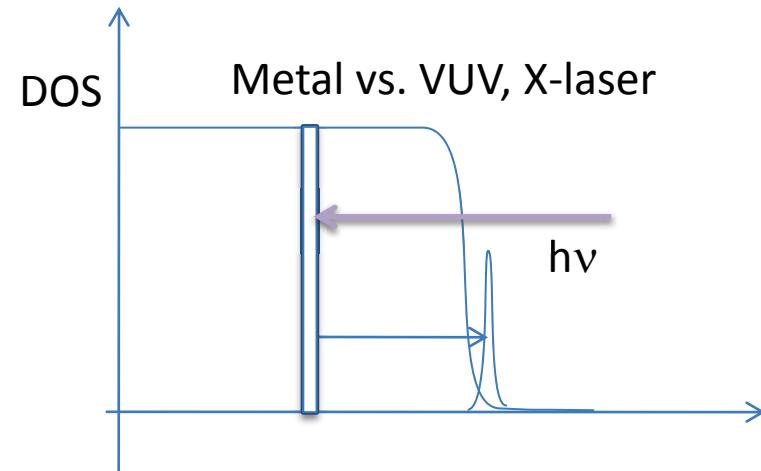
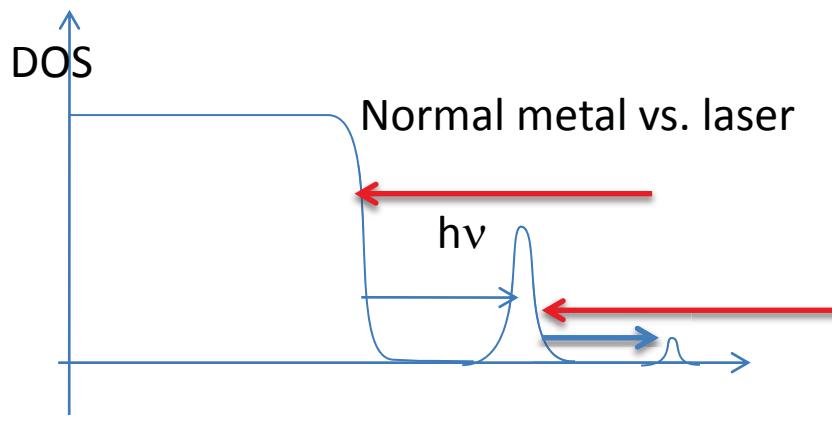
FIG. 4 (color). Density-temperature phase diagram along with the results of the x-ray scattering measurements and simulation using the ACTEX model [18], an averaged atom model [19], and radiation-hydrodynamic modeling [20]. The data and the ACTEX model include the free plus weakly bound (conduction) electrons. The averaged atom model and LASNEX only include free (ionized) electrons.



Creation of new states of matter

# What different in interaction physics with XFEL?

# absorption

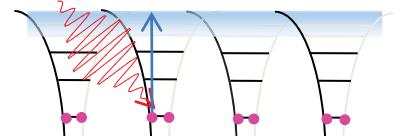
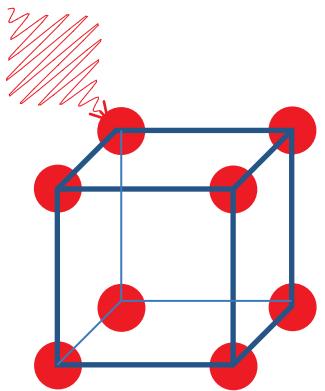


$$n_{e_{solid}} < n_c \text{ for EUV, Xray}$$

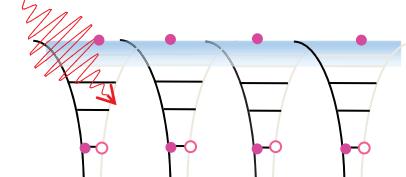
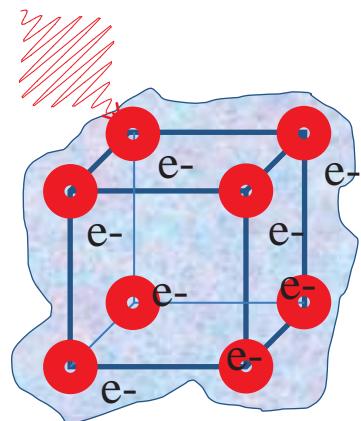
There is No effect of preplasma even in high intensity laser interaction.

It is easy to penetrate into solid density.

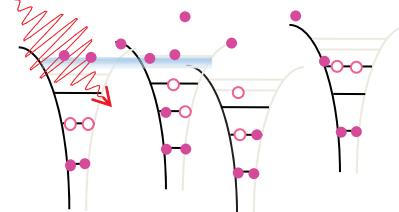
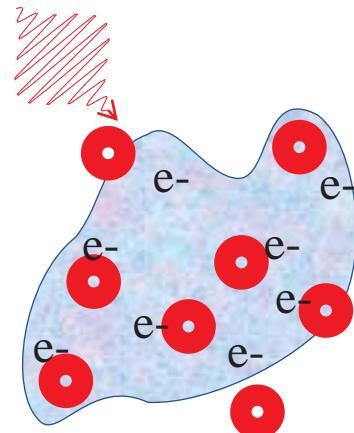
# Material condition



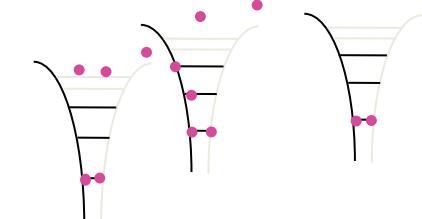
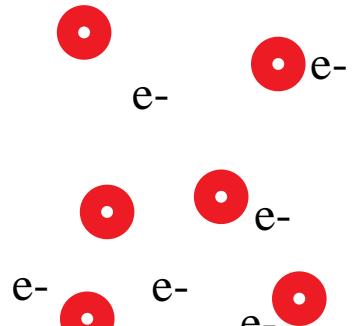
Condensed matter  
(1)



Hollow atom solid  
(2)



Warm dense matter  
(3)



Plasma  
(4)

→ time

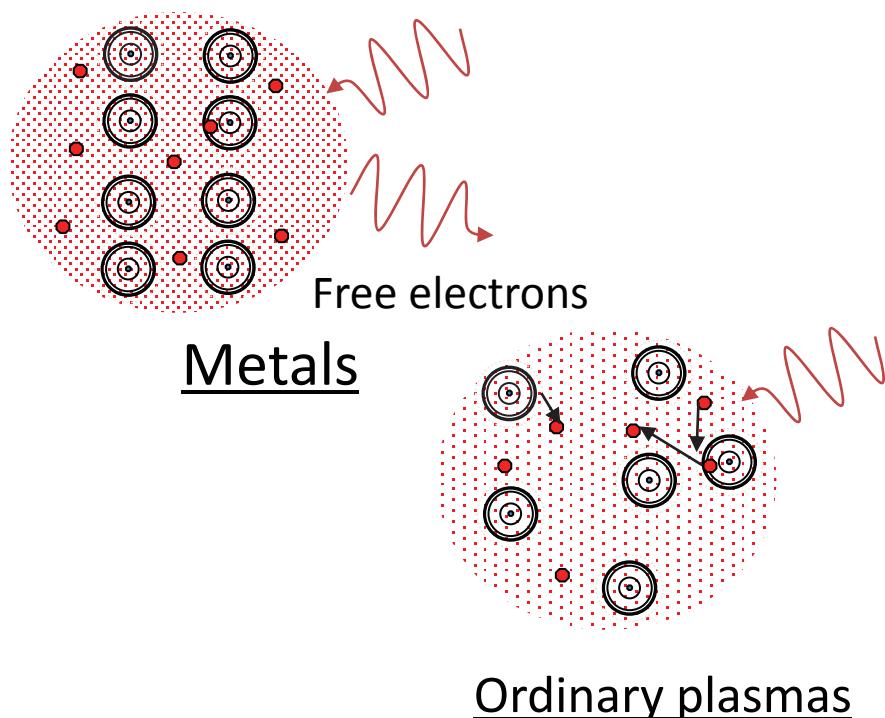
# High energy density matter

Xray & EUV laser interaction experiments

$$F[\rho, T] = E_{cold}[\rho] + F_{elec}[\rho, T] + F_{ion}[\rho, T] + E_{excitation}$$

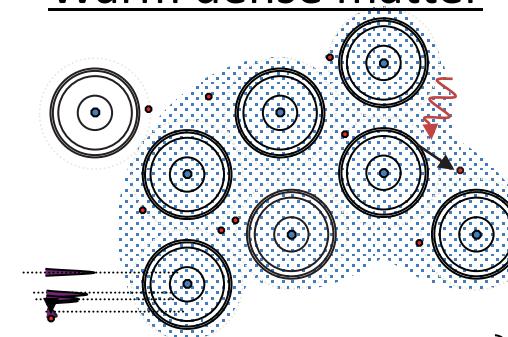
EOS internal energy    Cold curve    Electron-Thermal    Ion-Thermal

Ion excitation Energy

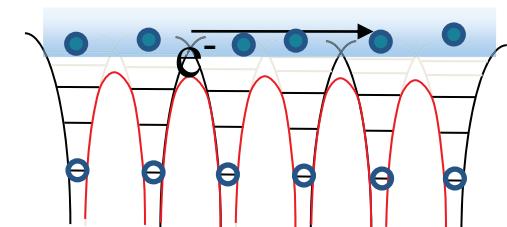


Disorder  $\Rightarrow$  Ordered

Warm dense matter

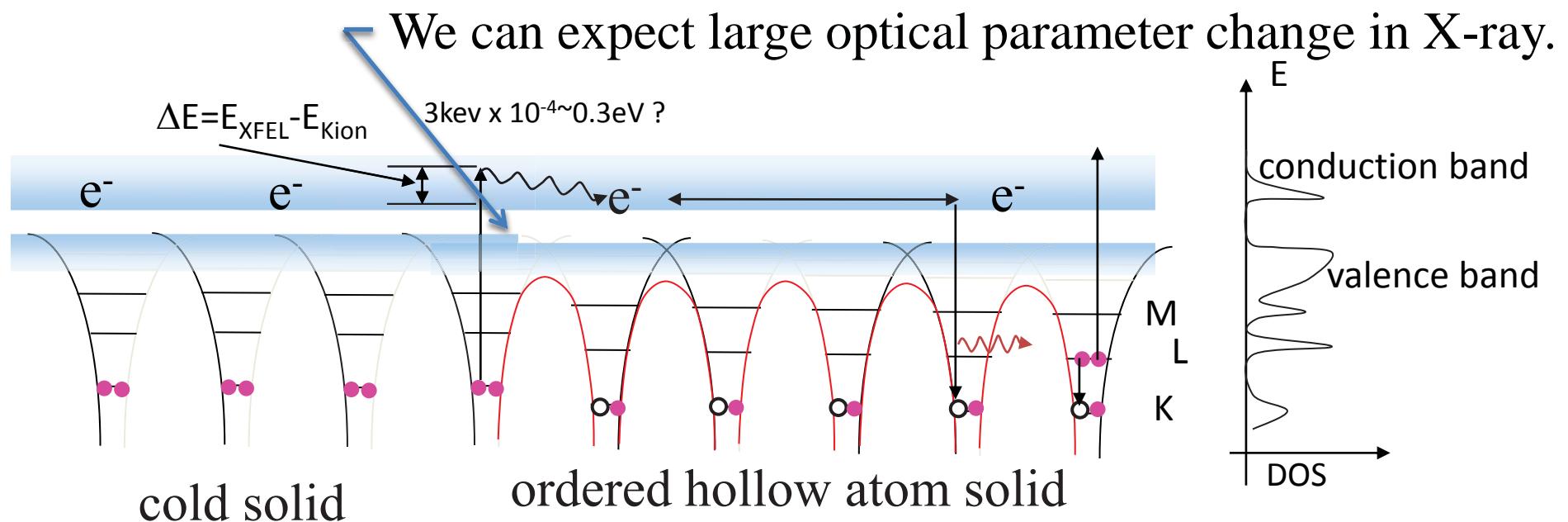


band structure  
+ionization

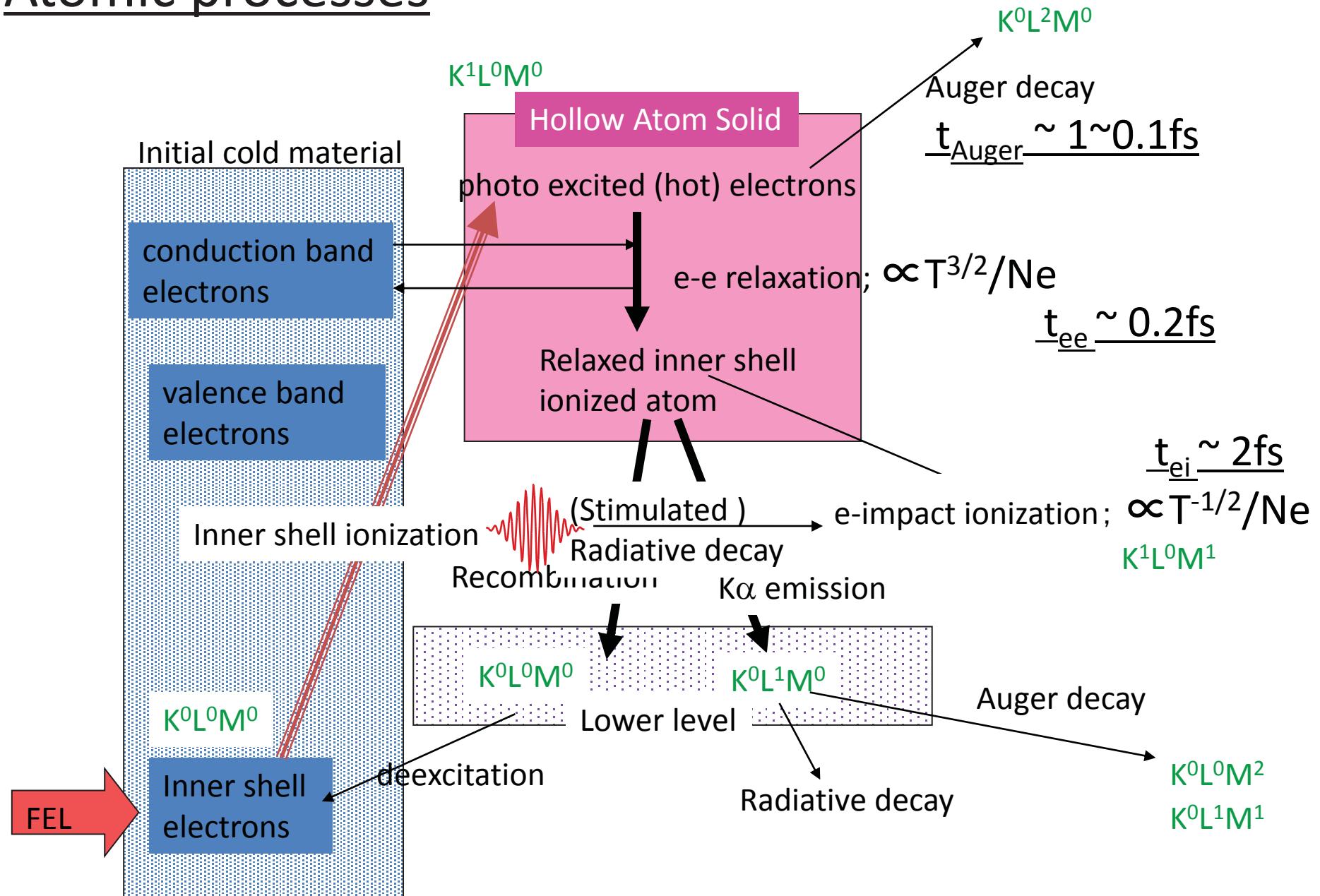


# Well-tuned X-ray and EUV laser photon

If  $t_{\text{pulse}} < t_{\text{disassemble}}$ , we create ordered solid with inner shell excited ions.



# Atomic processes



## What different in interaction physics with XFEL? (cont.)

Speed of excitation can be higher than  
relaxation rate of atomic process in solid.

$$\frac{dN}{dt} = \frac{\alpha I}{h\nu} \quad \Rightarrow \quad \tau_{excitation} = \frac{N}{dN/dt} = \frac{N_{solid}}{\alpha I} h\nu$$

# EUV interaction has big advantage because large  $\alpha$  and smaller  $h\nu$ .

(We can sustain inner shell hole under radiation.)

$\alpha$  at Ti K edge  $\sim 3350\text{cm}^{-1}$  at 5keV  $\Rightarrow \tau < 1\text{fs}$  for  $I > 10^{18}\text{W/cm}^2$

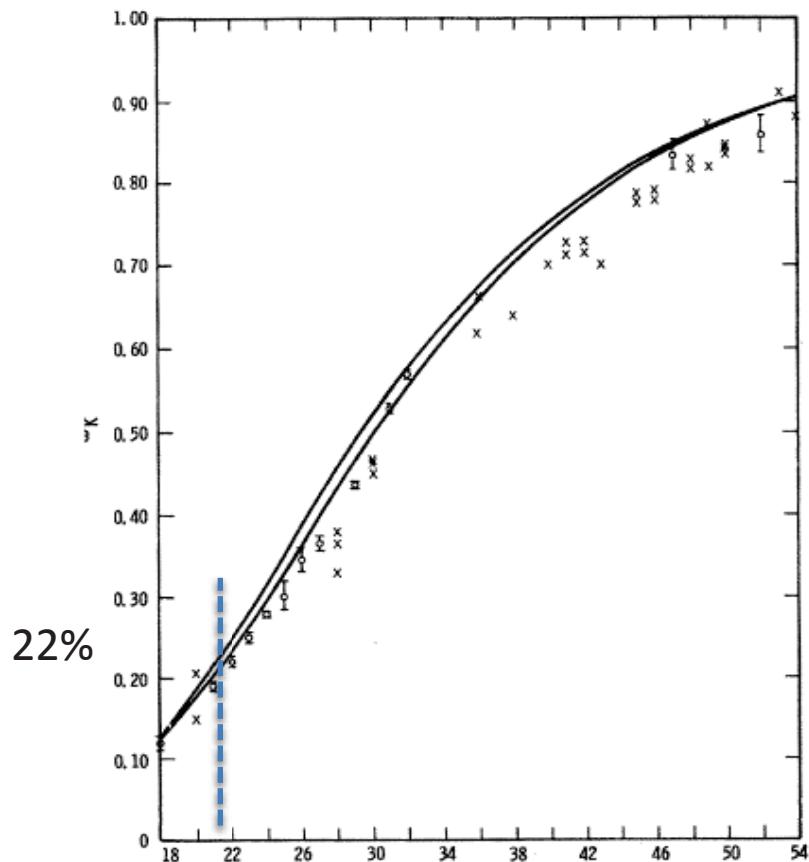


FIG. 1.  $K$ -shell fluorescence yield versus  $Z$ . The upper curve is the uncorrected; the lower the corrected calculations. The data points with error bars are from Refs. 5–11, and the crosses from Ref. 12.

$\alpha$  at Ti K edge  $\sim 3350\text{cm}^{-1}$

$$\frac{dN}{dt} = \frac{\alpha I}{h\nu}$$

If  $I = 10^{18}\text{W/cm}^2$ ,  $(dN/dt/N)^{-1} \sim \tau_{\text{Auger}}$

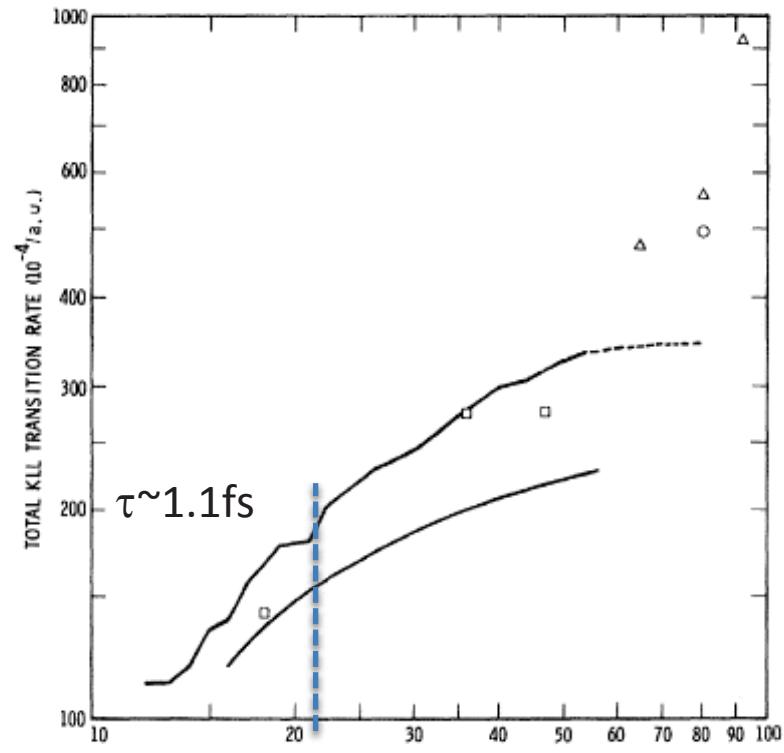
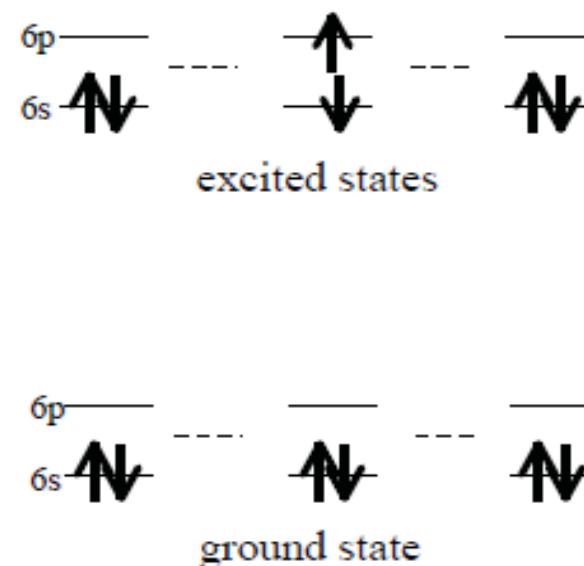
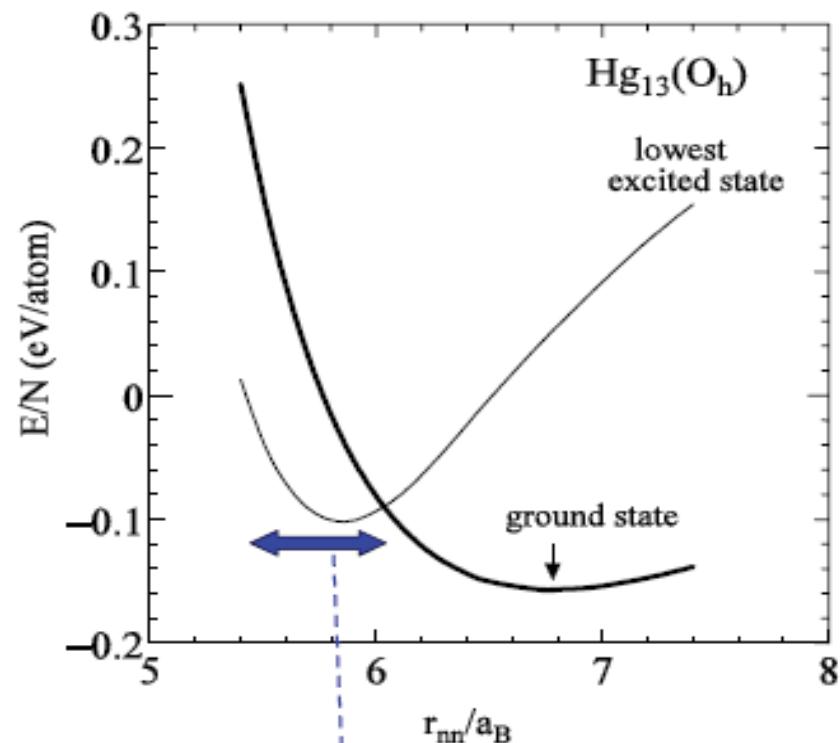


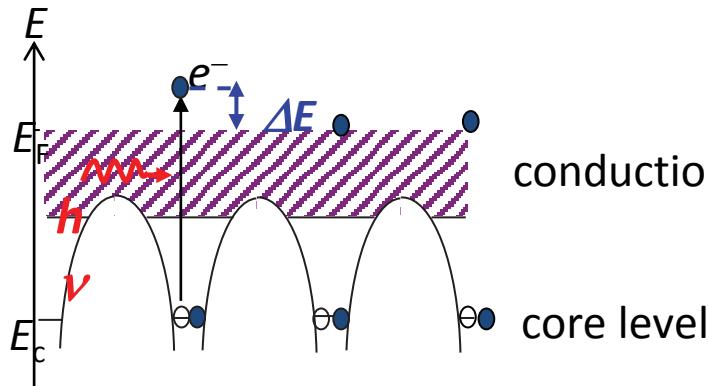
FIG. 3. Total  $KLL$  Auger rate versus  $Z$  in  $10^{-4}/\text{a.u.}$  (1 a.u. =  $2.42 \times 10^{-17}$  sec). The upper solid curve is from our calculations; the dashed portion was obtained for  $Z = 60, 65, 70$ , and  $80$  for which we did  $KLL$  calculations only. The lower solid curve is from Ref. 19; the squares from Ref. 36; the circle from Ref. 20; and the triangles from Ref. 21. E. J. McGuire, Phys. Rev. A 2, 273–278 (1970)

## potential energy curve of $\text{Hg}_{13}$ (cuboctahedron)



*“excitonic ground state” for  $r_{nn} < 6.0$  a.u  
--- large admixture of p-state*

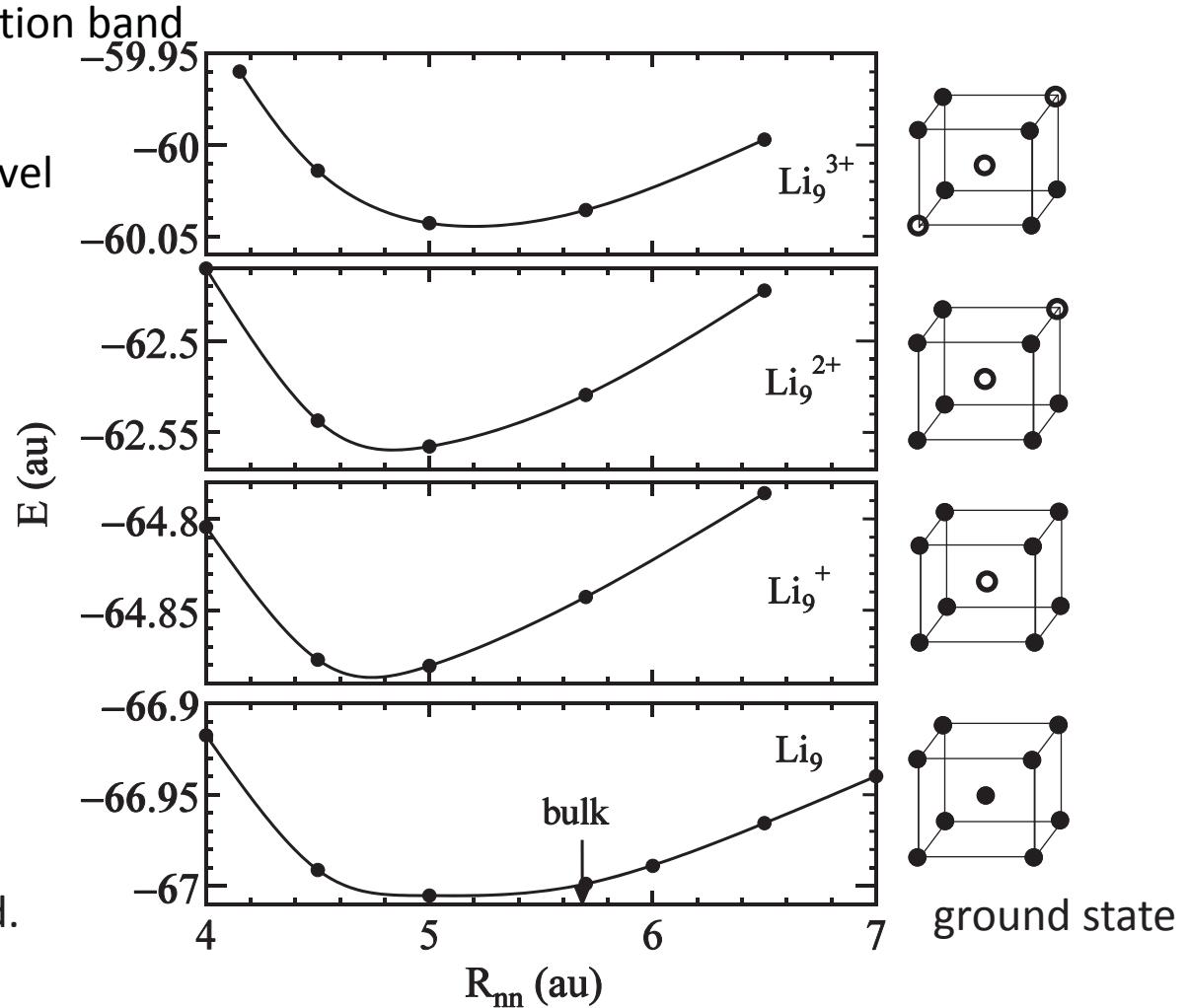
# Potential energy curves for $\text{Li}_9^{z+}$ ions



Clusters are bound  
at least for  $Z \leq 3$  !

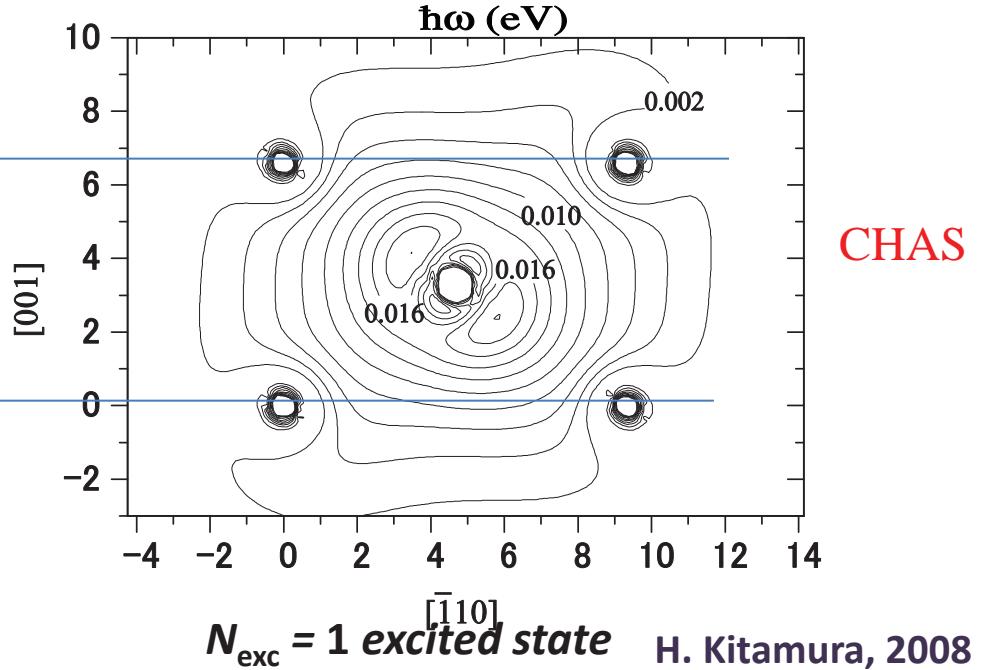
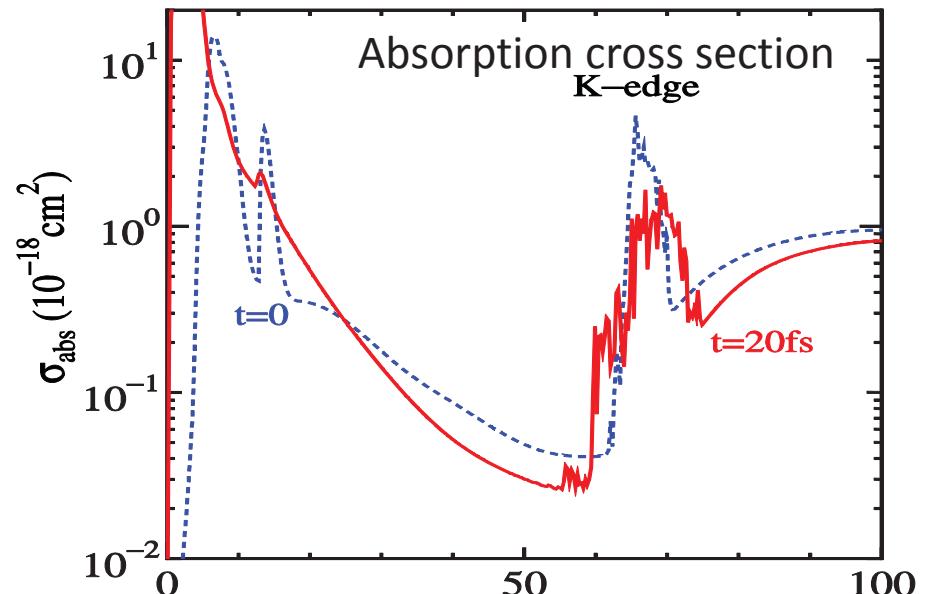
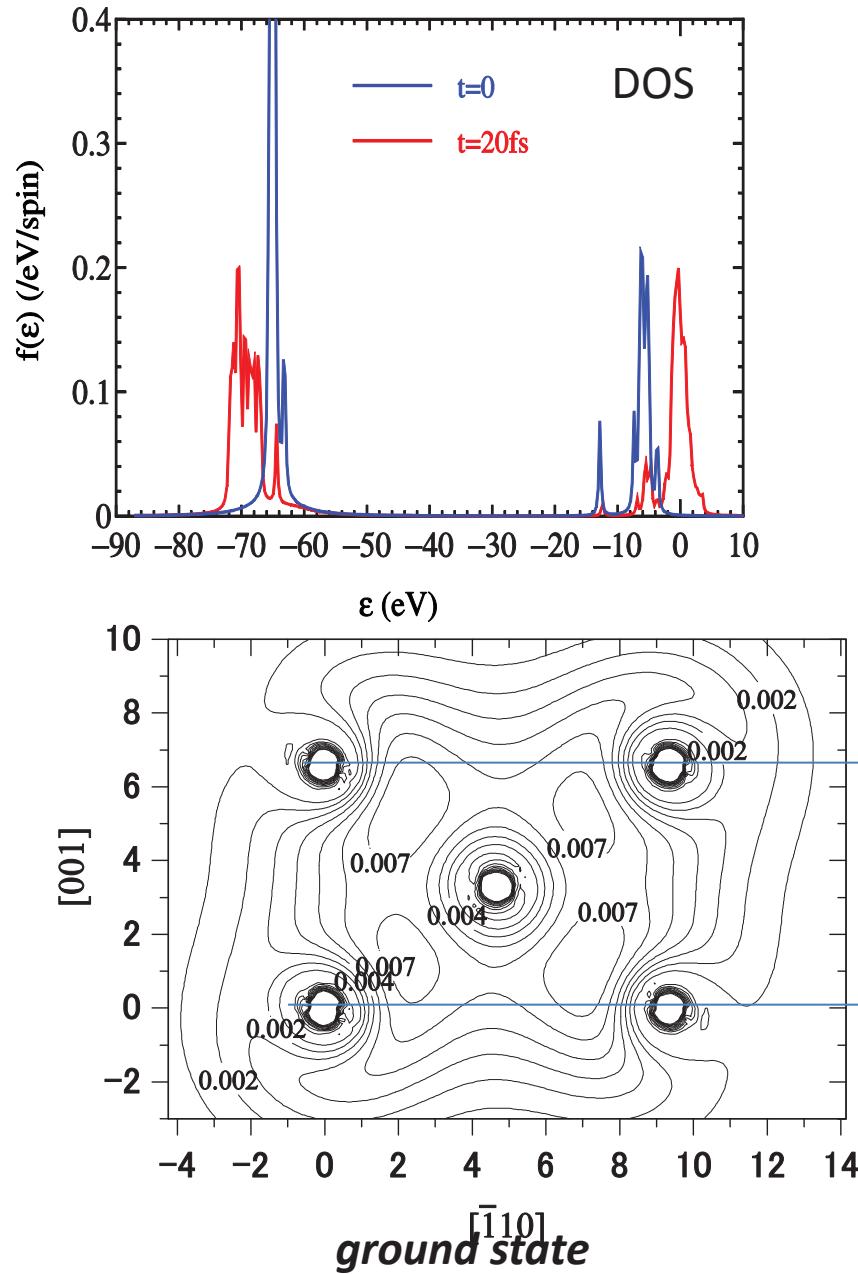
A key point is ionized electrons  
are trapped in the conduction band.

Unrestricted Hartree-Fock calculations



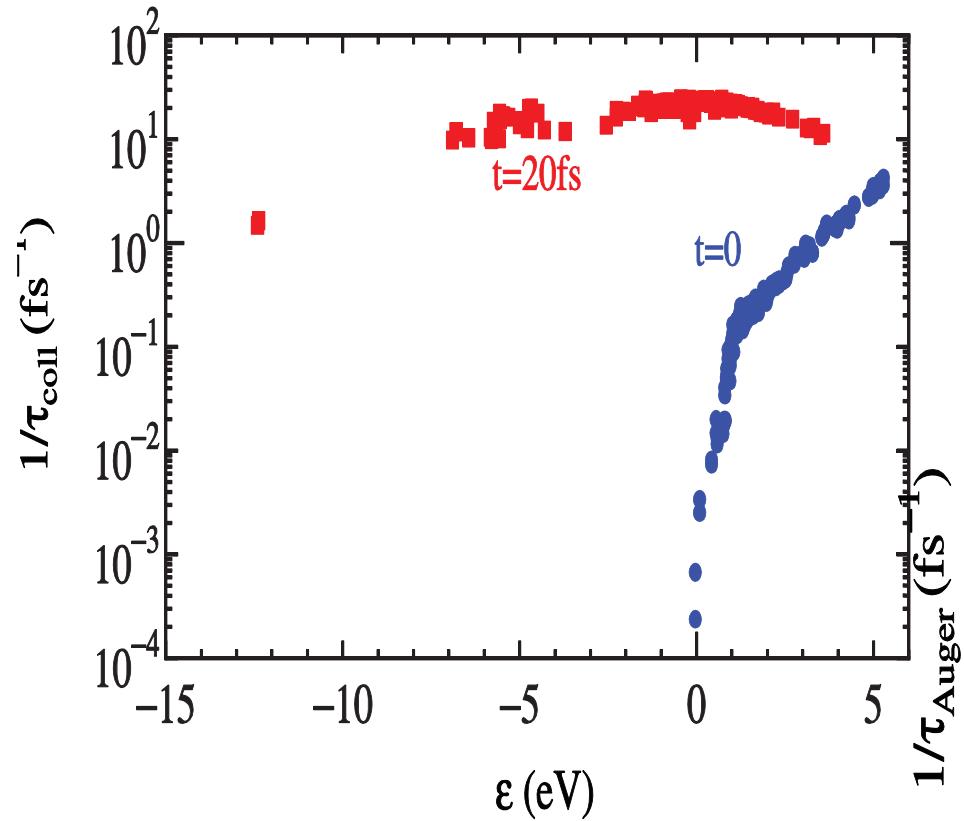
# Energy state after illumination of resonant Xray

Li<sub>24</sub> cluster calculation

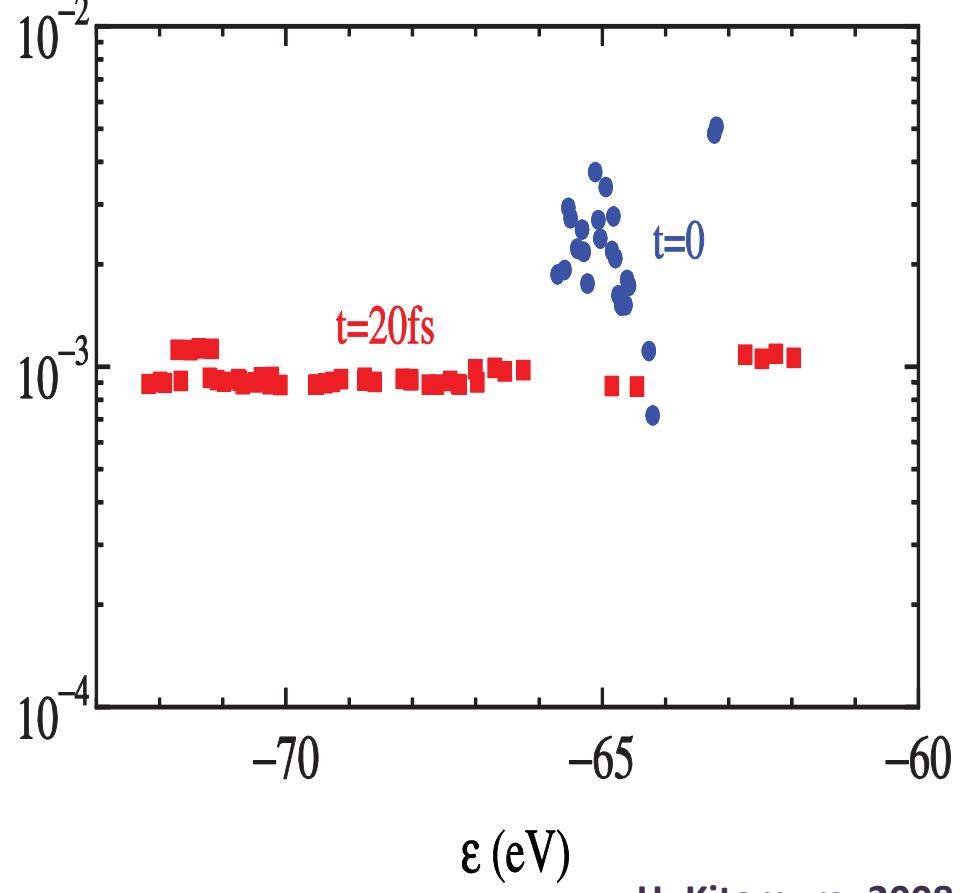


## 24 Li atoms cluster

Electron collision relaxation time



KVV Auger decay rate



H. Kitamura, 2008

## K edge shift due to L vacancy

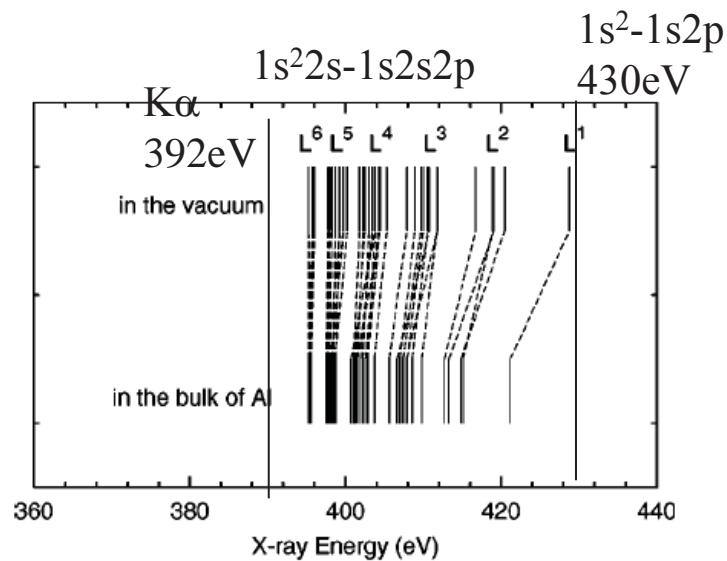


FIG. 3. The x-ray spectra emitted from  $N^{q+}$  hollow atoms and ions in the bulk of Al and in vacuum.

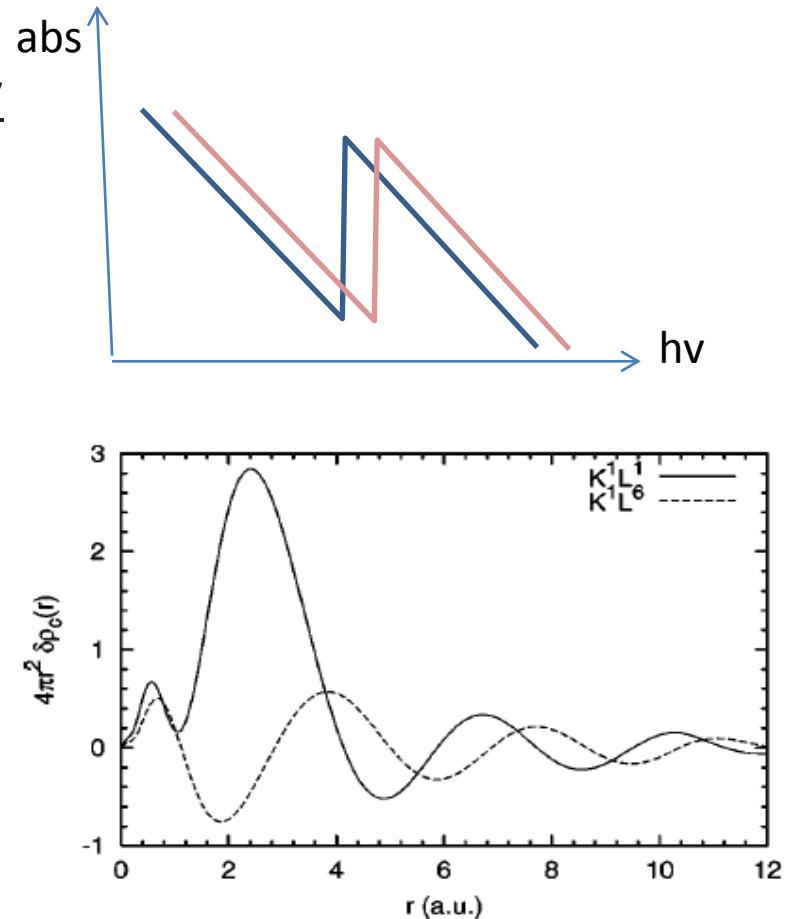
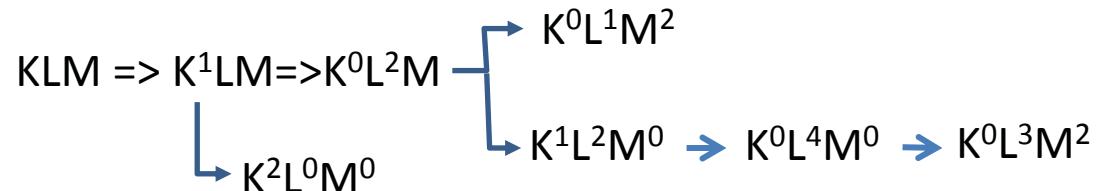
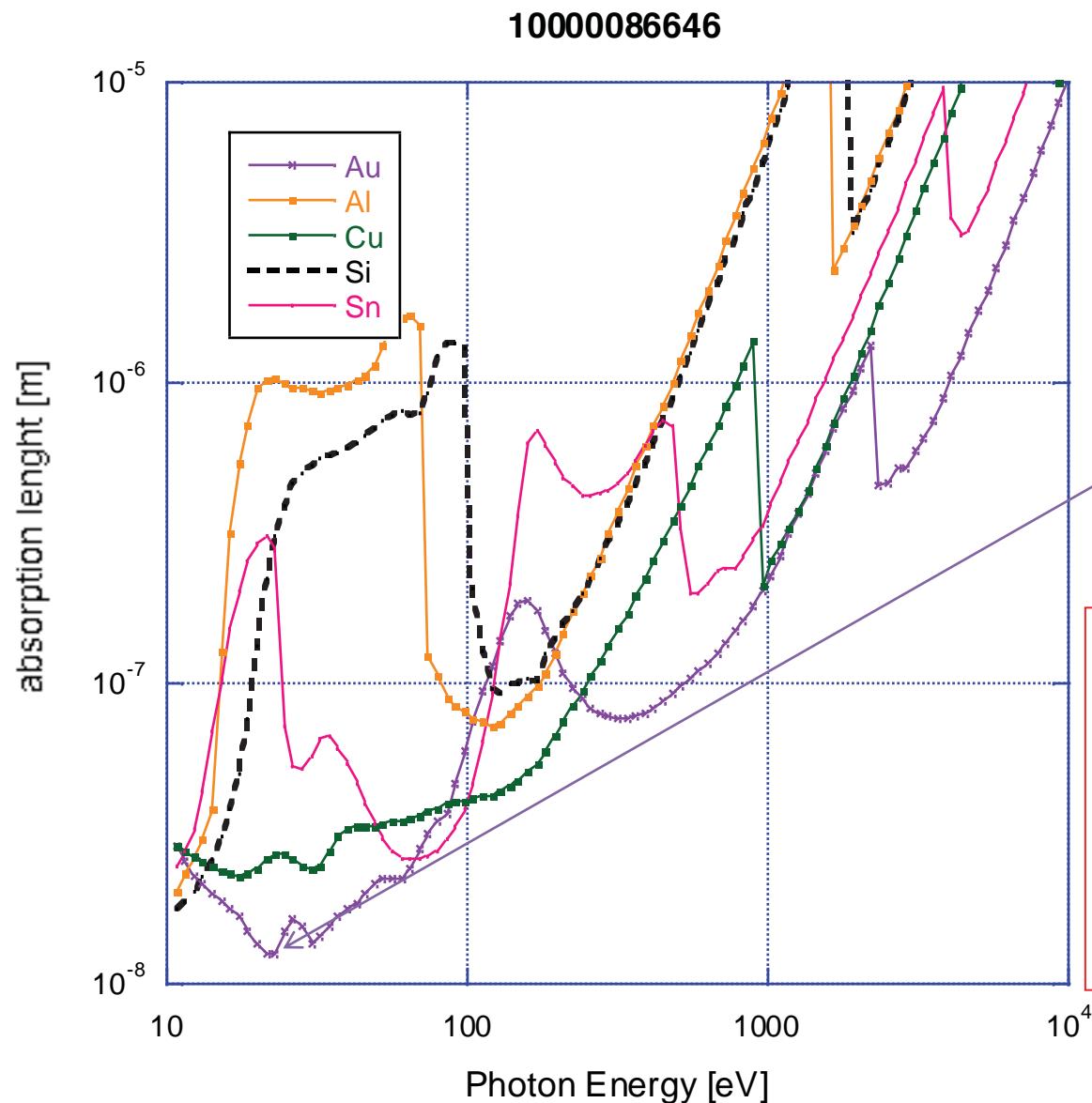


FIG. 2. Conduction electron density changes due to the perturbation of  $N^{q+}$  hollow atoms or ions in the bulk of Al.



# Demonstration experiment with EUV laser

# Absorption length in EUV region



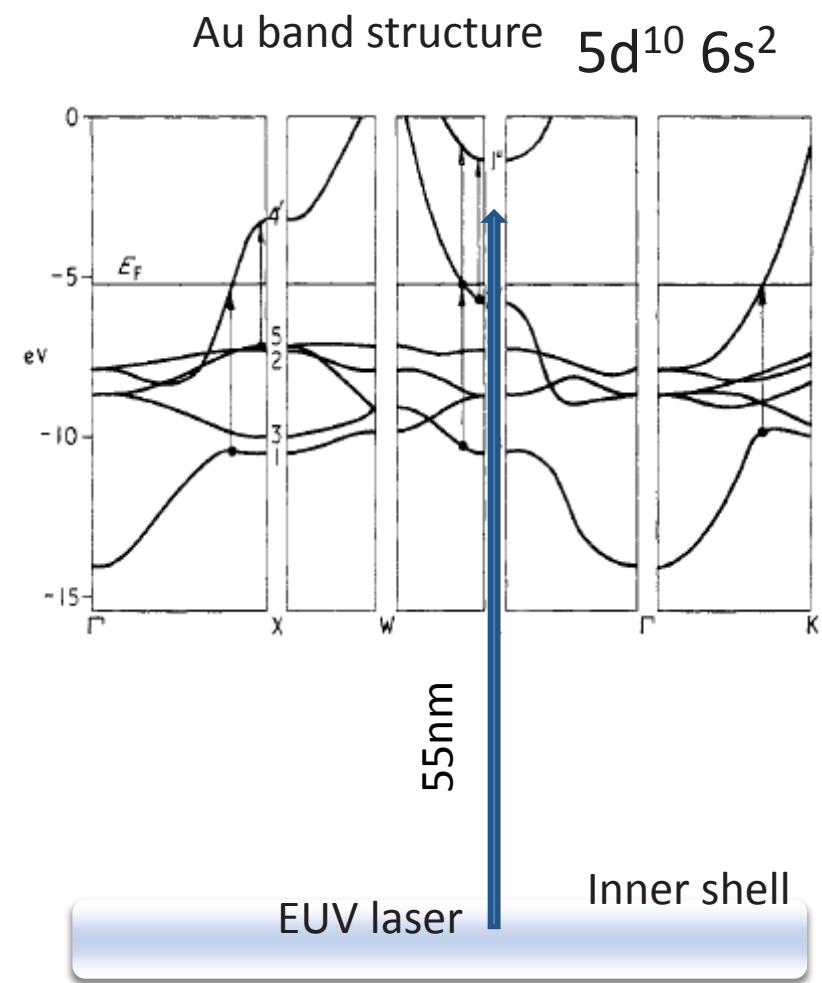
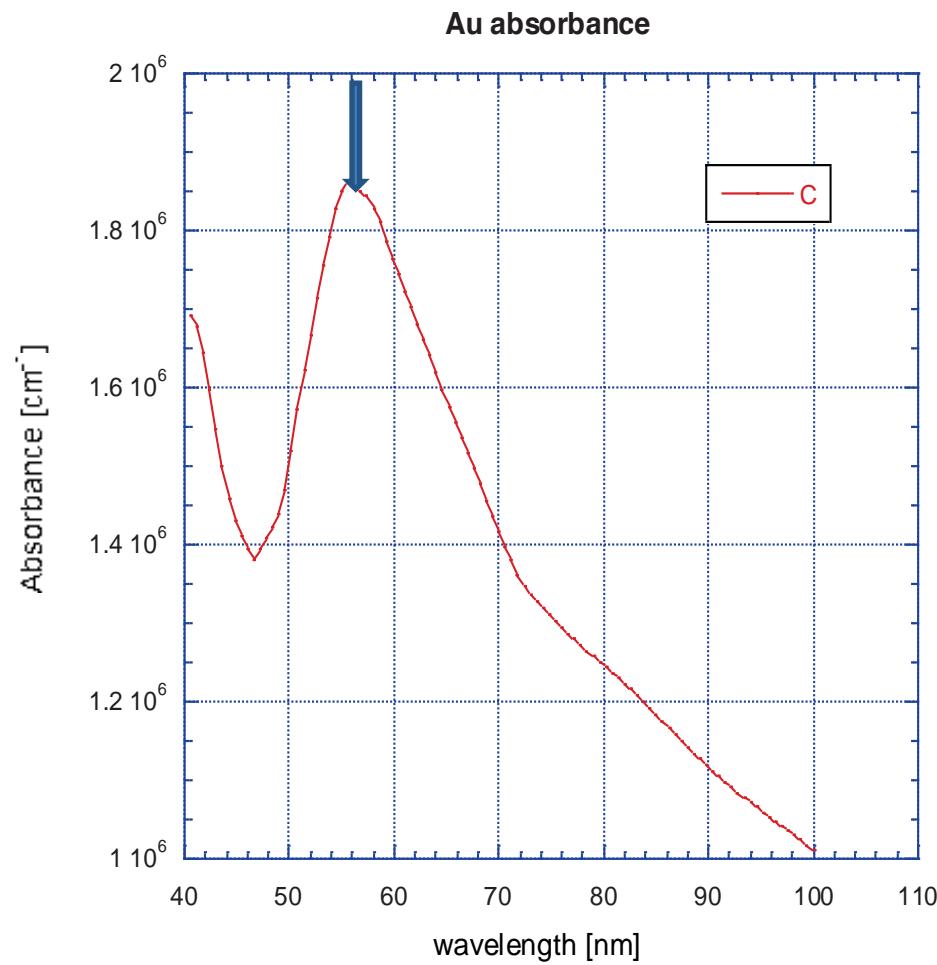
If  $I=10^{13}\text{W/cm}^2$   
Deposition power  
 $\Rightarrow 10^{19}\text{W/cm}^3$

Even within 10fs,  
 $E \sim 10^5\text{J/cm}^3$   
 $= 10^{11}\text{J/m}^3 = 1\text{Mbar}$

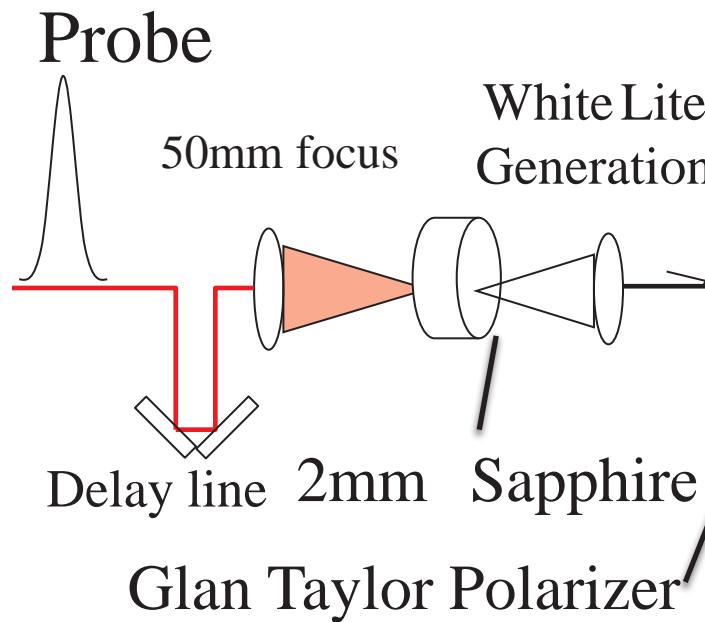
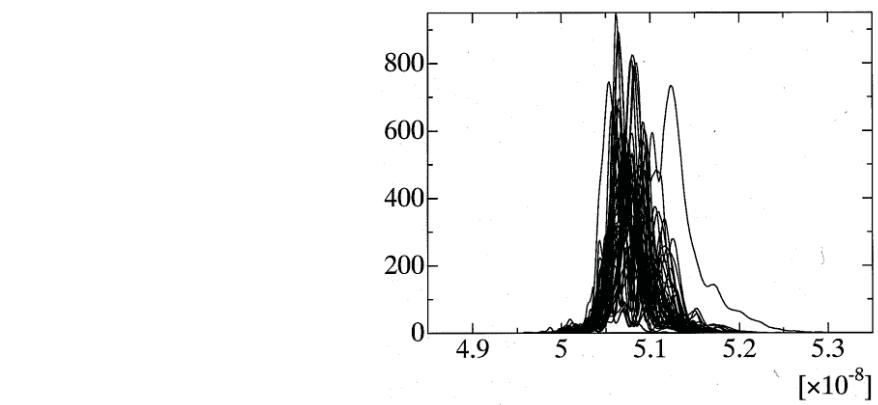


$\rho = 19.3 \text{ g/cc}$ ,  $h\nu = 25\text{eV}$   
 $n_i = 5 \times 10^{22}\text{cm}^{-3}$   
 $\alpha = 5 \times 10^6\text{cm}^{-1}$   
 $\tau_{\text{exp}} \sim 10\text{fs}$   
( $I = 3 \times 10^{13}\text{W/cm}^2$ )

What is excited state condition of electrons in solid with well-tuned high energy photons?



## Experimental system



## EUV FEL

Collection mirror

Pomp pulse Energy monitor

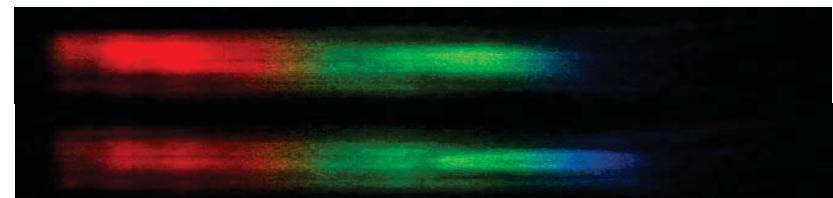
Synchronized

Spectrometer

Slit

CCD

Wollaston Polarizer

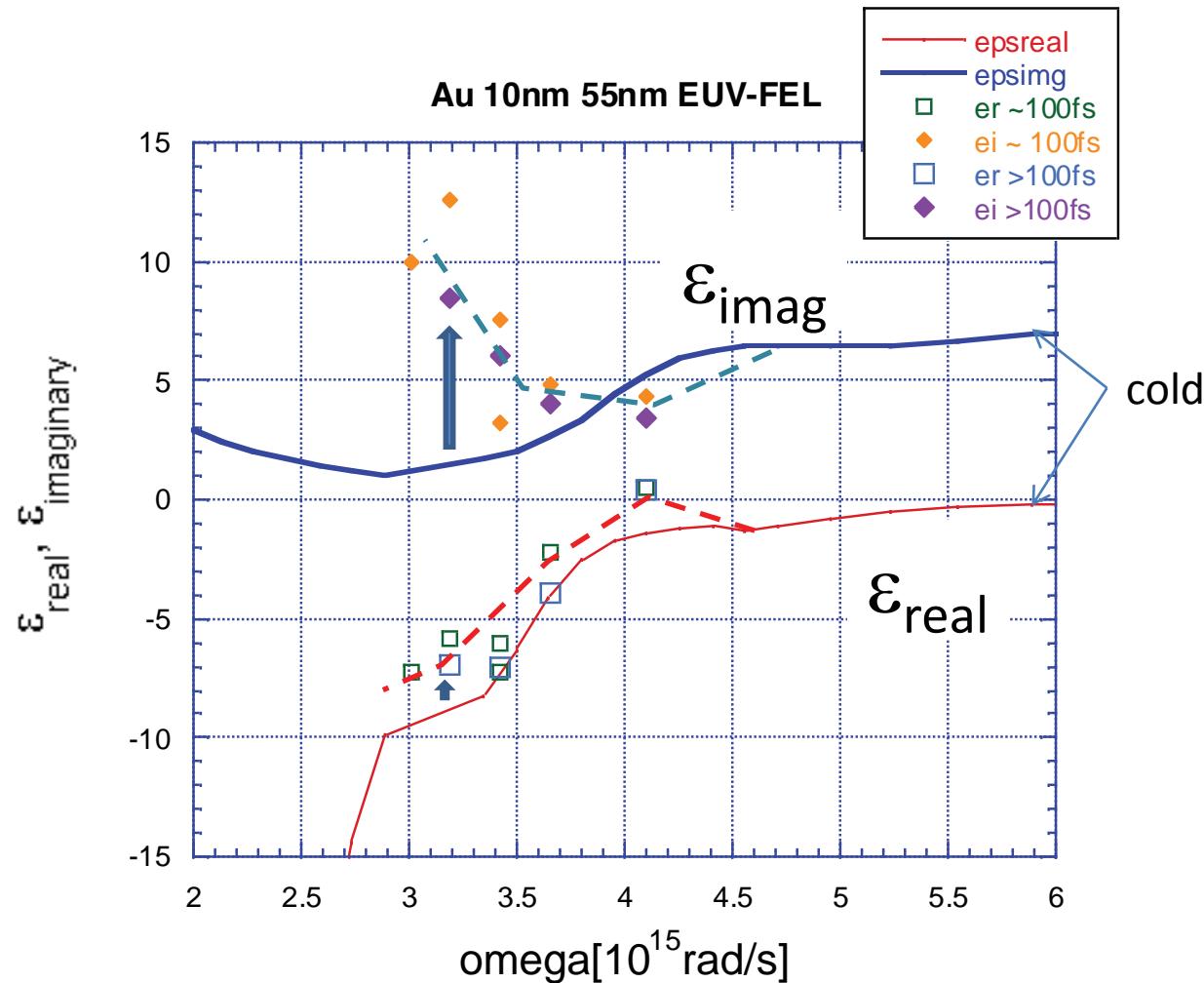


Drude ?

$$\epsilon_1 = n^2 - k^2 = \epsilon_\infty - \frac{\omega_p^2}{\omega^2 + \gamma^2}$$

$$\epsilon_2 = 2nk = \frac{\omega_p^2}{\omega(\omega^2 + \gamma^2)}$$

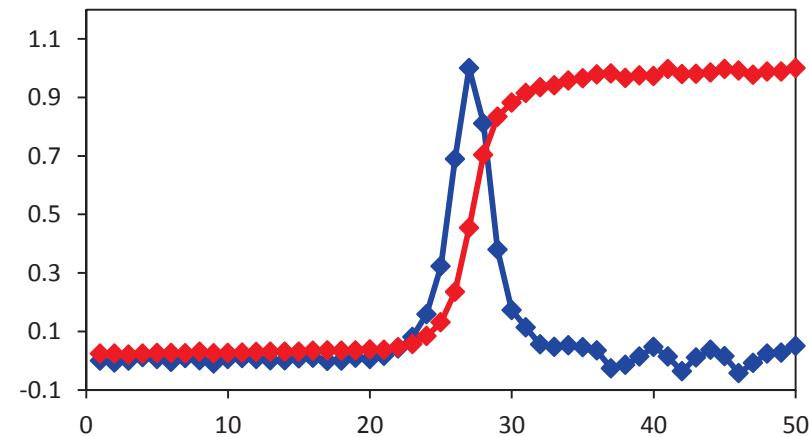
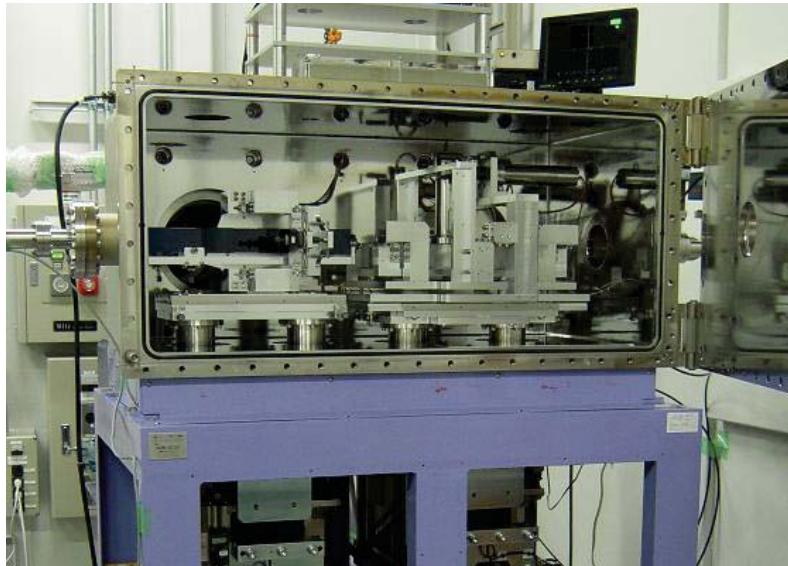
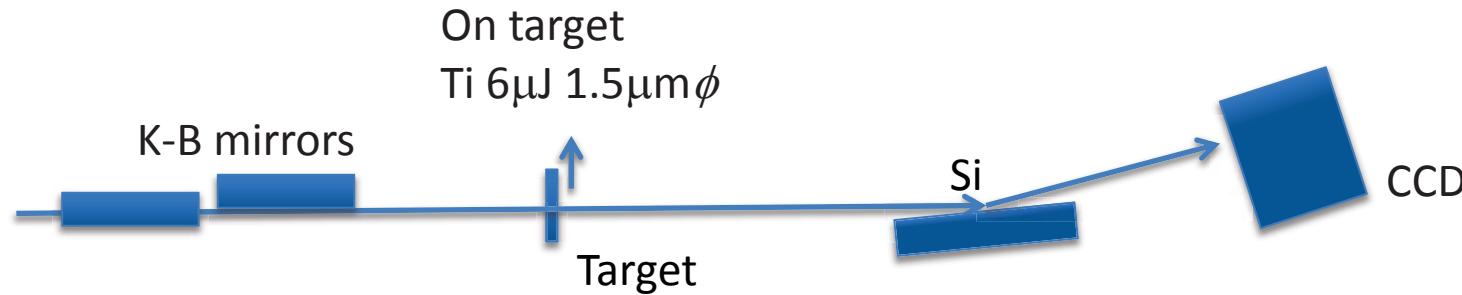
This phenomena is not observed at Ag.



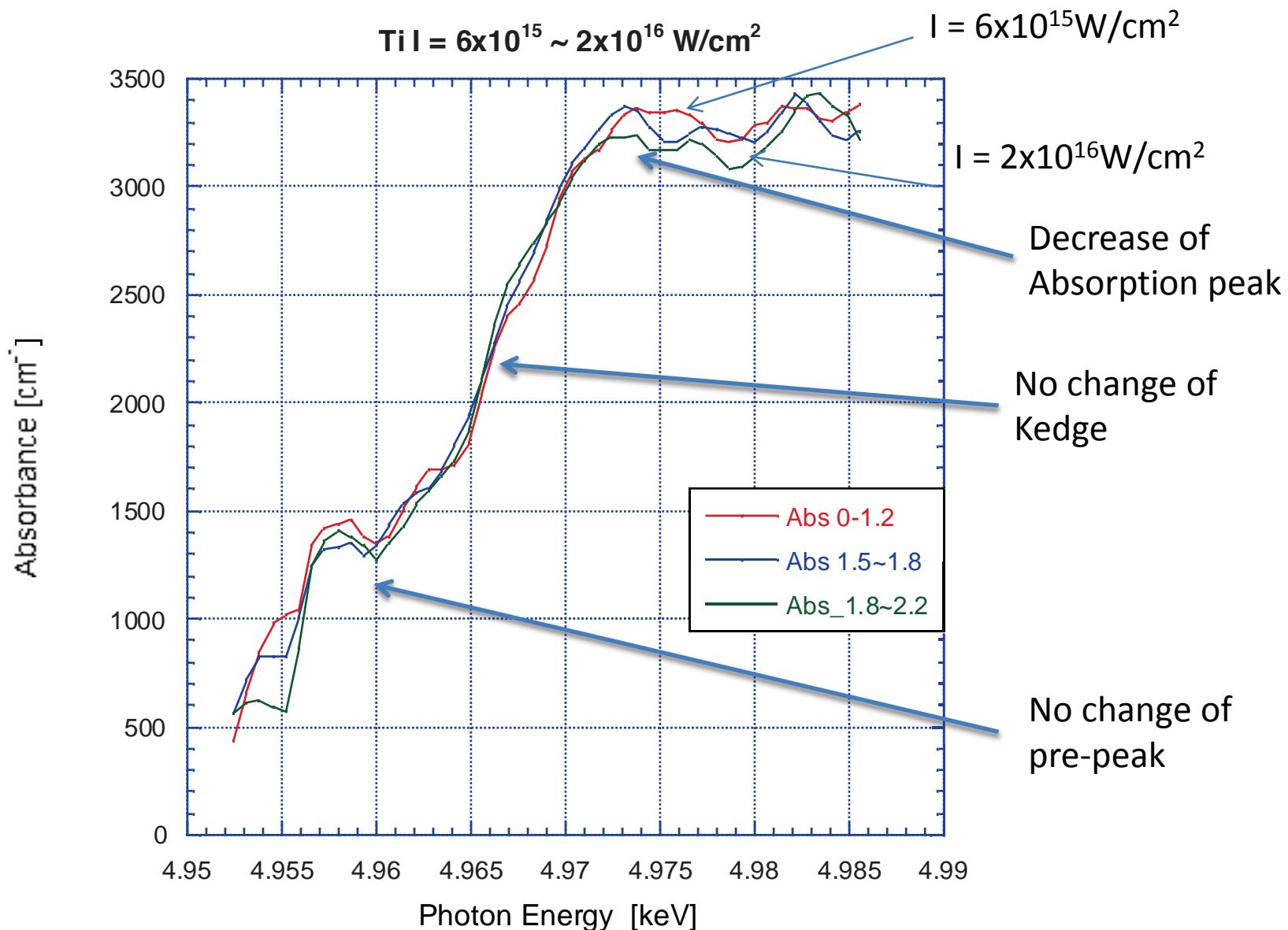
It is NOT free electrons. There are also resonant feature.

Can the photo excitation rate exceed Auger rate in K-shell electrons?

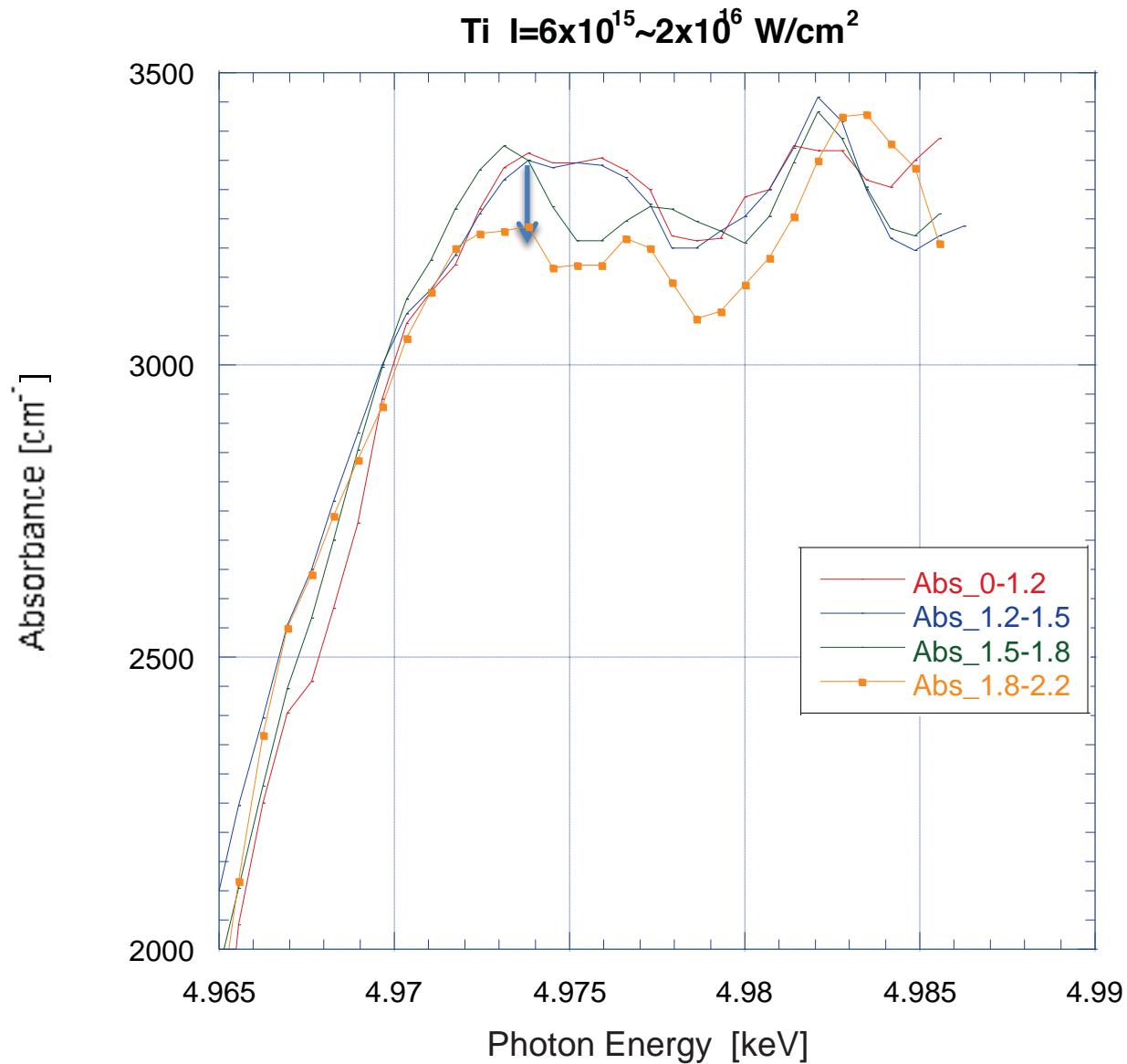
## XFEL SACLA experiments



Osaka U: Yamauchi et al, U Tokyo: Mimura et al  
SACLA/SP8: Ohashi, Yumoto, Koyama, Tono et al

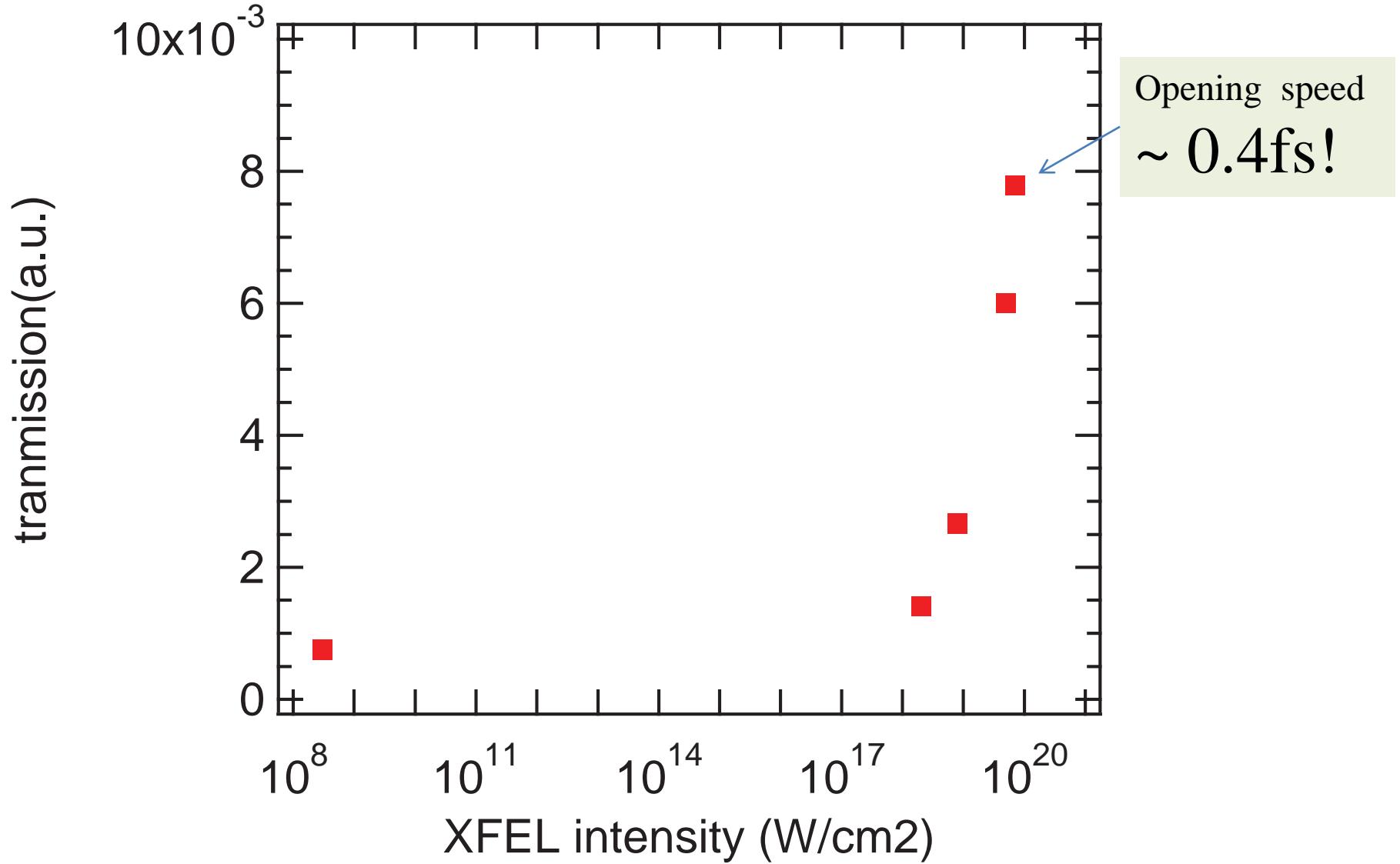


# Transmission spectrum under illumination of intense XFEL

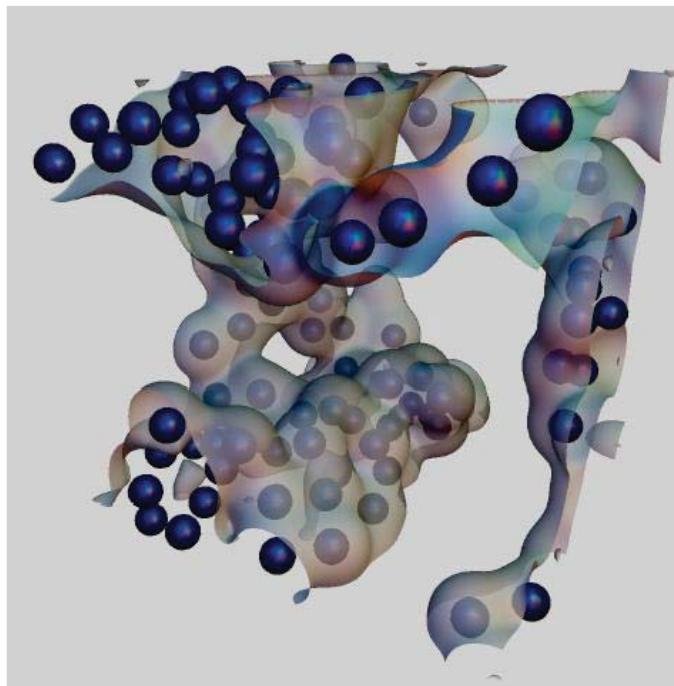


$\alpha$  at Fe K edge  $\sim 3216\text{cm}^{-1}$  at 7keV

Auger time is 0.9fs



We should measure directly the energy density state of electrons, disordering of ions, transport coefficient (electron, diffusion, thermal....), sound velocity, viscosity, and specific heat.



QMD simulation view of wdm Al, Michael Desjarlais, 2008 WDM school