



The Abdus Salam
International Centre
for Theoretical Physics



2451-2

Workshop on Interferometry and Interactions in Non-equilibrium Meso- and Nano-systems

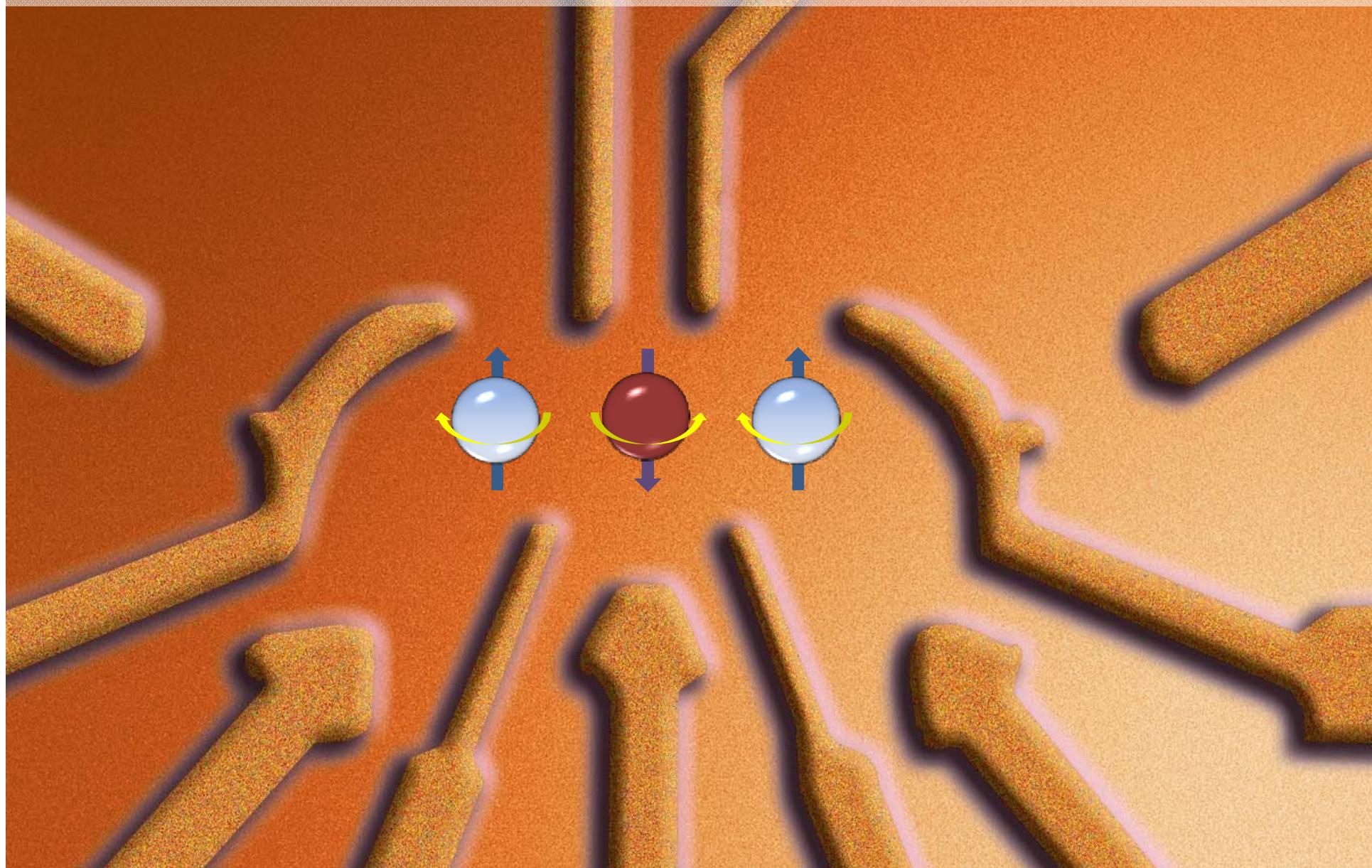
8 - 12 April 2013

Quantum Interference Phenomena in Triple Quantum Dot Circuits

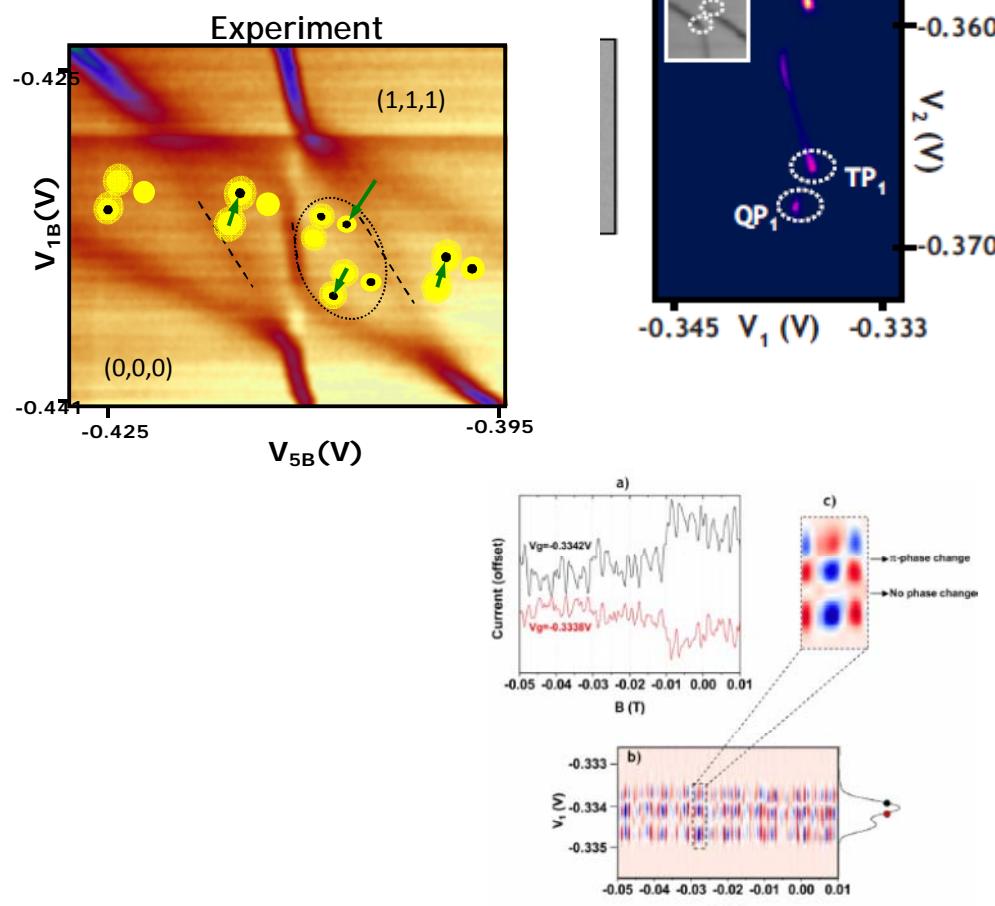
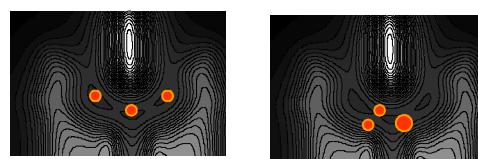
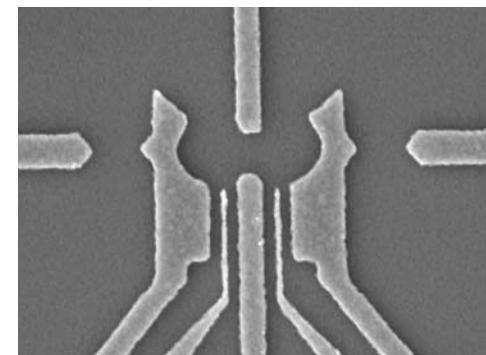
Andy SACHRAJDA
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CANADA*

Quantum Interference Phenomena in Triple Quantum Dot Circuits

Andy Sachrajda, NRC Canada

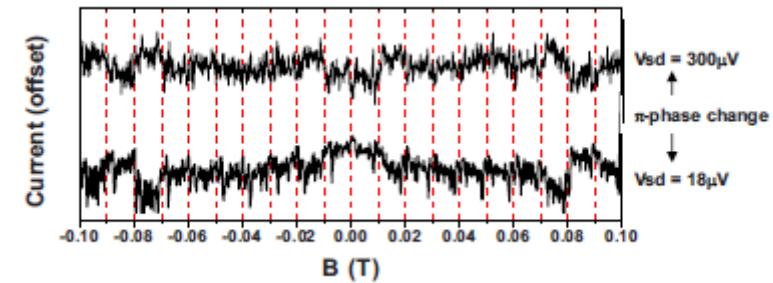
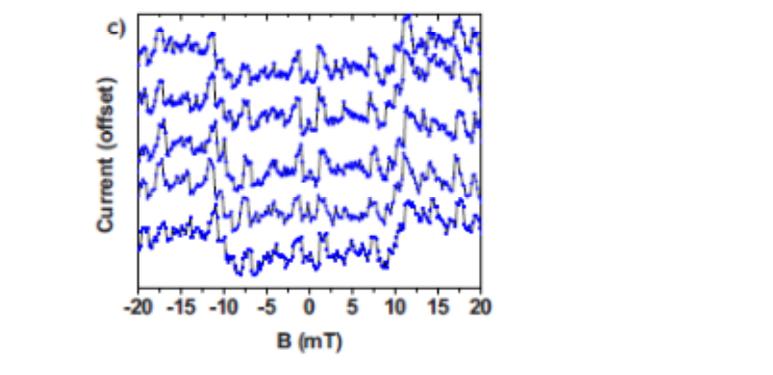
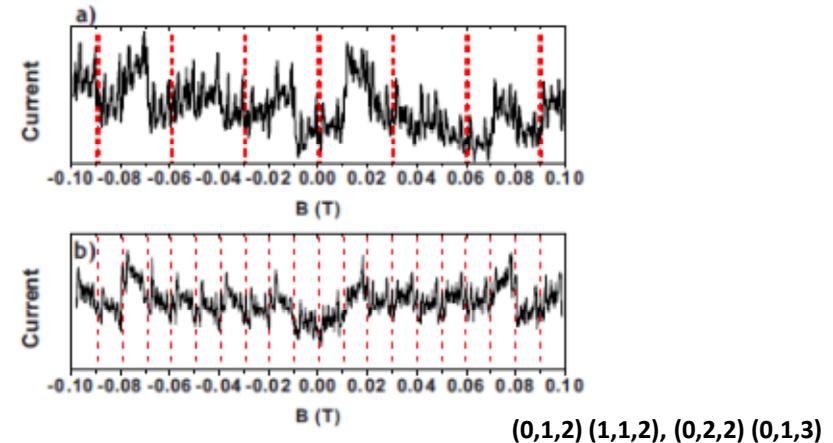


Transport Through an ‘incidental’ Ring of Quantum Dots



L. Gaudreau *et al.* Phys. Rev. Lett. 97, 036807 (2006)
and PRB 80 075415 (2009)

D.Schroer et al. PRB 76 075306 (2007)



Outline

*(1) Quantum Backaction due to Absorption of Phonons via
Quantum Interference - Single Phonon Detector (collaboration
with groups of Stefan Ludwig and Aash Clerk)*

*(2) Transferring electrons between distant qubits (collaboration
with group of Gloria Platero)*

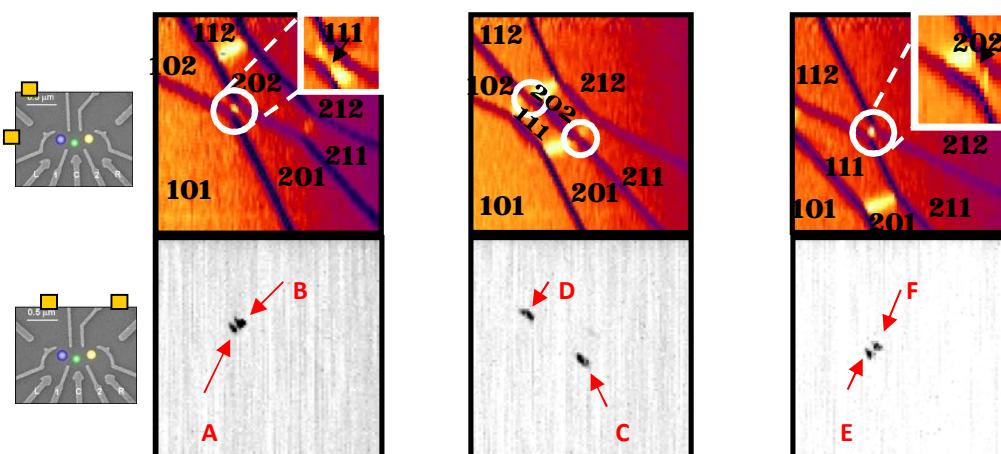
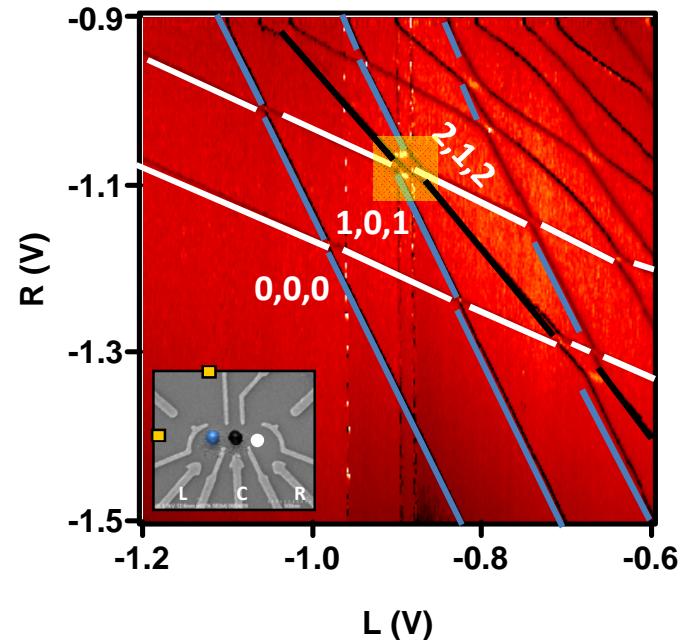
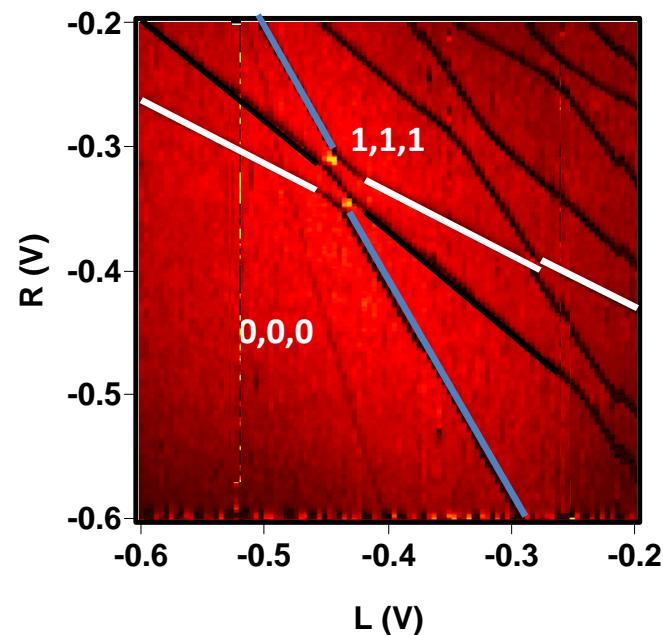
Quantum Superpositions overcoming Spin Blockade

*(3) Interplay between 3 spin Qubits
(LZS and all exchange)*



Stability Diagram / Transport Regimes

L. Gaudreau et al. APL 99 193101 (2009)



G.Granger et al., PRB 82, 075304 (2010).

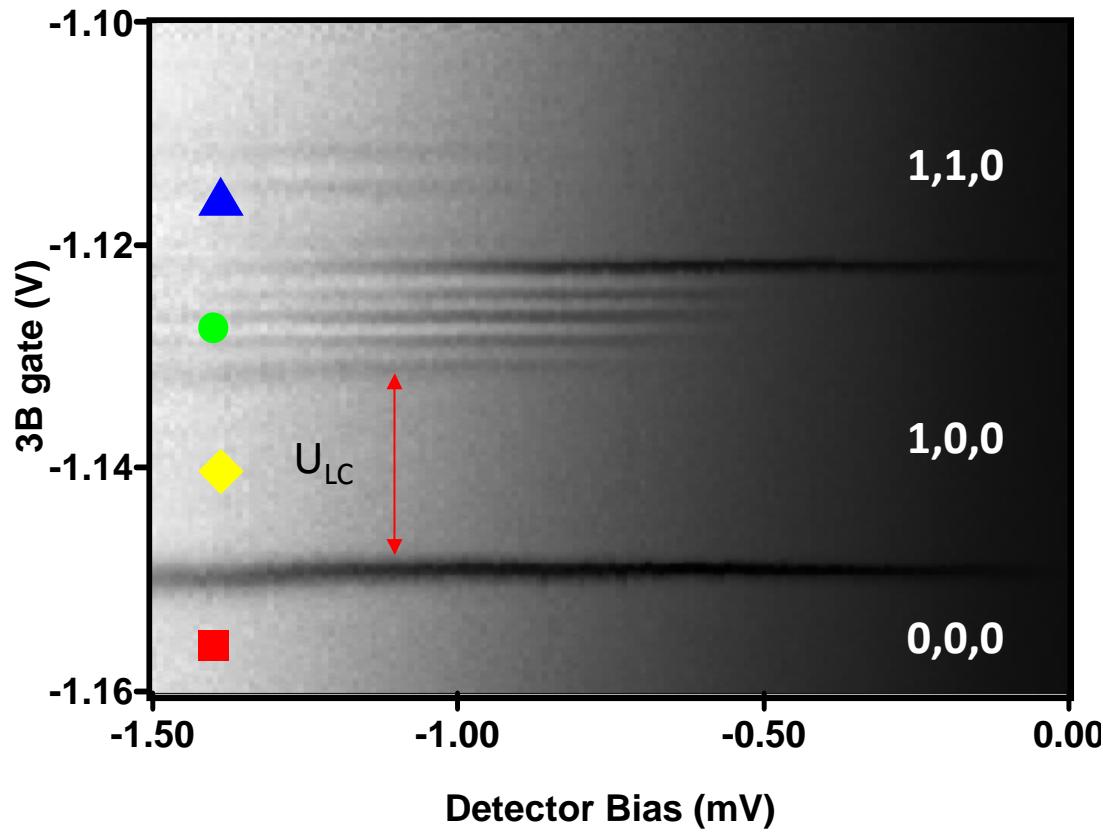
Predicted by M. Rogge, R. Haug. New Journal of Physics, 11 113037 (2009)

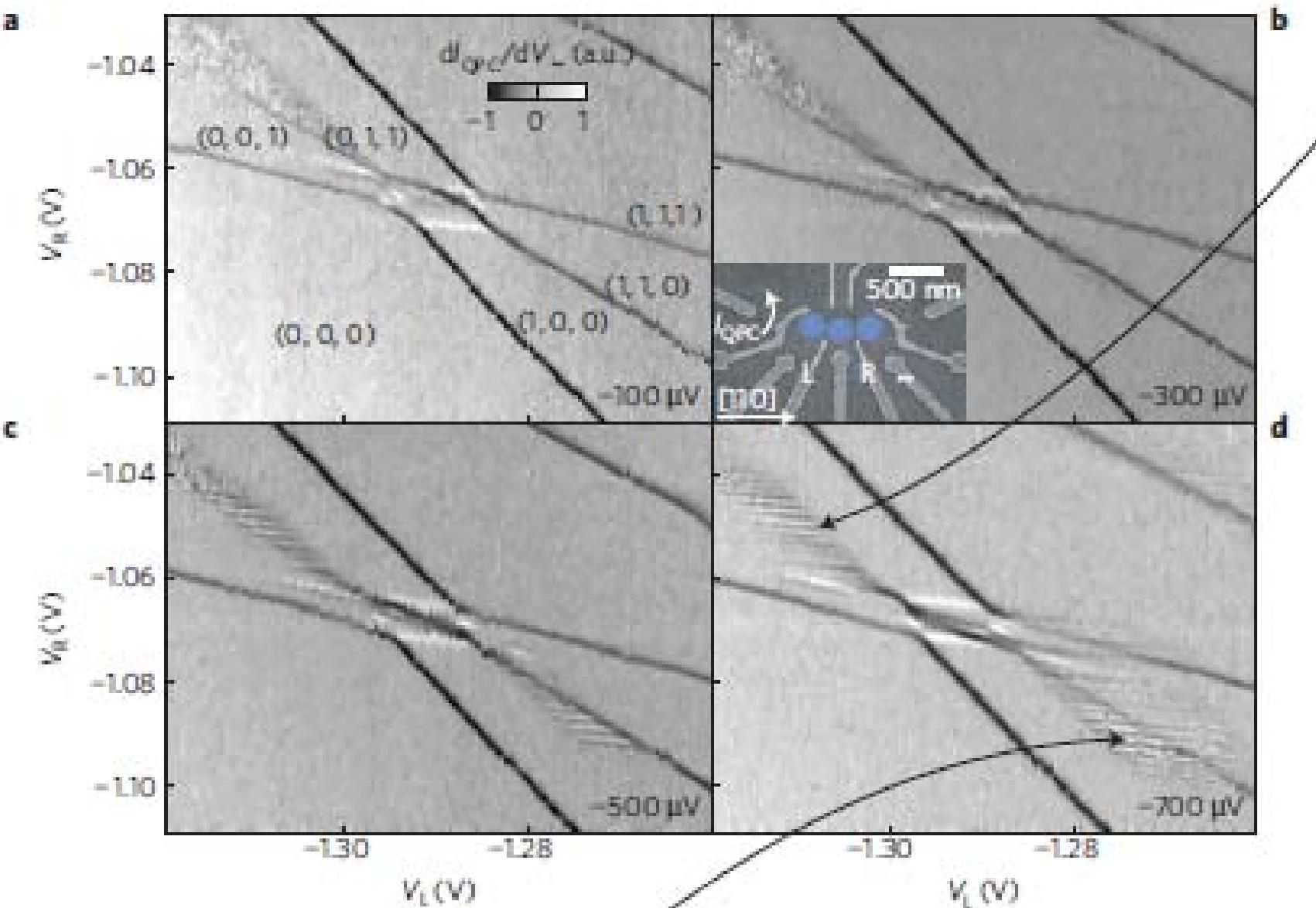
Why not keep increasing bias across QPC

Work in collaboration with Stefan Ludwig and Aash Clerk



Granger et al. Nature Physics 8, pp. 522-527 (2012)

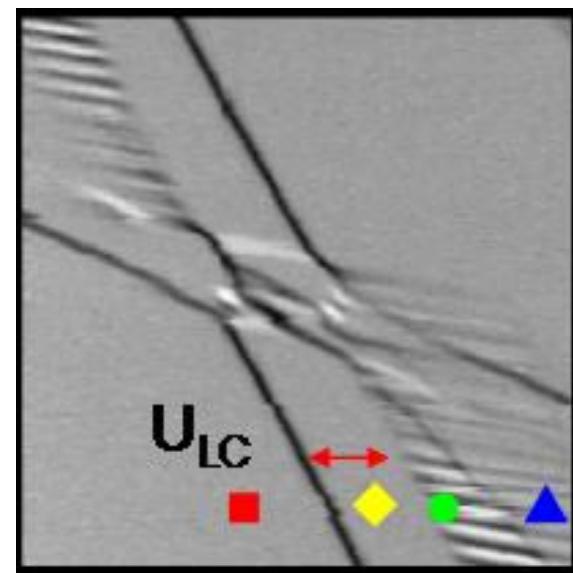
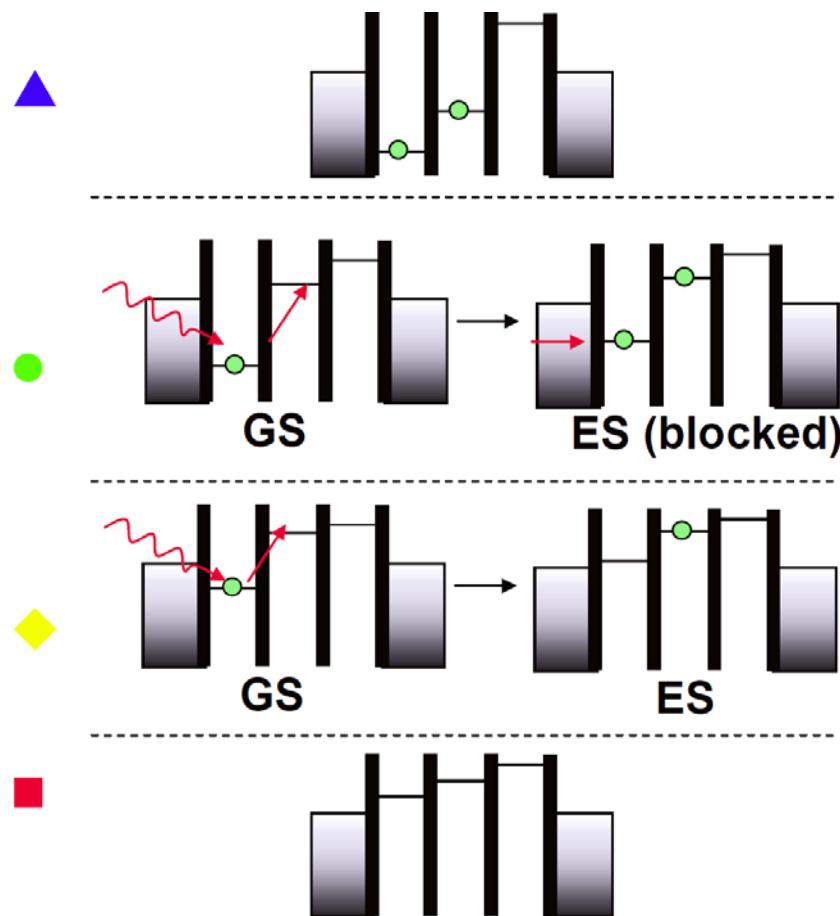


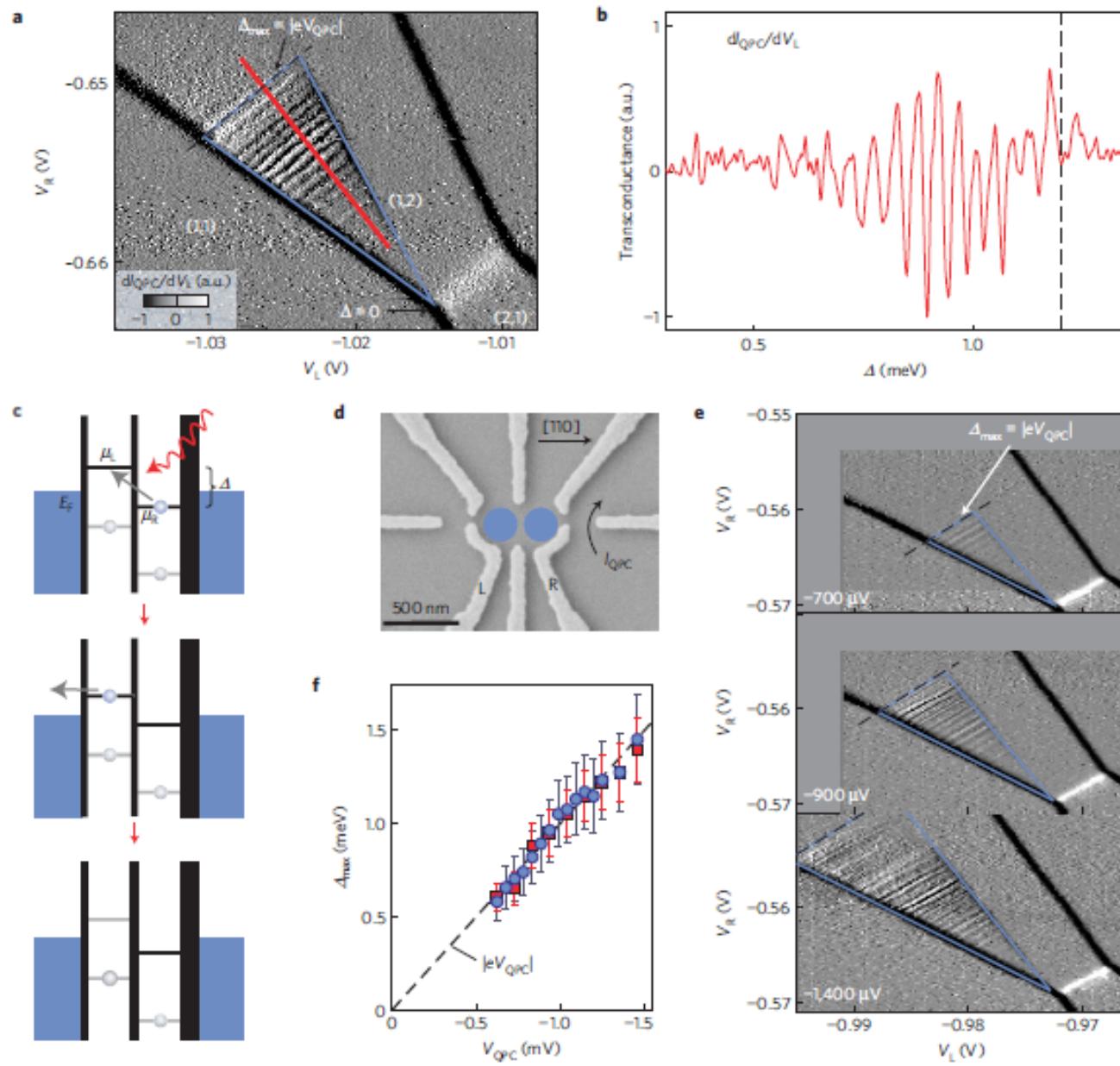


Harbusch, D et al. Phys. Rev. Lett. 104, 196801 (2010).

Taubert, D. et al. Phys. Rev. Lett. 100, 176805 (2008).

Young, C. E. & Clerk, A. A. Phys. Rev. Lett. 104, 186803 (2010).

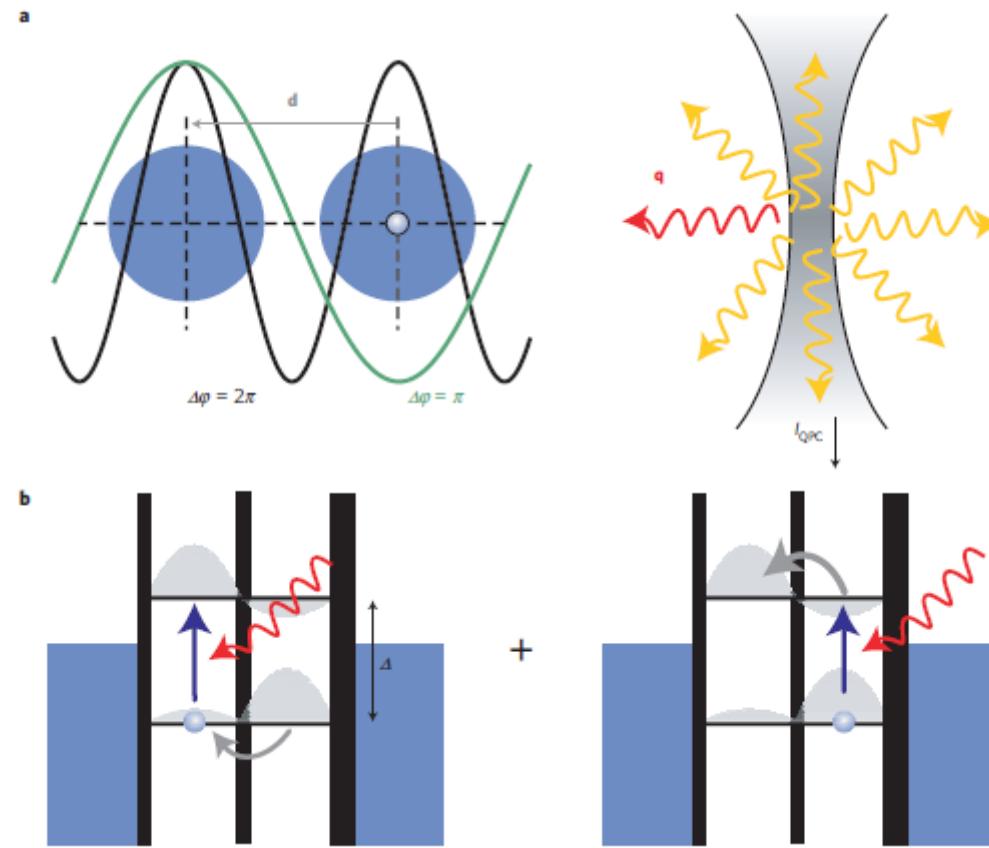




Why Periodic Pattern?

Why Stronger than emission equivalent?

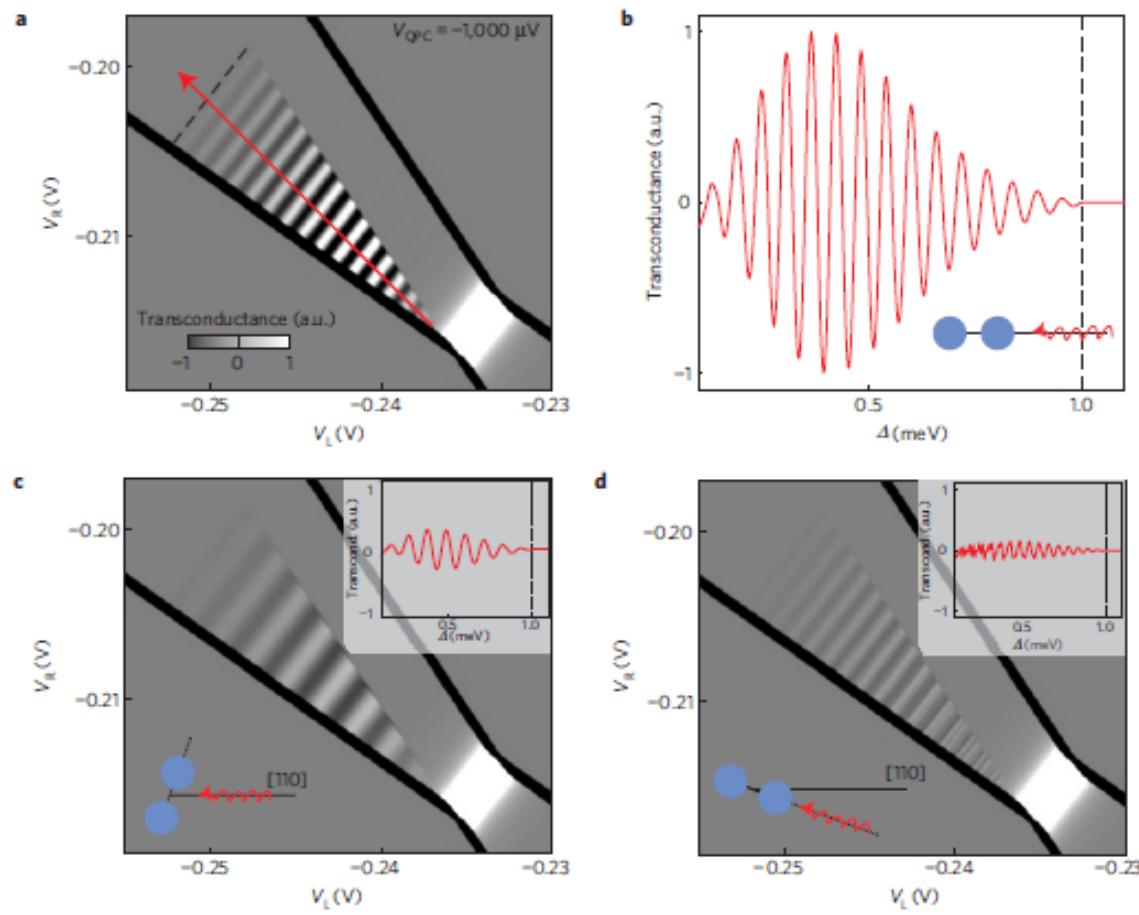
Equivalent Emission Mechanism
Fujisawa, T. et al. Science 282, 932935 (1998).
Rouleau, P. et al. Nature Commun. 2, 239 (2011)



$$\Delta = \hbar|q|v_{ph}$$

$$\Delta\phi = d \cdot q$$

$$\Delta\phi = (2n+1)\pi \quad \text{Constructive Interference}$$

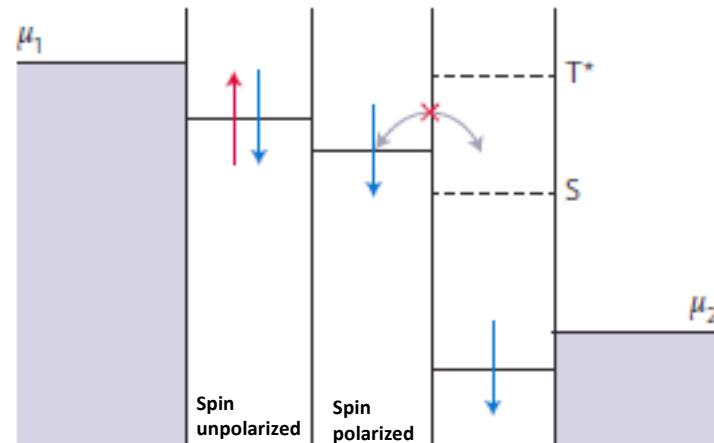


$$\hat{H}_{\text{int}} = \frac{t_c}{\Delta} \sum_{\mathbf{q}, \mu} \lambda_{\mathbf{q}, \mu} (e^{i\mathbf{q} \cdot \mathbf{r}_L} - e^{i\mathbf{q} \cdot \mathbf{r}_R}) (\hat{a}_{\mu, \mathbf{q}} + \hat{a}_{\mu, -\mathbf{q}}^\dagger) (|g\rangle \langle e| + \text{h.c.})$$

Bipolar spin blockade in the serial triple quantum dot

c

e.g. (1,1,1) to (2,1,1) to (2,0,2) to (1,1,2) to (1,1,1)

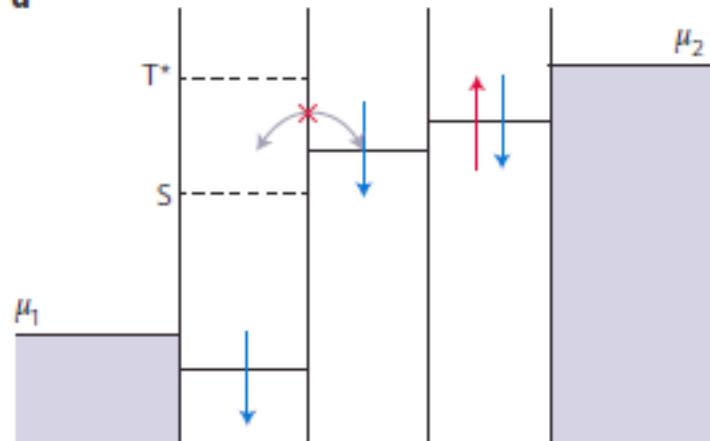


two spin blockade rectifiers placed back to back, forming a spinsulator

B=0

Due to the nuclear Overhauser fields that spins in neighboring dots experience, the singlet and triplet states of these spins are mixed, and spin blockade is lifted.

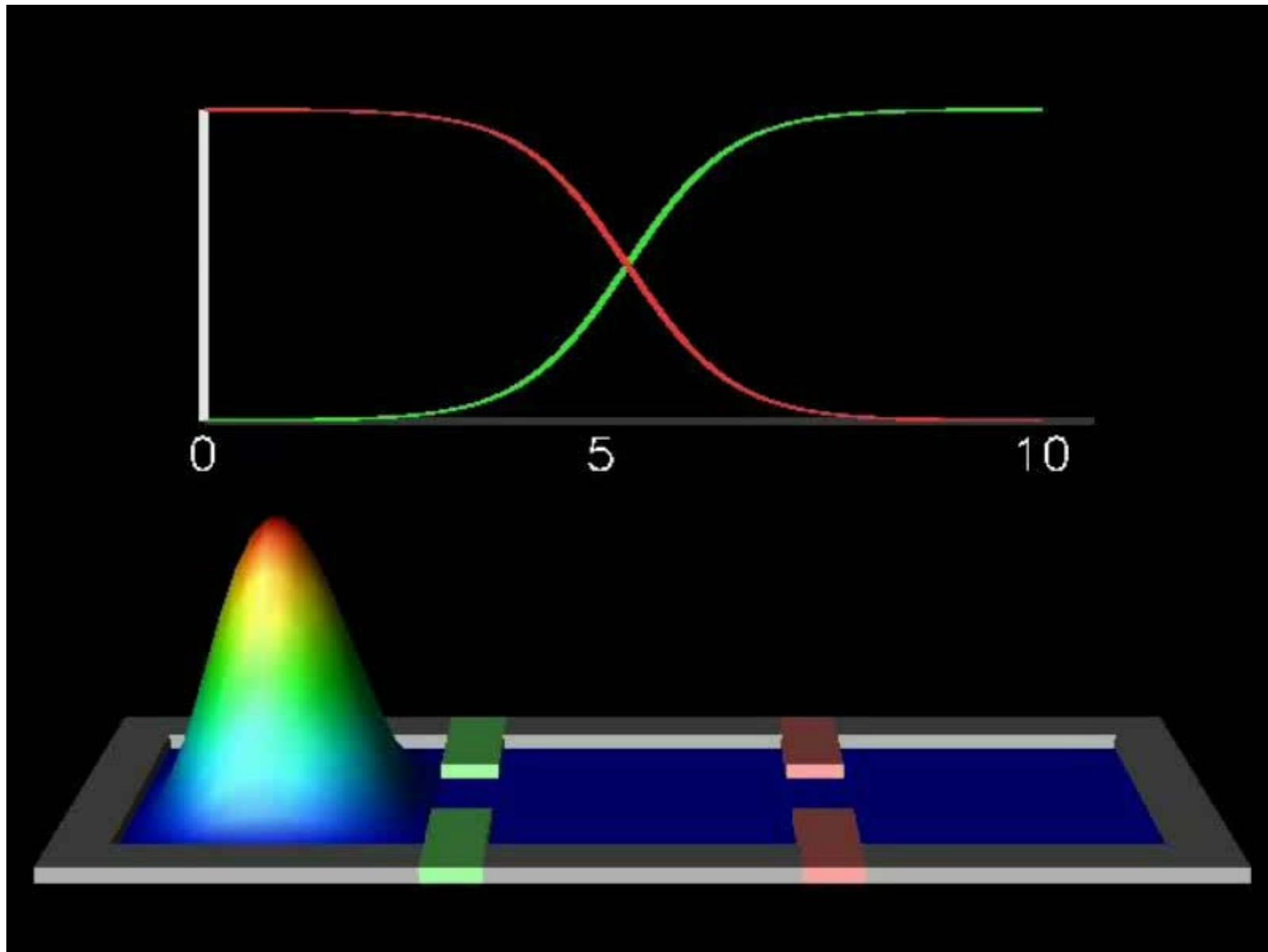
d



B>10mT

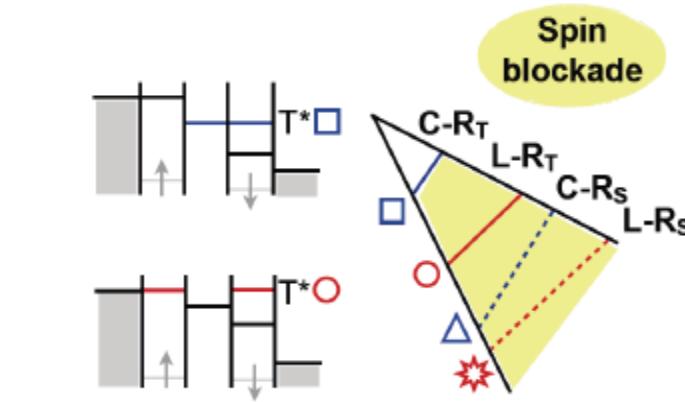
For an external magnetic field larger than the Overhauser field, triplet states with parallel spins provide unmixed spin blockade channels and the leakage current is suppressed.

A.D. Greentree, J.H. Cole, A.R. Hamilton, and L.C.L. Hollenberg, Phys. Rev. B 70, 235317 (2004).

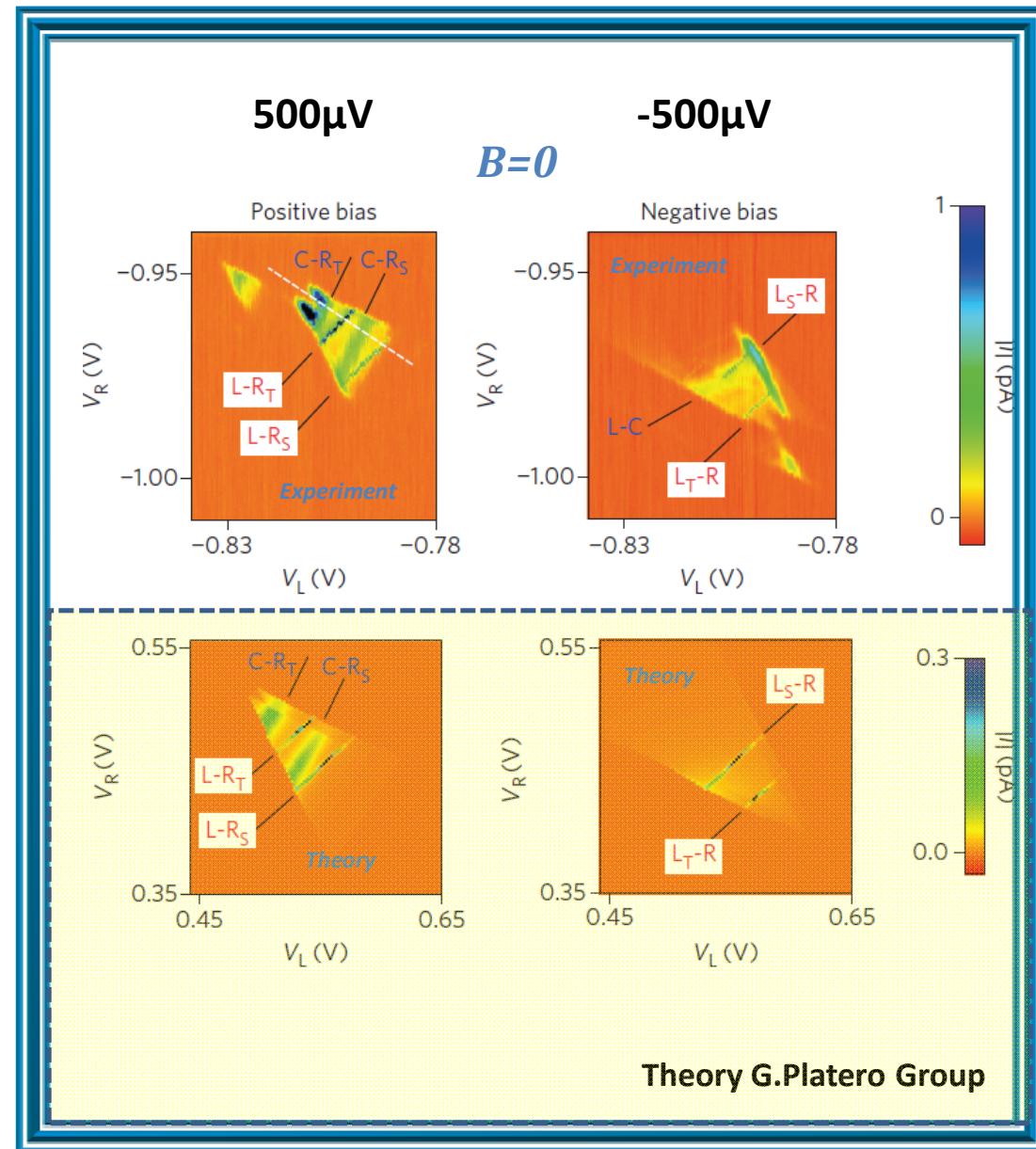
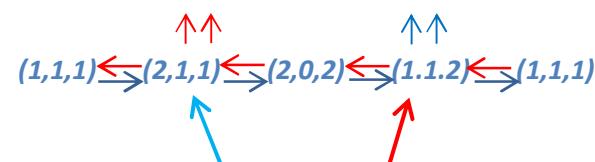


$$A|1\rangle + 0|2\rangle + C|3\rangle$$

Bipolar Spin Blockade (Spinsulator) and Leakage Paths at $B=0$



Sequential Transport Sequence

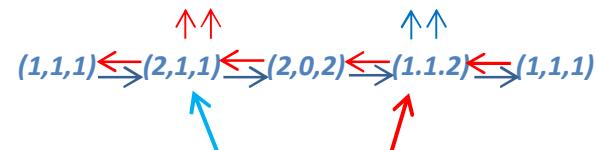


Bipolar Spin Blockade (Spinsulator) and Leakage Paths at $B=0.2T$

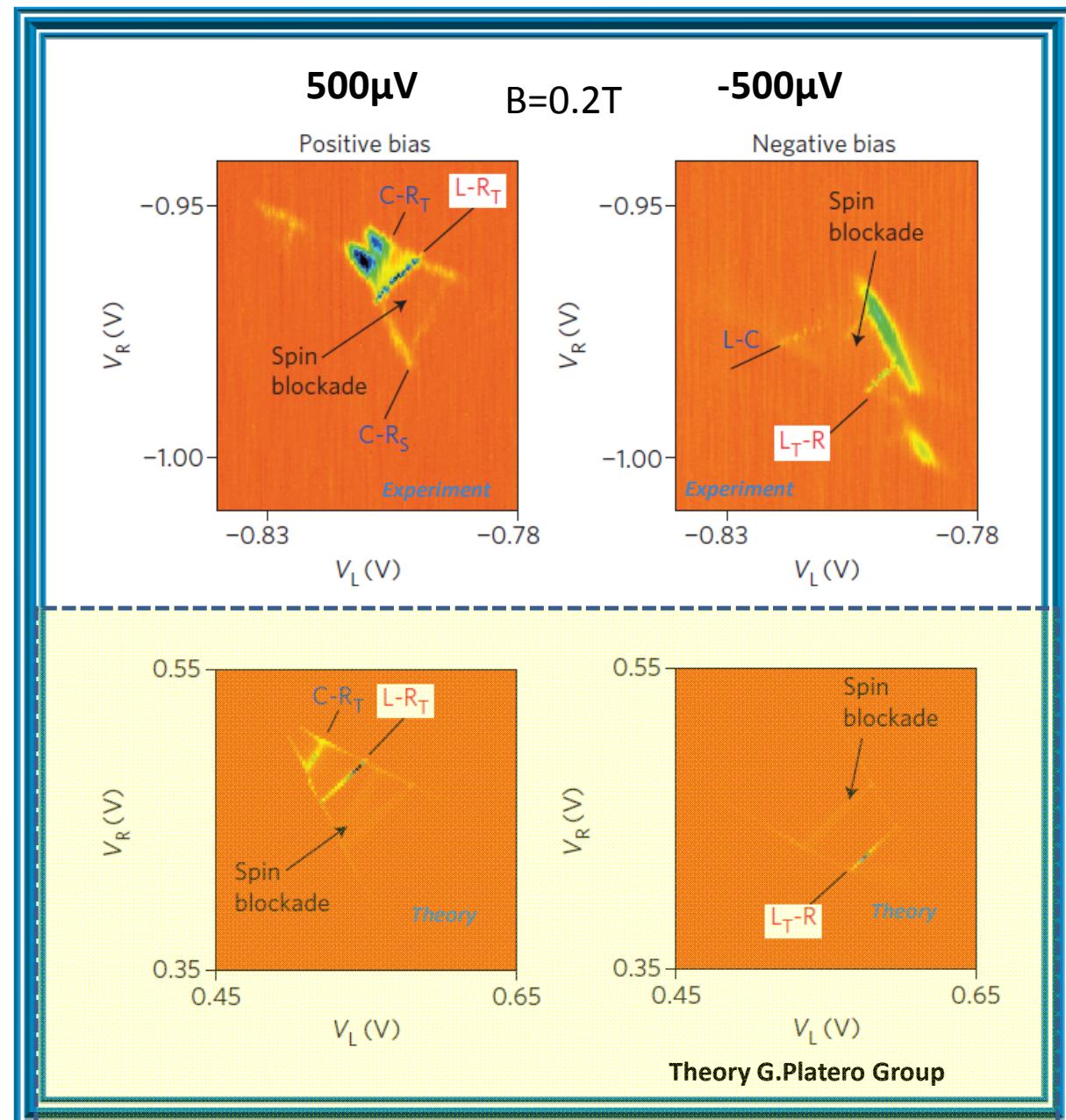
Quadruple Point

$(1,1,1), (2,1,1), (2,0,2) (1,1,2)$

Sequential Transport Sequence

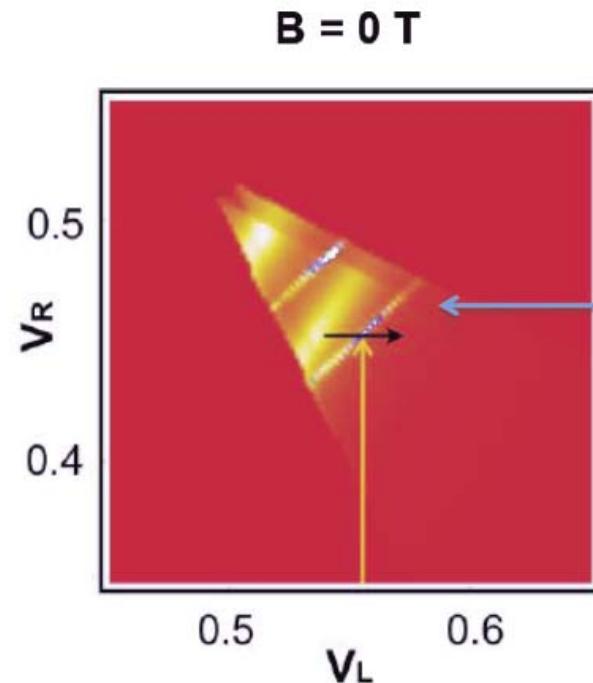
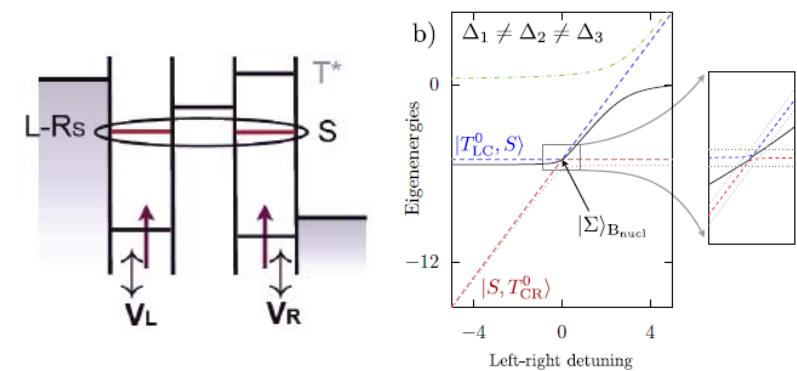


Spin Blockade

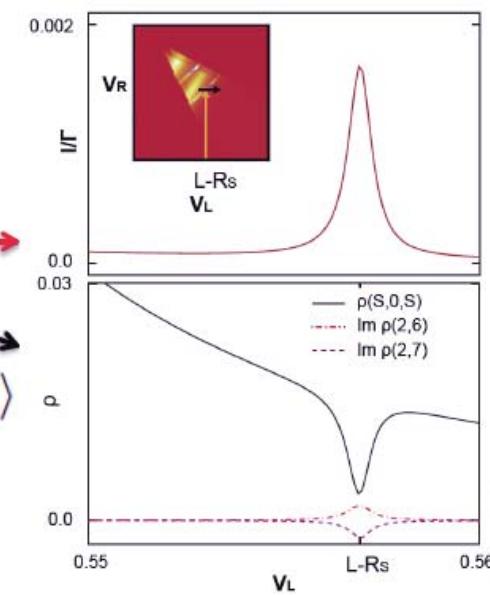


L-R_s resonance

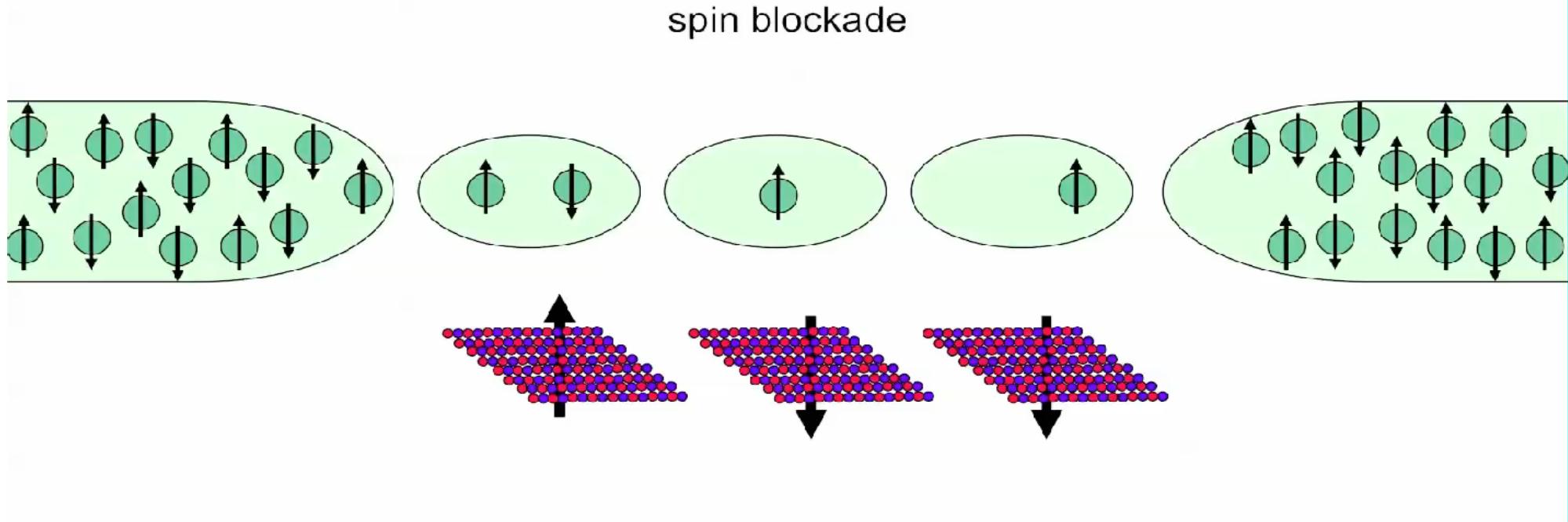
Off-resonance
 \downarrow
 $(2,1,1) \rightarrow (2,0,2) \rightarrow (1,1,2)$



High current peak →
without participation of
The intermediate state →
 $|\uparrow\downarrow, 0, \uparrow\downarrow\rangle = |S, 0, S\rangle$



Spin blockade and Coherent Superpositions

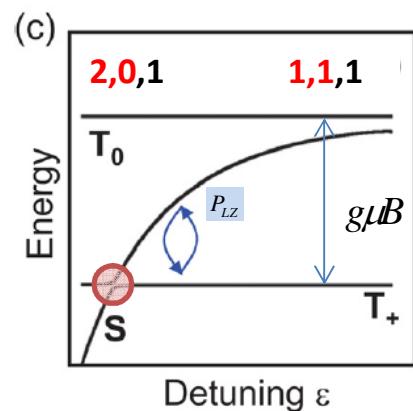


coherent superposition

$$|\Sigma_2\rangle = \frac{1}{2} (|\uparrow\downarrow, \downarrow, \uparrow\rangle - |\uparrow\downarrow, \uparrow, \downarrow\rangle - |\downarrow, \uparrow, \uparrow\downarrow\rangle + |\uparrow, \downarrow, \uparrow\downarrow\rangle)$$

$|\uparrow\downarrow, 0, \uparrow\downarrow\rangle = |S, 0, S\rangle$ not involved

Two Spin Landau-Zener-Stückelberg Oscillations



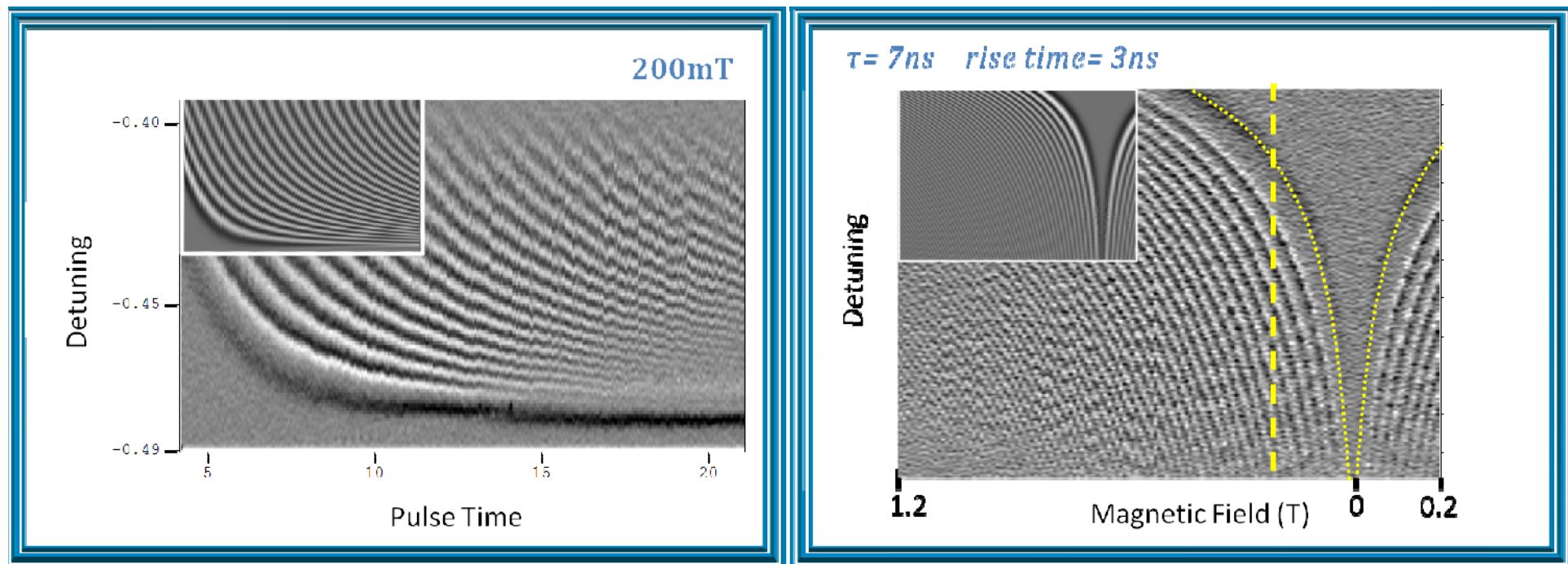
$$P_{LZ} = e^{-2\pi\Delta^2/\hbar v},$$

$$v = \frac{d(E_+ - E_-)}{dt}$$

$$|S\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$



$$|T+\rangle = |\uparrow\uparrow\rangle$$



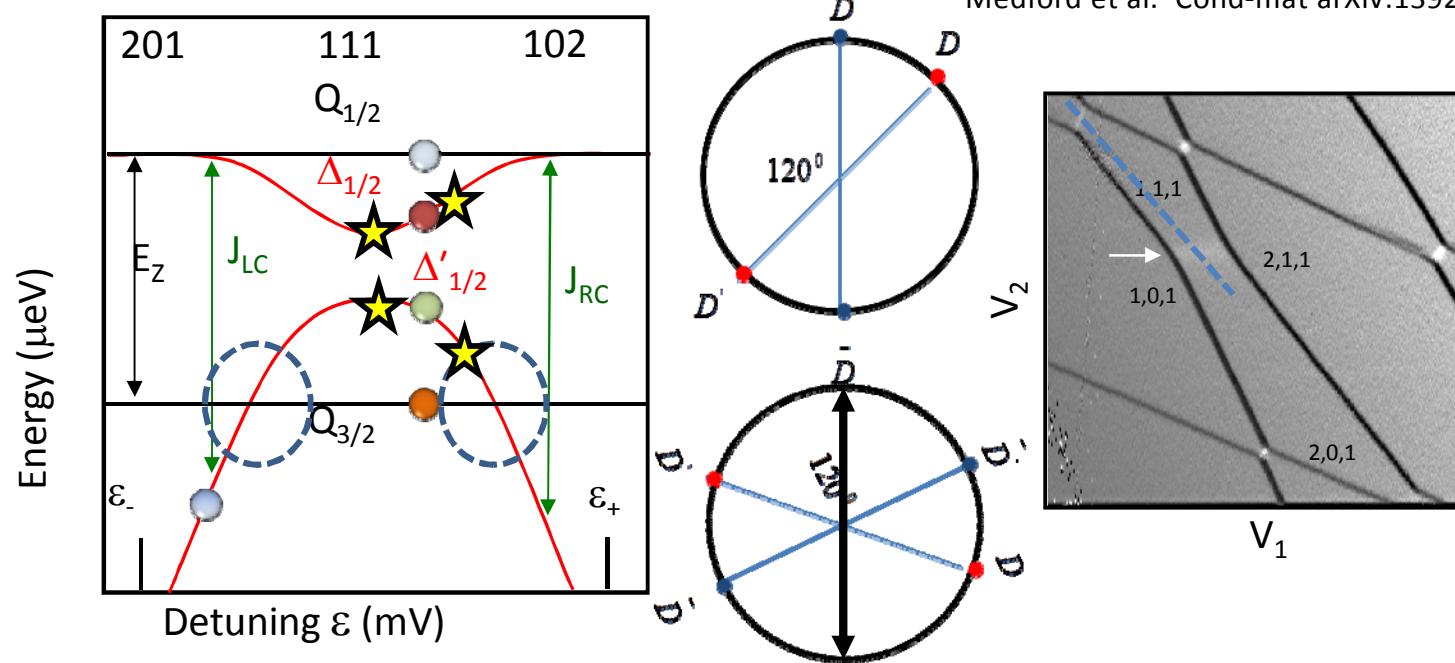
Measurements Made with the Harvard Protocol

Petta et al., Science 309, (2005) 2180

Petta et al., Science 327, (2010) 669

DiVincenzo et al. Science 408 (2000) 339

Medford et al. Cond-mat arXiv:1392.1933



$$|Q_{1/2}\rangle = \frac{1}{\sqrt{3}}(|\uparrow\uparrow\downarrow\rangle + |\uparrow\downarrow\uparrow\rangle + |\downarrow\uparrow\uparrow\rangle)$$

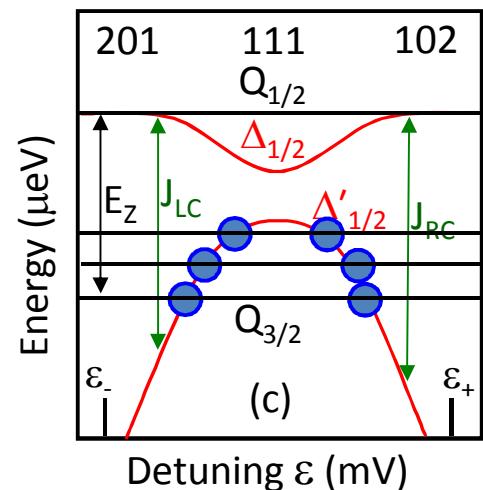
$$|\Delta_{1/2}\rangle = \frac{1}{\sqrt{4\Omega^2 + 2\Omega(J_{LC} - 2J_{RC})}}((J_{LC} - J_{RC} + \Omega)|\uparrow\uparrow\downarrow\rangle + (J_{RC} - \Omega)|\uparrow\downarrow\uparrow\rangle - J_{LC}|\downarrow\uparrow\uparrow\rangle)$$

$$|Q_{3/2}\rangle = |\uparrow\uparrow\uparrow\rangle$$

$$|\Delta'_{1/2}\rangle = \frac{1}{\sqrt{4\Omega^2 + 2\Omega(2J_{RC} - J_{LC})}}((-J_{LC} + J_{RC} + \Omega)|\uparrow\uparrow\downarrow\rangle - (J_{RC} + \Omega)|\uparrow\downarrow\uparrow\rangle + J_{LC}|\downarrow\uparrow\uparrow\rangle)$$

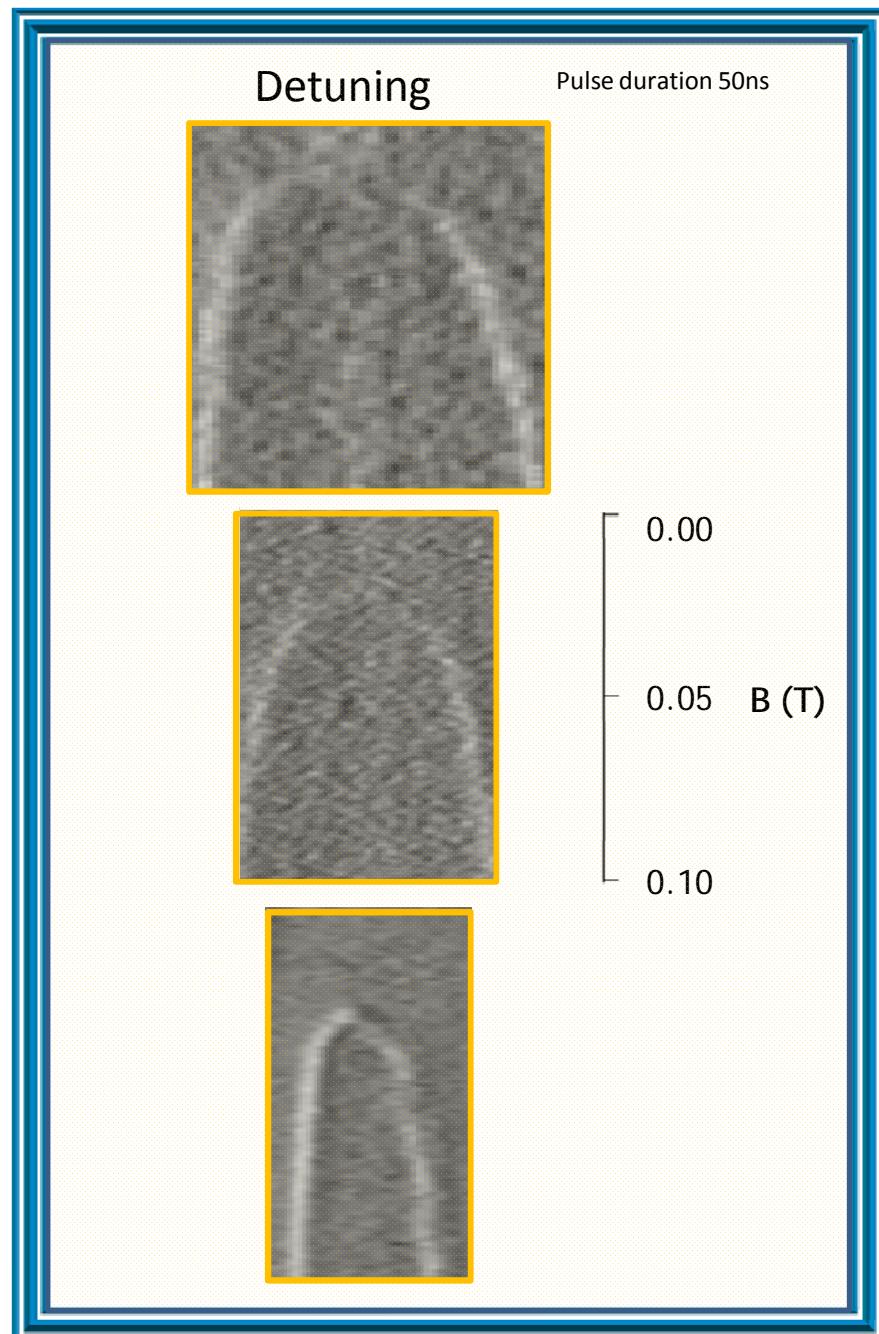
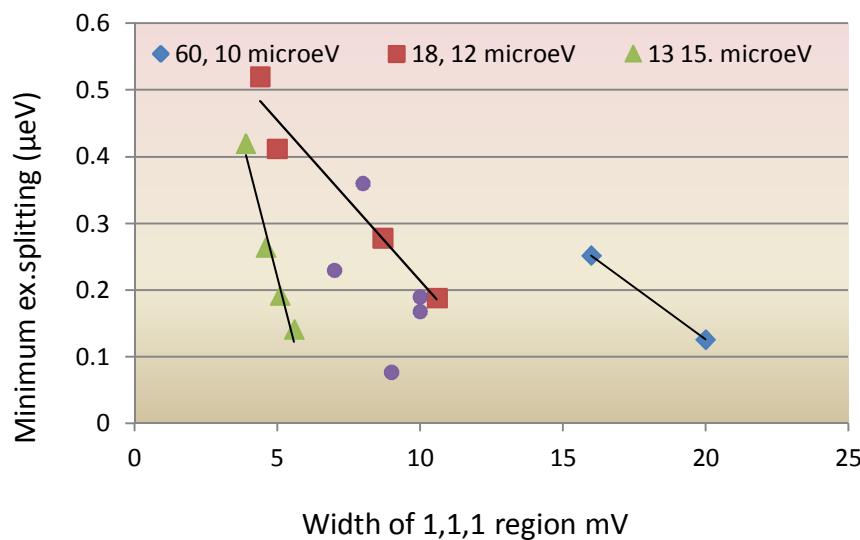
$$\Omega = \sqrt{J_{LC}^2 + J_{RC}^2 - J_{LC}J_{RC}}$$

Spin Arch – back to back spin funnels



$$T_{LC}(\epsilon) = \begin{cases} T_{LC}\exp[C_{LC}(\epsilon - \epsilon_+)], & \epsilon < \epsilon_+ \\ T_{LC}, & \epsilon \geq \epsilon_+ \end{cases}$$

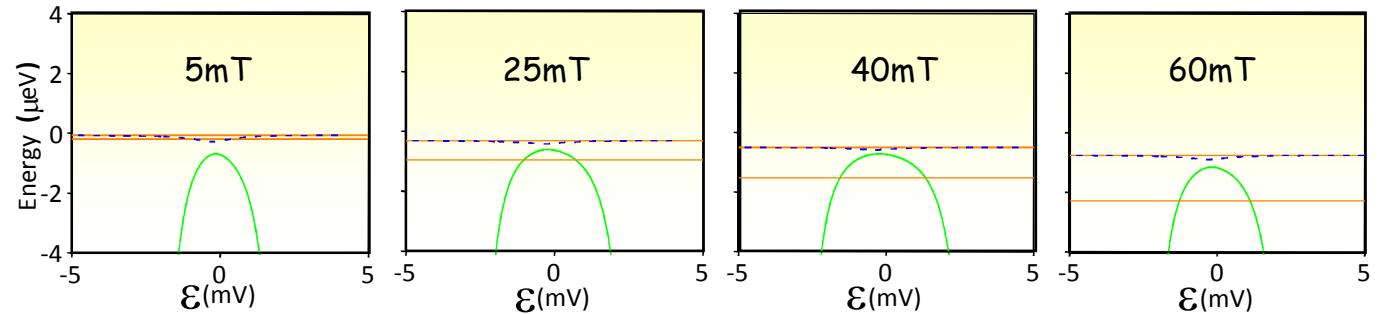
$$T_{RC}(\epsilon) = \begin{cases} T_{RC}\exp[C_{RC}(\epsilon_- - \epsilon)], & \epsilon > \epsilon_- \\ T_{RC}, & \epsilon \leq \epsilon_- \end{cases}$$



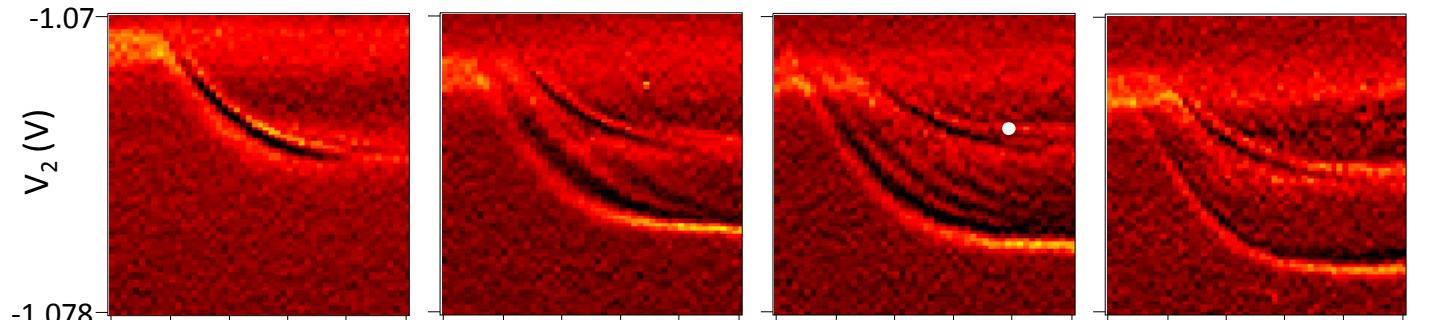
Three Spin LZS

Rise time ~ 6.6 ns

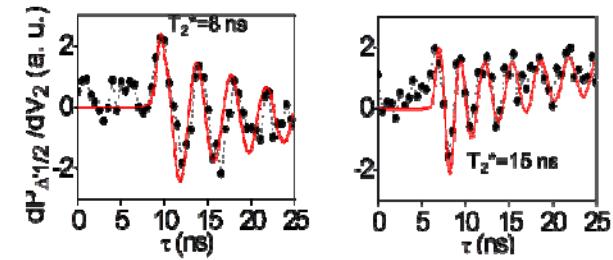
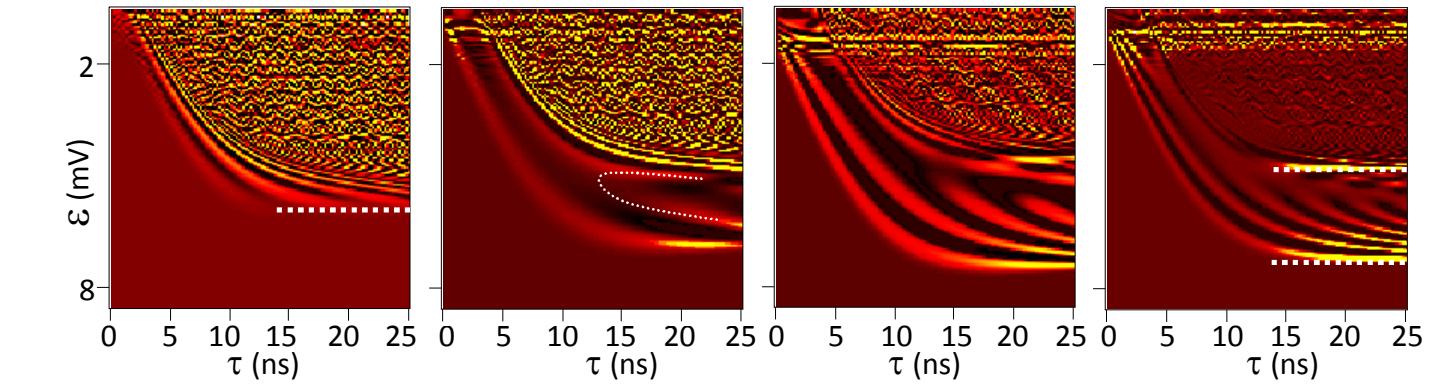
Level spectrum



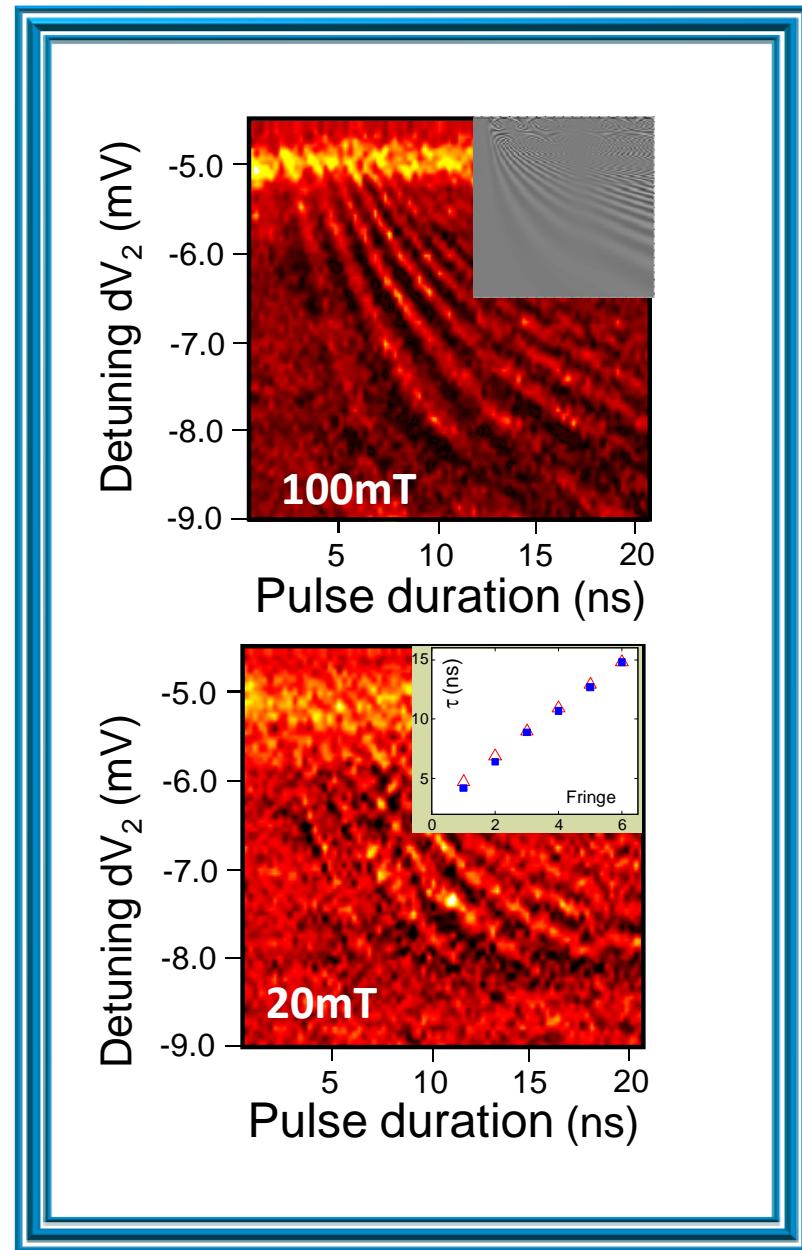
Experiment



Theory

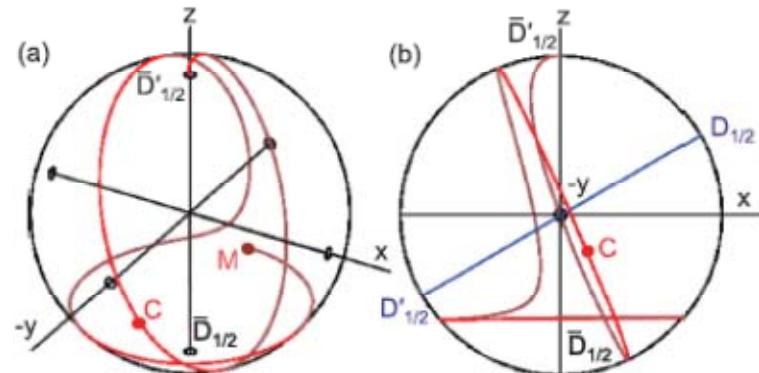
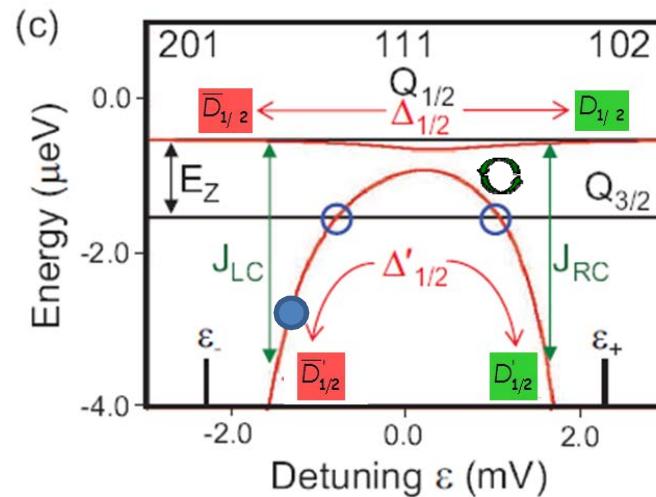


Exchange oscillations - 1ns rise time



$$|D_{1/2}\rangle = \frac{-1}{\sqrt{6}}(2|\uparrow\uparrow\downarrow\rangle - |\uparrow\downarrow\uparrow\rangle - |\downarrow\uparrow\uparrow\rangle),$$

$$|D'_{1/2}\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\uparrow\rangle - |\downarrow\uparrow\uparrow\rangle).$$



G.C.Aers et al. Phys. Rev. B 86, p. 45316 (2012)

DiVincenzo et al. Science 408 (2000) 339
 Laird et al. PRB 82 075403 (2010)
 Medford et al. Cond-mat arXiv:1392.1933

Summary

(1) Spin Blockade Leakage via Coherent Superpositions

M.Busl et al Nature Nanotechnology –on line February 2013

(2) Quantum Backaction related to coherent absorption of phonons

Granger et al. *Nature Physics* 8, pp. 522-527 (2012)

(3) Holistic View of Possible Coherent Behaviours with 3 spins in Triple Quantum Dots

Gaudreau et al. *Nature Physics* 8, 54–58 (2012)

Studenikin et al. *Phys. Rev. Lett.* 108, (2012) 226802

G.C.Aers et al. *Phys. Rev. B* 86, p. 45316 (2012)

Acknowledgements

