Lorentz-symmetry violation and neutrino physics



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Outline:

- A. Motivations
 - → Possible origins of LV
- B. Low-E effects; EFT
 - \rightarrow SME test framework for ν 's
- C. Phenomenology
 - \rightarrow qualitatively new effects for v's

A. Why test Spacetime Symmetries?

Motivation (i): philosophical necessity

physics is an experimental science

→ solid experimental confirmation of foundations of physics is crucial

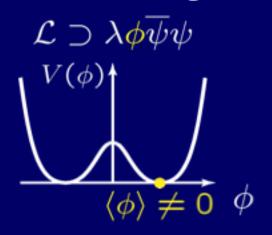
Motivation (ii): discovery potential

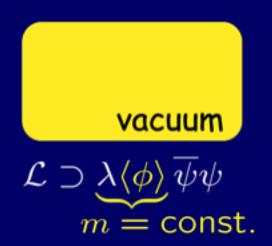
various approaches to physics beyond the Standard Model (e.g., "quantum gravity") can accommodate tiny violations of Relativity

Possible origins for Lorentz/CPT violation

(1) Spontaneous Lorentz breaking in string theory

conventional case: gauge symmet.





string theory: Lorentz symmetry

$$\mathcal{L} \supset \lambda B^{\mu} \overline{\psi} \gamma_5 \gamma_{\mu} \psi$$

$$\downarrow^{V(B^{\mu})} \qquad \downarrow^{B^{\mu}} \qquad \downarrow^{B^{\mu}}$$

 $\mathcal{L}\supset \underline{\lambda\langle B^{\mu}\rangle}\,\overline{\psi}\gamma_{5}\gamma_{\mu}\psi$ $b^{\mu}=\mathrm{const.}$

Kostelecký, Perry, Potting, Samuel '89; '90; '91; '95; '00

(2) Cosmol. varying scalars (e.g., fine-structure parameter)

intuitive argument:



gradient of the scalar <u>selects</u> pref. direction

mathematical argument:

$$\xi = \xi(x) \text{ ... varying coupling} \\ \phi, \Phi \qquad \text{... dynamical fields} \qquad \qquad \mathcal{L} \supset \xi \ \partial^{\mu} \phi \ \partial_{\mu} \Phi$$
 Integration by parts:
$$\mathcal{L}' \supset -(\partial^{\mu} \xi) \ \phi \ \partial_{\mu} \Phi$$
 slow variation of ξ :
$$K^{\mu} \equiv (\partial^{\mu} \xi) \simeq \text{const.} \qquad \qquad \mathcal{L}' \supset -K^{\mu} \phi \ \partial_{\mu} \Phi$$

Kostelecký, R.L., Perry '03; Arkani-Hamed et al. '03

Other mechanisms for Lorentz violation

Noncommutative field theory

$$[\hat{x}^{\mu}, \hat{x}^{\nu}] = i\theta^{\mu\nu}$$

Seiberg-Witten: $\hat{x}^{\mu} \rightarrow$ usual Minkowski coordinates x^{μ}

ightarrow SME terms emerge: $\mathcal{L}_{\mathrm{photon}}\supset \frac{1}{8}\,q\,\theta^{\alpha\beta}F_{\alpha\beta}F^{\mu\nu}F_{\mu\nu}$

e.g. Mocioiu et al. '00; Carroll et al. '01

Topology (in usual 4d spacetime: 1 large compact dim.)

Vacuum fluctuations along this dim.

have periodic boundary conditions

- → preferred direction in vacuum
- ightarrow calculation: $k^{\mu}A^{
 u} ilde{F}_{\mu
 u}\subset\mathcal{L}_{\mathsf{SME}}$

Klinkhamer '00

B. Effective field theory test framework

Why test framework?

- prediction of experimental effects
- analysis and comparison of tests
- theoretical insight

How to obtain test framework?

Purpose: broad experimental searches independent of details of underlying physics

→ construct gen. model compatible with key phys. principles

Ingredients for test framework?

- established physics → all feasible tests can be described
- effective field theory (well established & versatile tool)
- Lorentz/CPT violation (preferred directions; prev. Sec)

Actual construction of the effective field theory (SME)

$$\mathcal{L}_{\text{SME}} = \underbrace{\mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{EH}}}_{\text{EH}} + ek_{\mu}A_{\nu}\tilde{F}^{\mu\nu} + \frac{e}{2\kappa}s^{\mu\nu}R_{\mu\nu} + \dots$$
 present physics

- k^{μ} , $s^{\mu\nu}$, ... coefficients for Lorentz violation
- minimal SME → fermion 44, photon 23, ...
- generated by underlying physics (Sec A)
- amenable to ultrahigh-precision tests (Sec C)

Remarks:

- k^{μ} , $c^{\mu\nu}$, ... coefficients usually taken as spacetime constant
- situations involving gravity require more care
- many theor. SME studies → thus far no inconsistencies

Neutrinos with Lorentz/CPT violation

consider 6-dim. Hilbert space $\left(|e\rangle,|\mu\rangle,|\tau\rangle,|e\rangle^C,|\mu\rangle^C,|\tau\rangle^C\right)$:

q.m. ν propagation (incl. oscillations) governed by 6x6 Hamiltonian

$$H = \underbrace{\begin{pmatrix} |\vec{p}| + \frac{m_l \, m_l^{\dagger}}{2|\vec{p}|} & 0\\ 0 & |\vec{p}| + \frac{m_l^{\dagger} m_l}{2|\vec{p}|} \end{pmatrix}}_{\text{usual LI piece}} + \underbrace{\frac{1}{|\vec{p}|} \begin{pmatrix} \hat{a} - \hat{c} & \hat{H} - \hat{g}\\ \hat{H}^{\dagger} - \hat{g}^{\dagger} & -\hat{a}^T - \hat{c}^T \end{pmatrix}}_{\text{LV \& CPTV correction}}$$

Lorentz-/CPT-violating a, c, g, and H contain:

- (1) preferred directions arising from LV/CPTV background in $\delta \mathcal{L}$
- (2) powers $(p^{\mu})^{(d-3)}$ arising from operators of mass-dimension d
- (3) polarization vectors ε^{μ} arising from the underlying spin d.o.f.

sample structure:
$$\hat{H} \supset \sim \epsilon_{\mu} \epsilon_{\nu}^* H^{(d)\mu\nu}_{\alpha_1...\alpha_{d-3}} p^{\alpha_1} \dots p^{\alpha_{d-3}}$$

Kostelecký and Mewes, arXiv:1112.6395

Reasoning leading to this result:

define
$$\Psi_A = (\psi_e, \psi_\mu, \psi_\tau, \psi_e^C, \psi_\mu^C, \psi_\tau^C)^T, \quad A = 1, \dots, 6$$

most general free, local, unitary, transl.-inv. EFT Lagrangian

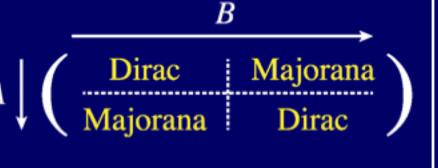
$$\delta \mathcal{L}_{\text{EFT}} = \overline{\Psi}_A (\hat{S}_{AB} + \hat{P}_{AB}\gamma_5 + \hat{V}^{\mu}_{AB}\gamma_{\mu} + \hat{A}^{\mu}_{AB}\gamma_5\gamma_{\mu} + \hat{T}^{\mu\nu}_{AB}\sigma_{\mu\nu})\Psi_B$$

S, P, V, A, T consist of LV/CPTV backgrounds & arbitrary # of ∂ 's

Example:
$$\hat{S}_{AB} \equiv \sum_{d=3}^{\infty} S_{AB}^{(d)}{}^{\alpha_1\alpha_2...\alpha_{d-3}} \partial_{\alpha_1}\partial_{\alpha_2} \ldots \partial_{\alpha_{d-3}}$$

Remarks:

- still some redundancies and LI pieces (removable) $_{A}$
- need to add h.c. for hermitian Lagrangian



assumptions: - LI masses are type-I seesaw compatible - treat LV/CPTV as perturbation

LI unitary transformation #1:

- separate spinors into light Ψ_L and heavy Ψ_R fields
- effect: linear mixture of LV/CPTV terms

heavy Ψ_R don't propagate at low E can be projected out

- may now use 2-component Weyl spinors ξ for remaining Ψ_{L}
- effect: entries of flavor-space H reduce from 4x4 $\gamma \rightarrow 2x2 \sigma$

LI unitary transformation (ultrarel. limit) #2:

- block-diagonalize to disentangle LI Dirac & Majorana m terms
- effects: linear mixture of LV/CPTV terms & p-1 dependence

remove C-conjugation redundancy by projecting onto positive E:

- effect 1: $\xi^+ \sigma^\mu \xi$ introduces polarizations ε^μ
- effect 2: entries of flavor-space H reduce from 2x2 $\sigma \rightarrow \mathbb{C}$

ightarrow prev. displayed most general q.m. v-propagation Hamiltonian in $\left(|e\rangle,|\mu\rangle,| au\rangle,|e\rangle^C,|\mu\rangle^C,| au\rangle^C\right)$ space compatible with EFT:

$$\delta H_{\rm SME}^{\rm propagation} = \frac{1}{|\vec{p}|} \begin{pmatrix} \hat{a} - \hat{c} & \hat{H} - \hat{g} \\ \hat{H}^{\dagger} - \hat{g}^{\dagger} & -\hat{a}^{T} - \hat{c}^{T} \end{pmatrix}$$

where a, c, g, H are lin. combinations of the LV/CPTV S, P, V^{μ} , A^{μ} , $T^{\mu\nu}$ quantities in the EFT Lagrangian

<u>note:</u> v and anti-v coefficients cannot be chosen independently

Ex. 1: no room for mass difference between ν and anti- ν in SME

$$\begin{array}{l} \underline{Ex.\ 2:} \text{ special choice: diagonalizable, isotropic, } \textbf{\textit{c}} = \textbf{0} \\ \rightarrow \textbf{\textit{g}} = \textbf{0}, \textbf{\textit{H}} = \textbf{0}, \text{ and } \textbf{\textit{a}} \text{ takes special form involving } \overset{\boldsymbol{\circ}}{a} \overset{\boldsymbol{\circ}}{e} \overset{\boldsymbol{\circ}}{e} \overset{\boldsymbol{\circ}}{e} \\ \text{neutrino disp. rel.} \qquad \overset{\boldsymbol{\circ}}{E}_{\nu_e} = |\vec{p}| + \frac{m_e}{2|\vec{p}|} + \sum_{d} |\vec{p}|^{d-3} \overset{\boldsymbol{\circ}}{a} \overset{\boldsymbol{\circ}}{e} \overset{\boldsymbol{\circ}}{e} \\ \text{antineutrino disp. rel.} \qquad \overset{\boldsymbol{\circ}}{E}_{\overline{\nu}_e} = |\vec{p}| + \frac{m_e}{2|\vec{p}|} - \sum_{d} |\vec{p}|^{d-3} \overset{\boldsymbol{\circ}}{a} \overset{\boldsymbol{\circ}}{e} \overset{\boldsymbol{\circ}}{e} \end{array}$$

C. Phenomenology and Tests

Novel qualitative features in neutrino oscillations

eff. Hamiltonian for 3 left-handed v's in flavor basis $(a,b \dots e,\mu,\tau)$: $p_{\mu} \dots \nu$ propagation direction \rightarrow directional dependence $(H + \delta H)_{ab} = \frac{1}{2|\vec{p}|} (m_l \, m_l^{\dagger})_{ab} + (a^{\mu})_{ab} \, \hat{p}_{\mu} - (c^{\mu\nu})_{ab} \, \hat{p}_{\mu} \hat{p}_{\nu} E + \dots$ $(a^{\mu})_{ab} \xrightarrow{\bar{\nu}} -(a^{\mu})^*_{ab}$ unconventional E dependence → CPT violation $\rightarrow U = U(E)$

→ great flexibility for matching observations

models based on this general effective Hamiltonian:

- "bicycle model" (Kostelecký et al., 2004)
- "tandem model" (Katori et al., 2006)
- "BMW model" (Barger et al., 2007)
- "puma model" (Díaz et al., 2010)
- "isotropic bicycle model" (Barger et al., 2011) ...

I	phenomenology: example "puma model"			
ı		3 <i>v</i> SM	"puma model"	
	established results (with θ_{13} = 0)	needs 5 parameters	needs only 2 parameters	
	LSND anomaly MiniBooNE low- <i>E</i> excess	its 6 parameters are insufficient	needs 6 parameters	

most recent exp. results: arXiv:0801.0287v5, January 2012 edition

- 12 pages of tables for d = 2, 3, 4, 5, 6, and d > 6
- analyses by MINOS, IceCube, MiniBooNE, LSND, ...
- sample section of table:

Table XIII. Neutrino sector, d = 3

Combination	Result	System	Ref.
å(3)	$<1.9\times10^{-7}~{\rm GeV}$	IceCube meson threshold	[13]*
$ \operatorname{Re}(a_L)_{\mu\tau}^X $	$<5.9\times10^{-23}~{\rm GeV}$	MINOS	[103]
$ \operatorname{Im}(a_L)_{u\tau}^X $	$<2.2\times10^{-20}~{\rm GeV}$	20	[103]
$ \operatorname{Re}(a_L)_{\mu\tau}^Y $	$<6.1\times10^{-23}~{\rm GeV}$	22	[103]
$ { m Im}(a_L)_{\mu au}^Y $	$<2.2\times10^{-20}~{\rm GeV}$	29	[103]
a_L^X, a_L^Y	$<1.8\times10^{-23}~{\rm GeV}$	IceCube	[104]
$ (a_L)_{u\tau}^X $	$<5.9\times10^{-23}~{\rm GeV}$	MINOS FD	[105]
$ (a_L)_{u au}^Y $	$<6.1\times10^{-23}~{\rm GeV}$	22	[105]
$ a_L^X , a_L^Y $	$<3.0\times10^{-20}~{\rm GeV}$	MINOS ND	[106]
$ \text{Re}(a_{\text{eff}}^{(3)})_{00}^{e\mu} $	$< 1.3 \times 10^{-18} \text{ GeV}$	LSND	[13]*
»	$<1.5\times10^{-19}~{\rm GeV}$	MiniBooNE	[13]*
27	$<9.2\times10^{-20}~{\rm GeV}$	MiniBooNE $\overline{\nu}$	[13]*
12 (8) seu	0.0 40-18 0 77	T COLID	[40]#

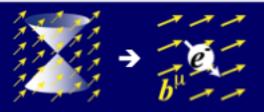
Summary

presently no compelling exp. evidence for Relativity violations, but:

(1) various theoretical approaches to quantum gravity can cause such violations



(2) at low E, such violations are described by SME test framework (eff. field theory + background fields)



- (3) for free v's, all Lorentz-/CPT-violating operators compatible with local, unitary EFT have been classified
- (4) Lorentz/CPT violation predicts a wealth of qualitatively novel features; provides ample ground for phenomenological model building