

### 3 Correlation functions in $\mathcal{N} = 4$ SYM

We shall consider the correlation function

$$G(1, \dots, n) = \langle \mathcal{O}(x_1, Y_1) \dots \mathcal{O}(x_n, Y_n) \rangle \quad (3.1)$$

of the operators  $\mathcal{O}(x_i, Y_i)$  built from six real scalar fields  $\Phi_{ij}^I$  (with  $I = 1, \dots, 6$ ) in the  $U(N_c)$  gauge group ( $i, j = 1, \dots, N_c^2$ )

$$\mathcal{O}(x, Y) = Y_I Y_J \text{tr} [\Phi^I(x) \Phi^J(x)] \quad (3.2)$$

Here  $Y_I$  is an auxiliary 6-dimensional null vector ( $Y^2 \equiv Y_I Y_I = 0$ ) and the scalar field is normalized as

$$\langle \Phi_{i_1 j_1}^{I_1}(x_1) \Phi_{i_2 j_2}^{I_2}(x_2) \rangle = \frac{1}{2} \delta^{I_1 I_2} \delta_{i_1 j_2} \delta_{j_1 i_2} \frac{1}{x_{12}^2} \quad (3.3)$$

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**Problem 1:** Write a Mathematica code to compute 2-, 3- and 4-point correlation functions at Born level

$$\begin{aligned} G^{(0)}(1, 2) &= \frac{1}{2} N_c^2 \frac{(y_{12}^2)^2}{(x_{12}^2)^2} \\ G^{(0)}(1, 2, 3) &= N_c^2 \frac{y_{12}^2 y_{23}^2 y_{31}^2}{x_{12}^2 x_{23}^2 x_{31}^2} \\ G^{(0)}(1, 2, 3, 4) &= \frac{N_c^4}{4} \left[ \left( \frac{y_{12}^2 y_{34}^2}{x_{12}^2 x_{34}^2} \right)^2 + \left( \frac{y_{13}^2 y_{24}^2}{x_{13}^2 x_{24}^2} \right)^2 + \left( \frac{y_{41}^2 y_{23}^2}{x_{41}^2 x_{23}^2} \right)^2 \right] \\ &\quad + N_c^2 \left( \frac{y_{12}^2 y_{23}^2 y_{34}^2 y_{41}^2}{x_{12}^2 x_{23}^2 x_{34}^2 x_{41}^2} + \frac{y_{12}^2 y_{24}^2 y_{34}^2 y_{13}^2}{x_{12}^2 x_{24}^2 x_{34}^2 x_{13}^2} + \frac{y_{13}^2 y_{23}^2 y_{24}^2 y_{41}^2}{x_{13}^2 x_{23}^2 x_{24}^2 x_{41}^2} \right) \end{aligned} \quad (3.4)$$

Here the notation is introduced for  $x_{ij}^2 = (x_i - x_j)^2$  and  $y_{ij}^2 = (Y_i \cdot Y_j)$ .

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The superconformal symmetry of the  $\mathcal{N} = 4$  SYM theory impose tight restrictions on the possible form of the loop corrections to the correlation function (3.1).  $G(1, 2)$  and  $G(1, 2, 3)$  are protected from loop corrections. The loop corrections to  $G(1, 2, 3, 4)$  take the following factorized form (the so-called ‘‘partial non-renormalization’’):

$$G_4 = G^{(0)}(1, 2, 3, 4) + R(1, 2, 3, 4) \times F(x_i; a) \quad (\text{for } a = g_{\text{YM}}^2 N_c / (8\pi^2)), \quad (3.5)$$

where  $R(1, 2, 3, 4)$  is a universal rational function of  $x_i$ - and  $y_i$ -coordinates at the four external points 1, 2, 3, 4:

$$\begin{aligned} R(1, 2, 3, 4) &= 2N_c^2 \left[ \frac{y_{13}^4 y_{24}^4}{x_{13}^2 x_{24}^2} + \frac{y_{14}^4 y_{23}^4}{x_{14}^2 x_{23}^2} + \frac{y_{12}^4 y_{34}^4}{x_{12}^2 x_{34}^2} + \frac{y_{12}^2 y_{23}^2 y_{34}^2 y_{14}^2}{x_{12}^2 x_{23}^2 x_{34}^2 x_{14}^2} (x_{13}^2 x_{24}^2 - x_{12}^2 x_{34}^2 - x_{14}^2 x_{23}^2) \right. \\ &\quad \left. + \frac{y_{12}^2 y_{13}^2 y_{24}^2 y_{34}^2}{x_{12}^2 x_{13}^2 x_{24}^2 x_{34}^2} (x_{14}^2 x_{23}^2 - x_{12}^2 x_{34}^2 - x_{13}^2 x_{24}^2) + \frac{y_{13}^2 y_{14}^2 y_{23}^2 y_{24}^2}{x_{13}^2 x_{14}^2 x_{23}^2 x_{24}^2} (x_{12}^2 x_{34}^2 - x_{14}^2 x_{23}^2 - x_{13}^2 x_{24}^2) \right] \end{aligned} \quad (3.6)$$

According to (3.5), the loop corrections to the four-point correlation function are determined by a single function with the perturbative expansion

$$F(x_i; a) = a F^{(1)}(x_i) + a^2 F^{(2)}(x_i) + O(a^3) \quad (3.7)$$

The goal of the next exercise is to determine the lowest order correction  $F^{(1)}(x_i)$

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**Problem 2:** Consider the correlation function  $G(1, 2, 3, 4)$ , Eq. (3.1), for the  $Y$ -variables of the form

$$Y_1 = (1, i, 0, 0, 0, 0), \quad Y_2 = (0, 0, 1, i, 0, 0), \quad Y_3 = Y_1^*, \quad Y_4 = Y_2^* \quad (3.8)$$

corresponding to the following choice of the operators (with  $Z = \Phi^1 + i\Phi^2$  and  $W = \Phi^3 + i\Phi^4$ )

$$O(1) = \text{tr}[Z^2(x_1)], \quad O(2) = \text{tr}[W^2(x_2)], \quad O(3) = \text{tr}[\overline{Z}^2(x_3)], \quad O(4) = \text{tr}[\overline{W}^2(x_4)] \quad (3.9)$$

Examine Feynman diagrams contributing to the correlation function to order  $a$  (with the interaction of the form  $\sim a \text{tr}([\Phi^I, \Phi^J][\Phi^I, \Phi^J])$ ) and compare them with the general expression for  $G(1, 2, 3, 4)$ , Eqs. (3.5) and (3.6),

$$G(1, 2, 3, 4) = \begin{array}{c} 1 \\ \diagup \\ \text{---} \\ \diagdown \\ 3 \end{array} + \begin{array}{c} 2 \\ \diagup \\ \text{---} \\ \diagdown \\ 4 \end{array} + \begin{array}{c} 1 & 2 \\ \diagdown & \diagup \\ & 5 \\ \diagup & \diagdown \\ 3 & 4 \end{array} = \frac{16N_c^2}{(x_{13}^2 x_{24}^2)^2} \left[ \frac{N_c^2}{4} + 2a x_{13}^2 x_{24}^2 F^{(1)}(x_i) + O(a^2) \right] \quad (3.10)$$

to obtain

$$F^{(1)}(x_i) = -\frac{1}{4\pi^2} \int \frac{d^4 x_5}{x_{15}^2 x_{25}^2 x_{35}^2 x_{45}^2} \quad (3.11)$$

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To any order  $\ell$  in the coupling, the function  $F^{(\ell)}(x_i)$  has the following properties. It admits the representation

$$F^{(\ell)}(x_i) = \frac{x_{12}^2 x_{13}^2 x_{14}^2 x_{23}^2 x_{24}^2 x_{34}^2}{\ell! (-4\pi^2)^\ell} \int d^4 x_5 \dots d^4 x_{4+\ell} f^{(\ell)}(x_1, \dots, x_{4+\ell}) \quad (3.12)$$

where  $f^{(\ell)}$  is a rational function depending on the four external coordinates  $x_1, \dots, x_4$  and the  $\ell$  additional (internal) coordinates  $x_5, \dots, x_{4+\ell}$  giving the positions of the interaction vertices. It has the following form

$$f^{(\ell)}(x_1, \dots, x_{4+\ell}) = \frac{P^{(\ell)}(x_1, \dots, x_{4+\ell})}{\prod_{1 \leq i < j \leq 4+\ell} x_{ij}^2} \quad (3.13)$$

where the denominator contains the product of all distances between the  $(4 + \ell)$  points. The function  $F^{(\ell)}(x_i)$  should satisfy the following two conditions:

- It should transform under inversions  $I[x_i^\mu] = x_i^\mu/x_i^2$  as

$$I[F^{(\ell)}(x_1, x_2, x_3, x_4)] = x_1^2 x_2^2 x_3^2 x_4^2 F^{(\ell)}(x_1, x_2, x_3, x_4) \quad (3.14)$$

- Its *integrand*  $f^{(\ell)}(x_1, \dots, x_{4+\ell})$  should be symmetric under the exchange of any pair of points (both external and internal)

$$f^{(\ell)}(\dots, x_i, \dots, x_j, \dots) = f^{(\ell)}(\dots, x_j, \dots, x_i, \dots) \quad (3.15)$$

**Problem 3:**

Show that the above two conditions imply that

- (a)  $P^{(\ell)}$  should also be invariant under the  $S_{4+\ell}$  permutations of  $x_1, \dots, x_{4+\ell}$
- (b)  $P^{(\ell)}$  should transform under inversions as

$$I[P^{(\ell)}(x_1, \dots, x_{4+\ell})] = (x_1^2 \dots x_{4+\ell}^2)^{-\ell+1} P^{(\ell)}(x_1, \dots, x_{4+\ell}) \quad (3.16)$$

- (c)  $P^{(\ell)}$  should be a homogeneous polynomial in  $x_{ij}^2$  of degree  $(\ell - 1)(\ell + 4)/2$

Verify that at one ( $\ell = 1$ ) and two ( $\ell = 2$ ) loops the most general expression for  $P^{(\ell)}$  is

$$P^{(1)} = 1, \quad P^{(2)} = \frac{c^{(2)}}{48} \sum_{\sigma \in S_6} x_{\sigma_1 \sigma_2}^2 x_{\sigma_3 \sigma_4}^2 x_{\sigma_5 \sigma_6}^2 = c^{(2)} [x_{12}^2 x_{34}^2 x_{56}^2 + \dots] \quad (3.17)$$

with  $c^{(2)}$  arbitrary constant.

Substitute (3.17) into (3.12) and (3.13) and obtain the following result

$$\begin{aligned} F^{(1)} &= g(1, 2, 3, 4), \\ F^{(2)}/c^{(2)} &= h(1, 2; 3, 4) + h(3, 4; 1, 2) + h(2, 3; 1, 4) + h(1, 4; 2, 3) \\ &\quad + h(1, 3; 2, 4) + h(2, 4; 1, 3) + \frac{1}{2} (x_{12}^2 x_{34}^2 + x_{13}^2 x_{24}^2 + x_{14}^2 x_{23}^2) [g(1, 2, 3, 4)]^2, \end{aligned} \quad (3.18)$$

where  $g$ - and  $h$ - integrals have the form

$$\begin{aligned} g(1, 2, 3, 4) &= -\frac{1}{4\pi^2} \int \frac{d^4 x_5}{x_{15}^2 x_{25}^2 x_{35}^2 x_{45}^2}, \\ h(1, 2; 3, 4) &= \frac{x_{34}^2}{(4\pi^2)^2} \int \frac{d^4 x_5 d^4 x_6}{(x_{15}^2 x_{35}^2 x_{45}^2) x_{56}^2 (x_{26}^2 x_{36}^2 x_{46}^2)}. \end{aligned} \quad (3.19)$$

The correlation function  $G(1, 2, 3, 4)$  is related in planar  $\mathcal{N} = 4$  SYM to (the square of) the scattering amplitude  $A_4$  in the light-cone limit  $x_{12}^2, x_{23}^2, x_{34}^2, x_{41}^2 \rightarrow 0$  through the duality relation

$$\lim_{x_{i,i+1}^2 \rightarrow 0} \frac{G(1, 2, 3, 4)}{G^{(0)}(1, 2, 3, 4)} \sim [A(p_1, p_2, p_3, p_4)]^2 \Big|_{p_i = x_i - x_{i+1}}, \quad (3.20)$$

or equivalently

$$\lim_{x_{i,i+1}^2 \rightarrow 0} \left( 1 + 2x_{13}^2 x_{24}^2 \sum_{\ell} a^{\ell} F^{(\ell)} \right) = \left( 1 + \sum_{\ell} a^{\ell} M^{(\ell)} \right)^2. \quad (3.21)$$

Here  $M^{(\ell)}$  defines  $\ell$ -loop correction to *planar* four-gluon scattering amplitude  $A_4 \sim 1 + \sum_{\ell} a^{\ell} M^{(\ell)}$  and it is given by the sum of *connected* (nonfactorizable) planar integrals.

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**Problem 4:**

- (a) Use the duality relation (3.21) together with (3.18) to show that the condition of absence of (disconnected)  $[g(1, 2, 3, 4)]^2$  integral inside  $M^{(2)}$  fixes uniquely the normalization coefficient in (3.18)

$$c^{(2)} = 1 \quad (3.22)$$

- (b) Use the duality relation (3.21) to compute two-loop corrections to the 4-particle scattering amplitude

$$\begin{aligned} M^{(1)} &= x_{13}^2 x_{24}^2 g(1, 2, 3, 4), \\ M^{(2)} &= x_{13}^2 x_{24}^2 [h(1, 3; 2, 4) + h(2, 4; 1, 3)]. \end{aligned} \quad (3.23)$$


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